

$K^-$

# $Y^*$ resonances investigation originated in $K$ -light nuclei absorption by AMADEUS

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on behalf of AMADEUS and in collaboration with  
Prof S. Wycech

Investigating strangeness: from accelerators to compact stellar  
14 May 2014, LNF INFN, Frascati, Italy

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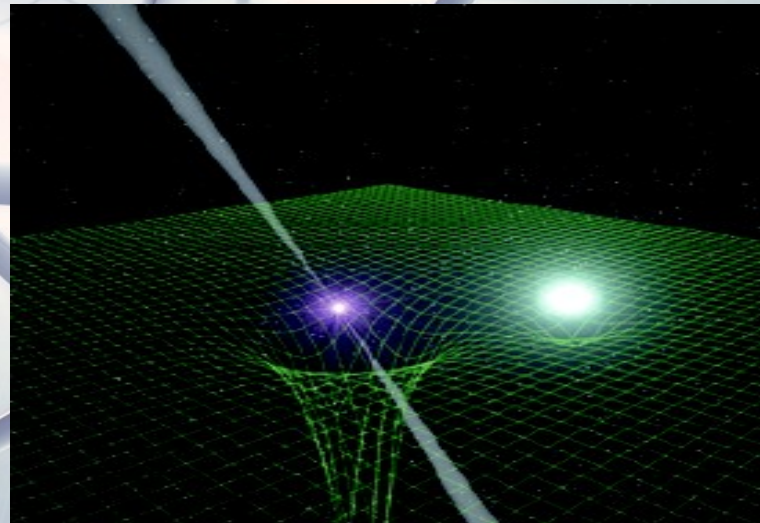
# Framework: Low-Energy QCD with Strange Quarks

$K^-$

Strangeness in baryonic matter:

- role of strangeness in **EoS of neutron stars**
- hyperon-nucleon and hyperon-hyperon interactions role in the investigation of dense baryonic matter
- new constraints from **2 solar masses neutron stars**, very stiff Equation of State required!

But



- the basic ingredient .. namely  **$\bar{K}N$  interaction still unclear** and mysterious from the experimental point of view.



$K^-$

## Framework: Low-Energy QCD with Strange Quarks

Approached by the investigation of the antikaon-nucleon interaction

Important constraints:

- $K^-N$  threshold physics (shift and width of kaonic atoms levels measured by SIDDHARTA)
  - $\Sigma\pi - \Lambda\pi$  mass spectra
- Nature and properties of the  $\Lambda(1405)$  considered as  $K^-N$  quasibound state embedded in the  $\Sigma\pi$  continuum

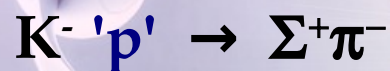
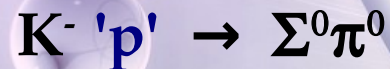
# Investigation of $K^-$ absorption on light nuclei

$K^-$

(H,  $^4\text{He}$ ,  $^9\text{Be}$ ,  $^{12}\text{C}$ )

AT-REST ( $K^-$  absorbed from atomic orbit) or IN-FLIGHT

Reactions:



$\Lambda p$  from 1NA or 2NA (single or multi-nucleon absorption)

$\Lambda d$  and  $\Lambda t$  channels (I. Tucakovic Talk)

**R&D for more refined setup: ScFi + SiPM (trigger system) TPC – GEM (inner tracker)**

Experimental tests of the trigger prototype for the AMADEUS experiment based on Sci-Fi read by MPPC,

Nucl.Instrum.Meth. A671 (2012) 125-128

Performances of a GEM-based TPC prototype for new high-rate particle experiments,

Nucl.Instrum.Meth. A617 (2010) 183-185

$K^-$

# Framework: Low-Energy QCD with Strange Quarks

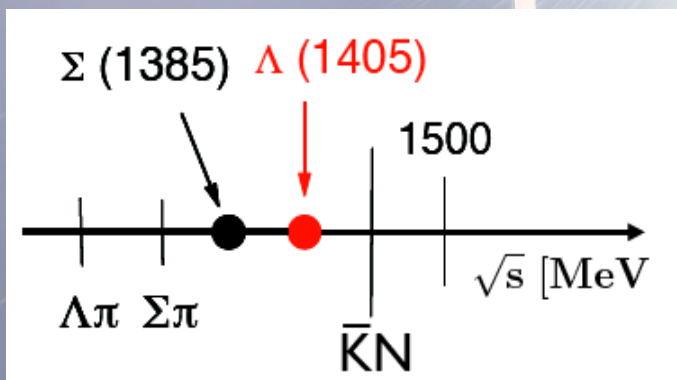
**CHIRAL PERTURBATION THEORY**  
Interacting systems of **NAMBU-GOLDSTONE BOSONS**  
(pions, kaons) coupled to **BARYONS**

$$\mathcal{L}_{eff} = \mathcal{L}_{mesons}(\Phi) + \mathcal{L}_B(\Phi, \Psi_B)$$

works well for low-energy pion-pion and pion-nucleon interactions

... but **NOT** for systems with strangeness  $S = -1$

**BECAUSE  $\Lambda(1405)$  just below  $K^-N$  threshold (1432 MeV)**

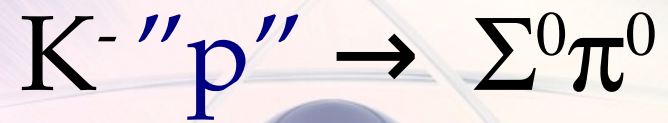


**Solutions:**

- Non-perturbative Coupled Channels approach based on Chiral SU(3) Dynamics
- phenomenological  $\bar{K}N$  and NN potentials



$K^-$



bound proton in  ${}^4\text{He}$  /  ${}^{12}\text{C}$

# Scientific case $\Lambda(1405)$

$K^-$   $\Lambda(1405)$  : mass =  $1405.1^{+1.3}_{-1.0}$  MeV, width =  $50 \pm 2$  MeV

$I = 0, S = -1, J^p = 1/2^-$ , Status: \*\*\*\*, strong decay into  $\Sigma\pi$

Its nature has been a puzzle for decades: three quark state, unstable  $\bar{K}N$  bound state, penta-quark, two poles??

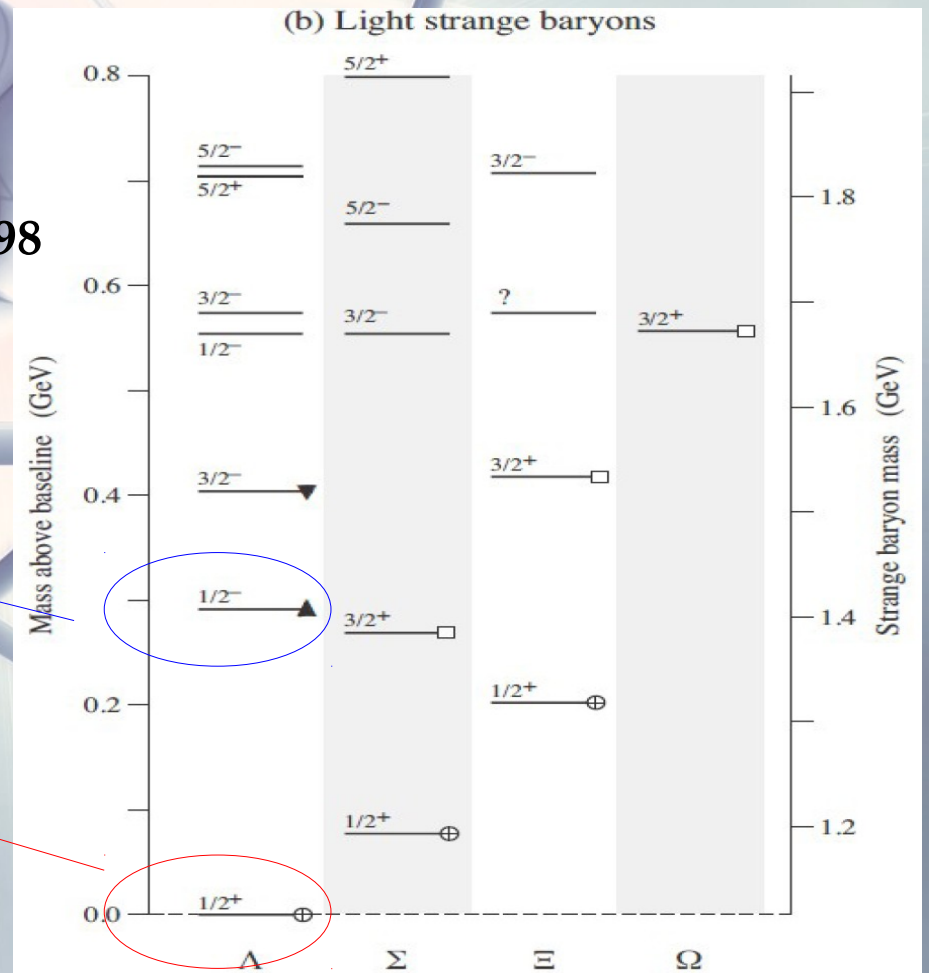
First experimental evidence:

M. H. Alston, et al., Phys. Rev. Lett. 6 (1961) 698

$K^- p \rightarrow \pi\pi\pi\Sigma$

$\Lambda(1405)$

$\Lambda(1116)$



## Scientific case $\Lambda(1405)$

$\Lambda(1405)$  is a negative parity baryon resonance (spin = 1/2, isospin = 0, strangeness = -1) located slightly below the  $\bar{K}N$  threshold, decaying into the  $\Sigma\pi$  channel through the strong interaction.

The **three quark model picture** has some difficulties to reproduce the  $\Lambda(1405)$ . According to its negative parity, one of the quarks has to be excited to the  $l = 1$  orbit. Similar to the nucleon sector, where one of the lowest negative parity baryon is the  $N(1535)$ , **the expected mass of the  $\Lambda^*$  is around 1700 MeV** (since it contains one strange quark). Another difficulty is the energy splitting observed between the  $\Lambda(1405)$  and the  $\Lambda(1520)$ , if is interpreted as the spin-orbit partner ( $J^p = 3/2^-$ ).

R. Dalitz and collaborators first suggested to interpret  $\Lambda(1405)$  as an  $\bar{K}N$  quasibound state.



# Scientific case $\Lambda(1405)$

- Chiral unitary models:  $\Lambda(1405)$  is an  $I = 0$  quasibound state emerging from the coupling between the  $\bar{K}N$  and the  $\Sigma\pi$  channels. Two poles in the neighborhood of the  $\Lambda(1405)$ :

4) *two poles*:  $(z_1 = 1424^{+7}_{-23} - i 26^{+3}_{-14} ; z_2 = 1381^{+18}_{-6} - i 81^{+19}_{-8})$  MeV (Nucl. Phys. A881, 98 (2012))

mainly coupled to  $\bar{K}N$

mainly coupled to  $\Sigma\pi$

→ line-shape depends on production mechanism

- Akaishi-Esmaili-Yamazaki phenomenological potential

Phys. Lett. B 686 (2010) 23-28 Confirmation of single pole ansatz?

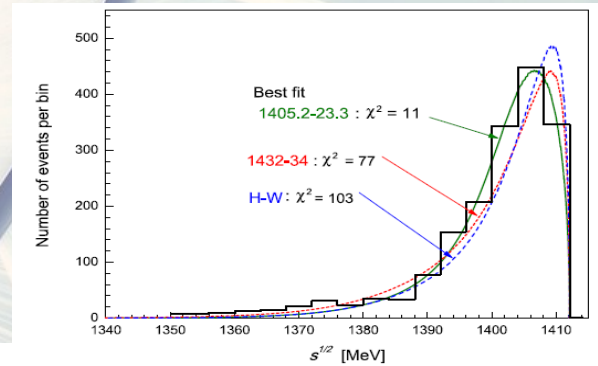
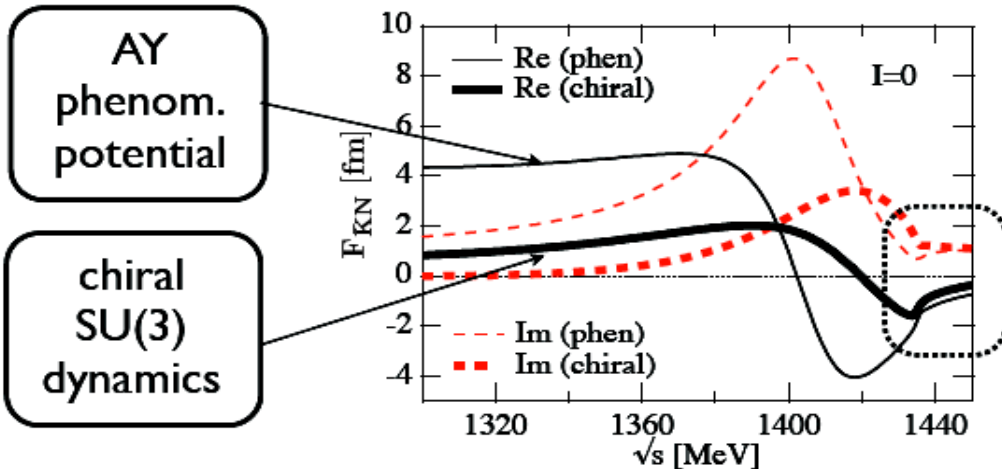
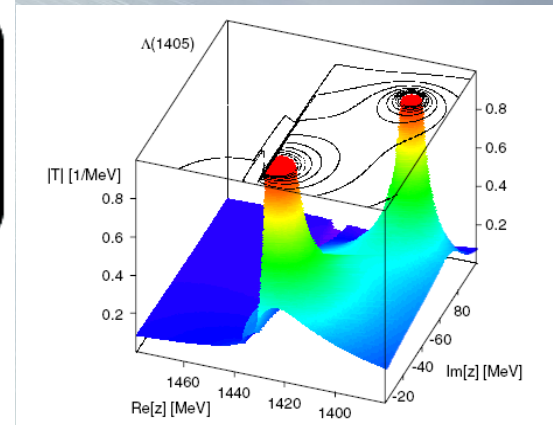


Fig. 6. Detailed differences in  $M_{\Sigma\pi}$  spectra among the Hyodo-Weise prediction and the present model predictions.



large differences in subthreshold extrapolations

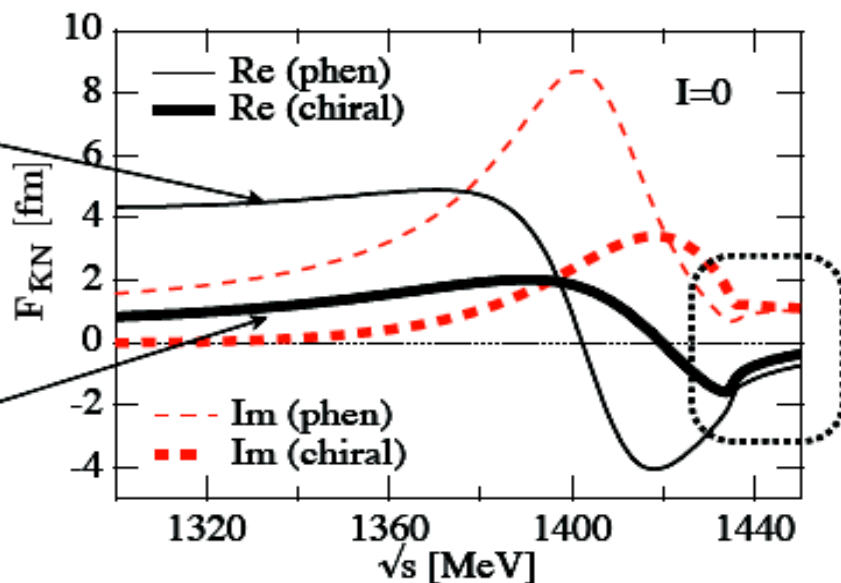


- Chiral dynamics predicts significantly weaker attraction than AY (local, energy independent) potential in far-subthreshold region

# Scientific case $\Lambda(1405)$

AY  
phenom.  
potential

chiral  
SU(3)  
dynamics



large differences  
in  
**subthreshold**  
extrapolations

- **Chiral dynamics** predicts significantly **weaker attraction** than AY (local, energy independent) potential in **far-subthreshold** region

Distribution shape depends

TO TEST THE HIGHER POLE:

- **production in  $\bar{K}N$  reactions** (only chance to observe the high mass pole)
- **decaying in  $\Sigma^0\pi^0$**  (free from  $\Sigma(1385)$  background I=1)

on the decay channel:

$$\frac{d\sigma(\Sigma^-\pi^+)}{dM} \propto \frac{1}{3} |T^0|^2 + \frac{1}{2} |T^1|^2 + \frac{2}{\sqrt{6}} \text{Re}(T^0 T^{1*})$$

$$\frac{d\sigma(\Sigma^+\pi^-)}{dM} \propto \frac{1}{3} |T^0|^2 + \frac{1}{2} |T^1|^2 - \frac{2}{\sqrt{6}} \text{Re}(T^0 T^{1*})$$

$$\frac{d\sigma(\Sigma^0\pi^0)}{dM} \propto \frac{1}{3} |T^0|^2$$

# Scientific case $\Lambda(1405)$

$K^-$  nuclear absorption experiments .. long history .. BUT

$K^-$

- 1)  $m_{\pi\Sigma}$  spectra **CUT AT THE ENERGY LIMIT AT-REST**    2)  $(\Sigma\pm\pi^\mp)$   **$\Sigma(1385)$  CONTAMINATION**

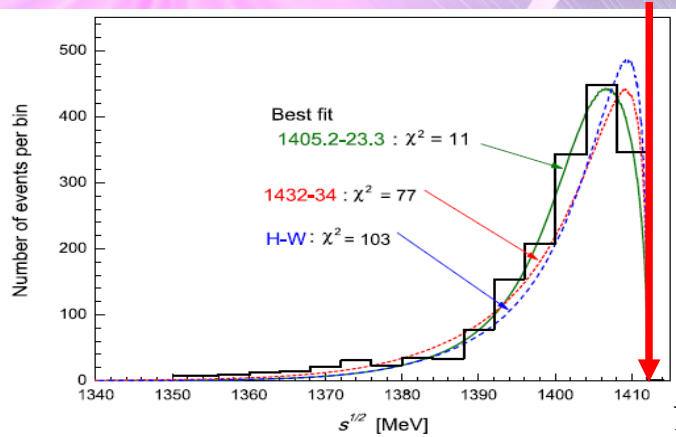
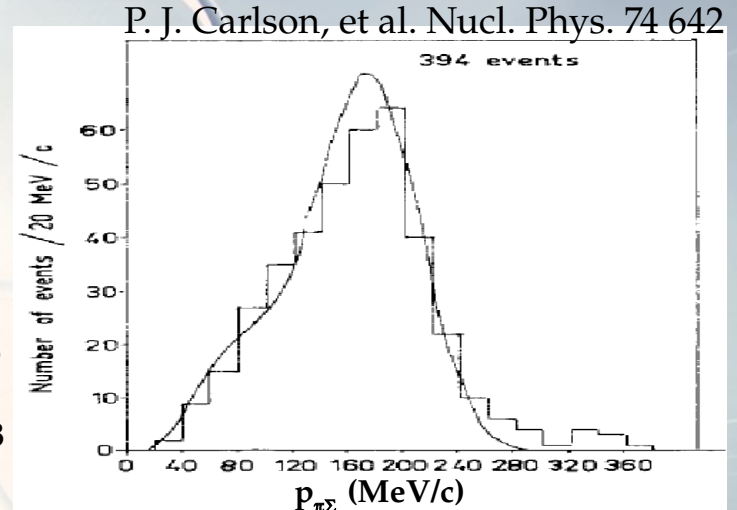


Fig. 6. Detailed differences in  $M_{\Sigma\pi}$  spectra among the Hyodo-Weise prediction and the present model predictions.

“A study of  $K^- \text{ } ^4\text{He} \rightarrow (\Sigma\pm\pi^\mp) + \text{}^3\text{H}$  using slow instead of stopping  $K^-$  would be very useful in eliminating some of the uncertainties in interpretation”

D. Riley, et al. Phys. Rev. D11 (1975) 3065

Esmaili et al., Phys.Lett. B686 (2010) 23-28



P. J. Carlson, et al. Nucl. Phys. 74 642

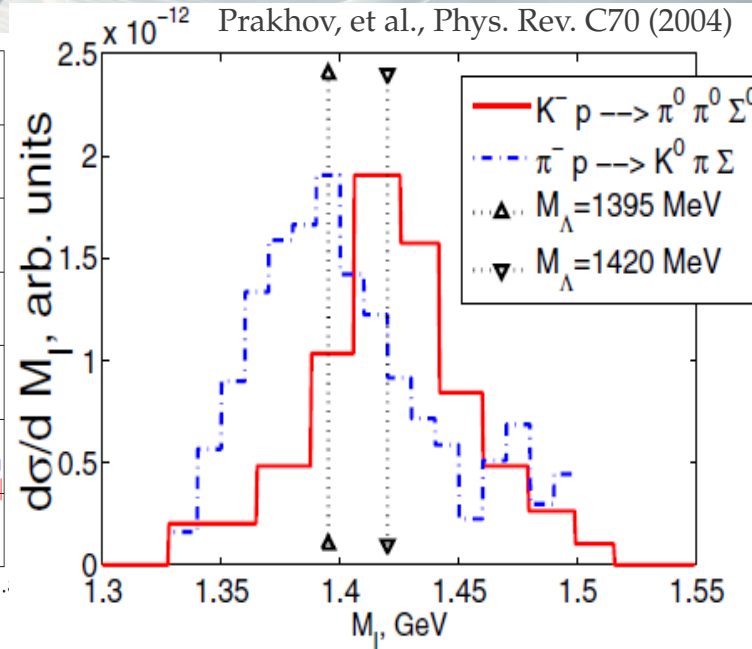
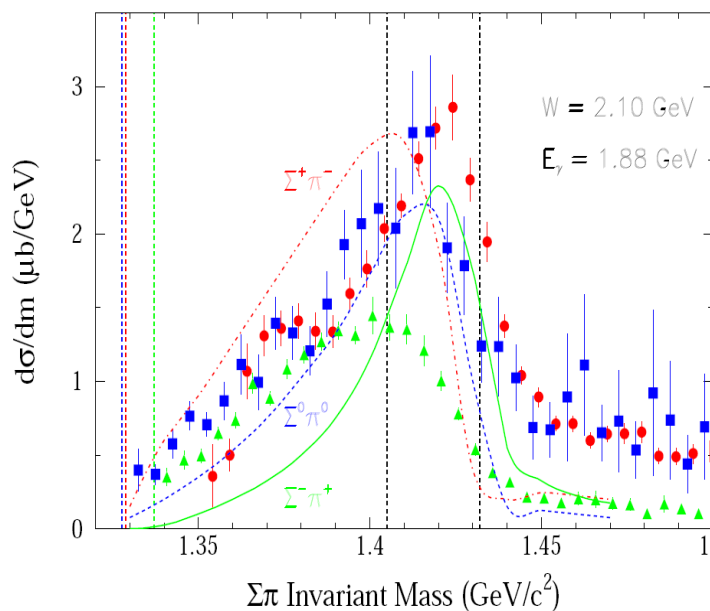
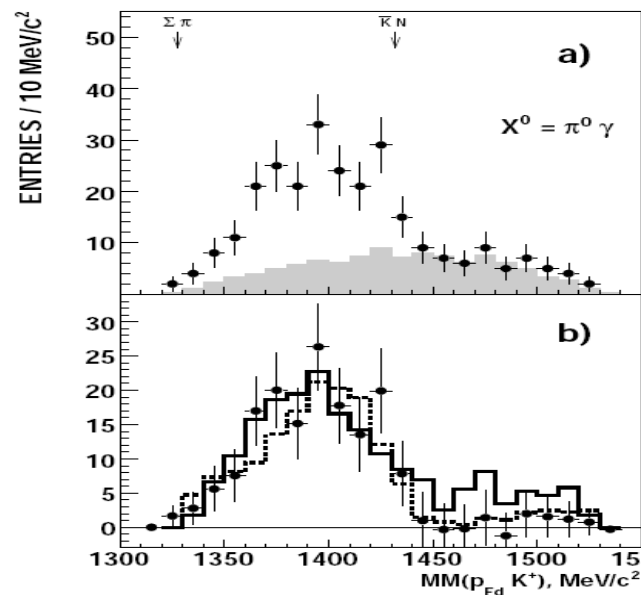
The  $\Sigma^0\pi^0$  spectrum was only observed in 3 experiments ... with different line-shapes !

I. Zychor et al., Phys. Lett. B 660 (2008) 167

K. Moriya, et al., (Clas Collaboration) Phys. Rev. C 87, 035206 (2013)

Magas et al. PRL 95, 052301 (2005) 034605 S.

