

# Probing the Spin Structure of the Neutron:

## New Experimental Results on $d_2$ and $g_2$ for the Neutron from JLab

**Brad Sawatzky**  
**Jefferson Lab**

**E06-014, E12-06-121, and Hall A Collaborations**

**HiX 2014 Conference**  
**November 17—21, 2013**

# SSF Measurements at JLab

- $g_1$  measured in all halls
  - $\text{NH}_3, \text{ND}_3$  in all Halls
  - $^3\text{He}$  in Hall A
- $g_2$  in C and A
- **Duality in  $g_1$**
- **Transverse structures  $A_2$  and  $g_T$**
- **Moments and twist-3**
- Sum Rules: GDH, B-C, Bjorken
- $n$  SSF's from  $^3\text{He}$  and from  $d - p$

Inclusive Program at 6 GeV				
Experiment	Hall	Target	Measured quantity	Kinematics $Q^2$ GeV <sup>2</sup>
94-010	A	$^3\text{He}$	$A_{\parallel}, A_{\perp}$	Resonances 0.1 – 0.9
CLAS eg1a-b	B	$p, d$	$A_{\parallel}$	DIS, Resonances 0.2 – 3.5
97-103	A	$^3\text{He}$	$A_{\perp}$	DIS 0.6 – 1.4
97-110	A	$^3\text{He}$	$A_{\parallel}, A_{\perp}$	Elastic, Resonances 0.02 – 0.5
99-117	A	$^3\text{He}$	$A_{\parallel}, A_{\perp}$	DIS 2.7, 3.5, 4.8
01-006 (RSS)	C	$p, d$	$A_{\parallel}, A_{\perp}$	Resonances 1.3
01-012	A	$^3\text{He}$	$A_{\parallel}, A_{\perp}$	Resonances 1 – 4
CLAS eg4	B	$p$	$A_{\parallel}$	Elastic, Resonances 0.01 – 0.5
07-003 (SANE)	C	$p$	$A_{\parallel}, A_{\perp}$	DIS, Resonances 1.6 – 6
06-014	A	$^3\text{He}$	$A_{\parallel}, A_{\perp}$	DIS <3>
08-027 (g2p)	A	$p$	$A_{\parallel}, A_{\perp}$	Resonances 0.03 – 0.3

Slide from O. Rondon

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Slide from O. Rondon

# Polarized DIS cross sections

$$\frac{d^2\sigma}{dE'd\Omega}(\downarrow\uparrow - \uparrow\uparrow) = \frac{4\alpha^2}{MQ^2} \frac{E'}{\nu E} \left[ (E + E' \cos \theta) g_1(x, Q^2) - \frac{Q^2}{\nu} g_2(x, Q^2) \right] = \Delta\sigma_{\parallel}$$

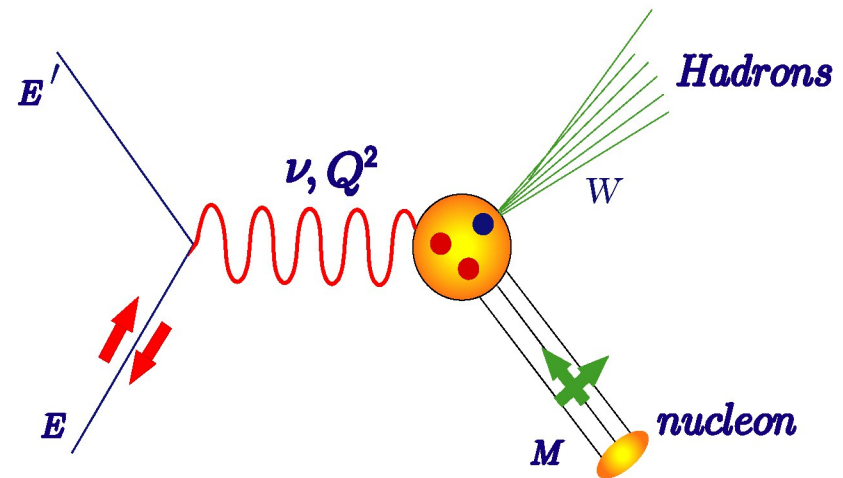
$$\frac{d^2\sigma}{dE'd\Omega}(\downarrow\Rightarrow - \uparrow\Rightarrow) = \frac{4\alpha^2 \sin \theta}{MQ^2} \frac{E'^2}{\nu^2 E} \left[ \nu g_1(x, Q^2) + 2E g_2(x, Q^2) \right] = \Delta\sigma_{\perp}$$

$Q^2$  = 4-momentum transfer squared of the virtual photon.

$\nu$  = energy transfer.

$\theta$  = scattering angle.

$x = \frac{Q^2}{2M\nu}$  fraction of nucleon momentum carried by the struck quark.

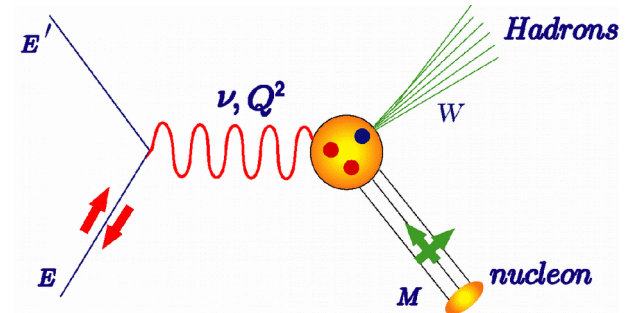


# What are $g_1$ and $g_2$ ?

- The “g's” play a role analogous to the “F's” in the unpolarized cross section

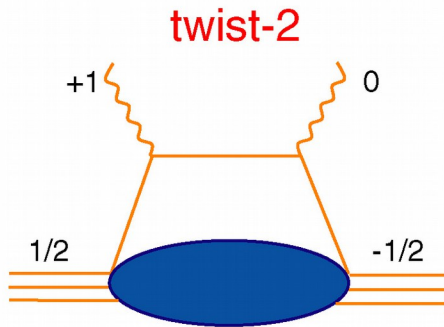
$$\frac{d^2\sigma}{d\Omega dE'} = \frac{\alpha^2}{4E^2 \sin^4 \frac{\theta}{2}} \left( \frac{2}{M} F_1(x, Q^2) \sin^2 \frac{\theta}{2} + \frac{1}{\nu} F_2(x, Q^2) \cos^2 \frac{\theta}{2} \right)$$

- $F$  encodes information about the **momentum structure** of the nucleon
- $g_1$  and  $g_2$  encode information about the **spin structure** of the target nucleon



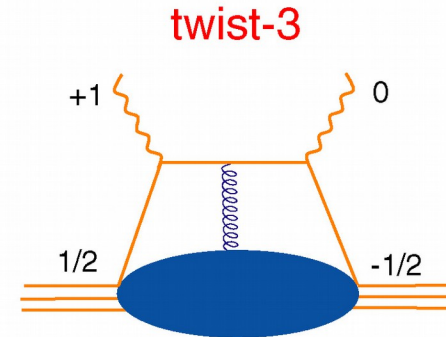
- The Parton Model
  - $g_1$  is a measure of the spin distribution among the individual constituent quarks (ie. aligned parallel and anti-parallel to the nucleon spin)
  - $g_2$  ???

# $g_2$ and Quark-Gluon Correlations



Carry one unit of orbital angular momentum

QCD allows the helicity exchange to occur in two principle ways



Couple to a gluon

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

- a twist-2 term (Wandzura & Wilczek, 1977):

$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 g_1(y, Q^2) \frac{dy}{y}$$

- a twist-3 term with a suppressed twist-2 piece (Cortes, Pire & Ralston, 92):

$$\bar{g}_2(x, Q^2) = -\int_x^1 \frac{\partial}{\partial y} \left( \frac{m_q}{M} h_T(y, Q^2) + \xi(y, Q^2) \right) \frac{dy}{y}$$

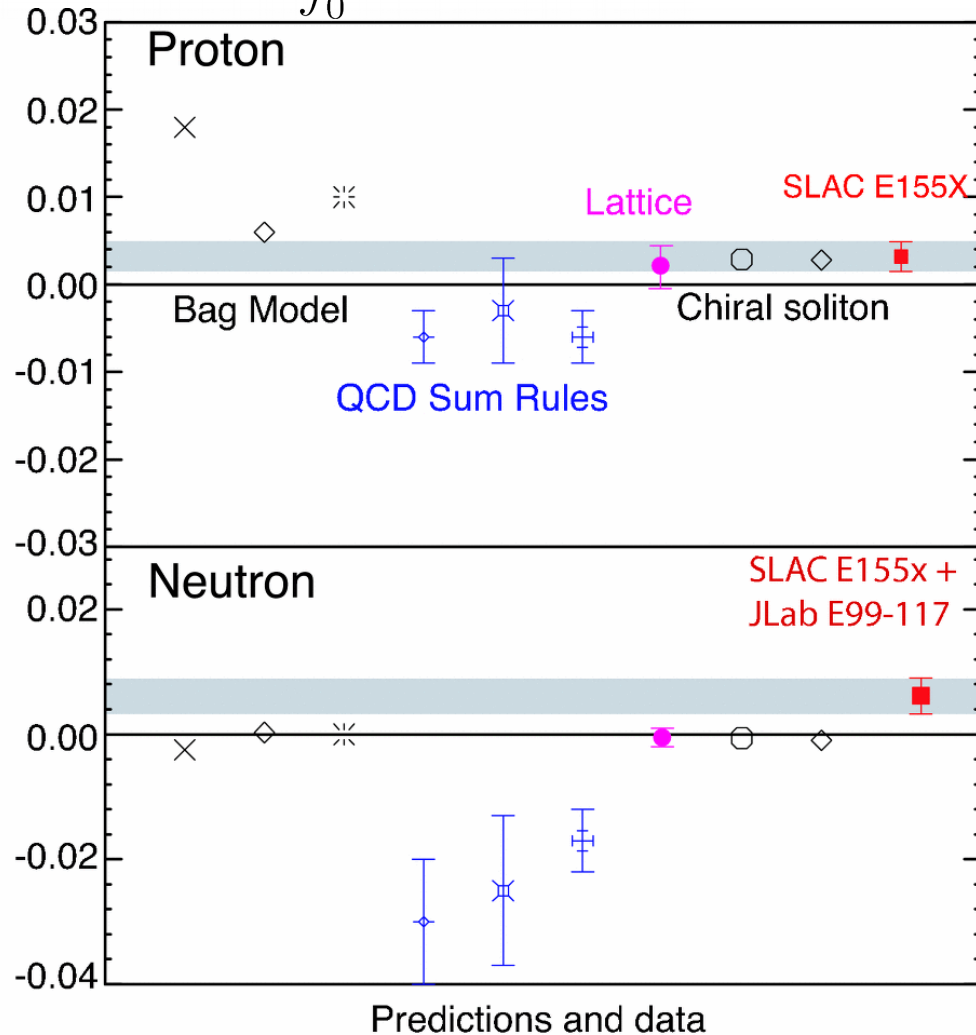
transversity

quark-gluon correlation

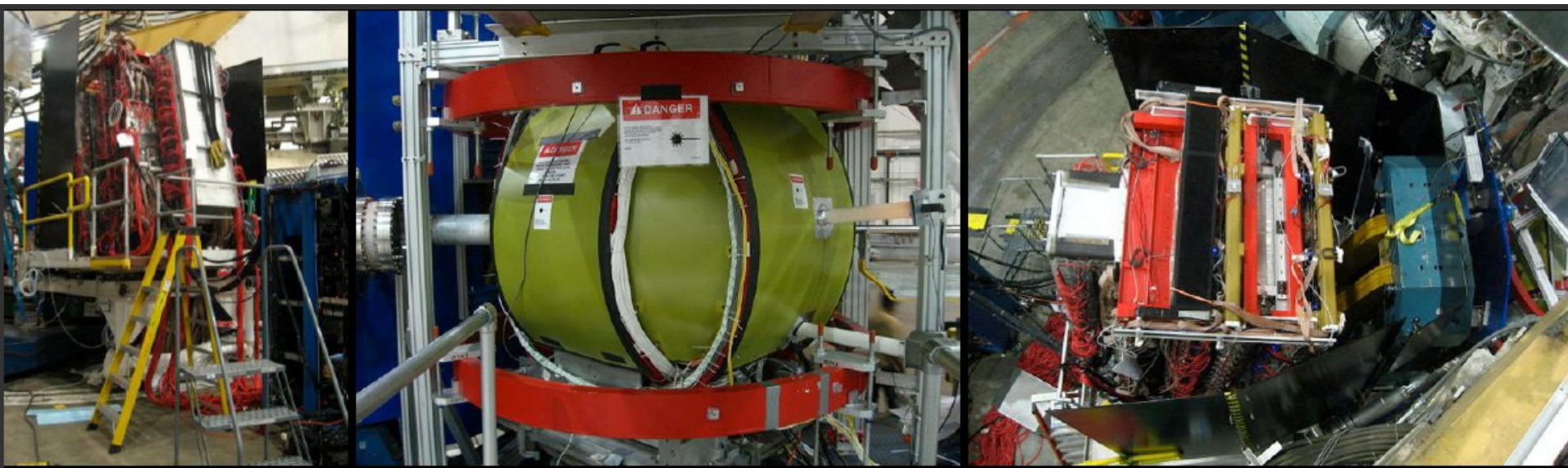
# $d_2$ : A clean probe of quark-gluon correlations

$$d_2(Q^2) = \int_0^1 x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)] dx = 3 \int_0^1 x^2 \bar{g}_2(x, Q^2) dx$$

