

Experimental considerations on emittance growth in the Drive Beam recombination at CTF3

EAAC Workshop 2015 - La Biodola (IT) - 14 September 2015
Davide Gamba (davide.gamba@cern.ch)

Outline

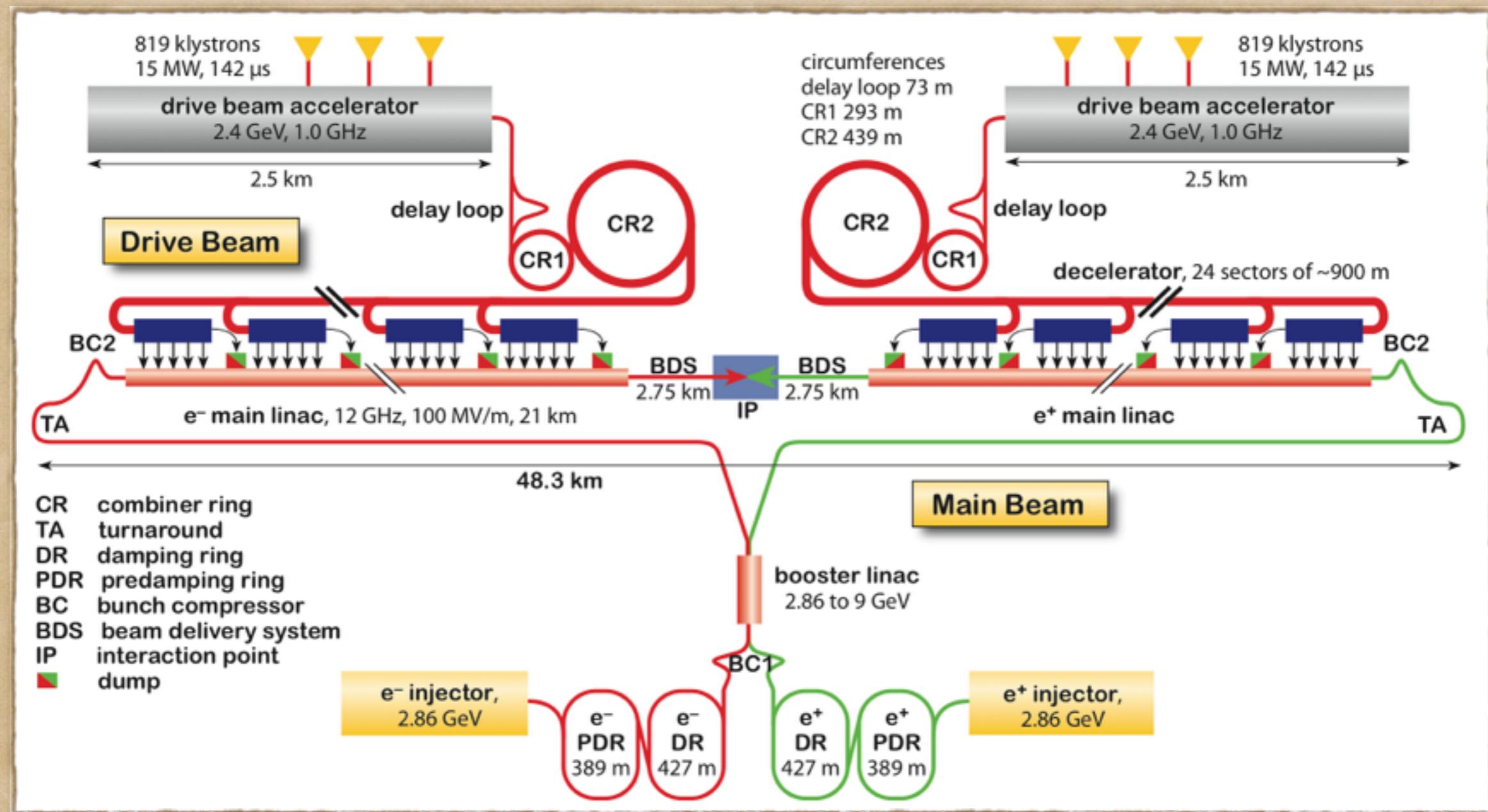
- CLIC and the CLIC Test Facility (CTF3)
- Drive Beam Recombination
 - Projected Emittance as a figure of merit
 - Sources of emittance growth
- Tools to control and optimise the beam
 - Dispersion measurements
 - Generic feedback application
 - Optics measurements
- Summary



CLIC - Design of the Compact Linear Collider

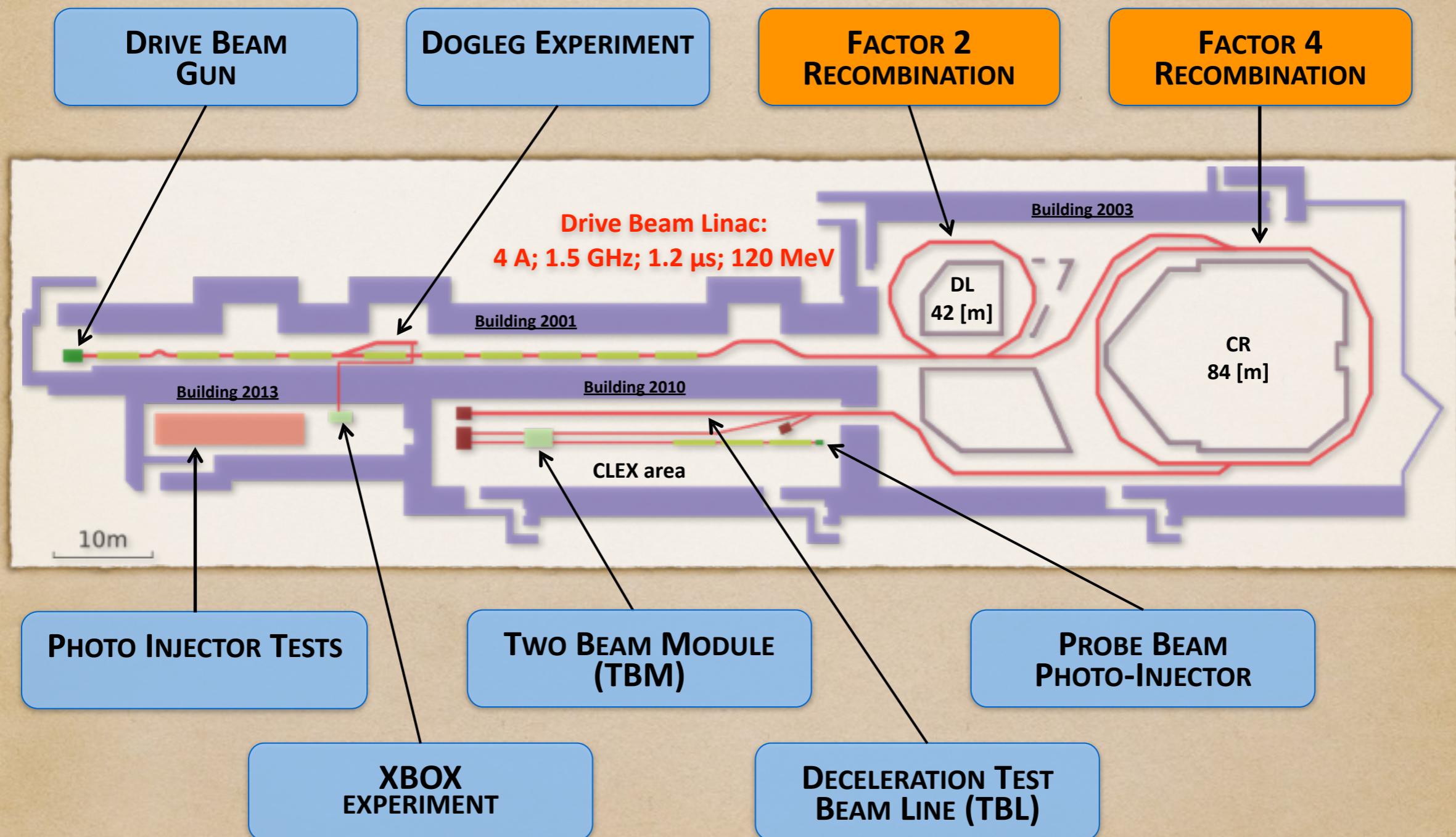
Based on two-beam acceleration scheme:

- Besides the two colliding beams, it requires two parallel high intensity Drive Beam for RF power production.
- Loaded accelerating gradient **100 MV/m**.
- 12 GHz RF generated from a **12x recombined** Drive Beam.



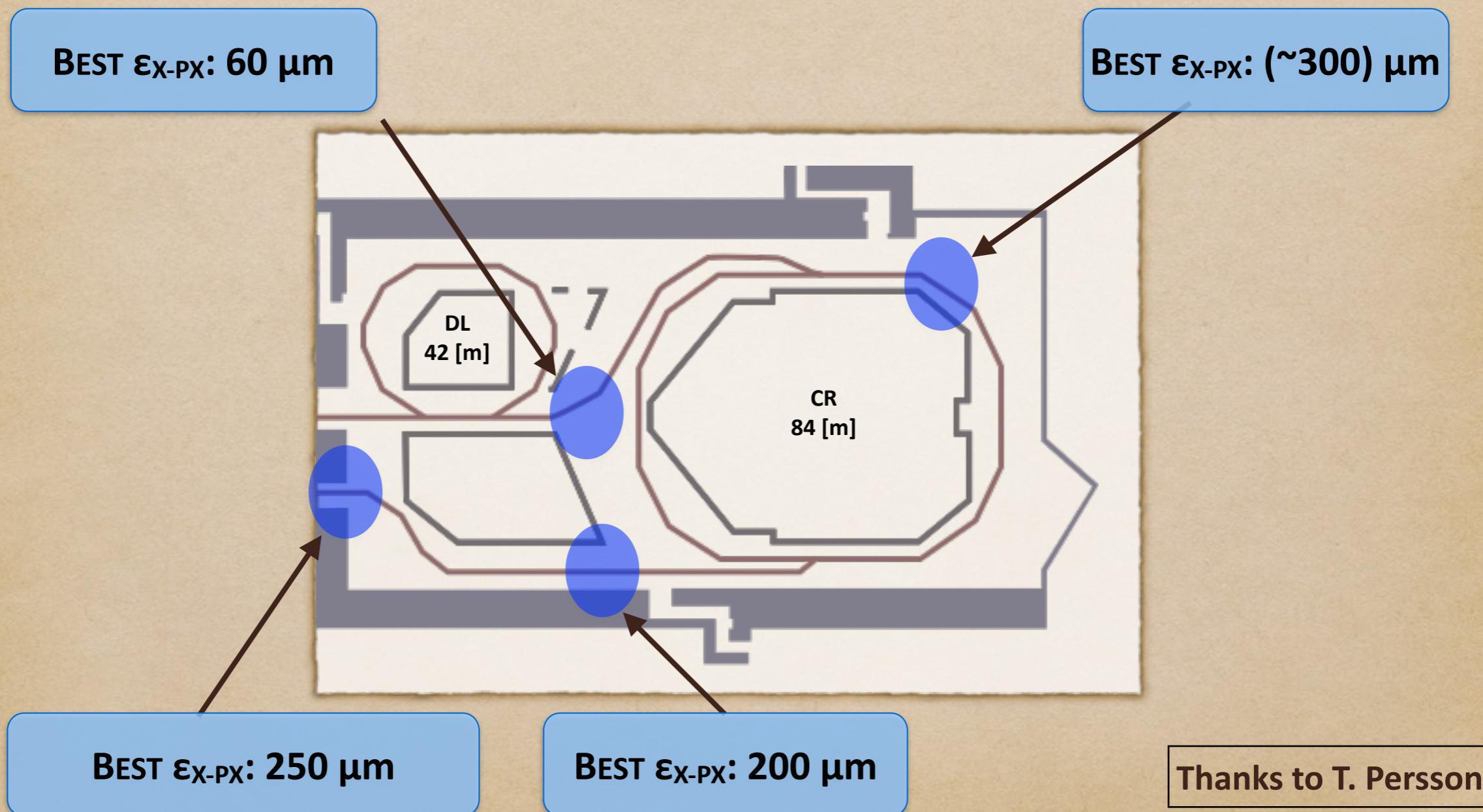


CTF3 - The CLIC Test Facility at CERN



Emittance along the Drive Beam Recombination.

- Experimentally measured by means of **quadrupole scan**: seen **inconsistencies...**
 - It might be some **optics/setup problems**, but also a **measurement related issue**.

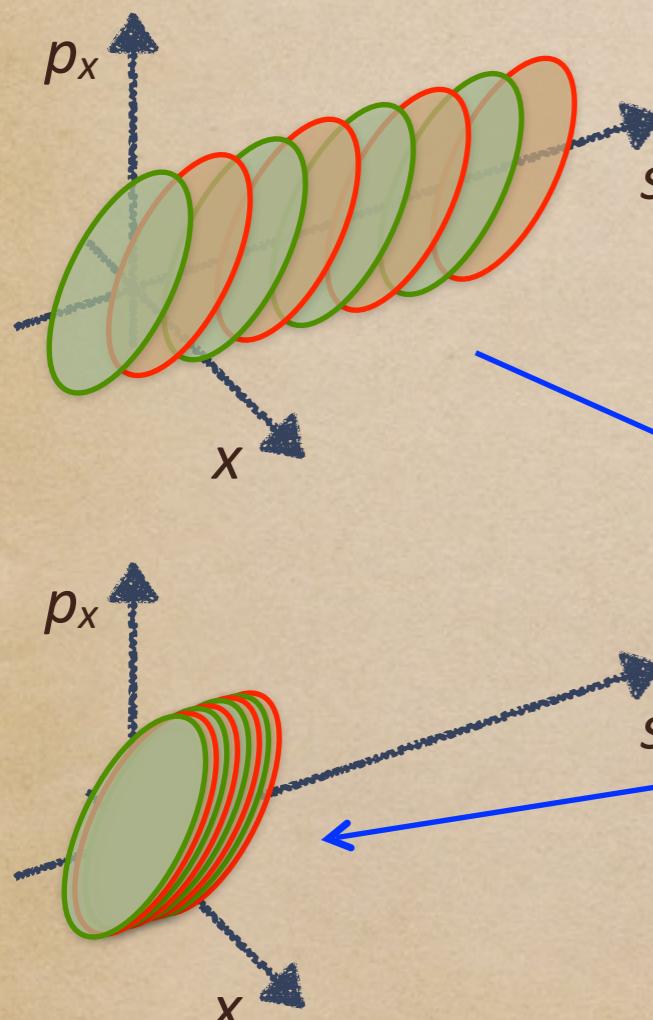


Sources of emittance growth

CTF3 Drive Beam Recombination

Incoming beam:

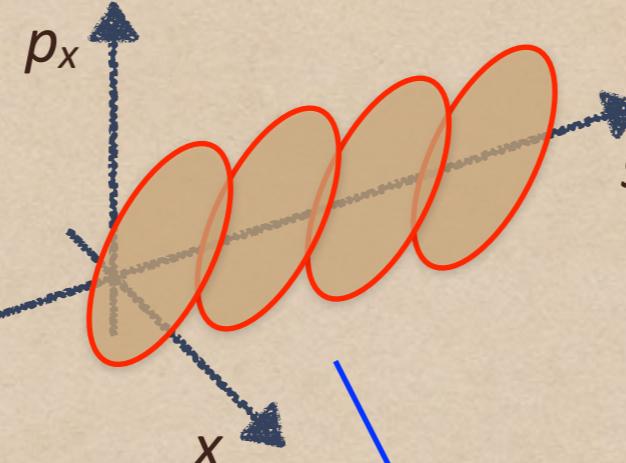
- 4 A; 1.5 GHz; 1.2 μ s; 120 MeV
- divided in eight 140 ns sub-trains



