

Odd moments of nucleon charge and magnetization distribution in baryon chiral perturbation theory

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Outline

I) Motivation:

a) Muonic hydrogen Lamb shift(LS), proton structure correctionsb) Hyperfine splitting

2) Moments of nucleon charge distribution in BChPT

3) Summary





A. Antognini et al., Science **339**, 417 (2013).

F=0







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 $[R_E^{\mu \rm H} = 0.84087(39)\,\rm{fm}]$



 $\Delta E_L^{\text{th}} \,[\text{meV}] = 206.0336(15) - 5.2275(10) \,R_E^2 \,[\text{fm}^2] + \Delta E_{\text{TPE}}$

Leading-order Finitesize correction, a⁴

$$\Delta E_{nS}(\text{LO}) = \frac{2(Z\alpha)^4 m_r^3}{3n^3} R_E^2$$

$$R_E^2 = -6 \lim_{Q^2 \to 0} \frac{\mathrm{d}}{\mathrm{d}Q^2} G_E(Q^2)$$

 $\Delta E_L^{\text{th}} [\text{meV}] = 206.0336(15) - 5.2275(10) R_E^2 [\text{fm}^2] + \Delta E_{\text{TPE}}$ Leading-order Finitesize correction, a^4 $\Delta E_{nS}(\text{LO}) = \frac{2(Z\alpha)^4 m_r^3}{3n^3} R_E^2$ $\Delta E_{nS}(\text{NLO}) = \frac{2(Z\alpha)^4 m_r^3}{3n^3} R_E^2$ $B_{E}^2 = -6 \lim_{Q^2 \to 0} \frac{d}{dQ^2} G_E(Q^2)$ $R_E^3 = -6 \lim_{Q^2 \to 0} \frac{d}{dQ^2} G_E(Q^2)$ $R_E^3 = \frac{48}{\pi} \int_0^\infty \frac{dQ}{Q^4} \left\{ G_E^2(Q^2) - 1 + \frac{1}{3} R_E^2 Q^2 \right\}$

 $\Delta E_L^{\text{th}} [\text{meV}] = 206.0336(15) - 5.2275(10) R_E^2 [\text{fm}^2] + \Delta E_{\text{TPE}}$ Two-photon exchange Leading-order Finitesize correction, a⁴ correction, a⁵ elastic: $\Delta E_{nS}(\text{NLO}) = -\frac{(Z\alpha)^3 m_r^4}{3n^3} R_{E(2)}^3$ $\Delta E_{nS}(\mathrm{LO}) = \frac{2(Z\alpha)^4 m_r^3}{2m^3} R_E^2$ 3rd Zemach moment: $R_{E(2)}^{3} = \frac{48}{\pi} \int_{0}^{\infty} \frac{\mathrm{d}Q}{Q^{4}} \left\{ G_{E}^{2}(Q^{2}) - 1 + \frac{1}{3}R_{E}^{2}Q^{2} \right\}$ $R_E^2 = -6 \lim_{Q^2 \to 0} \frac{\mathrm{d}}{\mathrm{d}O^2} G_E(Q^2)$ polarizability:

> BChPT: Alarcon, Lensky & Pascalutsa, EPJC (2014)

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$\Delta E_{HFS}^{\text{th}} \text{ [meV]} = 22.9763(15) - 0.1621(10) \langle r \rangle_Z \text{[fm]} + \Delta E_{\text{HFS}}^{\text{pol}}$

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Fermi energy + radiative and recoil corrections



Zemach radius:

$$\langle r \rangle_Z = \int_0^\infty \frac{\mathrm{d}Q}{Q^2} \left[\frac{G_E(Q^2) \, G_M(Q^2)}{1+\kappa} - 1 \right]$$



Moments of charge distribution

Charge distribution:

$$\rho_E(r) = \int \frac{d\vec{q}}{(2\pi)^3} G_E(\vec{q}^2) e^{-i\vec{q}\vec{r}}$$

$$G_E(Q^2) = 1 - \frac{Q^2}{\pi} \int dt \, \frac{\text{Im}G_E(t)}{(t+Q^2)t}$$
where $Q^2 = -q^2$

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$$\langle r^N \rangle_E \equiv \int \mathrm{d}\vec{r} \, r^N \rho_E(r)$$
$$= \frac{(N+1)!}{\pi} \int_{t_0}^{\infty} \mathrm{d}t \, \frac{\mathrm{Im} \, G_E(t)}{t^{1+N/2}}$$

