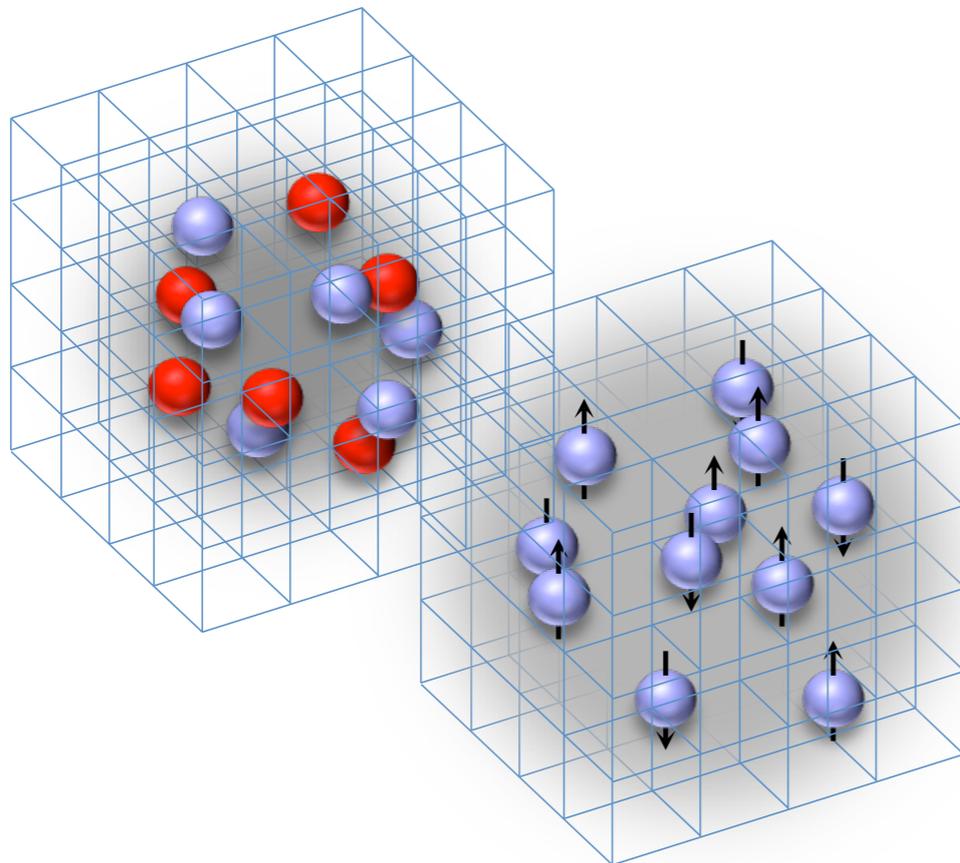
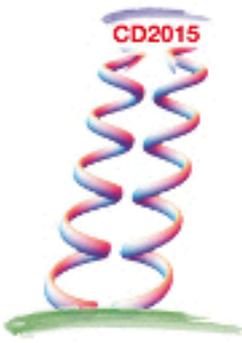


Scattering cluster wave functions on the lattice using the adiabatic projection method



Serdar Elhatisari - Bonn
Michelle Pine - NC State
Dean Lee - NC State
Evgeny Epelbaum - RUB
Hermann Krebs - RUB



Outline

- Introduction
- Adiabatic projection method
- Asymptotic cluster wave functions
- One and three dimensional examples
- Summary



- Objective: **ab initio** calculation of scattering and reactions involving two **clusters**. Processes with alpha-clusters are involved in stellar nucleosynthesis.



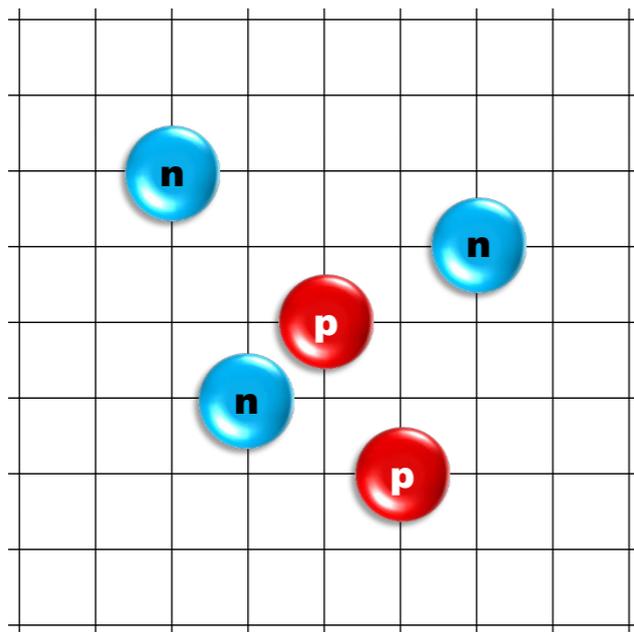
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 - Example: ${}^4\text{He} + {}^4\text{He} \rightarrow {}^4\text{He} + {}^4\text{He}$ (see next talk)



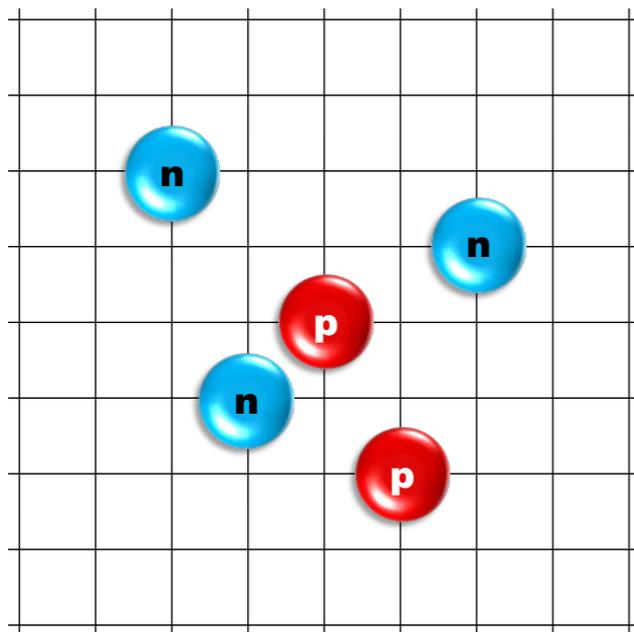
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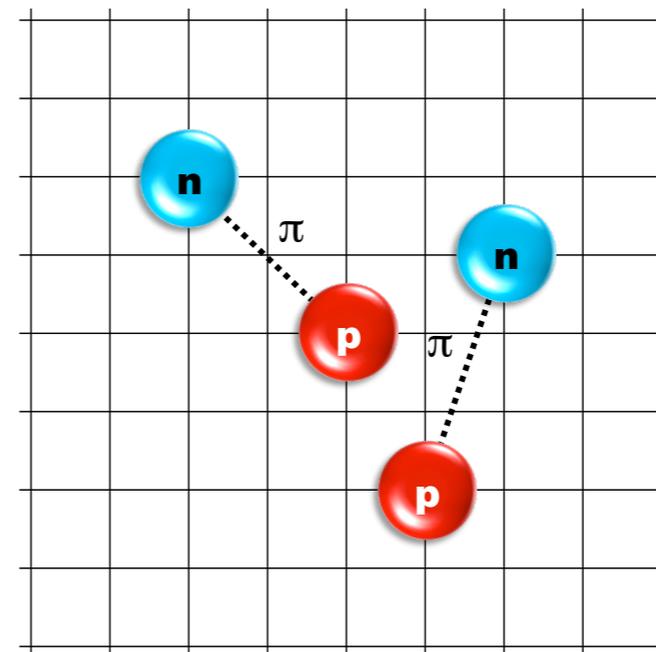
LEFT(\hbar)



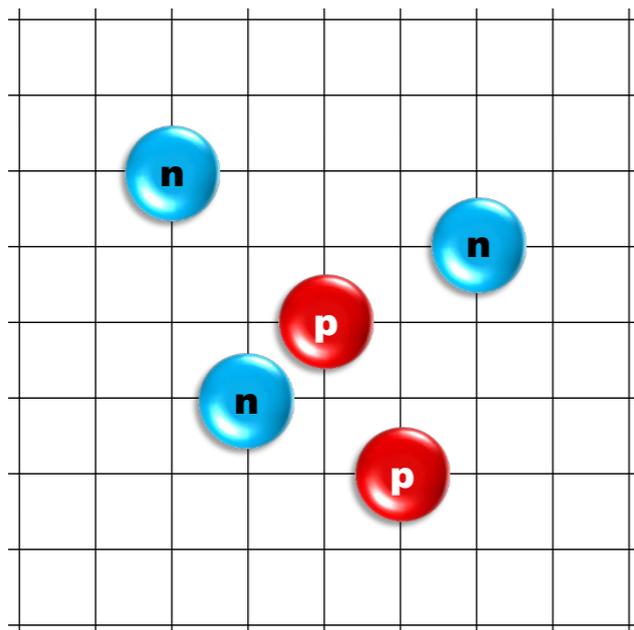
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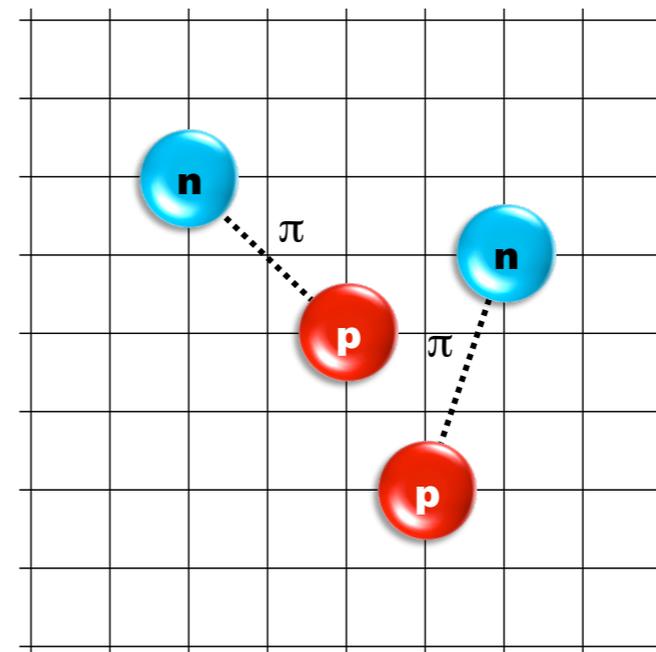
chirale LEFT



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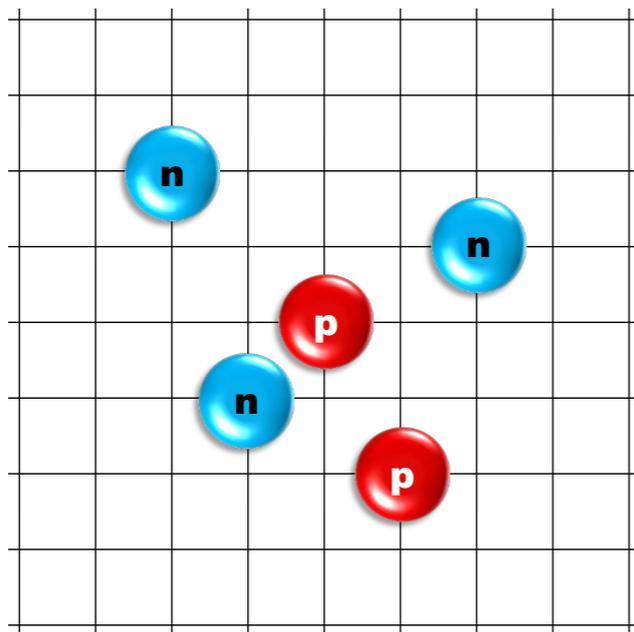
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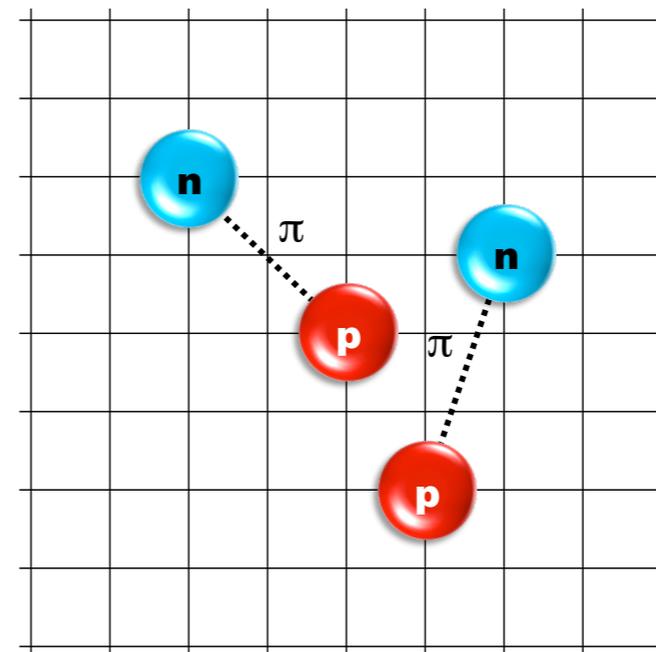


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M. Pine, D. Lee and G. Rupak: Eur. Phys. J. A (2013) 49: 151
 G. Rupak and D. Lee, Phys. Rev. Lett. 111, no. 3 (2013), 032502
 S. Elhatisari and D. Lee: Phys. Rev. C 90, no. 6 (2014), 064001