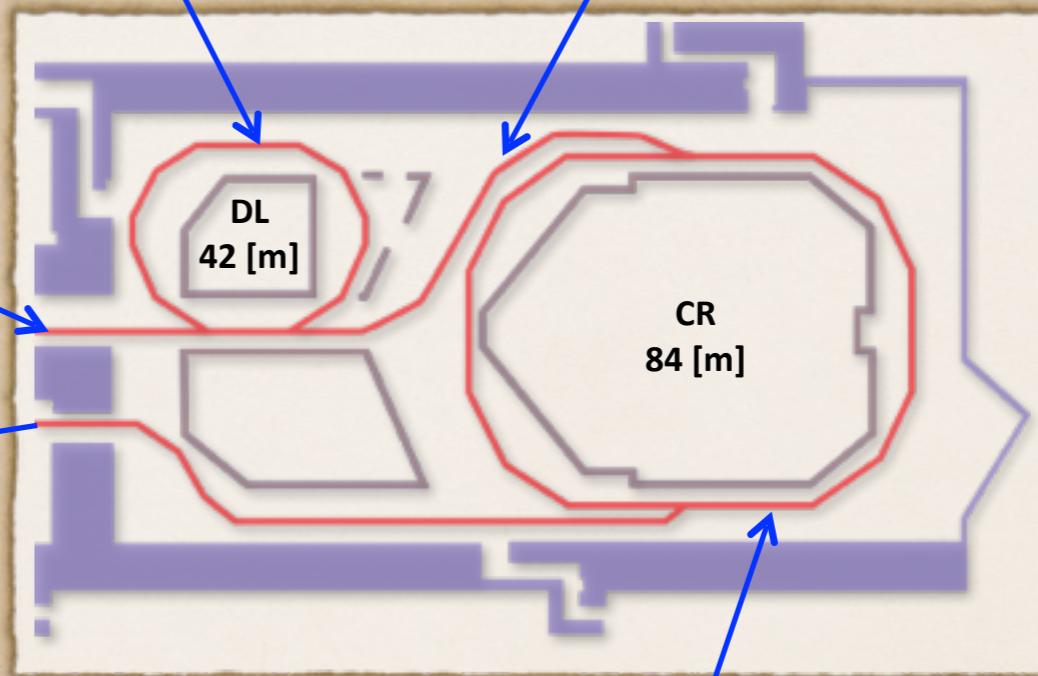
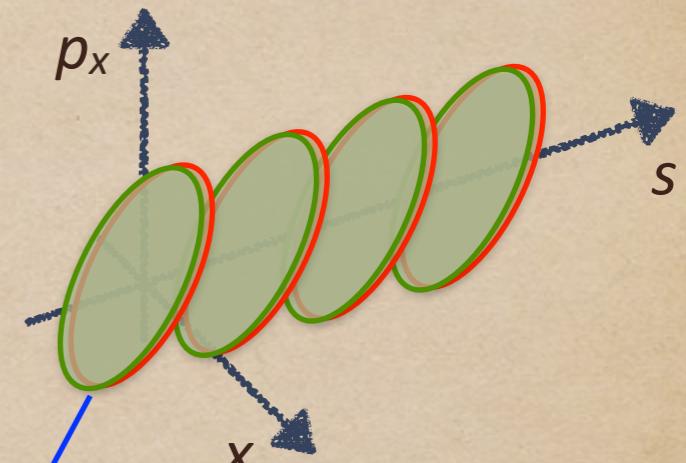
Outgoing beam:

- ~30 A; 12 GHz; 140 ns; 120 MeV

Every second sub-train is
'delayed' in the Delay Loop

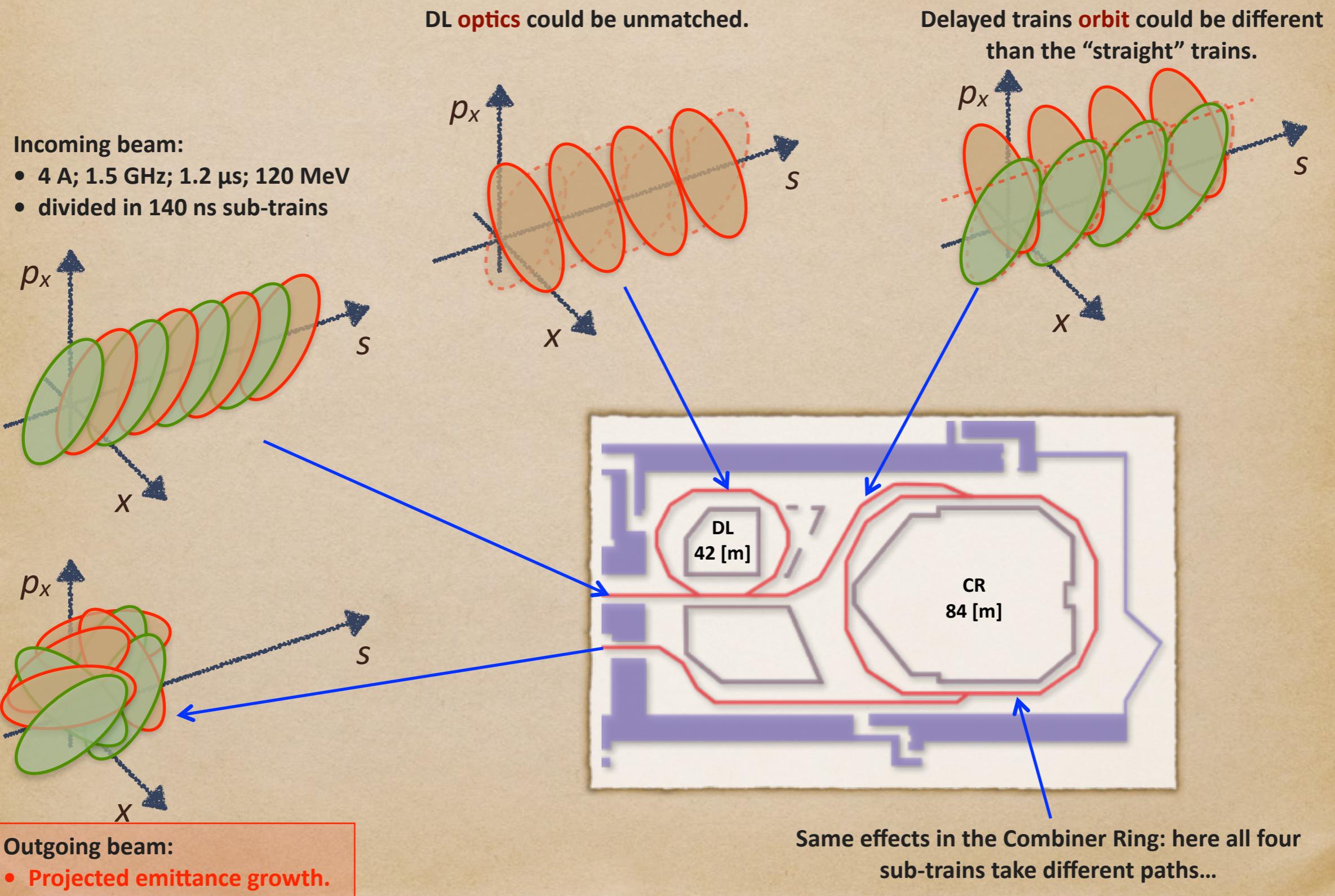


'Odd' and 'Even' sub-trains are
recombined into four 3 Ghz sub-trains

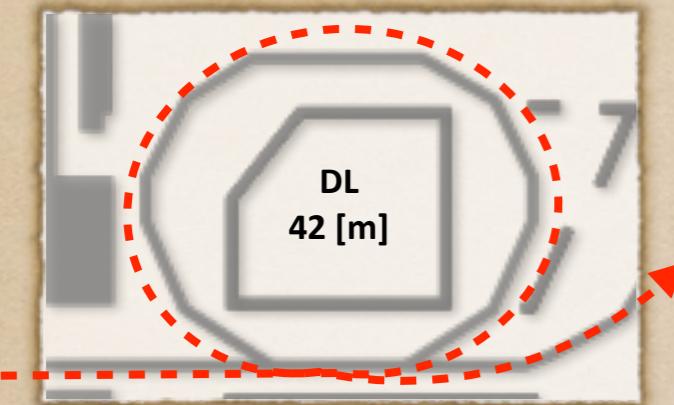


The four 3 Ghz sub-trains are interleaved by performing
different number of turns in the Combiner Ring.

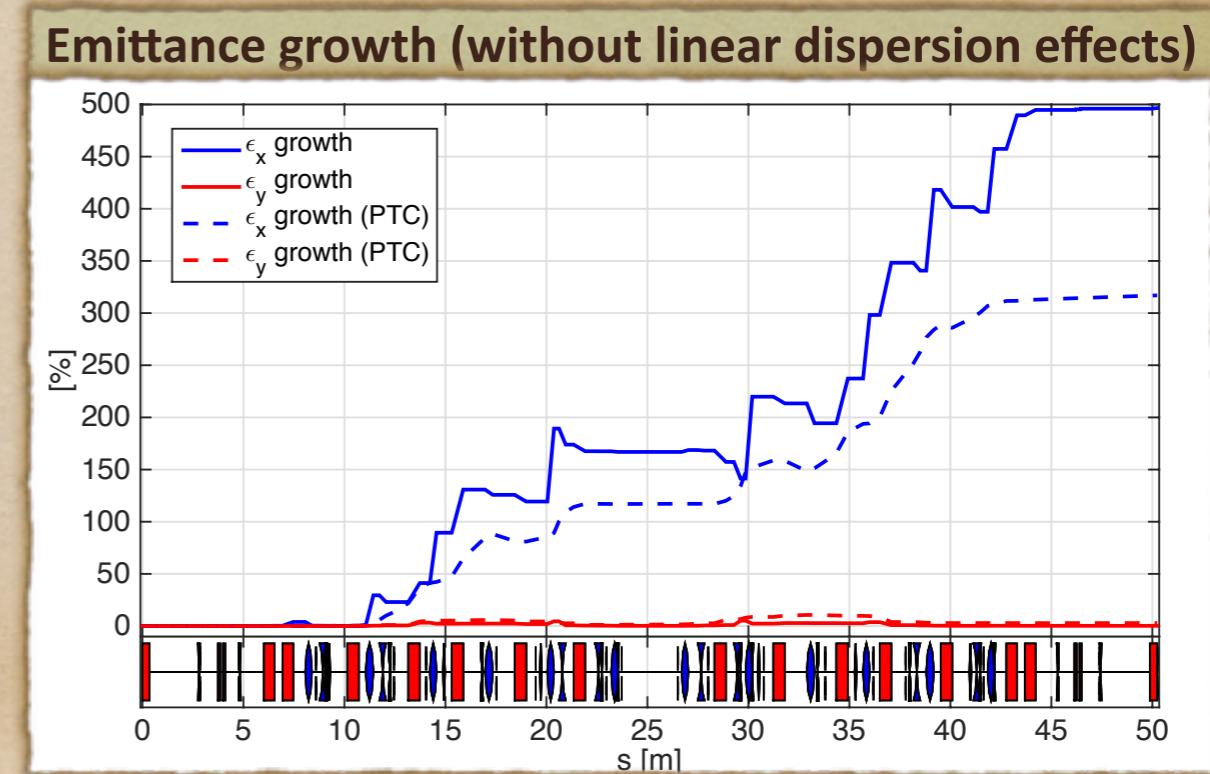
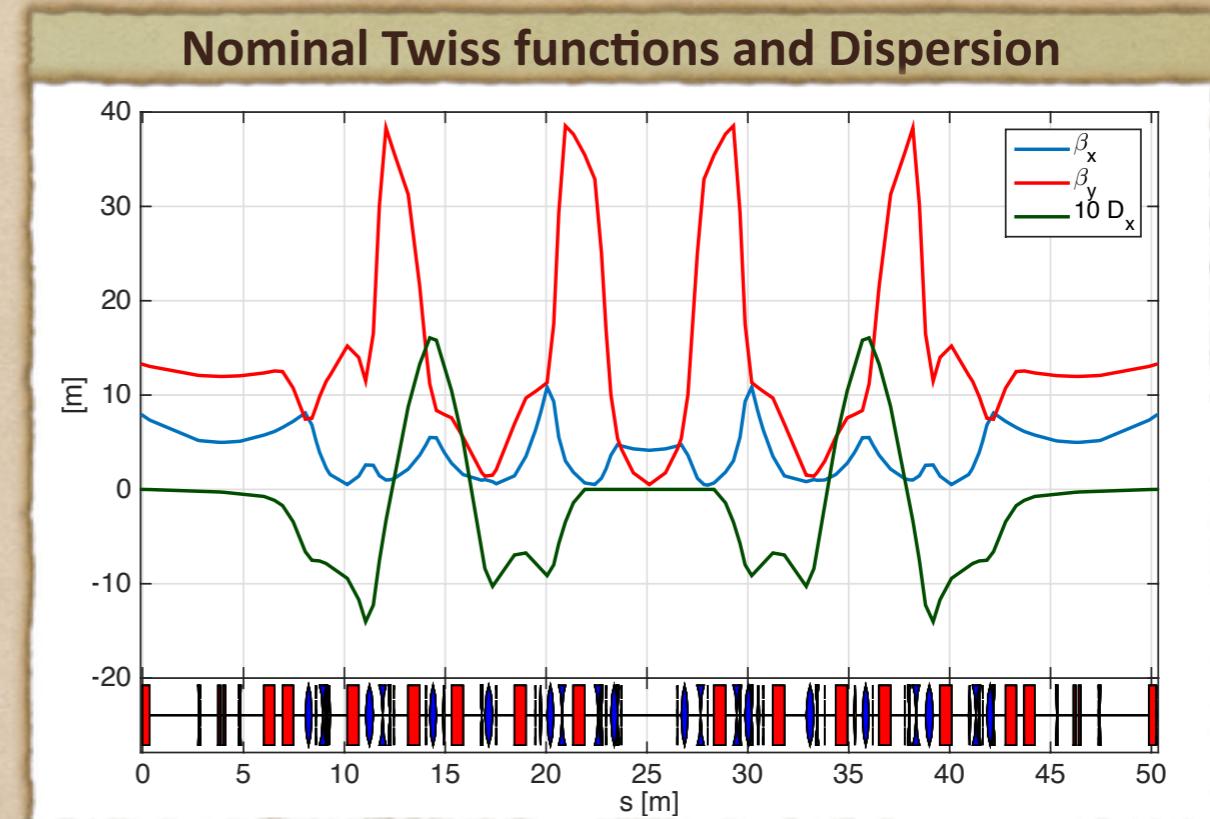
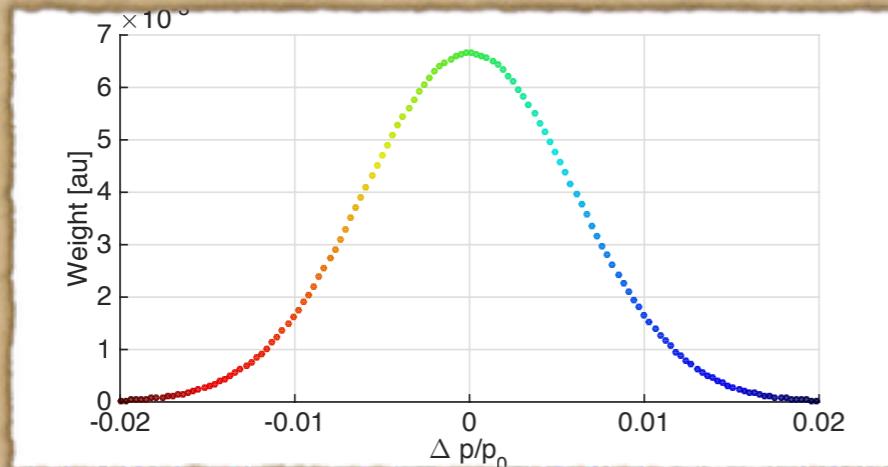
Recombination: everything can go wrong...



