



SHIP

SEARCH FOR HIDDEN PARTICLES

A New Experiment Proposal



Giovanni De Lellis

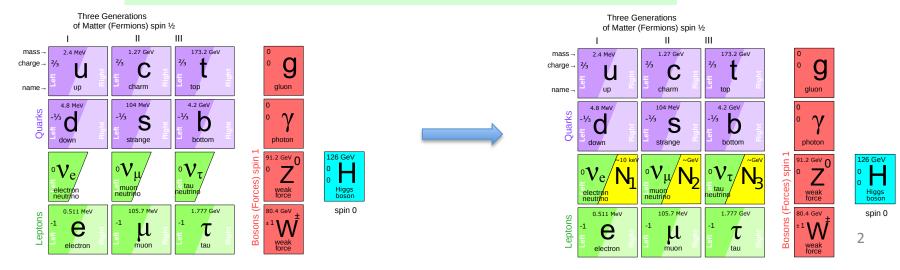
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Theoretical motivation

- Discovery of the 126 GeV Higgs boson

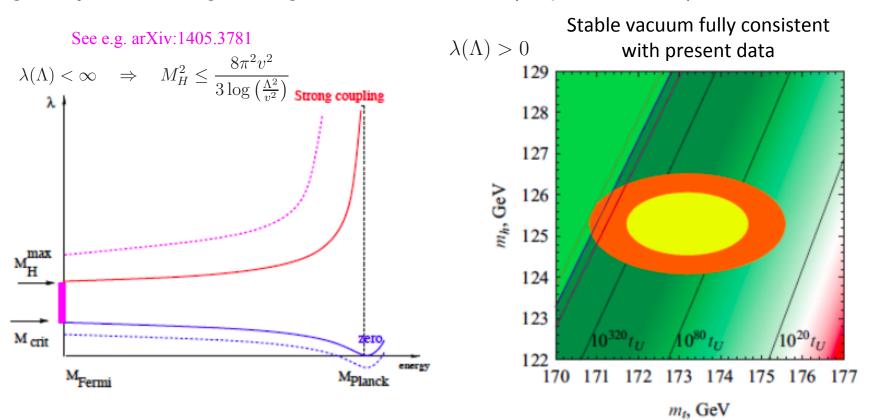
 Triumph of the Standard Model
 The SM may work successfully up to Planck scale!
- SM is unable to explain:
 - Neutrino masses
 - Excess of matter over antimatter in the Universe
 - The nature of non-baryonic Dark Matter
- All three issues can be solved by adding three new fundamental fermions, right-handed Majorana Heavy Neutral Leptons (HNL): N₁, N₂ and N₃

vMSM: T.Asaka, M.Shaposhnikov PL B620 (2005) 17



SM may well be a consistent effective theory all the way up to the Plank scale

- ✓ M_H < 175 GeV \rightarrow SM is a weakly coupled theory up to the Plank energies !
- ✓ $M_H > 111 \text{ GeV} \rightarrow EW$ vacuum is stable or metastable with a lifetime greatly exceeding the age of our Universe (Espinosa et al)



✓ No sign of New Physics seen

No sign of New Physics seen What is not found...

ATLAS SUSY Searches* - 95% CL Lower Limits

 $\sqrt{s} = 7 \text{ TeV}$

full data

 $\sqrt{s} = 8 \text{ TeV}$

 $\sqrt{s} = 8 \text{ TeV}$

ATLAS Preliminary

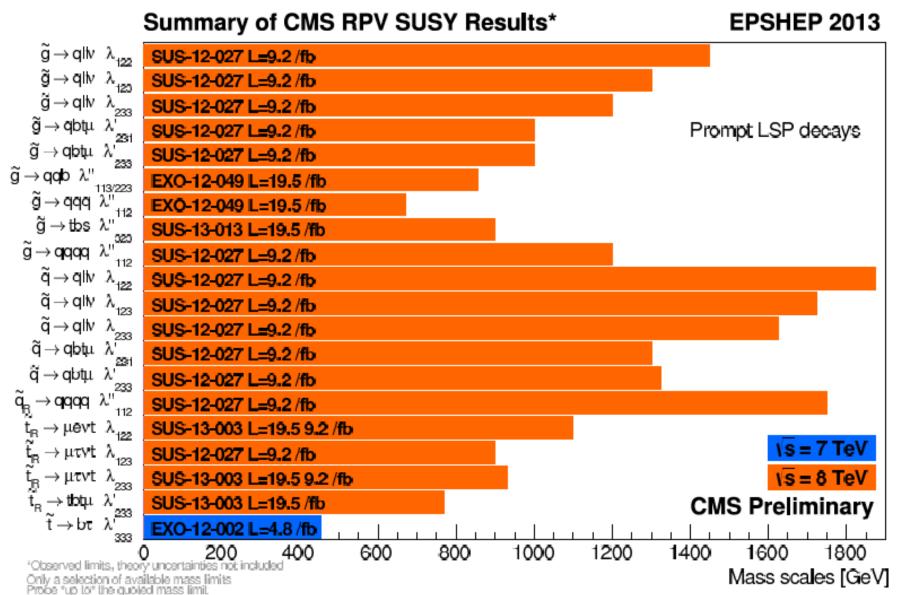
Mass scale [TeV]

Status: SUSY 2013 $\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$ $\sqrt{s} = 7.8 \text{ TeV}$ e, μ, τ, γ Jets $E_{\tau}^{miss} \int \mathcal{L} dt[fb^{-1}]$ Mass limit Model Reference 2-6 jets $m(\tilde{q})-m(\tilde{g})$ ATLAS-CONF-2013-047 MSUGRA/CMSSM 0 Yes 20.3 1.7 TeV 3-6 jets MSUGRA/CMSSM 1 e.u Yes 20.3 1.2 TeV any m(3) ATLAS-CONF-2013-062 MSUGRA/CMSSM 0 7-10 jets 20.3 any m(§) Yes 1.1 TeV 1308,1841 0 2-6 jets 20.3 740 GeV ATLAS-CONF-2013-047 $m(\hat{\xi}_1^0)=0$ GeV $\bar{q}\bar{q}, \bar{q} \rightarrow q\bar{q}_1$ ğğ, ğ→qōk 0 2-6 jets 20.3 $m(\hat{\xi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-047 $\vec{g}\vec{g}, \vec{g} \rightarrow qq\vec{\ell}_1^a \rightarrow qqW^{\pm}\vec{\ell}_1^c$ 3-6 jets 20.3 1.18 TeV $m(\tilde{k}_{1}^{2}) < 200 \text{ GeV}, m(\tilde{k}^{2}) = 0.5(m(\tilde{k}_{1}^{2}) + m(\tilde{g}))$ ATLAS-CONF-2013-062 $\tilde{R}\tilde{R}, \tilde{R} \rightarrow qq(\ell\ell/\ell\nu/\nu\nu)\tilde{X}_{1}^{0}$ $2e, \mu$ 0-3 jets 20.3 ATLAS-CONF-2013-089 1.12 TeV $m(\tilde{\ell}_1^0)=0 \text{ GeV}$ 2 e.u 2-4 jets tan/8<15 GMSB (NLSP) Yes 4.7 1.24 TeV 1208.4688 GMSB (/ NLSP) 1-2 T 0-2 jets Yes 20.7 1.4 TeV tan/6 > 18 ATLAS-CONF-2013-026 GGM (bine NLSP) 2 y Yes 4.8 1.07 TeV m(t) >50 GeV 1209.0753 GGM (wino NLSP) 619 GeV $1e, \mu + \gamma$ Yes 4.8 m(kg)>50 GeV ATLAS-CONF-2012-144 GGM (higgsino-bino NLSP) 4.8 1 b Yes m(£3)>220 GeV 1211,1167 GGM (higgsino NLSP) $2e, \mu(Z)$ 0-3 jets Yes 5.8 m(H)>200 GeV ATLAS-CONF-2012-152 Gravitino LSP mono-jet Yes 10.5 m(2)>10⁻⁴ eV ATLAS-CONF-2012-147 0 0 1.2 TeV ATLAS-CONF-2013-061 $\bar{g} \rightarrow b\bar{b}\bar{\chi}$ 3 b Yes 20.1 7-10 jets Yes 20.3 1.1 TeV $\tilde{g} \rightarrow t \tilde{t} \tilde{X}$ $m(\tilde{k}_1^3) < 350 \text{ GeV}$ 1308,1841 0-1 e, µ 3 b Yes 20.1 1.34 TeV ATLAS-CONF-2013-061 $\tilde{g} \rightarrow t \tilde{t} \tilde{X}_1^0$ 0-1 e, µ 20.1 1.3 TeV 3 b Yes m(£1)<300 GeV ATLAS-CONF-2013-061 $\tilde{g} \rightarrow b\tilde{t}\tilde{X}_1$ 0 100-620 GeV $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{k}_1^{\dagger}$ 2 Б Yes 20.1 1308.2631 2 e.u (SS) 275-430 GeV $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow t \tilde{\chi}_1$ 0-3 6 Yes 20.7 $m(\tilde{\mathcal{X}}_1^*)=2 m(\tilde{\mathcal{X}}_1^0)$ ATLAS-CONF-2013-007 110-167 GeV 1-2 e, µ 1208.4305, 1209.2102 $\tilde{t}_1 \tilde{t}_1(light), \tilde{t}_1 \rightarrow b\tilde{t}_1^2$ 1-2 b Yes 4.7 $m(\tilde{\ell}_1^3)=55 \text{ GeV}$ $2e,\mu$ 0-2 jets $\tilde{t}_1 \tilde{t}_1(light), \tilde{t}_1 \rightarrow Wb\tilde{t}_1^2$ Yes 20.3 130-220 GeV $m(\tilde{\mathcal{E}}_1^0) = m(\tilde{\mathfrak{T}}_1) \cdot m(W) \cdot 50 \text{ GeV. } m(\tilde{\mathfrak{T}}_1) << m(\tilde{\mathcal{E}}_1^+)$ ATLAS-CONF-2013-048 2 jets ATLAS-CONF-2013-065 $\tilde{t}_1 \tilde{t}_1 \text{(medium)}, \tilde{t}_1 \rightarrow t \tilde{k}_1^c$ $2e, \mu$ Yes 20.3 225-525 GeV $\tilde{t}_1 \tilde{t}_1 \text{ (medium)}, \tilde{t}_1 \rightarrow b \tilde{V}_1^*$ 0 2 b Yes 20.1 150-580 GeV $m(\tilde{\ell}_1^2)$ <200 GeV, $m(\tilde{\ell}_1^2)$ - $m(\tilde{\ell}_1^2)$ =5 GeV 1308,2631 $1e, \mu$ 20.7 200-610 GeV ATLAS-CONF-2013-037 1 b Yes $m(\tilde{\ell}_1^2)=0 \text{ GeV}$ $\tilde{t}_1 \tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t \tilde{V}_1$ 320-660 GeV 0 2 b Yes 20.5 $m(\tilde{k}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-024 $\tilde{t}_1 \tilde{t}_1 \text{(heavy)}, \tilde{t}_1 \rightarrow t \tilde{t}_1^{\prime}$ ō mono-jet/c-tag Yes 20.3 90-200 GeV m(t/)-m(t//)<85 GeV ATLAS-CONF-2013-068 $\bar{t}_1\bar{t}_1, \bar{t}_1 \rightarrow c k_1'$ t1 t1 (natural GMSB) 2 e.μ (Z) Yes 20.7 m(k³₁)>150 GeV ATLAS-CONF-2013-025 1 b 500 GeV $\tilde{t}_{2}\tilde{t}_{2}, \tilde{t}_{2} \rightarrow \tilde{t}_{1} + Z$ 3 e.μ (Z) 20.7 271-520 GeV ATLAS-CONF-2013-025 1 b Yes $m(\overline{t}_1)=m(\overline{X}_1^2)+180 \text{ GeV}$ $\tilde{l}_{LR}\tilde{l}_{LR}, \tilde{l} \rightarrow \ell \tilde{r}_1^0$ $2e,\mu$ 0 Yes 20.3 85-315 GeV $m(\tilde{k}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-049 $\tilde{X}_{1}^{-}\tilde{X}_{1}^{-}, \tilde{X}_{1}^{+} \rightarrow \tilde{\ell}\nu(\ell\tilde{\nu})$ 2 e.u 0 Yes 20.3 125-450 GeV $m(\tilde{\mathcal{X}}_1^0)=0$ GeV, $m(\tilde{\mathcal{X}}, \tilde{\mathcal{V}})=0.5(m(\tilde{\mathcal{X}}_1^0)+m(\tilde{\mathcal{X}}_1^0))$ ATLAS-CONF-2013-049 $\widetilde{X}_{1}^{1}\widetilde{X}_{1}^{1},\widetilde{X}_{1}^{1}\rightarrow\widetilde{\tau}\nu(\tau\overline{\nu})$ Yes 20.7 180-330 GeV $m(\bar{k}_1^3)=0$ GeV, $m(\bar{\tau}, \bar{\tau})=0.5(m(\bar{k}_1^3)+m(\bar{X}_1^3))$ 27 ATLAS-CONF-2013-028 20.7 $\tilde{\chi}_{1}^{*}\tilde{\chi}_{2}^{0} \rightarrow \tilde{\ell}_{L} \nu \tilde{\ell}_{L} \ell(\tilde{\nu}_{Y}), \ell \tilde{\nu} \tilde{\ell}_{L} \ell(\tilde{\nu}_{Y})$ 3 e, µ Ō Yes 600 GeV $m(\overline{\mathcal{K}}_{1}^{\pm})=m(\overline{\mathcal{K}}_{2}^{0}), m(\overline{\mathcal{K}}_{1}^{0})=0, m(\overline{\mathcal{E}}, \overline{\mathcal{V}})=0.5(m(\overline{\mathcal{K}}_{1}^{\pm})+m(\overline{\mathcal{K}}_{1}^{0}))$ ATLAS-CONF-2013-035 $\vec{x}_1 \vec{x}_2 \rightarrow W \vec{x}_1 Z \vec{x}_1$ 3 e, µ 20.7 Ō Yes 315 GeV $m(\tilde{\ell}_1^{\pm})=m(\tilde{k}_2^0), m(\tilde{\ell}_1^0)=0$, sleptons decoupled ATLAS-CONF-2013-035 $\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} h \tilde{\chi}_{1}^{0}$ $m(\tilde{\ell}_{1}^{\pm})=m(\tilde{\ell}_{2}^{0}), m(\tilde{\ell}_{1}^{0})=0$, sleptons decoupled $1e, \mu$ 2 b Yes 20.3 285 GeV ATLAS-CONF-2013-093 Disapp. trk Direct $\bar{X}_{1}^{+}\bar{X}_{1}^{-}$ prod., long-lived \bar{X}_{1}^{\pm} 1 jet Yes 20.3 270 GeV $m(\tilde{k}_1^z)$ - $m(\tilde{k}_1^0)$ =160 MeV, $\tau(\tilde{k}_1^z)$ =0.2 ns ATLAS-CONF-2013-069 Stable, stopped & R-hadron 22.9 822 GeV m(k⁰₁)=100 GeV, 10 με ντ(ḡ) ×1000 e 1-5 jets Yes ATLAS-CONF-2012-057 10<tan6<50 GMSB, stable $\bar{\tau}, \bar{X}_1^0 \rightarrow \bar{\tau}(\bar{e}, \bar{\mu}) + \tau(e, \mu) = 1 \cdot 2 \mu$ 15.9 475 GeV ATLAS-CONF-2013-058 2 y 230 GeV $0.4 < r(\tilde{\chi}_1^0) < 2 \text{ ns}$ GMSB, $\tilde{\chi}_1 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1$ Yes 4.7 1304,6310 1 µ, displ. vtx 20.3 1.0 TeV ATLAS-CONF-2013-082 ãã. X₁→qqμ (RPV) 1.5 < cr < 156 mm, BR(μ)=1, m(\overline{K}_{1}^{0})=108 GeV λ'_{in1} =0.10, λ_{132} =0.05 LFV $\rho p \rightarrow \bar{\nu}_{\tau} + X$, $\bar{\nu}_{\tau} \rightarrow e + \mu$ $2e, \mu$ 4.6 1.61 TeV 1212.1272 LFV $pp \rightarrow \bar{\nu}_{\tau} + X, \bar{\nu}_{\tau} \rightarrow e(\mu) + \tau$ $\lambda'_{101}=0.10, \lambda_{112100}=0.05$ $1e, \mu + \tau$ 4.6 1.1 TeV 1212,1272 Bilinear RPV CMSSM 1 e. µ 7 jets Yes 4.7 1.2 TeV $m(\tilde{q})=m(\tilde{g}), c\tau_{LSP}<1 mm$ ATLAS-CONF-2012-140 $\tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow W \tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow ee\bar{\nu}_{a}, e\mu\bar{\nu}_{a}$ $4 e. \mu$ Yes 20.7 760 GeV m(x2)>300 GeV, x22>0 ATLAS-CONF-2013-036 $\tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow W\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow \tau \tau \tilde{\nu}_{e}, e \tau \tilde{\nu}_{\tau}$ $3e, \mu + \tau$ Yes 20.7 350 GeV $m(\tilde{k}_1^9) > 80 \text{ GeV}, \lambda_{133} > 0$ ATLAS-CONF-2013-036 BR(c)=BR(b)=BR(c)=0% $\tilde{g} \rightarrow qqq$ 0 6-7 jets 20.3 916 GeV ATLAS-CONF-2013-091 2 e, µ (SS) 20.7 $\tilde{g} \rightarrow \tilde{t}_1 t$, $\tilde{t}_1 \rightarrow bs$ 0-3 b Yes 880 GeV ATLAS-CONF-2013-007 Scalar gluon pair, sgluon→gq 0 4 jets 4.6 agluon 100-287 GeV incl. limit from 1110,2693 1210.4826 Scalar gluon pair, sqluon $\rightarrow t\bar{t}$ 2 e, μ (SS) 1 b Yes 14.3 800 GeV ATLAS-CONF-2013-051 WIMP interaction (D5, Dirac y) 0 mono-jet Yes 10.5 m(x)<80 GeV, limit at<687 GeV for D8 ATLAS-CONF-2012-147

