

# Axion Physics in a Nut Shell

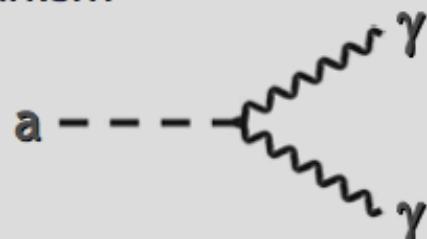
## Particle-Physics Motivation

CP conservation in QCD by Peccei-Quinn mechanism

→ Axions  $a \sim \pi^0$

$$m_\pi f_\pi \approx m_a f_a$$

For  $f_a \gg f_\pi$  axions are “invisible” and very light



## Solar and Stellar Axions

Axions thermally produced in stars, e.g. by Primakoff production



- Limits from avoiding excessive energy drain
- Search for solar axions (CAST)

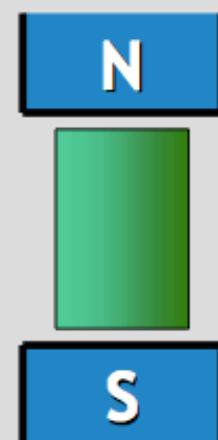
## Cosmology

In spite of small mass, axions are born non-relativistically (“non-thermal relics”)

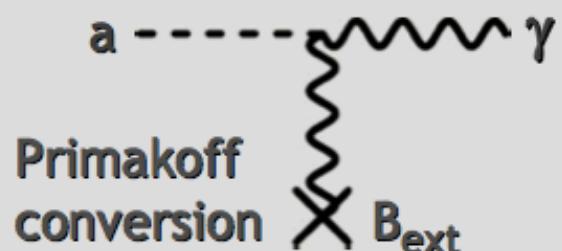
→ Cold dark matter candidate  
 $m_a \sim 1\text{-}1000 \mu\text{eV}$



## Search for Axion Dark Matter



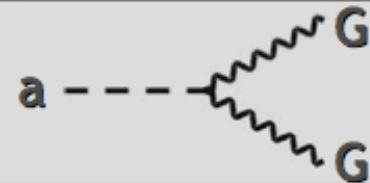
Microwave resonator  
(1 GHz = 4 μeV)



# Axion Properties

Gluon coupling  
(Generic property)

$$L_{aG} = \frac{\alpha_s}{8\pi f_a} G \bar{G} a$$



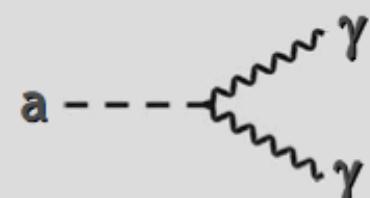
Mass

$$m_a = \frac{0.6 \text{ eV}}{f_a / 10^7 \text{ GeV}} \approx \frac{m_\pi f_\pi}{f_a}$$

Photon coupling

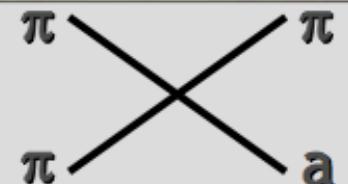
$$L_{a\gamma} = -\frac{g_{a\gamma}}{4} F \bar{F} a = g_{a\gamma} \vec{E} \cdot \vec{B} a$$

$$g_{a\gamma} = \frac{\alpha}{2\pi f_a} \left( \frac{E}{N} - 1.92 \right)$$



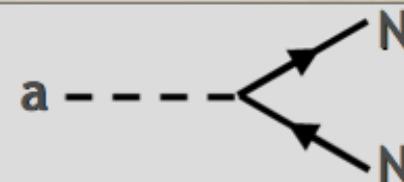
Pion coupling

$$L_{a\pi} = \frac{C_{a\pi}}{f_a f_\pi} (\pi^0 \pi^+ \partial_\mu \pi^- + \dots) \partial^\mu a$$



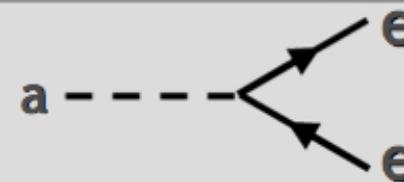
Nucleon coupling  
(axial vector)

$$L_{aN} = \frac{C_N}{2f_a} \bar{\Psi}_N \gamma^\mu \gamma_5 \Psi_N \partial_\mu a$$

