#### Recent HERMES results from inclusive and semi-inclusive hadron production with a transversely polarised target

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### Outline

- · Dihadron ( $\pi\pi$  and KK) production in TMD semiinclusive DIS on a transversely polarized proton target
- Transverse target single-spin asymmetry in inclusive electroproduction of charged pions and kaons
- Transverse polarization of Λ hyperons from quasi-real photoproduction on nuclei

# Dihadron production in semi-inclusive DIS



## Dihadron production





dihadrons h and h

guarks g and g

- new convention for FFs:
  - FFs entirely defined by quark spin  $\chi$ ,  $\chi'$
  - final-state polarisation of (di-)hadrons /  $\chi = \frac{1}{q'\chi'}$  $|l_1, m_1 >, |l_2, m_2 > \text{contained in partial-wave expansion}$
- exactly 2 FFs:
  - · unpolarised FF  $D_1$  with  $\chi = \chi'$
  - polarised (Collins) FF  $H_1^{\perp}$  with  $\chi \neq \chi$

(\*) S. Gliske, "Transverse target moments of dihadron production in semi-inclusive DIS at HERMES", PhD thesis, University of Michigan, 2011. 5

## Partial-wave expansion

 $\cdot \ \mbox{direct sum base} \mid l,m > \mbox{rather than}$ direct product base  $|l_1, m_1 \rangle, |l_2, m_2$ 

$$\frac{1}{2} \otimes \frac{1}{2} \otimes \frac{1}{2} \otimes \frac{1}{2} = \left(\frac{1}{2} \otimes \frac{1}{2}\right) \otimes \left(\frac{1}{2} \otimes \frac{1}{2}\right)$$
$$= (1 \oplus 0) \otimes (1 \oplus 0),$$
$$= 2 \oplus \underline{1} \oplus \underline{1} \oplus \underline{1} \oplus \underline{0} \oplus \underline{0}.$$
experimentally  $2 \oplus \underline{1} \oplus \underline{1} \oplus \underline{0} \oplus \underline{0}.$ 

$$|\ell_1, m_1\rangle$$
  $h' |\ell_2, m_2\rangle$   
 $q \chi$   $q' \chi'$ 

h

2 🕀 1 experimentally  $\oplus$ 

partial wave

$$D_{1} = \sum_{\ell=1}^{\infty} \sum_{m=-\ell}^{\ell} P_{\ell,m}(\cos\vartheta) e^{im(\phi_{R}-\phi_{k})} D_{1}^{|\ell,m\rangle}(z, M_{h}, |\boldsymbol{k}_{T}|),$$
  
$$H_{1}^{\perp} = \sum_{\ell=1}^{\infty} \sum_{m=-\ell}^{\ell} P_{\ell,m}(\cos\vartheta) e^{im(\phi_{R}-\phi_{k})} H_{1}^{\perp|\ell,m\rangle}(z, M_{h}, |\boldsymbol{k}_{T}|)$$

$$\begin{aligned} d\sigma_{UT} &= \frac{\alpha^2 M_h P_{h\perp}}{2\pi x y Q^2} \left( 1 + \frac{\gamma^2}{2x} \right) |\mathbf{S}_{\perp}| \\ &\times \sum_{\ell=0}^2 \sum_{m=-\ell}^{\ell} \left\{ A(x, y) \left[ P_{\ell,m} \sin((m+1)\phi_h - m\phi_R - \phi_S)) \right. \\ &\times \left( F_{UT,T}^{P_{\ell,m} \sin((m+1)\phi_h - m\phi_R - \phi_S)} + \epsilon F_{UT,L}^{P_{\ell,m} \sin((m+1)\phi_h - m\phi_R - \phi_S)} \right) \right] \\ &+ B(x, y) \left[ P_{\ell,m} \sin((1-m)\phi_h + m\phi_R + \phi_S) F_{UT}^{P_{\ell,m} \sin((1-m)\phi_h + m\phi_R + \phi_S)} \right. \\ &+ P_{\ell,m} \sin((3-m)\phi_h + m\phi_R - \phi_S) F_{UT}^{P_{\ell,m} \sin((-m\phi_h + m\phi_R - \phi_S))} \right] \\ &+ V(x, y) \left[ P_{\ell,m} \sin((-m\phi_h + m\phi_R + \phi_S) F_{UT}^{P_{\ell,m} \sin((-m\phi_h + m\phi_R - \phi_S))} \right] \\ &+ P_{\ell,m} \sin((2-m)\phi_h + m\phi_R - \phi_S) F_{UT}^{P_{\ell,m} \sin((2-m)\phi_h + m\phi_R - \phi_S)} \right] \right\}. \end{aligned}$$

and analogously for  $d\sigma_{UU}, d\sigma_{UL}, d\sigma_{LU}, d\sigma_{LL}, d\sigma_{LT}$ 