LEFT(π)

chirale LEFT

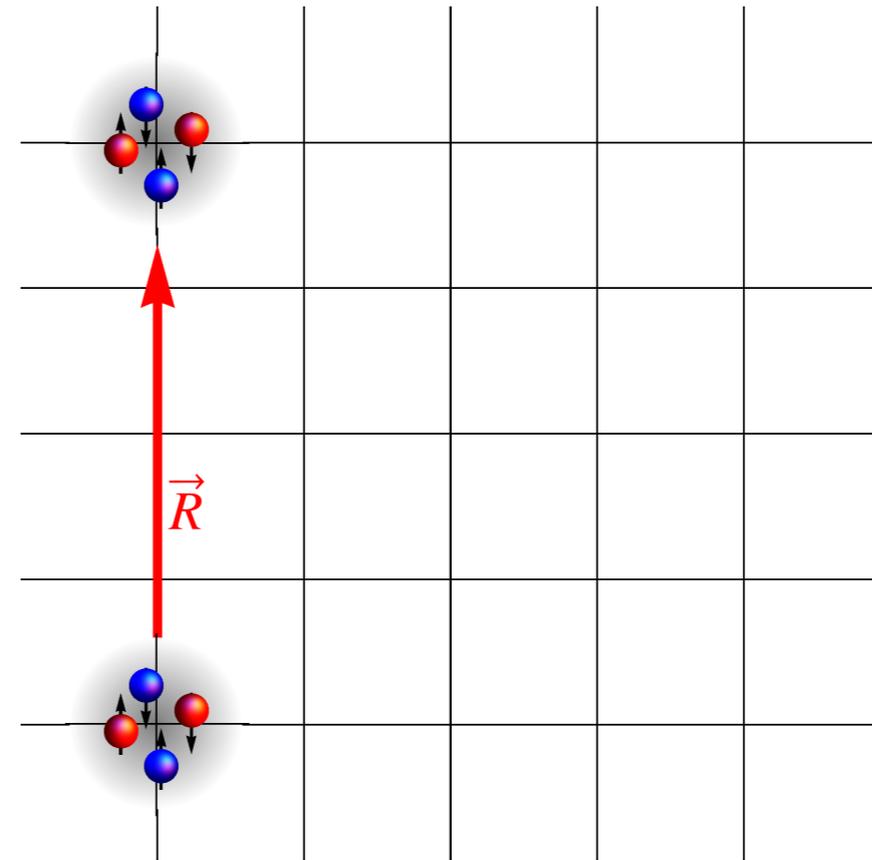




- **First Step:** low-energy cluster Hamiltonian

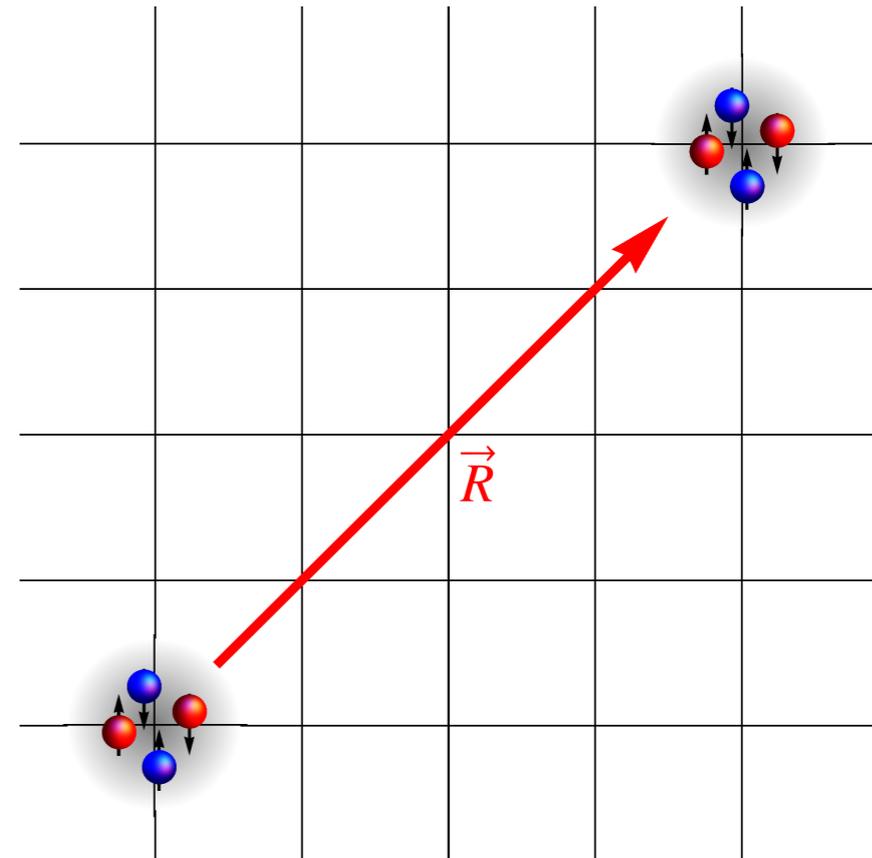
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$$|\vec{R}\rangle = \sum_{\vec{r}} |\vec{r} + \vec{R}\rangle_1 \otimes |\vec{r}\rangle_2$$



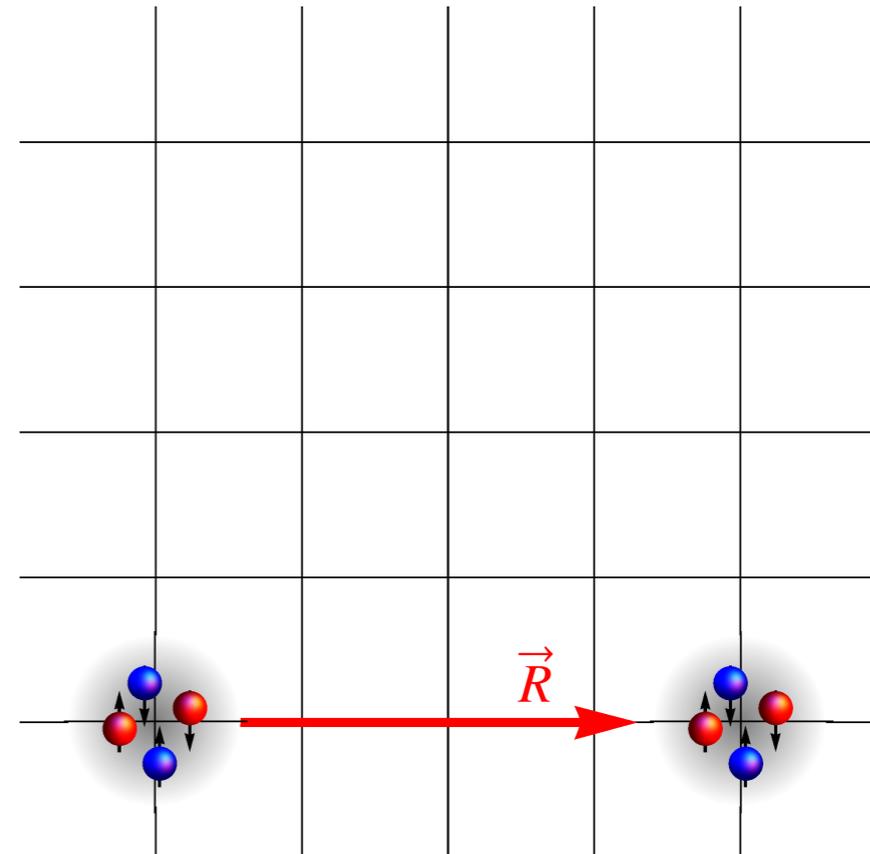
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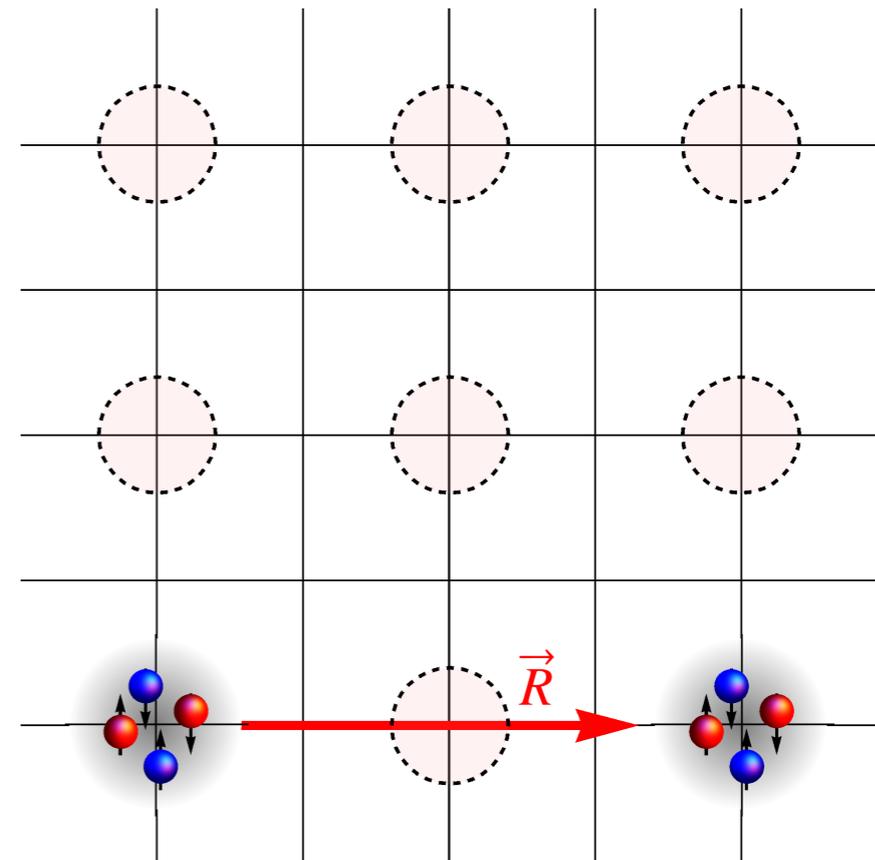
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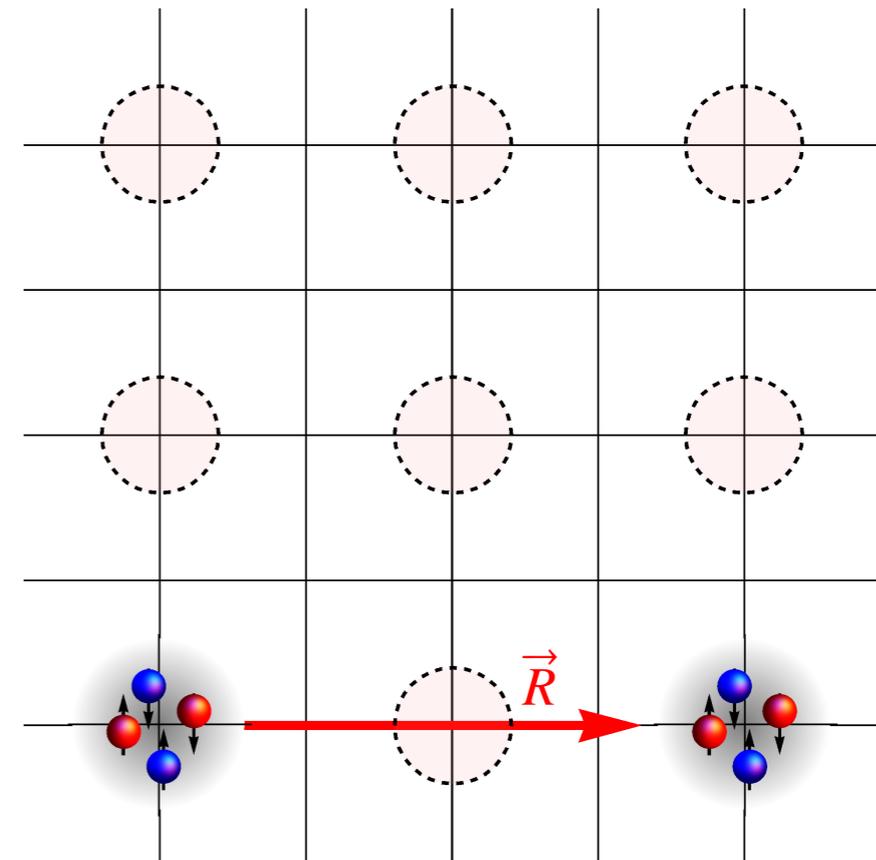


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$$|\vec{R}\rangle_\tau = \exp(-H\tau)|\vec{R}\rangle$$



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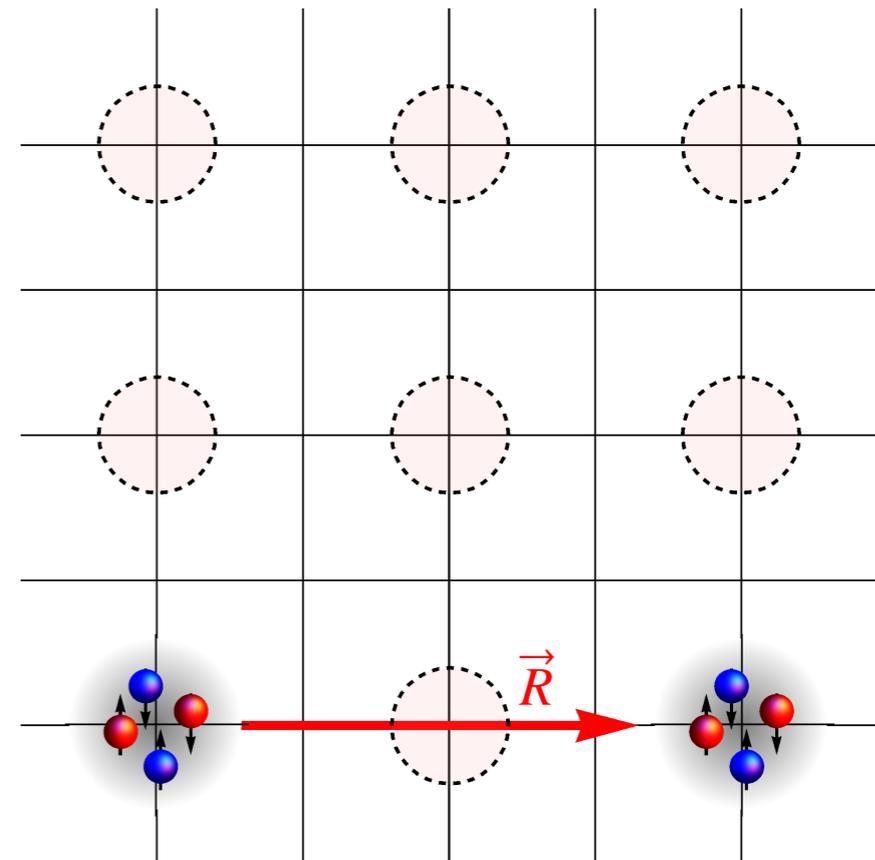
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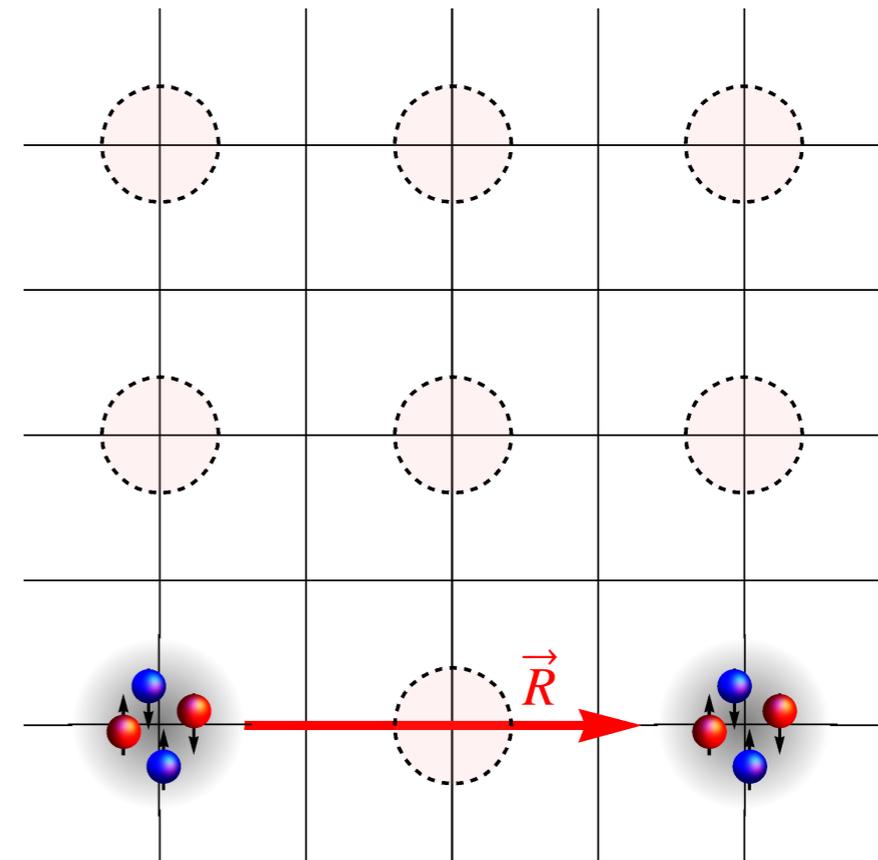
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Matrix elements of dressed cluster states