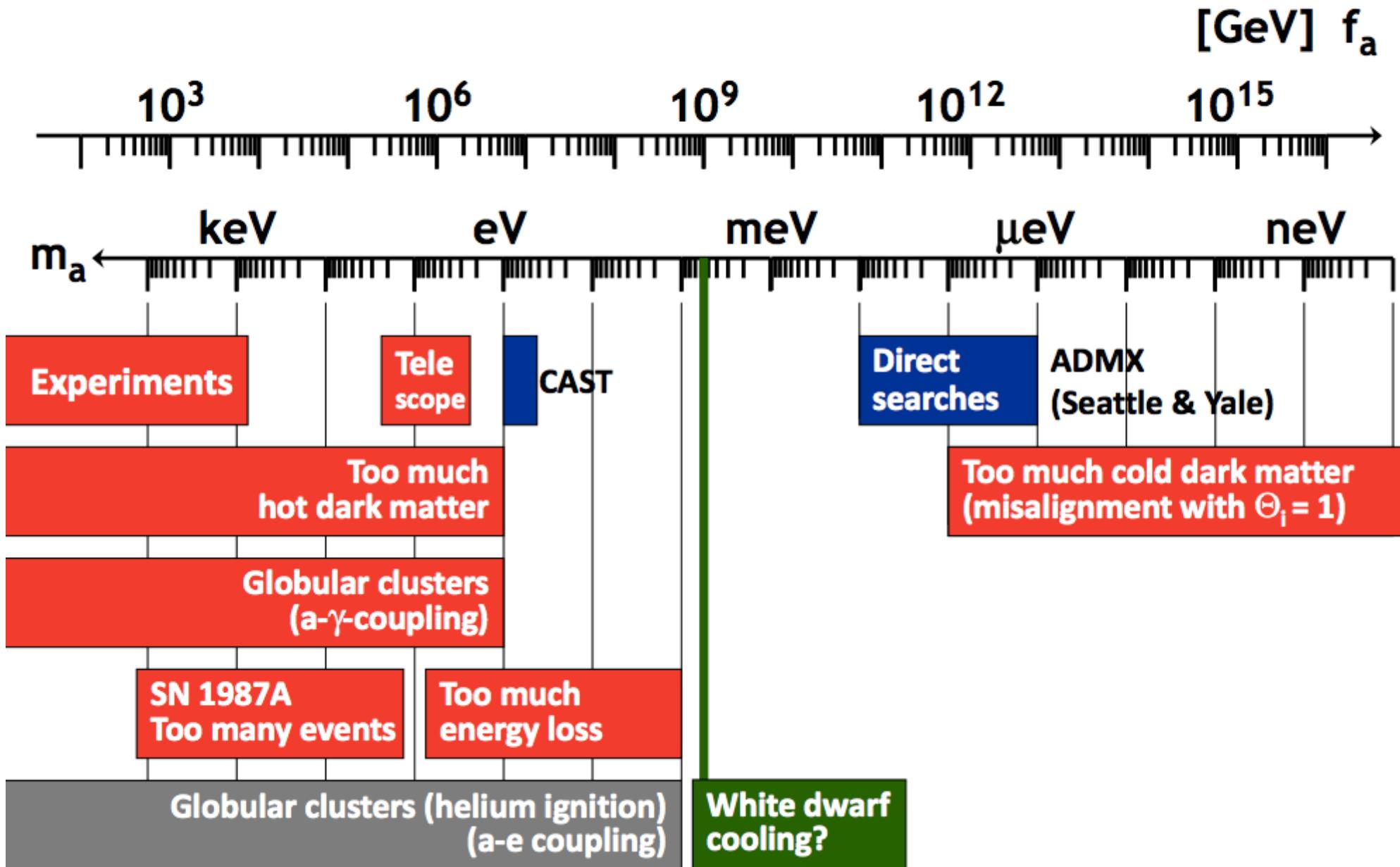


Electron coupling  
(optional)

$$L_{ae} = \frac{C_e}{2f_a} \bar{\Psi}_e \gamma^\mu \gamma_5 \Psi_e \partial_\mu a$$



# Axion Bounds and Searches

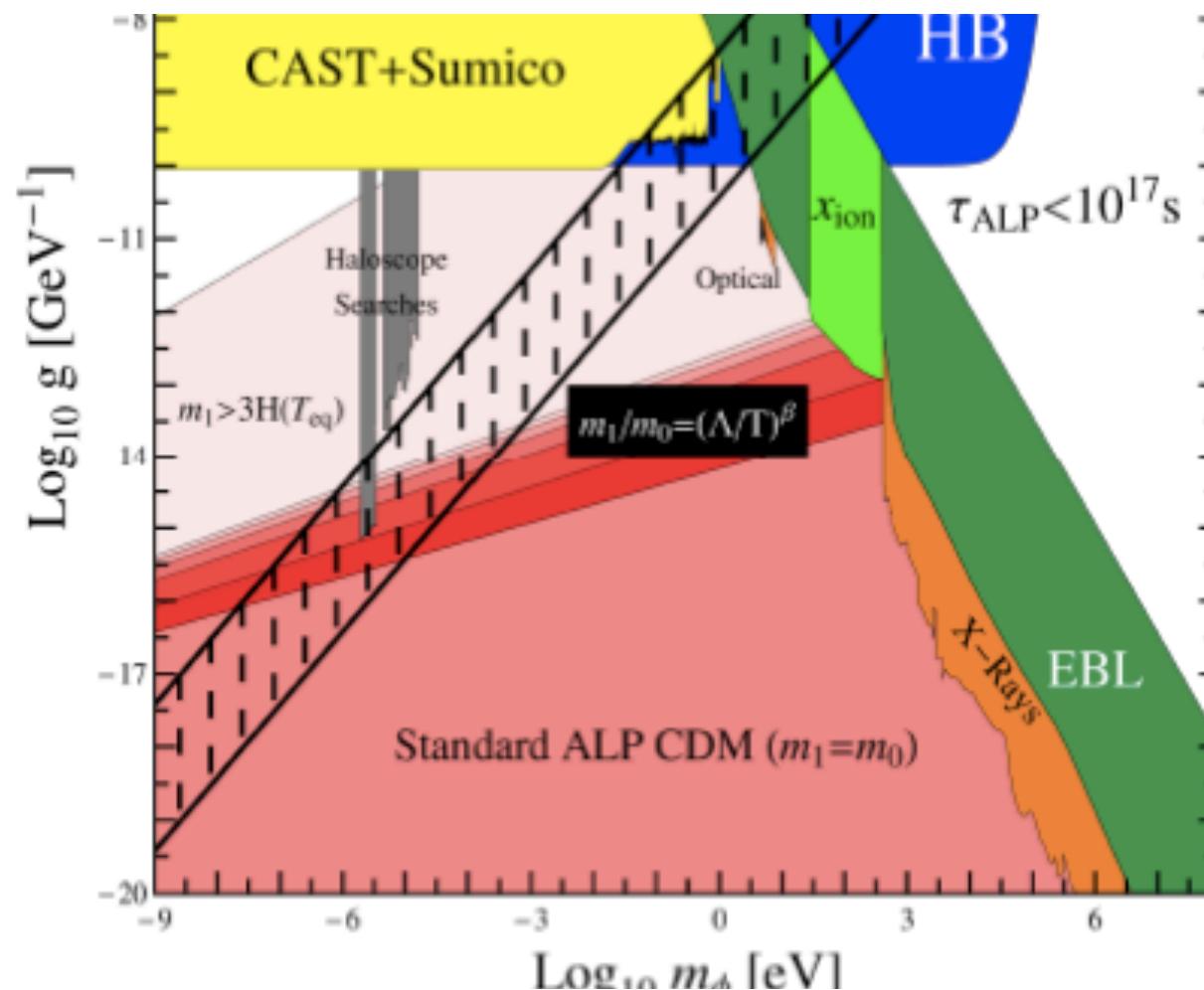


# Current experimental limits

$$L = -\frac{1}{4} g_W \varphi F_{\alpha\beta} \tilde{F}^{\alpha\beta}$$

$$g_W \approx \frac{\alpha}{m_e} g_p$$

Primakov Conversion



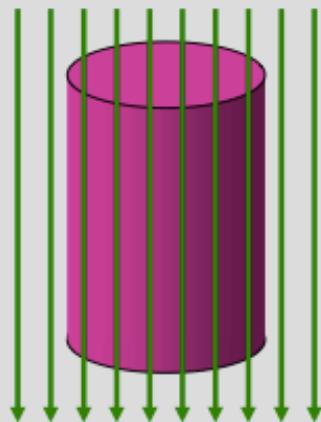
# Search for Galactic Axions (Cold Dark Matter)

