



Digital FDIRC

a Focused Differential
Internal Reflection Cherenkov
imaged by SiPM arrays

Pier Simone Marrocchesi

Univ. of Siena and INFN Pisa

FRONTIER DETECTORS FOR FRONTIER PHYSICS

13° Pisa Meeting on Advanced Detectors

24-30 May 2015 - *La Biodola, Isola d'Elba*

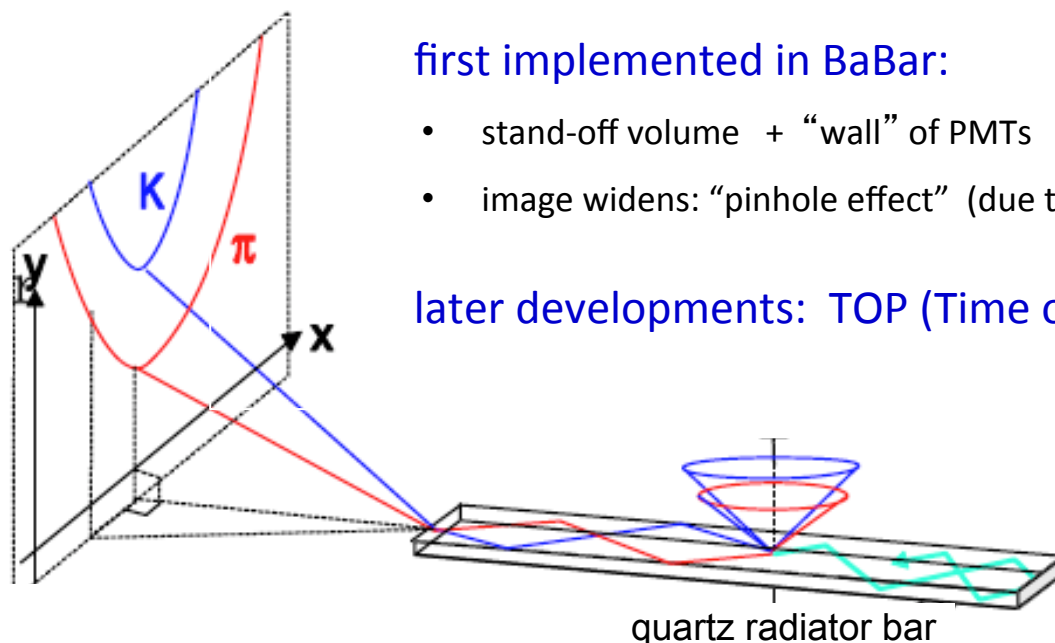
A

DIRC: Detection of Internally Reflected Cherenkov light

first implemented in BaBar:

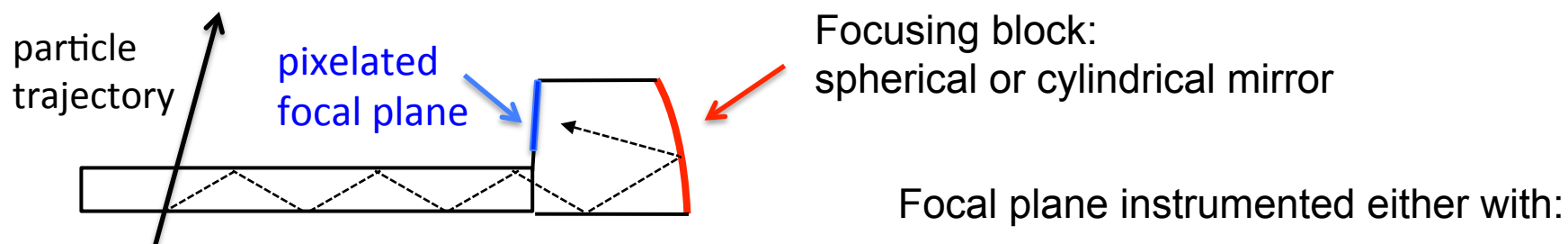
- stand-off volume + “wall” of PMTs
- image widens: “pinhole effect” (due to finite width of quartz bar)

later developments: TOP (Time of Propagation), etc...



B

A focalization scheme is introduced: FOCUSING DIRC (FDIRC)

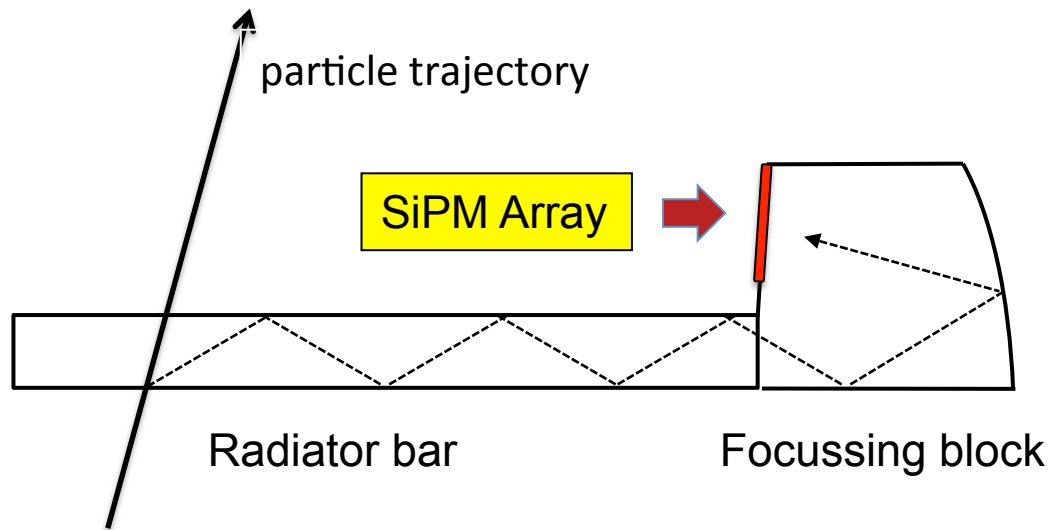


- ❑ pixelated photosensors (e.g. MAPMT, HPD, MCP):
 - Limited number of pixels
 - Large pixel size
- ❑ or fiber bundles coupled to pixelated photosensors → Low light trapping efficiency

this talk

Digital FDIRC

FDIRC with high resolution SiPM array



Focal plane instrumented with a **mosaic of SiPM arrays**:

- small pixels (mm^2 or sub- mm^2)
- large number of sensors ($\sim 10^3$)
- **photon counting**

Possible applications

High Energy Physics:

- PID @ a few GeV momenta (e.g.: π/K separation as in BABAR, BELLE, PANDA)

Astroparticle Physics:

- Charge identification of cosmic ray nuclei
- **isotopic separation** in cosmic rays (for space/balloon borne experiments)

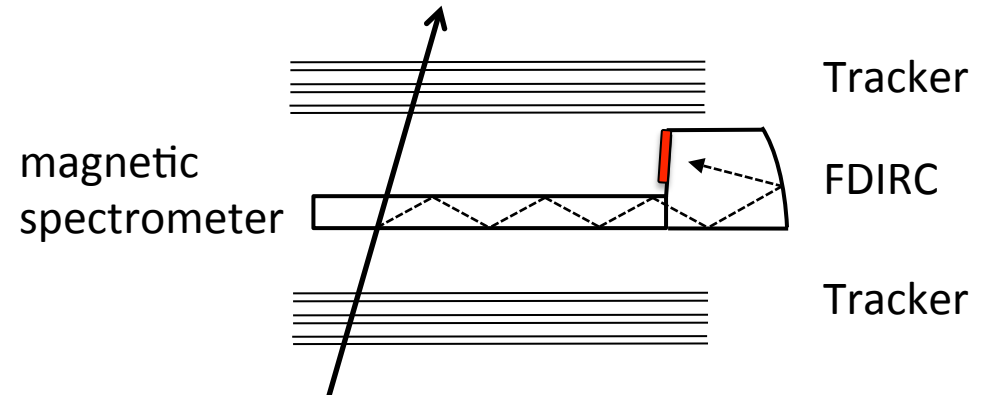
Isotopic separation by momentum + velocity (FDIRC) measurements

For **not too large** total-particle momenta P (\sim tens of GeV/c) an accurate measurement of β can provide an adequate **mass separation** for fully stripped ions of atomic number Z and mass A :
$$\left(\frac{\sigma_M}{M}\right)^2 = \left(\frac{\sigma_{P_T}}{P_T}\right)^2 + \left(\gamma^2 \frac{\sigma_\beta}{\beta}\right)^2$$

CONCEPT

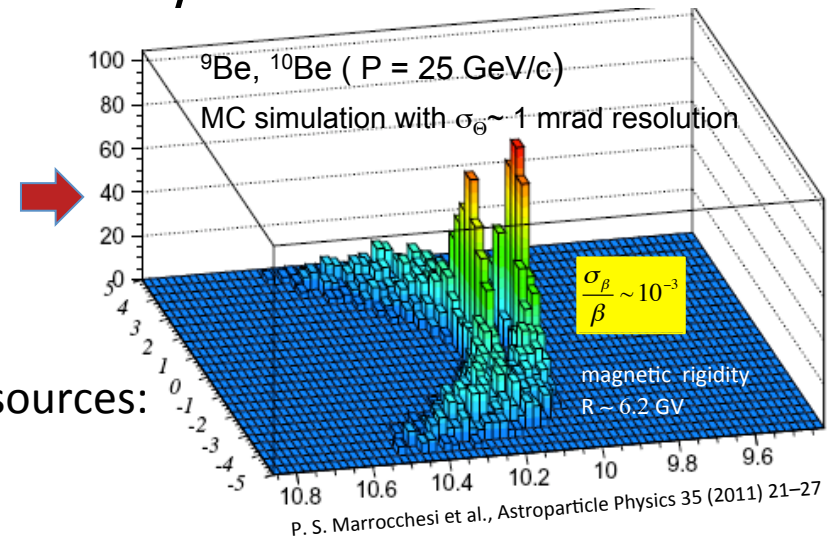
FDIRC embedded in a magnetic spectrometer:

- non destructive measurement of β
- track measured **before** and **after** FDIRC
- **large photostatistics** $\sim Z^2$ for **charged ions**



Example: $\Delta\Theta_c$ difference in Cherenkov angles for ${}^9\text{Be}$, ${}^{10}\text{Be}$ ions with $P = 25$ GeV/c is $\Delta\Theta_c \sim 11$ mrad

A mass separation better than 0.2 a.m.u. can be achieved with an angular resolution of $\sigma_\theta = 1.5$ mrad and $\sigma_p/p \sim 1\%$



Not easy to achieve: σ_θ is affected by several error sources:

- FDIRC {
- geometry/optics
 - **chromatic dispersion** along the optical path
 - bar imperfections (surface and angles)
 - pixel size, etc...



tracking angular error

FDIRC PROTOTYPE

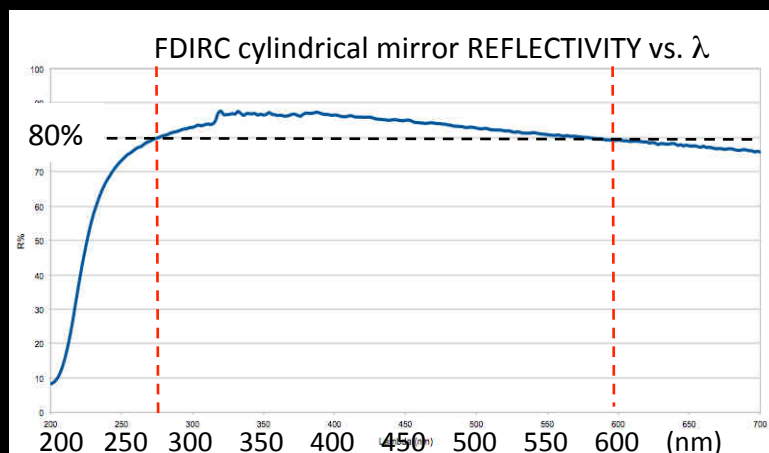
Thu Sep 4 14:04:13 2014

particle trajectory

Cylindrical mirror

radius = 26 cm

width = 14 cm; height ~ 16 cm

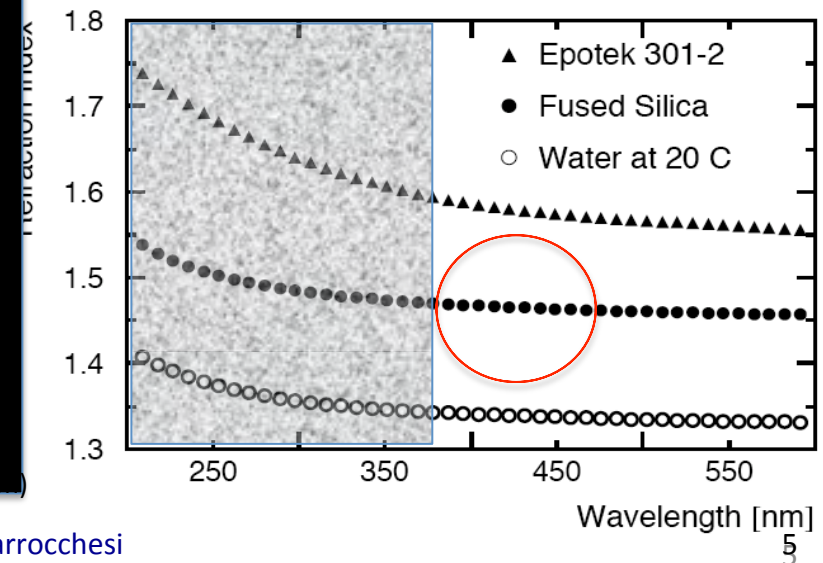


Quartz radiator bar

- Fused Silica (SiO_2) radiator bar
- 3 spare short bar segments from BaBar
- 17.25 mm (thickness) x 35 mm (width)
- 200 mm (long)

