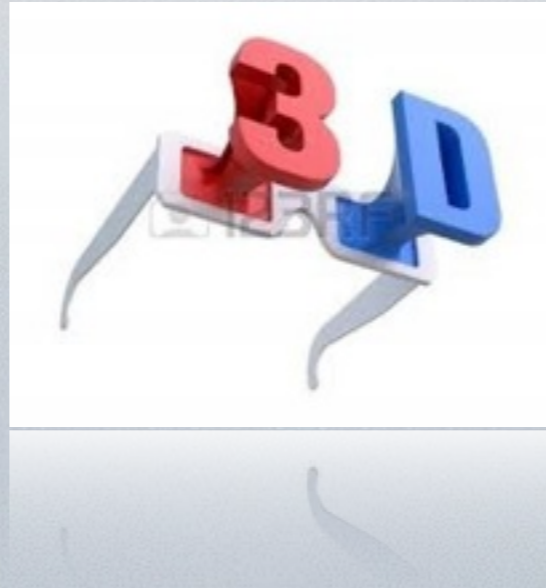


Nucleon



Structure theory

Marco Radici

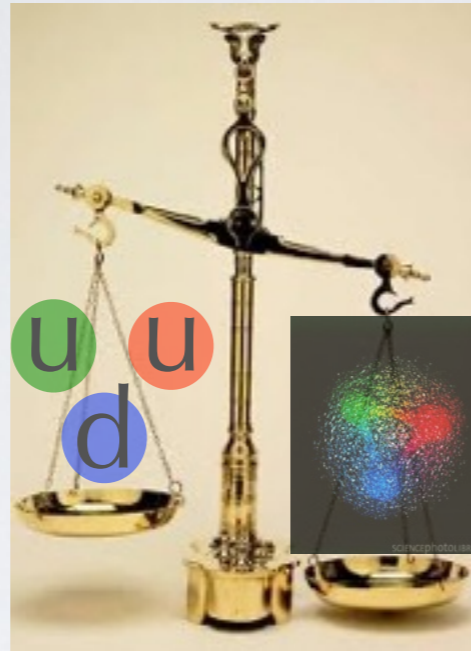


Pavia

understand the proton

quark-Higgs
coupling

$\sim 9 \text{ MeV}$

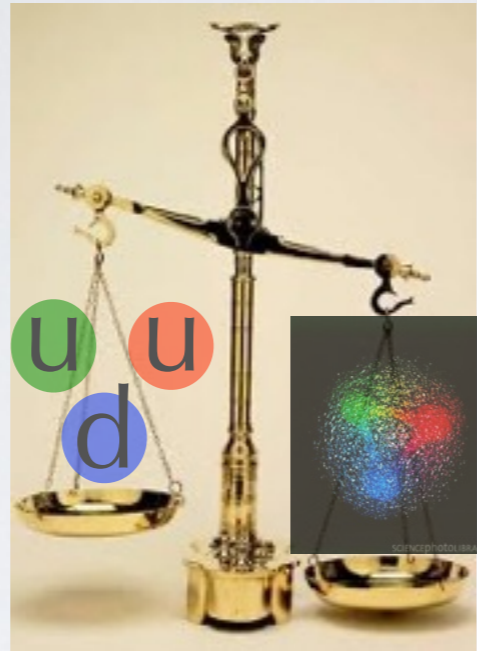


938 MeV

understand the proton

quark-Higgs
coupling

$\sim 9 \text{ MeV}$



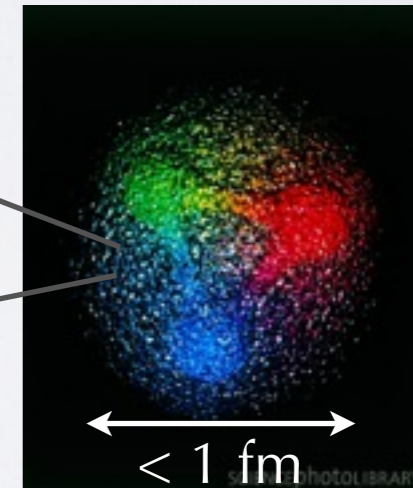
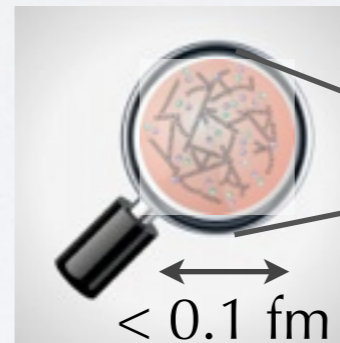
99% of proton mass is
generated by dynamics of
QCD confinement

938 MeV



lattice QCD

&

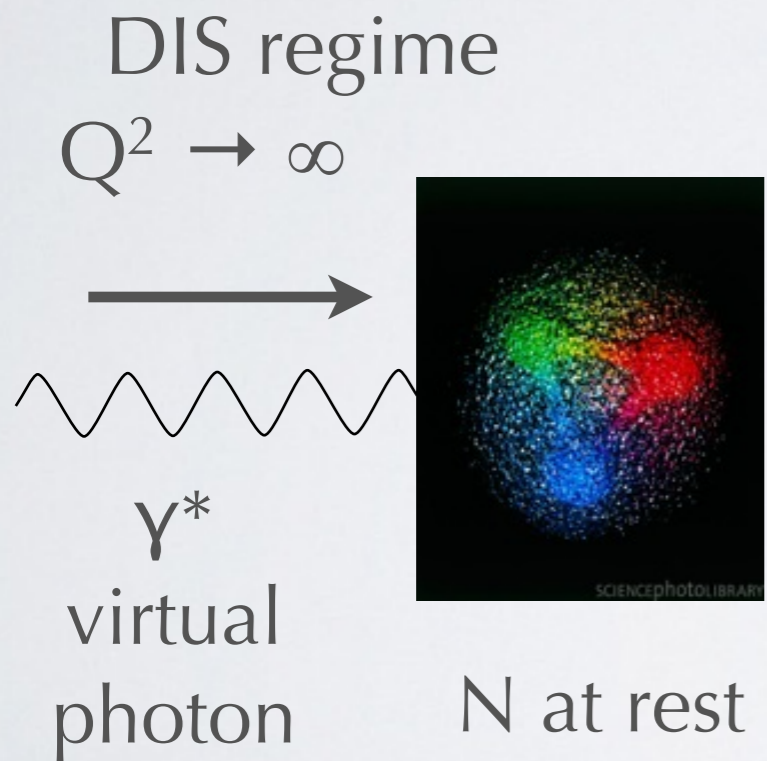
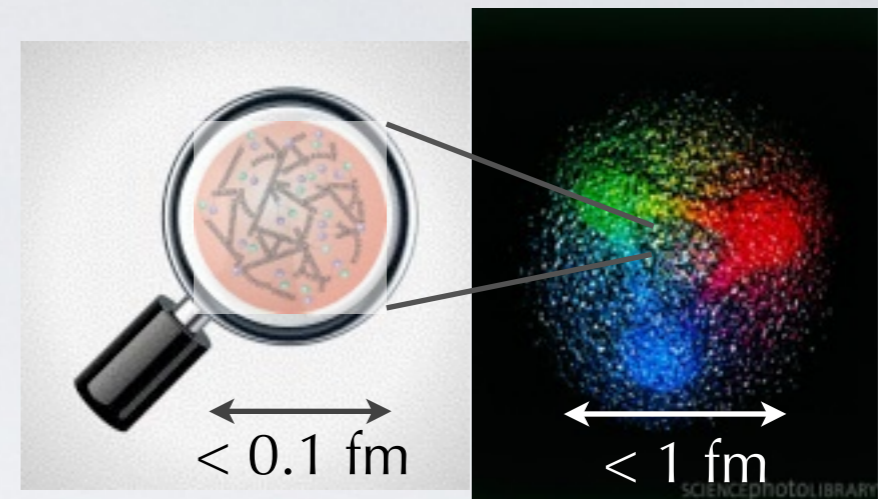


Hadron Physics

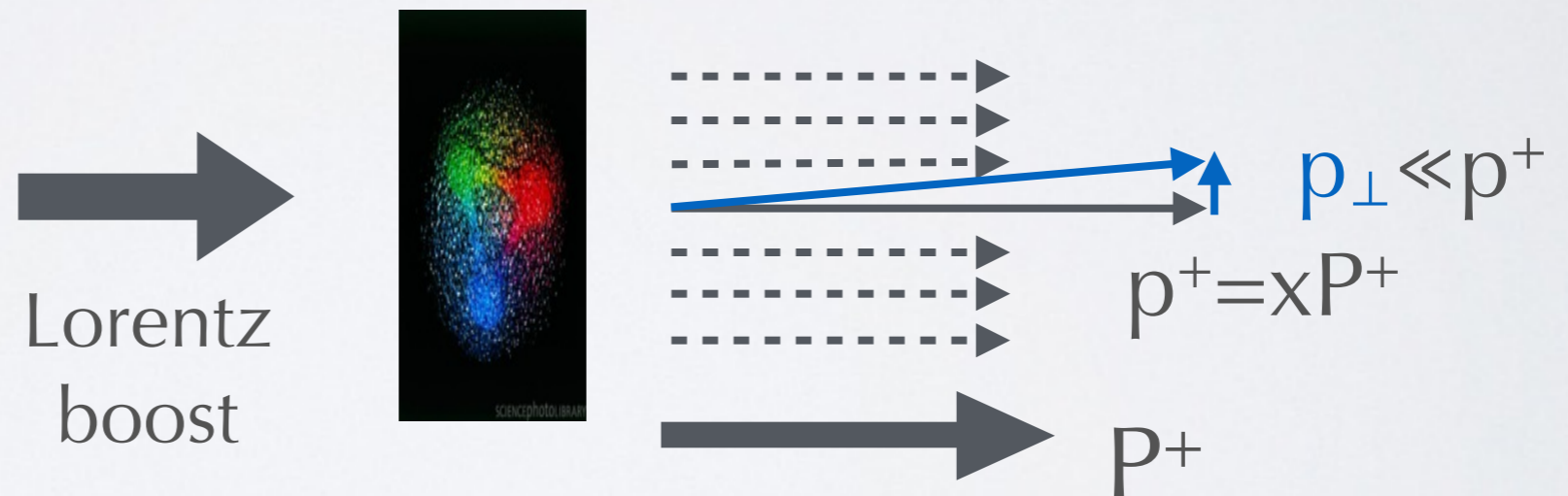
this talk

the Infinite Momentum Frame (IMF)

probe short distances
 \Rightarrow Deep-Inelastic (DIS) regime

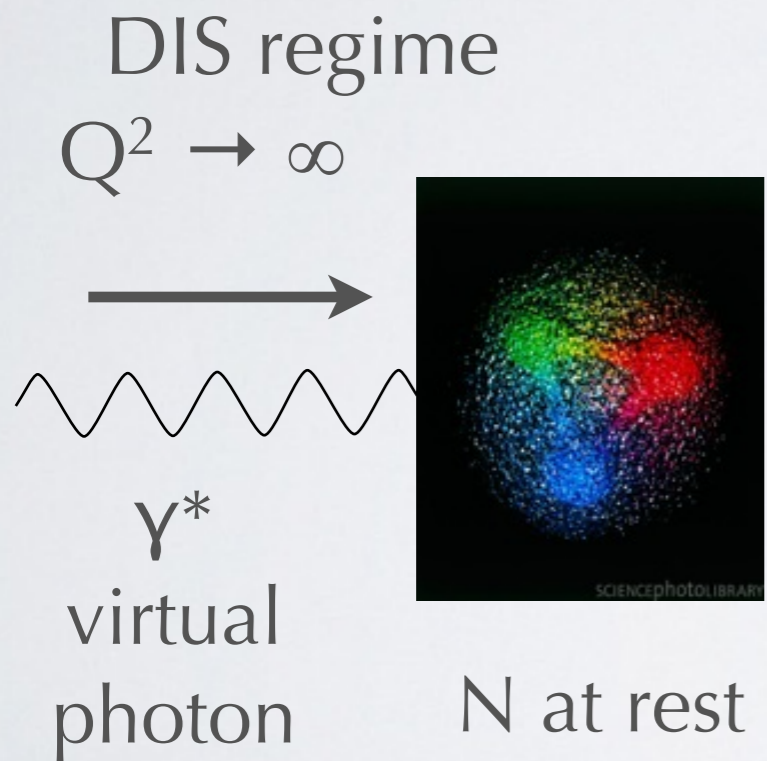
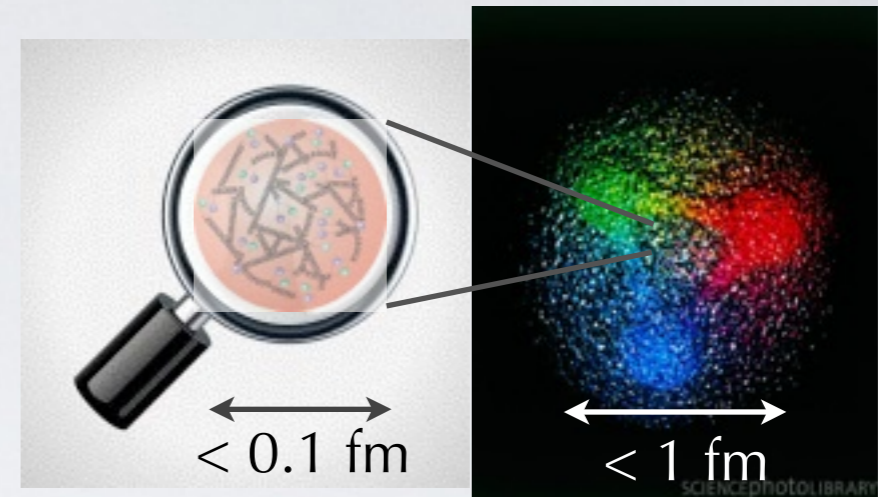


IMF \Leftrightarrow Light-Cone (LC) kin.



the Infinite Momentum Frame (IMF)

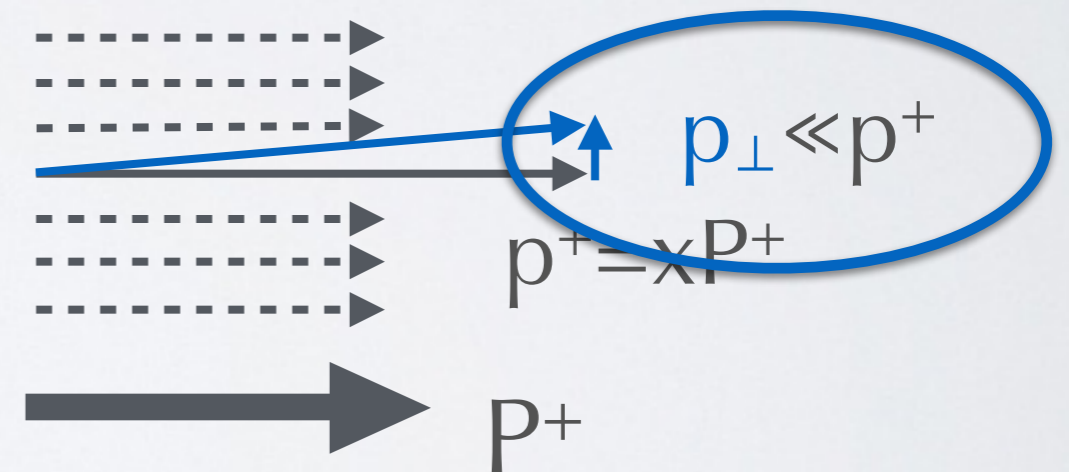
probe short distances
 \Rightarrow Deep-Inelastic (DIS) regime



\rightarrow
 Lorentz
 boost



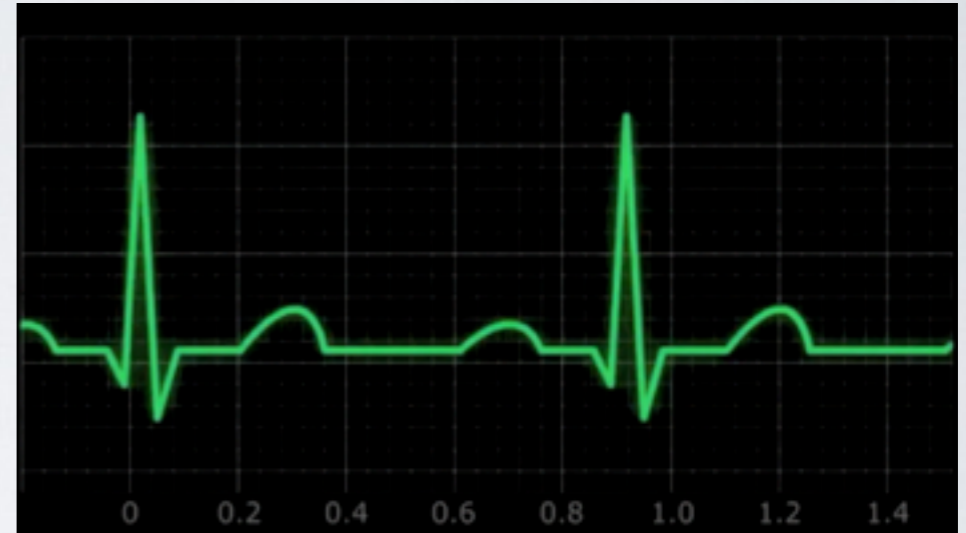
IMF \Leftrightarrow Light-Cone (LC) kin.



all partons \sim collinear
go beyond this approx.

main goal

the 3D-structure
of the Nucleon



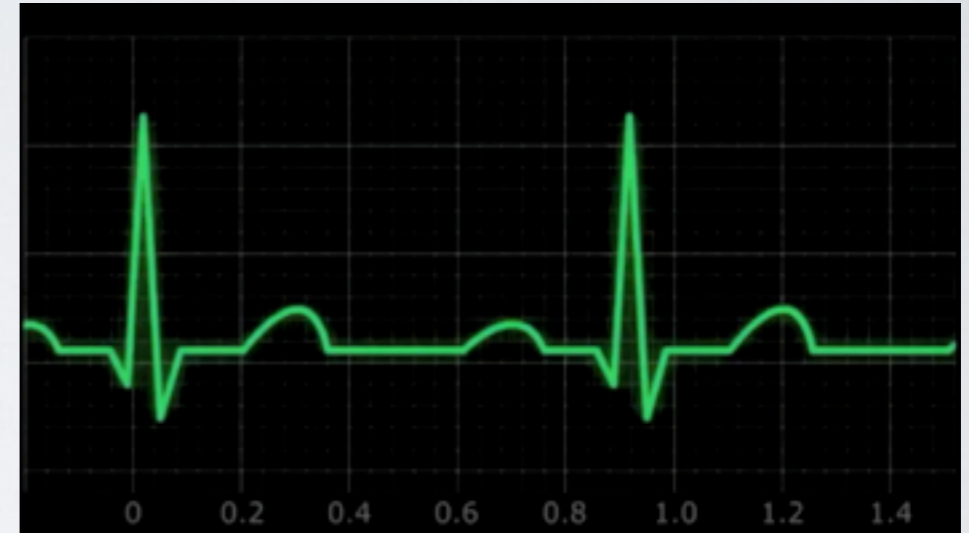
main goal

the 3D-structure
of the Nucleon



mono-dim. info
on heart activity

ECG



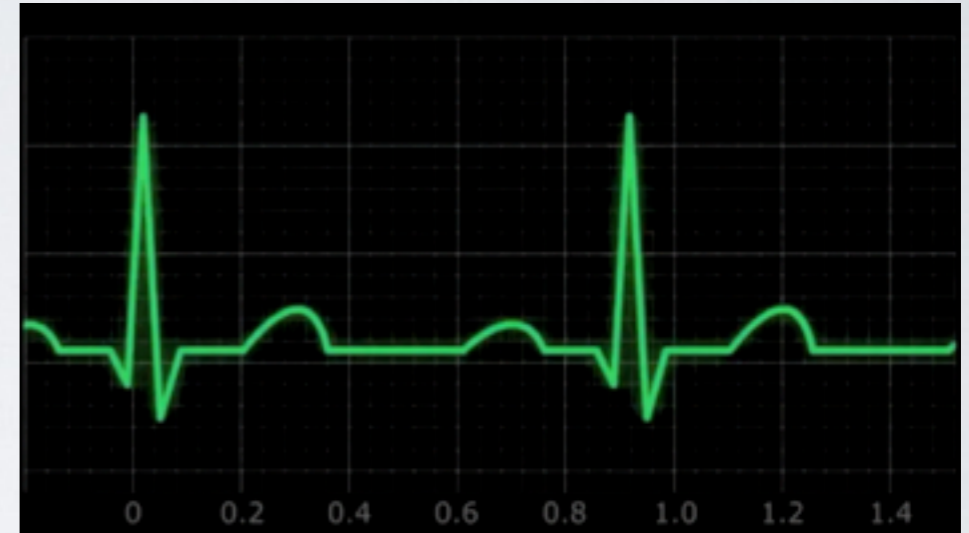
main goal

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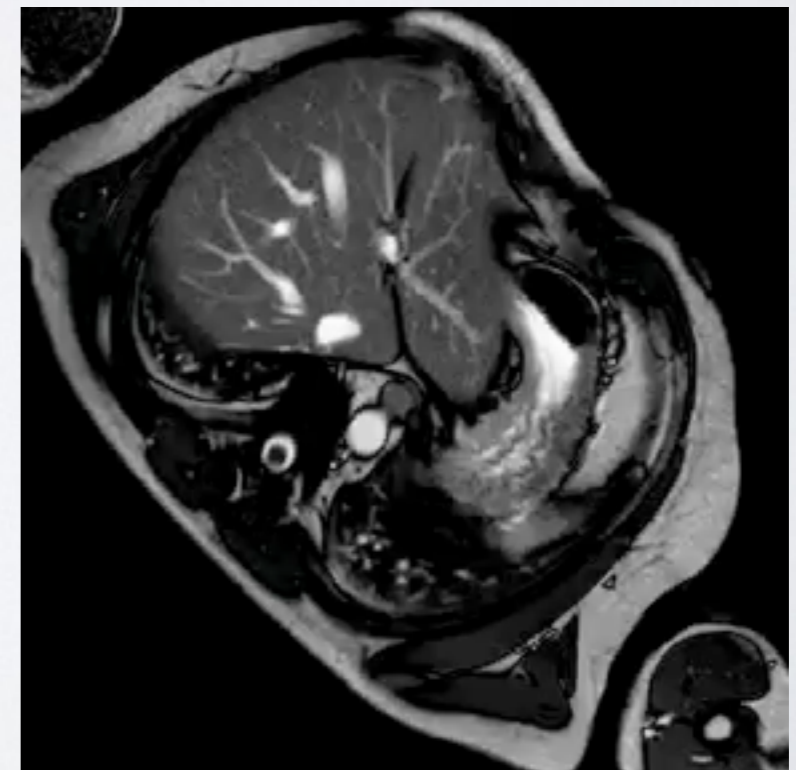
mono-dim. info
on heart activity

