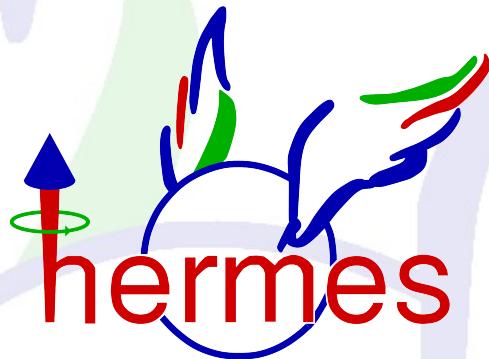


Recent HERMES results on TMDs from unpolarized targets

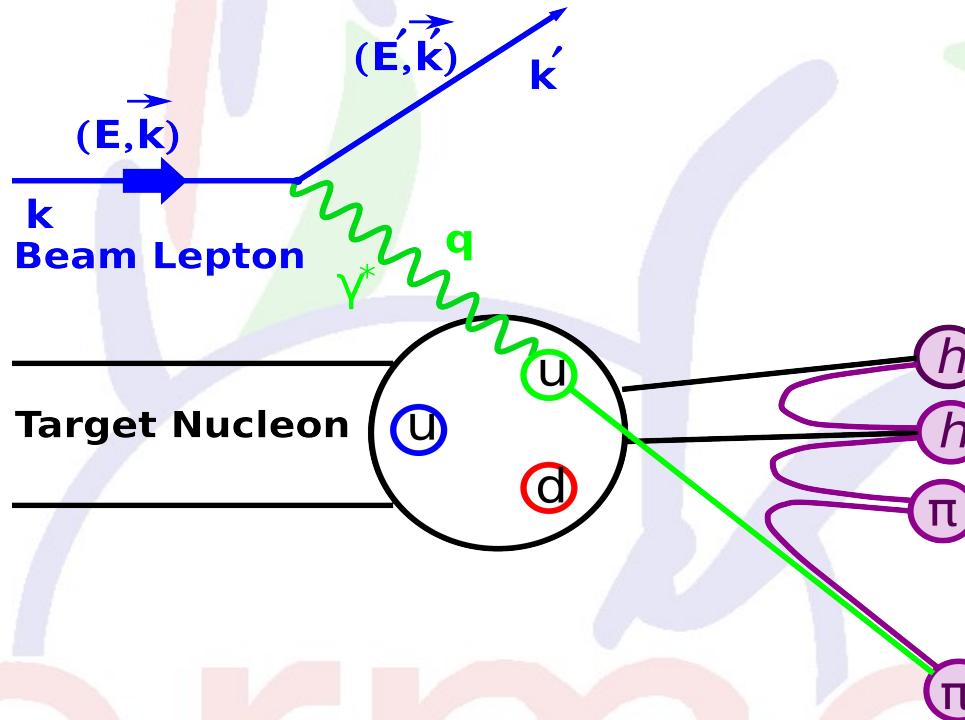


Gevorg Karyan

on behalf of the HERMES Collaboration

Alikhanyan National Science Laboratory
Yerevan, Armenia

Semi-Inclusive Deep-Inelastic Scattering

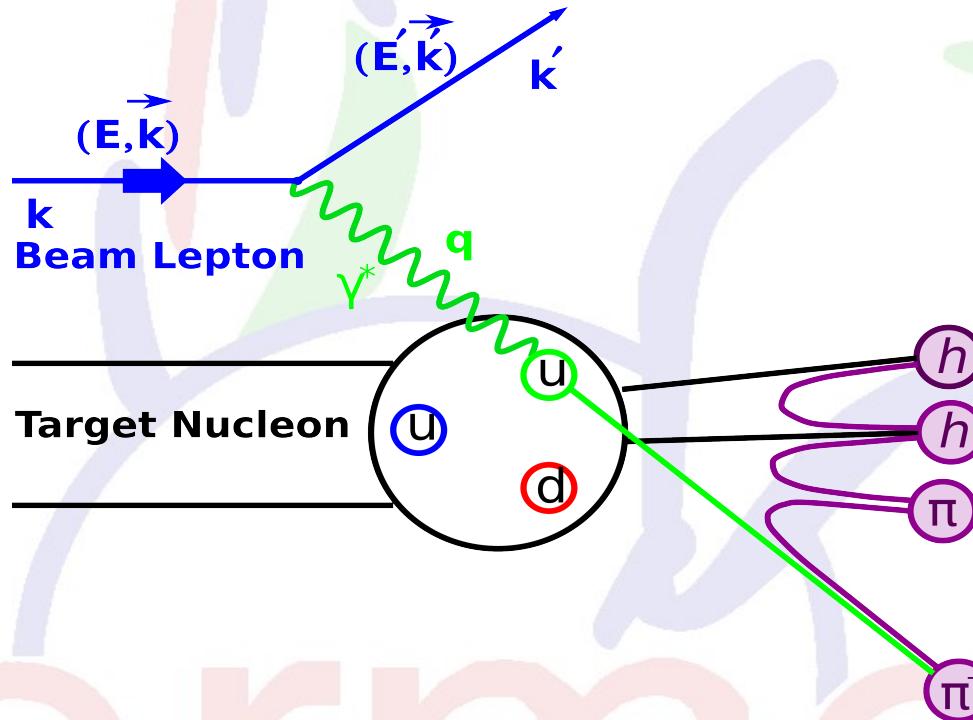


Semi-Inclusive Deep-Inelastic Scattering

$$Q^2 = -(\mathbf{k} - \mathbf{k}')^2$$

$$\nu = E - E'$$

$$W^2 = (\mathbf{P} + \mathbf{q})^2$$



$$x_B = \frac{Q^2}{2 \cdot M_N \cdot \nu}$$

$$z = E_h / \nu$$

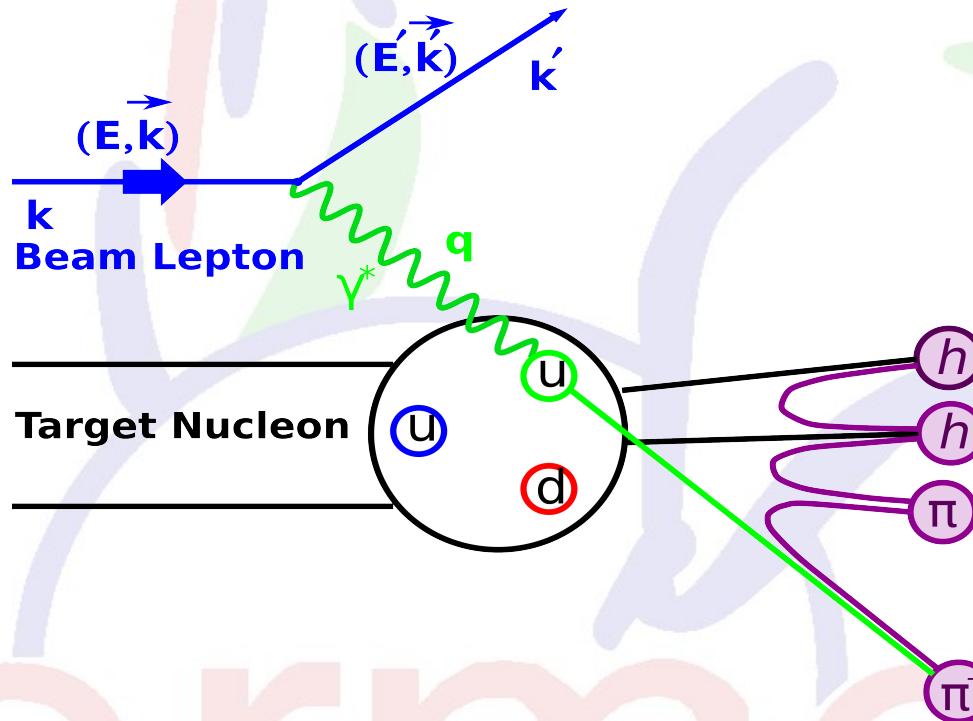
$$P_{h\perp} = \frac{|\vec{P}_h \times \vec{q}|}{|\vec{q}|}$$

Semi-Inclusive Deep-Inelastic Scattering

$$Q^2 = -(\mathbf{k} - \mathbf{k}')^2$$

$$\nu = E - E'$$

$$W^2 = (\mathbf{P} + \mathbf{q})^2$$



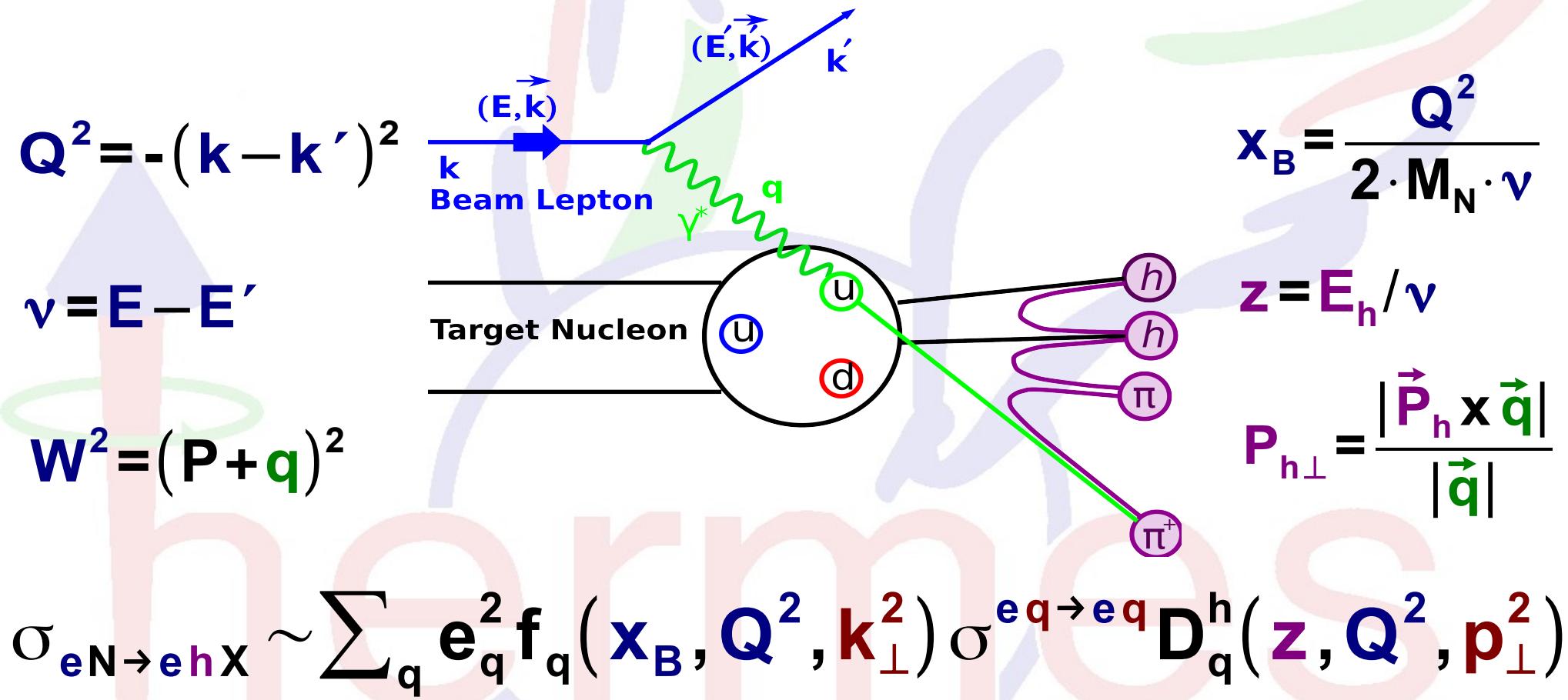
$$x_B = \frac{Q^2}{2 \cdot M_N \cdot \nu}$$