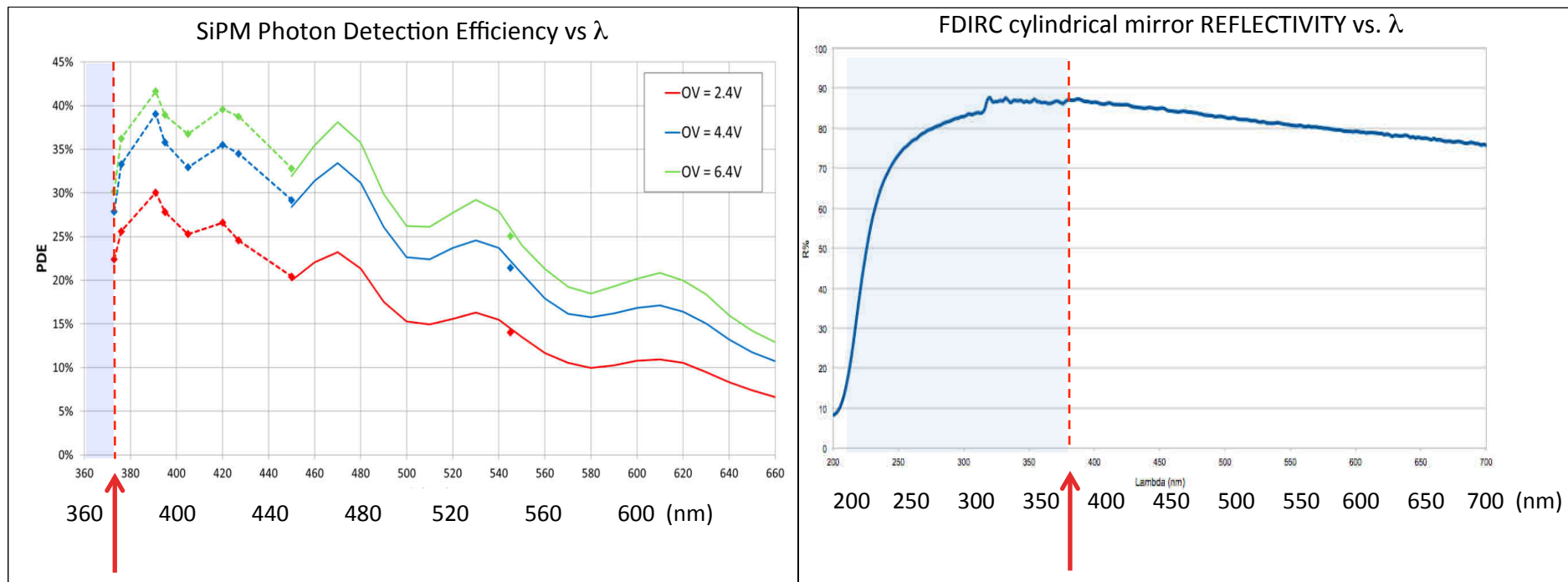
Dispersion relation in the radiator:

- above ~ 370 nm (**PDE cutoff**)
 $n(\lambda)$ is almost constant vs. λ
- $n \sim 1.47$ at 435 nm

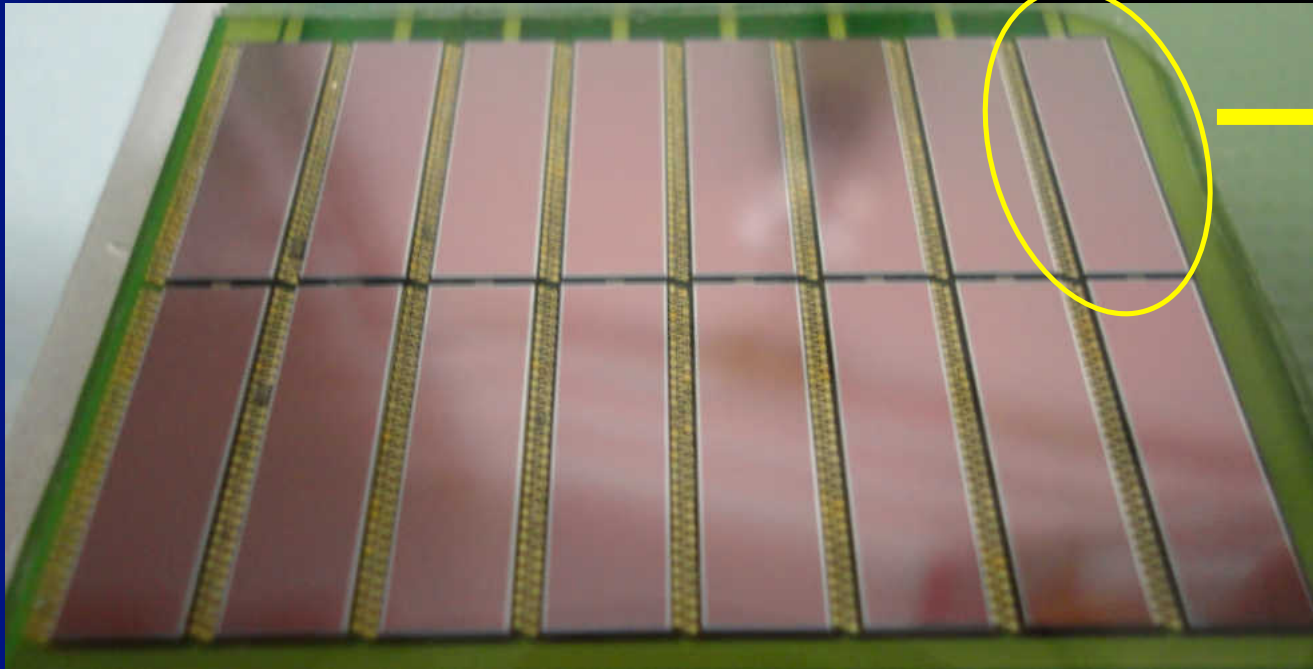


Chromatic dispersion vs. bandwidth

- Cherenkov light yield $\sim 1/\lambda^2 \rightarrow$ grows in the UV
- in SiO_2 radiator, the rate of change of the index of refraction vs. wavelength $dn(\lambda)/d\lambda$ is larger in the UV: angular resolution (chromatic term) gets worse
- tradeoff between light yield and bandwidth: in our case **mirror reflectivity** and **photon detection efficiency (PDE) of SIPM** limit the effective bandwidth to $\sim 370 < \lambda < 600 \text{ nm}$



FOCAL PLANE (4.3 cm x 2.7 cm)



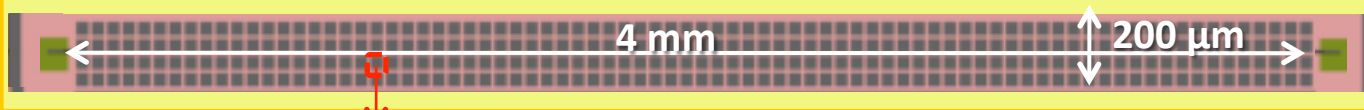
Array of 64 SiPMs

Array size
4.3 mm x 13.1 mm

Pixel size
4 mm x 0.2 mm

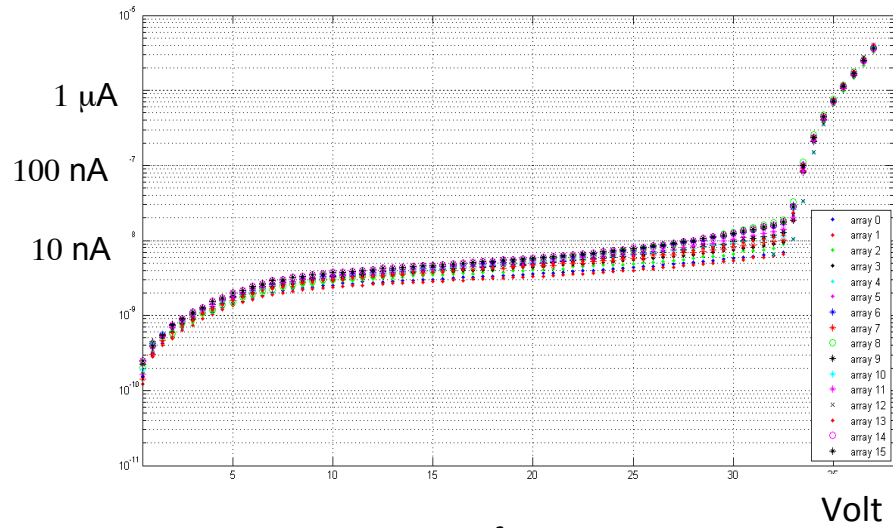
- Each **Array**: 64 pixels (SiPM sensors)
inter-pixel gap 10 μm
- Each **Pixel** : 4mm x 200 μm = 0.8 mm² SiPM
4 x 100 micro cells

**16 SiPM Arrays
(1024 SiPMs)**

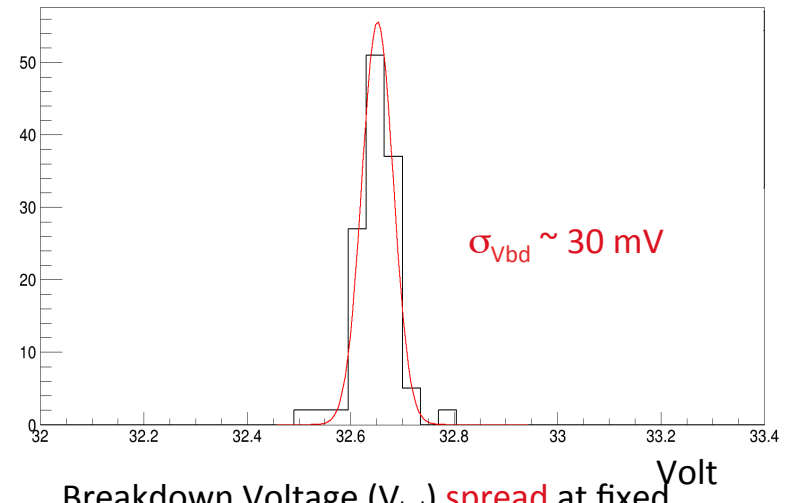


Micro cell
47.5 μm X 40 μm

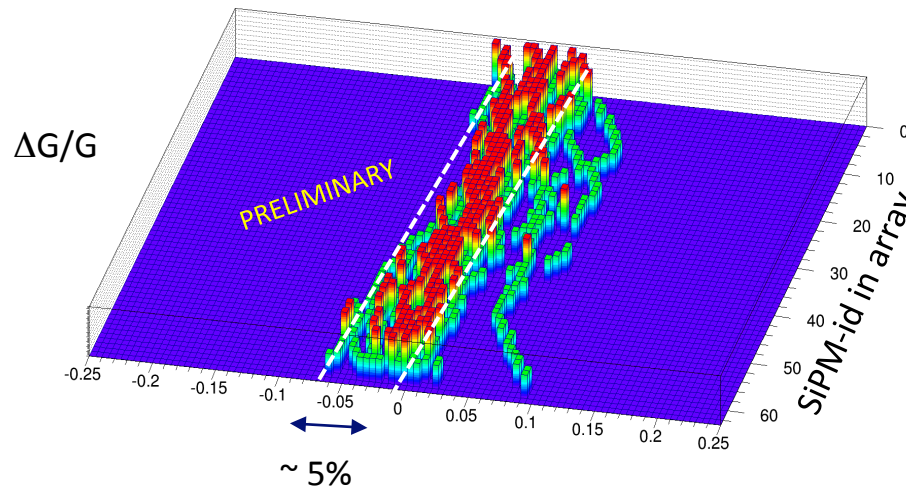
Characterization of 16 SiPM arrays (1024 SiPM)



