

# Optimized matching strategy for laser driven plasma booster

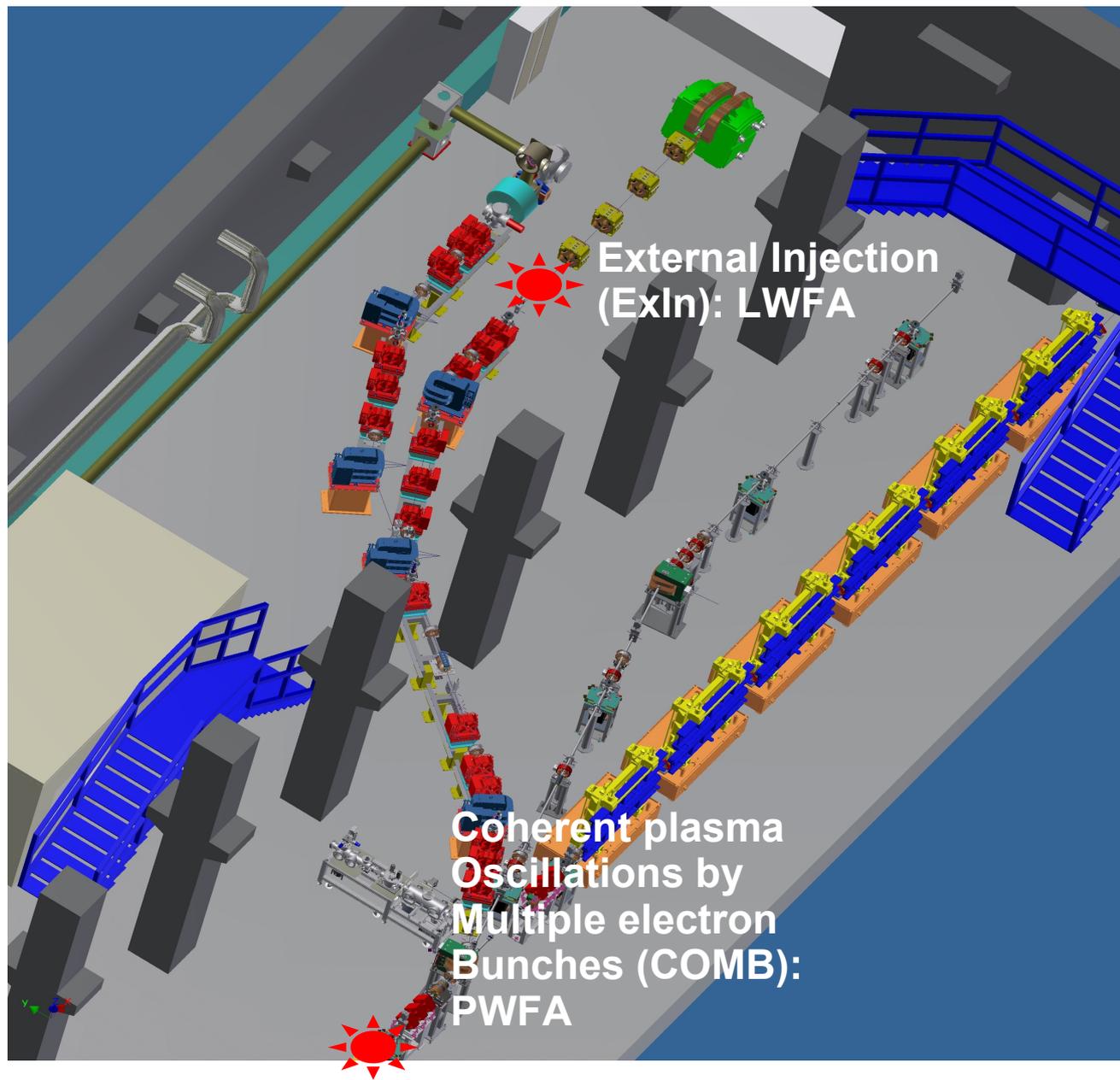
Andrea R. Rossi\*  
on behalf of the SPARC\_LAB collaboration



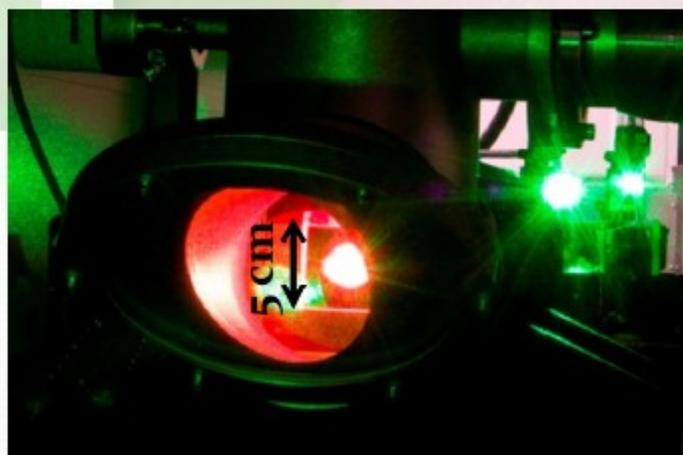
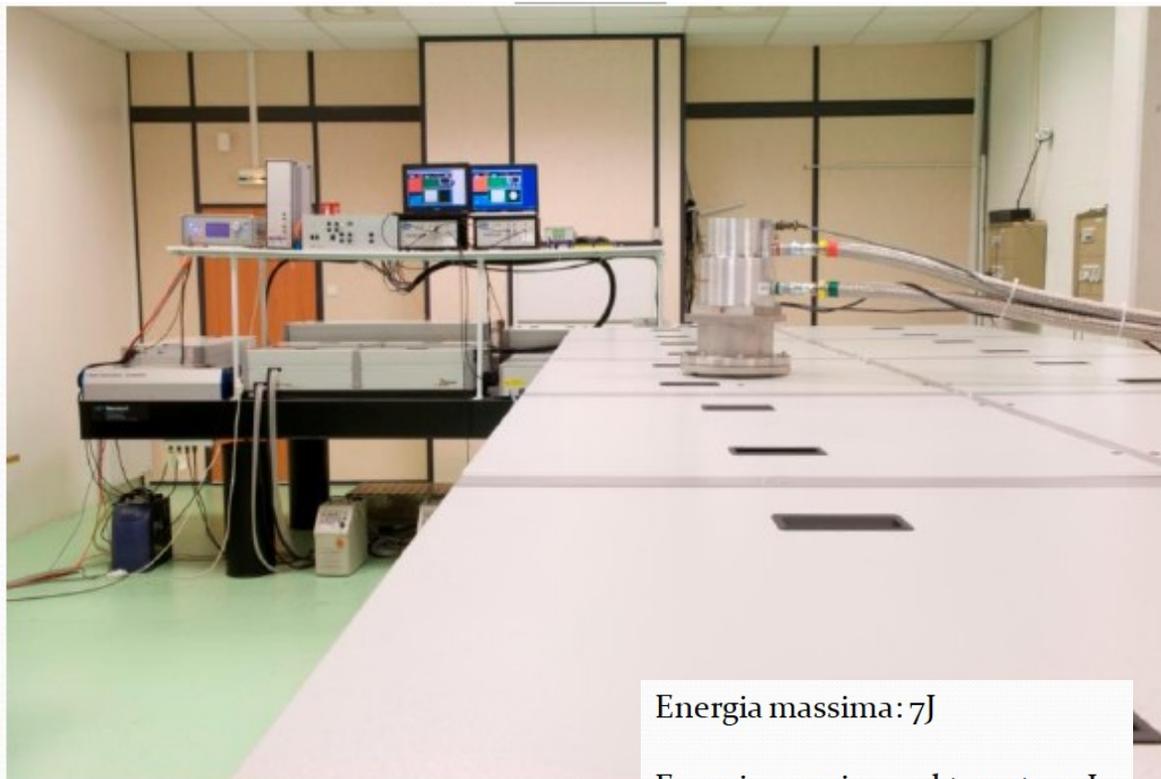
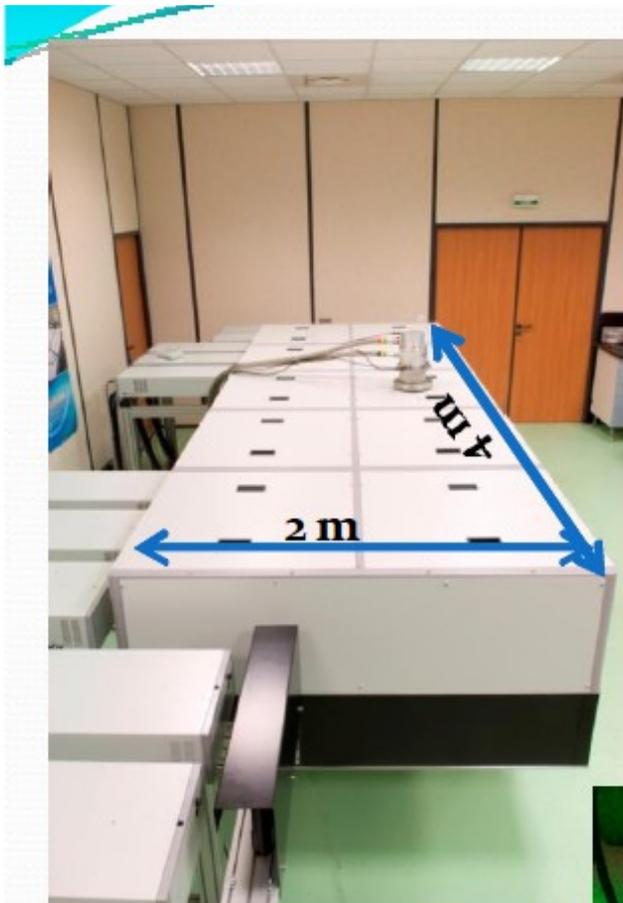
\* [andrea.rossi@mi.infn.it](mailto:andrea.rossi@mi.infn.it)



# External injection experiment @ SPARC\_LAB



# ExIn: the FLAME laser



Energia massima: 7J

Energia massima sul target: ~5J

Durata minima: 23 fs

Lunghezza d'onda: 800 nm

Larghezza di banda: 60/80 nm

Spot-size @ focus: 10  $\mu\text{m}$

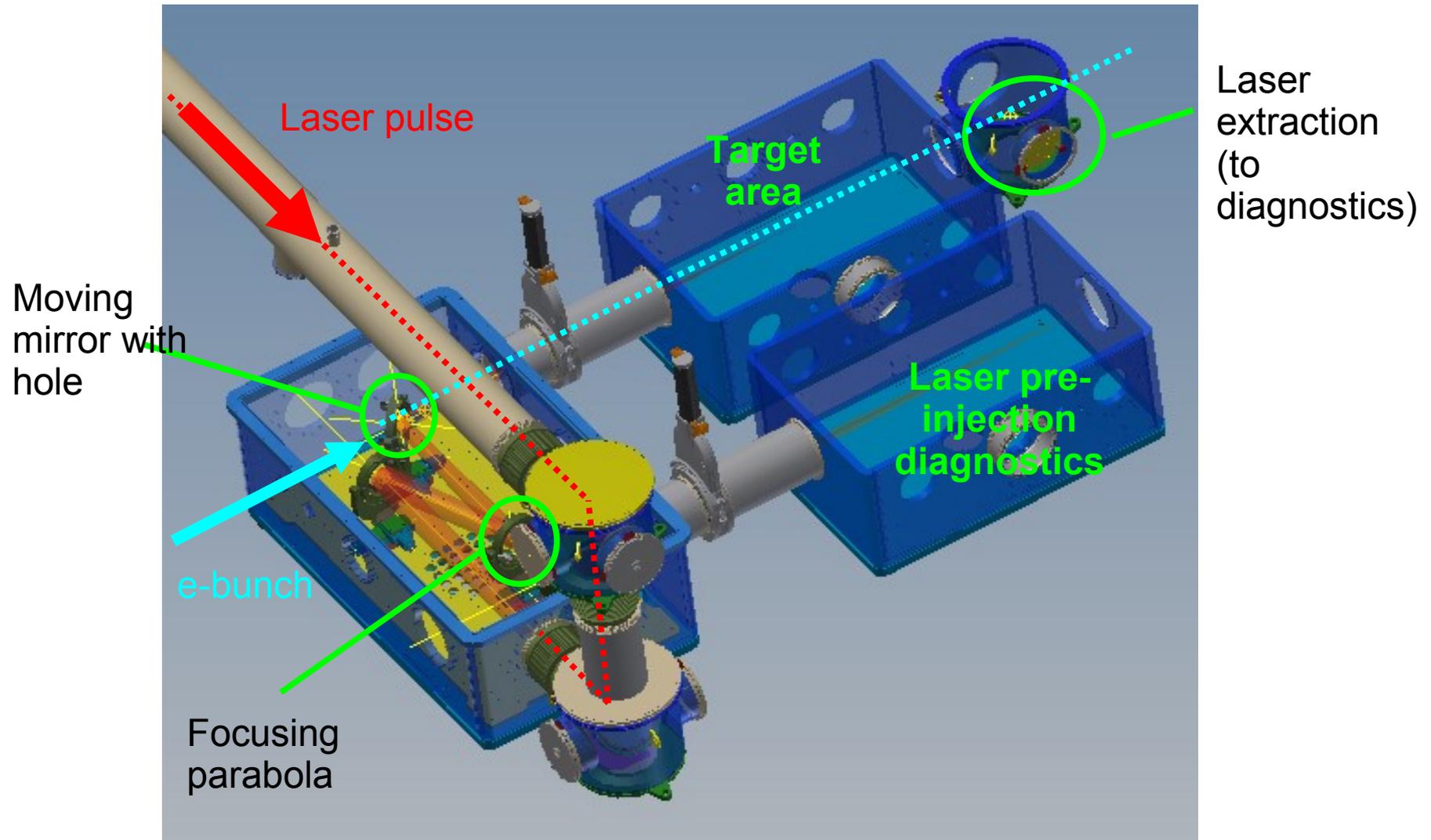
Potenza massima: ~300 TW

Contrasto:  $10^{10}$

## ExIn goals

- Demonstrate GV/m acceleration fields
- Stability.
- Reproducibility.
- Preserve beam brightness