$$z = E_h / \nu$$

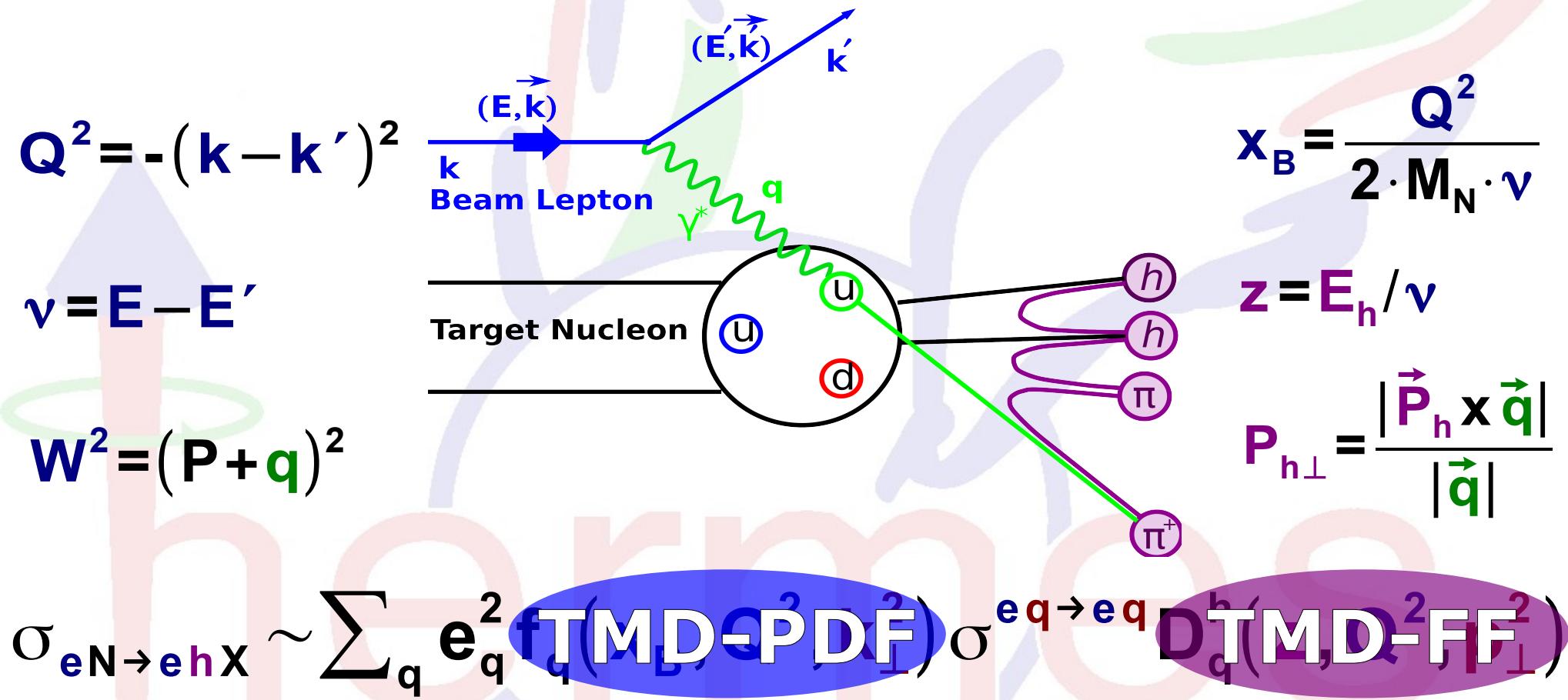
$$P_{h\perp} = \frac{|\vec{P}_h \times \vec{q}|}{|\vec{q}|}$$

TMD factorization at leading order in (k_\perp/Q) , $P_{h\perp} \simeq k_\perp \ll Q$

Semi-Inclusive Deep-Inelastic Scattering



Semi-Inclusive Deep-Inelastic Scattering



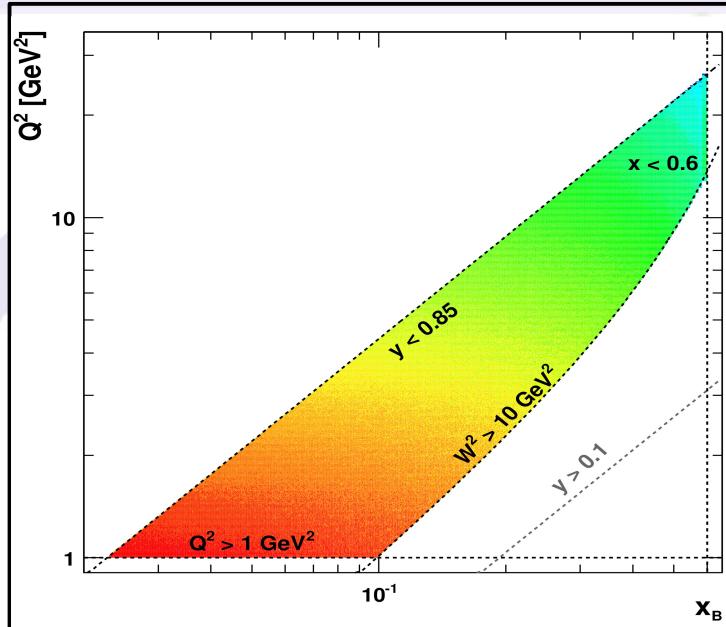
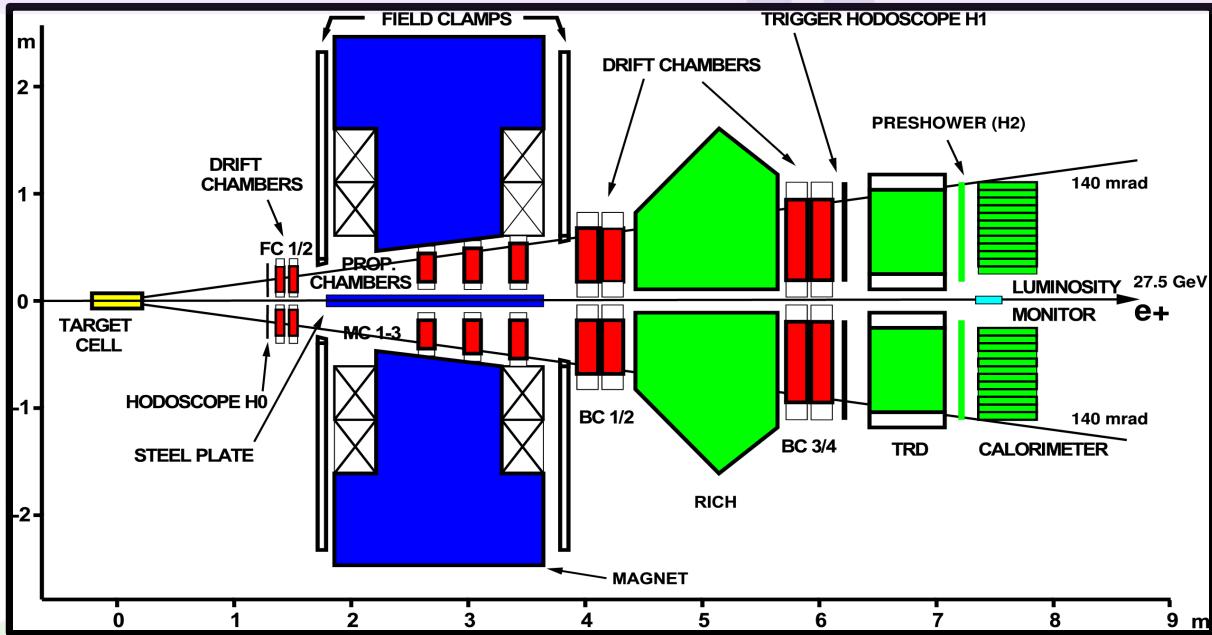
Experimental observable

SIDIS hadron yields

$$M^h(x_B, Q^2, z, P_{h\perp}, \phi_h) = \frac{N^h(x_B, Q^2, z, P_{h\perp}, \phi_h)}{N^e(x_B, Q^2)}$$

DIS event yields

Experiment



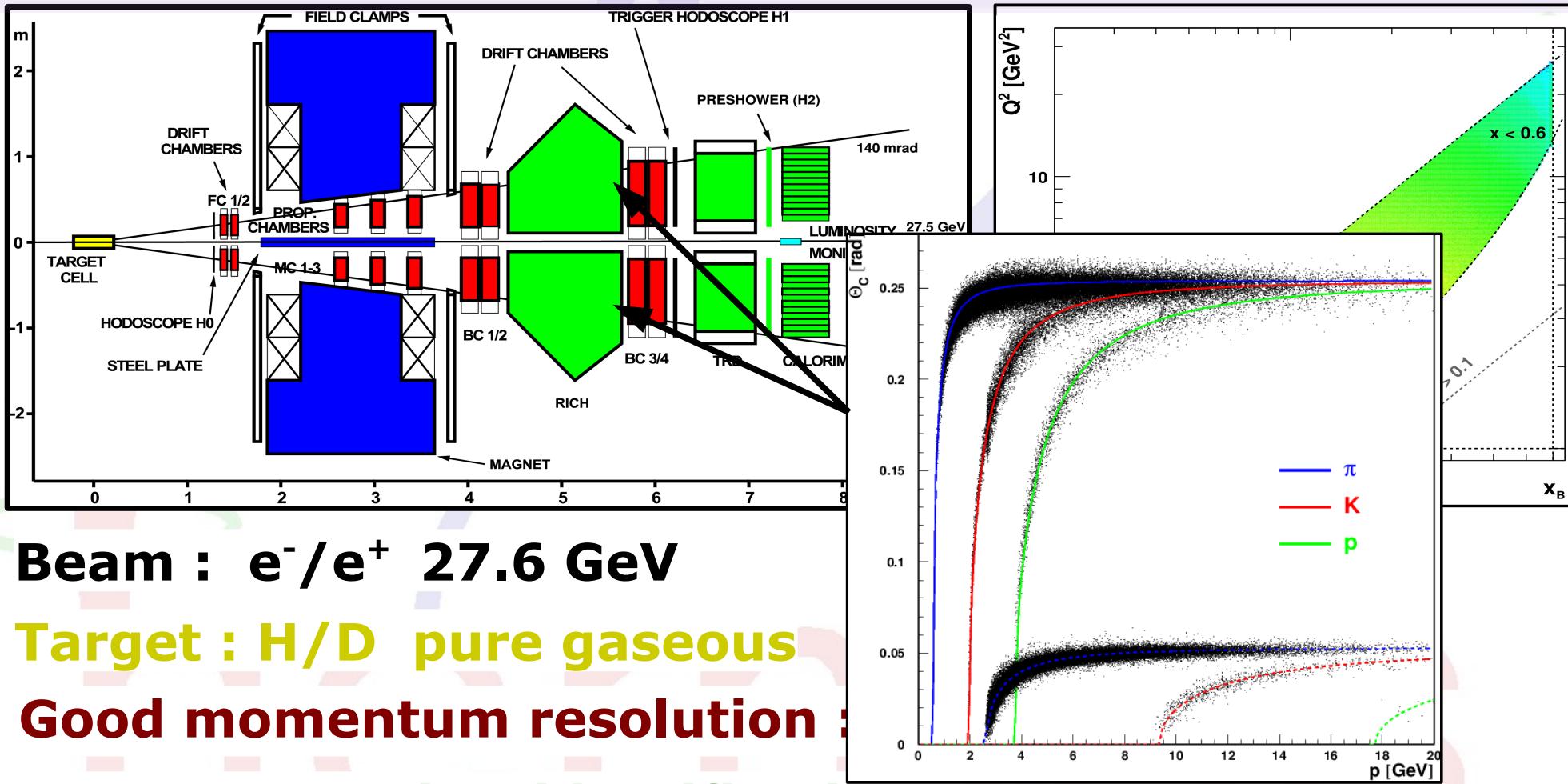
Beam : e^-/e^+ 27.6 GeV

Target : H/D pure gaseous

Good momentum resolution : $\frac{\delta p}{p} < 2 \%$

Excellent particle identification

Experiment



Data selection

DIS regime

- $Q^2 > 1 \text{ GeV}^2$
- $W^2 > 10 \text{ GeV}^2$
- $0.1 < v/E_{\text{beam}} < 0.85$