ECG



3-dim. tomography
of heart activity

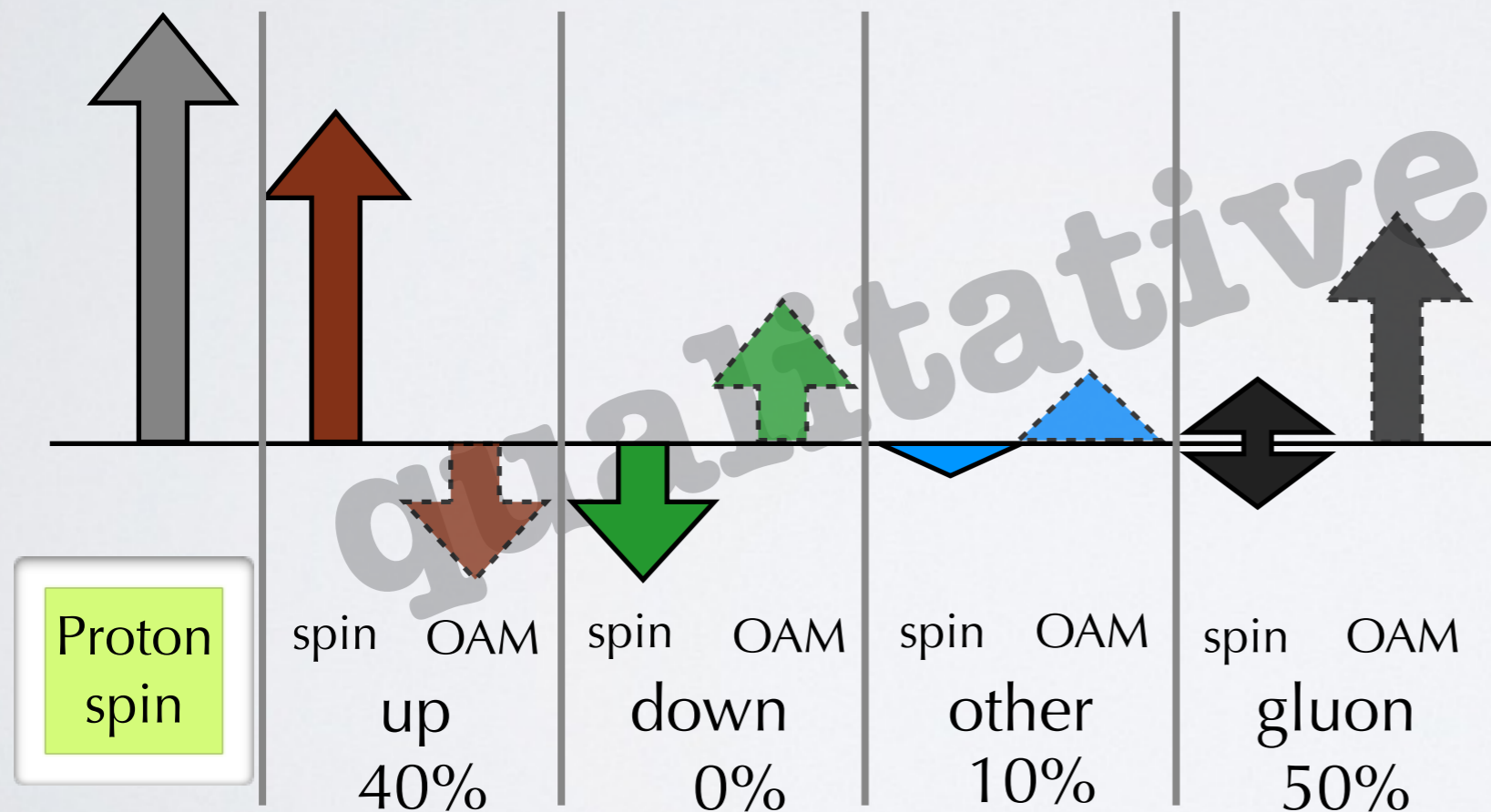
cardio
MR



the proton spin budget ?

since EMC (1988, the “spin crisis”)
we can't yet explain the proton spin
in terms of its constituents

$$1/2 = 1/2 \Delta\Sigma + \text{OAM}_q + \Delta g + \text{OAM}_g$$

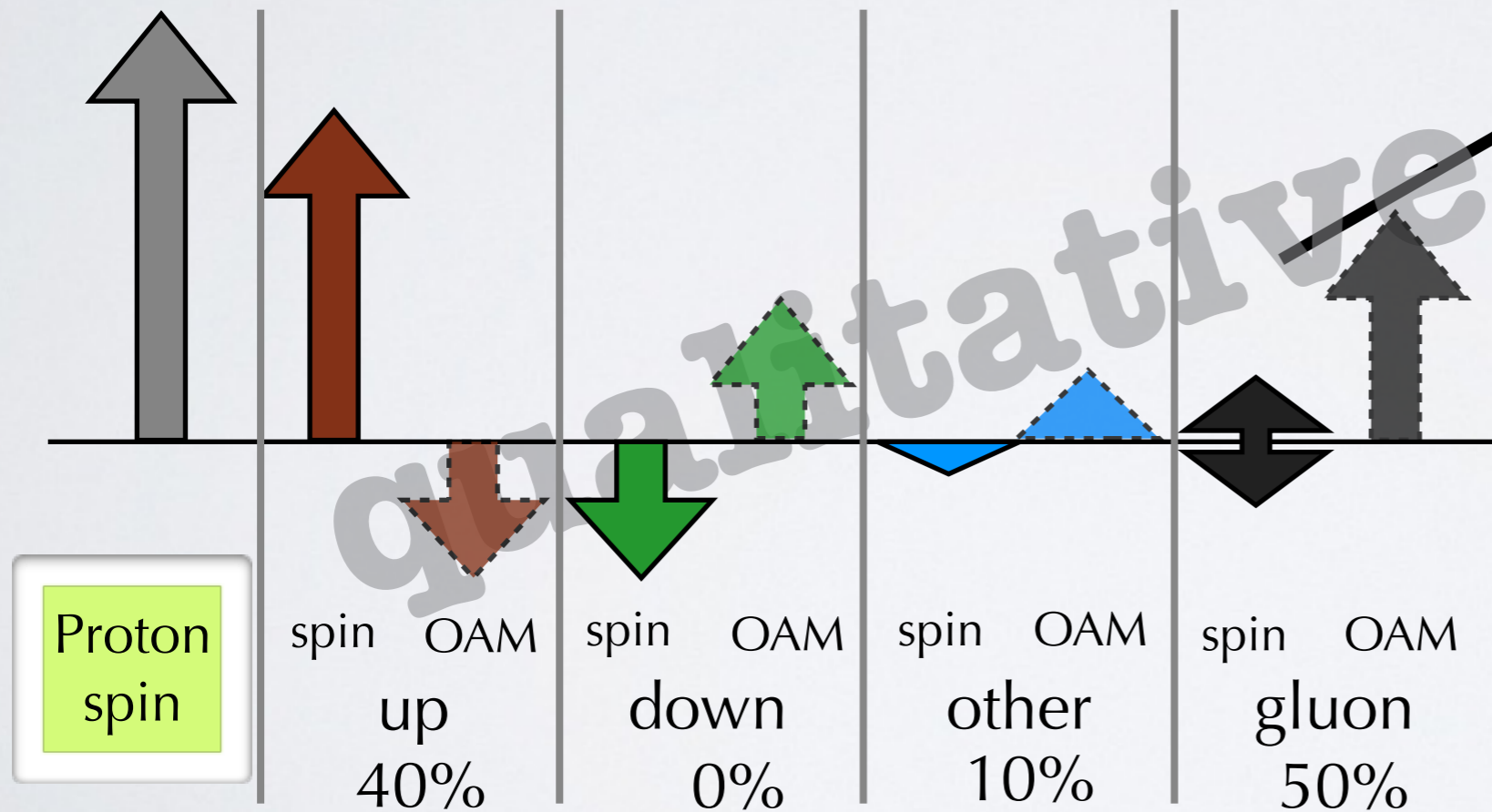


OAM = Orbital Angular Momentum

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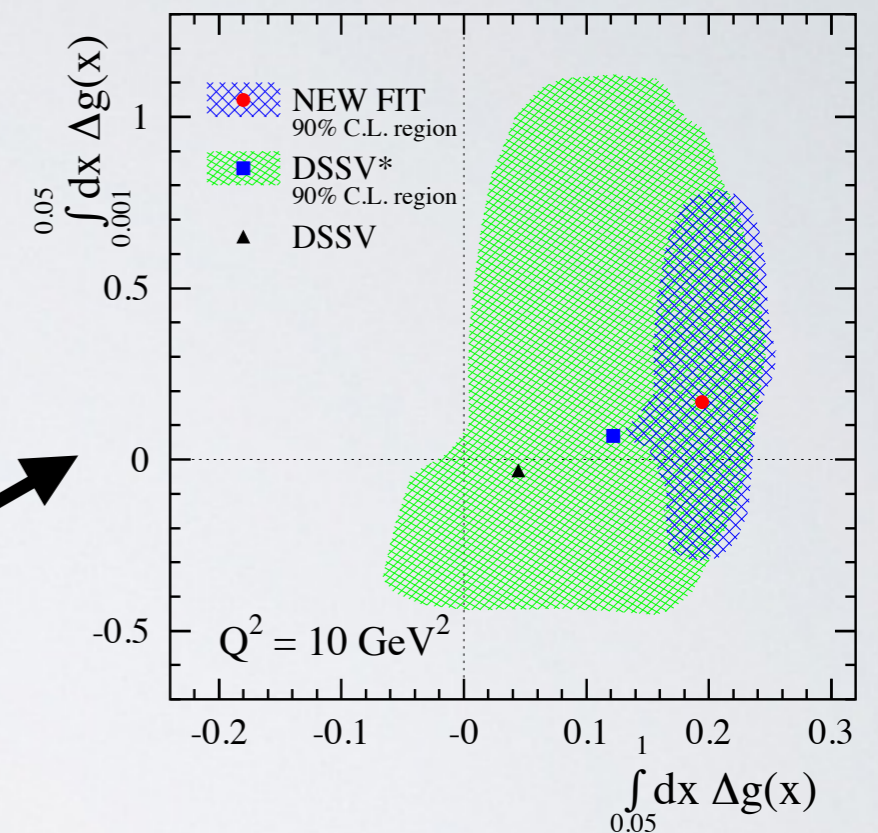
$$1/2 = 1/2 \Delta\Sigma + \text{OAM}_q + \Delta g + \text{OAM}_g$$



OAM = Orbital Angular Momentum

low x

*De Florian et al.,
arXiv:1404.4293*



valence

we don't even know
the gluon helicity
 $-0.15 \approx \Delta g \approx 1$

new tools needed

$$\bar{u}_{N'} \gamma^+ u_N F_1(t) + \bar{u}_{N'} \frac{i\sigma^{+\nu} \Delta_\nu}{2M} u_N F_2(t)$$

generalize to m-index operator

X. Ji, P.R.L. **78** (97) 610

the Ji's sum rule

$$J_i = \varepsilon_{ijk} (x_j T^{0k} - x_k T^{0j})$$

QCD
energy-momentum
tensor

$$J_z^q(Q^2) = \frac{1}{2} \int_0^1 dx x [H^q(x, 0, 0; Q^2) + E^q(x, 0, 0; Q^2)]$$

total angular
momentum
of parton q

Generalized Parton Distributions
GPD(x, ξ , t; Q^2)

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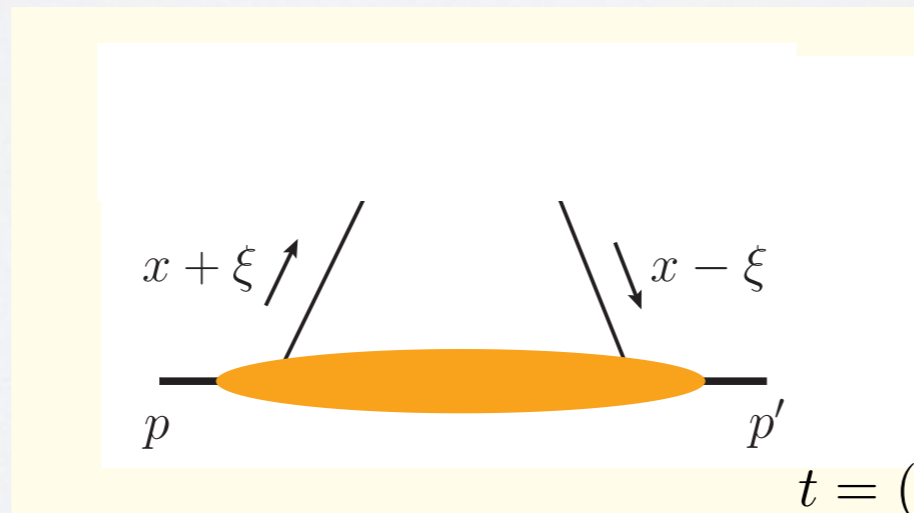
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Generalized Parton Distributions

GPD(x, ξ, t; Q²)



$$\xi = \frac{(P - P')^+}{(P + P')^+}$$

change in N
long. momentum

$$t = (P' - P)^2 = \Delta^2$$

non-diagonal (P' ≠ P) hadronic matrix element

new tools needed

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X. Ji, P.R.L. **78** (97) 610

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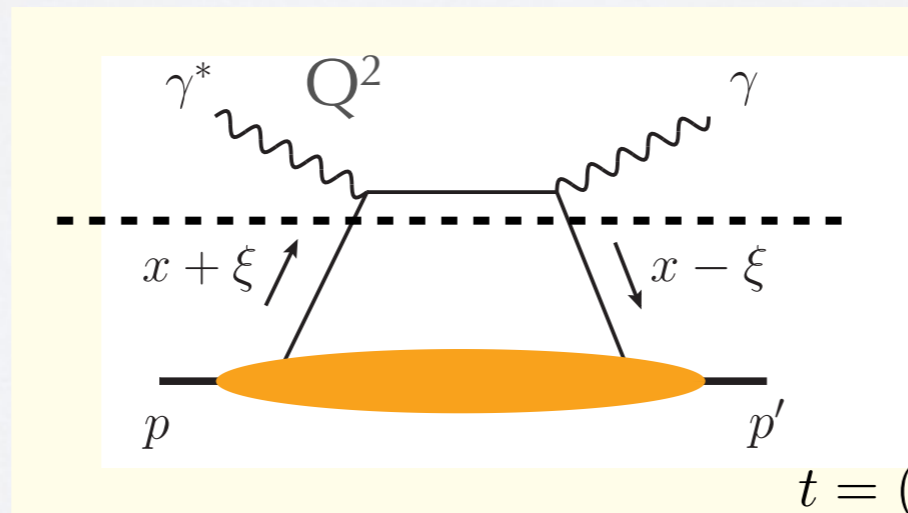
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momentum
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Generalized Parton Distributions

GPD(x, ξ , t; Q²)

DVCS
factoriz. th.
for Q² ≫ t



$$\xi = \frac{(P - P')^+}{(P + P')^+}$$

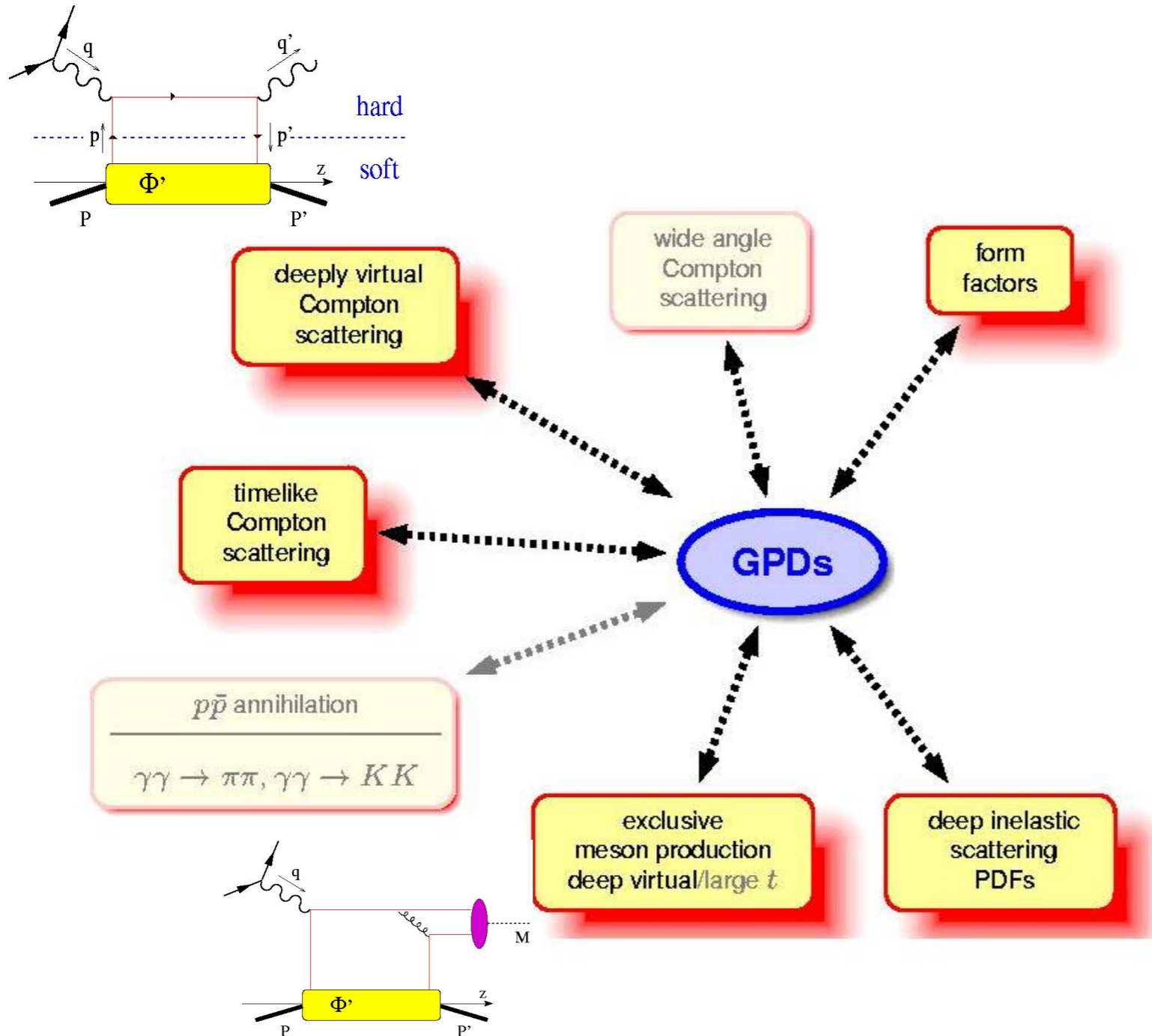
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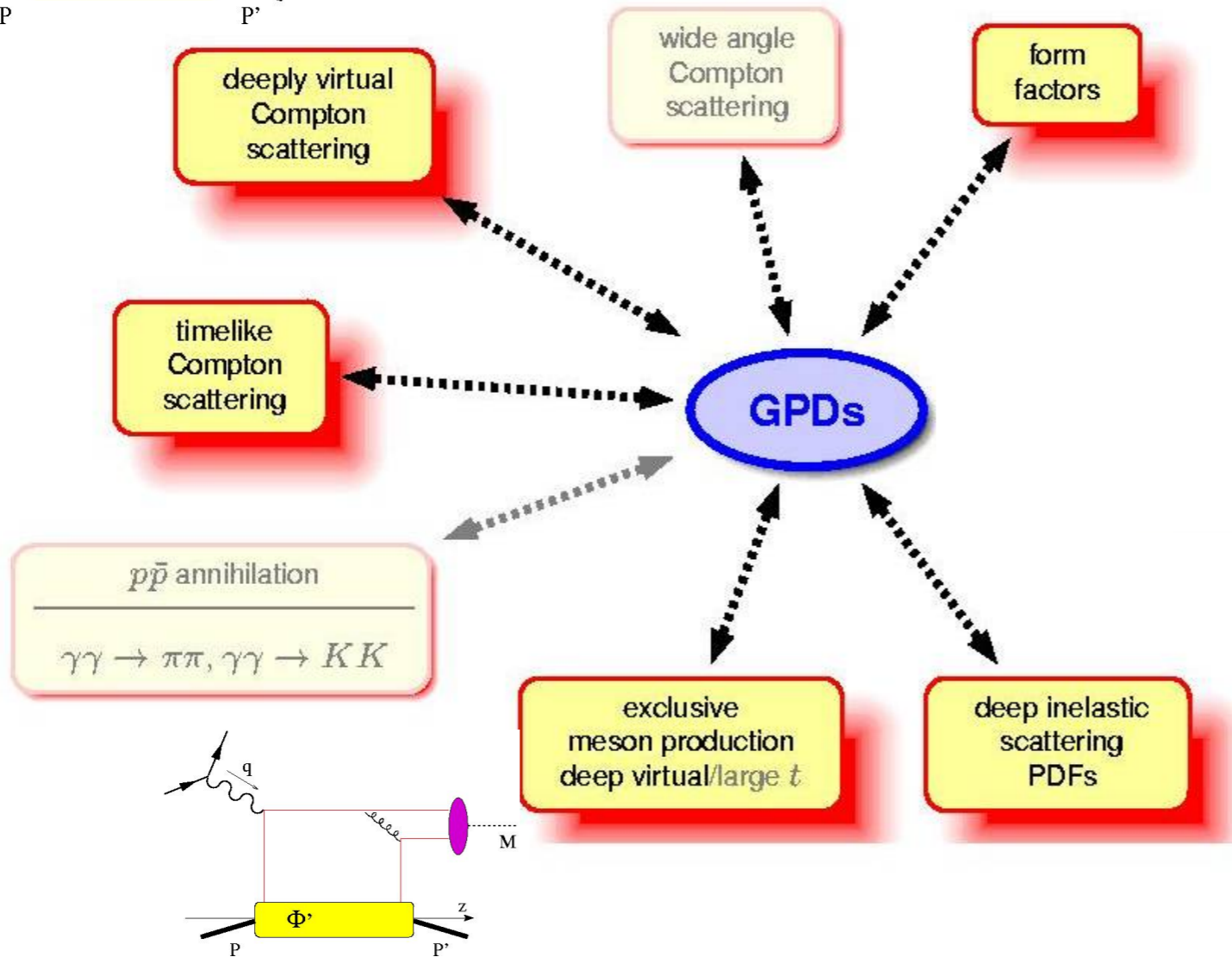
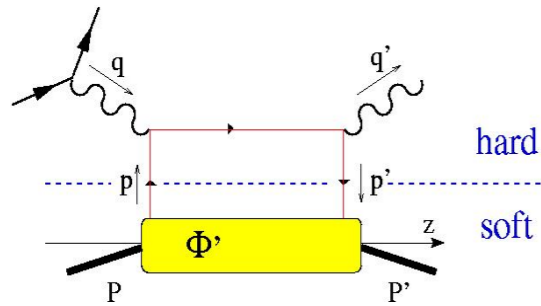
the GPD

