

Exclusive Processes at Jefferson Lab at 6 GeV

for Transversity 2014
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Generalized Parton Distributions

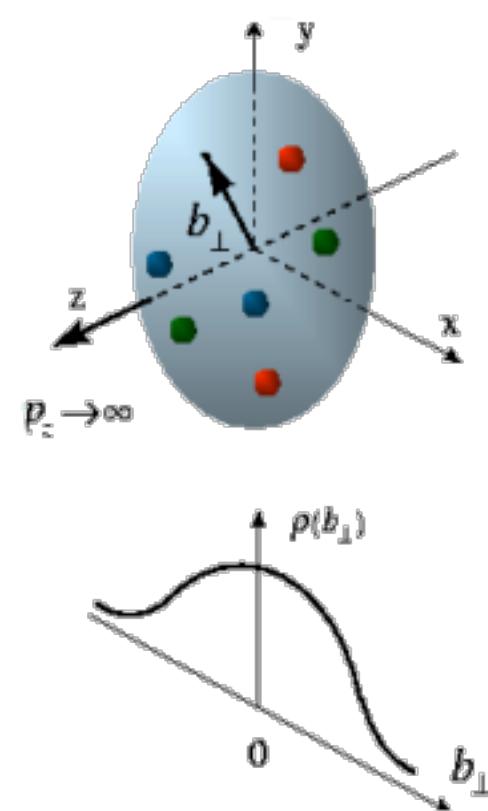
Elastic Scattering

$$\int dx \sum_q H^q(x, \xi, t) = F_1(t)$$

$$\int dx \sum_q E^q(x, \xi, t) = F_2(t)$$

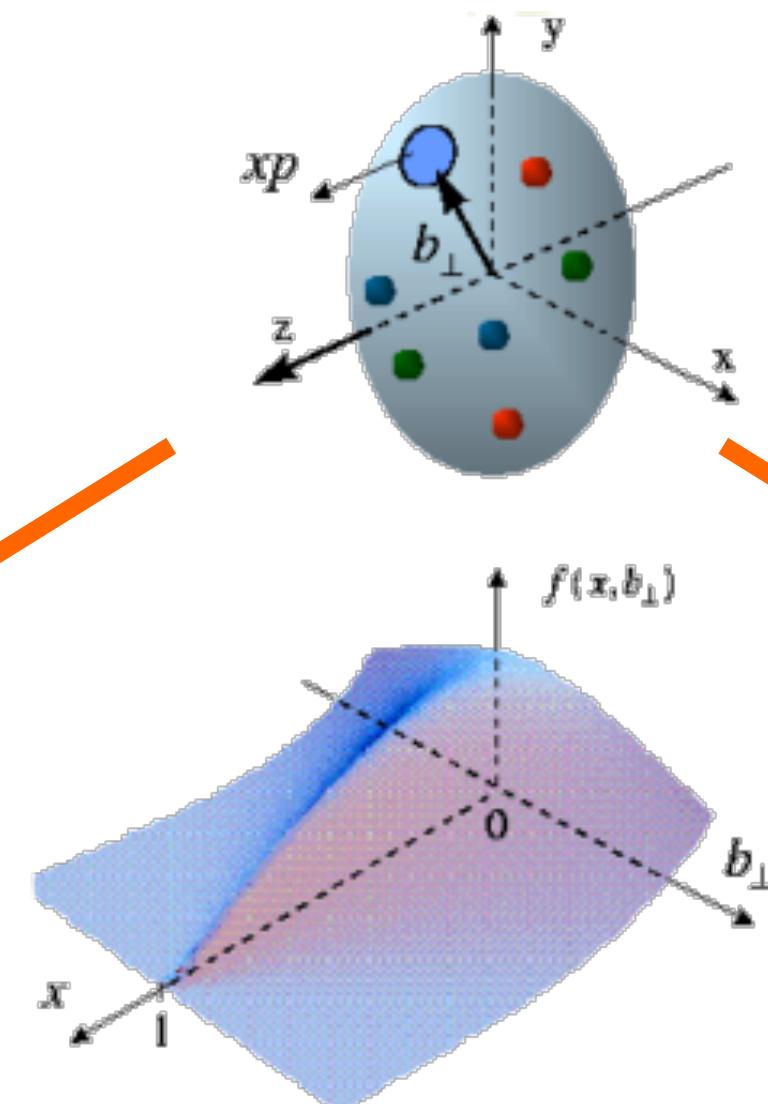
$$\int dx \tilde{H}^q(x, \xi, t) = G_A(t)$$

$$\int dx \tilde{E}^q(x, \xi, t) = G_P(t)$$



Form factors

HARD EXCLUSIVE PROCESSES

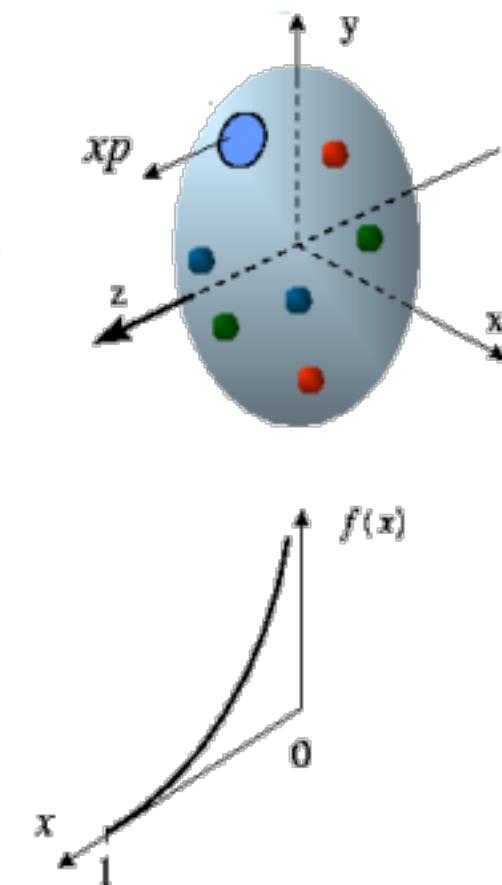


Correlation between quark longitudinal momentum and transverse spatial distributions