SIDIS selection

- $2 \text{ GeV} < p < 15 \text{ GeV}$
- $0.2 < z < 0.8$
- $x_F > 0.2$

Raw Data

Data analysis

Raw Data

**Charge Symmetric Background
(Dalitz decay, $\gamma \rightarrow e^+e^-$)**

RICH Unfolding

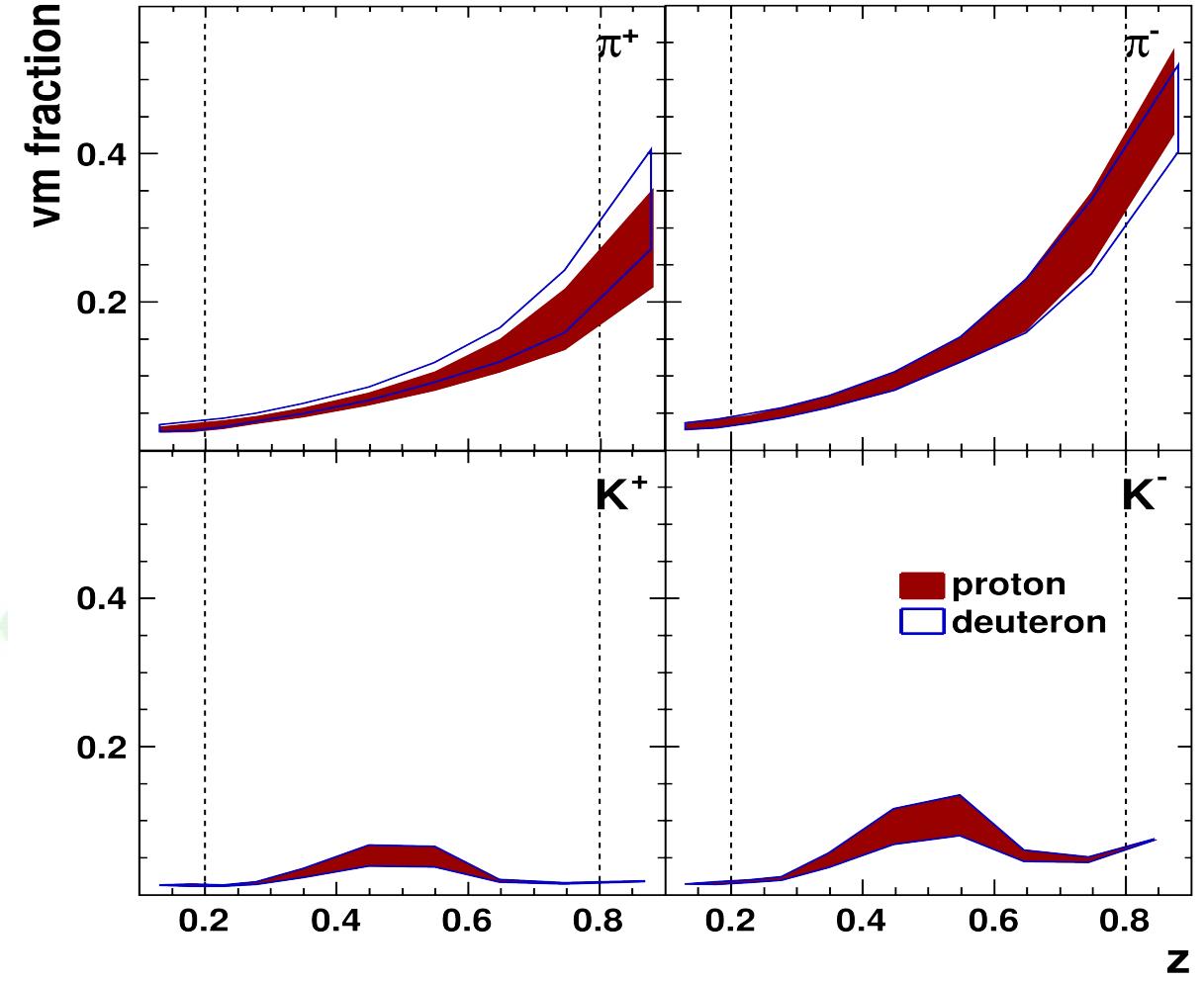
Trigger Efficiencies

**Diffractive Vector Meson
Contribution**

**Detector Smearing & QED Radiative
Effects**

Final Data

Diffractive vector meson contribution



Due to diffractively produced exclusive
 $\rho^0 \rightarrow \pi^+ \pi^-$ and $\phi \rightarrow K^+ K^-$

➤ Results with and without
VM subtraction.

Experimental observable

SIDIS hadron yields

$$M^h(x_B, Q^2, z, P_{h\perp}, \phi_h) = \frac{N^h(x_B, Q^2, z, P_{h\perp}, \phi_h)}{N^e(x_B, Q^2)}$$

DIS event yields

Experimental observable

SIDIS hadron yields

Collinear framework

$$M^h(x_B, Q^2) = \frac{N^h(x_B, Q^2, z)}{N^e(x_B, Q^2)}$$

DIS event yields

Underlying physics

Collinear framework

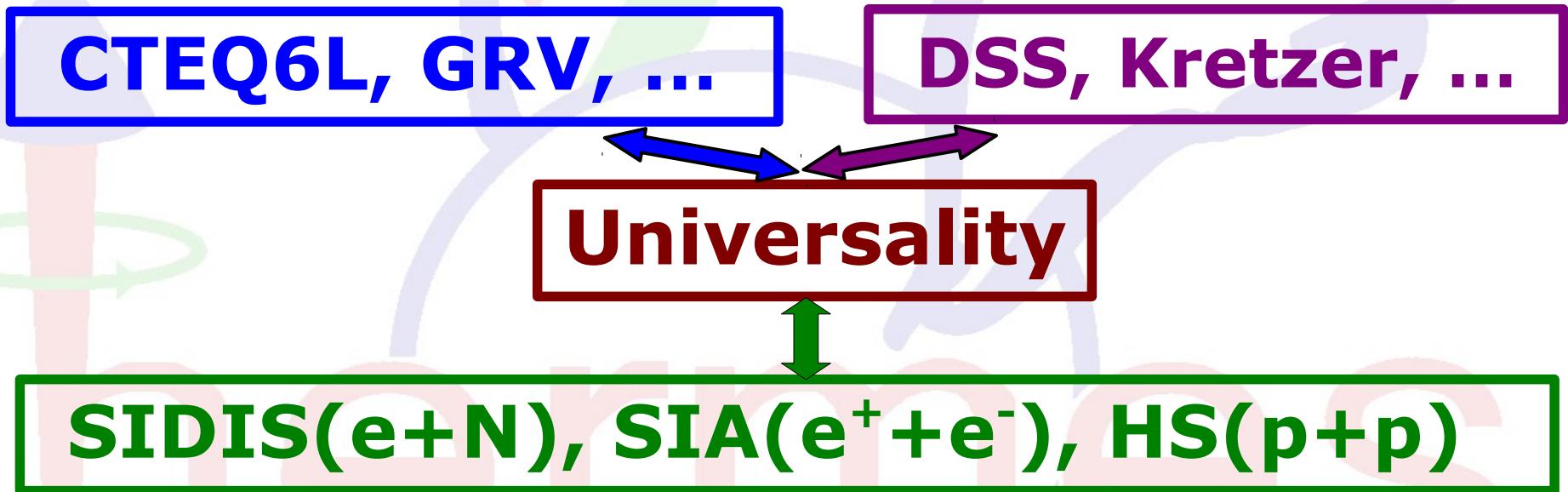
$$M^h \sim \frac{\sum_q e_q^2 f_q(PDF, Q^2) D_q^h(zFF, Q^2)}{\sum_q e_q^2 f_q(PDF, Q^2)}$$

using collinear PDFs (well known)

extract collinear FFs

Underlying physics

$$M^h \sim \frac{\sum_q e_q^2 f_q(PDF_B) Q^2 D_q^h(z, FF^2)}{\sum_q e_q^2 f_q(PDF_B) Q^2}$$



Advantage of SIDIS

Charge separated FFs



SIDIS

($D_u^\pi, D_{\bar{u}}^K, \dots$)

(D_u^π, D_s^K, \dots)