# ExIn interaction chamber

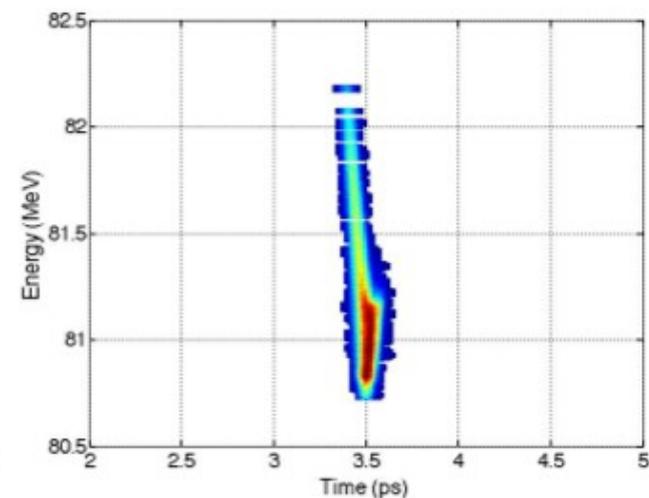
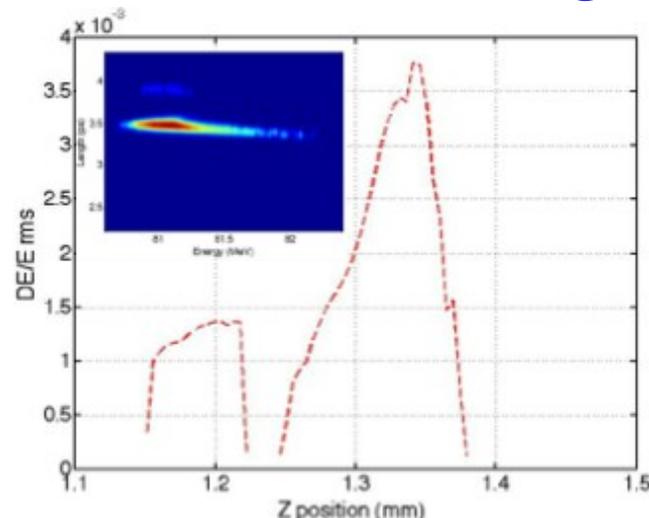
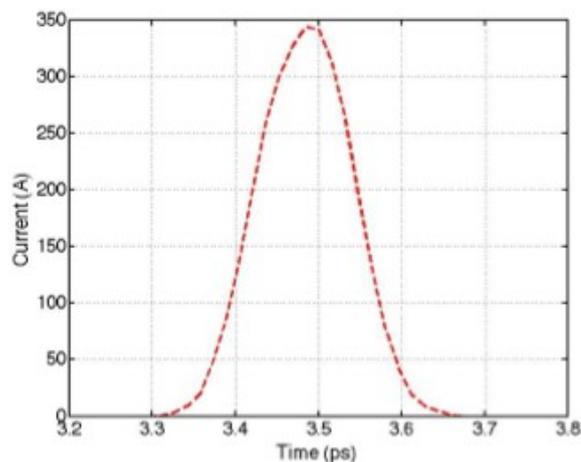


# ExIn experimental parameters

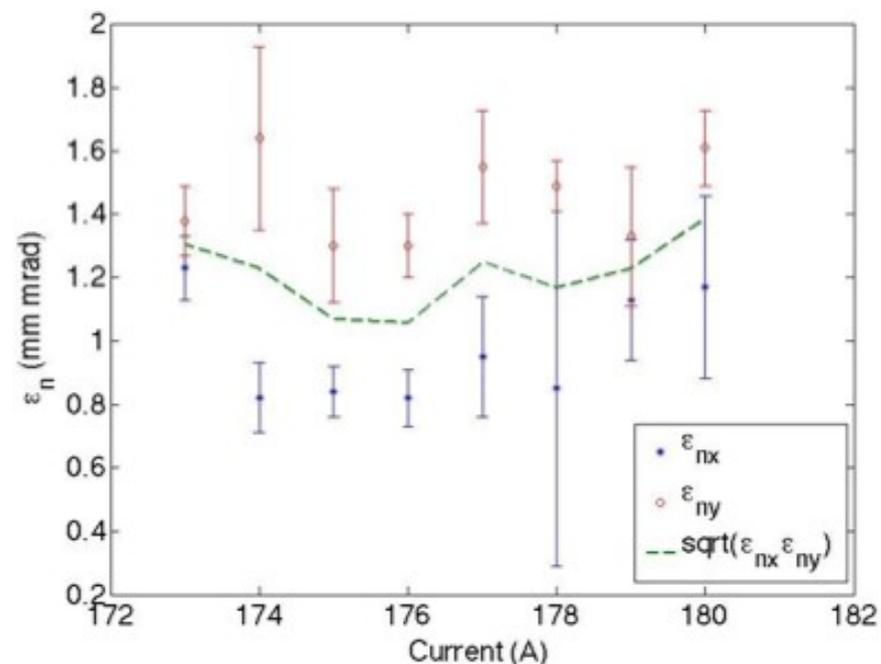
- Acceleration by LWFA of an electron beam produced by the SPARC photoinjector.
- Incoming bunch properties:  $E = 70 - 110$  MeV;  $L_{FWHM}$ : 20 – 30 fs;  $\epsilon_{nt}$ : few  $\mu\text{m}$ ;  $\sigma_{tr}$ : down to 4  $\mu\text{m}$  with permanent magnet quads;  $\delta E/E$ : some  $10^{-3}$ ; charge: 5 – 30 pC.
- Plasma wave in quasi-linear regime for best compromise between fields strength, acceleration length (pump depletion), de-phasing, tolerances to jitters, ecc...
- Laser is guided by dielectric capillaries exploited as hollow waveguide\*, though first preparatory runs will be performed with larger laser waist size and without any guiding.
- Assuming  $E_{laser} \approx 3.5$  J and  $q \approx 20$  pC the working point is chosen to be with  $R_{cap} = 60$   $\mu\text{m}$  and  $n_0 = 10^{17} \text{cm}^{-3}$  \*\*, with average accelerating field around 7 GV/m.
- Expected performances:  $E_f = 300 - 800$  MeV,  $\epsilon_f < 1.2 \epsilon_0$ ,  $\delta E/E < 10^{-2}$ .

\*B. Cross *et al.*, PRE 65, 026405 (2002).\*\*A.R. Rossi *et al.*, PRSTB 740, 60 (2014).

# Beam dynamics experimental results: EXIN setting



Energy (MeV)	81.23(0.03)
Energy spread (%)	0.35(0.03)
Bunch duration (ps)	0.048(0.001)
TransvEmit (mm mrad)	1.06(0.13)
Charge (pC)	55(4)
LongEmit (keV*mm)	3.37(0.29)



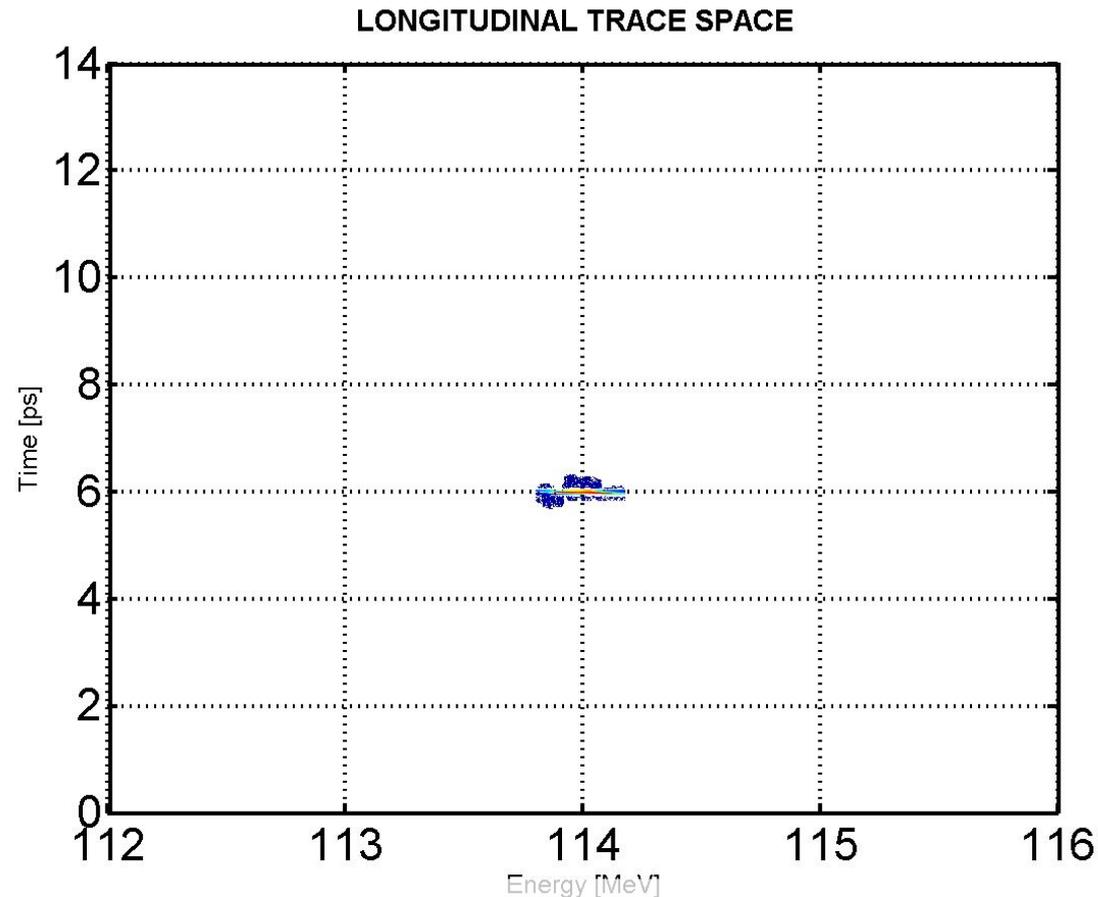
# Beam dynamics experimental results: EXIN setting

$$\sigma_t = 26 \text{ fs}$$

$$\varepsilon_{nx} = 1.2 \mu\text{m}$$

$$\sigma_E = 0.1\%$$

$$I = 400 \text{ A}$$



**Beam matching is critical for preserving  
beam brightness!!**

# Matching strategies

- Matching is critical for preserving beam quality, both when the bunch enters the plasma and when it leaves the interaction area.
- Up to now we know how the focusing strength should vary for adiabatic matching [1] and the theory is checked with simulations avoiding beam loading; a possible practical implementation consists in using plasma lenses [2].
- Plasma lenses are very effective but add to the overall complexity of the system if considered as stand alone optical elements; tapering is a collateral, unavoidable feature, that may act as a plasma lens. It is difficult but not impossible to control. However it may fail to have some needed degree of freedom (e.g. for compensating longitudinal fields variations, beam loading, etc...)
- Matching conditions are analytically established only in bubble regime (and linear regime, which is not so appealing).
- In our setting, simulations show that the matched transverse size is  $< 2 \mu\text{m}$ , while the focusing with permanent magnets quads reaches a size  $> 4 \mu\text{m}$ .

[1] K. Floettmann, PRSTAB 17, 054402 (2014); I. Dornmair et al., PRSTAB 18, 041302 (2015).

[2] R. Lehe *et al.*, PRSTAB 17, 121301 (2014).

# Matching strategies

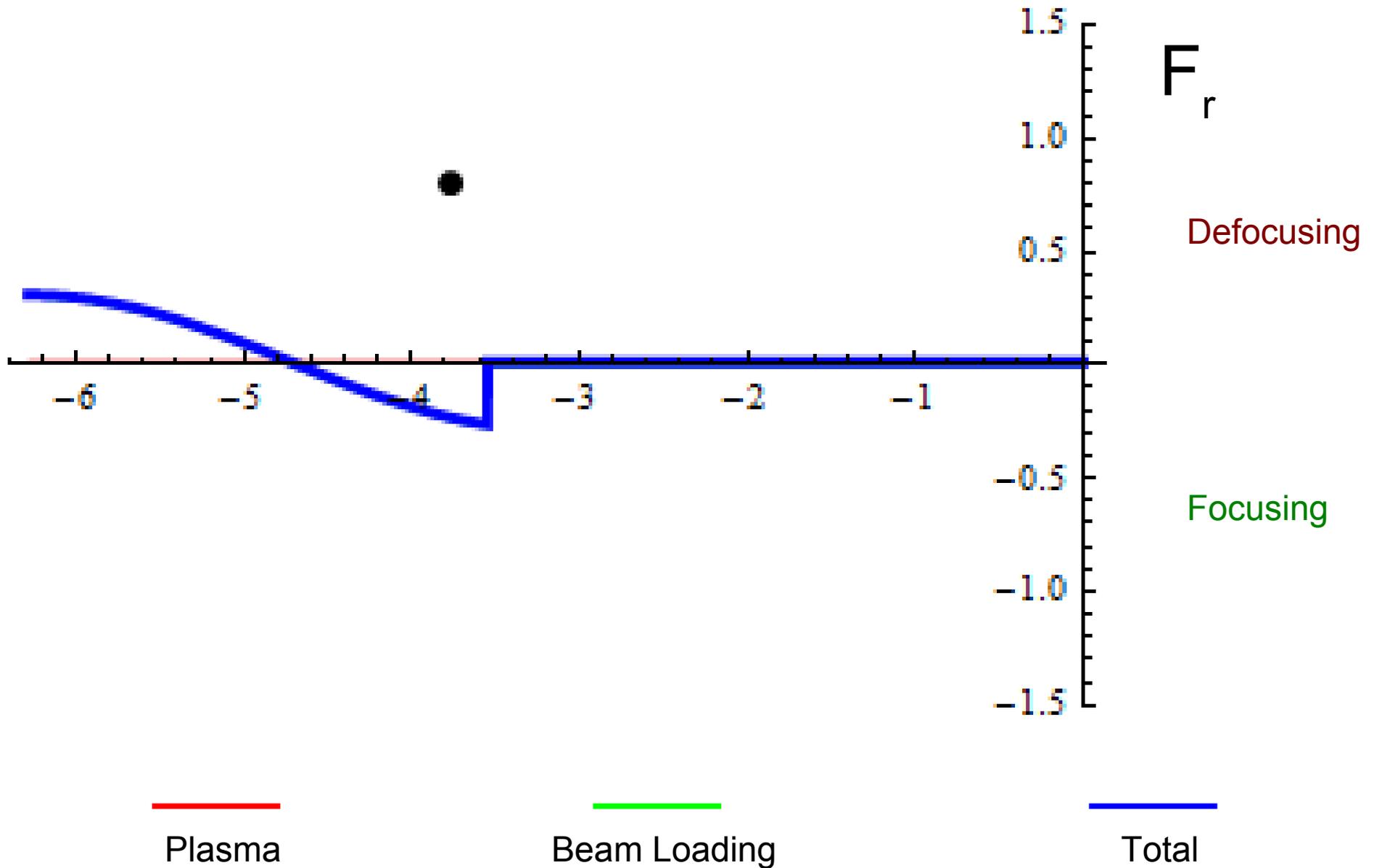
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Our proposal for ExIn is to exploit tapering while the plasma wave regime is varying due to laser focusing. This adds a degree of freedom that can be fruitfully exploited to preserve beam quality. The need for guiding also allows to somehow control the tapering.

[1] K. Floettmann, PRSTAB 17, 054402 (2014); I Dornmair et al., PRSTAB 18, 041302 (2015).

[2] R. Lehe *et al.*, PRSTAB 17, 121301 (2014).