- $d_2$  is a clean probe of **quark-gluon correlations / higher twist effects**
  - $d_2$  is the **2<sup>nd</sup> moment** of a sum of the spin structure functions
  - **matrix element** in the Operator Product Expansion  $d_2^2$ 
    - » *it is cleanly computable using Lattice QCD*
- Connected to the **color Lorentz (transverse) force** acting on the struck quark (Burkardt)
  - same underlying physics as in SIDIS  $k_\perp$  studies



# E06-014: The Neutron $d_2$ (Hall A)



- A measurement of the neutron  $d_2$ 
  - Polarized  $^3\text{He}$  target
  - Large acceptance detector to measure asymms (BigBite)
  - High-precision device to measure unpol. x-sec (HRS)
  - Focus:  $d_2$ ,  $g_2$  on the neutron
    - » extracted  $A_1$ ,  $g_1$  as well



# The Experiment

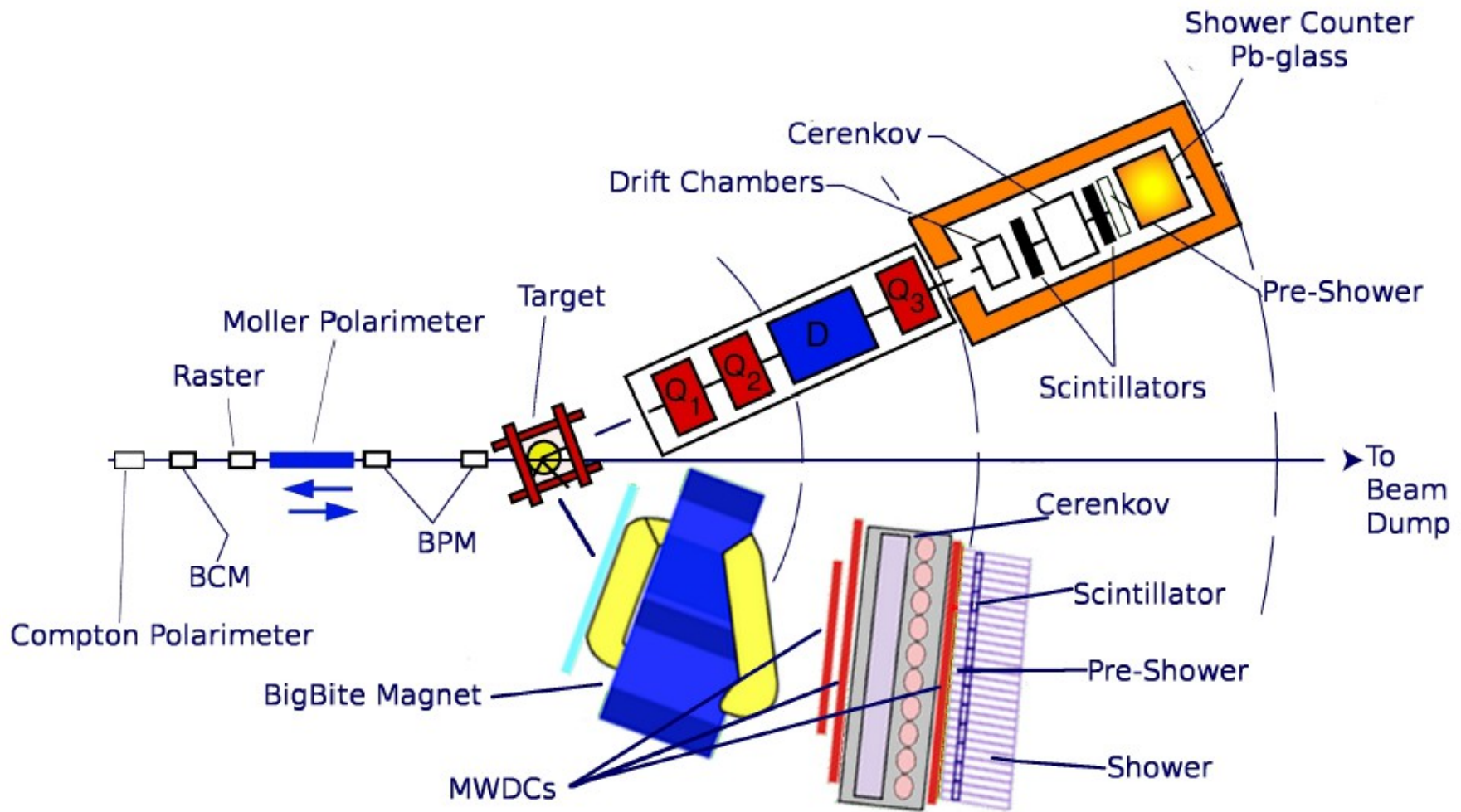
- A 4.75 and 5.9 GeV polarized electron beam scattering off a polarized  $^3\text{He}$  target
- Measure unpolarized cross section for  $^3\vec{\text{H}}\text{e}(\vec{e}, e')$  reaction  $\sigma_0^{^3\text{He}}$  in conjunction with the transverse asymmetry  $A_{\perp}^{^3\text{He}}$  and the parallel asymmetry  $A_{\parallel}^{^3\text{He}}$  for  $0.23 < x < 0.65$  with  $2 < Q^2 < 5 \text{ GeV}^2$ .
  - Asymmetries measured by BigBite
  - Absolute cross sections measured by L-HRS
- Determine  $d_2^n$  using the relation

$$\begin{aligned} \tilde{d}_2(x, Q^2) &= x^2[2g_1(x, Q^2) + 3g_2(x, Q^2)] \\ &= \frac{MQ^2}{4\alpha^2} \frac{x^2 y^2}{(1-y)(2-y)} \sigma_0 \left[ \left( 3 \frac{1 + (1-y)\cos\theta}{(1-y)\sin\theta} + \frac{4}{y} \tan\frac{\theta}{2} \right) A_{\perp} + \left( \frac{4}{y} - 3 \right) A_{\parallel} \right] \end{aligned}$$

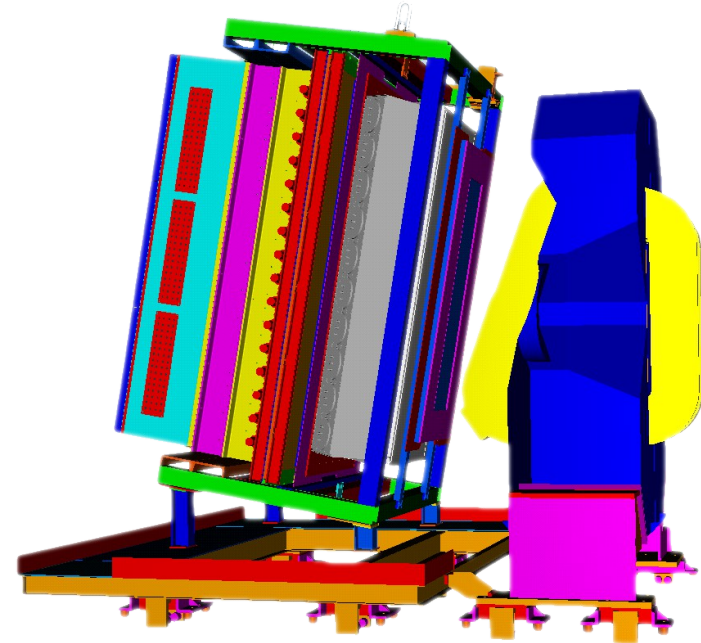
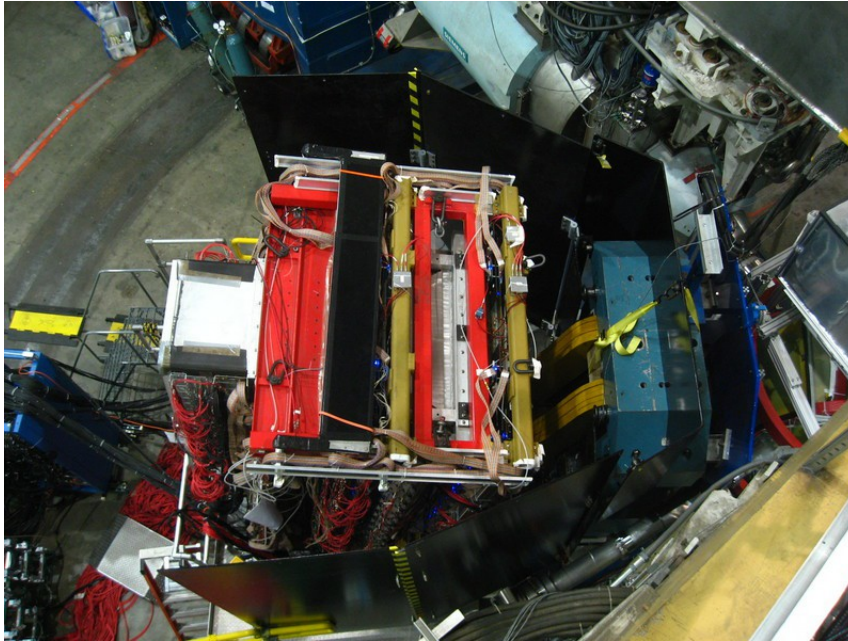
where,

$$\begin{aligned} A_{\perp} &= \frac{\sigma^{\downarrow\Rightarrow} - \sigma^{\uparrow\Rightarrow}}{2\sigma_0} & A_{\parallel} &= \frac{\sigma^{\downarrow\uparrow} - \sigma^{\uparrow\uparrow}}{2\sigma_0} \\ A_{\perp}^{^3\text{He}} &= \frac{\Delta_{\perp}}{P_b P_t \cos\phi} & A_{\parallel}^{^3\text{He}} &= \frac{\Delta_{\parallel}}{P_b P_t} \\ \Delta_{\perp} &= \frac{(N^{\downarrow\Rightarrow} - N^{\uparrow\Rightarrow})}{(N^{\downarrow\Rightarrow} + N^{\uparrow\Rightarrow})} & \Delta_{\parallel} &= \frac{(N^{\downarrow\uparrow} - N^{\uparrow\uparrow})}{(N^{\downarrow\uparrow} + N^{\uparrow\uparrow})} \end{aligned}$$

# Floor configuration for $d_2^n$



# BigBite Electron Stack



- **BigBite detector package:**

- 3 Multi-wire drift chambers (MWDC)
- Scintillator plane
- Pb-glass Calorimeter (Pre-shower + Shower)
- Gas Cherenkov