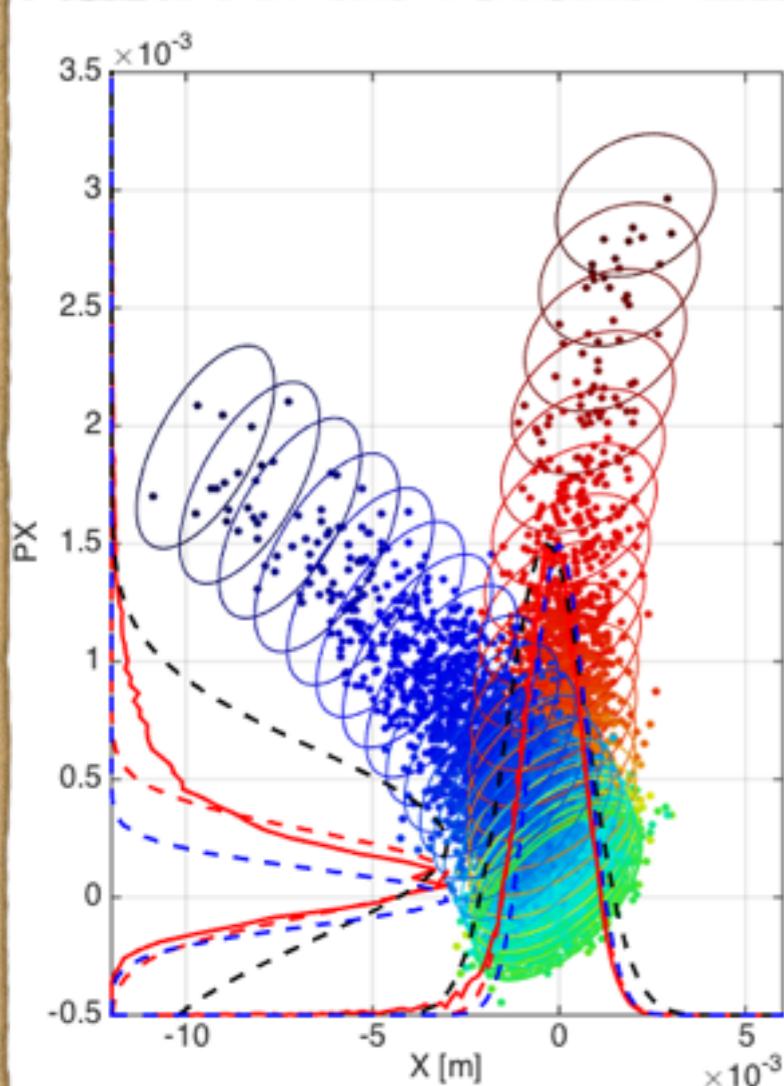
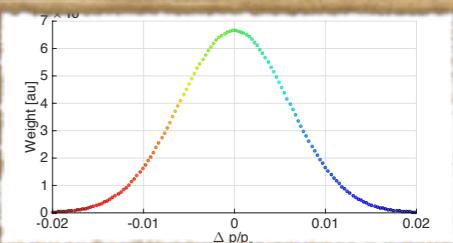
On top: energy spread...



- Single passage in the Delay Loop.
- Assume perfect transverse matching.
- Typical incoming beam:
 - $\epsilon_x = 60 \mu\text{m}$; $\epsilon_y = 100 \mu\text{m}$;
 - $\Delta p/p_0 = 0.6 \% \sigma$

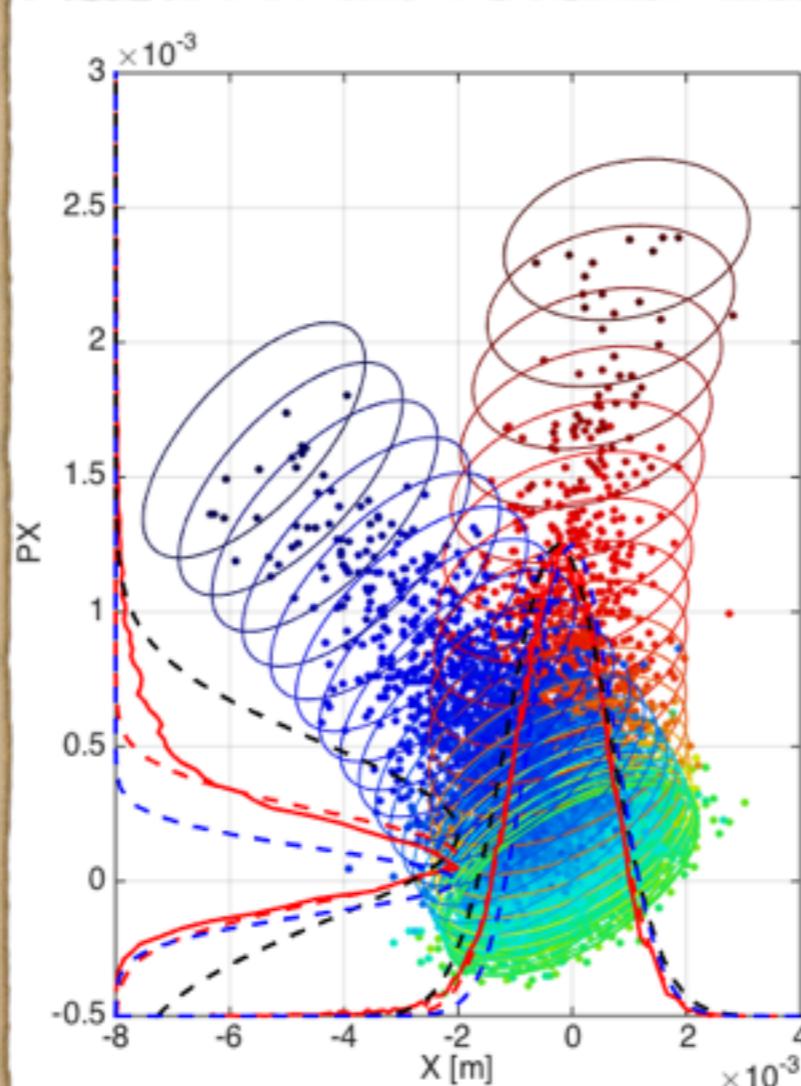


Energy spread simulations: X-PX



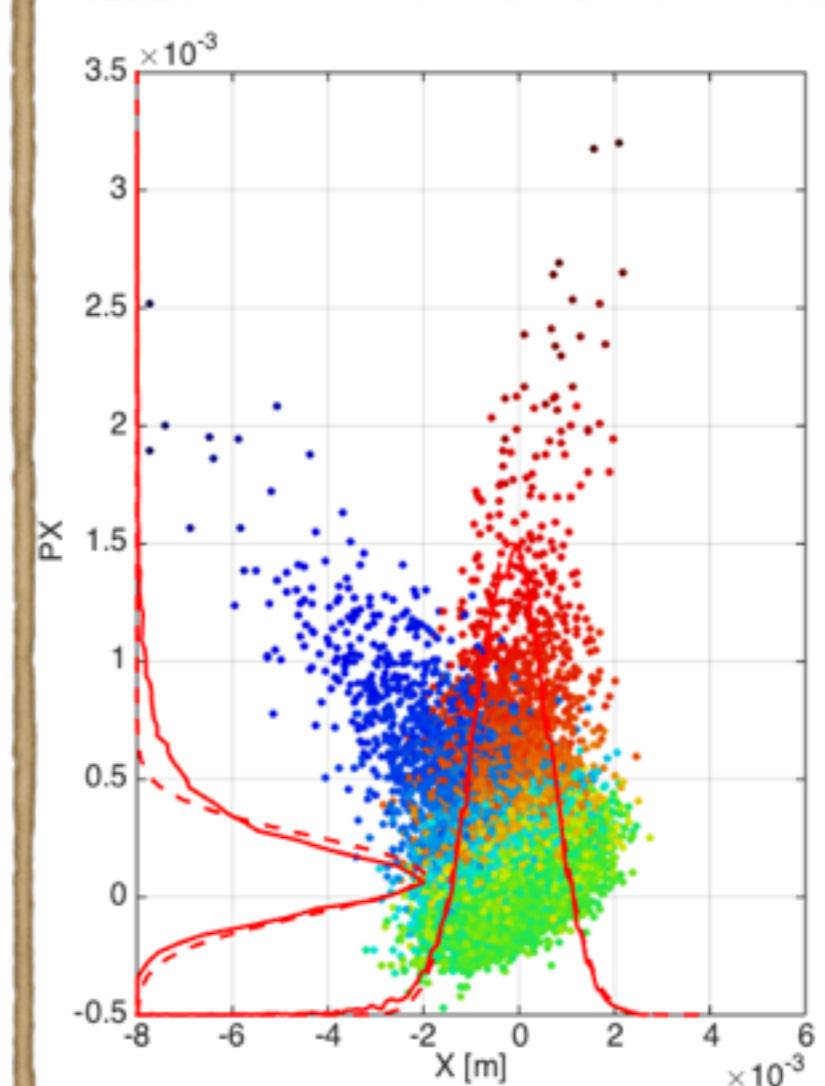
MAD-X simulation:

- ϵ_x growth = ~ 496.5 [%]
- nominal FWHM X = 1.7424 [mm]
- nominal FWHM PX = 0.2781e-3
- Gaussian fit FWHM X = 1.8925 [mm]
- Gaussian fit FWHM PX = 0.41397e-3



PTC TWISS simulation:

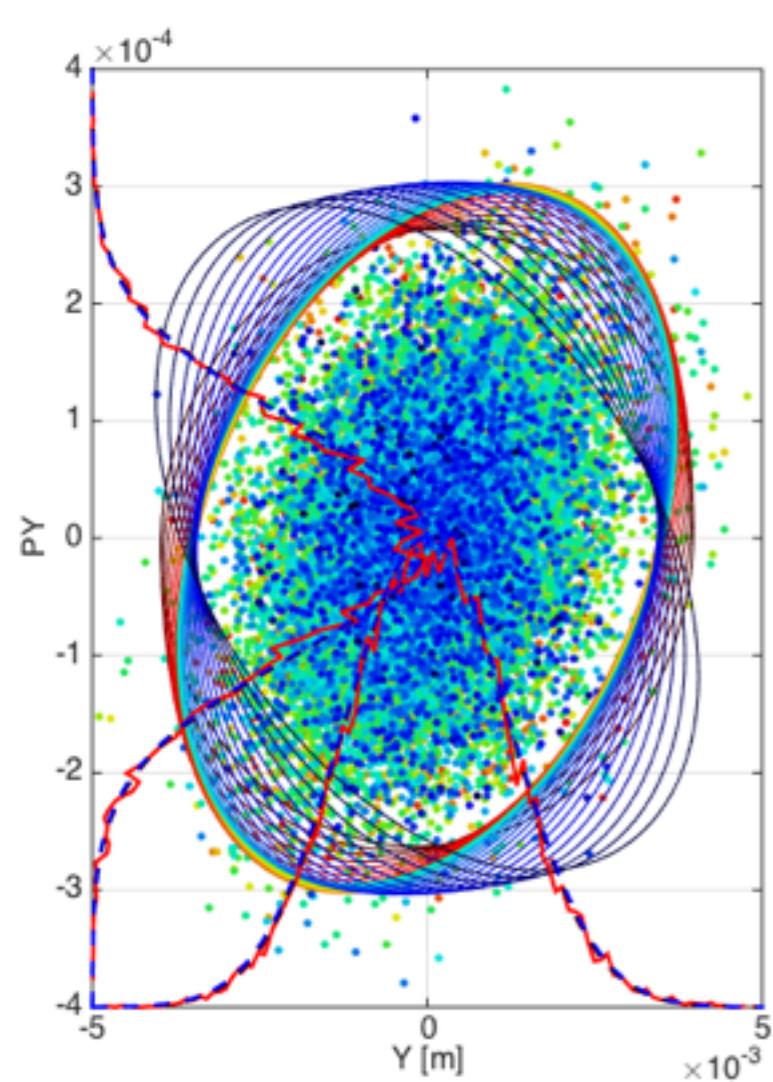
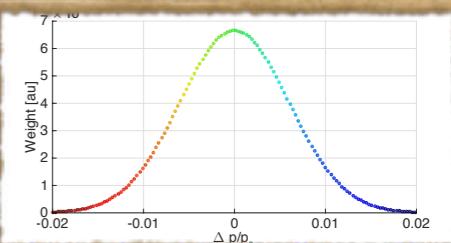
- ϵ_x growth = ~ 317 [%]
- **nominal FWHM X = 1.7424 [mm]**
- nominal FWHM PX = 0.2781e-3
- **Gaussian fit FWHM X = 1.8437 [mm]**
- Gaussian fit FWHM PX = 0.39237e-3



PTC TRACK simulation:

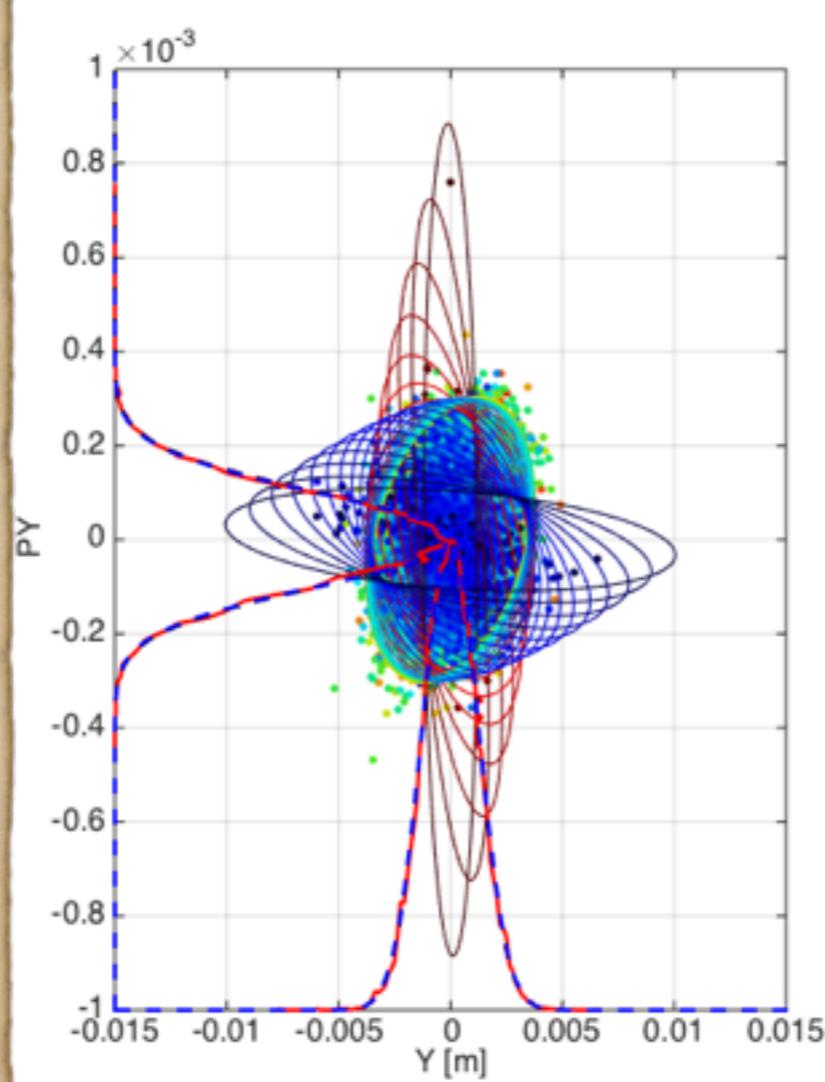
- **ϵ_x growth = ~ 257.8 [%]**
- nominal FWHM X = 1.7424 [mm]
- **nominal FWHM PX = 0.2781e-3**
- Gaussian fit FWHM X = 1.8595 [mm]
- **Gaussian fit FWHM PX = 0.39229e-3**