DM axions  
Velocities in galaxy  
Energies therefore

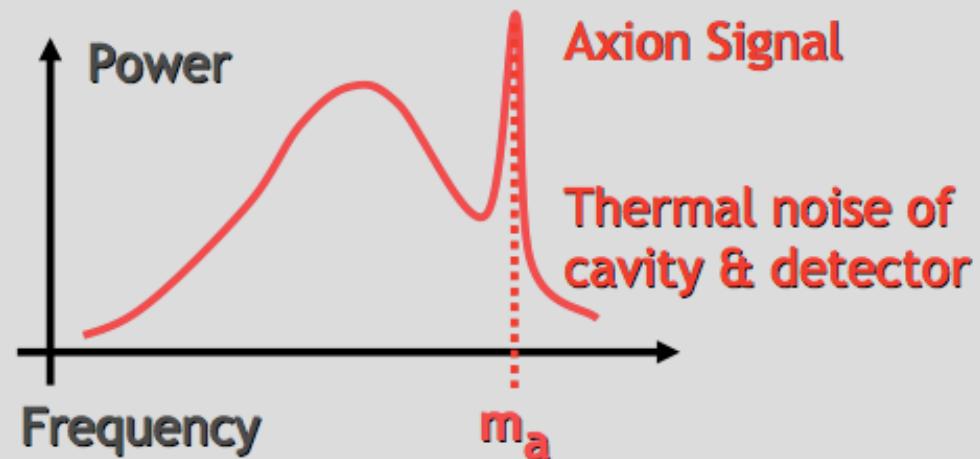
$$m_a = 10\text{-}3000 \mu\text{eV}$$
$$v_a \approx 10^{-3} c$$
$$E_a \approx (1 \pm 10^{-6}) m_a$$

Microwave Energies  
(1 GHz  $\approx 4 \mu\text{eV}$ )

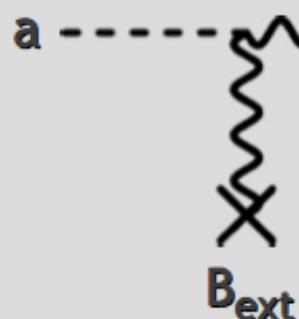
## Axion Haloscope (Sikivie 1983)



$B_{\text{ext}} \approx 8 \text{ Tesla}$   
Microwave Resonator  
 $Q \approx 10^5$



