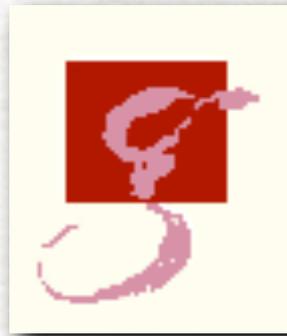


Binary source modelling for ET

Luciano Rezzolla



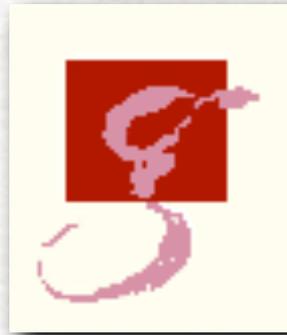
Albert Einstein Institute, Potsdam, Germany

Dept. of Physics and Astronomy, Louisiana State Univ. Louisiana, USA

ET-ILIAS meeting, Nov. 24-25 2008, Cascina, Italy

What NR can do for ET and viceversa

Luciano Rezzolla



Albert Einstein Institute, Potsdam, Germany

Dept. of Physics and Astronomy, Louisiana State Univ. Louisiana, USA

ET-ILIAS meeting, Nov. 24-25 2008, Cascina, Italy

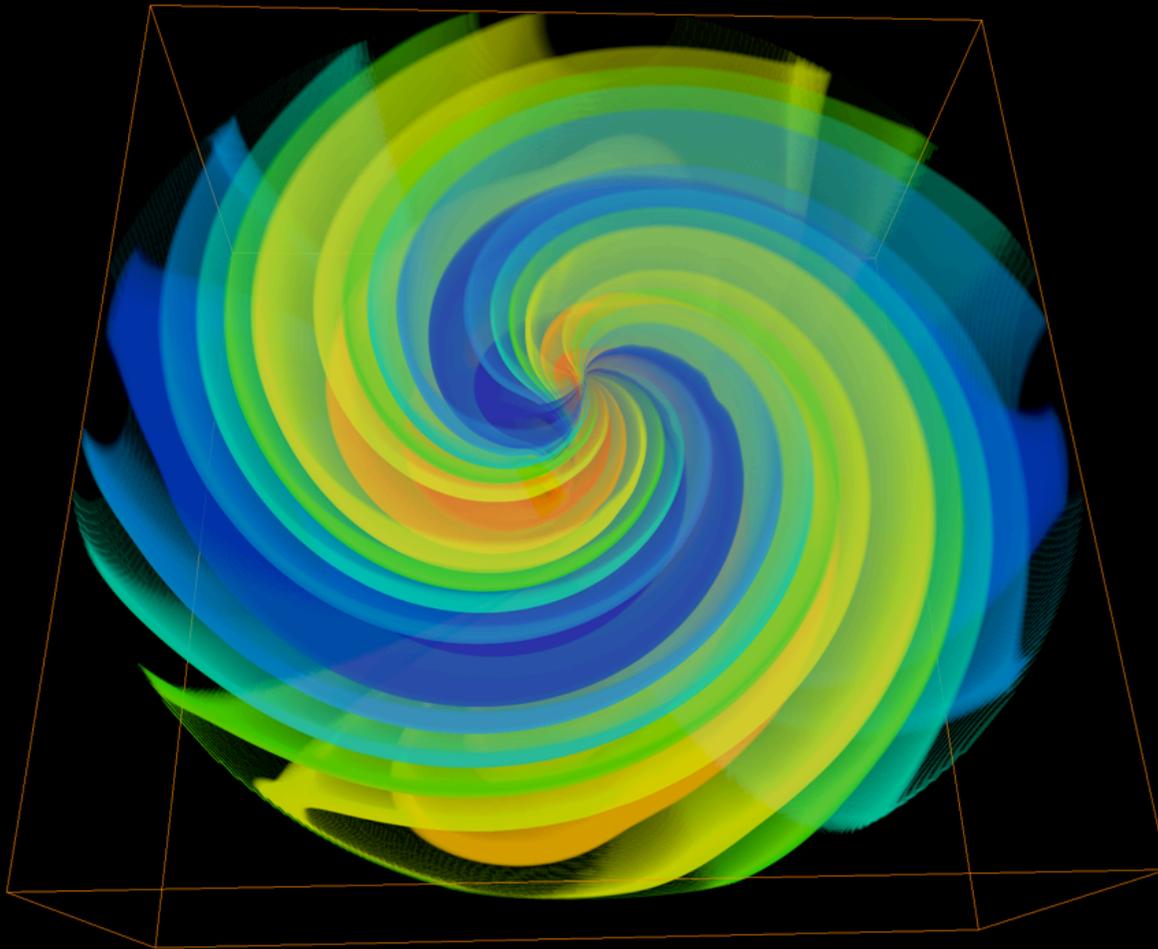


Plan of the talk

- What NR can do for ET
- Some representative examples
 - binary black holes:
 - hybrid waveforms for unequal-mass non-spinning
 - SNR for equal-mass spinning
 - binary neutron stars:
 - deciphering the EOS
 - connecting GW and EM emission from GRBs

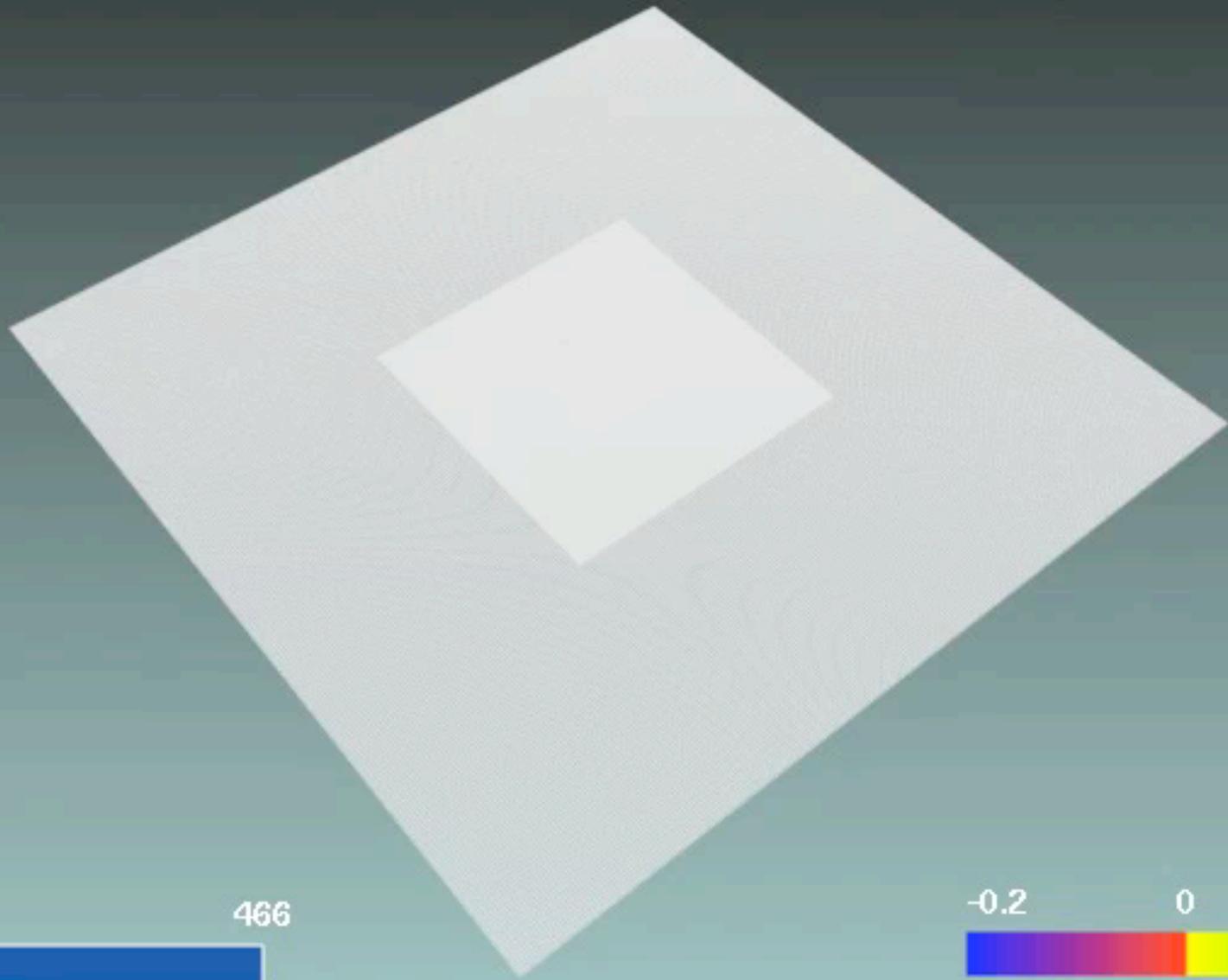
Modelling binary black holes

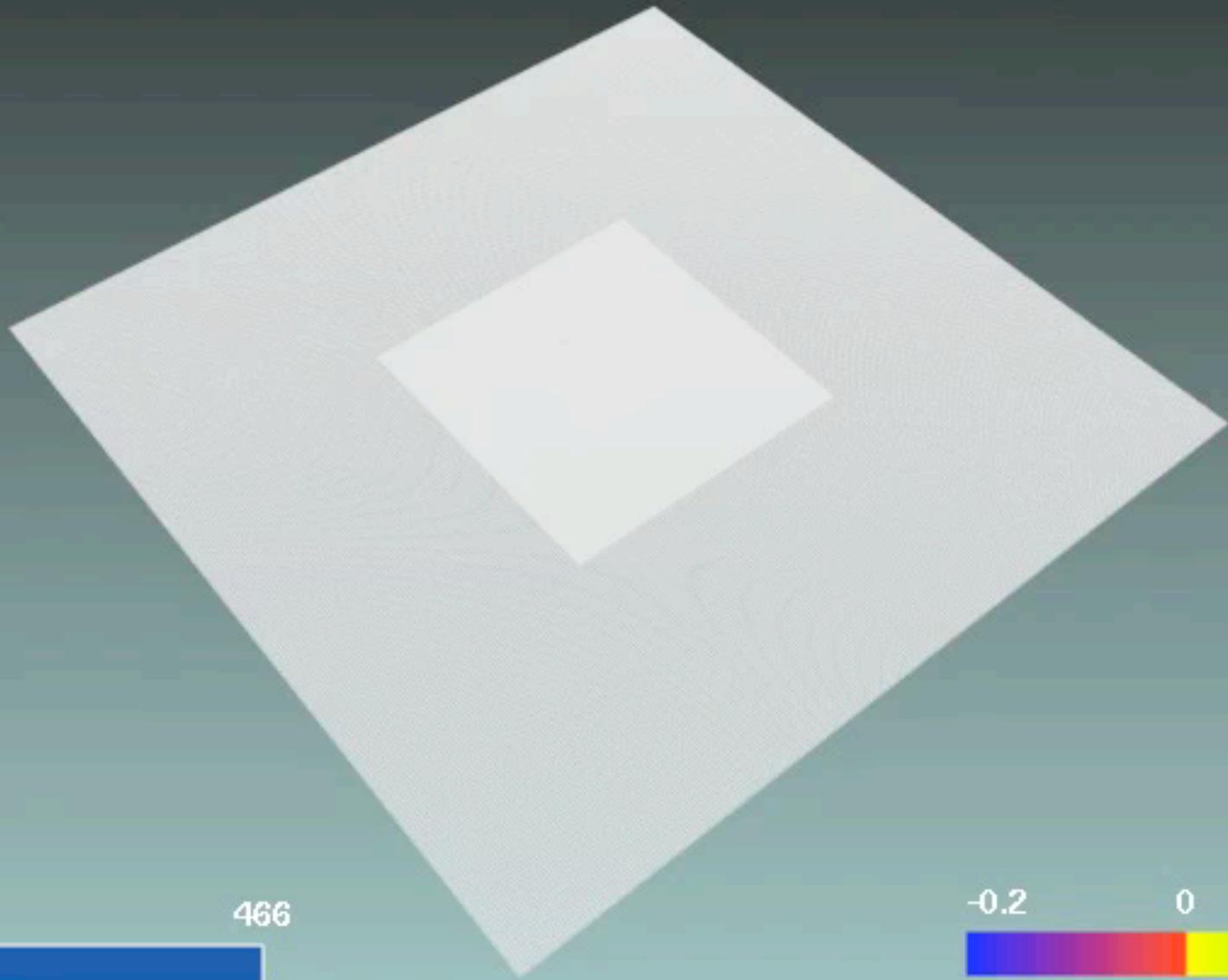
Ajith et al. 2007, CQG
Ajith et al. 2008, PRD
Riesswig et al. 2008



In collaboration with:

- GW group AEI Hannover
- GW group AEI Potsdam
- Numrel group Jena





Hybrid waveforms: the motivations

- There is no fundamental obstacle to long-term (i.e. covering $\sim 10+$ orbits) NR calculations of the three stages of the binary evolution: **inspiral**, **merger** and **ringdown**
- Yet, NR simulations are computationally expensive and building a template bank out of them is prohibitive and awkward
- Present data-analysis pipelines employ phenomenological template: analytic, fast, and collect most of the information
- In the past phenomenological templates have been built using approximations and a certain amount of heuristics
- Ideally phenomenological templates should be built upon modelling NR waveforms for the three stages.

The basic logic...

The basic logic...

PN waveform

+

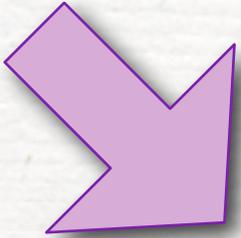
NR waveform

The basic logic...

PN waveform

+

NR waveform



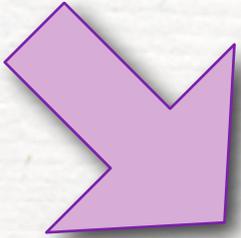
hybrid waveform

The basic logic...

PN waveform

+

NR waveform



hybrid waveform

+

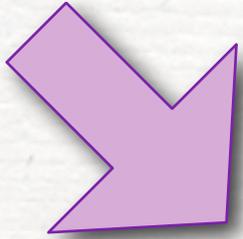
10D phenomenological waveform

The basic logic...

PN waveform

+

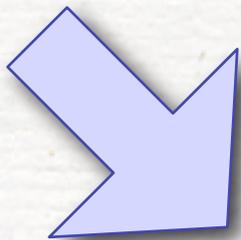
NR waveform



hybrid waveform

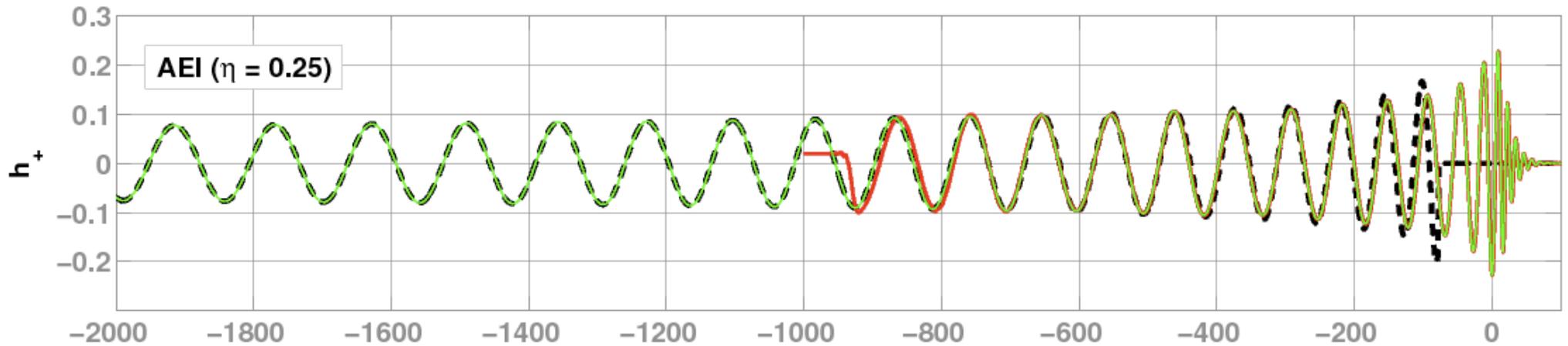
+

10D phenomenological waveform



2D phenomenological waveform

A hybrid waveform in practice...

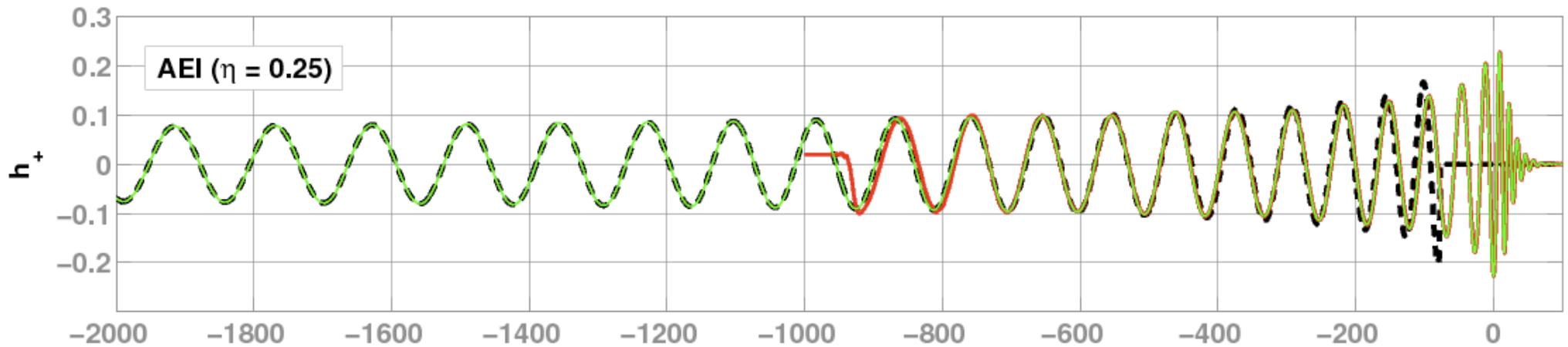


Red line is the numrel waveform

Black dashed line is the 3.5PN waveform

Green line is the hybrid waveform

A hybrid waveform in practice...