Prakhov, et al., Phys. Rev. C70 (2004)



# TWO SAMPLES OF DATA:

$K^-$

- **2004-2005** KLOE data (Analyzed luminosity of  $\sim 1.5 \text{ fb}^{-1}$ )

$K^-$  absorbed in KLOE materials (H,  $^4\text{He}$ ,  $^9\text{Be}$ ,  $^{12}\text{C}$ )

**At-rest + In-flight**

- Dedicated **2012** run with pure graphite Carbon target inside KLOE  
( $\sim 90 \text{ pb}^{-1}$ ; analyzed  $37 \text{ pb}^{-1}$ , x1.5 statistics)

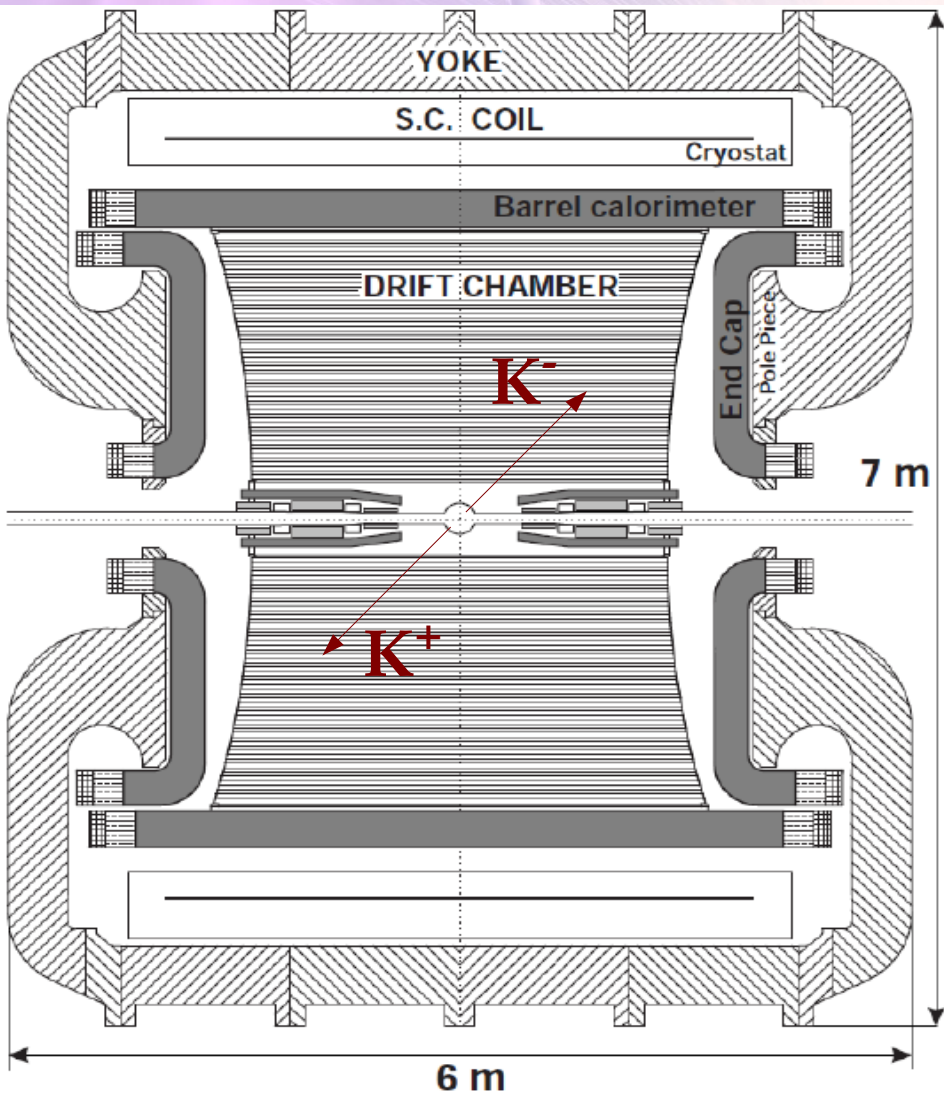
**$K^-$   $^{12}\text{C}$  absorptions At-rest**

# Low-energy $K^-$ hadronic interactions studies with KLOE, why?

$K^-$

MC simulations show that :

- $\sim 0.1$  of  $K^-$  stopped in the DC gas (90% He, 10%  $C_4H_{10}$ )
- $\sim 2\%$  of  $K^-$  stopped in the DC wall (750  $\mu\text{m}$  c. f. , 150  $\mu\text{m}$  Al foil).



Possibility to use KLOE materials as an active target

Advantage:  
excellent resolution ..

$$\sigma_{p\Lambda} = 0.49 \pm 0.01 \text{ MeV}/c \text{ in DC gas}$$

$$\sigma_{m\gamma\gamma} = 18.3 \pm 0.6 \text{ MeV}/c^2$$

Disadvantage:

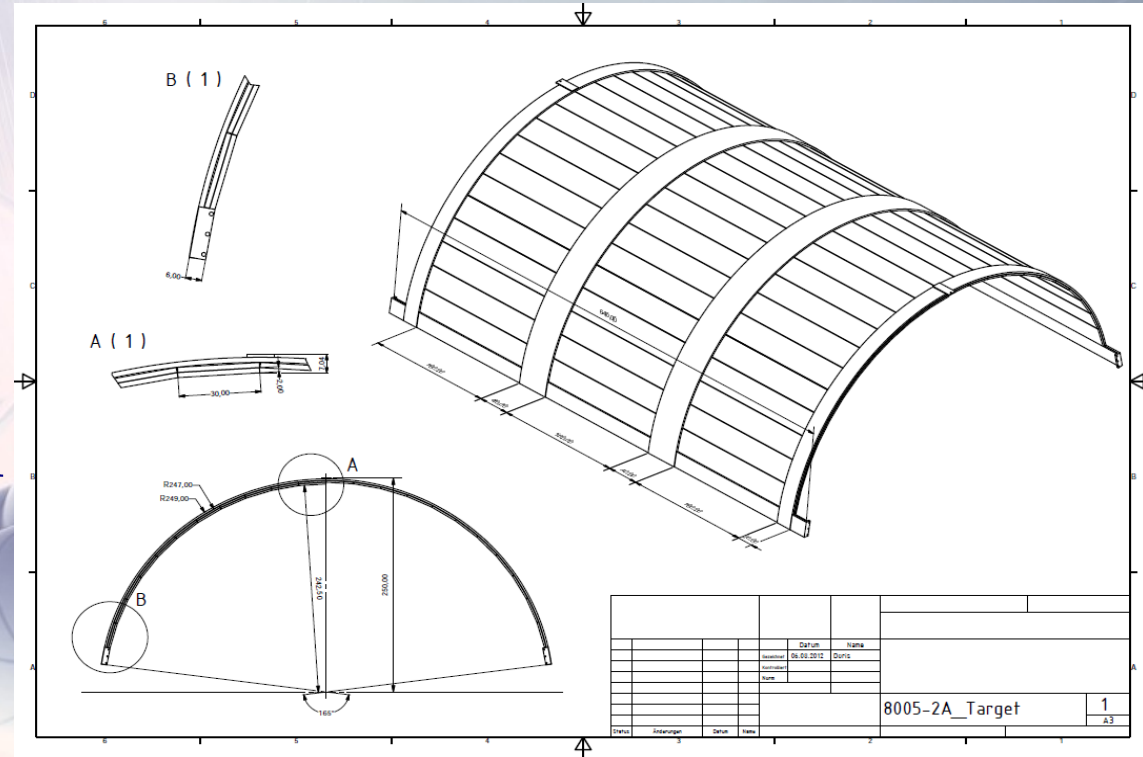
Non dedicated target  $\rightarrow$  different nuclei contamination  $\rightarrow$  complex interpretation ..  
but  $\rightarrow$  new features ..  $K^-$  in flight absorption.

# Carbon target inside KLOE

$K^-$

## Advantages:

- gain in statistics
- $K^-$  absorptions occur in Carbon
- absorptions at-rest.



- MC simulation: 26% of  $K^-$  stopped in C, 2% of  $K^-$  stopped in Al hence aluminium contamination from 19%  $\rightarrow$  7% !
- Thickness optimized (based on MC simulations) to maximize the number of stopping  $K^-$  in the target, minimizing the charged particles energy loss.

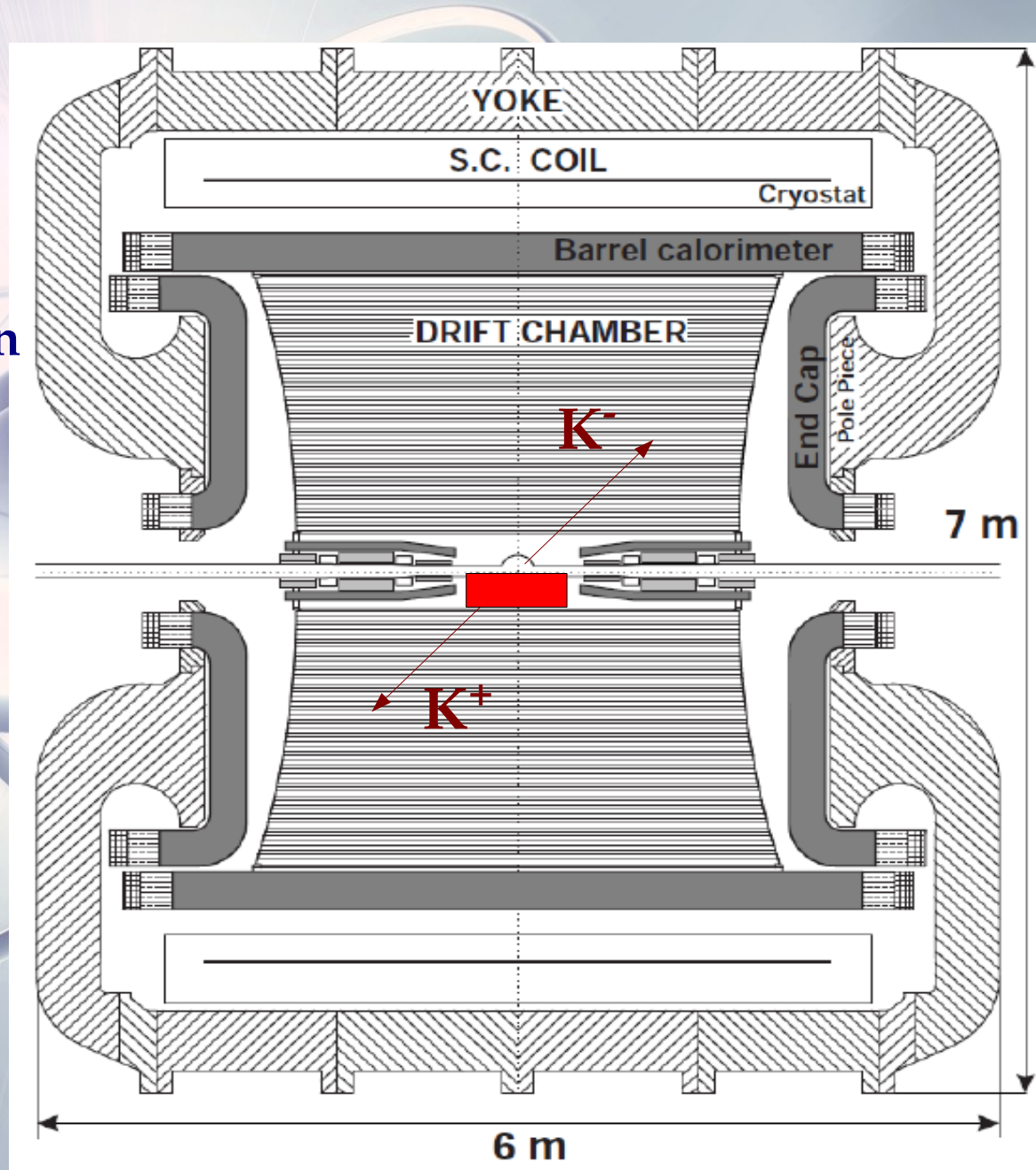
( $\sim 90 \text{ pb}^{-1}$ ; analyzed  $37 \text{ pb}^{-1}$ , x1.5 statistics)

# Carbon target inside KLOE

$K^-$

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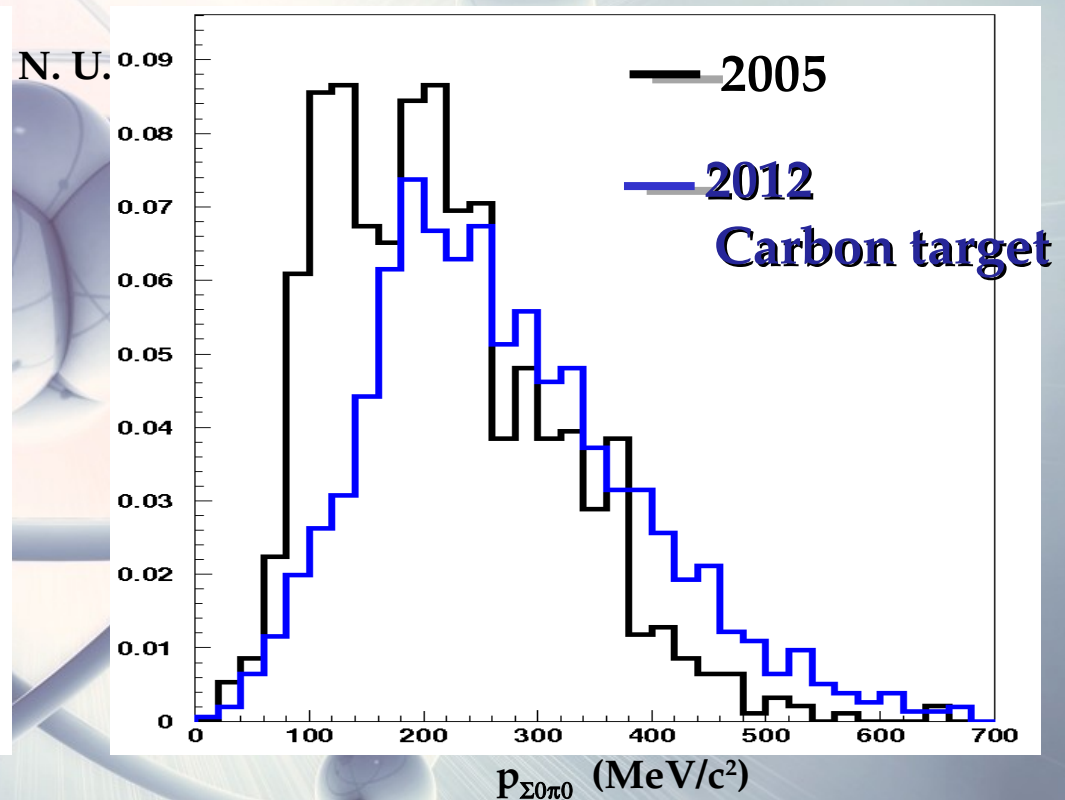
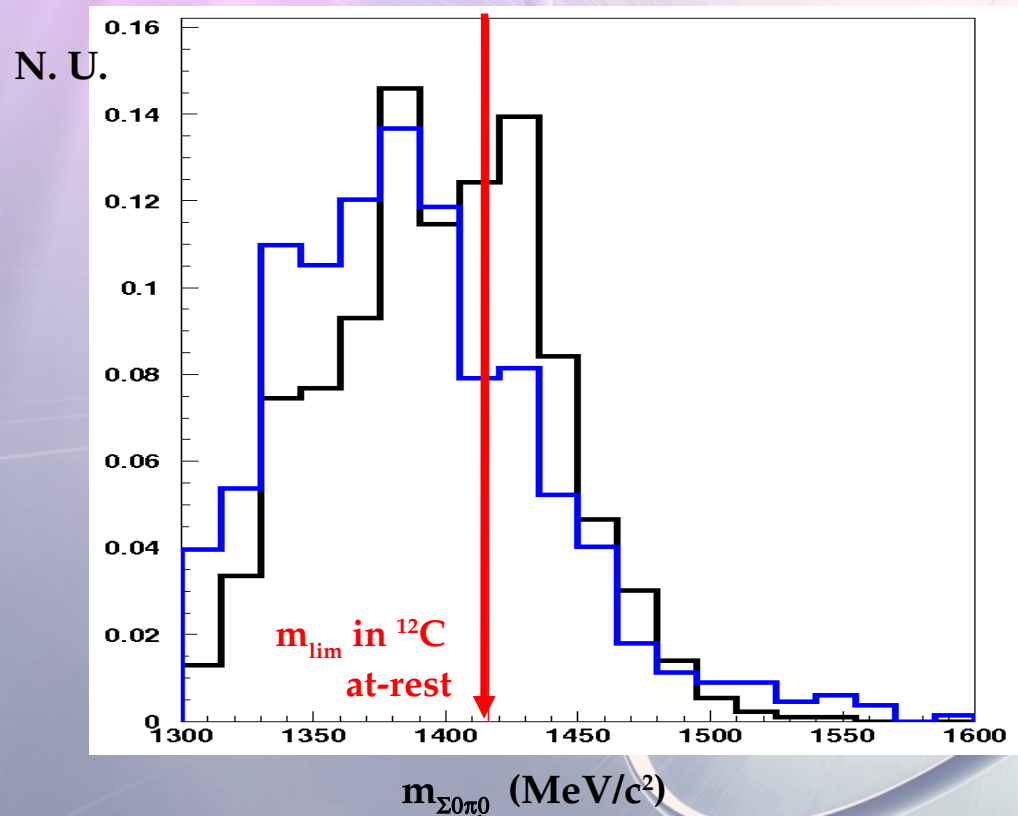


# $\Sigma^0 \pi^0$ channel

$K^- \Lambda(1405)$  signal searched by  $K^-$  interaction with a **bound proton** in Carbon

$K^- p \rightarrow \Sigma^0 \pi^0$  detected via:  $(\Lambda\gamma)$   $(\gamma\gamma)$

Strategy:  $K^-$  absorption in the DC entrance wall, mainly  $^{12}\text{C}$  with H contamination (epoxy)



$m_{\pi^0\Sigma^0}$  resolution  $\sigma_m \approx 32 \text{ MeV}/c^2$  ;  $p_{\pi^0\Sigma^0}$  resolution:  $\sigma_p \approx 20 \text{ MeV}/c$ .

Negligible  $(\Lambda \pi^0 + \text{internal conversion})$  background =  $(3 \pm 1) \%$   $\rightarrow$  no I=1 contamination



# $\Sigma^0 \pi^0$ channel

$K^-$  nuclear absorption experiments .. long history .. BUT

$K^-$

- 1)  $m_{\pi\Sigma}$  spectra always cut at the **at-rest limit** 2)  $(\Sigma^\pm\pi^\mp)$  spectra suffer  $\Sigma(1385)$  contamination

P. J. Carlson, et al. Nucl. Phys. 74 642

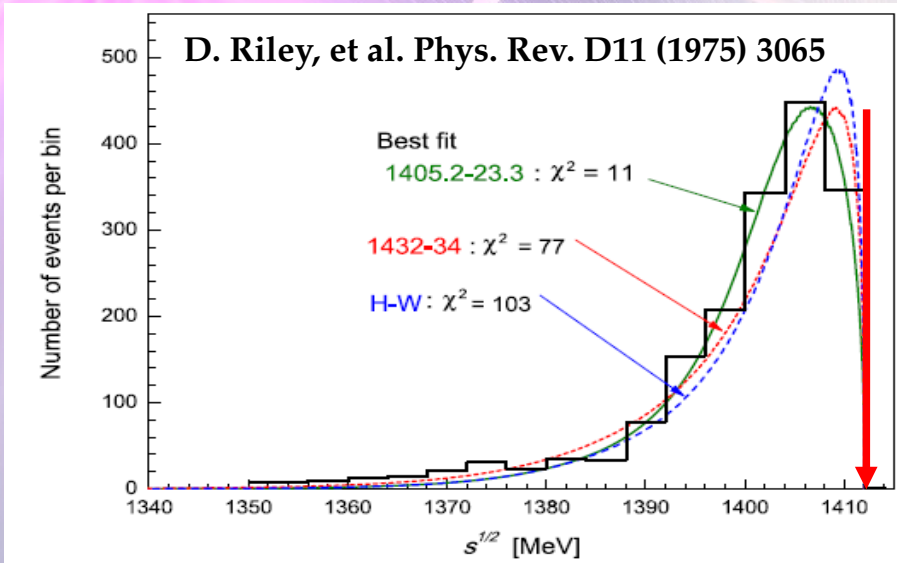
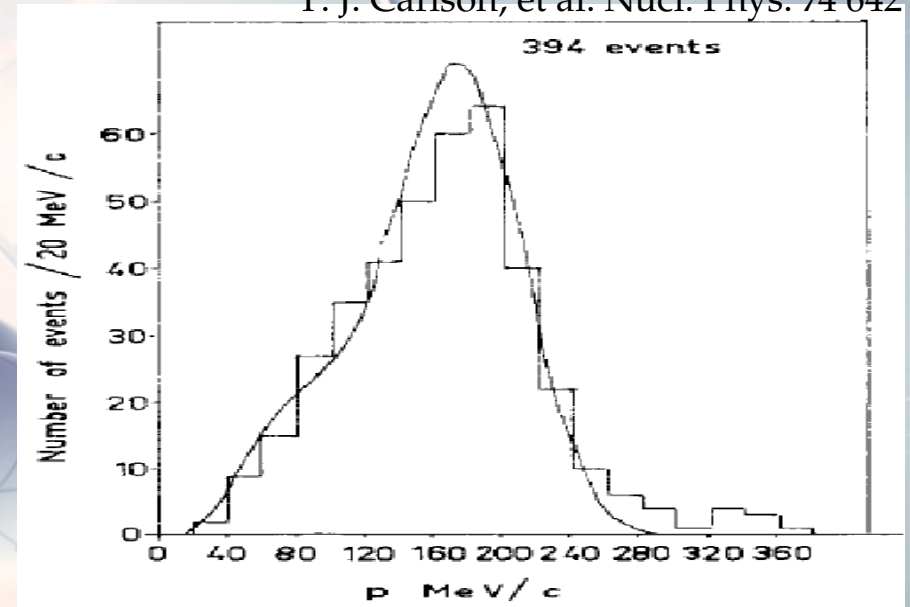
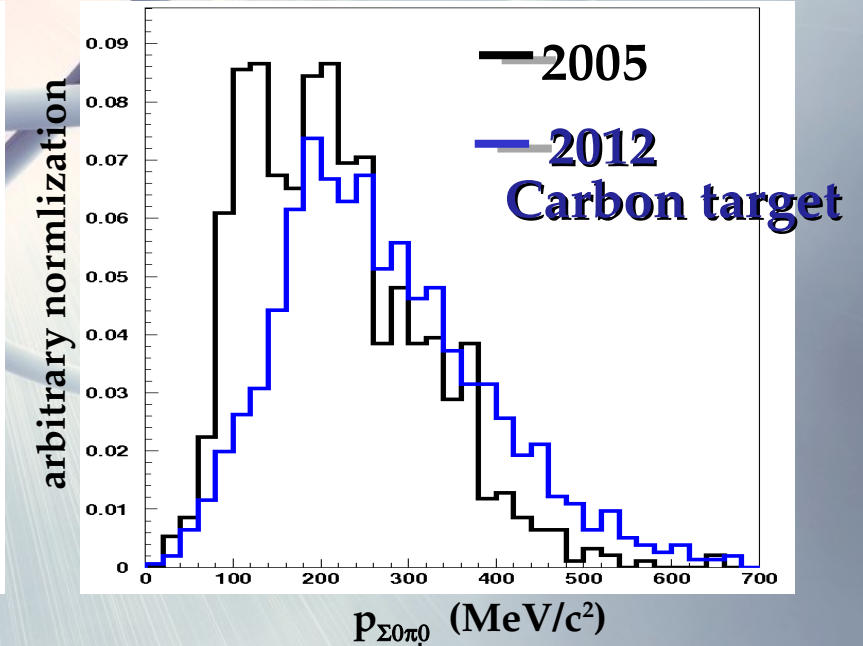
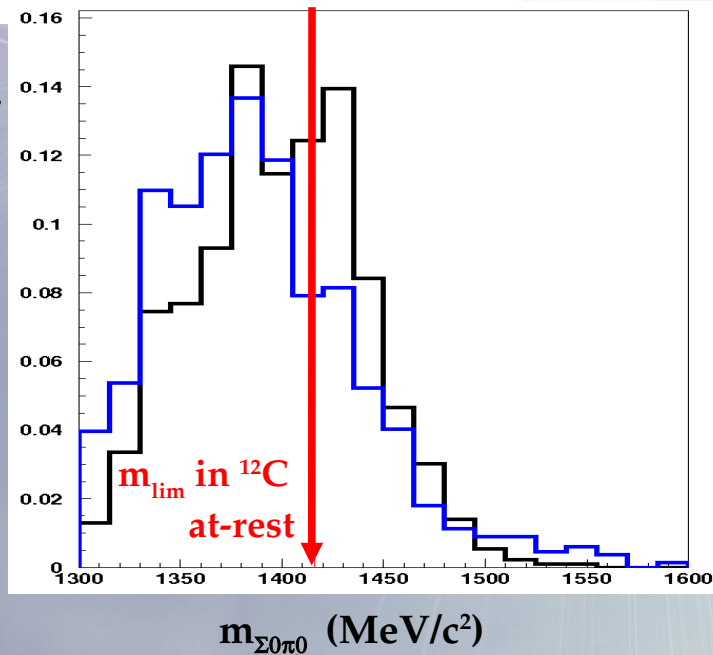


Fig. 6. Detailed differences in  $M_{\Sigma\pi}$  spectra among the Hyodo-Weise prediction and the present model predictions.