I-V curves for 16 arrays



Breakdown Voltage (V_{bd}) spread at fixed temperature $T \sim 22^\circ\text{C}$ for 16 arrays

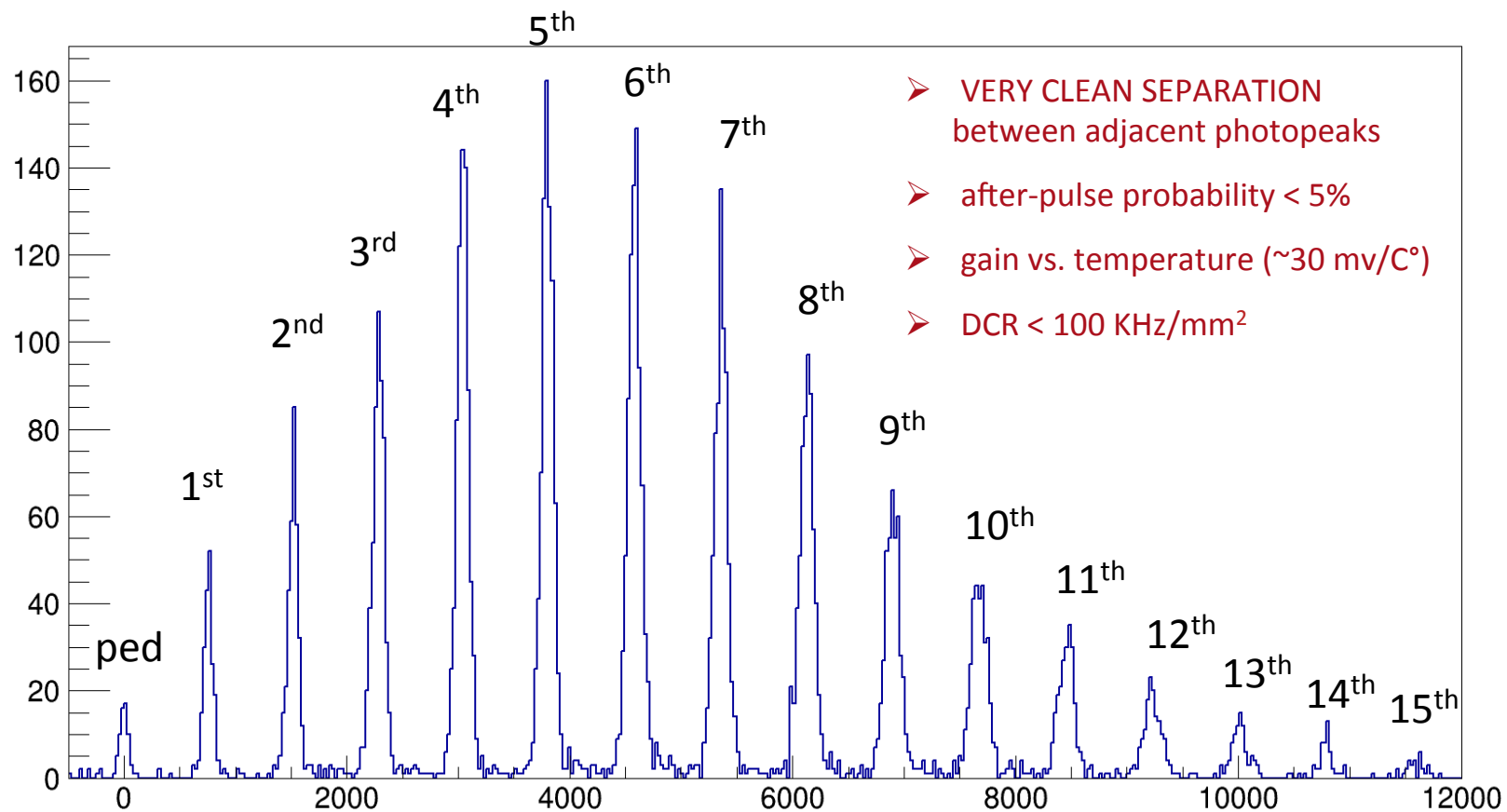


Gain spread $\Delta G/G$ at fixed temperature for $\sim 10^3$ channels (the gain spread of the FE electronics is FOLDED in)

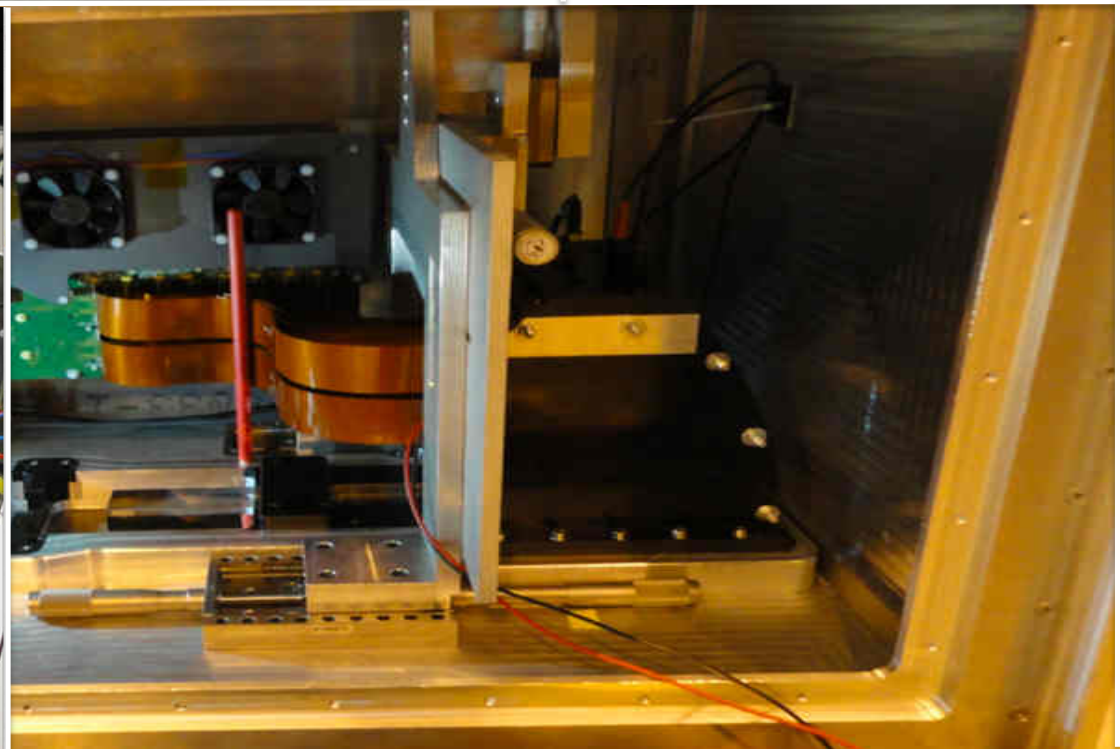
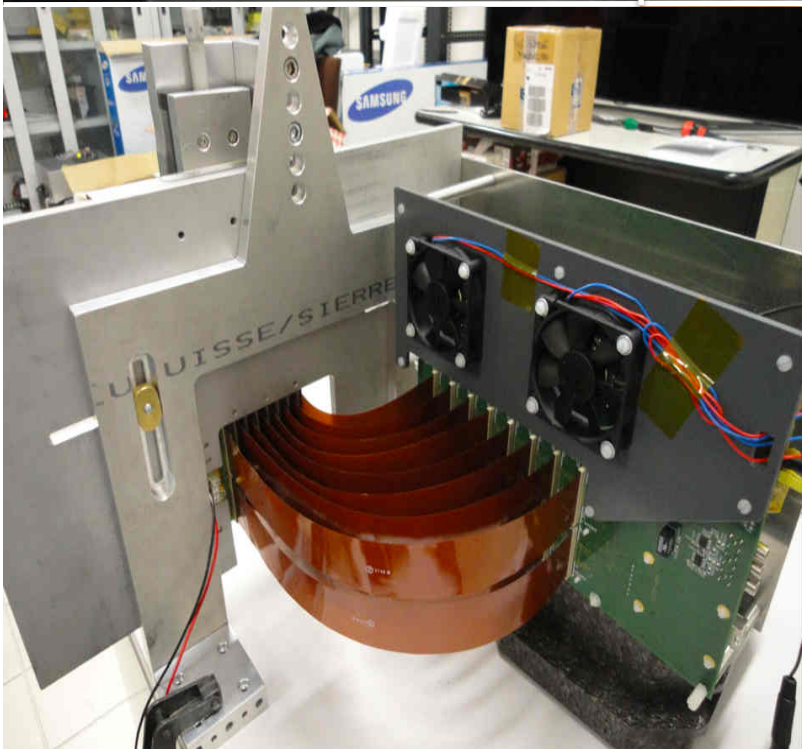
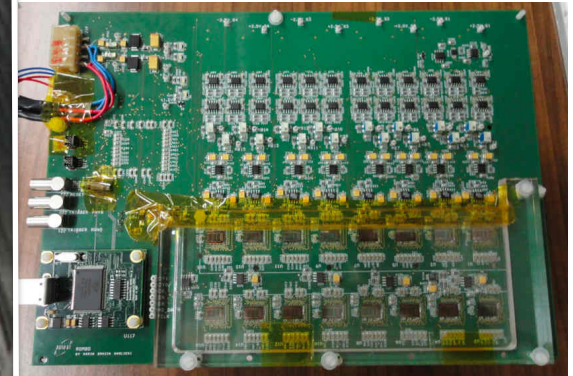
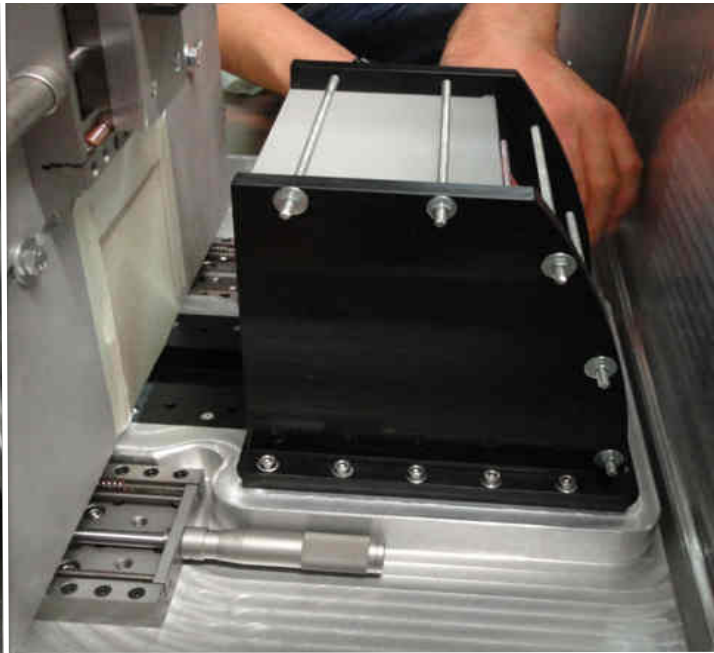
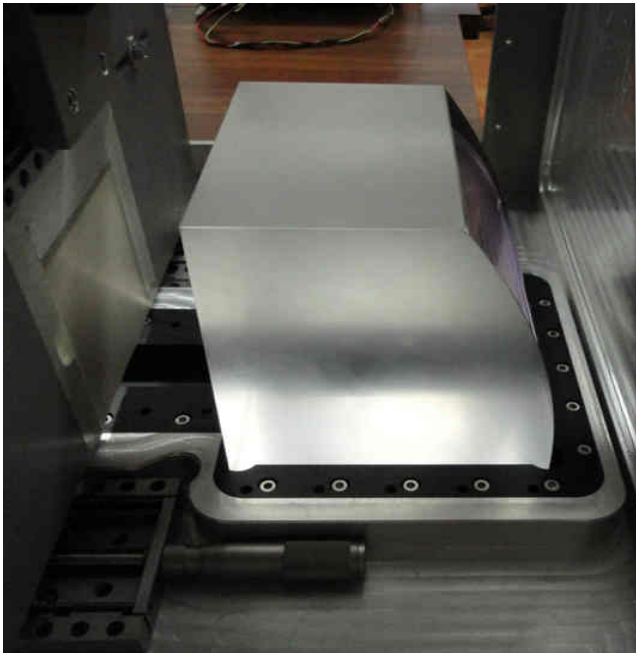
A real Photon Counting Device at room temperature !