Red line is the numrel waveform

Black dashed line is the 3.5PN waveform

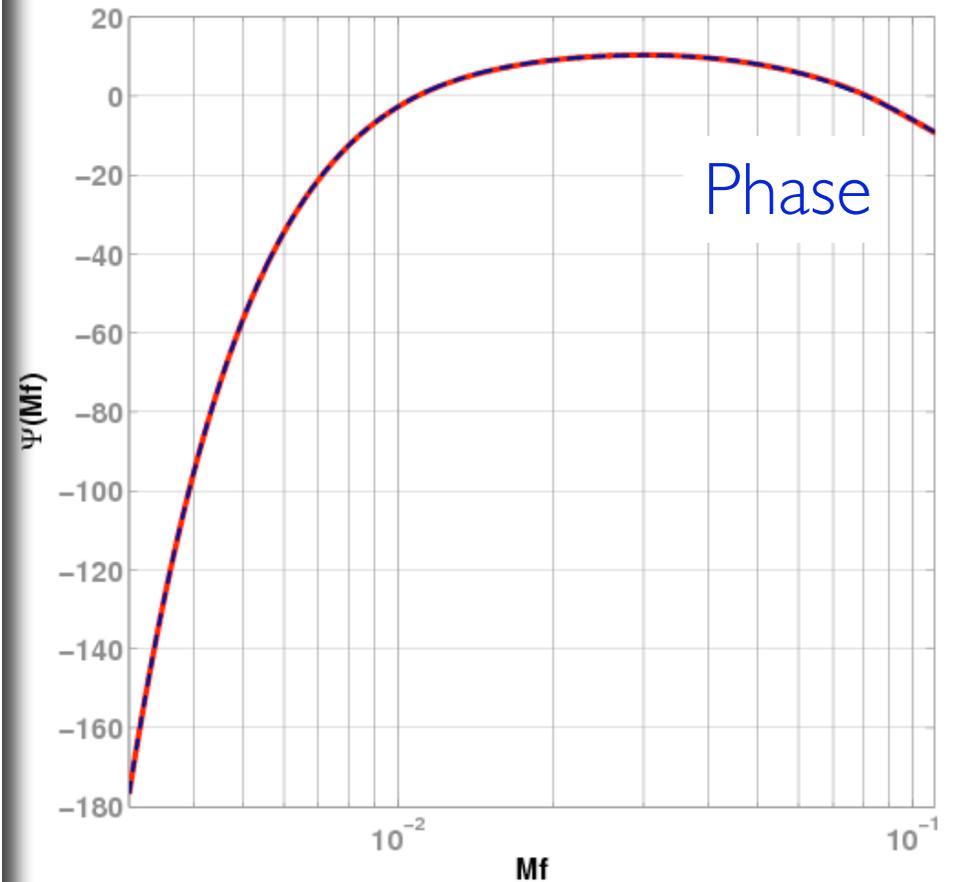
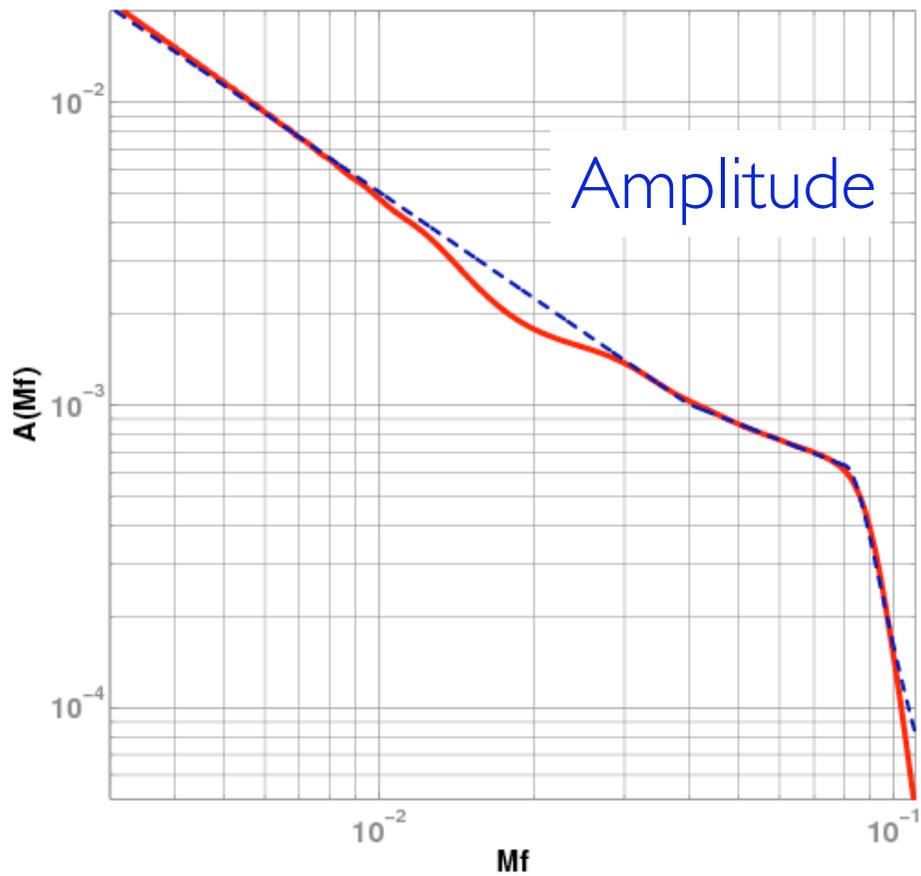
Green line is the hybrid waveform

Once the hybrid waveform is computed, it can be parametrized in the Fourier domain via 10 phenomenological parameters (4 for the amplitude, 6 for the phase).

A suitable fitting procedure allows one to reduce them to the 2 physical ones: M mass of the binary and the symmetric mass-ratio ν

Solid lines are hybrid waveforms

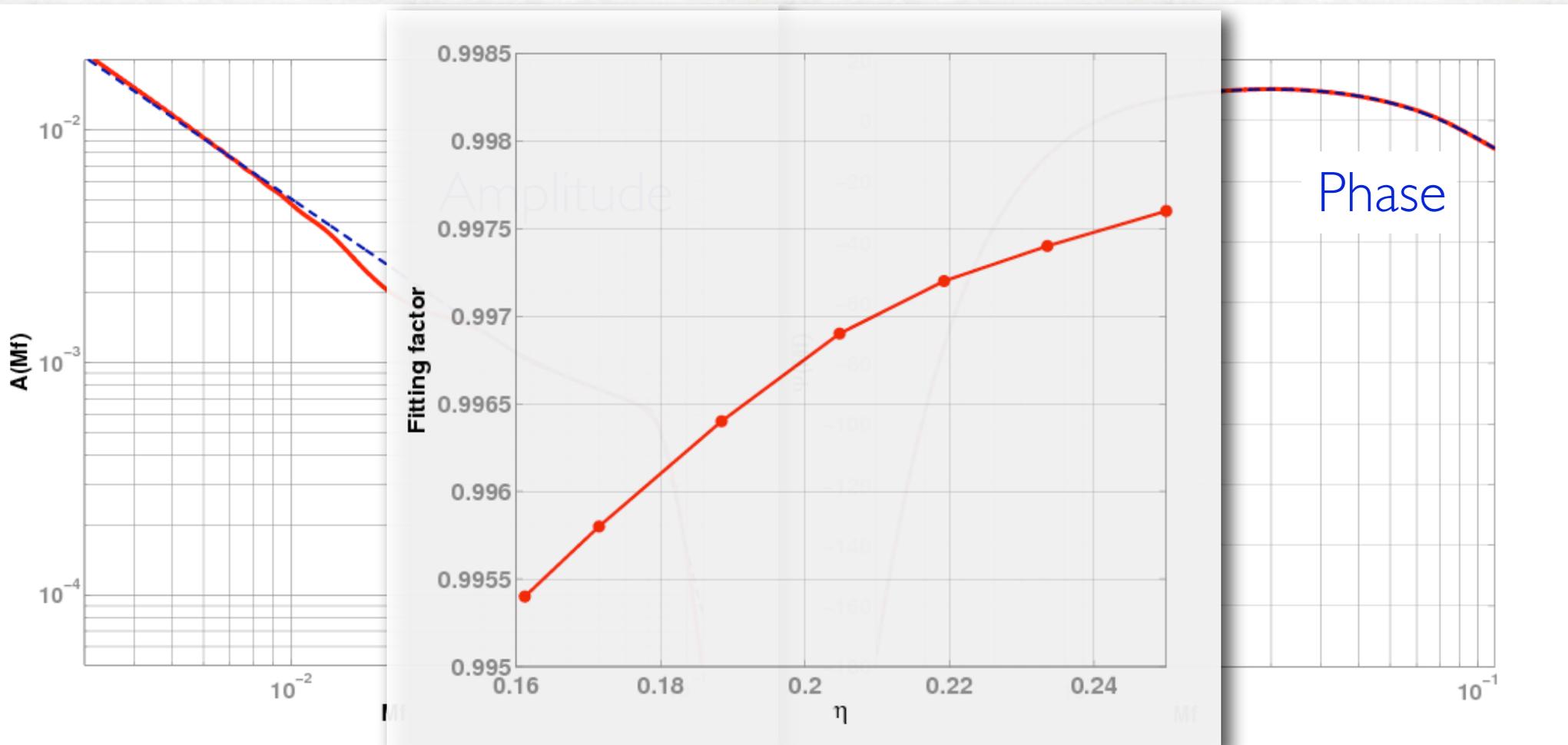
Dashed lines are phenomenological waveforms



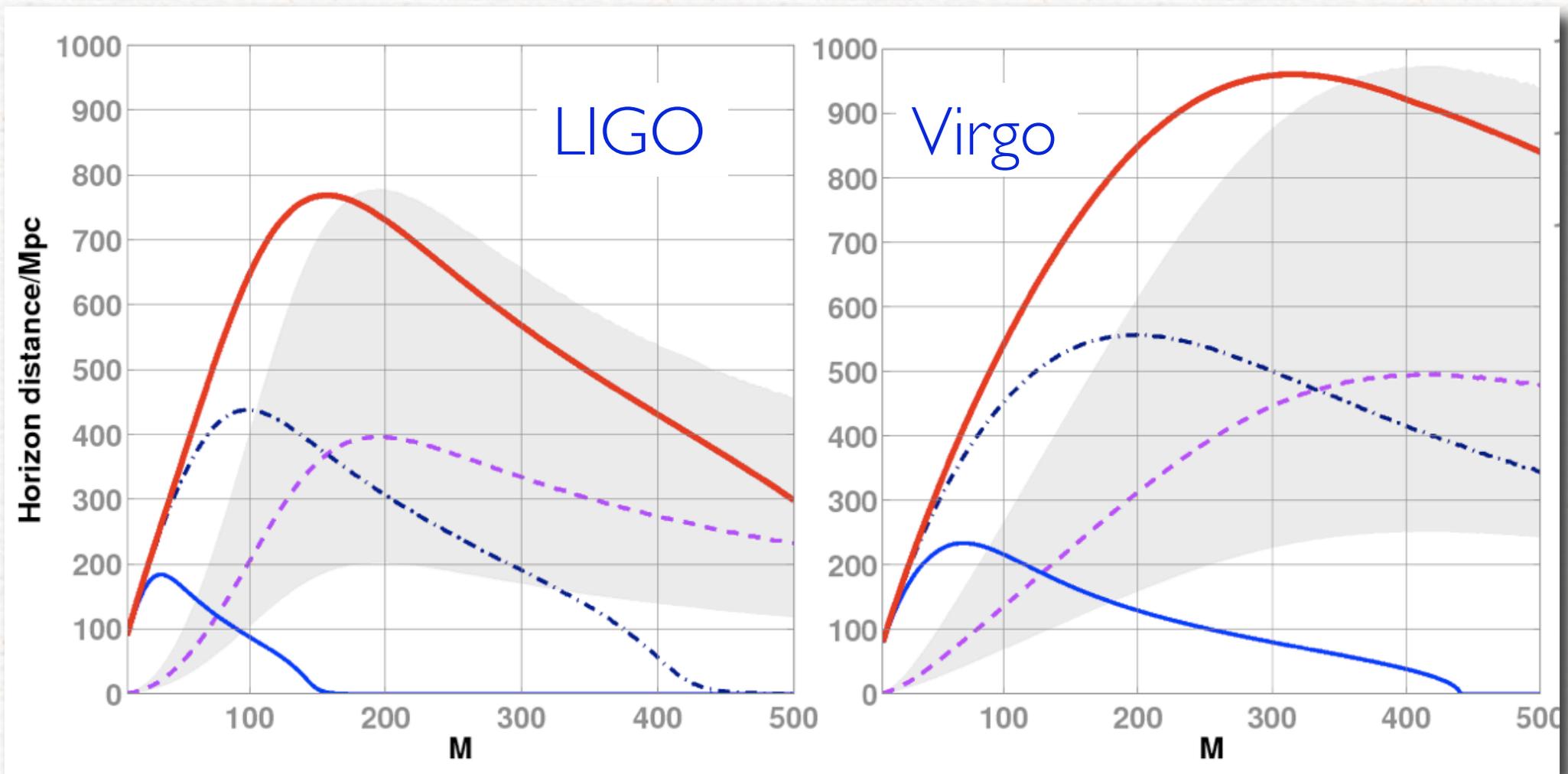
Solid lines are hybrid waveforms

Dashed lines are phenomenological waveforms

Fitting Factor: measures deviations between two waveforms; should be larger than 97%; here is $> 99\%$



What is this good for?



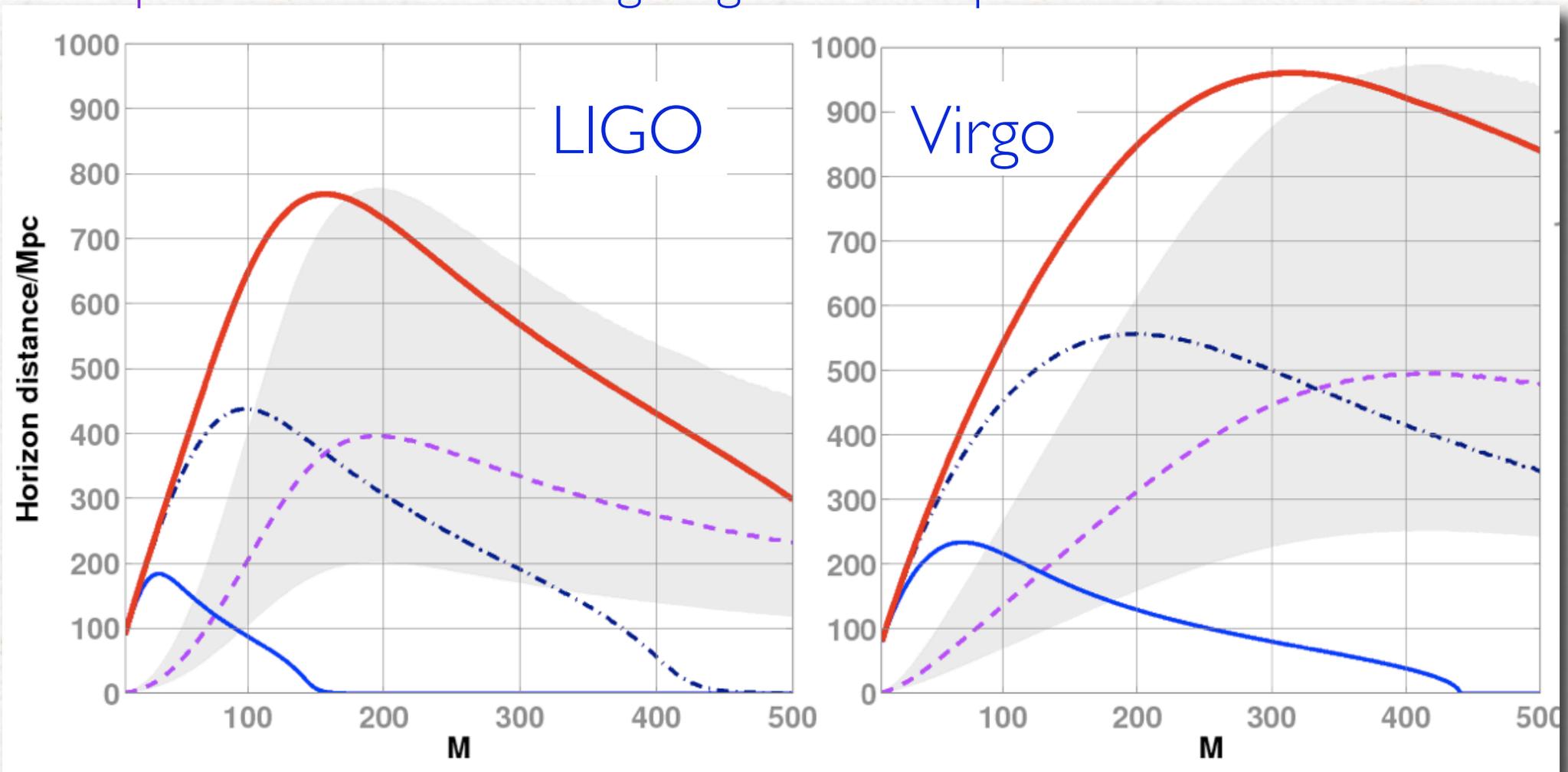
What is this good for?

Red line is the complete (inspiral, merger, ringdown) template

Blue line is the PN template truncated at ISCO

Black dot-dashed line is EOB template truncated at light-ring

Purple dashed line is using ringdown templates

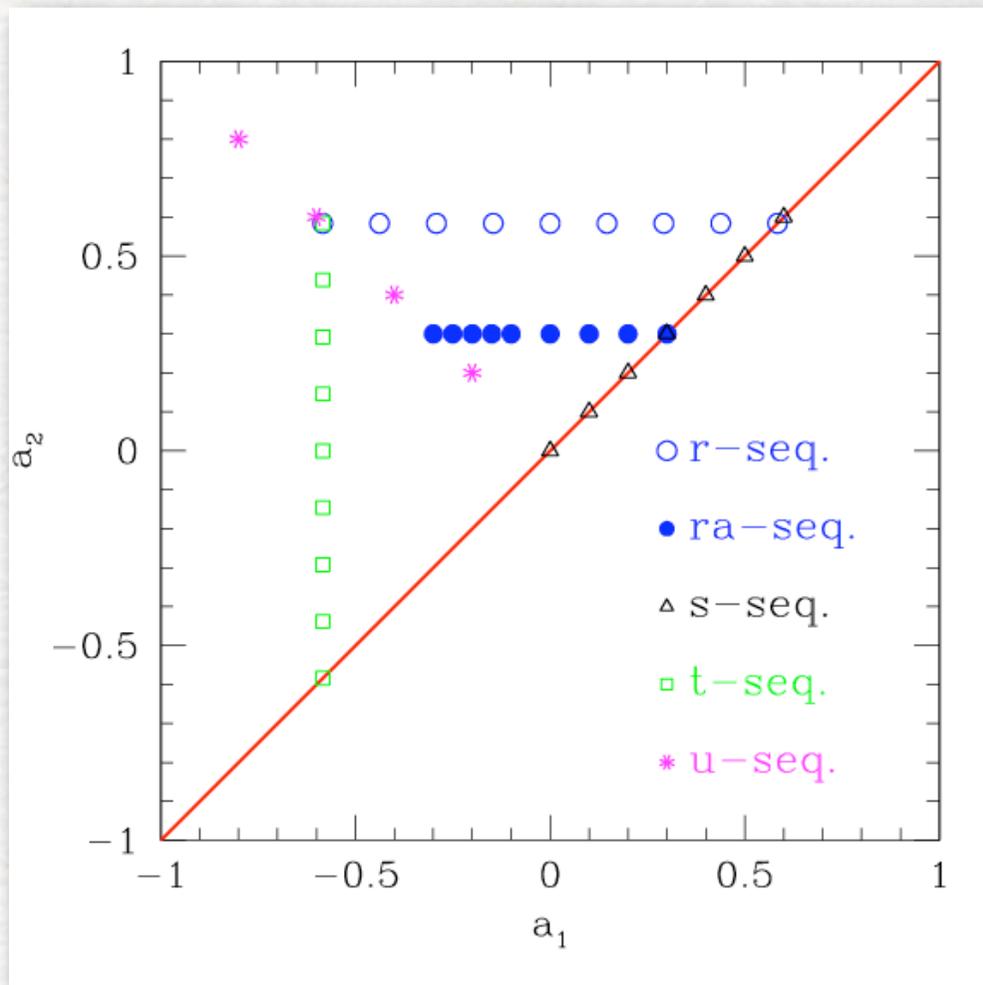


SNR for equal-mass **spinning** binaries

Given a precise portion of the space of parameters (e.g. equal-mass, aligned spinning binaries) it is possible to compute generic SNRs by “stitching” PN, and NR waveforms.