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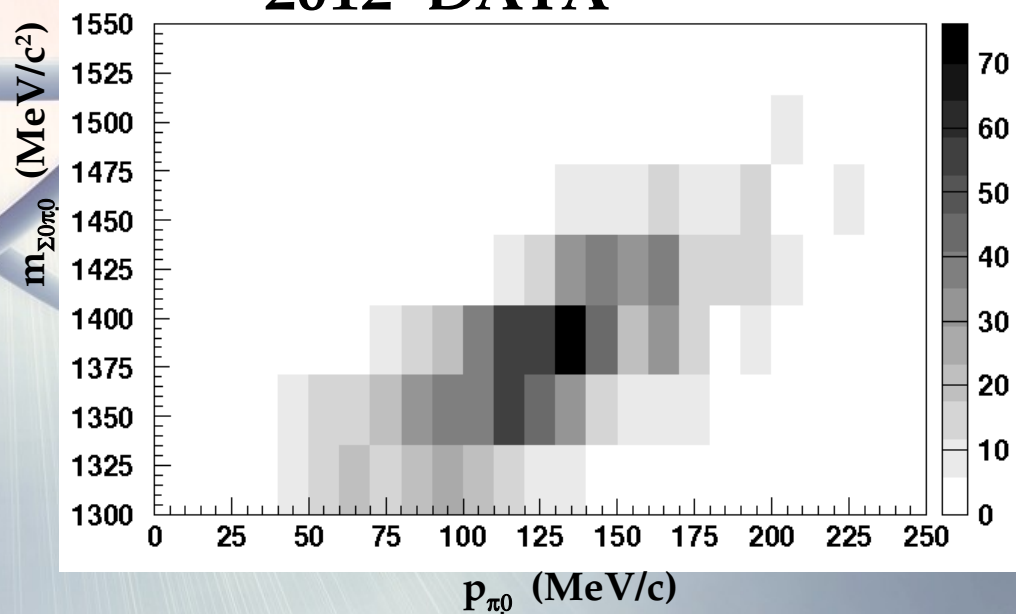
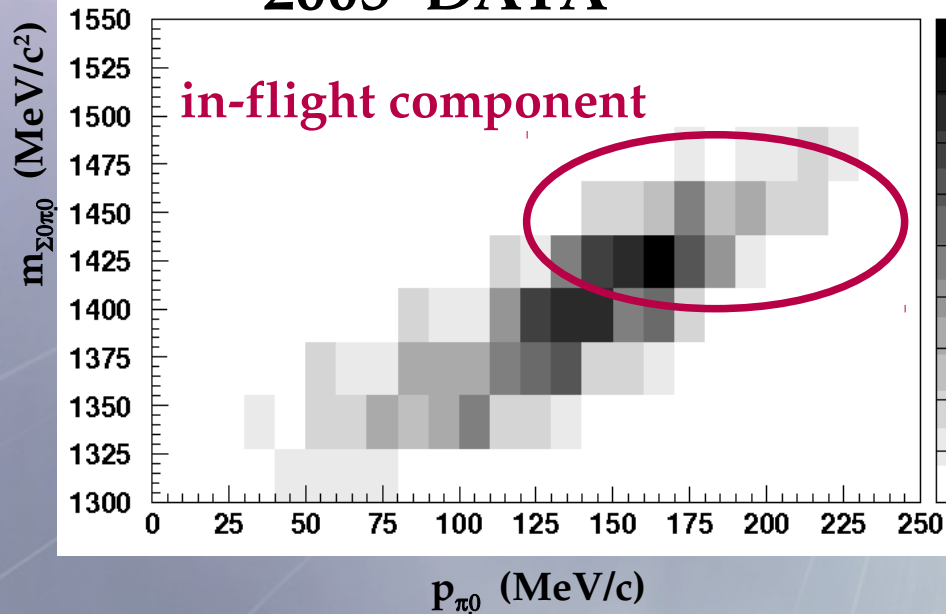
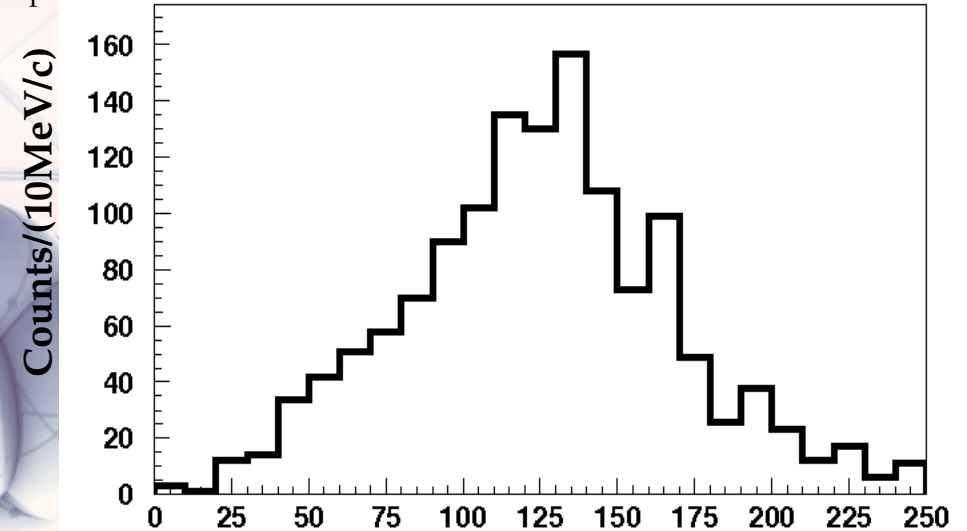
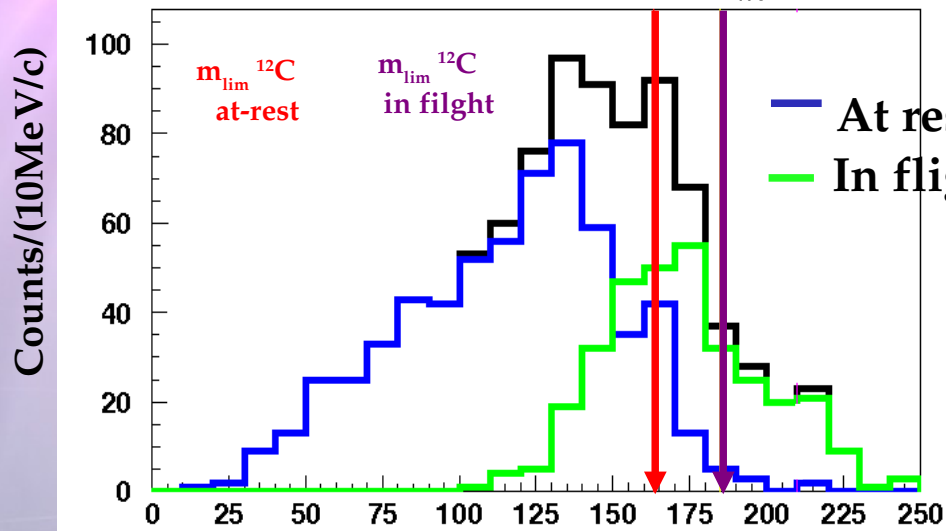
# $\Sigma^0 \pi^0$ channel

In-flight component ... **FIRST EVIDENCE IN  $K^-$  ABSORPTION MASS SPECTROSCOPY**

$K^-$

open a higher invariant mass region

$p_{\pi^0}$  resolution:  $\sigma_p \approx 12$  MeV/c



# $\Sigma^0 \pi^0$ channel

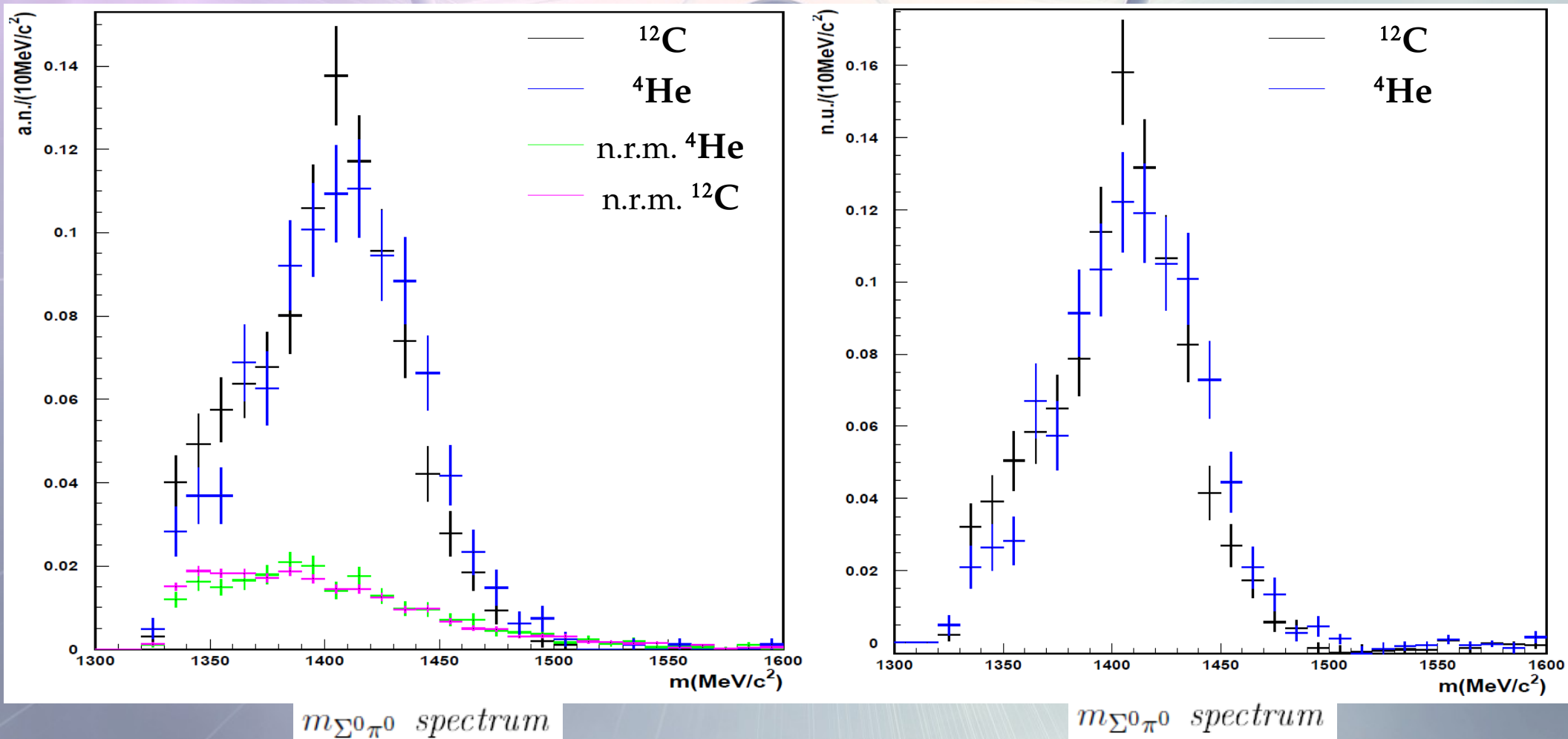
$K^-$

Invariant mass spectra with mass hypothesis on  $\Sigma^0$  and  $\pi^0$  *non resonant misidentification background subtracted (right)*

$$\sigma_m \approx 17 \text{ MeV}/c^2 \text{ } (^{12}\text{C}) \quad \sigma_m \approx 15 \text{ MeV}/c^2 \text{ } (^4\text{He})$$

Similar  $m_{\pi^0\Sigma^0}$  shapes due to the similar kinematical thresholds for  $^4\text{He}$  and  $^{12}\text{C}$ .

## 2005 DATA



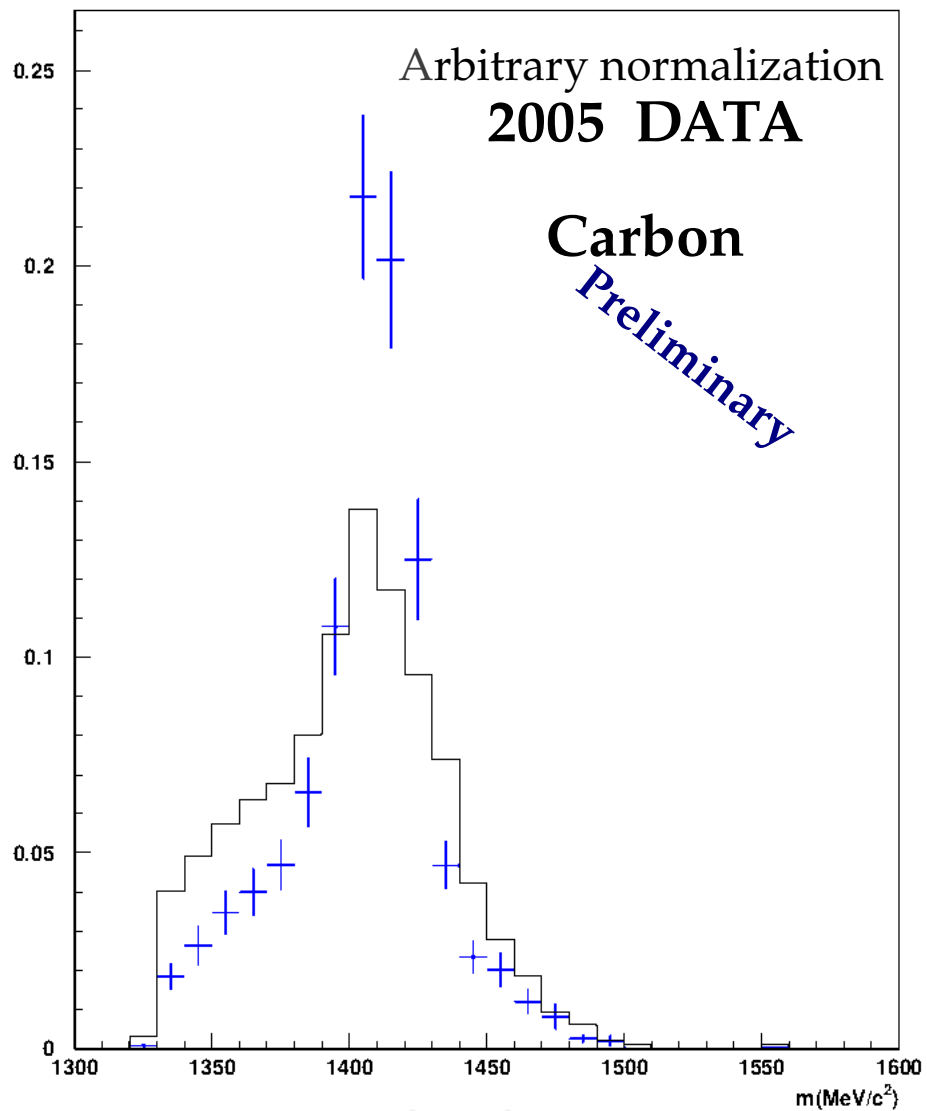
# $\Sigma^0 \pi^0$ channel

Acceptance corrected  $m_{\pi^0 \Sigma^0}$  spectra, DC wall (left) DC gas (right)

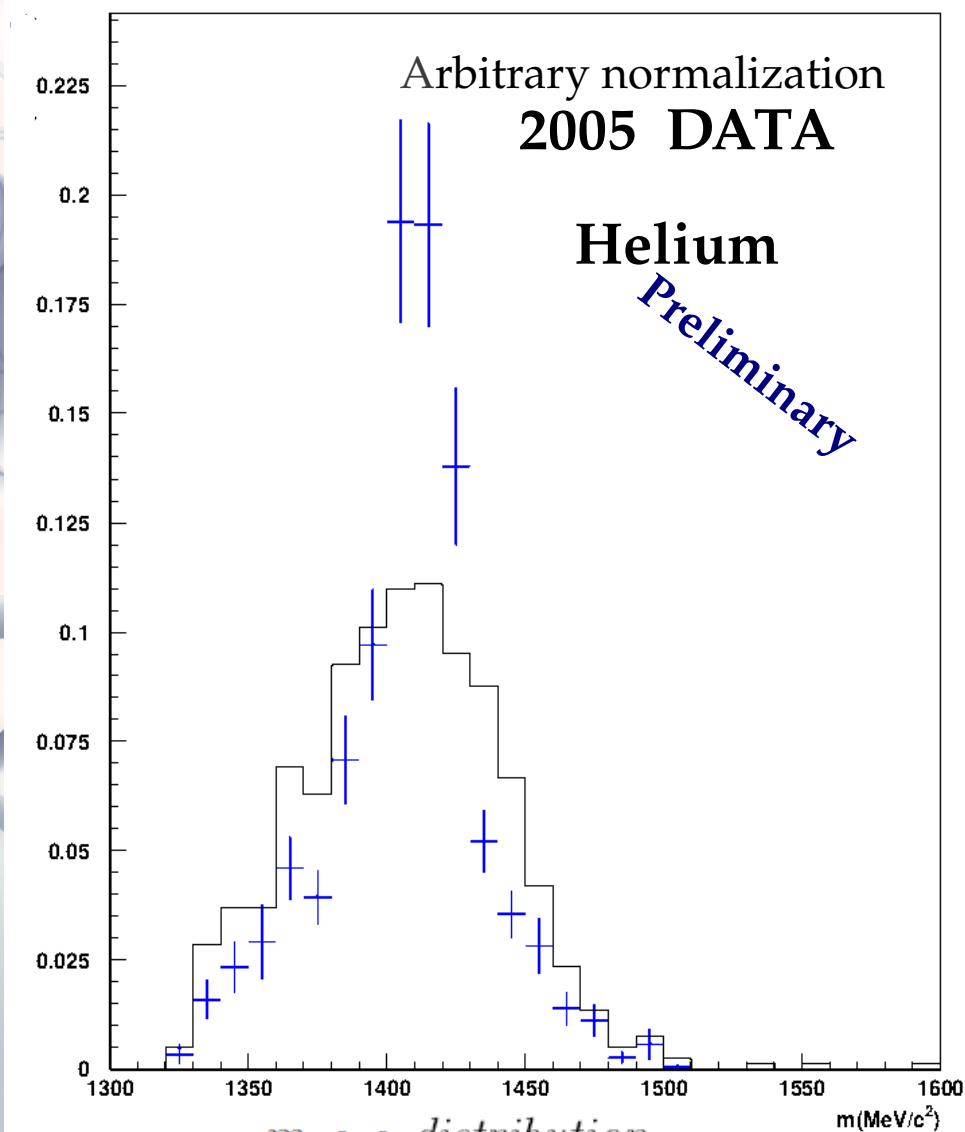
$K^-$

Acceptance function evaluated in 8 intervals of  $p_{\pi^0 \Sigma^0}$  (between 0 and 700 MeV/c) 8 intervals

of  $\theta_{\pi^0 \Sigma^0}$  (between 0 and 3.15 rad) 30 intervals of  $m_{\pi^0 \Sigma^0}$  (between 1300 and 1600 MeV/c<sup>2</sup>)



$m_{\Sigma^0 \pi^0}$  distribution

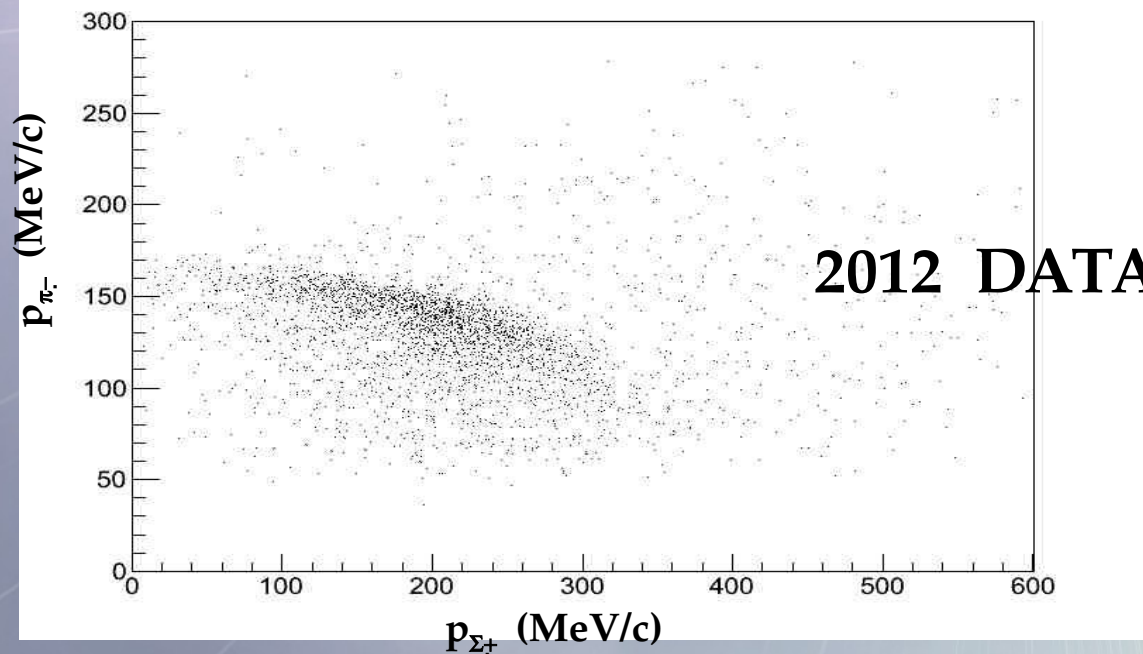
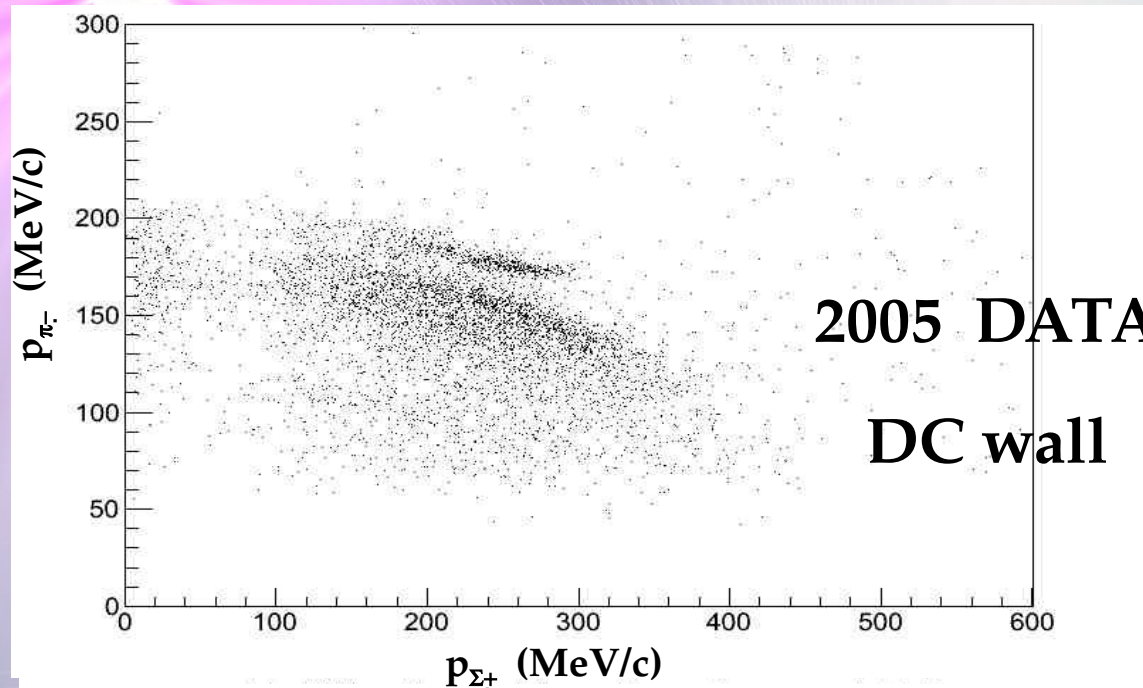


$m_{\Sigma^0 \pi^0}$  distribution

# $\Sigma^+ \pi^-$ channel ... A NEW POWERFUL TOOL (A. Scordo) !

$K^- p \rightarrow \Sigma^+ \pi^-$  detected via:  $(p\pi^0) \pi^-$

$K^-$



BEFORE ...

$K^-H$  interaction probability estimate  
based on  $K^-$  interaction AT-REST in  
hydrocarbons mixture data (Lett.  
Nuovo Cimento, C 1099 (1972))

order of 1% !!!

