

# Neutrinos as a probe of New Physics effects

Daide Meloni

Dipartimento di Matematica e Fisica  
RomaTre

FUTURO  
IN RICERCA



# Searching for New Physics

Neutrino oscillation physics is entering a precision era:

- Good knowledge of mixing angles and mass differences
- First hints of a non-vanishing  $CP$  phase
- Sum of the neutrino masses bounded from above



It is time to devote serious efforts to the search of New Physics in the lepton sector, using neutrinos

# Two (three) different items

Main message of this talk:

Neutrinos can be used as probes for some New Physics scenarios



New research field in theoretical particle physics

Not discussing new results in details, only two examples to show how New Physics shows up in the  $\nu$  sector

- Sterile neutrinos
- Large Extra Dimensions
- (Non Standard Interactions)

# Standard neutrino oscillations

- Neutrinos can also be described in terms of mass eigenstates  $\nu_i$



neutrino matrix  
matrix

$$|\nu_\alpha\rangle = \sum_{i=1}^3 U_{\alpha i} |\nu_i\rangle$$

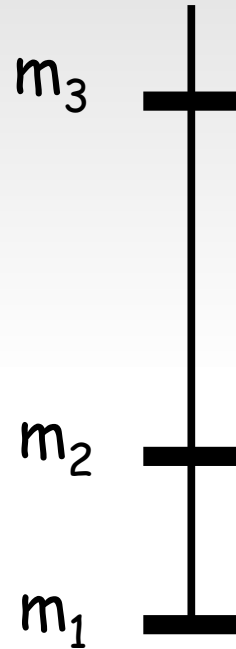
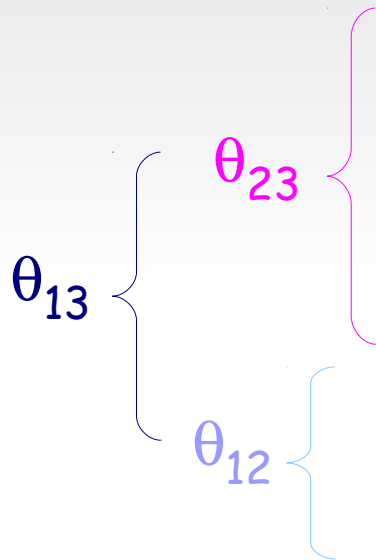
- Considering time evolutions of mass eigenstates:

$$P(\nu_\alpha \rightarrow \nu_\beta) = |\langle \nu_\beta | \nu_\alpha(t) \rangle|^2 = \left| \sum_j U_{\beta j} U_{\alpha j}^* e^{\frac{-im_j^2 L}{2E_\nu}} \right|^2$$

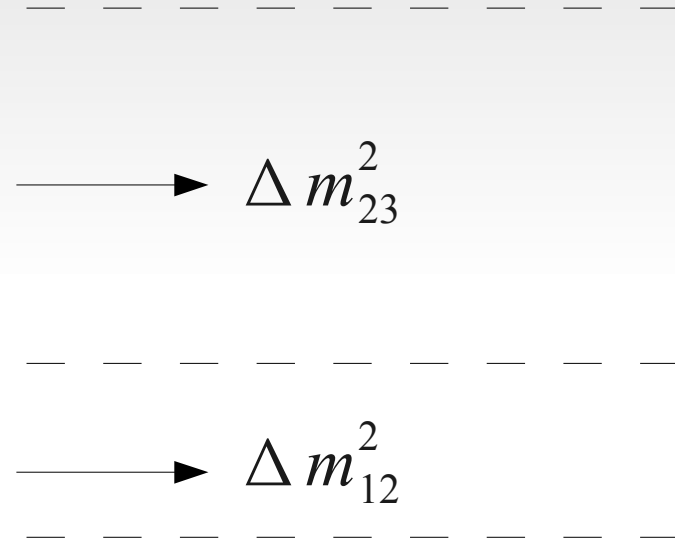
$$U_{PMNS} = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{atmospheric mixing}} \times \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13} e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13} e^{-i\delta} & 0 & c_{13} \end{pmatrix}}_{\text{reactor mixing}} \times \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{solar mixing}}$$