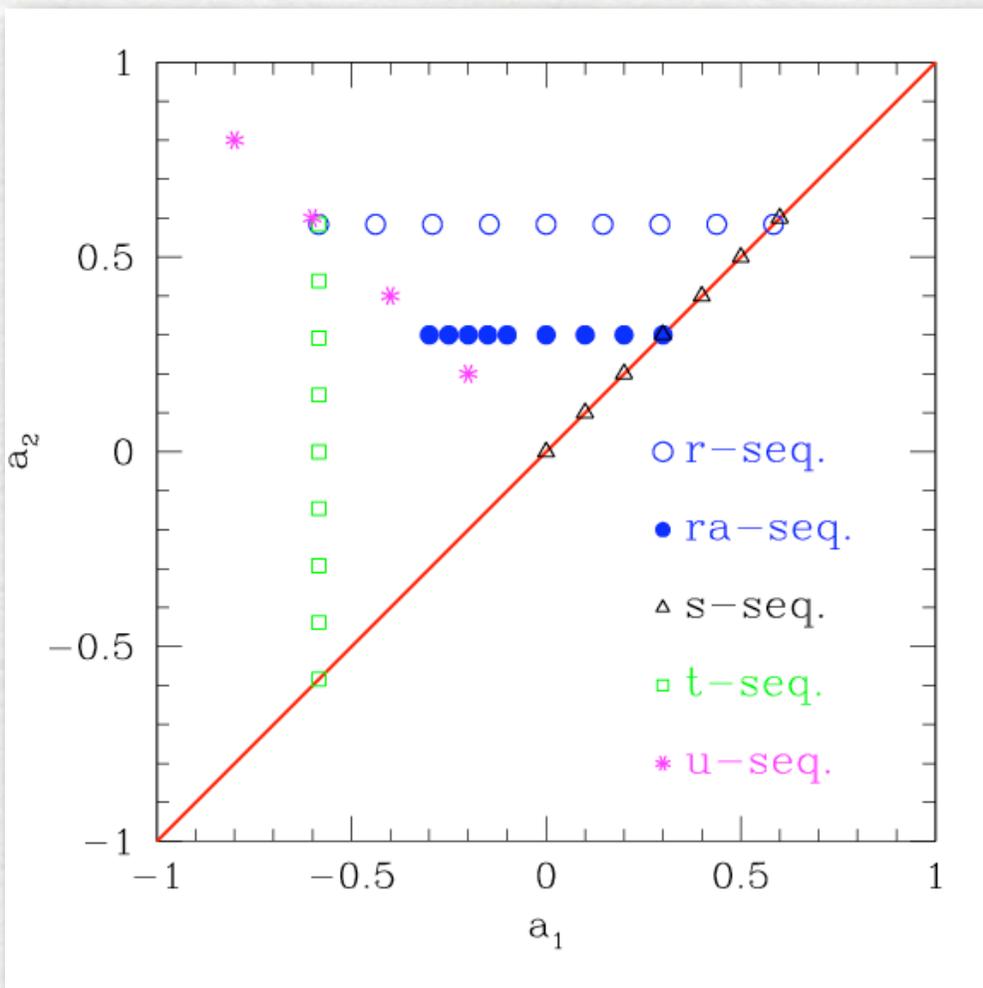
SNR for equal-mass **spinning** binaries

Given a precise portion of the space of parameters (e.g. equal-mass, aligned spinning binaries) it is possible to compute generic SNRs by “stitching” PN, and NR waveforms.



SNR for equal-mass **spinning** binaries

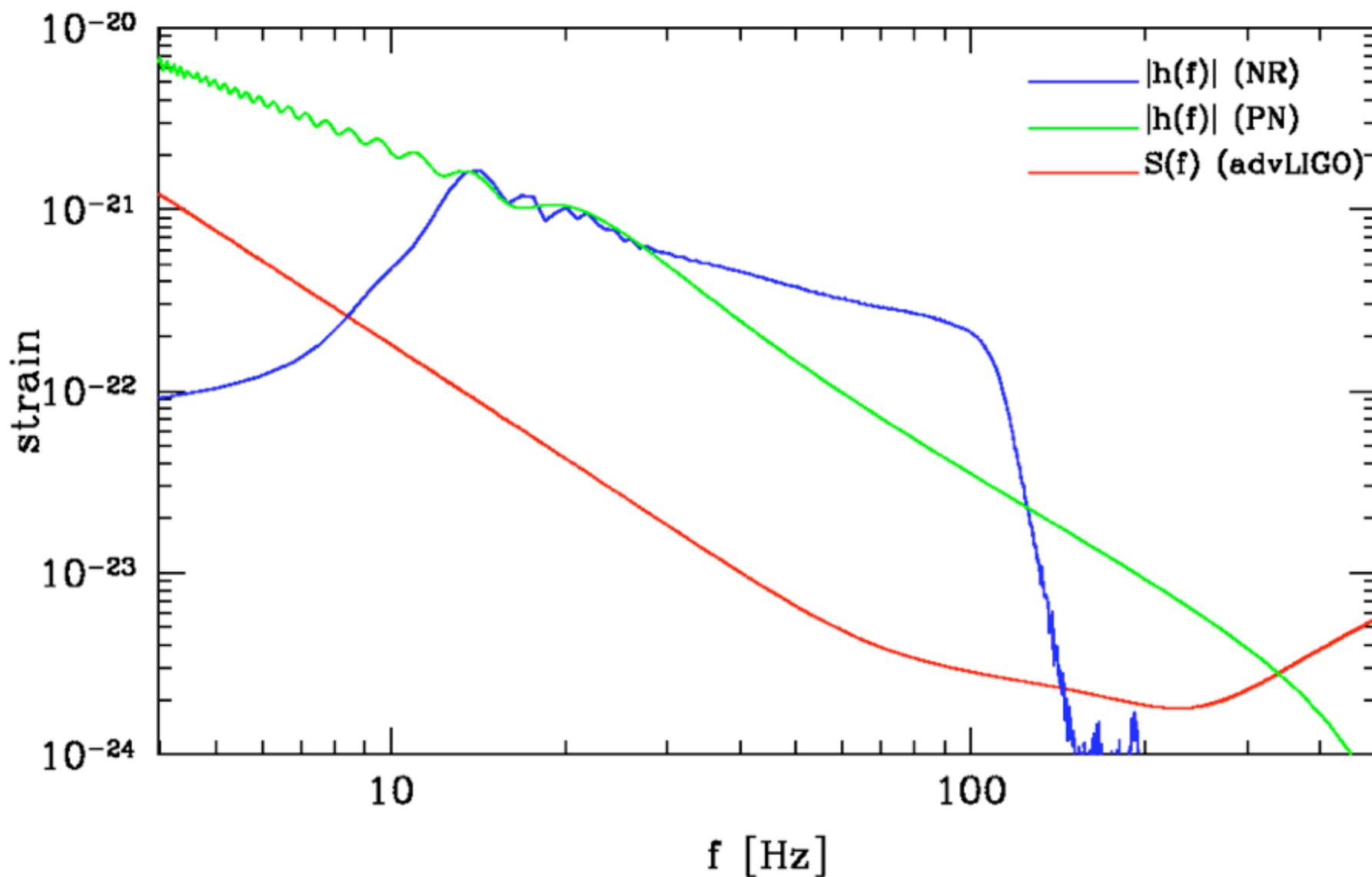
Given a precise portion of the space of parameters (e.g. equal-mass, aligned spinning binaries) it is possible to compute generic SNRs by “stitching” PN, and NR waveforms.



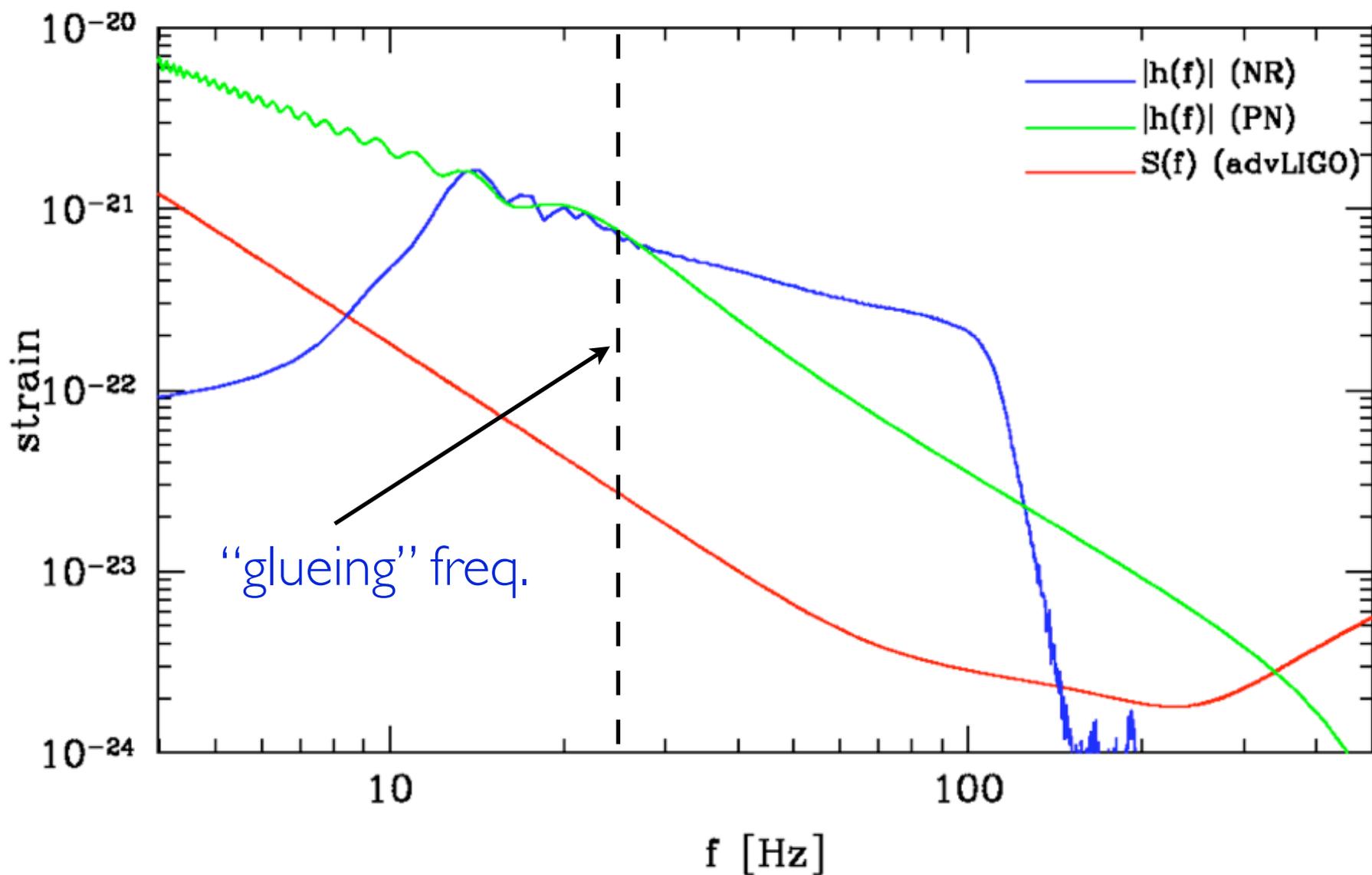
In practice, we cover the “**spin diagram**” with suitably chosen sequences and evaluate the SNR for different detectors and different masses.

Particularly relevant is the SW-NE diagonal: $a_1 = a_2$ and thus spanning the case of **accelerated inspiral** $a_1 = a_2 < 0$ and of **decelerated inspiral** (“hang-up”) $a_1 = a_2 > 0$

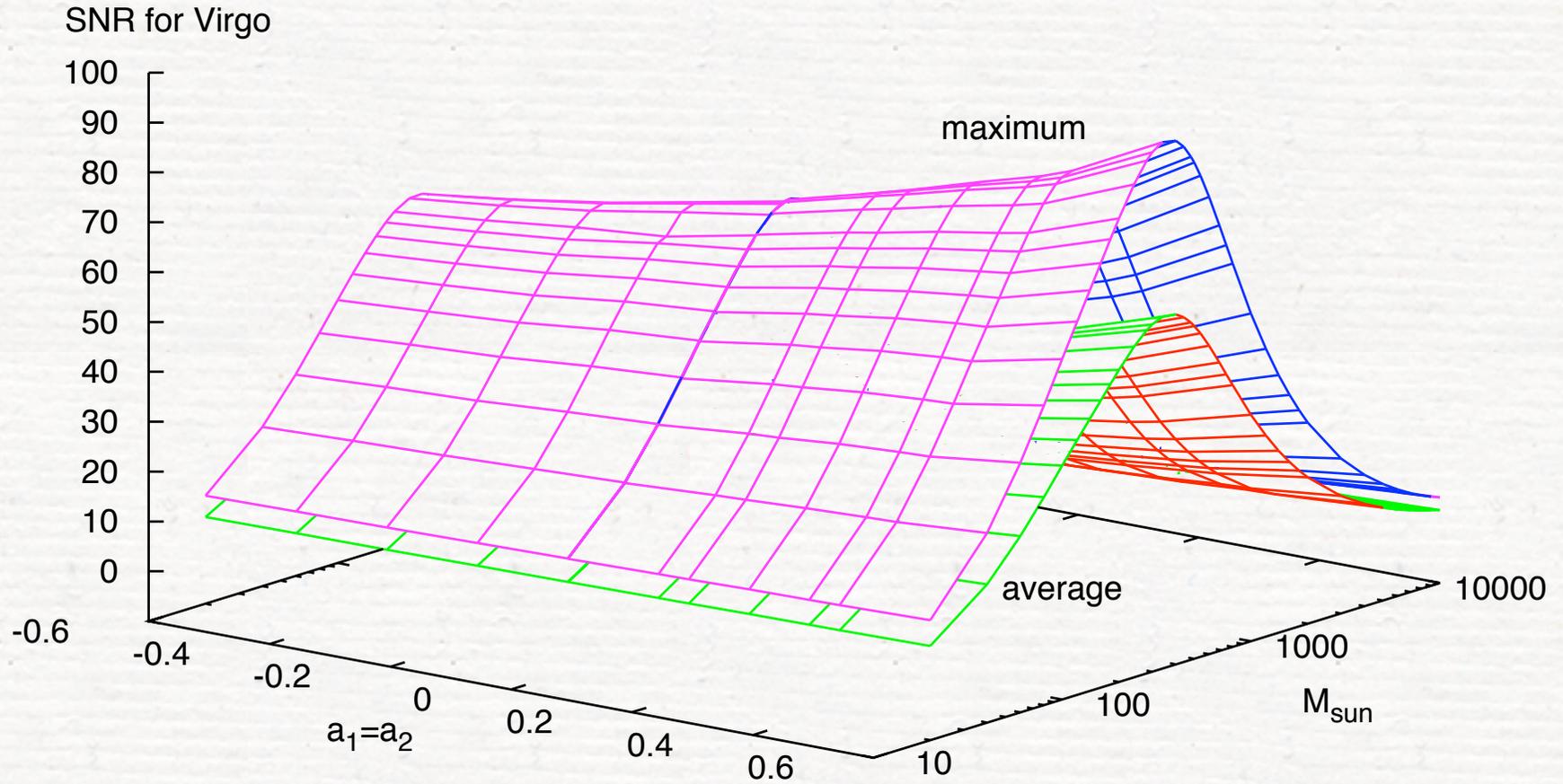
The "stitching" is made in **frequency space** and at a precise "glueing" frequency. Proper windowing removes additional high-and low-frequency artifacts from the NR waveform.



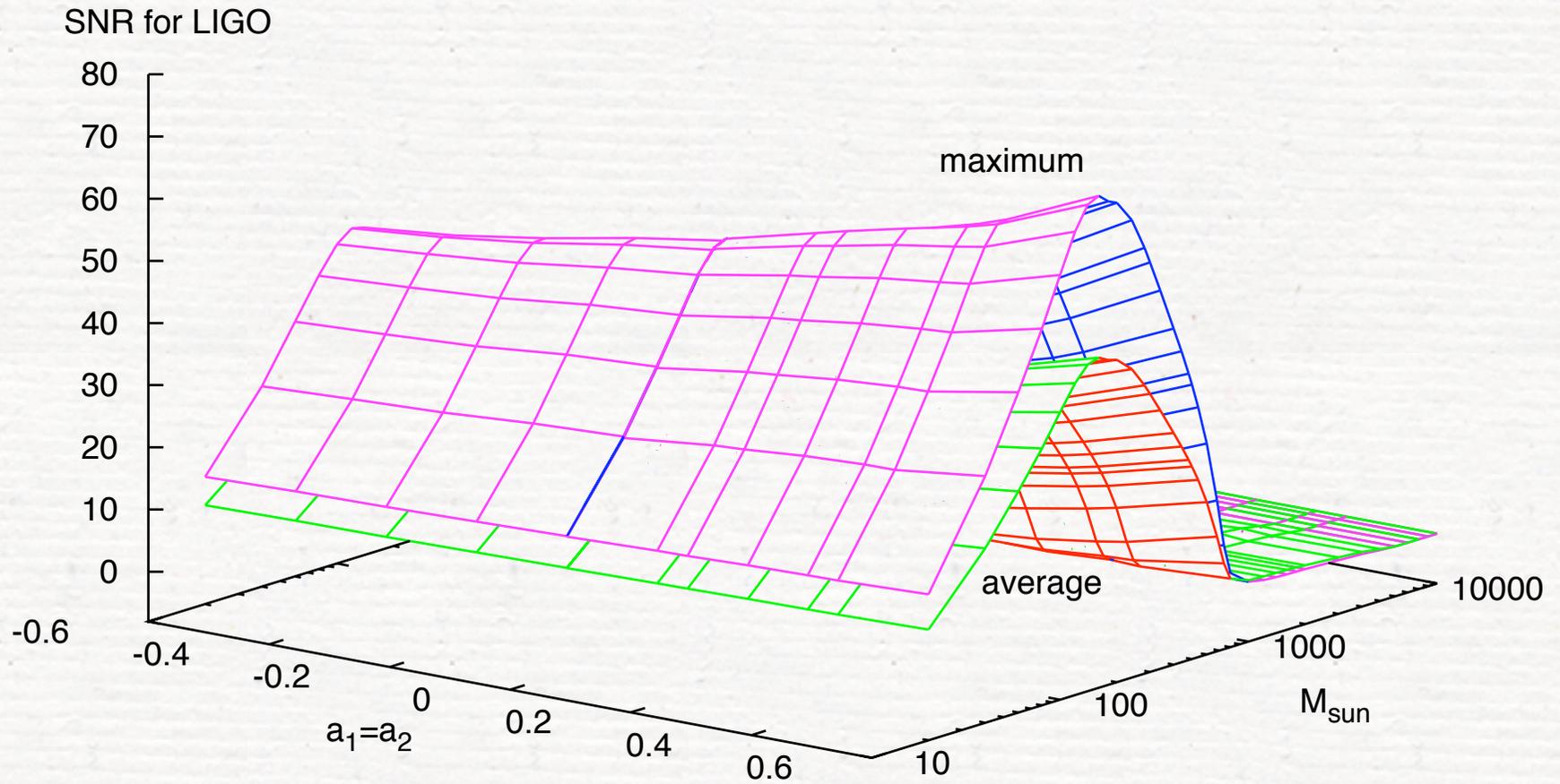
The "stitching" is made in frequency space and at a precise "glueing" frequency. Proper windowing removes additional high-and low-frequency artifacts from the NR waveform.



