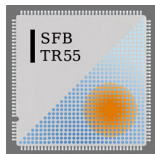


SU(3) flavor breaking in baryon octet light-cone distribution amplitudes

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DAs: Definition

Schematically:

$$|B\rangle = |qqq\rangle + |qqqg\rangle + |qqq\bar{q}q\rangle + \dots$$

- **Distribution Amplitudes** (DAs) describe the distribution of the lightcone-momentum within a specific Fock state
- in **hard exclusive processes** Fock states are increasingly power-suppressed with a rising number of partons

⇒ at high momentum transfer the 3-quark contribution plays the most important role

DAs: Definition

Three quark DAs are defined via baryon-to-vacuum matrix elements of 3-quark operators:

$$\begin{aligned} & \langle 0 | q_\alpha^a(a_1 n) q_\beta^b(a_2 n) q_\gamma^c(a_3 n) | B(p, s) \rangle \\ &= \int [dx] e^{-i n \cdot p} \sum_i a_i x_i \left[V_1^B(x_i) (\not{p} C)_{\alpha\beta} (\gamma_5 B^+(p, s))_\gamma + \dots \right] \end{aligned}$$

- color antisymmetrization and Wilson-lines are not written out explicitly
- a , b and c are **flavor indices**; α , β and γ are **Dirac indices**; n is a light-like vector
- on the r.h.s. one has 24 different structures and the same number of different DAs:

	twist-3	twist-4	twist-5	twist-6
scalar		S_1^B	S_2^B	
pseudoscalar		P_1^B	P_2^B	
vector	V_1^B	V_2^B, V_3^B	V_4^B, V_5^B	V_6^B
axialvector	A_1^B	A_2^B, A_3^B	A_4^B, A_5^B	A_6^B
tensor	T_1^B	T_2^B, T_3^B, T_7^B	T_4^B, T_5^B, T_8^B	T_6^B

→ ME decomposition by BFMS: **Nucl.Phys. B589 (2000) 381**

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- on the l.h.s. one has to **choose the correct flavor content**
- the order of flavors is relevant for the symmetry properties of the DAs
- a convenient choice is:

$$\begin{array}{llll} p \hat{=} uud, & n \hat{=} ddu, & \Sigma^+ \hat{=} uus, & \Sigma^0 \hat{=} uds, \\ \Sigma^- \hat{=} dds, & \Xi^0 \hat{=} ssu, & \Xi^- \hat{=} ssd, & \Lambda \hat{=} uds. \end{array}$$

Three quark operators

- split the quark fields in left- and righthanded parts to make use of chiral symmetry

$$\begin{aligned}
 q_\alpha^a(a_1 n) q_\beta^b(a_2 n) q_\gamma^c(a_3 n) &= \mathcal{O}_{RR, \alpha\beta\gamma}^{abc}(a_1, a_2, a_3) + \mathcal{O}_{LL, \alpha\beta\gamma}^{abc}(a_1, a_2, a_3) \\
 &+ \mathcal{O}_{RL, \alpha\beta\gamma}^{abc}(a_1, a_2, a_3) + \mathcal{O}_{LR, \alpha\beta\gamma}^{abc}(a_1, a_2, a_3) \\
 &+ \mathcal{O}_{RL, \gamma\alpha\beta}^{cab}(a_3, a_1, a_2) + \mathcal{O}_{LR, \gamma\alpha\beta}^{cab}(a_3, a_1, a_2) \\
 &+ \mathcal{O}_{RL, \beta\gamma\alpha}^{bca}(a_2, a_3, a_1) + \mathcal{O}_{LR, \beta\gamma\alpha}^{bca}(a_2, a_3, a_1)
 \end{aligned}$$

- i.e. there are **only two** different **types of operators**: chiral-even ($\mathcal{O}_{RR}/\mathcal{O}_{LL}$) and chiral-odd ($\mathcal{O}_{RL}/\mathcal{O}_{LR}$)

$$\mathcal{O}_{XY, \alpha\beta\gamma}^{abc}(a_1 n, a_2 n, a_3 n) = q_{X, \alpha}^a(a_1 n) q_{X, \beta}^b(a_2 n) q_{Y, \gamma}^c(a_3 n)$$

Ansatz for the effective operator

$$\mathcal{O}_{XY,\alpha\beta\gamma}^{abc}(a_1 n, a_2 n, a_3 n) = \int [dx] \sum_{i,j} \sum_{k=1}^{k_j} \mathcal{F}_{XY}^{i,j,k}(x_1, x_2, x_3) \Gamma_{\alpha\beta\gamma\delta}^{i,XXY} B_{\delta,abc}^{j,k,XXY}(z)$$