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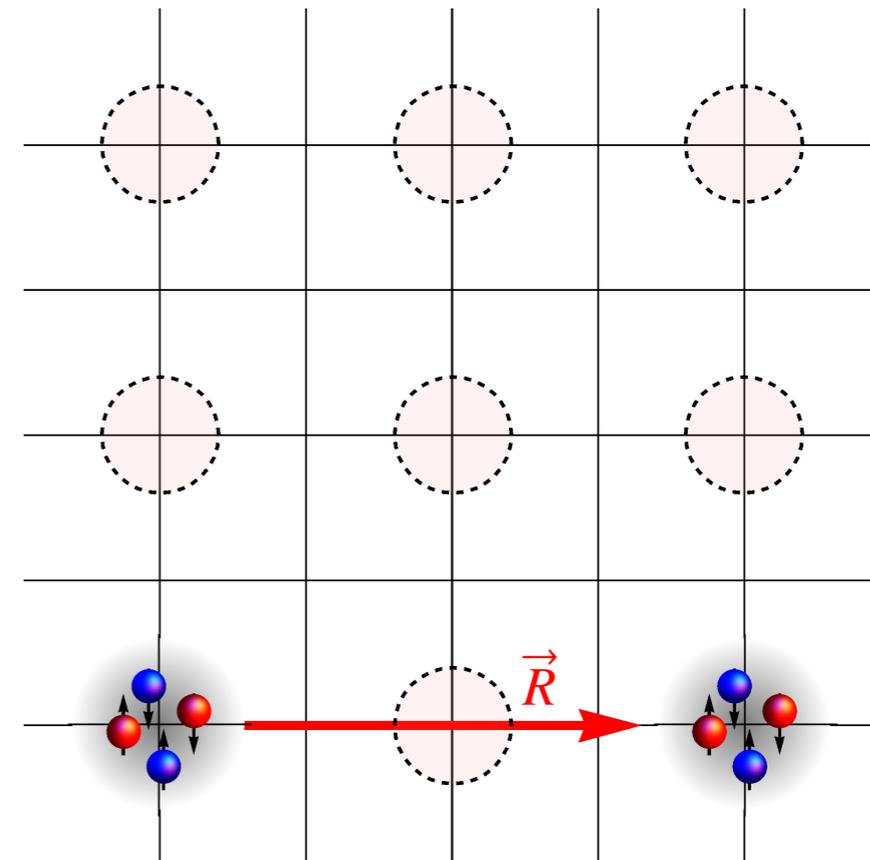
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Norm matrix

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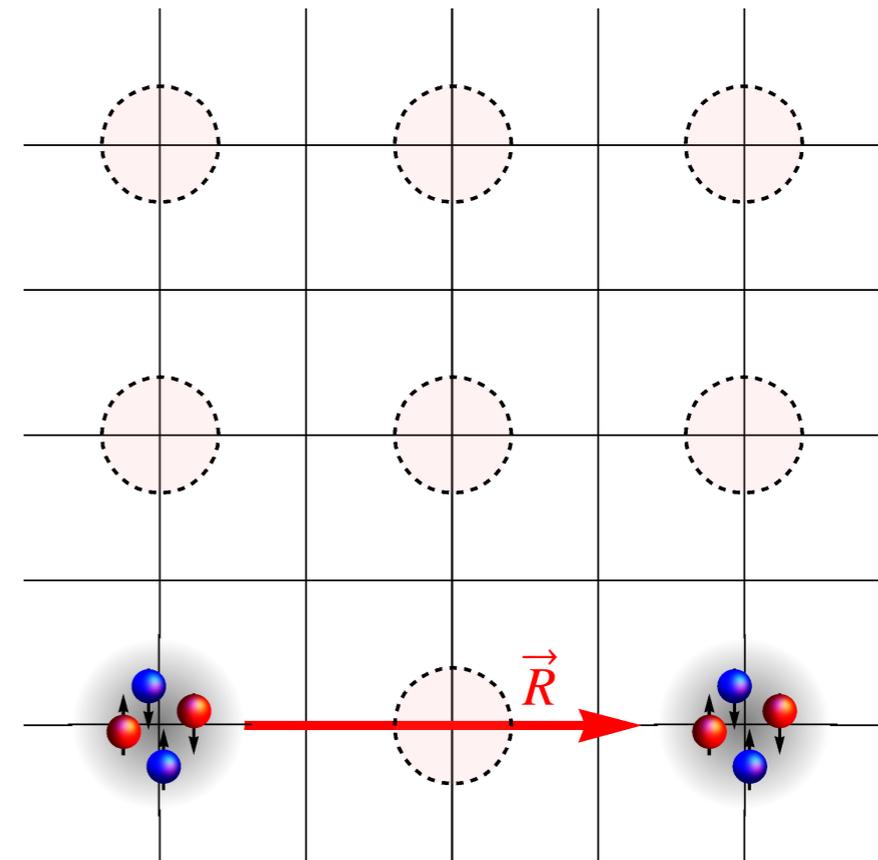
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similar to NCSM

Navratil, Quaglioni, Phys. Rev. C 83, 044609 (2011).

Norm matrix

Matrix elements of dressed cluster states





- **2 Step:** Extracting phase shifts

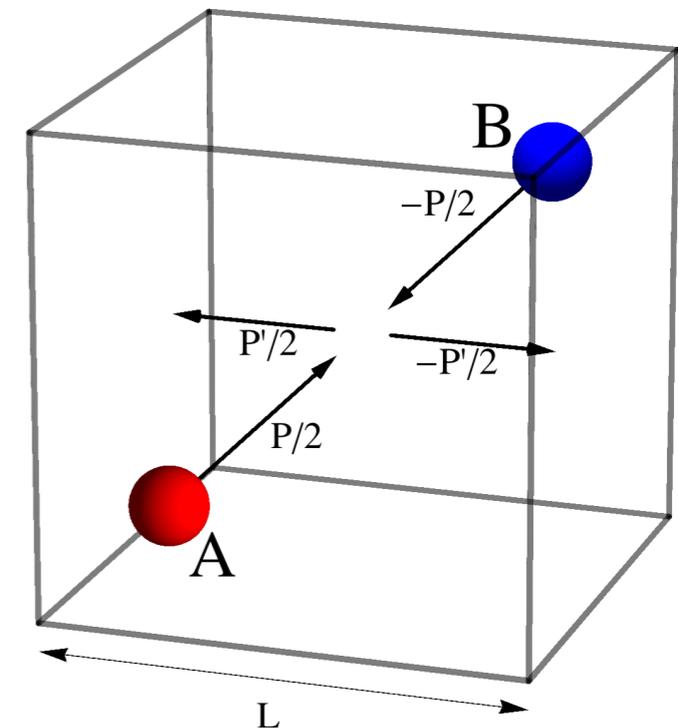


- **2 Step:** Extracting phase shifts
- **Lüscher's method:** Relation between energy levels in a finite periodic box and the infinite volume scattering phase shifts

M.Lüscher, *Commun. Math. Phys.* 105 (1986), 153

$$p \cot \delta_0(p) = \frac{1}{\pi L} S(\eta), \quad \eta = \frac{p(L)^2 L^2}{4\pi^2}$$

$$S(\eta) = \lim_{\Lambda \rightarrow \infty} \left[\sum_{\vec{k} \in \mathbb{Z}^3} \frac{\theta(\Lambda^2 - \vec{k}^2)}{\vec{k}^2 - \eta} - 4\pi \Lambda \right]$$

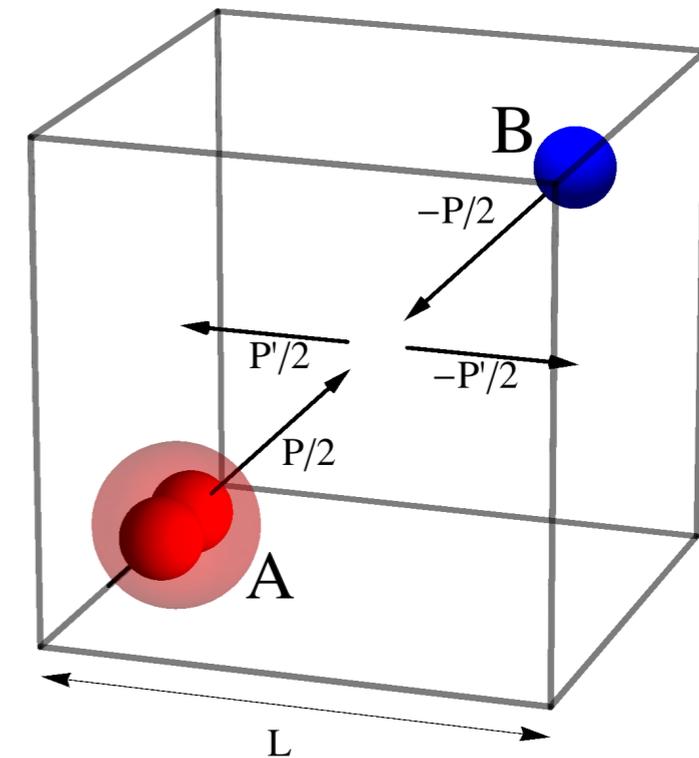


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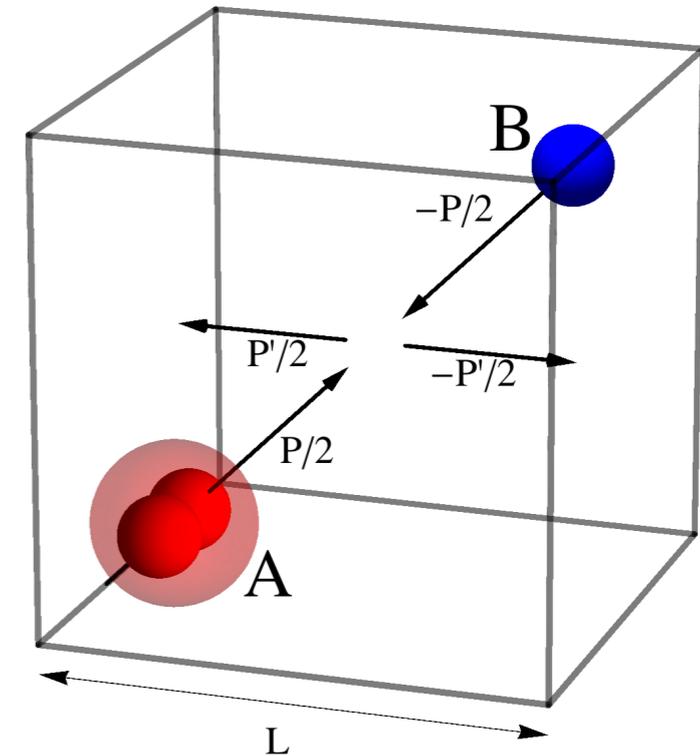


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- **Topological corrections** due to cluster character

$$E(p, L) = \frac{p^2}{2\mu} - B_1 - B_2 + \tau_1(\eta) \Delta E_1(L) + \tau_2(\eta) \Delta E_2(L)$$