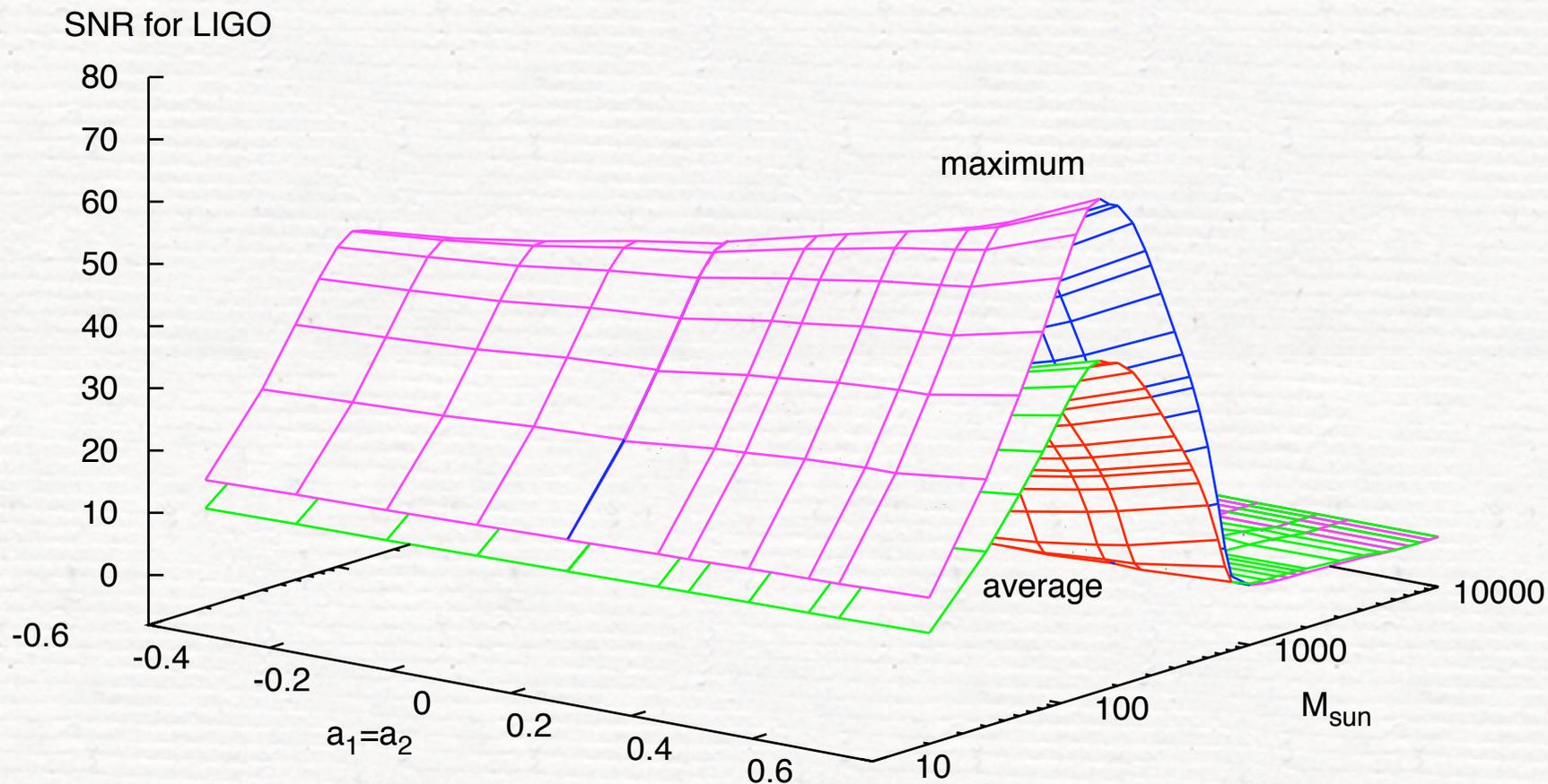
The end result can be shown as a function of the **initial spins** and **masses** for each detector



The end result can be shown as a function of the **initial spins** and **masses** for each detector



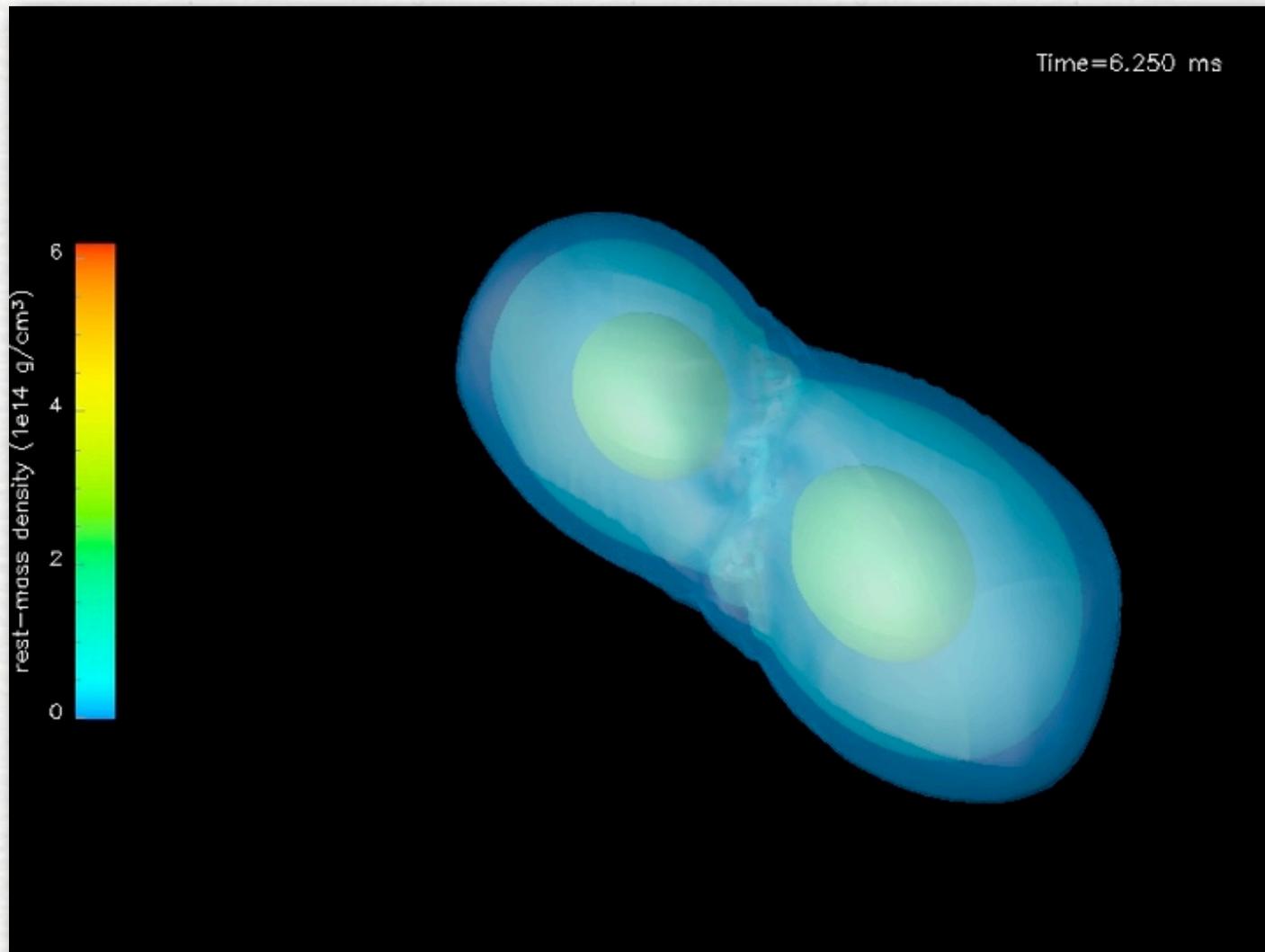
The end result can be shown as a function of the **initial spins** and **masses** for each detector



Overall, large **aligned spins** $a_1 = a_2 \sim 1$ have a SNR which is twice as large as the **antialigned spins** $a_1 = a_2 \sim -1$. The explanation is obvious: \sim four times the number of cycles

Modelling binary neutron stars

Baiotti, et al. 2008, PRD

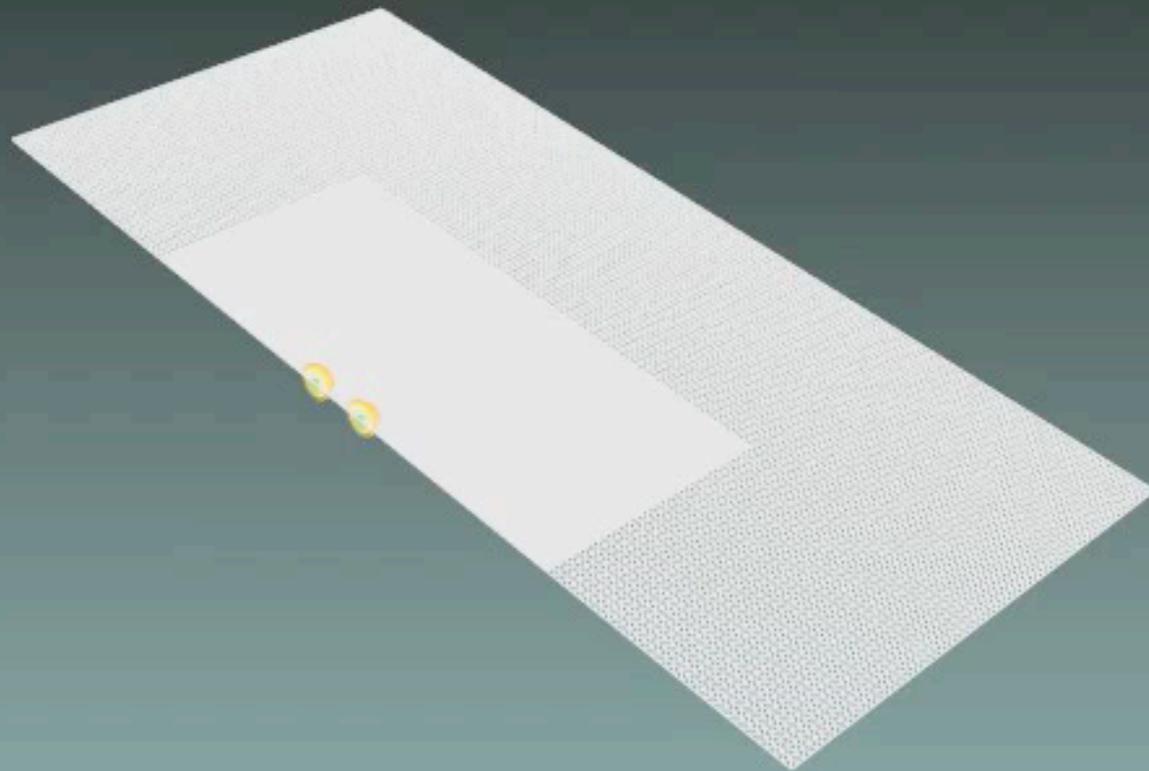


Polytropic EOS, 1.6 M_{\odot}

T[ms] = 0.00



T[M] = 0.00



0.0

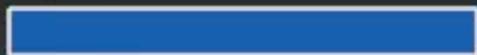
6.1E+14



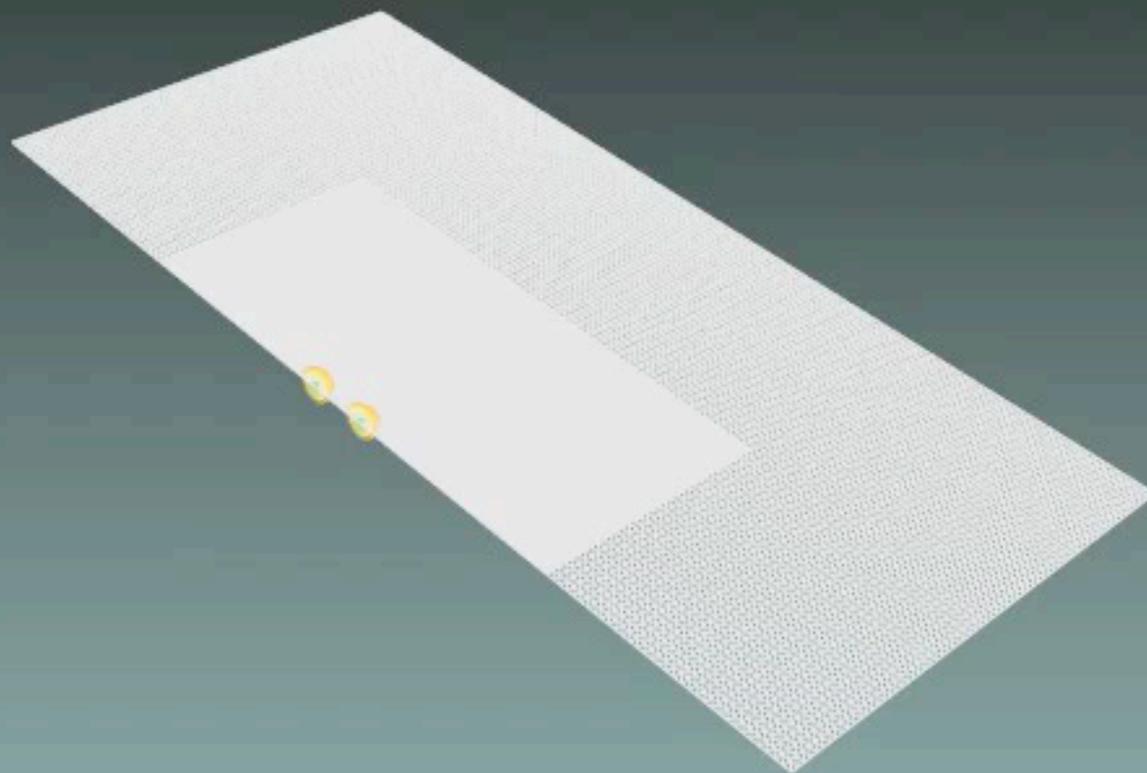
Density [g/cm³]

Polytropic EOS, 1.6 M_⊙

T[ms] = 0.00



T[M] = 0.00



0.0

6.1E+14



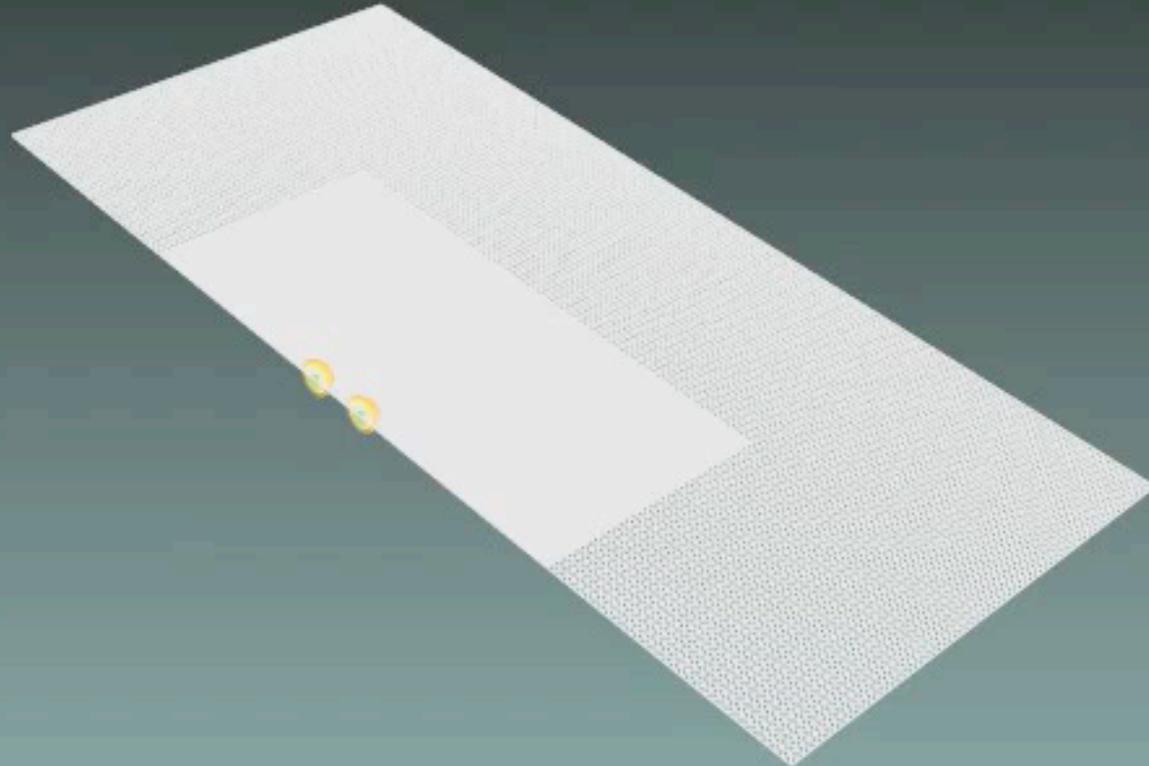
Density [g/cm³]

T[ms] = 0.00



T[M] = 0.00

A hot, low-density torus is produced orbiting around the bh. This is what expected in short GRBs.