- the functions \mathcal{F} are **distribution amplitudes**
 → play the role of **LECs**
- Γ : i labels the possible **Dirac structures** (6 for chiral-even and 6 for chiral-odd)
 → Lorentz indices are contracted with n or derivatives acting on B
- B : contains **hadron fields**; e.g. $B_{\delta,abc}^{1,1,RRL} = (u B_{\delta})_{aa'} (u)_{bb'} (u^{\dagger})_{cc'} \varepsilon_{a'b'c'}$
 $j = 2, 3 \rightarrow$ structures with quark mass matrix (contained in χ^+)
 $k = 2, 3, \dots \rightarrow$ different positions of baryon-octet field B_{δ} and χ^+

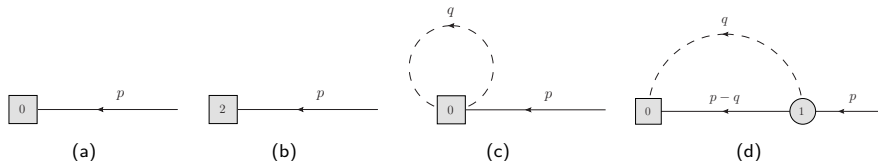
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- $z = n \sum_i x_i a_i$ and $[dx] = dx_1 dx_2 dx_3 \delta(1 - x_1 - x_2 - x_3)$
 → correct behaviour under **translations in n direction**
- **Parity**: relates left- and right-handed operators → $\mathcal{F}_{RR} = -\mathcal{F}_{LL}$ and $\mathcal{F}_{LR} = -\mathcal{F}_{RL}$
- **CP or Time-reversal**: yield “reality condition” → $\mathcal{F}^\dagger = \mathcal{F}$ up to an overall phase
- symmetry under **quark exchange**: relates different \mathcal{F} to each other

→ detailed description: **JHEP 1505 (2015) 073**

Leading one-loop calculation



- loop integrals are (more or less) straightforward
- multiply with \sqrt{Z} to take into account wave function renormalization
- we use IR regularization scheme
 - Becher, Leutwyler: **Eur. Phys. J. C9 (1999) 643**
- **main challenge:** handling the large number of different structures
- **last step:** obtain results for the 24 standard DAs $S_{1,2}^B$, $P_{1,2}^B$, V_{1-6}^B , A_{1-6}^B , T_{1-8}^B by matching to the general matrix element decomposition

Results: Overview

(a posteriori) we can define specific combinations of DAs with **remarkable features**:

- 1 **the Λ fits in** with the other octet baryons: same value for DAs at the flavor symmetric point (at the cost of a **different definition** for the Λ)
- 2 up to the fit parameters **extrapolation formulas** are **twist independent**
- 3 complete non-analytic structure contained in prefactors
- 4 constraints for $SU(3)_f$ symmetry breaking terms: e.g. $\Delta\Pi_3^\Sigma = -\frac{1}{2}\Delta\Phi_{+,3}^\Sigma - \frac{3}{2}\Delta\Phi_{+,3}^\Lambda$

Example: Leading twist DAs

$$\Phi_{\pm,3}^B(x_1, x_2, x_3) = \frac{c_B^\pm}{2} \left([V_1^B - A_1^B]_{(x_1, x_2, x_3)} \pm [V_1^B - A_1^B]_{(x_3, x_2, x_1)} \right)$$

$$\Pi_3^B(x_1, x_2, x_3) = c_B^- (-1)_B [T_1^B]_{(x_1, x_3, x_2)}$$

$$c_B^+ = \begin{cases} 1 & \text{if } B \neq \Lambda \\ \sqrt{\frac{2}{3}} & \text{if } B = \Lambda \end{cases}, \quad c_B^- = \begin{cases} 1 & \text{if } B \neq \Lambda \\ -\sqrt{6} & \text{if } B = \Lambda \end{cases}, \quad (-1)_B \equiv \begin{cases} +1 & \text{for } B \neq \Lambda \\ -1 & \text{for } B = \Lambda \end{cases}$$

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Example: Extrapolation for chiral-odd DAs

here: $\bar{m} = \bar{m}^{\text{phys}}$, $\delta m \propto m_s - m_l \propto m_K^2 - m_\pi^2$

$$\Phi_{\pm,i}^B = \sqrt{\frac{Z_B}{Z^*}} \left(1 + \Delta g_{\Phi\pm}^B\right) \left(\Phi_{\pm,i}^* + \delta m \Delta\Phi_{\pm,i}^B\right)$$

$$\Pi_i^B = \sqrt{\frac{Z_B}{Z^*}} \left(1 + \Delta g_\Pi^B\right) \times \begin{cases} \Phi_{+,i}^* + \delta m \Delta\Pi_i^B & , \text{ if } B \neq \Lambda \\ \Phi_{-,i}^* + \delta m \Delta\Pi_i^B & , \text{ if } B = \Lambda \end{cases}$$