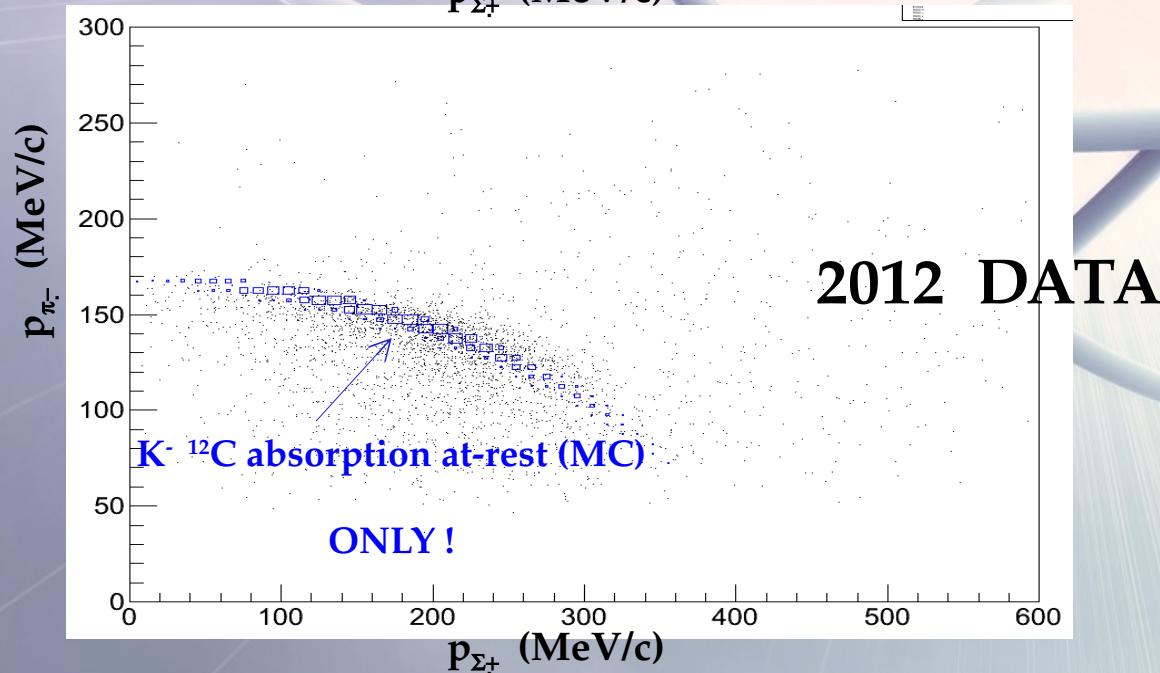
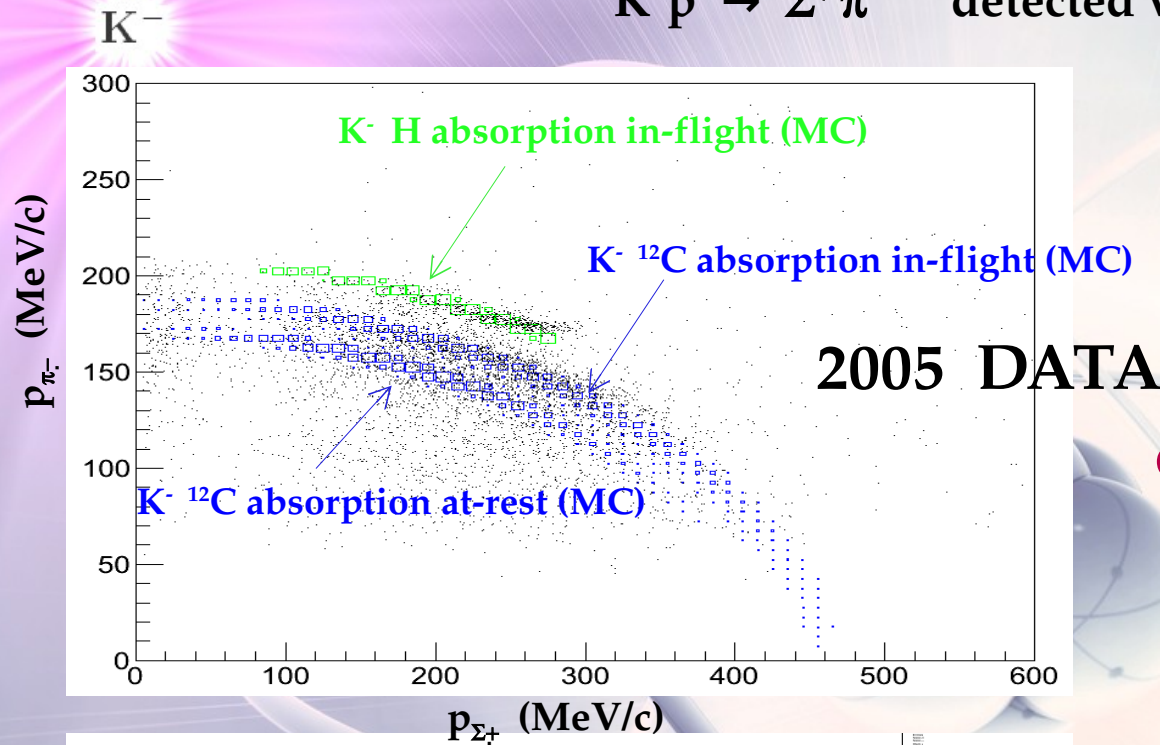
NOW

Thanks to the excellent  $p_{\pi^-}$  resolution

...

# $\Sigma^+ \pi^-$ channel ... A NEW POWERFUL TOOL (A. Scordo) !

$K^- p \rightarrow \Sigma^+ \pi^-$  detected via:  $(p\pi^0) \pi^-$



Complete understanding of different nuclear targets in different KLOE materials

AND

$K^-$  H contribution  $\sim 20\%$

... in-flight the situation is extremely different

# Fit of $\Sigma^0\pi^0$ spectrum in C

$K^-$

8 component fit, simultaneously  $m_{\Sigma^0\pi^0}$  &  $p_{\Sigma^0\pi^0}$ :

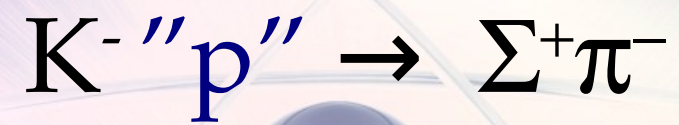
- Breit-Wigner resonant component  $K^- C$  at-rest/in-flight.  $(M,\Gamma) = (1405 \div 1430, 5 \div 52)$ 
  - Non resonant  $\Sigma^0\pi^0$   $K^- H$  production at-rest/in-flight
  - Non resonant  $\Sigma^0\pi^0$   $K^- C$  production at-rest/in-flight
    - $\Lambda\pi^0$  background ( $\Sigma(1385) + I.C.$ )
  - non resonant misidentification (*n.r.m.*) background

$K^- {}^{12}C \rightarrow \Sigma^0\pi^0 + {}^{11}B$  (Boron spectator, left in ground state)

secondary interactions not taken into account.



$K^-$



bound proton in  ${}^4\text{He}$  /  ${}^{12}\text{C}$



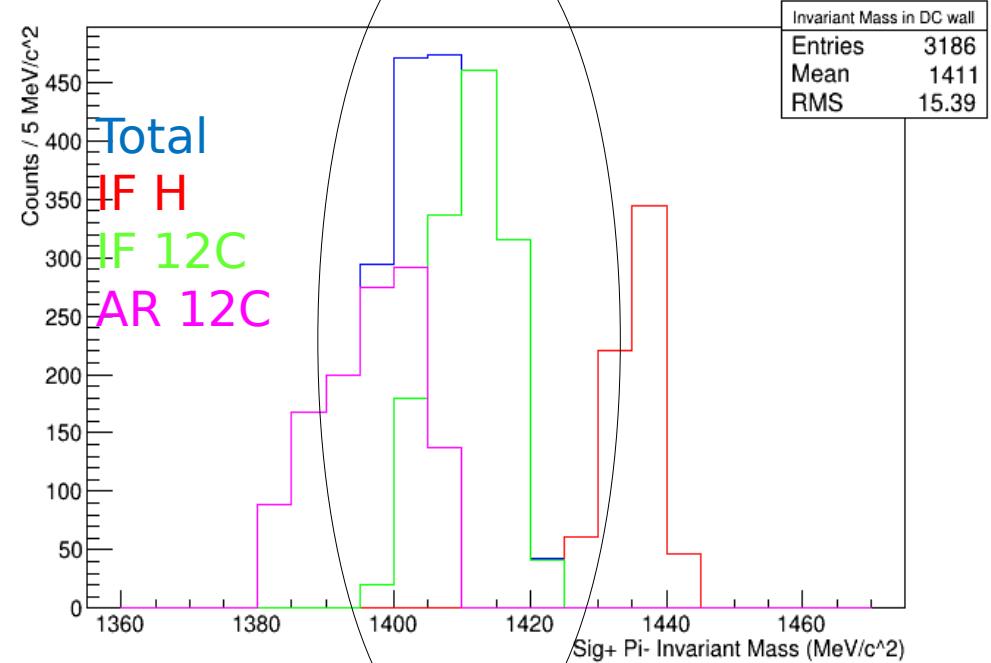
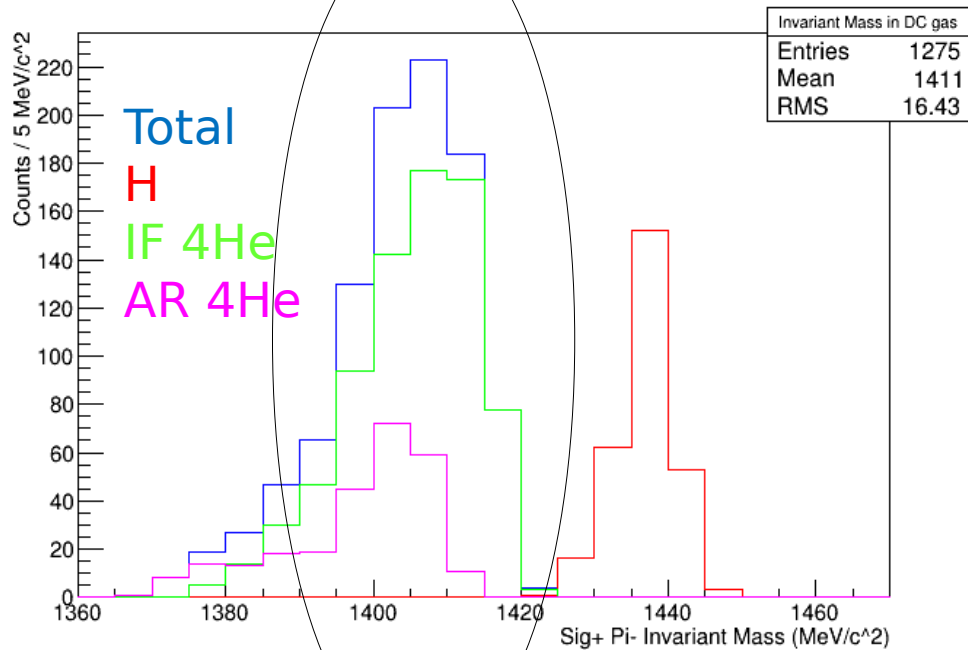
# $\Sigma^+\pi^-$ invariant mass spectra

$K^-$

$K^- p \rightarrow \Sigma^+ \pi^-$  detected via:  $(p\pi^0) \pi^-$

The excellent momentum resolution for  $\pi^-$  enables to **disentangle in-flight from at-rest  $K^-$  capture**

Hint: if resonant production contribution is important a high mass pole appears!



# $K^-$ Resonant VS non-resonant

Another unsolved question ..



how much comes from resonance ?

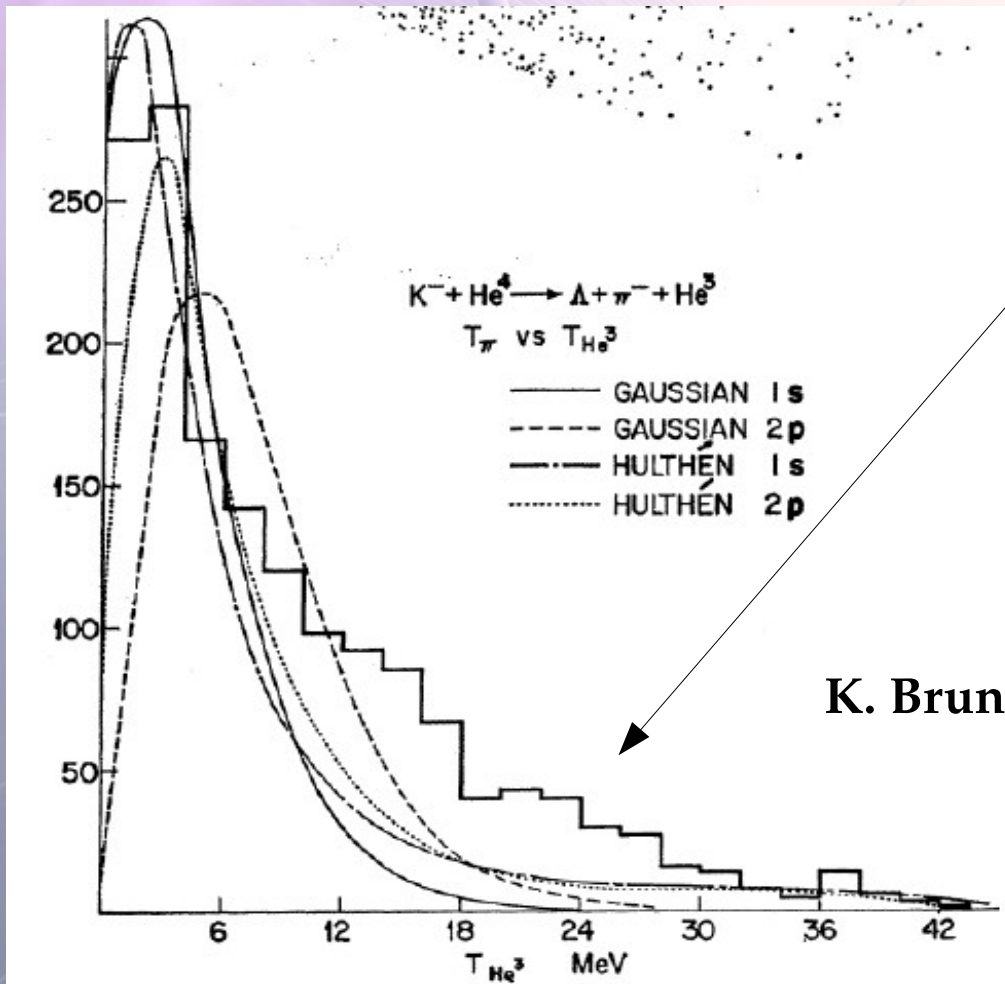
Investigated using:



# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ ... the idea

Bubble chamber experiments exhibit two components:

- Low momentum  $\Lambda \ \pi^-$  pair  $\rightarrow$  S-wave,  $I=1$ , **non-resonant** transition amplitude.
- High momentum  $\Lambda \ \pi^-$  pair  $\rightarrow$  P-wave **resonant formation** ?



Also exists in S-state K-mesic atom  
as a result of the  
**three body structure of the system**

( $K = 1, n=2, \ ^3\text{He} = 3$ )

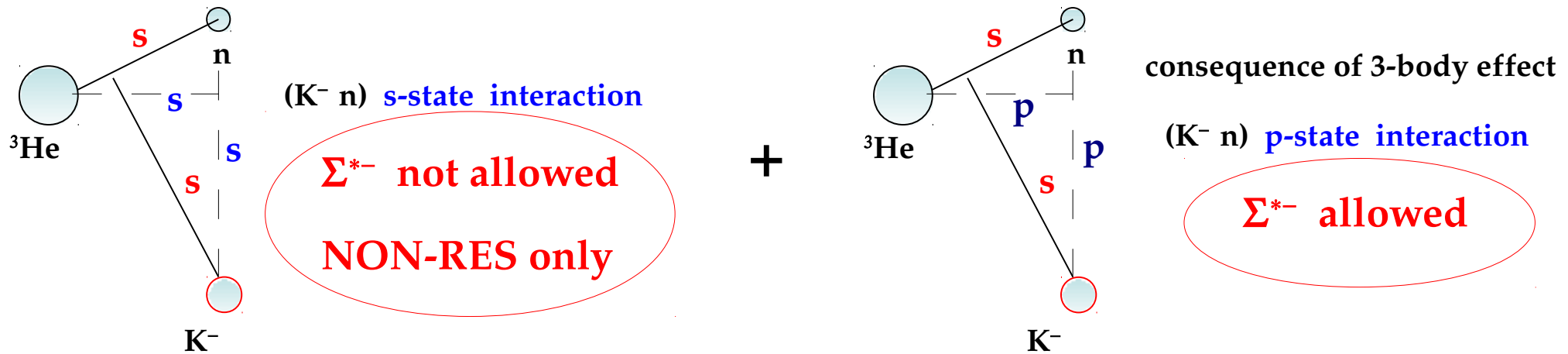
K. Brunnel et al., Phys.Rev. D2 (1970) 98

# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ ... the idea

$K^-$

$K^-(s=0) \ ^4\text{He}(s=0) \ n(s=1/2) \ \Sigma^{*-}(s=3/2) \rightarrow$  **resonance p-wave only**

## atomic s-state capture:



- $(K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He})$  absorptions from  $(n \ s)$  - atomic states are assumed  $\rightarrow$   $^4\text{He}$  bubble chamber data (Fetkovich, Riley interpreted by Uretsky, Wienke)

- Coordinates recoupling enables for P-wave resonance formation

# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ ... the strategy

$K^-$

- **Fit of the  $p_{\Lambda\pi^-}$  observed distribution** using calculated distributions :

$$P_s^s(p_{\Lambda\pi}) = |\Psi_N(p_{\Lambda\pi})|^2 |f^s(p_{\Lambda\pi})|^2 \rho \quad \text{non-resonant}$$

$$P_s^p(p_{\Lambda\pi}) = |\Psi_N(p_{\Lambda\pi})|^2 c^2 |2f^{\Sigma^*}(p_{\Lambda\pi})|^2 \rho/3 (kp_{\Lambda\pi})^2 \quad \text{resonant}$$

- **To determine *for the first time* the ratio resonant/non-res.**

↓

$$|f^{\text{N-R}}_{\Lambda\pi}| \text{ given the fairly well known } |f^{\Sigma^*}_{\Lambda\pi}|$$

# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ ... calculated reactions

Calculated primary hadronic interactions:

$K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$

At-rest: S-wave non-Res / P-wave  $\Sigma(1385)$  Res

In-flight: S-wave non-Res / P-wave  $\Sigma(1385)$  Res

$K^- \ ^4\text{He} \rightarrow \Sigma^0 \ \pi^- \ ^3\text{He}$

At-rest: S-wave non-Res / P-wave  $\Sigma(1385)$  Res

In-flight: S-wave non-Res / P-wave  $\Sigma(1385)$  Res

$K^- \ ^4\text{He} \rightarrow (\Sigma \ \pi)^0 \ ^3\text{H}$

At-rest: S-wave non-Res / S-wave  $\Lambda(1405)$  Res /  
P-wave  $\Sigma(1385)$  Res

In-flight: S-wave non-Res / S-wave  $\Lambda(1405)$  Res /  
P-wave  $\Sigma(1385)$  Res

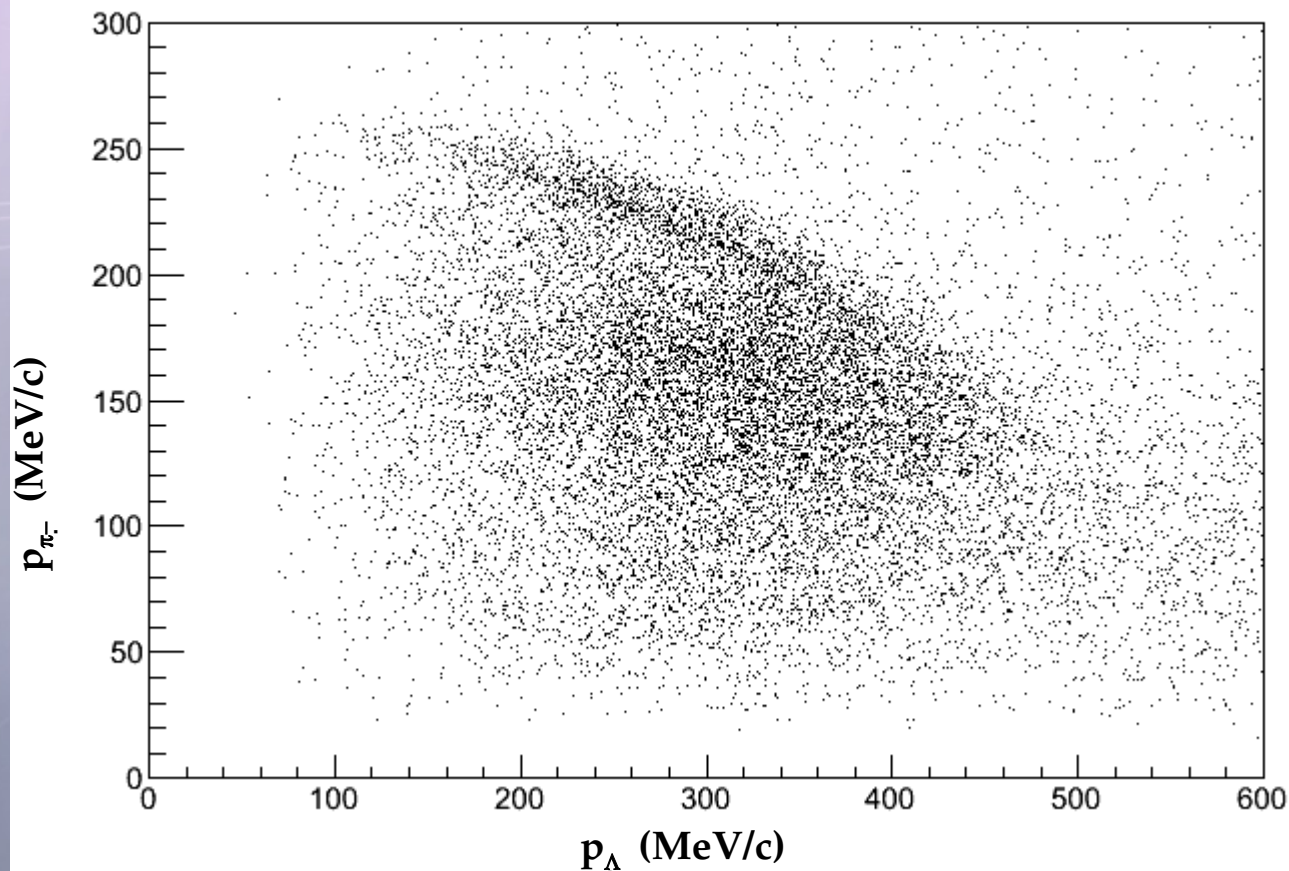
Channel:  $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$  ... calculated reactions

Calculated secondary hadronic interactions:

EACH INTERNAL CONVERSION PROCESS:

$$\Sigma \text{ p/n} \rightarrow \Lambda \text{ p/n}$$

was calculated for both P-wave and S-wave produced  $\Sigma$ s.



# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ ... calculated reactions

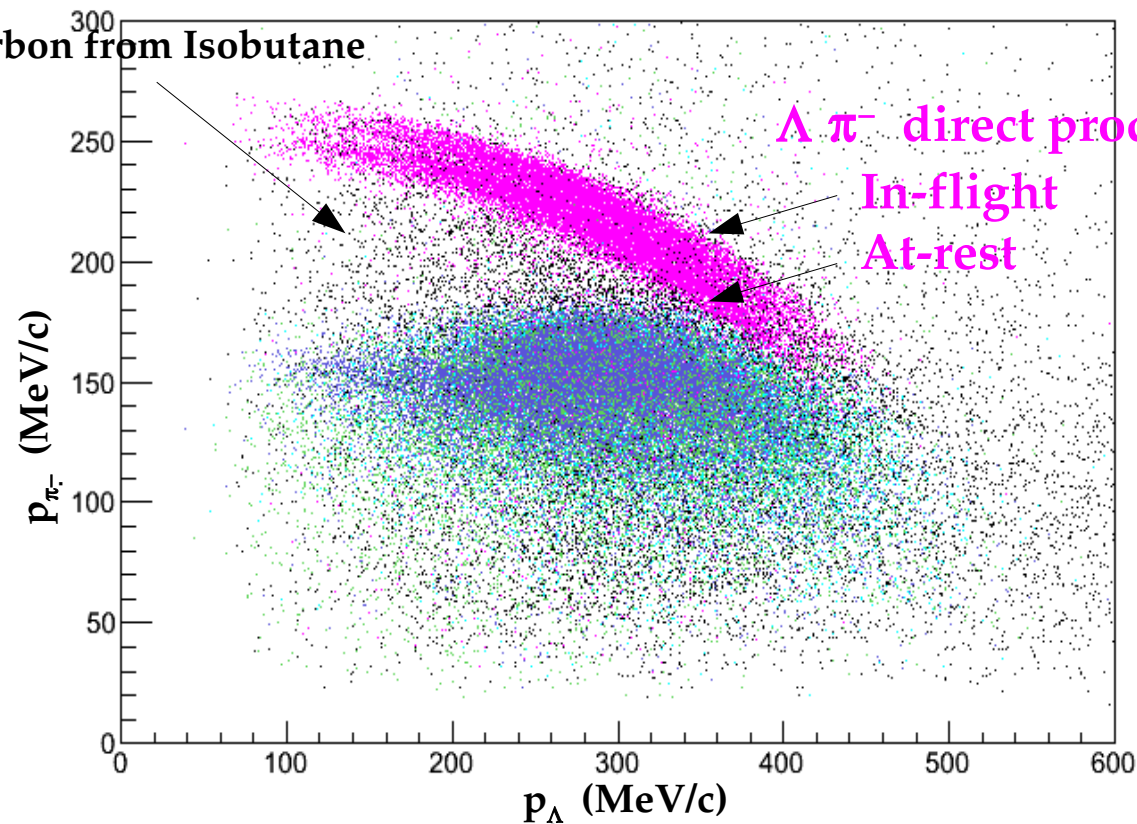
Calculated secondary hadronic interactions:

EACH INTERNAL CONVERSION PROCESS:

$$\Sigma \text{ p/n} \rightarrow \Lambda \text{ p/n}$$

was calculated for both P-wave and S-wave produced  $\Sigma$ s.