## Primakoff Conversion

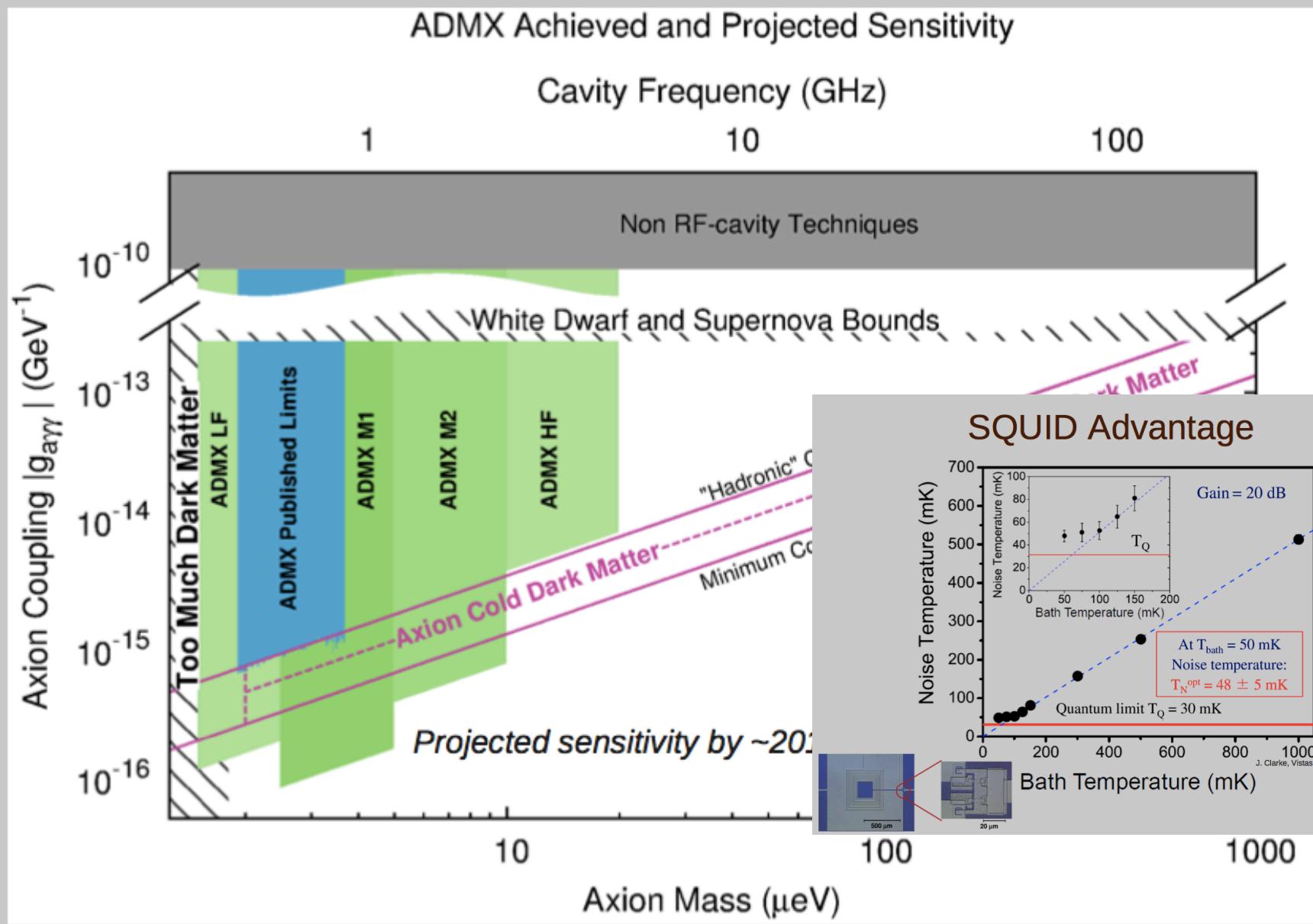


Cavity  
overcomes  
momentum  
mismatch

## Power of galactic axion signal

$$4 \times 10^{-21} W \frac{V}{0.22 \text{ m}^3} \left( \frac{B}{8.5 \text{ T}} \right)^2 \frac{Q}{10^5} \times \left( \frac{m_a}{2\pi \text{ GHz}} \right) \left( \frac{P_a}{5 \times 10^{-25} \text{ g/cm}^3} \right)$$

# ADMX Near Term Targets

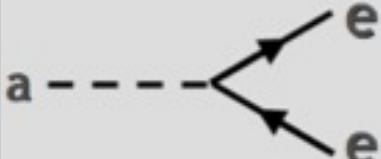


# Search for galactic axion – new proposal

Exploit the axion-electron coupling

$$L_{ae} = \frac{C_e}{2f_a} \bar{\Psi}_e \gamma^\mu \gamma_5 \Psi_e \partial_\mu a$$

$C_e \leq 10^{-13} \text{ GeV}^{-1}$



$$H_a = -\vec{S} \cdot \left[ \frac{g_p}{m_e} \nabla a \right]$$

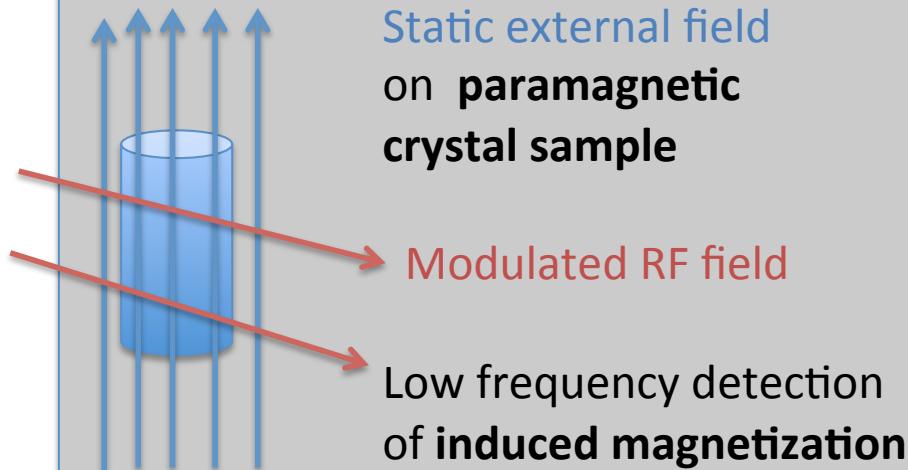
(only DFSZ axion)

Axion wind equivalent to magnetic field

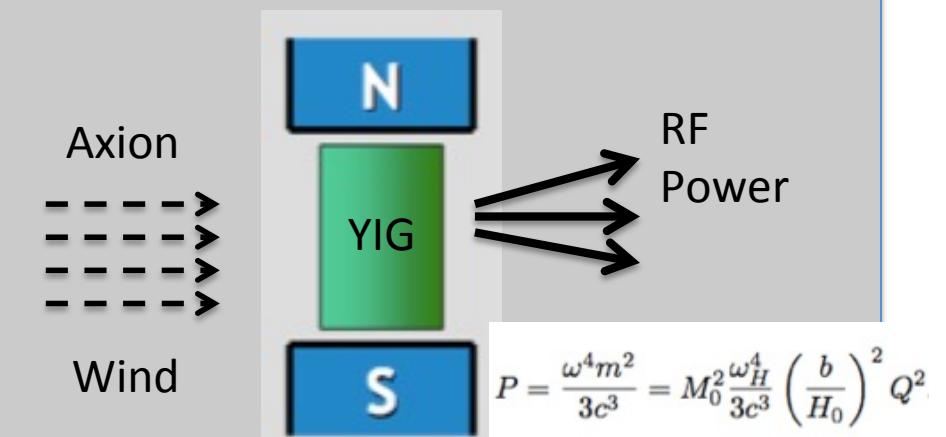
$$B_E = \frac{2g_p}{e} \frac{g_a}{g_J} \nabla_z a$$

