

Beam dynamics in resonant plasma wakefield acceleration @SPARC_LAB

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on behalf of SPARC_LAB collaboration











SPARC Layout

In response to the necessity of more compact accelerating structures one of the most promising ways is the plasma wakefield acceleration beam driven

Requirements: High accelerating field and <u>beam quality</u> (Energy spread, emittance)



Beam Driven Plasma Acceleration @SPARC_LAB

- > 1 Driver (High Charge) + Witness
- Resonant wakefield acceleration:
 N Drivers (Lower Charge)+ Witness
- Plasma driven FEL

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GOAL

- Matching conditions
 - Revision of longitudinal matching for resonant scheme

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Quasi Non Linear Regime



- Low charge high density bunches
- $\tilde{Q} = \frac{N_b k_p^3}{n_0} < 1$ $\tilde{A} = \frac{n_b}{n_0} > 1$
- Linear fields+resonant plasma response

Rosenzweig, J. B., et al. "Plasma wakefields in the quasi-nonlinear regime: Experiments at ATF." *ADVANCED ACCELERATOR CONCEPTS: 15th Advanced Accelerator Concepts Workshop*. Vol. 1507. No. 1. AIP Publishing, 2012.

Londrillo, P., C. Gatti, and M. Ferrario. "Numerical investigation of beam-driven PWFA in quasi-nonlinear regime." *Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment* 740 (2014): 236-241.

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Resonant Wakefield Acceleration



- Accelerating field increases with the number of drivers
- Different longitudinal matching for drivers following the first one

Hogan, M. J., et al. "Plasma wakefield acceleration experiments at FACET."*New Journal of Physics* 12.5 (2010): 055030.

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Hybrid Tool: Architect

- Bunch(es) are treated kinetically
- background plasma as a fluid
- systematic scan in no-time
- 1cm/1 hour!

Equations

$$d_{t}\mathbf{p}_{\text{particle}} = q(\mathbf{E} + c\boldsymbol{\beta}_{\text{particle}} \times \mathbf{B})$$

$$d_{t}\mathbf{x}_{\text{particle}} = \boldsymbol{\beta}_{\text{particle}}c$$

$$\partial_{t}n_{e} = -\nabla \cdot (\boldsymbol{\beta}_{e}c n_{e})$$

$$\partial_{t}\mathbf{p}_{e} = -\nabla \cdot (\mathbf{p}_{e} \otimes \boldsymbol{\beta}_{e}c) + q(\mathbf{E} + c\boldsymbol{\beta}_{e} \times \mathbf{B})$$

$$\partial_{t}\mathbf{B} = -\nabla \times \mathbf{E}$$

$$\partial_{t}\mathbf{E} = c^{2}\nabla \times \mathbf{B} - q\mu_{0}c^{3} (n_{e}\boldsymbol{\beta}_{e} + n_{b}\boldsymbol{\beta}_{b})$$



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Beam Parameters

SIMULATED BUNCH PARAMETERS



MEASURED DRIVER BUNCH PARAMETERS

 $\succ \epsilon_r \cong 4 \ mm \ mrad \quad \succ \sigma_z = 30 - 60 \ \mu m$ \succ $E \cong 114 MeV$ $\succ \sigma_E \cong 0.5\%$

 $\blacktriangleright Q \cong 200 - 210 pC$

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Longitudinal Matching Conditions for 1st Driver

$$\succ k_p \sigma_z \cong \sqrt{2}$$

Joshi, C., et al. "High energy density plasma science with an ultrarelativistic electron beam." *Physics of Plasmas (1994-present)* 9.5 (2002): 1845-1855.
Lu, W., et al. "Limits of linear plasma wakefield theory for electron or positron beams." *Physics of Plasmas (1994-present)* 12.6 (2005): 063101.

Theoretical result for linear theory that grants the best coupling between the beam and the plasma

$$\succ$$
 In our case $\sigma_z \cong 75 \mu m$

- We perform a scan keeping all the bunch characteristics constant except the bunch length in order to find optimal conditions in our quasi non linear regime with \approx 2cm accelerating length
- As a figure of merit for longitudinal matching we choose the peak accelerating field inside the first bubble
- We look for a matching that grants the highest field that keeps constant during all accelerating length

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Single Driver Peak Field vs.Accelerating LengthPre ionized plasma $n = 10^{16} cm^{-3}$



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Best Longitudinal Conditions with Short Accelerating Length



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Head Erosion Mitigation by Longitudinal Matching

- Head erosion effects lead to a lack in accelerating field
- Longitudinal matching allows to mitigate head erosion consequences
- Most stable fields are guaranteed for $k_p \sigma_z \cong \sqrt{2}$



Blumenfeld, Ian, et al. "Energy doubling of 42 GeV electrons in a metre-scale plasma wakefield accelerator." *Nature* 445.7129 (2007): 741-744.

Zhou, M., et al. "Beam head erosion in self-ionized plasma wakefield accelerators." *Particle Accelerator Conference, 2007. PAC. IEEE*. IEEE, 2007.

An, W., et al. "Strategies for mitigating the ionization-induced beam head erosion problem in an electronbeam-driven plasma wakefield accelerator. "*Physical Review Special Topics-Accelerators and Beams* 16.10 (2013): 101301.

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Longitudinal Matching

 $k_p\sigma_z < \sqrt{2}$

The beam is best matched when the driver fills completely the first bubble

$$\succ k_p \sigma_z \cong \sqrt{2}$$

With short accelerating lengths longitudinal matching conditions are more relaxed





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Longitudinal Matching Conditions for 2nd Driver

- → $\sigma_z \cong 60 75 \mu m$ -> Driver tails on witness bunch
- Lower second bubble quality
- The presence of the field caused by the first driver suggest different longitudinal matching conditions for the 2nd driver
- We perform the same scan we performed for the first driver in order to look for possible advantages using resonant scheme
- > We plot the **Maximum accelerating field in the second bubble** vs. accelerating length
- During the scan, the 1st driver characteristics and the injection distance are kept constant

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Peak Field with 2 drivers



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Conclusions

- Well known longitudinal matching condition are still valid for the first driver
- For multiple driver configurations the longitudinal matching for the drivers following the first one is different
 - Bunches can be even shorter allowing to avoid the presence of driver tails in witness accelerating region
- > 2D+1W resonant accelerating scheme:
 - 1st driver follows the longitudinal matching condition $k_p \sigma_z \leq \sqrt{2}$
 - Short 2nd driver inside the second bubble

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 $k_p \sigma_z \leq \sqrt{2}$





No Longitudinal matching conditions for drivers injected within the second bubble