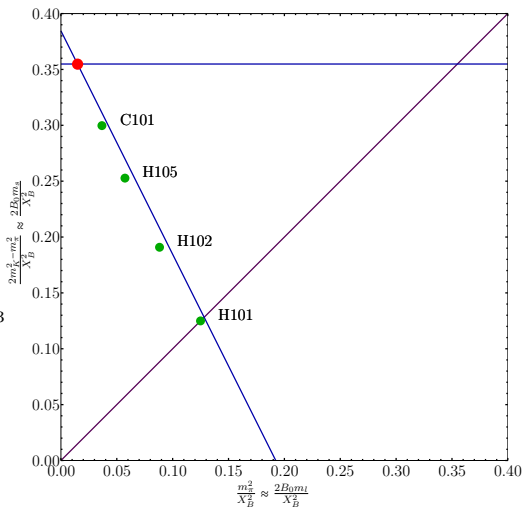
→ all results: **JHEP 1505 (2015) 073**

Results: Lattice QCD

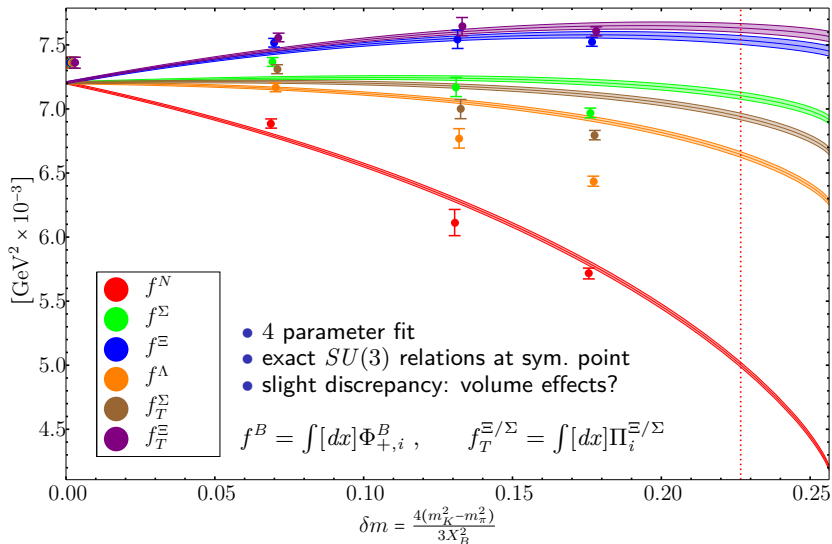
- Normalization and moments of DAs can be extracted from 2-point-correlation functions
- data points from **C**oordinated **L**attice **S**imulations effort:
 - 2+1 dynamical Wilson fermions
 - lattice spacing: $a \approx 0.086$ fm
 - volume: H: $(2.75 \text{ fm})^3$, C: $(4.13 \text{ fm})^3$

Disclaimer: preliminary!

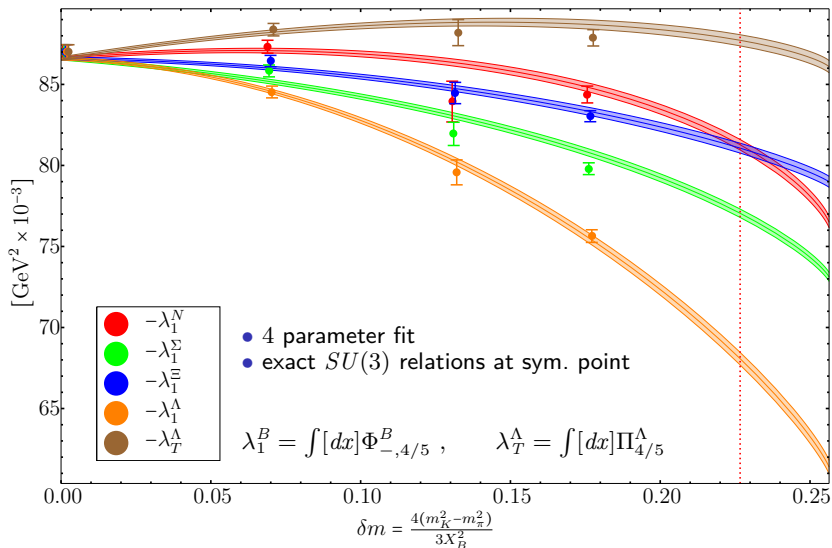
- data not renormalized yet
- purely statistical errors