0.0

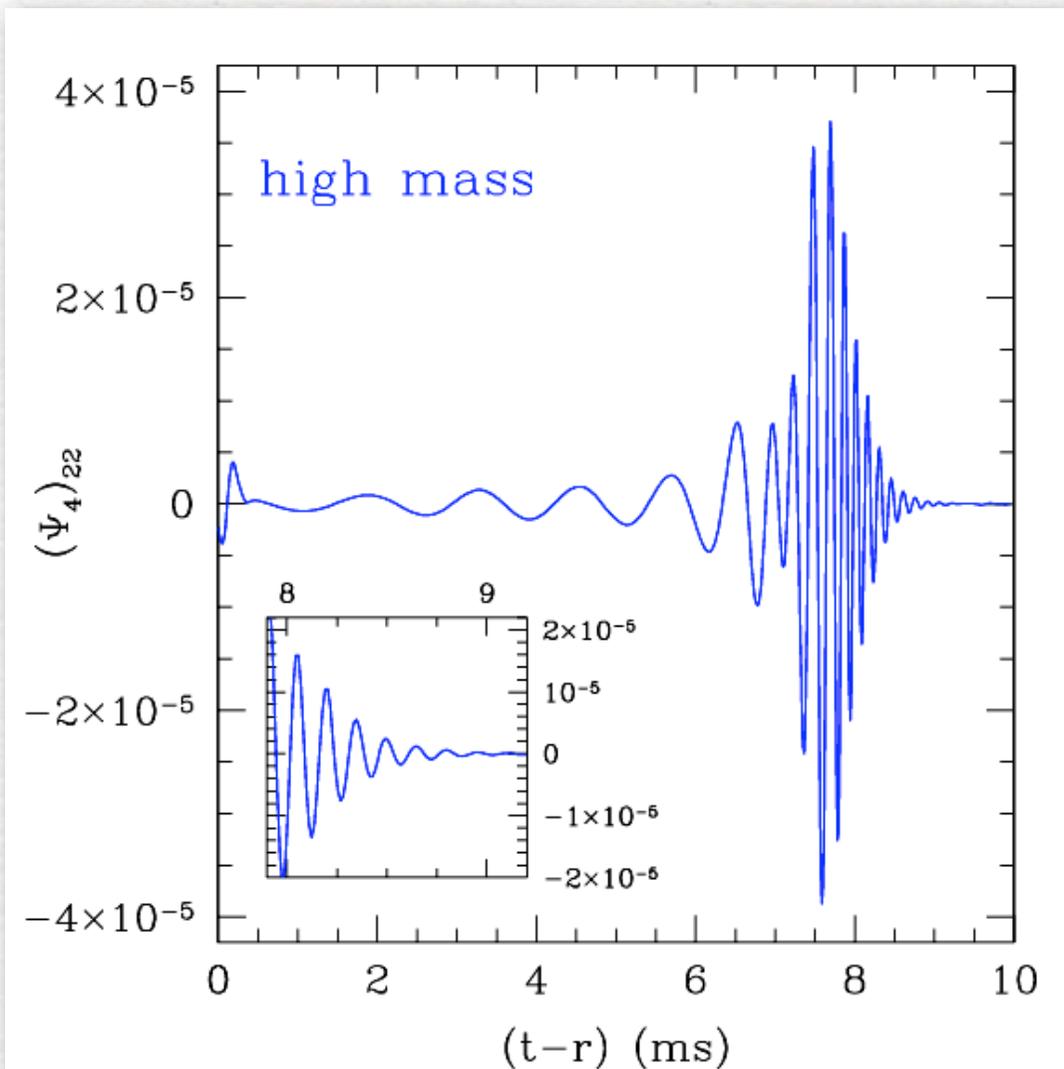
6.1E+14



Density [g/cm³]

Waveforms: polytropic EOS

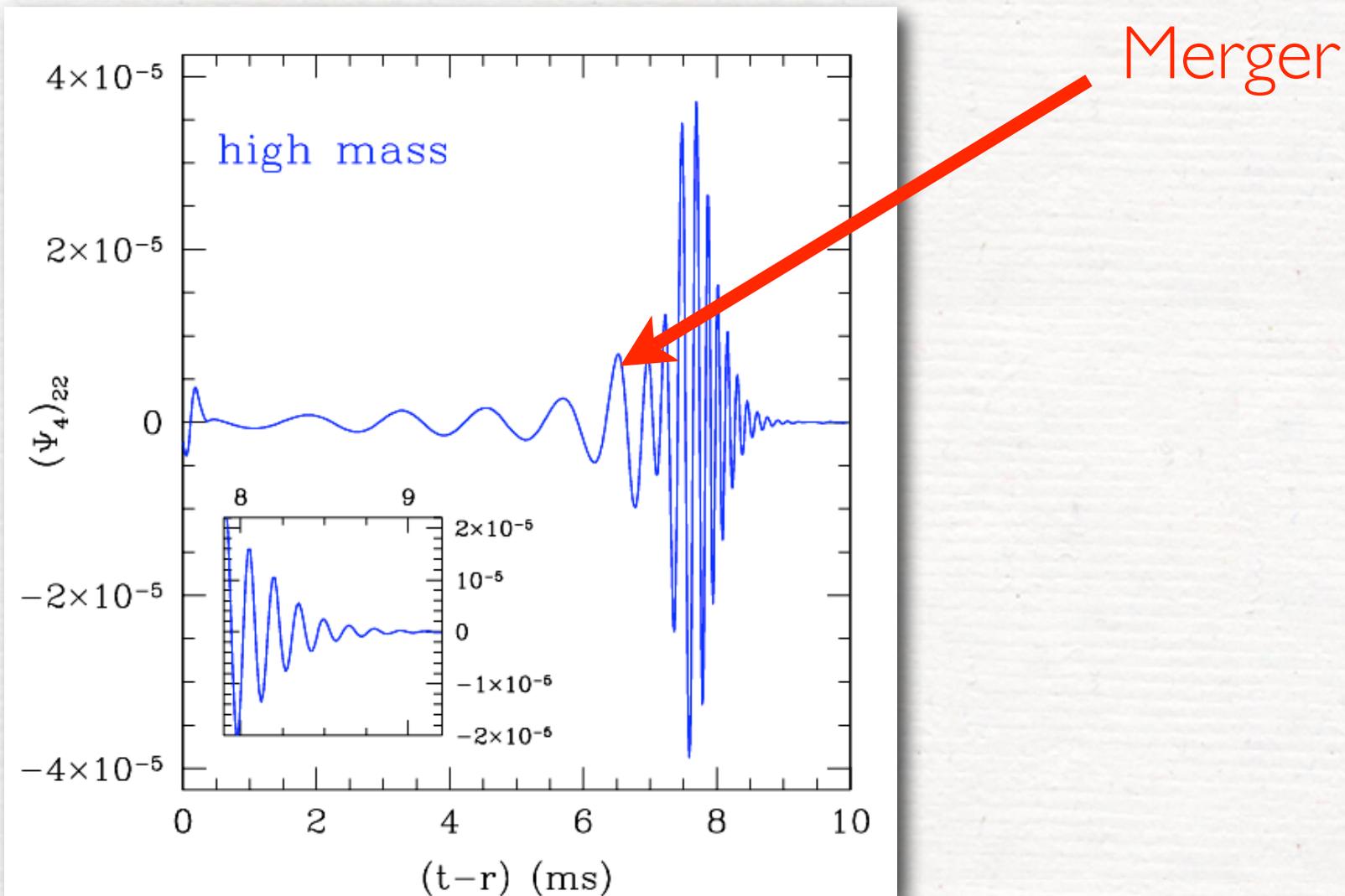
high-mass binary



first time the full signal from the formation to a bh has been computed

Waveforms: polytropic EOS

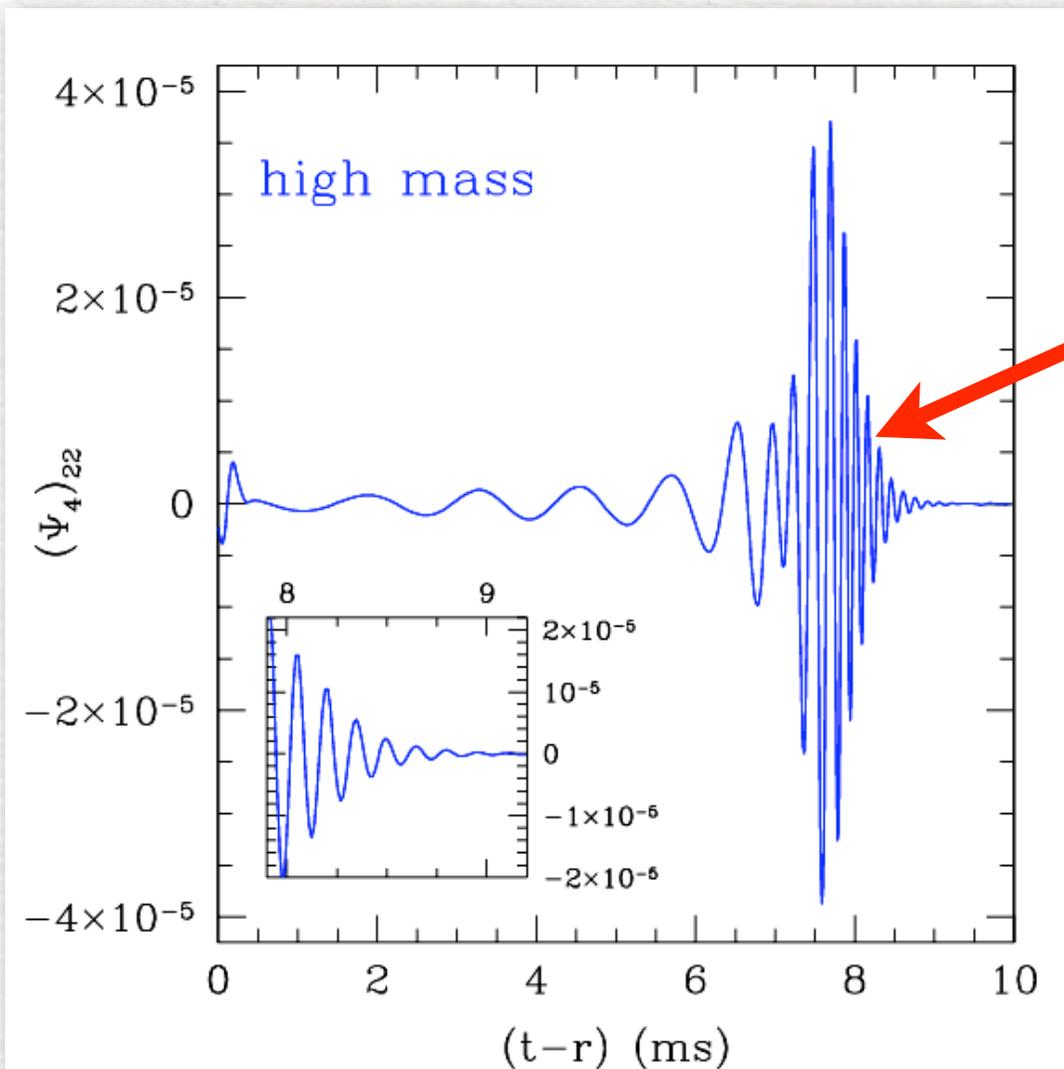
high-mass binary



first time the full signal from the formation to a bh has been computed

Waveforms: polytropic EOS

high-mass binary



Collapse
to BH

first time the full signal from the formation to a bh has been computed

The behaviour:

“merger  *HMNS*  *BH + torus”*

is general but only qualitatively

The behaviour:

“merger → HMNS → BH + torus”

is general but only *qualitatively*

Quantitative differences are produced by:

- **differences in the mass for the same EOS:**

a binary with smaller mass will produce a HMNS which is further away from the stability threshold and will collapse at a later time

The behaviour:

“merger \longrightarrow HMNS \longrightarrow BH + torus”

is general but only qualitatively

Quantitative differences are produced by:

- differences in the **mass** for the same **EOS**:

a binary with smaller mass will produce a HMNS which is further away from the stability threshold and will collapse at a later time

- differences in the **EOS** for the same **mass**:

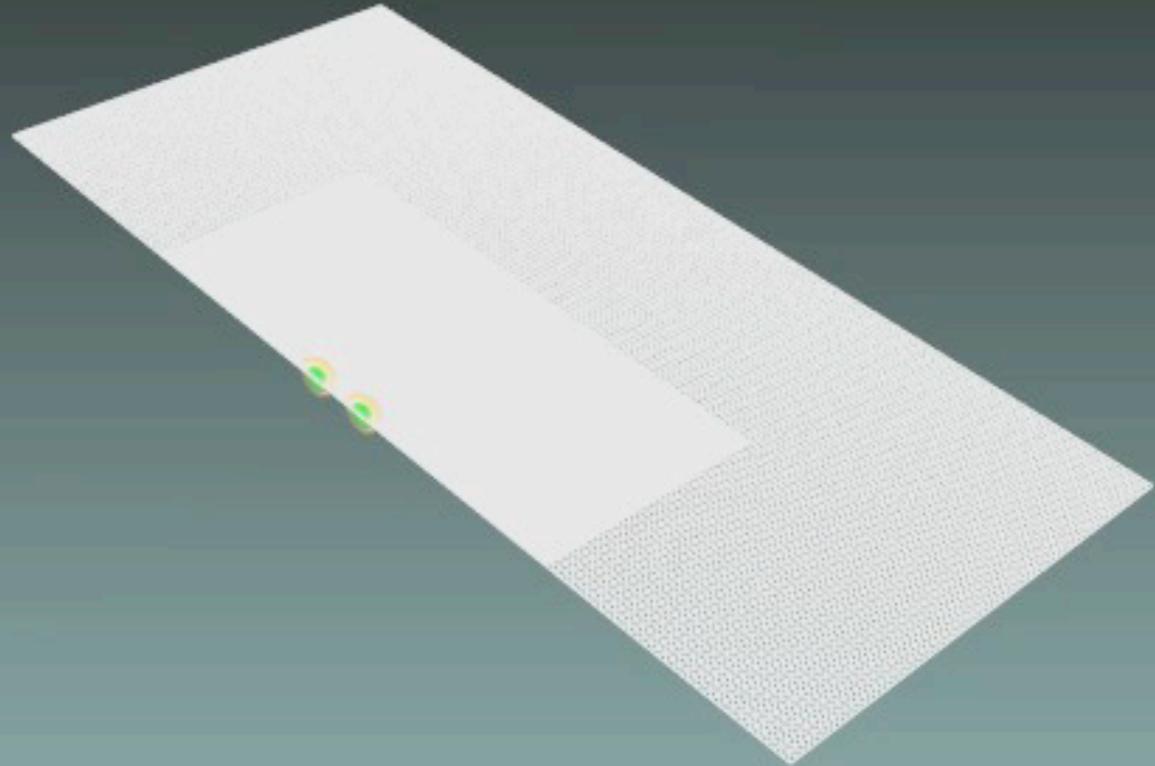
a binary with an EOS allowing for a larger thermal internal energy (ie hotter after merger) will have an increased pressure support and will collapse at a later time

T[ms] = 0.00



T[M] = 0.00

The HMNS is far from the instability threshold and survives for a longer time while losing energy and angular momentum



0.0

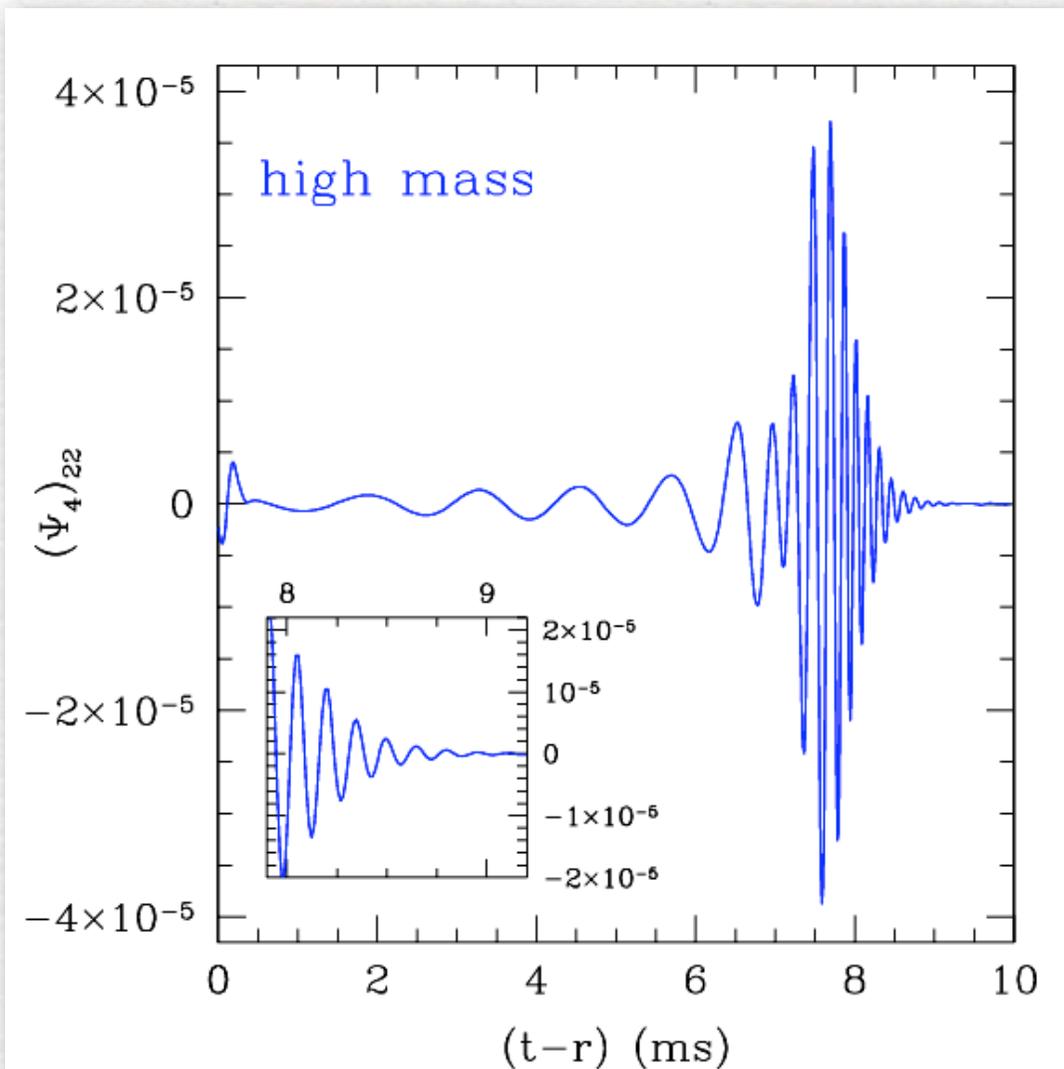
6.1E+14



Density [g/cm³]