NUV-SiPM arrays developed at FBK-Trento

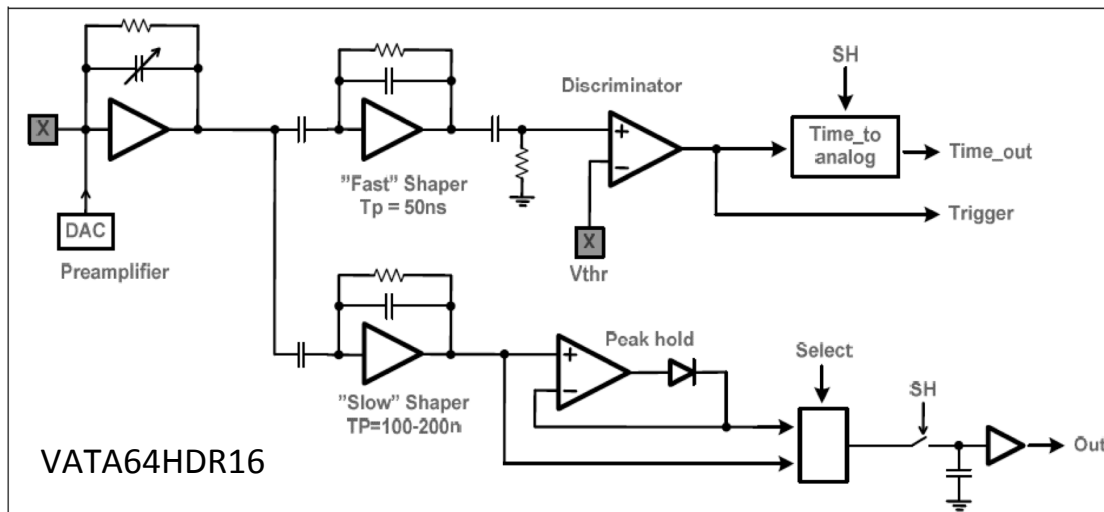
- designed for SPIDER2 R&D project under INFN funding
- delivered in December 2014



Prototype Construction in Siena/Pisa



Front End electronics and readout of 1024 SiPMs



M.G. Bagliesi et al., Nucl.Phys.B (Proc.Suppl.) 215:344-348, 2011

Custom ASIC VATA64HDR16:

- autotrigger mode / external trigger
- adjustable threshold / channel
- programmable slow shaper 50 – 300 ns
- Peak&Hold device → pulse height
- fast shaper + TAC → time measurement (~ 160 ps resolution)

2014: FDIRC readout board (1024 chns) developed in Siena/Pisa

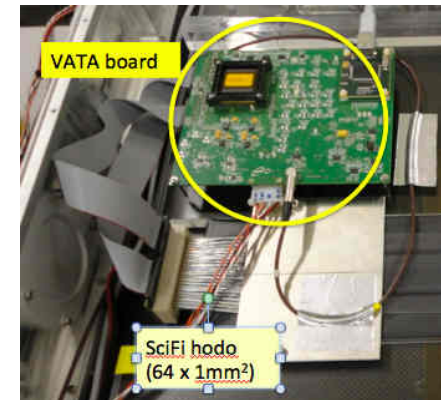
- 16 VATA chips
- 16 bit ADCs
- 1024 SiPM digitized signals
- 1024 time stamps
- autotrigger + 2 external triggers: random/physics
- USB-2 connection to PC

2009/10: development of custom ASIC for SiPM readout (64 channels)

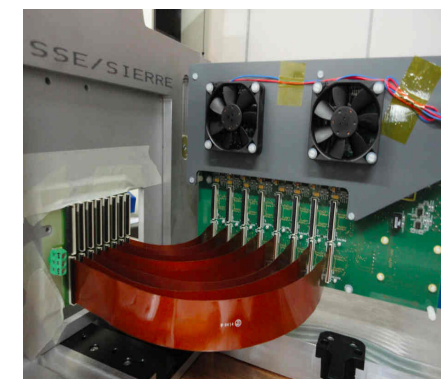
jointly funded by INFN (SPIDER project) and GM-IDEAS

- 64 pulse height measurements
- 64 time measurements
- **custom SiPM r/o board (64 chns)**

2011 SciFi hodoscope with SiPMs

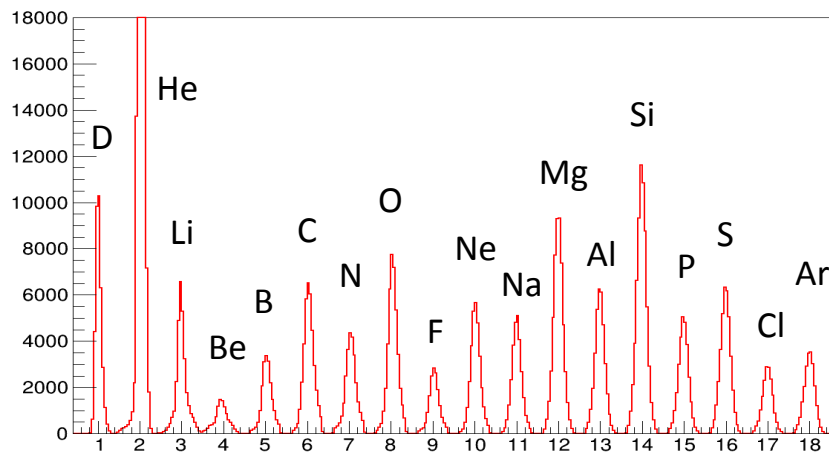


2013:
Scintillation Fiber
hodoscope
(1 VATA)
@CERN beam test

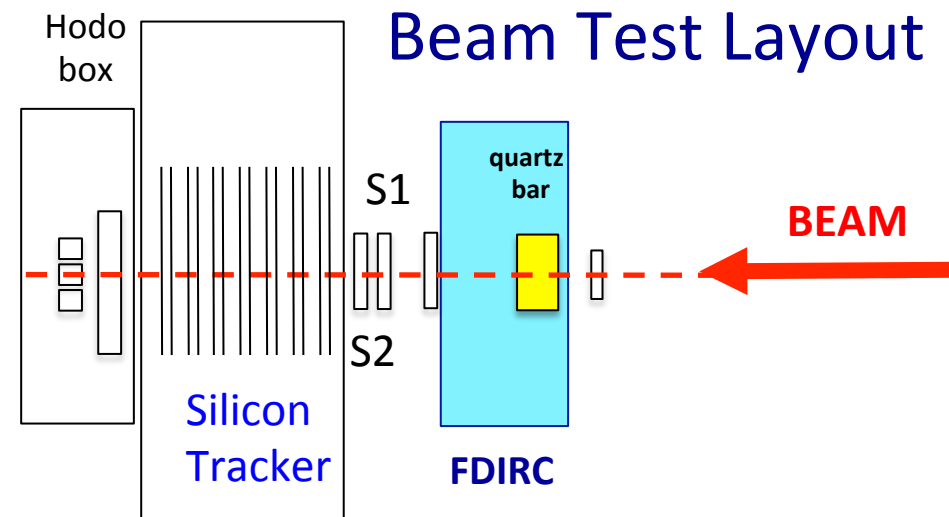


2014:
FDIRC prototype
(16 VATA)

- to test isotopic separation → need a LOW ENERGY ion beam ($P < \text{few tens of GeV}/c$)
- However, in March 2015 we had the opportunity of a parasitic beam test at CERN SPS (H8) primary Ar beam → internal target → ion fragments with $A/Z=2$: ${}^2\text{H}$, ${}^4\text{He}$, ..., ${}^{34}\text{Cl}$, ${}^{36}\text{Ar}$
- available beam energies (13, 19, 30 GeV/n) too large: no chance for isotopic separation by $\Delta\beta$ measurements. Nevertheless we decided to test the performance of the FDIRC prototype:

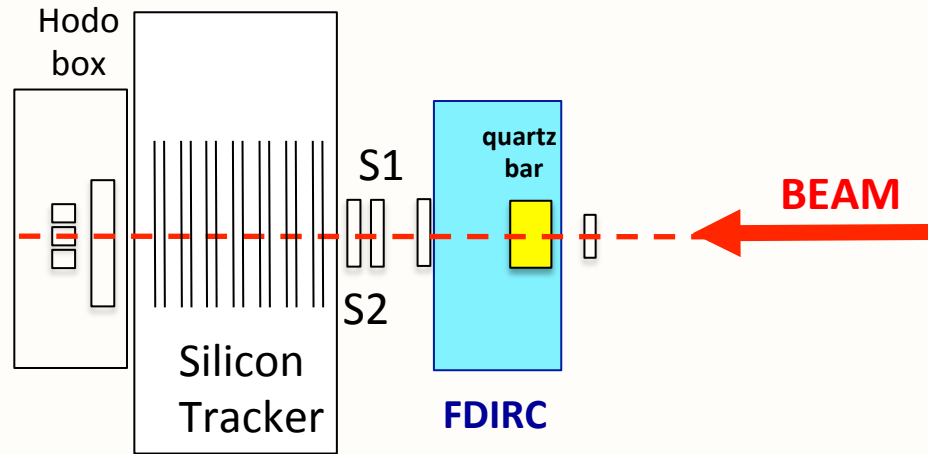


charge tagging by Silicon Tracker



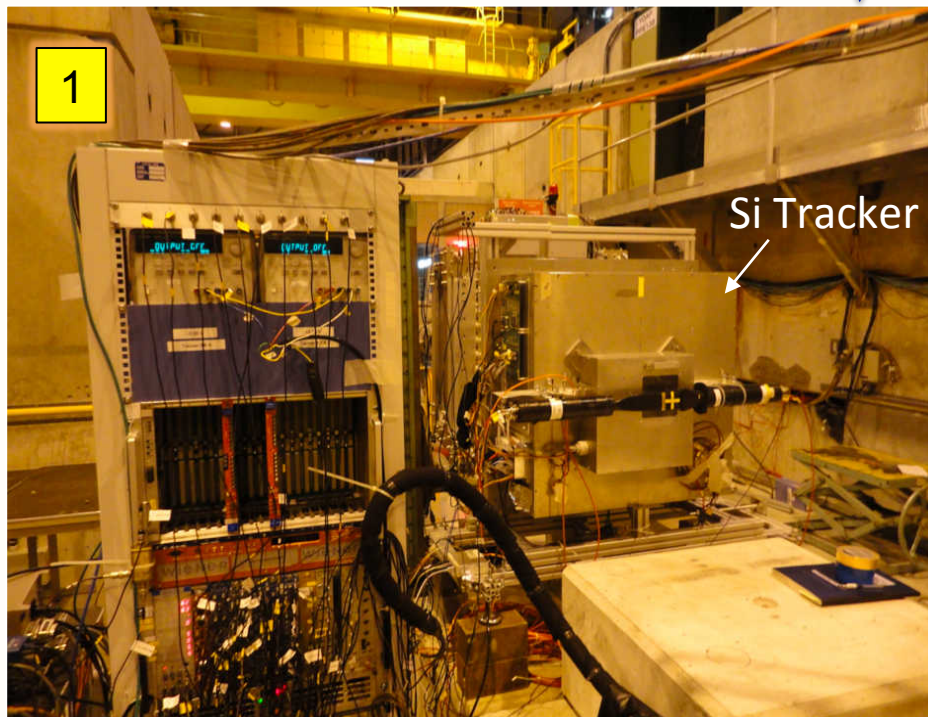
- Silicon Tracker identifies beam particles (+ interact. products in FDIRC)
- CHARGE-TAGGING: up to 14 independent dE/dx samples: 4 Si pixel layers + 10 Si strip layers