Deep Inelastic Scattering

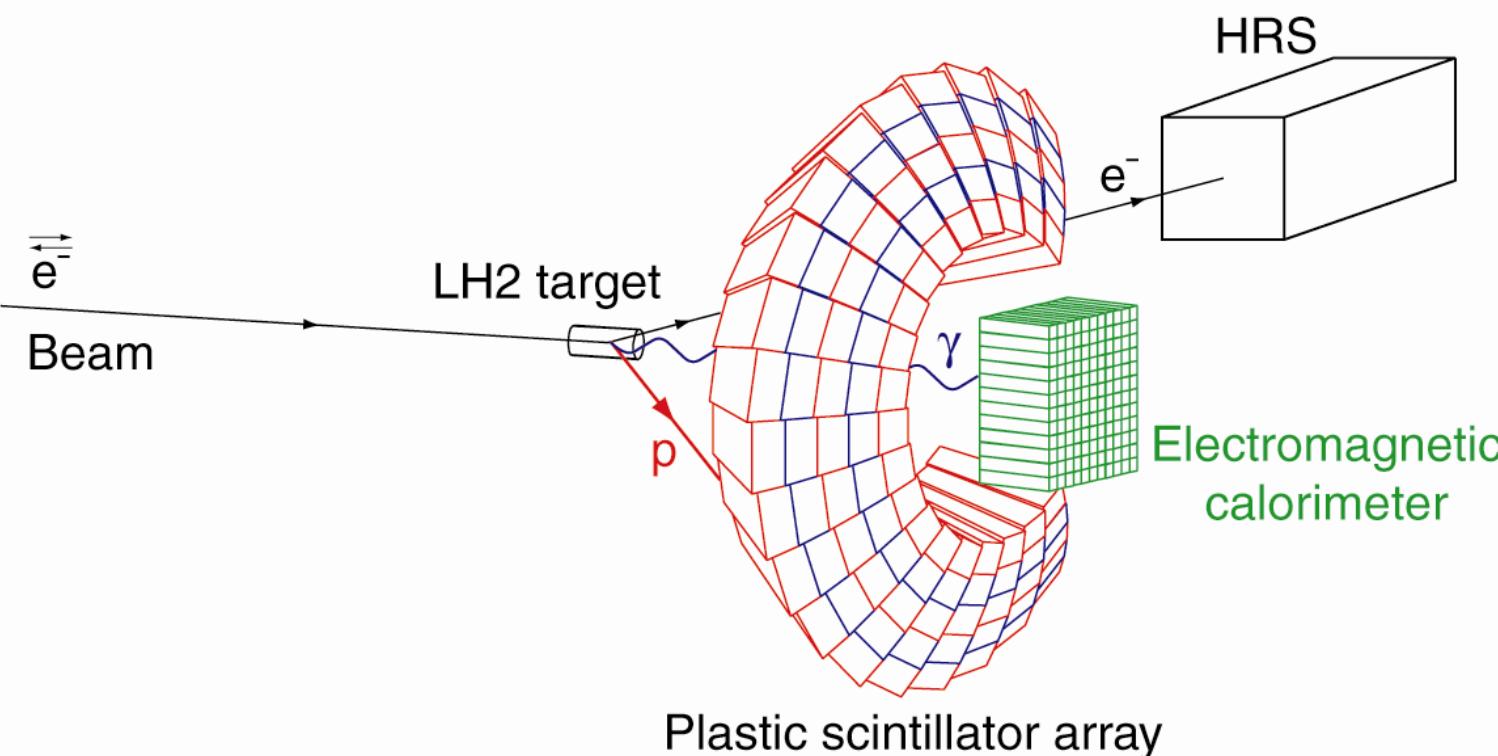
$$H^q(x, \xi = 0, t = 0) = q(x)$$

$$\tilde{H}^q(x, \xi = 0, t = 0) = \Delta q(x)$$

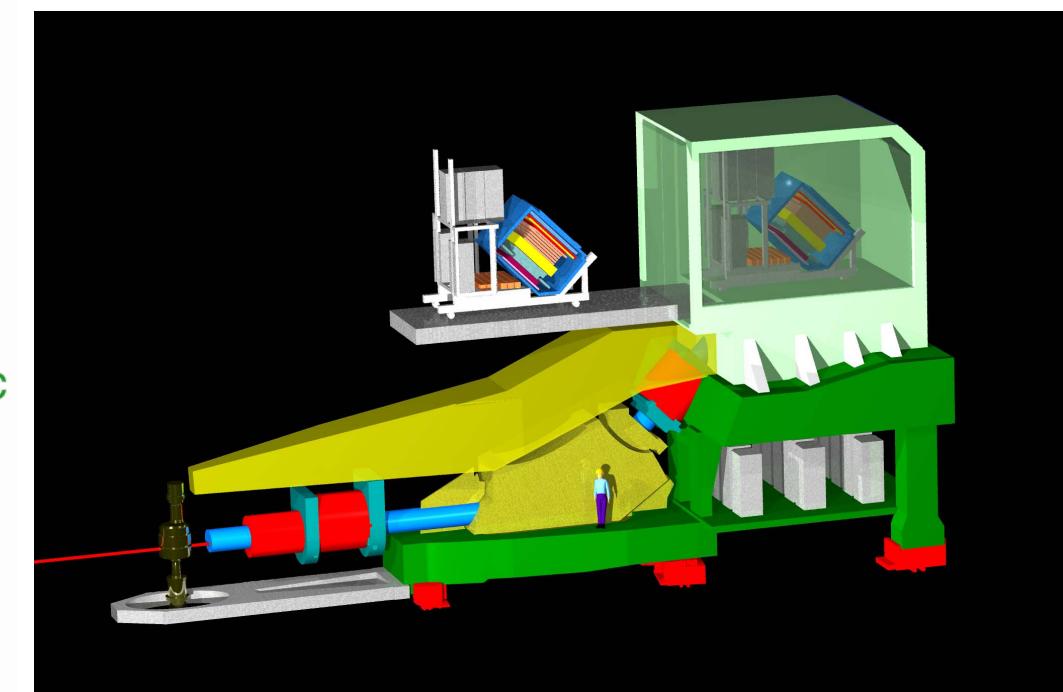


Parton distributions

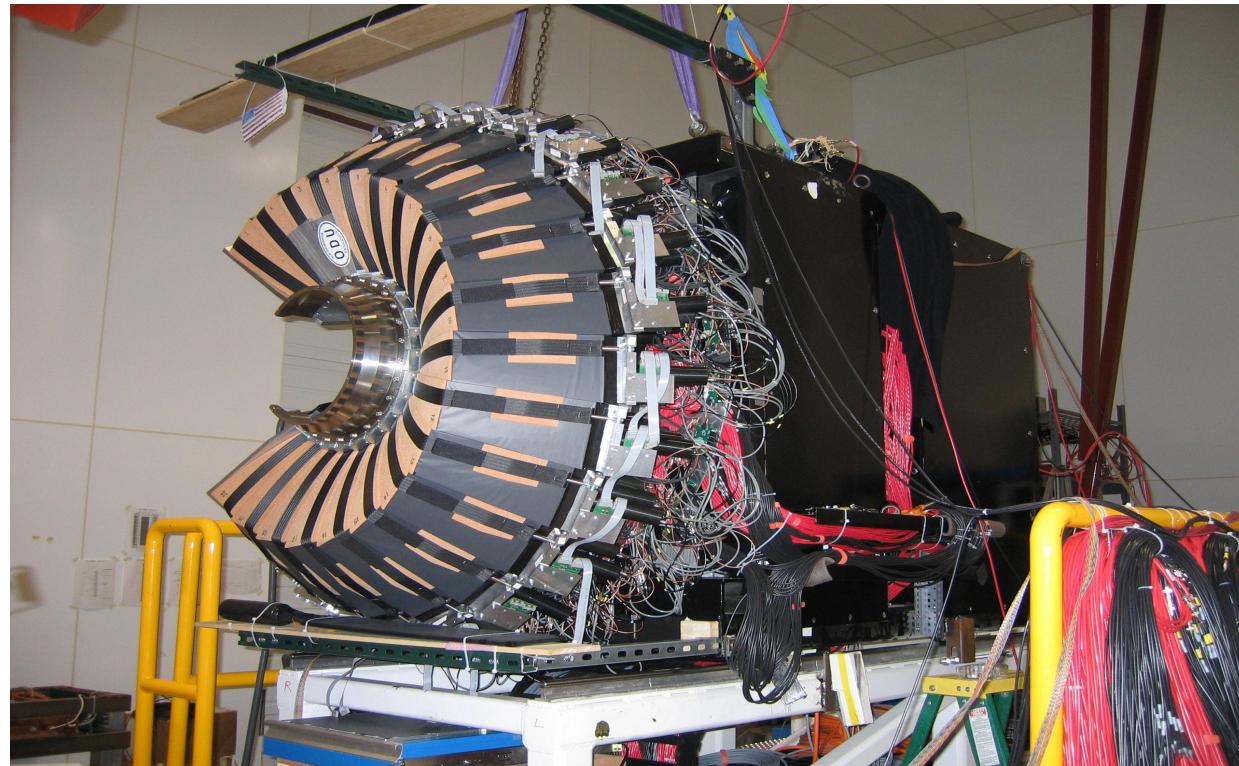
Hall A Experimental Setup at JLab



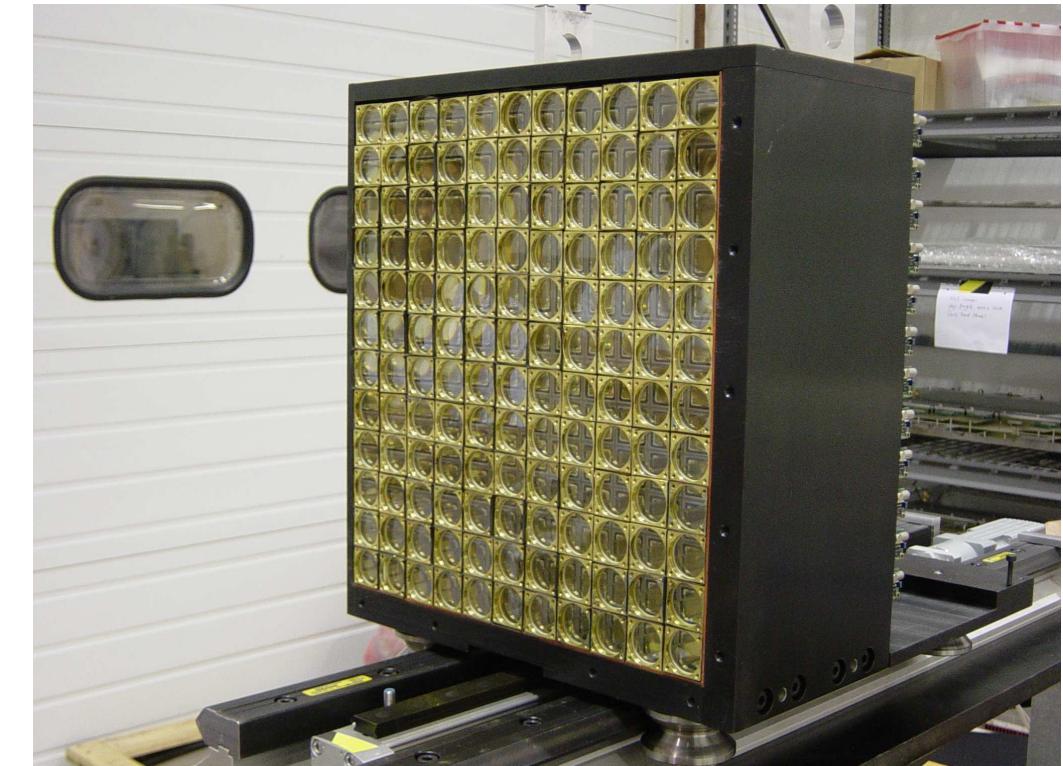
HIGH RESOLUTION SPECTROMETER



100-CHANNEL SCINTILLATOR ARRAY



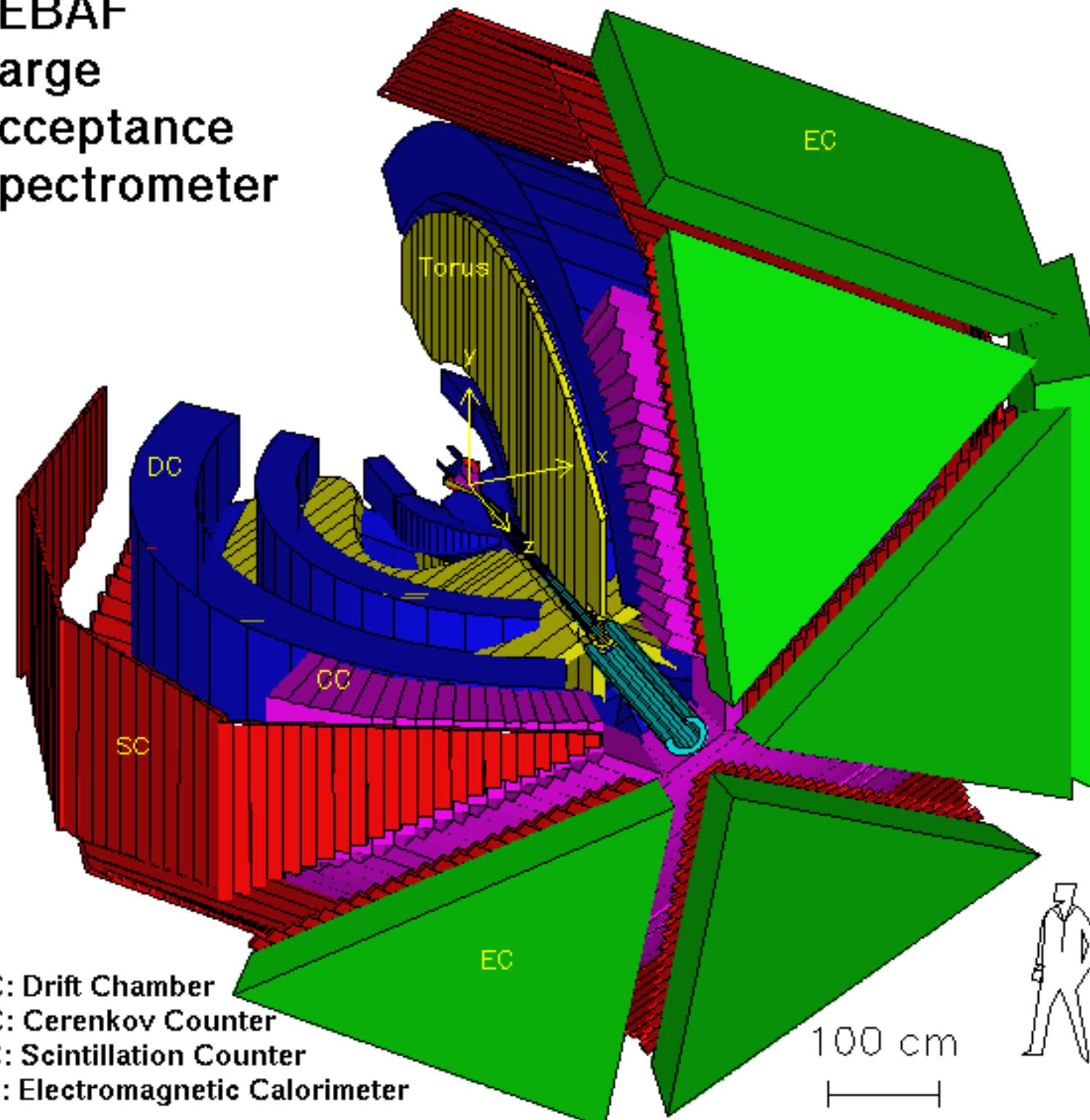
ELECTROMAGNETIC CALORIMETER



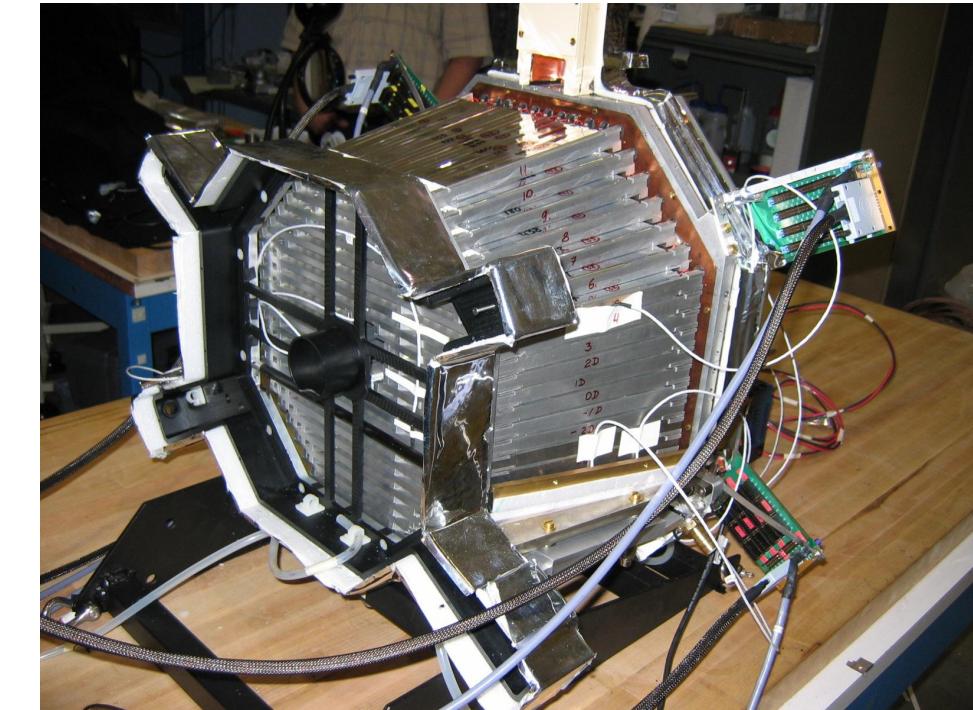


Hall B Experimental Setup at JLab

**CEBAF
Large
Acceptance
Spectrometer**

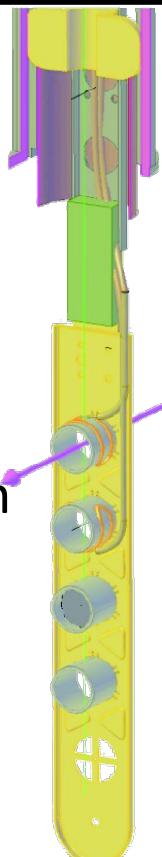


IC: Inner Calorimeter



**Longitudinally
Polarized Target**

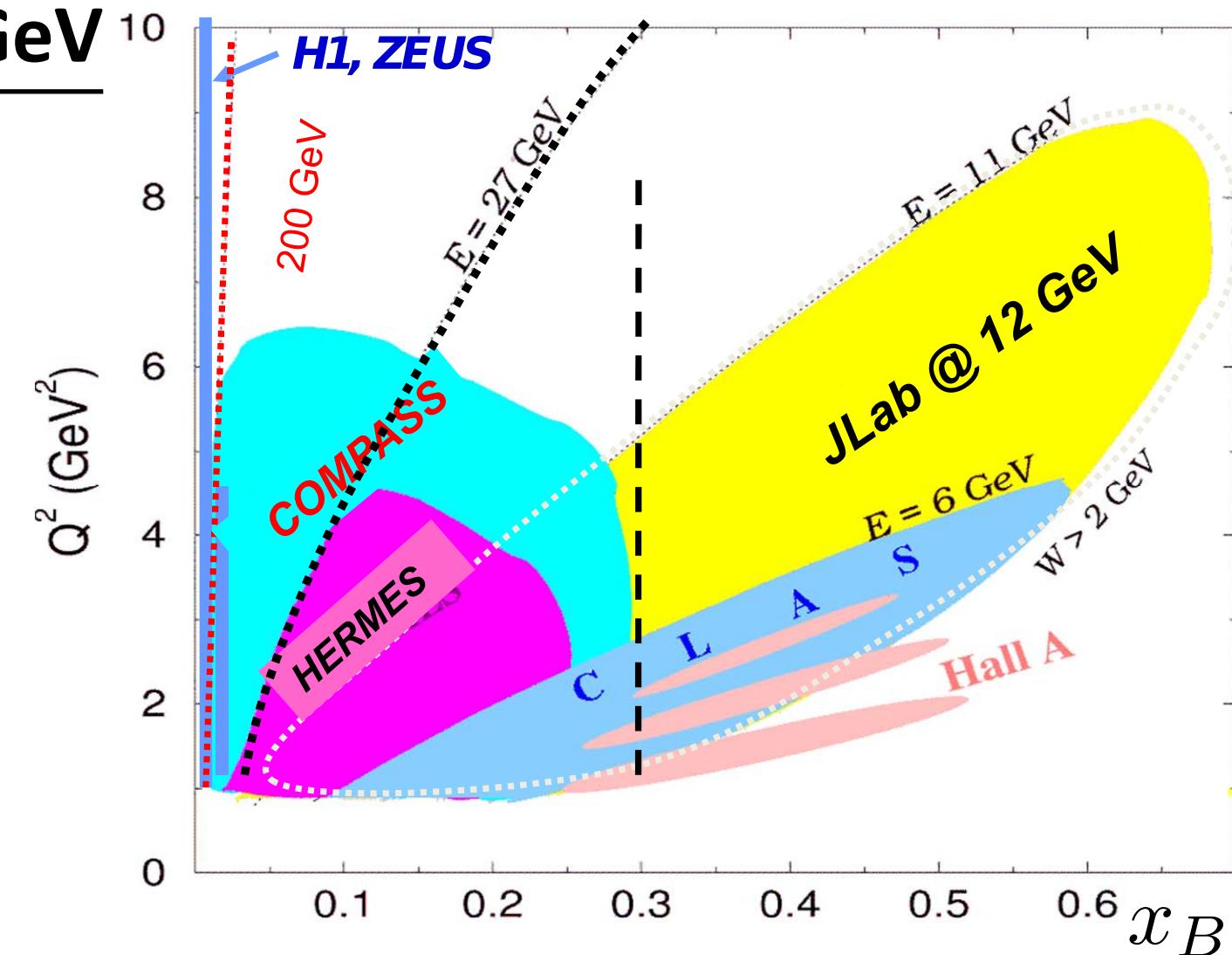
- ◆ Frozen ammonia as a target material
- ◆ polarized by Dynamic Nuclear Polarization in a 5 Tesla homogeneous magnetic field
- ◆ monitored using a Nuclear Magnetic Resonance system



Exclusive Processes at JLab at 6 GeV

♦ Hall A and Hall B (CLAS) at 6 GeV

- ♦ partially: complimentary, overlapping
- ♦ Hall A: high accuracy limited kinematics
- ♦ Hall B: wide kinematic range limited accuracy



♦ The roadmap:

- ♦ Early results (2001) from non-dedicated experiment (CLAS)
- ♦ First round of dedicated experiments in Halls A/B in 2004/2005
- ♦ Second round in 2008-2010
- ♦ Compelling exclusive program in Halls A/B at 12 GeV

DVCS Experimental Observables

DVCS:

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi_\pi}$$

unpolarized terms

$$\propto c_0^{BH} + \sum_{n=1}^2 c_n^{BH} \cos(n\phi) + s_1^{BH} \sin \phi$$

$$\propto c_0^{DVCS} + \sum_{n=1}^2 [c_n^{DVCS} \cos(n\phi) + s_n^{DVCS} \sin(n\phi)]$$

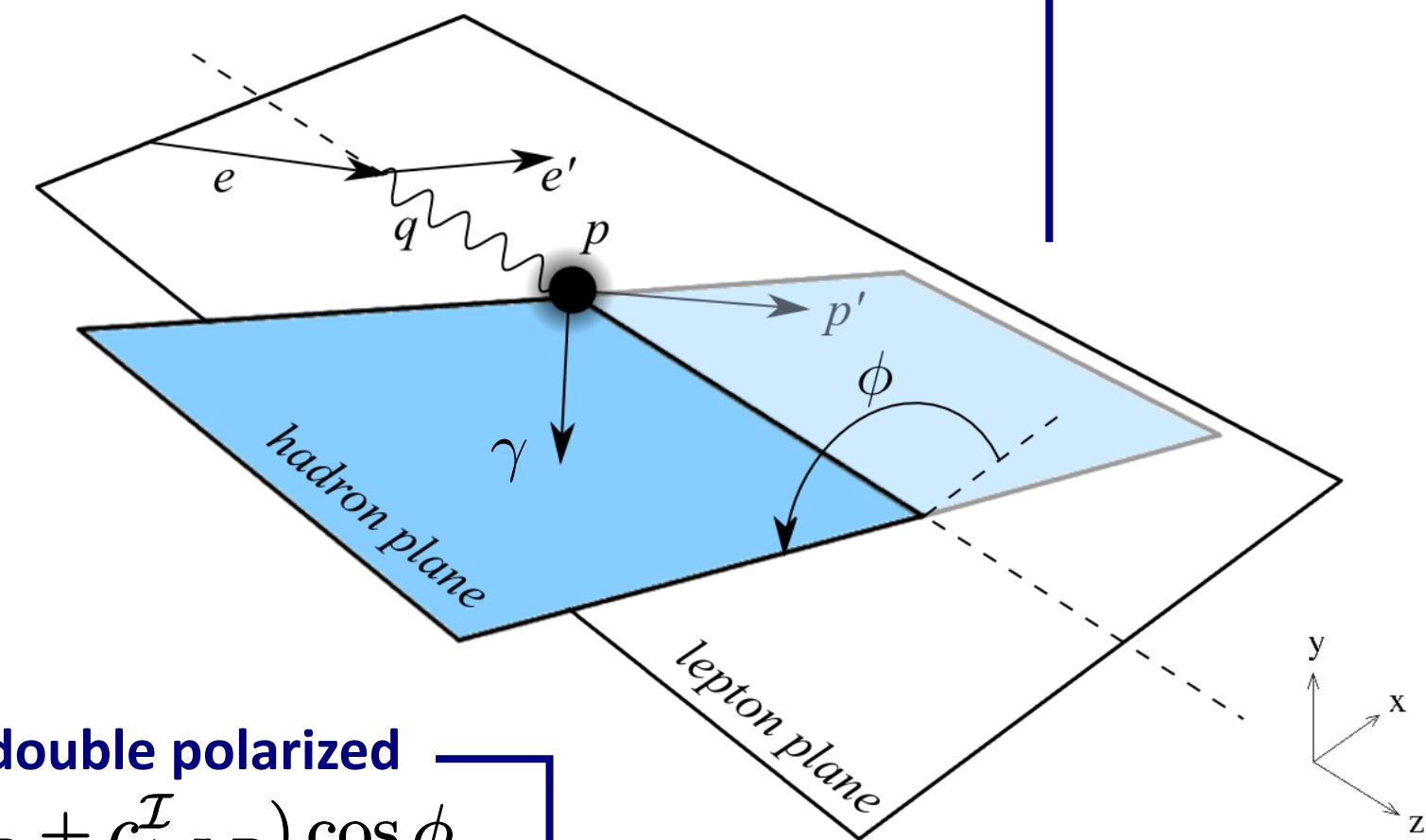
$$\propto c_0^{\mathcal{I}} + \sum_{n=1}^3 [c_n^{\mathcal{I}} \cos(n\phi) + s_n^{\mathcal{I}} \sin(n\phi)]$$

beam polarized

$$\propto s_{1,unp}^{\mathcal{I}} \sin \phi$$

target polarized

$$\propto s_{1,LP}^{\mathcal{I}} \sin \phi$$



double polarized

$$\propto c_{0,LP}^{BH} + c_{0,LP}^{\mathcal{I}} + (c_{1,LP}^{BH} + c_{1,LP}^{\mathcal{I}}) \cos \phi$$

Deeply Virtual Compton Scattering

Unpolarized cross sections, unpolarized beam and target

$$\sigma_0 \sim \text{Re} \left\{ F_1 H - \frac{t}{4M^2} F_2 E - \xi^2 (F_1 + F_2)(H + E) \right\}$$

Polarized beam and Unpolarized target (BSA)