Waveforms: polytropic EOS

high-mass binary

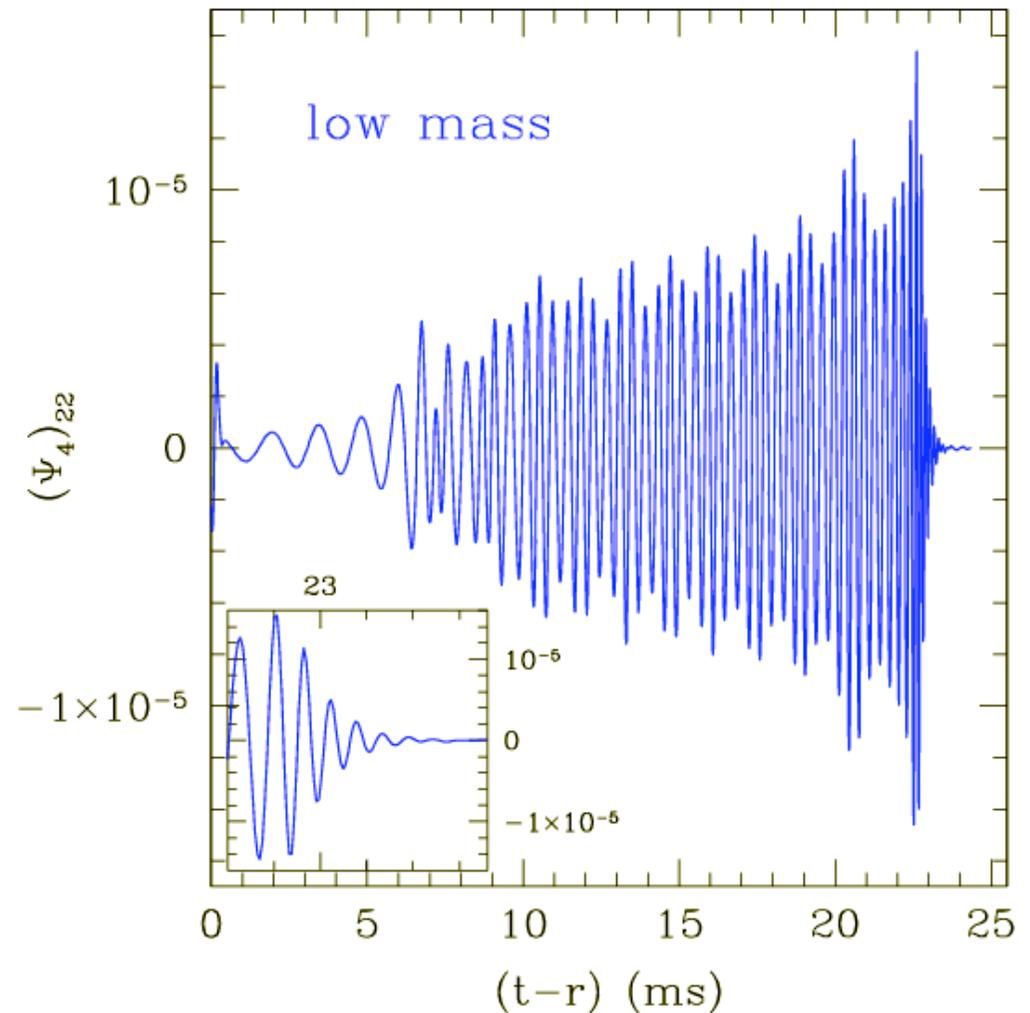
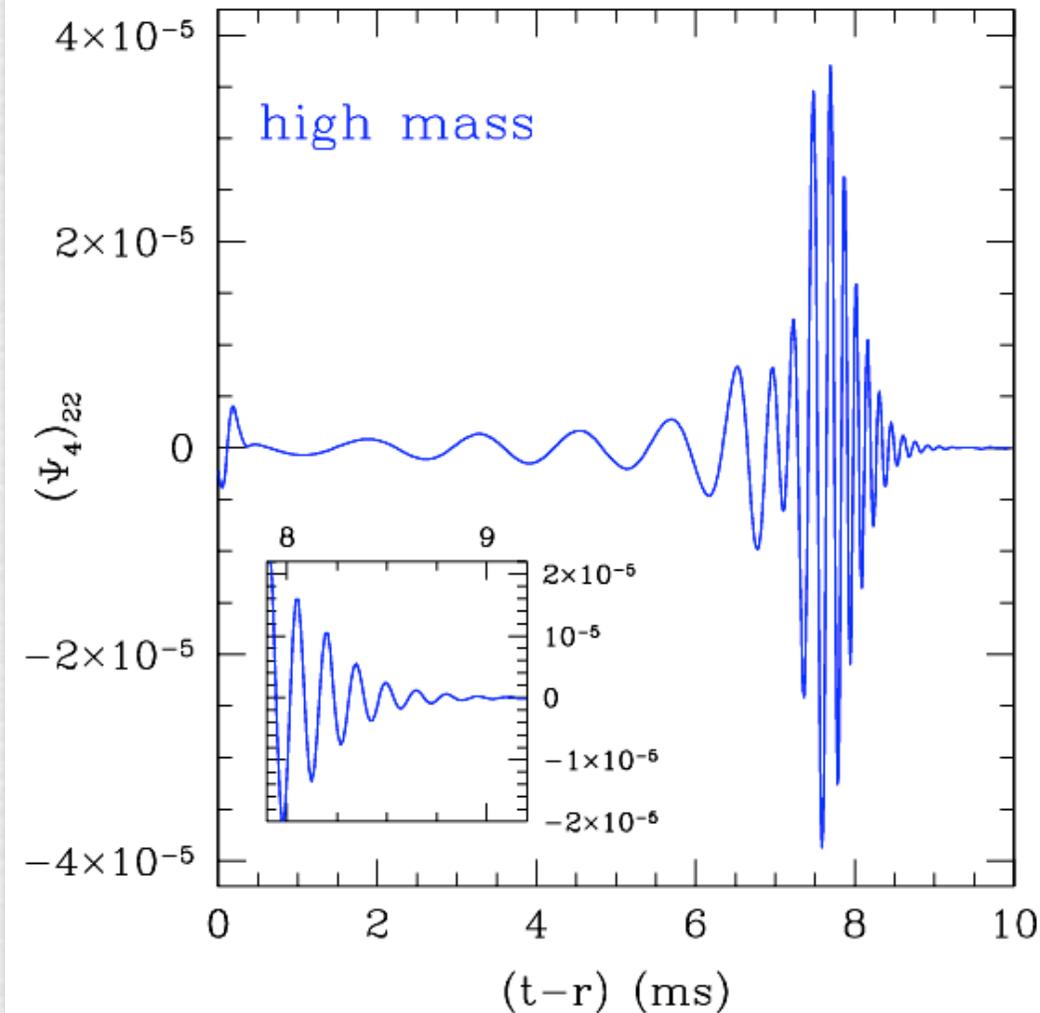


first time the full signal from the formation to a bh has been computed

Waveforms: polytropic EOS

high-mass binary

low-mass binary



first time the full signal from the formation to a bh has been computed

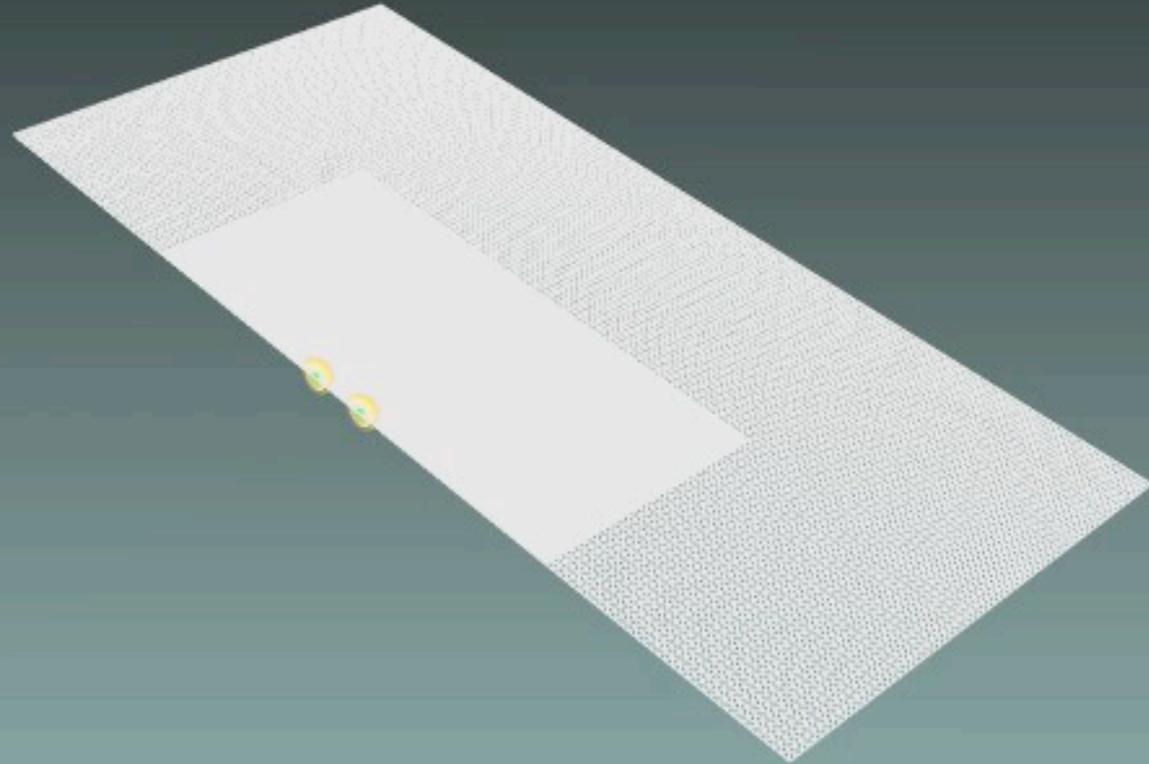
development of a bar-deformed NS leads to a long gw signal

T[ms] = 0.00



T[M] = 0.00

The HMNS is not close to the instability threshold and survives for a much longer time



0.0

6.1E+14

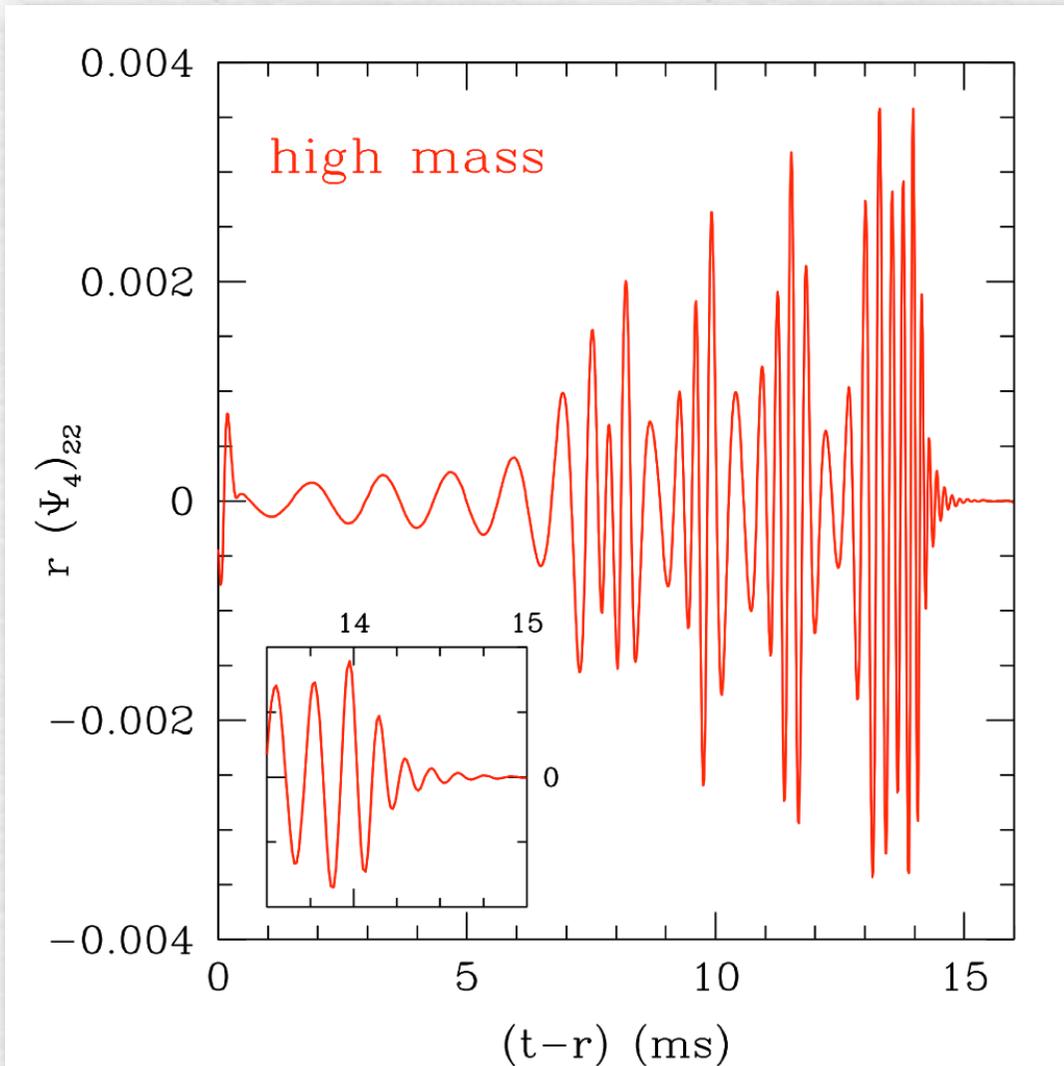


Density [g/cm³]

Waveforms: ideal-fluid EOS

Waveforms: ideal-fluid EOS

high-mass binary

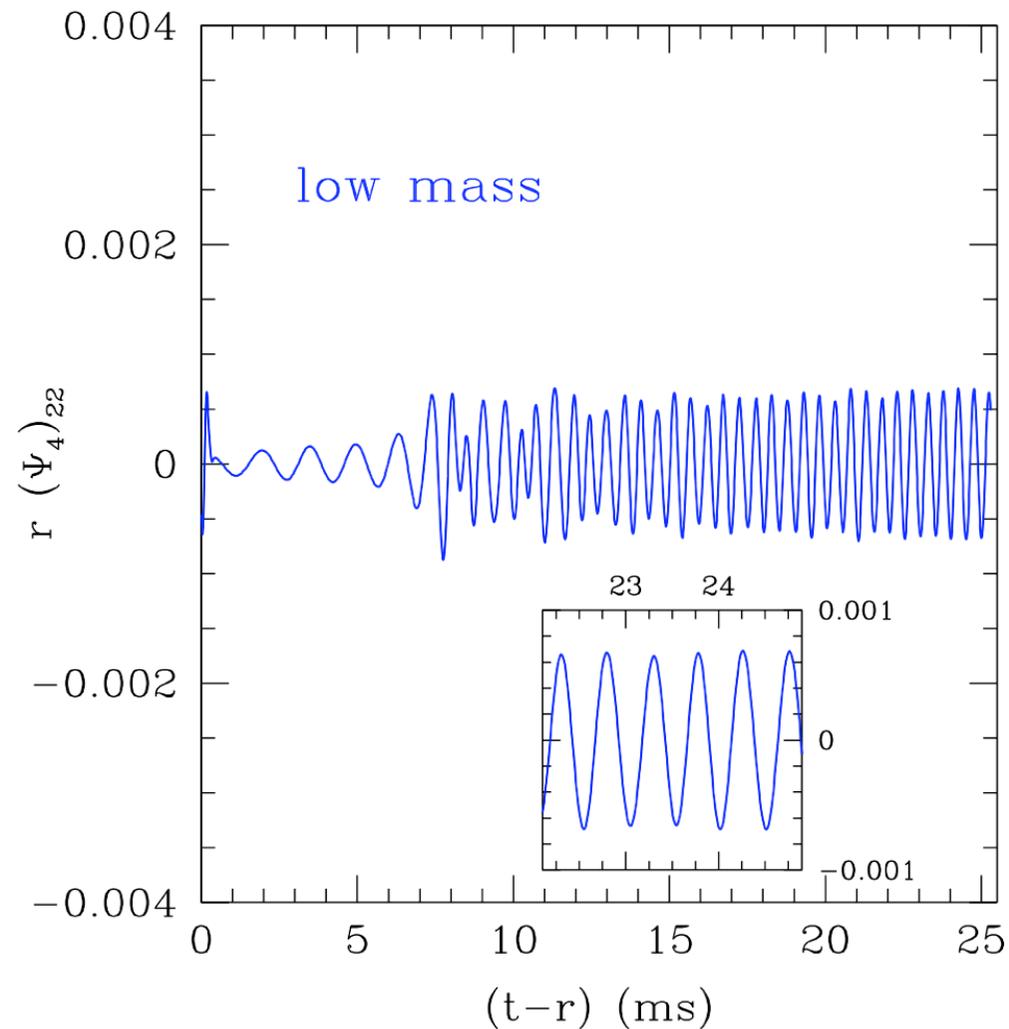
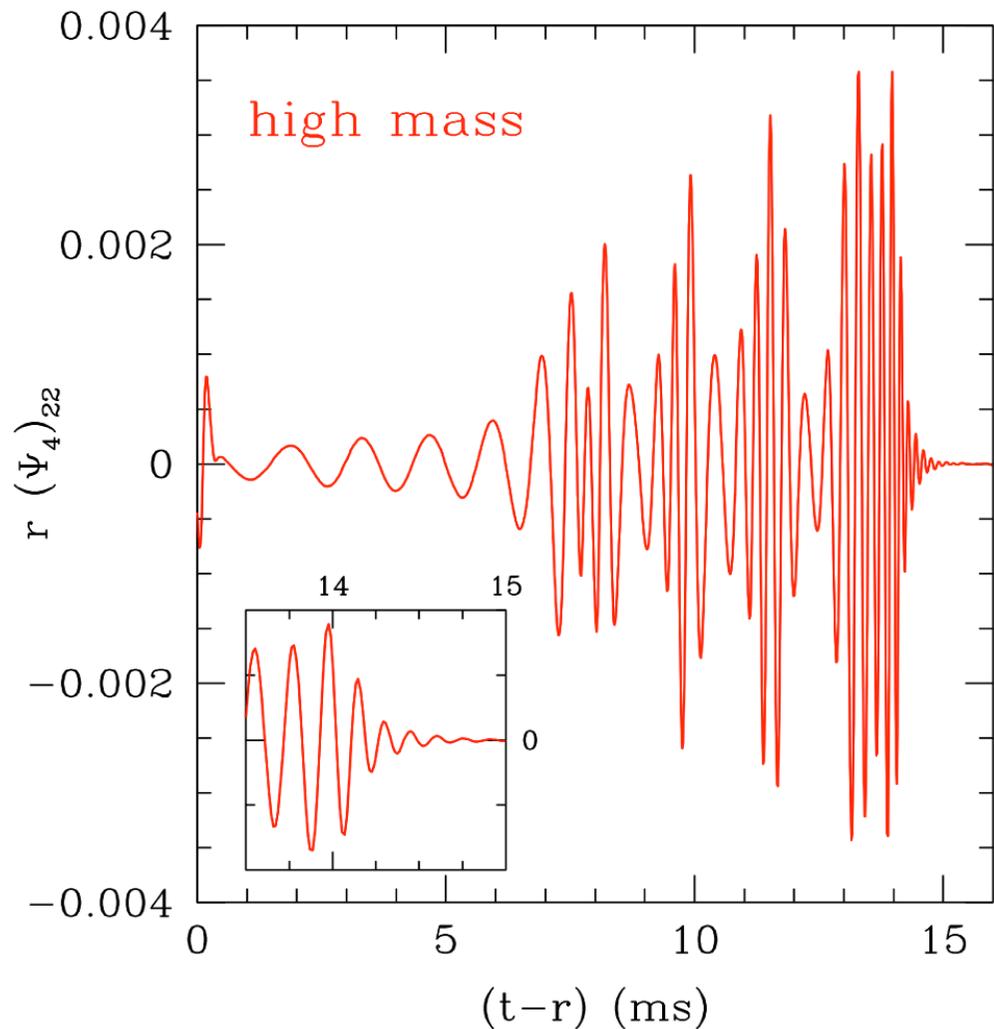