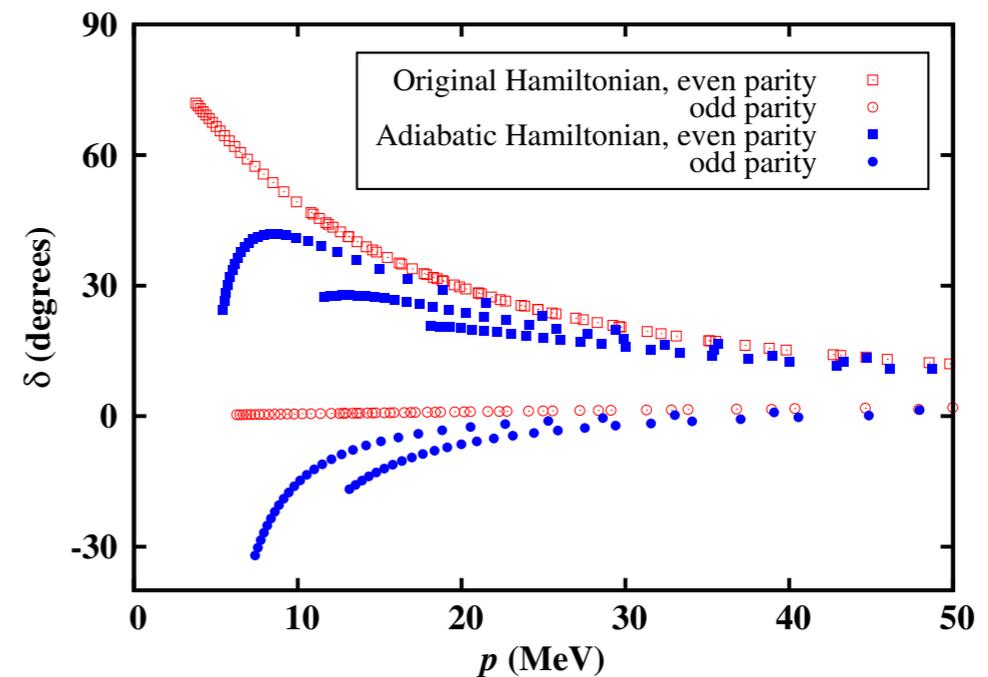
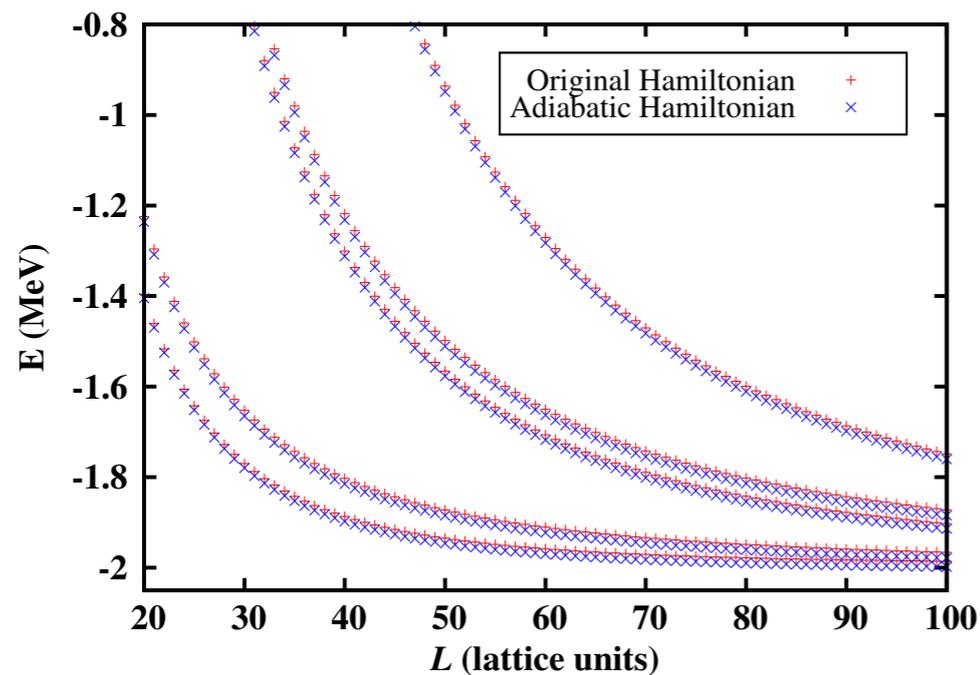
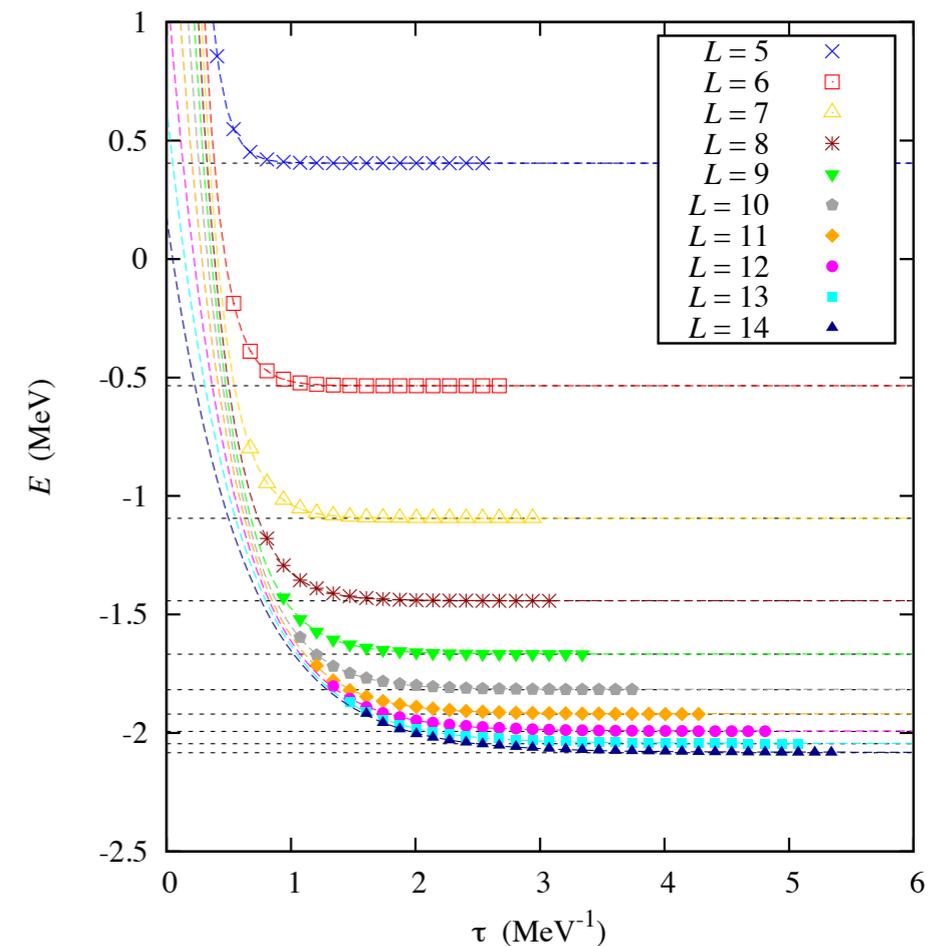
$$\tau(\eta) = \frac{1}{\sum_{\vec{k}} (\vec{k}^2 - \eta)^{-2}} \sum_{\vec{k}} \frac{\sum_{i=1}^3 \cos(2\pi k_i \alpha)}{3(\vec{k}^2 - \eta)^2}$$

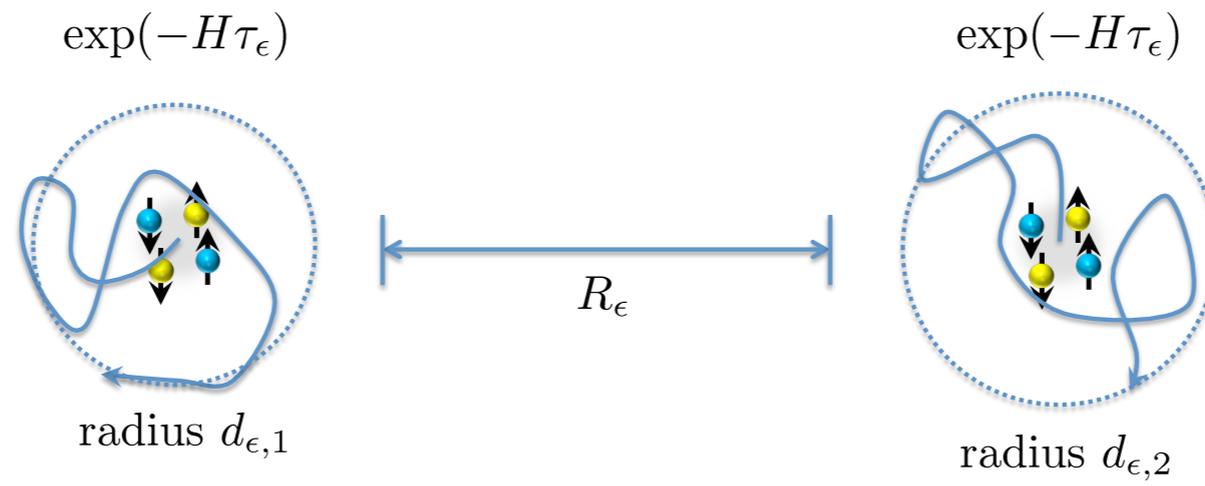
S. Bour, S. König, D. Lee, H.-W. Hammer and U.-G. Meißner, *Phys. Rev. D* 84 (2011), 091503

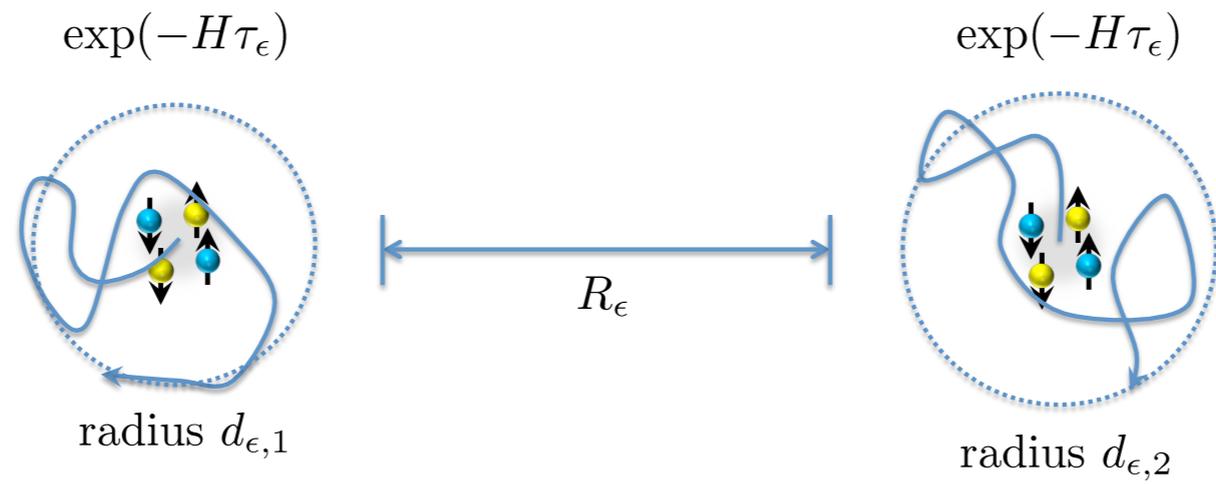
S. Bour, H.-W. Hammer, D. Lee and U.-G. Meißner, *Phys. Rev. C* 86 (2012), 034003

- There is an exponentially small error in energy levels due to Euclidean time projection
- In larger systems there is a statistical error due to Monte Carlo methods

Lüscher's method is unfortunately very sensitive to small errors in energy levels!

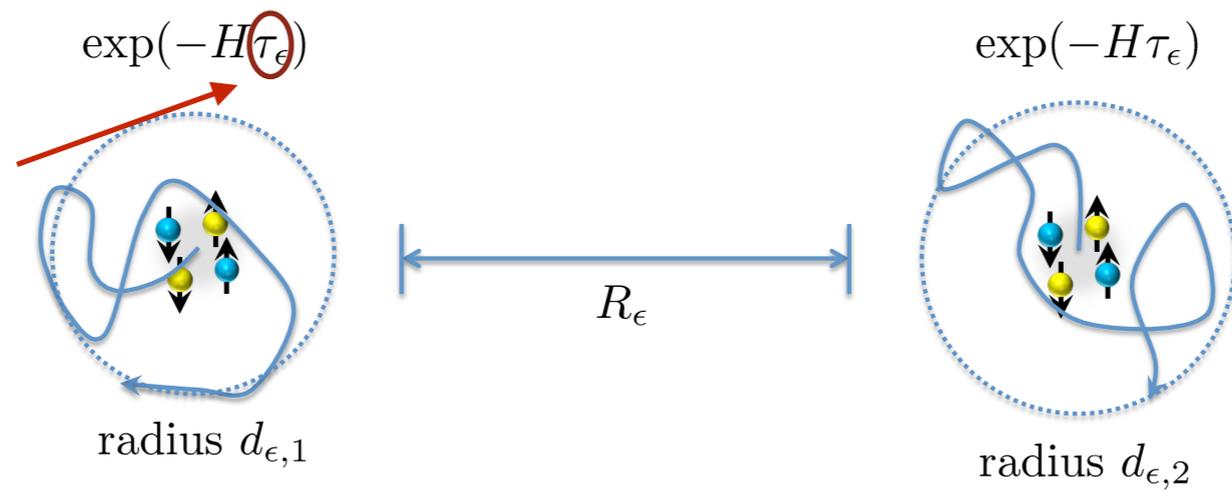






ϵ : relative error

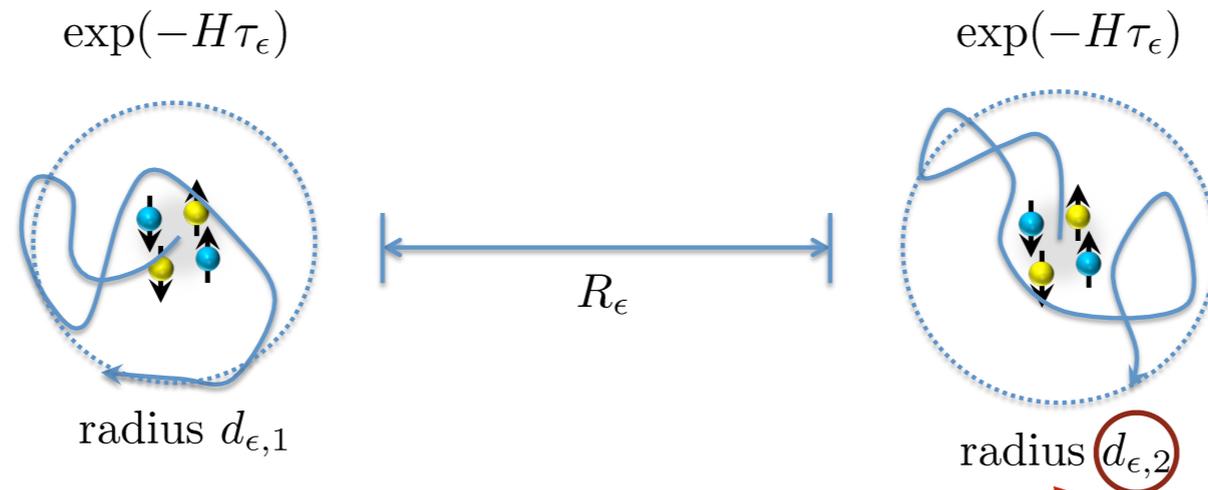
time interval such that
the relative contamination
from excited cluster states
is less than ϵ