$$B_{Ef} \approx \left( \frac{m_a}{10^{-4} \text{ eV}} \right) 9.4 \times 10^{-23} T$$

Detection using EPR magnetometry



- Excess Noise in Magnetized Crystal  
- Zeeman Transition in Optical Crystals



For both schemes **annual modulation expected**

# ELECTRON – AXION interaction through SPIN

F. Wilczek, J.M. Martín, J. Leon, R. Barbieri, I.V. Kolokolov, G. Raffelt

$$L = \bar{\psi}(x) (i\hbar \partial_x - mc) \psi(x) - a(x) \bar{\psi}(x) (g_s + ig_p \gamma_5) \psi(x) \quad (1)$$

The energy of interaction  
between spin and axion-field is

$$H_a = -\vec{S} \cdot \left[ \frac{g_p}{m_e} \nabla a \right]$$

This looks like a **magnetic coupling** through electron spin

$$\Delta E_a = \frac{g_p}{m_e} \hbar g_a \nabla_z a$$

$$B_E = \frac{2g_p}{e} \frac{g_a}{g_J} \nabla_z a \quad \text{Effective magnetic field for cosmological axion} \quad B_{Ef} \approx \left( \frac{m_a}{10^{-4} eV} \right) 9.4 \times 10^{-23} T$$