the high internal energy (temperature) of the HMNS prevents a prompt collapse

Waveforms: ideal-fluid EOS

high-mass binary

low-mass binary



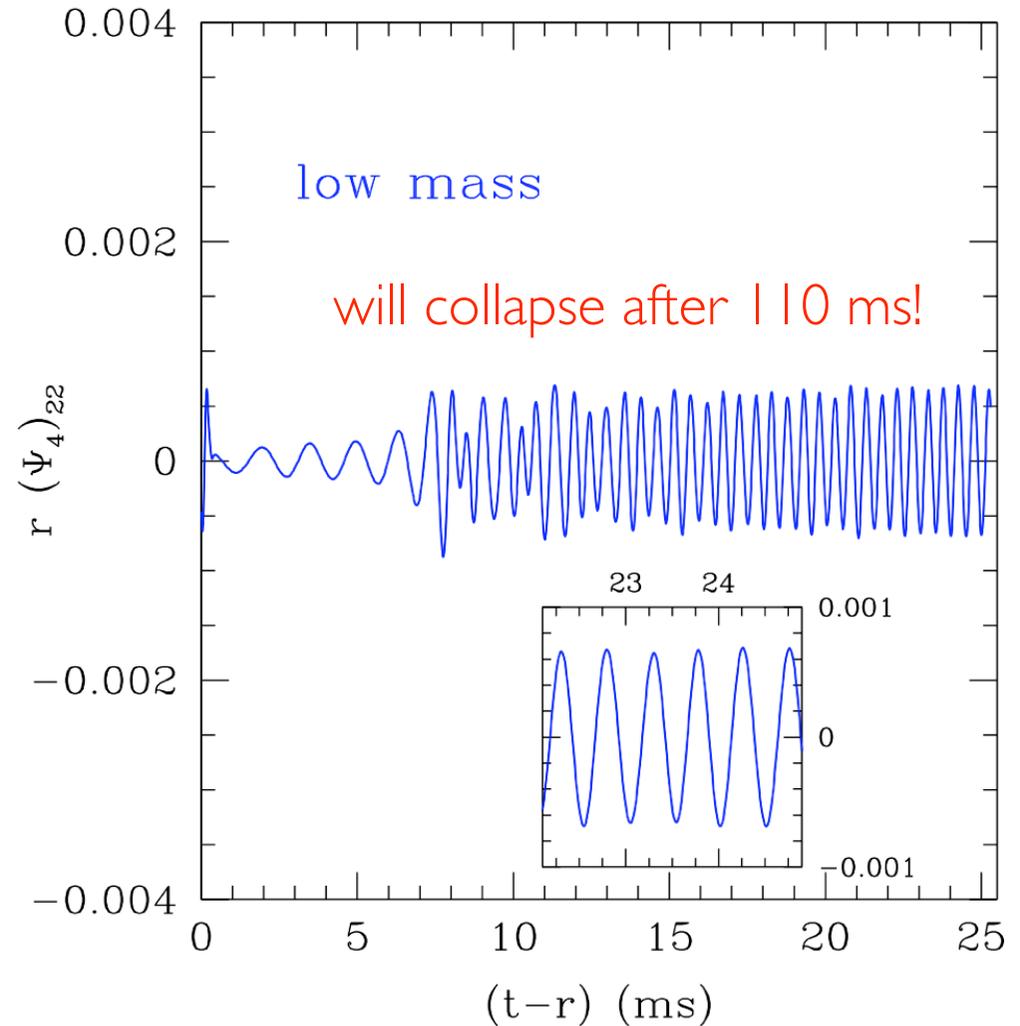
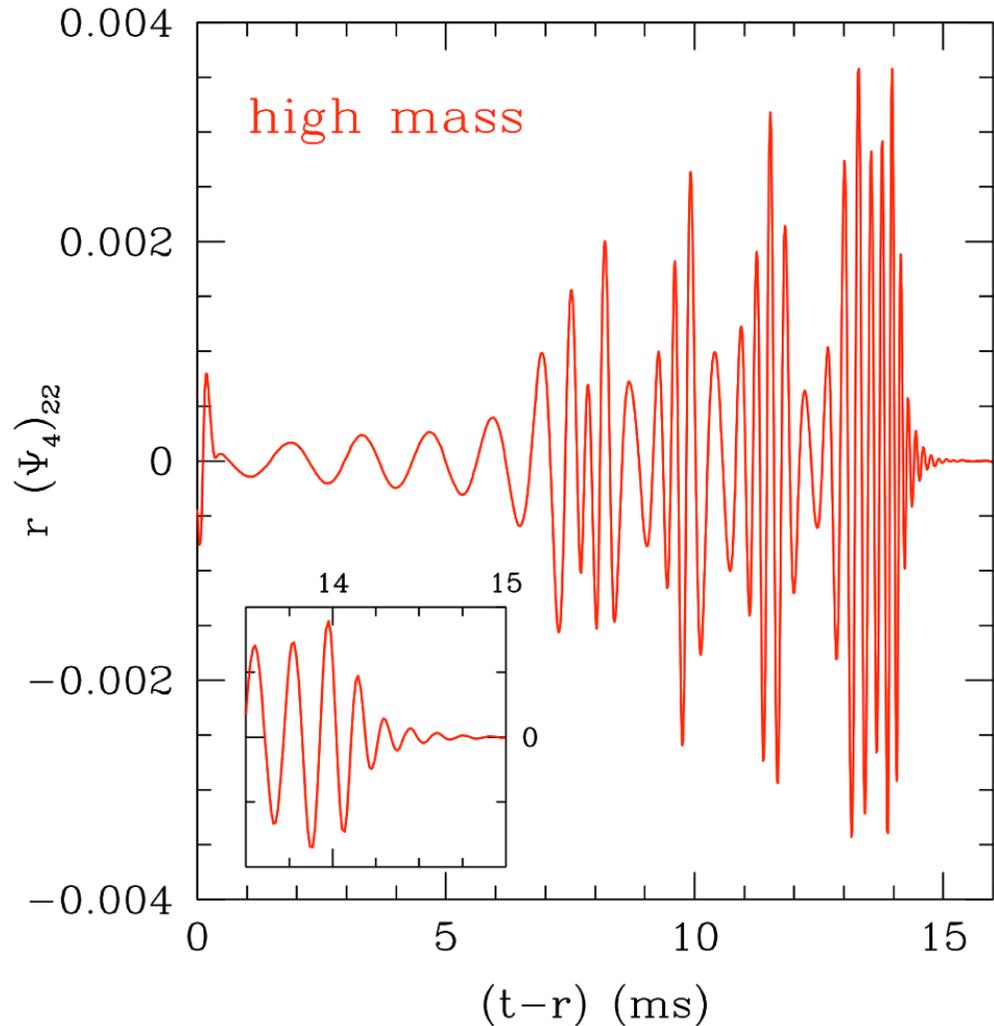
the high internal energy (temperature) of the HMNS prevents a prompt collapse

the HMNS evolves on longer (radiation-reaction) timescale

Waveforms: ideal-fluid EOS

high-mass binary

low-mass binary

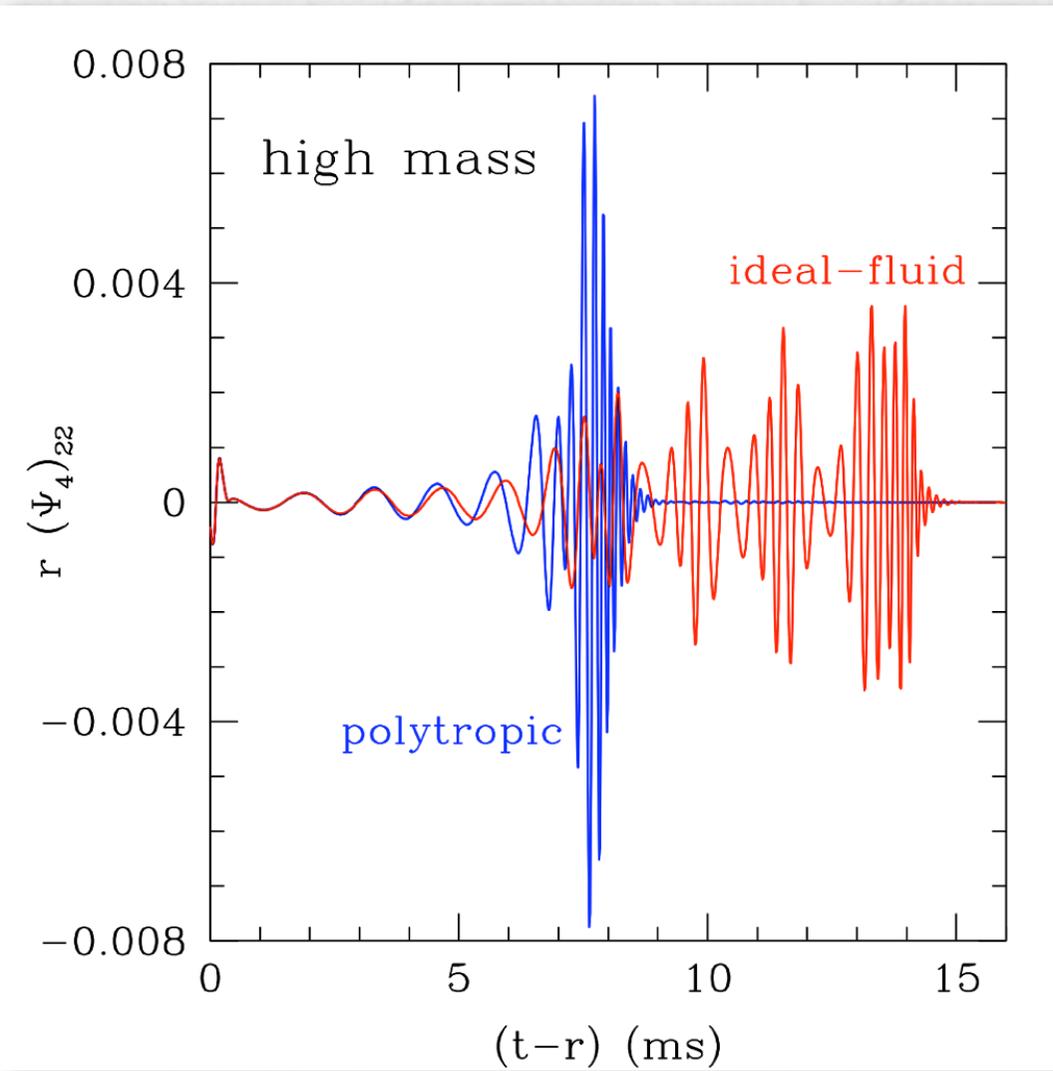


the high internal energy (temperature) of the HMNS prevents a prompt collapse

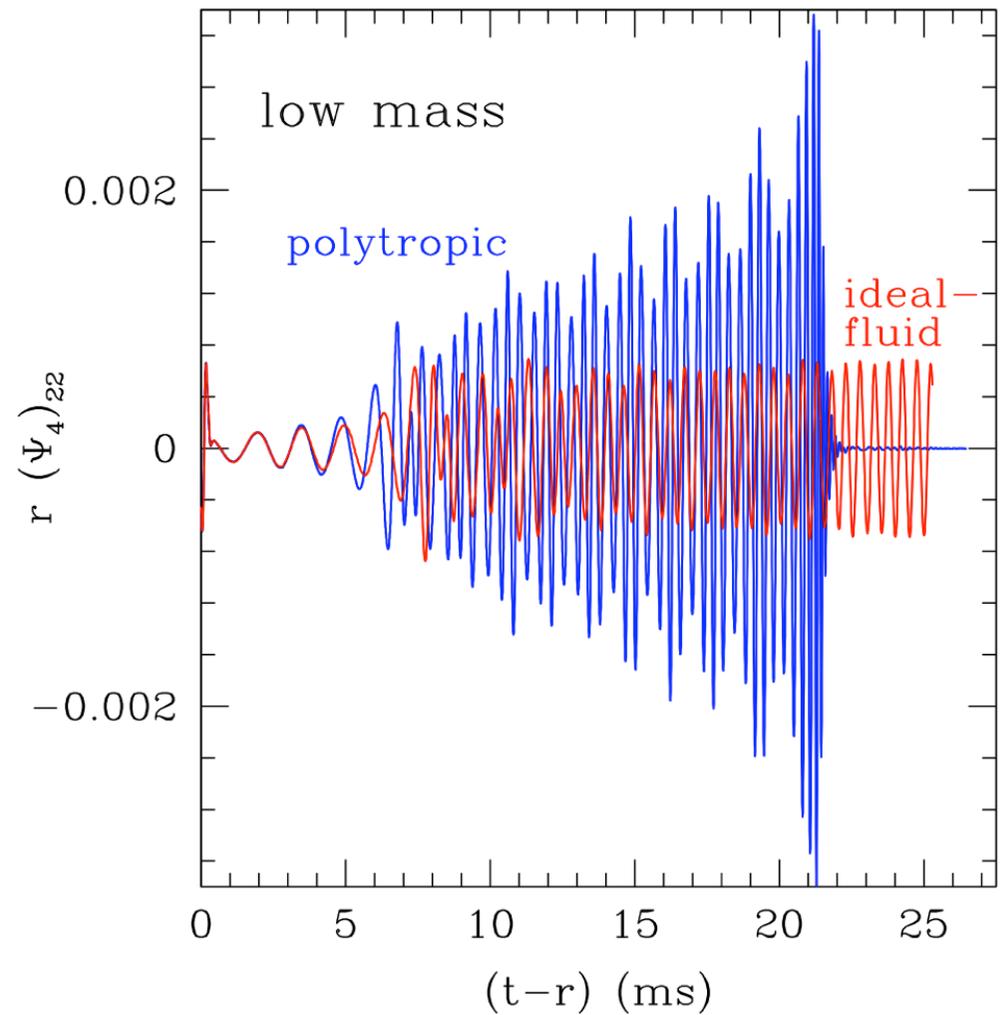
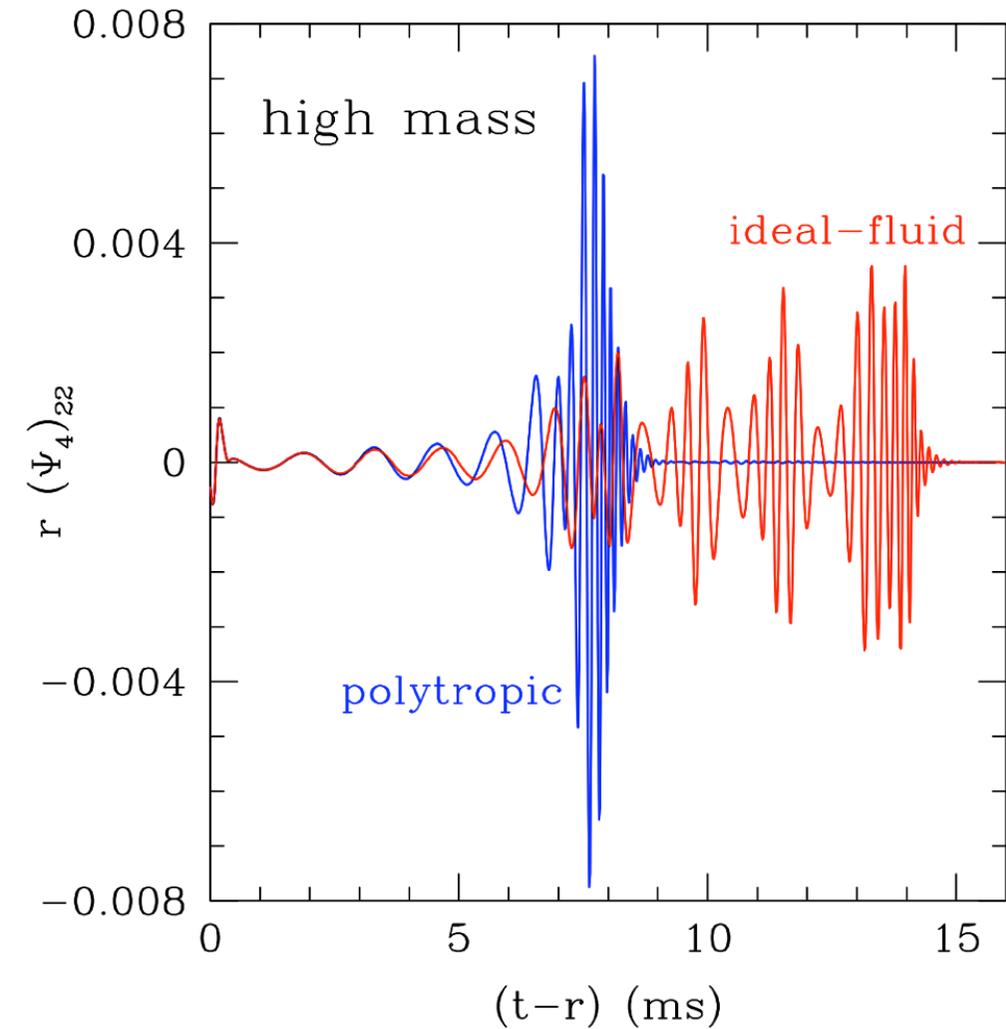
the HMNS evolves on longer (radiation-reaction) timescale

