

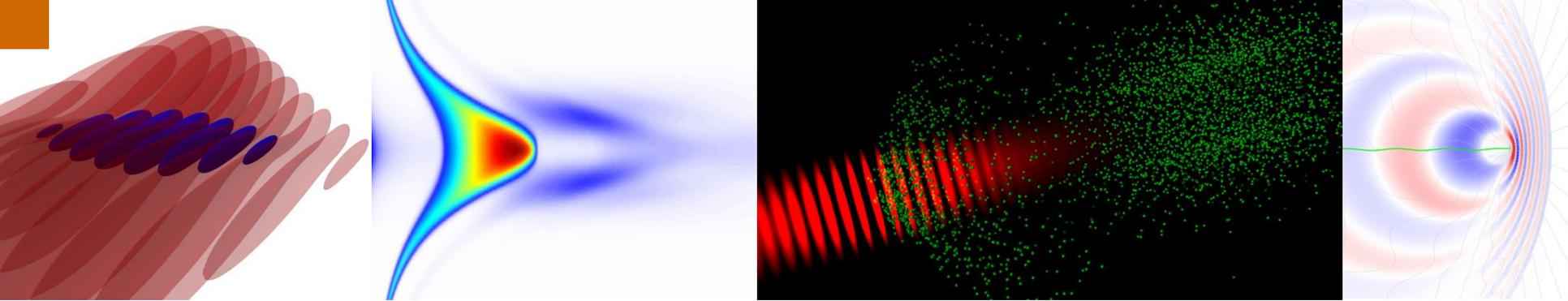
Ultra-compact all-optical FELs and Compton sources

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EAAC, La Biodola, Isola d'Elba
September 13-19th, 2015





Optical free-electron lasers (OFELs)

- EUV to X-rays : **100eV - 10 keV**
- **Coherent** light source
- $10^{10} - 10^{12}$ photons per pulse at sub-% bandwidth
- Vision of a **compact (X-ray) FEL, no experimental realization yet.**

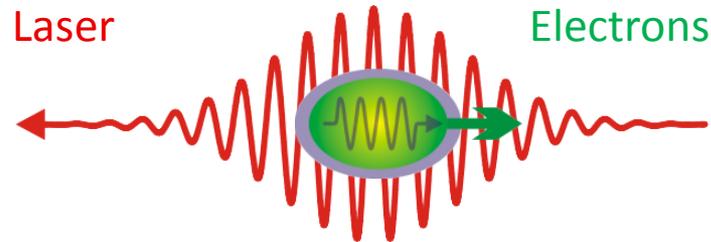
Compton sources

- Potentially most brilliant source at hard X-rays : **keV - MeV**
- **Incoherent** source
- Narrow bandwidth @ %-level, or broad bandwidth, tunable in wavelength
- **Experimental realizations** (linacs and LWFA) are available, but **more photons per shot desired.**



Optical Free Electron Lasers

Optical free-electron laser (OFELs) – shorter FEL wavelengths while using less electron energy on table-top?

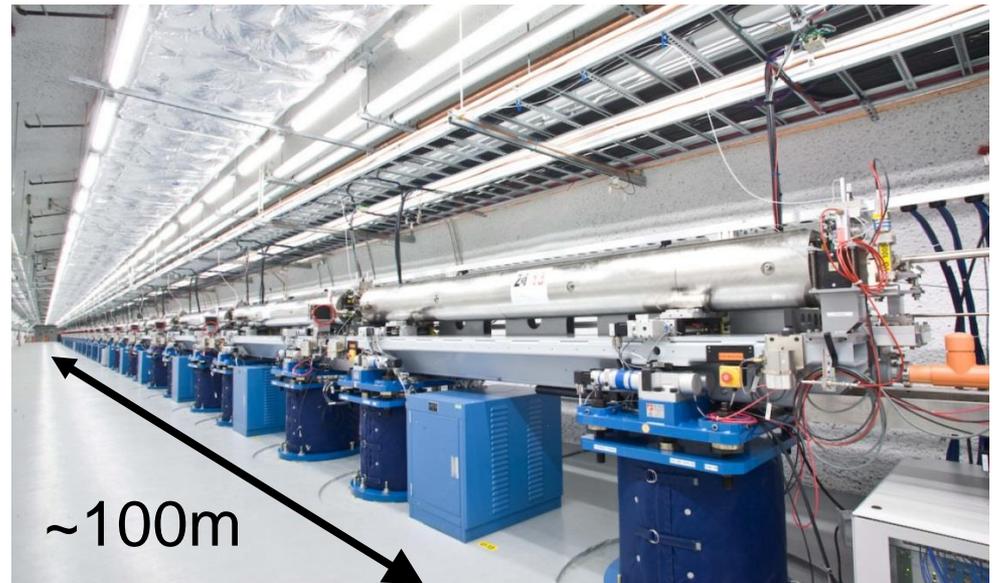


Even if GeV-scale LWFA accelerators are table-top, X-ray FEL undulators usually are not !

- **Optical undulators** provide μm -scale undulator lengths at high fields.

$$\lambda_{FEL} = \frac{\lambda_{laser}(1 + a^2/2)}{2\gamma^2}$$

- **Optical free-electron lasers** additionally drive an electron beam instability and thus have to meet very different requirements.

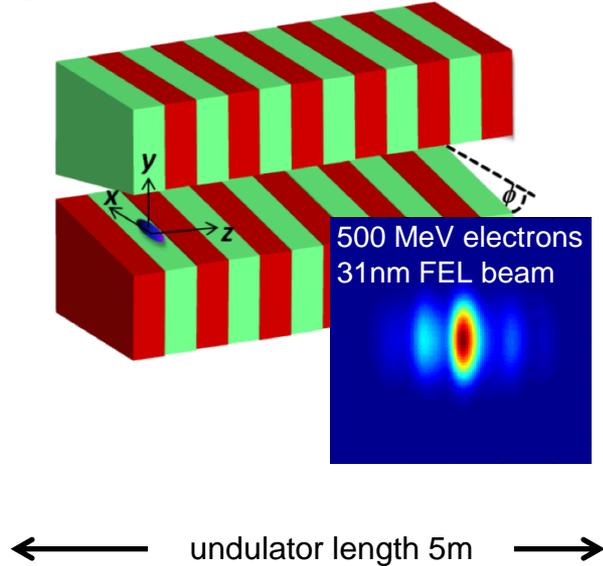


The 0.1nm FEL at SLAC (<http://lcls.slac.stanford.edu>)

Laser-plasma driven free-electron lasers using magnetic undulators can be compact at EUV to soft x-rays.

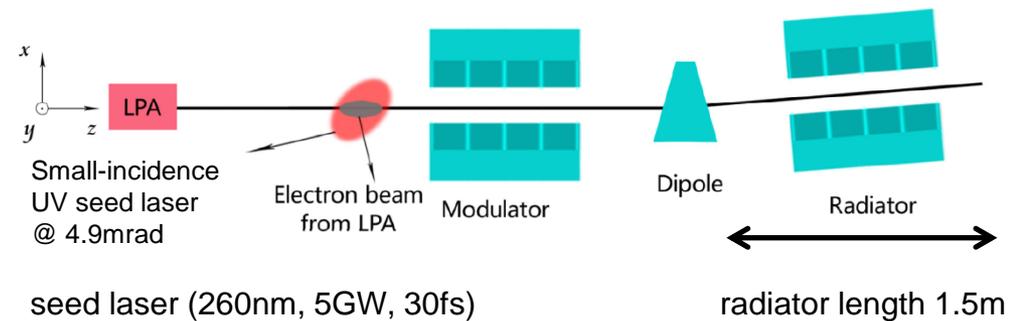
SASE FEL using transverse-gradient undulators

Huang *et. al.*, Phys. Rev. Lett. **109**, 204801 (2012)

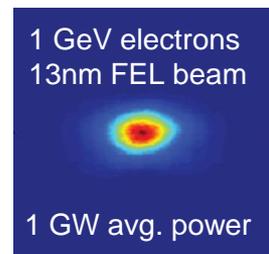


Coherent harmonic generation

Chao Feng *et. al.*, Opt. Express **23**, 14993-15002 (June 2015)



- Introduce **angular modulation** onto LPA beam and transform to **density modulation**
- Not sensitive to energy spread
- Good transverse coherence

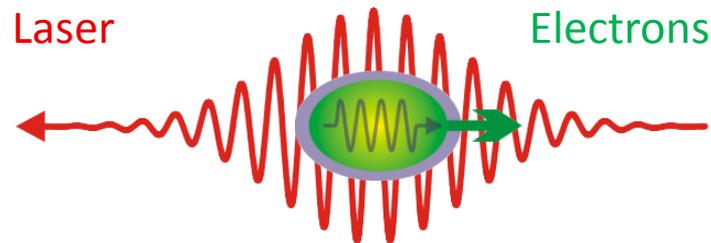


When scaling to sub-nm wavelengths magnetic undulators cease to be compact!

$$L_{\text{gain}} = \frac{\lambda_u}{4\pi\sqrt{3}\rho}$$

$$\rho \propto \gamma^{-1} \lambda_u^{2/3}$$

Optical free-electron laser (OFELs) – shorter FEL wavelengths while using less electron energy on table-top?