$$\text{GPD}(x, \xi, t; Q^2)$$

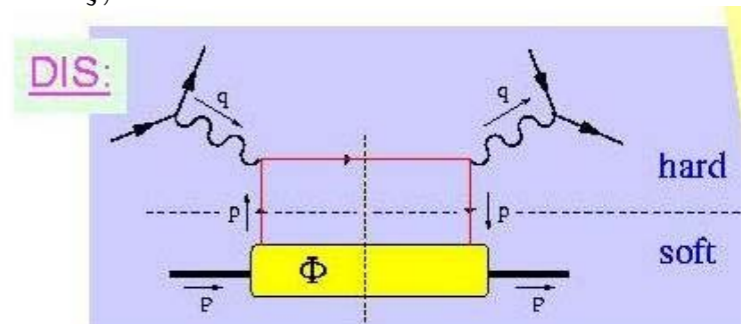


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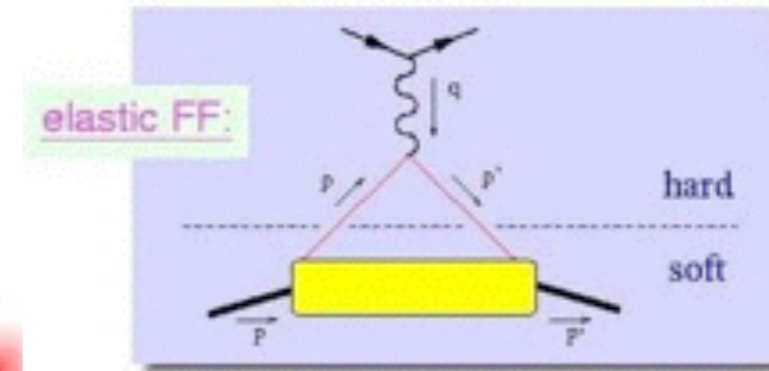
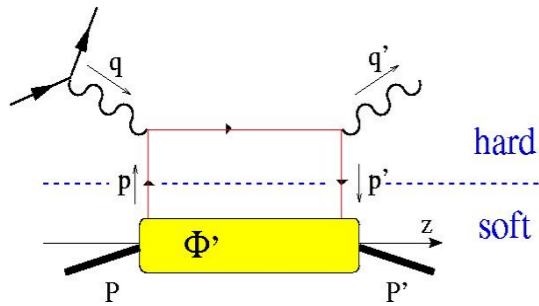


$$\lim_{\xi, t \rightarrow 0} \text{GPD}(x, \xi, t) = \text{PDF}(x)$$

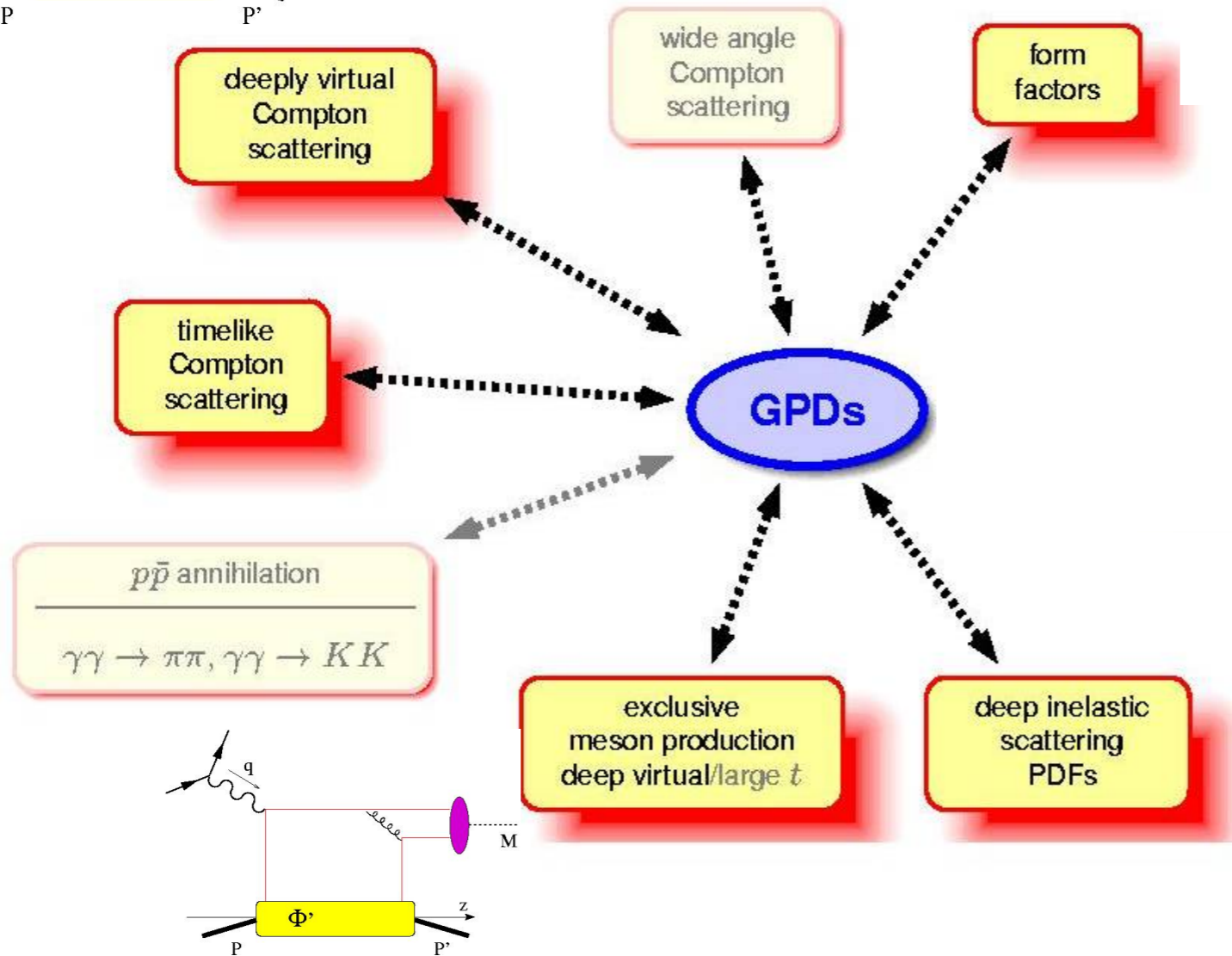


the GPD

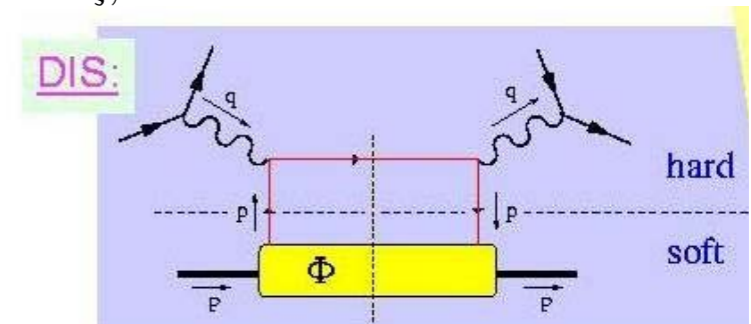
$$\text{GPD}(x, \xi, t; Q^2)$$



$$\int dx \text{GPD}(x, \xi, t) = \text{FF}(t)$$



$$\lim_{\xi, t \rightarrow 0} \text{GPD}(x, \xi, t) = \text{PDF}(x)$$



the GPD

$$\text{GPD}(x, \xi, t; Q^2)$$

$$\lim_{\xi, t \rightarrow 0} \text{GPD}(x, \xi, t) = \text{PDF}(x)$$

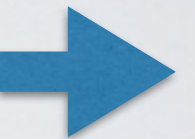
$$H^q(x, \xi \rightarrow 0, t \rightarrow 0) \Rightarrow f_1^q(x)$$

not directly accessible

($E^q \rightarrow N$ spin flip)

need model extrapolation

$$J_z^q = \frac{1}{2} \int dx x [H^q(x, 0, 0) + E^q(x, 0, 0)]$$



the GPD

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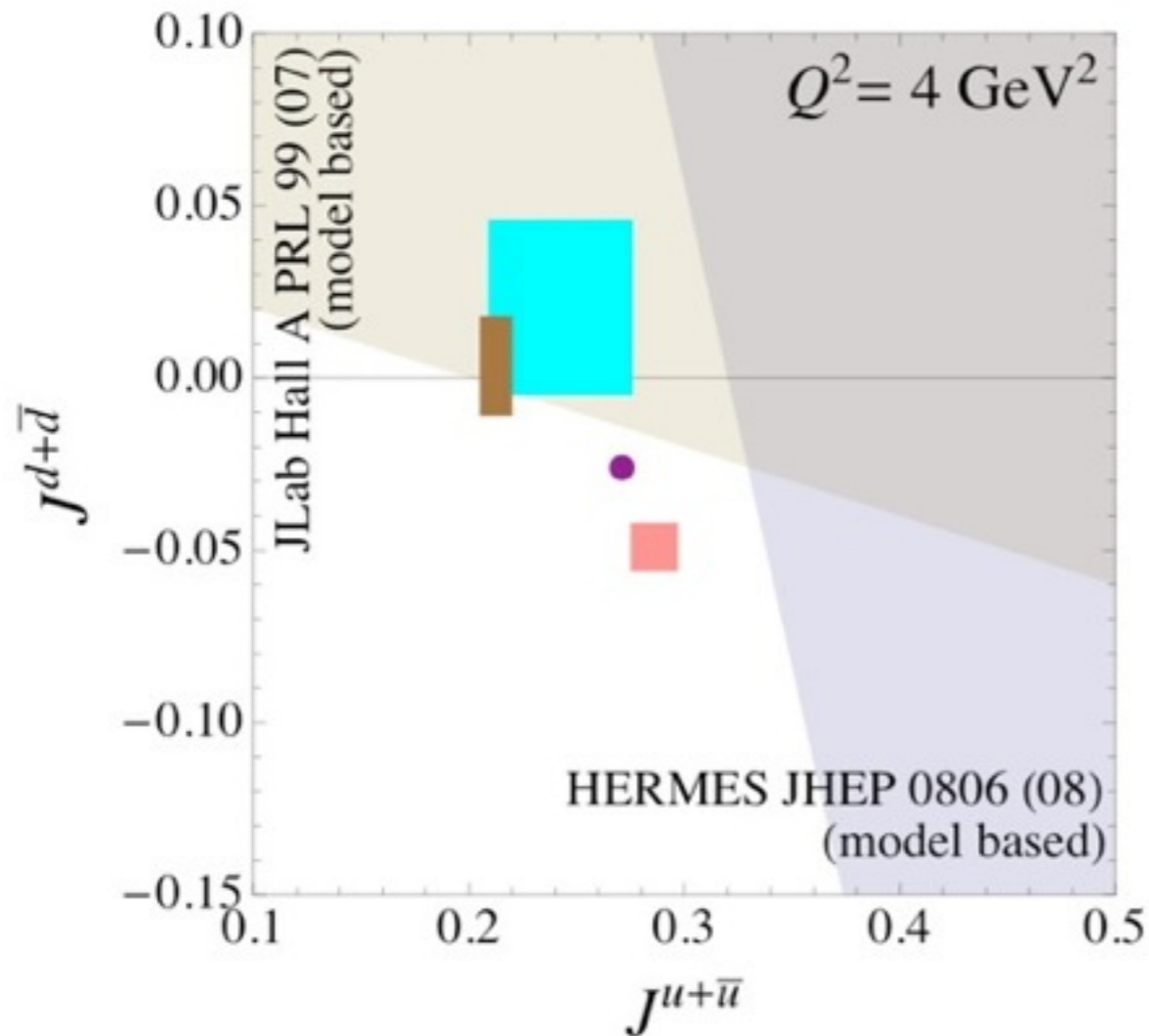
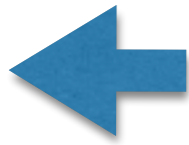
$$\begin{aligned} J_z^q &= \frac{1}{2} \int dx x [H^q(x, 0, 0) + E^q(x, 0, 0)] \\ &= \frac{1}{2} [A_{2,0}^q(0) + B_{2,0}^q(0)] \end{aligned}$$

moments of GPD

Generalized Form Factors calculable on lattice

$$A_{1,0} (\equiv F_1), B_{1,0} (\equiv F_2), \quad \mathbf{A_{2,0}}, \mathbf{B_{2,0}}, \quad A_{3,0}, A_{3,2} \dots B_{3,0}, B_{3,2} \dots$$

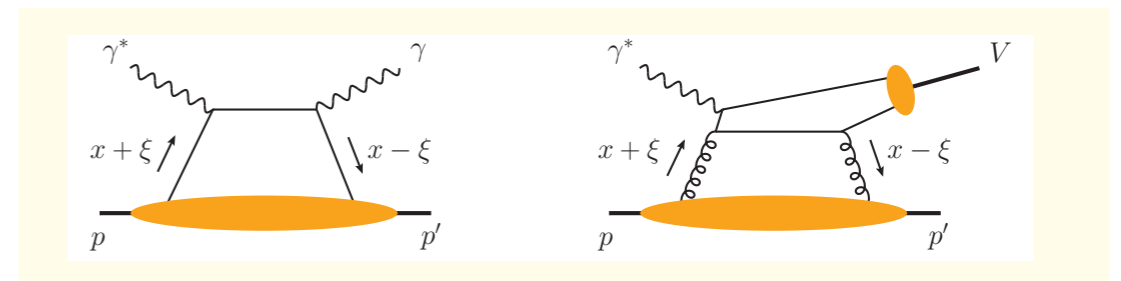
J^q results (model) params. of GPD



- Goloskokov & Kroll, EPJ C59 (09) 809
- Diehl et al., EPJ C39 (05) 1
- Guidal et al., PR D72 (05) 054013
- Liuti et al., PRD 84 (11) 034007

DVCS

DVMP



GPD convoluted in **C**ompton **F**orm **F**actors

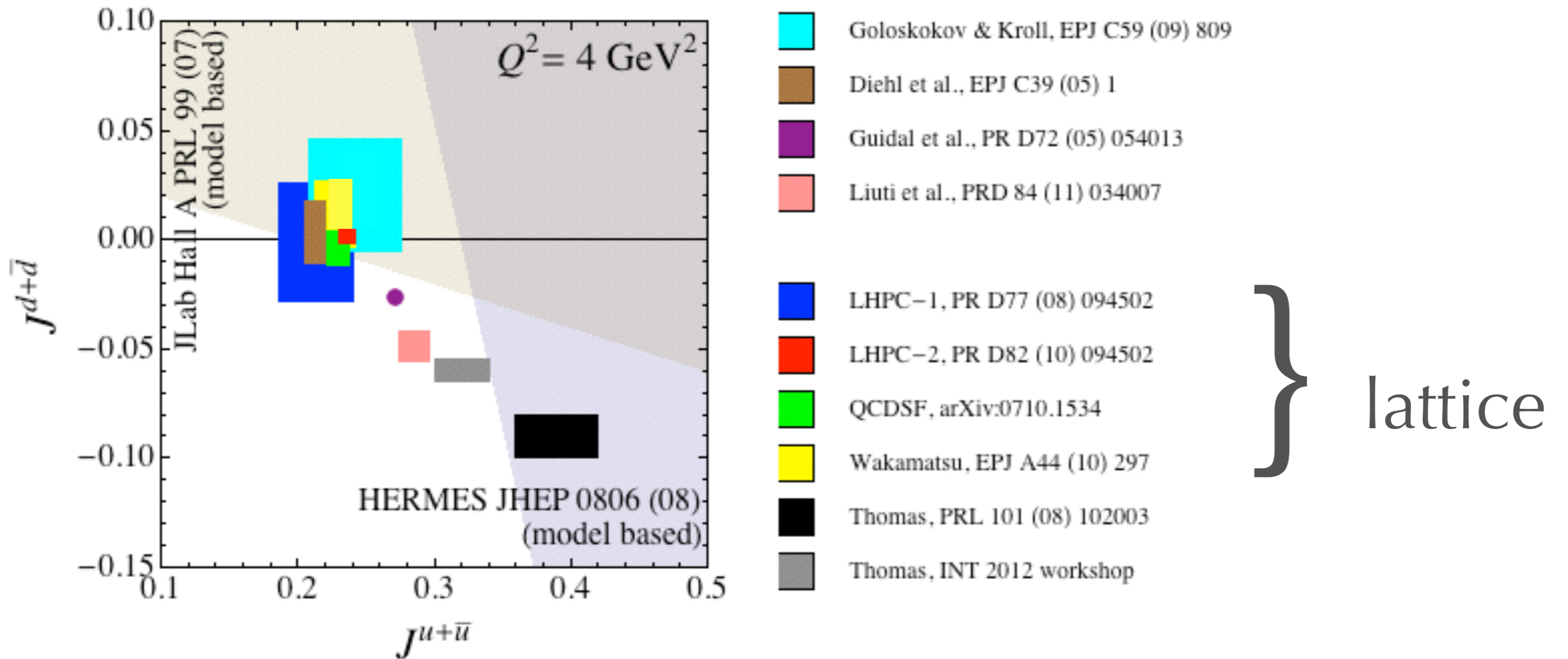
$$d\sigma \sim \sum_n C_n(\text{Re}[CFF], \text{Im}[CFF]) \cos(n\Phi) + S_n(\text{Re}[CFF], \text{Im}[CFF]) \sin(n\Phi)$$

Φ azimuthal angle of γ / M

A. Bacchetta & M. Radici, arXiv:1206.2565 [hep-ph]

“Physics Opportunities with the 12 GeV Upgrade at Jefferson Lab”, E.P.J. A48 (12) 187

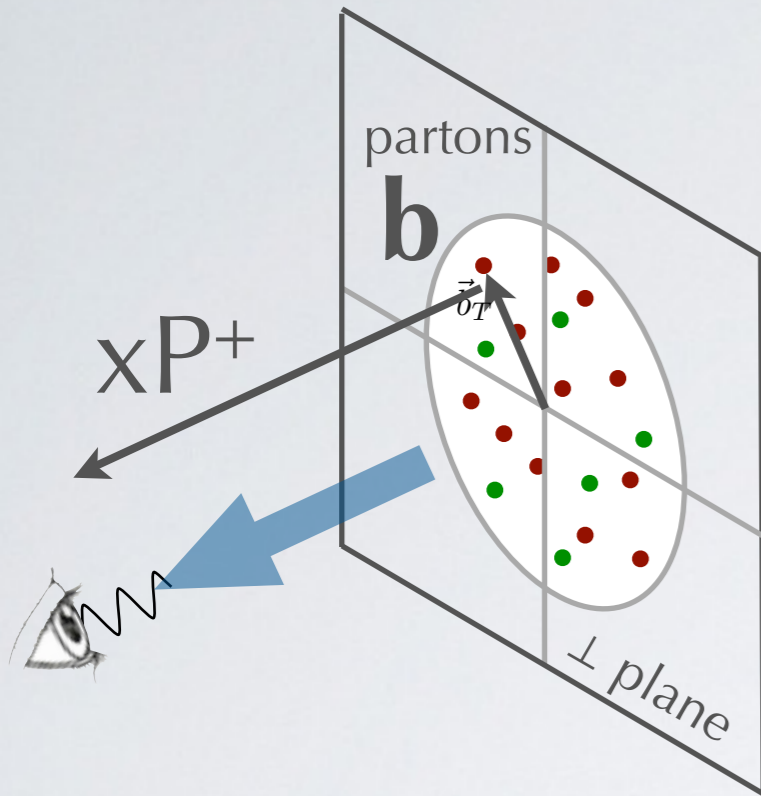
J^q results compare with lattice QCD



A. Bacchetta & M. Radici, arXiv:1206.2565 [hep-ph]

“Physics Opportunities with the 12 GeV Upgrade at Jefferson Lab”, E.P.J. A**48** (12) 187

tomography of the Nucleon



GPD limit : $\xi \rightarrow 0$ ($P^+ = P'^+$) ; $t \rightarrow -(\mathbf{P}'_{\perp} - \mathbf{P}_{\perp})^2 = -\mathbf{q}^2$

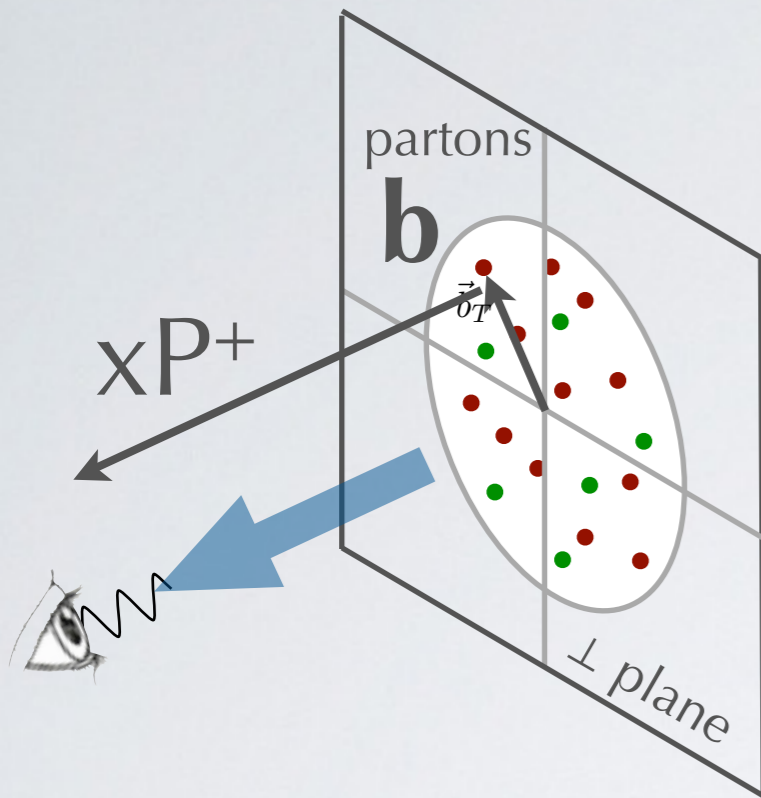
$$q(x, \mathbf{b}) = \int \frac{d\mathbf{q}}{(2\pi)^2} e^{i\mathbf{q} \cdot \mathbf{b}} H(x, 0, t = -\mathbf{q}^2)$$