ϵ : relative error



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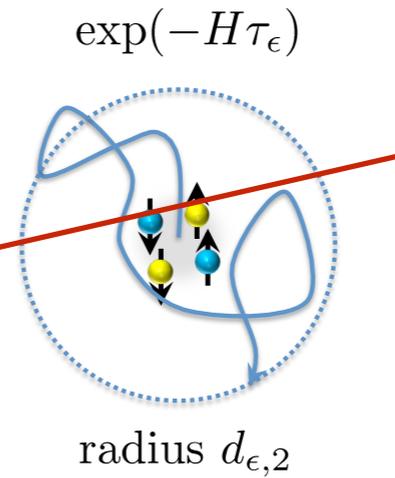
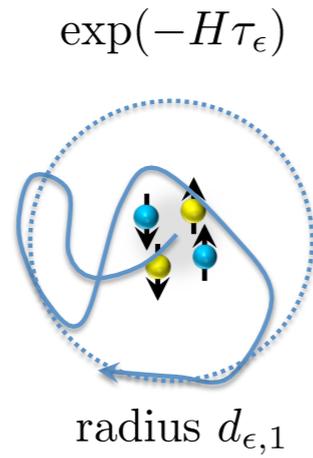
spatial diffusion that
the clusters undergo in time τ_ϵ

ϵ : relative error

Asymptotic cluster wave functions (I)



time interval such that the relative contamination from excited cluster states is less than ϵ

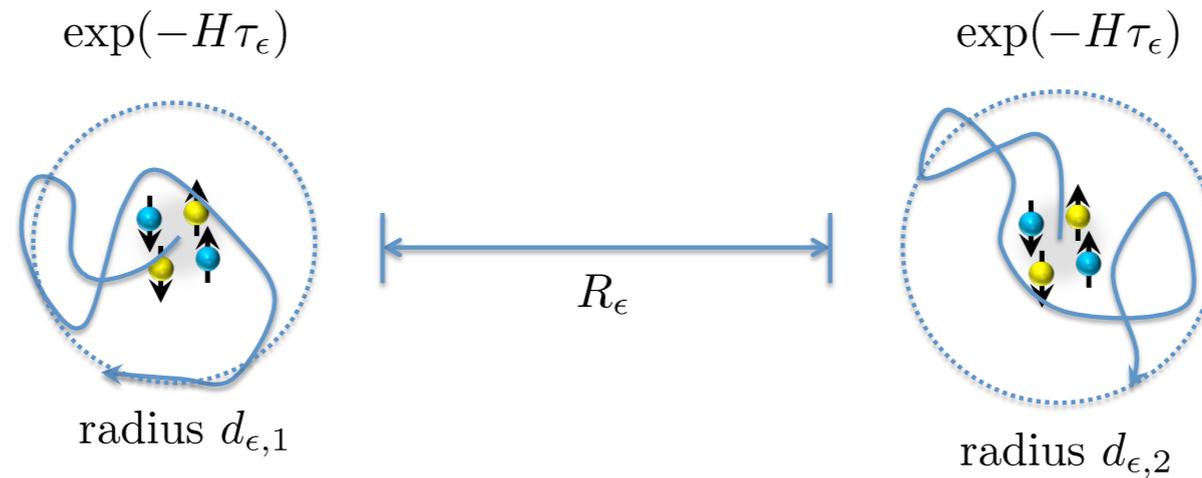


asymptotic distance such that the overlap between clusters is less than ϵ

spatial diffusion that the clusters undergo in time τ_ϵ

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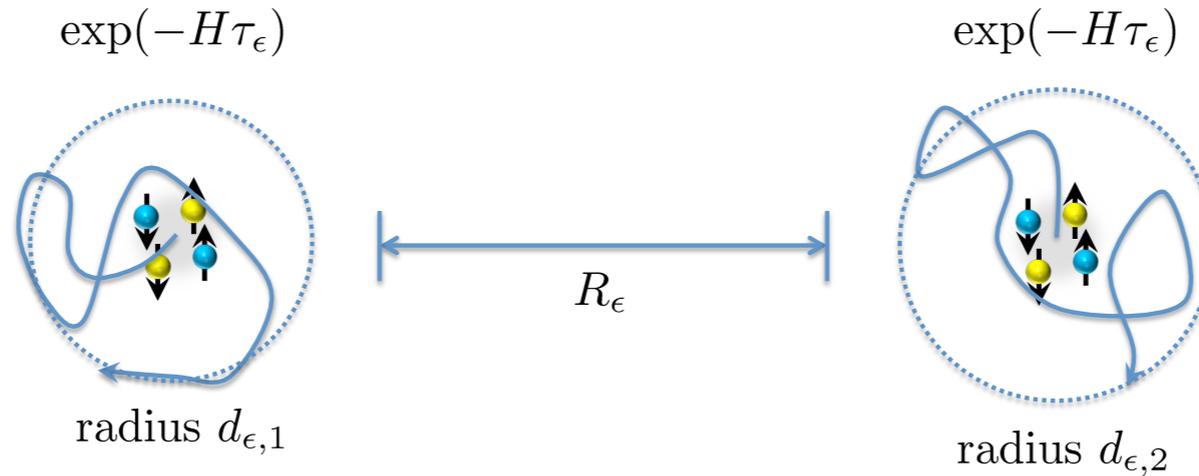
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the Hamiltonian is similar to a free lattice Hamiltonian H_{eff}

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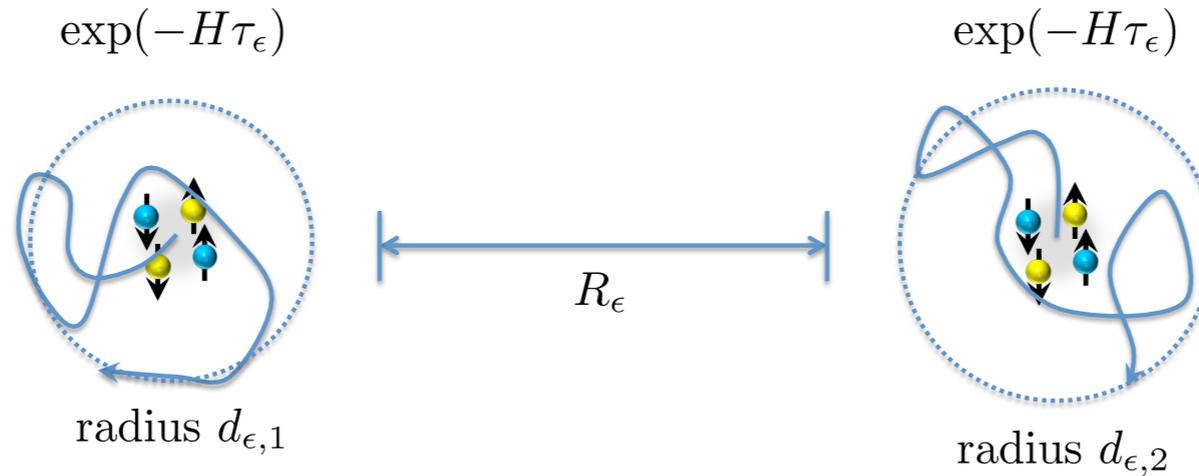
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\uparrow
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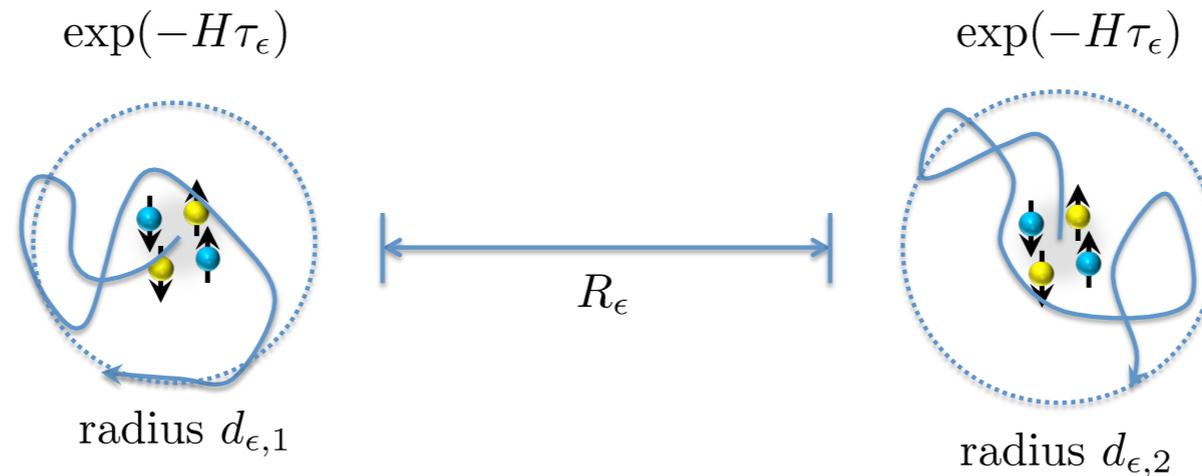
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- The asymptotic cluster wave function can also be used to extract phase shifts



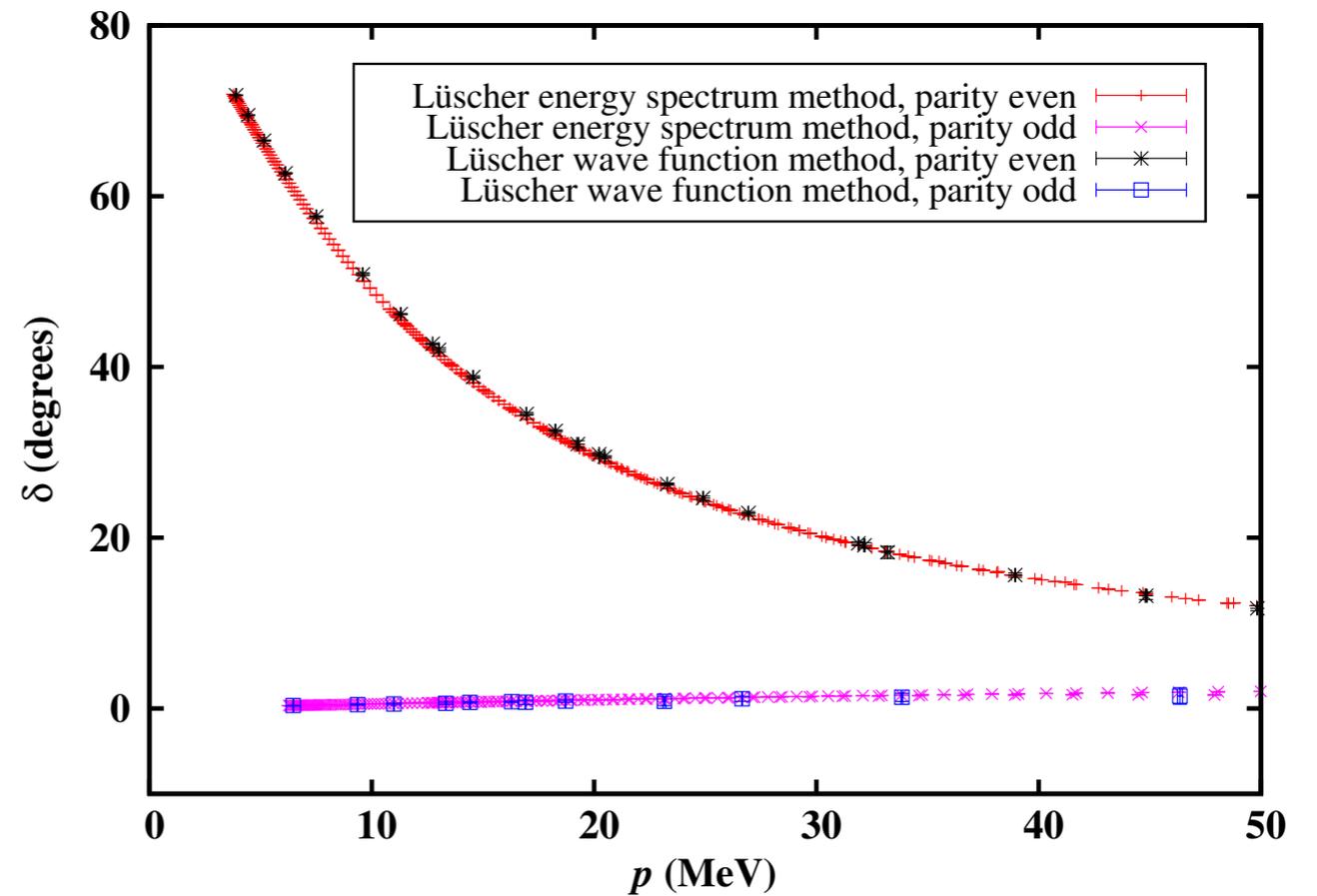
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check with Lüscher's method applied
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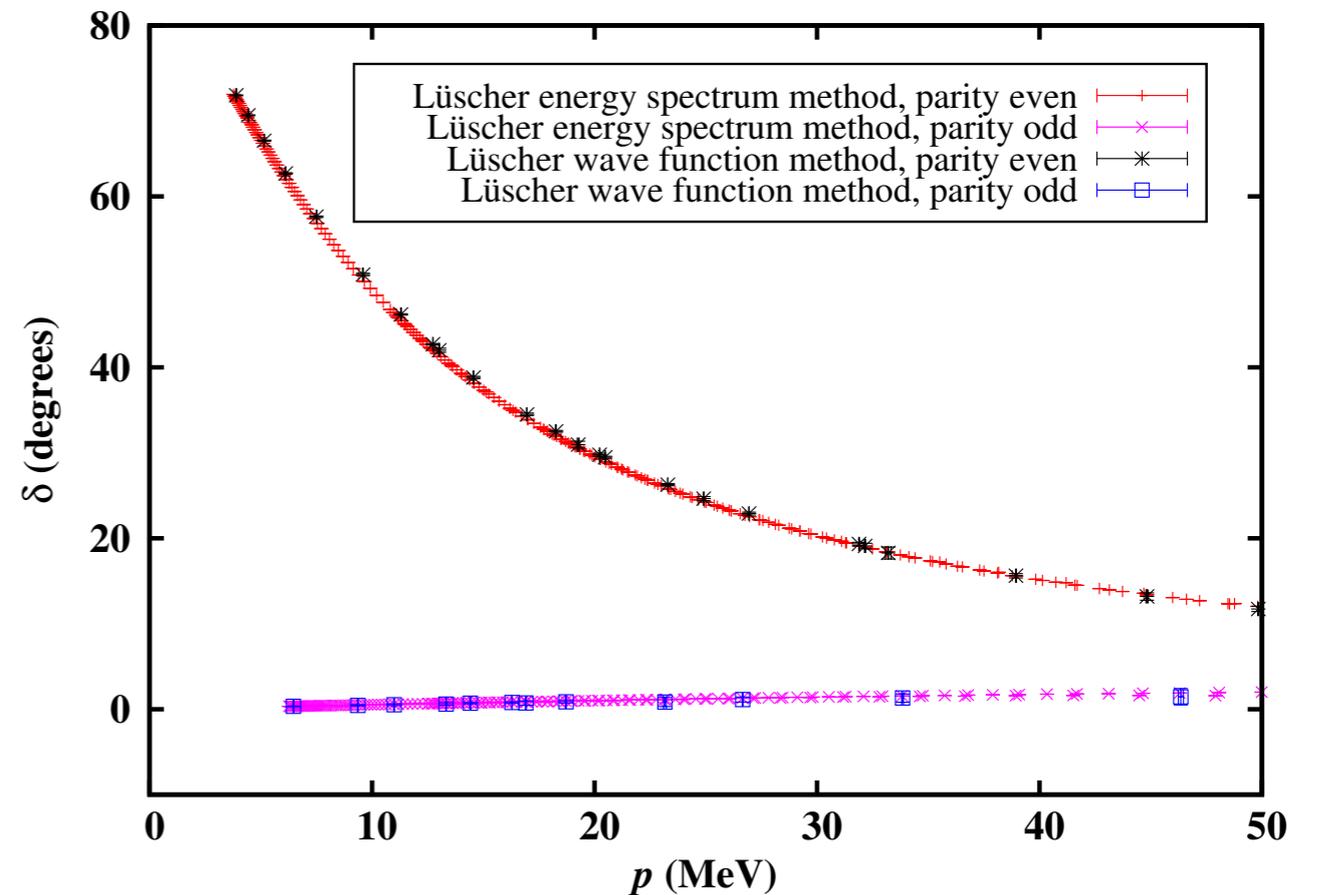