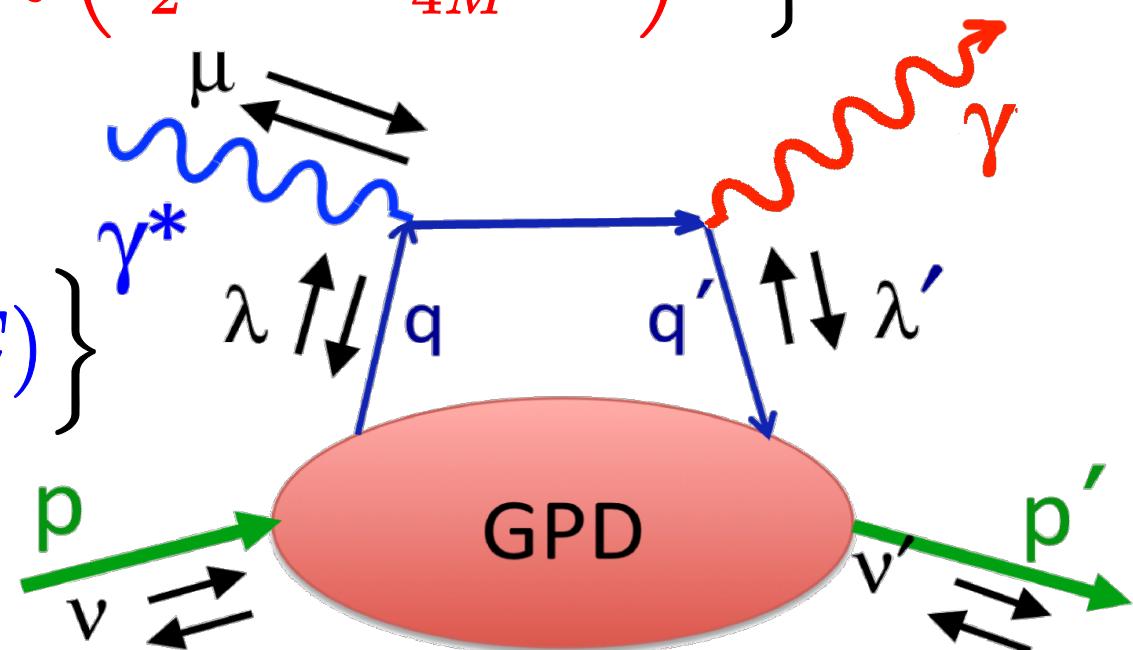
$$A_{LU} \sim \text{Im} \left\{ F_1 H + \xi(F_1 + F_2) \tilde{H} + \frac{t}{4M^2} F_2 E \right\}$$

Unpolarized beam and Polarized target (TSA)

$$A_{UL} \sim \text{Im} \left\{ F_1 \tilde{H} + \xi(F_1 + F_2) \left(H - \frac{x_B}{2} E \right) - \xi \left(\frac{x_B}{2} F_1 - \frac{t}{4M^2} F_2 \right) \tilde{E} \right\}$$

Polarized beam and Polarized target (DSA)

$$A_{LL} \sim \text{Re} \left\{ F_1 \tilde{H} + \xi(F_1 + F_2) \left(H + \frac{x_B}{2} E \right) \right\}$$



DVCS Cross Sections

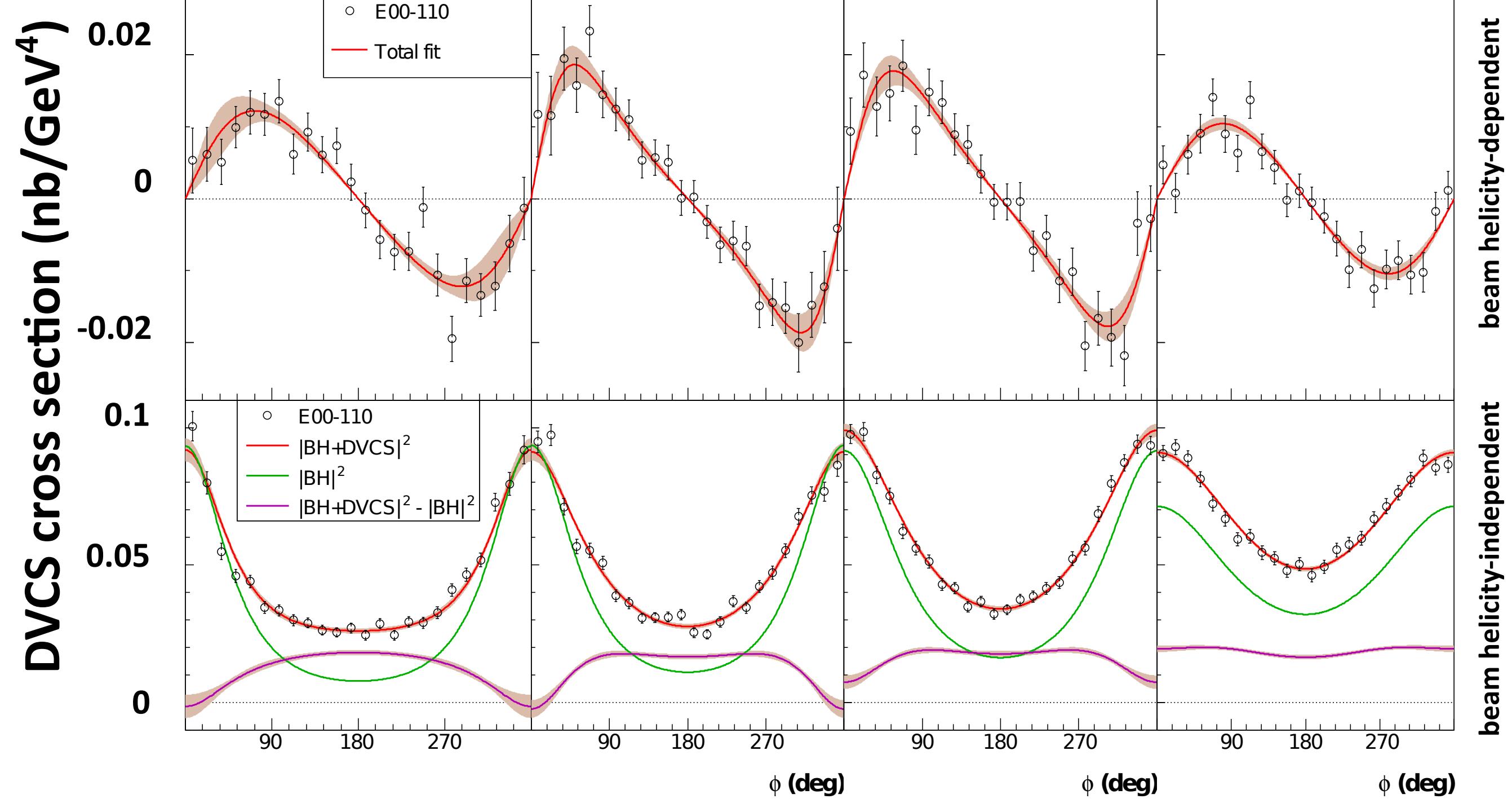
First round experiment at Hall A: E00-110

$t = -0.33 \text{ GeV}^2$

$t = -0.28 \text{ GeV}^2$

$t = -0.23 \text{ GeV}^2$

$t = -0.17 \text{ GeV}^2$



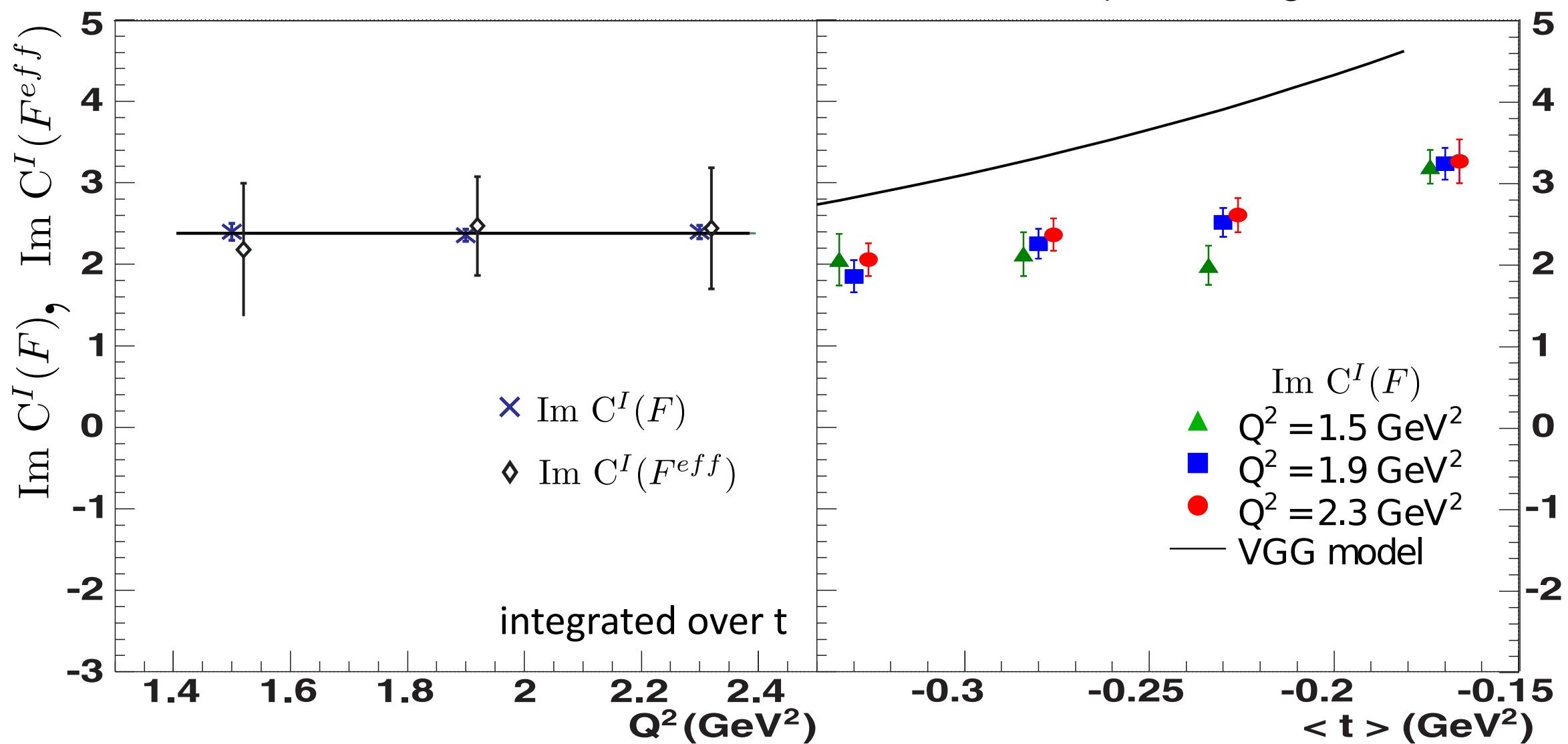
DVCS Scaling Test

First round experiment at Hall A: E00-110

♦ **Q^2 independence:**

♦ Twist-2 dominance (GPDs)

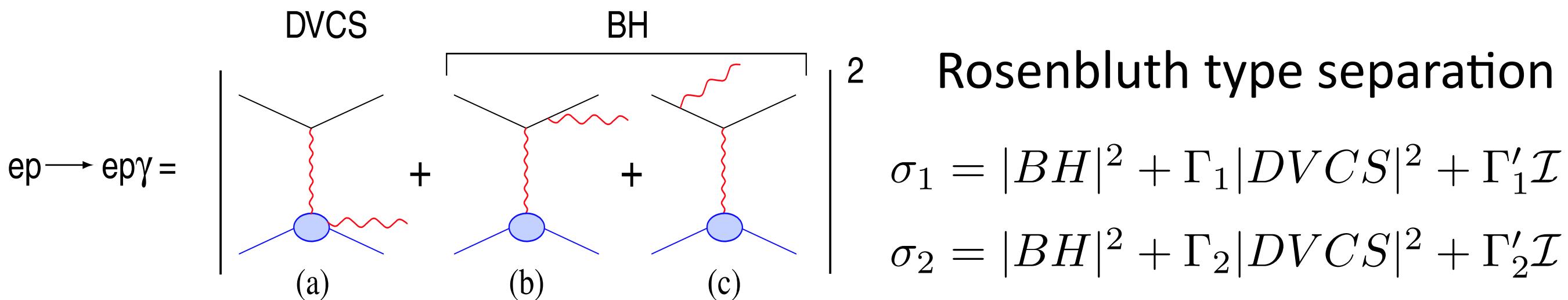
curve is from VGG model calculations
(Vanderhaeghen, Guichon, Guidal)



Separation of DVCS and Interference Terms

Second round experiment at Hall A: E07-007

$$\sigma(ep \rightarrow ep\gamma) = |BH|^2 + \underbrace{\mathcal{I}(BH \cdot DVCS)}_{\text{Linear combination of GPDs}} + \underbrace{|DVCS|^2}_{\text{Bilinear combination}}$$



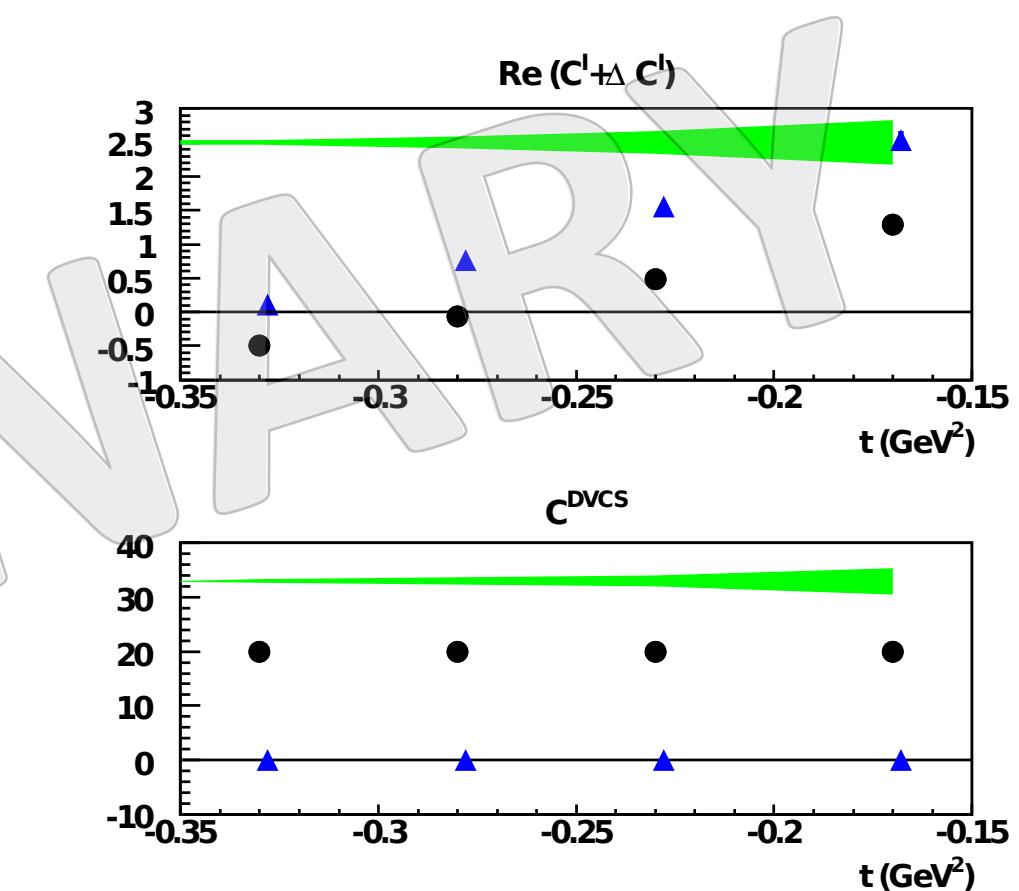
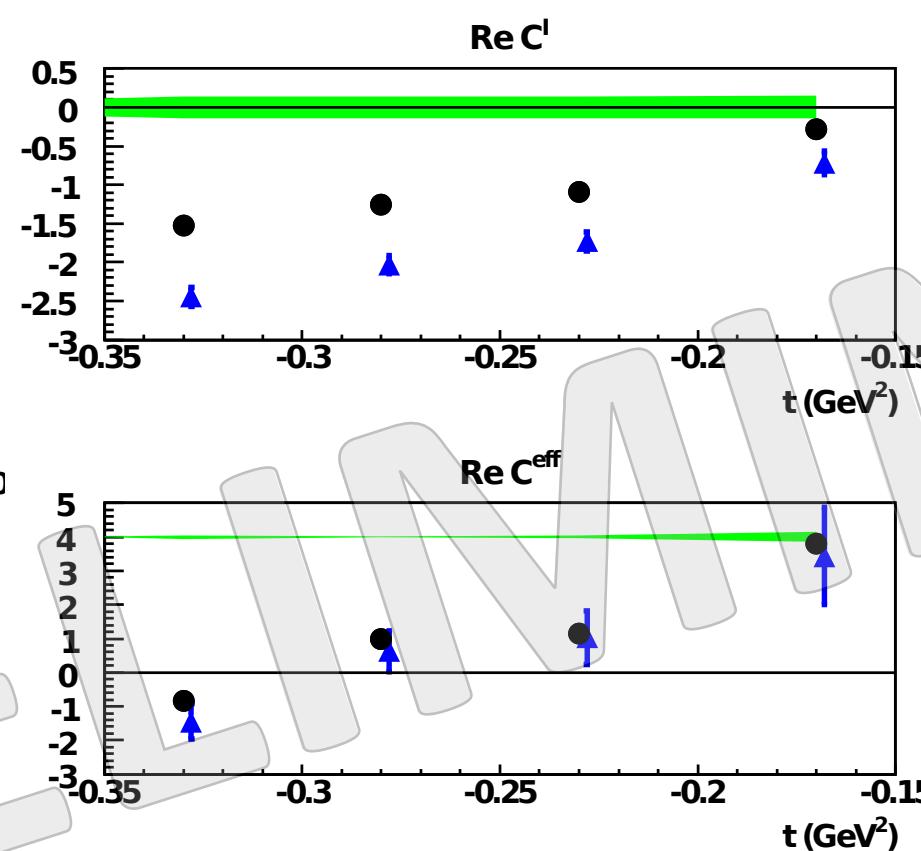
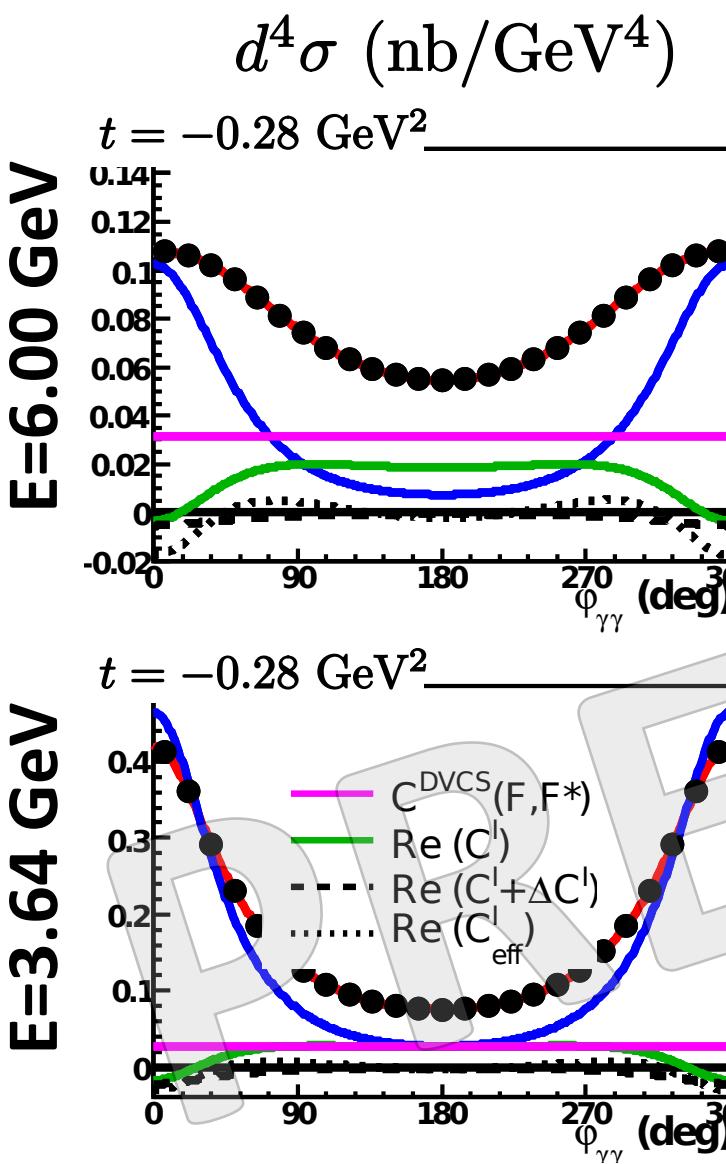
DVCS cross section has a very rich structure:

- ◆ Azimuthal analysis allows separation of different contributions
- ◆ Interference and DVCS terms mix in an azimuthal analysis
- ◆ The different energy dependence allows a full separation

Rosenbluth-like separation

Second round experiment at Hall A: E07-007

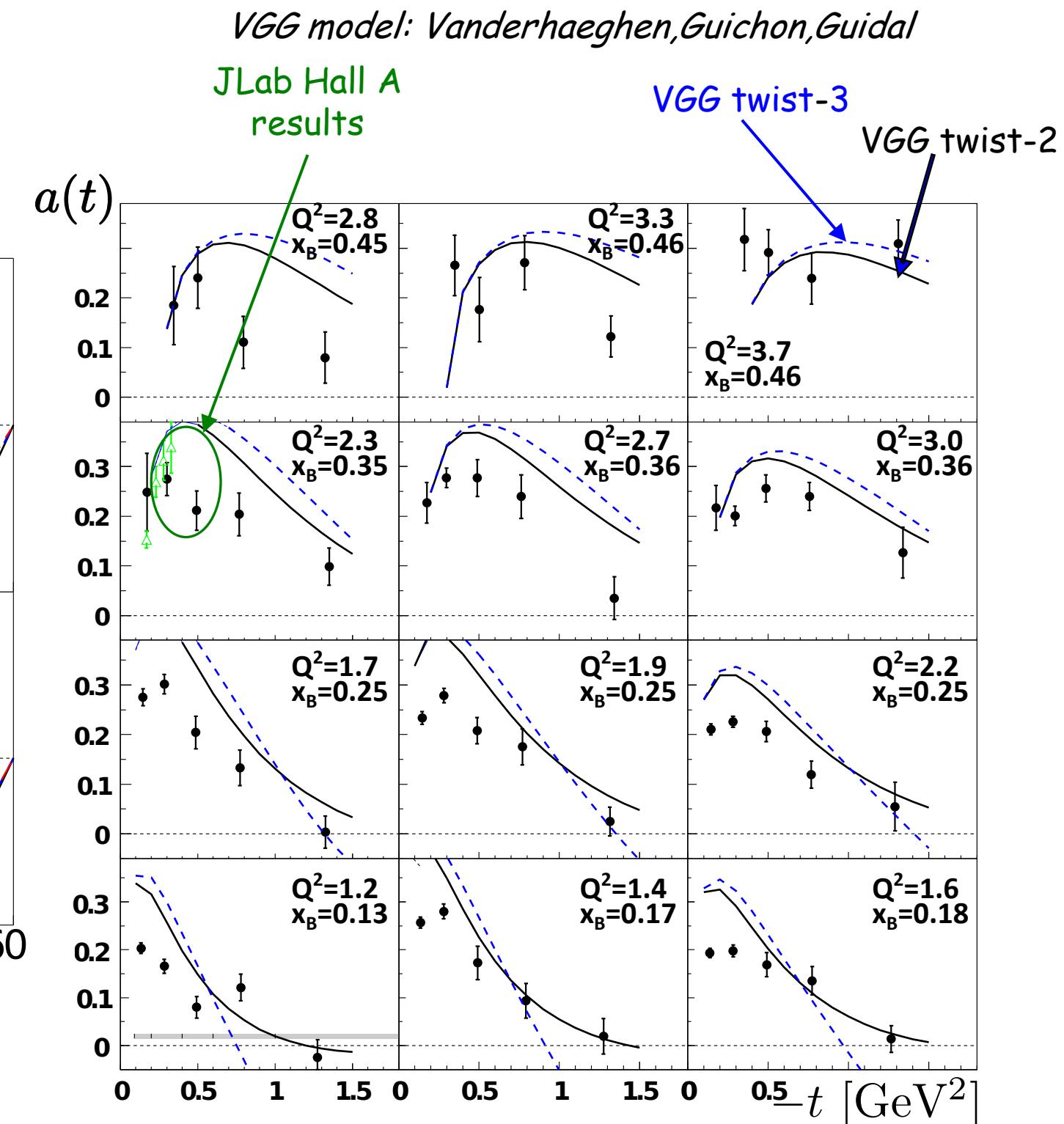
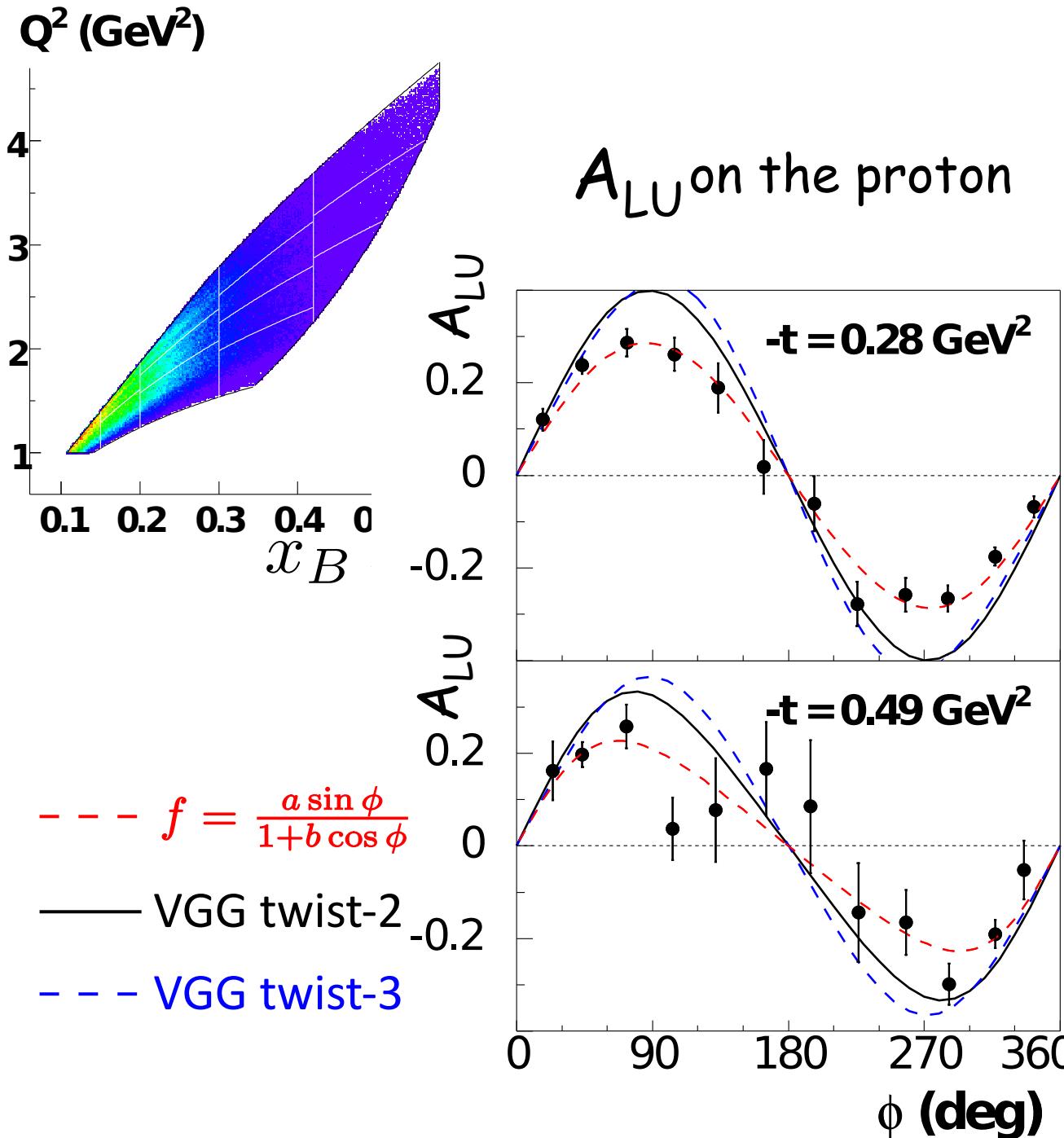
- ♦ Clean separation of BH-DVCS Interference term from pure DVCS²



- ▲ First round experiment E00-110: assuming DVCS²=0
- Second round experiment E07-007: separation of DVCS² and \mathcal{I}

DVCS Beam Spin Asymmetries

First round experiment at Hall B: E01-113

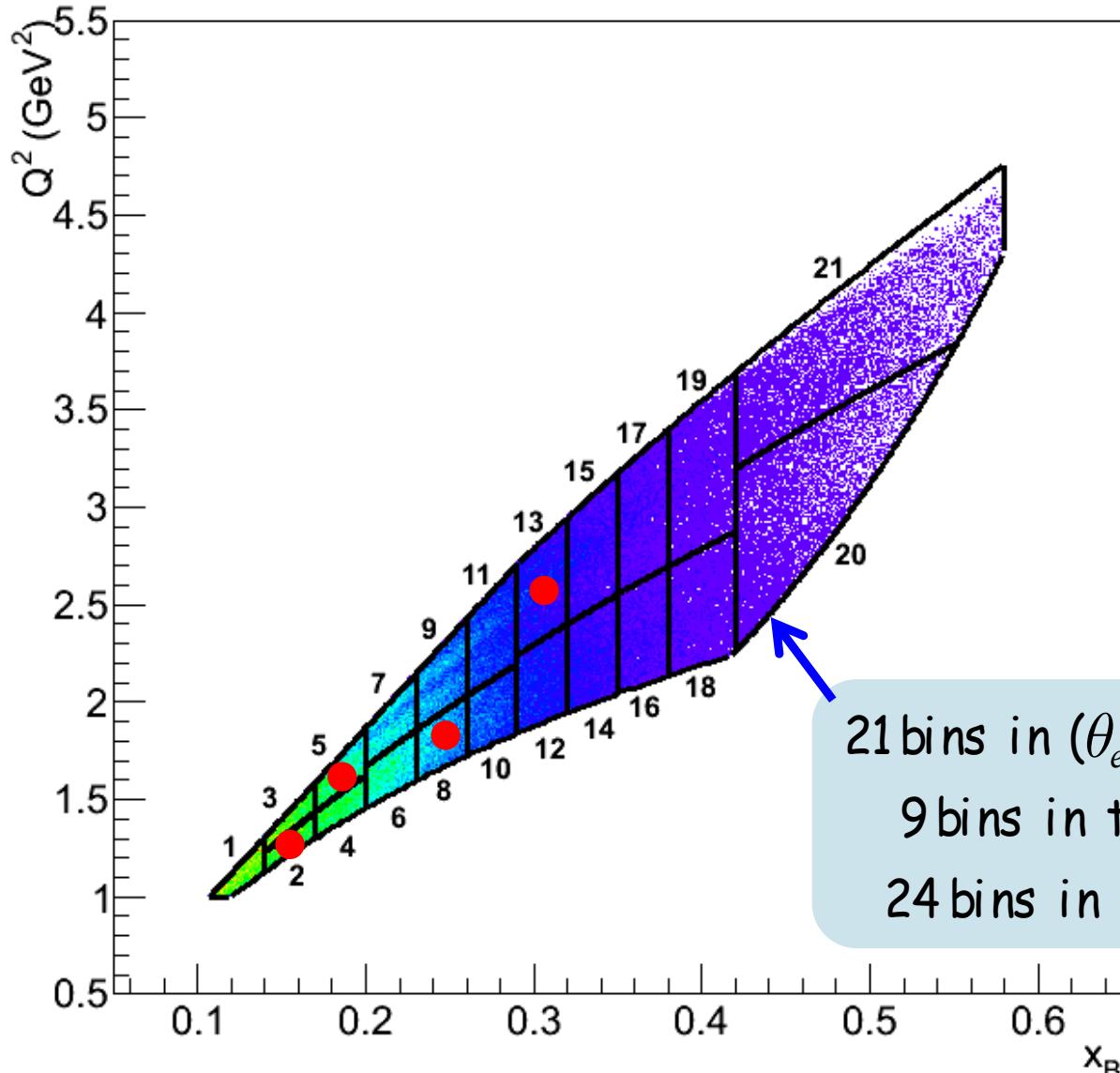


DVCS Kinematic Coverage at CLAS

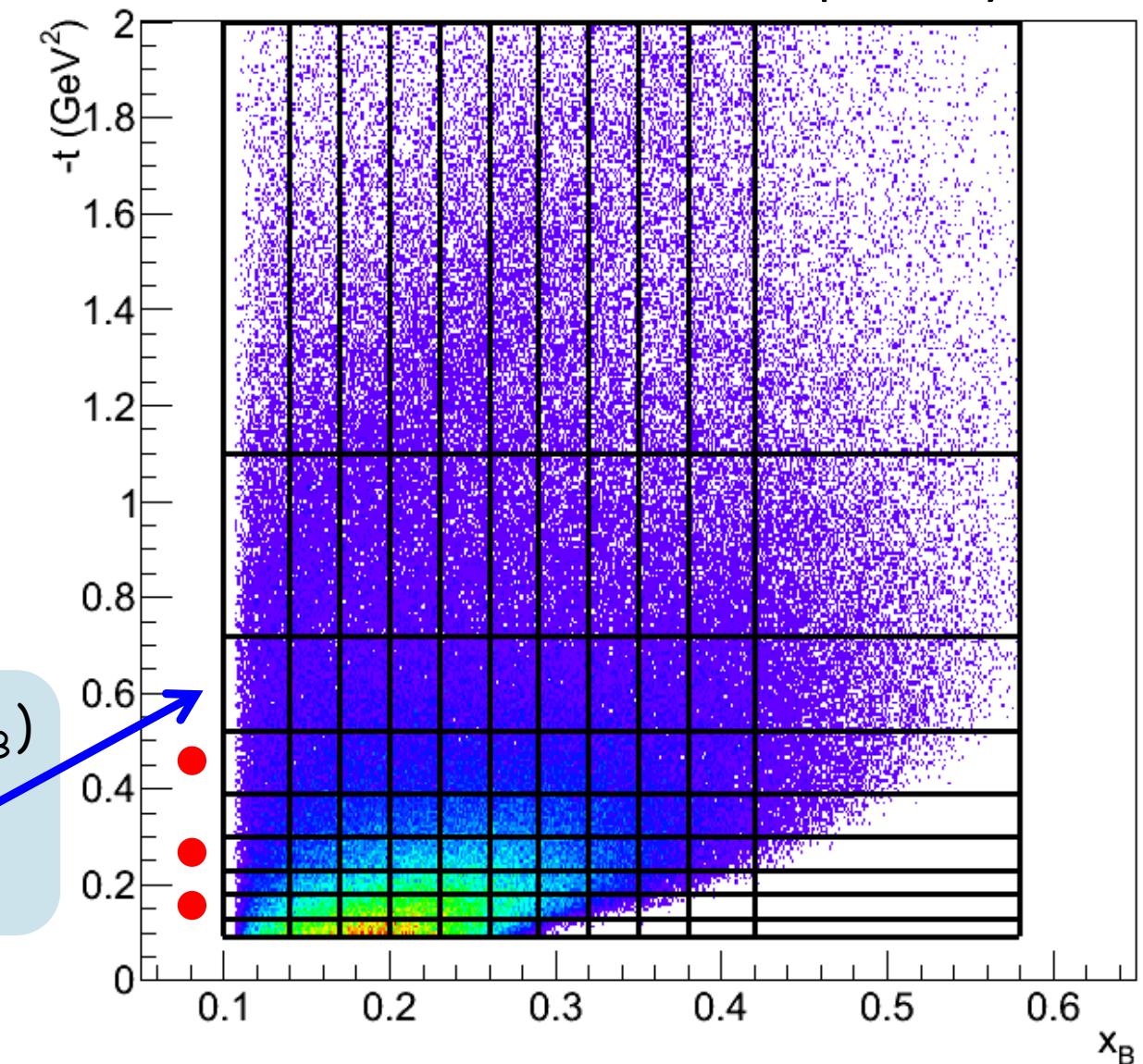
First round experiment at Hall B: E01-113