CERN 2015 Ion Beam Test Layout

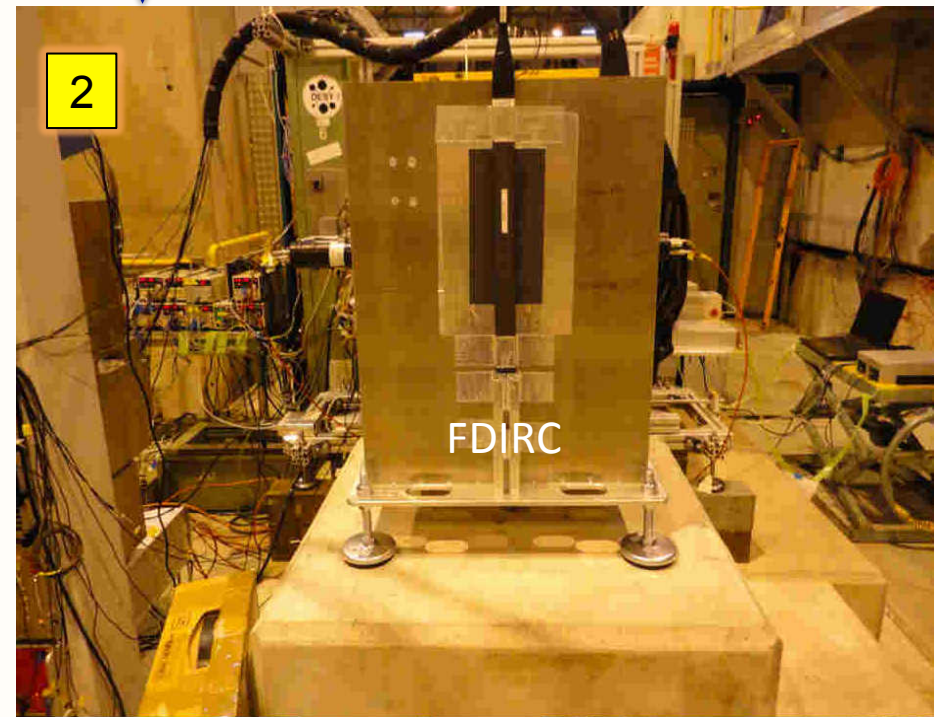


March 2015, SPS H8 beam line

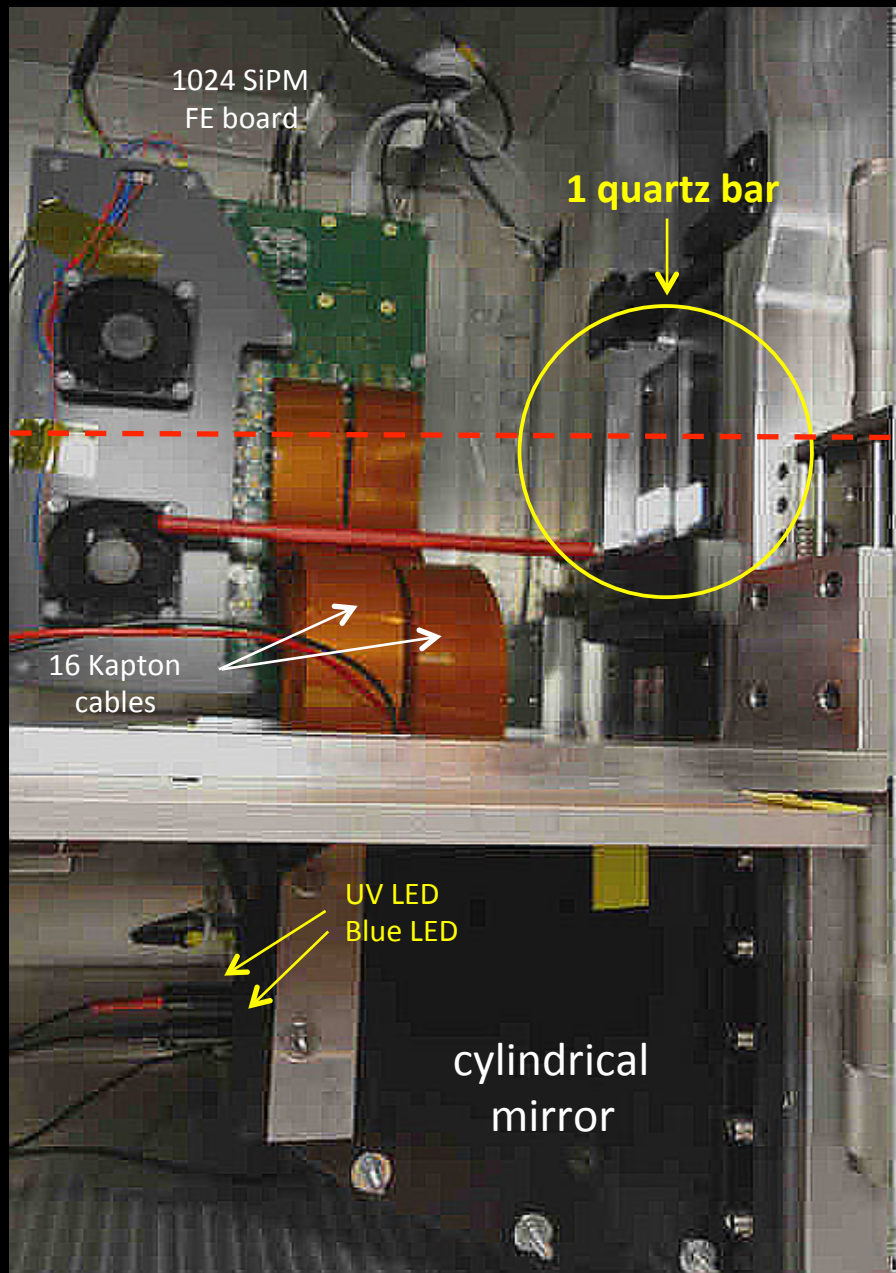
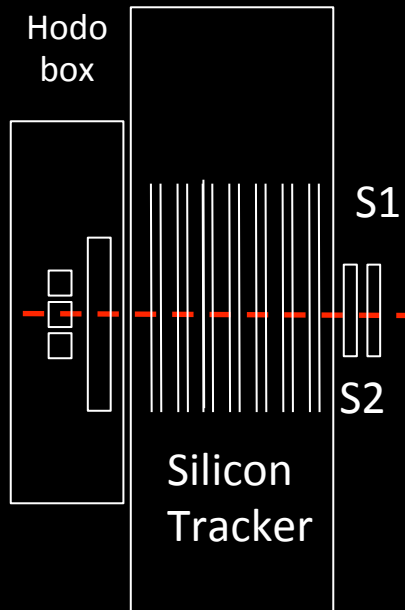
data: $\sim 9 \cdot 10^6$ triggers collected



beam line after TRACKER installation



beam line after FDIRC installation



BEAM ←

CERN-H8
BEAM TEST
LAYOUT
(March 2015)

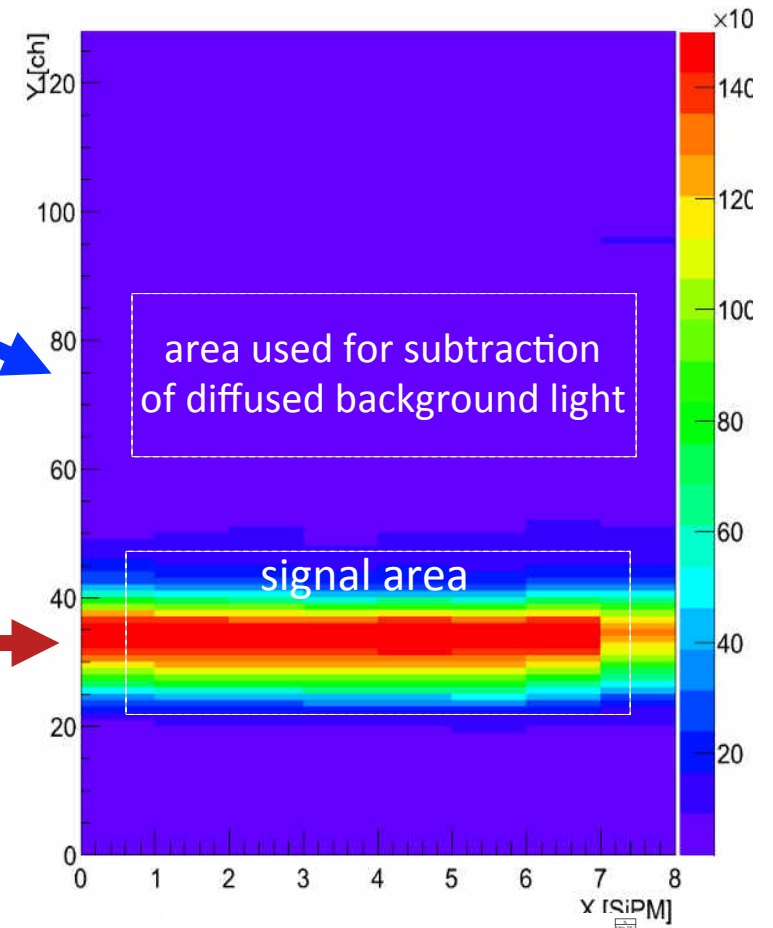
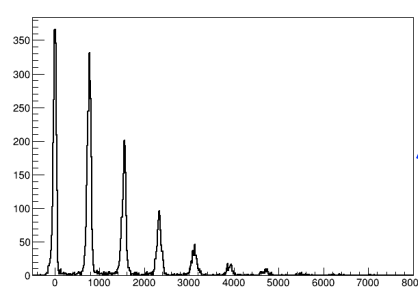
FDIRC Signal and Background subtraction

□ Diffused light background:

reflected light: mainly from the lateral walls of the mirror

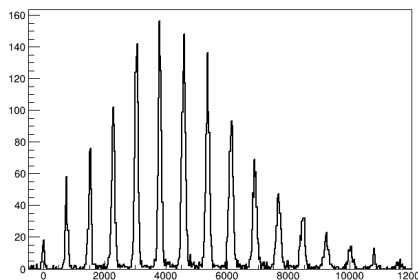
- small ($\sim 5\%$)
- from Cherenkov light
→ proportional to Z^2
- measured with beam triggers

in a region with **same area** but **far away** signal region.



□ Cherenkov signal:

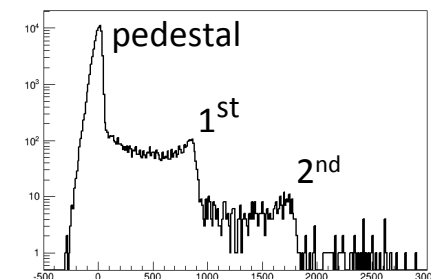
- from signal illuminated band
- well below SiPM saturation



□ Dark count background:

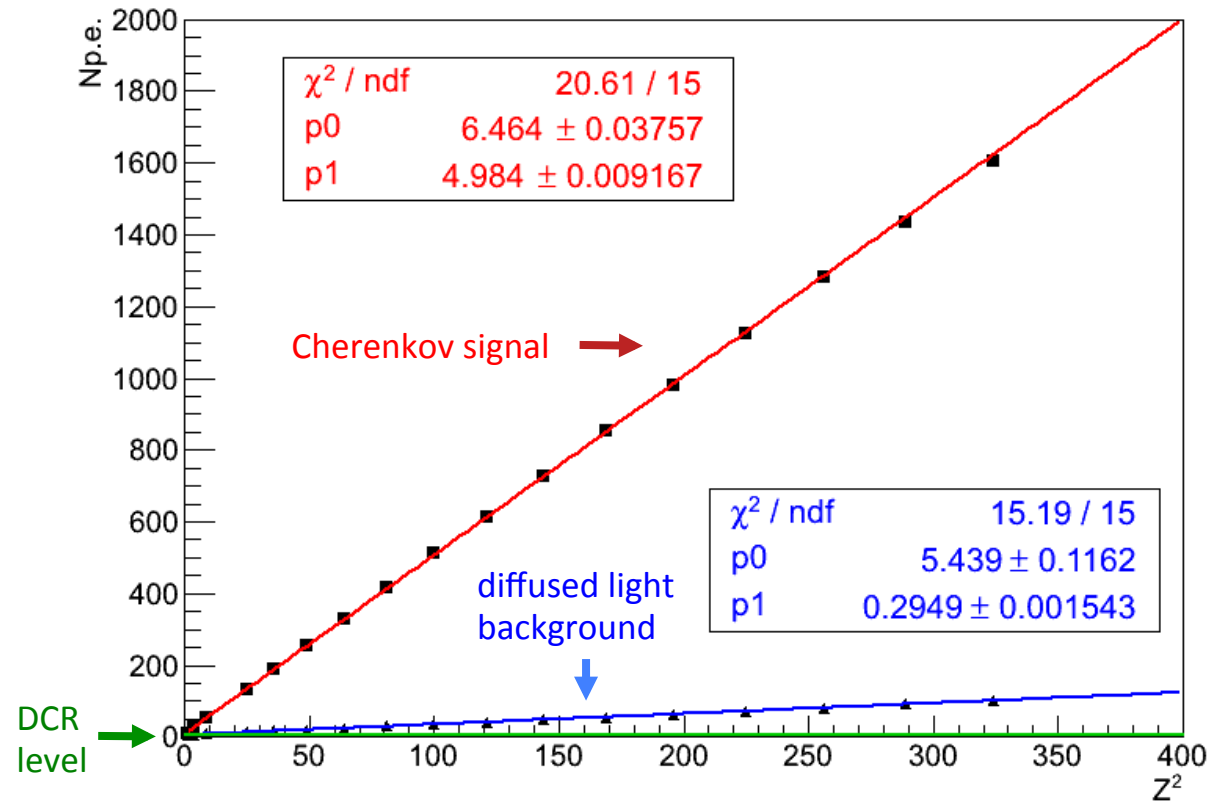
due to **SiPM dark count rate (DCR)**

- does not depend on atomic number Z
- measured with random triggers (off spill: no beam triggers)
- very small (%) due to excellent performance of SiPM arrays