Energy spread simulations: Y-PY



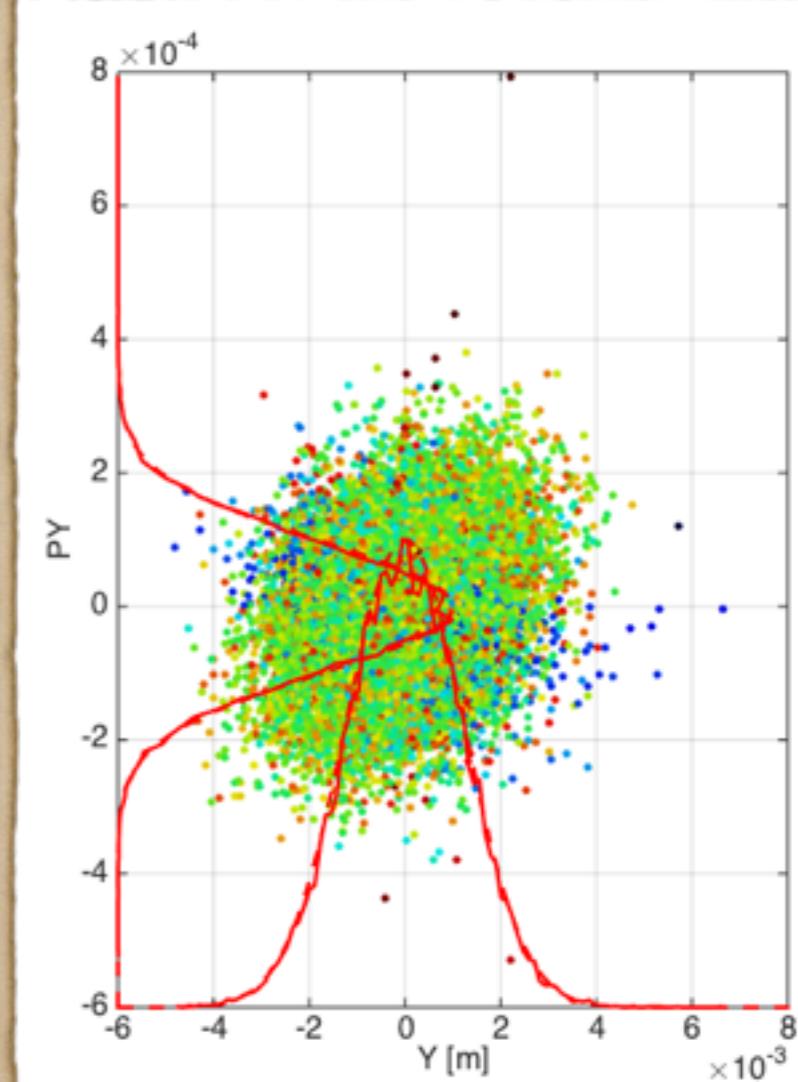
MAD-X simulation:

- ϵ_y growth = ~ 0.2 [%]
- nominal FWHM Y= 2.9228 [mm]
- nominal FWHM PY = 0.23193e-3
- Gaussian fit FWHM Y = 2.8701 [mm]
- Gaussian fit FWHM PY = 0.23448e-3



PTC TWISS simulation:

- ϵ_y growth = ~ 2.8 [%]
- nominal FWHM Y= 2.9228 [mm]
- nominal FWHM PY = 0.23193e-3
- Gaussian fit FWHM Y = 2.8947 [mm]
- Gaussian fit FWHM PY = 0.23522e-3

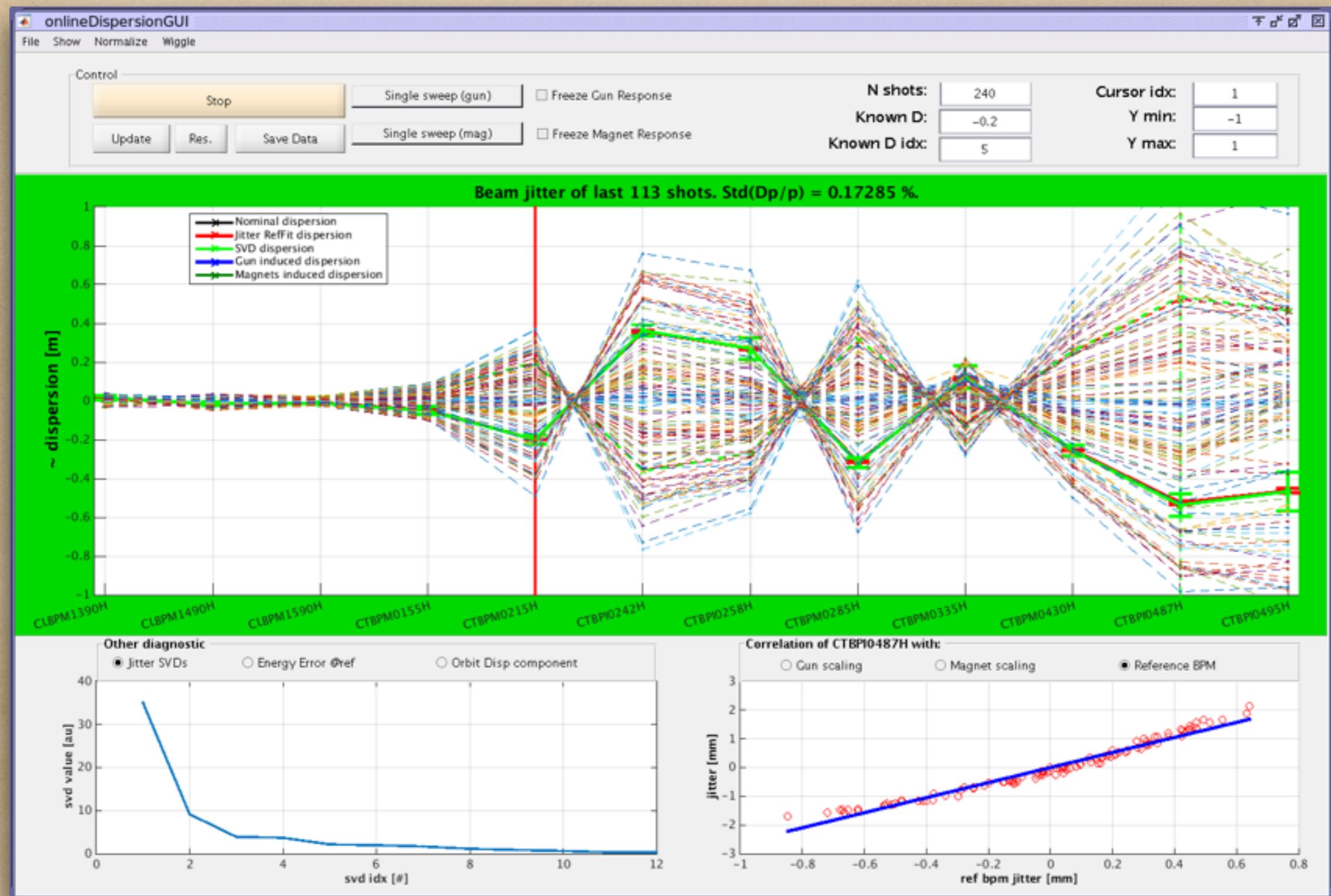


PTC TRACK simulation:

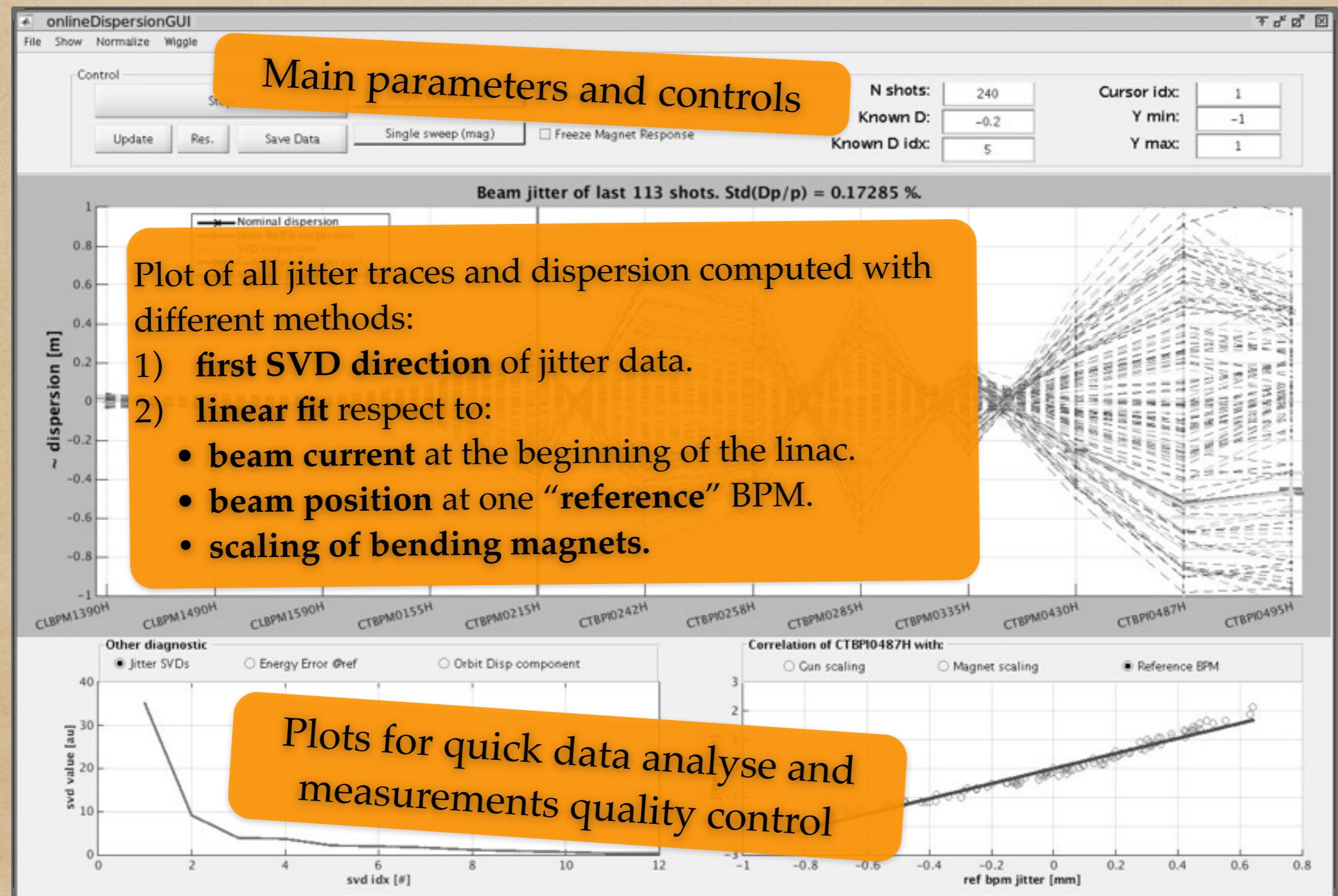
- ϵ_y growth = ~ 3 [%]
- nominal FWHM Y = 2.9228 [mm]
- nominal FWHM PY = 0.23193e-3
- Gaussian fit FWHM Y = 2.9099 [mm]
- Gaussian fit FWHM PY = 0.23317e-3

Implemented tools to diagnose and optimise the beam.
Some experimental results.

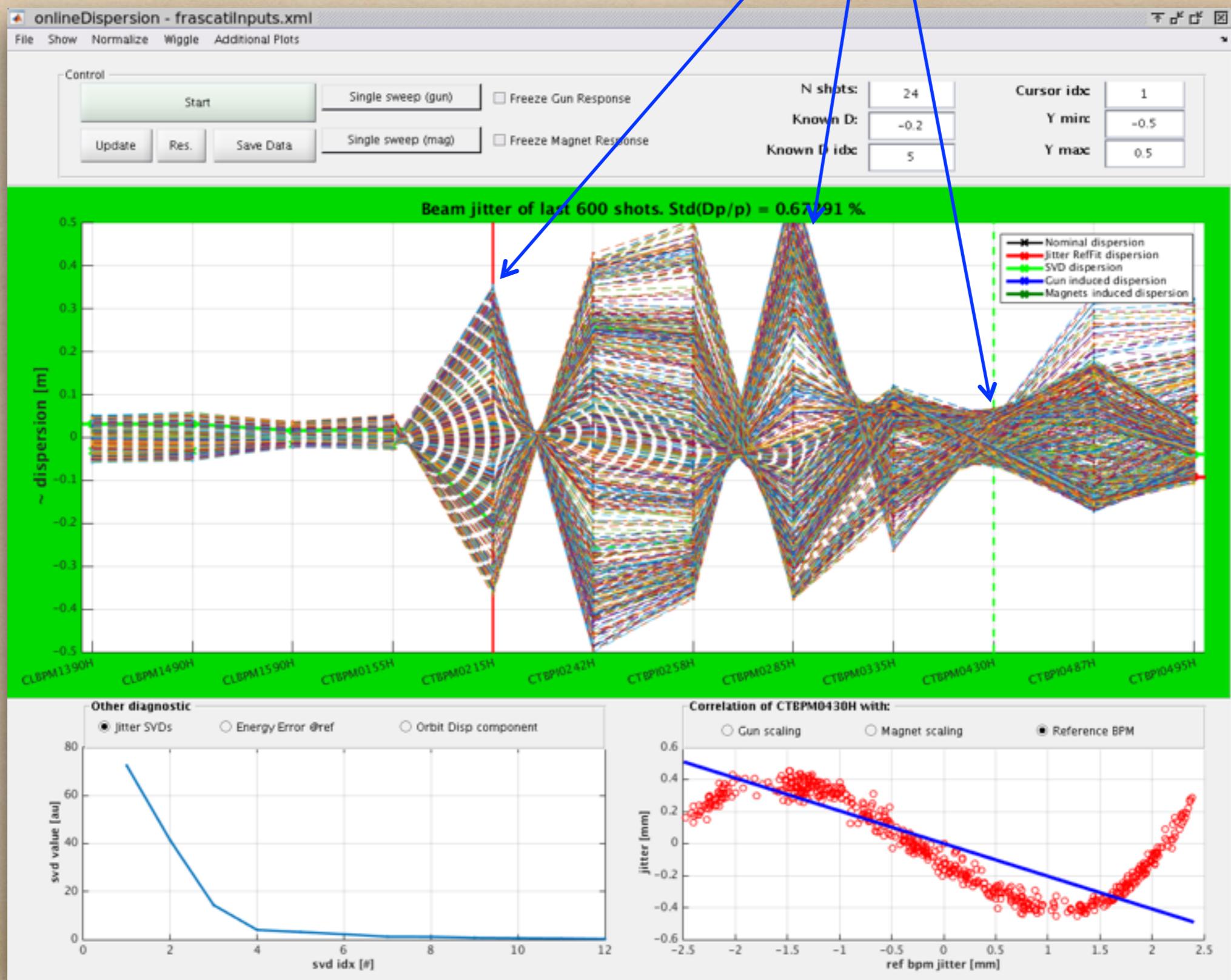
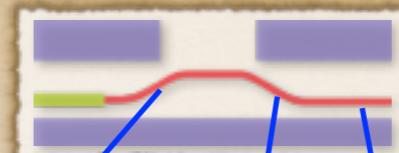
Dispersion measurements