 10^{-1}

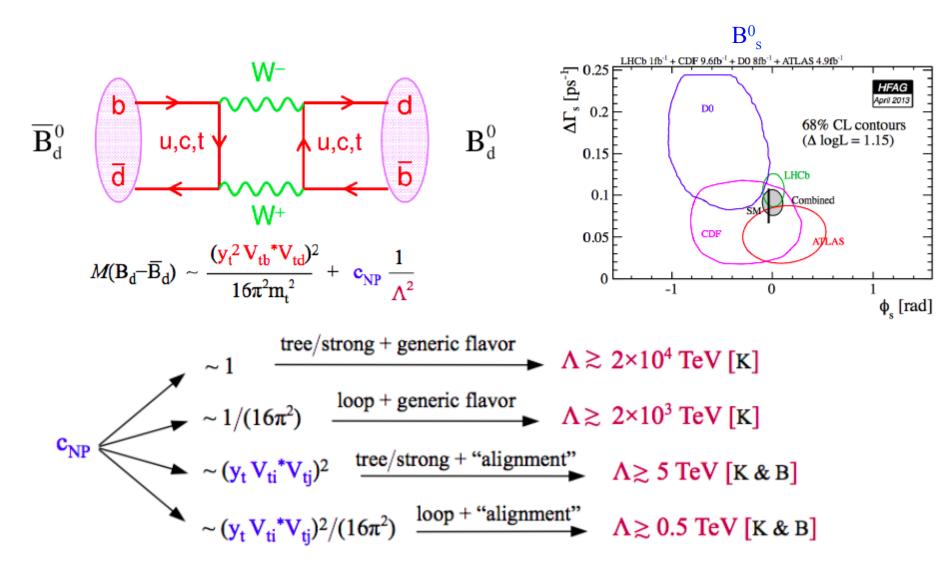
No sign of New Physics seen

What is not found...



Bounds on the scale of New Physics

Most stringent limits come from observables in BB mixing

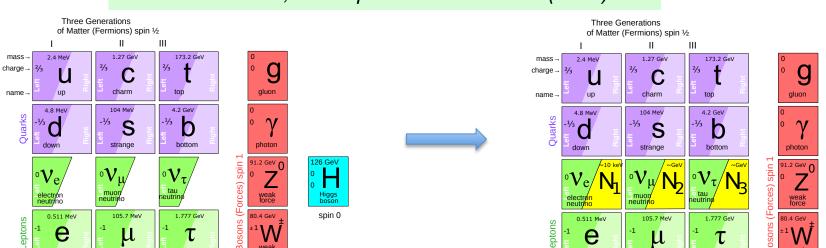


Theoretical motivation

- Discovery of the 126 GeV Higgs boson

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 The SM may work successfully up to Planck scale!
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See-saw generation of neutrino masses

Most general renormalisable Lagrangian of all SM particles (+3 singlets wrt the SM gauge group):

$$L_{singlet} = i\bar{N}_I \partial_\mu \gamma^\mu N_I - Y_{I\alpha} \bar{N}_I^c \tilde{H} L_\alpha - M_I \bar{N}_I^c N_I + h.c.$$

Yukawa term: mixing of N_I with active neutrinos to explain oscillations

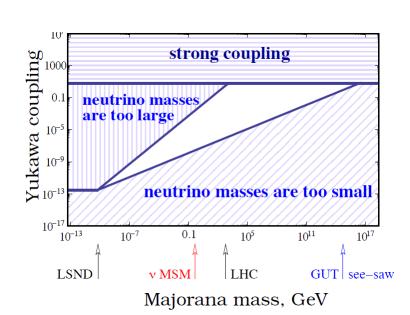
Majorana term which carries no gauge charge

The scale of the active neutrino mass is given by the see-saw formula: $m_
u \sim \frac{m_D^2}{M}$ where $m_D \sim Y_{I\alpha} v$ - typical value of the Dirac mass term

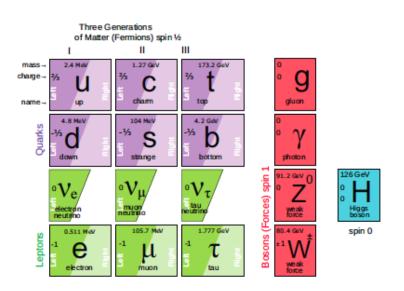
$$v \sim 246 \text{ GeV}$$

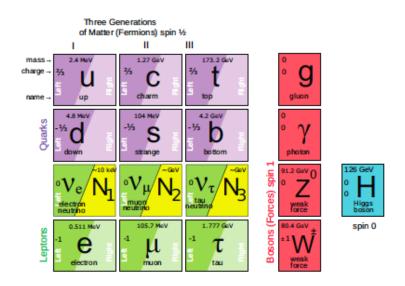
Example:

For $M \sim 1$ GeV and $m_v \sim 0.05$ eV it results in $m_D \sim 10$ keV and Yukawa coupling $\sim 10^{-7}$



The vMSM model





N = Heavy Neutral Lepton - HNL

Role of N_1 with mass in keV region: dark matter

Role of N_2 , N_3 with mass in 100 MeV – GeV region: "give" masses to neutrinos and produce baryon asymmetry of the Universe

Masses and couplings of HNLs

N₁ can be sufficiently stable to be a DM candidate, M(N₁)~10keV

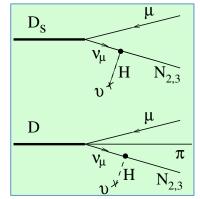


• $M(N_2) \approx M(N_3) \sim a$ few GeV \rightarrow CPV can be increased dramatically to explain Baryon Asymmetry of the Universe (BAU)

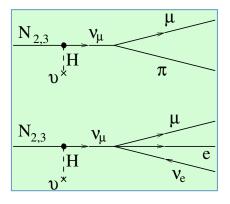
Very weak $N_{2,3}$ -to-v mixing (~ U^2) $\rightarrow N_{2,3}$ are much longer-lived than the SM particles

Example:

*N*_{2,3} production in charm

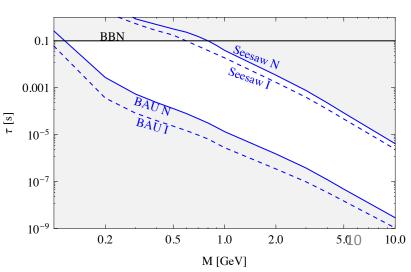


and subsequent decays



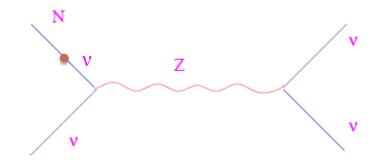
- Typical lifetimes > 10 μ s for $M(N_{2,3}) \sim 1$ GeV Decay distance O(km)
- Typical BRs (depending on the flavour mixing):

Br(N $\rightarrow \mu/e^{-} \rho^{+}$) ~ 0.1 – 50% Br(N $\rightarrow \mu/e^{-} \rho^{+}$) ~ 0.5 – 20% Br(N $\rightarrow \nu \mu e$) ~ 1 – 10%

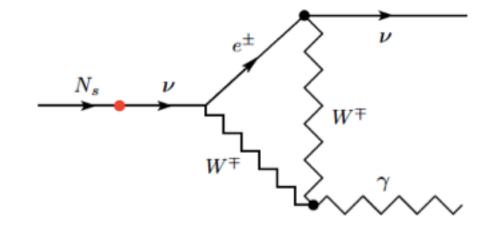


Dark Matter candidate HNL N₁

Yukawa couplings are small \rightarrow N can be very stable.