Some Carbon from Isobutane



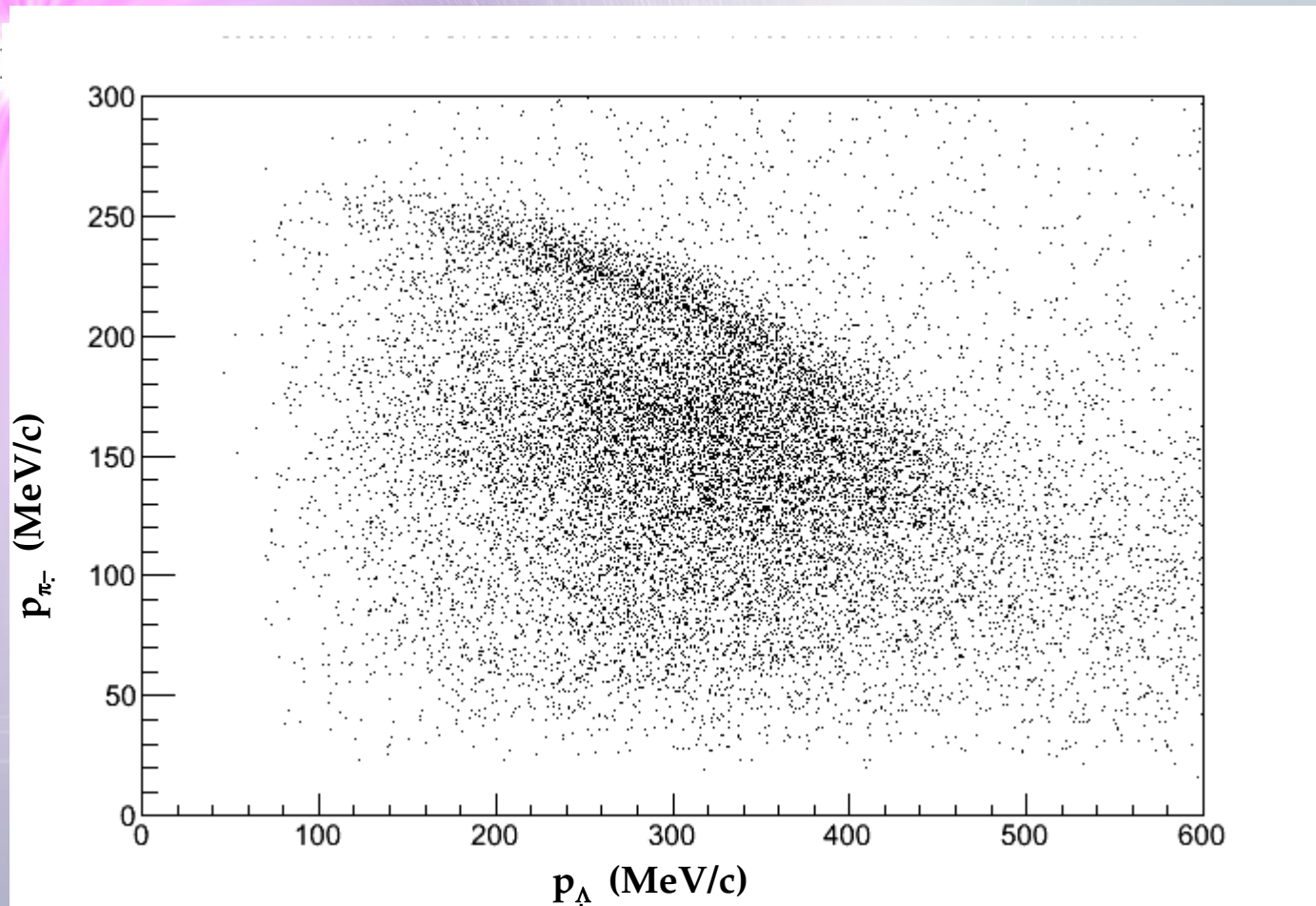
$\Sigma^0$  p conversion

$\Sigma^0$  n conversion

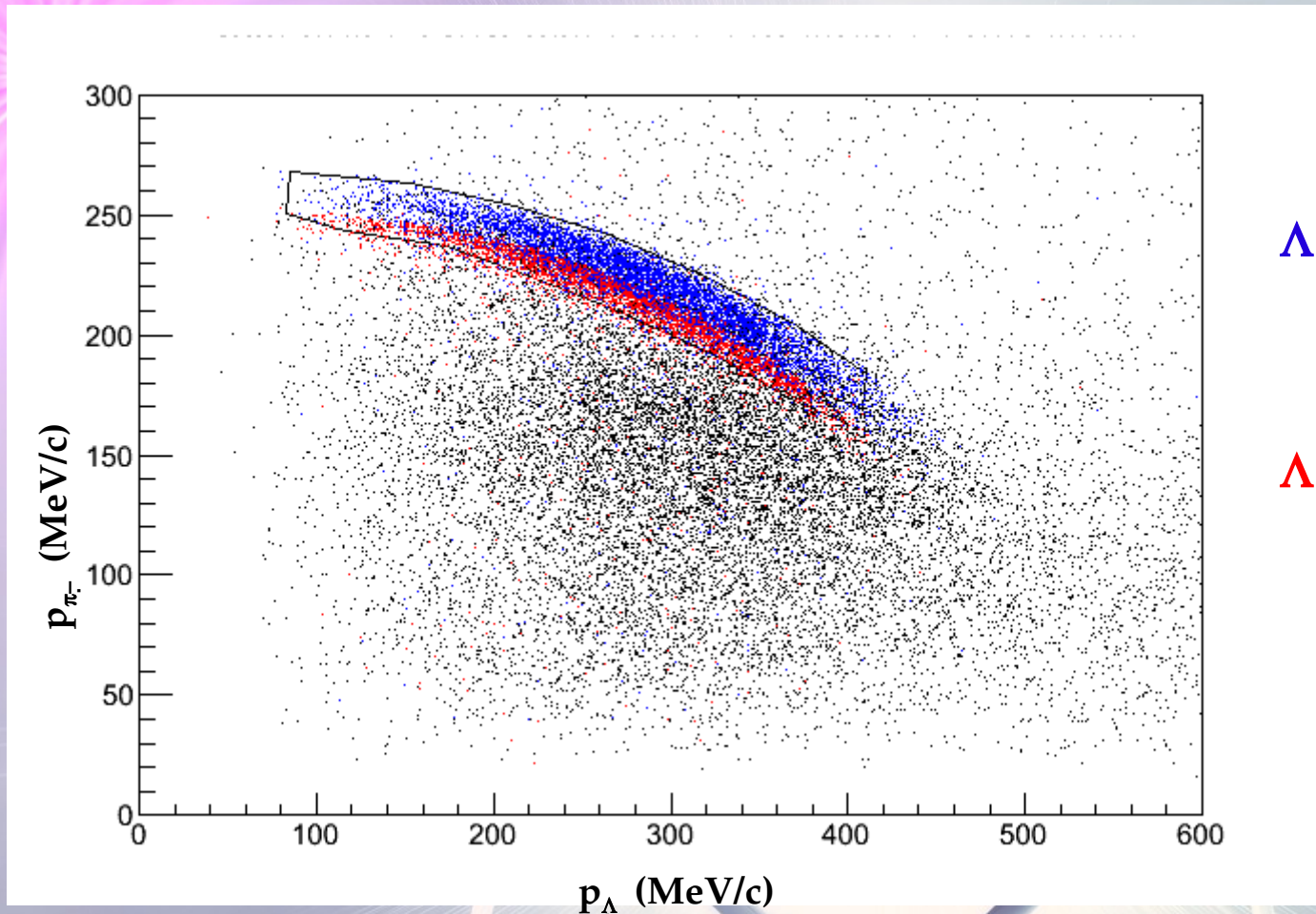
$\Sigma^+$  n conversion



# $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ events selection



# $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ events selection



$\Lambda \pi^-$  direct production  
In-flight RES + N-R

$\Lambda \pi^-$  direct production  
At-rest RES + N-R

- **CUT** based on MC simulations used to select  $\Lambda \pi^-$  direct production events
- At-rest **CAN NOT** be separated from In-flight  $\rightarrow$  global fit performed
- Background sources:
  - $\Lambda \pi^-$  events from  $\Sigma$  p/n  $\rightarrow$   $\Lambda$  p/n conversion
  - $\Lambda \pi^-$  events from  $K^- \ ^{12}\text{C}$  absorptions in Isobutane

# $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ background

- $\Sigma$  p/n  $\rightarrow$   $\Lambda$  p/n conversion:

Each possible conversion channel was simulated

$\Sigma^0$  p /  $\Sigma^0$  n /  $\Sigma^+$  n / At-rest / In-flight / from RES and N-R produced  $\Sigma$ s

- $\Lambda \pi^-$  events from  $K^- \ ^{12}\text{C}$  absorptions in Isobutane (90% He, 10%  $\text{C}_4\text{H}_{10}$ ):

$K^- \ ^{12}\text{C}$  DATA in the KLOE DC wall are used

estimated contribution:

$$N_{\text{KC}}/N_{\text{KHe}} = (n_{\text{KC}}/n_{\text{KHe}}) \cdot (\sigma_{\text{KC}}/\sigma_{\text{KHe}}) \cdot (\text{BR}_{\text{KC}}(\Lambda \pi^-)/\text{BR}_{\text{KHe}}(\Lambda \pi^-)) \sim 1/3$$

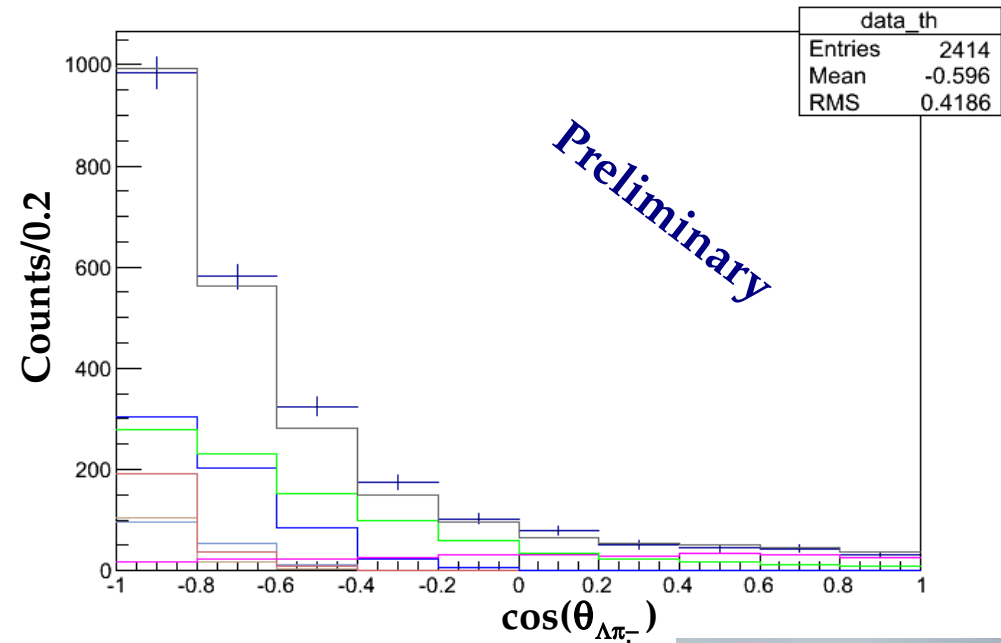
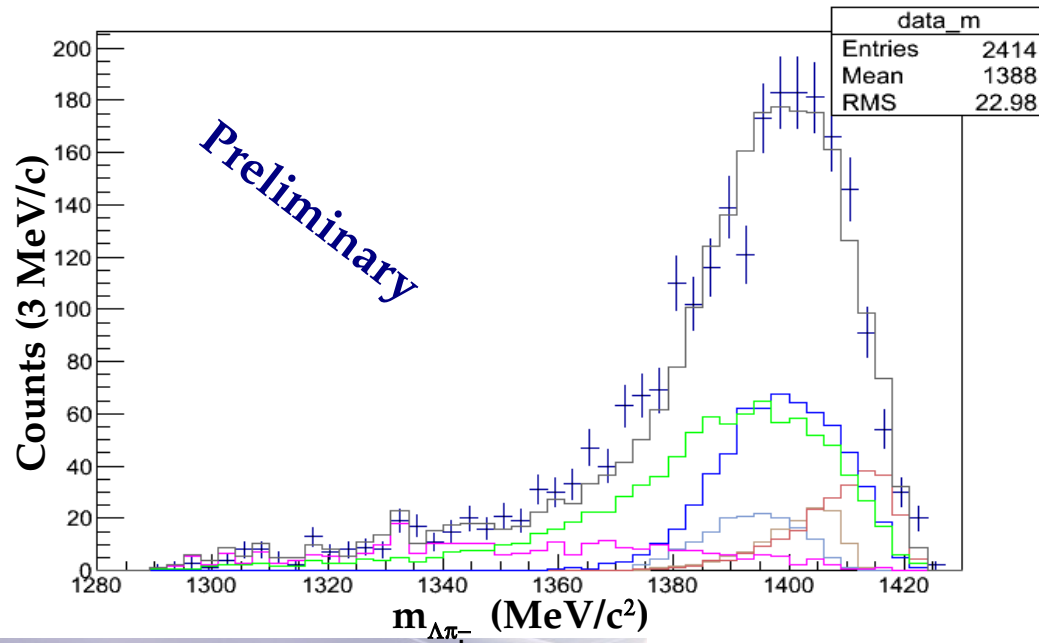
Nuovo Cimento 39 A 338-347 (1977)

$K^- \ ^{12}\text{C}$  still not calculated:

- uncertain initial state of K meson  $l_{\text{K}} = 1, 2, 3$
- 4 nucleons in s-orbit, 8 nucleons in p-orbit
- final state hyperon interactions

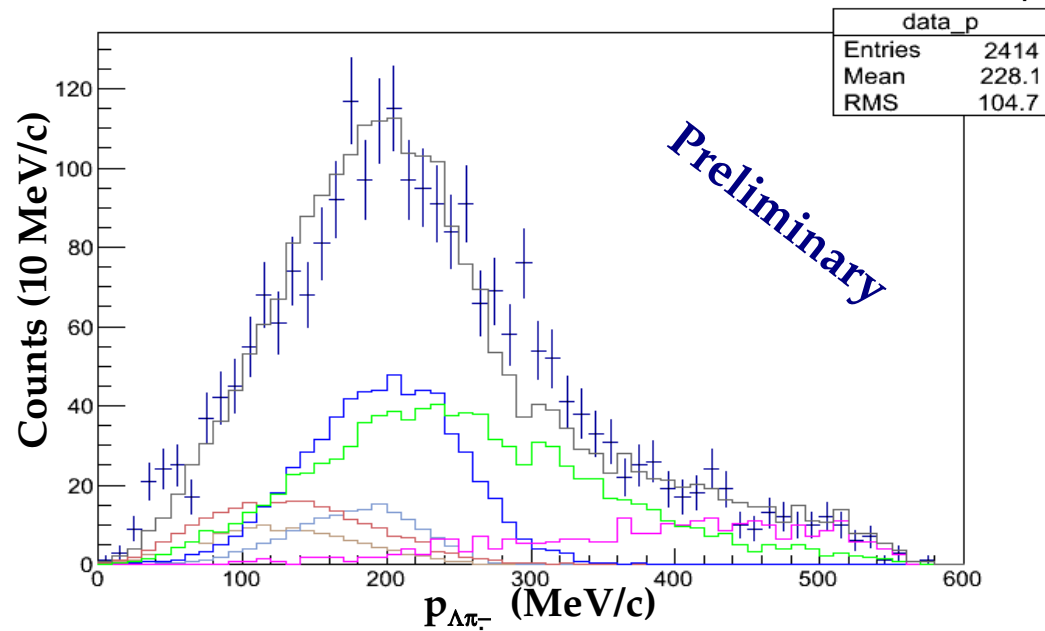
# $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ fit

**Simultaneous fit** ( $p_{\Lambda\pi^-} - m_{\Lambda\pi^-} - \theta_{\Lambda\pi^-}$ ) leaving the ratio At-rest /In-flight and  $^{12}\text{C}$  contamination to vary around the estimated values within errors:



Global fit

- $\Lambda \pi^-$  At-rest N-R
- $\Lambda \pi^-$  At-rest RES
- $\Lambda \pi^-$  In-flight N-R
- $\Lambda \pi^-$  In-flight RES
- $\Lambda \pi^-$  events from  $K^- \ ^{12}\text{C}$
- $\Sigma \text{ p/n} \rightarrow \Lambda \text{ p/n}$  conversion



# $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ fit

**Simultaneous fit** ( $p_{\Lambda\pi^-} - m_{\Lambda\pi^-} - \theta_{\Lambda\pi^-}$ ) leaving the ratio At-rest /In-flight and  $^{12}\text{C}$  contamination to vary around the estimated values within errors:

- $\chi^2 / (\text{ndf} - \text{np}) = 1.2$
- **(At-rest RES)/(At-rest N-R) =  $1.26 \pm 0.06$**
- **(In-flight RES)/(In-flight N-R) =  $2.59 \pm 0.3$**
  
- **(In-flight) / (At-rest) =  $2.9 \pm 0.5$  → consistent with the estimate in  $\Sigma^+\pi^-$**
- **$\Sigma$  p/n →  $\Lambda$  p/n conversion =  $(12.2 \pm 0.8)\%$**
- **$\Lambda \pi^-$  events from  $K^- \ ^{12}\text{C} = (38 \pm 1)\%$**

**Preliminary**

# Conclusions and perspectives ..

- $m_{\Sigma\pi^-}$  spectra show a **high invariant mass component** → associated to in-flight  $K^-$  capture
- PRELIMINARY  $\Lambda\pi^-$  first measurement of RES/N-R ratio in nuclear  $K^-$  absorption.

Next steps ...

- Same analysis is ongoing for  $\Sigma^0\pi^-$  → extraction of  $|f^{N-R}_{\Sigma^0\pi^-}(I=1)|$
- Similar description of  $\Sigma^+\pi^-$  and  $\Sigma\pi^+$  production → extraction of  $|f^{N-R}_{\Sigma^+\pi^-}|$  and  $|f^{N-R}_{\Sigma-\pi^+}|$ , a comparison of these could give an estimate of  $|f^{N-R}_{\Sigma^+\pi^-}(I=0) + f^{N-R}_{\Sigma^+\pi^-}(I=1)|$  against  $|f^{N-R}_{\Sigma^+\pi^-}(I=0) - f^{N-R}_{\Sigma^+\pi^-}(I=1)|$
- Branching ratio modifications in different targets (see A. Ohnishi et al., Phys. Rev. C 56 5 (1997) 2767) & **Density dependence of  $m_{\Sigma\pi}$  and  $p_{\Sigma\pi}$**  (see L. R. Staronski, S. Wycech, Nucl. Phys. 13 (1987) 1361 / A. Cieplý, E. Friedman, A. Gal, V. Krejčířík - Phys.Lett.B698 (2011) 226-230)

**Shedding light on the nature of the  $\Lambda(1405)$  and its behaviour in nuclear matter is crucial to understand the role of strangeness in our universe.**

**K<sup>-</sup>**



**Thanks**

$K^-$

**SPARE SLIDES ...**





# AMADEUS & DAΦNE

$K^-$

Completely neutral channel:

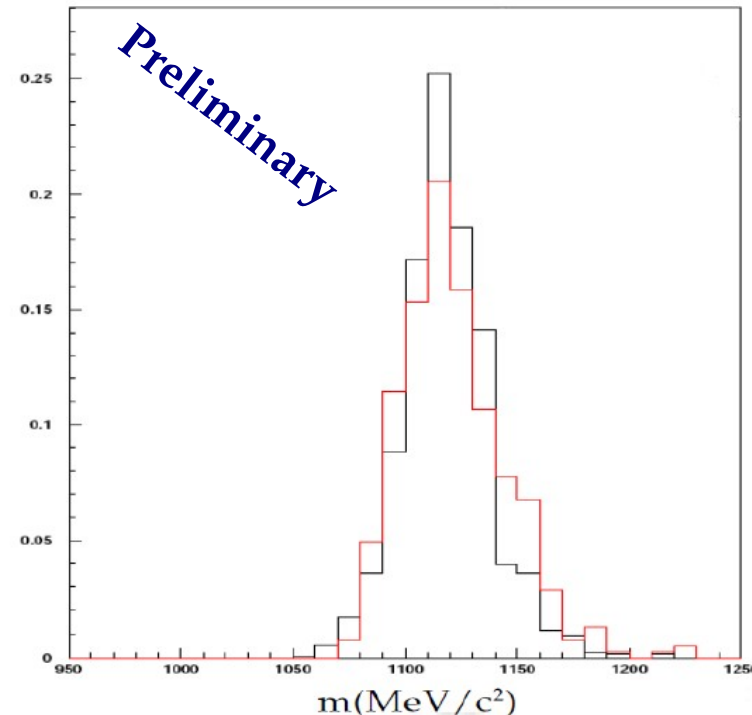


Possibility to detect neutrons!

black MC      red data

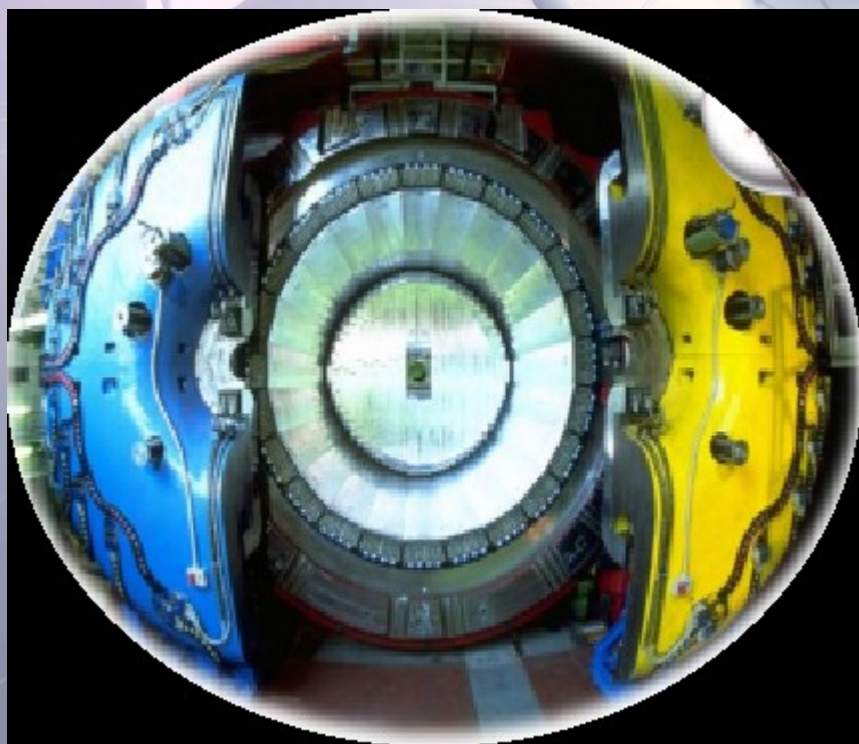
Perspective:  $\Sigma^-\pi^+ \rightarrow (n\pi^-)\pi^+$

a. u. / (10MeV/c<sup>2</sup>)



## KLOE

- 96% acceptance,
- optimized in the energy range of all charged particles involved
- good performance in detecting photons (and neutrons checked by kloNe group (M. Anelli et al., Nucl Inst. Meth. A 581, 368 (2007)))

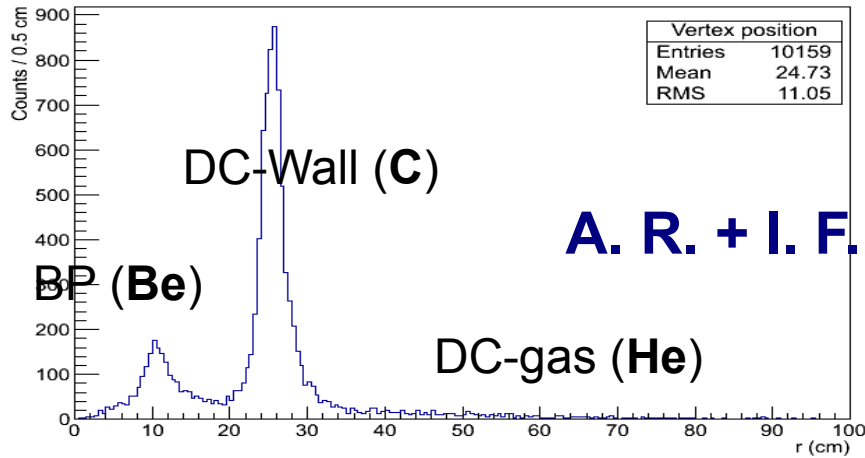


# KLOE data on $K^-$ nuclear absorption

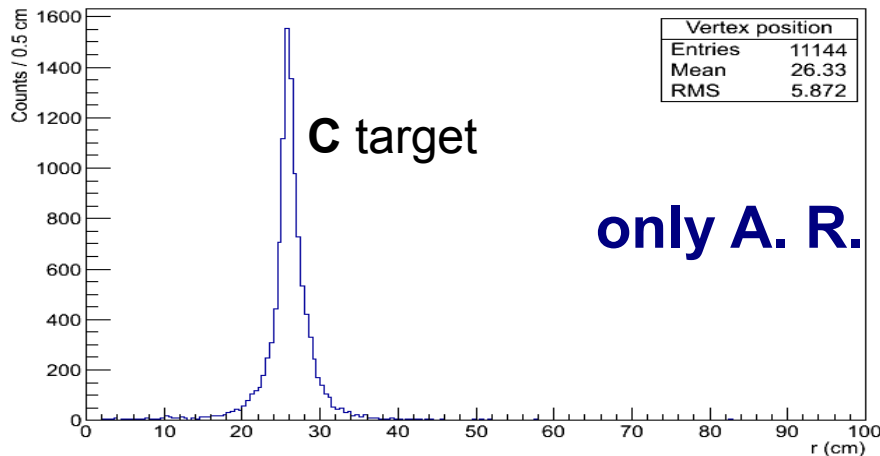
Use of two different data samples:

- KLOE data from 2004/2005 (2.2 fb<sup>-1</sup> total, 1.5fb<sup>-1</sup> analyzed)
- Dedicated run in november/december 2012 with a Carbon target 4/6 mm thickness (~90 pb<sup>-1</sup>; analyzed 37 pb<sup>-1</sup>, x1.5 statistics)

Position of the  $K^-$  hadronic interaction inside KLOE:



2005 data

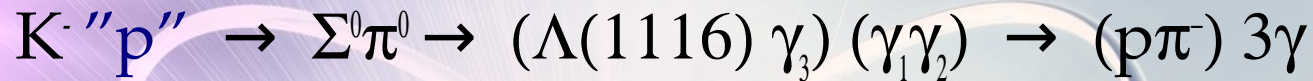


2012 with Carbon target



# Photon clusters identification

$K^-$



1) 3 neutral clusters selection ( $E_{cl} > 20$  MeV) not from  $K^+$  decay ( $K^+ \rightarrow \pi^+ \pi^0$ )

2) photon clusters selection:  $\chi_t^2 = t^2/\sigma_t^2$  where  $t = t_i - t_j$

time of flights in light speed hypothesis.