$q(x, \mathbf{b})$ is a density in $\mathbf{b} \leftrightarrow \mathbf{q} = \mathbf{P}'_{\perp} - \mathbf{P}_{\perp}$

density of partons with momentum x
and position \mathbf{b}

tomography of N

tomography of the Nucleon



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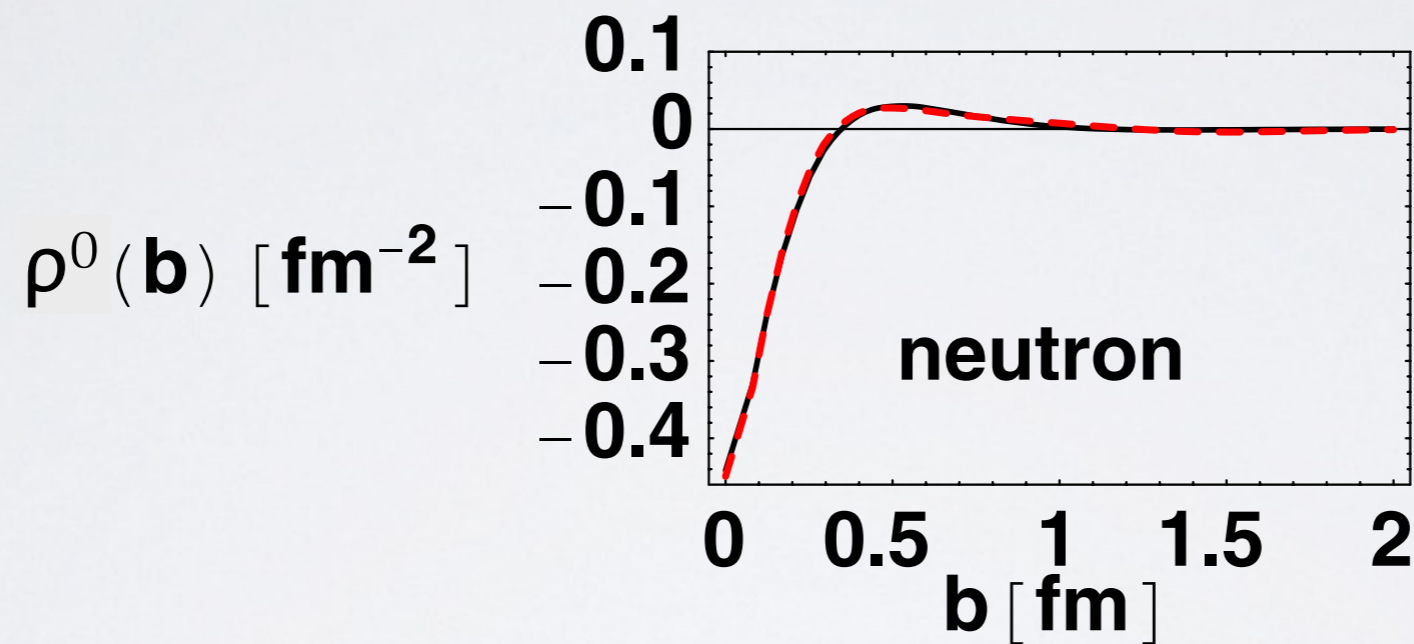
valid for all $x \Rightarrow$

$$\begin{aligned} \rho^0(\mathbf{b}) &= \int dx \int \frac{d\mathbf{q}}{(2\pi)^2} e^{i\mathbf{q} \cdot \mathbf{b}} H(x, 0, t = -\mathbf{q}^2) \\ &= \int \frac{d\mathbf{q}}{(2\pi)^2} e^{i\mathbf{q} \cdot \mathbf{b}} F_1(t = -\mathbf{q}^2) \end{aligned}$$

Dirac form factor

revolutionize the neutron

inside neutron



*G.A. Miller,
PRL99 (07) 112001*

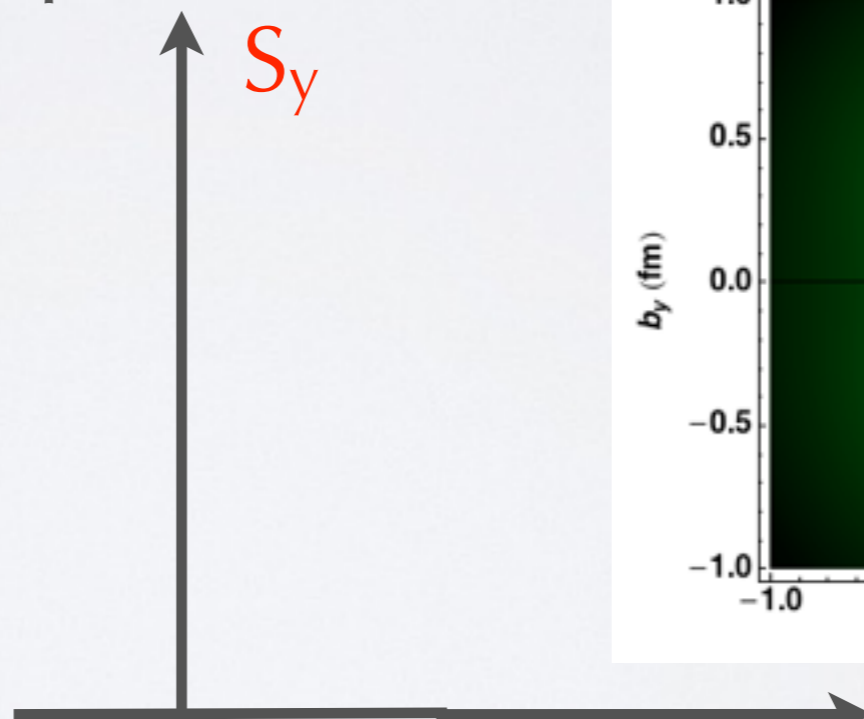
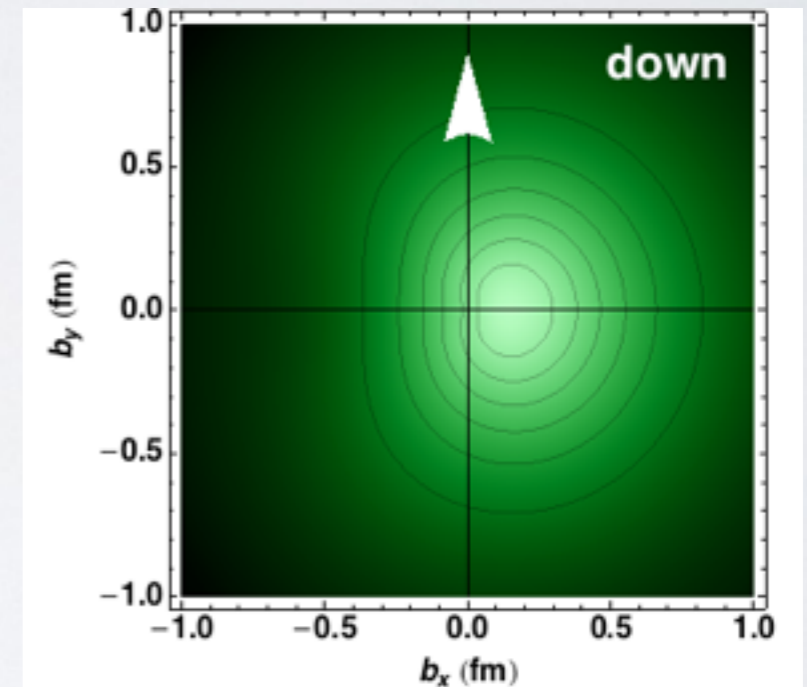
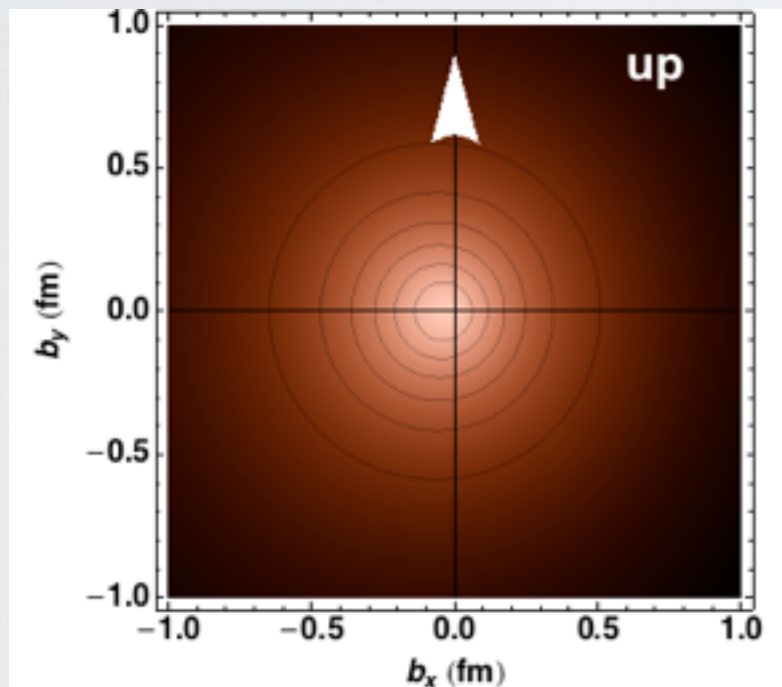
neutron core with negative charge
plus π cloud with positive charge !

polarized N \rightarrow deformation

polarization $S_y \rightarrow$ spin-flip $E(x,0,-q^2) \rightarrow$ b_x deformation
 $\mathbf{b} = b (\cos\Phi_b, \sin\Phi_b)$

$$\rho(\mathbf{b}) = \rho^0(\mathbf{b}) + \cos\phi_b \int_0^\infty \frac{d|\mathbf{q}|}{2\pi} \frac{q^2}{2M} J_1(|\mathbf{q}|b) F_2(Q^2 = \mathbf{q}^2)$$

proton
polarization



$E_x \sim$ dipole deformation

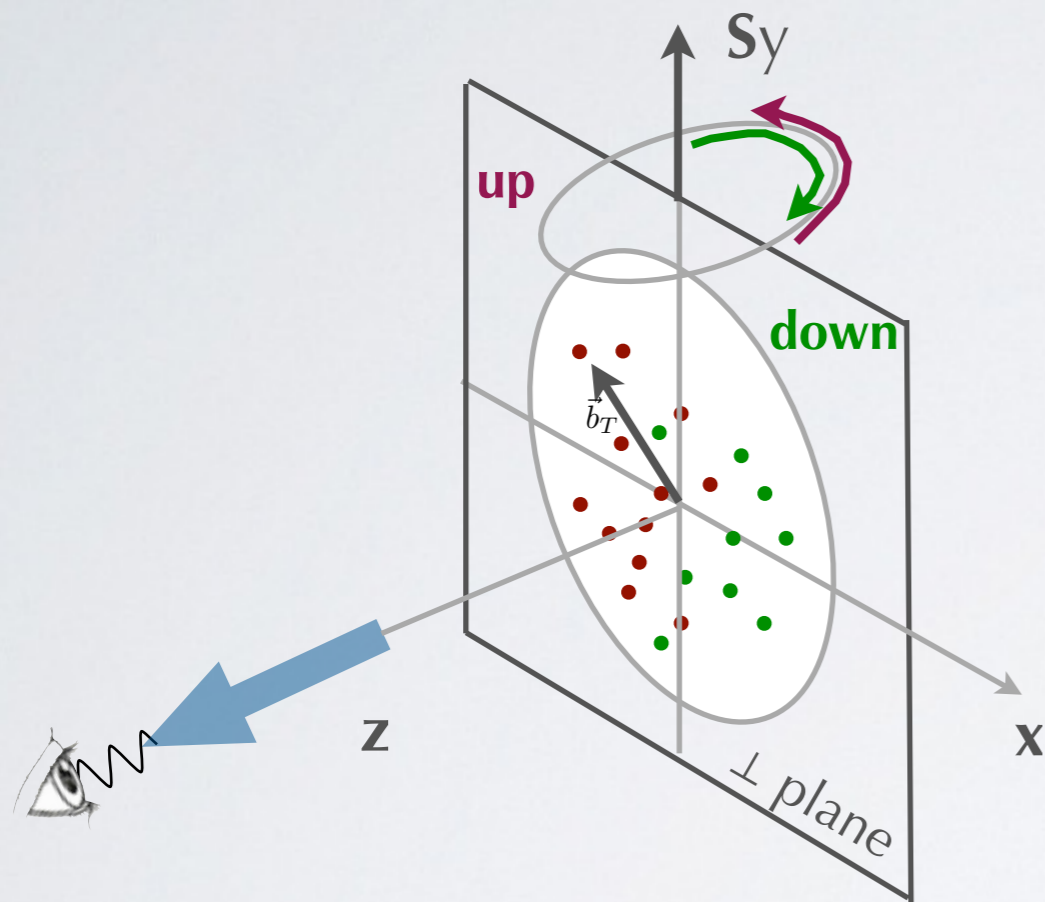
A. Bacchetta & M. Contalbrigo,
Il Nuovo Saggiatore **28** (12) n.1,2

and



C. Carlson & M. Vanderhaeghen,
P.R.L. **100** (08) 032004

parton **O**rbital **A**ngular **M**omentum



N^\uparrow polarization along **y**
gives a twist along **x**
to parton charge densities
because of their
Orbital **A**ngular **M**omentum
(OAM)

how to define it ?

(gauge-inv. definition is
common problem for
gauge field th.'s)

definition #1 of **OAM**

from Ji's sum rule :

OAM = **total J** – **helicity**

$$L_z^q(Q^2) \equiv J_z^q(Q^2) \left\{ = \frac{1}{2} \int dx x [f_1^q(x; Q^2) + E^q(x, 0, 0; Q^2)] \right\} \\ - S_z^q(Q^2) \left\{ = \int dx g_1(x; Q^2) \right\}$$

gauge invariant

measurable (DIS \rightarrow f_1, g_1 ; DVCS \rightarrow E)



definition #1 of **OAM**

from **Ji's sum rule** :

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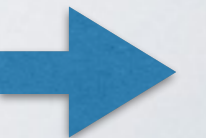
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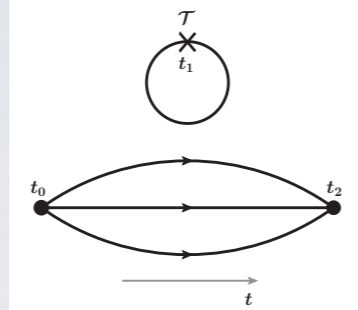
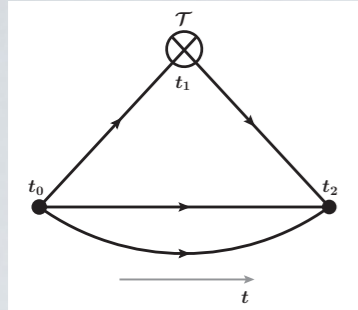
but L^q does not satisfy canonical relations

alternatives?...



the latest scenario from lattice

M. Deka et al. (χ QCD), arXiv:1312.4816 [hep-lat]



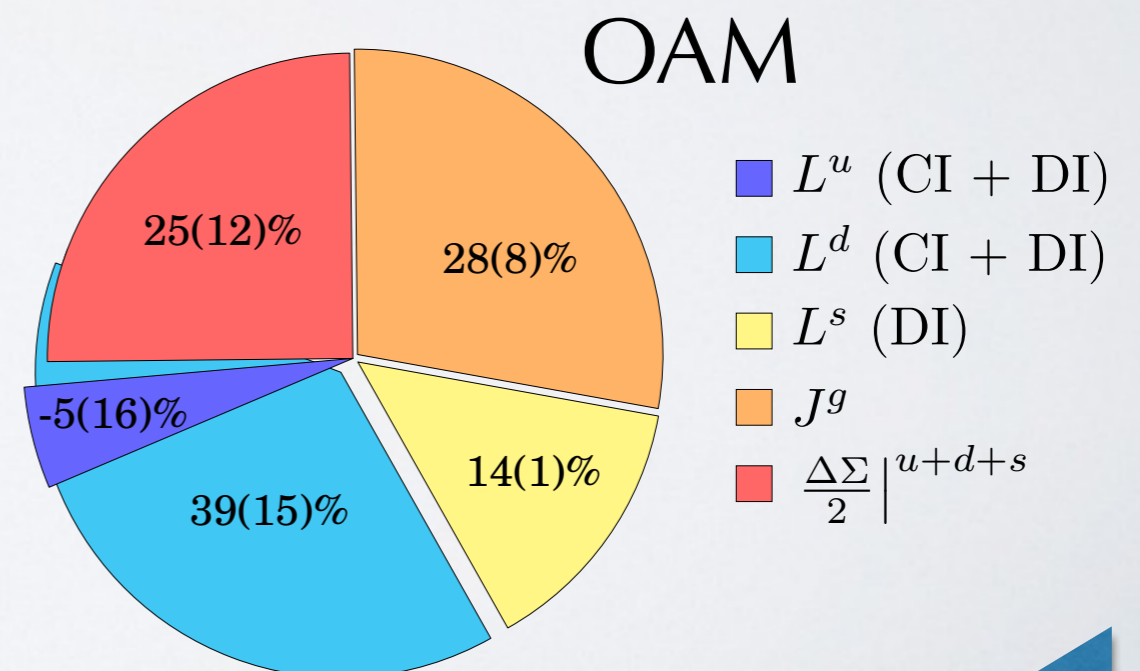
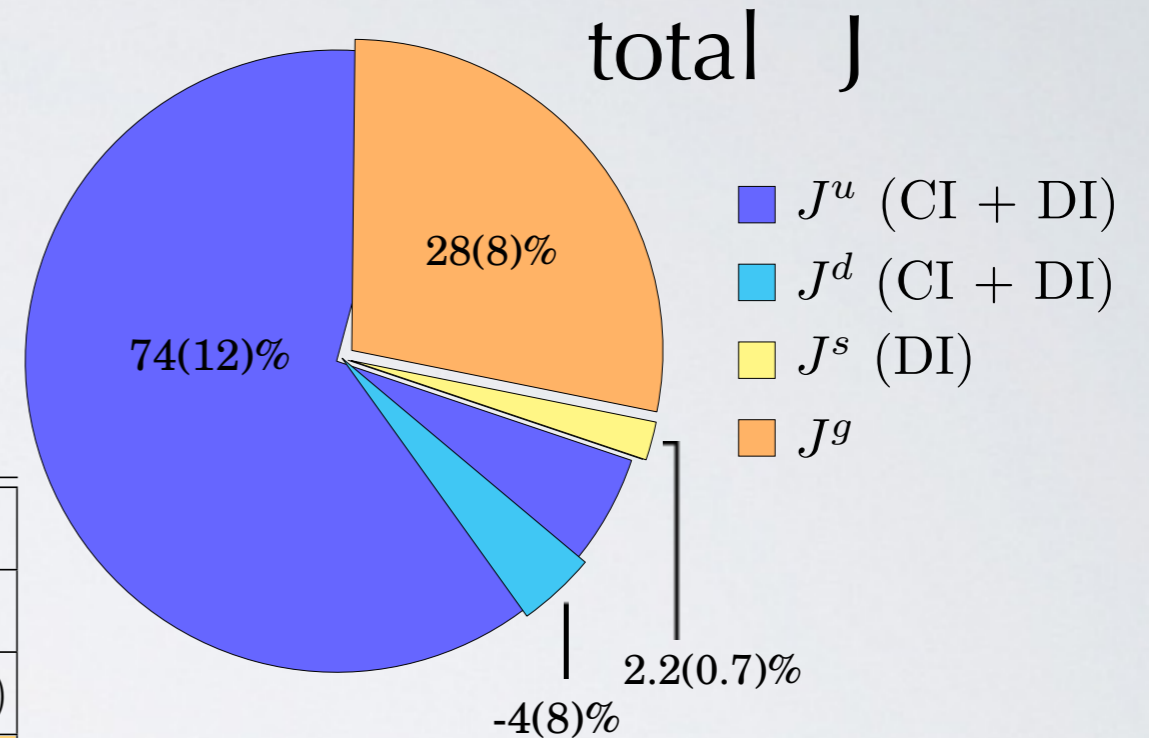
Connected Insertions Disconnected

	CI(u)	CI(d)	CI(u+d)	DI(u/d)	DI(s)	Glue
$\langle x \rangle$	0.416(40)	0.151(20)	0.567(45)	0.037(7)	0.023(6)	0.334(56)
$T_2(0)$	0.283(112)	-0.217(80)	0.061(22)	-0.002(2)	-0.001(3)	-0.056(52)
$2J$	0.704(118)	-0.070(82)	0.629(51)	0.035(7)	0.022(7)	0.278(76)
g_A	0.91(11)	-0.30(12)	0.62(9)	-0.12(1)	-0.12(1)	-
$2L$	-0.21(16)	0.23(15)	0.01(10)	0.16(1)	0.14(1)	-

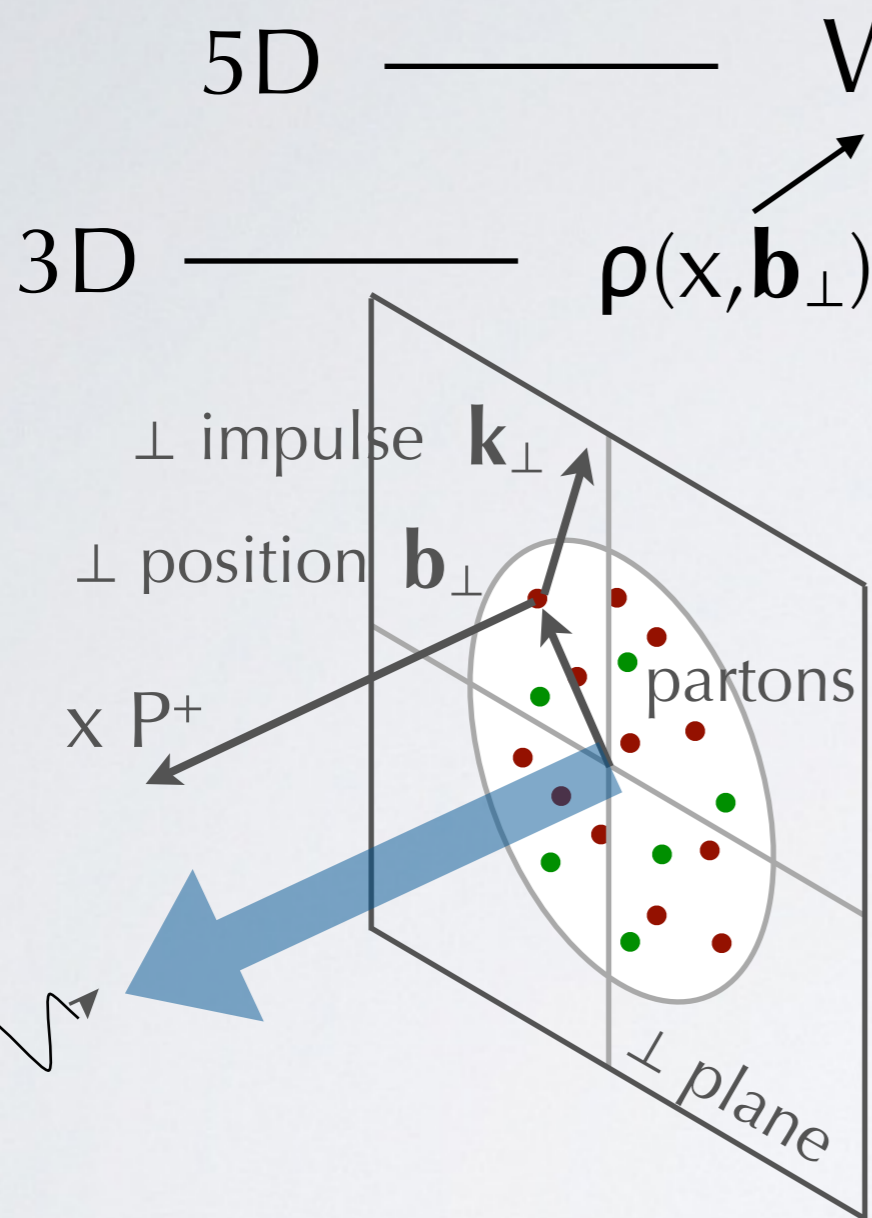
TABLE III. Renormalized values in \overline{MS} scheme at $\mu = 2$ GeV.