Predictive orders (no new LECs)

Leading order (p^3):



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At order O(p^4/Δ):



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$$\langle R^3 \rangle_E = \frac{48}{\pi} \int_0^\infty \frac{\mathrm{d}Q}{Q^4} \left\{ G_E(Q^2) - 1 + \frac{1}{6} \langle r^2 \rangle_E Q^2 \right\} = \frac{24}{\pi} \int_{t_0}^\infty \mathrm{d}t \, \frac{\mathrm{Im}G_E(t)}{t^{5/2}}$$

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LO heavy baryon ChPT:

$$\langle R^3 \rangle_E = \frac{21g_A^2}{128\pi f_\pi^2 m_\pi} + O\left(\frac{m_\pi}{M_p}\right)$$

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$$\langle R^3 \rangle_E \simeq 0.276 + 0.109 = 0.385 \,[\text{fm}^3]$$

 $1 \qquad 1 \qquad 1 \qquad 10$

Empirical value from FF fitted to ep data:

$$\langle R^3 \rangle_E = 1.16(4) \, [\text{fm}^3]$$

Distler, Bernauer, Walcher, Phys. Lett. B 696, 343 (2011).

Third Zemach moment

$$\langle R^3 \rangle_{E(2)} = 2 \langle R^3 \rangle_E + \frac{24}{\pi^2} \int_{t_0}^{\infty} dt \int_{t_0}^{\infty} dt' \frac{\text{Im}G_E(t) \,\text{Im}G_E(t')}{(t\,t')^{3/2}(\sqrt{t}+\sqrt{t'})}$$

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Numerical result:

$$\langle R^3 \rangle_{E(2)} \simeq 0.554 + 0.218 = 0.772 \,[{\rm fm}^3]$$

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m fm}^3]$$
 (Distler, NLO

Empirical value from FF fitted to ep data:

$$\langle R^3 \rangle_{E(2)} = 2.85(8) \, [\text{fm}^3]$$

Distler, Bernauer, Walcher, Phys. Lett. B **696**, 343 (2011).

Zemach radius

$$\langle r \rangle_{Z} = -\frac{2}{\pi(1+\kappa)} \Big\{ \frac{1}{\pi} \int_{t_{0}}^{\infty} \mathrm{d}t \int_{t_{0}}^{\infty} \mathrm{d}t' \frac{\mathrm{Im}G_{E}(t) \,\mathrm{Im}G_{M}(t')}{t \, t'(\sqrt{t} + \sqrt{t'})} \\ -(1+\kappa) \int_{t_{0}}^{\infty} \mathrm{d}t \, \frac{\mathrm{Im}G_{E}(t)}{t^{3/2}} - \int_{t_{0}}^{\infty} \mathrm{d}t \, \frac{\mathrm{Im}G_{M}(t)}{t^{3/2}} \Big\}$$

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Numerical result: $\langle r \rangle_Z \simeq 1.537 + 1.439 = 2.976 \, [\text{fm}]$ \hat{f} $\hat{\chi}$ LO NLO Empirical value from FF fitted to ep data: $\langle r \rangle_Z = 1.045(4) \, [\text{fm}]$

Distler, Bernauer, Walcher, Phys. Lett. B 696, 343 (2011).

Summary of the results

Odd moment	р^3	p^4 /∆	p^3+p^4/∆	Empirical value
$\langle R^3 \rangle_E$	0.276	0.109	0.385	1.16(4)
$\langle R^3 \rangle_{E(2)}$	0.554	0.218	0.772	2.85(8)
$\langle r \rangle_Z$	I.537	I.439	2.976	I.045(4)

Possible explanations:

i) ChTP does not work, new physics needs to be implemented

ii) Empirical values are wrong

iii) ...

Summary

(i) Calculation of the 3rd Zemach moment allows us to extract the proton radius independently on the elastic ep-scattering data.

(ii) The 3rd Zemach moment and Zemach radius can be predicted by ChPT (no new LECs).

(iii) We present the results of the leading and next-to-leading order contributions, O(p^3) and O(p^4/ Δ), to the third moment of the charge distribution, the 3rd Zemach moment, and Zemach radius.

(iv) More work is required to reproduce the empirical values.