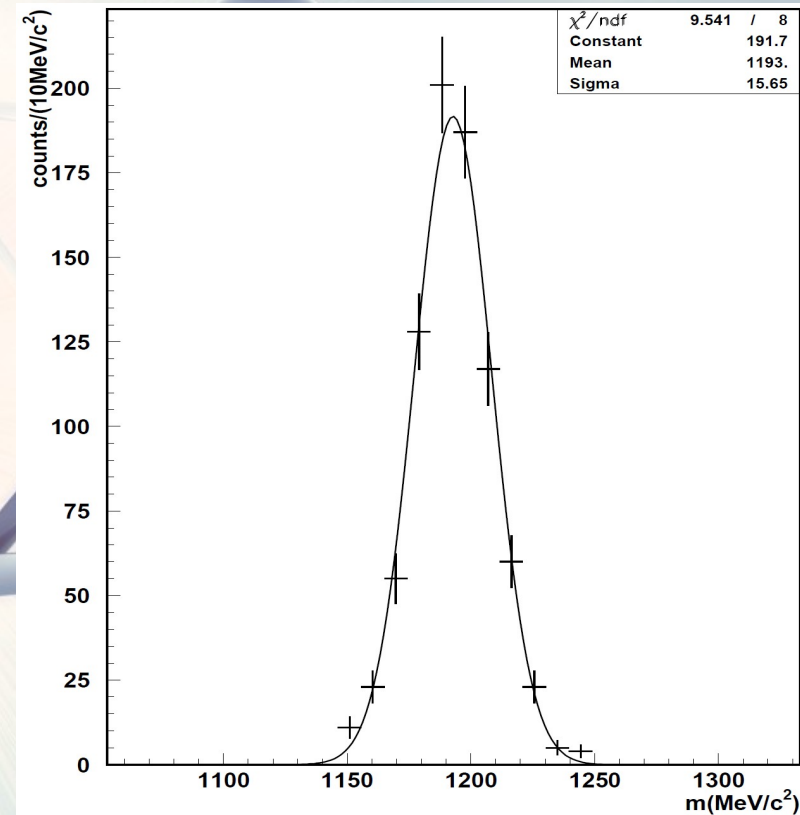
Three photons in time from the  $\Lambda$  decay vertex  $r_\Lambda$

3) photon clusters identification:  $\gamma_3$  from  $\pi^0 \rightarrow \gamma_1 \gamma_2$

$$\chi_{\pi\Sigma}^2 = \frac{(m_{\pi^0} - m_{ij})^2}{\sigma_{ij}^2} + \frac{(m_{\Sigma^0} - m_{k\Lambda})^2}{\sigma_{k\Lambda}^2}$$

$i, j$  and  $k$  candidate photon cluster.

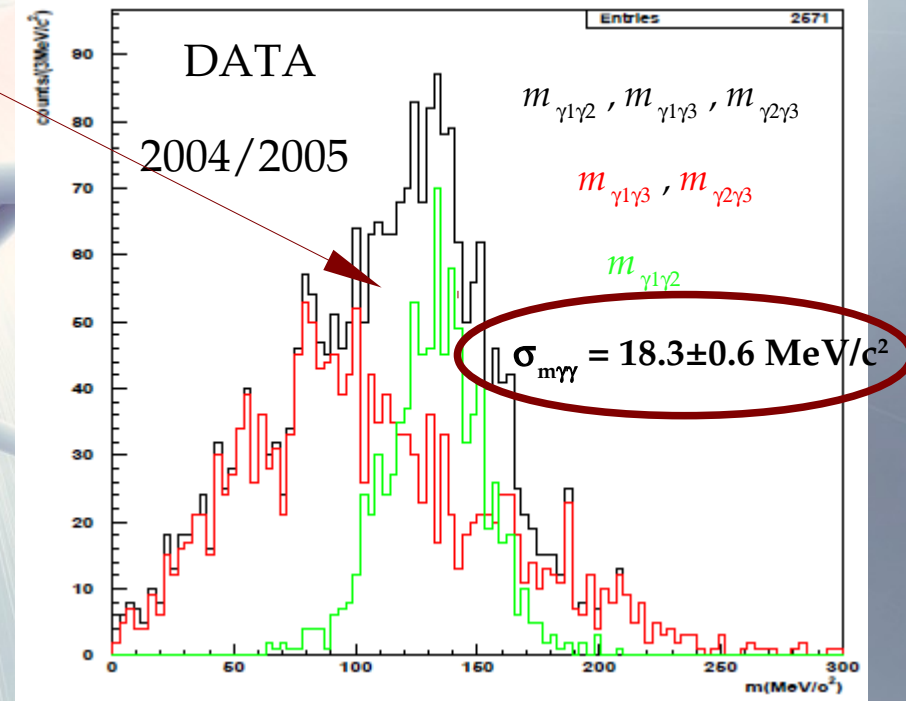
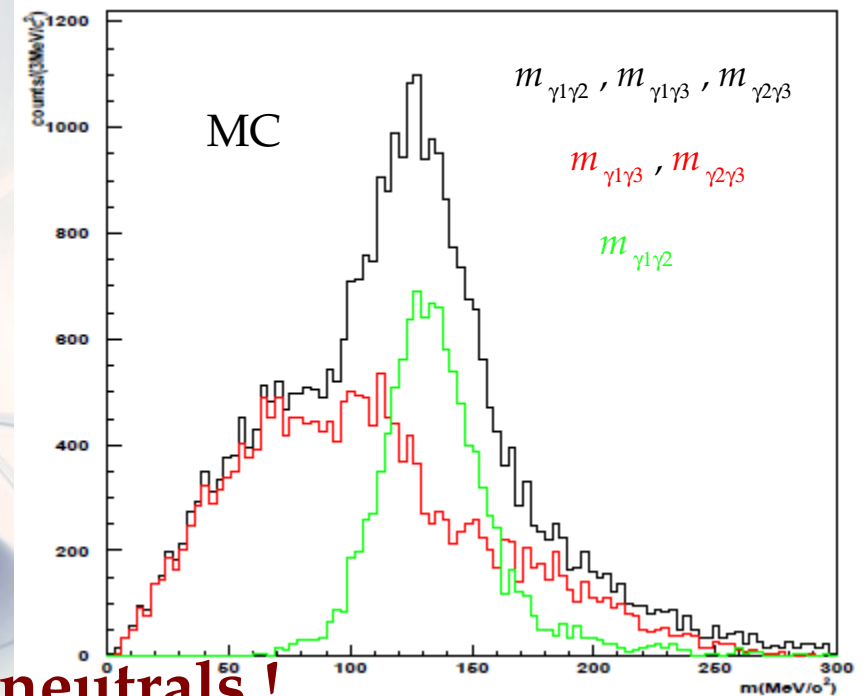
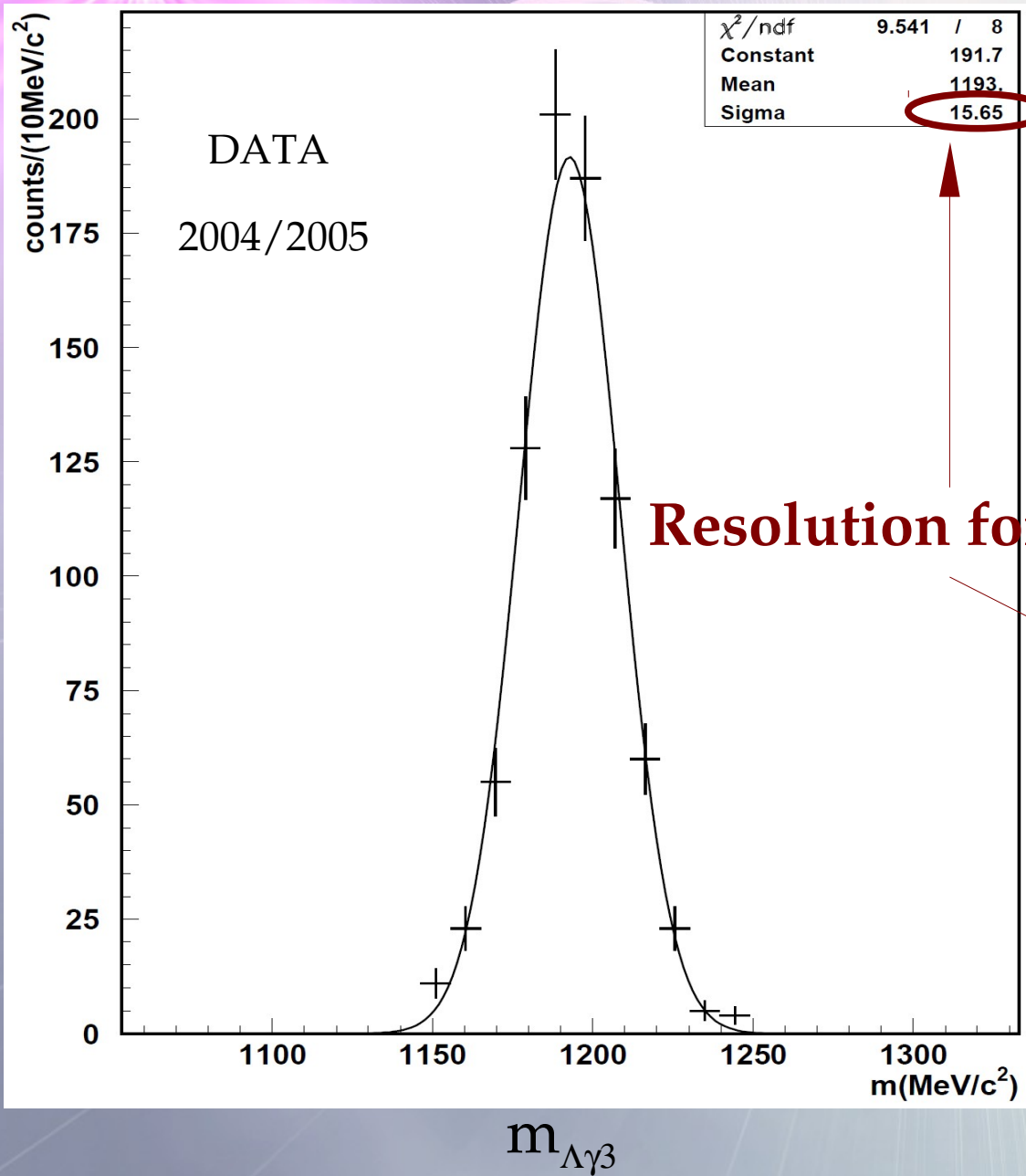
4) Cuts on  $\chi_i^2$  and  $\chi_{\pi\Sigma}^2$  optimized on MC simulations & splitted clusters rejection



**Efficiency (98±1)% to identify photons and (78±2)% to select the correct triple of neutral clusters.**

# Resolution for neutral clusters

$K^-$



Concluding, a fit of  $\Sigma^0\pi^0$  spectrum requires ...

$K^-$

9 components:

- Resonant component  $K^- C$  at-rest/in-flight.
- Non resonant  $\Sigma^0\pi^0 K^- H$  production at-rest/in-flight
- Non resonant  $\Sigma^0\pi^0 K^- C$  production at-rest/in-flight
- $\Lambda\pi^0$  background ( $\Sigma(1385) + I.C.$ )
- non resonant misidentification (*n.r.m.*) background

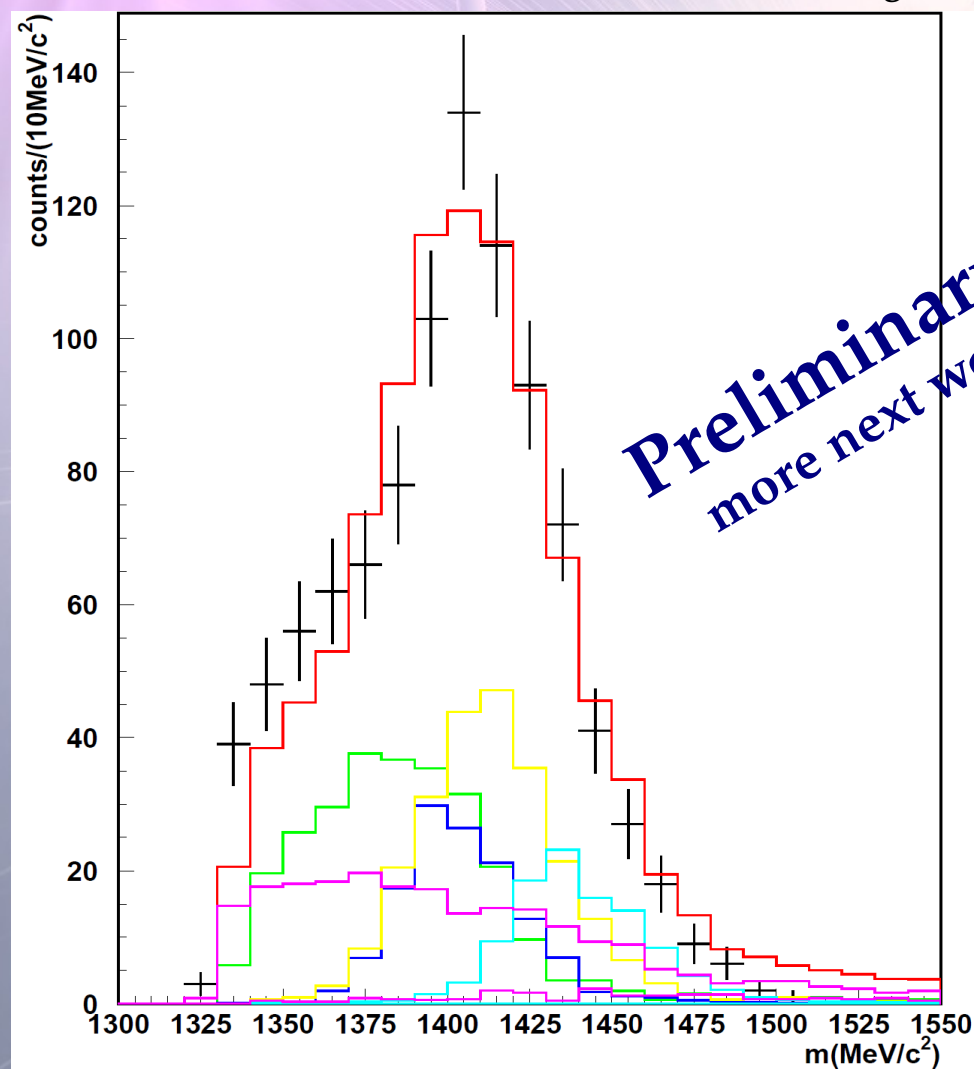
a careful model of  $K^-$  - nuclear interaction is required  
in Helium and Carbon!

# Fit of $\Sigma^0\pi^0$ spectrum in C

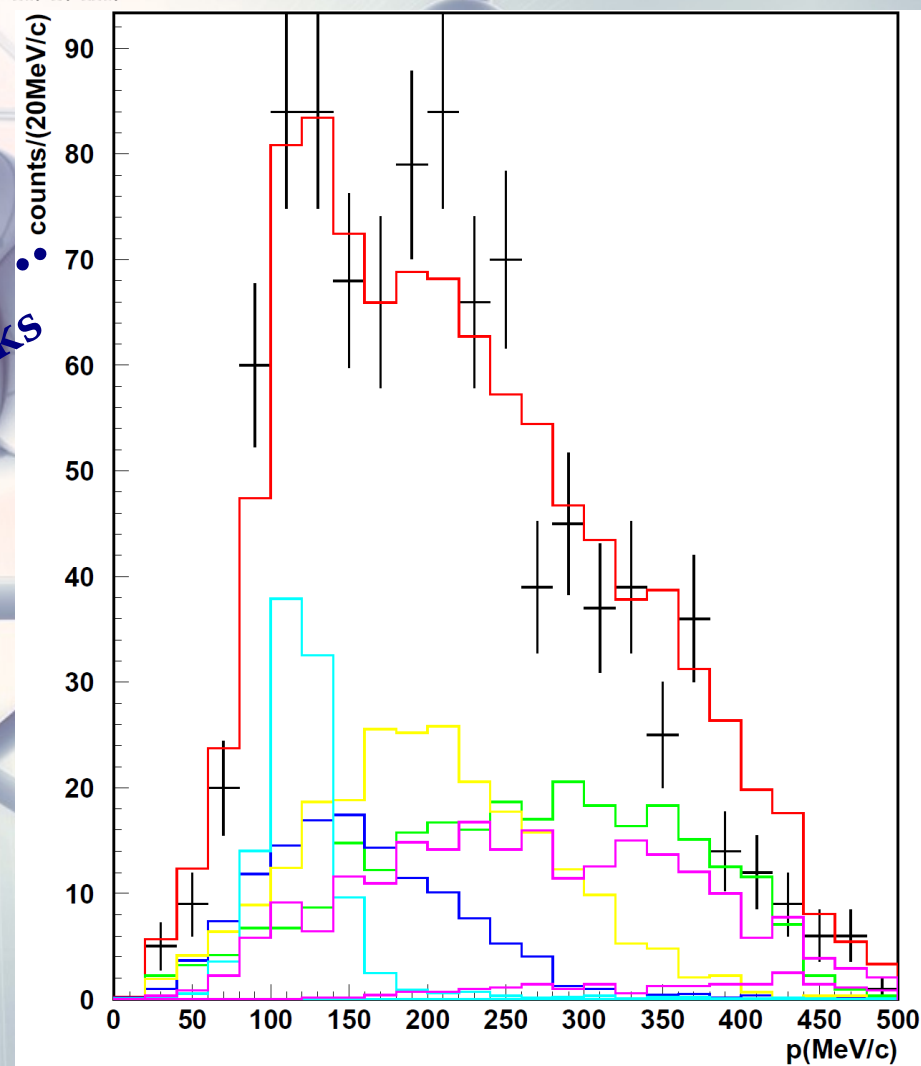
$\chi^2_{\min} / \text{ndf} \sim 1.7$  corresponding to  $(M_{\min}, \Gamma_{\min}) = (1426, 52) \text{ MeV}/c^2$

$K^-$

- Global fit ———
- Resonant component  $K^- C$  at-rest ———
- n. r.  $K^- C$  at-rest ———
- n. r.  $K^- C$  in-flight ———
- n. r.  $K^- H$  in-flight ———
- $\Lambda^0\pi^0$  background + n. r. m. ———



$m_{\Sigma^0\pi^0}$

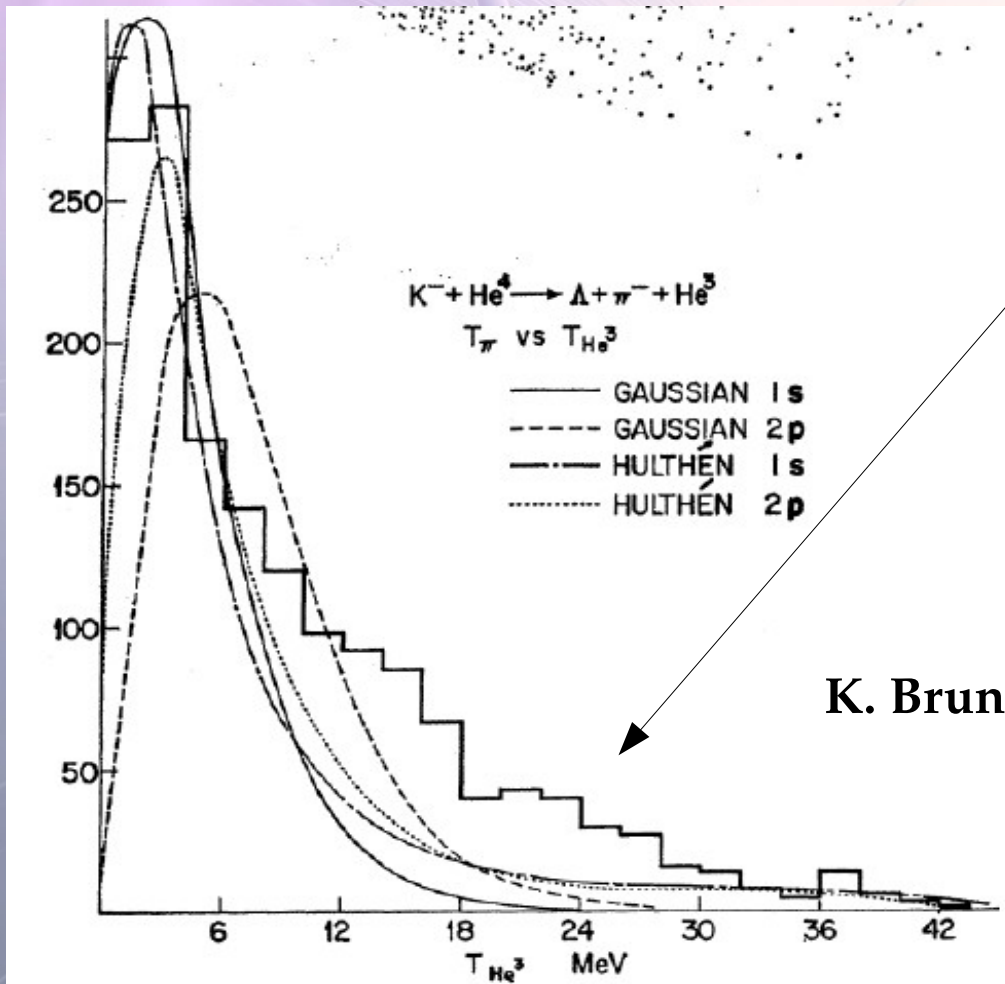


$p_{\Sigma^0\pi^0}$

# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ ... the idea

Bubble chamber experiments exhibit two components:

- Low momentum  $\Lambda \ \pi^-$  pair  $\rightarrow$  S-wave,  $I=1$ , non-resonant transition amplitude.
- High momentum  $\Lambda \ \pi^-$  pair  $\rightarrow$  P-wave resonant formation ?