DVCS cross section measurements are underway

plots by H.S. Jo



21 bins in (θ_e, x_B)
9 bins in t
24 bins in ϕ



Extraction
of 4-fold
cross sections

$$\frac{d^4\sigma_{ep \rightarrow ep\gamma}}{dQ^2 dx_B dt d\Phi}$$

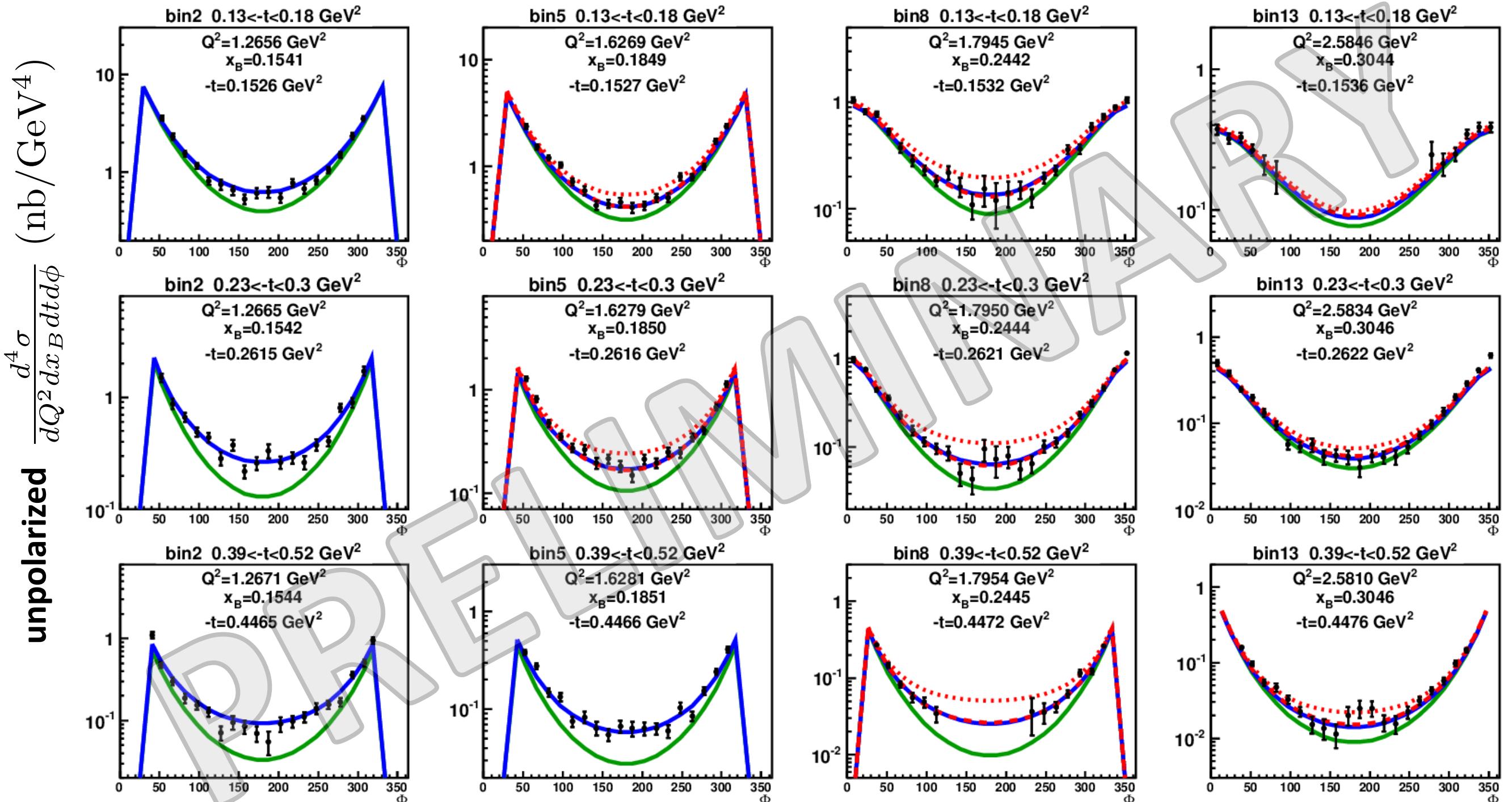
$$\frac{1}{2} \left(\frac{d^4\vec{\sigma}_{ep \rightarrow ep\gamma}}{dQ^2 dx_B dt d\Phi} - \frac{d^4\leftarrow\vec{\sigma}_{ep \rightarrow ep\gamma}}{dQ^2 dx_B dt d\Phi} \right)$$

DVCS Cross Sections

plots by H.S. Jo

First round experiment at Hall B: E01-113
Measurements in a **LARGE kinematic domain**

- BH only
- VGG (H only)
- - KM10 (Kumericki, Mueller)
- - - KM10a



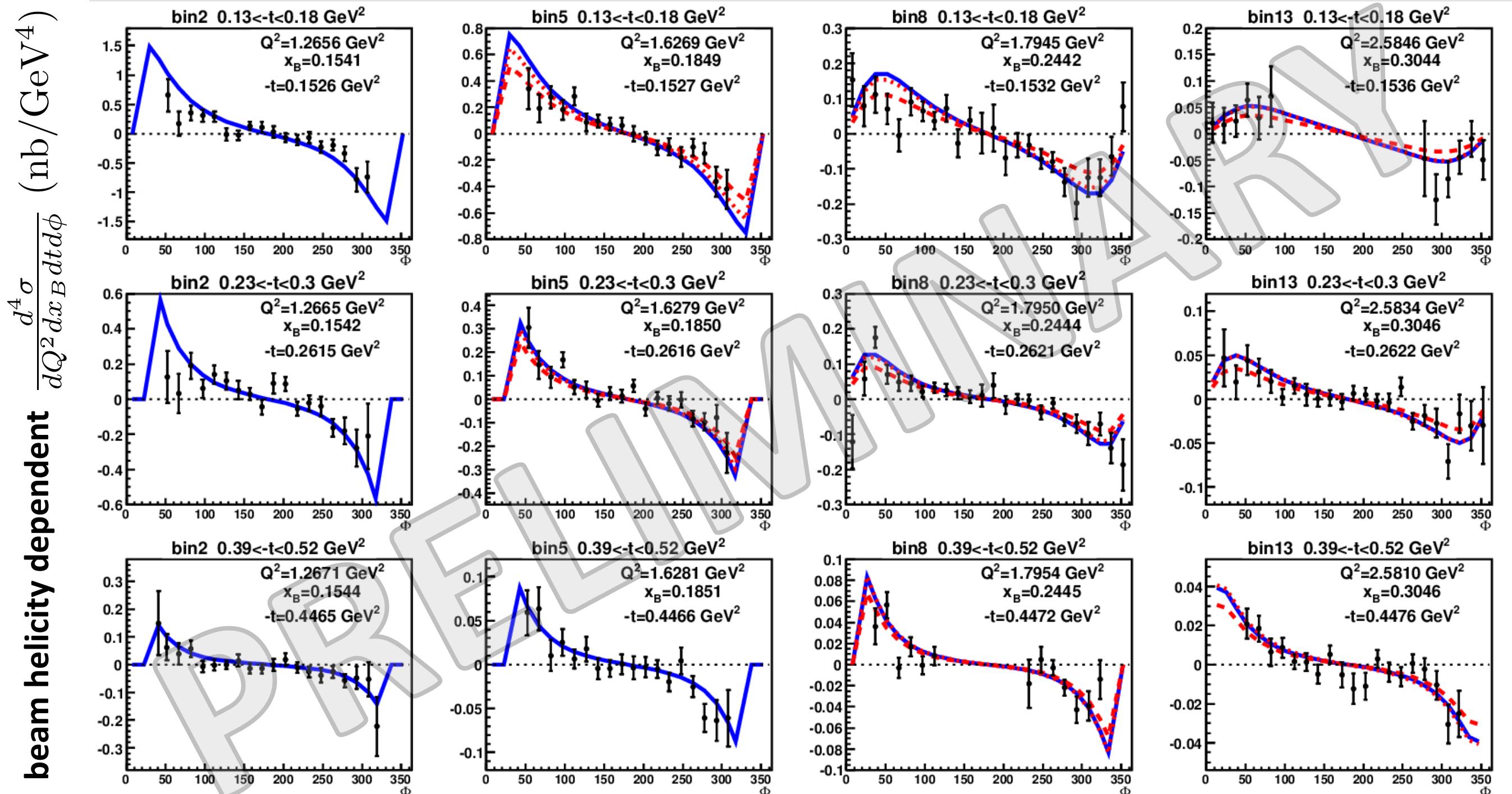
DVCS Cross Sections

plots by H.S. Jo

First round experiment at Hall B: E01-113

Measurements in a **LARGE** kinematic domain

- BH only
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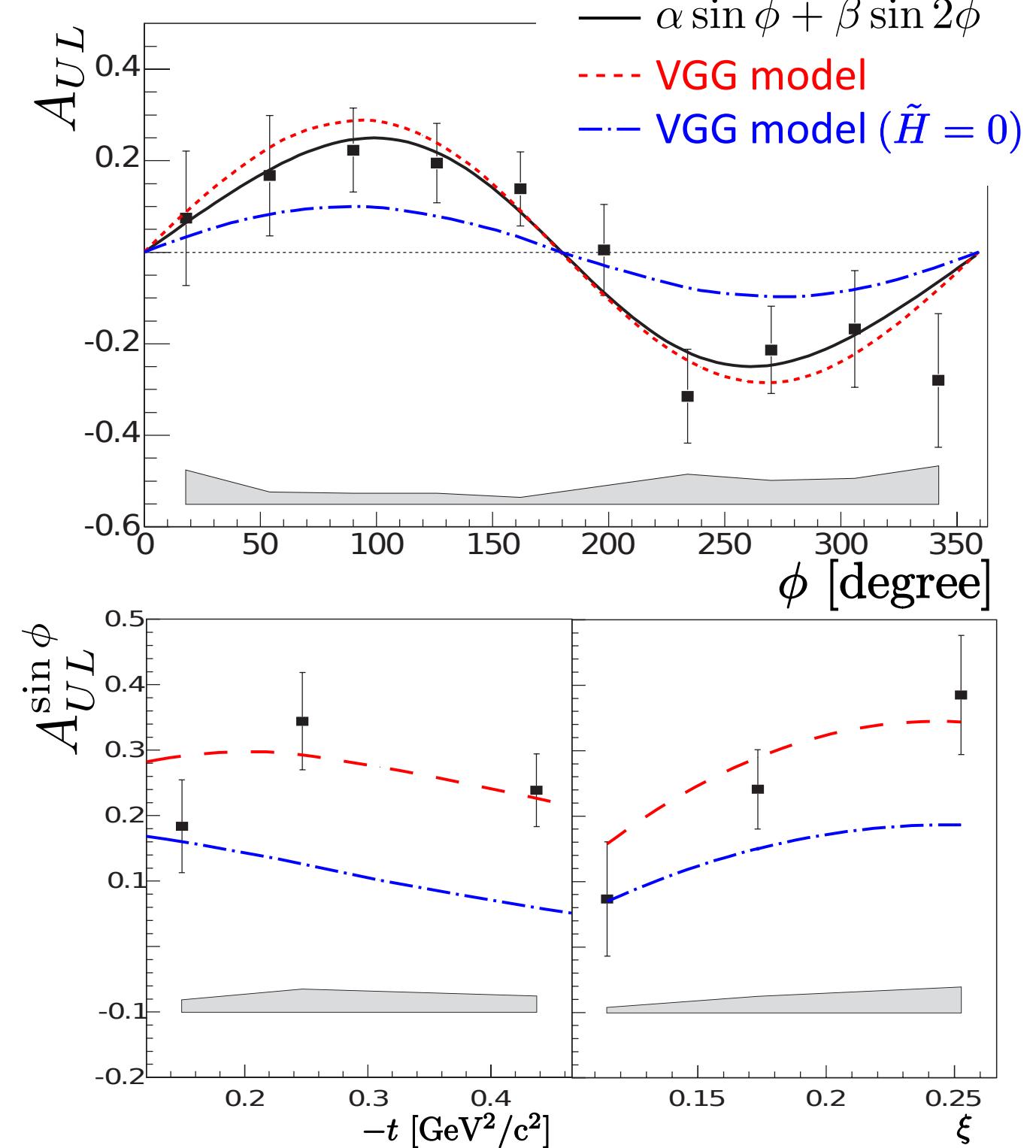
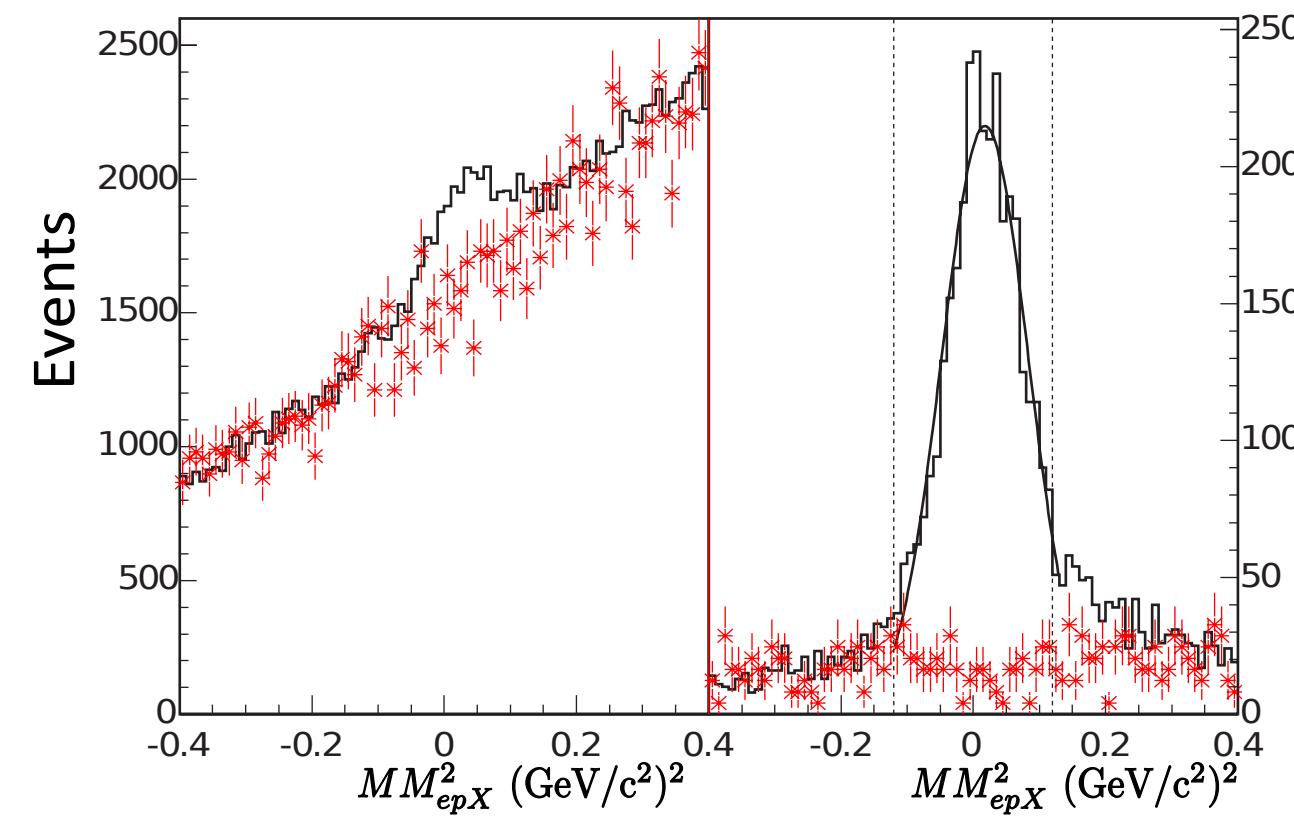
DVCS Target Spin Asymmetries