# To be determined by oscillation experiments

Mixing angles

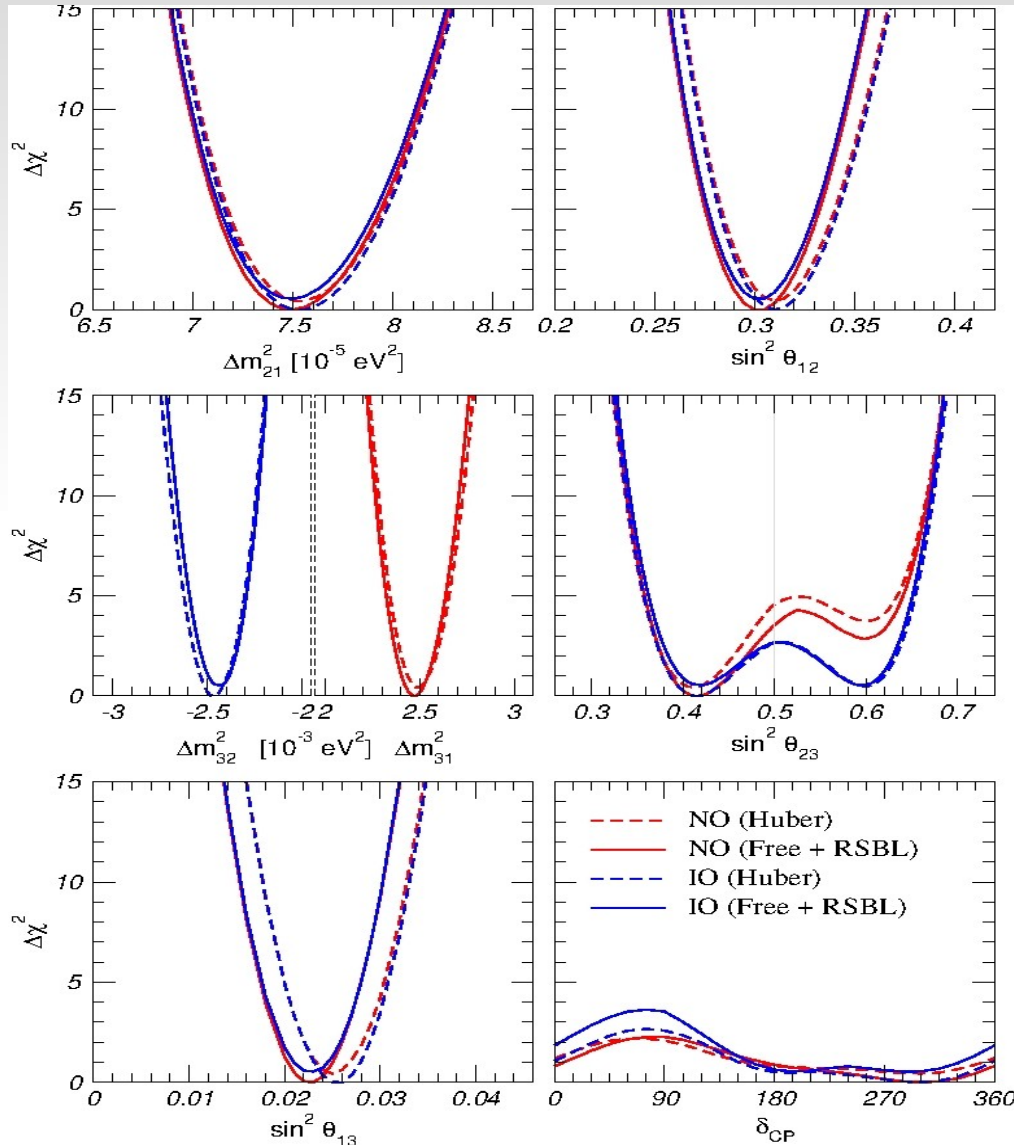


Mass differences



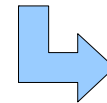
- And: a possible CP phase  $\delta$  and the absolute order of the mass eigenstates (normal or inverted hierarchy)