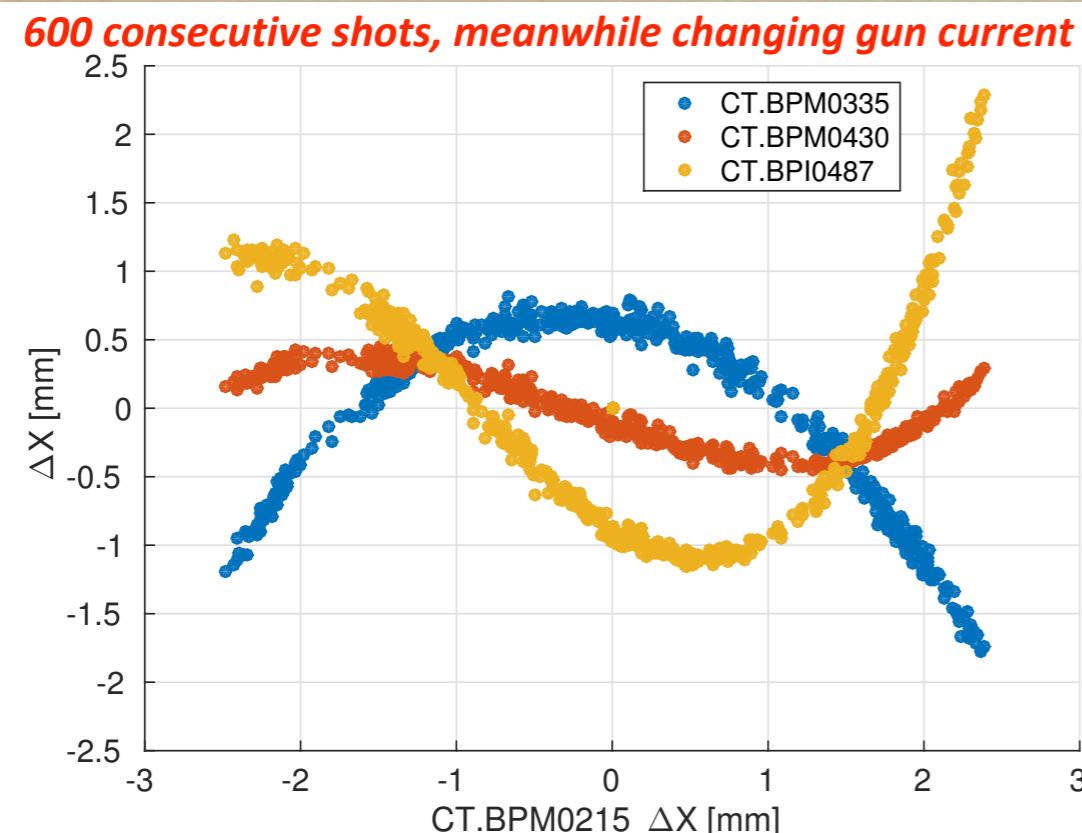
Dispersion measurements



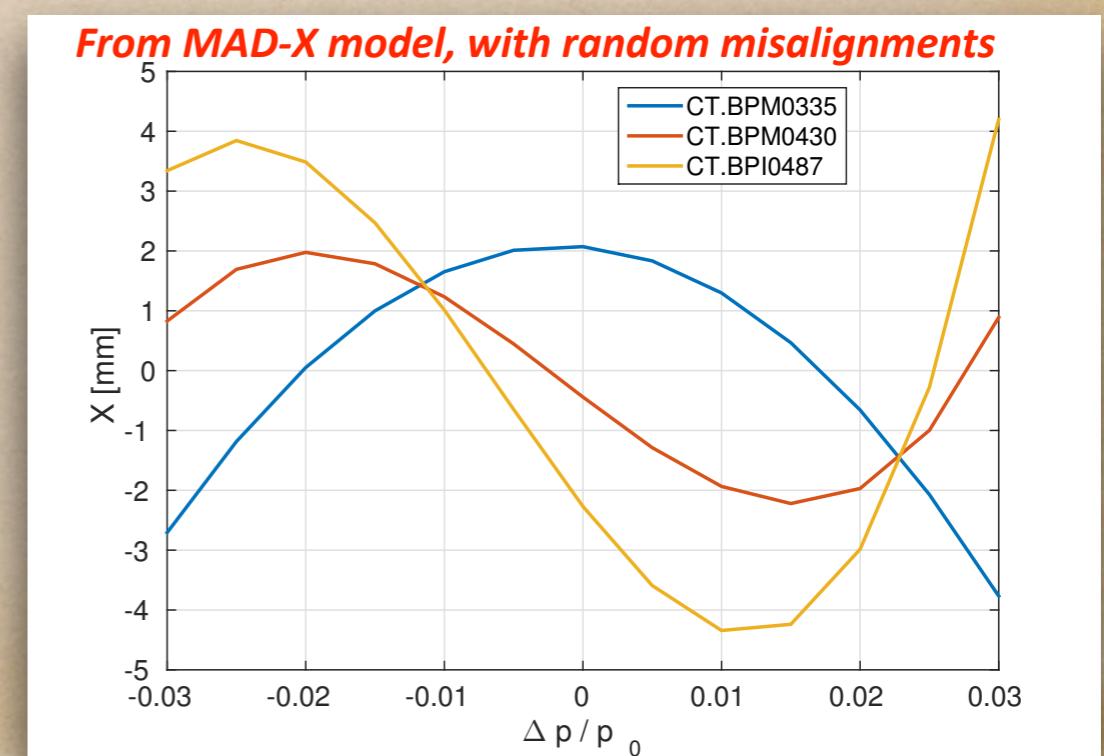
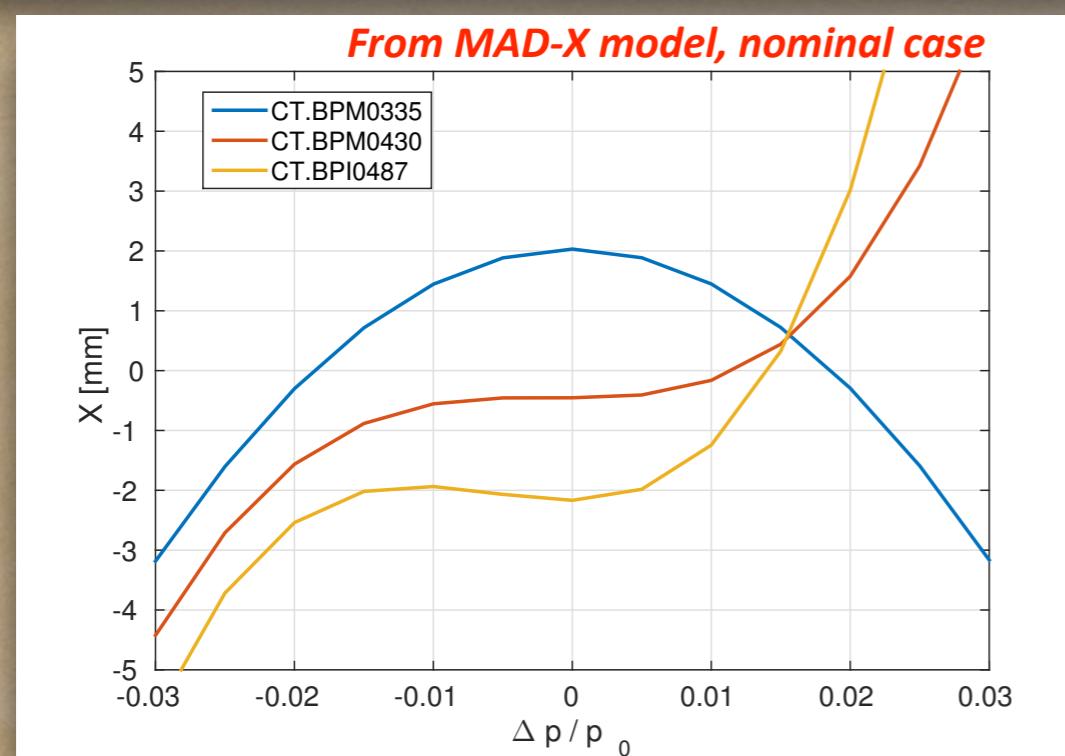
Dispersion in Frascati Chicane.



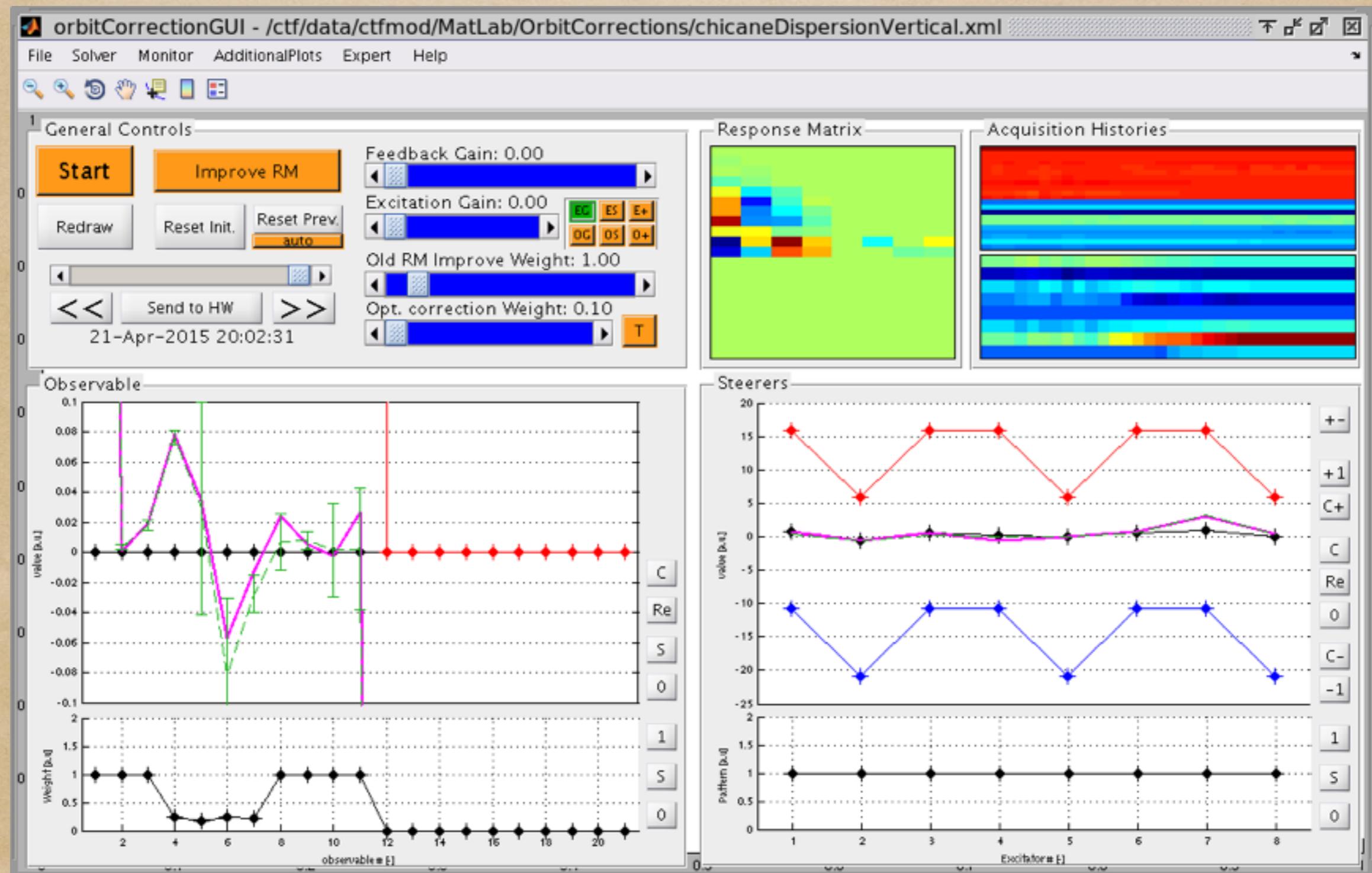
Dispersion in Frascati Chicane: a closer look.



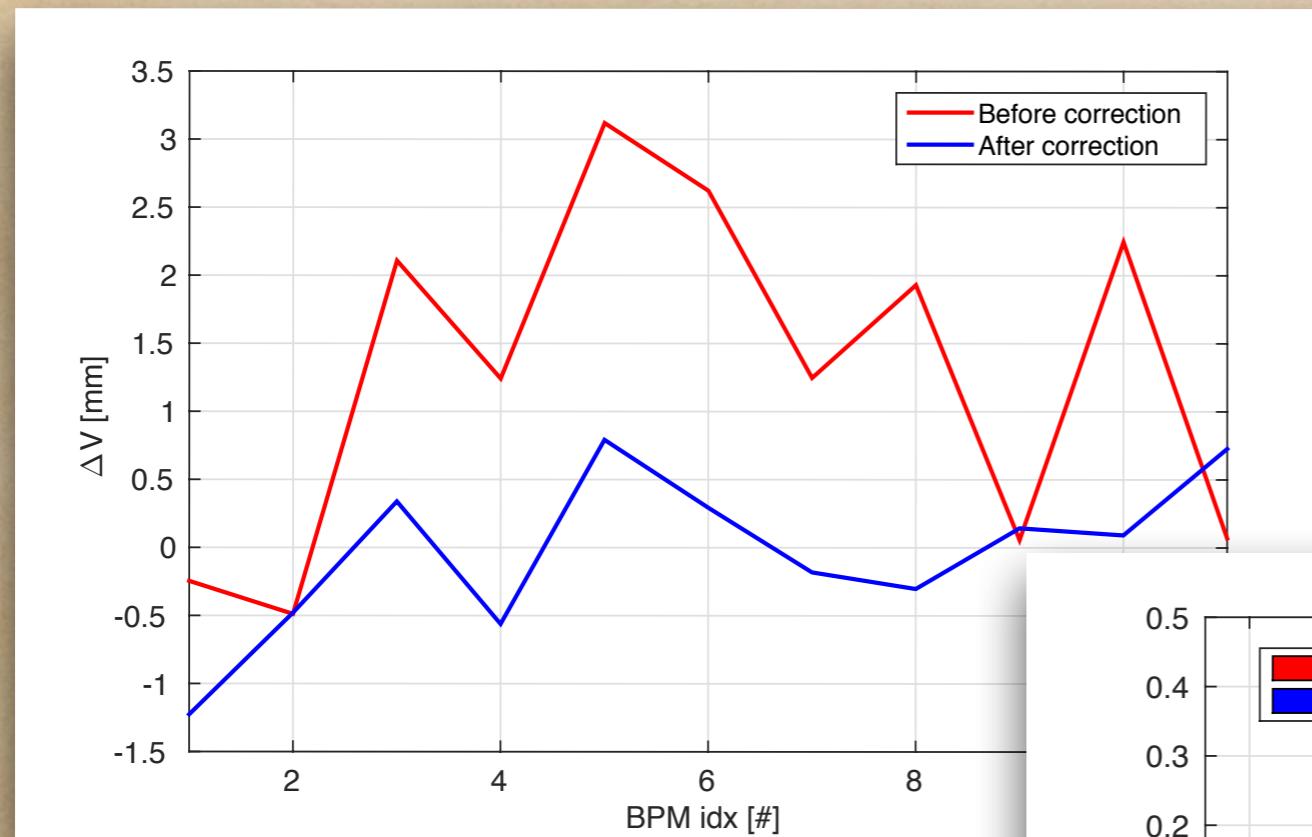
- Clear non-linearities visible on the dispersion free BPMs after the chicane.
- **Similar results** can be obtained by simulating different beam energies with MAD-X in the **nominal case** and with (*lucky*) random quadrupoles **misalignments** ($\sigma=0.2$ [mm]).