- Optical lasers provide μm -scale undulator lengths at high fields.
- However, yields of existing Compton sources

$$10^7 - 10^8 \frac{\text{photons}}{\text{shot}}$$

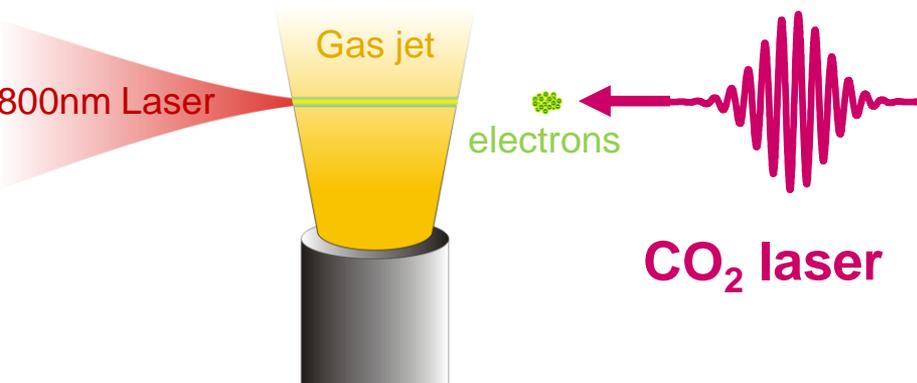
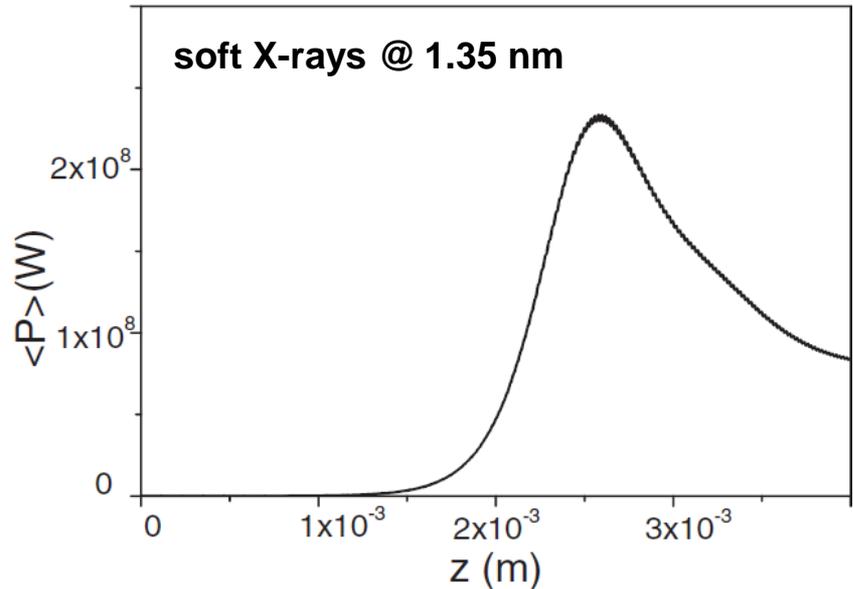
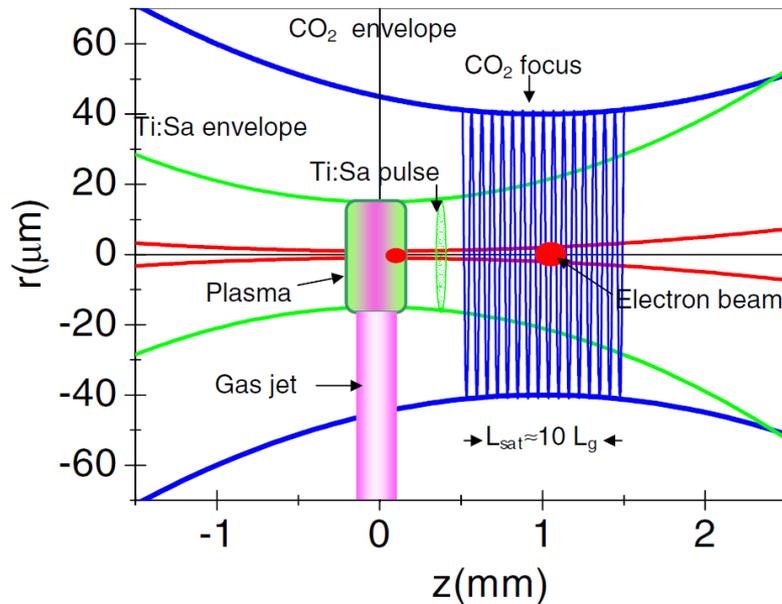
indicate difficulties in realization.

Challenges for OFELs

- **Interaction distances needs to be long**
Else the FEL instability cannot develop.
- **High energy efficiency**
Else too much (laser) energy required.
- **Avoid ultra-high intensities $a_0 \gg 1$**
Else impractical for real-world lasers.
- **One needs:** An optical undulator, which is both long and narrow.

Early „head-on“ OFEL designs show undulator wavelengths beyond optical wavelengths

Petrillo *et. al.*, PRSTAB **11**, 070703 (2008)



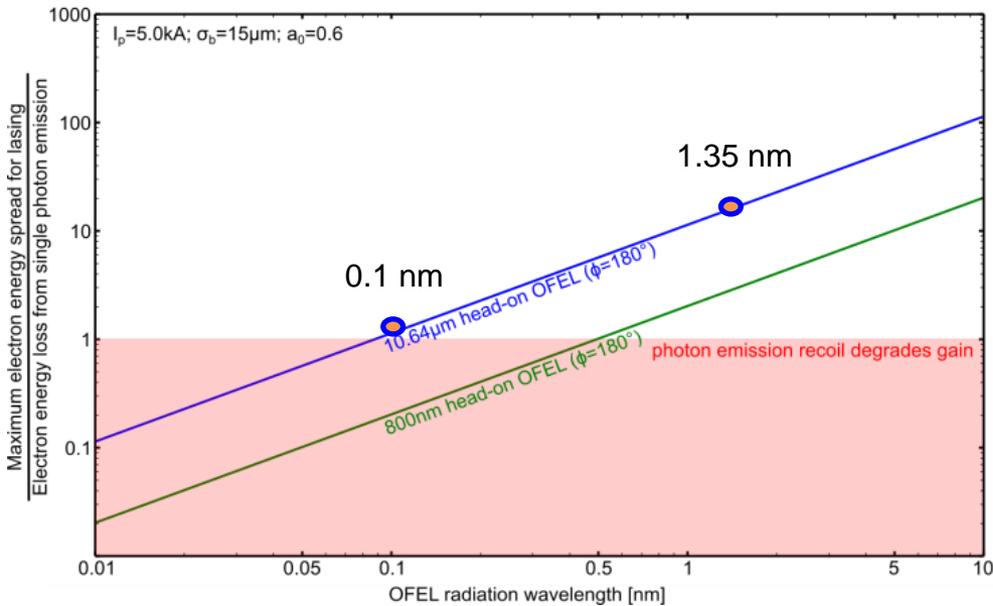
undulator laser

CO₂ laser, $\lambda=10\mu\text{m}$, $\sim 18.7\text{ps}$, $\sim 17\text{J}$, $40\mu\text{m}$ spot size, $a_0=0.8$

electrons

28 MeV ; >20 kA LWFA bunch ; 0.4% energy spread
 $\sim 10\text{fs}$ duration ; <0.3 mm mrad transverse emittance slices

„Head-on“ OFELs do not scale to the hard X-ray range



- High-power, flattop CO₂ laser on ps-scale at <1.2% field stability is challenging.
 - High photon emission recoil instantly kicks electrons out of FEL resonance.
- **Going to sub-nm X-rays not possible for „head-on“ designs.**

Petrillo *et. al.*, PRSTAB 11, 070703 (2008)