# Summary of the experimental results



Gonzalez-Garcia et al. JHEP1212,(2012)123

Parameter	Fit results
$\theta_{12}$	$33.36^{+0.81}_{-0.78}$
$\theta_{13}$	$8.66^{+0.44}_{-0.46}$
$\theta_{23}$	$40.0^{+2-1}_{-1.5}$
$\delta$ (?)	$300^{+66}_{-138}$
$\Delta m_{23}^2$ ( $10^{-3} \text{ eV}^2$ )	$2.47^{+0.07}_{-0.07}$
$\Delta m_{12}^2$ ( $10^{-5} \text{ eV}^2$ )	$7.50^{+0.18}_{-0.19}$



- precision era in the determination of mixing parameters

# Perturbative approach

- $\sin^2 2\theta_{13}$  is a small parameter



- New effects in oscillations must be even smaller

$$A(\nu_\alpha \rightarrow \nu_\beta) = \underbrace{A^{SM}(\nu_\alpha \rightarrow \nu_\beta)}_{\text{interference}} + \delta A_{\alpha\beta}$$

From the interference  
term one can:  
( $\delta^2$  is generally too small)

- set strong bounds on  $\delta$  if the data are precise and very well described by SM physics
- "measure"  $\delta$  if the data are precise and NOT well described by SM physics

$$\text{Ex: } \sin^2 2\theta_{13}^{\text{eff}} = \sin^2 2\theta_{13}^{\text{SM}} (1+\delta)$$

# The case of sterile neutrinos

- The number of active neutrino species is fixed by the Z-boson invisible decay width:  $N_\nu = 2.994 \pm 0.011$
- Extra families (if they exist) must have either very heavy neutrinos ( $m_N > m_Z/2$ ), or no neutrinos at all

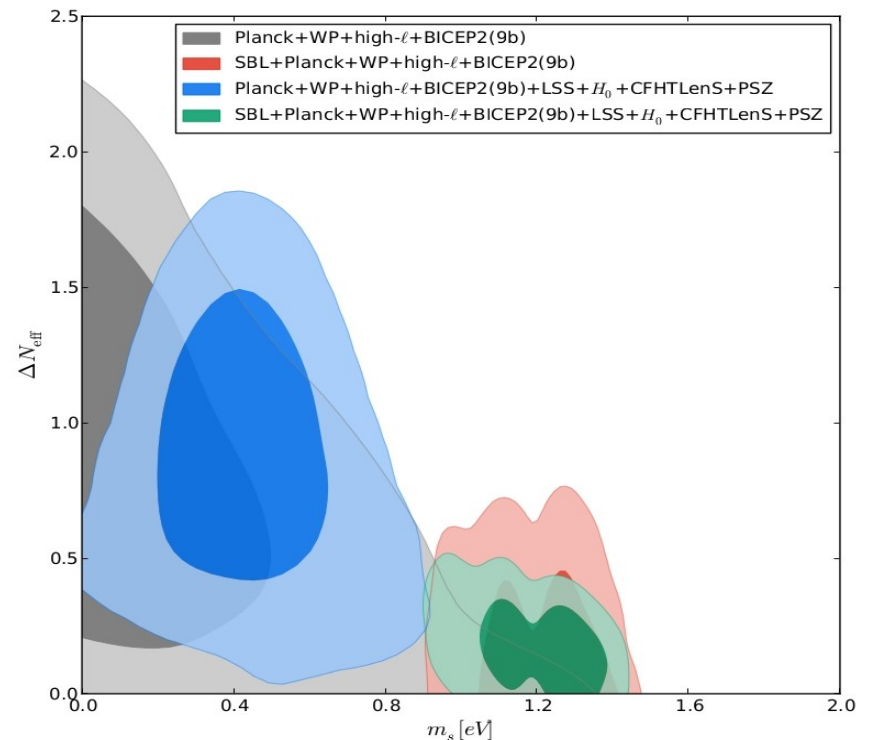


Standard Model singlets allowed:  
sterile neutrinos

Renewed interest after  
the recent BICEP2 results

extra species

M.Archidiacono et al., arXiv:1404.1794 [astro-ph.CO]