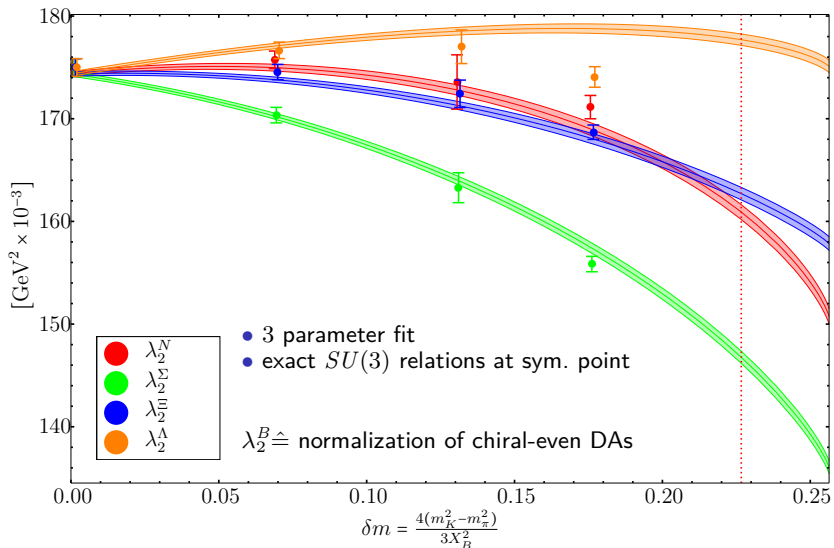
Leading twist normalization constants f^B and f_T^B (unrenormalized)



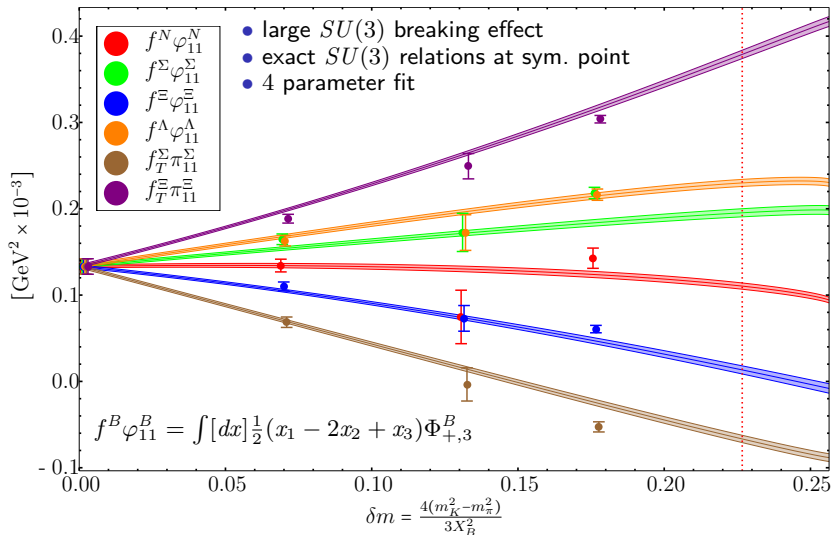
Higher twist normalization constants λ_1^B and λ_T^Λ (unrenormalized)



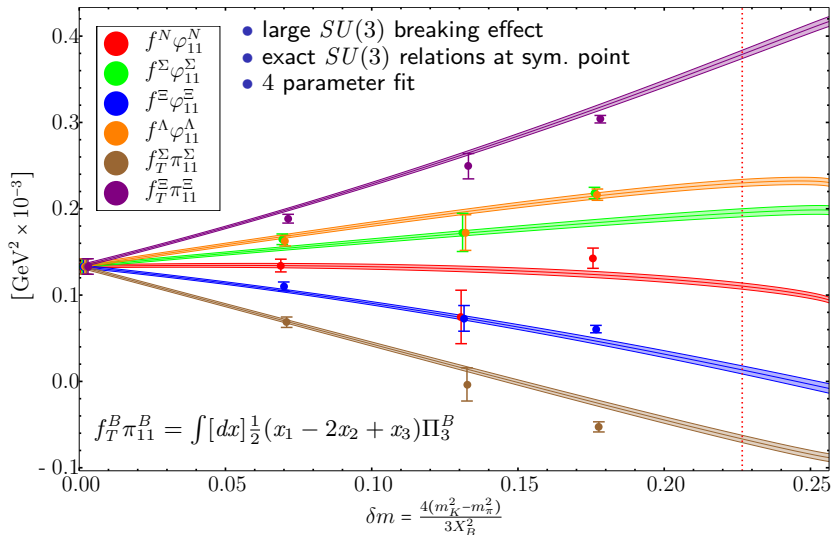
Higher twist normalization constants λ_2^B (unrenormalized)