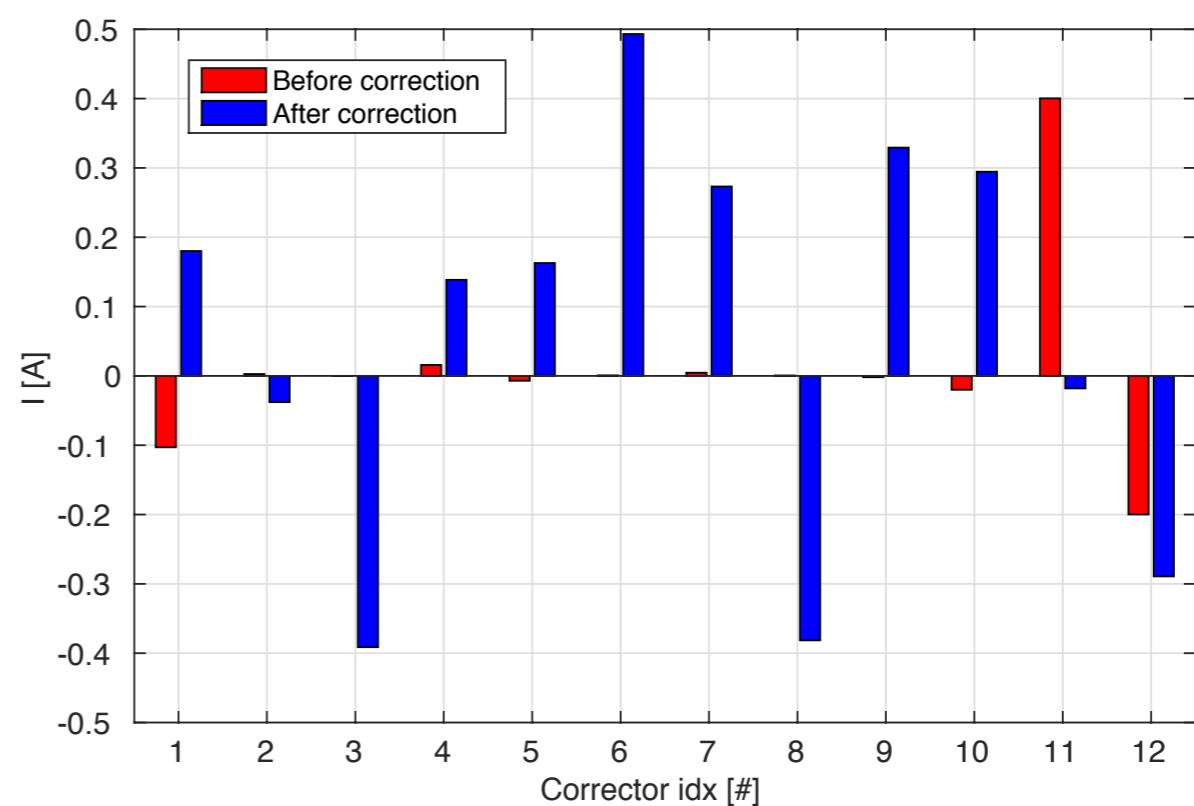
Slow feedback



Orbit correction: vertical closure of CR.



Strength of the vertical correctors in the CR before and after the correction.

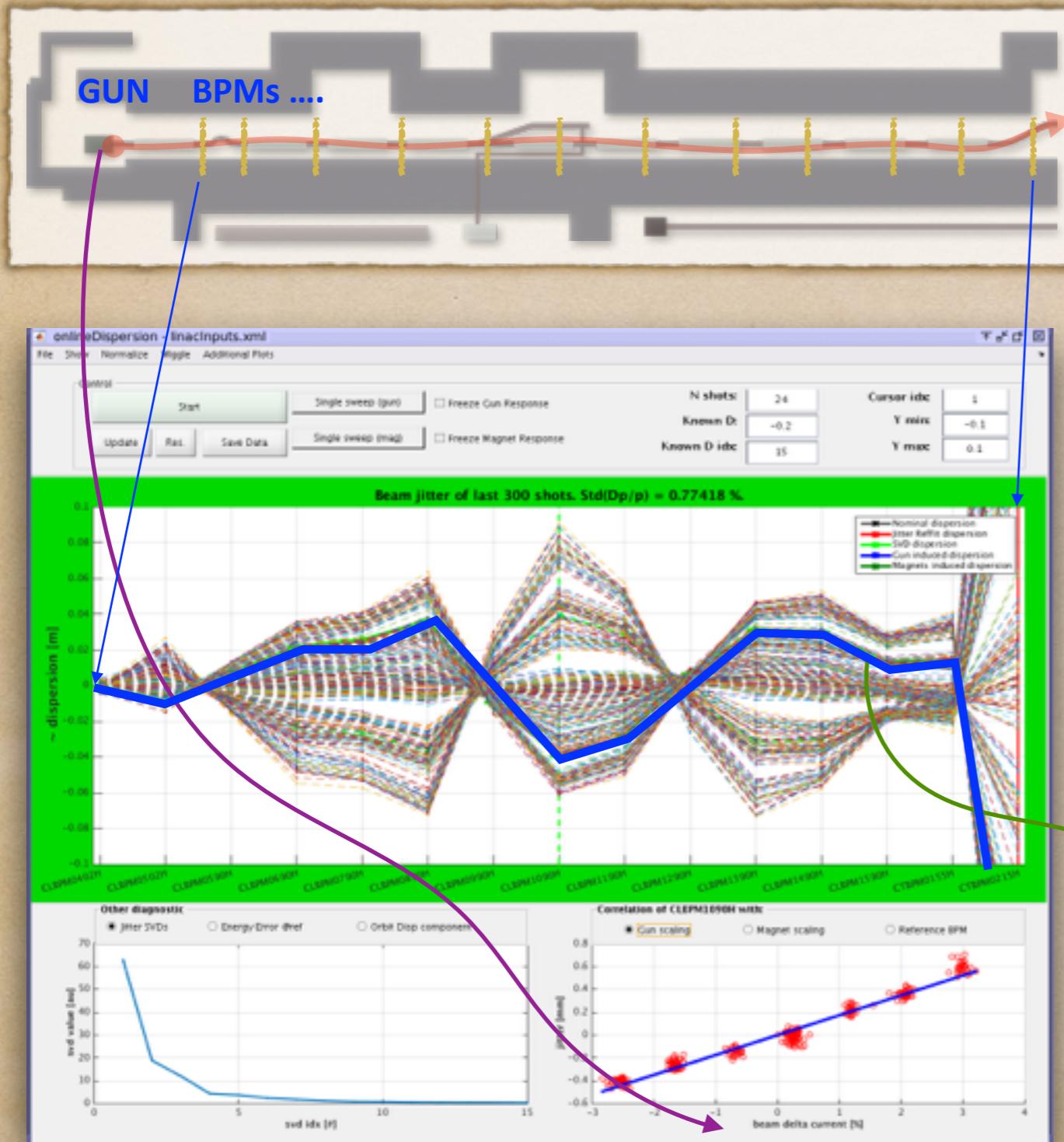


Observables (i.e. differences between one turn and the next one on the mostly trusted BPMs of CR) before and after correction.

Further improvements mainly limited by losses.

- Similar results has been obtained for horizontal closure, as well as general orbit correction in other parts of CTF3.

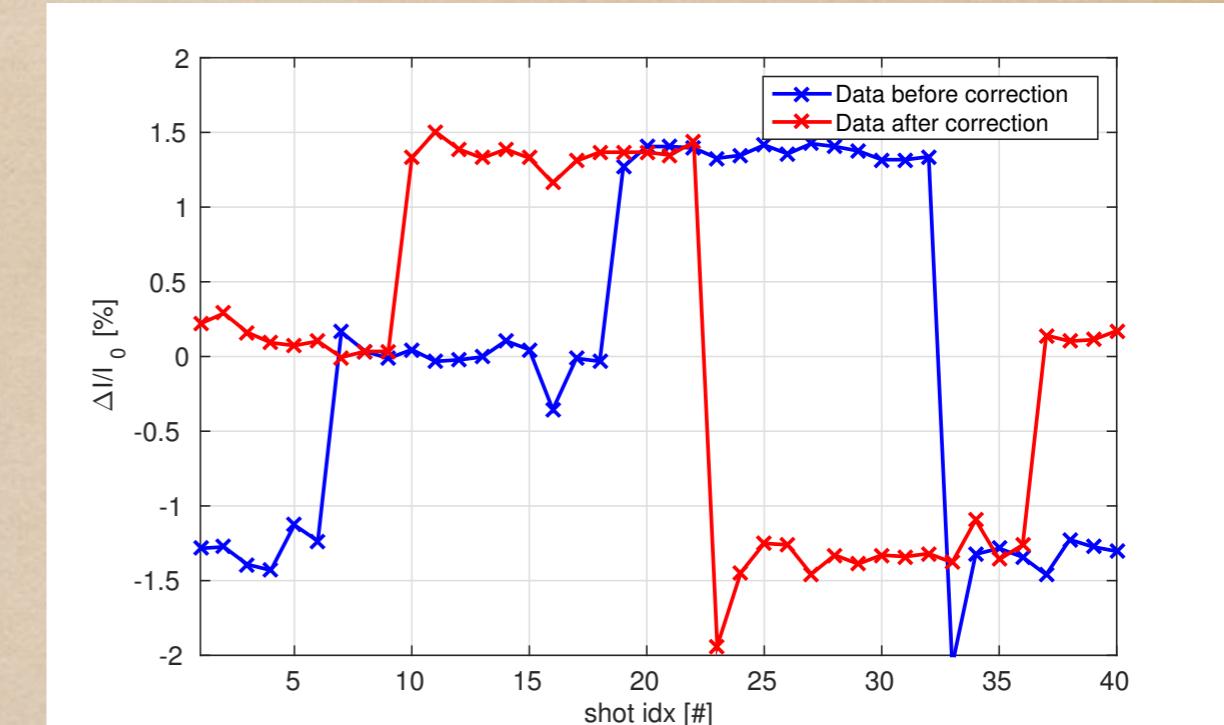
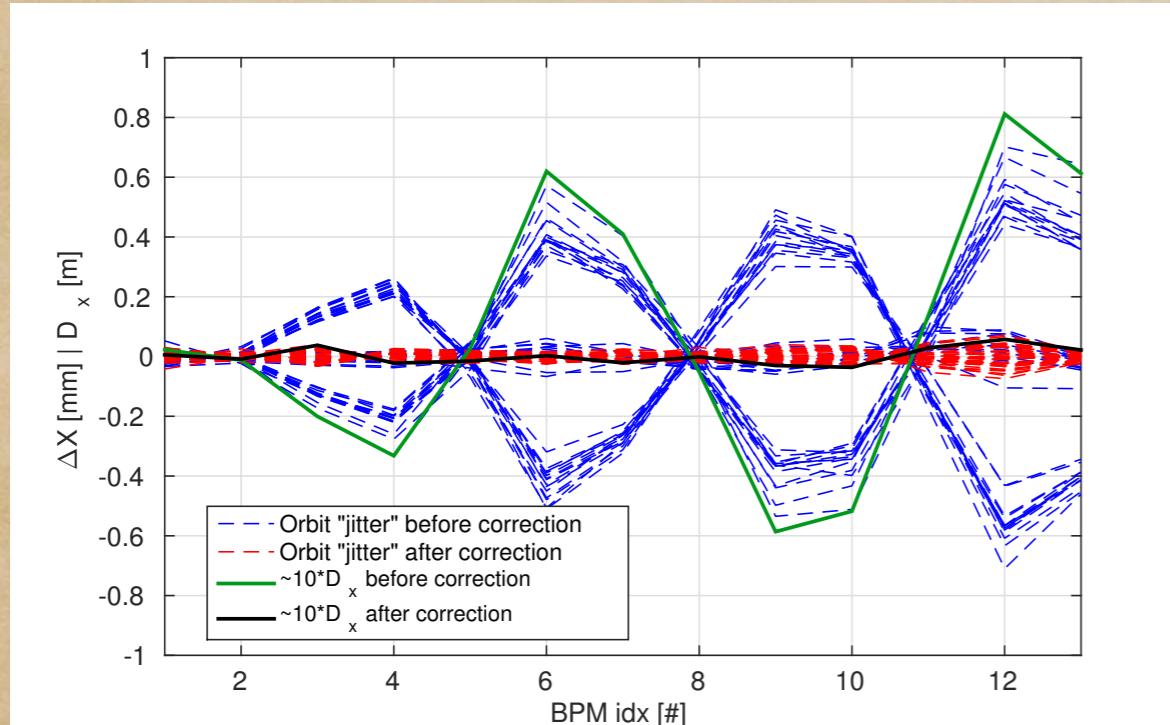
Dispersion free steering in the LINAC