Also exists in S-state K-mesic atom  
as a result of the  
three body structure of the system

( $K = 1, n=2, \ ^3\text{He} = 3$ )

K. Brunnel et al., Phys.Rev. D2 (1970) 98

# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ ... the strategy

$K^-$

- **Fit of the  $p_{\Lambda\pi^-}$  observed distribution** using calculated distributions :

$$P_s^s(p_{\Lambda\pi}) = |\Psi_N(p_{\Lambda\pi})|^2 |f^s(p_{\Lambda\pi})|^2 \rho$$

where  $\rho = k p_{\Lambda\pi}^2$

$$P_s^p(p_{\Lambda\pi}) = |\Psi_N(p_{\Lambda\pi})|^2 c^2 |2f^{\Sigma^*}(p_{\Lambda\pi})|^2 \rho/3 (k p_{\Lambda\pi})^2$$

The constant  $c = M_K/(M_K + M_n) = 0.345$  re-couples the S x S waves to P x P waves

- **To determine the ratio resonant/non-res.**

$|f_{\Lambda\pi}^{N-R}|$  given the fairly well known  $|f_{\Lambda\pi}^{\Sigma^*}|$



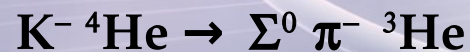
# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ ... calculated reactions

Calculated primary hadronic interactions:



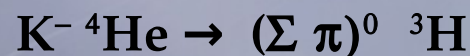
At-rest: S-wave non-Res / P-wave  $\Sigma(1385)$  Res

In-flight: S-wave non-Res / P-wave  $\Sigma(1385)$  Res



At-rest: S-wave non-Res / P-wave  $\Sigma(1385)$  Res

In-flight: S-wave non-Res / P-wave  $\Sigma(1385)$  Res



At-rest: S-wave non-Res / S-wave  $\Lambda(1405)$  Res /  
P-wave  $\Sigma(1385)$  Res

In-flight: S-wave non-Res / S-wave  $\Lambda(1405)$  Res /  
P-wave  $\Sigma(1385)$  Res

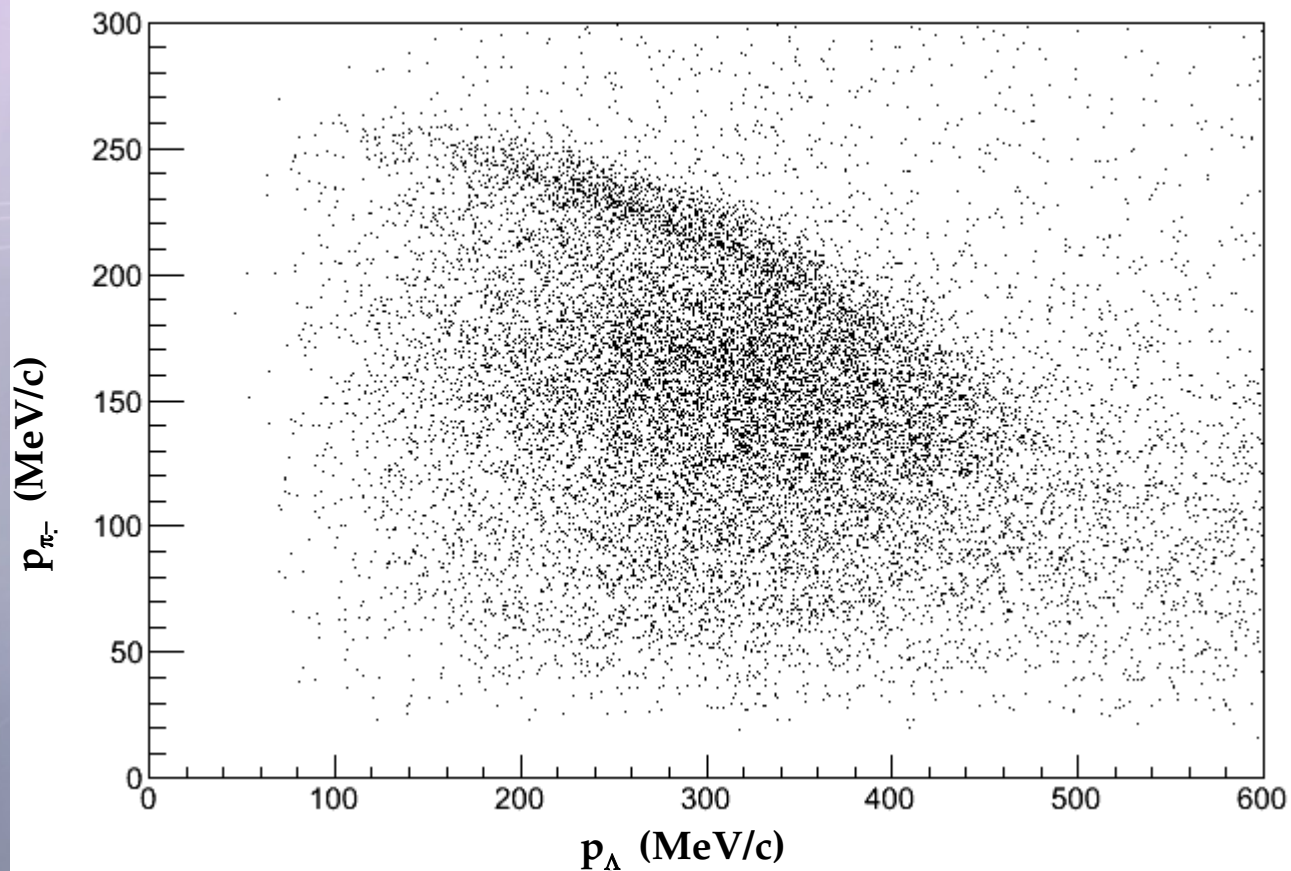
Channel:  $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$  ... calculated reactions

Calculated secondary hadronic interactions:

EACH INTERNAL CONVERSION PROCESS:

$$\Sigma \text{ p/n} \rightarrow \Lambda \text{ p/n}$$

was calculated for both P-wave and S-wave produced  $\Sigma$ s.



# Channel: $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ ... calculated reactions

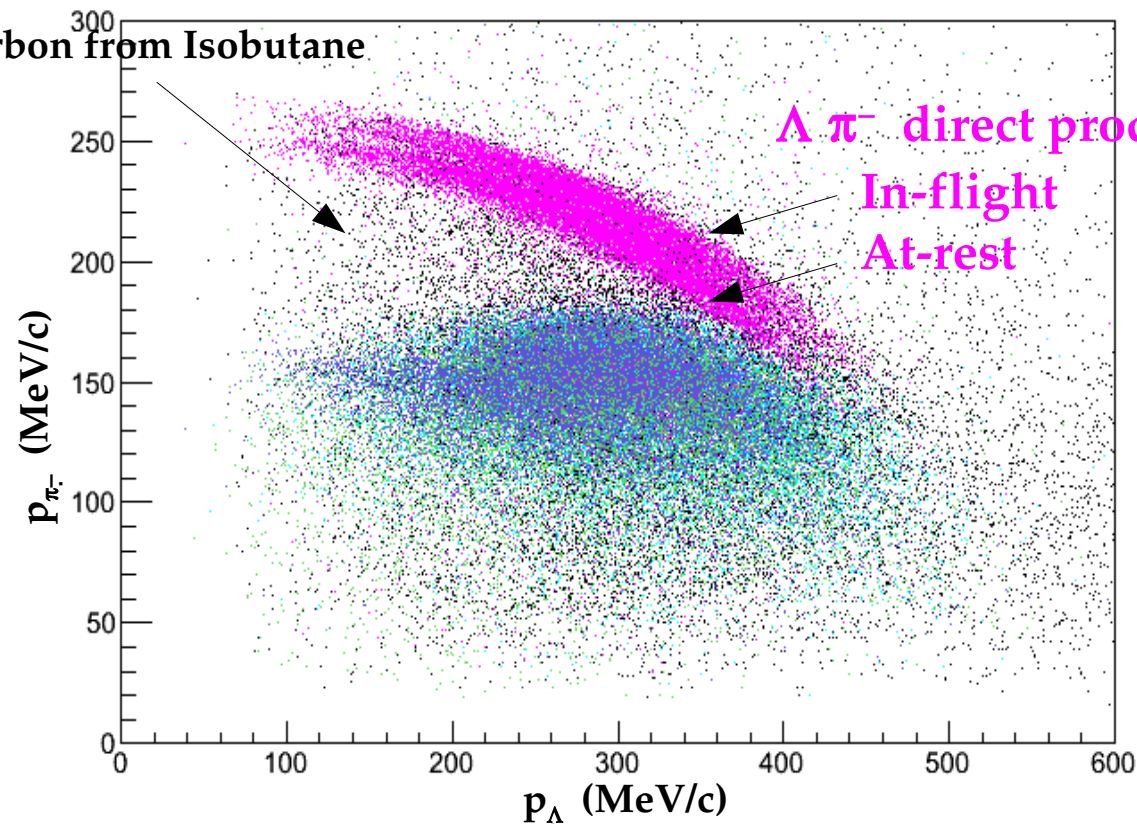
Calculated secondary hadronic interactions:

EACH INTERNAL CONVERSION PROCESS:

$$\Sigma \text{ p/n} \rightarrow \Lambda \text{ p/n}$$

was calculated for both P-wave and S-wave produced  $\Sigma$ s.

Some Carbon from Isobutane



$\Lambda \pi^-$  direct production

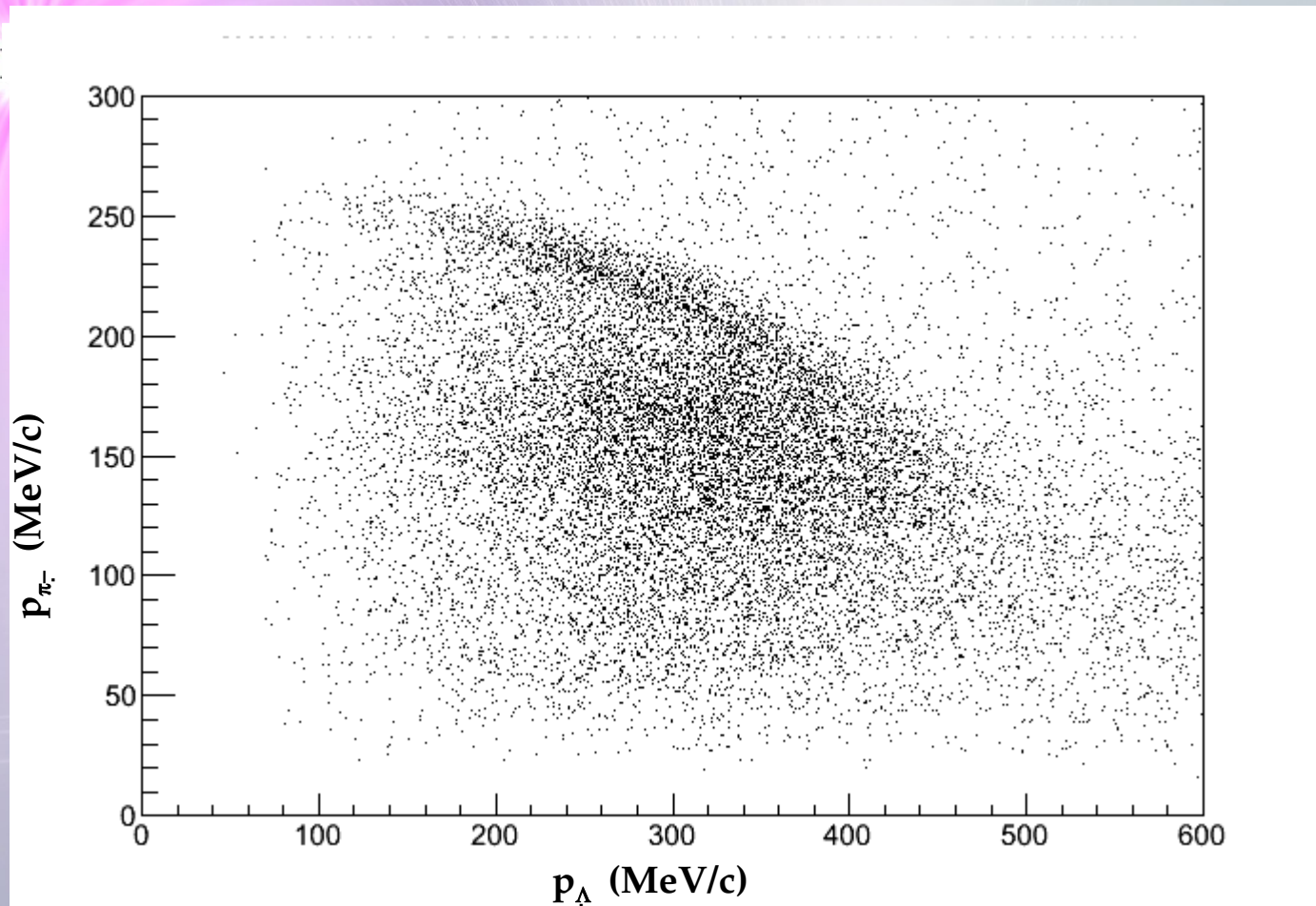
In-flight  
At-rest

$\Sigma^0$  p conversion

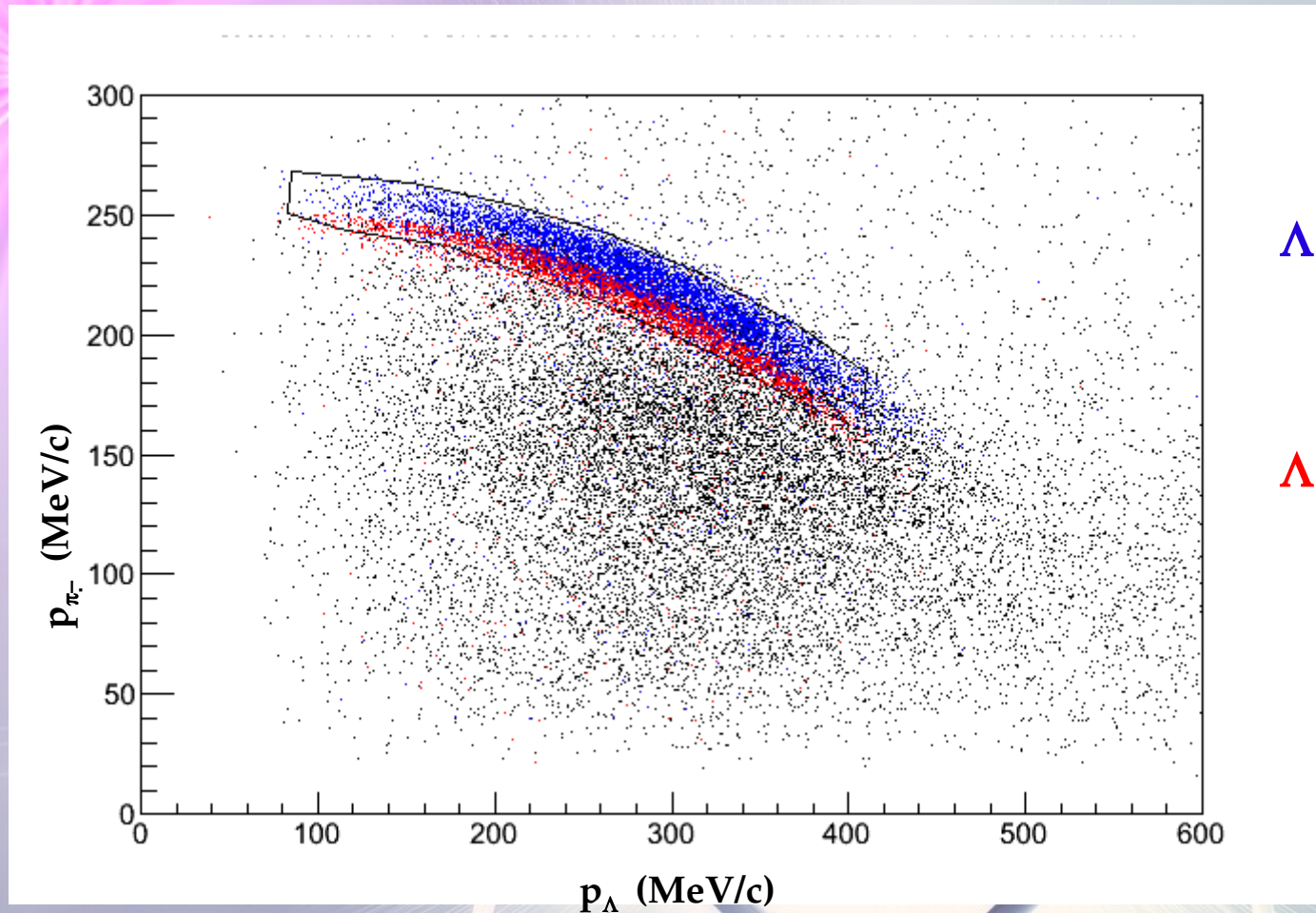
$\Sigma^0$  n conversion

$\Sigma^+$  n conversion

# $K^- \ ^4\text{He} \rightarrow \Lambda \ \pi^- \ ^3\text{He}$ events selection



# $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ events selection



$\Lambda \pi^-$  direct production  
In-flight RES + N-R

$\Lambda \pi^-$  direct production  
At-rest RES + N-R

- **CUT** based on MC simulations used to select  $\Lambda \pi^-$  direct production events
- At-rest **CAN NOT** be separated from In-flight  $\rightarrow$  global fit performed
- Background sources:
  - $\Lambda \pi^-$  events from  $\Sigma$  p/n  $\rightarrow$   $\Lambda$  p/n conversion
  - $\Lambda \pi^-$  events from  $K^- \ ^{12}\text{C}$  absorptions in Isobutane

# $K^- \ ^4\text{He} \rightarrow \Lambda \pi^- \ ^3\text{He}$ background

- $\Sigma$  p/n  $\rightarrow$   $\Lambda$  p/n conversion:

Each possible conversion channel was simulated

$\Sigma^0$  p /  $\Sigma^0$  n /  $\Sigma^+$  n / At-rest / In-flight / from RES and N-R produced  $\Sigma$ s

- $\Lambda \pi^-$  events from  $K^- \ ^{12}\text{C}$  absorptions in Isobutane (90% He, 10%  $\text{C}_4\text{H}_{10}$ ):

$K^- \ ^{12}\text{C}$  DATA in the KLOE DC wall are used

estimated contribution:

$$N_{\text{KHe}}/N_{\text{KC}} = (n_{\text{KHe}}/n_{\text{KC}}) \cdot (\sigma_{\text{KHe}}/\sigma_{\text{KC}}) \cdot (\text{BR}_{\text{KHe}}(\Lambda \pi^-)/\text{BR}_{\text{KC}}(\Lambda \pi^-)) \sim 1.3 \pm 0.3$$

Nuovo Cimento 39 A 338-347 (1977)

$K^- \ ^{12}\text{C}$  still not calculated:

- uncertain initial state of K meson  $l_K = 1, 2, 3$
- 4 nucleons in s-orbit, 8 nucleons in p-orbit
- final state hyperon interactions

# Study of the background

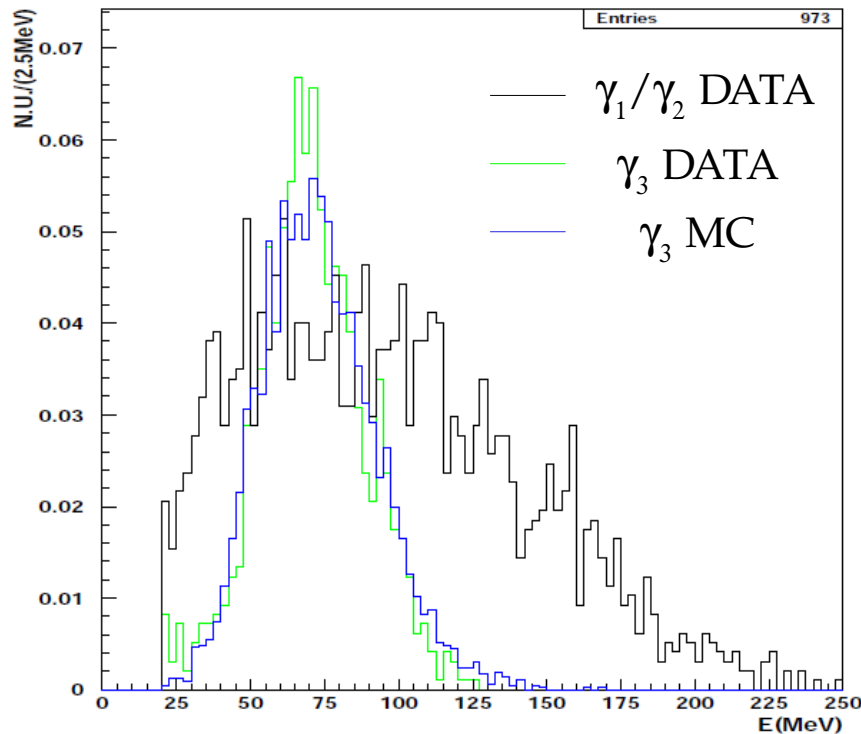
The main background sources for this channel are (example in  $^{12}\text{C}$ ):



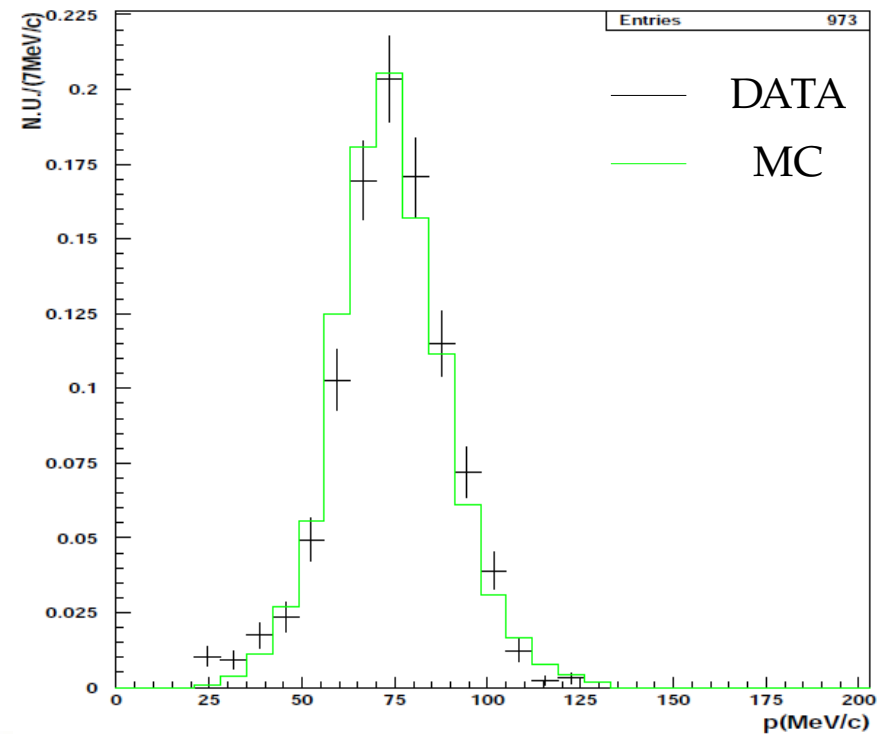
$\Sigma^0(1385)$  can not decay in  $\Sigma^0 \pi^0$  for isospin conservation.

- Internal conversion**  $\text{K}^- \text{ } ^{12}\text{C} \rightarrow \Lambda(1405) + \text{}^{11}\text{B} \rightarrow \Sigma^0\pi^0 + \text{}^{11}\text{B}$  ,  $\Sigma^0 \text{N} \rightarrow \Lambda \text{N}$  competes with the decay  $\Sigma^0 \rightarrow \Lambda \gamma$ .

Both background sources were analyzed by different methods:



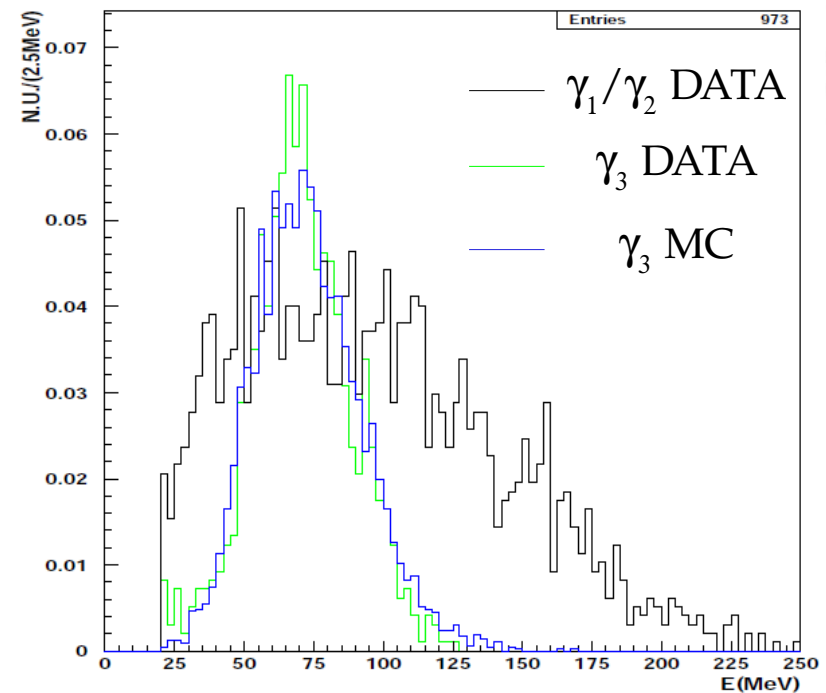
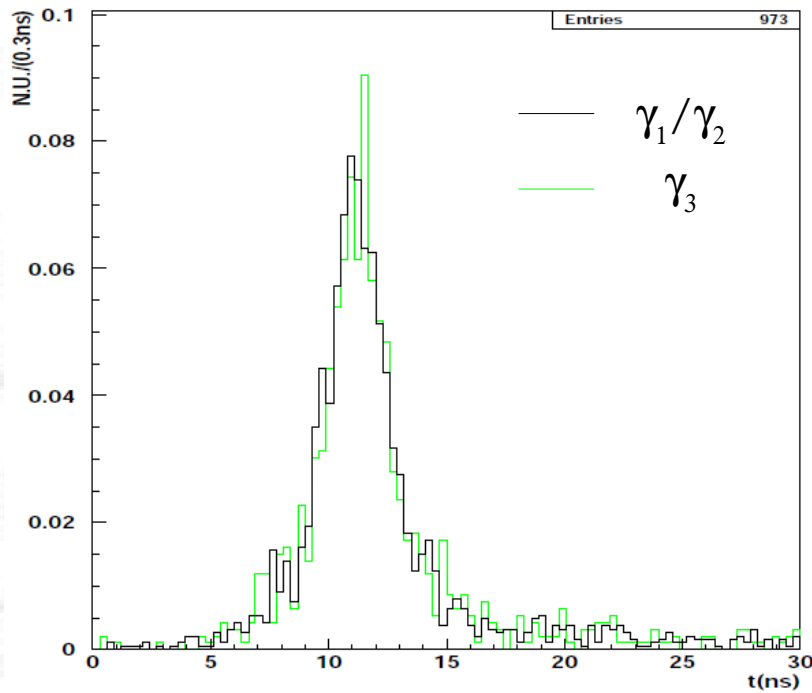
photons energy distribution



$\Lambda$  momentum in the  $\Sigma^0$  rest frame

# Study of the background

In both cases  $\gamma_3$  is not present, if a contamination is present, the neutral cluster which is associated to  $\gamma_3$  by reconstruction should show differences.