Flavor separated FFs

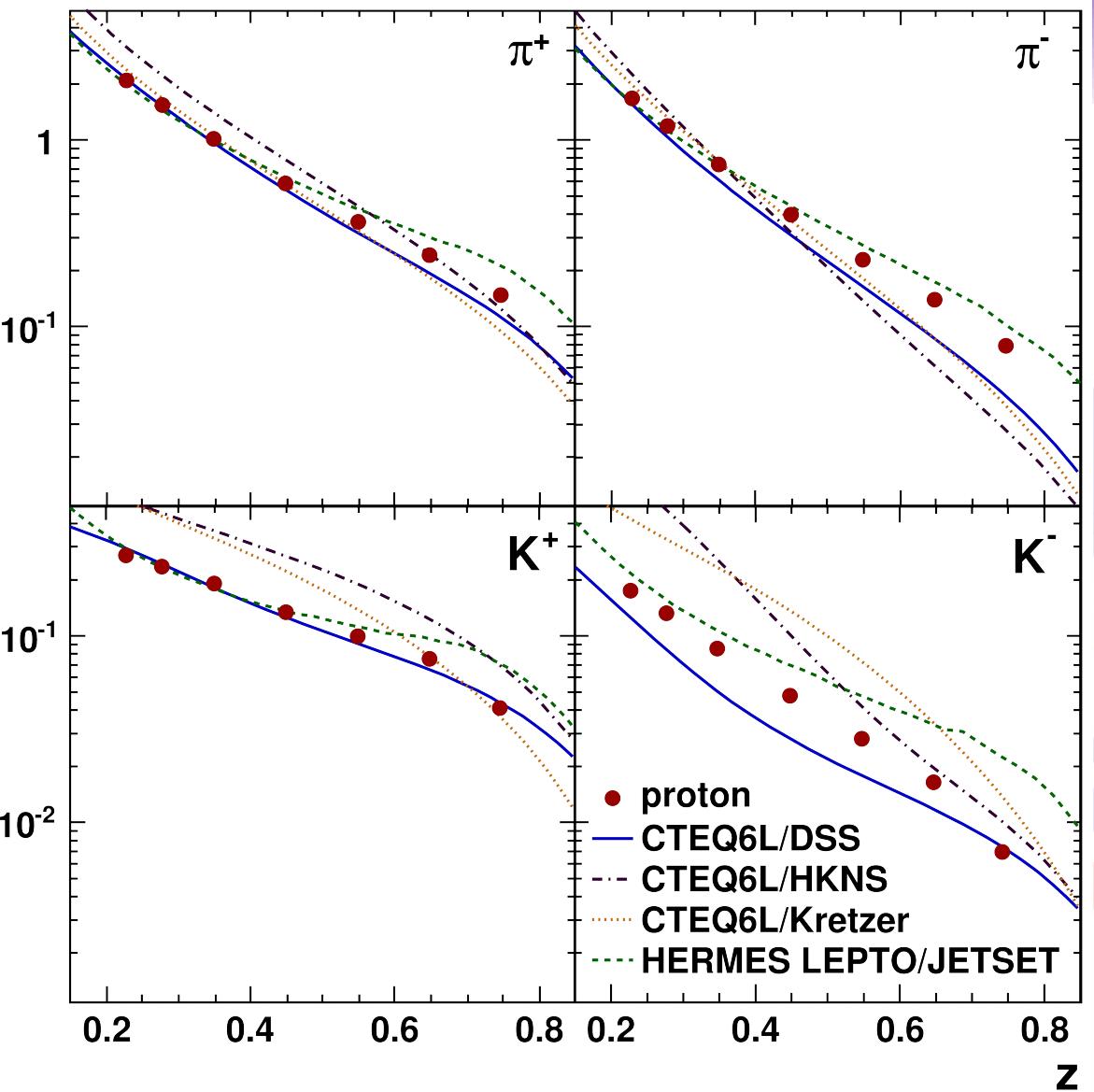
$$M^h(x_B, Q^2, z, P_{h\perp})$$

Proton target

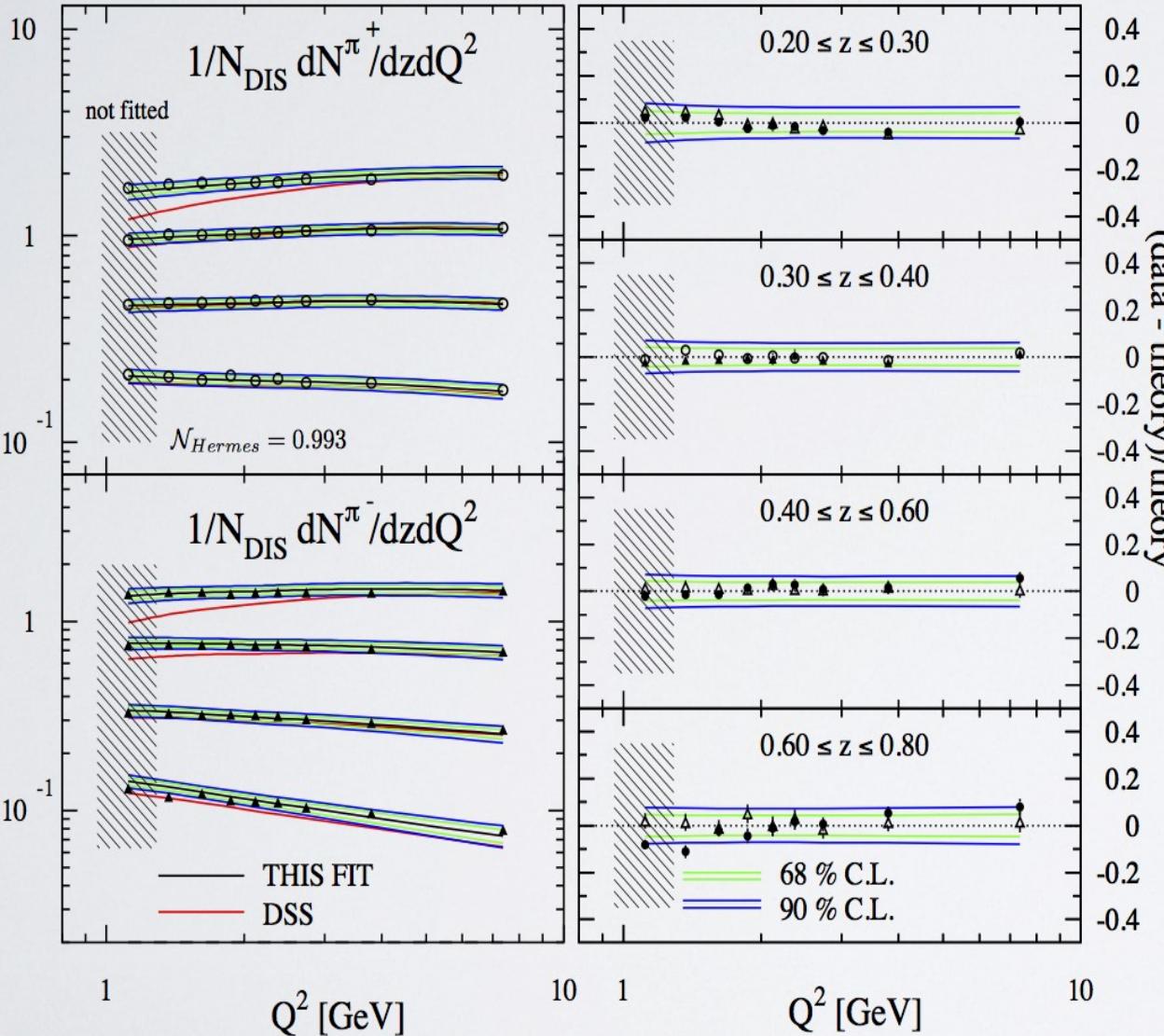
LO calculations

- Reasonable agreement between **DSS FFs** and **Data** for positively charged pions and kaons.

- Substantial differences between all FFs and **Data** for negatively charged kaons.



HERMES proton



$$M^h(x_B, Q^2, z, P_{h\perp})$$

Proton target

NLO calculations(DSS+)

- ◆ Much better agreement for both π^+ and π^- .

◆ Workshop on fragmentation functions, Bloomington, 2013

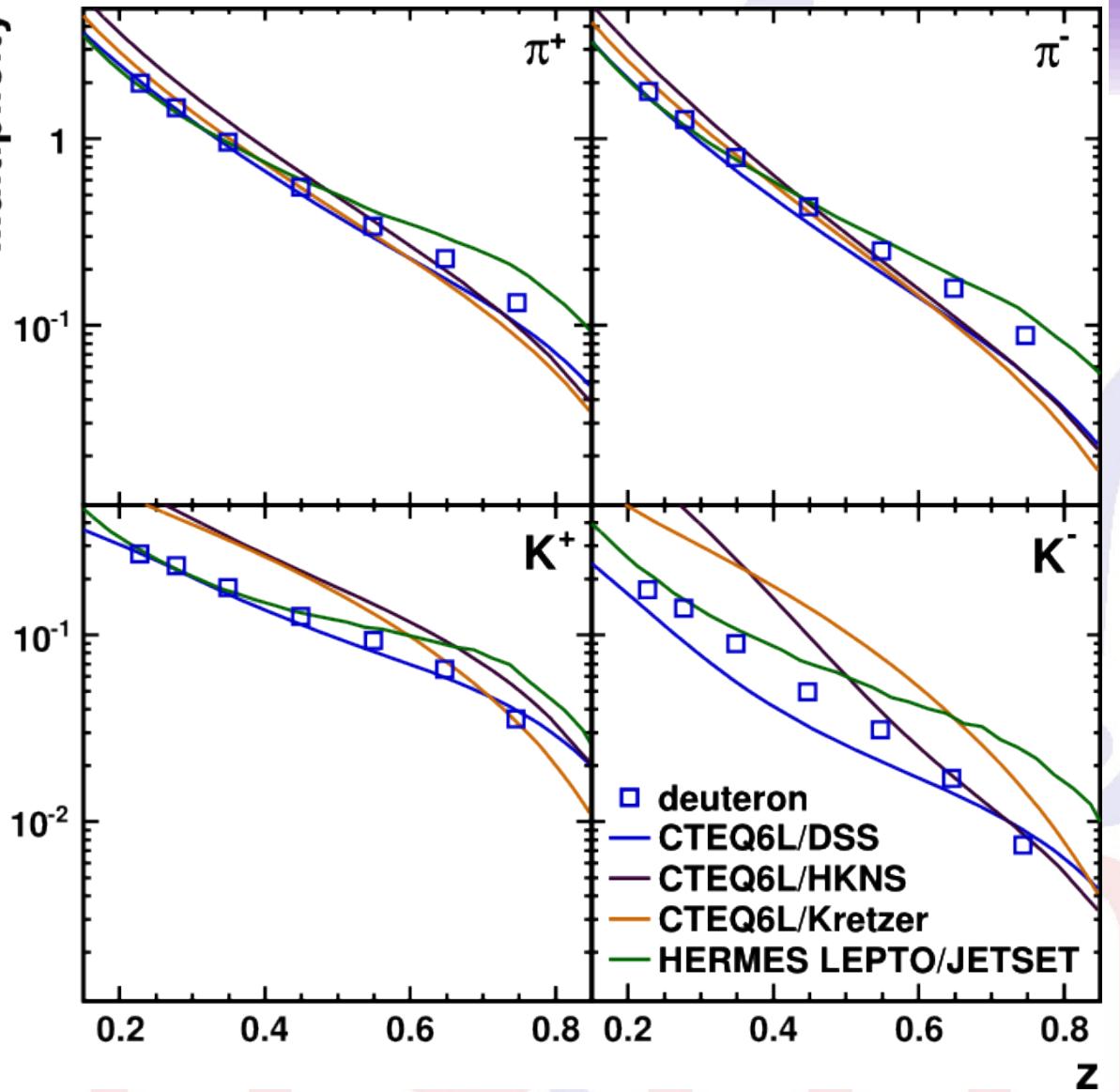
$$M^h(x_B, Q^2, z, P_{h\perp})$$

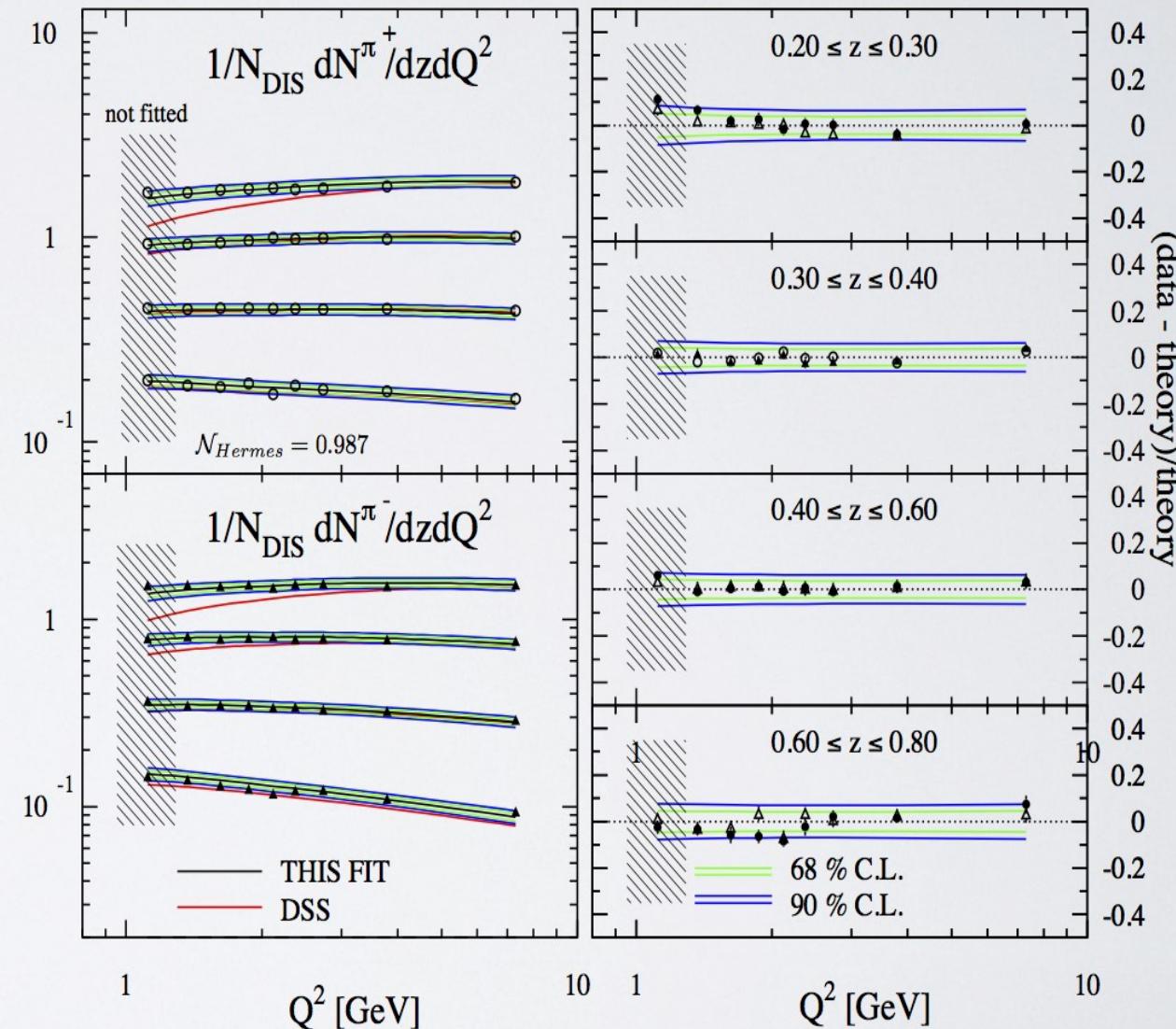
Deuteron target

LO calculations

- ◆ Reasonable agreement between **DSS FFs** and **Data** for positively charged pions and kaons.