# Matching strategy: the idea



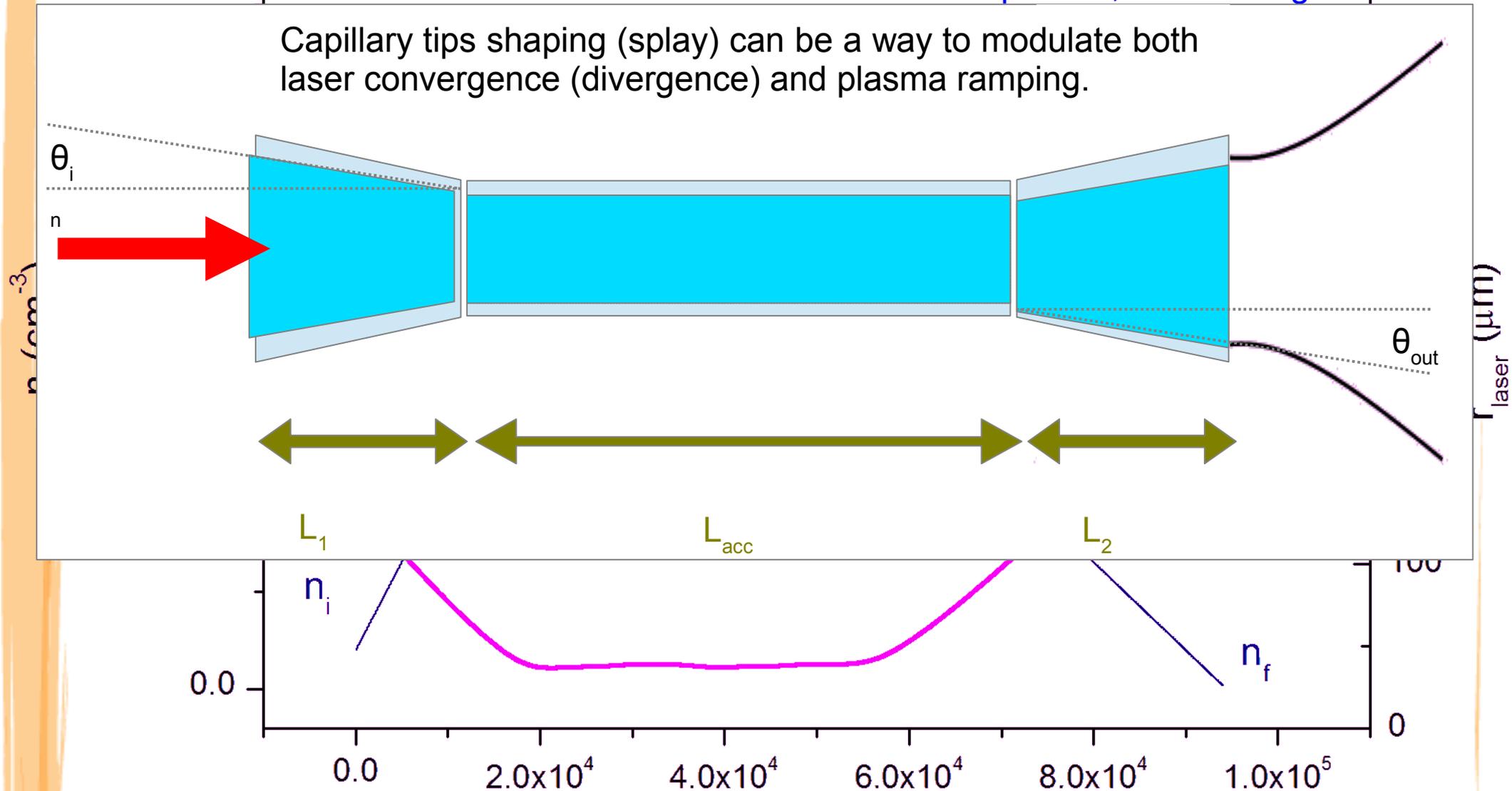
# Simulation settings

 $2.0 \times 10^{17}$ 

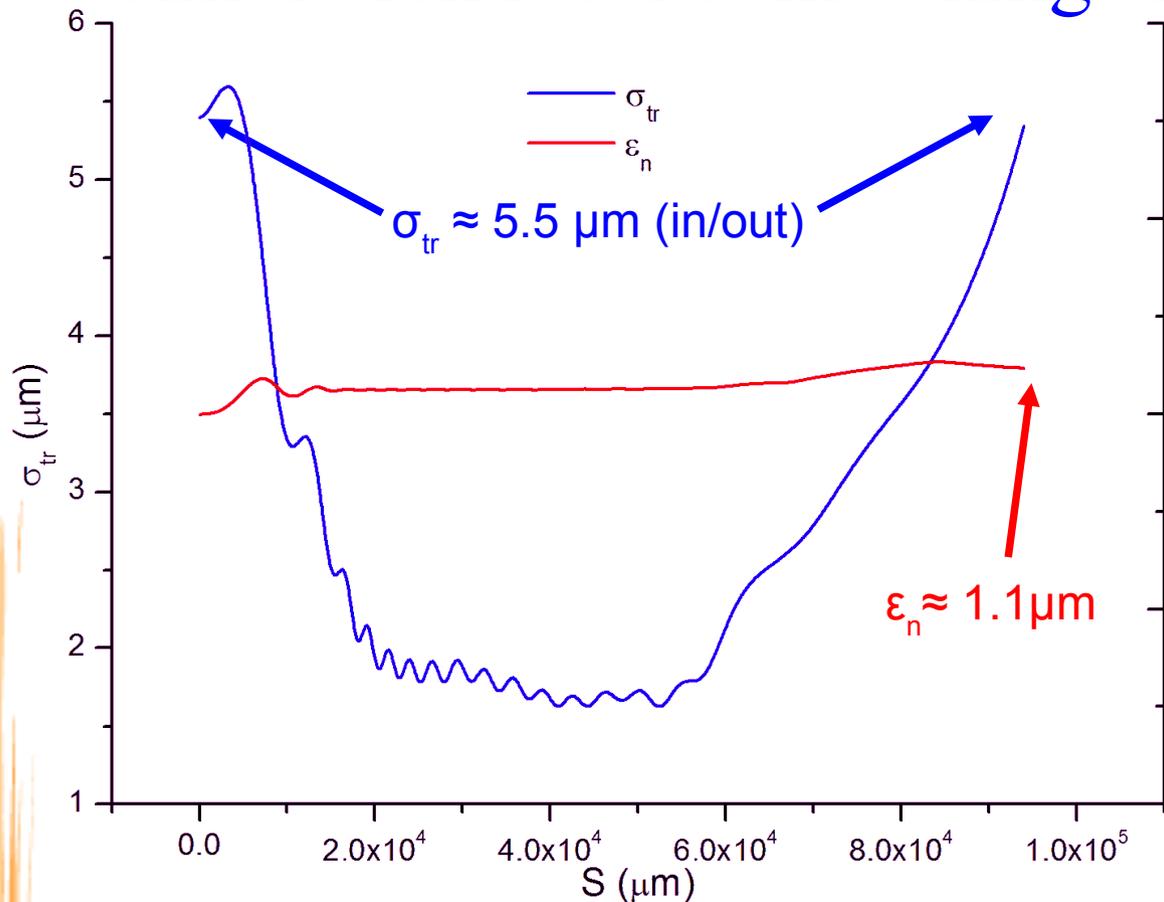
We assume linear tapering for sake of simplicity. The incoming bunch is an **ideal** one with bi-Gaussian profiles, 17 fs long

500

Capillary tips shaping (splay) can be a way to modulate both laser convergence (divergence) and plasma ramping.



# Simulation results: matching and focusing/defocusing



Energy gain is in excess of 200 MeV in a 3 cm acceleration length, which means an average electric field in excess of

**7 GV/m**

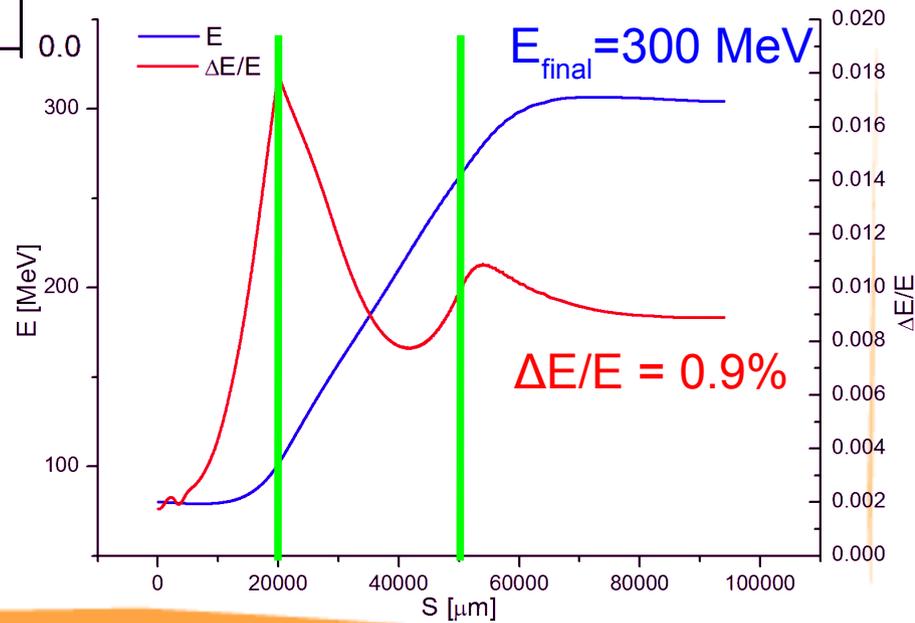
Energy spread is within 1% (a safe threshold for subsequent transport\*) and **REDUCED IN THE FINAL RAMP** due to beam loading and plasma wavelength increase.

\* P. Antici et al., J. App. Phys. 112, 044902 (2012).

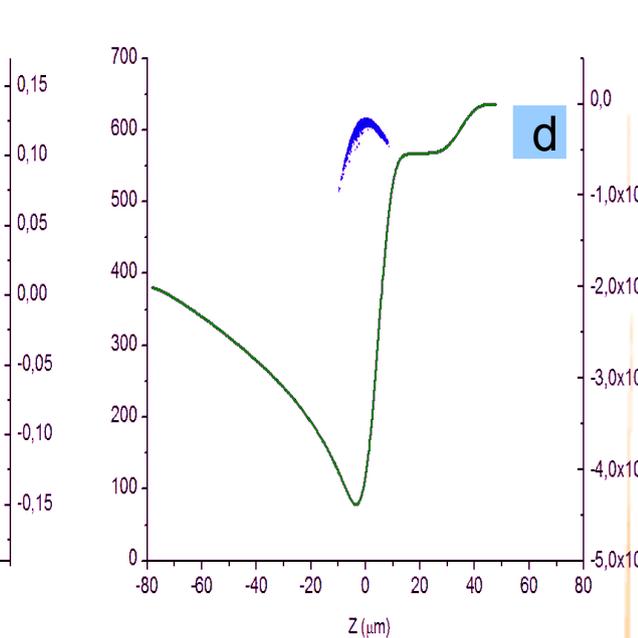
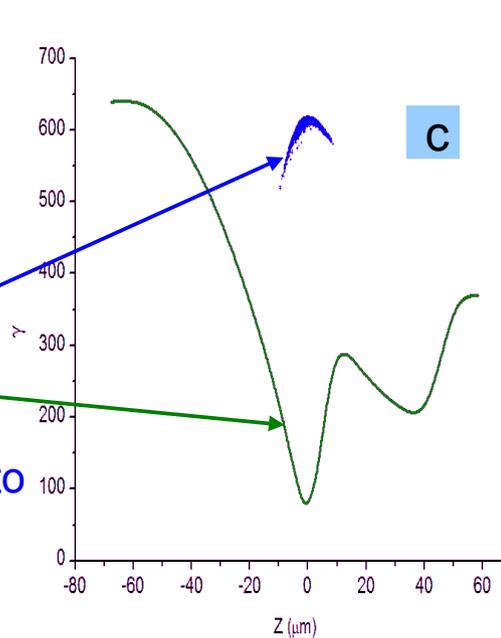
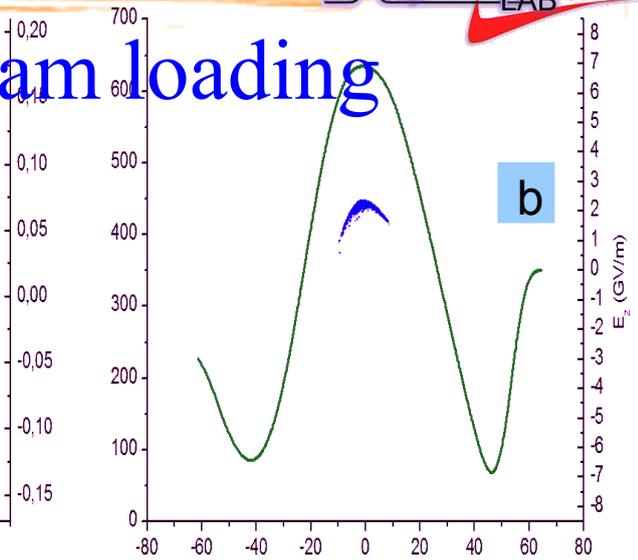
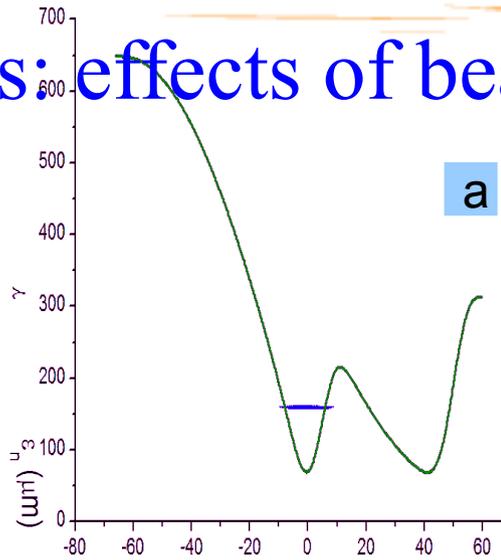
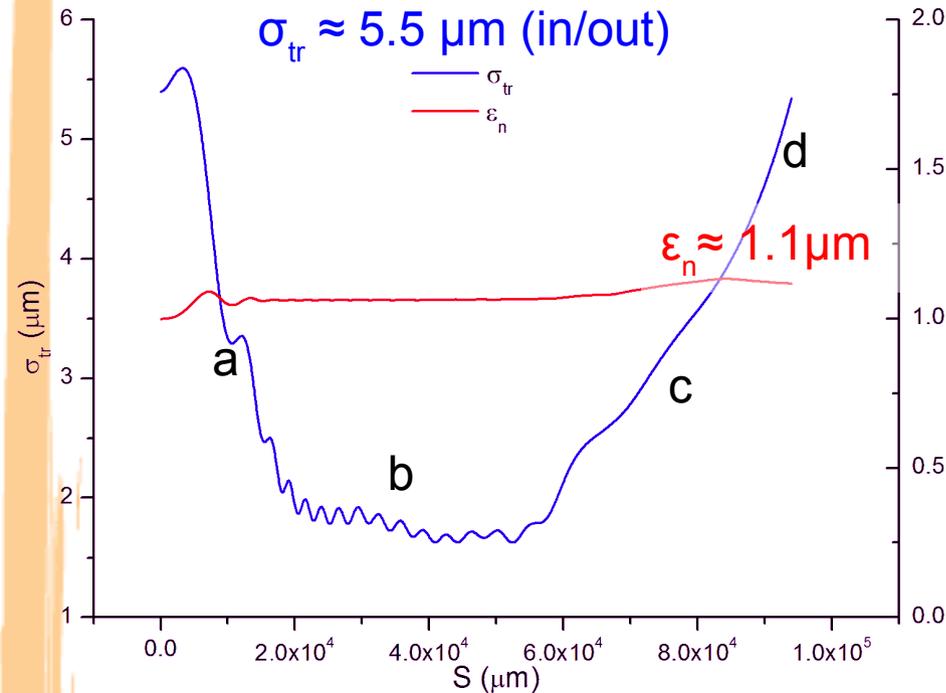
A very nice matching is obtained by laser focusing method.