Tracking

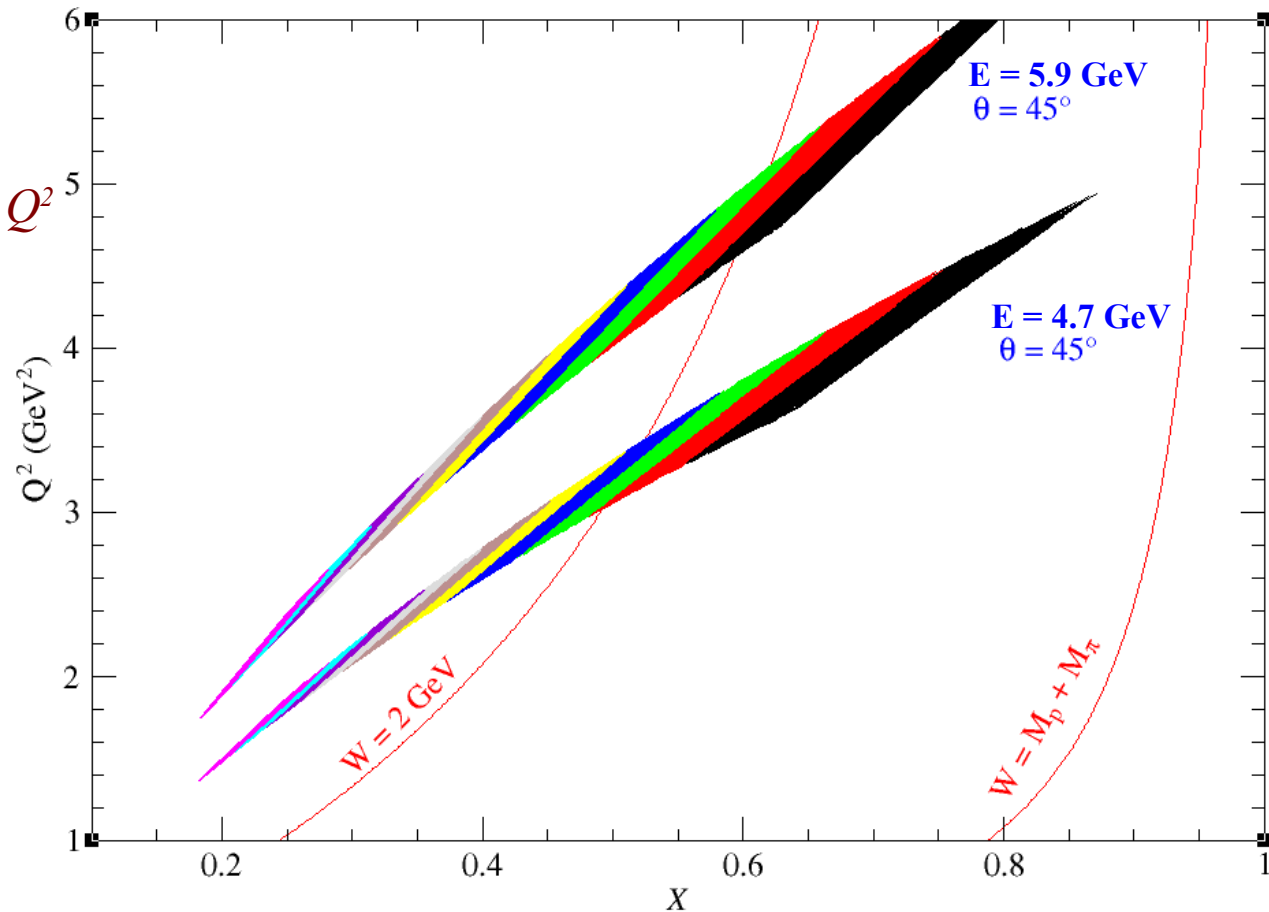
Timing

Energy/PID

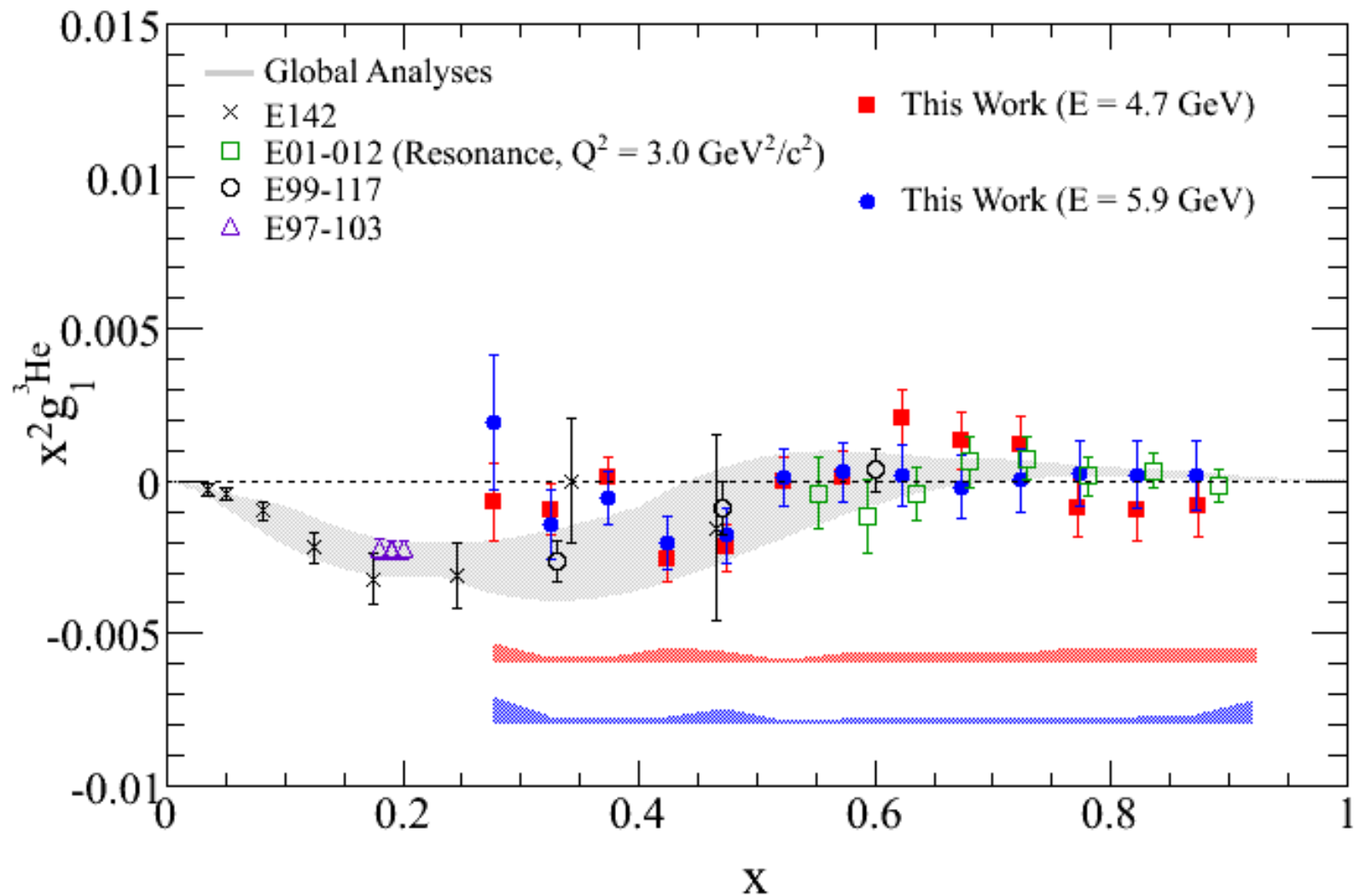
PID

# Kinematics of the measurement

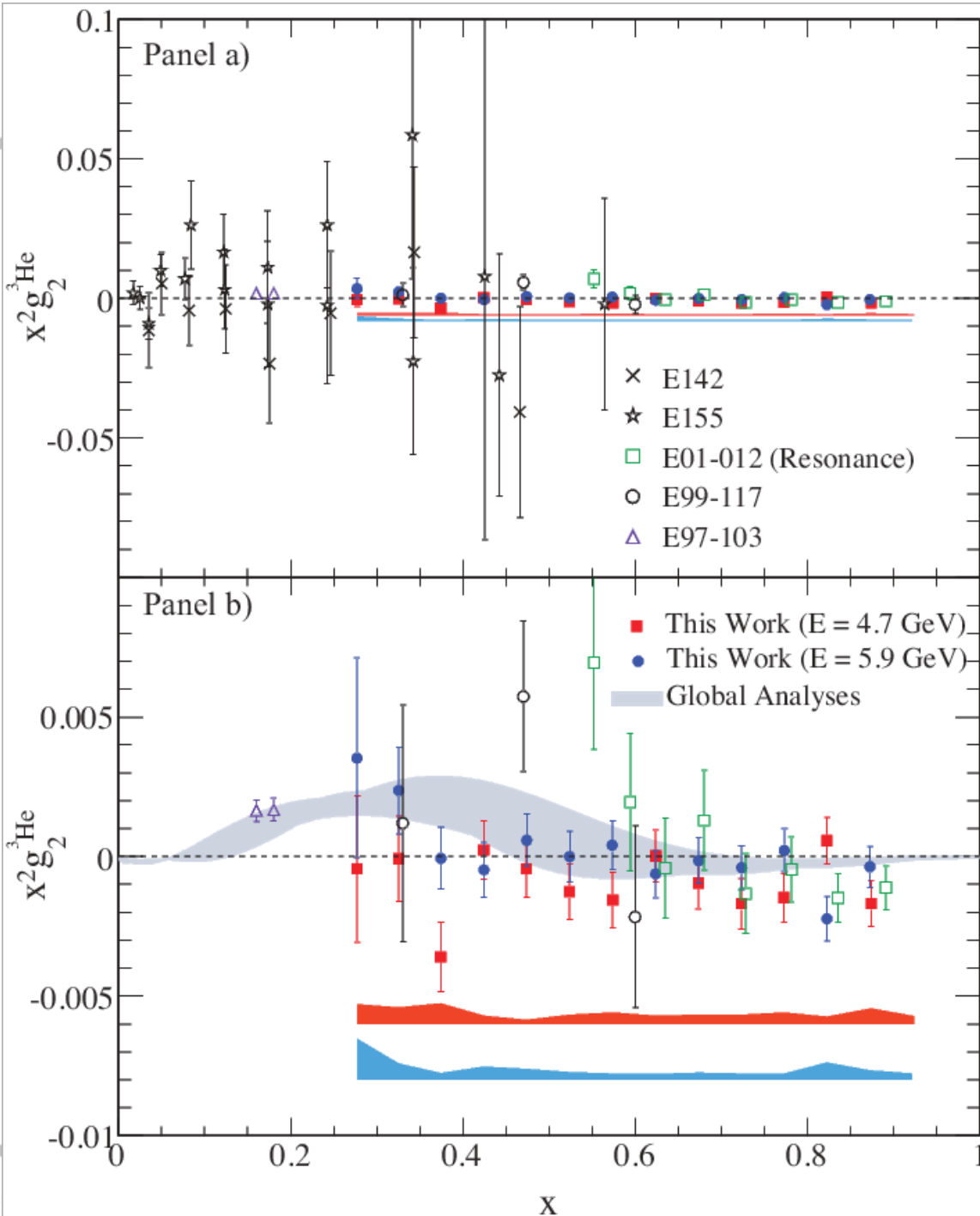
- Two beam energies  
4.75 and 5.9 GeV  
(4 pass, 5 pass)
  - provides a handle on the  $Q^2$  dependence of  $g_2$
  - supports rad. correction calculations
- BigBite fixed at single scattering angle ( $\theta=45^\circ$ )  
(data divided into bins during analysis)
- Avoid resonance region as much as possible.