Non-dedicated experiment at Hall B (longitudinally polarized target)

First measurements of DVCS target spin asymmetry:

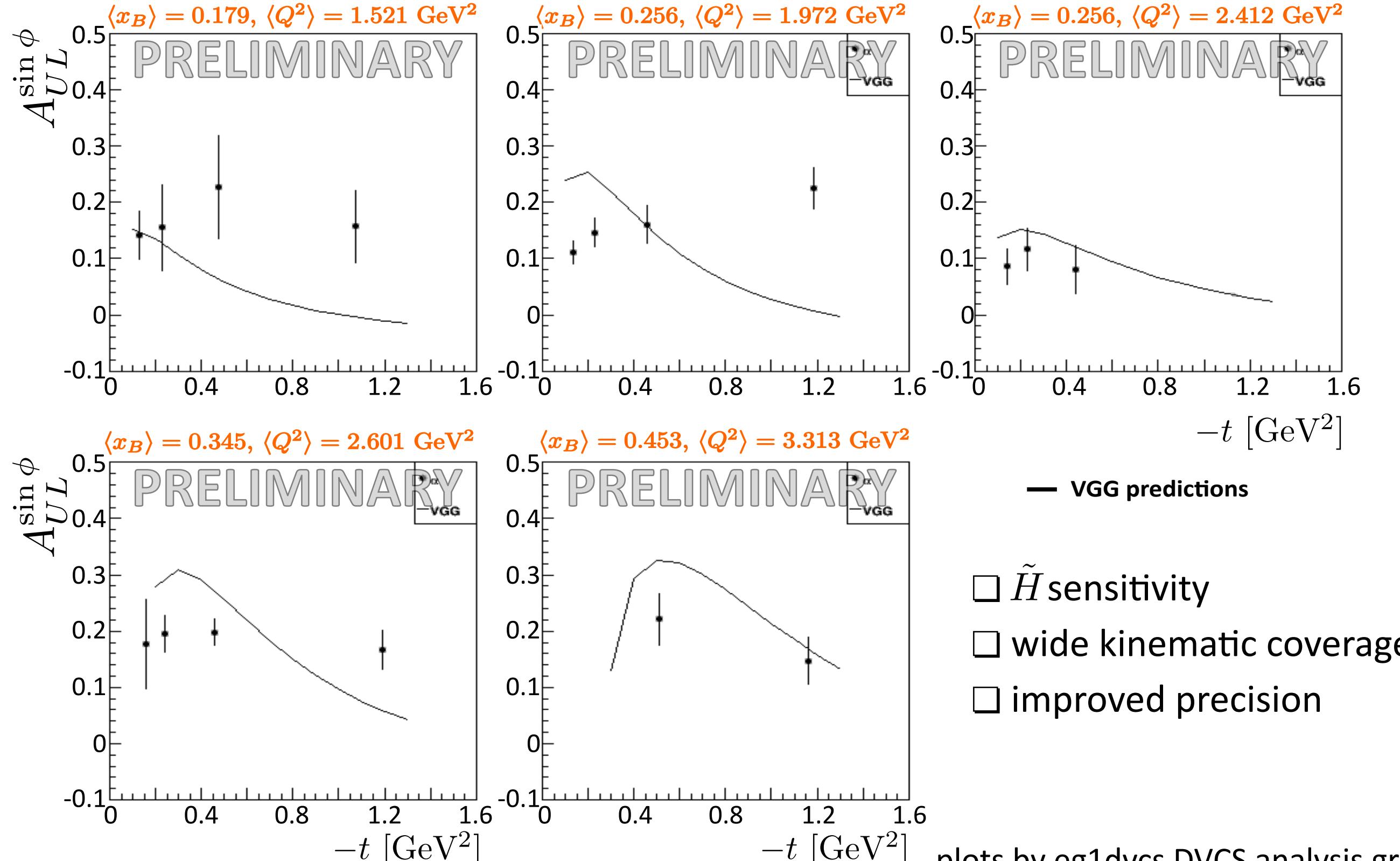
- ◆ non-dedicated experiment
- ◆ no Inner Calorimeter

$H(e, e'\gamma p)$ exclusivity



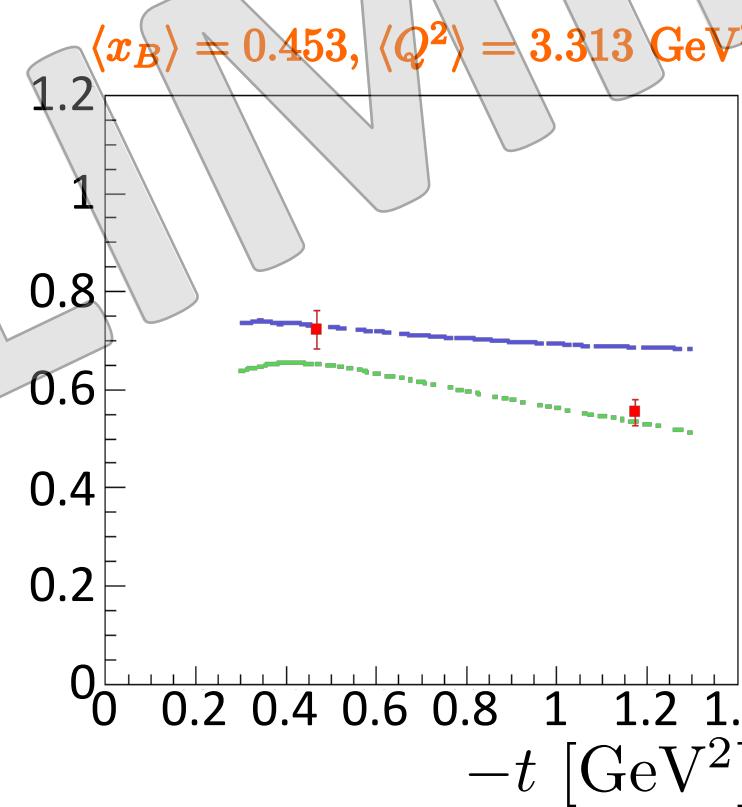
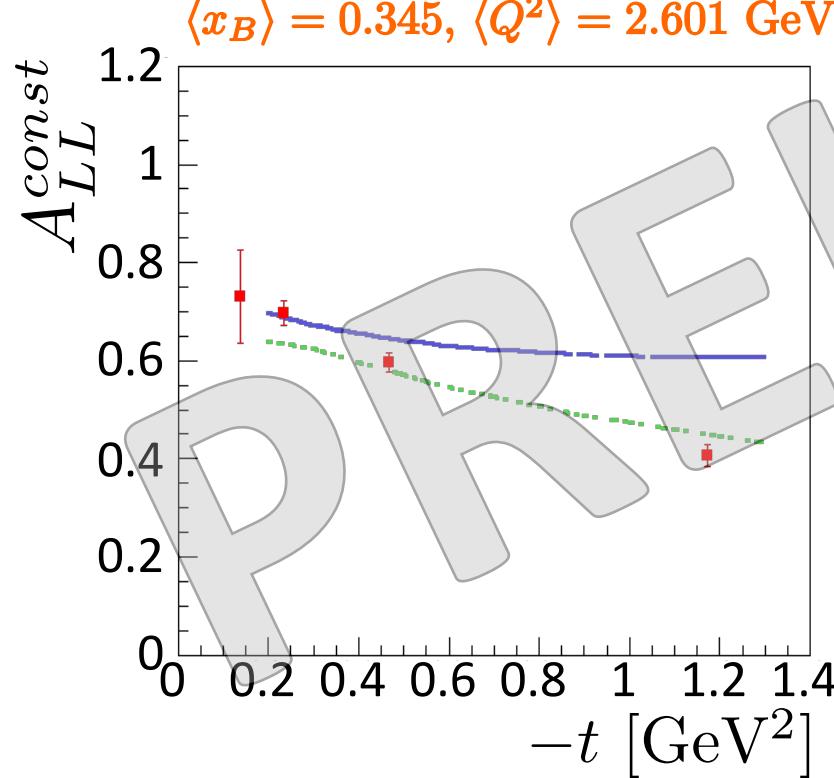
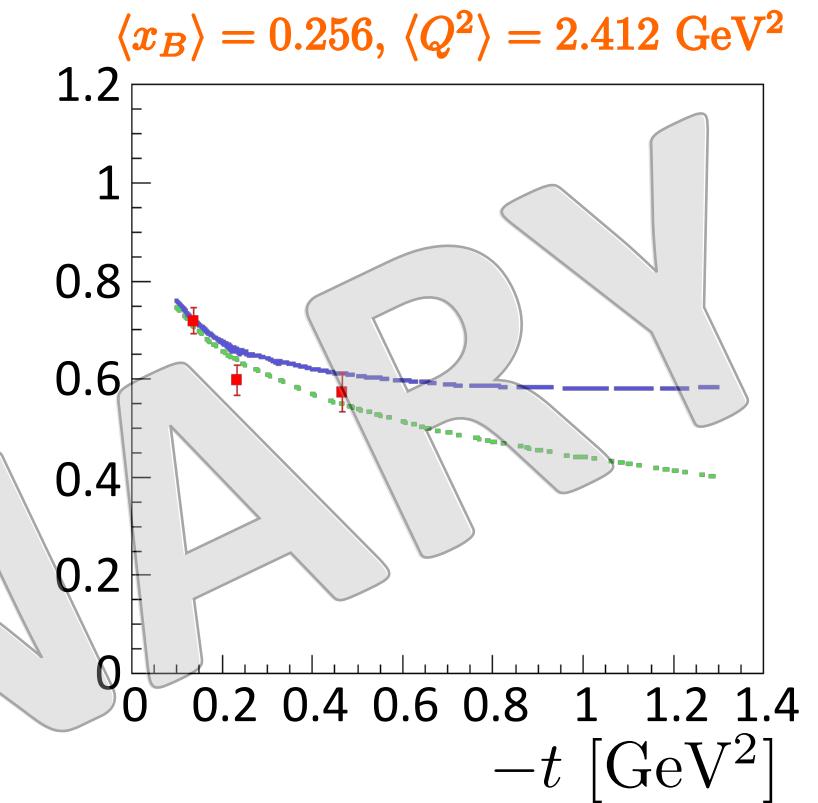
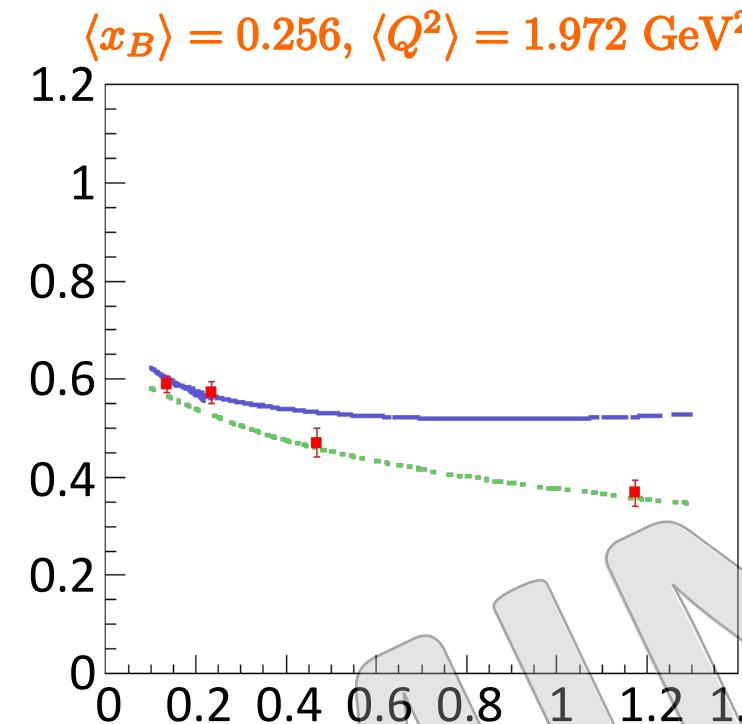
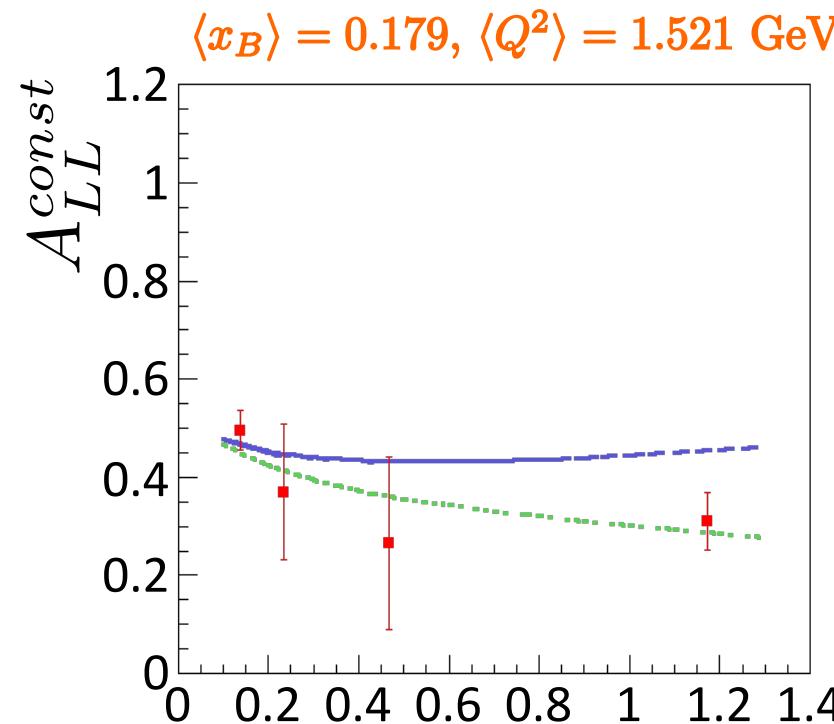
DVCS Target Spin Asymmetries

Second round experiment at Hall B: E05-114 (longitudinally polarized target)



DVCS Double Spin Asymmetries

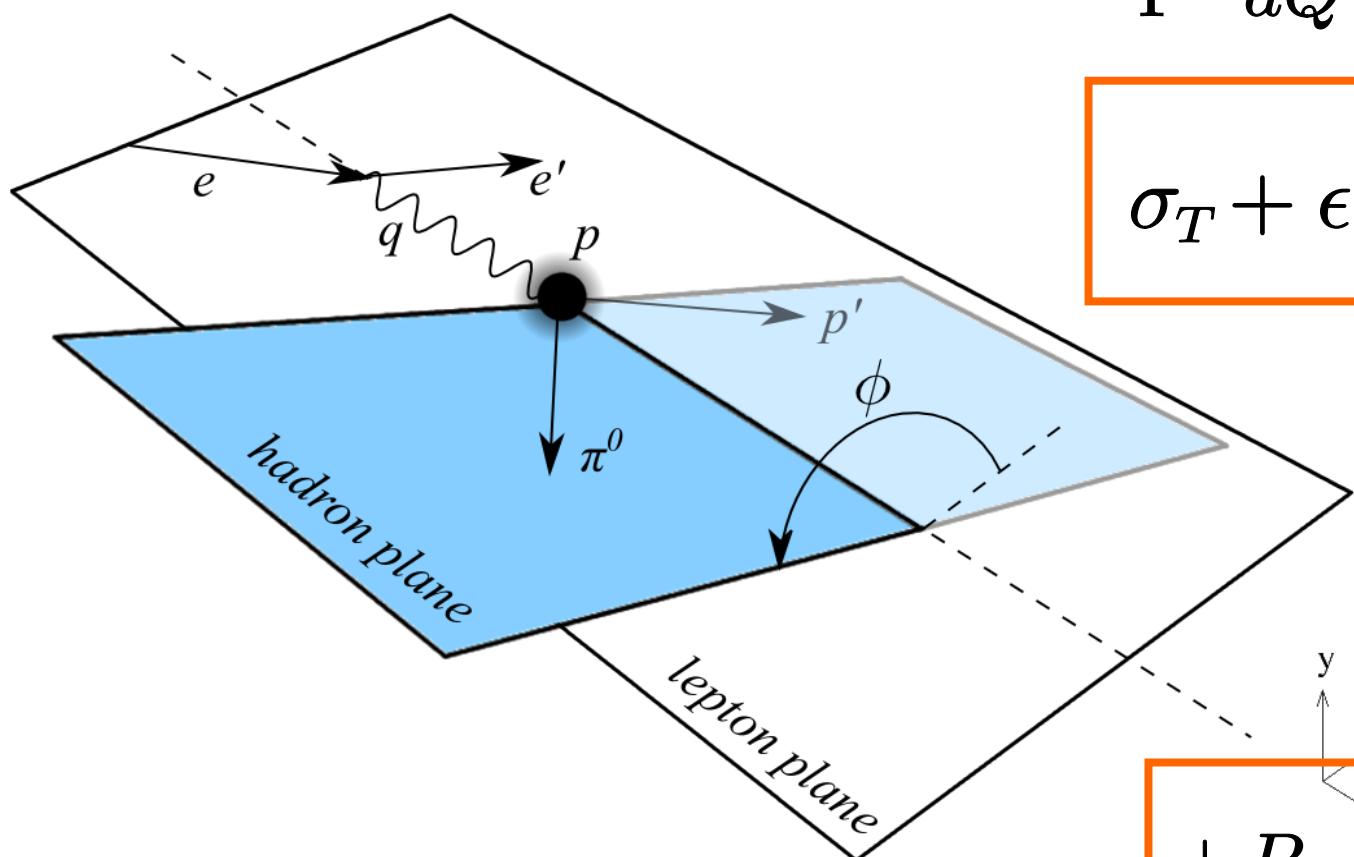
Second round experiment at Hall B: E05-114 (longitudinally polarized target)



- VGG predictions: BH term
- VGG predictions: BH+DVCS

- \tilde{H} sensitivity
- access to CFF real parts
- wide kinematic coverage
- improved precision

DV π^0 P Experimental Observables



$$\frac{2\pi}{\Gamma} \frac{d^4\sigma}{dQ^2 dx_B dt d\phi_\pi} =$$

unpolarized terms

$$\sigma_T + \epsilon\sigma_L + \epsilon\sigma_{TT} \cos 2\phi + \sqrt{\epsilon(1+\epsilon)}\sigma_{LT} \cos \phi$$

longitudinally polarized beam

$$+ P_b \sqrt{\epsilon(1-\epsilon)}\sigma_{LT'} \sin \phi$$

$$+ P_{tg} \left(\sqrt{\epsilon(1+\epsilon)}\sigma_{UL}^{\sin \phi} \sin \phi + \epsilon\sigma_{UL}^{\sin 2\phi} \sin 2\phi \right)$$

longitudinally polarized target

$$+ P_b P_{tg} \left(\sqrt{1-\epsilon^2}\sigma_{LL} + \sqrt{\epsilon(1-\epsilon)}\sigma_{LL}^{\cos \phi} \cos \phi \right)$$

longitudinally polarized beam and longitudinally polarized target

Deeply Virtual π^0 Production

Phys.Rev., D84:034007, 2011: G. R. Goldstein, J. O. Gonzalez, S. Liuti.
 Eur.Phys.J., A47:112, 2011: S.V. Goloskokov, P. Kroll.