Main decay mode: $N \to 3\nu$. Subdominant radiative decay channel: $N \to \nu \gamma$.



Photon energy:

$$E_{\gamma}=rac{M}{2}$$

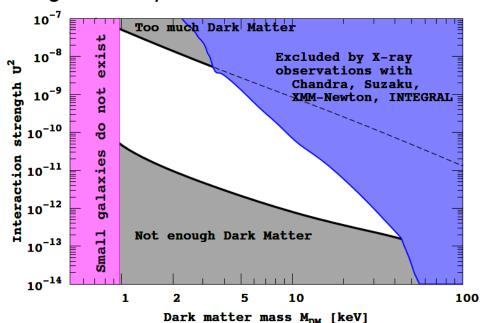
Radiative decay width:

$$\Gamma_{
m rad} = rac{9\,lpha_{
m EM}\,G_F^2}{256\cdot 4\pi^4}\,\sin^2(2 heta)\,{\it M_N}^5$$

Interaction strength

Constraints on DM HNL N₁

- ✓ **Stability** \rightarrow N_1 must have a lifetime larger than that of the Universe
- ✓ **Production** → N_1 are created in the early Universe in reactions $l \bar{l} \to \nu N_1 \quad q \bar{q} \to \nu N_1$ etc. Need to provide correct DM abundance
- ✓ Structure formation → N_1 should be heavy enough! Otherwise its free streaming length would erase structure non-uniformities at small scales (Lyman- α forest spectra of distant quasars and structure of dwarf galaxies)
- ✓ **X-ray spectra** → Radiative decays $N_1 \rightarrow \gamma v$ produce a mono-line in photon galaxies spectrum.



Searches for DM HNL N₁ in space

- Has been previously searched with XMM-Newton, Chandra, Suzaku, INTEGRAL
- Spectral resolution is not enough (required $\Delta E/E \sim 10^{-3}$)
- Proposed/planned X-ray missions with sufficient spectral resolution:

Astro-H



Athena+



LOFT



Origin/Xenia

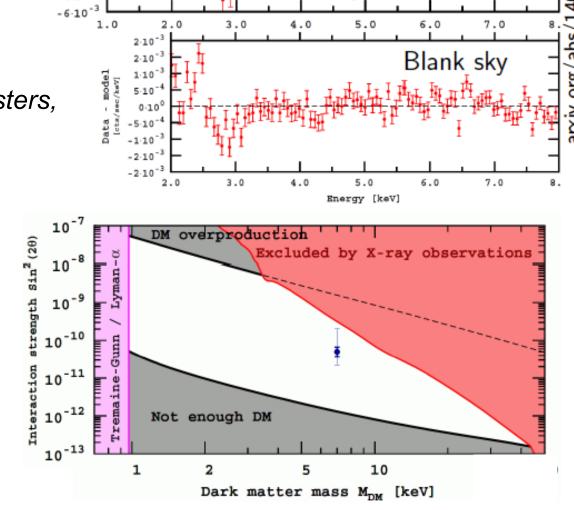


New line in photon galaxy spectrum ???

Two recent publications in arXiv:

- arXiv 1402.2301
 Detection of an unidentified
 emission line in the stacked
 X-ray spectrum of Galaxy Clusters,
 E, ~ 3.56 keV
- arXiv 1402.4119 An unidentified line in X-ray spectra of the Andromeda galaxy and Perseus galaxy cluster, E, ~ 3.5 keV

Will soon be checked by Astro-H with higher energy resolution



Andromeda

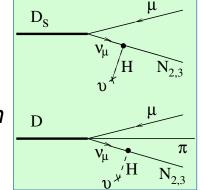
Masses and couplings of HNLs

M(N₂) ≈ M(N₃) ~ a few GeV→ CPV can be increased dramatically to explain
 Baryon Asymmetry of the Universe (BAU)

Very weak $N_{2,3}$ -to- ν mixing (~ U^2) $\rightarrow N_{2,3}$ are much longer-lived than the SM particles

Produced in semi-leptonic decays,

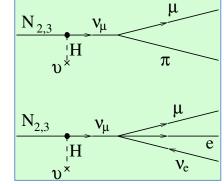
$$K \to \mu\nu$$
, $D \to \mu\pi\nu$, $B \to D\mu\nu$



$$\bullet \propto \sigma_D \times U^2$$

$$\bullet \ U_2^2 = U_{2,\nu_e}^2 + U_{2,\nu_\mu}^2 + U_{2,\nu_\tau}^2$$

and subsequent decays

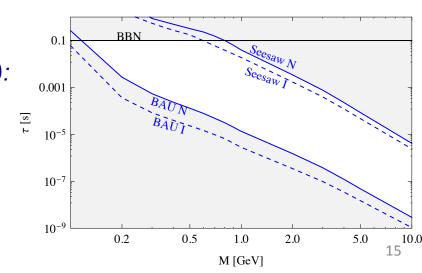


Example:

*N*_{2.3} production in charm

- Typical lifetimes > 10 μ s for $M(N_{2,3}) \sim 1$ GeV Decay distance O(km)
- Typical BRs (depending on the flavour mixing):

Br(N
$$\rightarrow \mu$$
/e π) ~ 0.1 – 50%
Br(N $\rightarrow \mu$ -/e⁻ ρ ⁺) ~ 0.5 – 20%
Br(N $\rightarrow \nu\mu$ e) ~ 1 – 10%



Baryon asymmetry

Sakharov conditions:

C and CP are not conserved in vMSM

6 CPV phases in the lepton sector and 1 CKM phase in the quark sector (to be compared with only one CKM phase in the SM)

Deviations from thermal equilibrium



- ✓ HNL are created in the early Universe
- ✓ CPV in the interference of HNL production and decay
- ✓ Lepton number goes from HNL to active neutrinos
- ✓ Then lepton number transfers to baryons in the equilibrium sphaleron processes

PS Explanation of DM with N_1 reduces a number of free parameters \rightarrow Degeneracy of $N_{2,3}$ masses is required to ensure sufficient CPV

Masses and couplings of HNLs

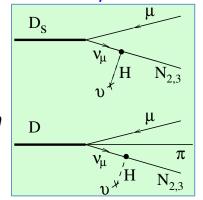
Very weak $N_{2,3}$ -to- ν mixing (~ U^2) $\rightarrow N_{2,3}$ are much longer-lived than the SM particles

• Produced in semi-leptonic decays,

$$K \to \mu\nu$$
, $D \to \mu\pi\nu$, $B \to D\mu\nu$

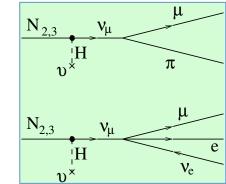


*N*_{2,3} production in charm



- $\bullet \propto \sigma_D \times U^2$
- $\bullet \ U_2^2 = U_{2,\nu_e}^2 + U_{2,\nu_\mu}^2 + U_{2,\nu_\tau}^2$

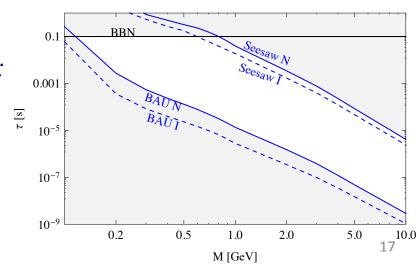
and subsequent decays



- Typical lifetimes > 10 μ s for $M(N_{2,3}) \sim 1$ GeV Decay distance O(km)
- Typical BRs (depending on the flavour mixing):

$$Br(N \to \mu/e \pi) \sim 0.1 - 50\%$$

 $Br(N \to \mu/e^- \rho^+) \sim 0.5 - 20\%$
 $Br(N \to \nu\mu e) \sim 1 - 10\%$

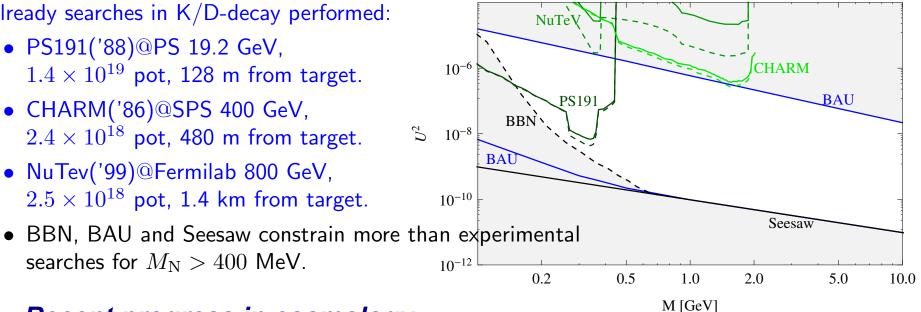


Experimental and cosmological constraints

Already searches in K/D-decay performed:

- PS191('88)@PS 19.2 GeV, 1.4×10^{19} pot, 128 m from target.
- CHARM('86)@SPS 400 GeV, 2.4×10^{18} pot, 480 m from target.
- NuTev('99)@Fermilab 800 GeV, 2.5×10^{18} pot, 1.4 km from target.

searches for $M_{\rm N} > 400$ MeV.

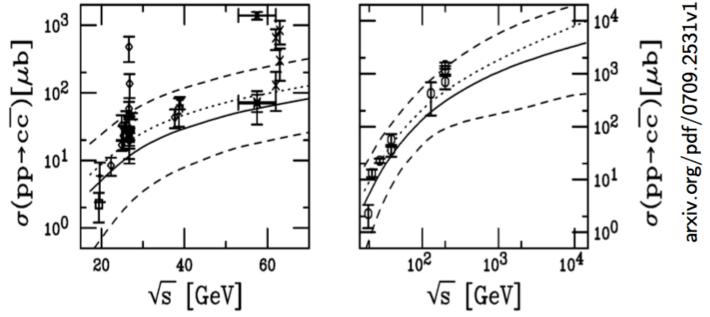


- Recent progress in cosmology
 - The sensitivity of previous experiments did not probe the interesting region for HNL masses above the kaon mass
 - Strong motivation to explore cosmologically allowed parameter space

Proposal for a new experiment at the SPS, SHIP to search for new long-lived particles produced in charm decays (see http://ship.web.cern.ch/ship)

Sensitivity for $N_{2.3} \propto U^4$

- PS-191: Used K decays → limited to 500 MeV (PLB 203 (1988) 332)
- Goal: Extend mass range to ~ 2 GeV by using charmed hadron decays
- B-decays: 20÷100 smaller σ, and B \rightarrow Dµv, i.e. limited to 3 GeV still



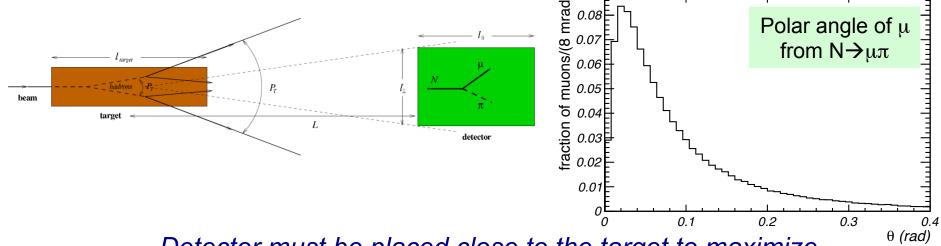
Where to produce charmed hadrons?