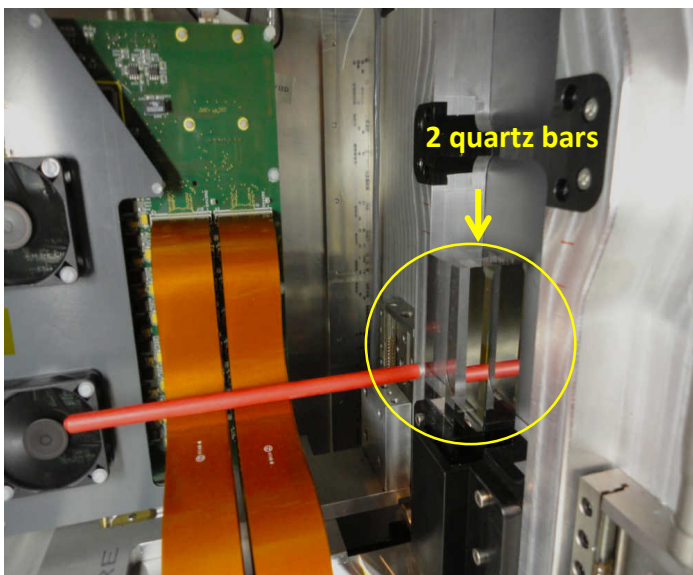
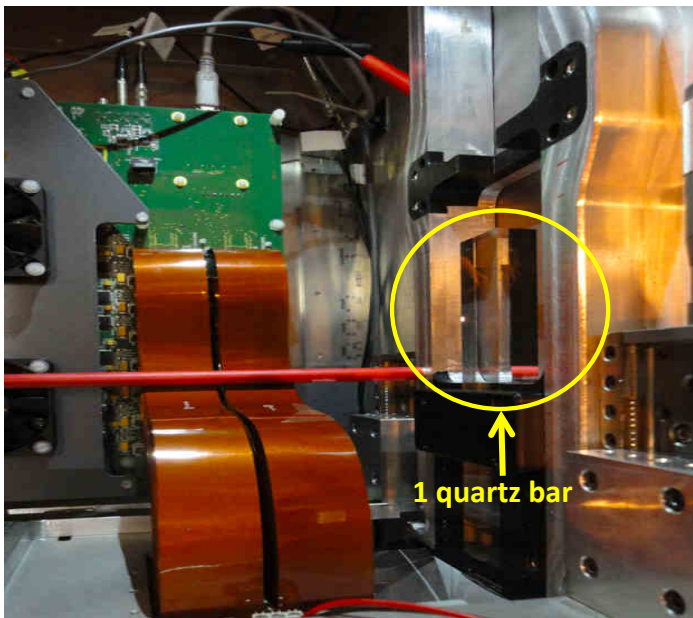


Z² dependence of Cherenkov (integrated) signal

- INTEGRAL measurement: use **Cherenkov light yield** $\sim Z^2$ to identify chemical elements by their CHARGE Z

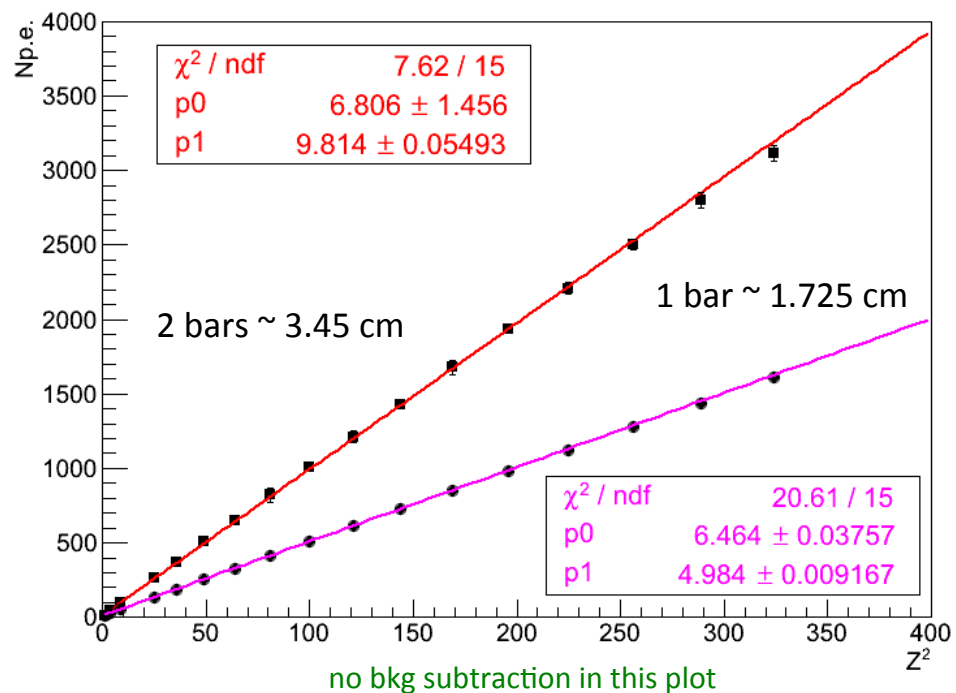


The mosaic of 16 SiPM arrays covers $\sim 1/3$ of the focal plane. Active area is $\sim 68\%$ of instrumented area (cracks among arrays). **For Z=1 particle:** ~ 5 p.e. in region covered by SiPM arrays $\rightarrow \sim 22$ p.e. on the whole focal plane (ideal seamless mosaic).



Integrated Cherenkov light from 2 radiators of different thickness

Z^2 dependence of Cherenkov signal

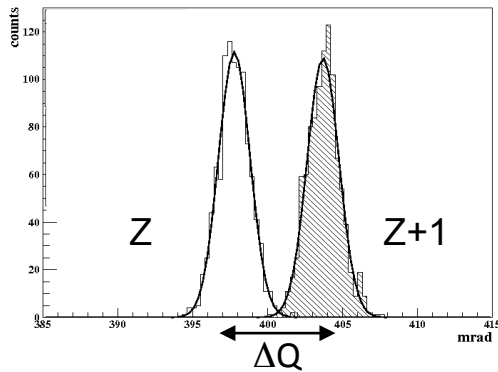


- scales linearly with radiator thickness
- NO SATURATION effects in the radiator as well as in the photosensors: each SiPM has 400 micro-cells and max photopeaks with 2 bars for Argon are < 50
→ well below SiPM saturation

Charge separation: Photon counting

- Integrated Cherenkov signal $Q(Z) \sim N_{pe} Z^2$
- $\sigma_Q^2 \sim N_{pe} Z^2 + \sigma_{electronics}^2 + \dots$
- σ_Q dominated by Poissonian photostatistics:
scales $\propto N_{pe}^{1/2} Z$ i.e. **linear in Z**

□ Charge separation $\equiv \sigma_Q / \Delta Q$

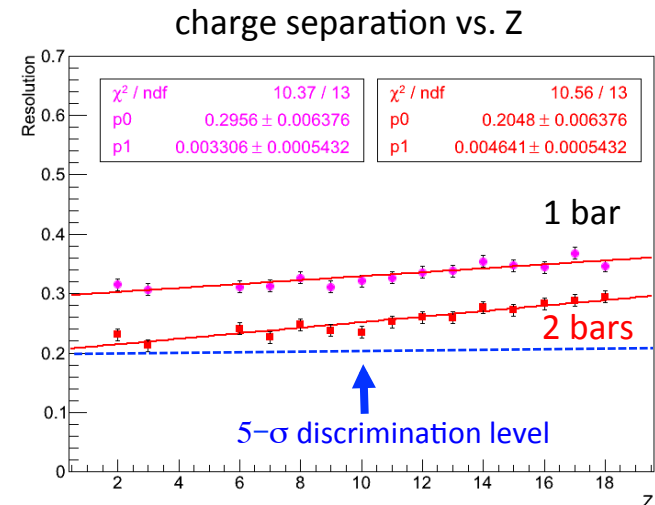
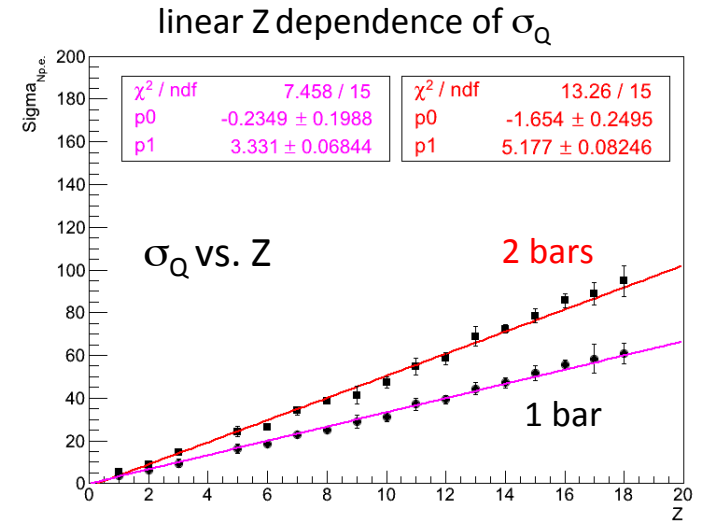


For two adjacent ions with $\Delta Z = 1$:
 $\Delta Q = Q(Z+1) - Q(Z) = (2Z+1) N_{pe}$

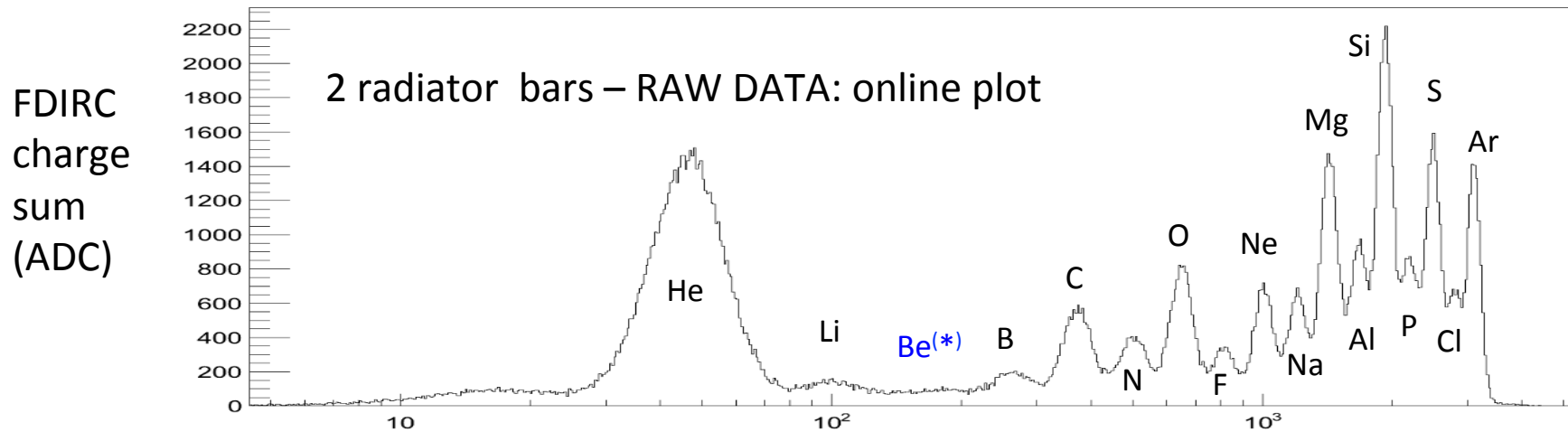
$$\frac{\sigma_Q}{\Delta Q} \propto \frac{1}{2\sqrt{N_{pe}}}$$

If purely Poissonian then charge separation should be independent of Z.
 Residual linear Z dependence is observed in the data.