- ◆ Substantial differences between all FFs and **Data** for negatively charged kaons.





$$M^h(x_B, Q^2, z, P_{h\perp})$$

Deuteron target

NLO calculations(DSS+)

- ◆ Much better agreement for both π^+ and π^- .

◆ Workshop on fragmentation functions, Bloomington, 2013

Experimental observable

SIDIS hadron yields

$$M^h(x_B, Q^2, z, P_{h\perp}, \phi_h) = \frac{N^h(x_B, Q^2, z, P_{h\perp}, \phi_h)}{N^e(x_B, Q^2)}$$

DIS event yields

Experimental observable

SIDIS hadron yields

$$M^h(x_B, Q^2, TMDs) \text{ via } \frac{N^h(x_B, Q^2, z, P_{h\perp})}{N^{eh}(x_B, Q^2)}$$

DIS event yields

Flavor-dependent Gaussian ansatz

$$f_q(x_B, Q^2, k_{\perp}^2) = f_q(x_B, Q^2) \frac{1}{\pi \langle k_{\perp, q}^2 \rangle} e^{-k_{\perp}^2 / \langle k_{\perp, q}^2 \rangle}$$

$$D_q^h(z, Q^2, p_{\perp}^2) = D_q^h(z, Q^2) \frac{1}{\pi \langle p_{\perp, q \rightarrow h}^2 \rangle} e^{-p_{\perp}^2 / \langle p_{\perp, q \rightarrow h}^2 \rangle}$$

$$\langle P_{h \perp, q}^2 \rangle = \langle p_{\perp, q \rightarrow h}^2 \rangle + z^2 \langle k_{\perp, q}^2 \rangle$$

A. Signori, A. Bacchetta, M. Radici and G. Schnell(JHEP, 2013)

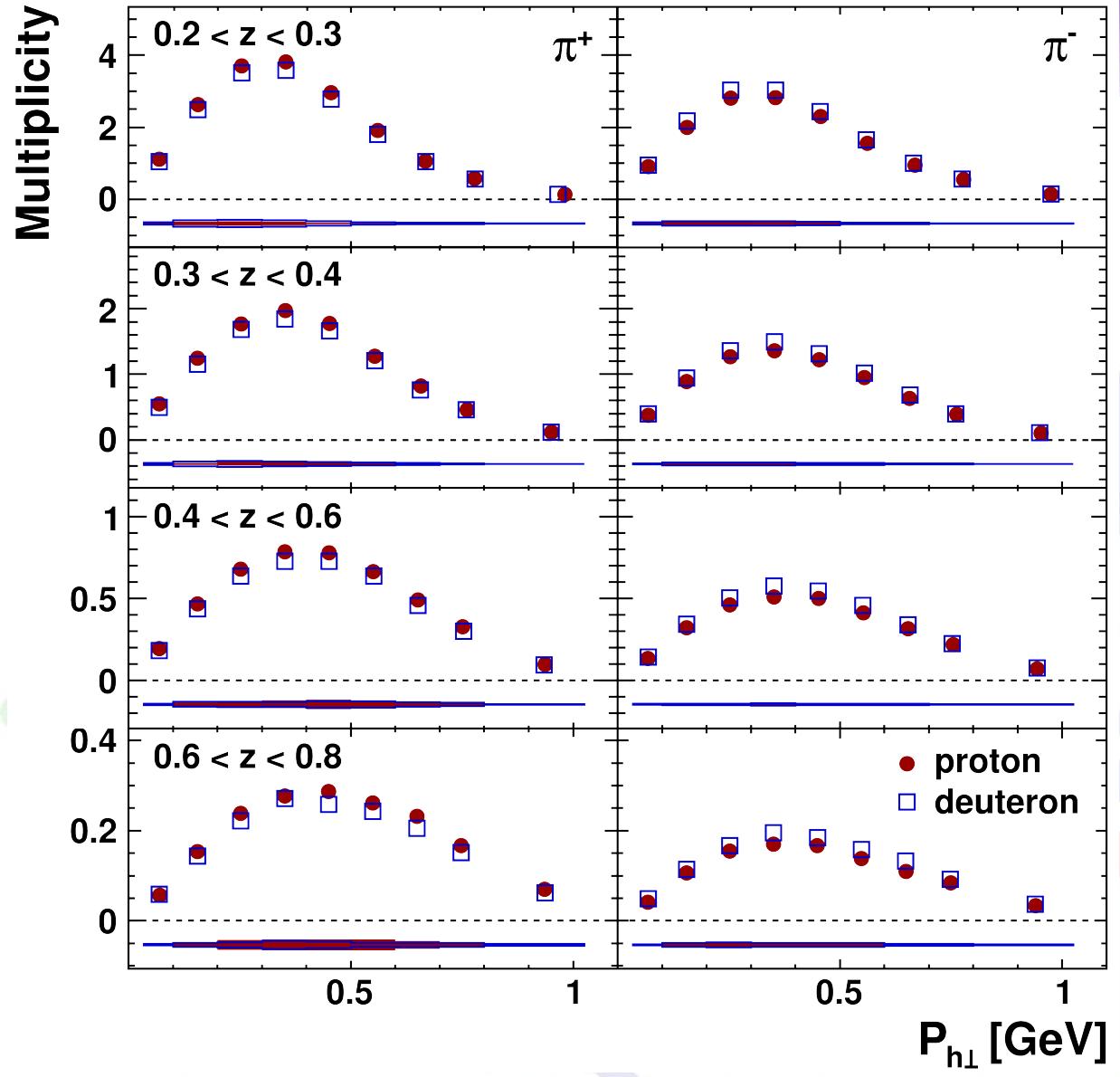
Flavor-dependent Gaussian ansatz

$$f_q(x_B, Q^2, k_\perp^2) = f_q(\text{PDF}, Q^2) \frac{1}{\pi \langle k_{\perp, q}^2 \rangle} e^{-k_\perp^2 / \langle k_{\perp, q}^2 \rangle}$$

$$\int f_q(x_B, Q^2, k_\perp^2) d^2 k_\perp$$

$$\int D_q^h(z, Q^2, p_\perp^2) d^2 p_\perp$$

$$D_q^h(z, Q^2, p_\perp^2) = D_q^h(\text{FF}, Q^2) \frac{1}{\pi \langle p_{\perp, q \rightarrow h}^2 \rangle} e^{-p_\perp^2 / \langle p_{\perp, q \rightarrow h}^2 \rangle}$$

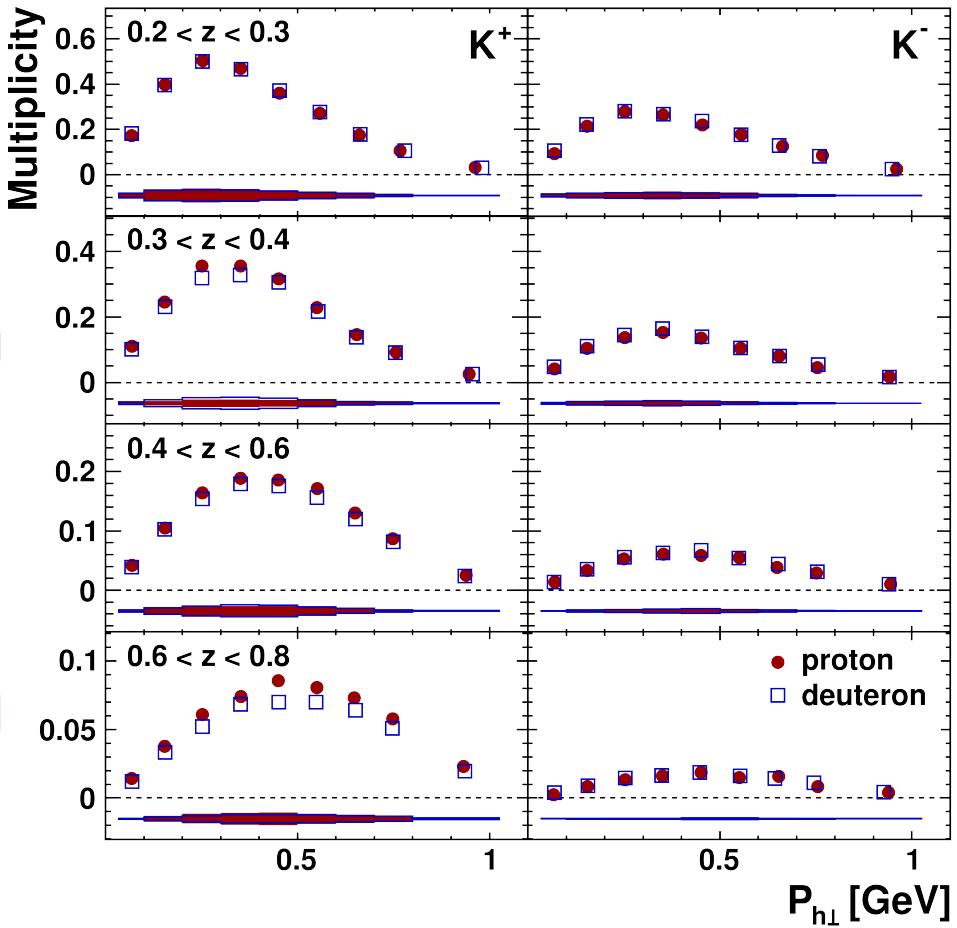
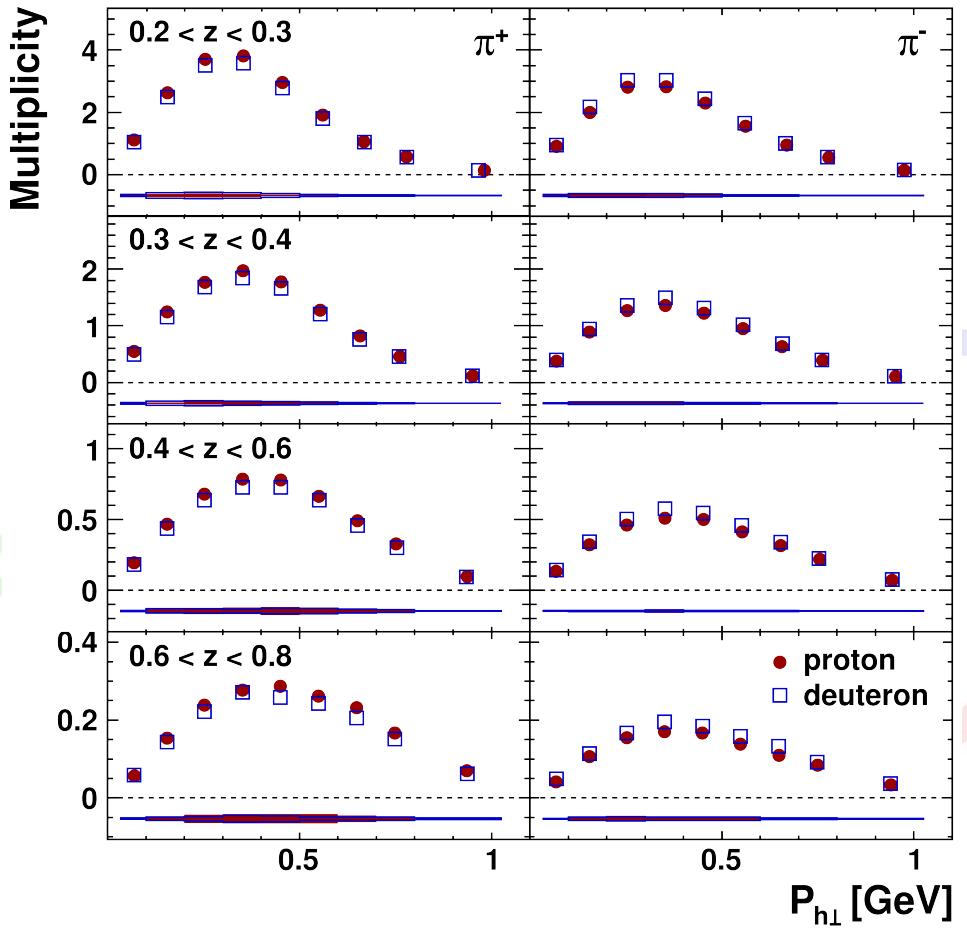