LHC ($\sqrt{s} = 14 \text{ TeV}$): with $1 \ ab^{-1} \ (\sim 3\text{-}4 \text{ years})$: $\sim 2 \times 10^{16} \text{ in } 4\pi$ SPS (400 GeV p-on-target (pot) $\sqrt{s}=27$ GeV): with 2×10^{20} pot (~3 -4 years): $\sim 2 \times 10^{17}$

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Experimental requirements

- Search for HNL in Heavy Flavour decays
 - Beam dump experiment at the SPS with a total of 2×10²⁰ protons on target (pot) to produce a large number of charmed hadrons CNGS: 1.8 x 10²⁰ pot, 2011 run: 4.8 x 10¹⁹ pot
- HNLs produced in charm decays have significant P_T



Detector must be placed close to the target to maximize geometrical acceptance

Effective (and "short") muon shield is essential to reduce
 → muon-induced backgrounds (mainly from short-lived resonances accompanying charm production)

Secondary beam-line

(different from a conventional neutrino facility)

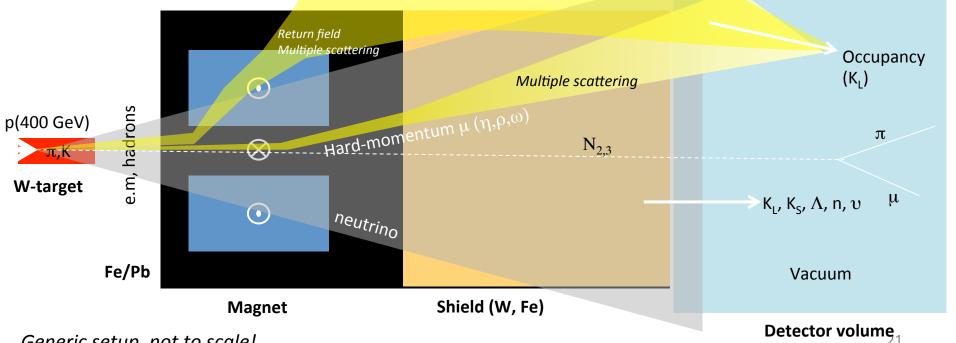
Initial reduction of beam induced backgrounds

- Heavy target (50 cm of W)
- Hadron absorber (5m Fe)
- Muon shield: optimization of active and passive shields is underway

Acceptable occupancy <1% per spill of 5×10¹³ p.o.t.

spill duration 1s \rightarrow < 50×10⁶ muons spill duration 1ms \rightarrow < 50×10³ muons spill duration 10 μ s \rightarrow < 500 muons

Low-mid-momentum μ from fast decays of π ,K

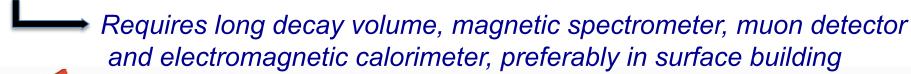


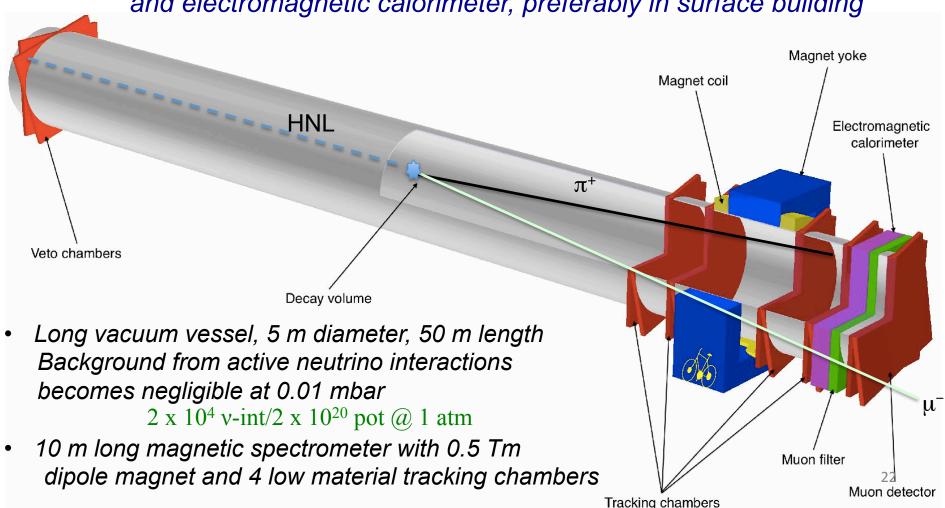
Generic setup, not to scale!

Detector concept

(based on existing technologies)

• Reconstruction of the HNL decays in the final states: $\mu^-\pi^+$, $\mu^-\rho^+$ & $e^-\pi^+$

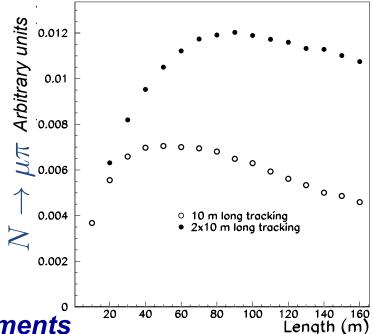




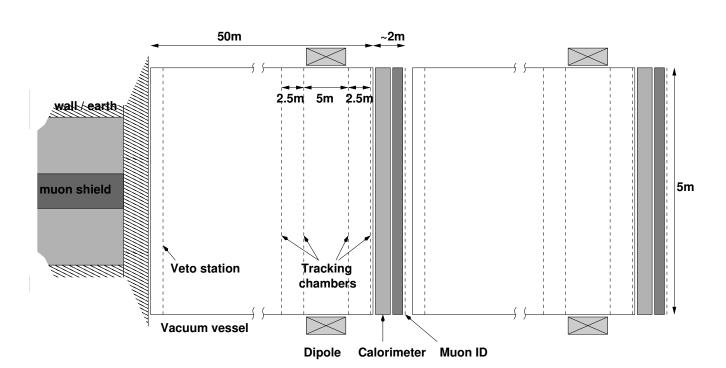
Detector concept (cont.)

Geometrical acceptance

- Saturates for a given HNL lifetime as a function of detector length
- The use of two magnetic spectrometers increases the acceptance by 70%



Detector has two almost identical elements

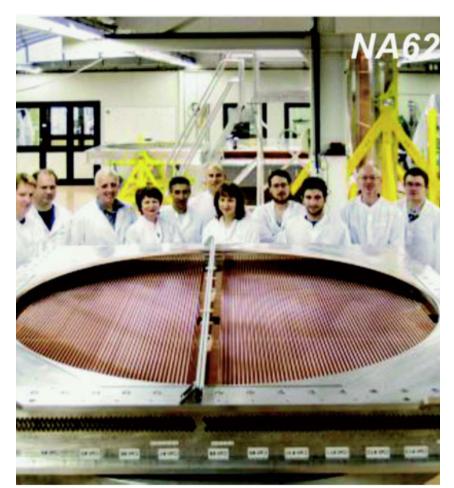


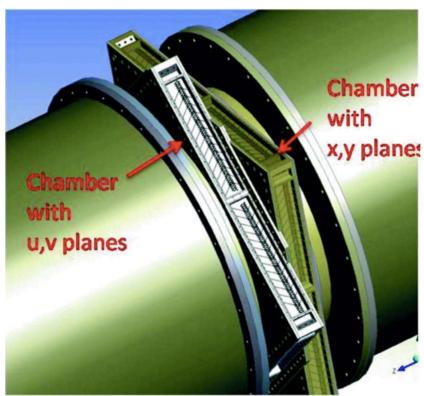
23

Tracking chambers

Same as NA62 $(K^+ \to \pi^+ \nu \bar{\nu})$ 2m diameter vessel at 0.01 µbar 10 mm diameter straws made of PET \to working well in vacuum $X/X_0 = 0.5\%$ for 4 view stations

120 μm resolution/straw



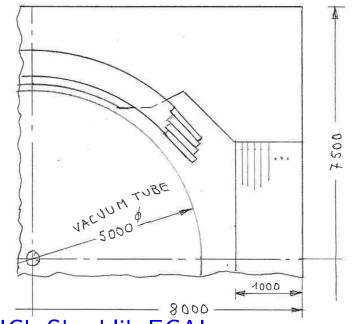


Magnet and e.m. calo

- With X/X0=0.5~% chambers: modest 0.5 Tm
- Need $\sim 20~{\rm m}^2$ aperture.

LHCb magnet: 4 Tm, 16 m^2 aperture Preliminary calculations (W.Flegel):

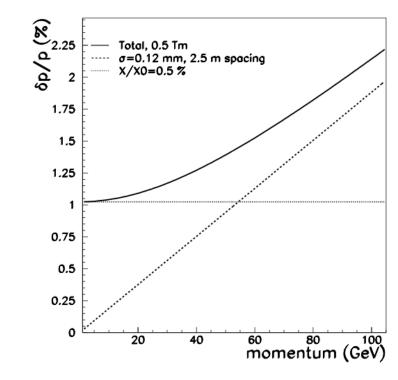
- Needs 30 % less iron/yoke than LHCb.
- Consumes 3 times less power.



LHCb Shashlik ECAL:

• $6.3 \times 7.8 \text{ m}^2$

• $\frac{\sigma(E)}{E} < 10\% / \sqrt{E} \oplus 1.5\%$





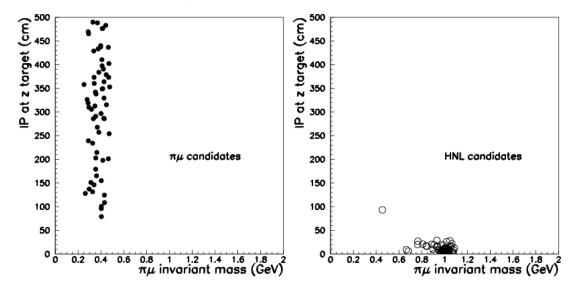
Mass resolution

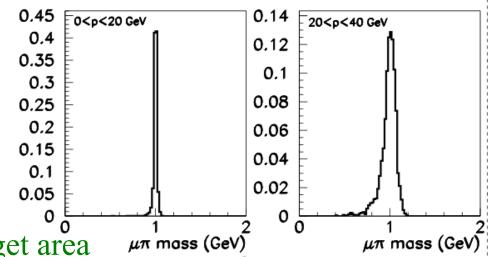
Expected resolution for 1 GeV N $\rightarrow \mu\pi$ $< p_N > = 25 \text{ GeV}$

• $\sigma_{mass} \sim 40 \text{ MeV for } P < 20 \text{ GeV}$

K₁ background suppression

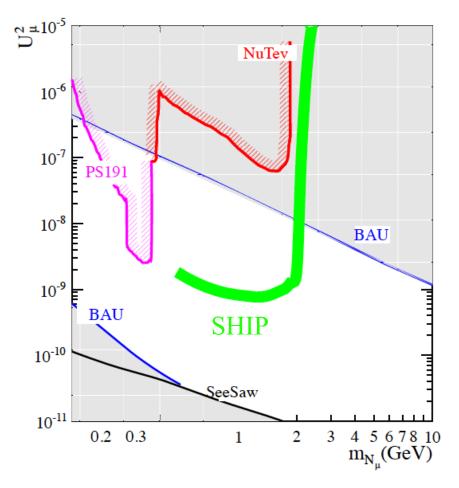
- Use pointing of candidates to target area μπ moss (GeV)
 Detect CC via extra μ in coincidence with μπ, upstream tagger
- Instrument μ -filter to tag CC/NC shower