240 MW soft X-rays @ 1.35 nm

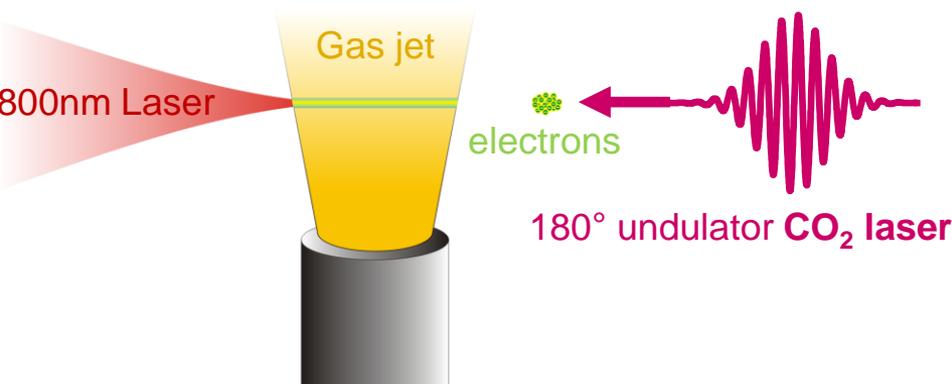
undulator laser

CO₂ laser, $\lambda=10\mu\text{m}$, $\sim 18.7\text{ps}$, $\sim 17\text{J}$, $40\mu\text{m}$ spot size, $a_0=0.8$

electrons

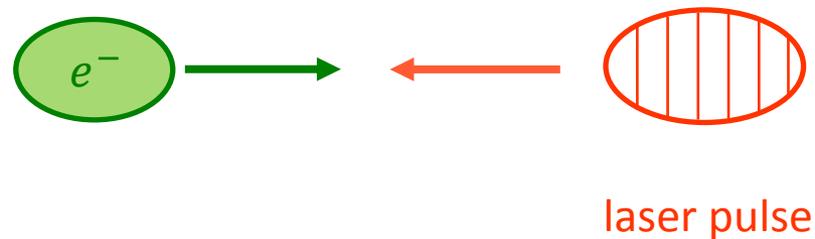
28 MeV ; >20 kA LWFA bunch ; 0.4% energy spread

$\sim 10\text{fs}$ duration ; <0.3 mm mrad transverse emittance slices



Traveling-Wave Thomson-Scattering (TWTS) in a nutshell

Using pulse-front tilted petawatt lasers in side scattering geometries for arbitrarily long interaction distances

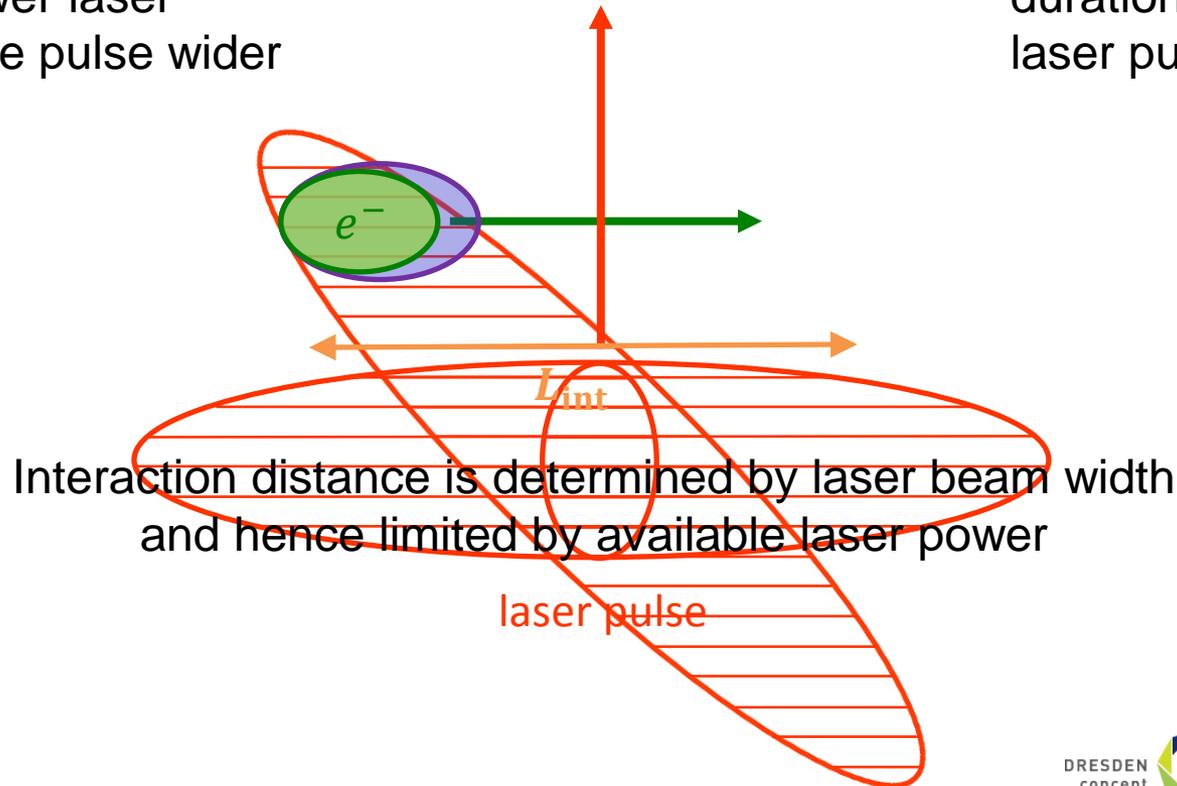


Traveling-Wave Thomson-Scattering (TWTS) in a nutshell

Using pulse-front tilted petawatt lasers in side scattering geometries for arbitrarily long interaction distances

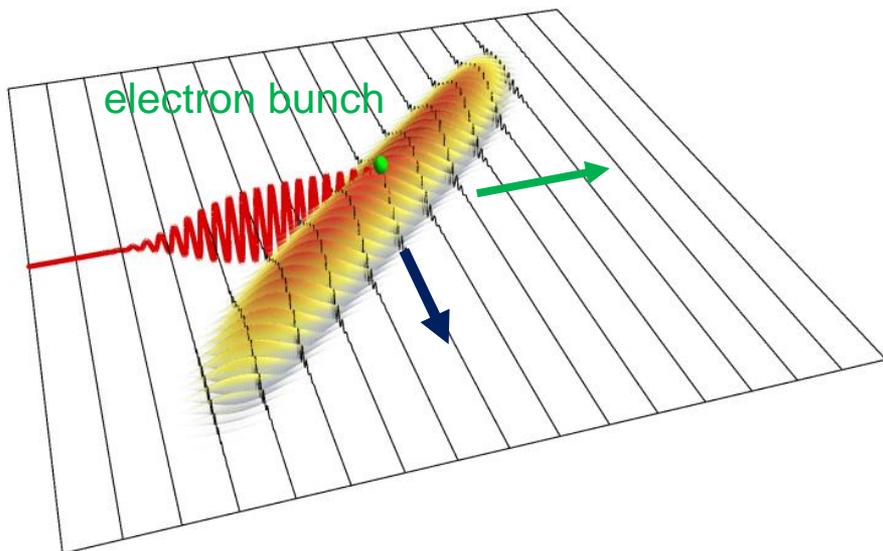
1) Reduce the local intensity of a high power laser by making the pulse wider

2) Increase interaction duration by tilting the laser pulse front



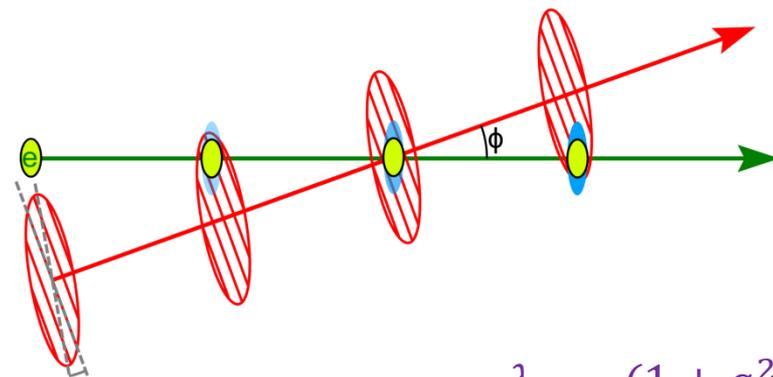
The undulator wavelength of a TWTS OFEL is scalable

Pulse front tilt of half the interaction angle ensures continuous overlap independent of interaction angle.