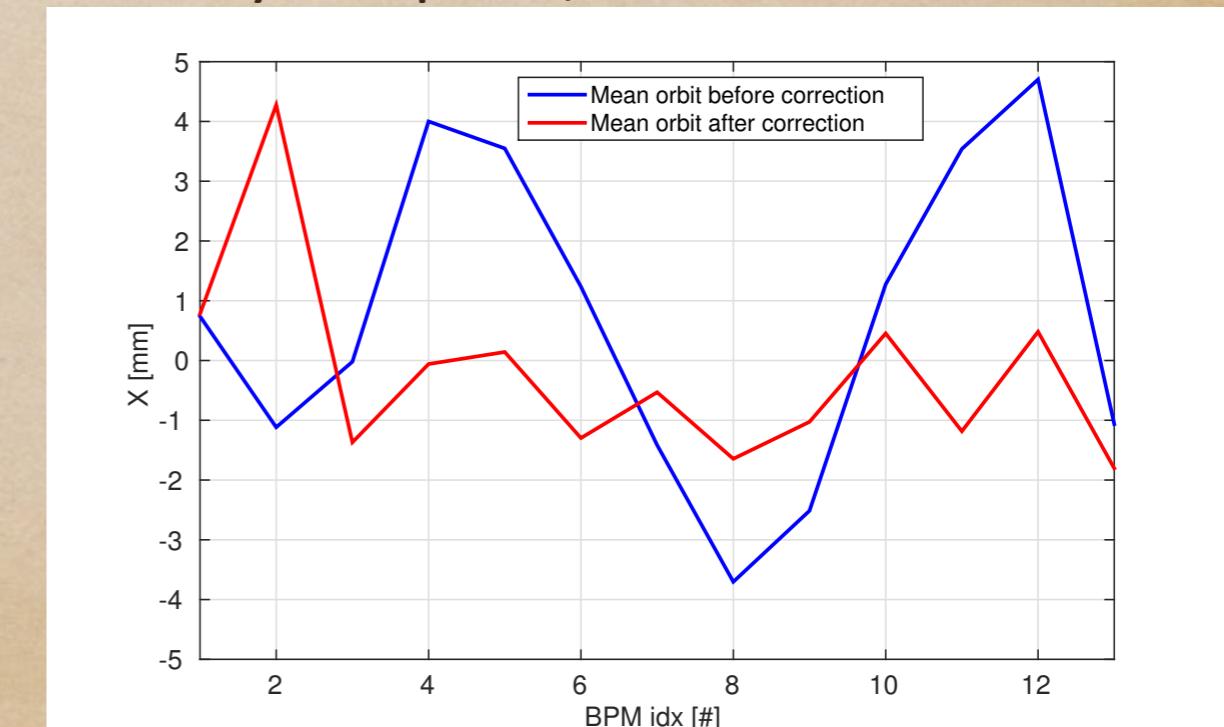
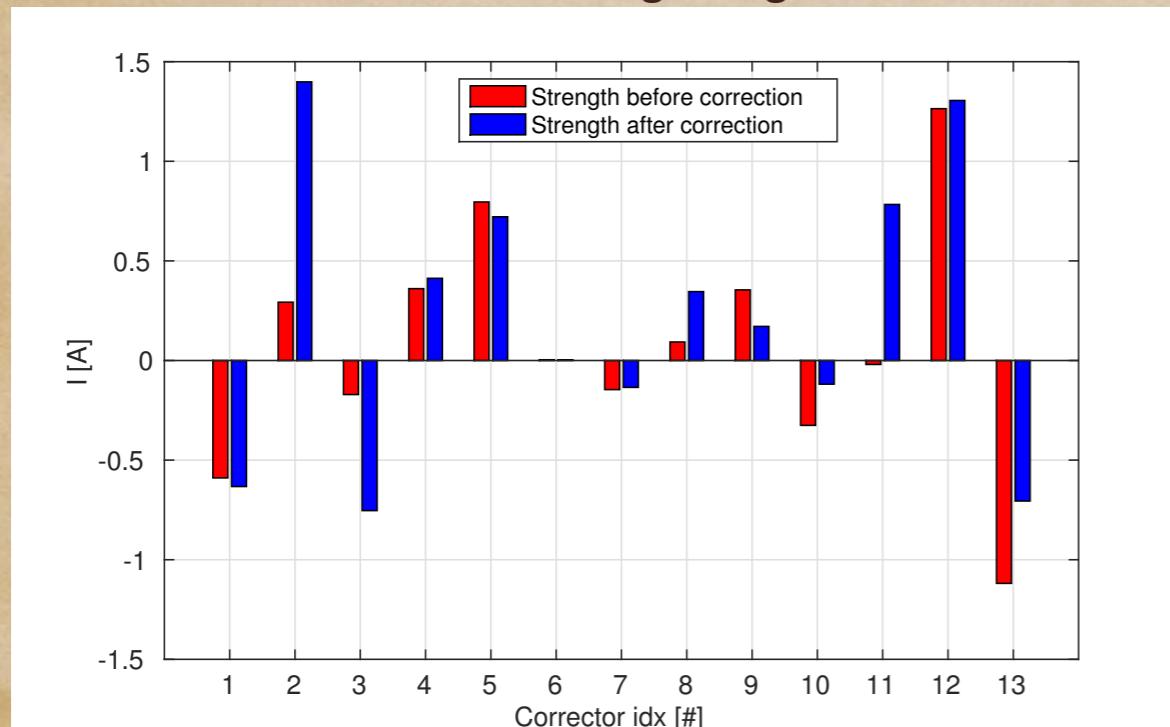
- Zero dispersion expected.
- The small chicane after the injector might leave some unclosed dispersion.
- We have big aperture, so we can afford to heavily scale the beam current without losses.
- **If we can measure something reasonable, we can correct it...**

Dispersion free steering in the LINAC

- Observables: 40 consecutive shots at different GUN current settings (scaling up/down of ~1.5% than nominal).

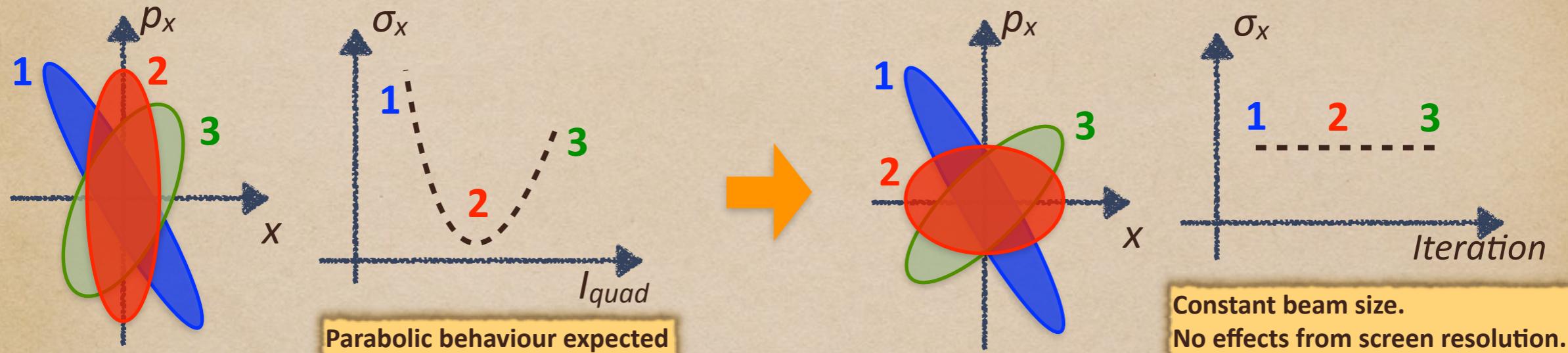


- Main correction is at the beginning of the linac. This reduces not only the dispersion, but also the orbit.



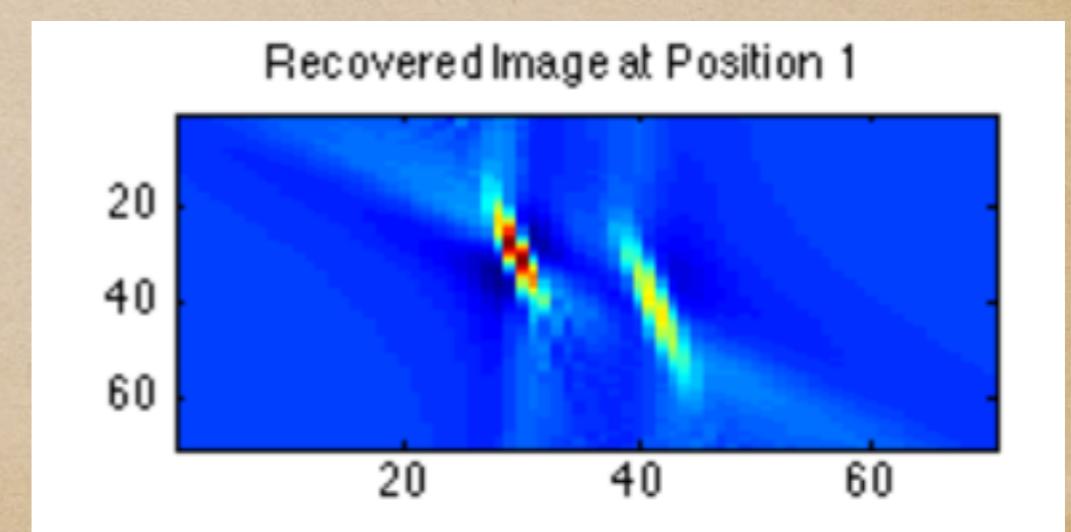
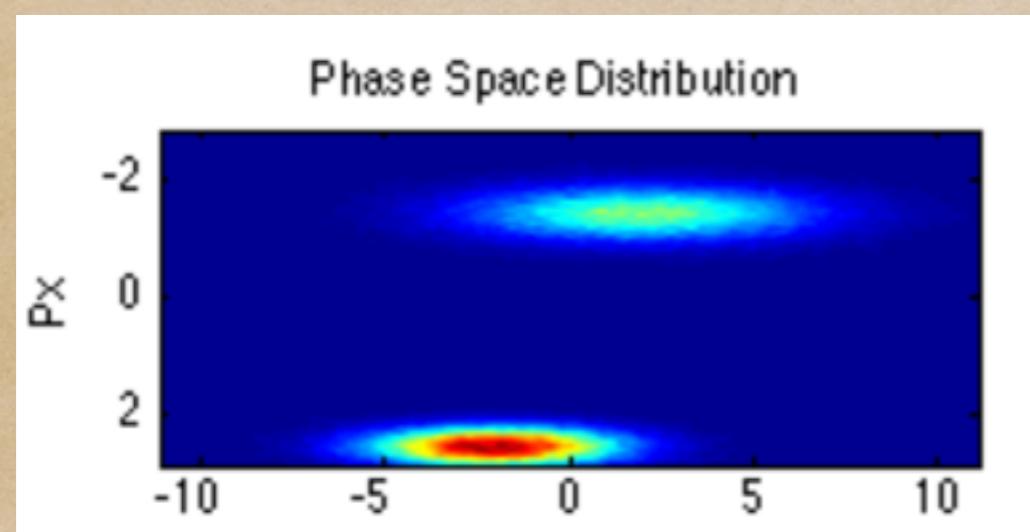
Work in progress: constant beam size Quadscan:

During a quadscan normally one quadrupole strength is varied and the beam size recorded on a screen...
... but one could use two quadrupoles and still “rotate” the beam, but keep the **beam size constant!**



Work in progress: phase-space tomography

In a similar way to quadscan, we can reconstruct the phase-space using the quadscan projections and applying **tomography** algorithms.



A. Rollings;
L. Martin

Summary

- The CTF3 Drive Beam recombination is affected by **emmitance growth** due to:
 - **Orbit mismatch** of the single sub-trains once recombined.
 - Transverse optics mismatch (less critical).
 - High beam energy spread and **non-linear dispersion**.
- We have a series of implemented tools to diagnose and correct (some) effects:
 - Dispersion measurements.
 - Generic **slow feedback** tool that has been demonstrated over a wide range of corrections. (*here only discussed orbit and dispersion, but also energy, bunching, gun current,...*)
- We are working on new methods to enhance our understanding of the leading sources of emittance growth:
 - Constant beam size based quadscan.
 - Transverse beam size tomography.

Thank you for your attention



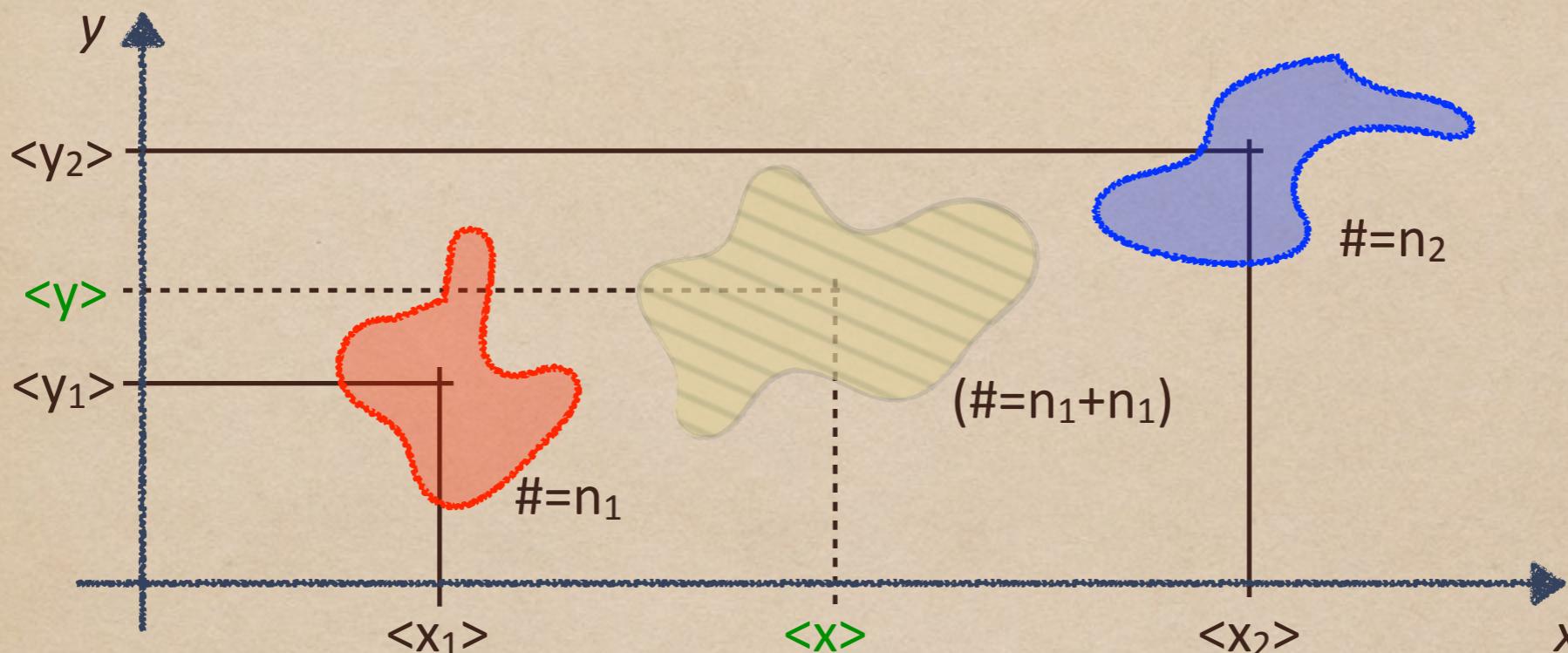
CLIC - A possible Compact Linear Collider

- Some parameters:

| | | | |
|---|----------------|----------------------|--------------------------------|
| Center of mass energy | E_{cm} | 3000 | GeV |
| Main Linac RF Frequency | f_{RF} | 11.994 | GHz |
| Luminosity | L | 5.9×10^{34} | $\text{cm}^{-2} \text{s}^{-1}$ |
| Linac repetition rate | f_{rep} | 50 | Hz |
| No. of particles / bunch | N_b | 3.72×10^9 | |
| No. of bunches / pulse | k_b | 312 | |
| Bunch separation | Δt_b | 0.5 (6 periods) | ns |
| Bunch train length | τ_{train} | 156 | ns |
| Beam power / beam | P_b | 14 | MW |
| Unloaded / loaded gradient | $G_{unl/l}$ | 120 / 100 | MV/m |
| Overall two linac length | L_{linac} | 42.16 | km |
| Total site AC power | P_{tot} | 415 | MW |
| Length of PETS | L_{PETS} | 0.213 | m |
| Nominal output RF Power / PETS | P_{out} | 136 | MW |
| Wall plug -> RF efficiency | η_{ACRF} | 58.6 | % |
| RF -> drive beam efficiency | η_{bRF} | 93 | % |
| drive beam -> RF efficiency (HDS input) | η_{decRF} | 65 | % |
| RF -> main beam efficiency | η_{bRF} | 27.7 | % |
| Wall plug to main beam power efficiency | η_{tot} | 7 | % |