# $x^2g_1$ for $^3\text{He}$



# $x^2g_2$ for $^3\text{He}$



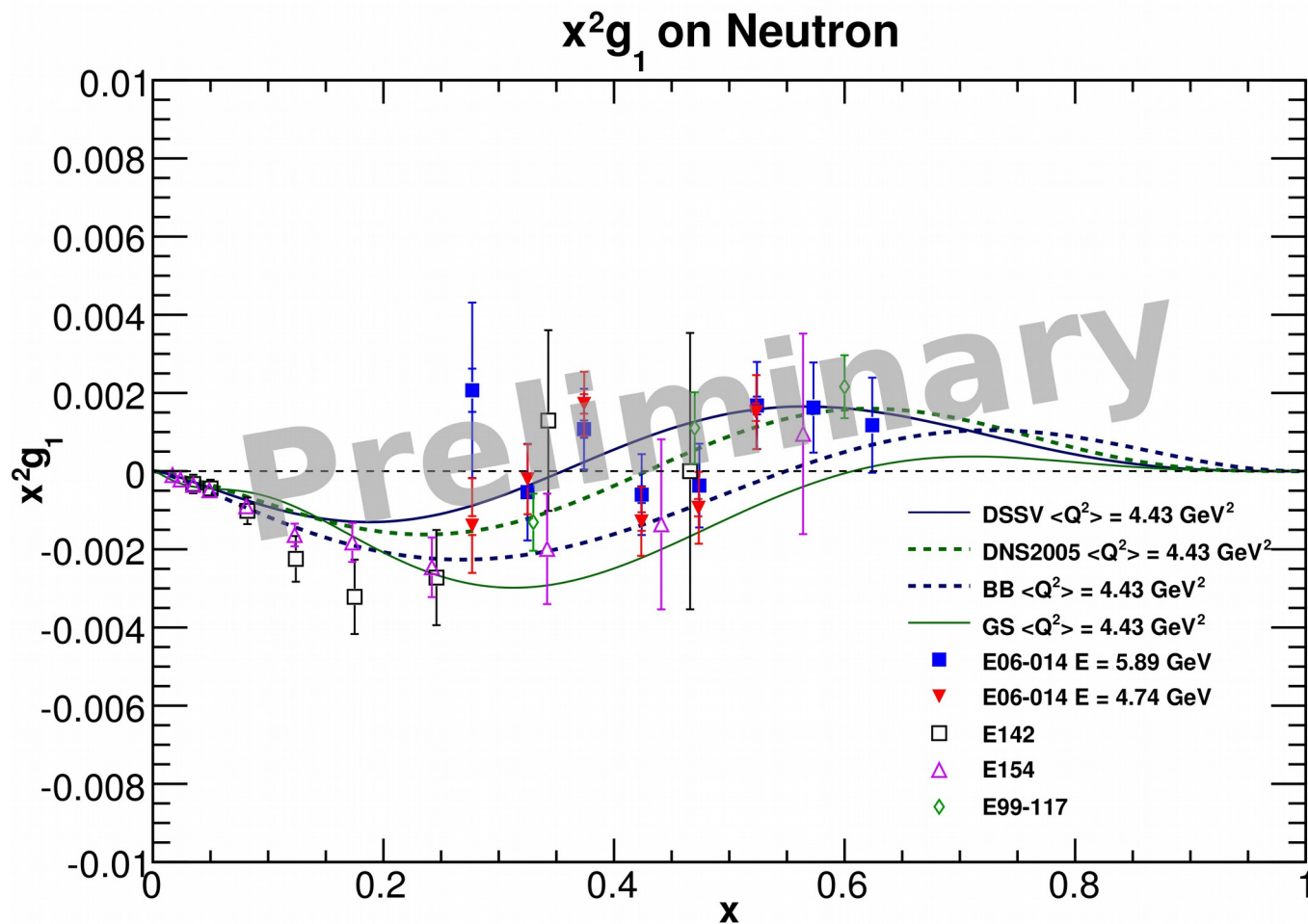
- Panel (a) shows comparison to world data

- Panel (b) is zoomed by 10x on y-axis to show error bars

- Global analyses envelope includes  $g_2^{ww}$  by

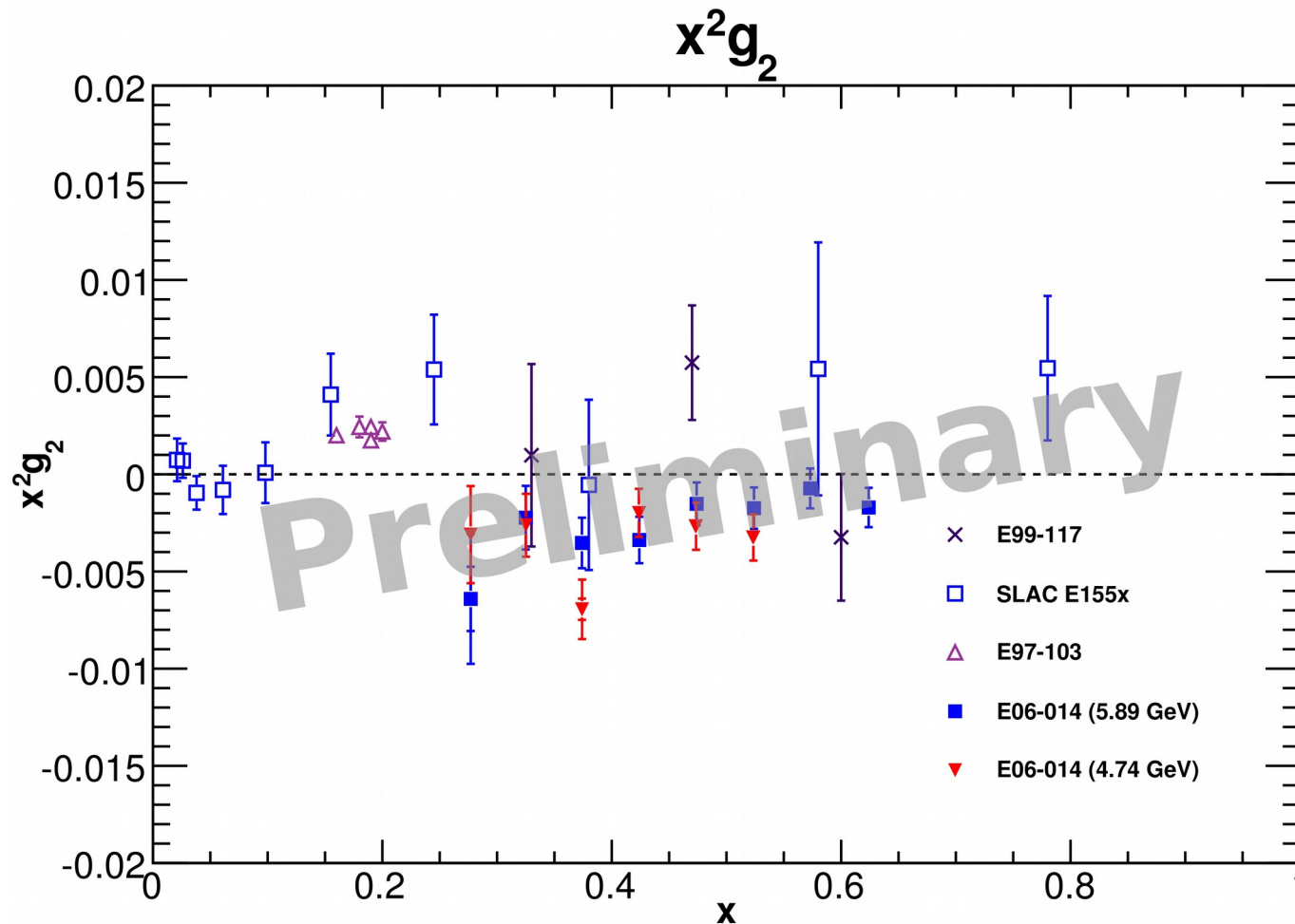
→ Bourelly, Soffer, Buccella; Stratmann, Thomas, Bissey, de Florian, Navarro, Sassot; Song

# $x^2g_1$ for the Neutron



- $3\text{He} \rightarrow$  neutron extraction done using effective polarization model (*very preliminary*)  
→ more sophisticated deconvolution method in progress (Melnitchouk, *et al*)

# $x^2g_2$ for the Neutron



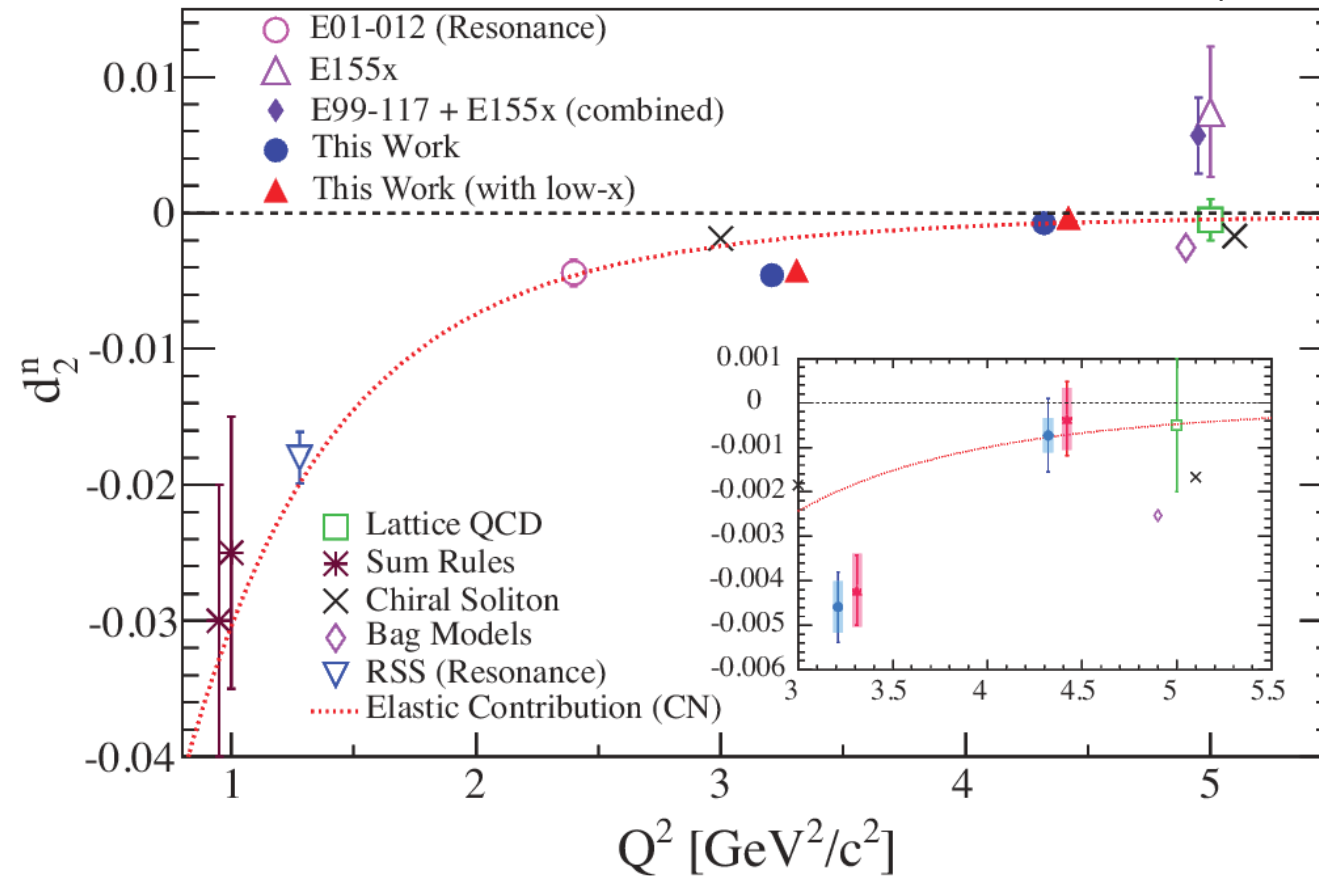
- $^3\text{He} \rightarrow$  neutron extraction done using effective polarization model (*very preliminary*)  
→ more sophisticated deconvolution method in progress (Melnitchouk, *et al*)



# $d_2$ for the Neutron

$$d_2(Q^2) = \int_0^1 x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)] dx = 3 \int_0^1 x^2 \bar{g}_2(x, Q^2) dx$$

Posik et al., PRL 113 022002 (2014)



- Our results are consistent with Lattice QCD prediction

- $d_2^n$  extracted at
  - $\langle Q^2 \rangle \sim 3.3 \text{ GeV}^2$  (E=4.7 GeV data)
  - $\langle Q^2 \rangle \sim 4.3 \text{ GeV}^2$  (E=5.9 GeV data)

- Shaded boxes in inset are systematic uncertainties

- Low-x contribution ( $0.02 < x < 0.25$ ) is provided by fits to world data (small impact)

- $^3\text{He} \rightarrow$  neutron correction using eff. polarization method applied to  $d_2$

→ (Bissey et al. Phys Rev C, 65:064137, 2002)

Archival paper in progress (David Flay ~ Temple U.)