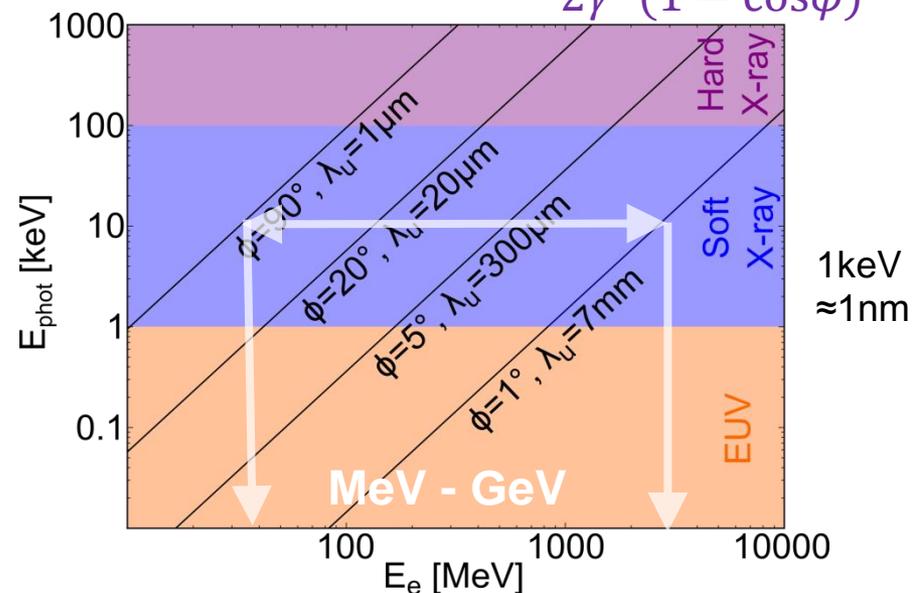


Note

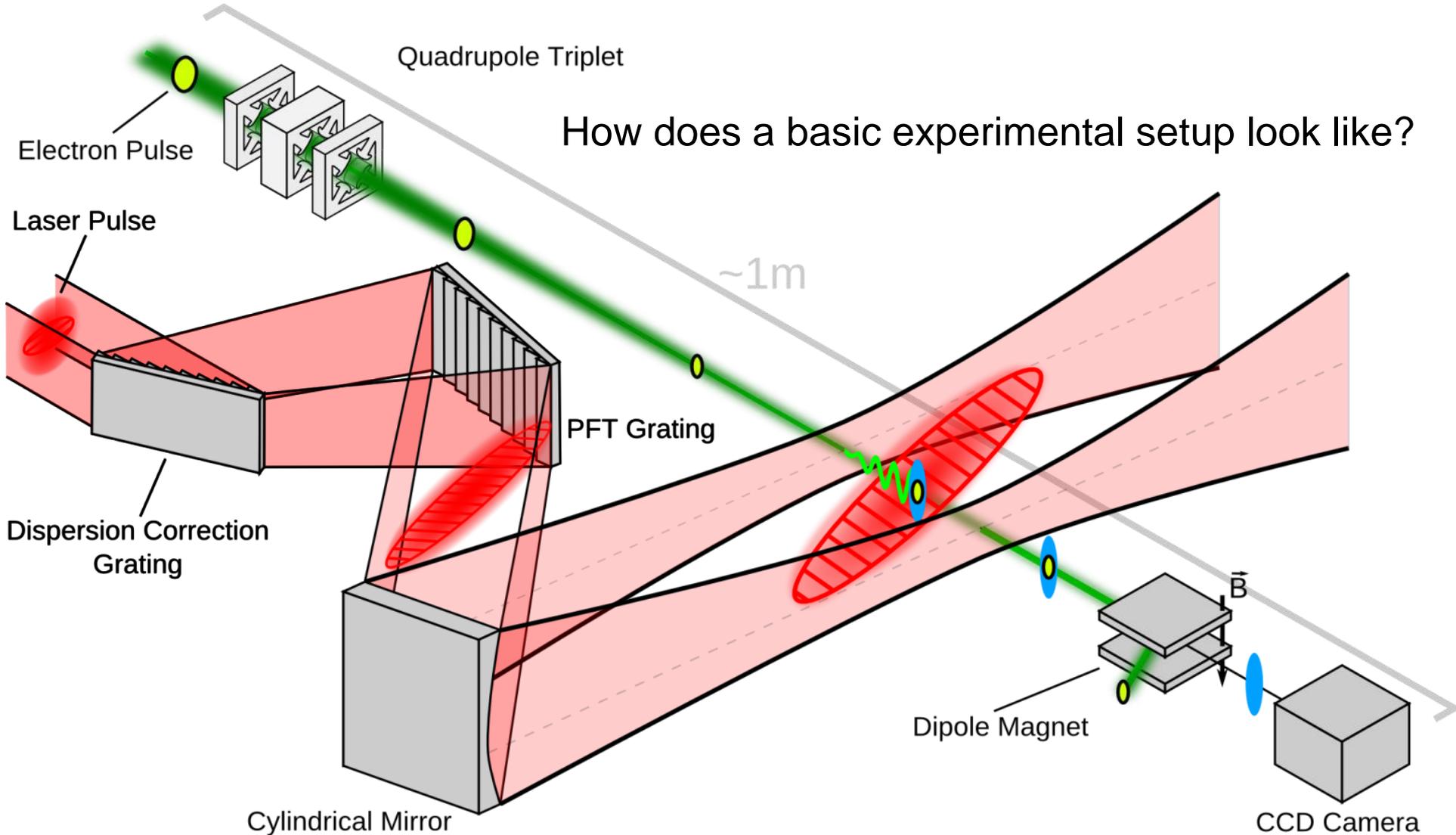
The number of laser periods seen by an electron exceeds the periods of the laser pulse duration.



$$\alpha_{tilt} = \phi/2 \quad \lambda_{FEL} = \frac{\lambda_{laser}(1 + a^2/2)}{2\gamma^2(1 - \cos\phi)}$$



Traveling Wave-Thomson Scattering OFEL in experiment

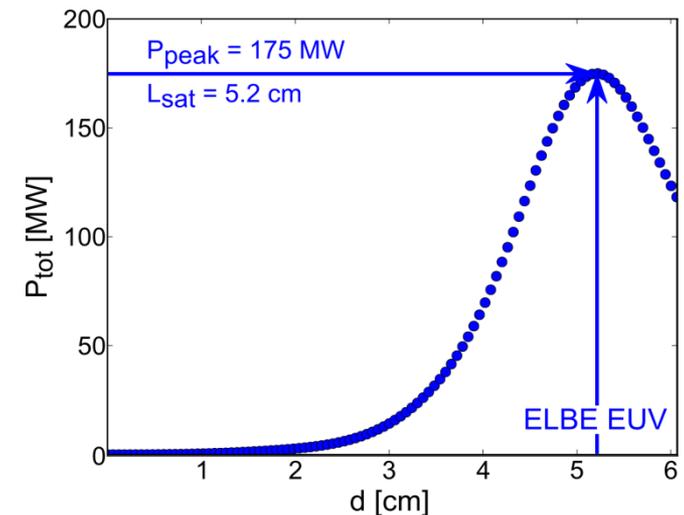


Steiniger, Debus *et. al.*, J. Phys. B **47** (2014) 234011

Design 1

13.5nm Flash-type TWTS OFEL at ELBE energy

Parameter	EUV TWTS OFEL ELBE@HZDR	FLASH conventional FEL
Scattered wavelength [nm]	13.5	13.5
Interaction angle [deg]	6.6	-
Undulator period [mm]	0.15	27.3
Electron energy [MeV]	40	700
Peak current [kA]	2.5	2.5
Norm. emittance [mm mrad]	1.0	1.45
Rel. energy spread	0.2%	0.2%
Undulator parameter a_0/K	0.5	1.2
Laser power [TW]	995	-
Gain length [cm]	0.35	250
Interaction distance [m]	0.052	27
Peak power [MW]	175	10000



Simulation of ELBE EUV
TWTS OFEL based on
analytic 1.5D TWTS OFEL
equations

Traveling-wave Thomson Scattering (**TWTS**) OFELs can be realized with state-of-the-art accelerators and lasers.