Beam size is reduced to matched size and then increased back to initial value in a quasi-adiabatic transformation.

Both initial and final sizes seems to be within manageability of permanent magnet quads.



# Simulation results: effects of beam loading



$E_{\text{final}} = 300 \text{ MeV}$  Long. phase space

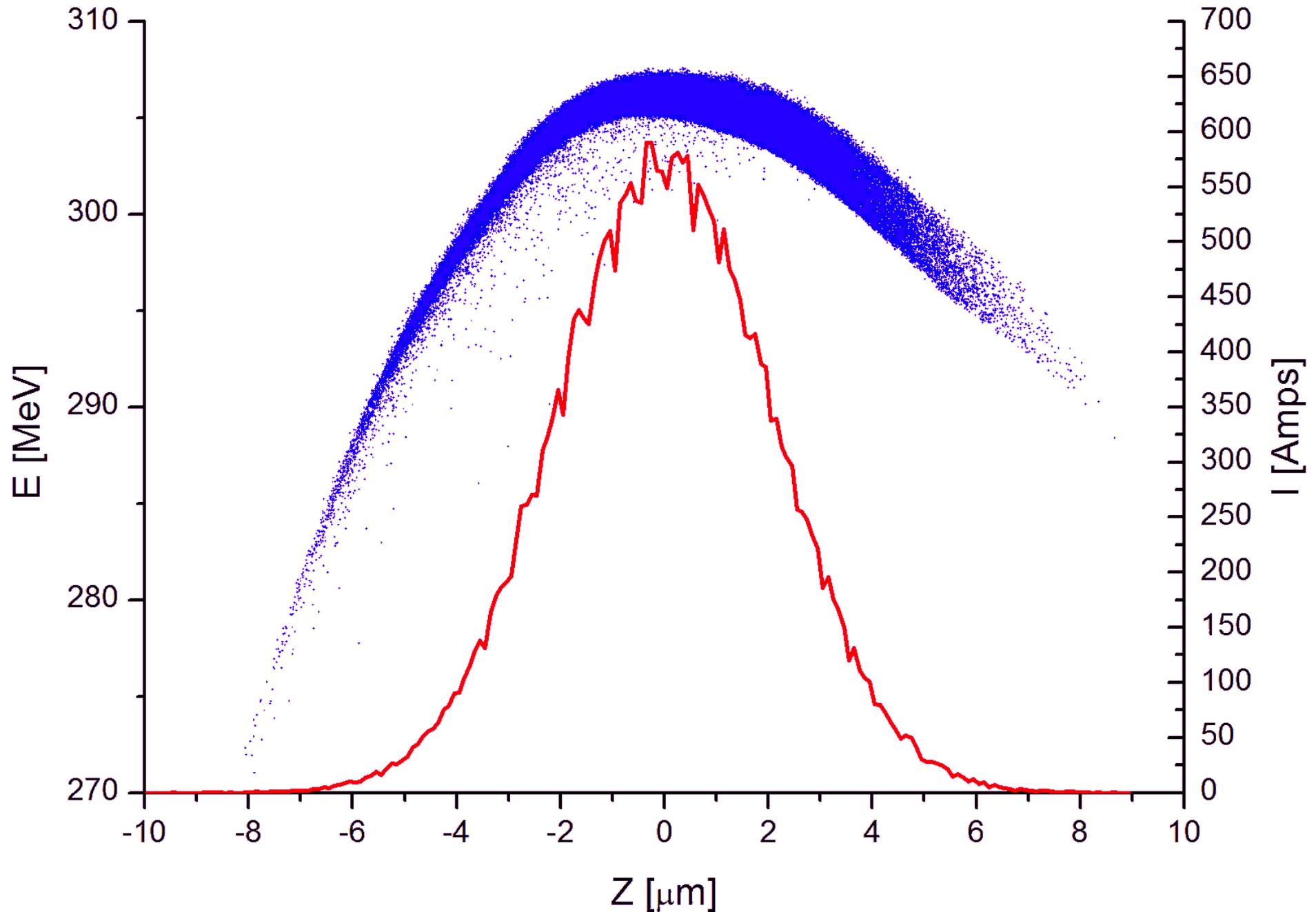
$\Delta E/E < 1\%$

$E_z$

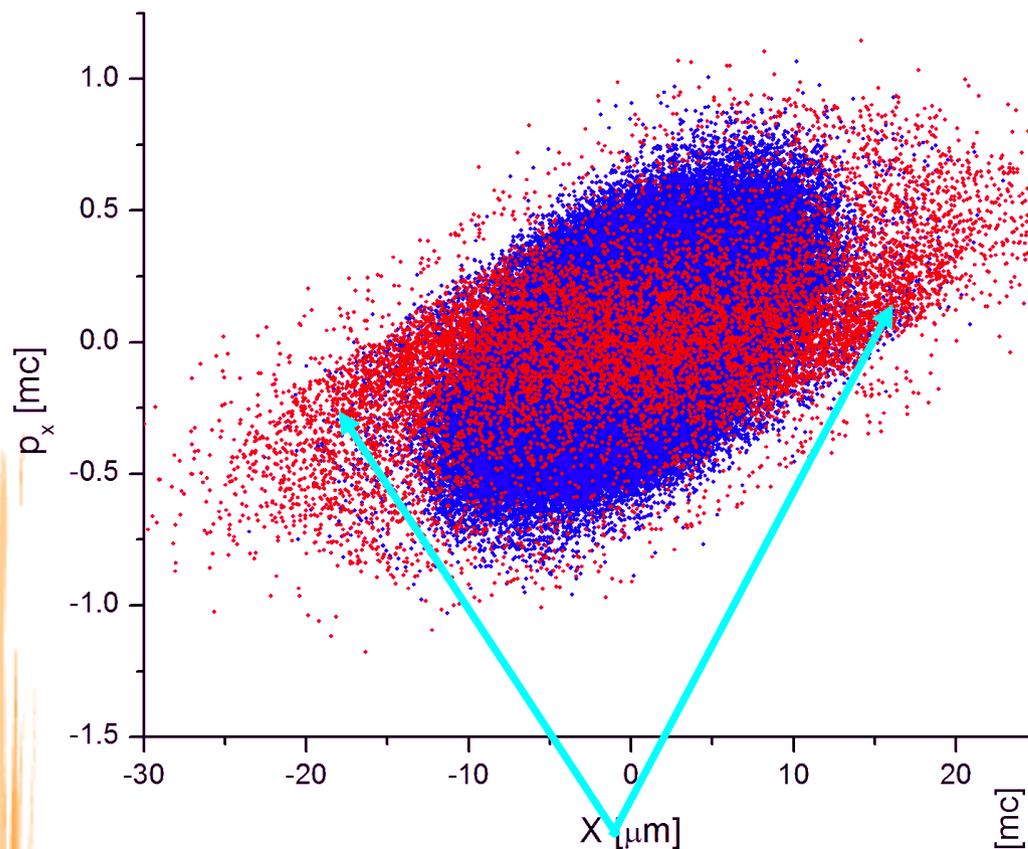
The plasma wave regime goes from linear to quasi linear and back to linear.

Beam loading has a varying impact on beam dynamics and helps to reduce final emittance and energy spread, being relevant when the plasma wavelength is such that the bunch would experience defocusing fields.

## Simulation results: final beam



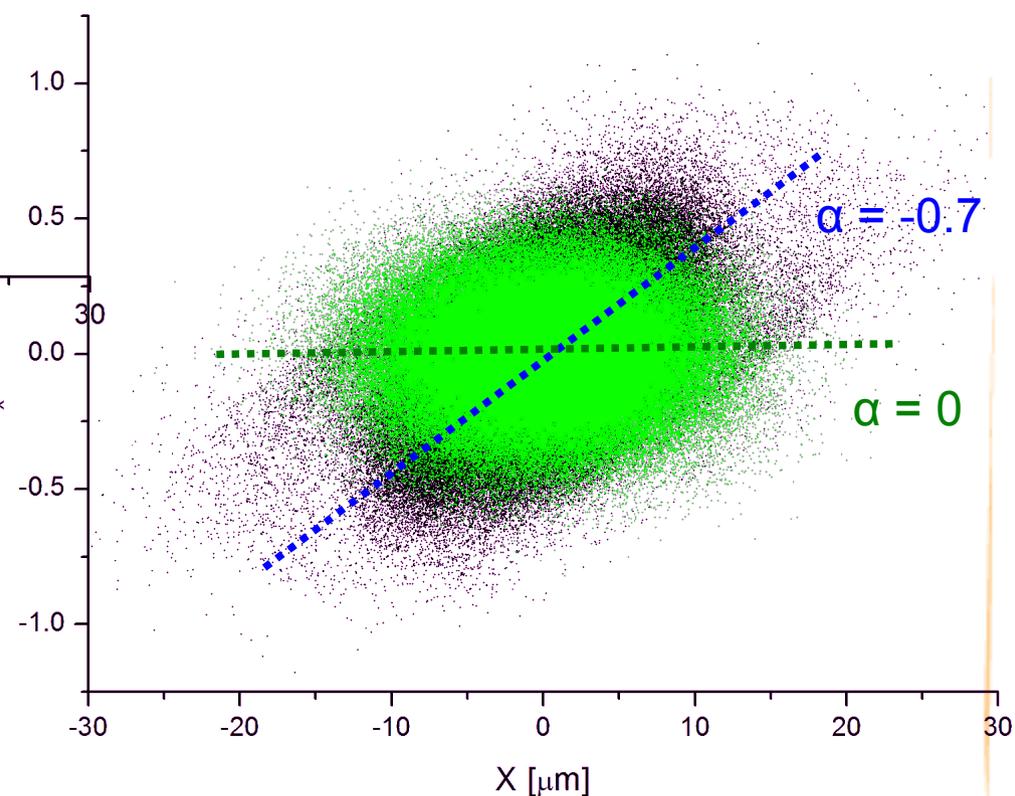
# Simulation results: final beam



A slight non-linearity in transverse phase space builds up since initial focusing, but only 5% of charge is involved, even if it is responsible for most of the emittance increase.

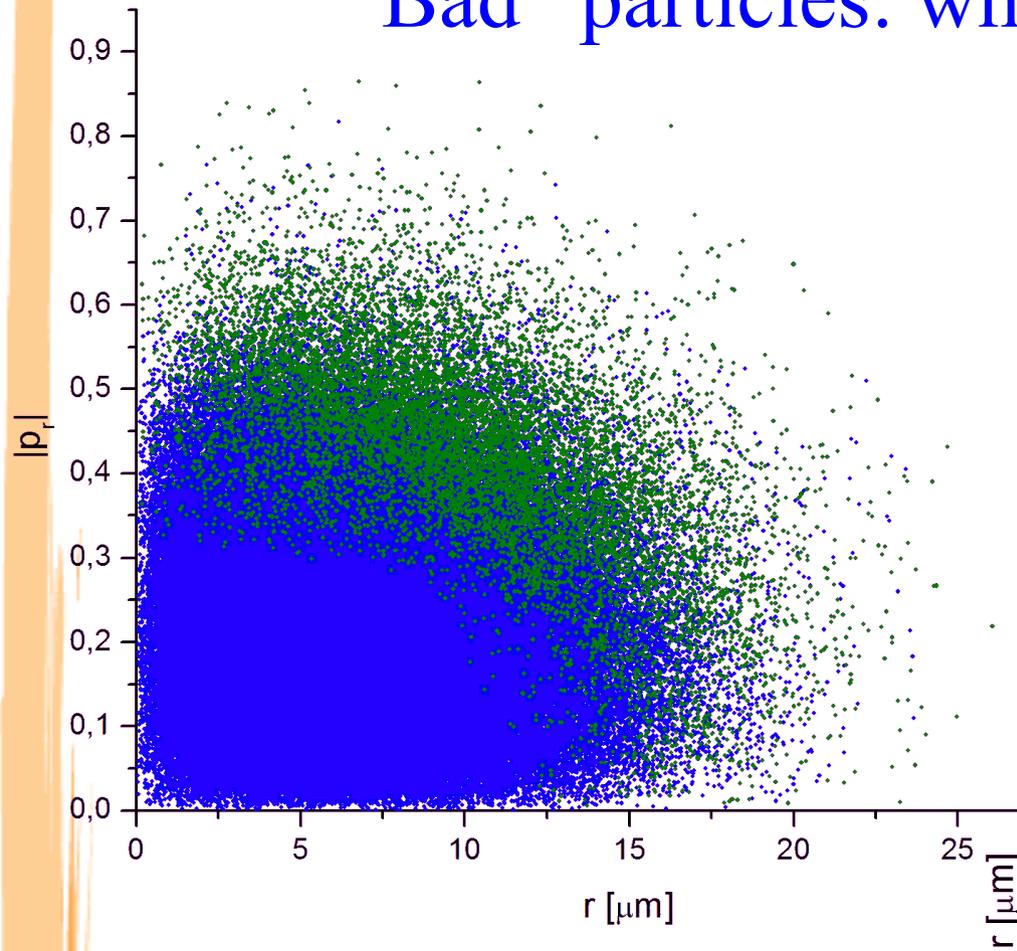
The final Twiss parameter  $\alpha$  has a value of  $-0.7$ . With a little more effort should be brought to  $-1^*$

\* K. Floettmann, PRSTAB 17, 054402 (2014).