# Moments of Structure Functions

$$\Gamma_1(Q^2) \equiv \int_0^1 g_1(x, Q^2) dx$$

$$= \mu_2 + \frac{M^2}{9Q^2} (a_2 + 4d_2 + 4f_2) + \frac{\mu_6}{Q^4} + \dots$$

$$\mu_2(Q^2) = C_{ns}(Q^2) \left( -\frac{1}{12} g_A + \frac{1}{36} a_8 \right) + C_s(Q^2) \frac{1}{9} \Delta\Sigma$$

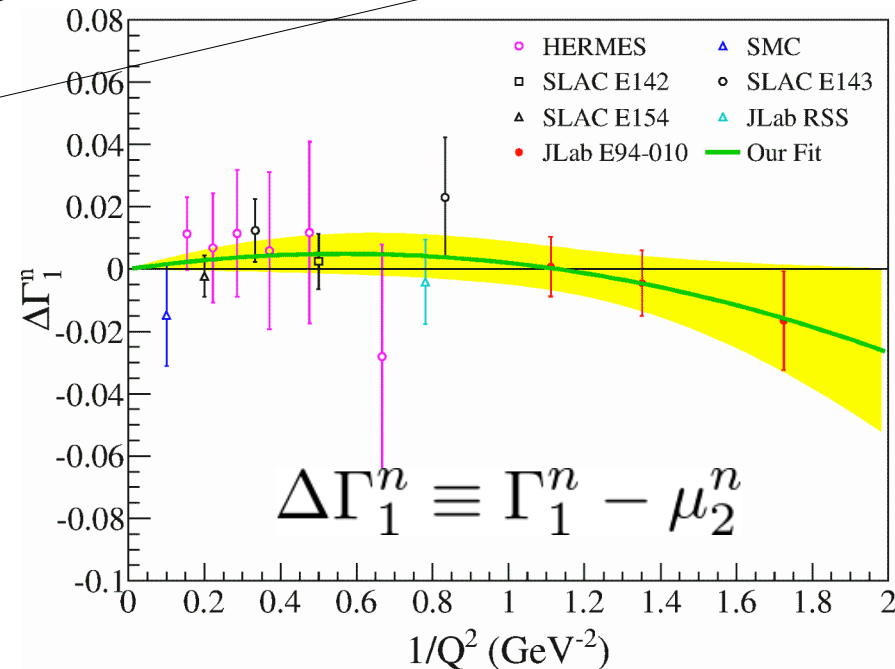
Wilson coefficients

Triplet axial charge

Octet axial charge

Singlet axial charge

(Extracted from  
neutron B decay and  
weak hyperon decay  
measurements.)



# Moments of Structure Functions

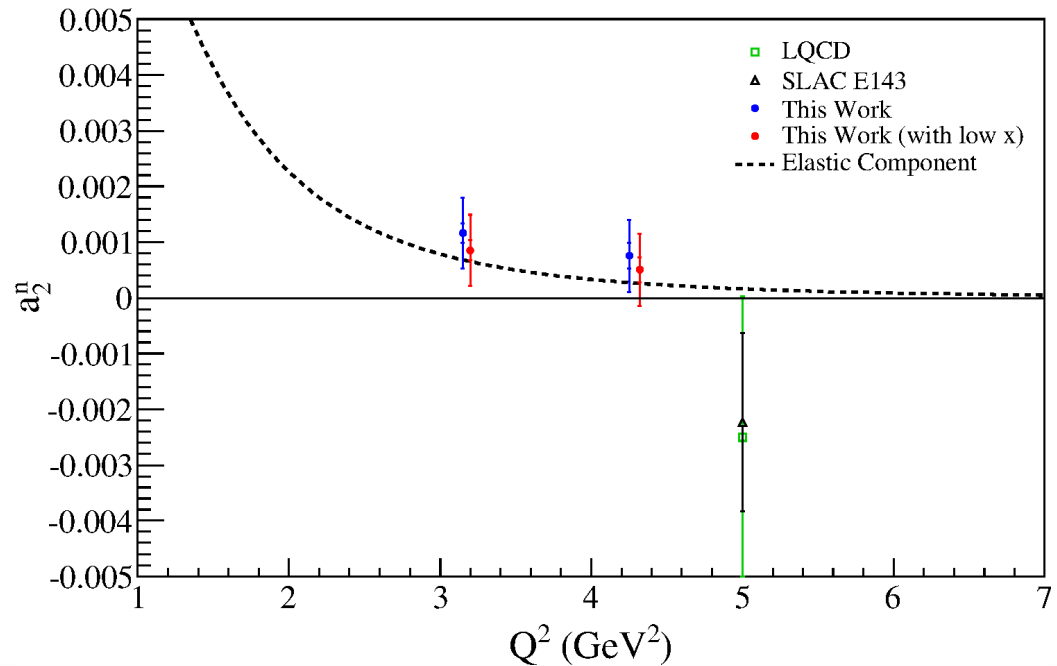
$$\Gamma_1(Q^2) \equiv \int_0^1 g_1(x, Q^2) dx$$

$$= \mu_2 + \frac{M^2}{9Q^2} (a_2 + 4d_2 + 4f_2) + \frac{\mu_6}{Q^4} + \dots$$

$\downarrow$   
 Moments of  $g_1$  and  $g_2$ !

$$a_2 = \int_0^1 x^2 g_1 dx$$

Twist-2 matrix element  
connected to target mass  
corrections.

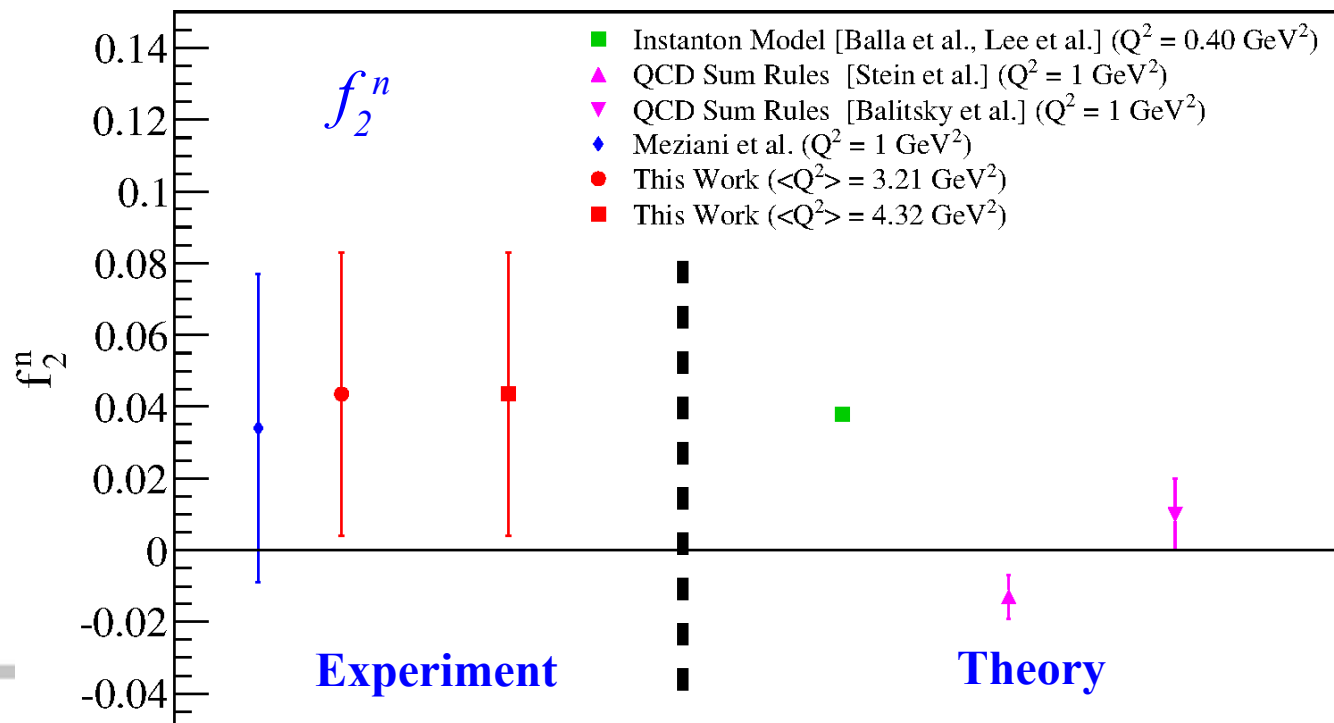


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$$\Gamma_1(Q^2) \equiv \int_0^1 g_1(x, Q^2) dx$$

$$= \mu_2 + \frac{M^2}{9Q^2} (a_2 + 4d_2 + 4f_2) + \frac{\mu_6}{Q^4} + \dots$$

- Rearrange OPE expansion and solve for  $f_2$



# Color Electric and Magnetic Forces

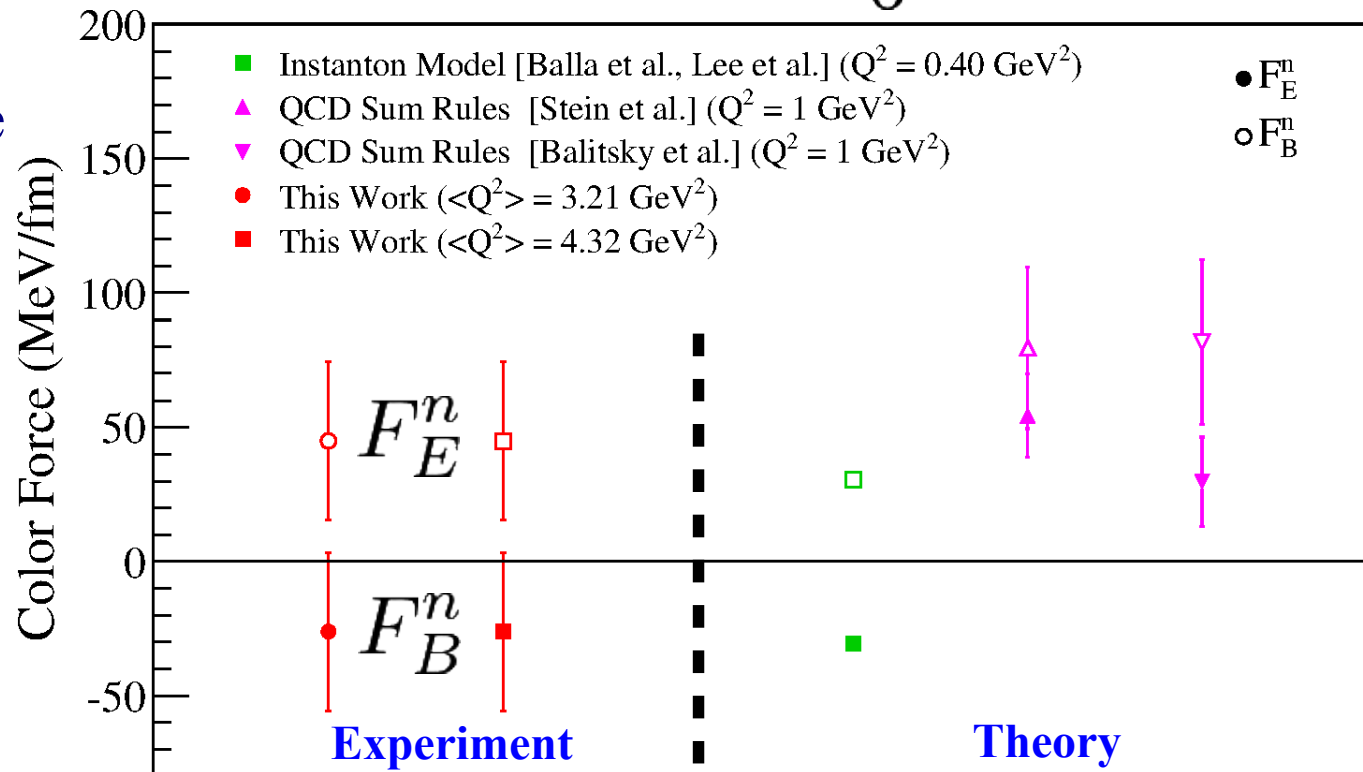
- $d_2$  and  $f_2$  are twist-3 and twist-4 matrix elements; both connected to quark-gluon correlations

$$F_E^n = -\frac{M_n^2}{6} (2d_2^n + f_2^n)$$

$$F_B^n = -\frac{M_n^2}{6} (4d_2^n - f_2^n)$$

- Together, they give information on the *Color Lorentz Electric and Magnetic forces*

Burkardt, PRD 88 114502 (2013)