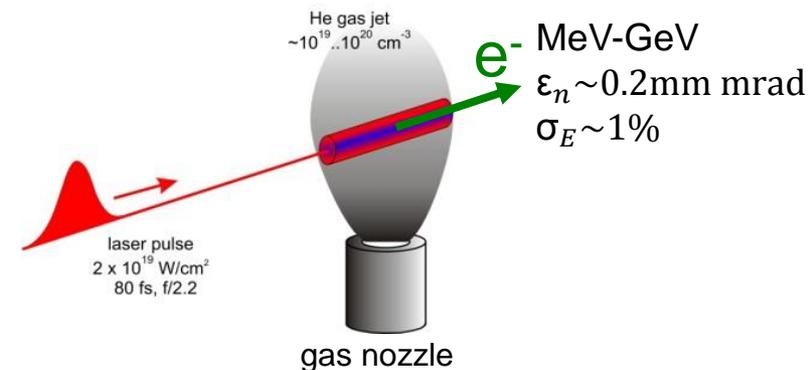
Design 2

Ultra compact TWTS OFEL using LWFA electrons

Parameter	EUV TWTS OFEL LWFA
Scattered wavelength [nm]	13.5
Interaction angle [deg]	3.0
Undulator period [mm]	0.578
Electron energy [MeV]	98
Peak current [kA]	5 (50pC/10fs)
Norm. emittance [mm mrad]	0.2
Rel. energy spread	1%
Undulator parameter a_0/K	1.2
Laser power [TW]	865
Gain length [cm]	0.27
Interaction distance [m]	0.043
Peak EUV power [MW]	4888

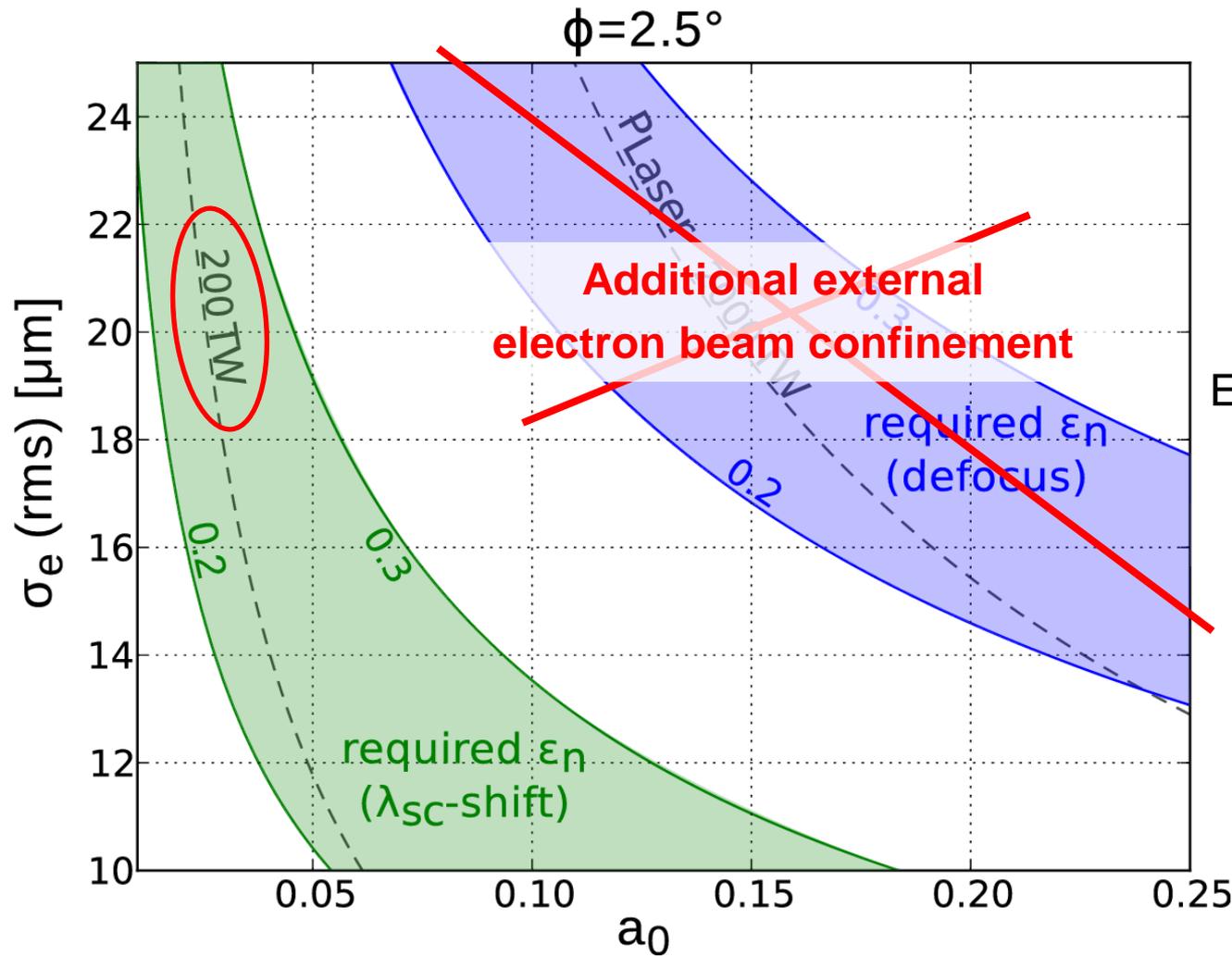
Use ultra-low emittance of LWFA electrons for tight focusing to $5\mu\text{m}$ (rms).

This maximizes radiation-electron coupling for LWFA electrons with 1% level energy spread.



- TWTS+LWFA offers inherent synchronization of laser and electrons.

Electron beam trapping could drastically reduce laser power requirements

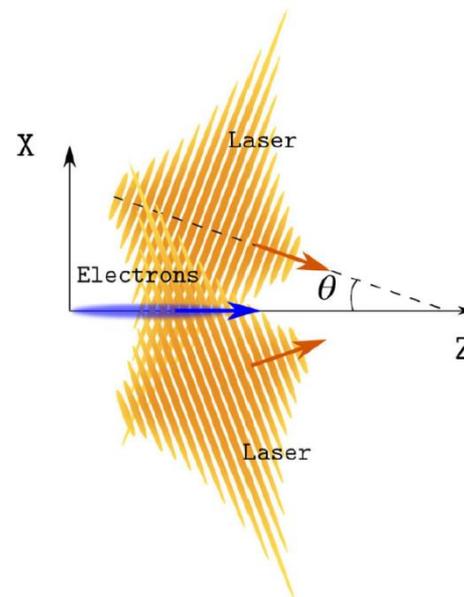
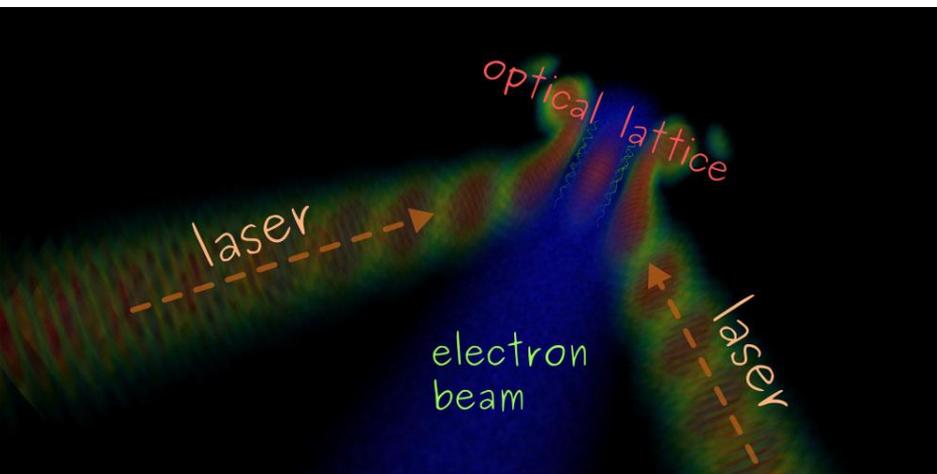


ϵ_n for FEL instability

ϵ_n for maintaining focus

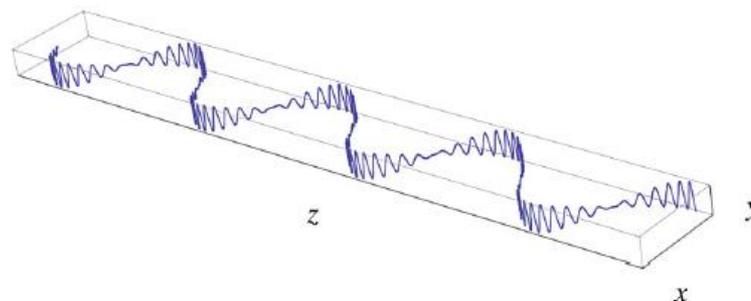
Emittance and laser power landscape for an 0.15nm X-ray OFEL