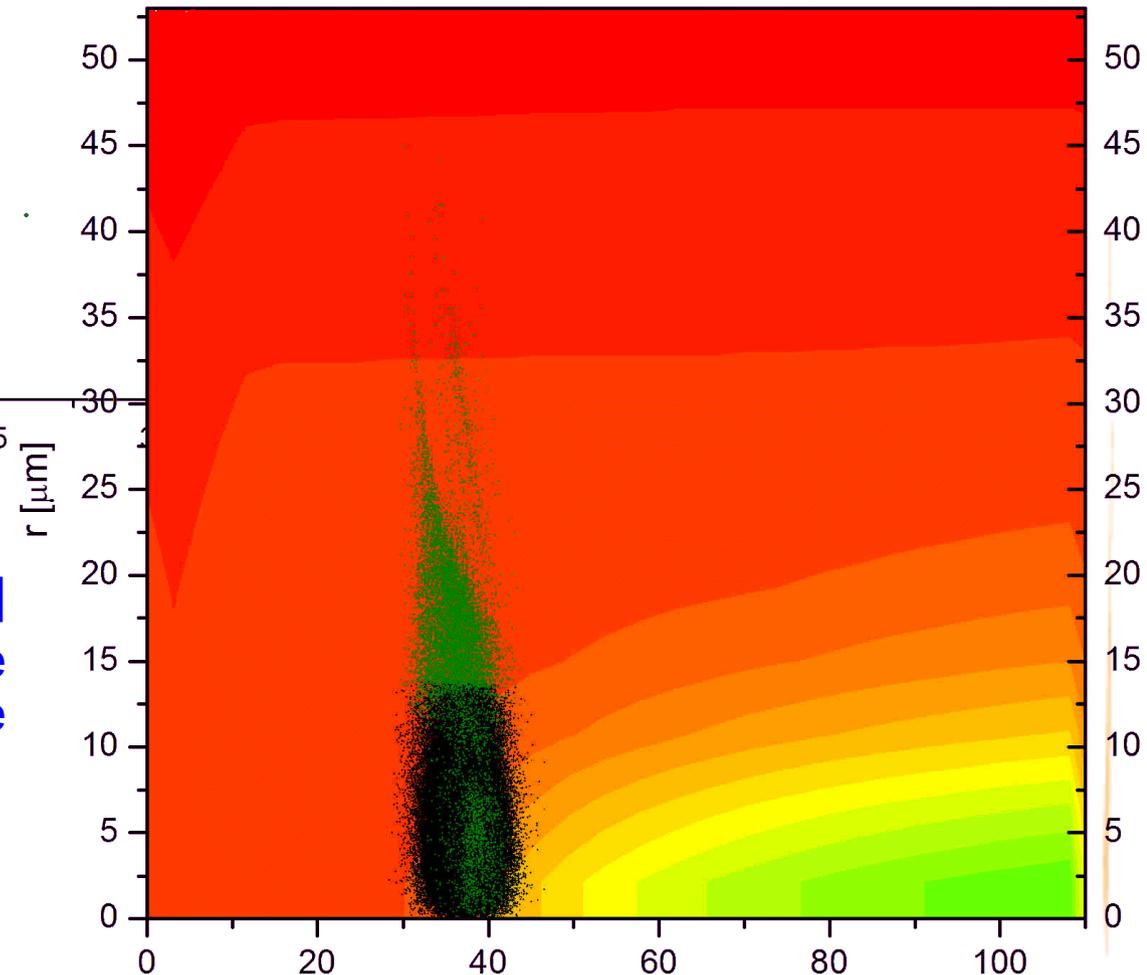


# “Bad” particles: where do they come from?

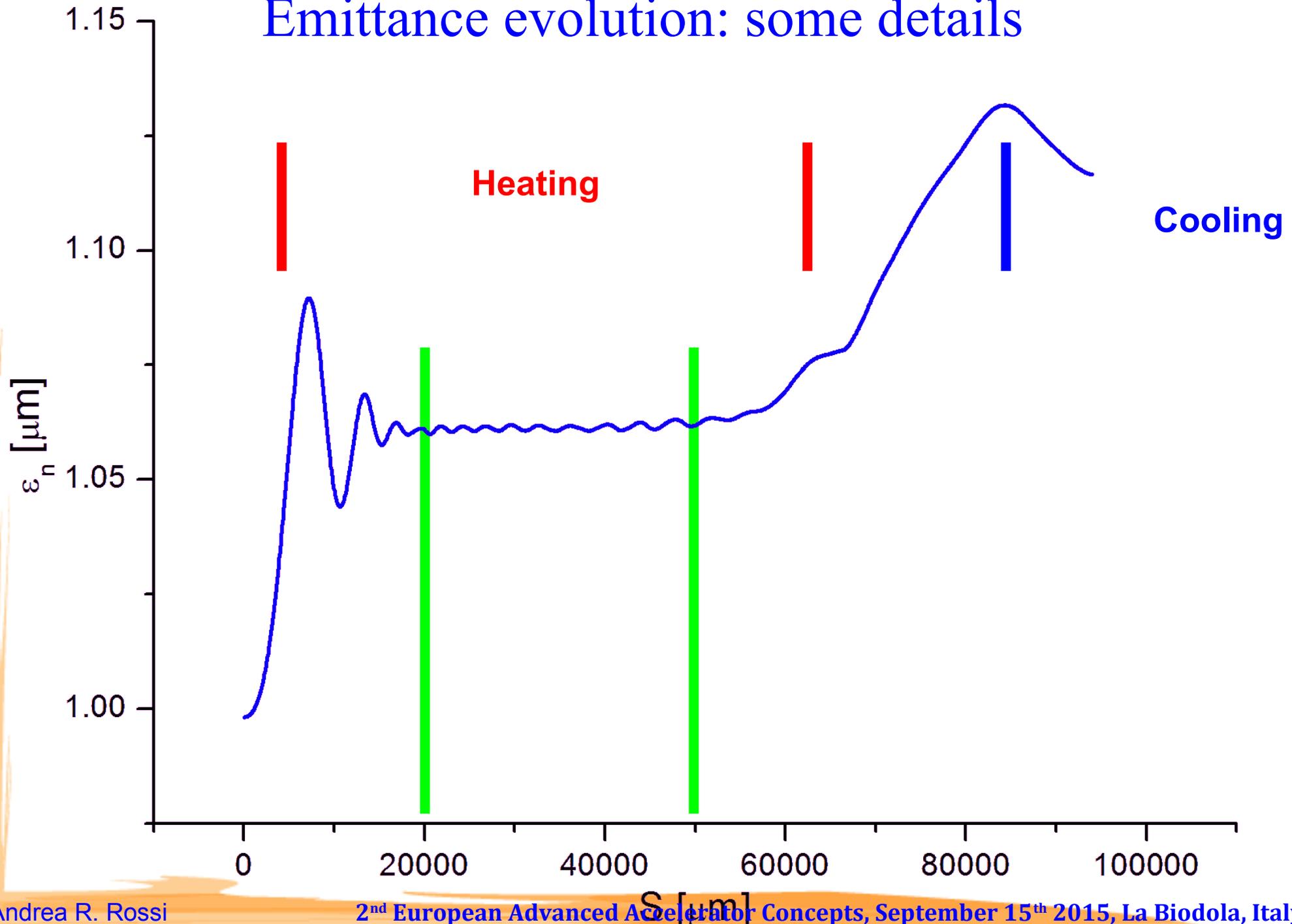
Why are they bad? Because, when beam loading dominates the plasma wave, they are eventually expelled from the (beam-driven) wave and start to drift in the neutral plasma.



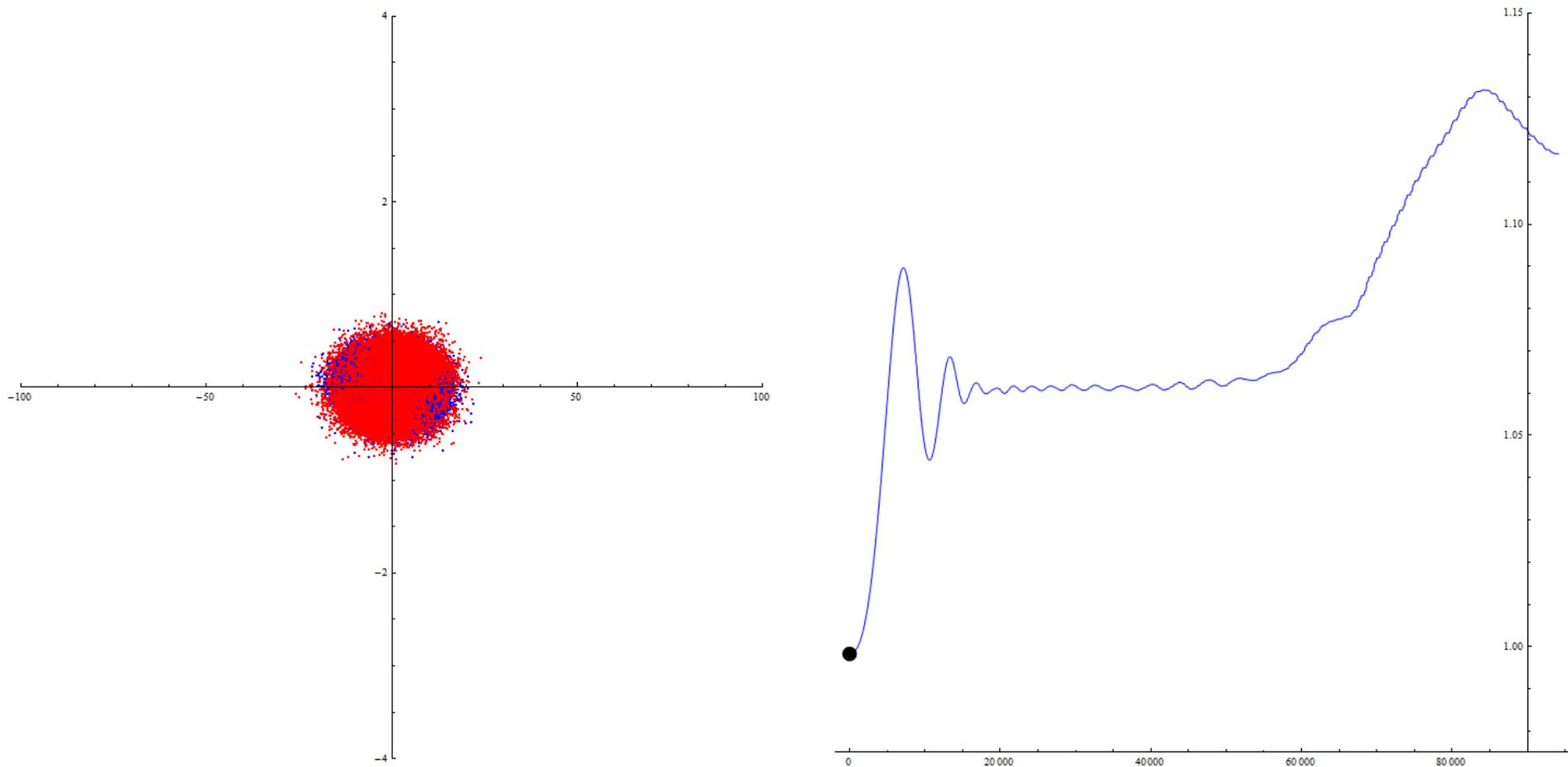
The “bad” particles, in the initial  $r$ - $|p_r|$  phase space, are around the outer edge of the distribution, i.e. large transverse position and/or transverse momentum.