$$g_A^0 = \Delta u + \Delta d + \Delta s$$

$$2J - g_A^0 = 2L$$



Wigner Distribution



The “mother”
distribution

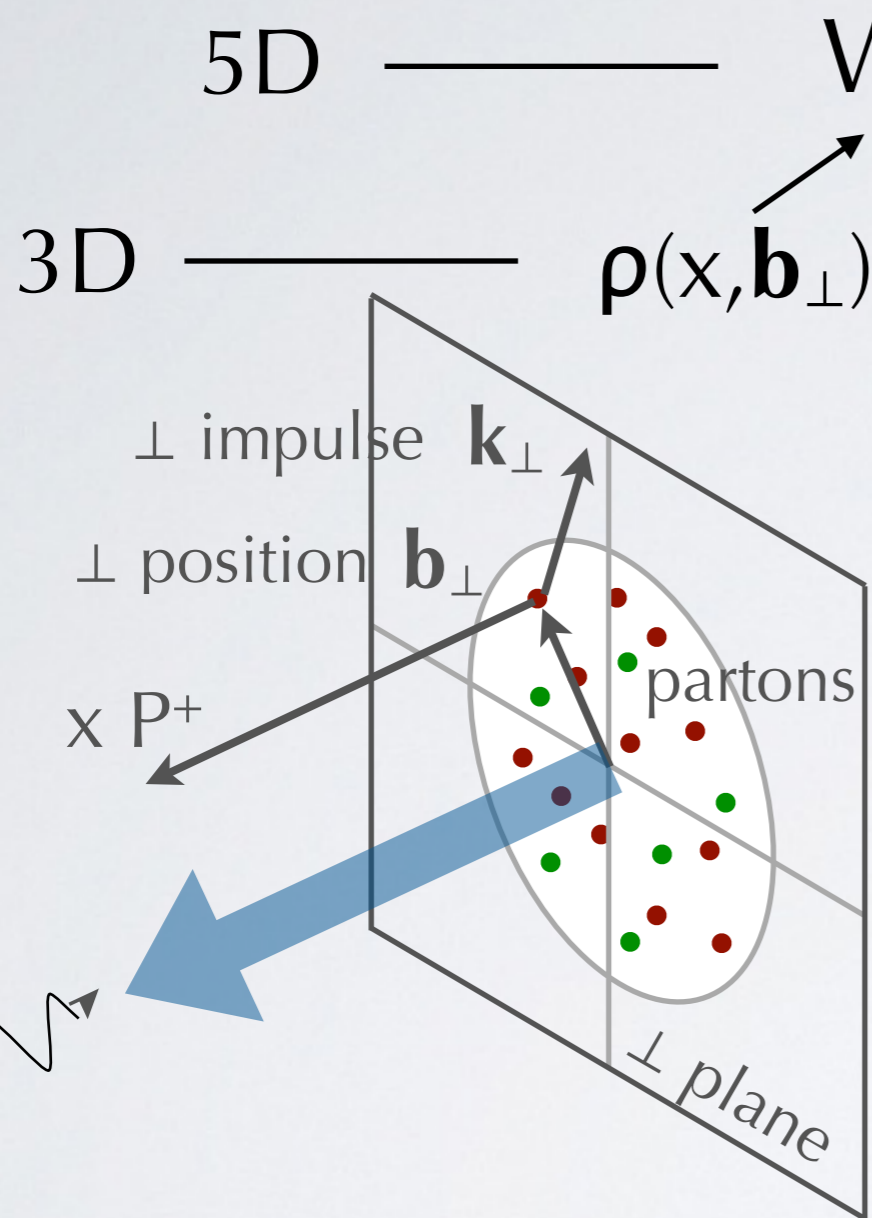
correlation of quark
 \perp **momentum** and **position**
for **S_N and S_q polarizations**

not positive-definite
but $\mathbf{b} \leftrightarrow \mathbf{q} = \mathbf{P}'_\perp - \mathbf{P}_\perp$

no constraint from
Heisenberg principle

*C. Lorcé, B. Pasquini, M. Vanderhaeghen,
JHEP 1105 (11) 041*

Wigner Distribution



The “mother”
distribution

correlation of quark
 \perp **momentum** and **position**
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no constraint from
Heisenberg principle

*C. Lorcé, B. Pasquini, M. Vanderhaeghen,
JHEP 1105 (11) 041*

$$\int d\mathbf{k}_\perp W(x, \mathbf{k}_\perp, \mathbf{b}_\perp) \rightarrow q(x, \mathbf{b}_\perp) \rightarrow \text{GPD}$$

Transverse Mom. Distributions (TMD)

5D ————— $\int d\mathbf{b}_\perp W(x, \mathbf{k}_\perp, \mathbf{b}_\perp)$

3D —————

$q(x, \mathbf{k}_\perp)$ **TMD**

parton density in k -space
is not the F.T. of $q(x, \mathbf{b}_\perp)$

Transverse Mom. Distributions (TMD)

5D ————— $\int d\mathbf{b}_\perp W(x, \mathbf{k}_\perp, \mathbf{b}_\perp)$

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parton density in k -space
is not the F.T. of $q(x, \mathbf{b}_\perp)$

leading twist: 8 **TMDs**

quark pol.

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Twist-2 TMDs

Transverse Mom. Distributions (TMD)

$$5D \text{ ————— } \int d\mathbf{b}_\perp W(x, \mathbf{k}_\perp, \mathbf{b}_\perp)$$

$$3D \text{ ————— } q(x, \mathbf{k}_\perp) \text{ TMD}$$

parton density in k -space
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leading twist: 8 TMDs

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	U	L	T
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T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Twist-2 TMDs

* Anselmino et al., P.R. D87 (13) 094019

* A. Bacchetta, A. Courtoy, M. Radici, JHEP 03 (13) 119

$$\int d\mathbf{k}_\perp \text{TMD}(x, \mathbf{k}_\perp) \rightarrow \text{PDF}(x)$$

Transverse Mom. Distributions (TMD)

$$5D \text{ ————— } \int d\mathbf{b}_\perp W(x, \mathbf{k}_\perp, \mathbf{b}_\perp)$$

$$3D \text{ ————— } q(x, \mathbf{k}_\perp) \text{ TMD} \quad \text{parton density in } k\text{-space}$$

is not the F.T. of $q(x, \mathbf{b}_\perp)$

leading twist: 8 TMDs

quark pol.

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

$$\int dx h_1(x) = \text{tensor charge}$$

Twist-2 TMDs

* Anselmino et al., P.R. D87 (13) 094019

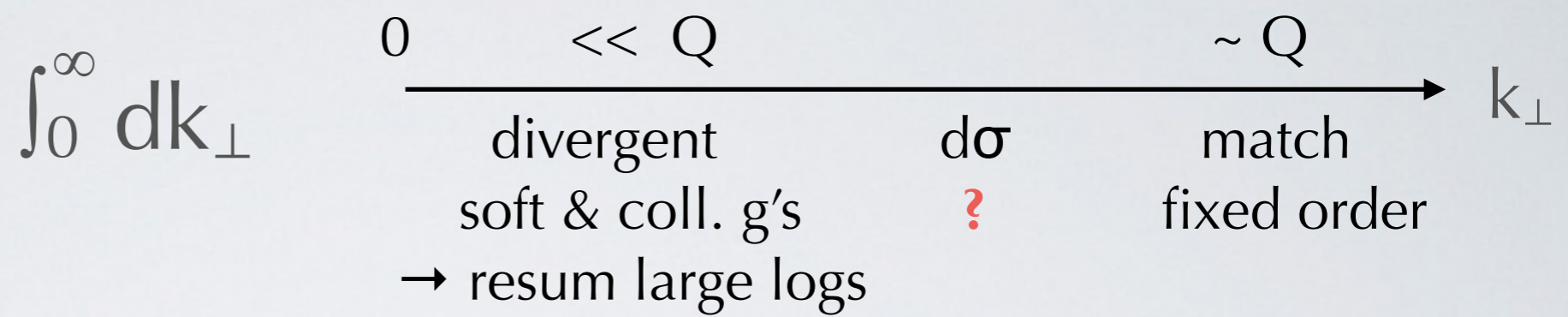
* A. Bacchetta, A. Courtoy, M. Radici, JHEP 03 (13) 119

$$\int d\mathbf{k}_\perp \text{TMD}(x, \mathbf{k}_\perp) \rightarrow \text{PDF}(x)$$

$$\int dk_{\perp} \text{TMD}(x, k_{\perp}) \rightarrow \text{PDF}(x) ?$$

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^{\perp}
	L		g_{1L}	h_{1L}^{\perp}
	T	f_{1T}^{\perp}	g_{1T}	h_1, h_{1T}^{\perp}

Twist-2 TMDs

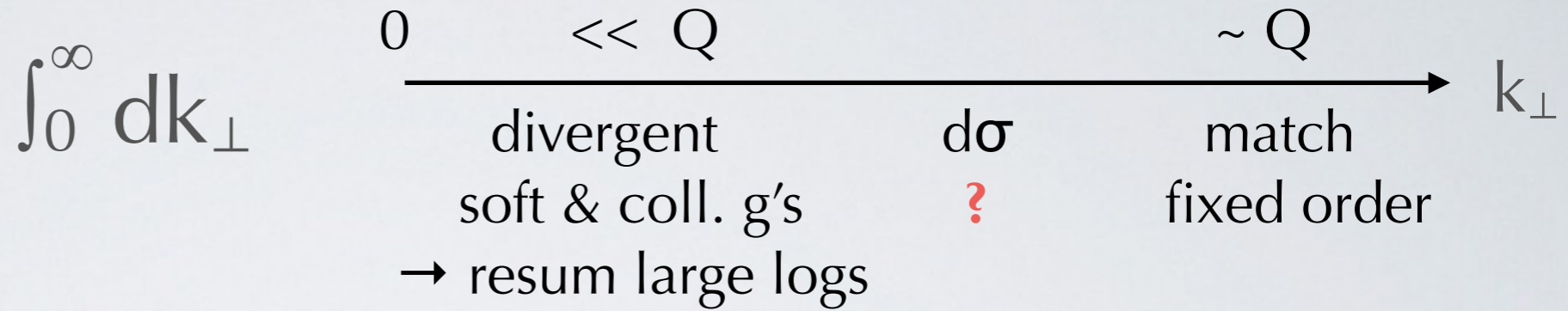


*Collins, Soper, Sterman, N.P. **B250** (85) 199*
Collins, "Foundations of perturb. QCD" (C.U.P., 11)
*Echevarria et al., E.P.J. **C73** (13) 2636*

$\int dk_{\perp} \text{TMD}(x, k_{\perp}) \rightarrow \text{PDF}(x) ?$

	quark pol.		
	U	L	T
nucleon pol.	U	f_1	h_1^{\perp}
	L		h_{1L}^{\perp}
	T	f_{1T}^{\perp}	h_1, h_{1T}^{\perp}

Twist-2 TMDs



Collins, Soper, Sterman, N.P. B250 (85) 199
Collins, "Foundations of perturb. QCD" (C.U.P., 11)
Echevarria et al., E.P.J. C73 (13) 2636

in \mathbf{b}_{\perp} space

$$f_1^q(x, \mathbf{b}_{\perp}; Q^2) = \sum_i [C_{qi} \otimes f_1^i] \left(x; \frac{c_0^2}{b_*^2} \right) e^{S_P(b_*; Q)} e^{S_{NP}(\mathbf{b}_{\perp}) \log Q/Q_0} f_1^q(x, \mathbf{b}_{\perp}; Q_0^2)$$

hard coeffs.

PDF

perturb.
Sudakov

non-perturb.
Sudakov

non-perturb.
input TMD

all divergent for $b_{\perp} \rightarrow \infty$ ($k_{\perp} \rightarrow 0$)

prescription: $b_{\perp} \Rightarrow b_* = \frac{b_{\perp}}{\sqrt{1 + \frac{b_{\perp}^2}{b_{\max}^2}}}$

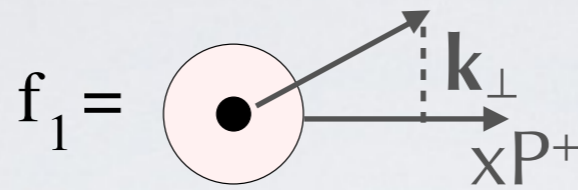
depend
on **parameters**
to be fitted

$f_1^q(x, \mathbf{k}_\perp) \rightarrow \text{LHC}$

quark pol.

	U	L	T
nucleon pol.	f_1		h_1^\perp
		g_{1L}	h_{1L}^\perp
	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Twist-2 TMDs



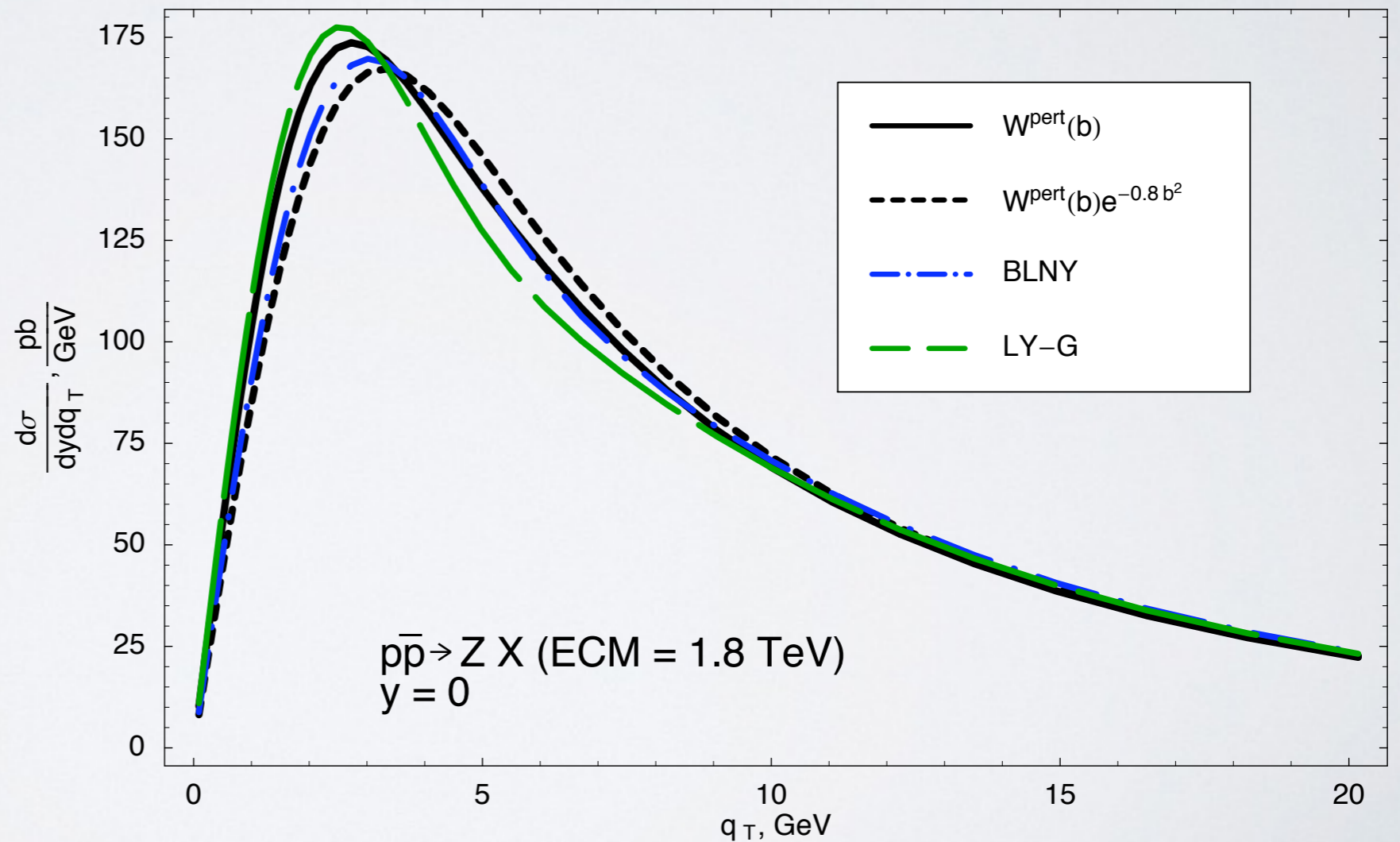
impact of TMD on Z^0 peak \rightarrow W mass

P. Nadolski, hep-ph/0412146

7.5%

30%

uncertainty

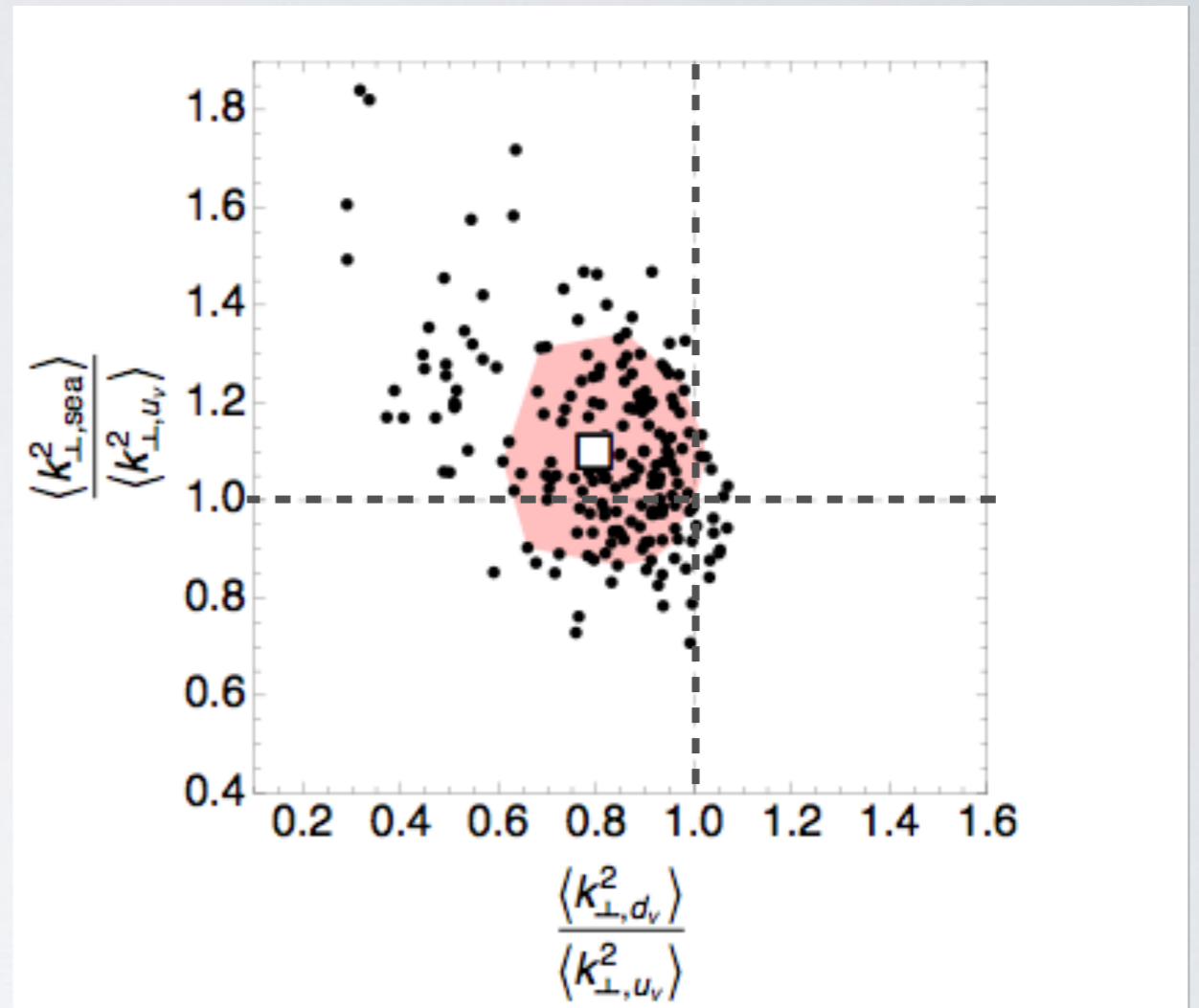


flavor analysis of TMD(x, \mathbf{k}_\perp)