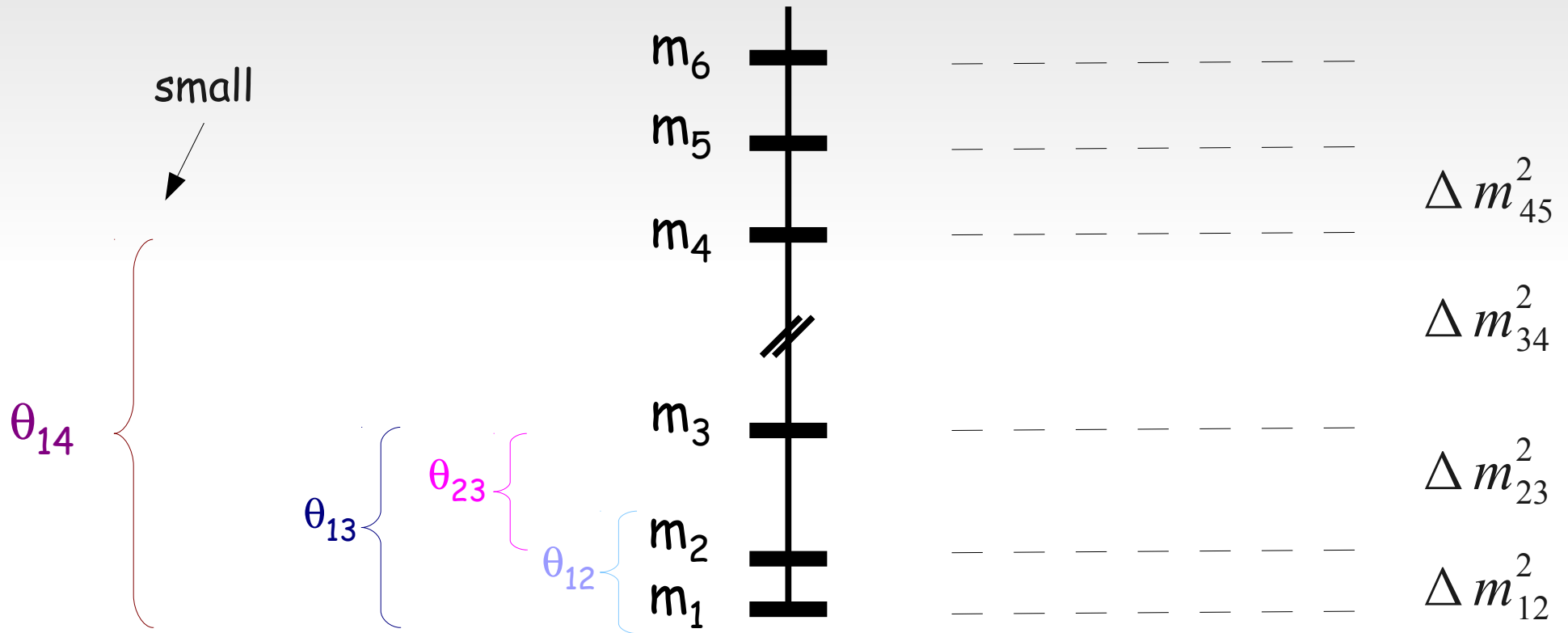
Mass of the extra species



# The case of sterile neutrinos

## Mixing angles

## Mass differences



# The case of sterile neutrinos

$$U = R_{34} R_{24} R_{14} R_{23} R_{13} R_{12}$$

perturbations

Example:  $\nu_e \rightarrow \nu_e$  transition

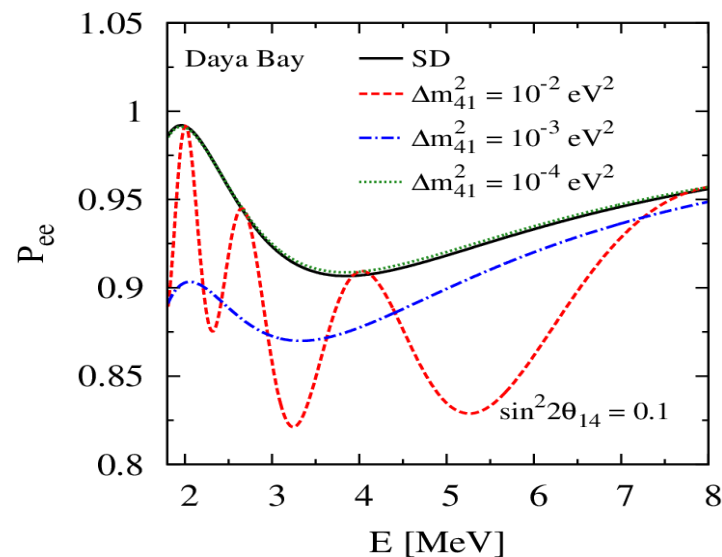
$$P_{ee} \sim 1 - s_{12} \sin^2 2\theta_{13} \sin^2\left(\frac{\Delta m_{32}^2 L}{4 E_\nu}\right) - c_{12}^2 \sin^2 2\theta_{13} \sin^2\left(\frac{\Delta m_{31}^2 L}{4 E_\nu}\right) \Rightarrow \text{Standard}$$

$$- s_{12} \sin^2 2\theta_{14} \sin^2\left(\frac{\Delta m_{42}^2 L}{4 E_\nu}\right) - c_{12}^2 \sin^2 