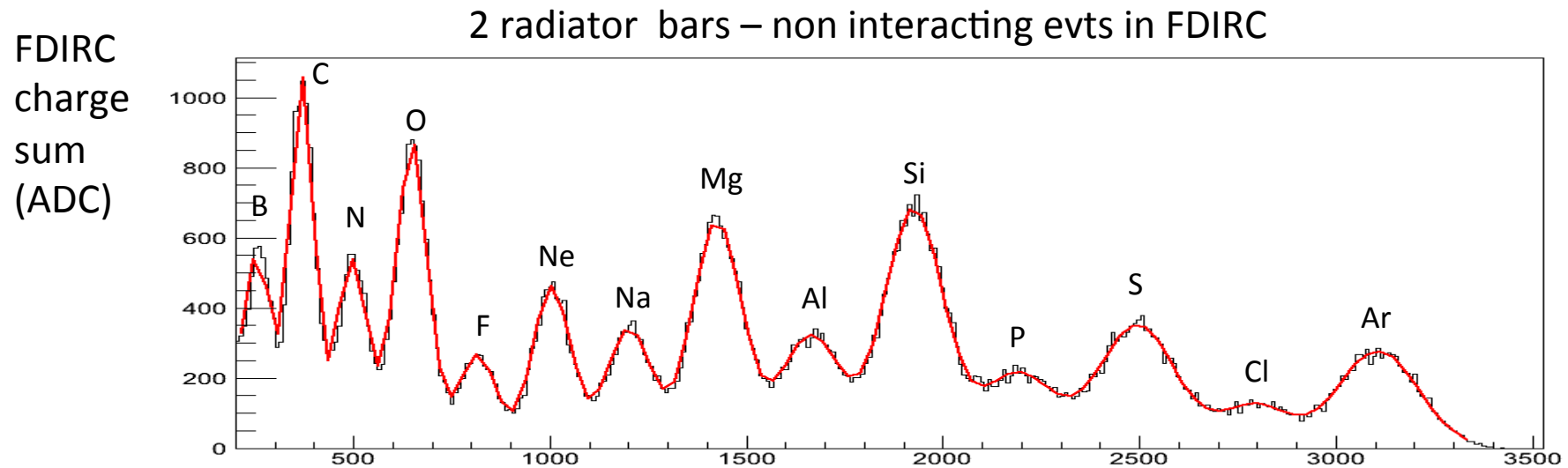
LEGENDA: Z = atomic number = charge of fully stripped ion
 N_{pe} = number of p.e. for Z=1 charge



Charge identification by FDIRC (integral measurement)

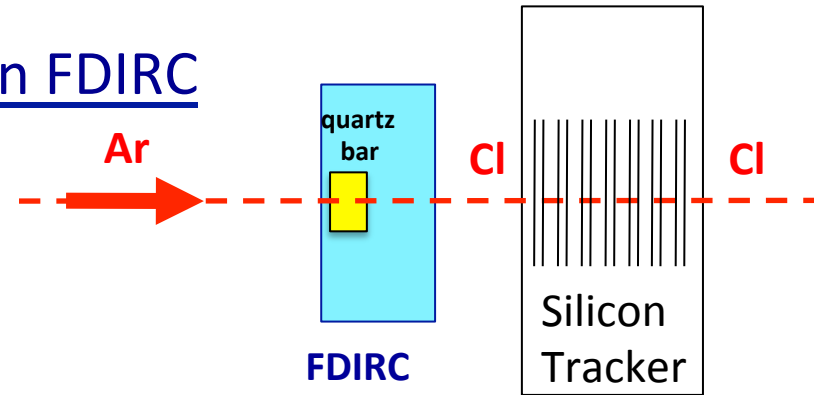
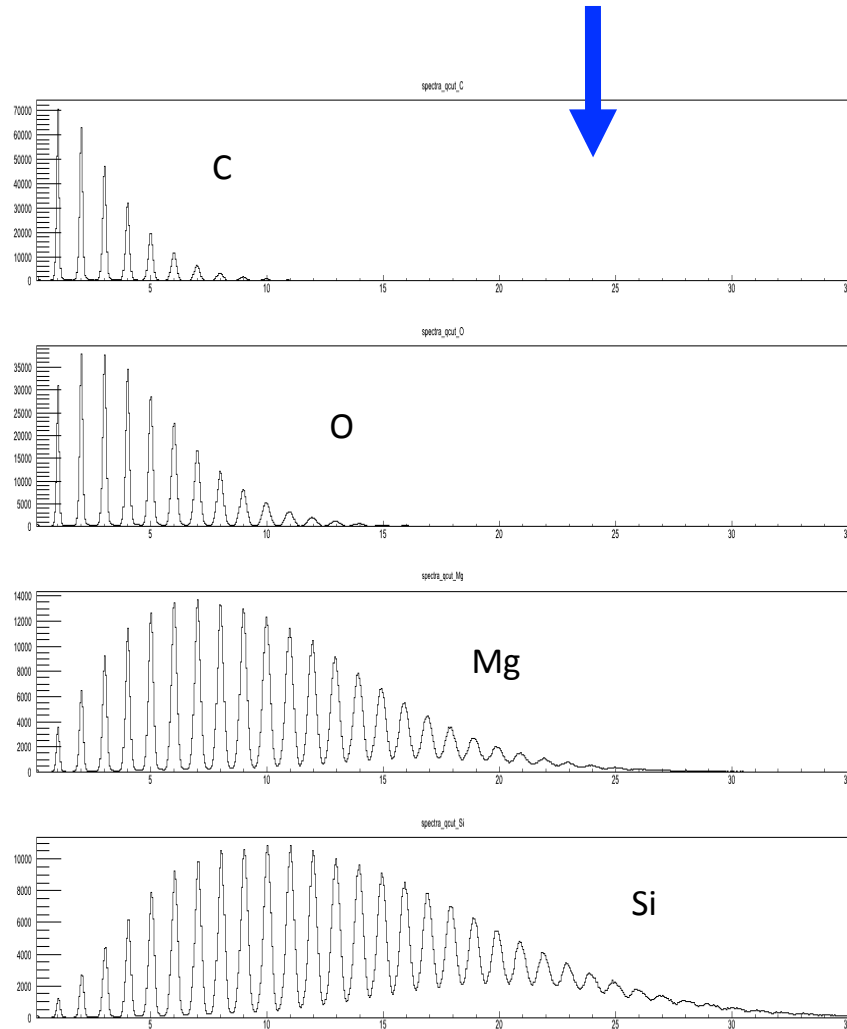


(*) practically no Be with $A/Z = 2$ beam setting: isotopes are ${}^7\text{Be}$, ${}^9\text{Be}$, ${}^{10}\text{Be}$ but not ${}^8\text{Be}$

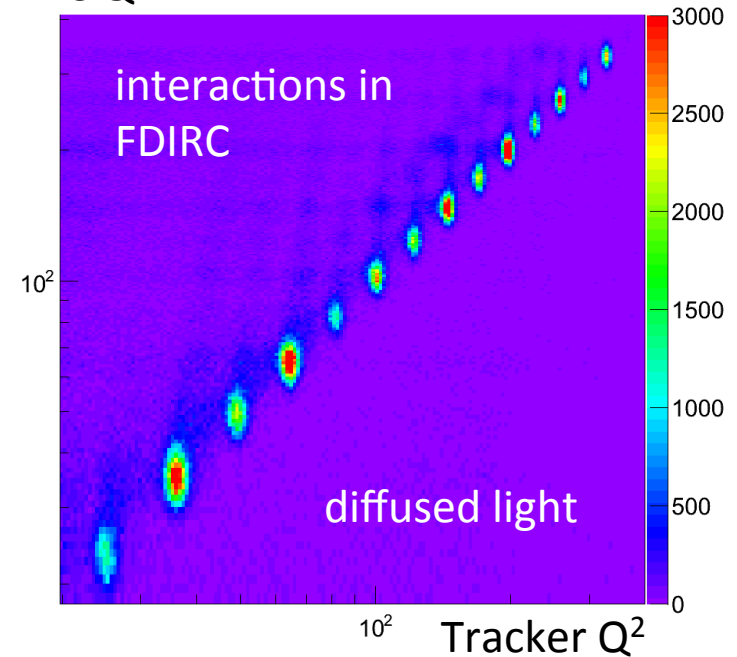


Selection of non-interacting nuclei in FDIRC

Photopeak spectra for non-interacting ions:
C, O, Mg, Si with 2 radiator bars



FDIRC Q^2

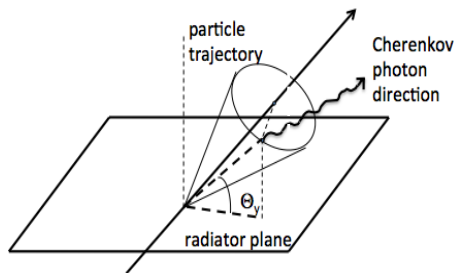
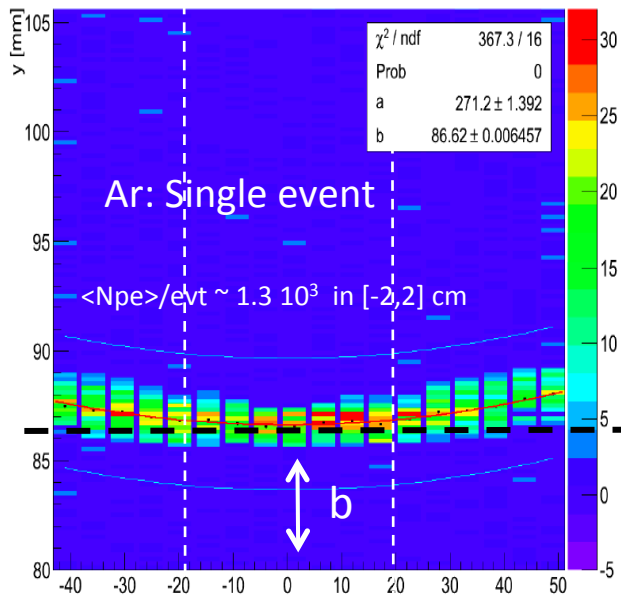


Use charge tagging from Tracker (downstream FDIRC): impose the same charge.

Cherenkov pattern fit: MC simulation vs beam test data

- DIFFERENTIAL measurement : fit Cherenkov pattern and measure FDIRC angular resolution.
- For tracks at normal incidence on the radiator the focal plane (FP) pattern is hyperbolic.

Geant4 simulation: 30 GeV/n Ar
(beam center shifted on the left as at beam test)



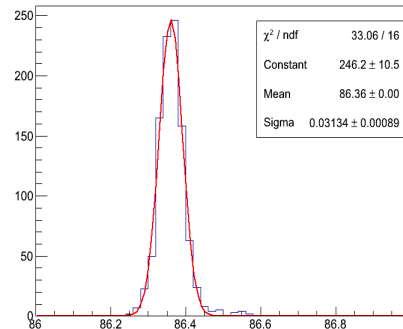
Θ_y is conserved during propagation along the bar

- Dead areas among SiPM arrays implemented in the simulation. In total 68% of instrumented FP area is active.

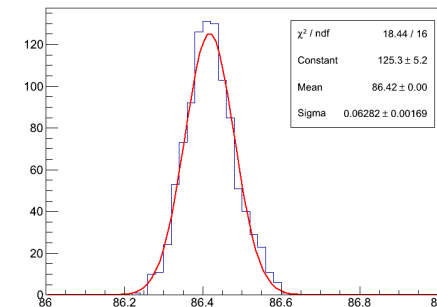
- fit FP pattern event-by-event:

→ Cherenkov angle Θ_y

→ apex of hyperbola (b parameter)



MC simulation:
fitted b parameter
 $\sigma_b \sim 31 \mu\text{m}$ for Ar

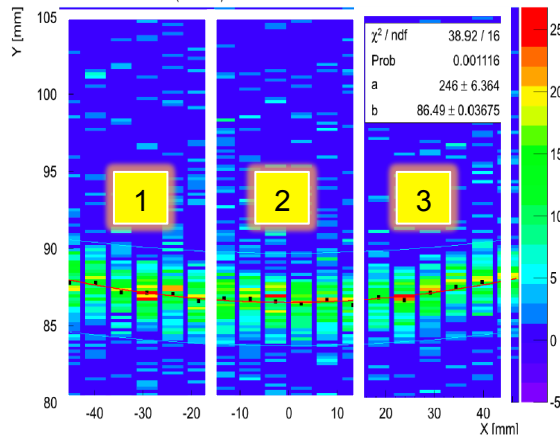


beam test data:
fitted b parameter
 $\sigma_b \sim 63 \mu\text{m}$ for Ar

Beam data fit to b-parameter is ~ 2 times larger than MC.