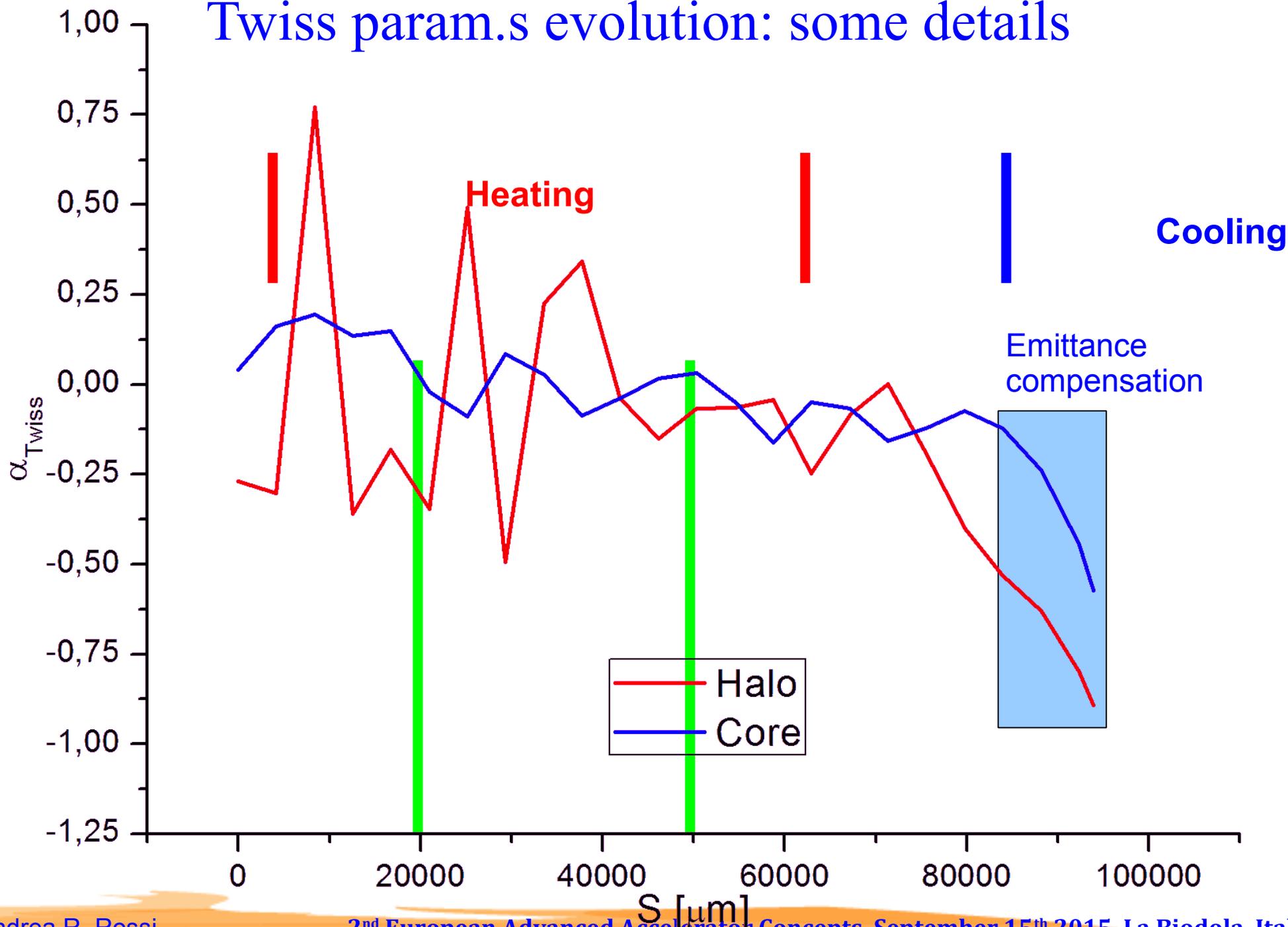
## Emittance evolution: some details



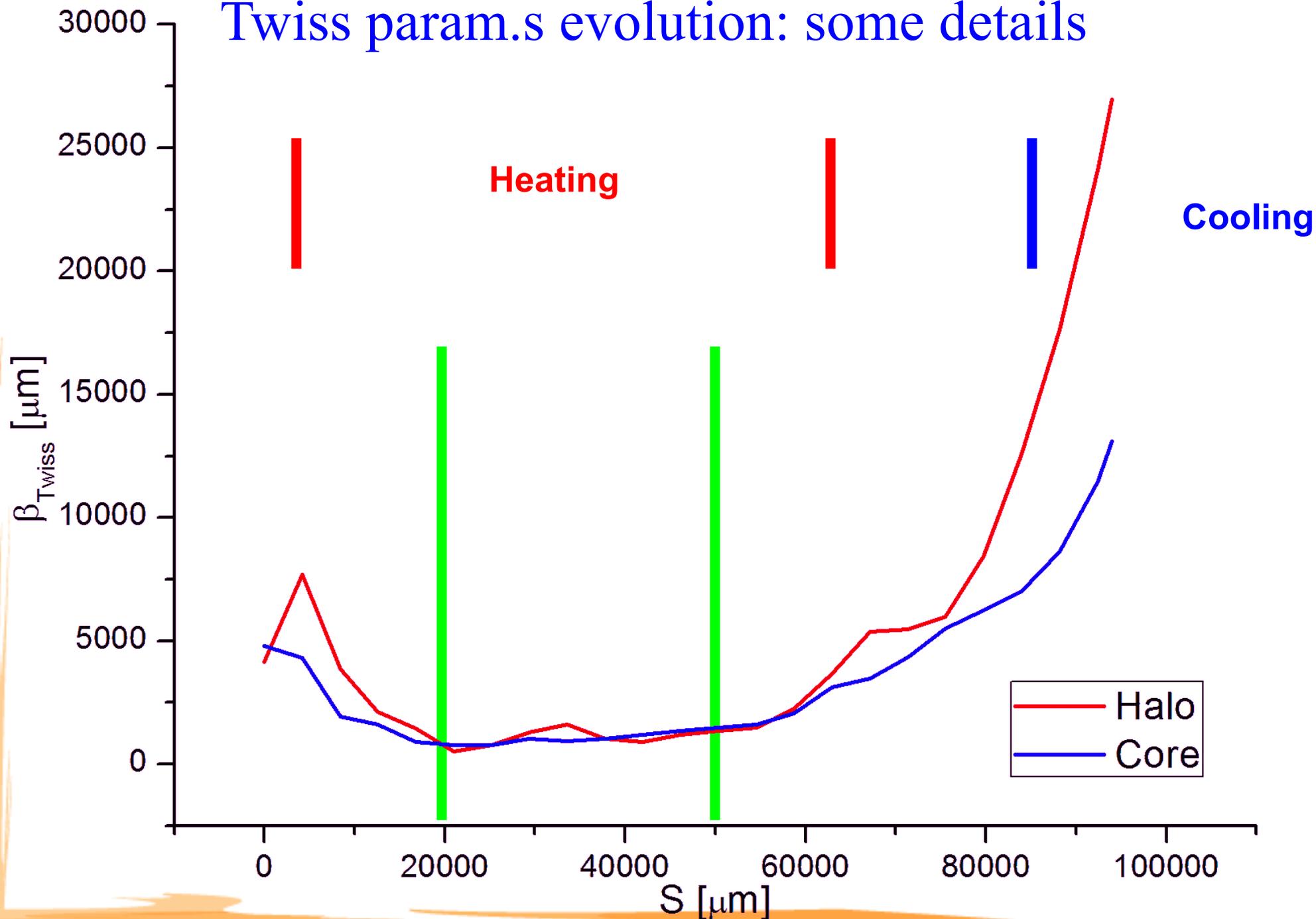
# Emittance evolution: what's going on?



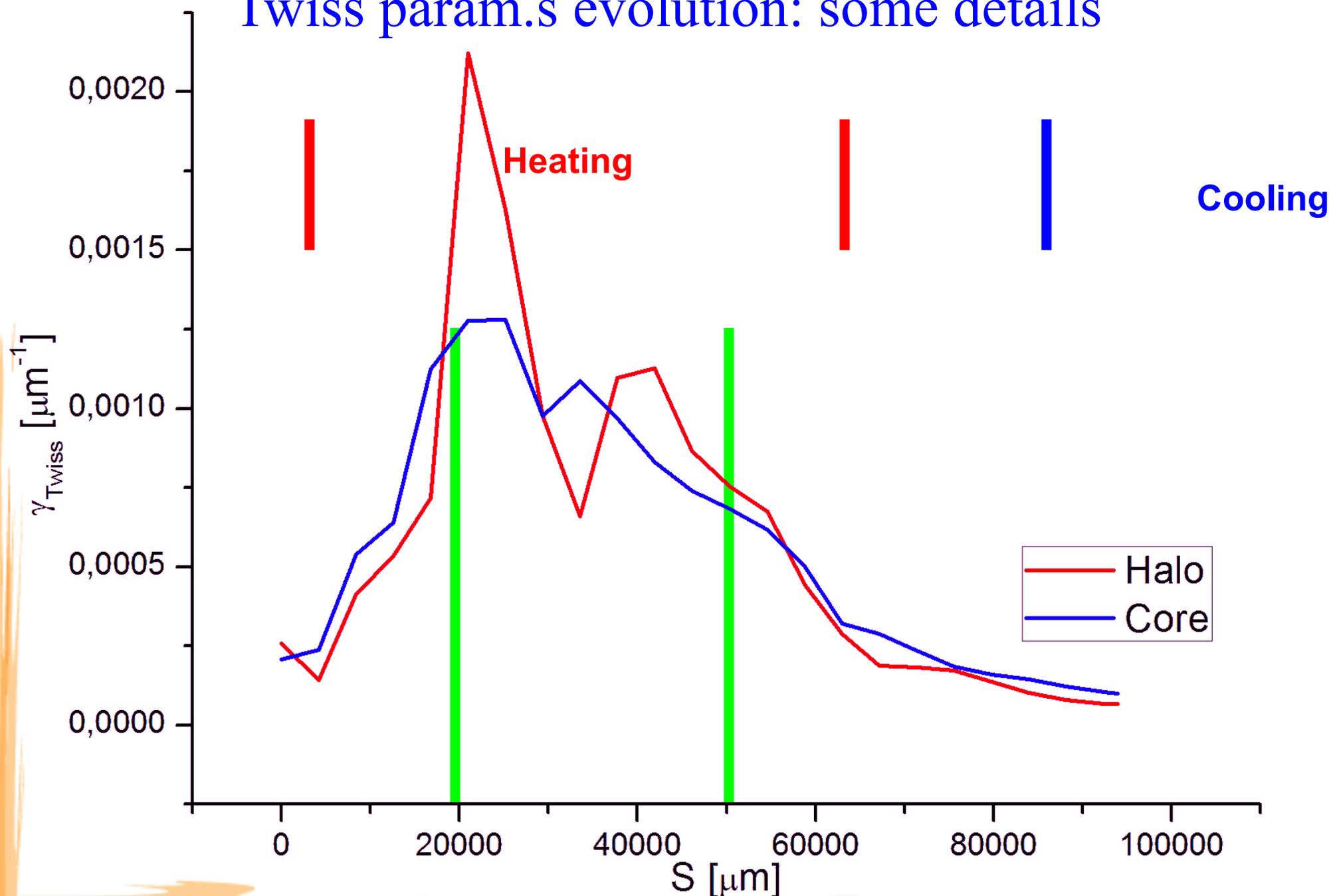
# Twiss param.s evolution: some details



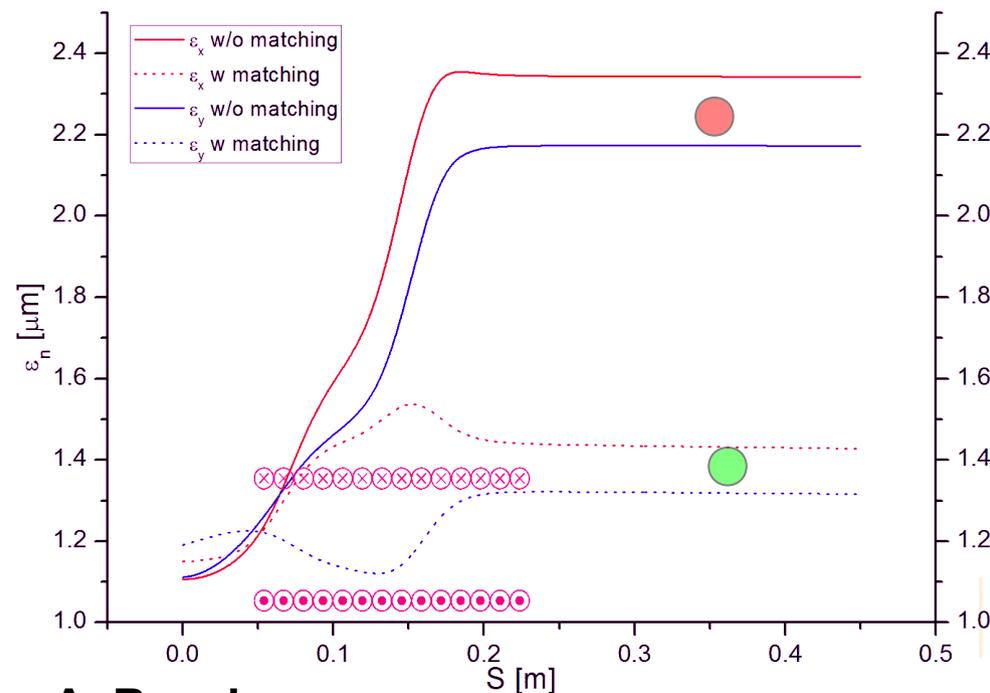
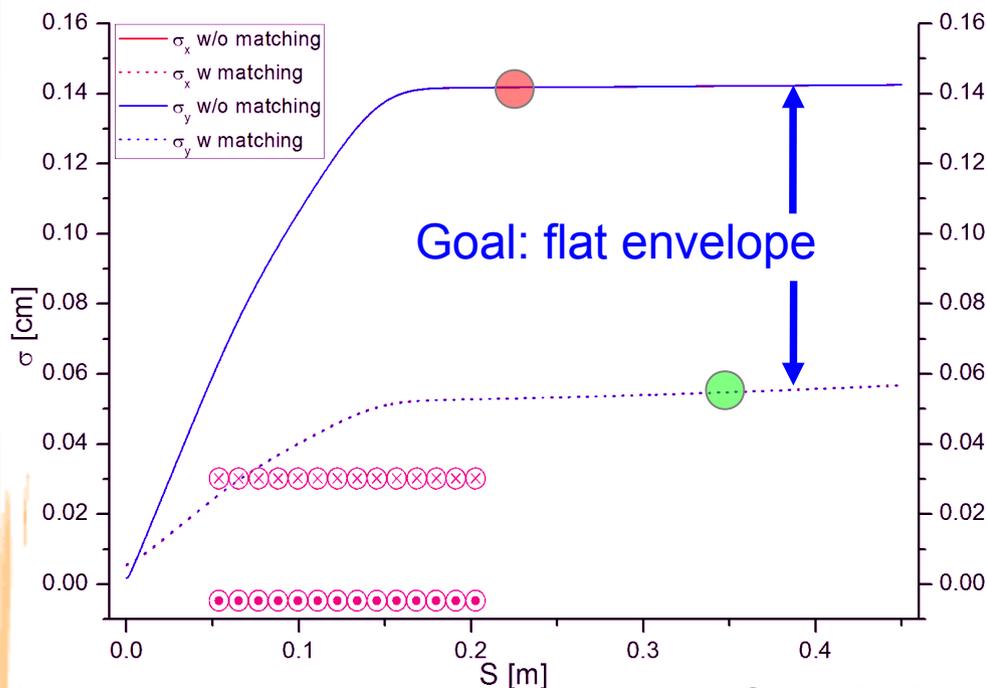
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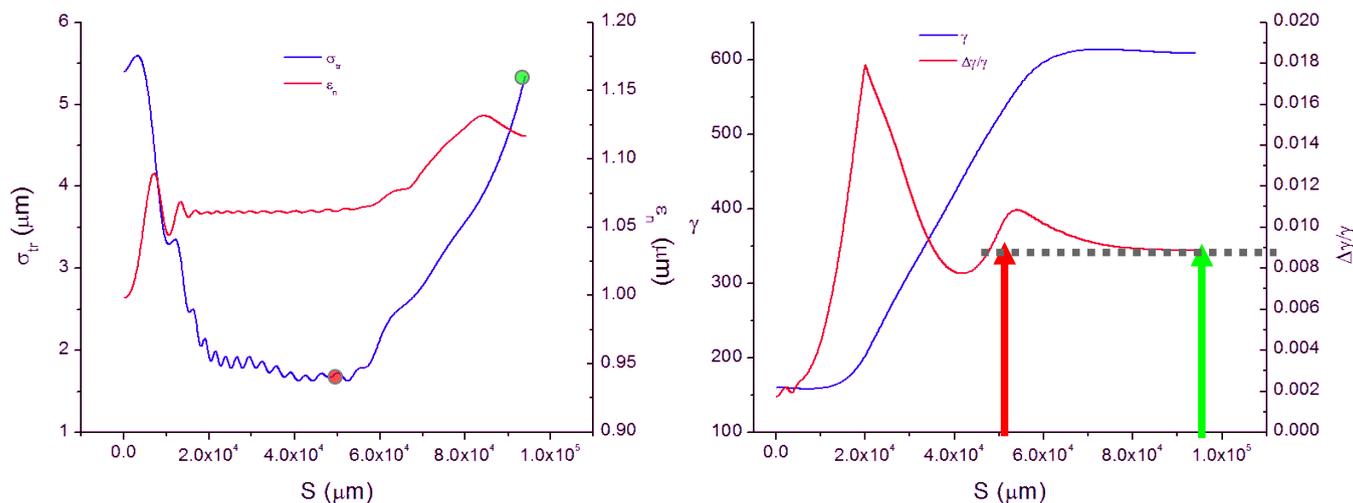
# Simulation results: post acceleration transport



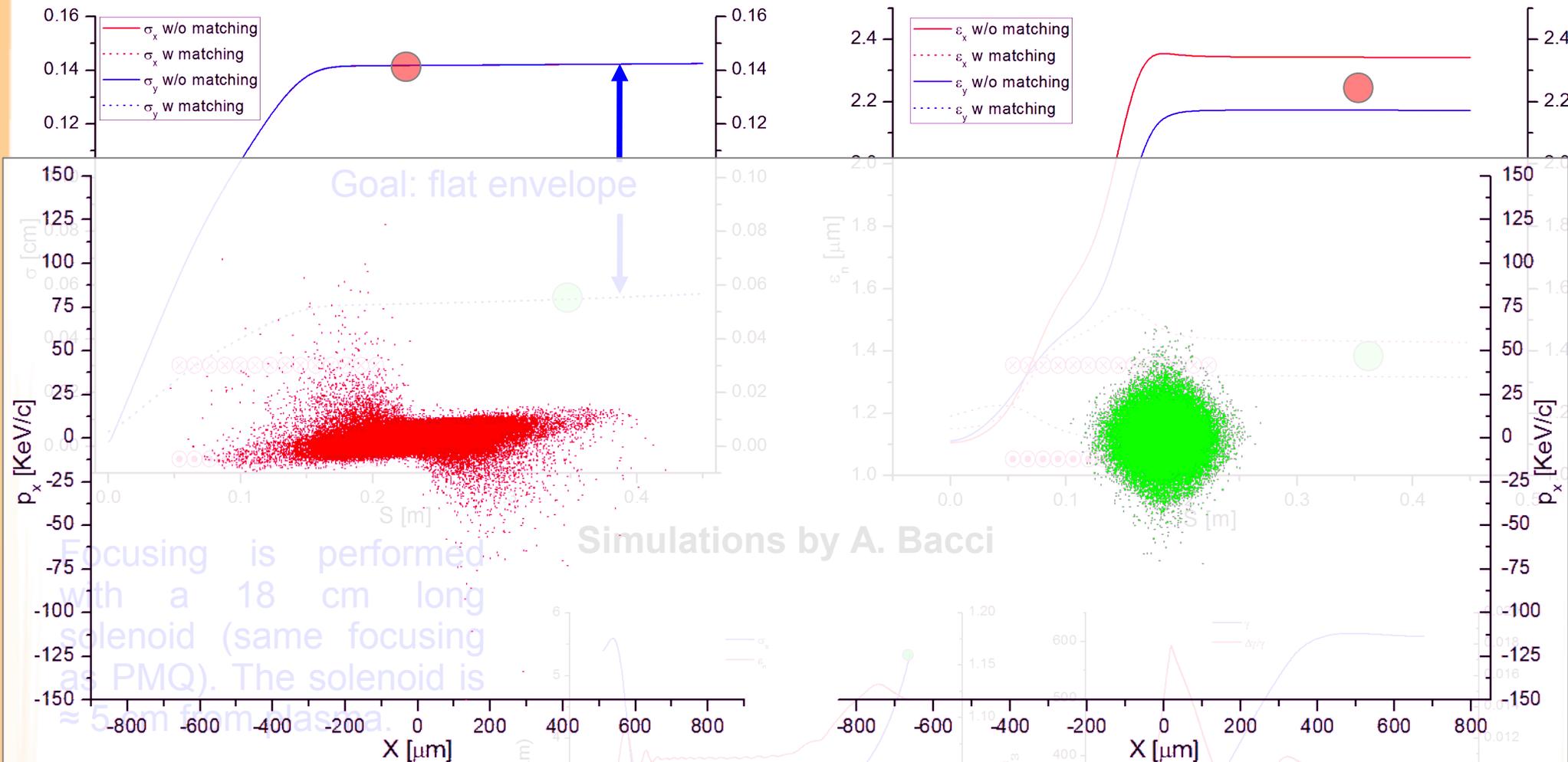
Simulations by A. Bacci

Focusing is performed with a 18 cm long solenoid (same focusing as PMQ). The solenoid is  $\approx 5$  cm from plasma.