2\theta_{14} \sin^2\left(\frac{\Delta m_{41}^2 L}{4 E_\nu}\right) \Rightarrow \text{Interference}$$

Current upper limit:

$$\sin^2 2\theta_{14} \sim 0.1$$

M.C.Gonzalez-Garcia et al.,  
JHEP 1212, 123 (2012)



# The case of sterile neutrinos

Bounds on new parameters require experiments with very low systematic uncertainties

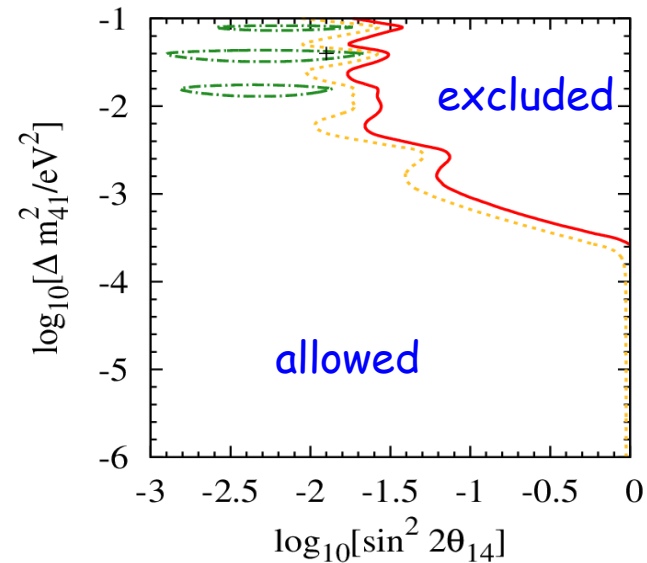
Daya Bay in China:

$\bar{\nu}_e \rightarrow \bar{\nu}_e$  transition

neutrinos from reactor plants



1,2 and 3 $\sigma$  confidence level



$$\sin^2 2\theta_{14} \lesssim 10^{-2}$$

# Large Extra Dimensions

Hierarchy problem:

there exist two fundamental energy scales:

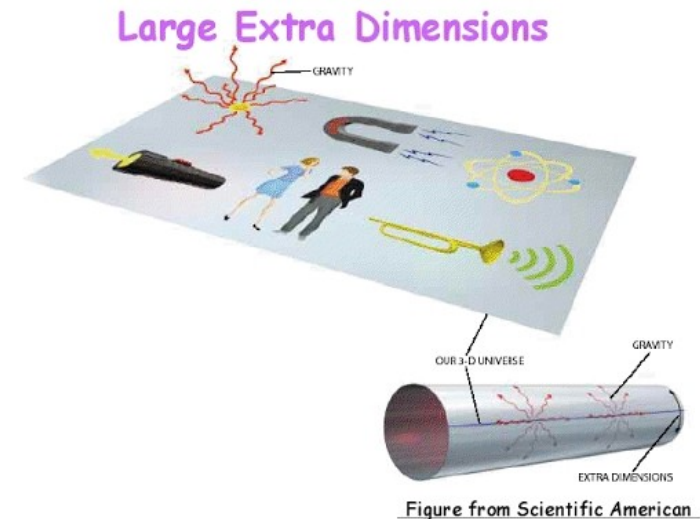
- the electroweak scale  $M_D \sim 1 \text{ TeV}$
- quantum gravity  $M_{PL} \sim 10^{18} \text{ GeV}$



A possible way-out:

- there are  $\delta$  compact extra dimensions of radius  $R$  and  $M_D$  is the only fundamental energy scale:  $M_{PL}^2 = R^\delta M_D^{\delta+2}$
- at distances less than  $R$  gravity propagates in all  $4+\delta$  dimensions
- Standard Model fields are confined in our 4D world

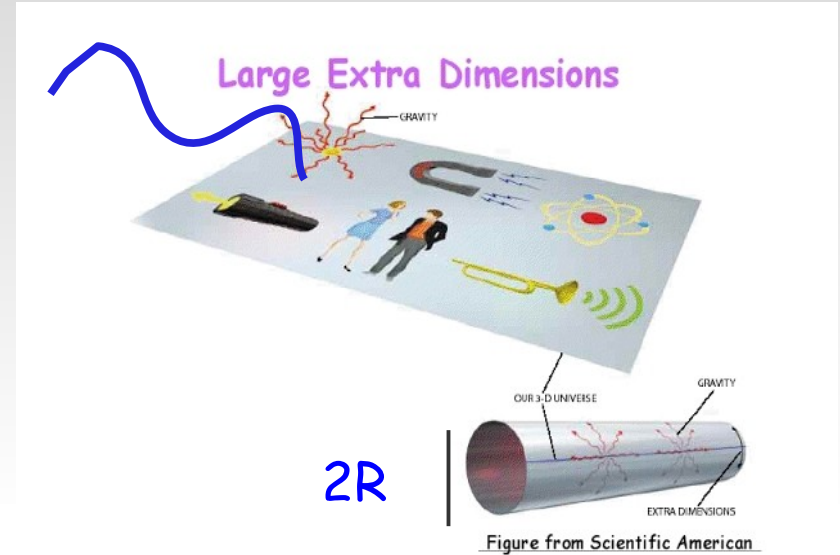
05/06/2014



# Right-handed neutrinos in extra dimensions

Since neutrinos are special: massless right-handed neutrinos  $\nu_R$  also feel the whole 5D space

- $\nu_R$  propagated in the whole 5D
- $\nu_R = \nu_R(x^\mu, y)$



Imposing  $y = y + 2\pi R$  on the  $\nu_R$  wave function generates infinite replica of the field

$$\nu_R(x, y) \sim \sum_{n=-\infty}^{+\infty} \nu_R^{(n)}(x, y) e^{\frac{iny}{R}}$$