## DETECTION TECHNIQUES :

Electron Spin Resonance Magnetometry with Paramagnetic Materials

Optical Spectroscopy in Paramagnetic Crystals

Effective B RF field @  $10^{-22}$  T seems feasible to attain

# ESR Magnetometry/Zeeman Spectroscopy:

from  $10\mu\text{eV}$  to  $\text{meV}$  ( GHz to THz)

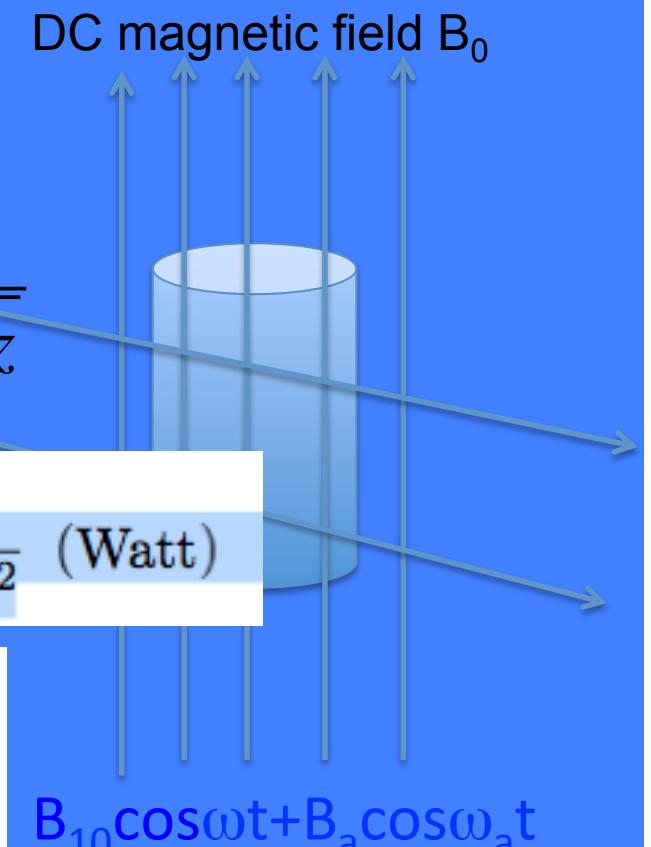
SAMPLE: DPPH , TbF<sub>3</sub> , YIG, Ruby

$$m = M_0 \frac{b}{H_0} Q,$$

$$P_{\min RF} = 10^{-22} \frac{W}{\sqrt{\text{Hz}}}$$

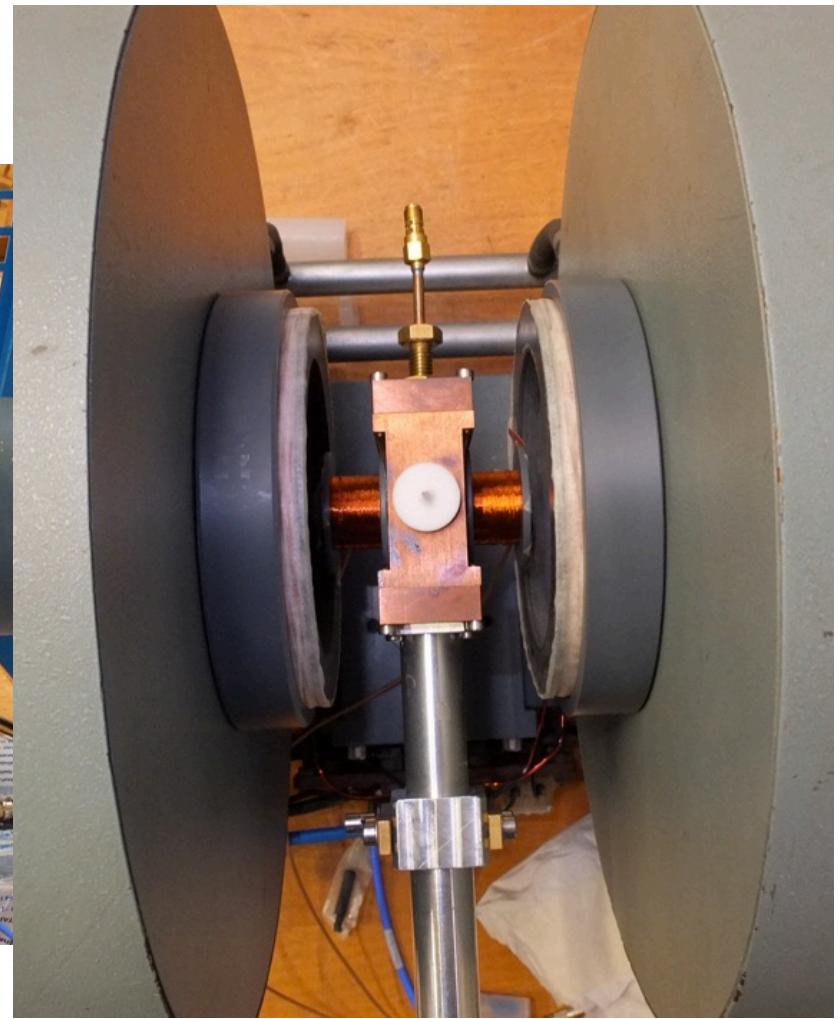
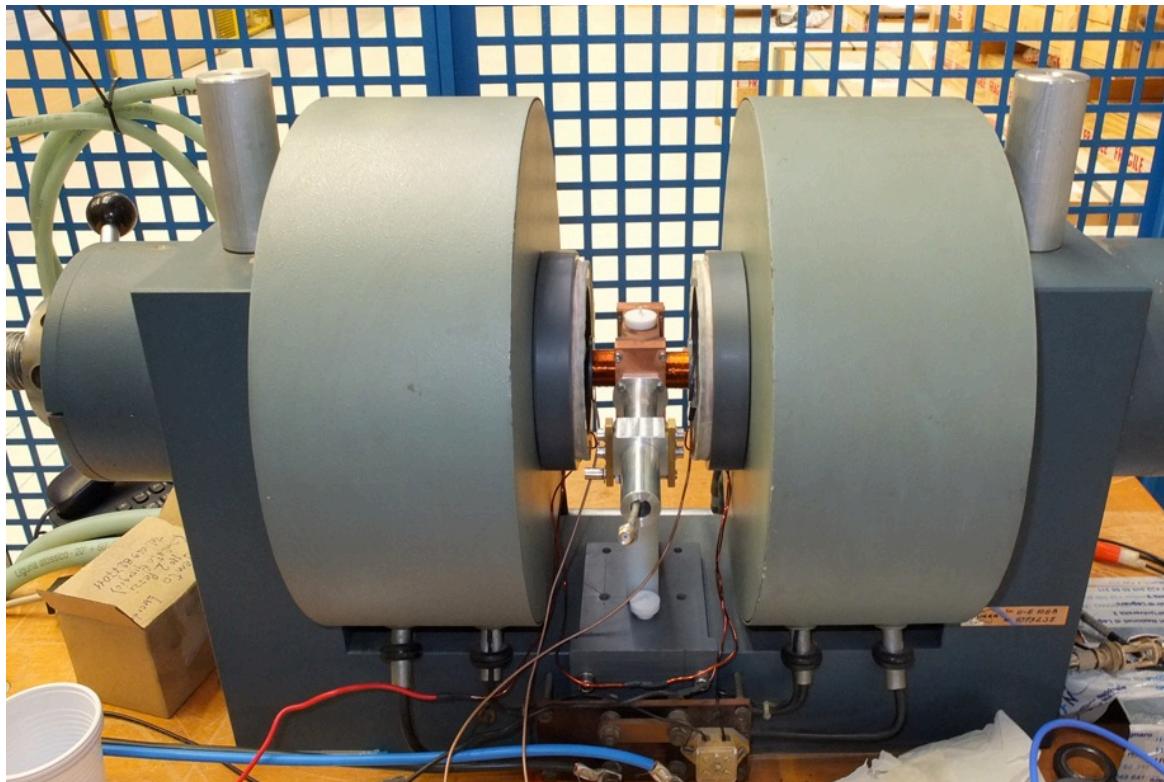
$$P_a(\omega) = \frac{\omega^2}{12\pi c^3} (\dot{m})^2 = \frac{1}{4\pi} \frac{\omega^4}{3c^3} \frac{B_a^2 V^2}{\mu_0} \frac{\gamma^2 M_0^2 \tau^2}{1 + \tau^2(\omega - \omega_0)^2} \text{ (Watt)}$$