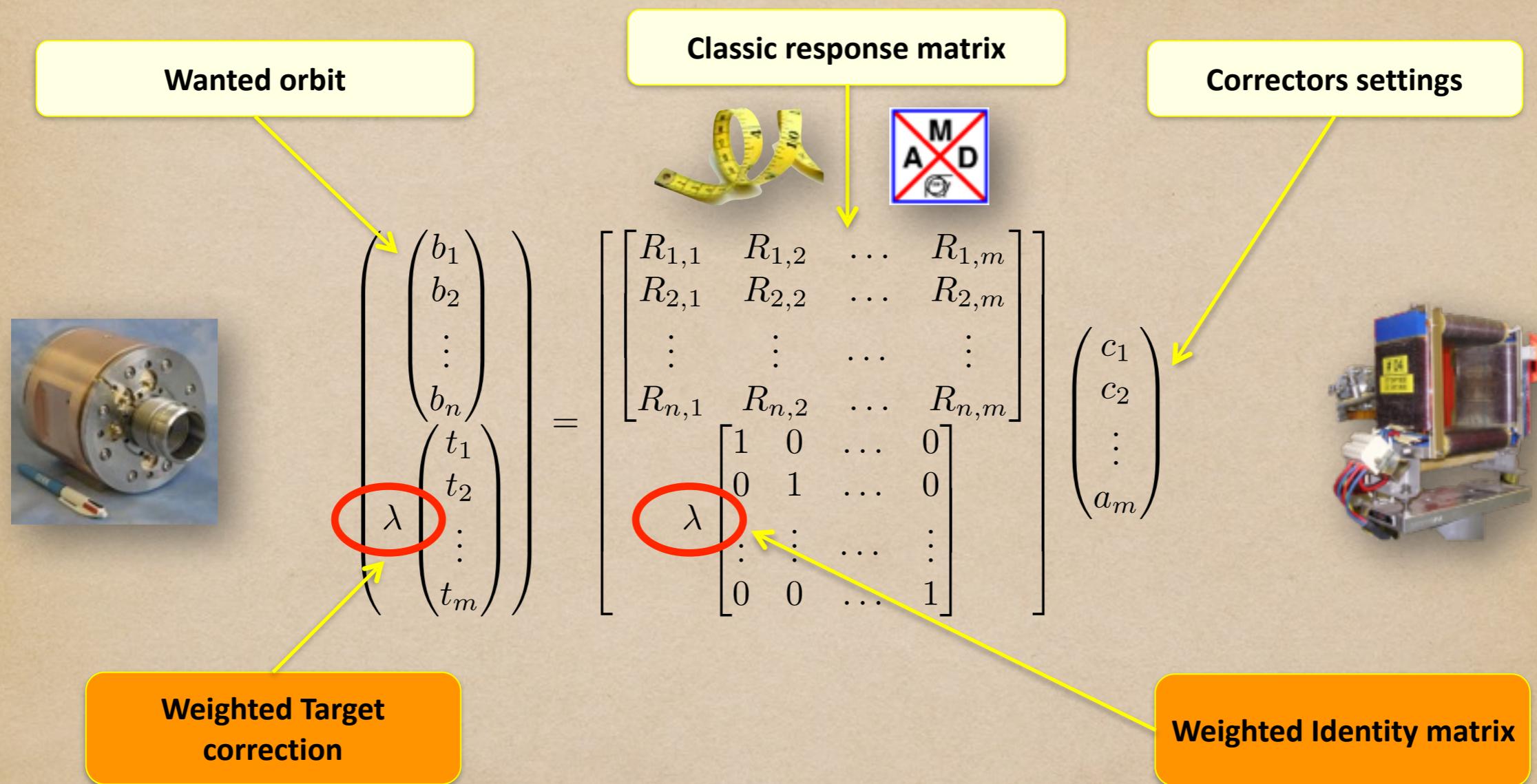
Simulation method: few details.

- Combine covariance matrices at one location:



$$\left. \begin{array}{l} \Sigma_1 = \begin{bmatrix} \sigma(x_1, x_1) & \sigma(x_1, y_1) \\ \sigma(y_1, x_1) & \sigma(y_1, y_1) \end{bmatrix} \\ \Sigma_2 = \begin{bmatrix} \sigma(x_2, x_2) & \sigma(x_2, y_2) \\ \sigma(y_2, x_2) & \sigma(y_2, y_2) \end{bmatrix} \end{array} \right\} \left. \begin{array}{l} \sigma(x_1, y_1) = \langle x_1 y_1 \rangle - \langle x_1 \rangle \langle y_1 \rangle \\ \sigma(x_2, y_2) = \langle x_2 y_2 \rangle - \langle x_2 \rangle \langle y_2 \rangle \\ \sigma(x, y) = \langle xy \rangle - \langle x \rangle \langle y \rangle \\ \langle x \rangle = \frac{n_1 \langle x_1 \rangle + n_2 \langle x_2 \rangle}{n_1 + n_2} \\ \langle y \rangle = \frac{n_1 \langle y_1 \rangle + n_2 \langle y_2 \rangle}{n_1 + n_2} \\ \langle xy \rangle = \frac{n_1 (\sigma(x_1, y_1) + \langle x_1 \rangle \langle y_1 \rangle) + n_2 (\sigma(x_2, y_2) + \langle x_2 \rangle \langle y_2 \rangle)}{n_1 + n_2} \end{array} \right\} \Sigma = \begin{bmatrix} \sigma(x, x) & \sigma(x, y) \\ \sigma(y, x) & \sigma(y, y) \end{bmatrix}$$

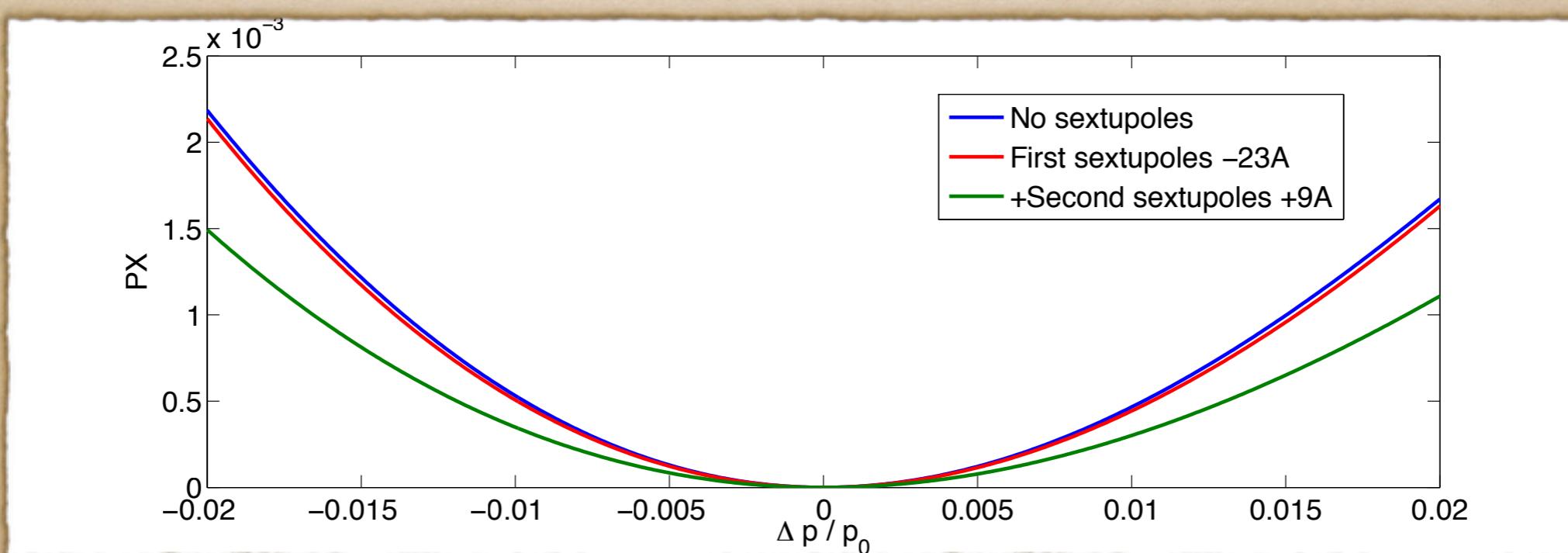
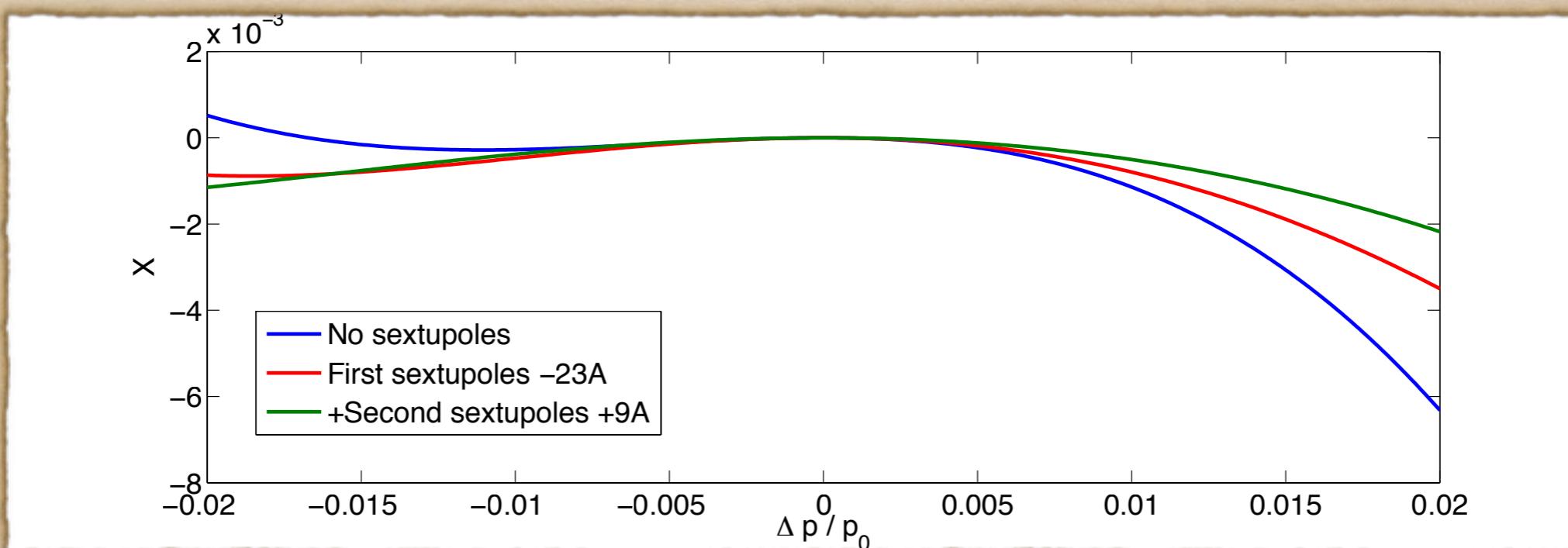
Feedback tool: what is behind



- If $\lambda = 0$, we are back in the classical situation.
- If $\lambda \neq 0$, then the full Response Matrix is **always over-constrained**.
 - Solution cannot explode.
- If $\lambda \gg 0$, then the classic response matrix is just “noise”. Solution will be identical to the target correction required.

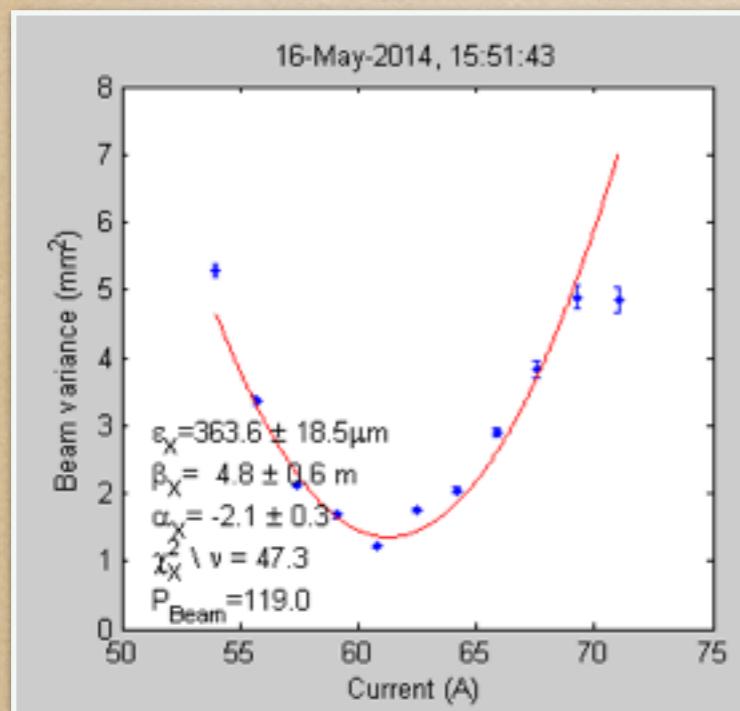
Experimental test: sextupolar dispersion correction

- Simulated end of DL X/PX as a function of DELTAP (MAD-X simulations)

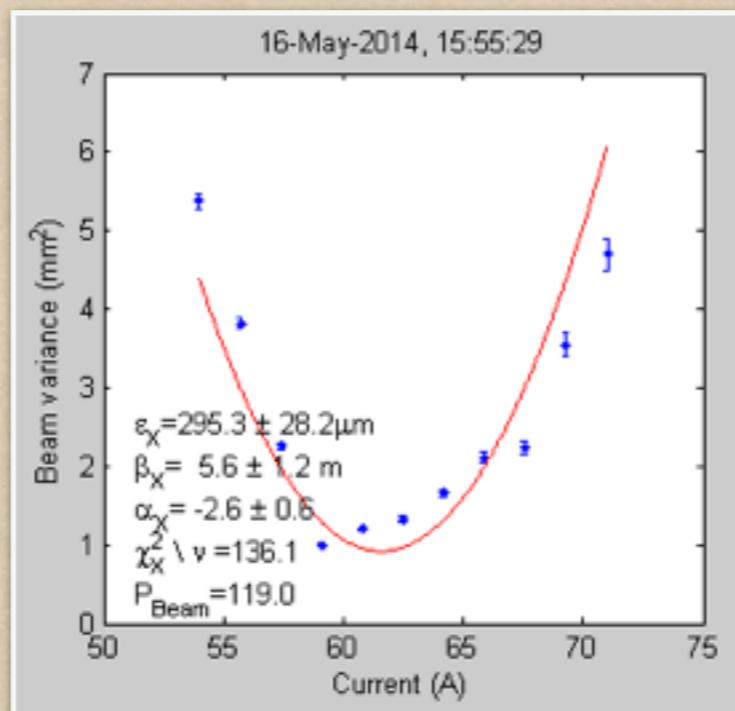


Experimental results

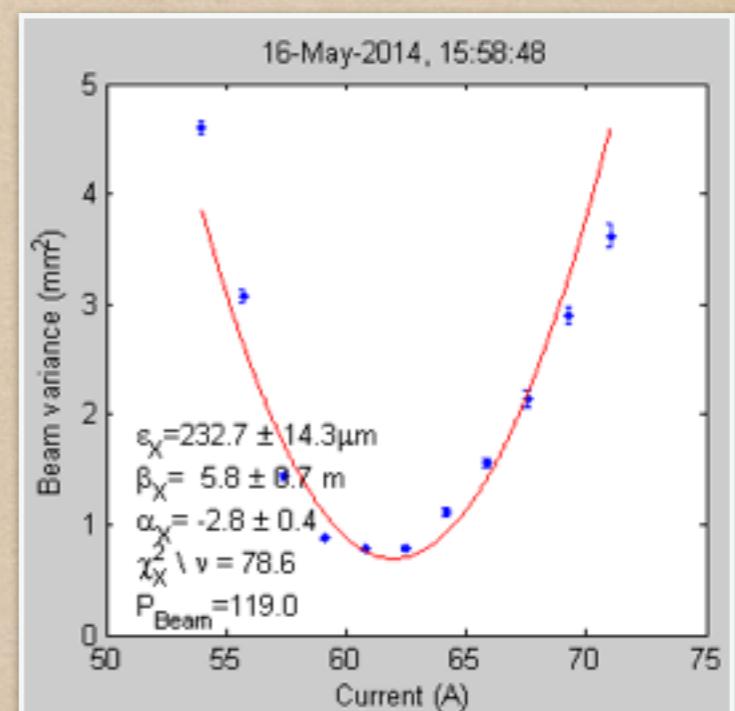
- Implementing sextupolar correction of second-order dispersion.
- Measuring emittance in downstream measurement line at different stages.



No sextupoles on.



First family at -23 A



First family at -23 A
Second family at +9 A



Thanks to P. Skowronski
and T. Persson