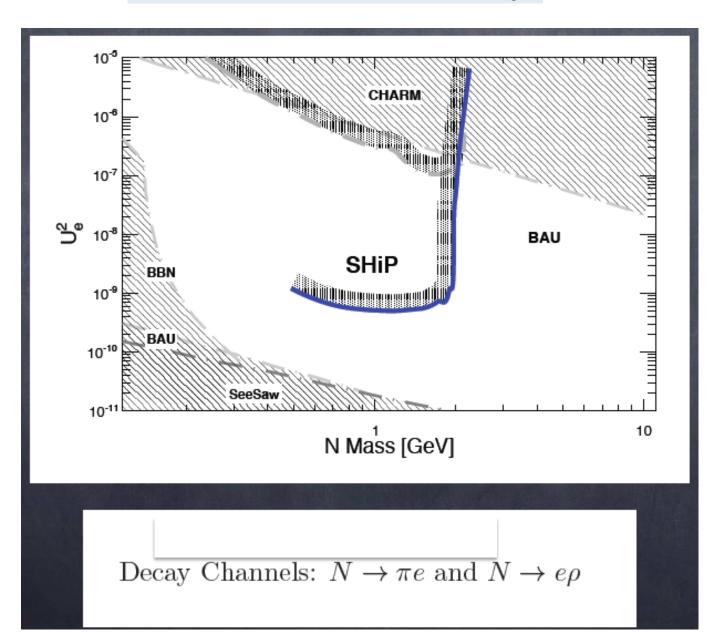
Expected event yield

Assuming $U_{\mu}^{2} = 10^{-7}$ (corresponding to the strongest experimental limit currently for $M_{N} \sim 1$ GeV) and $\tau_{N} = 1.8 \times 10^{-5}$ s $\sim 12k$ fully reconstructed $N \rightarrow \mu^{-}\pi^{+}$ events are expected for $M_{N} = 1$ GeV

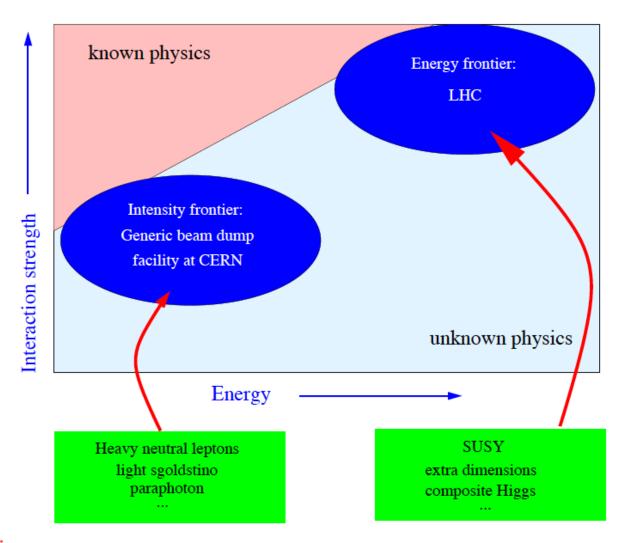


120 events for cosmologically favoured region: $U_{\mu}^{2} = 10^{-8} \& \tau_{N} = 1.8 \times 10^{-4} \text{s}$

Sensitivity to HNL: U_e



Physics case for general beam dump facility



hidden sector:

HNL: baryon asymmetry of the Universe, dark matter, neutrino masses sgoldstino, light neutralino: SUSY

paraphoton: mirror matter, dark matter

Physics case for general beam dump facility

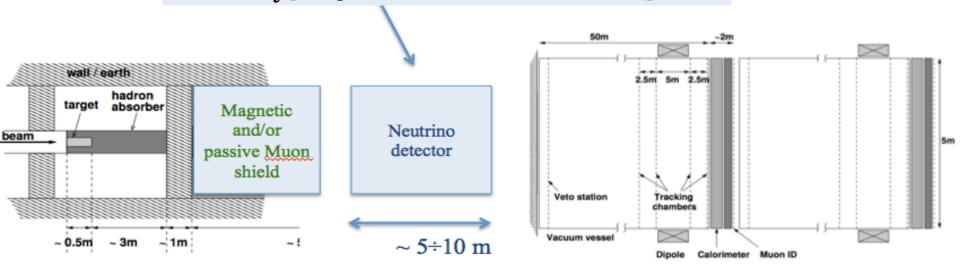
 \checkmark Study of v_{τ} interactions (guarantied SM physics)

Ideally suited since v_{τ} is produced in $D_s \rightarrow \tau v_{\tau}$ with similar to HNL kinematics. DONUT 9 events, OPERA 4 events (from oscillations)

- ✓ Search for any weakly interacting yet unstable particles produced in very rare charm (or hyperon) decays such as low mass SUSY or paraphotons or ...
- \checkmark Review of the SHIP sensitivities for v_{τ} physics and wide class of models with hidden portals is ongoing

Expect significant improvement of currently available measurements and constrains everywhere!

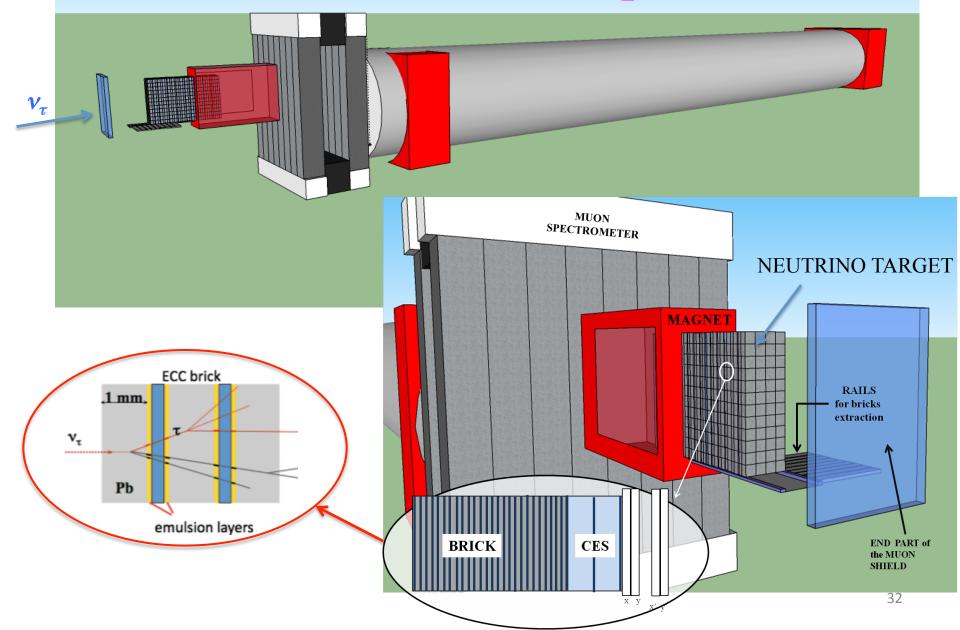
SM: v_{τ} physics with 2×10²⁰ pot



A powerful source of neutrinos originated by charm decays, $D_s \rightarrow \tau \, v_\tau$ Present configuration, flux rates: $v_e \sim 7.1\%$, $v_\mu \sim 92.5\%$, $v_\tau \, 0.4\%$ Anti- $v_\mu / v_\mu \sim 62\%$, Anti- $v_e / v_e \sim 1$, Anti- $v_\tau / v_\tau = 1$

- Expect $\sim 3400 \, v_{\tau}$ interactions in 6 tons emulsion target ($\sim 5\%$ of the OPERA emulsion films)
- Tau neutrino and anti-neutrino (never seen) physics
- Charm physics with neutrinos and anti-neutrinos
- v_{τ} : a probe of non-standard interactions PRD 87 (2013) 013002

A neutrino detector upstream

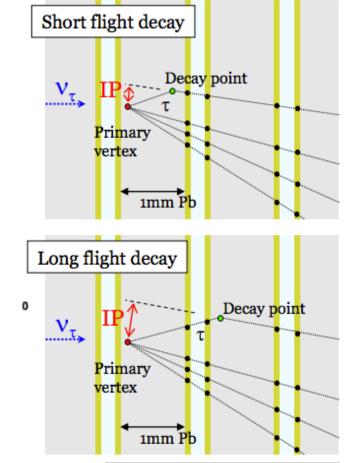


Detector choice

- Neutrino target: high density material (lead) alternated with emulsions (Emulsion Cloud Chamber).
- Vertex detector: emulsions as micrometric tracking device to see both the neutrino interaction and τ decay vertices.
- Nuclear fragments \rightarrow reduce background from hadronic int. (track angles up to $\vartheta \sim \pi/2$)
- Momentum measurement of hadrons (several GeV) by Multiple Coulomb Scattering with $\sim 20\%$ accuracy in the several GeV range
- High resolution tracker to provide the time stamp and reconstruct the μ track in the target ($\sim 100 \ \mu m$)
- Muon spectrometer to measure charge and momentum. Focus on the acceptance to low energy and large angle μ (as μ from ν_{μ} int., charmed interactions particularly important)

01/07/14

Detector Principle: hybrid detector with modular structure **TRACKERS** COMPACT EMULSION **SPECTROMETER** rohacell **ECC BRICK** lead τ ν_{τ} 1 mm emulsion 15 mm $\otimes B$ layers



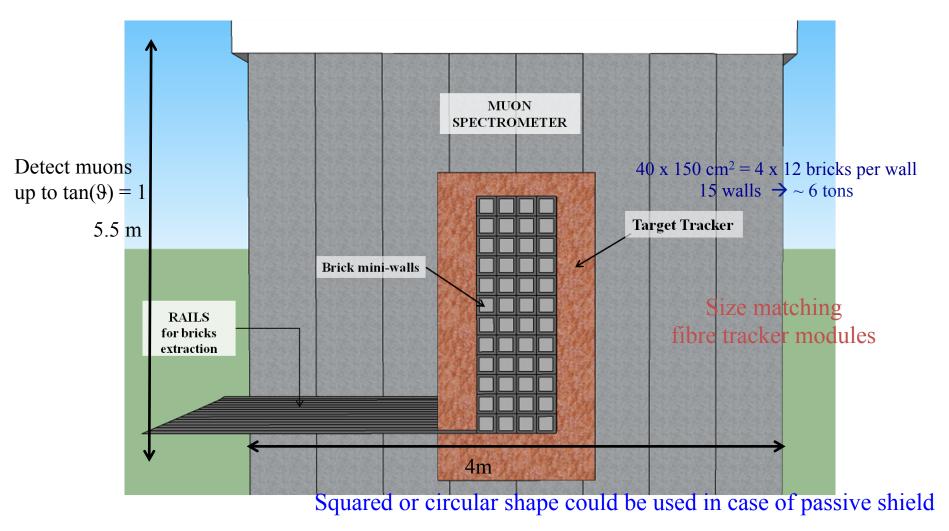
- Nuclear emulsions acting as trackers with micrometric resolution
- Detect τ-lepton production and decay vertices
- An option: Compact emulsion spectrometer to measure charge
- Electronic detectors to provide the "time stamp" and reconstruct μ charge/momentum

τ DECAY CHANNEL	BR (%)
$\tau\!\to\!\!\mu$	17.7
$\tau \rightarrow e$	17.8
$\tau \mathop{\to} h$	49.5
τ →3h	15.0

01/07/14

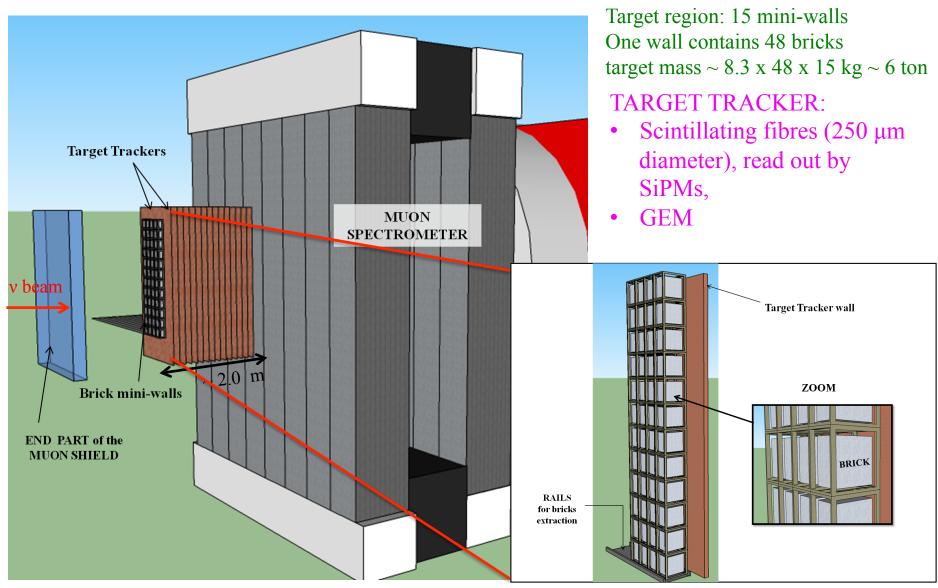
Detector design, beam view

Magnetic muon shield → elongate the target along the magnetic field → rectangular shape