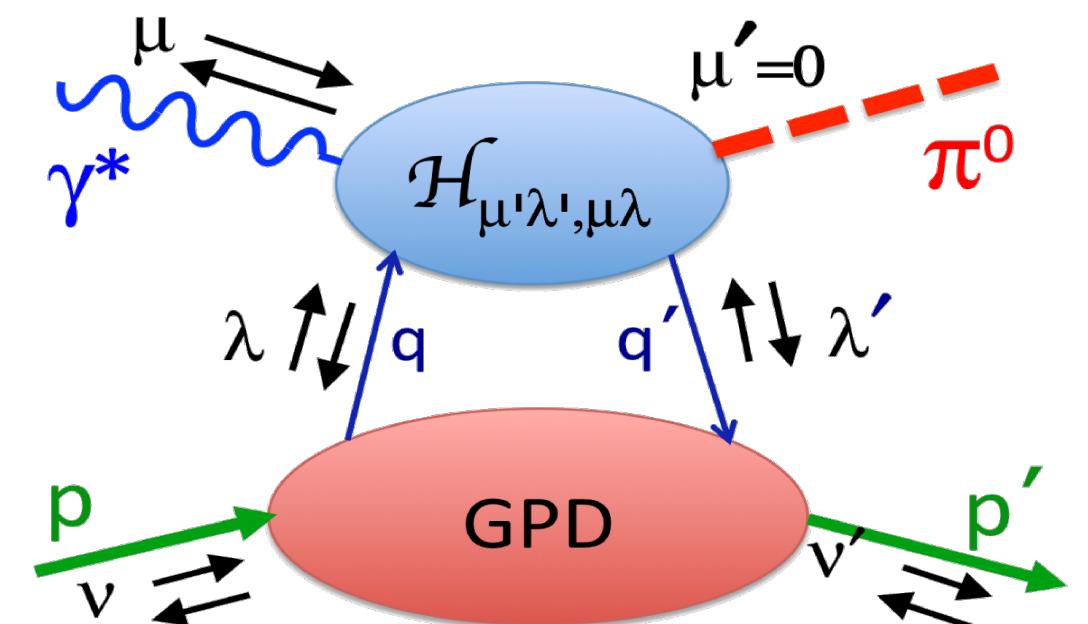
$$\langle F \rangle = \sum_{\lambda} \int_{-1}^1 dx \mathcal{H}_{0\lambda,\mu\lambda}(x, \xi, Q^2, t) F(x, \xi, t)$$

$$A_{UL}^{\sin \phi} \sigma_0 \sim \text{Im} \left[\langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle + \xi \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

$$A_{LU}^{\sin \phi} \sigma_0 \sim \text{Im} \left[\langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

$$A_{LL}^{\cos 0\phi} \sigma_0 \sim |\langle H_T \rangle|^2$$

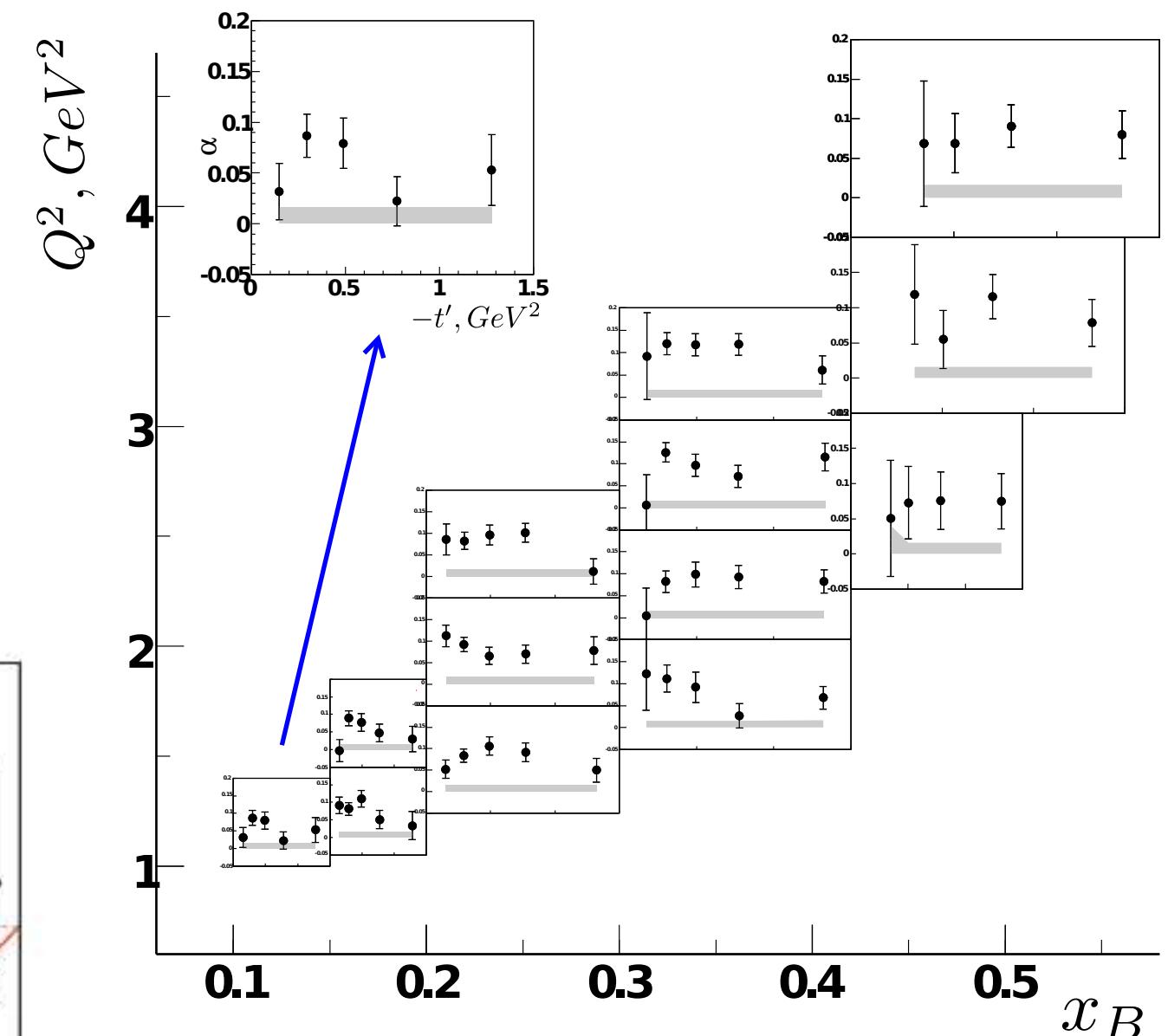
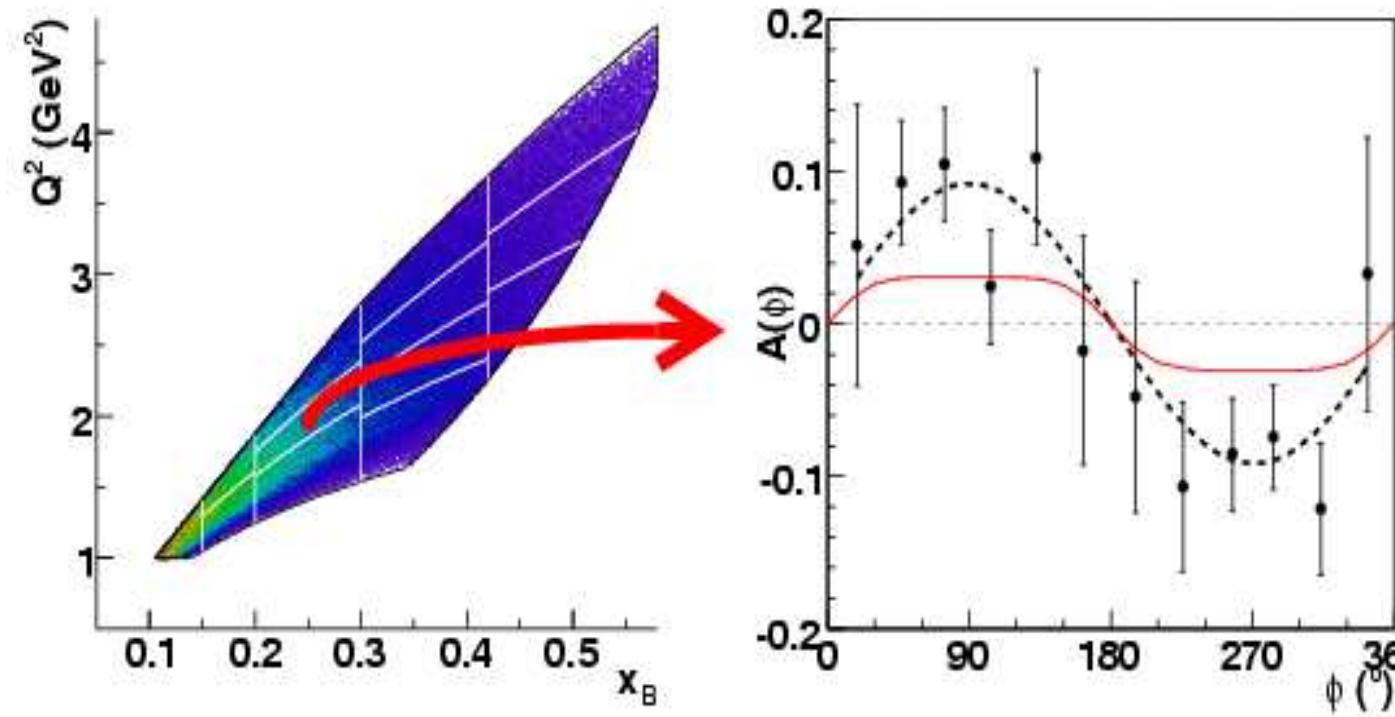
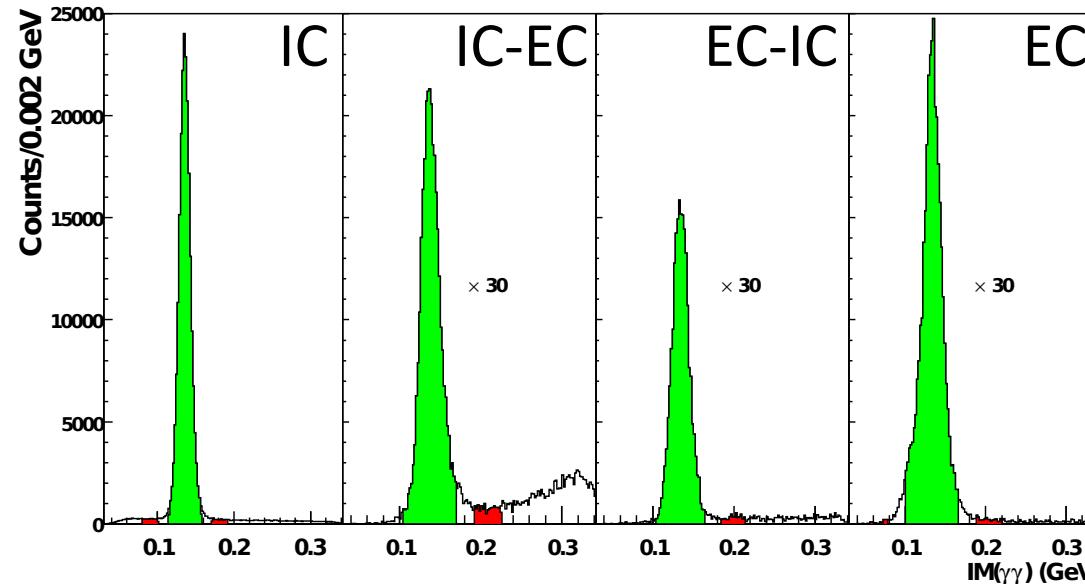
$$A_{LL}^{\cos \phi} \sigma_0 \sim \text{Re} \left[\langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle + \xi \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$