- Right: the energy distribution of  $\gamma_3$  (green) is in perfect agreement with MC simulations of pure signal events (blue) (energy spectrum of  $\gamma_1\gamma_2$  is shown in black).
- Left: the time distribution of  $\gamma_3$  (green) is in agreement with the time distributions of the two photons coming from  $\pi^0$  decay (black).



# Study of the background

The numbers of pure background  $\Sigma(1385)$  and  $\Sigma^0 N \rightarrow \Lambda N$  events passing the analysis cuts are normalized to pure signal  $\Lambda(1405)$  events, then weighted to the BRs for  $\Lambda\pi^0$  direct production (D), internal conversion (IC) and  $\Sigma^0\pi^0$  production due to  $K^-$  interaction in  $^4\text{He}$  and C respectively :

P. A. Katz et al., Phys.Rev. D1 (1970) 1267

C. Vander Velde-Wilquet et al., Nuovo Cimento 39 A, (1977) 538

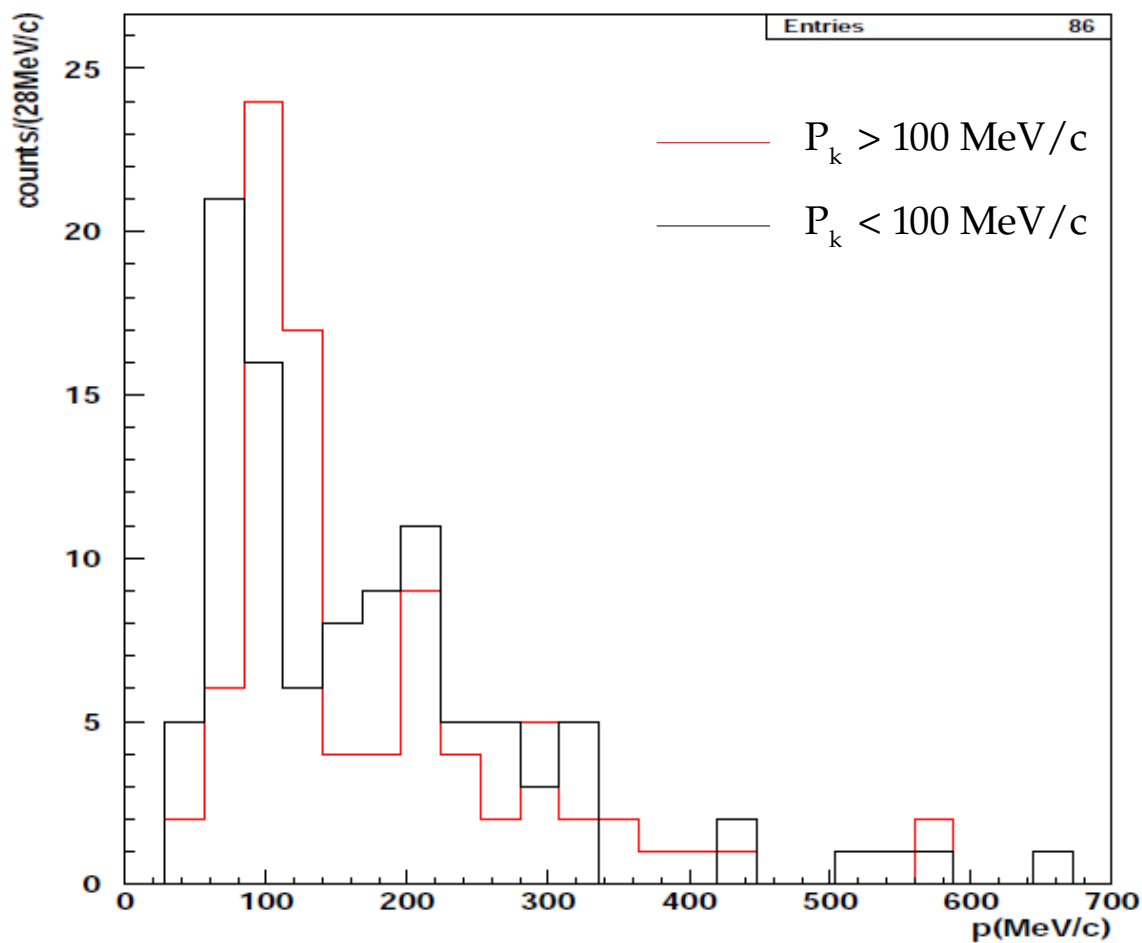
The percentages of background events entering the final selected samples are:

$$\frac{n_{\Lambda\pi^0 D \text{ norm}} + n_{\Lambda\pi^0 IC \text{ norm}}}{n_{\Sigma^0\pi^0} + n_{\Lambda\pi^0 D \text{ norm}} + n_{\Lambda\pi^0 IC \text{ norm}}} = 0.03 \pm 0.01 \quad \text{in DC wall} \quad (0.03 \pm 0.02 \text{ in DC gas})$$

# $p_{\pi^0\Sigma^0}$ spectrum for boost and anti-boost events

37

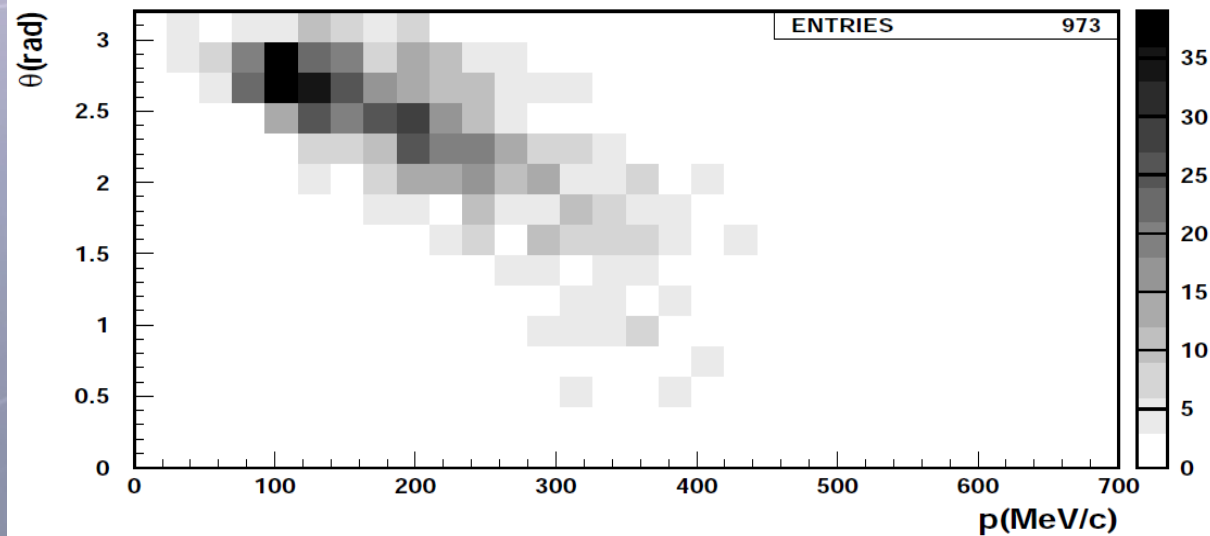
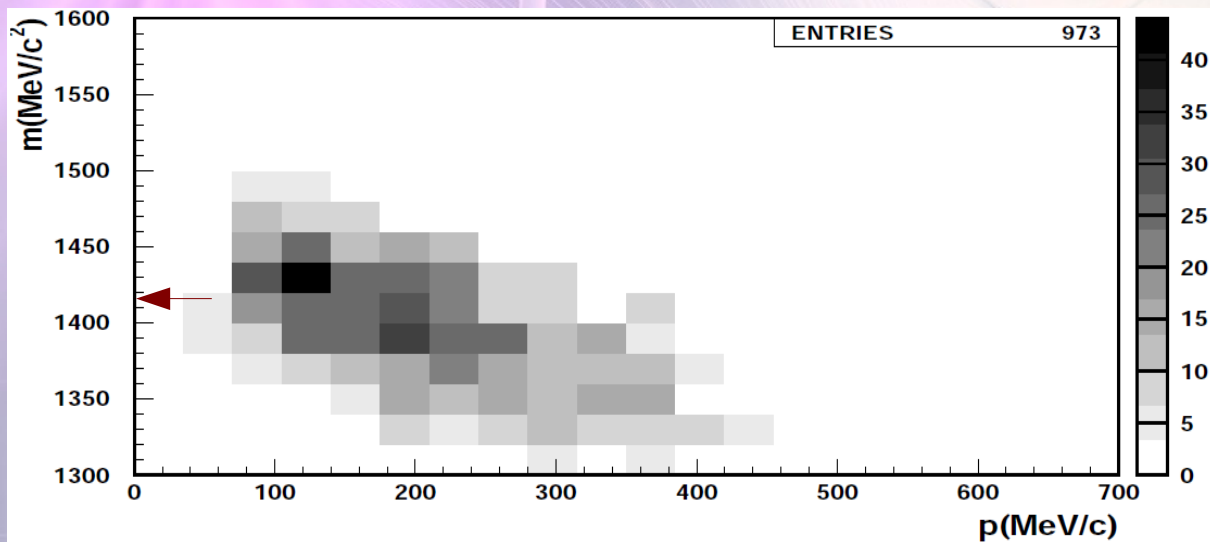
$p_{\Sigma^0\pi^0}$  distribution for lower (black) and higher (red)  $p_k$  values



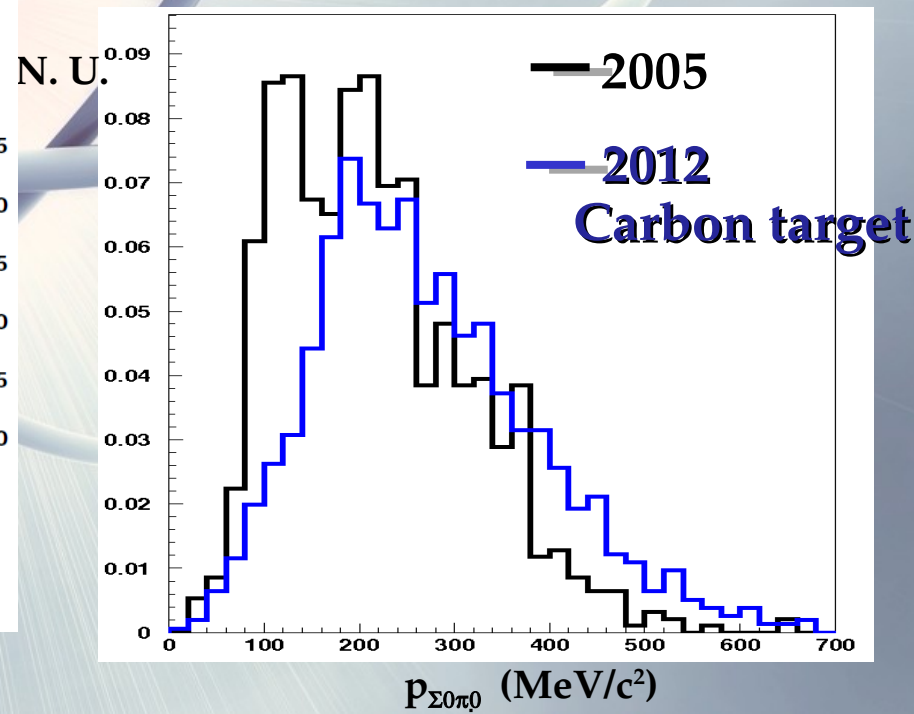
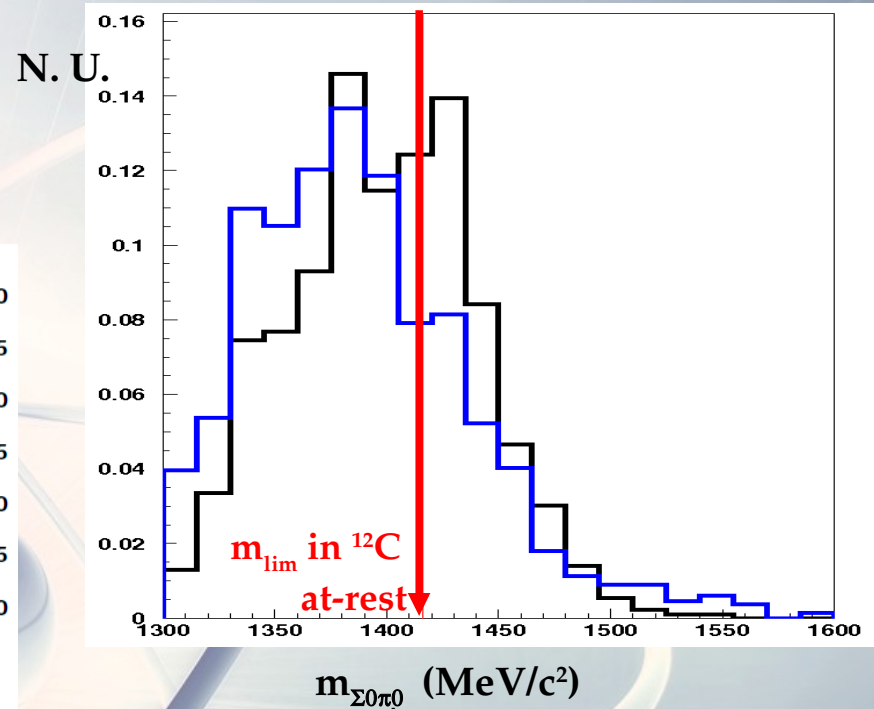
# $\Sigma^0 \pi^0$ channel

$K^-$

## Mass momentum correlation



Top  $m_{\Sigma^0 \pi^0}$  vs  $p_{\Sigma^0 \pi^0}$ , bottom  $\theta_{\Sigma^0 \pi^0}$  vs  $p_{\Sigma^0 \pi^0}$ .



# Fit of $\Sigma^0\pi^0$ spectrum in C

$K^-$

8 component fit, simultaneously  $m_{\Sigma^0\pi^0}$  &  $p_{\Sigma^0\pi^0}$ :

- Breit-Wigner resonant component  $K^- C$  at-rest/in-flight.  $(M,\Gamma) = (1405 \div 1430, 5 \div 52)$ 
  - Non resonant  $\Sigma^0\pi^0$   $K^- H$  production at-rest/in-flight
  - Non resonant  $\Sigma^0\pi^0$   $K^- C$  production at-rest/in-flight
    - $\Lambda\pi^0$  background ( $\Sigma(1385) + I.C.$ )
  - non resonant misidentification (*n.r.m.*) background

$K^- {}^{12}C \rightarrow \Sigma^0\pi^0 + {}^{11}B$  (Boron spectator, left in ground state)

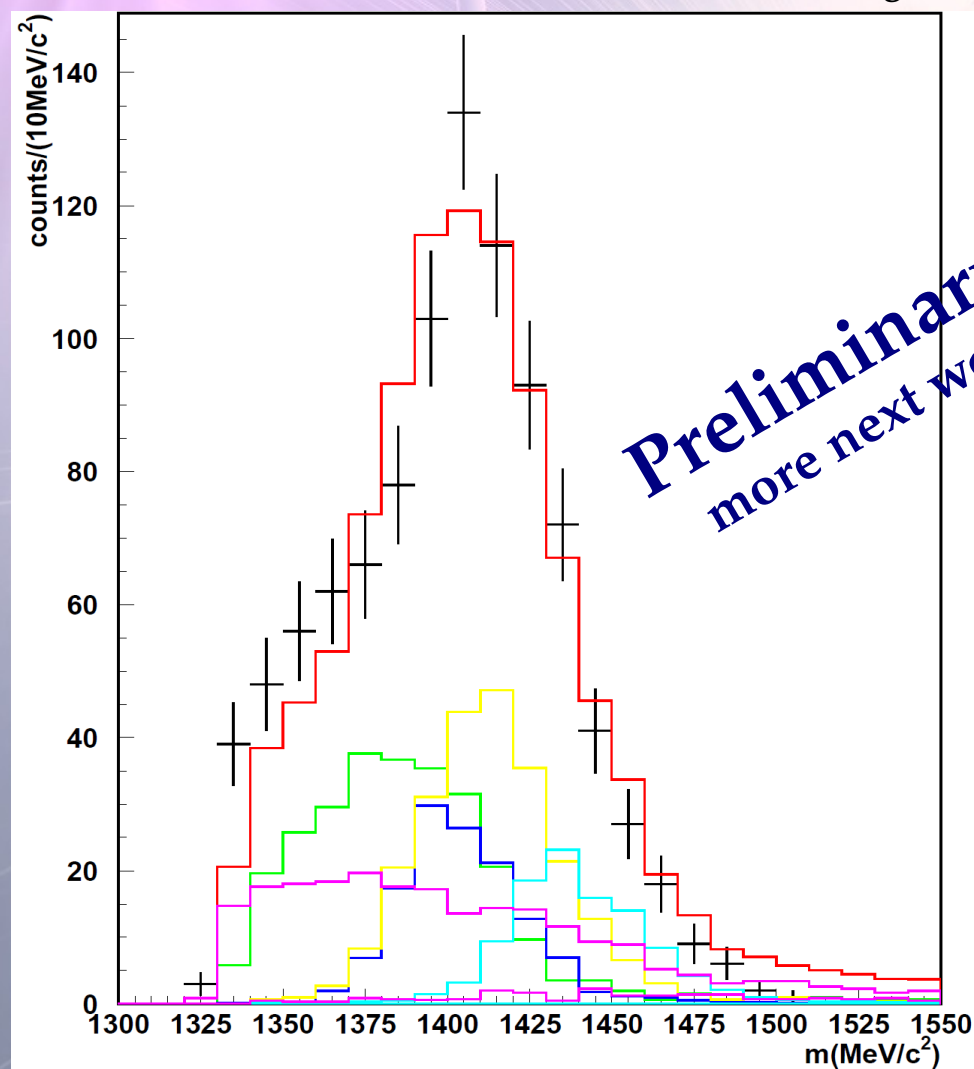
secondary interactions not taken into account.

# Fit of $\Sigma^0\pi^0$ spectrum in C

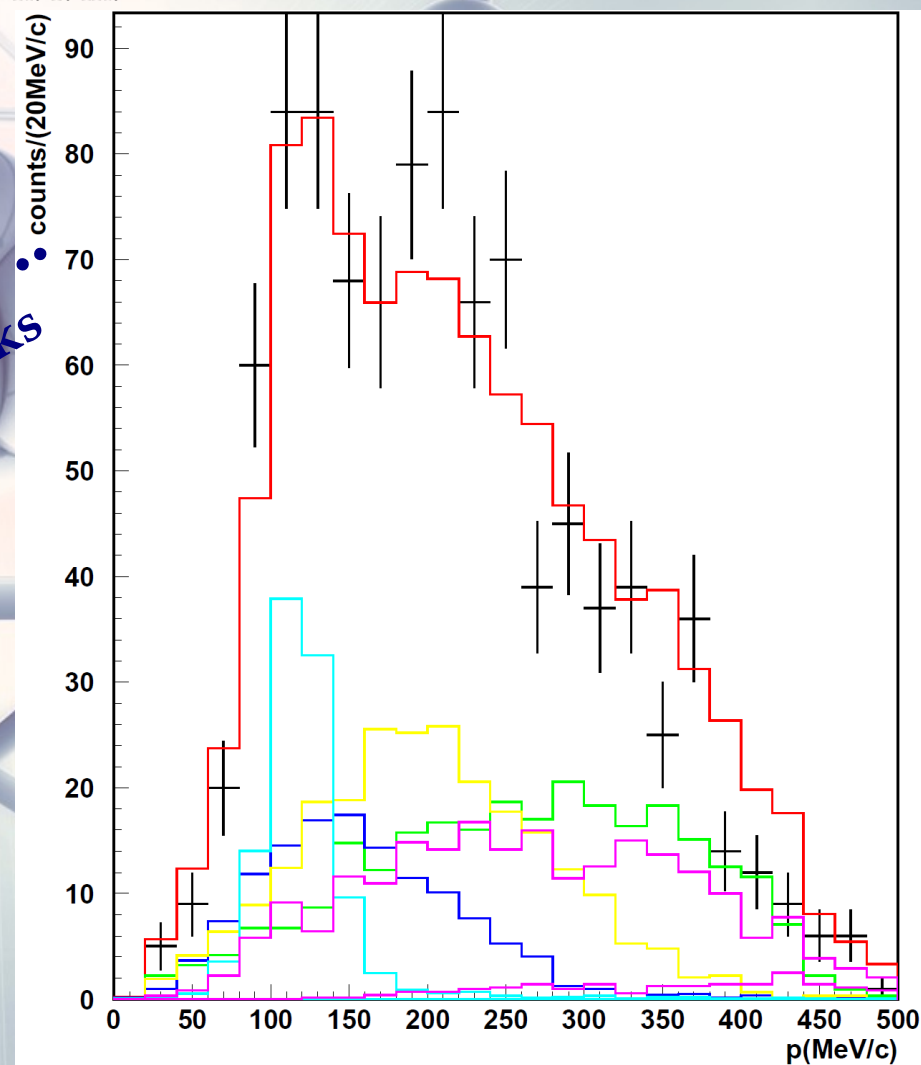
$\chi^2_{\min} / \text{ndf} \sim 1.7$  corresponding to  $(M_{\min}, \Gamma_{\min}) = (1426, 52) \text{ MeV}/c^2$

$K^-$

- Global fit — (red line)
- Resonant component  $K^- C$  at-rest — (green line)
- n. r.  $K^- C$  at-rest — (blue line)
- n. r.  $K^- C$  in-flight — (yellow line)
- n. r.  $K^- H$  in-flight — (cyan line)
- $\Lambda^0\pi^0$  background + n. r. m. — (magenta line)

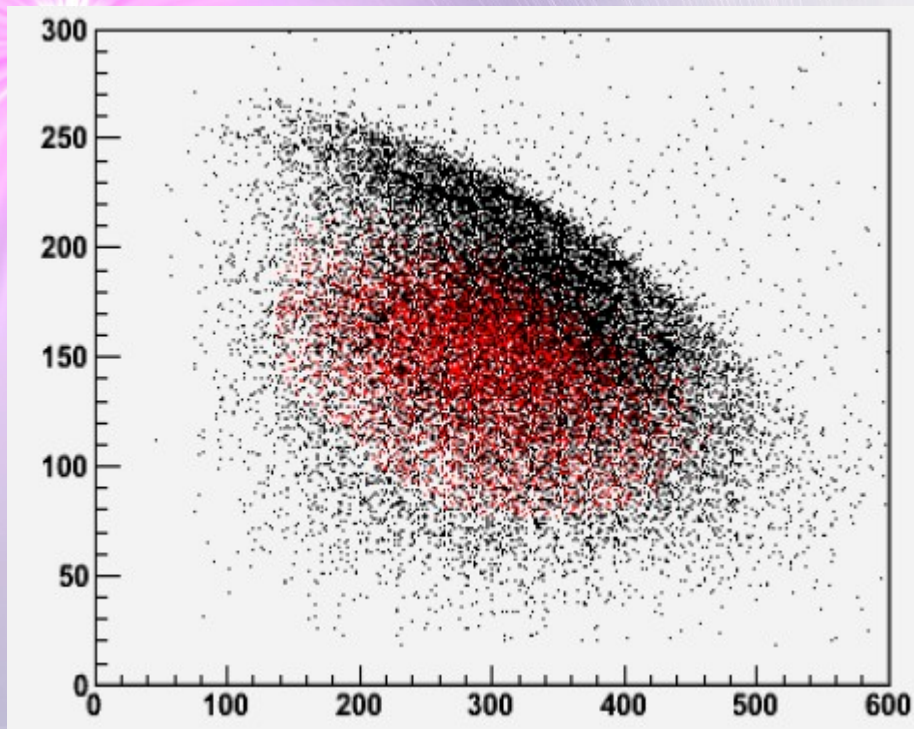


$m_{\Sigma^0\pi^0}$



$p_{\Sigma^0\pi^0}$

# $\Lambda\pi^-$ + extra-proton



Black-> lambda + pi-  
Red-> lambda + pi- + proton

The **extra-p** indicates  
**fragmentation of the residual  
nucleus**

( $\Sigma / \Lambda$  conversion,  $\Sigma / \Lambda / \pi$   
secondary interactions,  
multi-nucleon absorption)  
but ...

...

- if ( $\Lambda\pi^-$  direct production) then 3% extra-p
- if ( $\Lambda\pi^-$  in-direct production) then 25% extra-p