Both beams start with about the same energy spread. The **matched** one behaves in a much better way than the **unmatched**

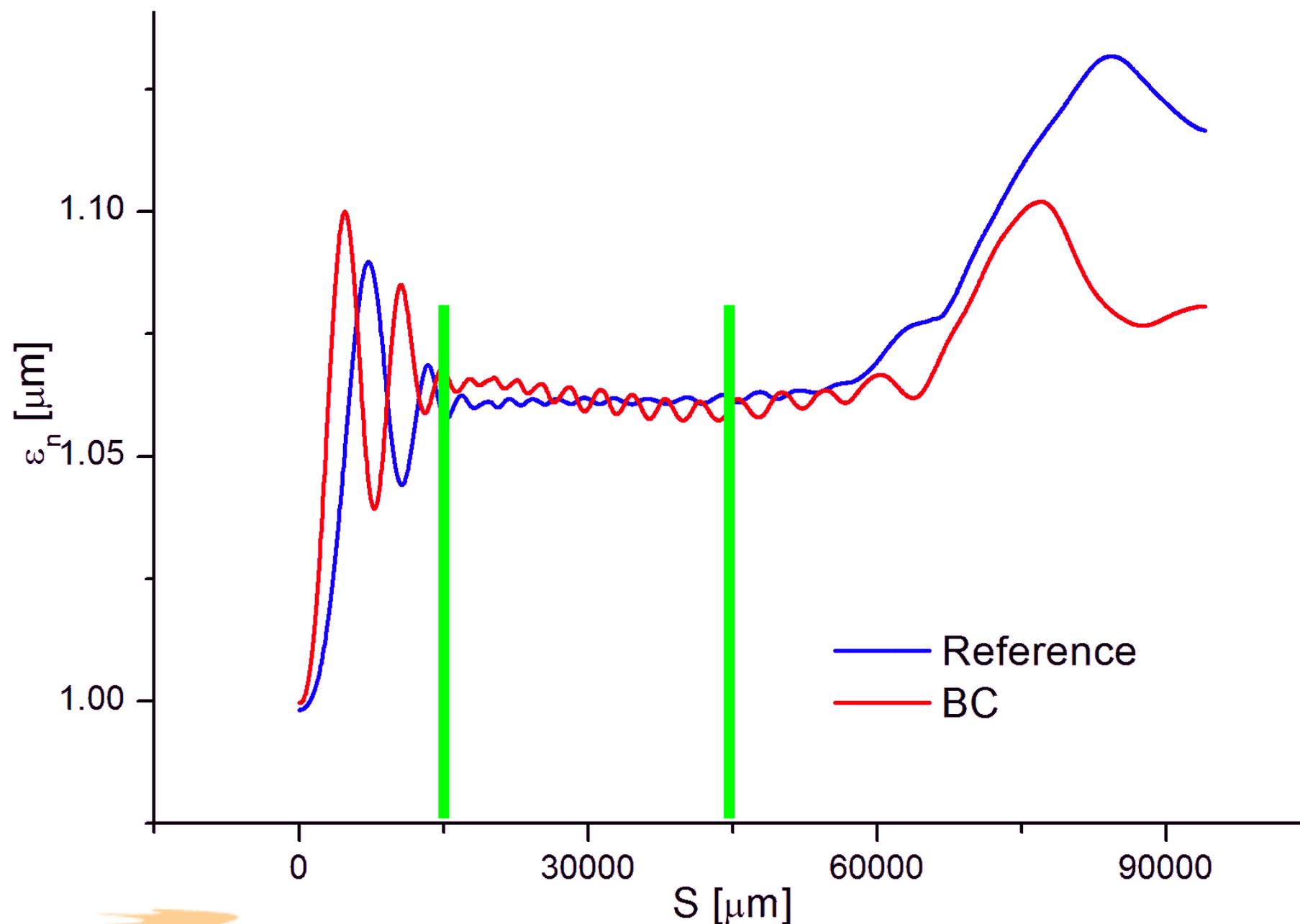


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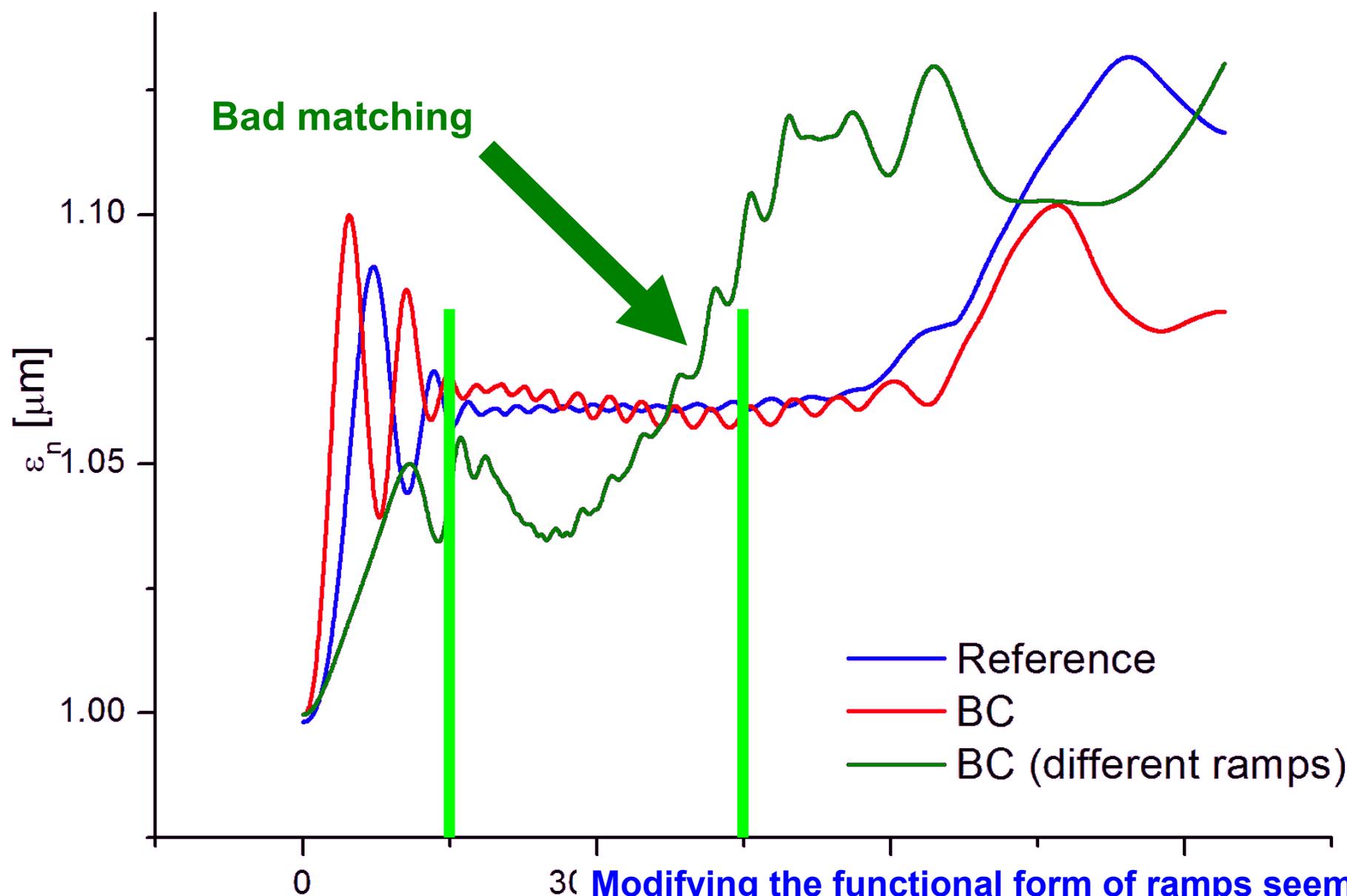


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# Stability: Reference vs Beer Can (BC, *ceteris paribus*)