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Twist-2 TMDs

fit SIDIS
multiplicity
from HERMES



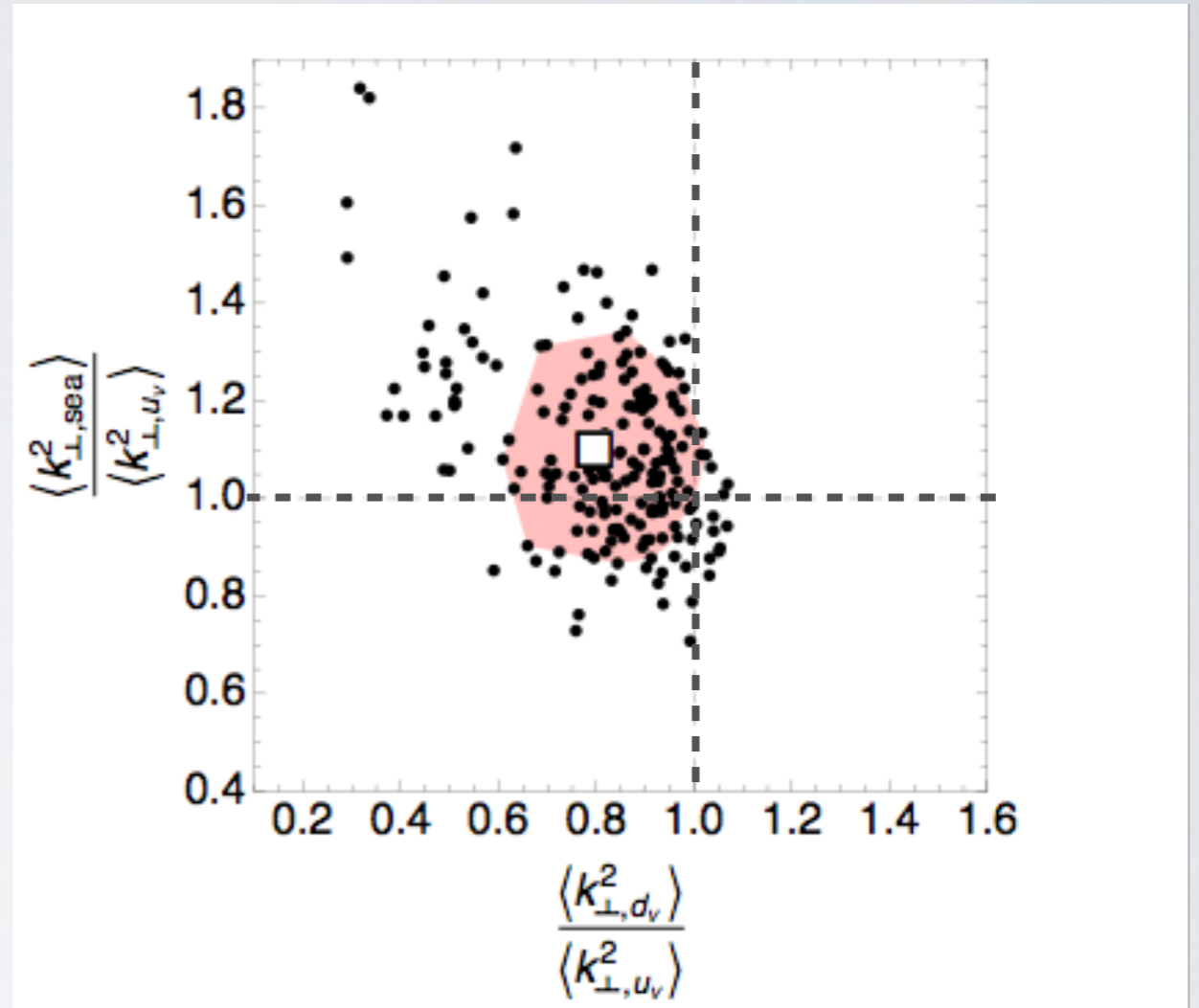
A. Signori et al., JHEP1311 (13) 194

flavor analysis of TMD(x, \mathbf{k}_\perp)

	quark pol.		
	U	L	T
nucleon pol.	f_1		h_1^\perp
		g_{1L}	h_{1L}^\perp
	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

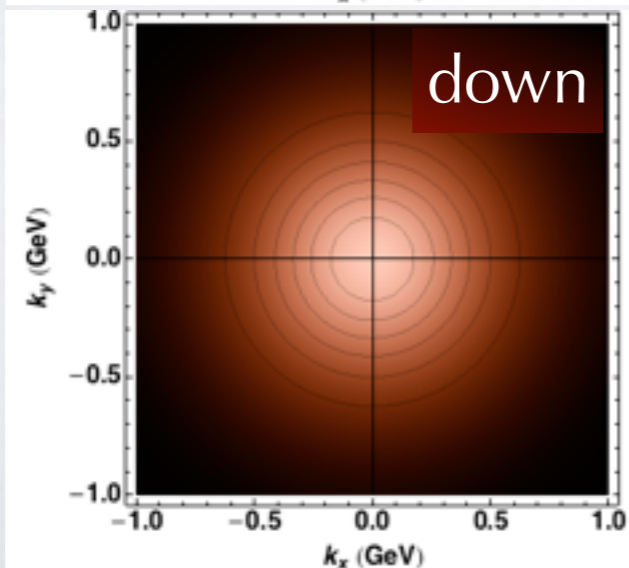
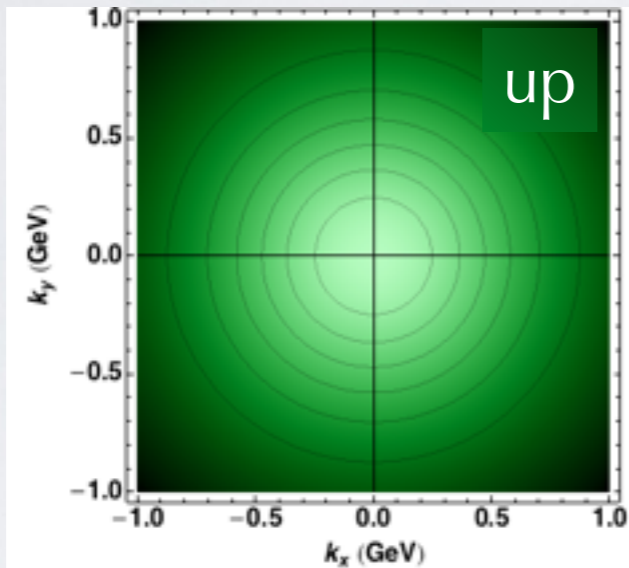
Twist-2 TMDs

fit SIDIS
multiplicity
from HERMES



A. Signori et al., JHEP1311 (13) 194

$x=0.1$



down < up < sea ?

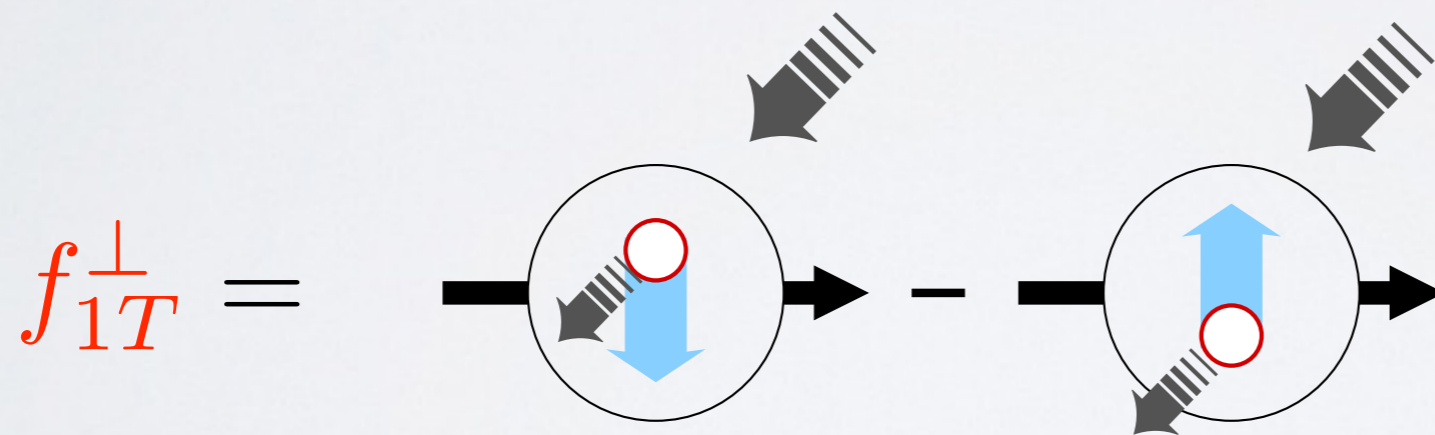
the Sivers effect

leading twist: 8 TMDs

	quark pol.		
	U	L	T
nucleon pol.	U	f_1	h_1^\perp
	L		g_{1L}
	T	f_{1T}^\perp	g_{1T}
			h_1, h_{1T}^\perp

Sivers function

Twist-2 TMDs

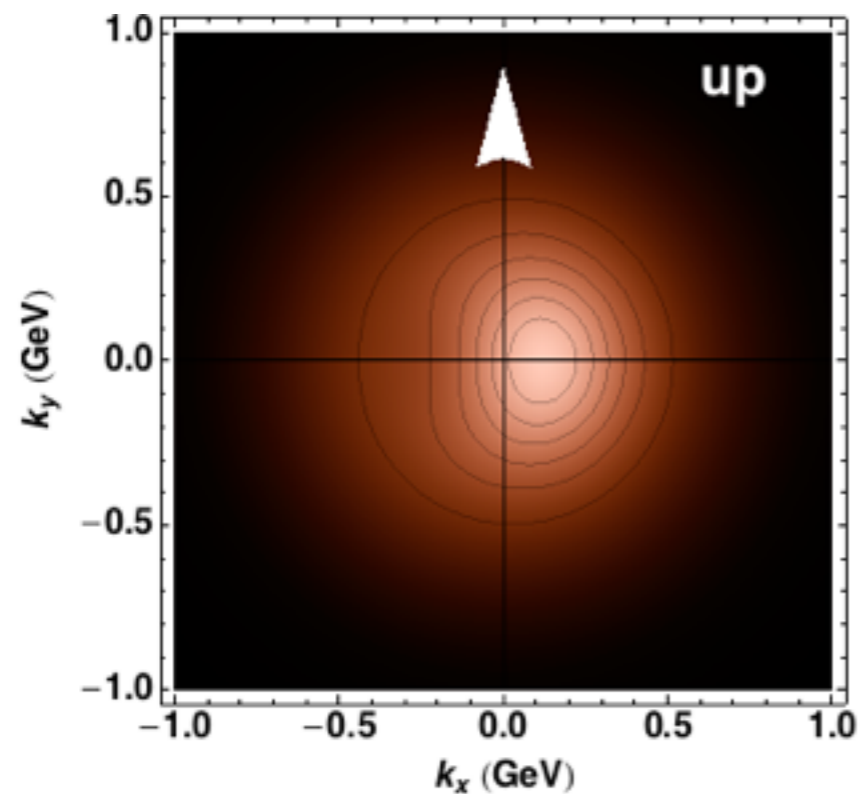


distortion of quark distribution
because of N^\uparrow polarization

flavor dependence of Sivers effect

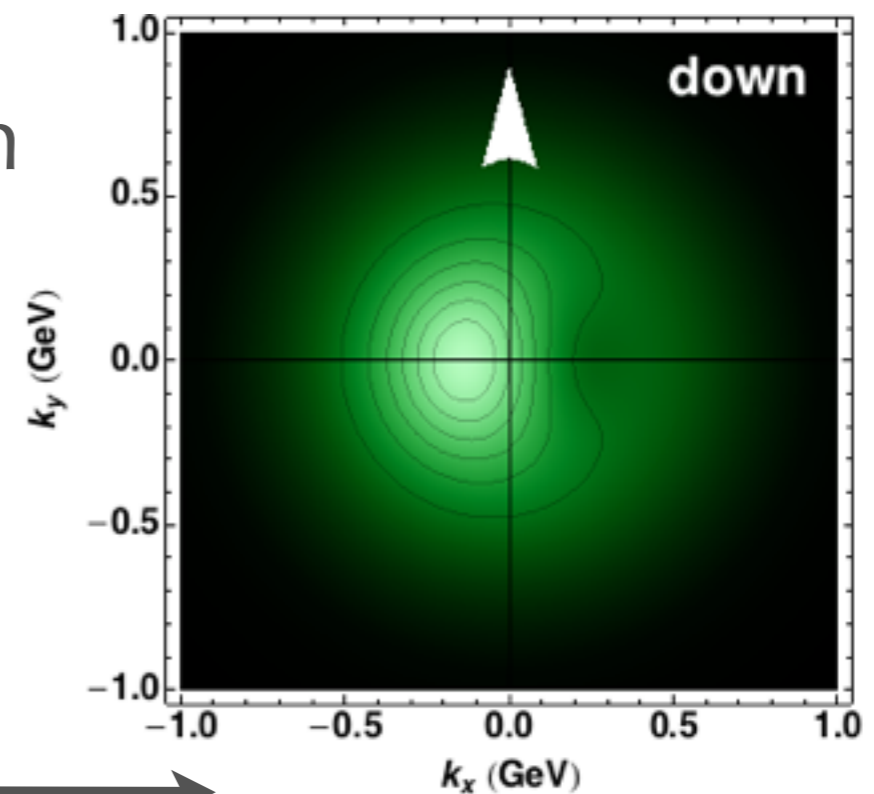
distribution of unpolarized q in polarized P^\uparrow

$$f_{q/p^\uparrow}(x, \mathbf{k}_\perp) = f_1^q(x, \mathbf{k}_\perp^2) - f_{1T}^{\perp q}(x, \mathbf{k}_\perp^2) \frac{(\hat{\mathbf{P}} \times \mathbf{k}_\perp) \cdot \mathbf{S}}{M}$$



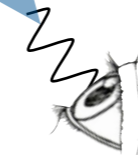
polarization

S_y



deformation along x

*A. Bacchetta & M. Contalbrigo,
Il Nuovo Saggiatore* **28** (12) n.1,2



the Sivers effect in semi-incl. DIS (SIDIS)



the Sivers effect in semi-incl. DIS (SIDIS)

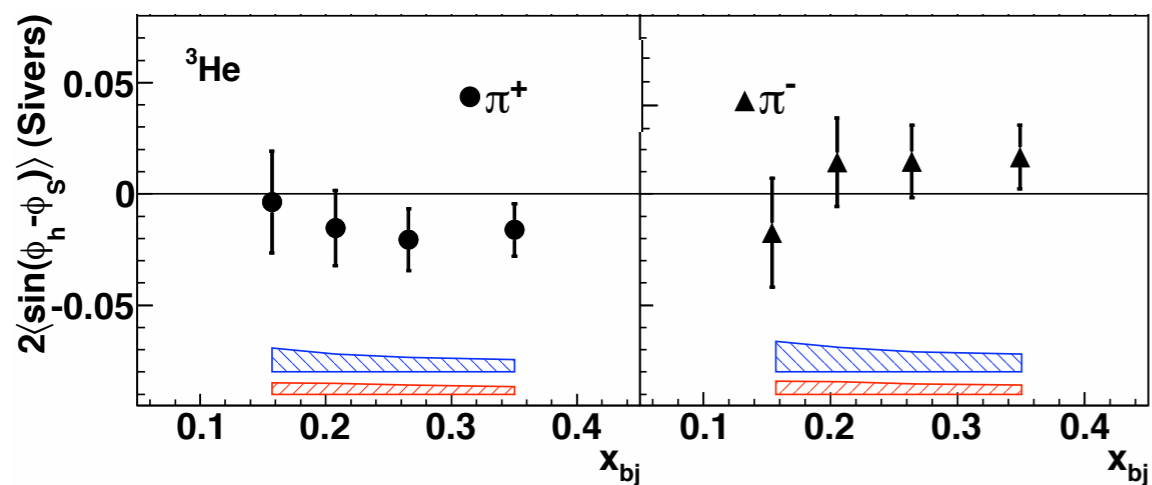
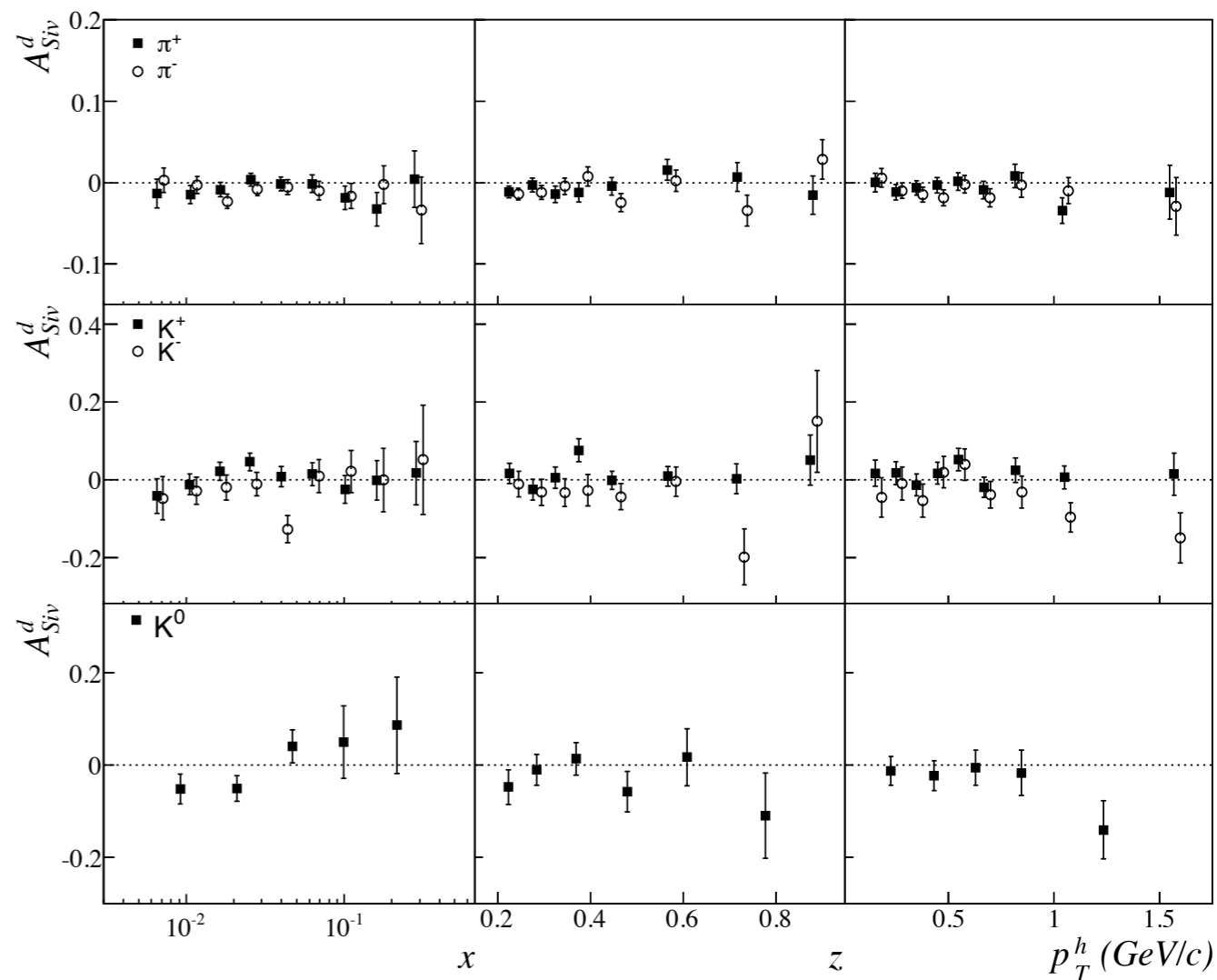
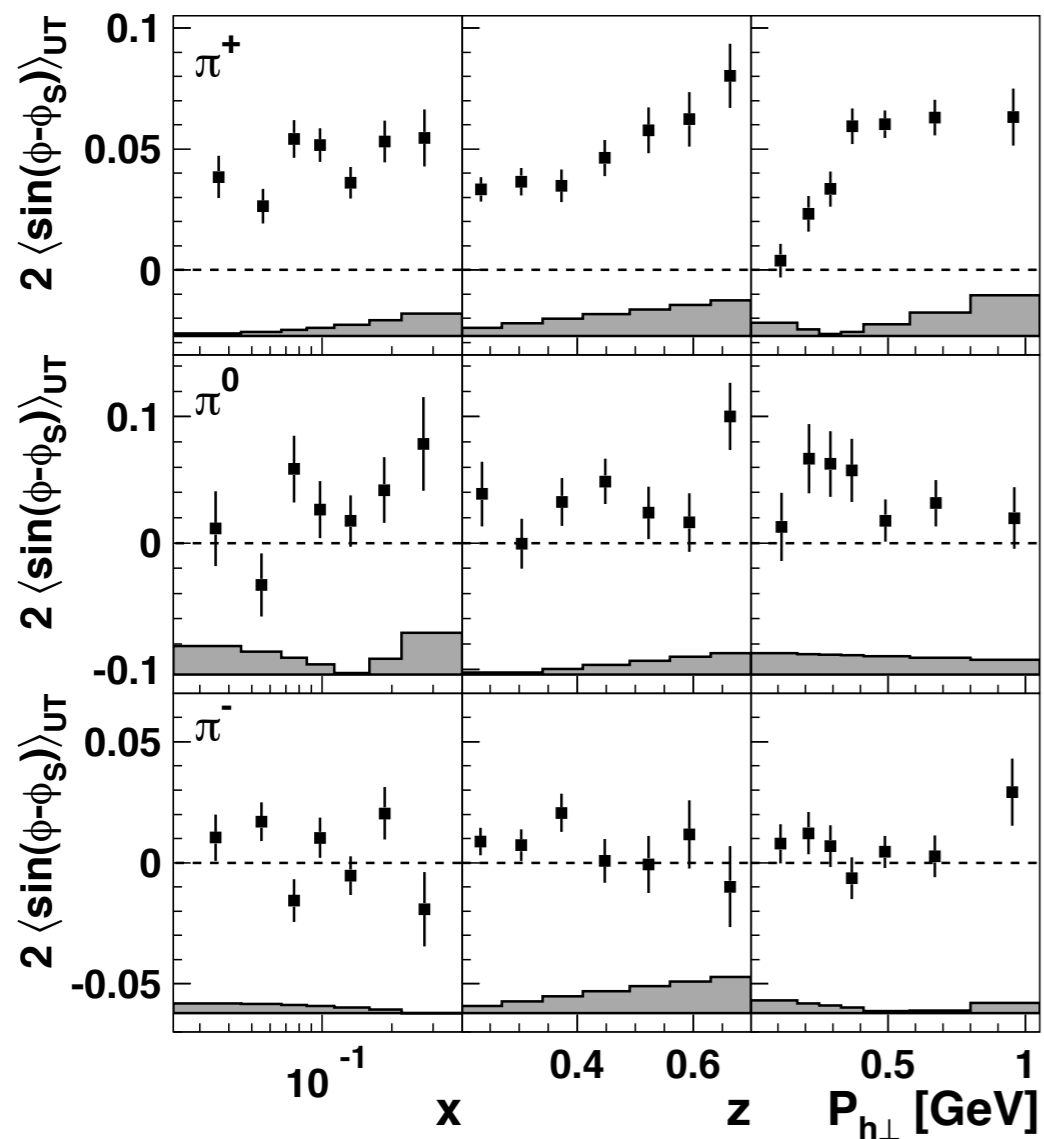




PRL103 (09) 152002



PL B673 (09) 127



Jefferson Lab
Hall A

PRL107 (11) 072003

Sivers effect has
been measured
in $N^\uparrow(e, e'\pi)$!

parametrizations of Sivers function

leading twist: 8 TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Twist-2 TMDs