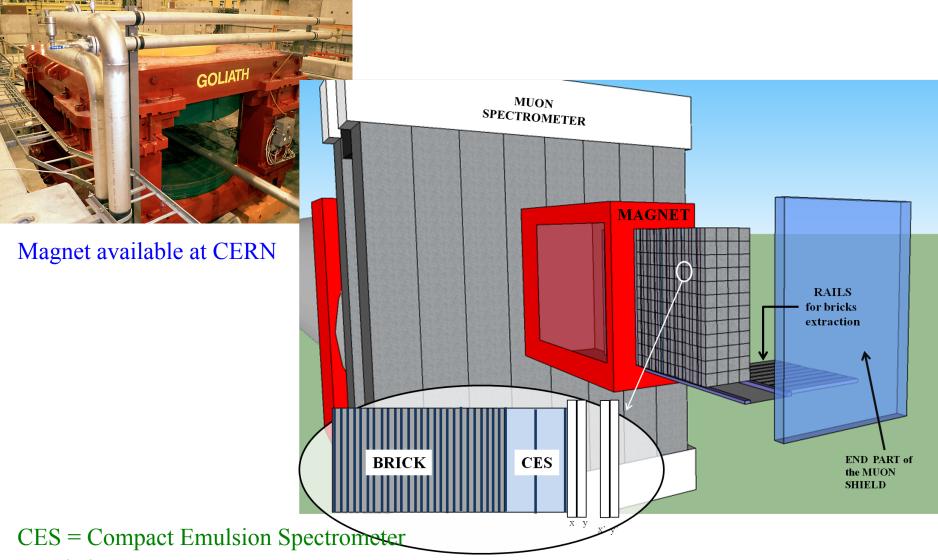
01/07/14

Detector design: first option



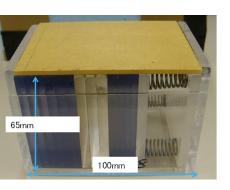
Wall thickness ~ 13 cm: 8 cm brick + 5 cm tracker plane

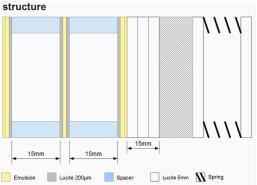
Detector design: second option with a magnetized target



01/07/14

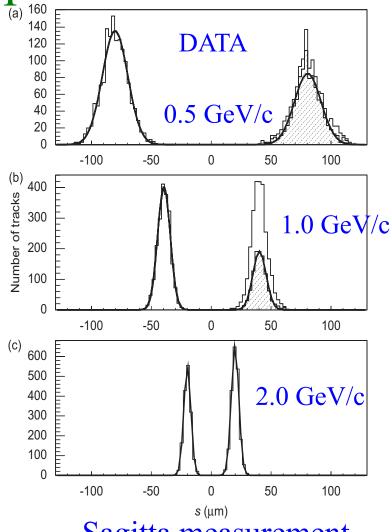
Compact emulsion spectrometer CES





3 emulsion films interleaved with 1.5 cm air gap in a magnetic field (~1T), 3cm thick device, H. Shibuya et al NIM A592 (2008) 56

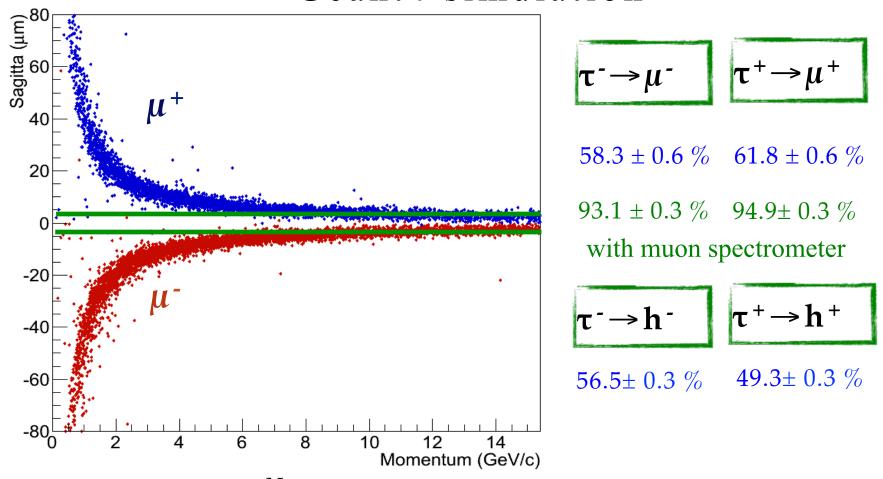
- Emulsion films alternated by low density material (Rohacell, 30÷100 kg/m³)
- With a good alignment, by means of high energy μ , the charge of 10 GeV muons detectable ($\pm 4~\mu m$ displacement).



Sagitta measurement

01/07/14

Charge Measurement with CES Geant4 simulation



Statistical gain due to
$$\frac{\sum_{i=1}^{N}br_{i}\varepsilon_{i}}{br_{\mu}\varepsilon_{\mu}}\simeq\frac{18\cdot0.95+50\cdot0.53+15\cdot0.53^{2}}{18\cdot0.9}\simeq 3$$

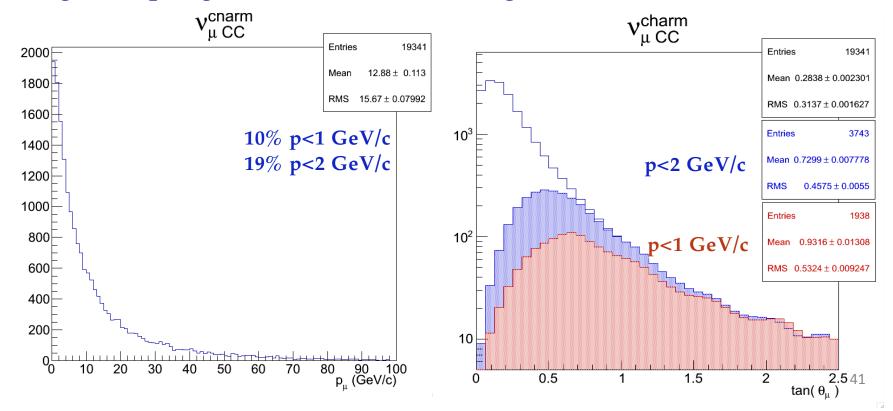
Real-time electronic detectors

- Time stamp of the event, association of tracks in emulsions with tracks in the target tracker
- In an environment with high interaction rate, a high resolution (~ 100 μm) is needed. Two identified options: 250 μm scintillating fibre tracker and GEM
- The muon spectrometer has to identify μ with high efficiency
- Charm background from μ mis-identification.
- Charge measurement is less important than μ identification in this respect
- Charge measurement important for v_{τ} and anti- v_{τ} separation

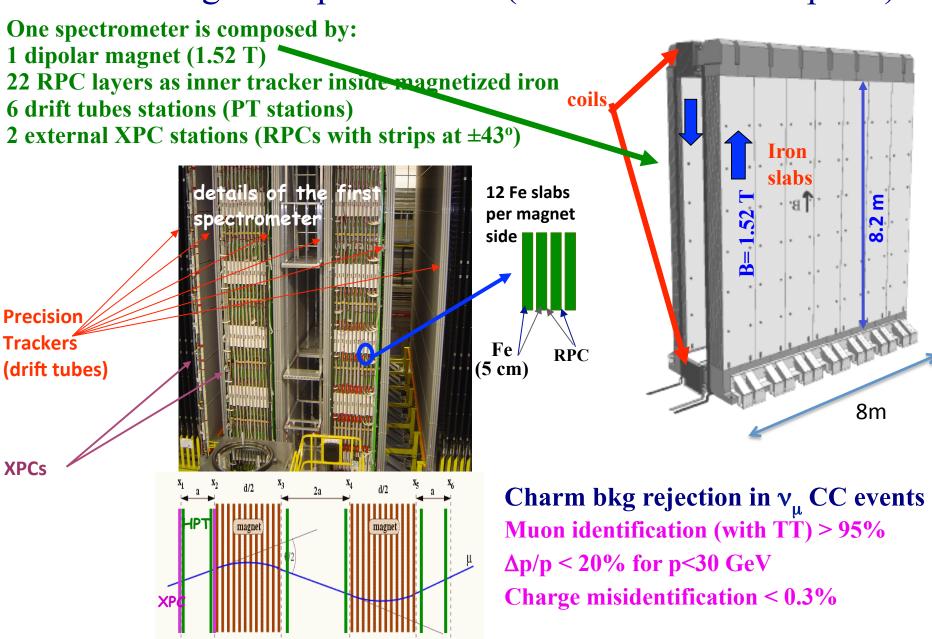
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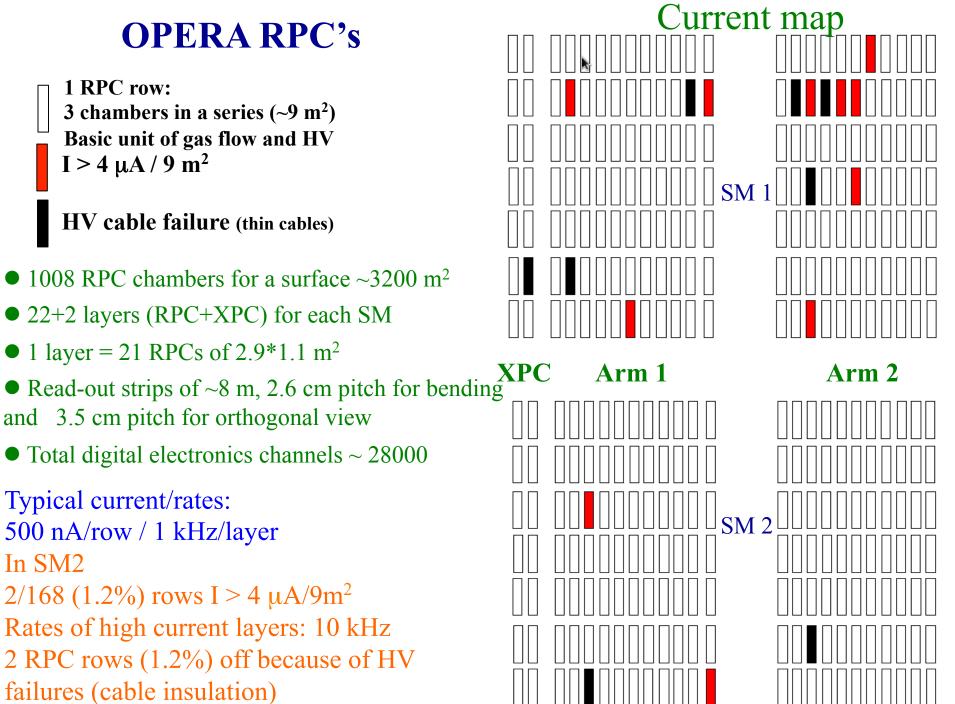
Muon detector requirements

- Detector performances driven by the background rejection → minimise muon misidentification
- Soft and large angle muons → difficult to be identified
- Large acceptance and fine graining to identify P < 2 GeV
- $4 \times 5.5 \text{ m}^2$ to detect angles up to $\tan(9) \le 1$
- High sampling to use momentum/range correlation



The magnetic spectrometer (OPERA one is an option)