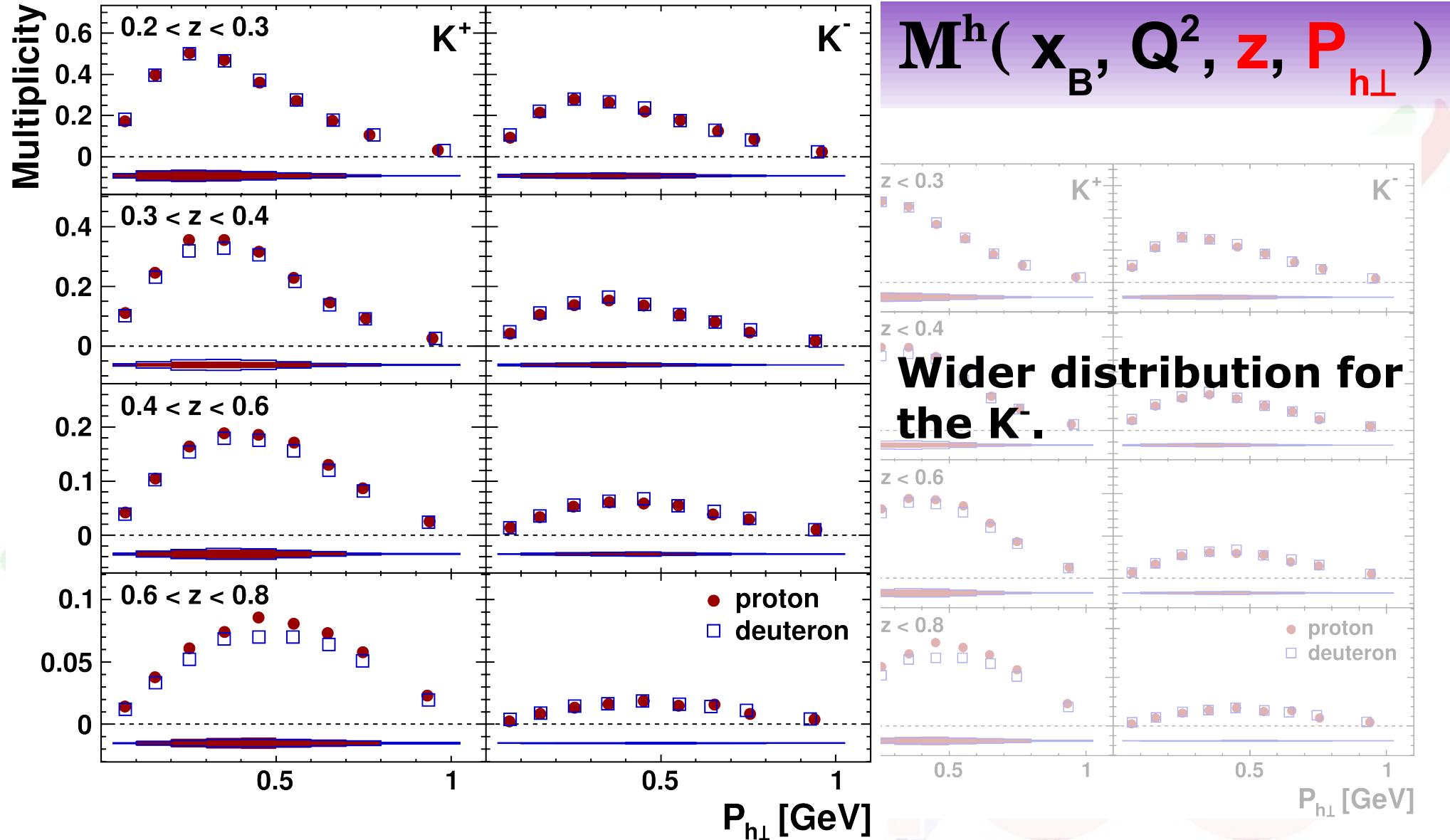
$$M^h(x_B, Q^2, z, P_{h\perp})$$

Arises from combined effect : initial transverse motion of the struck quark and the transverse momentum component generated by the fragmentation process :

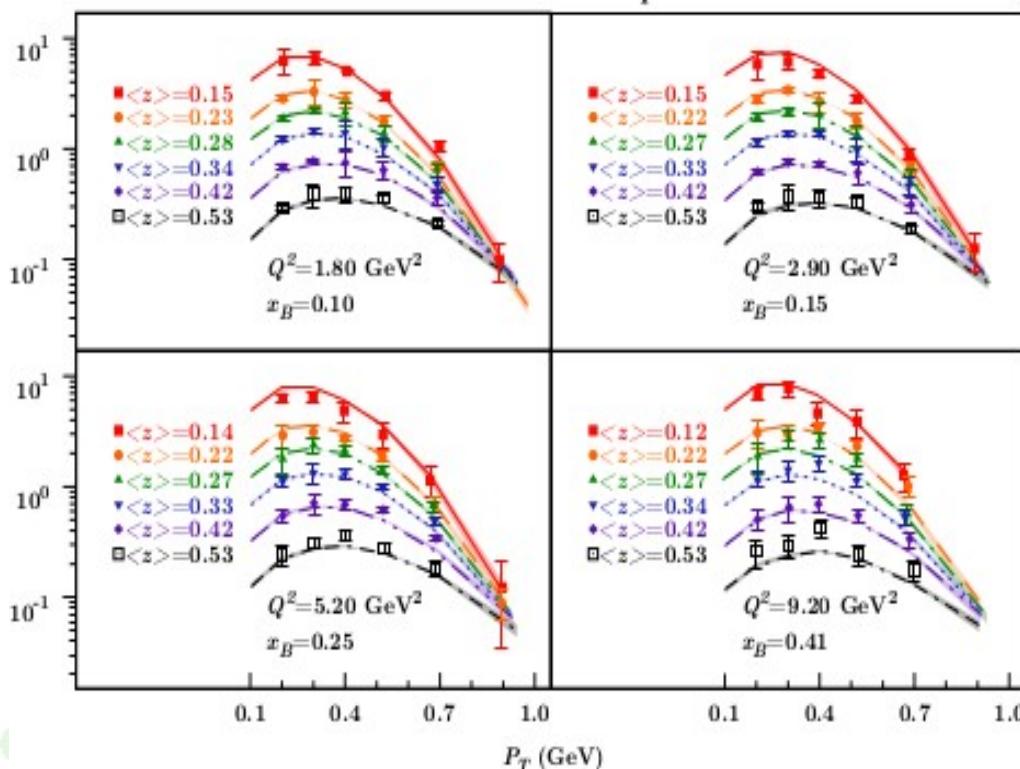
$$\vec{P}_{h\perp} = z \vec{k}_\perp + \vec{p}_\perp - O\left(\frac{\vec{k}_\perp^2}{Q^2}\right)$$

$$M^h(x_B, Q^2, z, P_{h\perp})$$



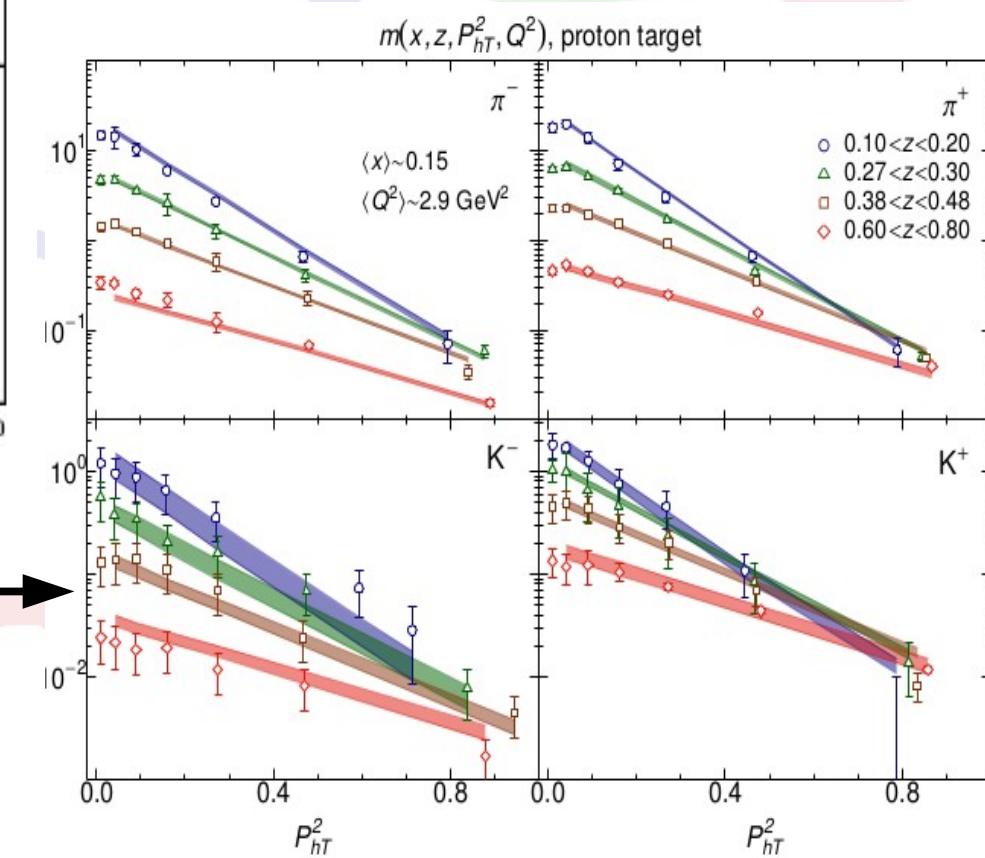


HERMES M_p^π



M. Anselmino, M. Boglione, J.O.
Gonzalez H., S. Melis, A. Prokudin
JHEP (2014)

A. Signori, A. Bacchetta, M. Radici
and G. Schnell JHEP (2013)



Experimental observable

SIDIS hadron yields

$$M^h(x_B, Q^2, z, P_{h\perp}, \phi_h) = \frac{N^h(x_B, Q^2, z, P_{h\perp}, \phi_h)}{N^e(x_B, Q^2)}$$

DIS event yields

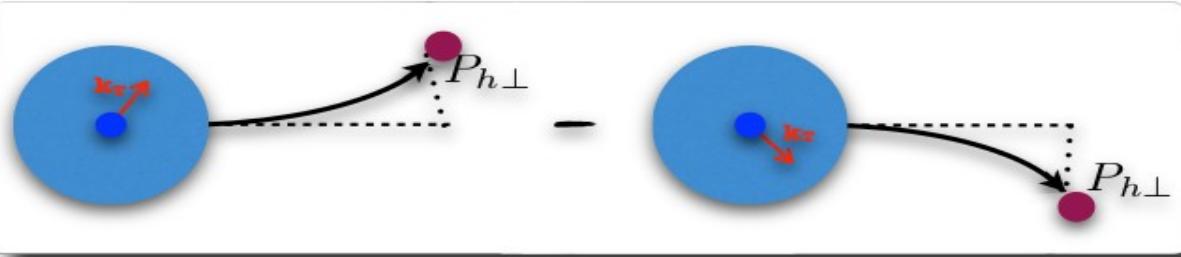
Experimental observable

SIDIS hadron yields

$$\frac{M^h(x_B, Q^2, z, \text{TMDs via } \phi_h)}{N^e(x_B, Q^2)}$$

DIS event yields

Azimuthal modulations



Cahn effect

kinematic effect caused by quark intrinsic transverse momentum.