Optical lattices provide electron beam trapping



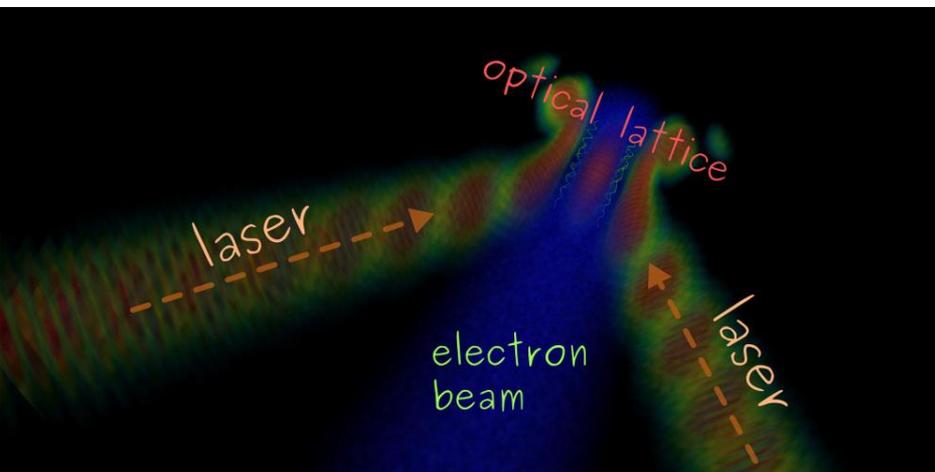
- Two counter propagating TWTS laser beams provide an optical lattice.
- Electrons experience
 - fast laser oscillation
 - slow ponderomotive oscillation in the optical lattice potential

Elliptical trajectory alters FEL dynamics



Andriyash *et. al.*, PRSTAB **18**, 050704 (2015)

Optical lattices provide electron beam trapping



laser

800nm; $2 \cdot 10^{16}$ W/cm² ($a_0=0.1$); ~300mJ

$\theta=10^\circ \rightarrow \lambda_u=52\mu\text{m}$

electron beam

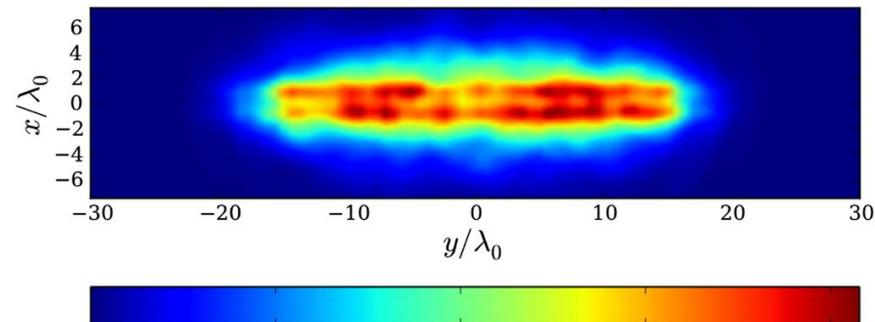
40 MeV; 100pC; 10fs; 1 μm diameter

1mrad (x) and 0.36mrad (y) divergence

x-ray beam

4.4nm; 200MW; 2cm saturation length; $\rho=4.4 \cdot 10^{-3}$

seeded FEL, transverse mode

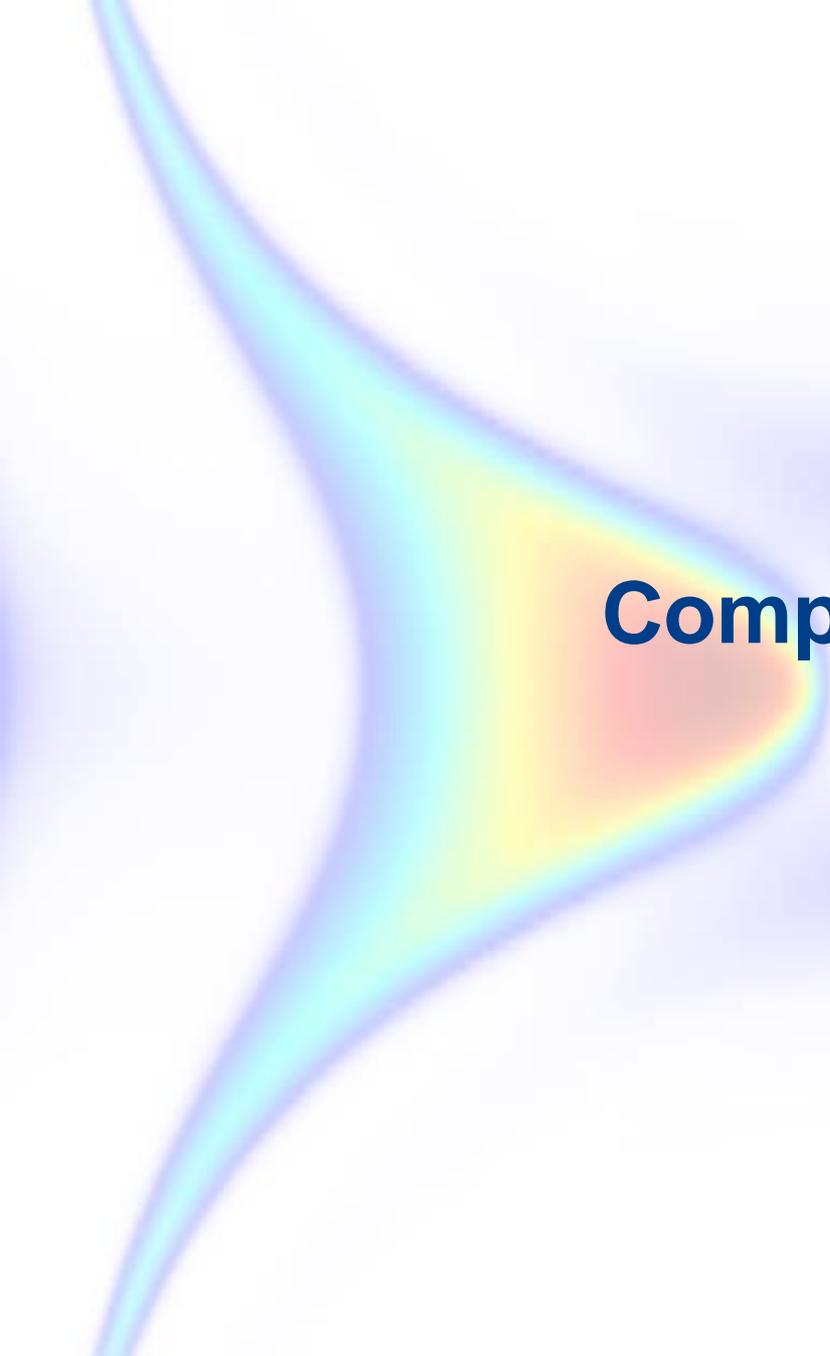


- Electron beam trapping in one dimension

Andriyash *et. al.*, PRSTAB **18**, 050704 (2015)

Andriyash *et. al.*, PRSTAB **16**, 100703 (2013)

Andriyash *et. al.*, PRL **109**, 244802 (2012)



Compton sources

RF accelerator-driven Compton source

PHOENIX source @ HZDR

Fresh results from last week

See **Axel Jochmann's** talk
on the **PHOENIX source**
at WG4 – **today 17:00**

**measured spectrum
photondensity
over energy & observation angle
in log. scale**

Removed unpublished data in Web-version.

ELBE RF accelerator (thermionic gun)

23.5 MeV ; 3ps; 80pC ; 0.4% energy spread
norm. transverse emittance 11 mm mrad

X-ray beam

$5.5 \cdot 10^5$ photons/shot total yield
30% BW @ 11.7 keV

scattering DRACO laser

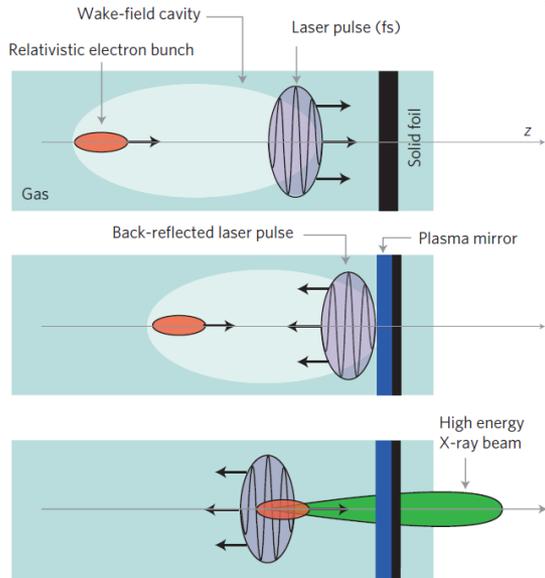
2J ; 27fs ; $a_0=1.5$; 35 μ m spot size

Peak brilliance

$1.5 \cdot 10^{17} \text{ s}^{-1} \text{ mm}^{-2} \text{ mrad}^{-2}$ [0.1% BW]

State-of-the-art LWFA Compton sources provide either broadband high photon yield or low-bandwidth