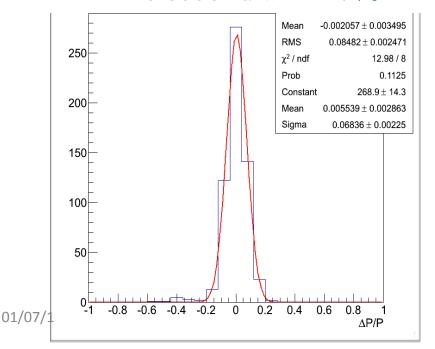




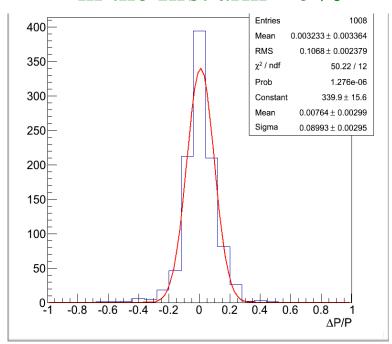
Muon identification

- Same configuration as OPERA, 5 cm iron slabs interleaved with active (RPC) trackers, 4 x 5.5 m²
- Exploit the momentum-range correlation
- $\varepsilon_{\mu} \sim 92\%$ for charm v_{μ} events (~3% out of acceptance)
- ε_{μ} ~ 94% for $\tau \rightarrow \mu$ events

Momentum resolution by range in the second arm $\sim 7\%$



Momentum resolution by range in the first arm $\sim 9\%$



Signal to noise

In v_{τ} interactions

$$\frac{\nu_{\tau}}{\bar{\nu}_{\mu_{CC}}^{charm}} = \frac{\int \phi_{\nu_{\tau}} \cdot \sigma_{\nu_{\tau_{CC}}} dE}{\int \phi_{\bar{\nu}_{\mu}} \cdot \sigma_{\bar{\nu}_{\mu_{CC}}^{charm}} \cdot \eta_{mis} dE} \sim \frac{6 \times 10^{-3} \times 300}{8 \times 10^{-2}} \sim 20$$

In anti- v_{τ} interactions

$$\frac{\bar{\nu}_{\tau}}{\nu_{\mu_{CC}}^{charm}} = \frac{\int \phi_{\bar{\nu}_{\tau}} \cdot \sigma_{\bar{\nu}_{\tau_{CC}}} dE}{\int \phi_{\nu_{\mu}} \cdot \sigma_{\nu_{\mu_{CC}}^{charm}} \cdot \eta_{mis} dE} \sim \frac{3.7 \times 10^{-3} \times 90}{8 \times 10^{-2}} \sim 4$$

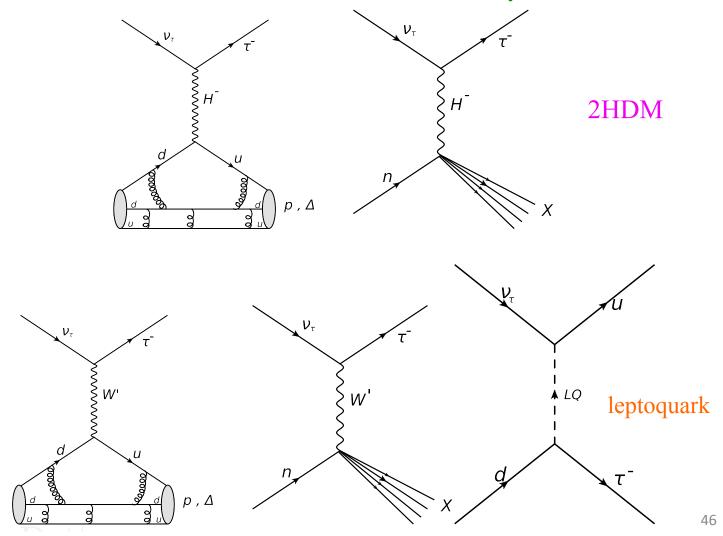
Effect of the angular acceptance not accounted Kinematical selection not included

In OPERA Signal/noise ~ 10 on average Signal/noise ~ 30 in the leptonic channels ($\epsilon_{\mu} \sim 98\%$)

New physics in v_{τ} scattering A. Datta et al., PRD 87 (2013) 013002

Several new physics models may contribute to v_{τ} scattering

01/07/14



W' in deep-inelastic scattering

$$\mathcal{L} = \frac{g}{\sqrt{2}} V_{f'f} \bar{f}' \gamma^{\mu} (g_L^{f'f} P_L + g_R^{f'f} P_R) f W'_{\mu} + h.c.$$

$$r_{W'} = \frac{\sigma_{NP}}{\sigma_{SM}}$$

$$E_{\nu} = 17 \text{ GeV}$$

$$E_{\nu} = 17 \text{ GeV}$$

$$M_{W'} = 200 \text{ GeV}$$

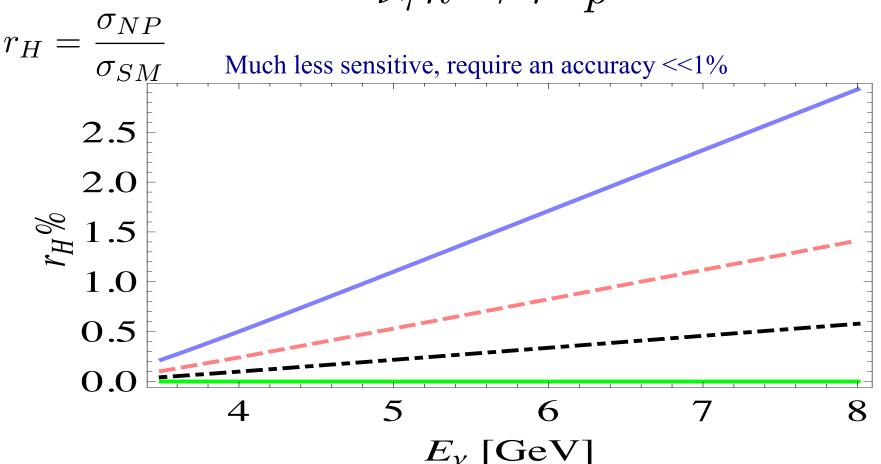
$$E_{\nu} = 1600 \text{ GeV}$$

$$E_{\nu} = 1600 \text{ GeV}$$

The lines correspond to different sets of the couplings $(g_L^{\tau\nu_\tau}, g_L^{ud}, g_R^{ud})$.

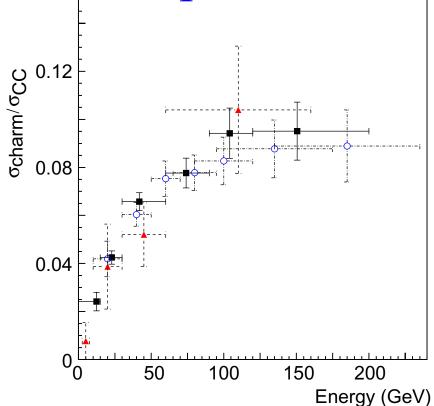
Charged Higgs in quasi-elastic scattering

$$\nu_{\tau} n \to \tau^- p$$



The lines correspond to different values of $\tan \beta$ (ratio of VEV's)

Expected v-induced charm events



CHORUS, New Journal of Physics 13 (2011) 093002

$$f(charm) = \frac{\int \Phi_{\nu_{\mu}} \sigma_{\nu_{\mu}}^{CC} \left(\frac{\sigma_{charm}}{\sigma_{\nu_{\mu}}^{CC}}\right) dE}{\int \Phi_{\nu_{\mu}} \sigma_{\nu_{\mu}}^{CC} dE} \approx 2\%$$

In v_{μ} interactions: $\sigma_{charm} \sim 2\%$, CC/(NC+CC) = 0.7 $\epsilon_{decay} \sim 40\% \Rightarrow 2 \times 10^6 \times 0.02 \times 0.7 \times 0.4 \sim 11000$ charm

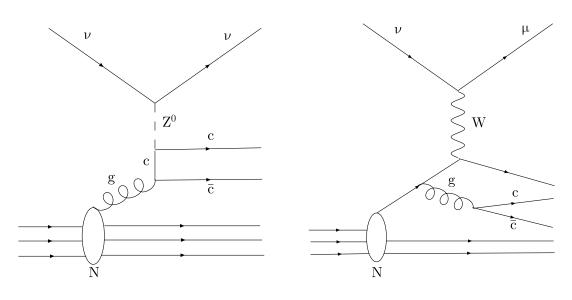
In anti- ν_{μ} interactions: anti- ν_{μ} / ν_{μ} ~ 63%, $\sigma_{\nu\text{-bar}}$ / σ_{ν} = 0.5 ~ 3500 events expected, 32 observed by CHORUS

Charm production in v_e interactions

$$f(charm) = \frac{\int \Phi_{\nu_e} \sigma_{\nu_e}^{CC} \left(\frac{\sigma_{charm}}{\sigma_{\nu_e}^{CC}}\right) dE}{\int \Phi_{\nu_e} \sigma_{\nu_e}^{CC} dE} \approx 4.6\%$$

Charm background in the v_{τ} search from electron mis-identification ($v_{e} \sim 10\%$)

Associated charm production



Only gluon bremsstrahlung in CC interactions
Both processes in NC

Fig. 1. Feynman diagrams for boson-gluon fusion (left) and gluon bremsstrahlung processes (right)

3 events in NC and 1 in CC processes, another CC event observed with a different search

$$\frac{\sigma(c\bar{c}\nu)}{\sigma_{\rm NC}^{\rm DIS}} = (3.62^{+2.95}_{-2.42}({\rm stat}) \pm 0.54({\rm syst})) \times 10^{-3}. \qquad \qquad \frac{\sigma(c\bar{c}\mu^-)}{\sigma_{\rm CC}} < 9.69 \times 10^{-4} \,,$$

~ 30 events expected

Double charm in NC interactions sensitive to the existence of Z' bosons

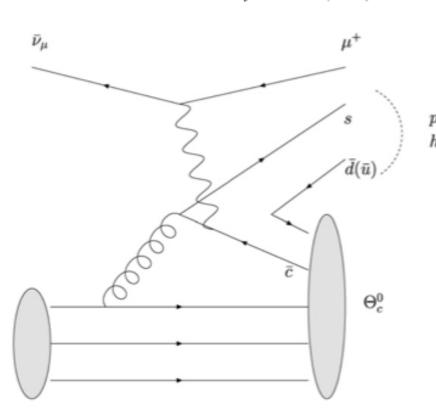
Charm production induced by v_{τ} interactions is competing Same topology, kinematics to be exploited

Search for multi-quark states in v interactions: charmed pentaquarks

Weakly decaying charmed hadron (below 2.8 GeV)

Unlike other processes like e+ e- scattering, the θ^0_c production in anti-neutrino interactions is favoured by the presence of three valence quarks

G. De Lellis et al. / Nuclear Physics B 763 (2007) 268-282



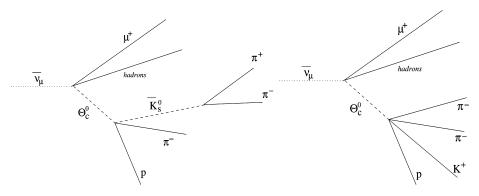


Fig. 2. Decay topology of Θ_c^0 events produced in $\bar{\nu}_{\mu}$ interactions with two and four prongs.

$$\sigma_{\Theta_c^0}/\sigma_{\bar{\nu}} < 0.039$$
 at 90% C.L.

lifetime equal to $0.5\tau_{D^0}$

Not a tight bound, larger than D⁰ prod, Limited by the anti-nu statistics

Fig. 1. Θ_c^0 production in $\bar{\nu}_\mu$ interactions.