# Summary for E06-014

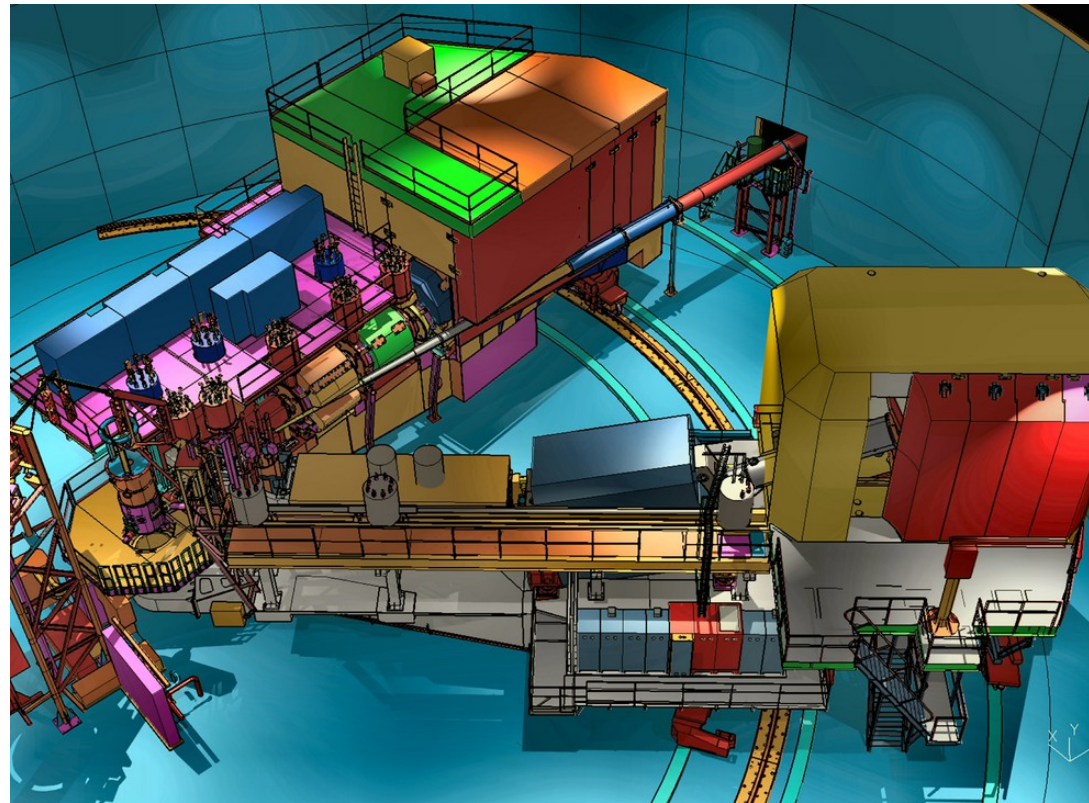
- “d2n” (E06-014) Neutron,  $^3\text{He}$  at 6 GeV
  - $A_1(^3\text{He}$  and neutron) in good agreement with world data
  - $d_2^n$  is small, negative, and consistent with LQCD prediction at  $Q_2 = 5$  (GeV/c) $^2$
  - $g_1$  and  $g_2$  in agreement with world data for both  $^3\text{He}$
  - *Color* Electric and Magnetic Forces on struck quark extracted
    - » roughly equal and opposite
- The “Future”
  - Major JLab upgrades
    - » 12 GeV beam, major new detector apparatus in Hall C (SHMS)
  - 12 GeV dedicated  $d_2^n$ ,  $g_2^n$  (and A1n) measurements are approved in both Halls A and C (tentatively scheduled to be in first “non-commissioning” run group)
    - » focus on high-x and  $Q_2$  evolution
      - E06-014 Special thanks to:
        - D. Flay, M. Posik, D. Parno, ZE Mezziani, G. Franklin

[This work is supported in part by  
DOE Award from Temple University #DE-FG02-94ER40844.]

# Hall C after 12 GeV Upgrade

- Beam Energy: 2 – 11 GeV/c
- **Super High Momentum Spectrometer (SHMS)**
  - Horizontal Bender, 3 Quads, Dipole
  - $P \rightarrow 11 \text{ GeV/c}$
  - $dP/P \text{ } 0.5\text{--}1.0 \times 10^{-3}$
  - Acceptance: 5msr, 30%
  - $5.5^\circ < \theta < 40^\circ$
- **High Momentum Spectrometer (HMS)**
  - $P \rightarrow 7.5 \text{ GeV/c}$
  - $dP/P \text{ } 0.5 - 1.0 \times 10^{-3}$
  - Acceptance: 6.5msr, 18%
  - $10.5^\circ < \theta < 90^\circ$
- Minimum opening angle:  $17^\circ$
- Well shielded detector huts

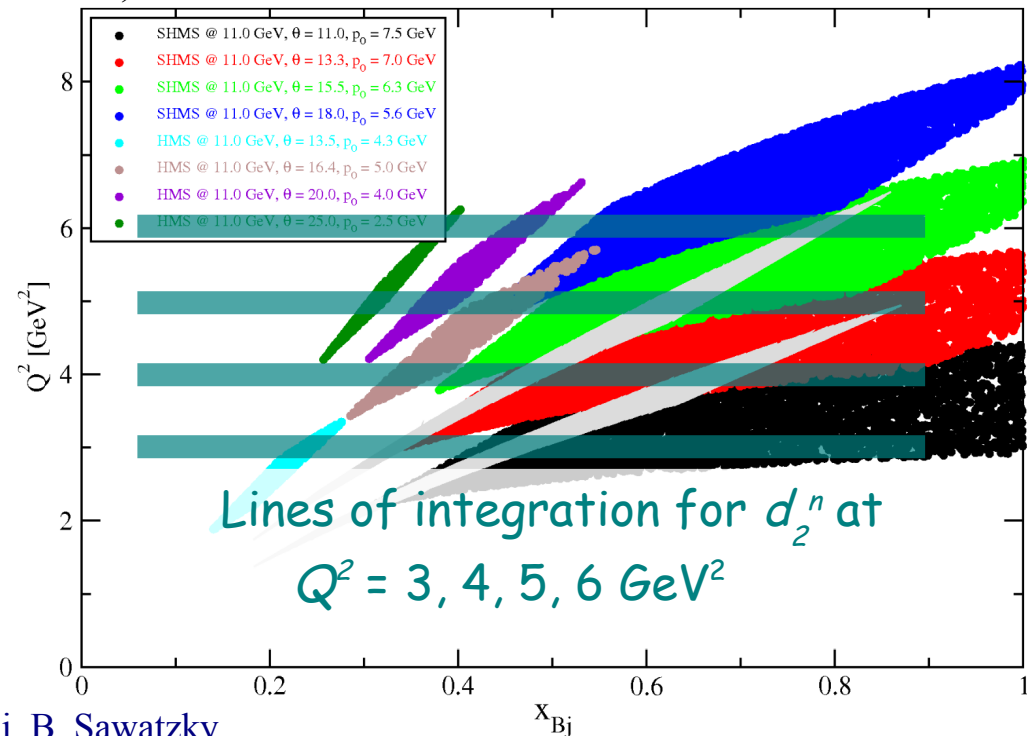
- Ideal facility for:
  - Rosenbluth (L/T) separations
  - Exclusive reactions
  - Low cross sections (neutrino level)



# E12-06-121: $d_2^n, g_2^n$

- Directly measure the  $Q^2$  dependence of the neutron  $d_2^n(Q^2)$  at  $Q^2 \approx 3, 4, 5, 6 \text{ GeV}^2$  with the new polarized  $^3\text{He}$  target.
  - The new Hall C SHMS is ideally suited to this task!
- Doubles number of precision data points for  $g_2^n(x, Q^2)$  in DIS region.
  - $Q^2$  evolution of  $g_2^n$  over  $(0.23 < x < 0.85)$

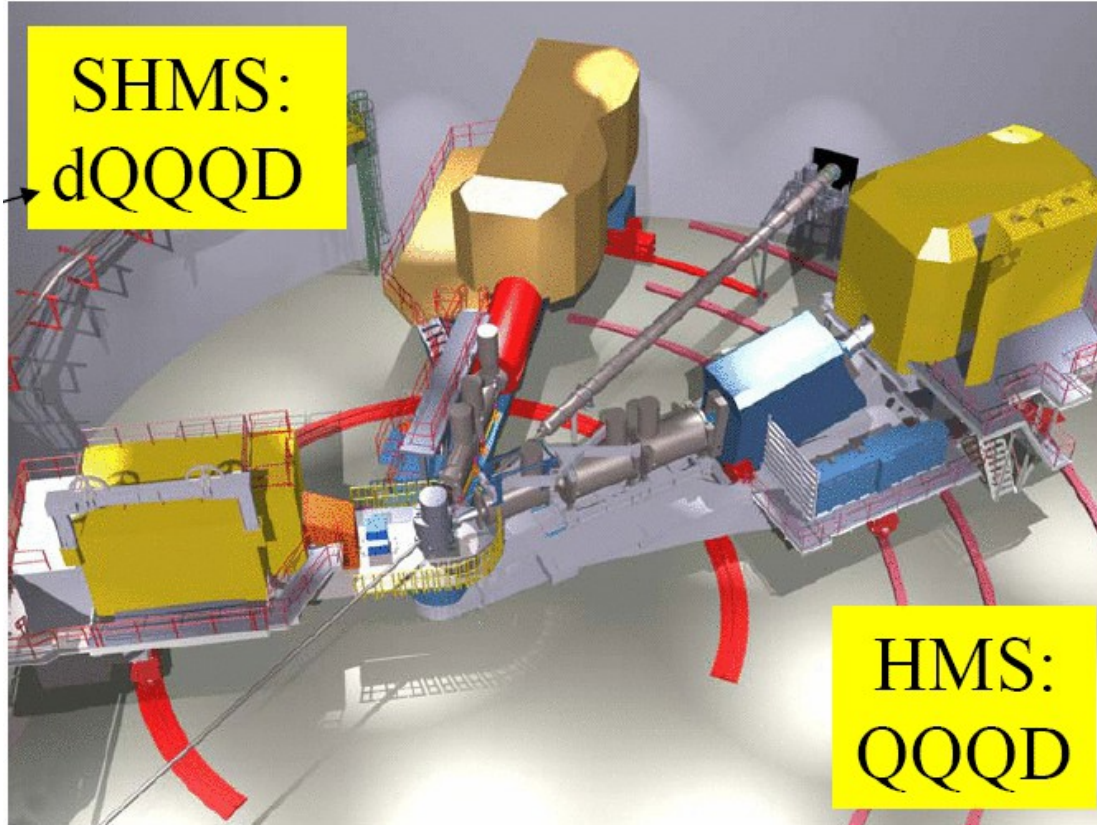
- $d_2$  is a clean probe of quark-gluon correlations / higher twist effects
- Connected to the color Lorentz force acting on the struck quark (Burkardt)
  - same underlying physics as in SIDIS  $k_\perp$  studies
- Investigate the present discrepancy between data and theories.
  - Theory calcs consistent but have wrong sign, wrong value.



- Spokespeople: T. Averett, W. Korsch, Z.E. Meziani, B. Sawatzky

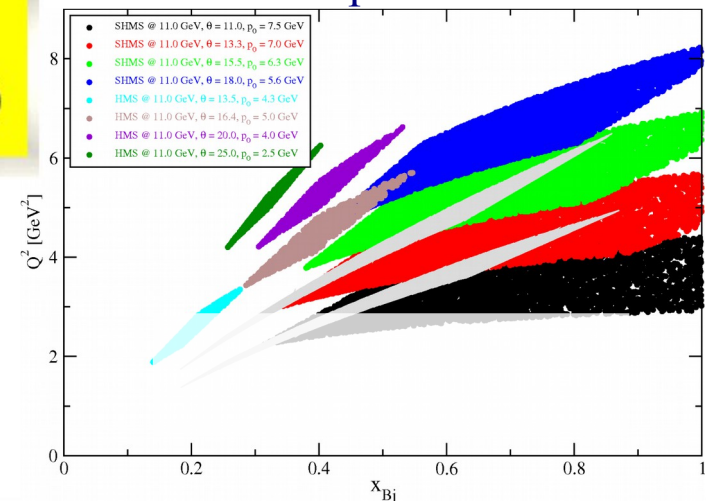


# E12-06-121: $d_2^n, g_2^n$

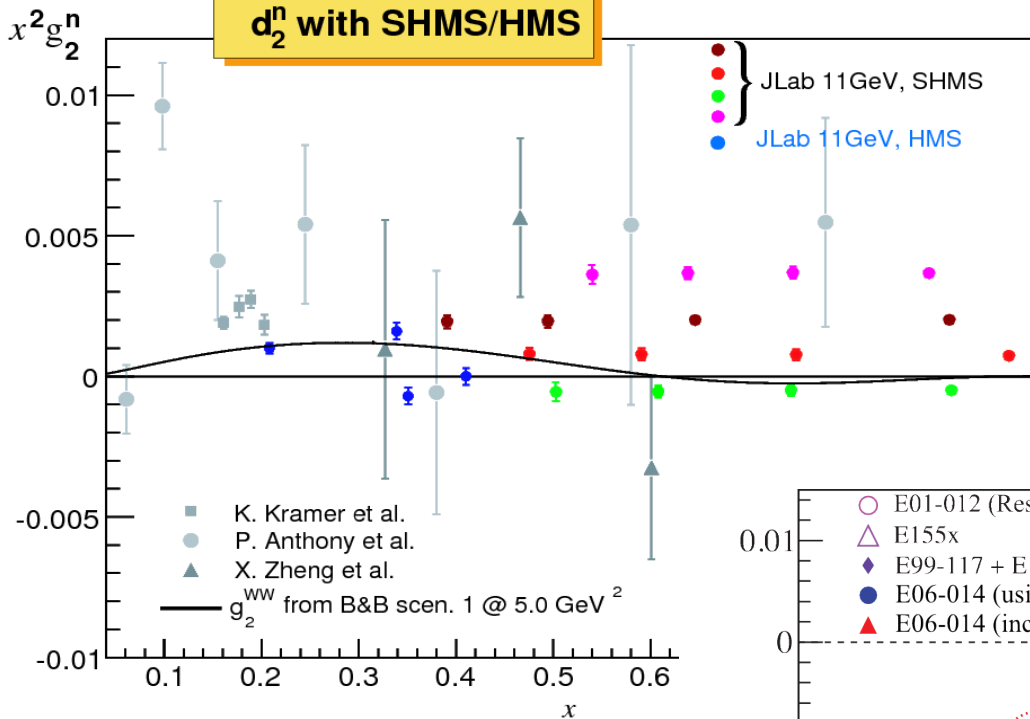


- SHMS collects data at  $\Theta = 11^\circ, 13.3^\circ, 15.5^\circ$  and  $18.0^\circ$  for 125 hrs each  
→ data from each setting divided into 4 bins
- HMS collects data at  $\Theta = 13.5^\circ, 16.4^\circ, 20.0^\circ$  and  $25.0^\circ$  for 125 hrs each