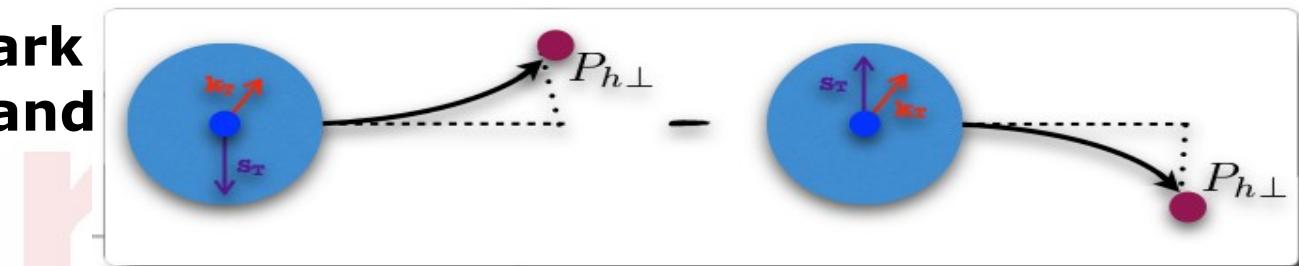
$$(\cos \phi_h)$$

R.N. Cahn, Phys. Lett. B78, (1978)

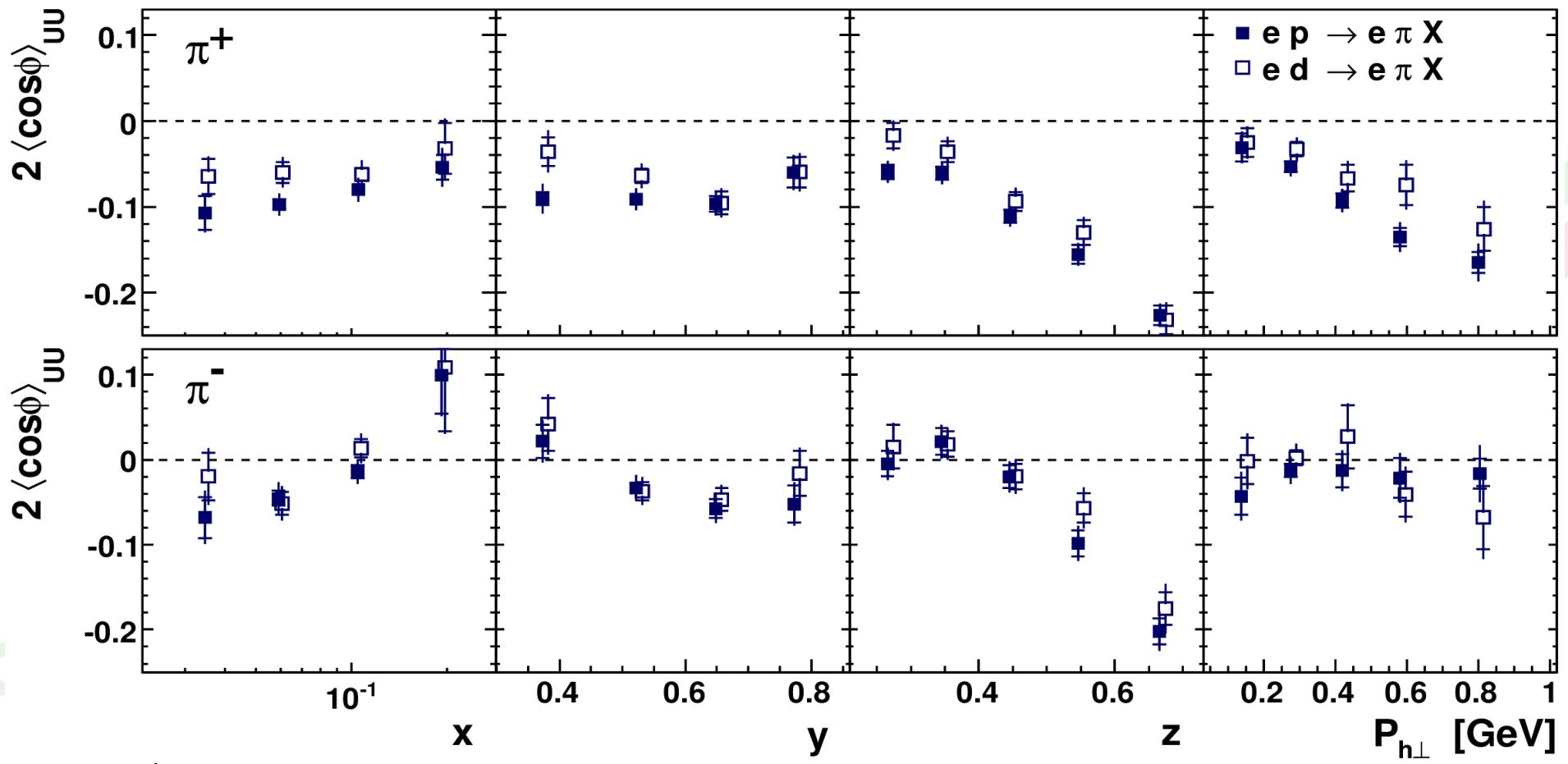
Boer-Mulders effect

correlation between quark transverse momentum and quark transverse spin.

$$(\cos 2 \phi_h)$$



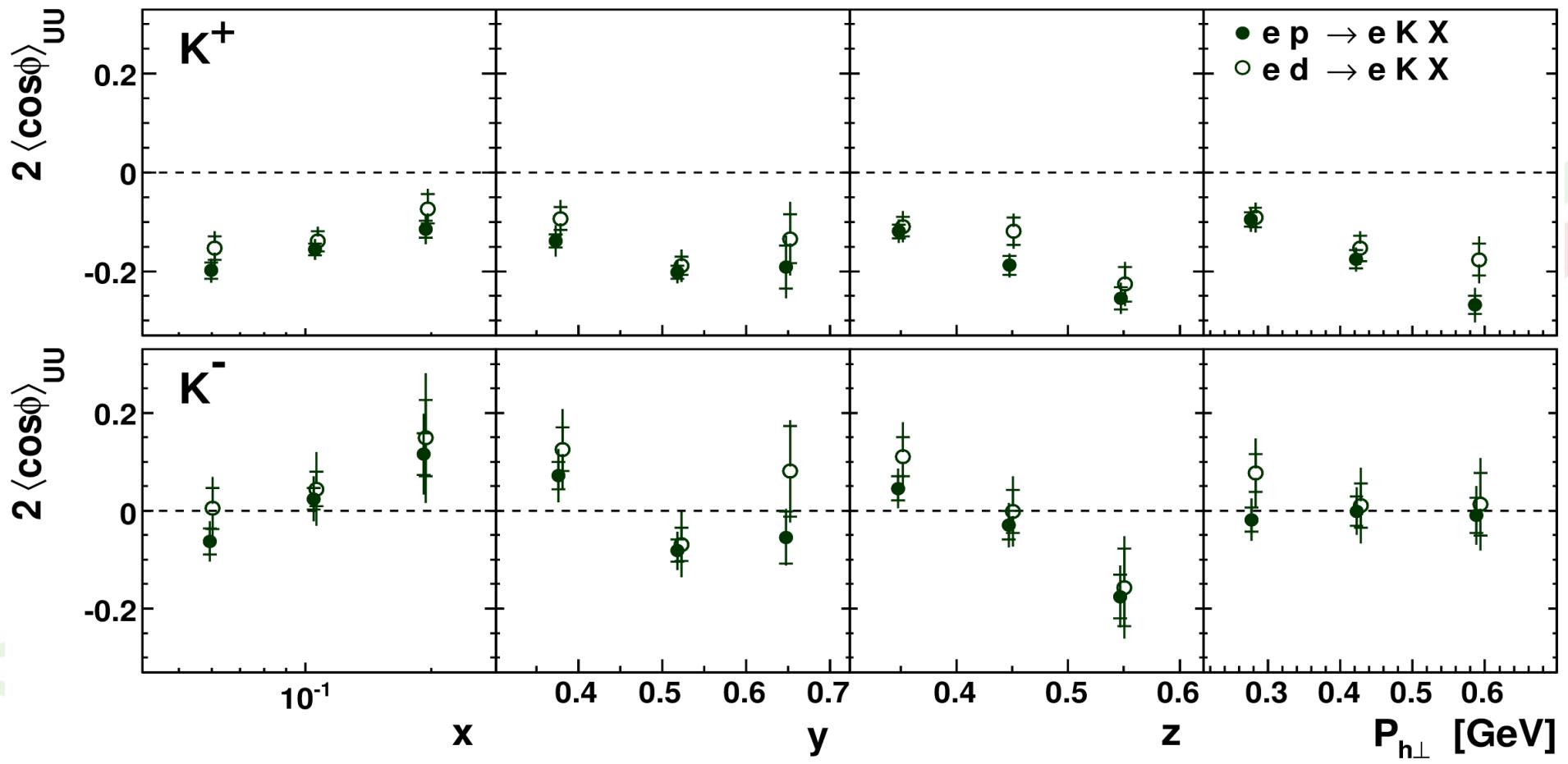
D. Boer and P.J. Mulders, Phys. Rev. D57, (1998)



Both π^+ and π^- have negative asymmetry amplitudes.