* M. Anselmino et al, *E.P.J. A***39** (09) 89

* M. Anselmino, M.E. Boglione, S. Melis,
*P.R. D***86** (12) 014028

* A. Bacchetta & M. Radici,
P.R.L. **107** (11) 212001

* W. Vogelsang & F. Yuan,
*P.R. D***72** (05) 054028

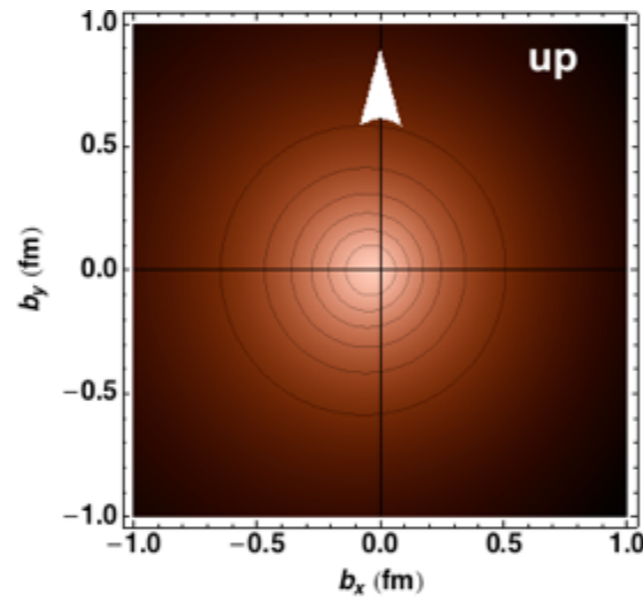
* J.C. Collins et al.,
*P.R. D***73** (06) 014021

* S.M. Aybat, A. Prokudin, T.C. Rogers,
P.R.L. **108** (12) 242003

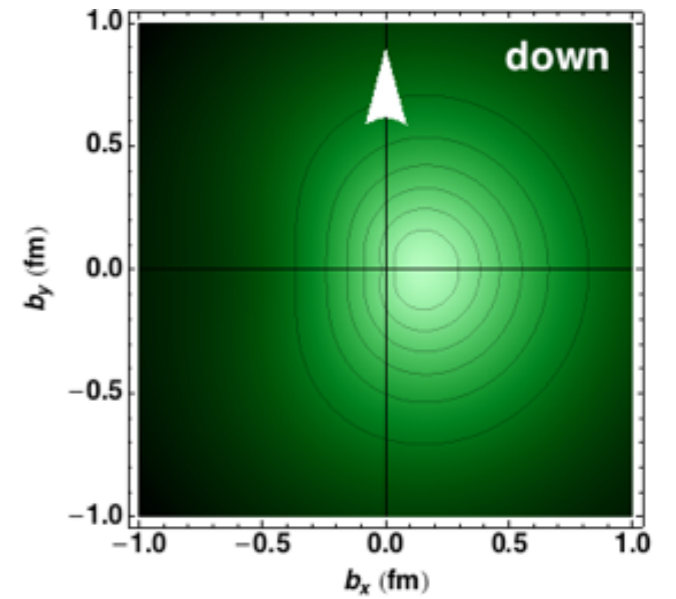
* P. Sun & F. Yuan,
*P.R. D***88** (13) 034016

* D. Boer,
N.P. **B874** (13) 217

GPD E



N^\uparrow
polarization



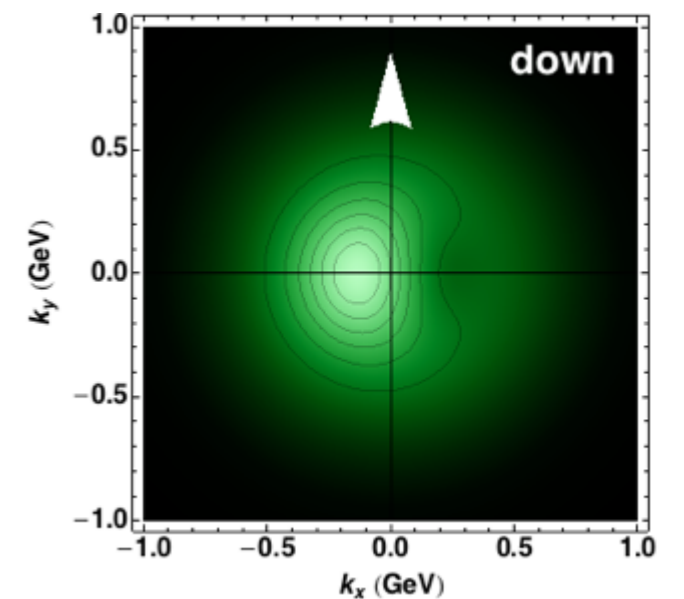
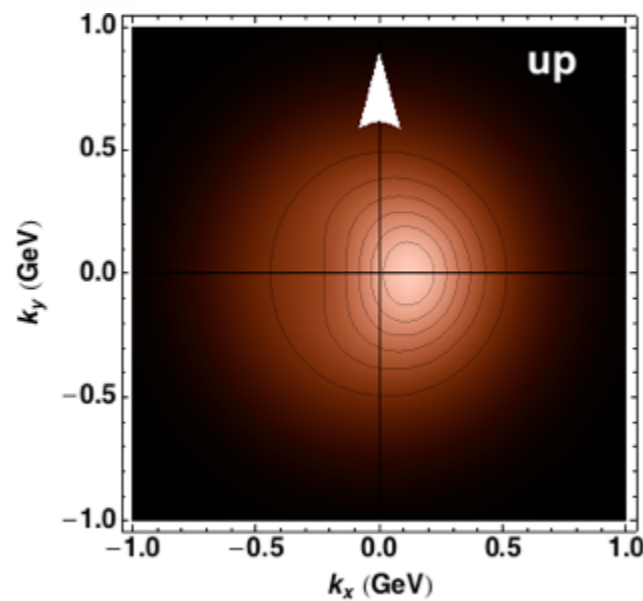
in **b** space

deformation

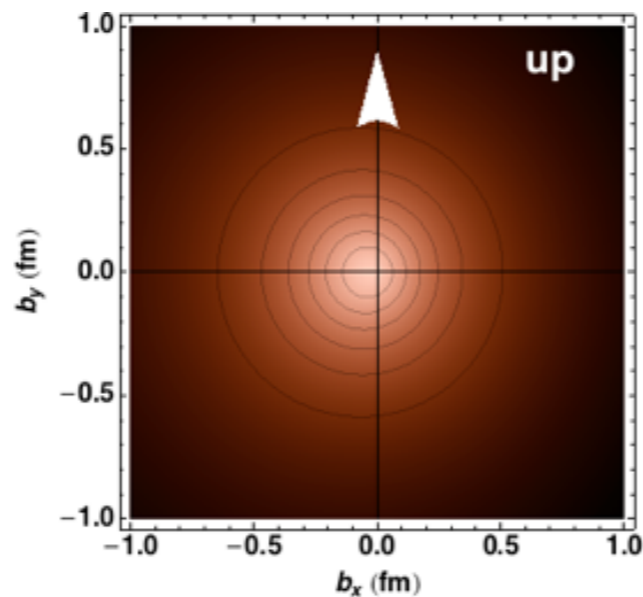
in **k** space

TMD f_{1T}^\perp

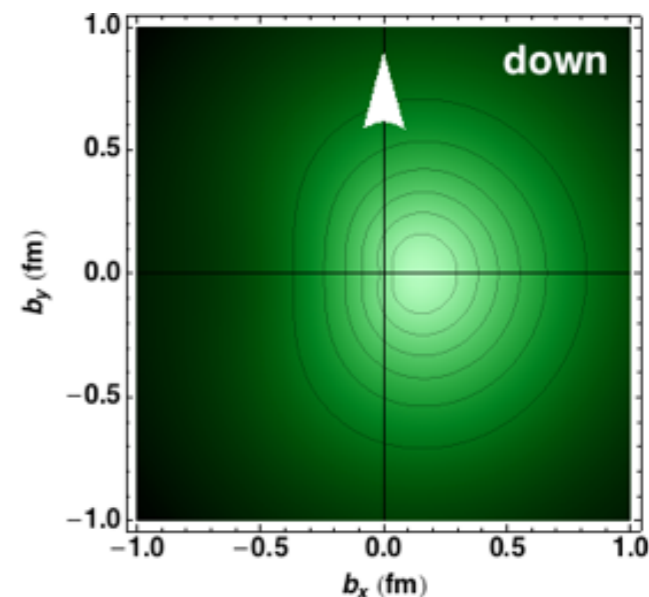
Sivers
effect



GPD E



N^\uparrow
polarization



in **b** space

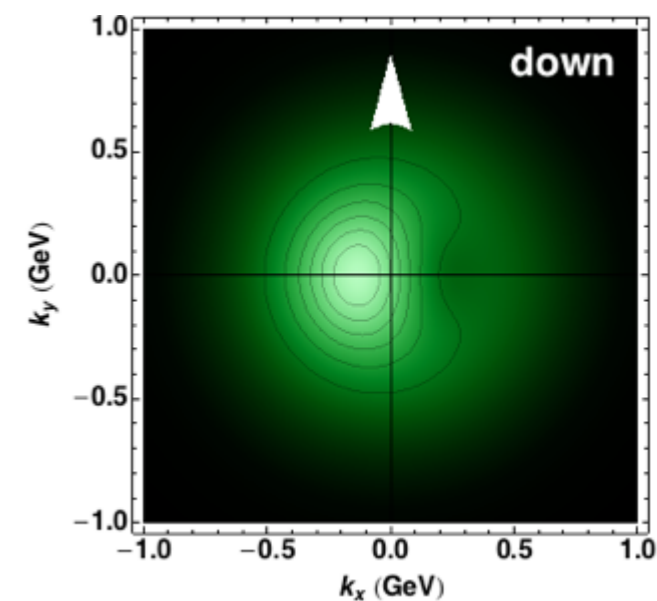
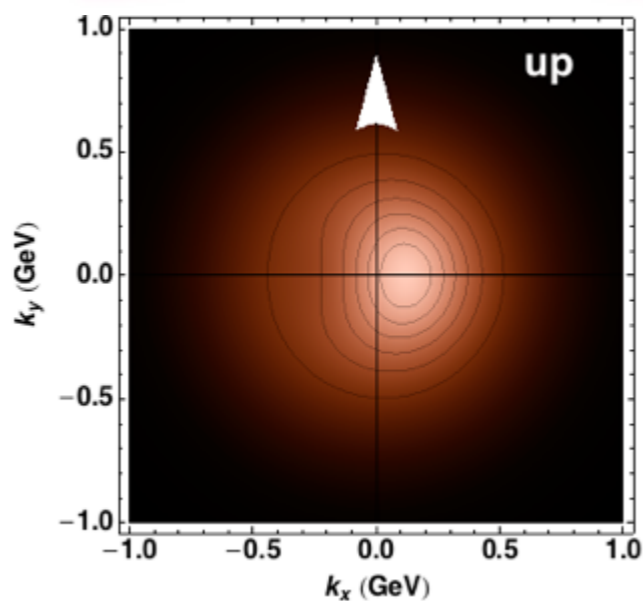
in **k** space

deformation

lensing function color FSI

TMD f_{1T}^\perp

Sivers effect



Ji's sum rule

$$J_z^q(Q^2) = \frac{1}{2} \int_0^1 dx x [H^q(x, 0, 0; Q^2) + E^q(x, 0, 0; Q^2)]$$

not accessible

assumption

$$f_{1T}^{\perp(0)q}(x; Q_L^2) = -L(x)E^q(x, 0, 0; Q_L^2)$$

lensing funct.

(at some Q_L)

*A. Bacchetta, F. Conti, M. Radici,
P.R. D78 (08) 074010*

Ji's sum rule

$$J_z^q(Q^2) = \frac{1}{2} \int_0^1 dx x [H^q(x, 0, 0; Q^2) + E^q(x, 0, 0; Q^2)]$$

not accessible

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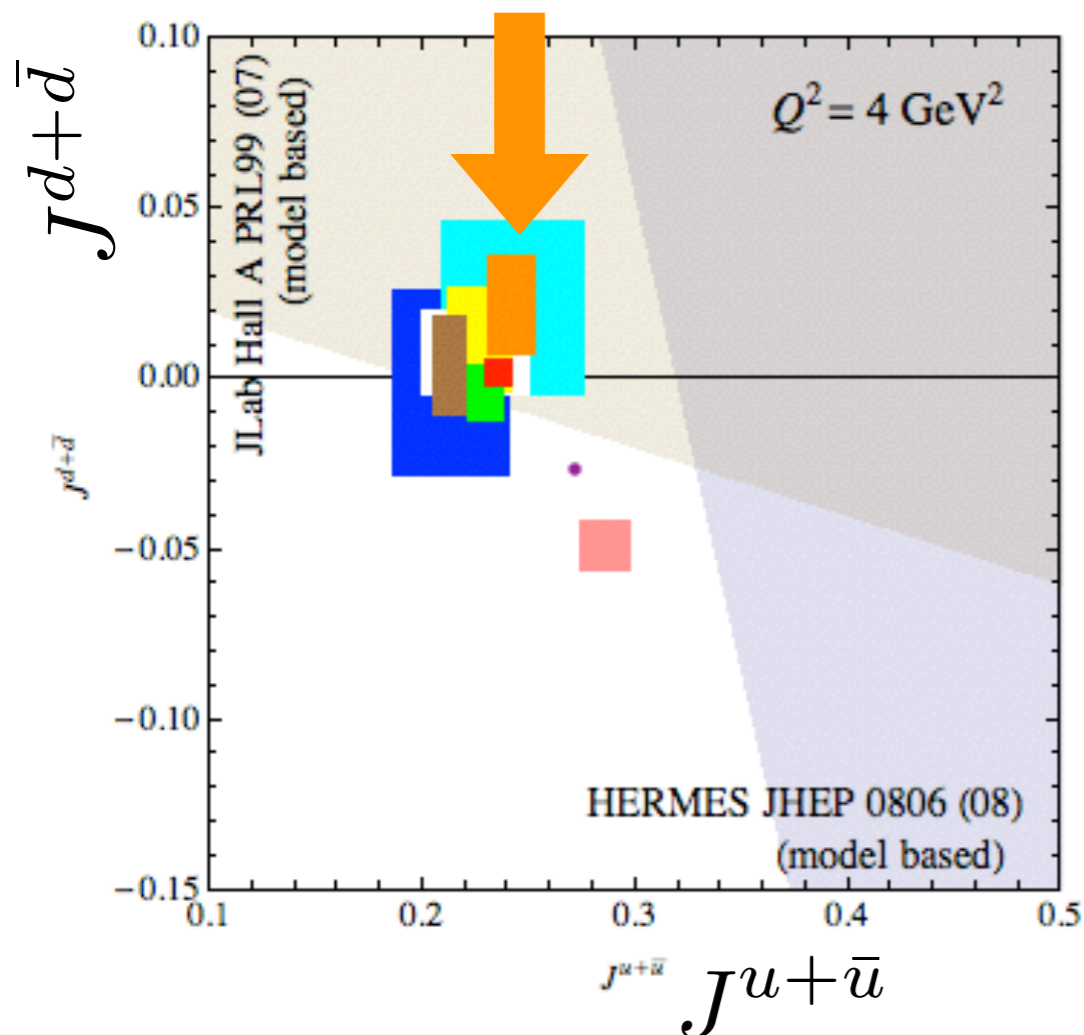
$$f_{1T}^{\perp(0)q}(x; Q_L^2) = -L(x)E^q(x, 0, 0; Q_L^2)$$

lensing funct.

(at some Q_L)

A. Bacchetta, F. Conti, M. Radici,
P.R. D78 (08) 074010

comparison with other GPD extractions and lattice results



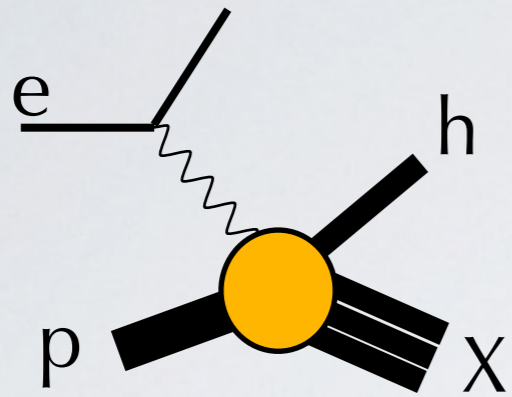
- Goloskokov & Kroll, EPJ C59 (09) 809
- Diehl & Kroll, E.P.J. C73 (13) 2397
- Diehl et al., EPJ C39 (05) 1
- Guidal et al., PR D72 (05) 054013
- Liuti et al., PRD 84 (11) 034007
- Bacchetta & Radici, PRL 107 (11) 212001
- LHPC-1, PR D77 (08) 094502
- LHPC-2, PR D82 (10) 094502
- QCDSF, arXiv:0710.1534
- Wakamatsu, EPJ A44 (10) 297

$$J^{u-\bar{u}} = 0.214^{+0.009}_{-0.013} \quad J^{d-\bar{d}} = -0.029^{+0.021}_{-0.008}$$

$$J^{u-\bar{u}} = 0.230^{+0.009}_{-0.024} \quad J^{d-\bar{d}} = -0.004^{+0.010}_{-0.016}$$

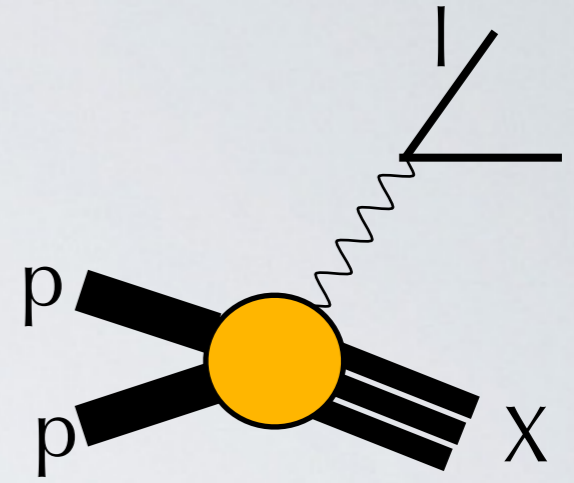


Reason #2 for the Sivers function

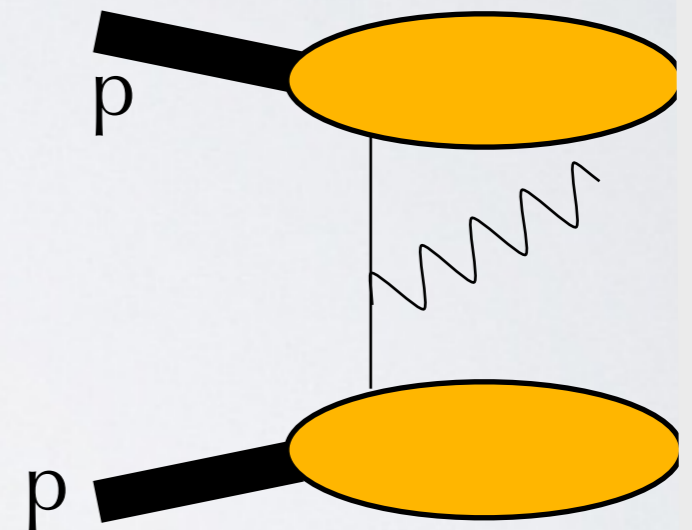
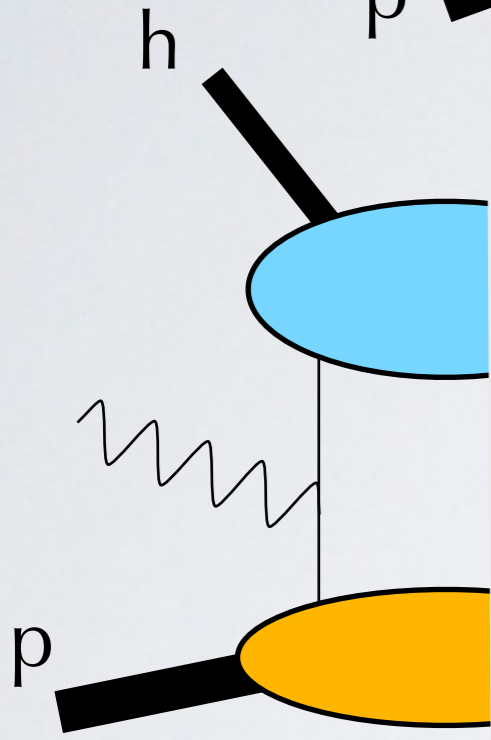