- Hall C: SHMS + HMS
- One beam energy  
→ 11 GeV
- Each arm measures a total cross section independent of the other arm.
- Experiment split into four pairs of 125 hour runs with spectrometer

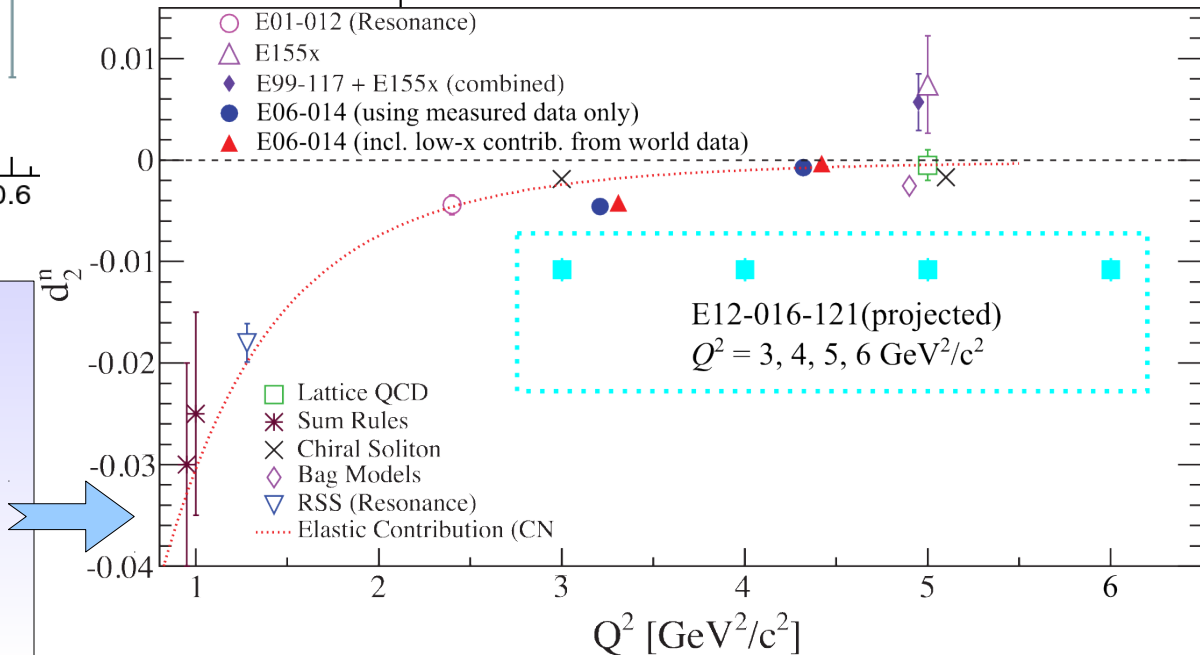


# Projected results for E12-06-121



Projected  $g_2^n$  points are vertically offset from zero along lines that reflect different (roughly) constant  $Q^2$  values from  $2.5-7 \text{ GeV}^2$ .

- $Q^2$  evolution of  $d_2^n$  in a region where models are thought to be accurate.
- Direct overlap with  $6 \text{ GeV}$  Hall A measurements.



# The E06-014 Collaboration (Hall A)

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A. Camsonne	R. Gilman	E. Long	R. D. Ransome	Y. Zhang
M. Canan	O. Glamazdin	A. Lukhanin	S. Riordan	Y.-W. Zhang
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C. Chen	J. Gomez	D. McNulty	<b>B. Sawatzky</b>	X. Zheng
J.-P. Chen	L. Guo	<b>Z.-E. Meziani</b>	M. H. Shabestari	
<b>S. Choi</b>	O. Hansen	R. Michaels	A. Shahinyan	
E. Chudakov	D. W. Higinbotham	M. Mihovilovič	S. Širca	
F. Cusanno	T. Holmstrom	B. Moffit	P. Solvignon	
M. M. Dalton	J. Huang	N. Muangma	R. Subedi	
W. Deconinck	C. Hyde	S. Nanda	V. Sulkosky	
C. W. de Jager	H. F. Ibrahim	A. Narayan	A. Tobias	
X. Deng	<b>X. Jiang</b>	V. Nelyubin	W. Troth	
A. Deur	G. Jin	B. Norum	D. Wang	
C. Dutta	J. Katich	Nuruzzaman	Y. Wang	
L. El Fassi	A. Kelleher	Y. Oh	B. Wojtsekhowski	
<b>D. Flay</b>	A. Kolarkar	<b>D. S. Parno</b>	X. Yan	

**Co-spokesperson**

**PhD (complete)**

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# BACKUP

# Other Interesting Observables

- Often experimentally simpler to measure Asymmetries
  - target polarized parallel to beam polarization
  - target polarized transverse to beam polarization
- Asymmetry is formed by measuring the difference in yields when you flip polarization of the beam
- Virtual photon asymmetries  $A_1$  and  $A_2$  can be expressed in terms of these experimental asymmetries
  - these are also connected to  $g_1$ ,  $g_2$ , of course

$$A_{\parallel} = \frac{N^{\uparrow\downarrow} - N^{\downarrow\downarrow}}{N^{\uparrow\downarrow} + N^{\downarrow\downarrow}}$$

$$A_{\perp} = \frac{N^{\uparrow\Rightarrow} - N^{\downarrow\Rightarrow}}{N^{\uparrow\Rightarrow} + N^{\downarrow\Rightarrow}}$$

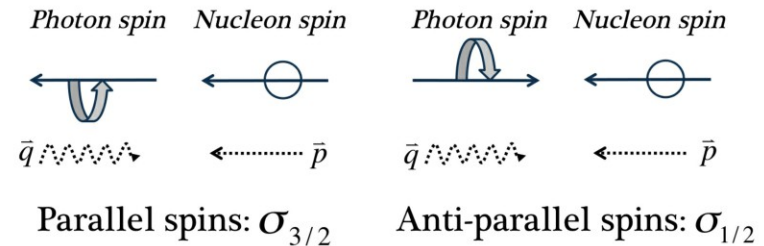
# Virtual Photon Asymmetries

$$A_1 = \frac{1}{(E + E')D'} \left( (E - E' \cos \theta) A_{\parallel} - \frac{E' \sin \theta}{\cos \phi} A_{\perp} \right)$$

$$A_2 = \frac{\sqrt{Q^2}}{2ED'} \left( A_{\parallel} + \frac{E - E' \cos \theta}{E' \sin \theta \cos \phi} A_{\perp} \right)$$