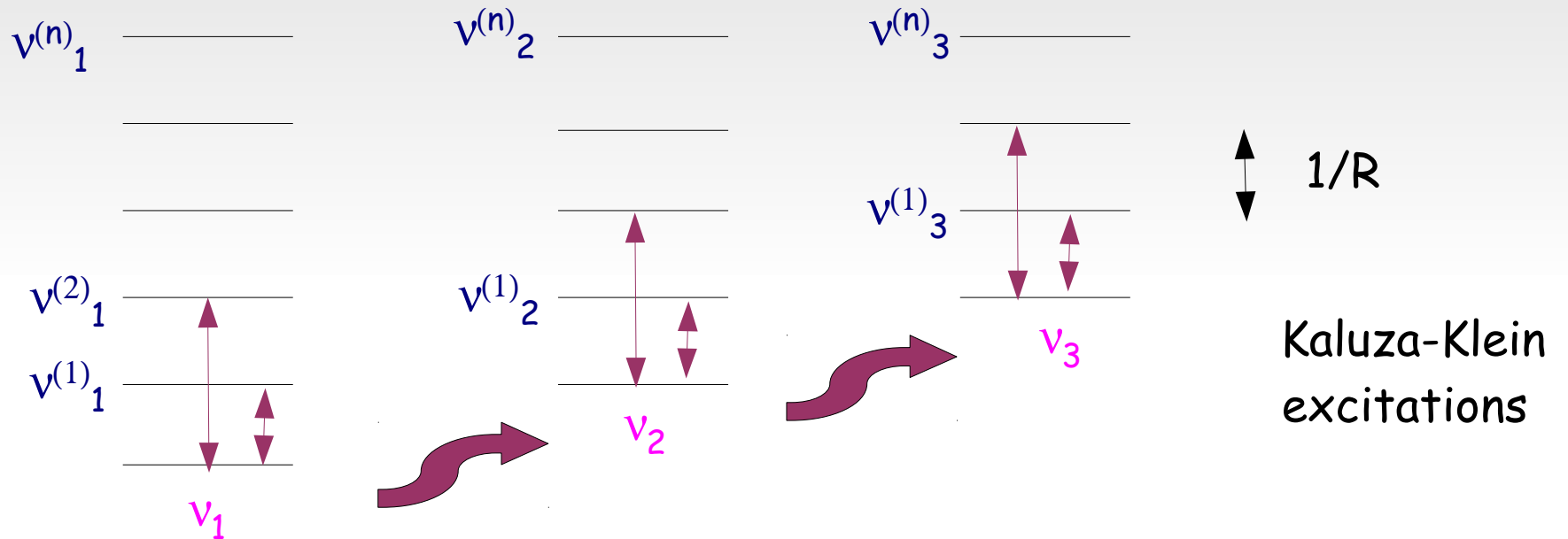
Interaction term:

$$\lambda_{\alpha\beta} \nu_L^\alpha(x) \nu_R^\beta(x, 0) H(x)$$



# Right-handed neutrinos in extra dimensions

For the mass eigenstates (normal ordering):



Oscillations can take place among active-active and active-KK (infinite) states

# Right-handed neutrinos in extra dimensions

$$A(\nu_\alpha \rightarrow \nu_\beta) = \sum_j U_{\beta j} U_{\alpha j}^* e^{\frac{-im_j^2 L}{2E_\nu}} \rightarrow \sum_{j=1}^3 \sum_{k=0}^{\infty} U_{\beta j} U_{\alpha j}^* |W_j^{(k)}|^2 e^{\frac{-i\lambda_j^{(k)2} L}{2E_\nu R^2}}$$