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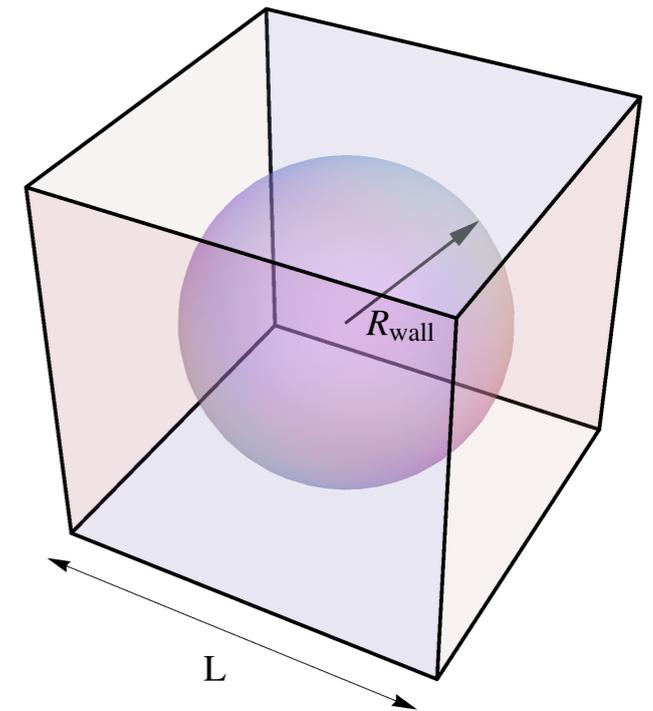
- 3-dim asymptotic form:

$$\Psi_{\ell, m_\ell}^{(p)}(\mathbf{r}) = A_\ell Y_{\ell, m_\ell}(\theta, \phi) [\cos \delta_\ell(p) j_\ell(pr) - \sin \delta_\ell(p) n_\ell(pr)]$$



- Impose a hard wall on the relative separation of two point-like particles and fit it to the asymptotic form

[Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 \(2007\) 185](#)



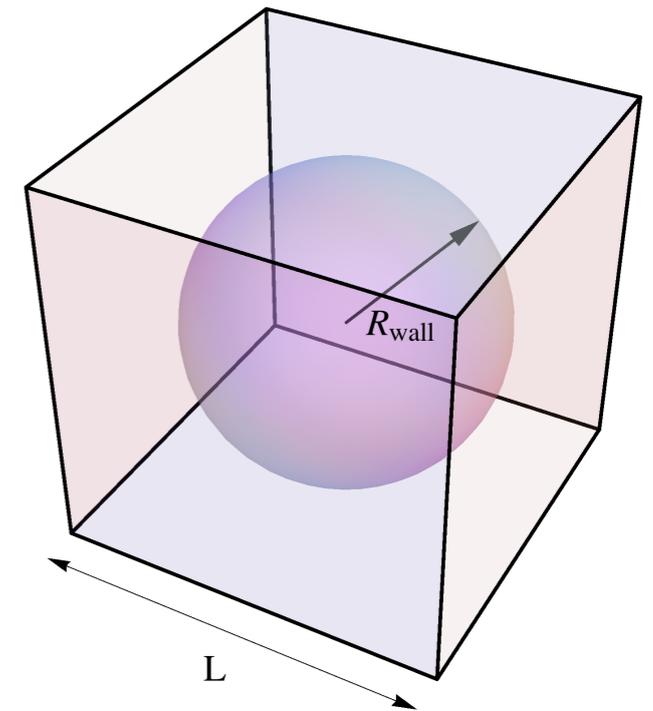


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Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

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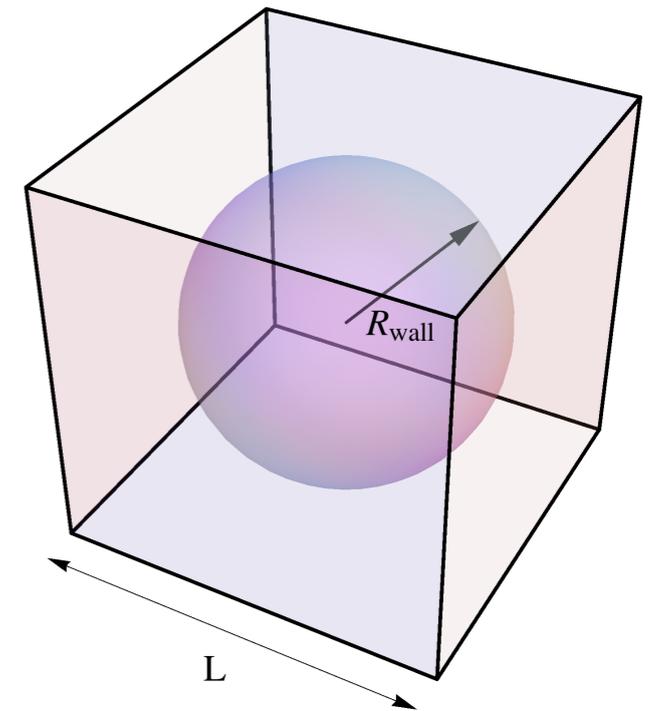


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Borasoy, Epelbaum, Krebs, Lee, Meißner, EPJA 34 (2007) 185