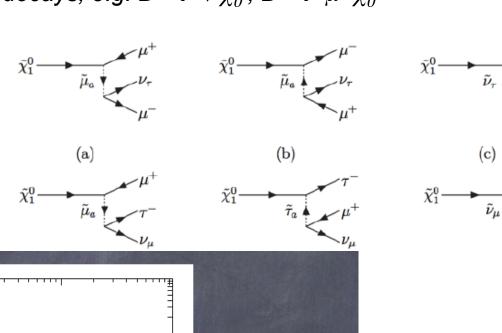
New physics portals

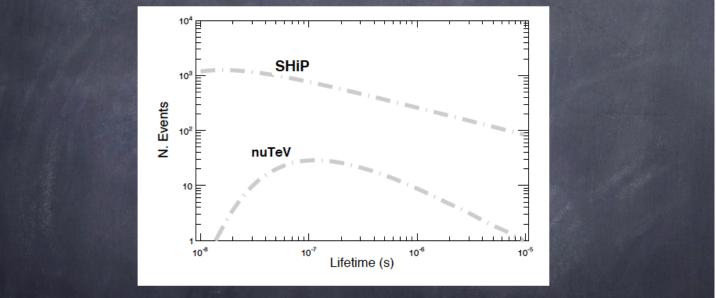
- Fermion portal
 - HNL in general
 - SUSY: Low mass neutralinos
- Scalar portal
 - Light inflaton
- Vector portal
 - Paraphoton
- Axion portal
 - SUSY: Light sgoldinos

In general: sensitive to very weakly interacting and long lived particles

Light neutralinos

- ✓ Can be produced in charm decays, e.g. $D^0 \rightarrow v \chi_0$, $D^+ \rightarrow \mu^+ \chi_0$
- ✓ Decay final states:





Number of events for $BR(B^+ \to \ell^+ \chi_0) = 10^{-7}$ which is equivalent to $BR(D^+ \to \ell^+ \chi_0) = 10^{-9}$ and $\chi \to \ell^+ \ell^- \nu$

Hidden photons

• Assuming no lighter hidden particles, γ' decay into SM particles through a virtual photon:

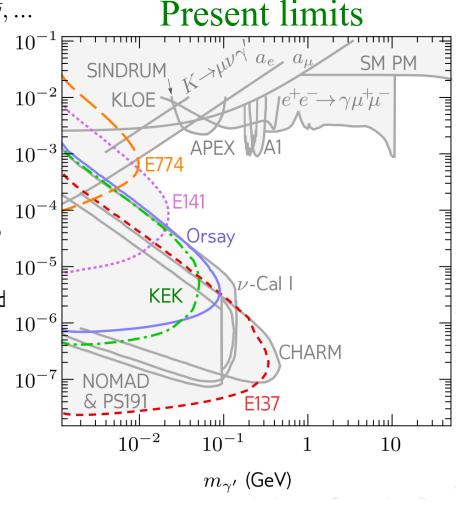
$$\gamma' \to e^+ e^-, \quad \mu^+ \mu^-, \quad q\bar{q}, \dots$$

- decay length $c au\sim arepsilon^{-2}m_{\gamma'}^{-1}$
- cosmological constraints (nucleo-synthesis): $\tau < 0.1 \text{ s} \Rightarrow \varepsilon^2 m_{\gamma'} > 10^{-21} \text{ GeV}$

γ' production

- proton bremsstrahlung:
 - initial-state radiation from the incoming proton, followed by a hard proton-nucleus interaction
- secondary particles decay:

Mass interval (GeV)	Process	$n_{\gamma'}/p.o.t$
$m_{\gamma'} < 0.135$	$\pi^0 \to \gamma \gamma'$	$\varepsilon^2 \times 5.41$
$0.135 < m_{\gamma'} < 0.548$	$\eta o \gamma \gamma'$	$\varepsilon^2 \times 0.23$
$0.548 < m_{\gamma'} < 0.648$	$\omega o \pi^0 \gamma'$	$\varepsilon^2 \times 0.07$
$0.648 < m_{\gamma'} < 0.958$	$\eta' o \gamma \gamma'$	$\varepsilon^2 \times 10^{-3}$



54

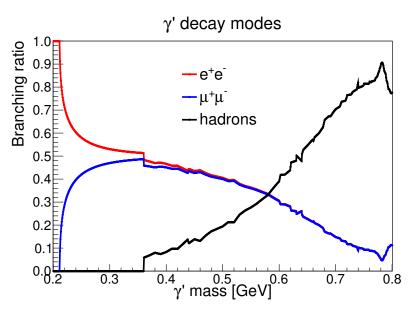
$$\gamma'$$
 decays

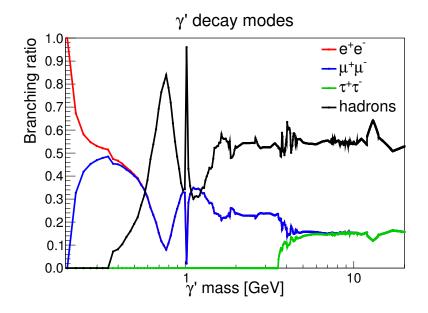
Phys. Lett. B731 (2014) 320

$$\Gamma\left(\gamma' \to \ell^+ \ell^-\right) = \varepsilon^2 \frac{\alpha_{QED}}{3} m_{\gamma'} \times \sqrt{1 - \frac{4m_\ell^2}{m_{\gamma'}^2}} \left(1 + \frac{2m_\ell^2}{m_{\gamma'}^2}\right)$$

$$\Gamma\left(\gamma' \to q\bar{q}\right) = \varepsilon^2 \frac{\alpha_{QED}}{3} m_{\gamma'} \times R\left(m_{\gamma'}\right)$$

$$R\left(\sqrt{s}\right) = \frac{\sigma\left(e^{+}e^{-} \to \text{hadrons}\right)}{\sigma\left(e^{+}e^{-} \to \mu^{+}\mu^{-}\right)}$$



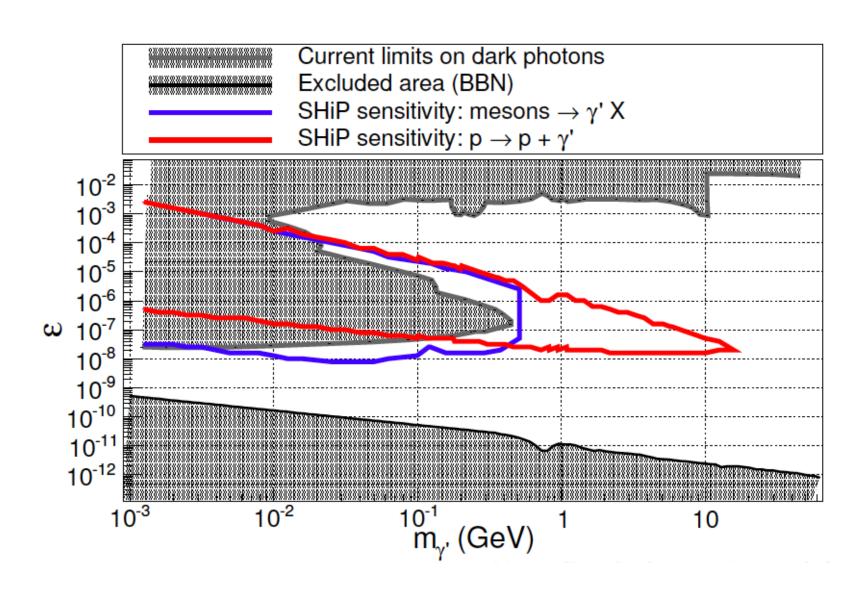


Only e^+e^- and $\mu^+\mu^-$ decays considered here

R used to compute the dark photon lifetime

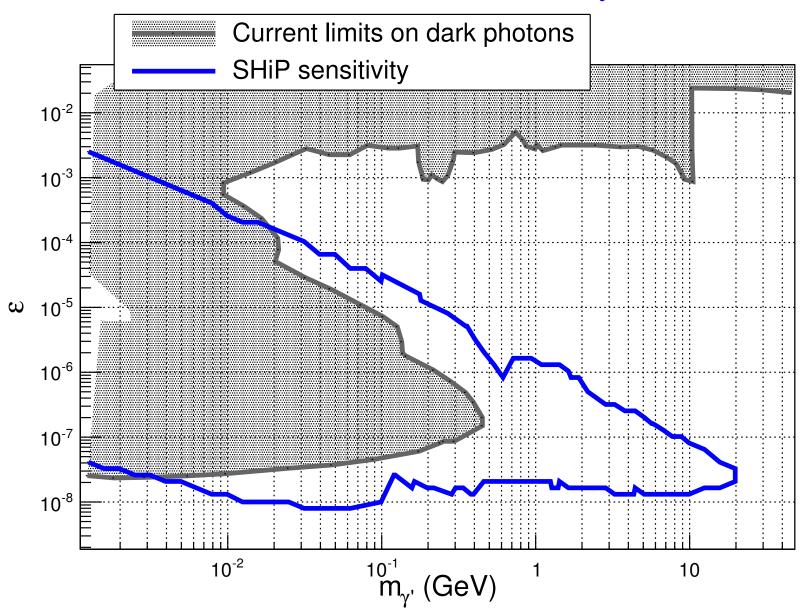
Sensitivity to Dark photons

Preliminary result: 90% C.L. limit



Sensitivity to Dark photons

90% C.L. Preliminary



Status of the SPSC review

- Oct 2013: submitted our EOI: CERN-SPSC-2013-024; arXiv:1310.1762; SPSC-EOI-010. 2013
- SPSC assigned 4 referees, who came with a list of questions.
- 3/1/2014: answers to questions: snoopy.web.cern.ch/snoopy/EOI/SPSC-EOI-010_ResponseToReferees.pdf
- 15/1/2014: SPSC discussed our proposal.

17/1/2014: The official feedback from the Committee is as follows :

"The Committee **received with interest** the response of the proponents to the questions raised in its review of EOI010.

The SPSC **recognises** the interesting physics potential of searching for heavy neutral leptons and investigating the properties of neutrinos.

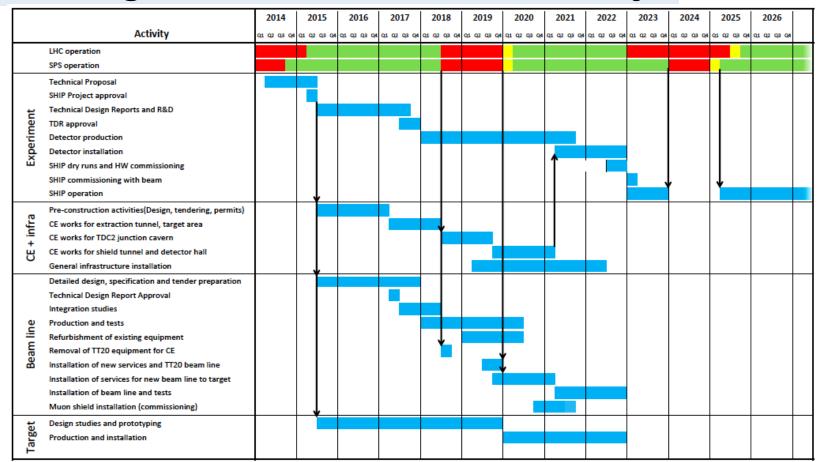
Considering the large cost and complexity of the required beam infrastructure as well as the significant associated beam intensity, such a project should be designed as a general purpose beam dump facility with the broadest possible physics programme, including maximum reach in the investigation of the hidden sector.

To further review the project the Committee **would need** an extended proposal with further developed physics goals, a more detailed technical design and a stronger collaboration."

Cheers,

Gavin, Lau, Matthew and Thierry (for the SPS Committee).

Planning schedule of the SHIP facility



A few milestones:

- √ Form SHIP collaboration
- √ Technical proposal
- ✓ Technical Design Report
- ✓ Construction and installation
- ✓ Commissioning
- ✓ Data taking and analysis of 2×10²⁰ pot → 2023 2027

- → June-September 2014
- **→** 2015
- → 2018
- → 2018 2022
- → 2022

Conclusion and Next steps

- The proposed experiment will search for NP in the largely unexplored domain of new, very weakly interacting particles with masses below the Fermi scale
- Detector is based on existing technologies
 Ongoing discussions of the beam lines with experts
- The impact of HNL discovery on particle physics is difficult to be overestimated!
- The proposed experiment complements the searches for NP at the LHC and in neutrino physics
- Tau neutrino physics and charm physics with neutrinos

A Collaboration is currently being setup. The first collaboration meeting was held in June. Let us know if you are interested to join!