SIDIS

Drell-Yan



factorization theorems

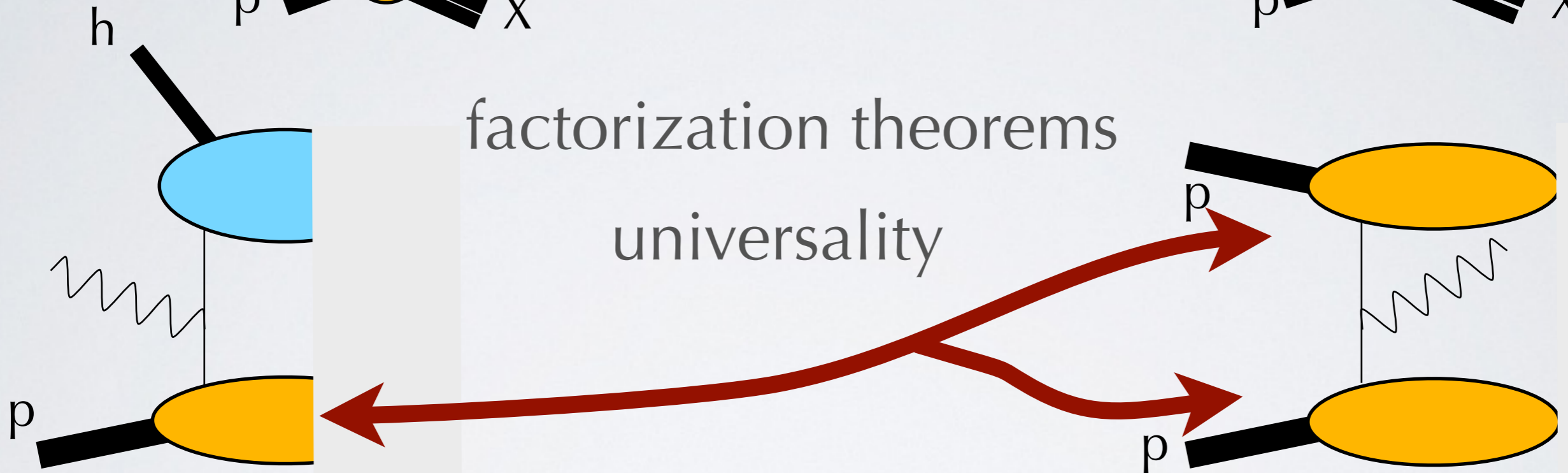


Reason #2 for the Sivers function



factorization theorems

universality

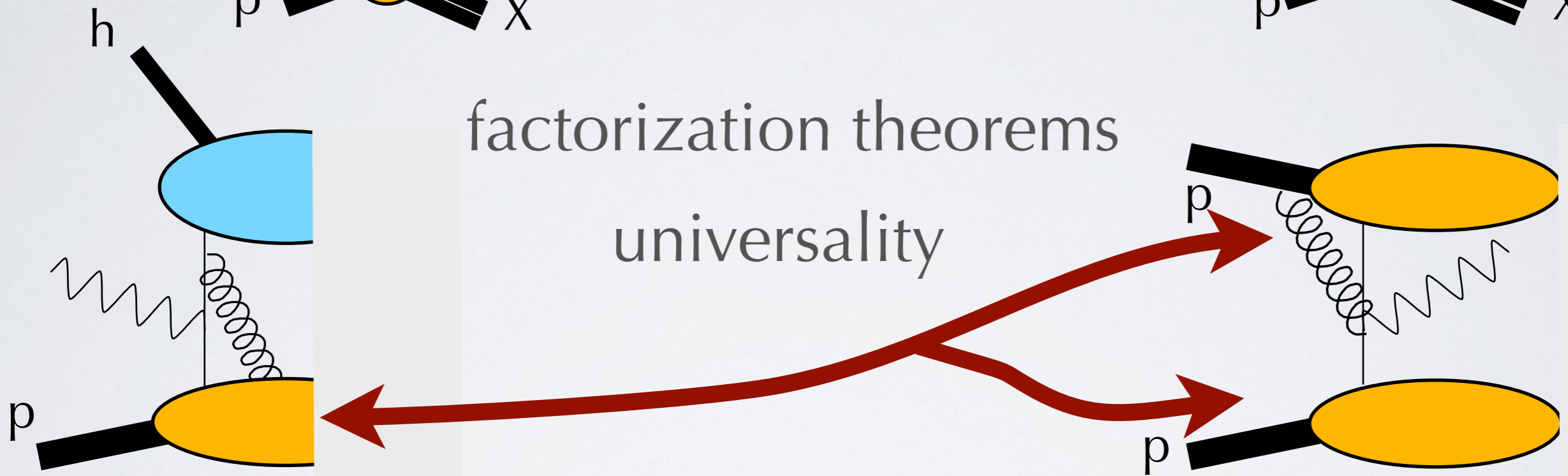


Reason #2 for the Sivers function



factorization theorems

universality



QCD prediction to be tested:
(at COMPASS)

$$\text{Sivers} \Big|_{\text{SIDIS}} = -\text{Sivers} \Big|_{\text{D-Y}}$$

J.C. Collins, P.L. B536 (02)

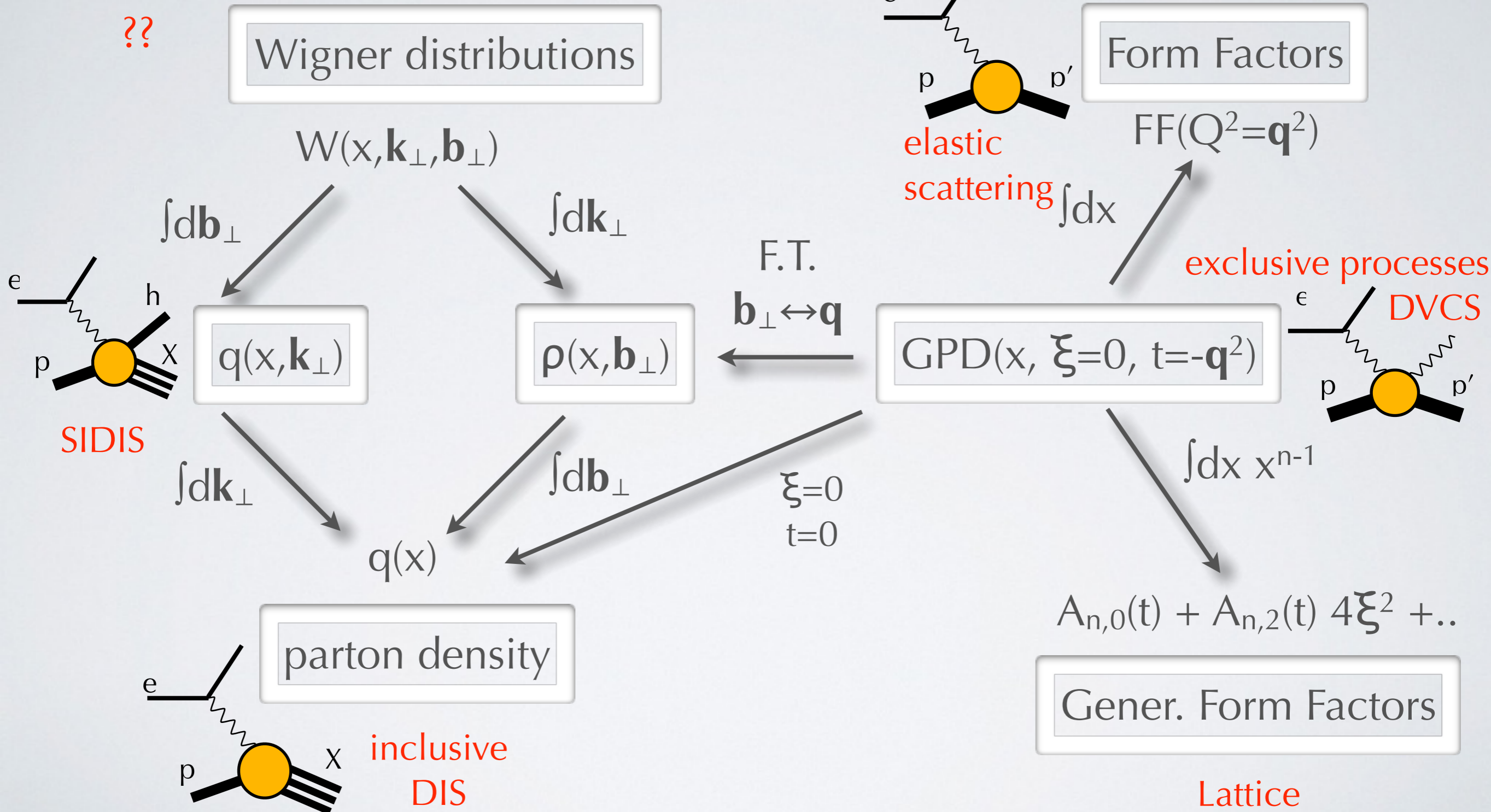
“roadmap” to a multi-dim. picture of N

*“Electron Ion Collider: the Next QCD Frontier”
arXiv:1212.1701 [nucl-ex]*

$$\text{GPD}(x, \xi=0, t=-\mathbf{q}^2)$$

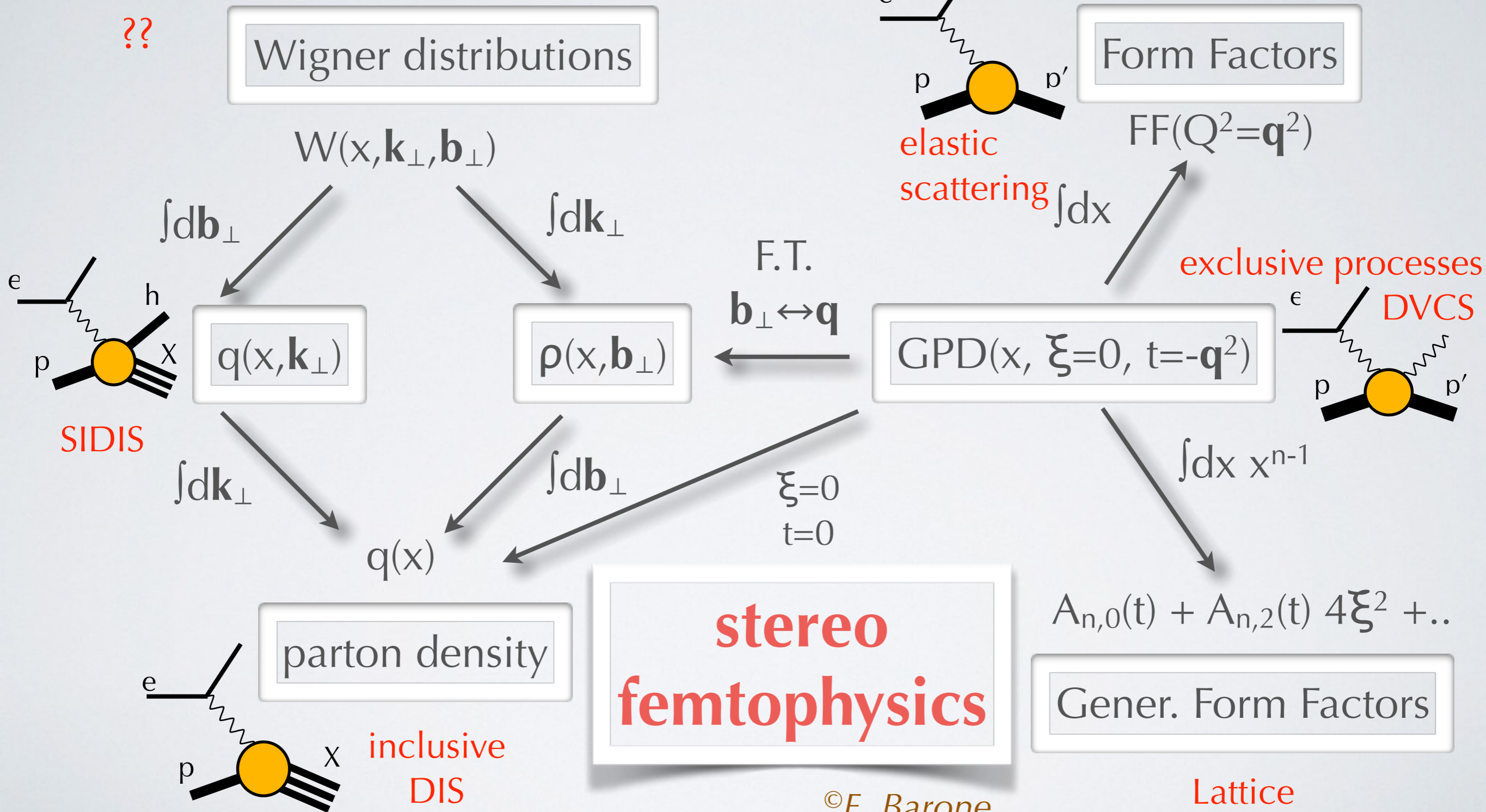
“roadmap” to a multi-dim. picture of N

“Electron Ion Collider: the Next QCD Frontier”
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“roadmap” to a multi-dim. picture of N

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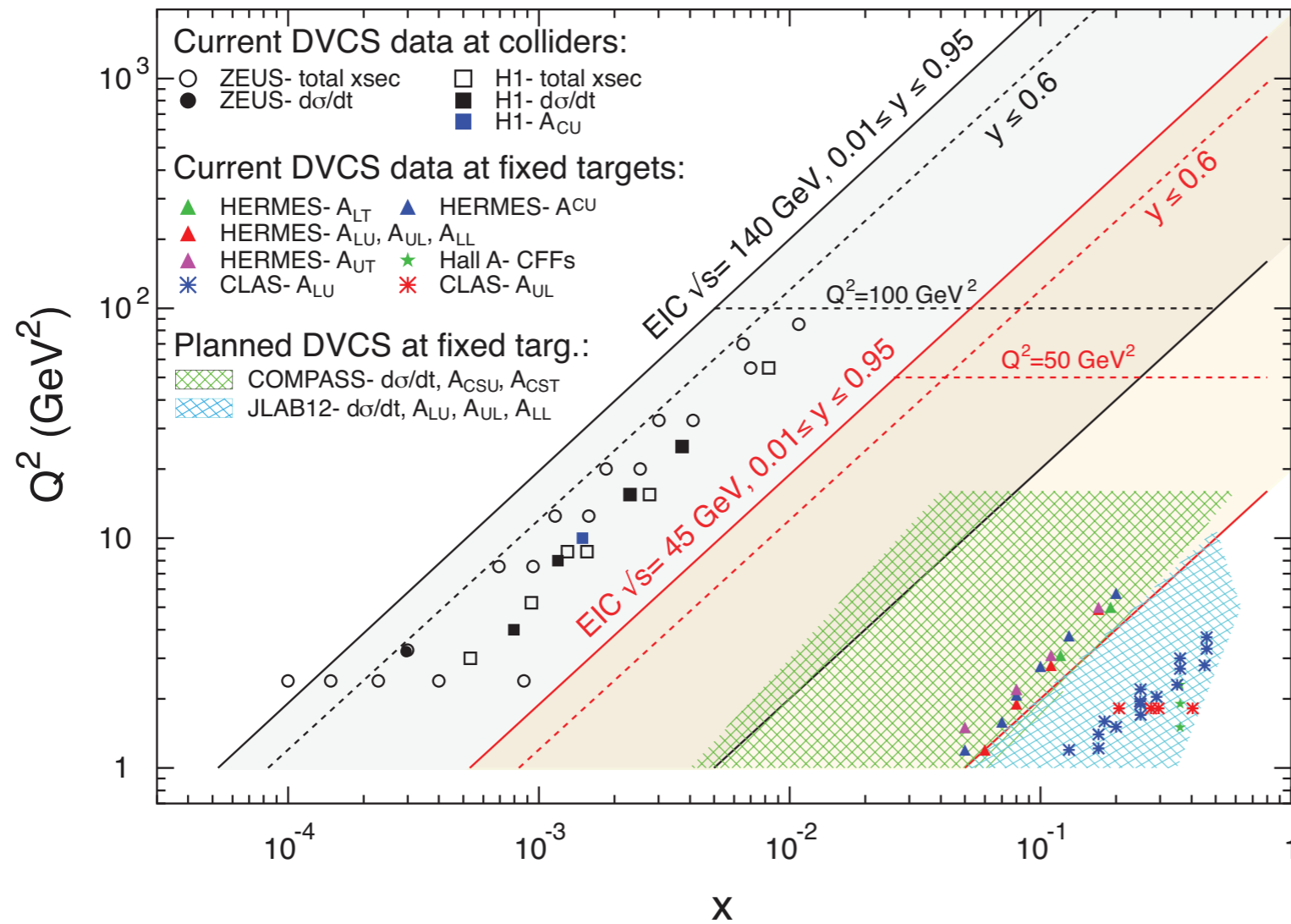


©E. Barone

future directions ?

“Electron Ion Collider:
the Next QCD Frontier”
arXiv:1212.1701 [nucl-ex]

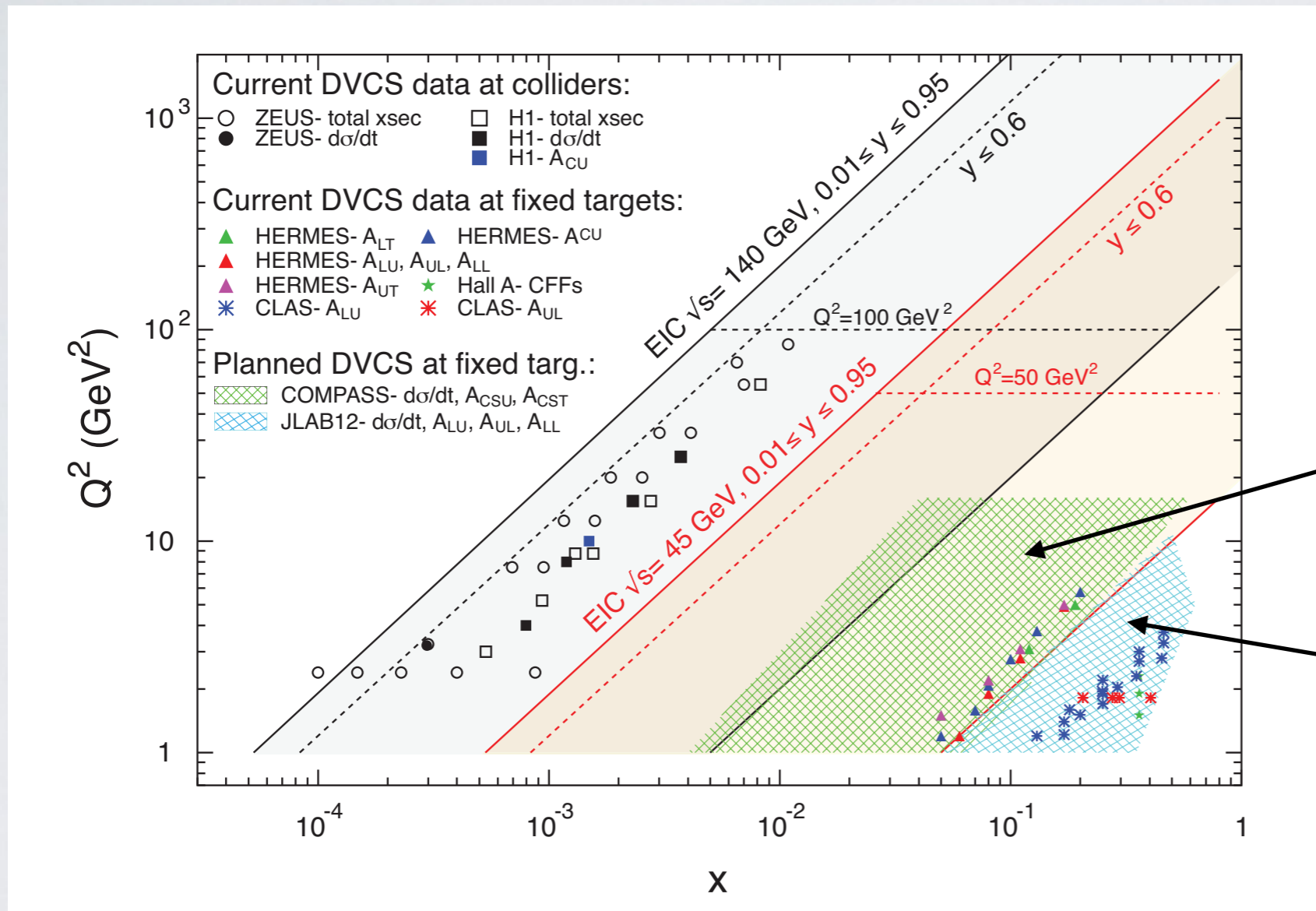
example : DVCS data



future directions ?

“Electron Ion Collider:
the Next QCD Frontier”
arXiv:1212.1701 [nucl-ex]

example : DVCS data



also
polarized
Drell-Yan

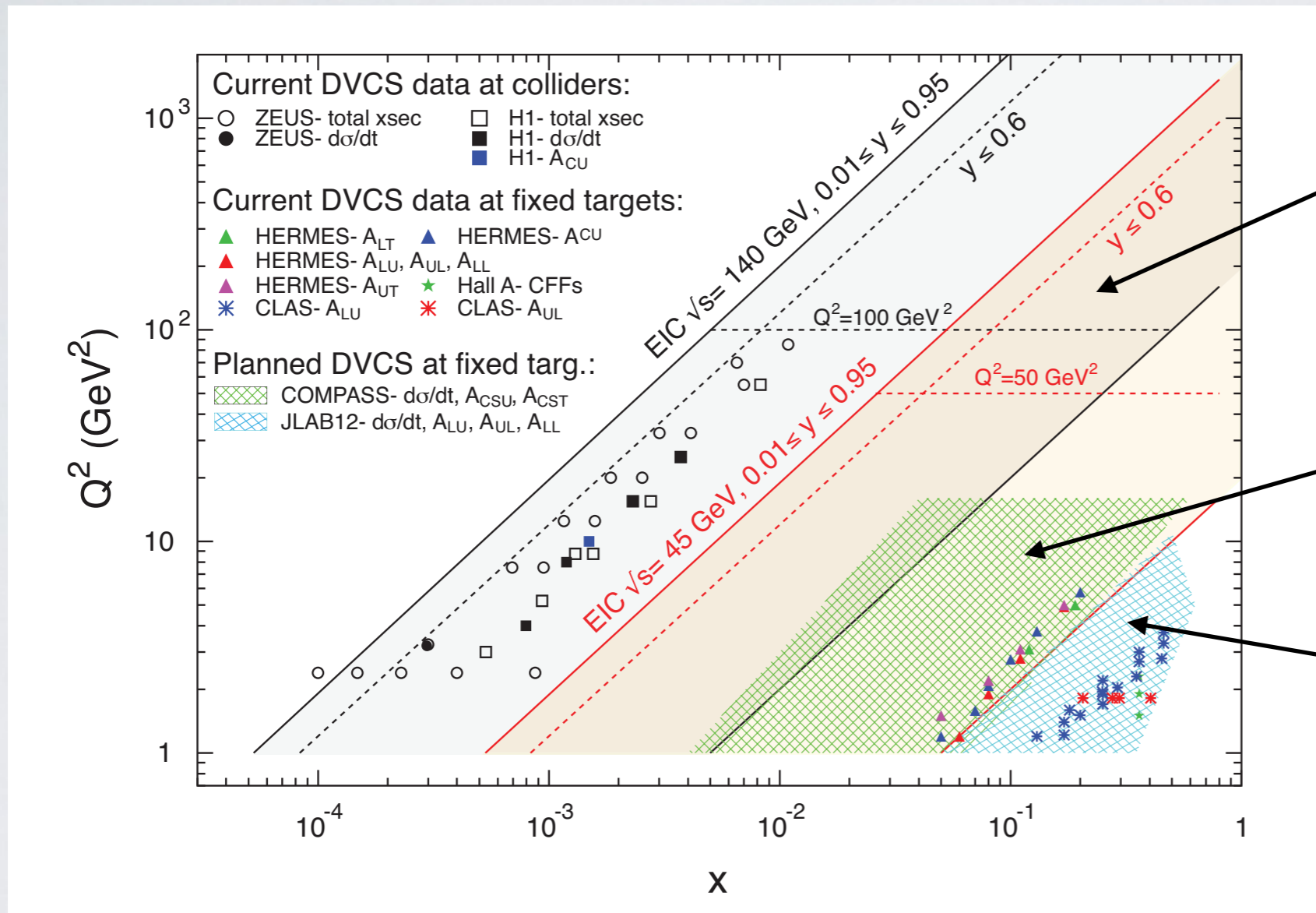


valence x

future directions ?

“Electron Ion Collider:
the Next QCD Frontier”
arXiv:1212.1701 [nucl-ex]

example : DVCS data



EIC
small x



also
polarized
Drell-Yan

Jefferson Lab @12 GeV

valence x

LHeC even smaller x, but no polarization...

“ With 3D projections, we will be entering a new age. Something which was never technically possible before: a stunning visual experience which ‘turbocharges’ the viewing. ”

James Cameron