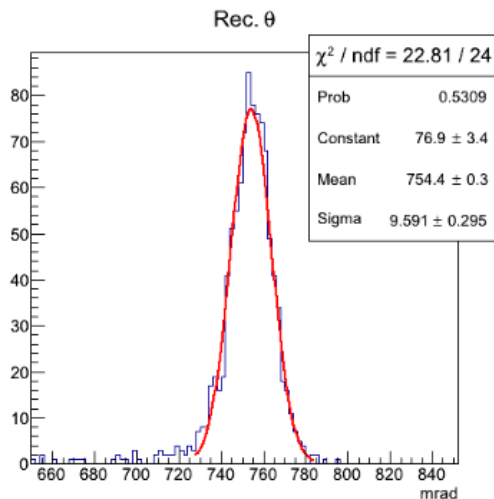
However: geometry + bar imperfections affecting light propagation along the bar are not simulated as well as diffused light background.

Cherenkov angle fit: MC simulation vs beam test data

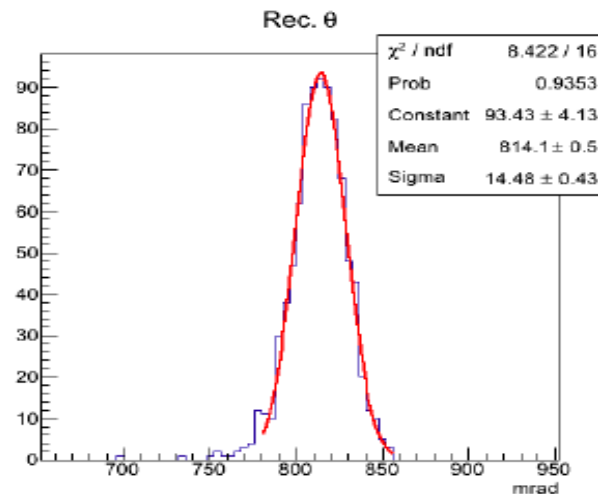


The SiPM instrumented area (4.3 cm x 2.7 cm) covers only $\sim 1/3$ of the Cherenkov pattern for tracks at normal incidence.

- Take data at 3 different positions moving the focal plane along the X-axis
- select events within a “narrow beam” (8 mm x 10 mm) using tracking info
- data stitching from 3 evts at different positions \rightarrow cover x: [-5, +4] cm
- angular fit “event-by-event” \rightarrow reconstruct Cherenkov angle Θ



Ar: MC simulation
fitted Cherenkov angle
 $\sigma_{\Theta} \sim 9.6$ mrad for Ar



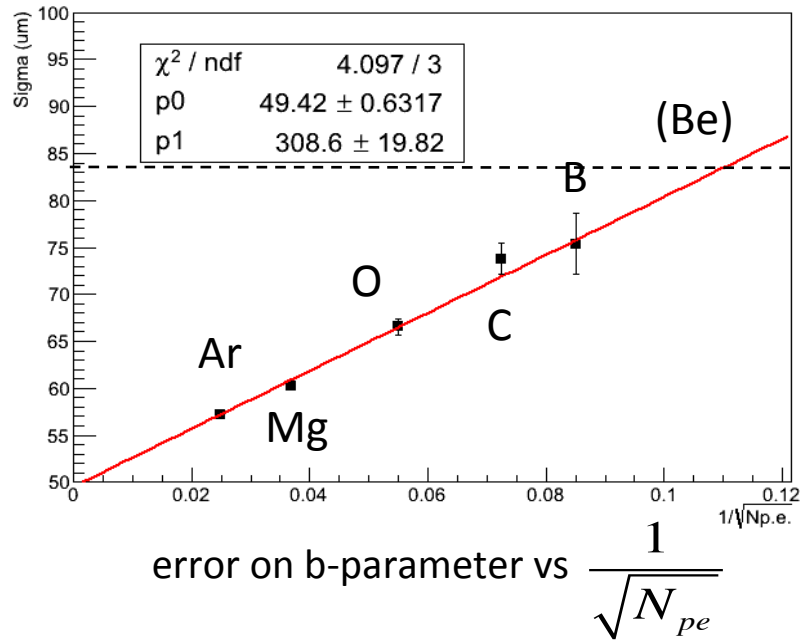
Ar: beam test data
fitted Cherenkov angle
 $\sigma_{\Theta} \sim 14.5$ mrad for Ar

Beam data angular resolution is ~ 1.5 larger than MC.

Of course, “stitching” 3 different events is just a poor approx. of having a larger instrumented area.

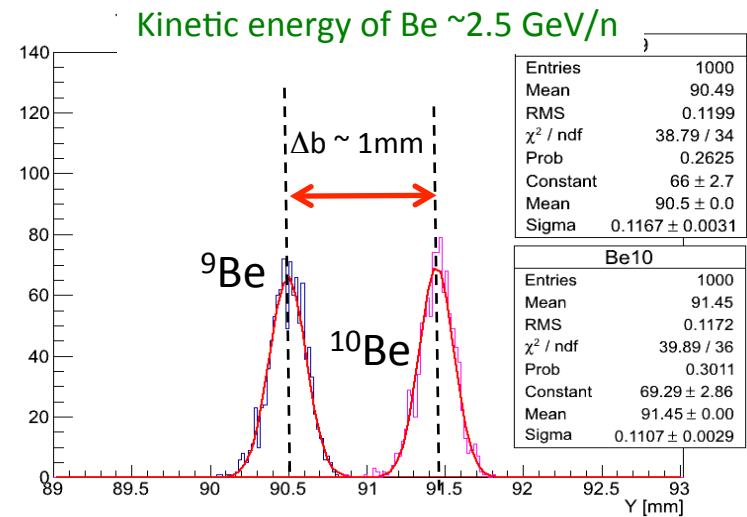
Be isotopes mass separation for $P < 30 \text{ GeV}/c$?

- ❑ To achieve $\sigma_M = 0.2 \text{ amu}$ (i.e. 5- σ mass separation) for $^{10}\text{Be}/^9\text{Be} \rightarrow \sigma_M/M < 0.02 \rightarrow \sigma_\theta \sim 1.5 \text{ mrad}$
- ❑ **Difficult** \rightarrow needs larger mirror aperture and larger FP coverage. How about apex parameter b?



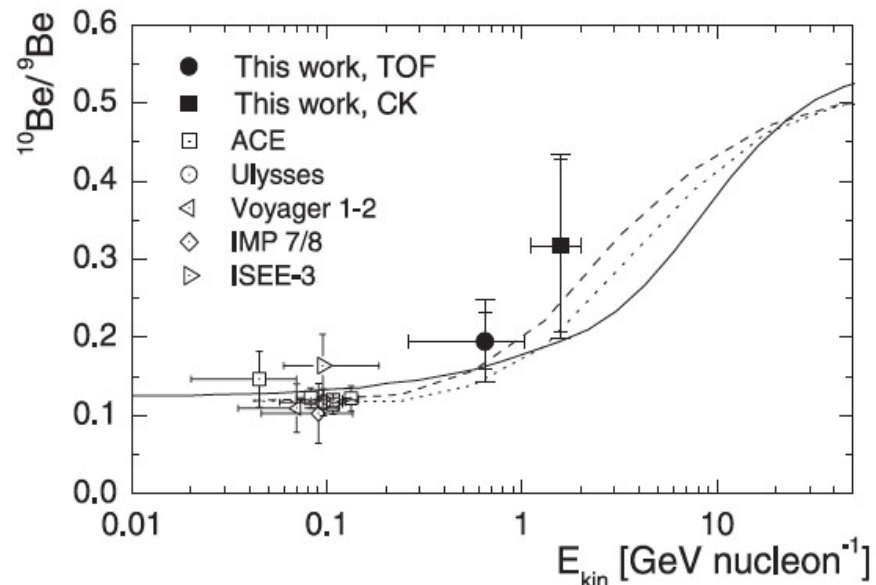
- Extrapolation of beam test data to Be gives $\sigma_b \sim 85 \mu\text{m}$
- This is possible because SiPM size is small ($200 \mu\text{m}$) along y.
- $\Delta b \sim 1 \text{ mm}$ distance on the focal plane for the two Be isotopes with $31.5 \text{ GeV}/c$ total particle momentum (at normal incidence).

- Simulated b-parameter distributions for ^9Be and ^{10}Be at $31.5 \text{ GeV}/n$ total momentum (7.9 GV).
- The difference in Cherenkov angles is $\sim 7.7 \text{ mrad}$
- The difference in b-parameter $\Delta b \sim 1 \text{ mm}$
- $[-2, +2] \text{ cm}$ instrumented area as in beam test.
- **track at normal incidence**



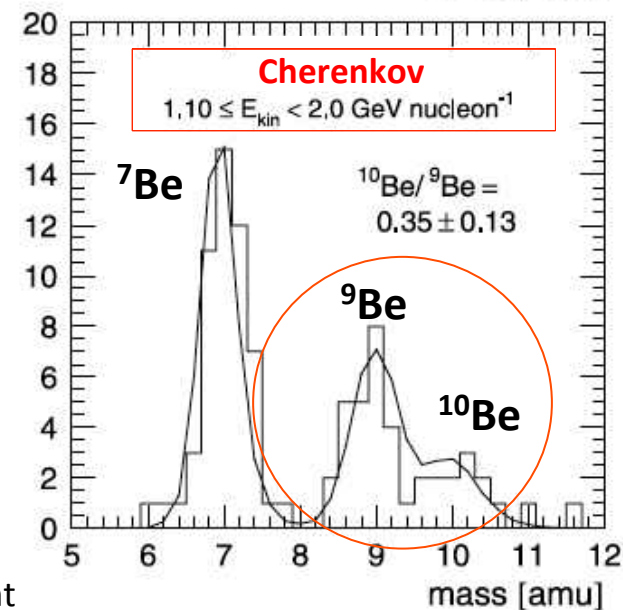
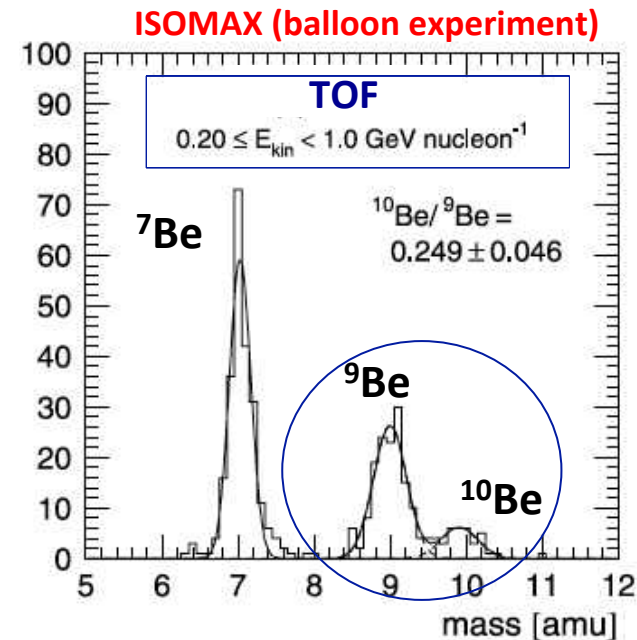
Cosmic $^{10}\text{Be}/^9\text{Be}$ Isotopic Ratio

- ^{10}Be decays to stable ^9Be isotope
- “radioactive clock” to measure propagation time in the galaxy: very important for astrophysics models
- few measured points vs kinetic energy per nucleon
- need better accuracy + extension above 2 GeV/n



T. Hams et al., The Astrophysical Journal, 611:892–905, 2004

$^{10}\text{Be}/^9\text{Be}$ Be isotopic ratio measurement by ISOMAX (balloon experiment with magnetic spectrometer) in the energy interval $0.2 < E_{\text{kin}} < 2.0$ GeV/n



Conclusions

- Good beam test performance of our first prototype of FDIRC with SiPM readout
 - Outstanding performance of the SiPM arrays
 - Ion beam test data are very useful to study detector's full dynamic range
 - Preliminary analysis presented (2 months after beam test) → more to come...
 - Proof-of-principle for isotope separation requires lower beam momentum
-
- $^{10}\text{Be}/^9\text{Be}$ measurement is not easy.
 - Effect of the incident track angle vs. mass resolution has to be studied.
 - Requires 1-2% momentum resolution from magnetic spectrometer at 30 GeV/c.
 - Quartz bars + coverage of focal plane with SiPM array are expensive.

(no meal is for free...)

Authors

P.S. Marrocchesi^(a), M.G. Bagliesi^(a), G. Bigongiari^(a), S. Bonechi^(a),
P. Brogi^(a), P. Maestro^(a), F. Morsani^(b), C. Piemonte^(c), F. Stolzi^(a),
J.E. Suh^(a), A. Sulaj^{a)}

(a) University of Siena and INFN Pisa

(b) INFN Sezione di Pisa

(c) Fondazione Bruno Kessler (FBK), Trento