#### Structure functions at leading twist

$$\begin{split} F_{UT,L}^{P_{\ell,m}\sin((m+1)\phi_h - m\phi_R - \phi_S)} &= 0 \\ F_{UT,T}^{P_{\ell,m}\sin((m+1)\phi_h - m\phi_R - \phi_S)} &= -\mathcal{I}\bigg[\frac{|\boldsymbol{p}_T|}{M}\cos\left((m+1)\phi_h - \phi_p - m\phi_k\right) \\ & \text{"Sivers"} \\ & \times \left(f_{1T}^{\perp}D_1^{|\ell,m\rangle +} + \operatorname{signum}[m]g_{1T}D_1^{|\ell,m\rangle -}\right)\bigg], \\ F_{UT}^{P_{\ell,m}\sin((1-m)\phi_h + m\phi_R + \phi_S)} &= -\mathcal{I}\bigg[\frac{|\boldsymbol{k}_T|}{M_h}\cos\left((m-1)\phi_h - \phi_p - m\phi_k\right)h_1H_1^{\perp|\ell,m\rangle}\bigg], \end{split}$$

$$F_{UT}^{P_{\ell,m}\sin((3-m)\phi_h+m\phi_R-\phi_S)} = \mathcal{I}\left[\frac{|\boldsymbol{p}_T|^2|\boldsymbol{\kappa}_T|}{M^2M_h}\cos\left((m-3)\phi_h+2\phi_p-(m-1)\phi_k\right)\right] \times h_{1T}^{\perp}H_1^{\perp|\ell,m\rangle}.$$

usual IFF related to  $H_1^{\perp|1,1>}$ 

 $\vec{p}_T, \phi_p$  struck quark  $\vec{k}_T, \phi_k$  fragmenting quark

### Results

- Collins moments for + + 0 = 0
  - $\pi^+\pi^-, \pi^+\pi^0, \pi^-\pi^0$
- Collins and Sivers moments for  $K^+K^-$  in  $\phi$  resonance region
- Collins, Sivers and pretzelocity for |0,0> moments for  $K^+K^-$  outside  $\phi$  resonance region since l>0,m>0 are zero (as expected)

### |1, 1> Collins moments for $\pi\pi$



allows collinear access to transversity

#### |1, 1 > Collins moments for $\pi\pi$



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#### $|2,\pm 2>$ Collins moments for $\pi\pi$ $|2,\pm 2>=|1,\pm 1>|1\pm 1>$



### Collins moments for $K K in \phi$ resonance region

sensitive to transversity s-quark distribution



no indication for different signal in and outside  $\phi$ -resonance region

### Sivers moments for $K^{+}K^{-}$ in $\phi$ resonance region

sensitive to Sivers s-quark distribution



no indication for different signal in and outside  $\phi$ -resonance region

### Moments for K K outside $\phi$ resonance region @ leading twist



consistent with small positive value

consistent with small positive value

### Moments for K K outside $\phi$ resonance region @ leading twist



consistent with zero

#### Moments for K K outside $\phi$ resonance region @ sub-leading twist



consistent with zero

### Aut inclusive



# Transverse target single-spin asymmetry in inclusive electroproduction of pions and kaons

• various polarized pp scattering experiments consistently observe since 35 years large A asymmetries, with  $\sqrt{s}$  from 5 to 200 GeV



not interpretable in leading-twist based on collinear factorisation

# Transverse target single-spin asymmetry in inclusive electroproduction of pions and kaons

• various polarized pp scattering experiments consistently observe since 35 years large A asymmetries, with  $\sqrt{s}$  from 5 to 200 GeV



- not interpretable in leading-twist based on collinear factorisation
- HERMES measurement of inclusive transverse target spin asymmetry  $A_{UT}^{\sin(\psi)}$  :

$$d\sigma = d\sigma_{UU} [1 + s_{\perp} A_{UT}^{\sin(\psi)} \sin(\psi)]$$

•  $A_{UT}^{\sin(\psi)} = \frac{\pi}{2}A_N$ 

Left

• **@HERMES**  $\sin(\psi) \sim \sin(\phi - \phi_S)$ 



## Results: xF dependence

 $\pi^+$ 

 $x_F = 2P_L/\sqrt{s}$ 



 compatible with zero, with small variations over x<sub>F</sub>

### Results: disentangle $x_F$ and $P_T$ dependence



 $\pi^{\dashv}$ 

- increase with  $P_{T}$  up to  $P_{T} \approx 0.8 \text{ GeV}$
- $P_{T}$  dependence independent of  $x_{F}$
- $\rightarrow x_{F}$  increase from  $P_{T}$  dependence

 $\pi$ 

small amplitudes,

varyingly positive and negative with  $\mathrm{P}_{_{\mathrm{T}}}$ 

• decrease with increasing  $x_{_{\rm F}}$ 

# Results: disentangle $x_F$ and $P_T$ dependence





- small amplitudes
- decrease with increasing  $x_{_{\rm F}}$



#### Contribution of various subsamples



3 sub samples:

- anti-tagged: no e<sup>±</sup> detected (mostly Q<sup>2</sup> ≈0)
- DIS with 0.2<z<0.7
- DIS with z>0.7
- anti-tagged results ~ overall results, majority of statistics
- 0.2<z<0.7 results: similar to Sivers amplitudes
- z>0.7 results: large asymmetries

# Transverse A polarization in inclusive measurement





the other inclusive SSA



atomic-mass number A

- clearly positive for light target nuclei
- consistent with zero for heavy targets

#### the other inclusive SSA



larger in backward direction w.r.t. incoming lepton

• consistent with  $x_F$  dependence of twist-3 calculation (opposite sign conventions for  $x_F$ !)

#### the other inclusive SSA



larger in backward direction w.r.t. incoming lepton

• distinct  $p_T$  dependences in forward and backward directions: rising with  $p_T$  in backward direction as in pp

## Summary

- SIDIS dihadron moments (in new partial wave expansion) provide potentially rich information on various distribution and fragmentation functions
- $\cdot$  inclusive  $A_{UT}$  provides information that can contribute to understanding of  $A_N$  in pp data
- · inclusive production of  $\Lambda$  in ep can provide complementary information to pp data on the mechanism to generate  $\Lambda$  polarization