Imprint of the EOS: Ideal-fluid vs polytropic

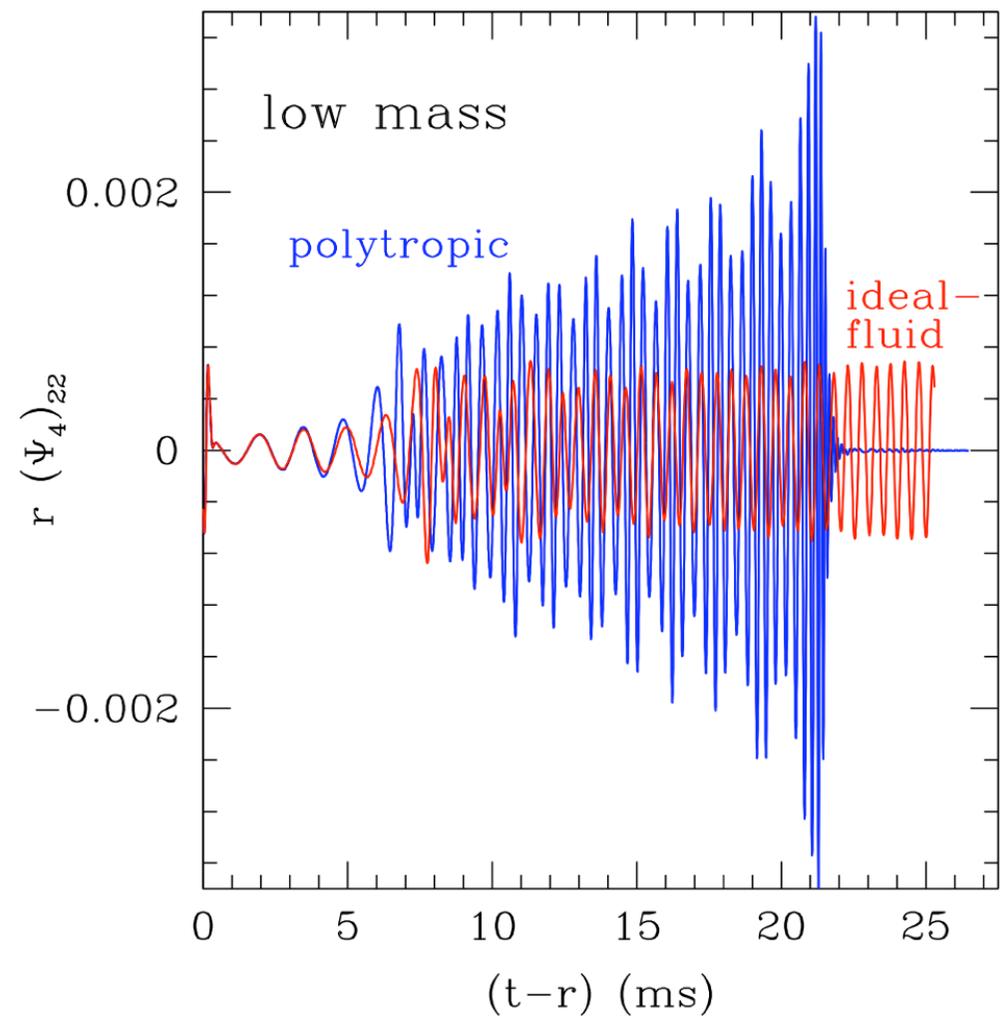
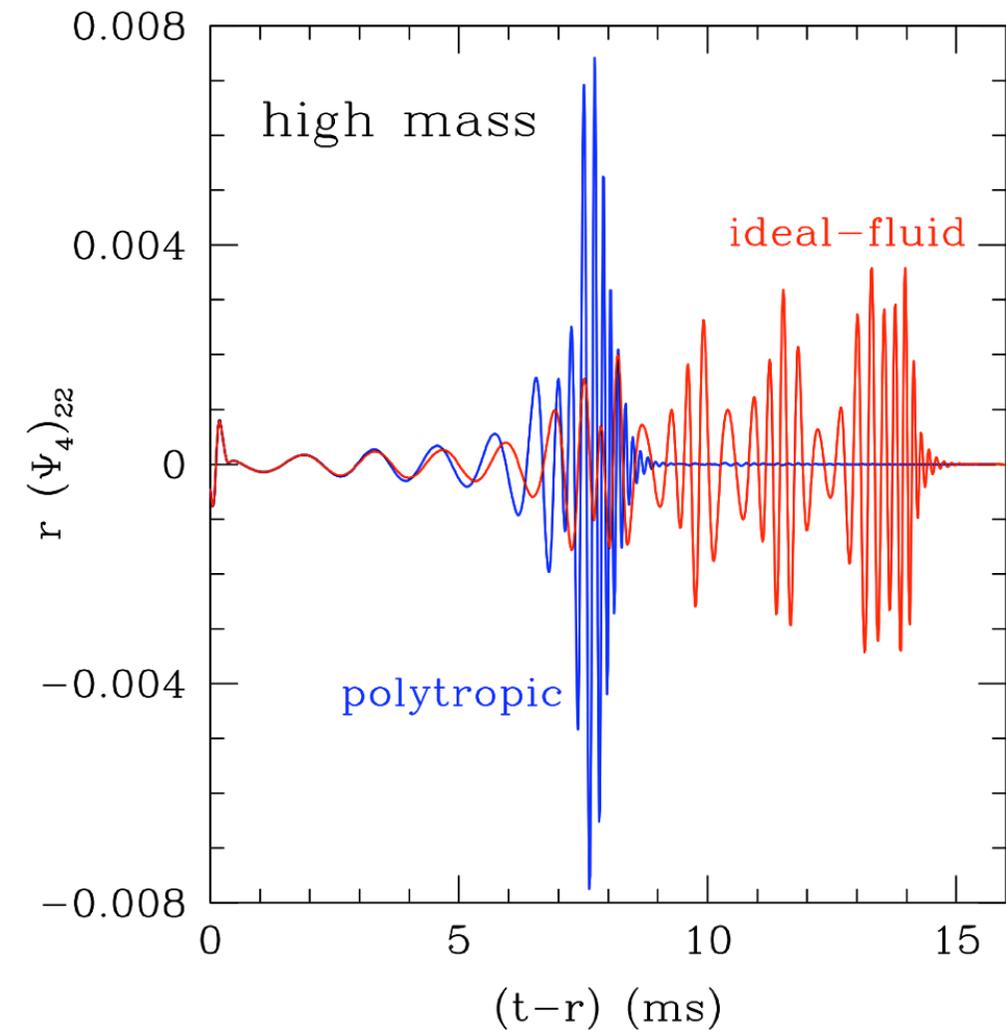
Imprint of the EOS: Ideal-fluid vs polytropic



Imprint of the EOS: Ideal-fluid vs polytropic

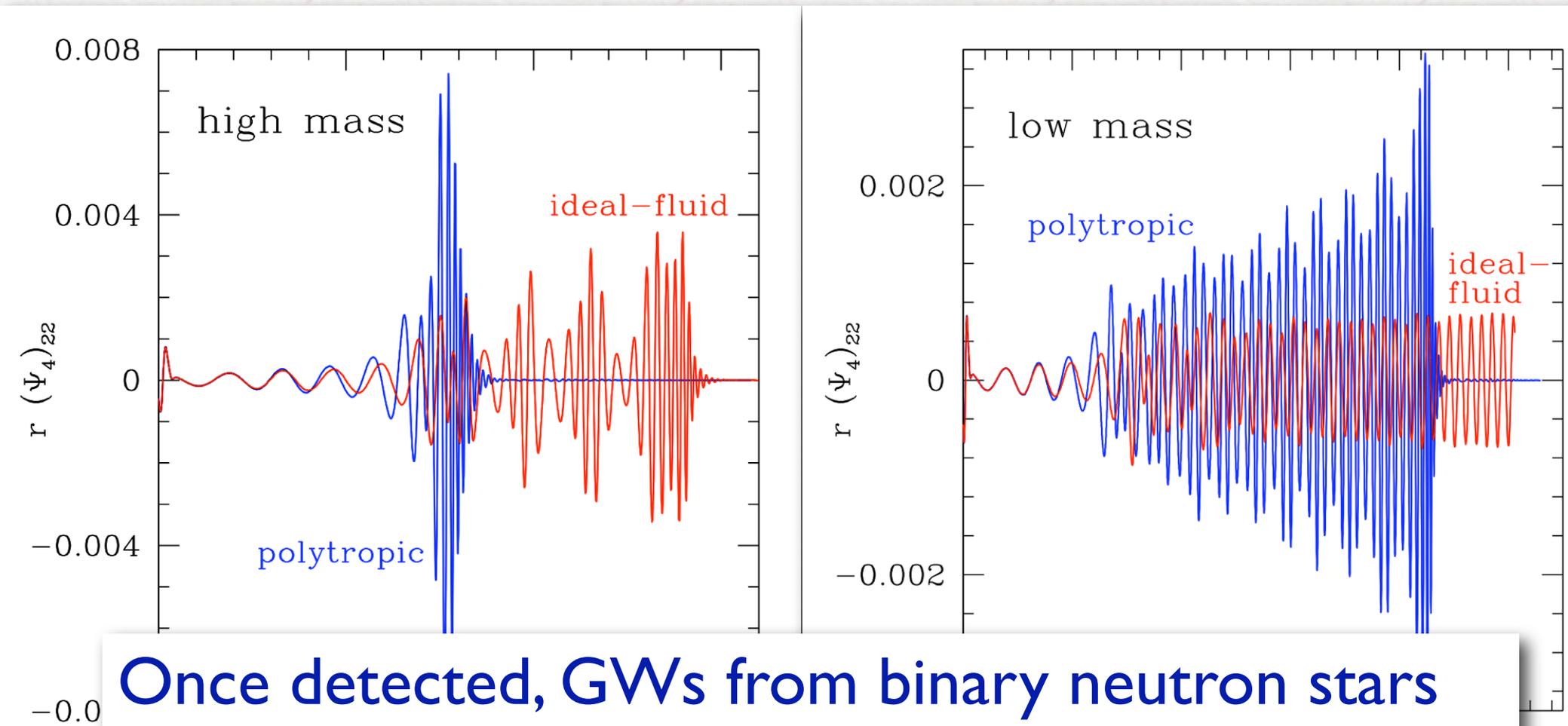


Imprint of the EOS: Ideal-fluid vs polytropic



It is reasonable to expect that for any realistic EOS the waveforms will be in between these two extreme cases

Imprint of the EOS: Ideal-fluid vs polytropic



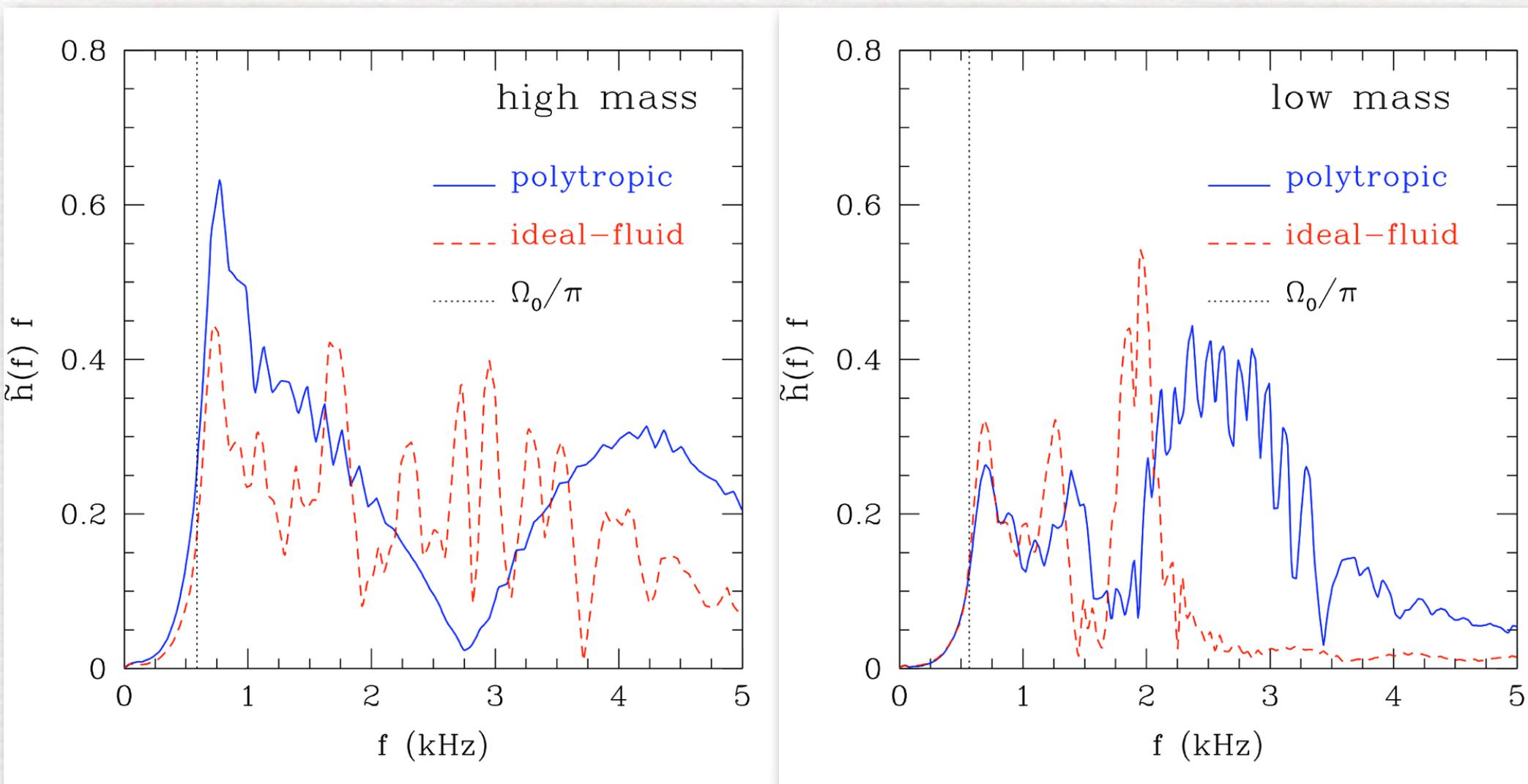
Once detected, GWs from binary neutron stars will be the **Rosetta stone** to decipher the NS EOS

It is reasonable to expect that for any realistic EOS the waveforms will be in between these two extreme cases

Conclusions and prospects

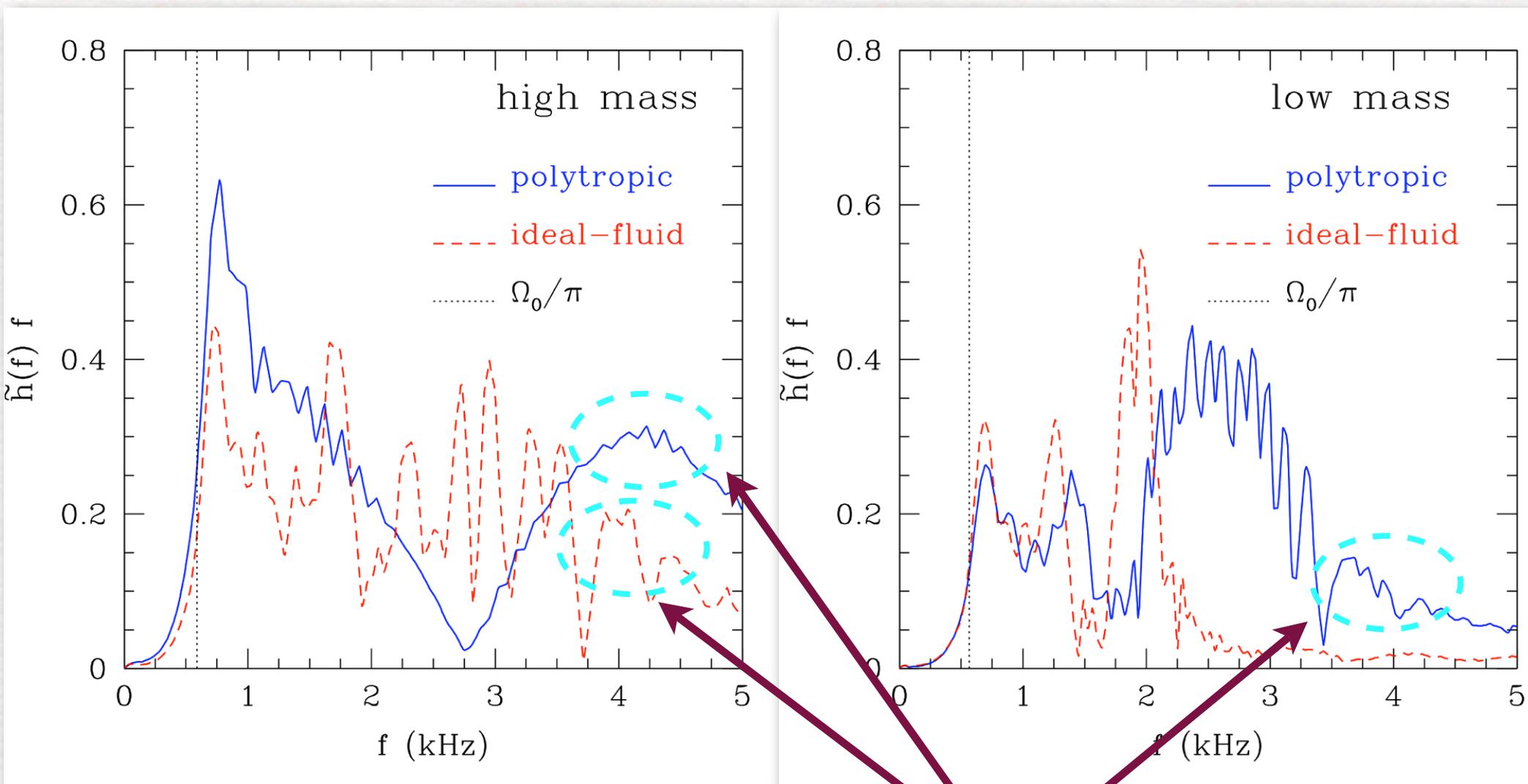
- * With improved numerical techniques and increased computational power, NR is passing from being a “tool-to-be-developed” to a “tool-to-be-used”.
- * NR can act as a bridge between GR and other research areas: astrophysics (GRBs), cosmology (BBHs), GWDA (templates).
- * The work at the AEI is aligned with this philosophy and often leading the way
- * In summary: there is a lot that NR can do for ET and a lot that ET can do for NR/GR!

Imprint of the EOS: frequency domain



The pre-merger dynamics is **very similar**; the post-merger phase is **very different**

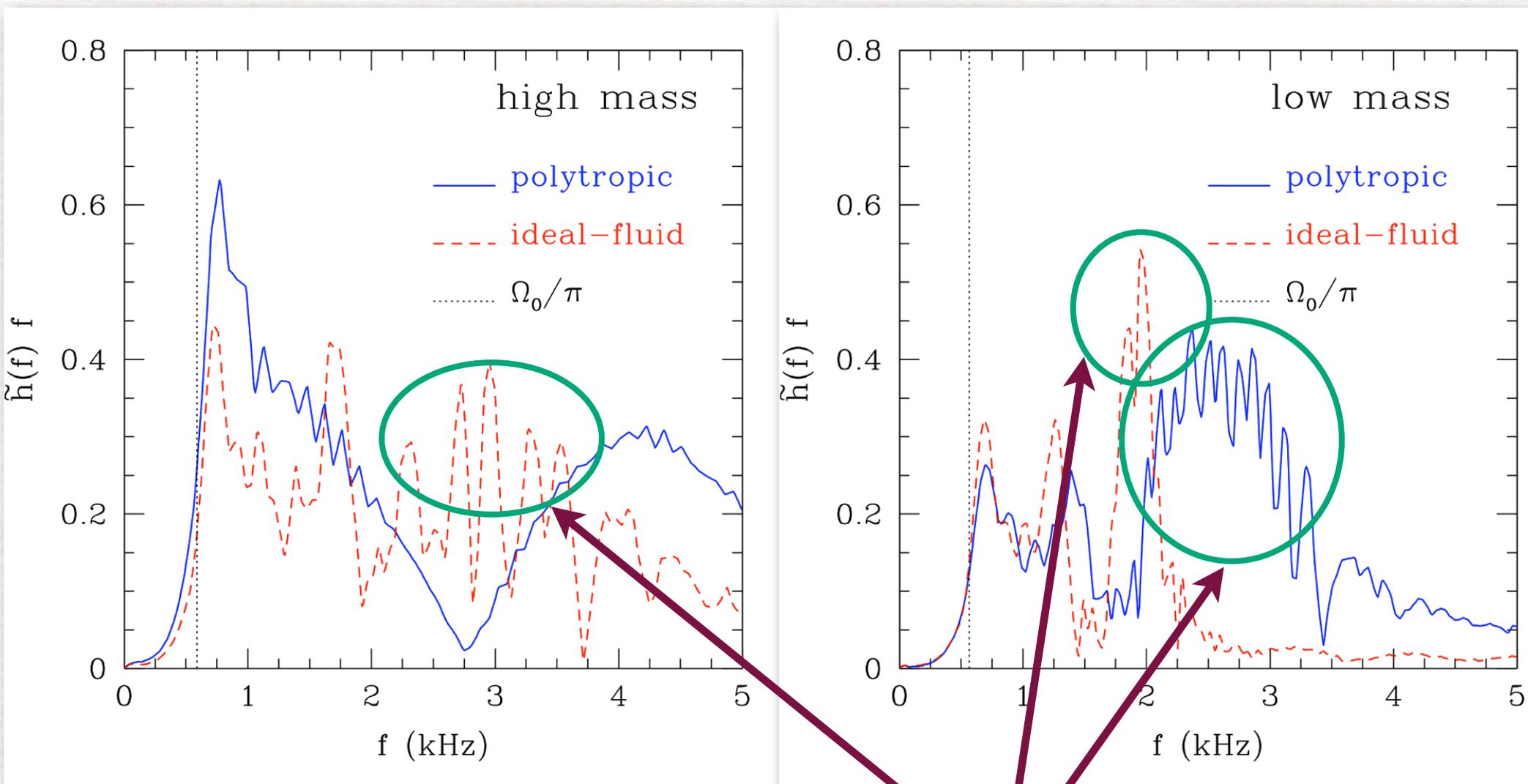
Imprint of the EOS: frequency domain



The pre-merger dynamics is very similar; the post-merger phase is very different

Contributions from the collapse to BH

Imprint of the EOS: frequency domain



The pre-merger dynamics is very similar; the post-merger phase is very different

Contributions from the bar-deformed HMNS