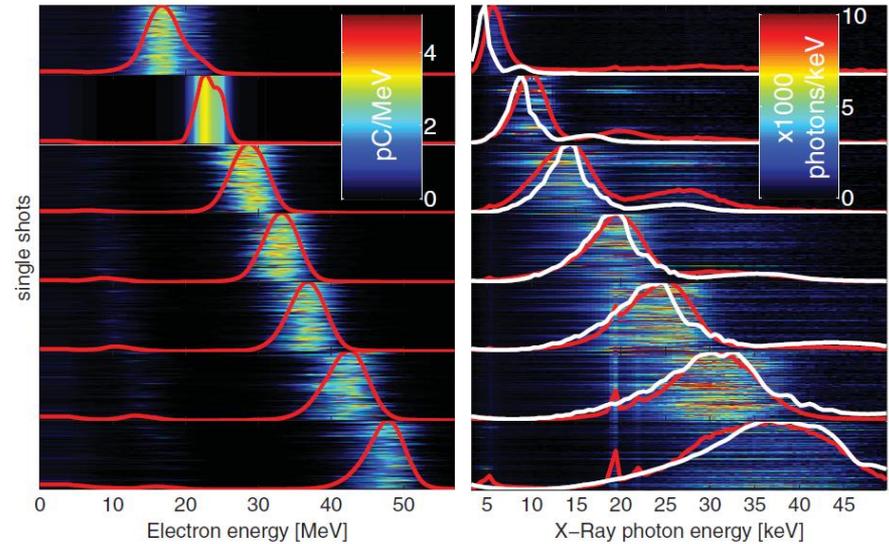
Ta Phuoc *et. al.*, PRL **114**, 195003 (2015)



tight overlap, but high intensity; $a_0 \sim 1.5$
 $17\mu\text{m}$ laser spot size

Broadband Compton spectra

K. Khrennikov *et. al.*, PRL **113**, 224801 (2014)

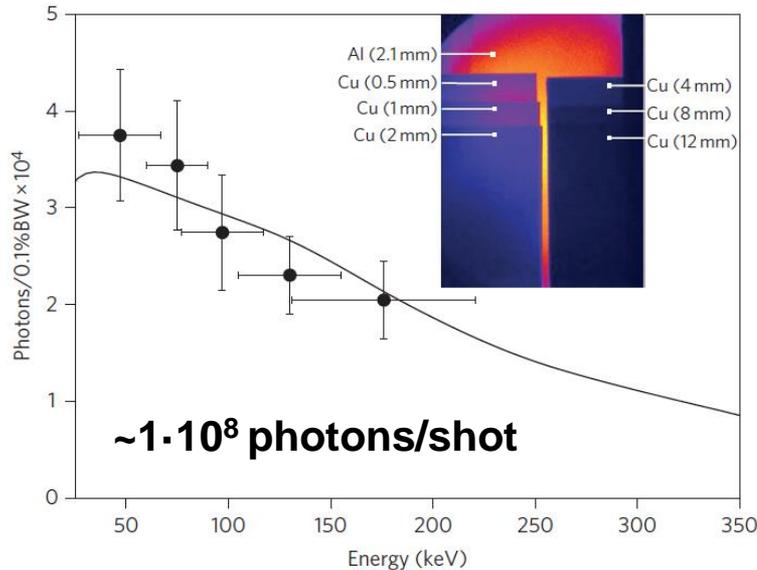


$25\mu\text{m}$ laser spot size
 max. ~ 40000 photons detectable

Tunable narrowband Compton spectra

State-of-the-art LWFA Compton sources provide either broadband high photon yield or low-bandwidth

Ta Phuoc *et. al.*, PRL **114**, 195003 (2015)



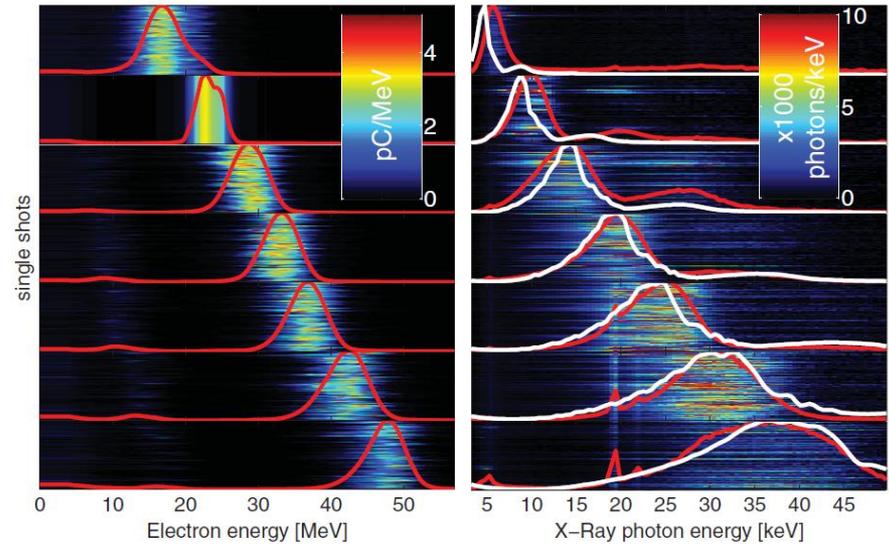
scattering laser

Ti:Sa, 1J, 35 fs, $a_0=1.5$

peak brilliance

$1 \cdot 10^{21} \text{ s}^{-1} \text{ mm}^{-2} \text{ mrad}^{-2}$
[0.1% BW]

K. Khrennikov *et. al.*, PRL **113**, 224801 (2014)



scattering laser

Ti:Sa, 300mJ, 28 fs, $a_0=0.9$

max. ~40000 photons detectable

peak brilliance

$0.2..15 \cdot 10^{18} \text{ s}^{-1} \text{ mm}^{-2} \text{ mrad}^{-2}$
[0.1% BW]

Sarri *et. al.*, PRL **114**, 195003 (2015)

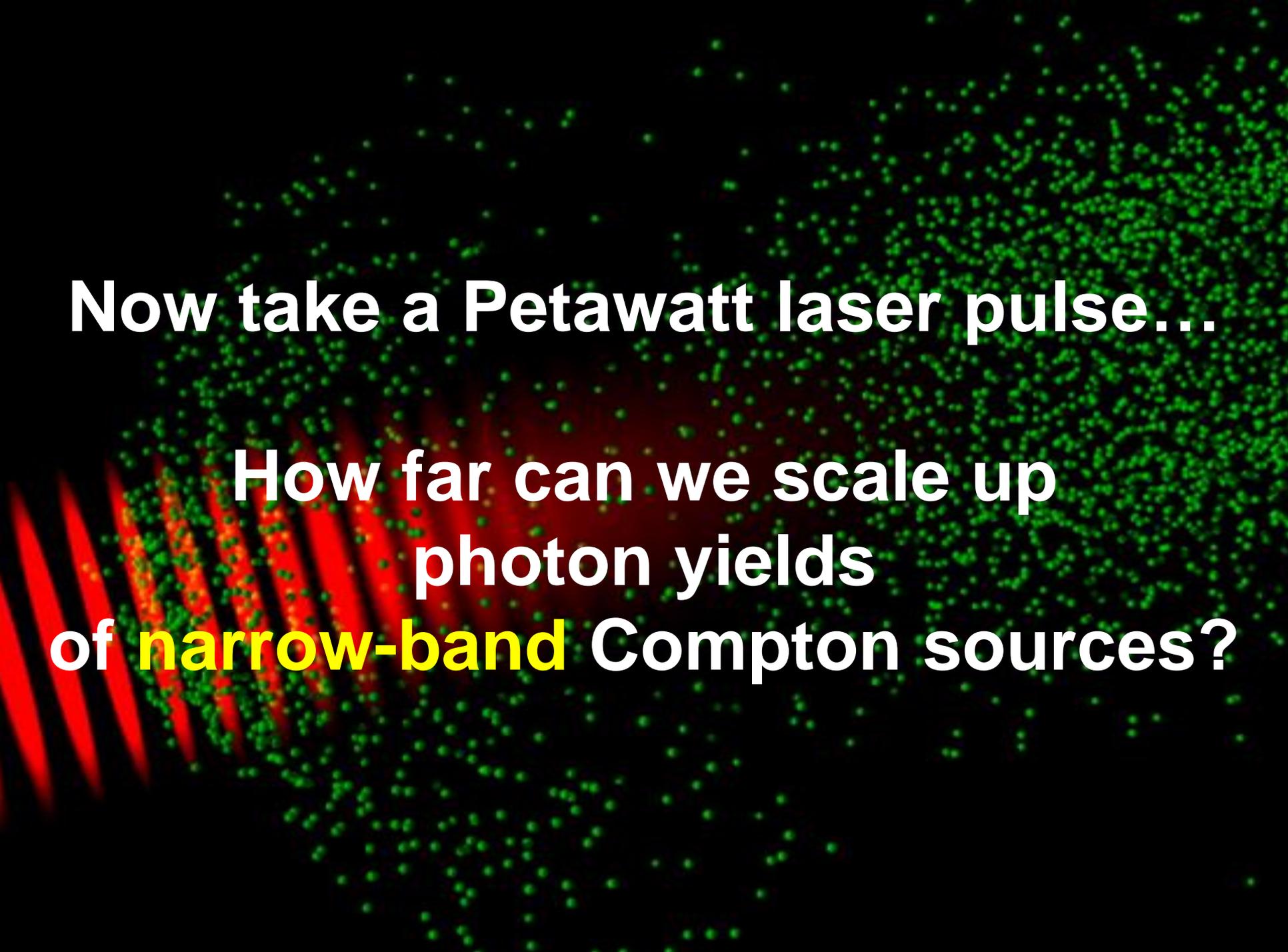
peak brilliance

$1.8 \cdot 10^{20} \text{ s}^{-1} \text{ mm}^{-2} \text{ mrad}^{-2}$ [0.1% BW] @ 6-18 MeV