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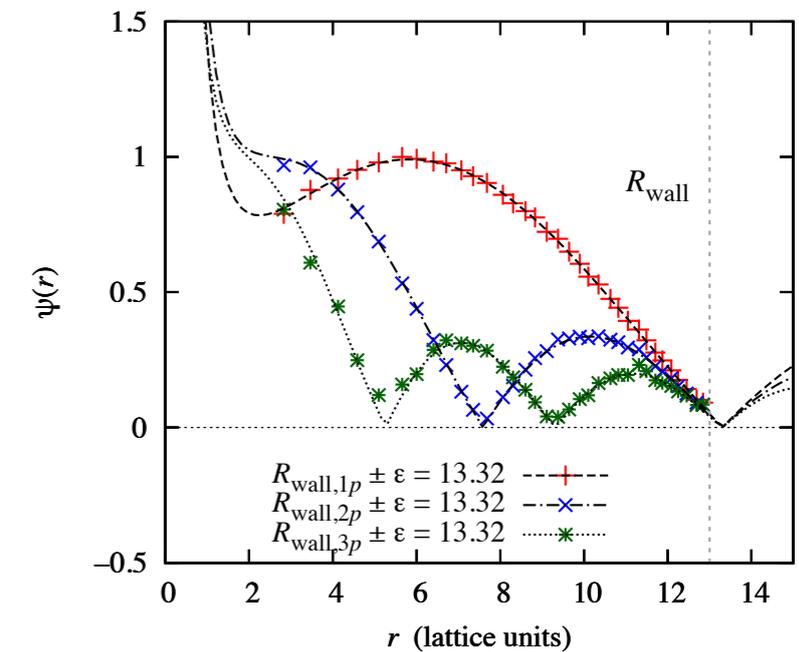
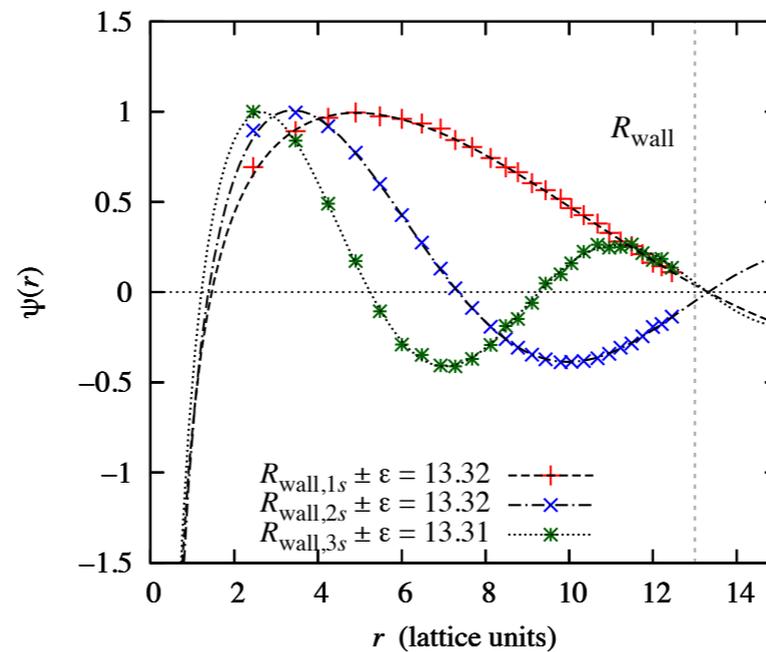
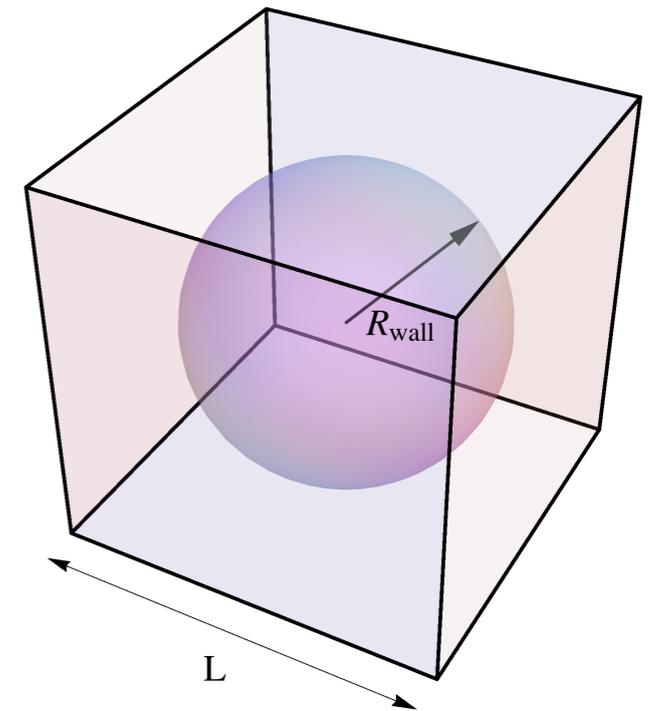
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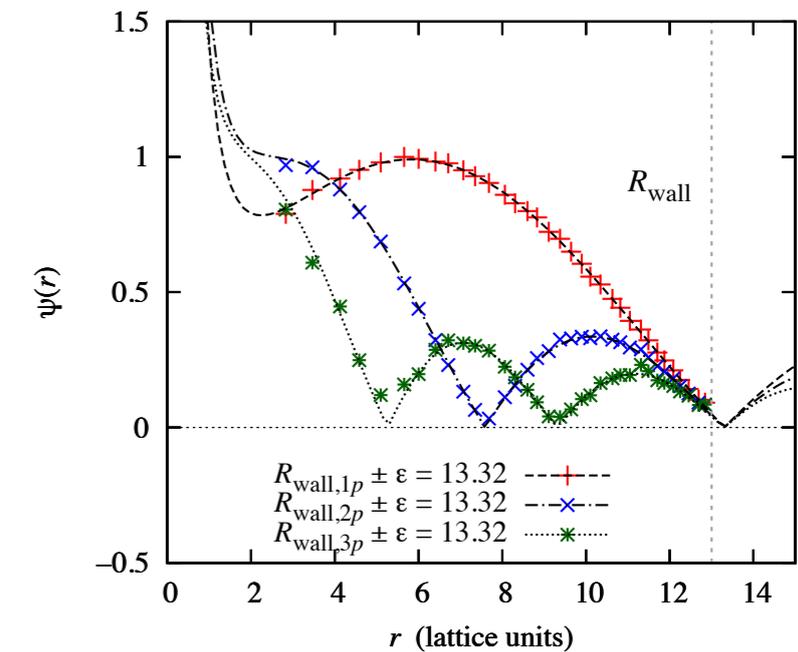
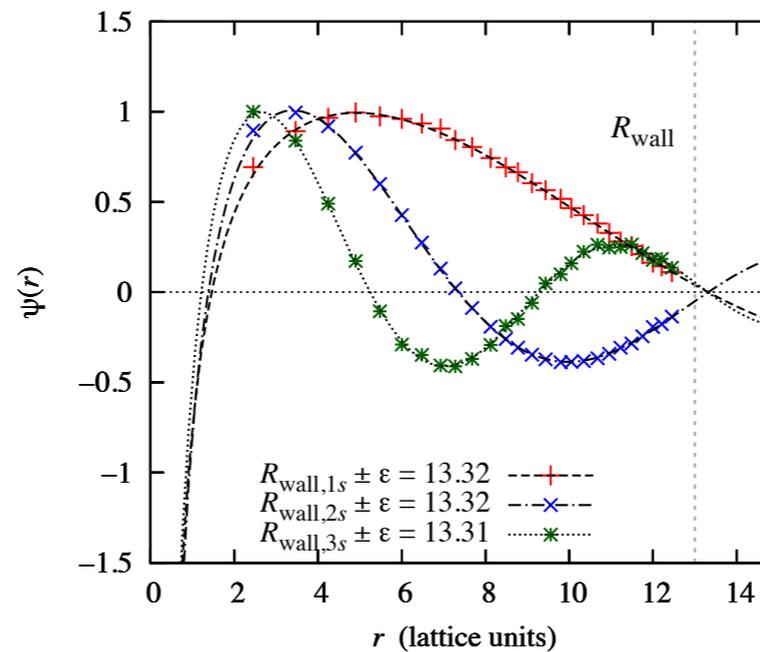
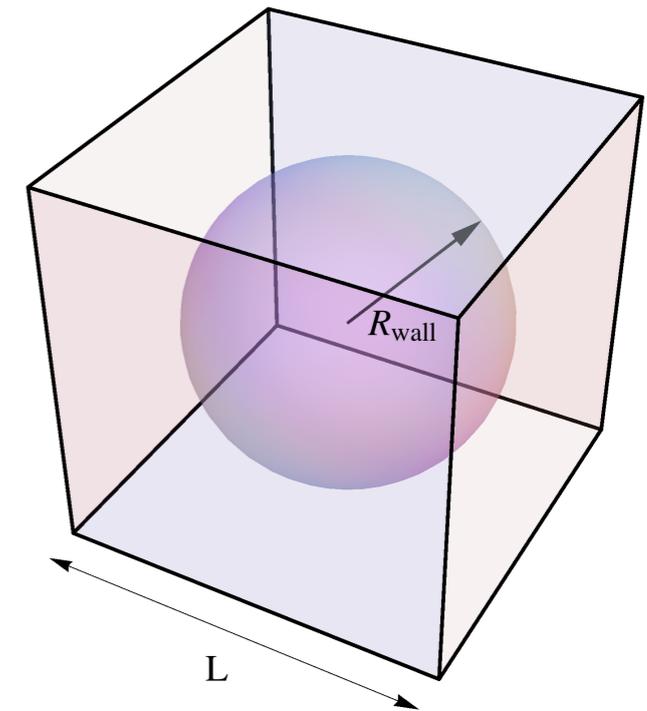
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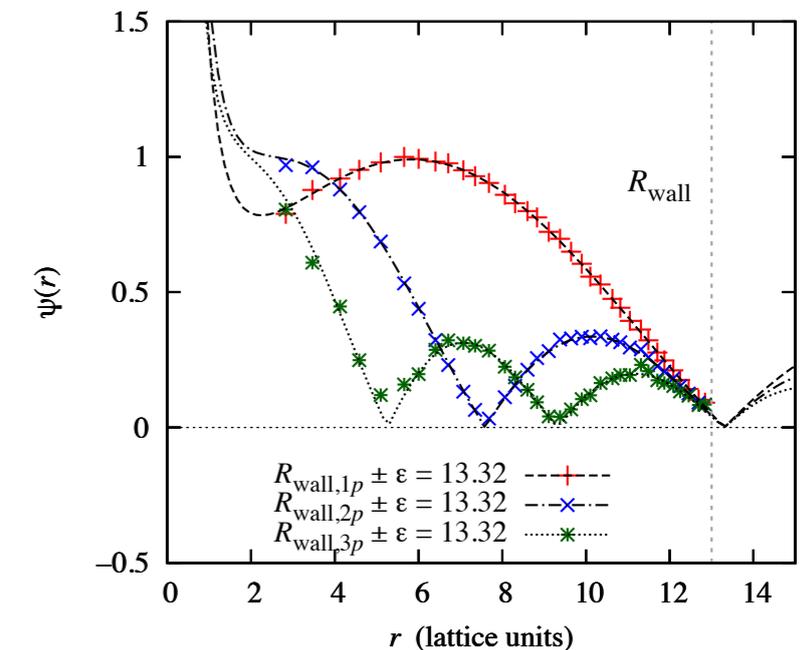
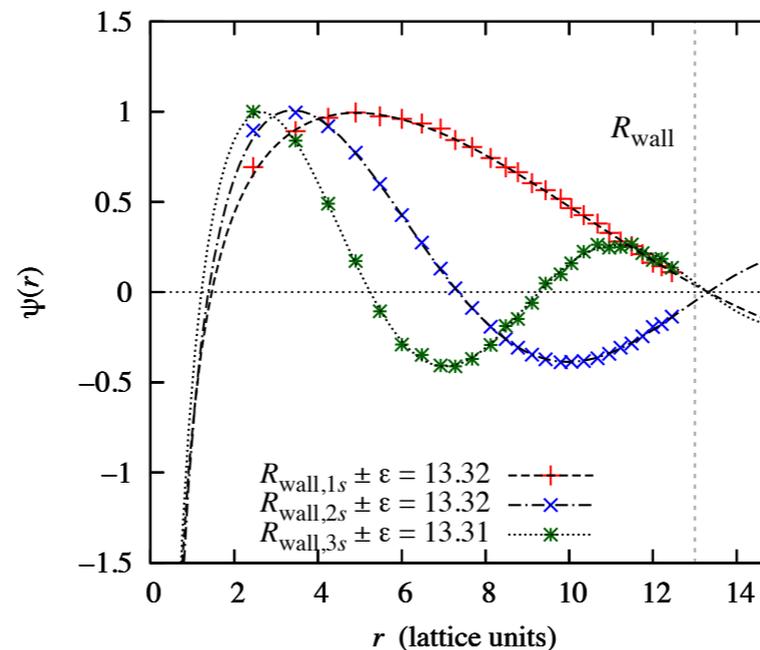
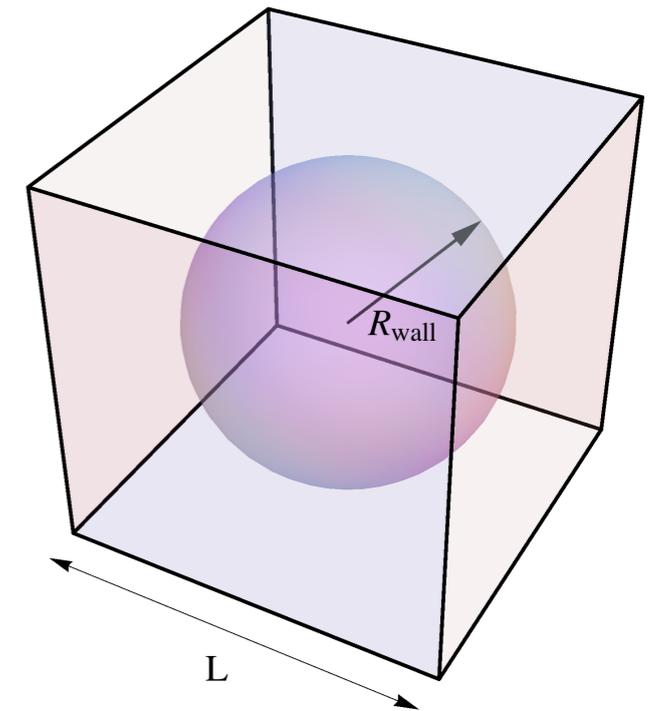
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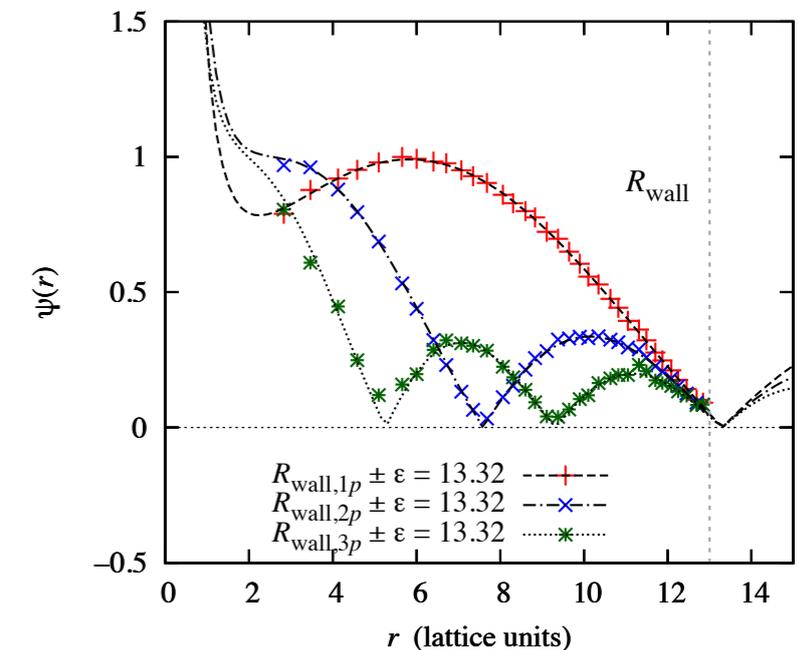
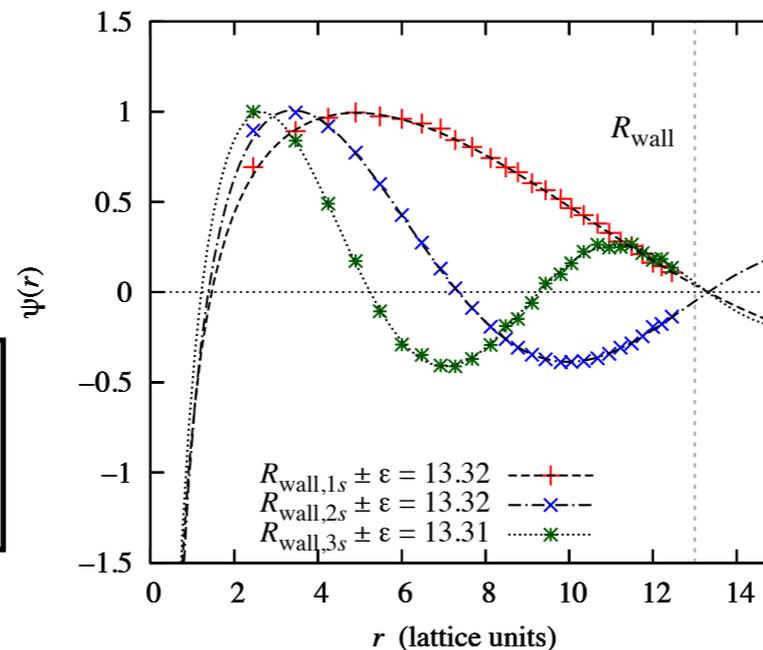
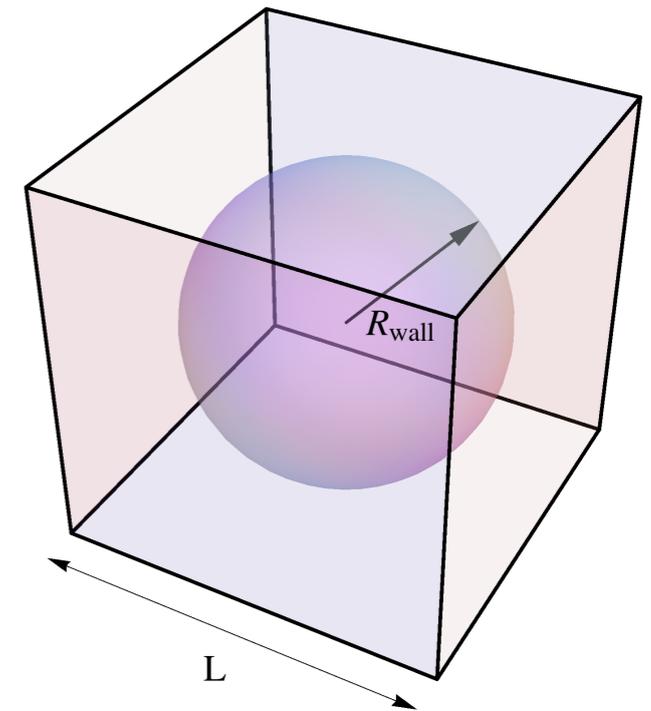
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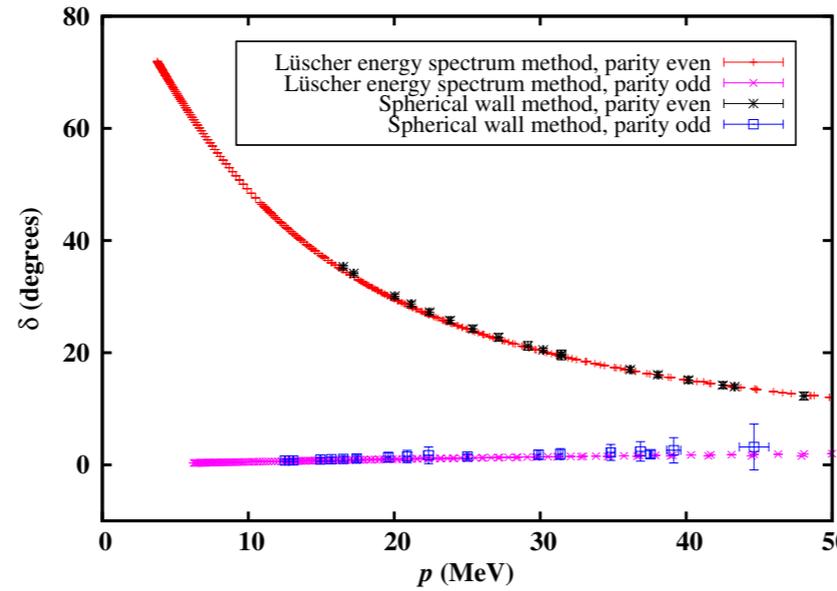
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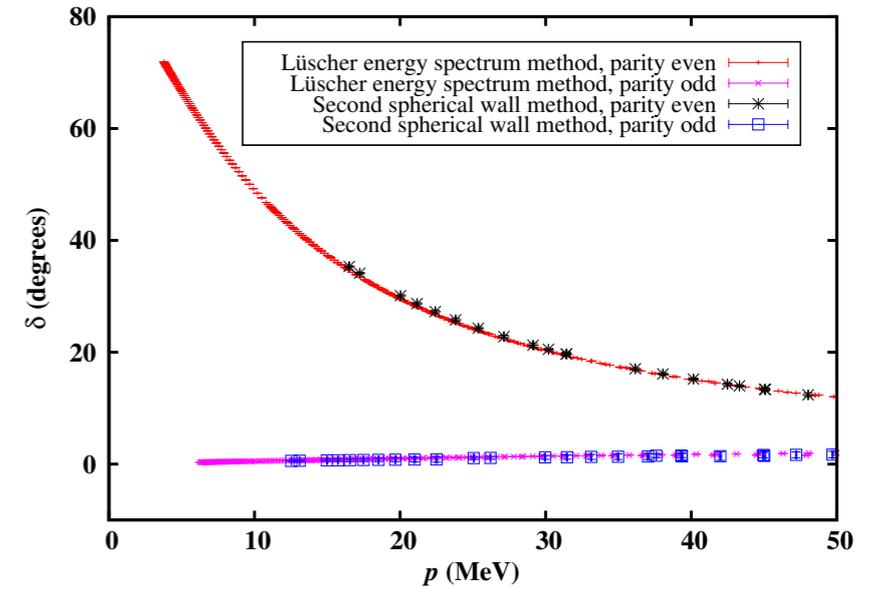
$$\delta_\ell(p) = \begin{cases} -pR'_{\text{wall}} + \frac{\pi(\ell+1)}{2} \text{ mod } \pi \\ \tan^{-1} \left[\frac{j_\ell(pR'_{\text{wall}}/a)}{n_\ell(pR'_{\text{wall}}/a)} \right] \end{cases}$$