D. Davoudiasl, P. Langacker and M. Perelstein,  
Phys. Rev. D 65, 105015 (2002)

$W_j^{(k)}$  = Transition between zero modes and KK

$\lambda_j^{(k)}$  = absolute neutrino masses

Depend on the lightest  
absolute neutrino mass  
 $m_0$  and  $R$

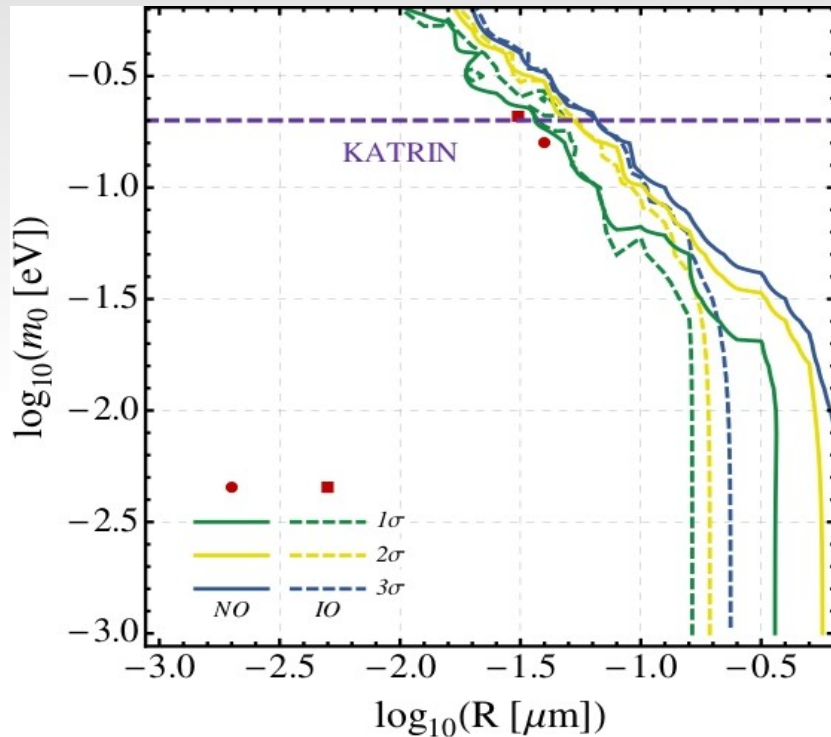
- Only upper limit on  $m_0$  (from  $\sum_i m_i < 0.2$  eV)
- Limits on  $R$  from experiments based on the torsion pendulum:  
 $R < 37 \times 10^{-6}$  m (95% CL) or equivalently  $M_D > 3.6$  TeV for  $\delta=2$  (LHC gives comparable bounds)

J. Beringer et al. (Particle Data Group  
Collaboration), Phys. Rev. D 86, 010001 (2012)

G.Aad et al. [ATLAS Collaboration],  
Phys. Rev. Lett. 110, 011802 (2013)

# First results...

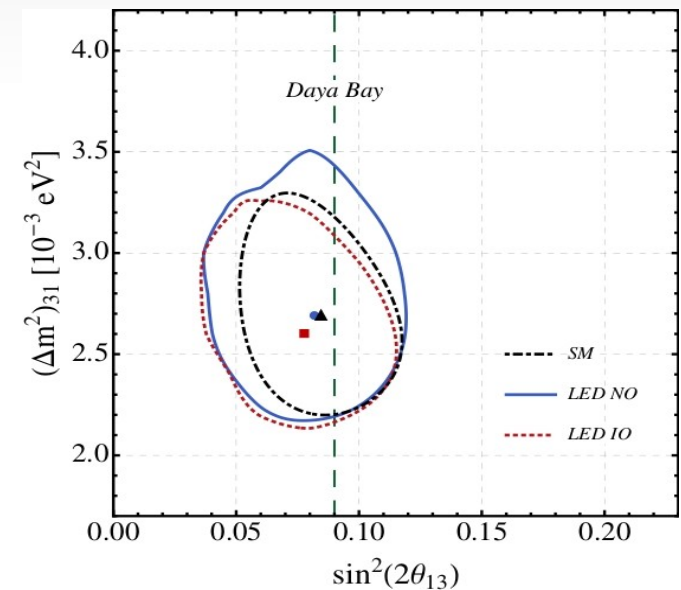
new limits from Daya Bay



$R < 0.2$  ( $0.6$ )  $\times 10^{-6}$  m @  $2\sigma$   
for Inverted (Normal) Ordering



Effects of New Physics



I. Girardi and D. Meloni, 1403.5507



# Conclusions

- Neutrino physics is an active field, from both experimental and theoretical point of views
- Many and precise data are now available, which are very well described in the context of the SM theory of neutrino oscillation
- As for the hadronic sector, New Physics must pop-up as perturbations of the standard picture
- We started to investigate such tiny effects for a variety of New Physics scenarios...

[Join the program !](#)