DV π^0 P Beam Spin Asymmetries

First round experiment at Hall B: E01-113

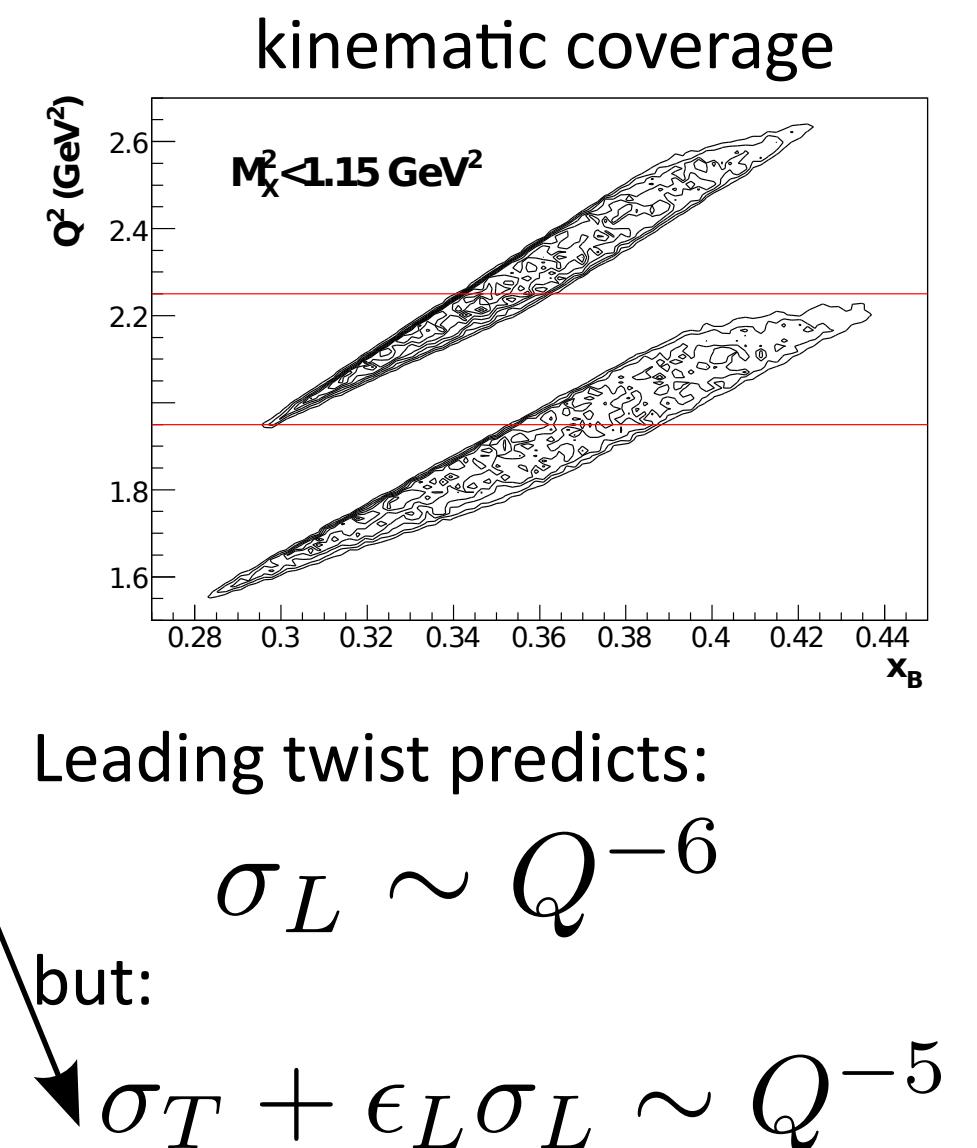
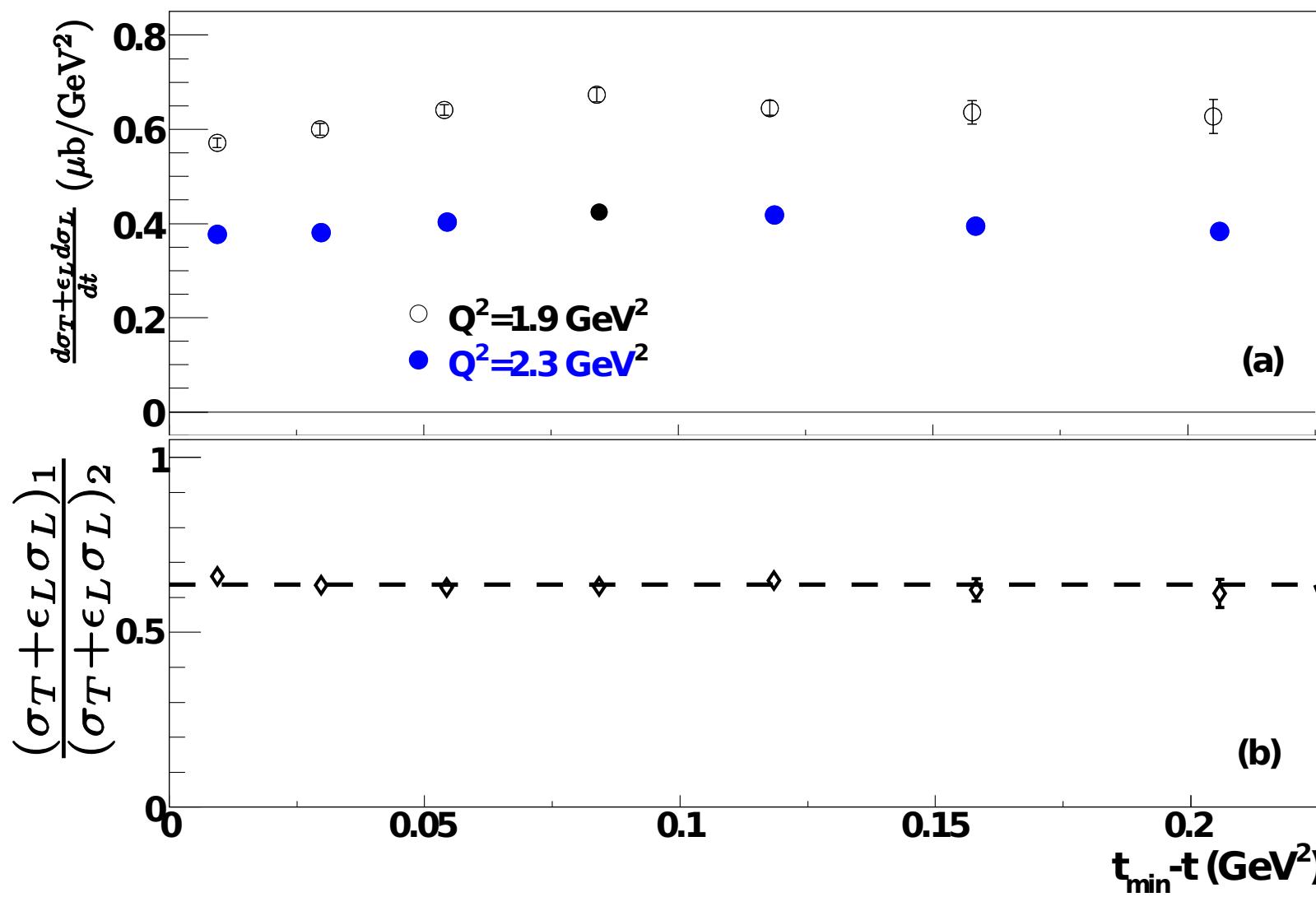
Two-photon invariant mass



R. De Masi et al. Measurement of $ep \rightarrow ep\pi^0$ beam spin asymmetries above the resonance region.
Phys. Rev., C77:042201, 2008.

DV π^0 P Cross Sections

First round experiment at Hall A: E00-110

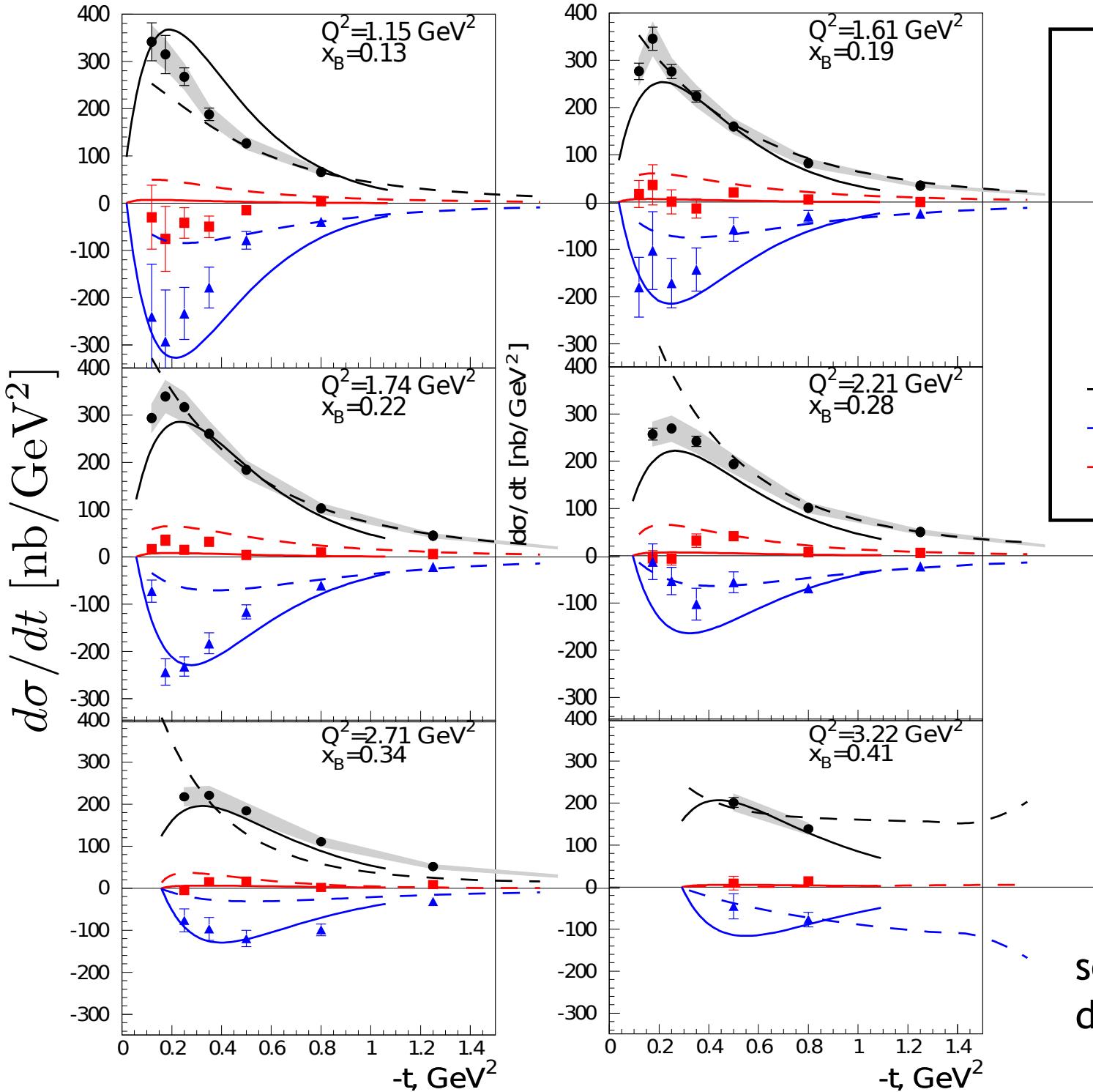


σ_T is likely to dominate at these Q^2

but L/T separation is necessary

DV π^0 P Cross Sections

First round experiment at Hall B: E01-113



Inclusion of the Chiral-odd GPDs
brings theoretical calculations into
moderate agreement with the data.

$$\sigma_T \sim (1 - \xi^2) |H_T|^2 - \frac{t'}{8m^2} |\bar{E}_T|^2$$

— $\sigma_0 = \sigma_T + \epsilon \sigma_L$
 — σ_{TT}
 — σ_{LT}

$$\sigma_{TT} \sim \frac{t'}{8m^2} |\bar{E}_T|^2$$

π^0 electroproduction
is uniquely sensitive process to access
transversity GPDs.

solid: P.Kroll & S.Goloskokov
dashed: G.R. Goldstein, J.O. Gonzalez & S.Liuti

DV π^0 P Spin Asymmetries

Second round experiment at Hall B: E05-114 (longitudinally polarized target)

DV π^0 P single and double spin asymmetry measurements are underway

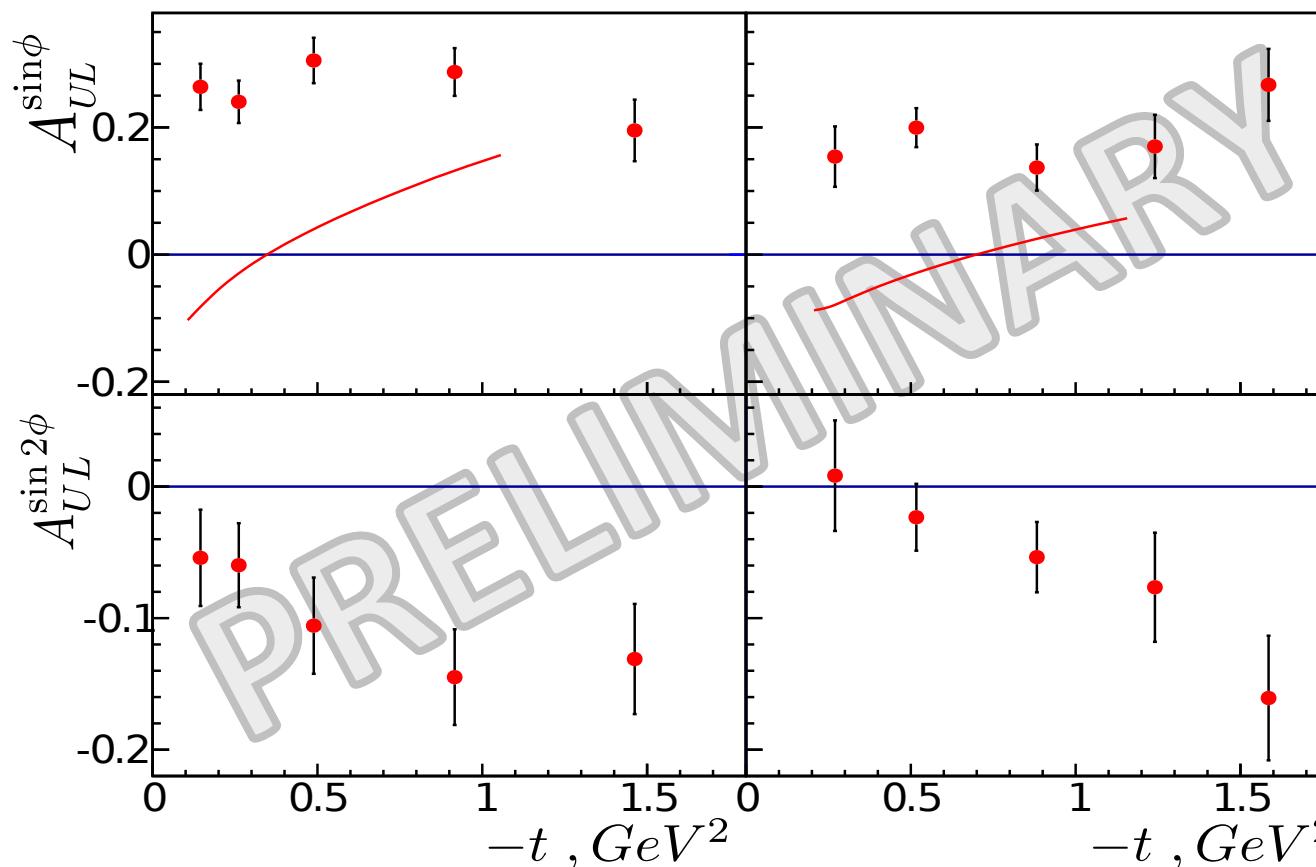
TARGET SPIN ASYMMETRY

$$\langle Q^2 \rangle = 1.94 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.25$$

$$\langle Q^2 \rangle = 2.83 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.40$$



$$A_{UL}^{\sin\phi} \sigma_0 \sim \text{Im} \left[\langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle + \xi \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

— Goloskokov-Kroll
- - - Goldstein-Liuti

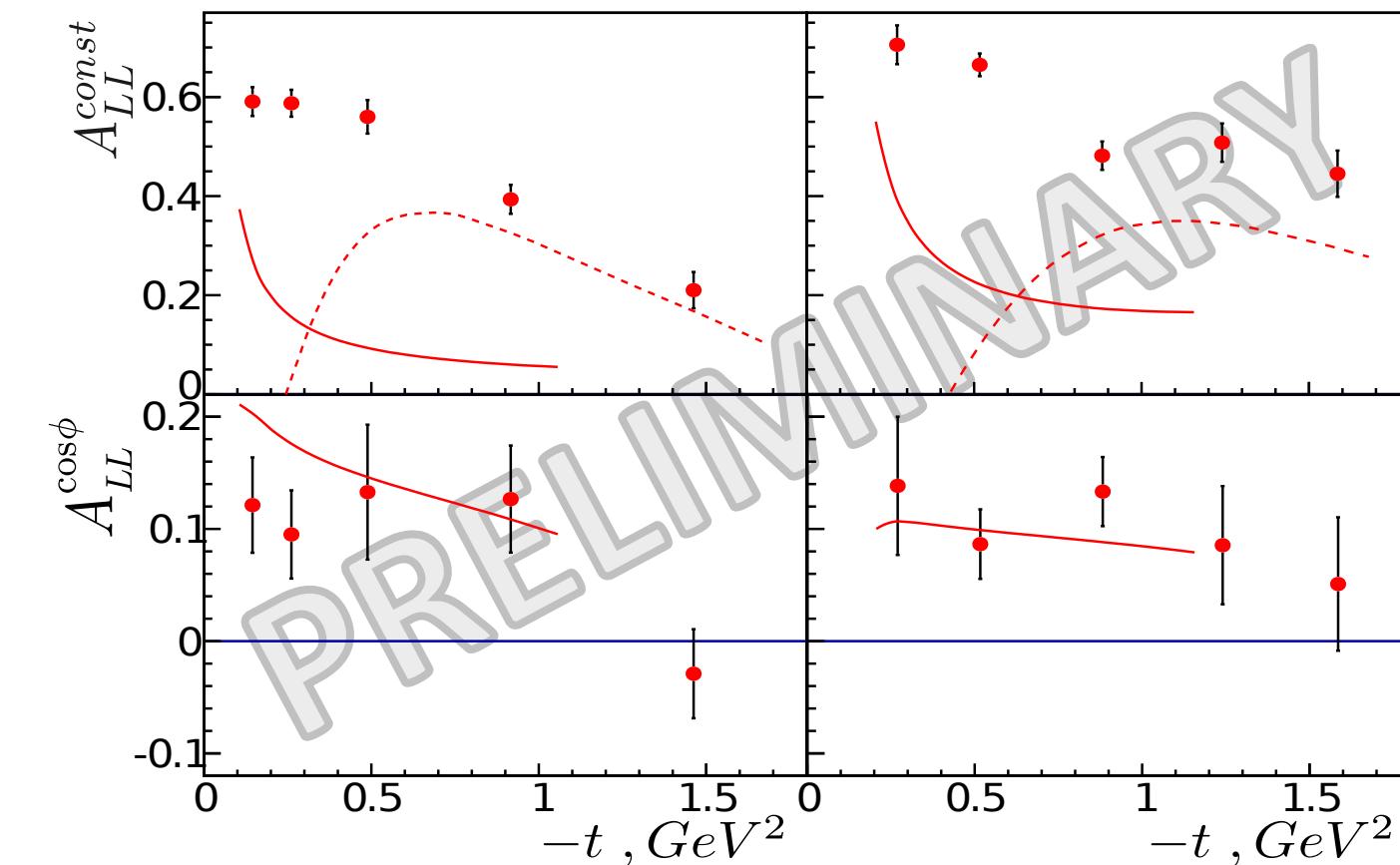
DOUBLE SPIN ASYMMETRY

$$\langle Q^2 \rangle = 1.94 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.25$$

$$\langle Q^2 \rangle = 2.83 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.40$$



$$A_{LL}^{\cos 0\phi} \sigma_0 \sim |\langle H_T \rangle|^2$$

$$A_{LL}^{\cos\phi} \sigma_0 \sim \text{Re} \left[\langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle + \xi \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

Summary

- ◆ Large data set (cross sections, single and double spin asymmetries) in the wide kinematic region is available
- ◆ DVCS and DV $\pi^0 P$ provide access to the extensive set of chiral-even and chiral-odd GPDs
- ◆ DVCS reaction is a clean probe of the GPDs
- ◆ Combination of polarized and unpolarized observables provide constraints for t dependence on underlying GPDs and will help to establish the role of transversity of pion electroproduction
- ◆ Compeling GPD program in the future at JLab at 12 GeV...