- particle-dimer in one dimension

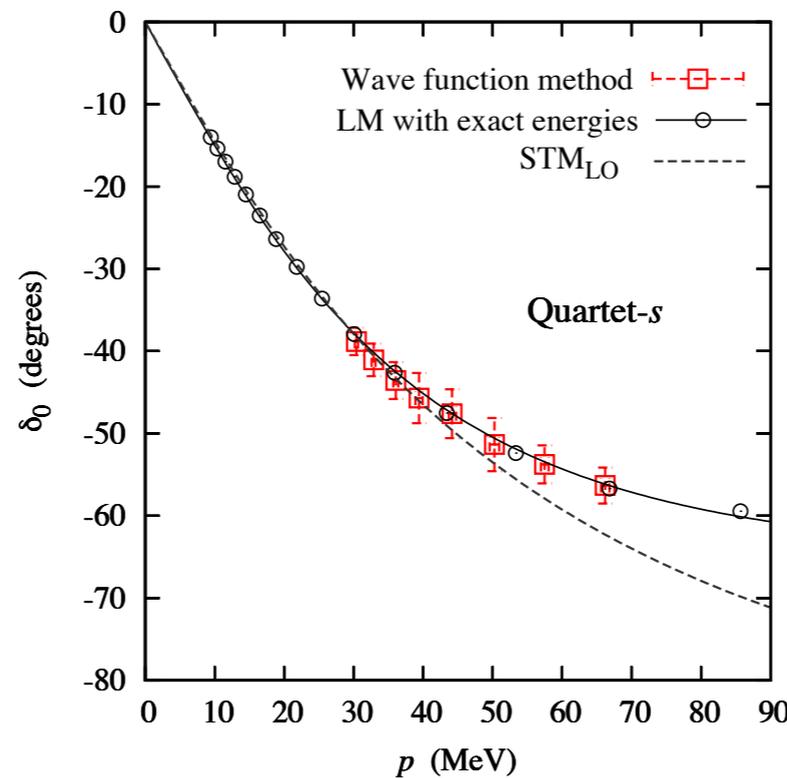


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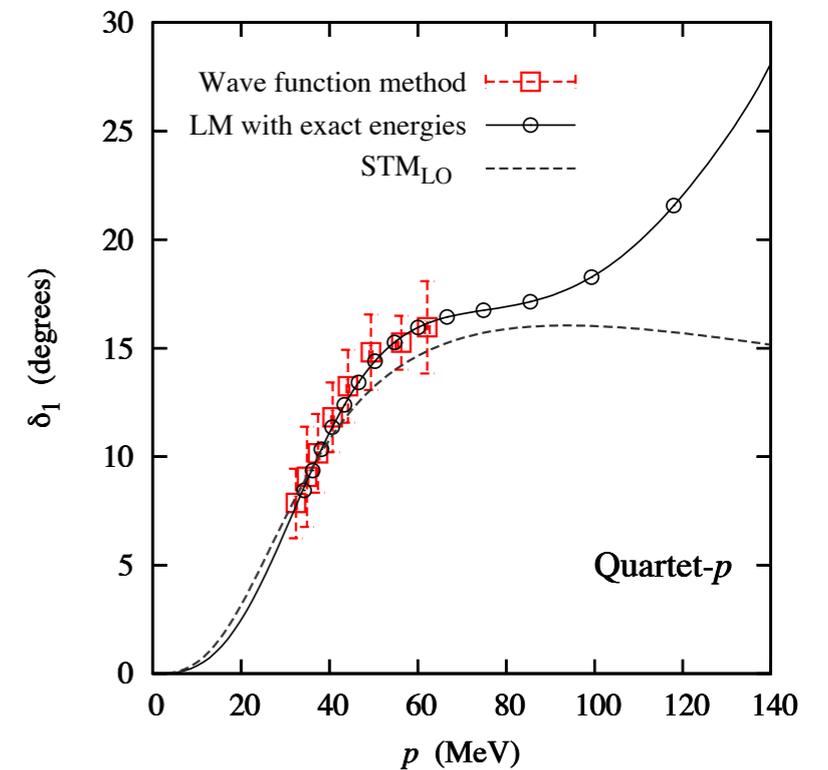


2-parameter fit

- fermion-dimer in three dimensions



2-parameter fit

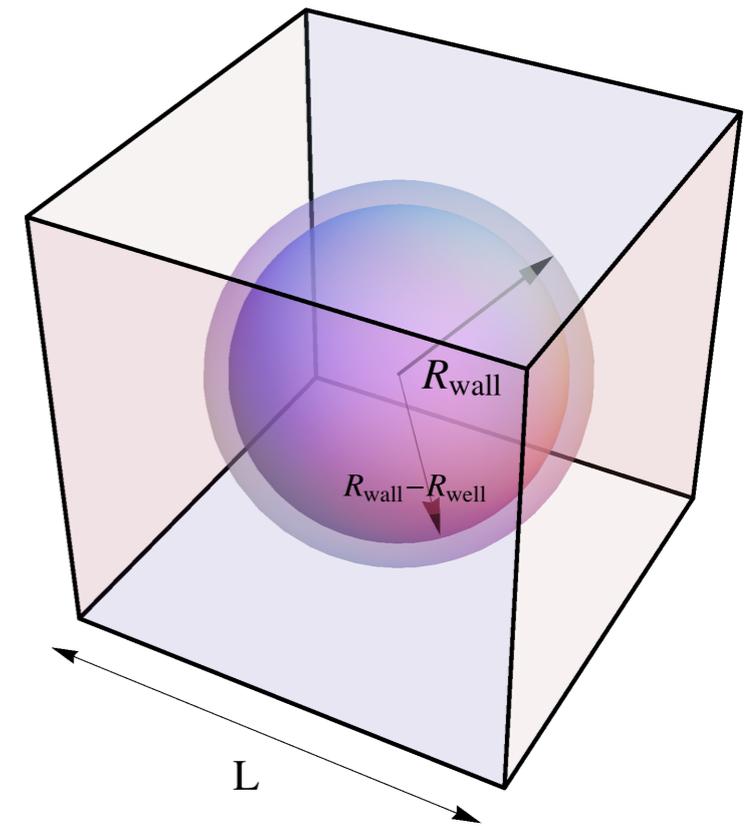


Gabbiani, Bedaque, Grißhammer, NPA 675 (2000) 601

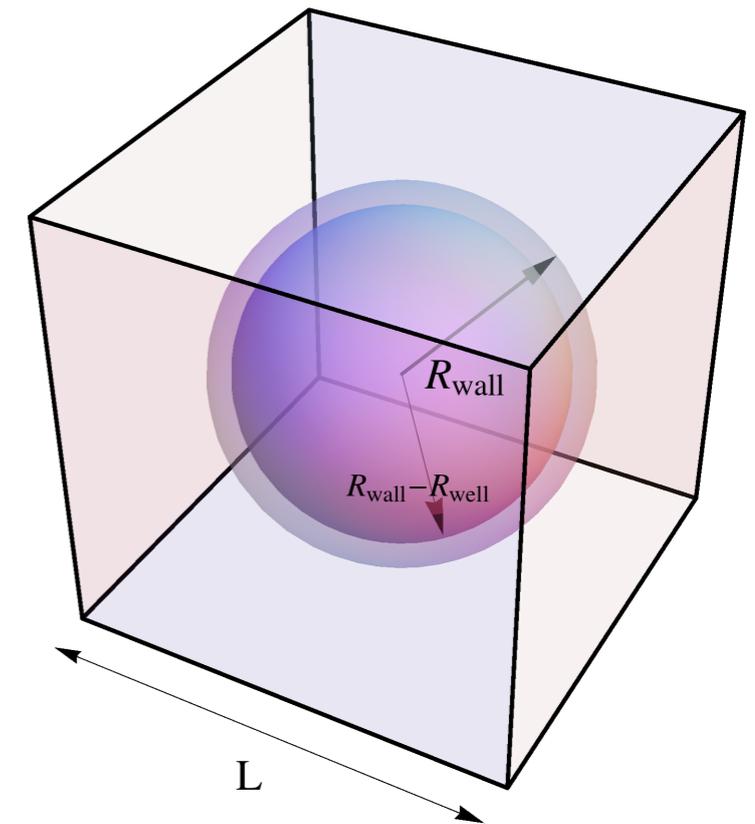


- one disadvantage of the spherical wall method are large R_{wall} and L for low energies

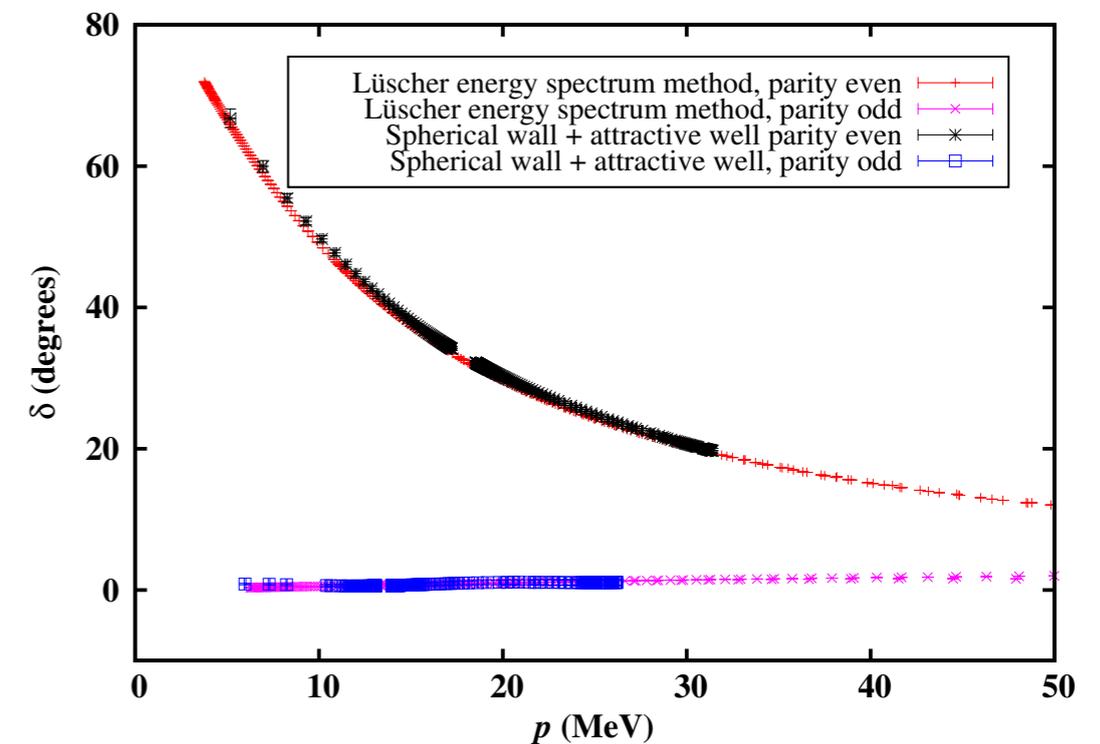
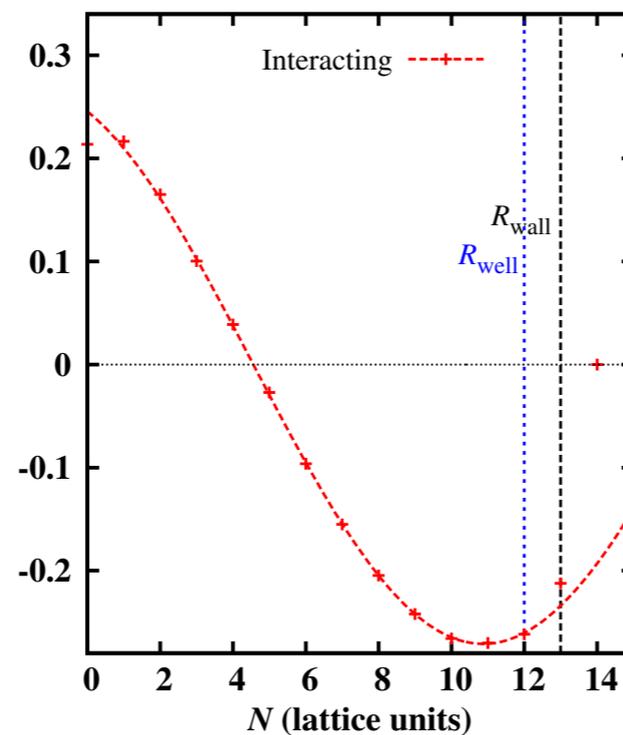
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Summary and Outlook





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Thank you for your attention!