$$\text{SNR}^2 = \frac{M_s^2(\nu)}{\mathcal{M}_n^2(\nu)} \frac{1}{\Delta\nu} = \frac{\gamma M_0 \tau B_a^2 V}{2\mu_0 \hbar \coth\left(\frac{\hbar\omega}{2k_B T}\right)} \frac{1}{\Delta\nu}$$



$$P_{\text{noise}} = \hbar\omega \sqrt{\frac{\Delta\nu}{t_{\text{int}}}} = 1.58 \cdot 10^{-23} \left( \frac{m_a}{10^{-4} \text{eV}} \right) \sqrt{\frac{\Delta\nu}{t_{\text{int}}}} \text{ W}$$

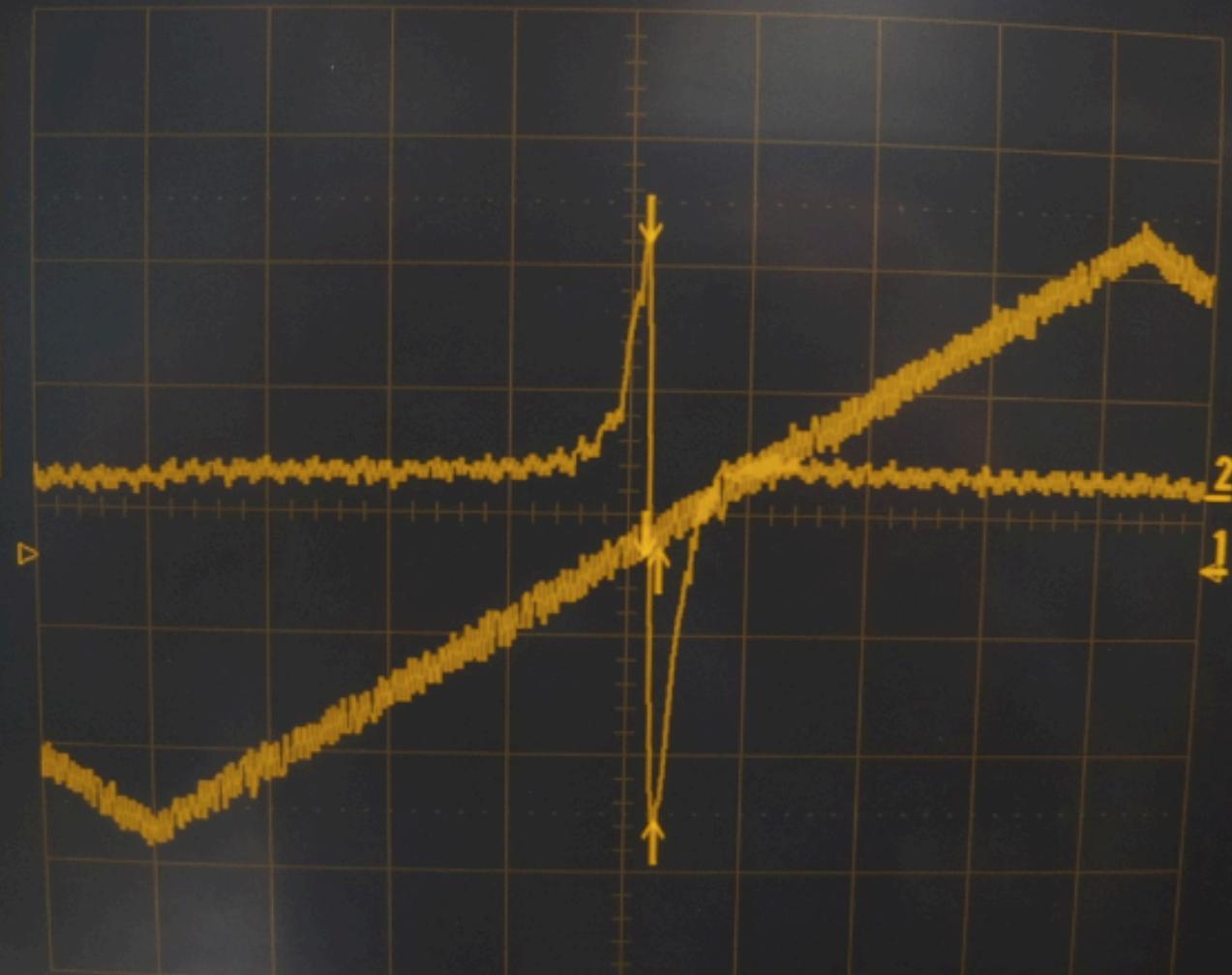
# LNL EPR SET UP



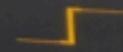
24-Jun-14  
13:27:26

1  
.1 s  
20.0 mV  
1.2 mV

2  
.1 s  
10.0 mV  
-47.50 mV



.1 s  
1 20 mV DC  
2 10 mV AC



1 HFREJ 0.0 mV

$\Delta t$

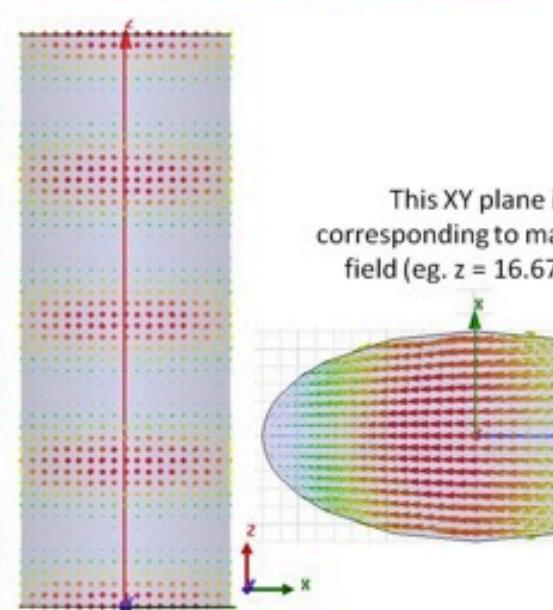
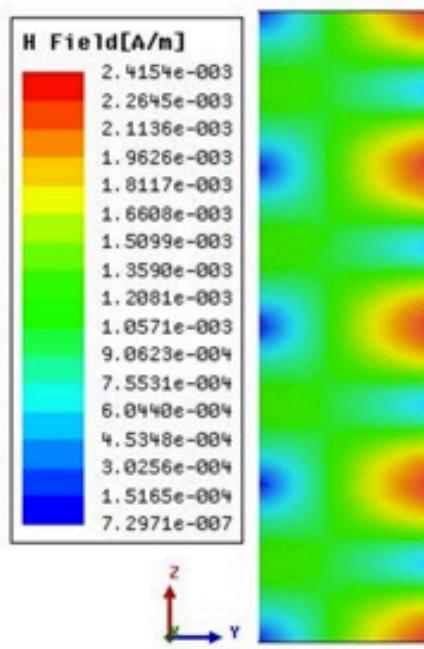
12.0 ms