Powers *et. al.*, Nat. Phot. **8**, 28-31 (2014)

peak brilliance

$1 \cdot 10^{21} \text{ s}^{-1} \text{ mm}^{-2} \text{ mrad}^{-2}$ [0.1% BW] @ 70-100 keV



Now take a Petawatt laser pulse...

How far can we scale up
photon yields
of **narrow-band** Compton sources?

Traveling-wave Thomson scattering improves photons yield per shot by 2-3 orders of magnitudes

	PW-TWTS ($\phi=120^\circ$)	PW-TWTS (LWFA) ($\phi=120^\circ$)
wavelength [μm]	1	1
W_{pulse} [J]	200	200
pulse length [fs]	150	150
I_0 [W/cm^2]	5.5×10^{16}	5.7×10^{16}
L_{int} [mm]	57.6	200
photon energy [keV]	23	3500
$N_{\text{phot},5\%}$ per pulse	4.9×10^{10}	3.4×10^{10}

@ 5% bandwidth



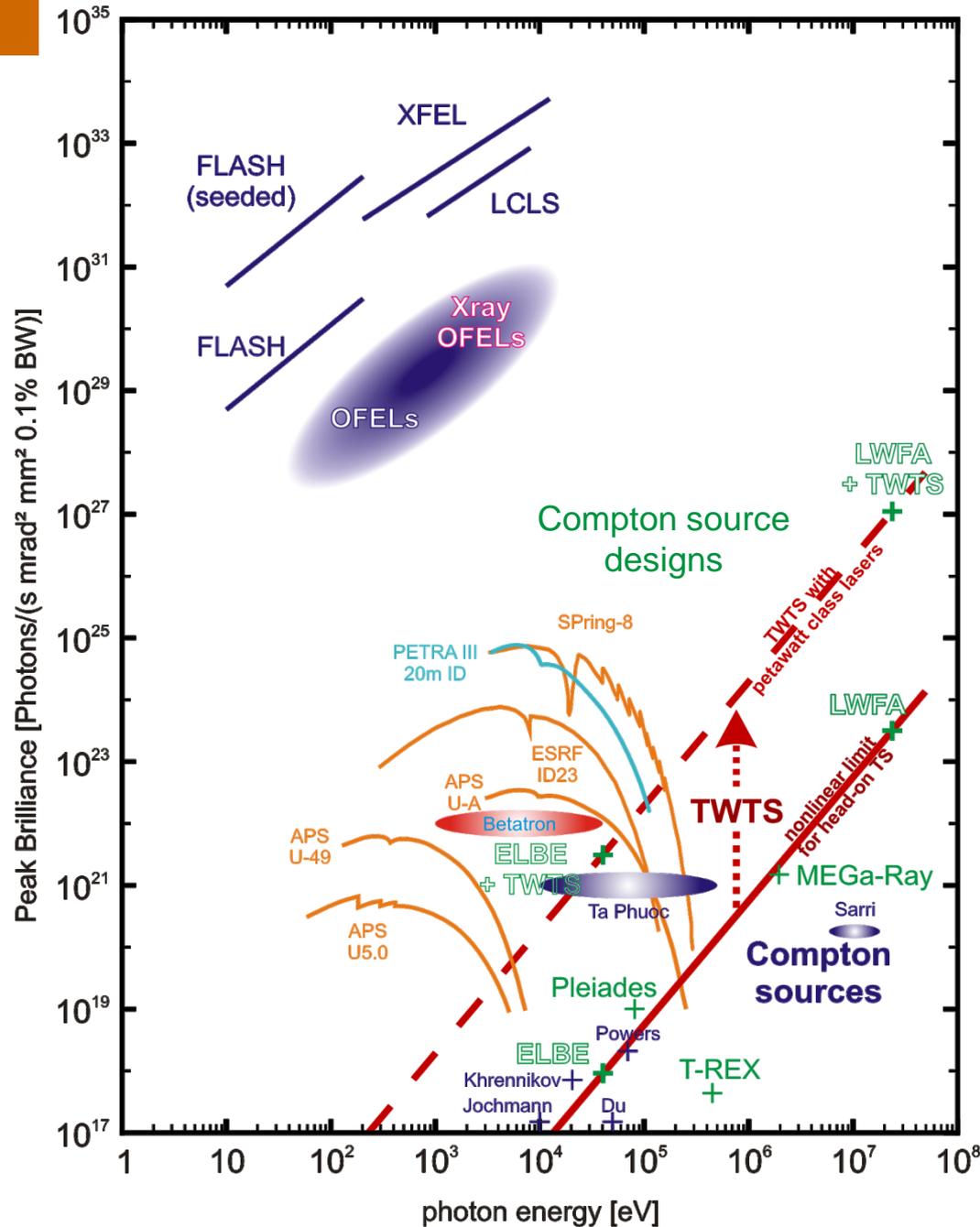
Debus *et. al.*
Appl. Phys. B (2010) **100**(1) 61-76

Electrons

ELBE electrons: 40 MeV, 1nC, 2π mm mrad (norm. trans. emittance)

LWFA electrons: 500 MeV, 200 pC, 0.1π mm mrad (norm. trans. emittance)

Towards brilliant Compton sources



- „Head-on“ Compton sources are limited in peak brilliance by nonlinear Thomson scattering.

- TWTS sources can increase peak brilliances in the X-ray range by 2 - 3 orders of magnitudes.

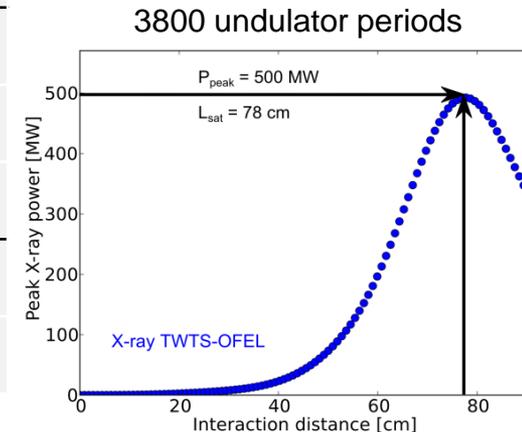
- The brilliance limit using TWTS is purely technical.

It arises from electron divergence, size of optics and available laser power.

Outlook

How would a LCLS-type TWTS-OFEL at 1.5Å look like?

Parameter	TWTS	LCLS
scattered wavelength [Å]	1.5	1.5
electron energy [GeV]	0.46	13.6
peak current [kA]	5	3.5
norm. emittance [mm mrad]	0.24	0.5
rel. energy spread	0.02 %	0.01 %
undulator parameter (a_0 or K)	0.2	3.5
laser power [PW]	1.1	-
interaction angle [°]	5	-
saturation length [m]	0.8	132
peak power [GW]	0.5	40



Required laser intensity uniformity within one gain length \rightarrow 2.4%

Conclusions

- Optical Free-Electron Lasers (OFELs) can be realized with existing electron sources and laser systems.
- In contrast to head-on OFEL designs, TWTS OFELs have favorable properties with regard to *transverse coherence*, *space charge* and *photon recoil*.
- TWTS OFELs are scalable to higher photon energies. An X-ray TWTS OFEL would remain compact.
- Narrowband, peak-brilliance Compton sources scale beyond 10^{10} photons per shot @ 5% bandwidth.



Thank you for your attention!

Acknowledgements

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Jakob Krämer, Axel Jochmann, Arie Irman