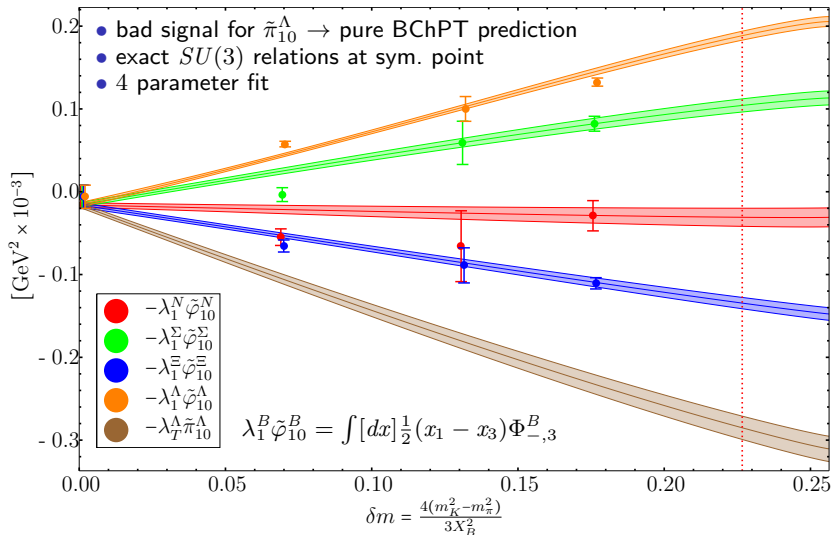
First moments of the leading twist DA φ_{11}^B and π_{11}^B (unrenormalized)



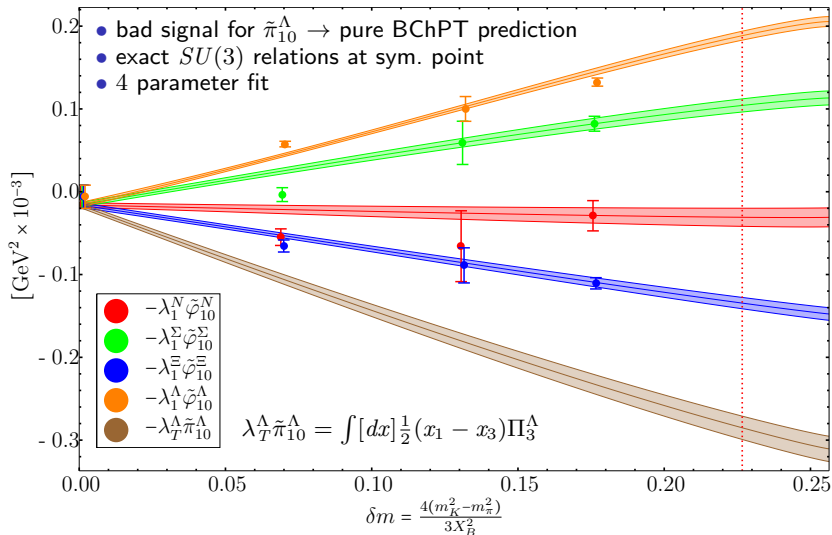
First moments of the leading twist DA φ_{11}^B and π_{11}^B (unrenormalized)



First moments of the leading twist DA $\tilde{\varphi}_{10}^B$ and $\tilde{\pi}_{10}^\Lambda$ (unrenormalized)



First moments of the leading twist DA $\tilde{\varphi}_{10}^B$ and $\tilde{\pi}_{10}^\Lambda$ (unrenormalized)



Summary

- we have performed a **NLO 3 flavor BChPT** calculation for baryon DAs
- → quark-mass dependence for all baryon octet DAs (including **higher twist**)
- we used non-local operators → **all moments** adressed simultaneously

investigation of preliminary lattice data:

- reasonable description of lattice QCD results at physical average quark mass (slight discrepancies due to volume effects?)
- normalization constants: $SU(3)$ breaking effects $\sim 15\%$
- deviations from the asymptotic shape exhibit large $SU(3)$ breaking effects