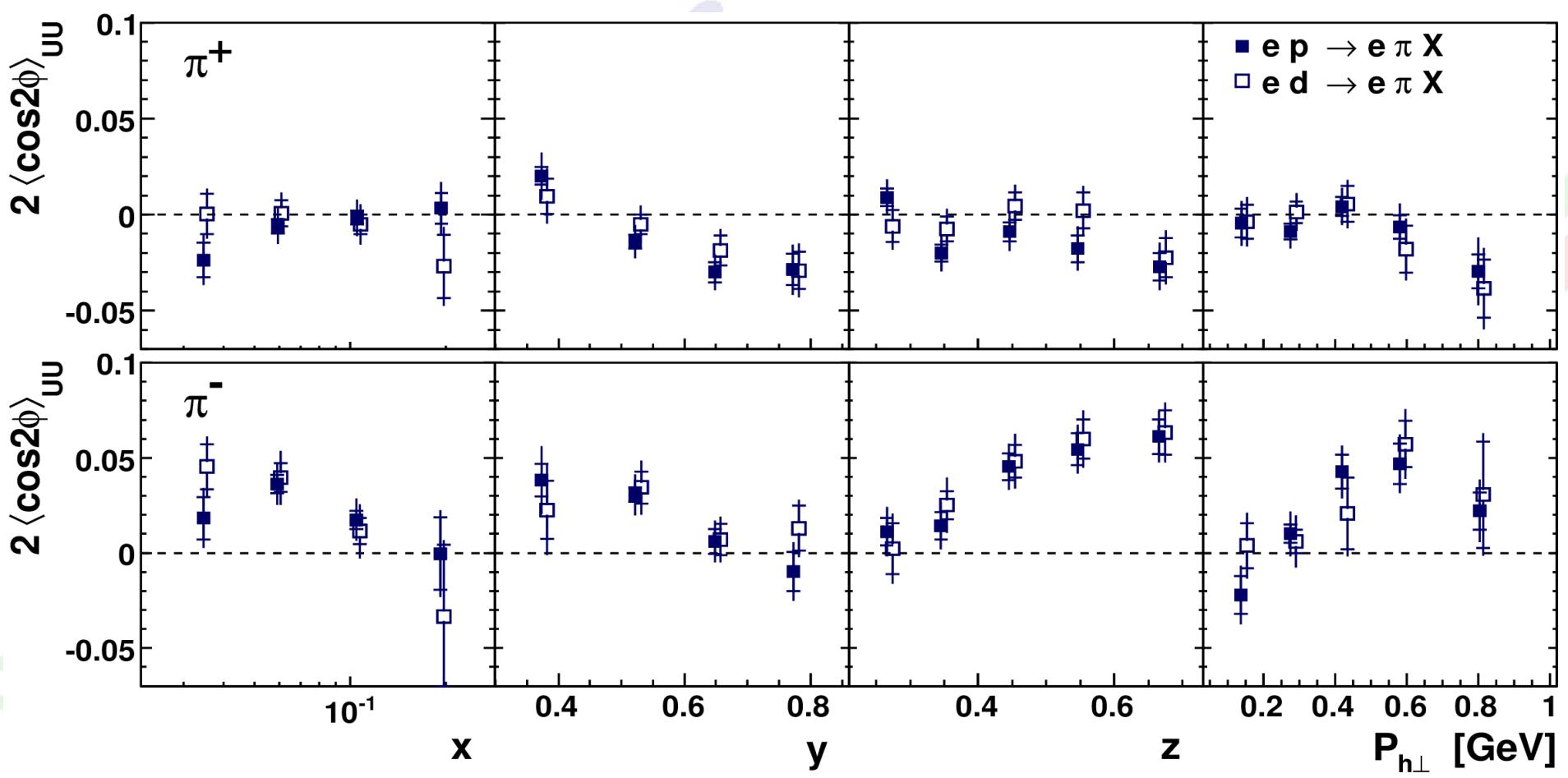
The magnitude of π^+ asymmetry amplitude is smaller on deuterium than on hydrogen target.

For both particles the asymmetry amplitude is the largest in magnitude at high z .



Large negative amplitudes for positively charged kaons.

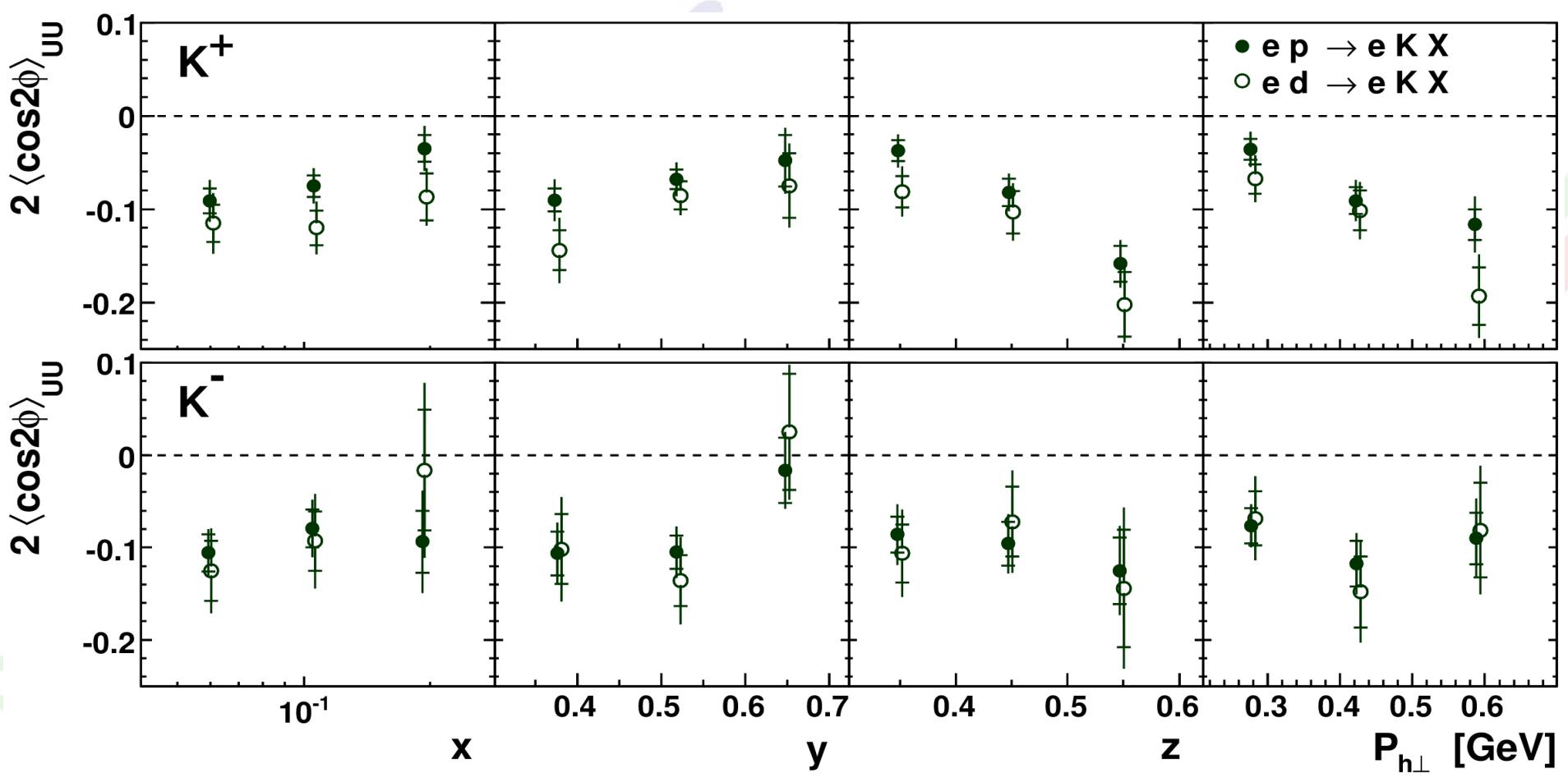
Similar results for hydrogen and deuterium targets.



The amplitudes for oppositely charged pions have the opposite signs.

π^+ and π^- asymmetry amplitudes have different magnitudes.

Similar results for hydrogen and deuterium targets.



The asymmetry amplitudes have negative sign for both K⁺ and K⁻.

Both kaons have large asymmetry amplitudes in magnitude.

Similar results for hydrogen and deuterium targets.

(hadron multiplicities)

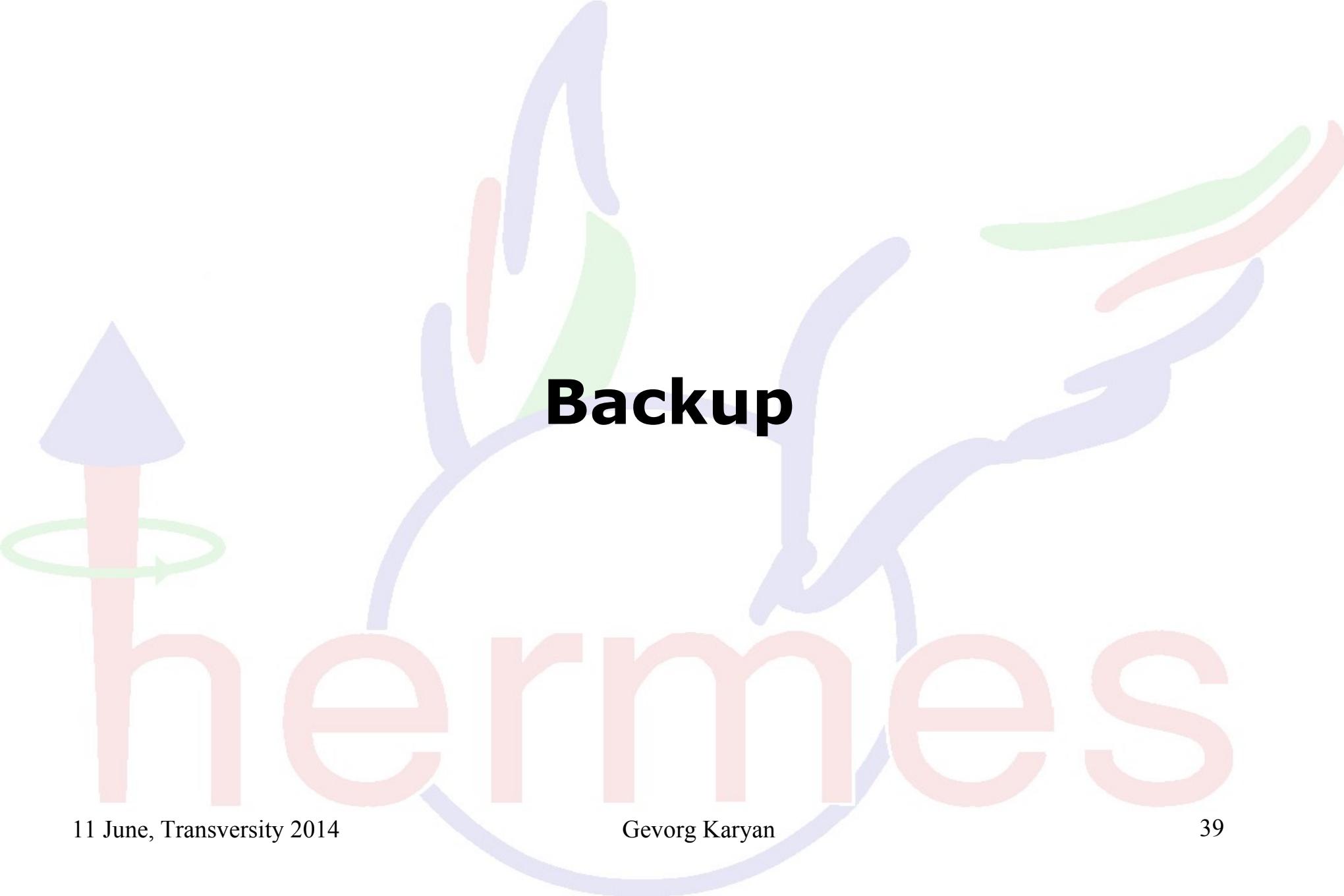
A. Airapetian et. al, Phys. Rev. D87 (2013) 074029
<http://www-hermes.desy.de/multiplicities>

(hadron azimuthal modulations)

A. Airapetian et. al, Phys. Rev. D87 (2013) 012010
<http://www-hermes.desy.de/cosnphi>

Summary

- High statistical data set for positively/negatively charged pion and kaon multiplicities on proton and deuteron.
- The extracted multiplicities integrated over hadron transverse momentum give an access to collinear fragmentation functions.
- The azimuthal modulations of produced hadrons indicate the presence of non-vanishing intrinsic transverse momentum of quarks inside an unpolarized nucleon.
- Dependence of multiplicities on hadron transverse momentum provides constraints on transverse momentum dependent distribution and fragmentation functions.
- A_{LU} results are coming soon!



hermes

Backup