$\frac{1}{\Delta t}$

83.3 Hz

□ STOPPED

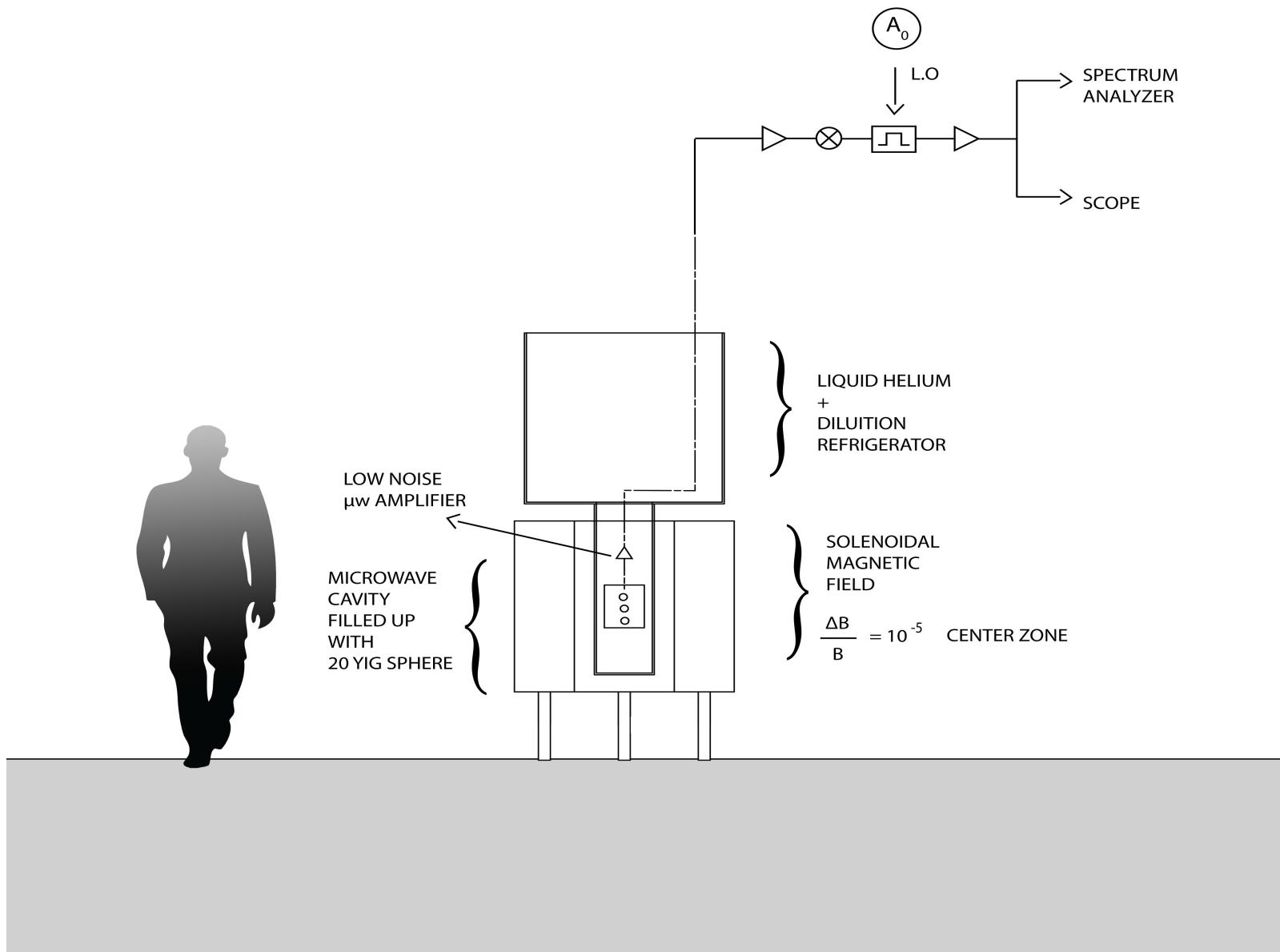
## Magnetic field distribution



This XY plane is  
corresponding to ma  
field (eg.  $z = 16.67$ )

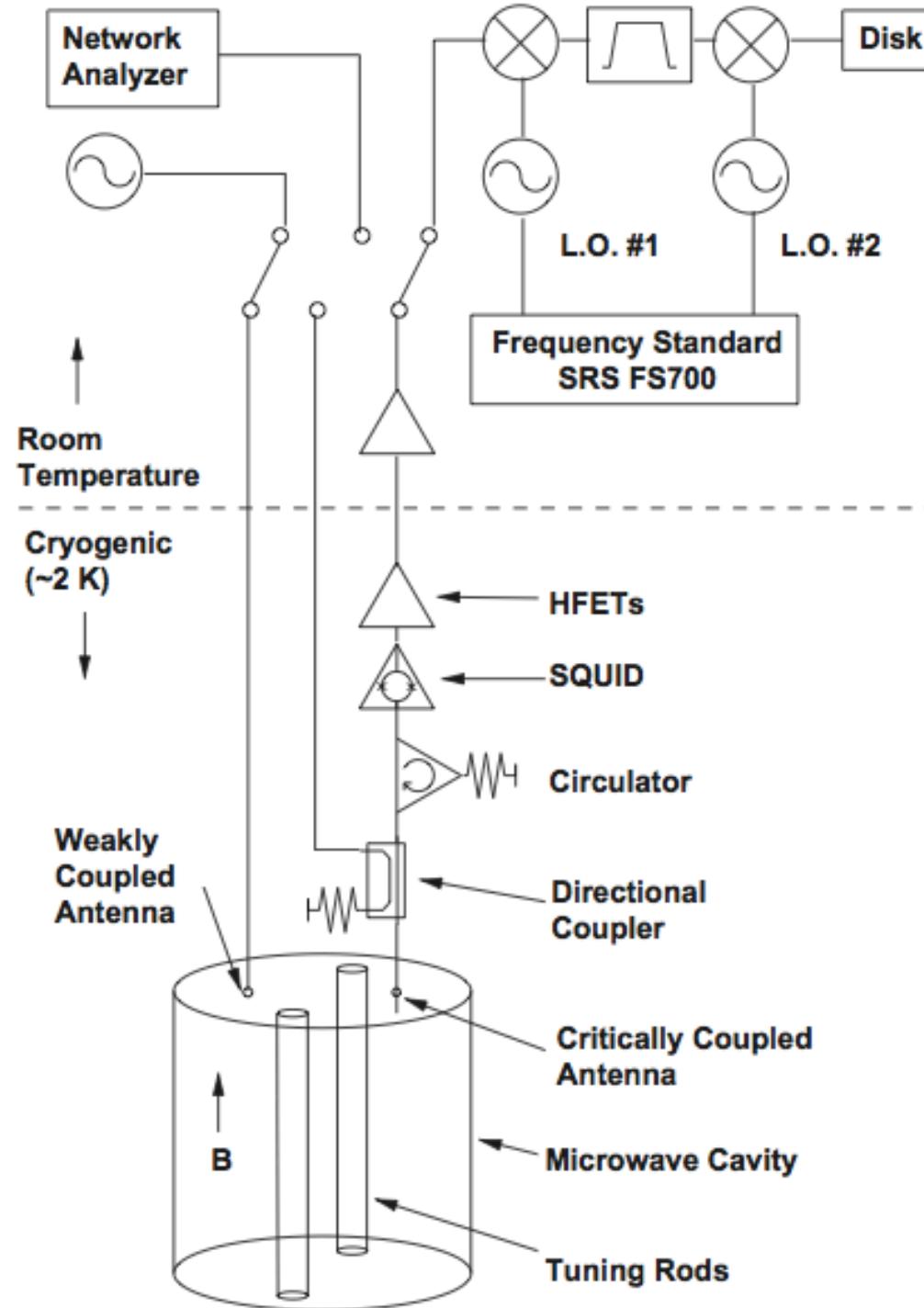


$Q_{UAERERE}$  -  $A_{XION}$  EXPERIMENTAL SET - UP



## **QUAX Activities ( 2015 – 2017/8 ) :**

- 1) High Susceptivity Material , Low Dissipation (LNL-INRIM.TO )
- 2) Low Noise Receiver : Linear / Quantum @ near QL( PD-LNL)
- 3) High Frequency High Q Microwave Cavity  $10^5$ : (PD-LNL)
- 4) High Magnetic Field : 05-1 Tesla ( GE-LNL )
- 5) Cryogenics @ 100 milliKelvin ( TN-PD )
- 6) Optical Spectroscopy on Paramagnetic Single Crystals  
Zeeman



“Hybrid” superconducting cavities may increase Q,  
and thus increase signal power



Open resonators may be the key to explore axion masses up to meV