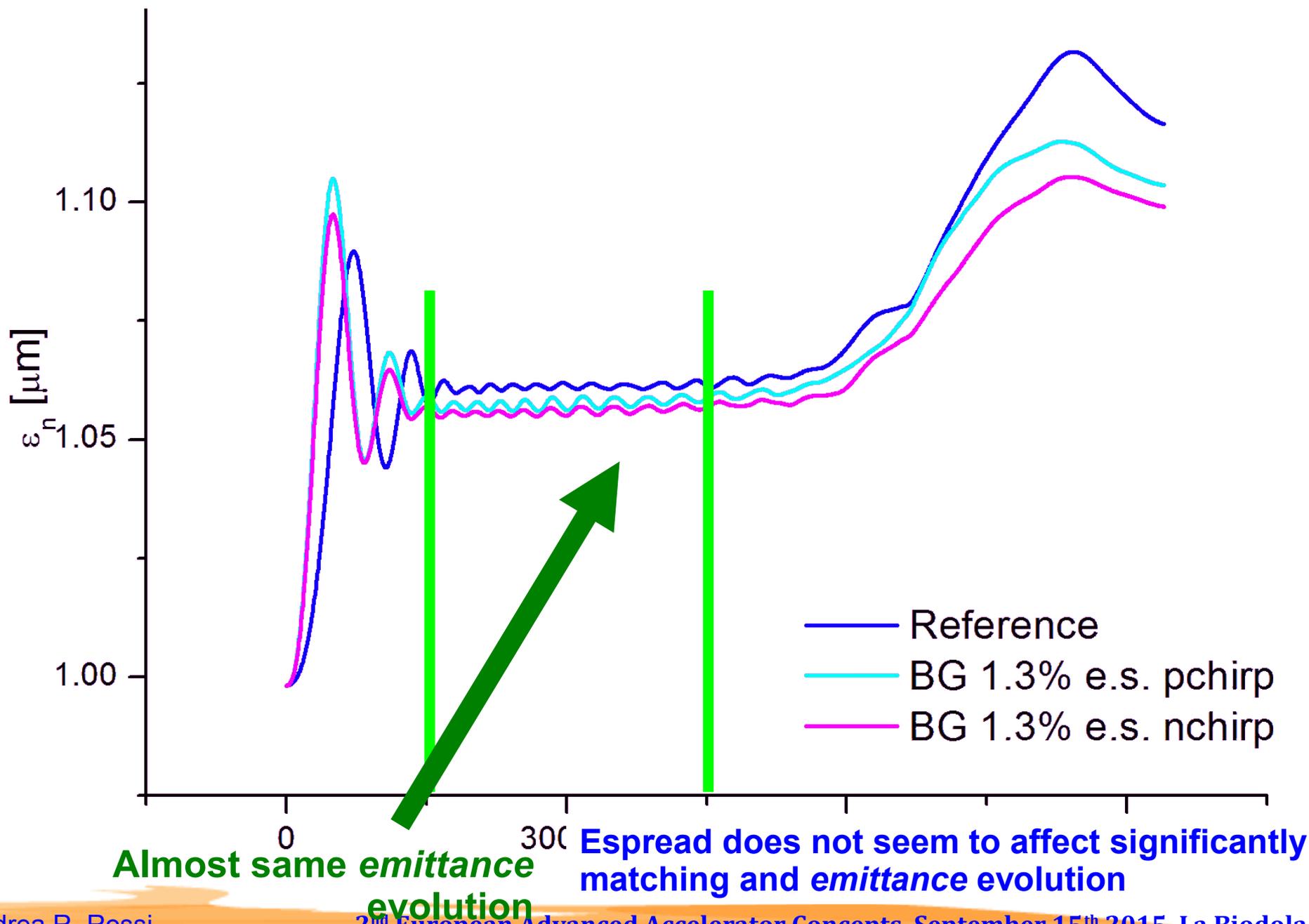


# Stability: Reference vs BC vs BC with different ramps

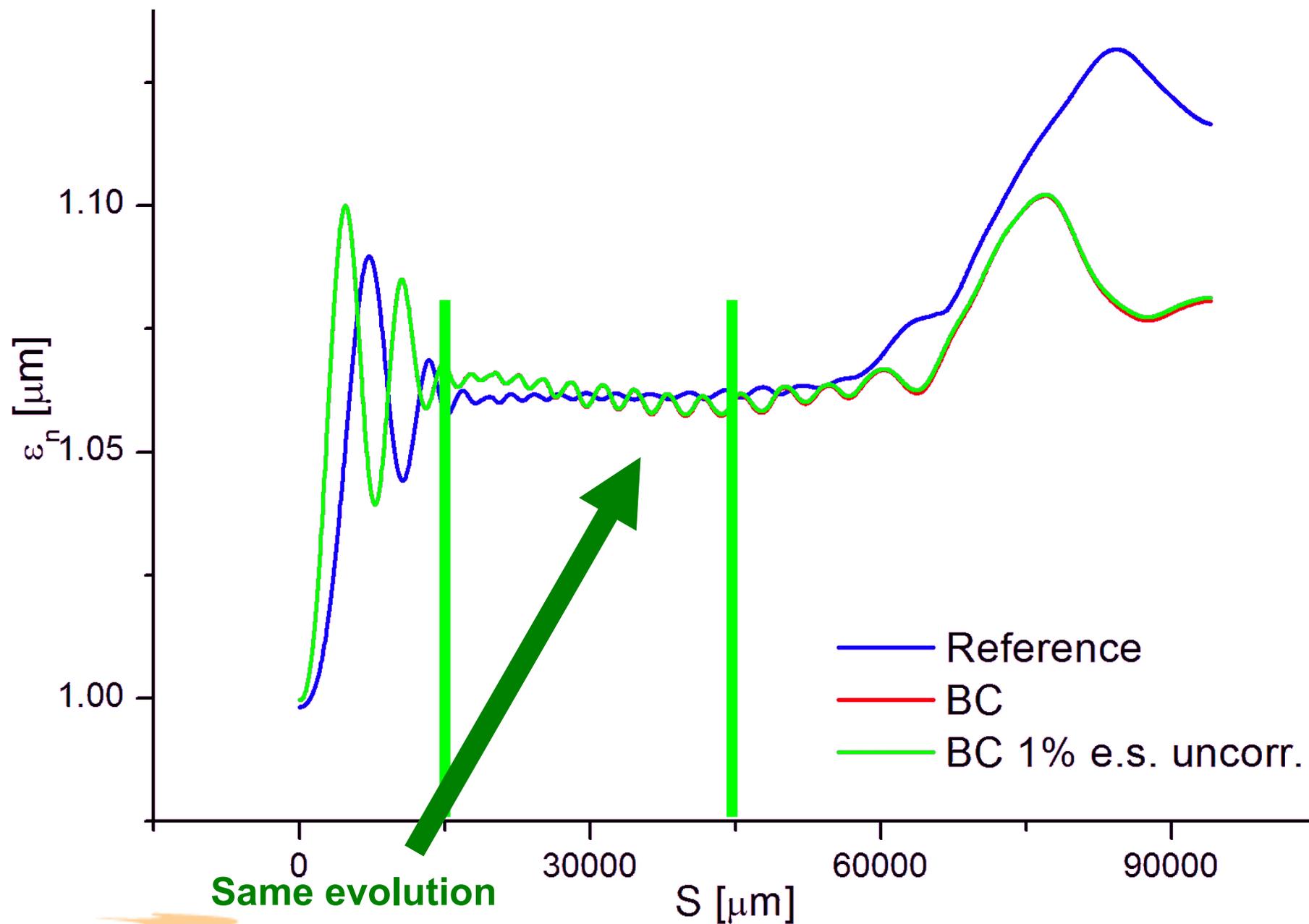


Modifying the functional form of ramps seems to be a way to improve results, as expected

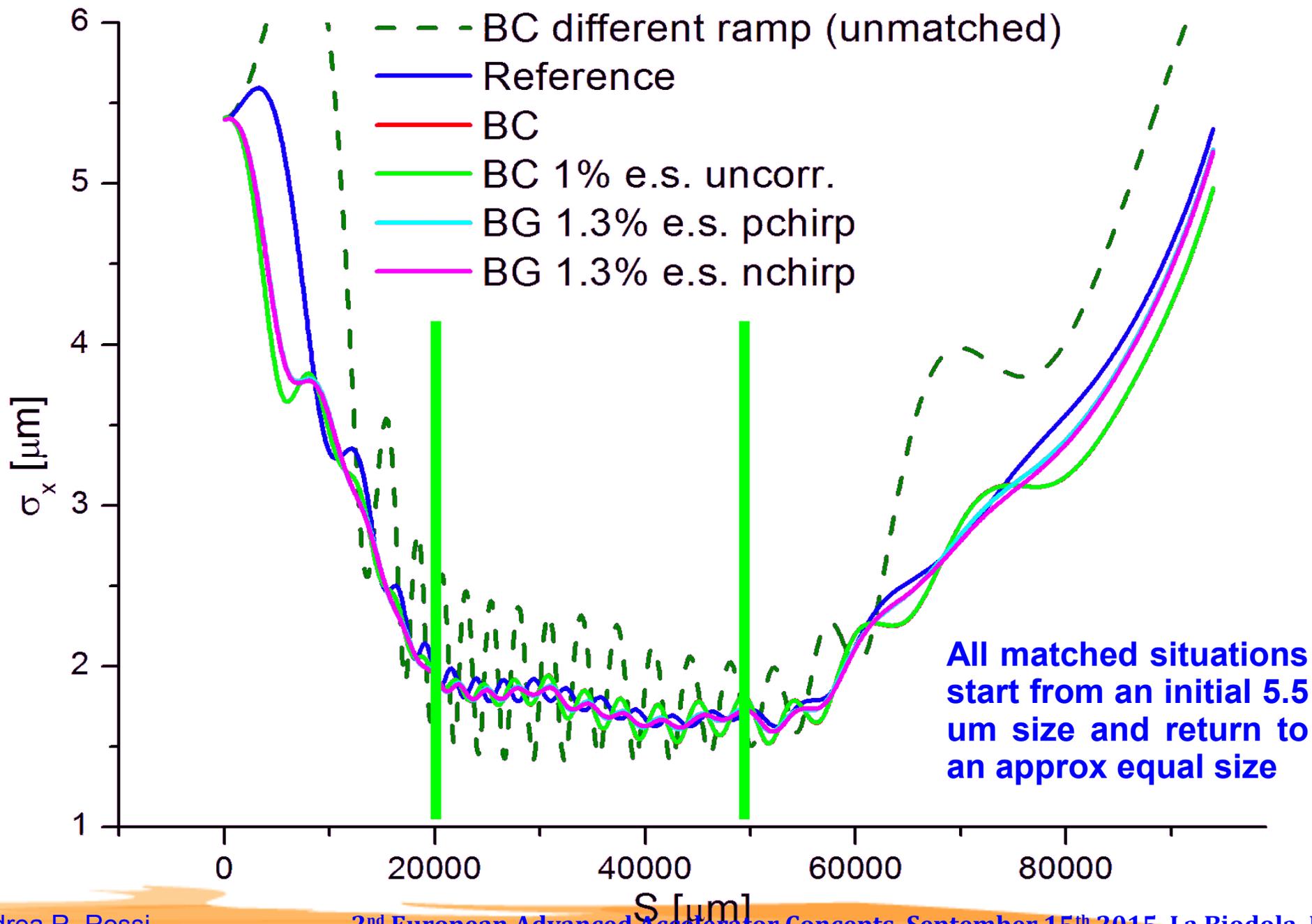
# Stability: Reference vs BG with 1.3% e.s. pchirp/nchirp



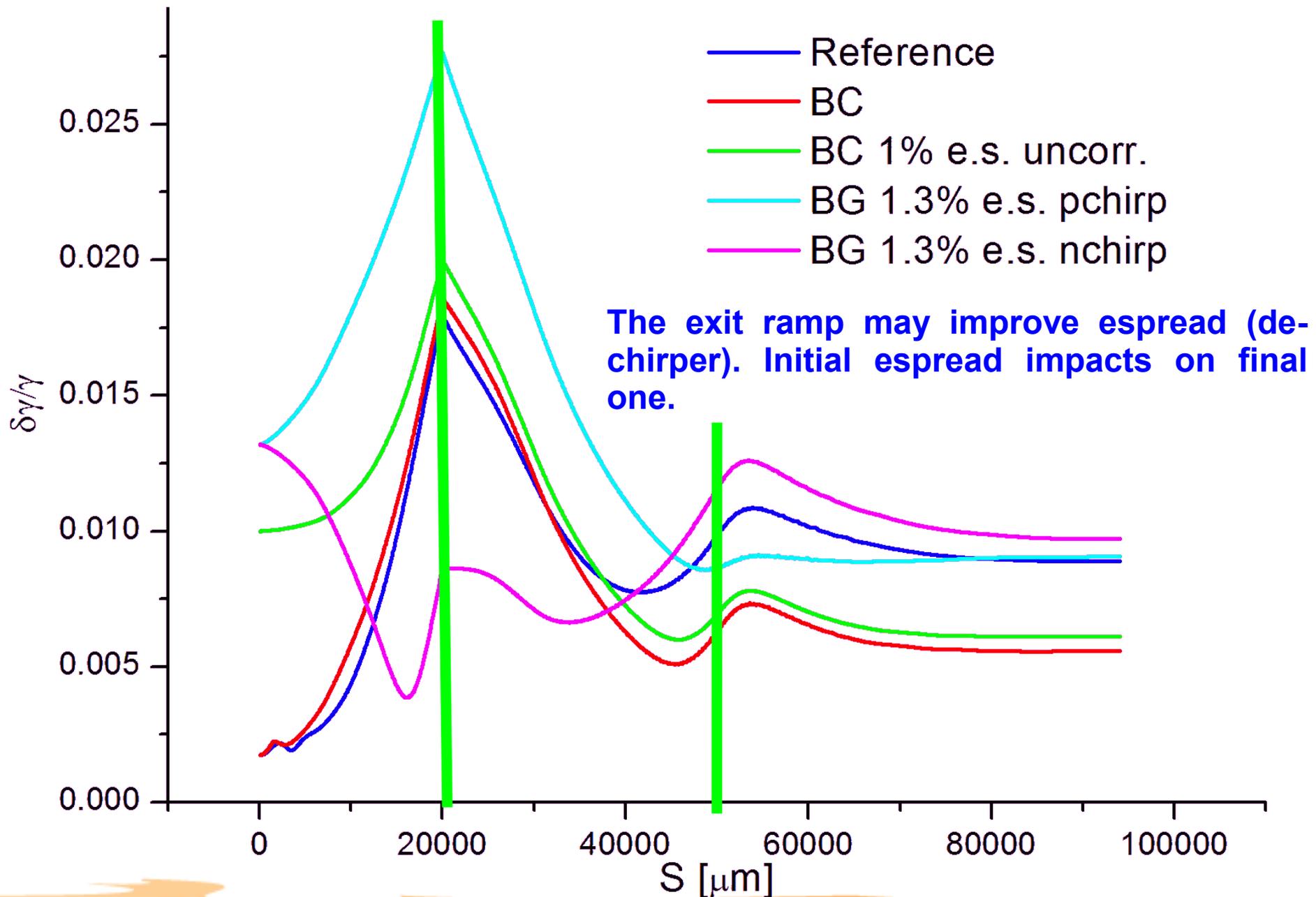
## Stability: Reference vs BC vs BC 1% uncorr. e.s.



# Transverse size evolution



# Espread evolution



# Conclusions

- ◆ We showed a method to match an electron beam in/out a laser driven plasma accelerating stage.
- ◆ The method consist in exploiting both plasma density ramps and laser (de)focusing switching from a beam driven wave to a laser driven wave and back to a beam driven wave; this allows to avoid problems from varying plasma wavelength at density ramps.
- ◆ The method, applied to an ideal beam, showed to be effective and robust against some changes in beam distribution and allows to relax requirements on conventional focusing devices.
- ◆ The whole process can be further optimized by changing splay angles and/or shapes.
- ◆ The effect of the splays and their shape on plasma ramps profiles must be checked by dedicated simulations.

Thanks for your  
attention!

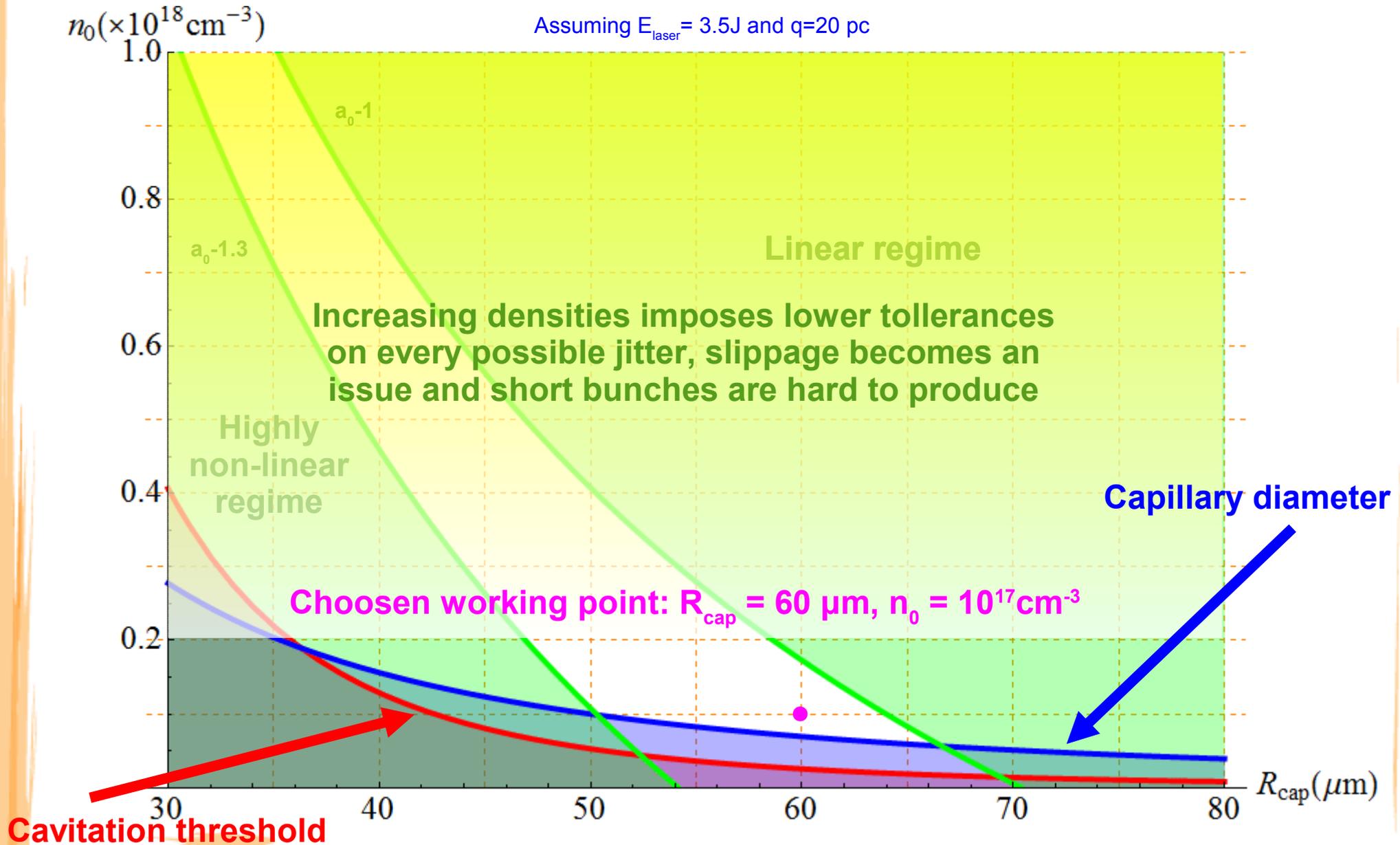
# Backup slides

# S2E simulation: plasma acceleration

Simulation tool: QFLUID2 by P. Tomassini and A.R. Rossi

- 2D cylindrical.
- Fluid approximation for plasma.
- Bunch macroparticles are fully kinetic.
- Supports quasi linear regimes.
- Laser evolution is self consistent and uses envelope approximation.
- Beam loading effects are included.

# Choice of parameters: physical and practical constraints



# Random particle trajectory