$$A_1 = \frac{1}{F_1} \left( g_1 - \frac{(2Mx)^2}{Q^2} g_2 \right)$$

$$= \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}}$$

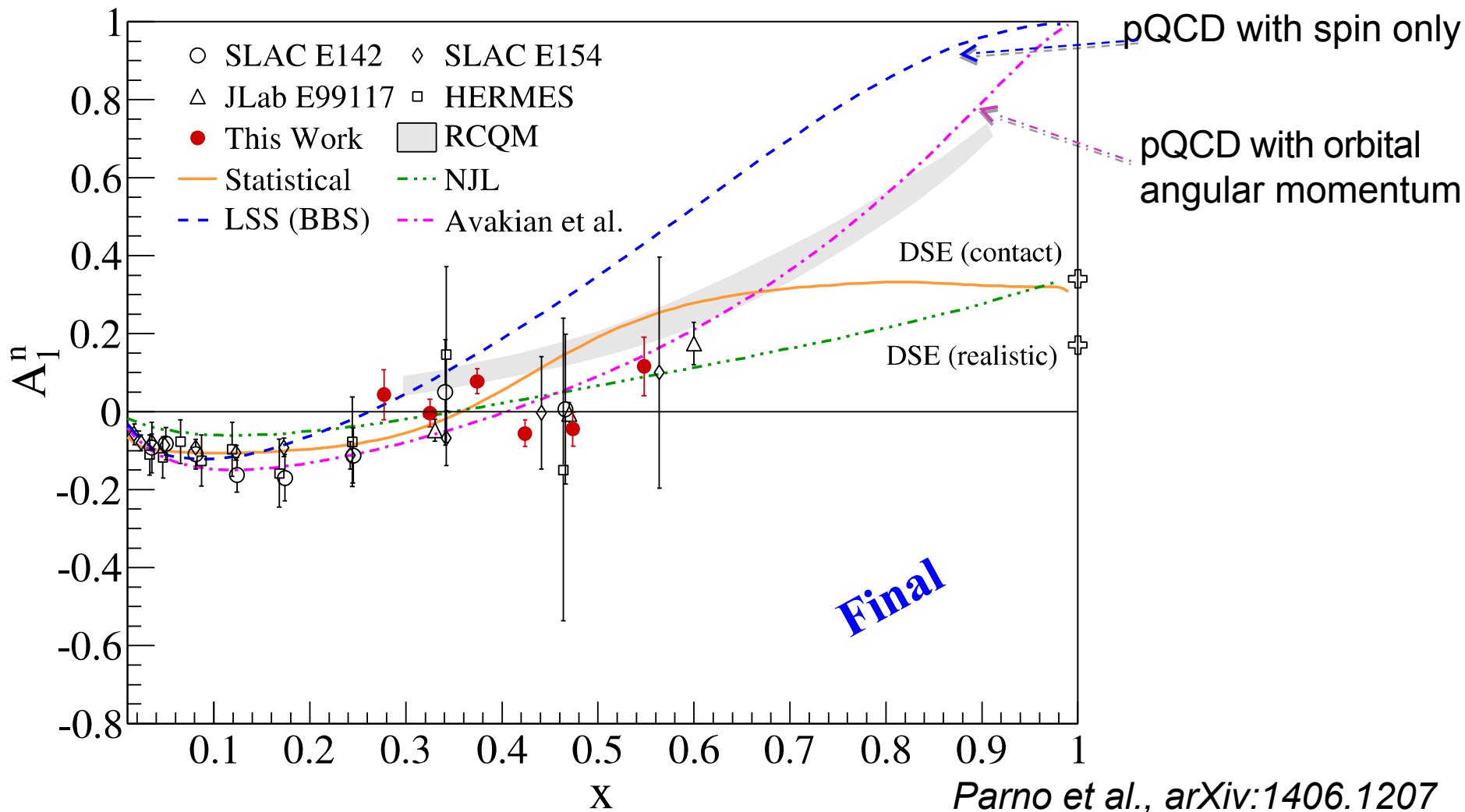


$$A_2 = \frac{\sqrt{Q^2}}{\nu} \frac{g_T}{F_1}$$

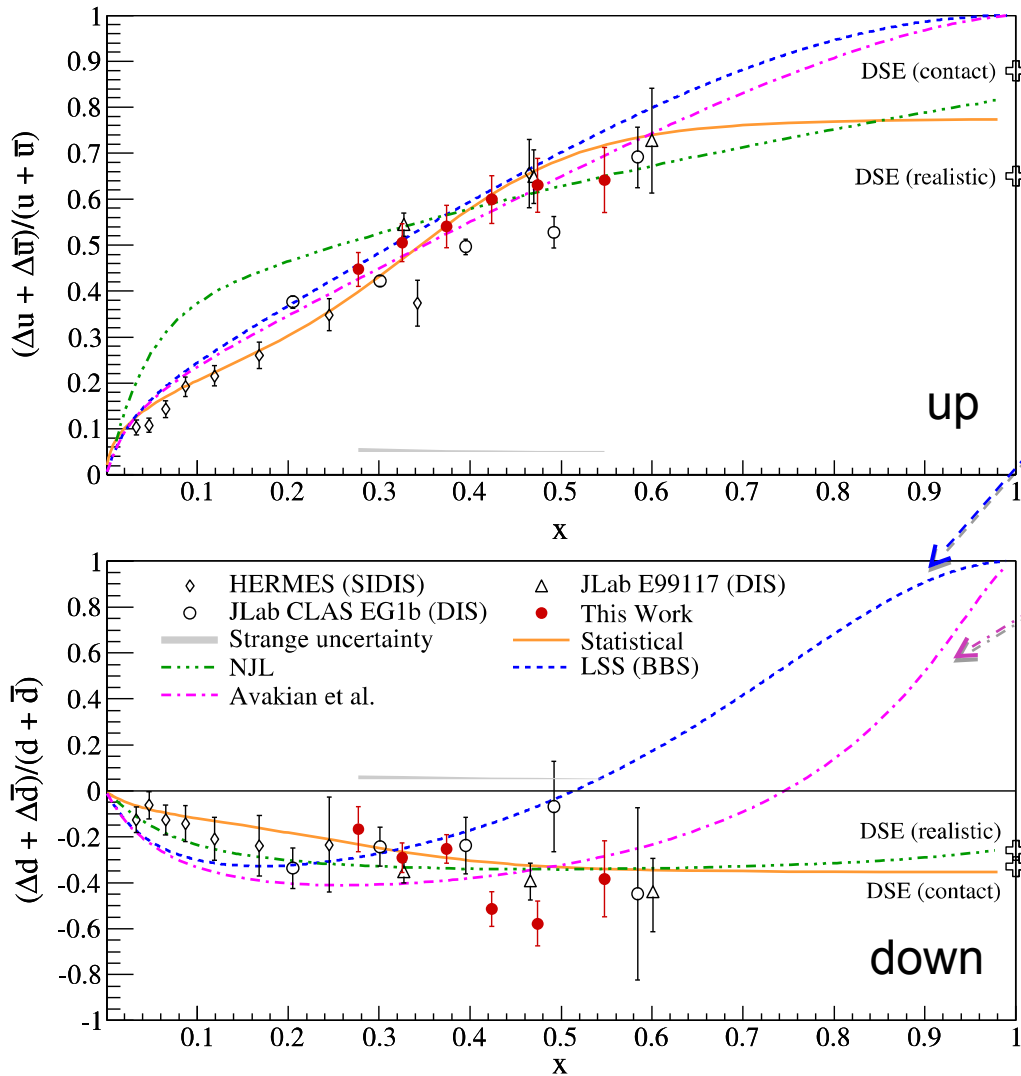
$$g_T = g_1 + g_2$$

- $g_T$  measures spin distribution normal to the virtual photon

# $A_1^n$ for Neutron



# $(\Delta u + \Delta \bar{u})/(u + \bar{u})$ and $(\Delta d + \Delta \bar{d})/(d + \bar{d})$



pQCD with spin only

pQCD with orbital angular momentum

- Spin-only pQCD is strongly disfavored
- Must go higher in  $x$  to distinguish between other models

Parno et al., arXiv:1406.1207



# $A_1^n$ vs. CQM Theory

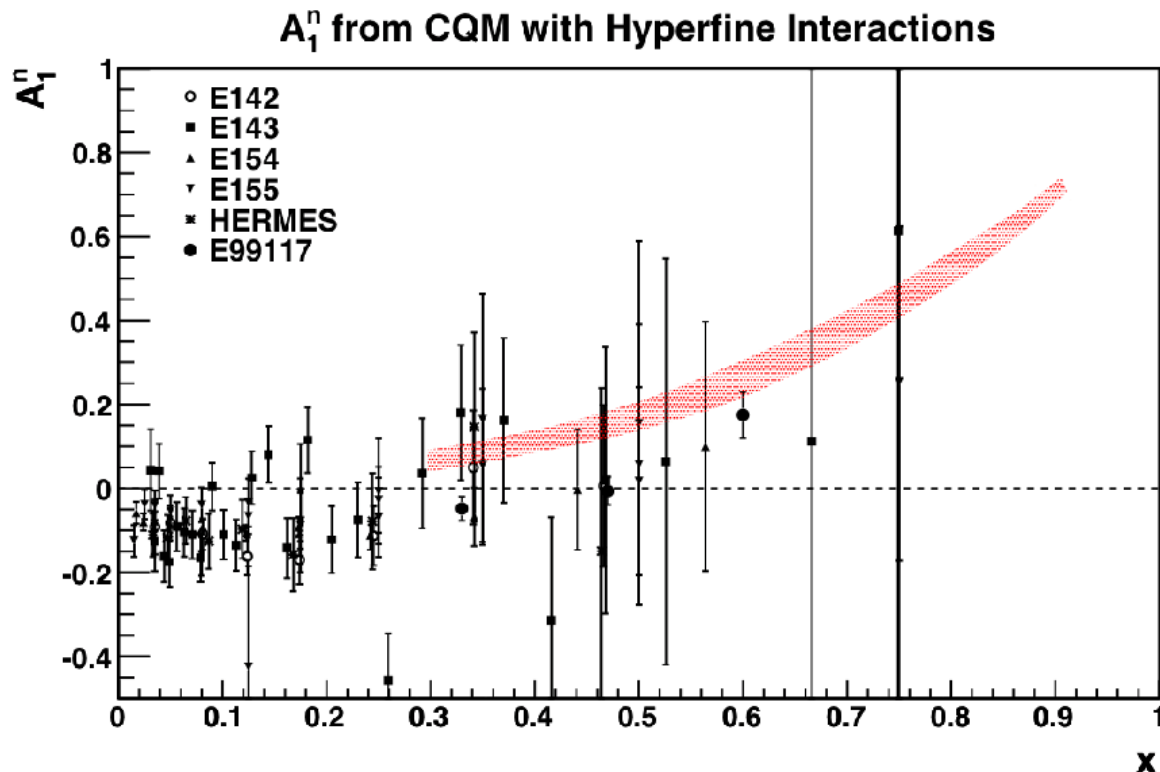


Figure 2.9: Predictions of  $A_1^n$  from constituent quark model, compared to world data. The dashed line on the x-axis marks  $A_1^n = 0$ , the prediction from unbroken SU(6) symmetry. The shaded red band shows the range of  $A_1^n$  values allowed in a model where SU(6) symmetry is broken by hyperfine interactions between quarks [70]. We have made use of parameterizations compiled by X. Zheng [71].

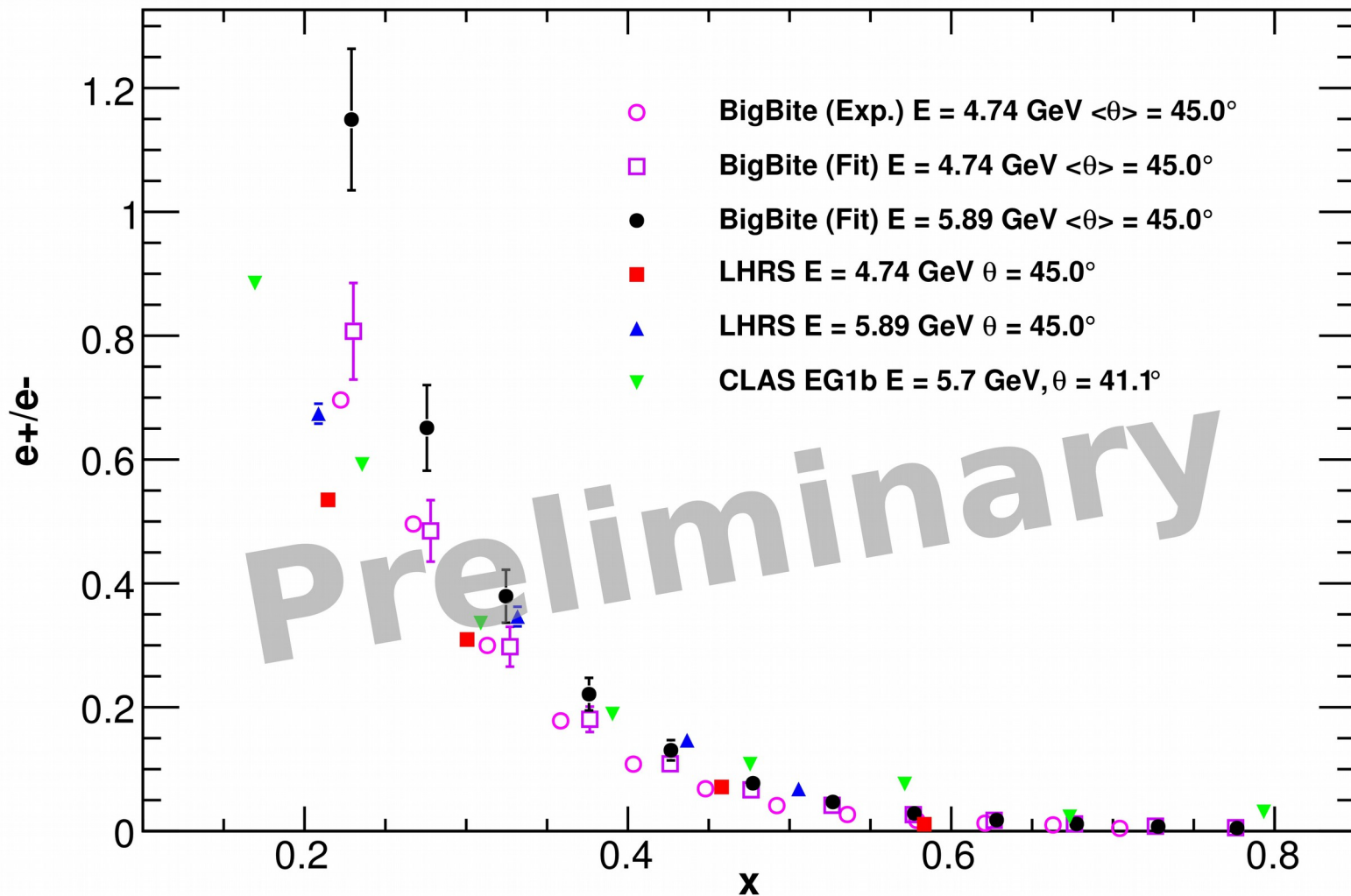
# Systematic Error Contributions to $g_2^n$ and $d_2^n$

Item description	Subitem description	Relative uncertainty
Target polarization		3 %
Beam polarization		1.5%
Asymmetry (raw)	<ul style="list-style-type: none"> <li>• Target spin direction (<math>0.1^\circ</math>)</li> <li>• Beam charge asymmetry</li> </ul>	$< 5 \times 10^{-4}$ $< 50$ ppm
Cross section (raw)	<ul style="list-style-type: none"> <li>• PID efficiency</li> <li>• Background Rejection efficiency</li> <li>• Beam charge</li> <li>• Beam position</li> <li>• Acceptance cut</li> <li>• Target density</li> <li>• Nitrogen dilution</li> <li>• Dead time</li> <li>• Finite Acceptance cut</li> </ul>	$< 1$ % $\approx 1$ % $< 1$ % $< 1$ % 2-3 % $< 2$ % $< 1$ % $< 1$ % $< 1$ %
Radiative corrections		$\leq 5$ %
From $^3\text{He}$ to Neutron correction		5 %
Total systematic uncertainty (for both $g_2^n(x, Q^2)$ and $d_2(Q^2)$ )		$\leq 10$ %
Estimate of contributions to $d_2$ from unmeasured region	$\int_{0.003}^{0.23} \tilde{d}_2^n dx$	$4.8 \times 10^{-4}$
Projected absolute statistical uncertainty on $d_2$		$\Delta d_2 \approx 5 \times 10^{-4}$
Projected absolute systematic uncertainty on $d_2$ (assuming $d_2 = 5 \times 10^{-3}$ )		$\Delta d_2 \approx 5 \times 10^{-4}$

- Radiative correction uncertainty cross-checked with E01-012 (Spin Duality) experiment  
 ☒ worst case: 4.4%
- Pion rejection ratio of  $\sim 10^4:1$  should be achievable with standard SHMS/HMS detectors ( $\sim 10^3:1$  would be adequate)

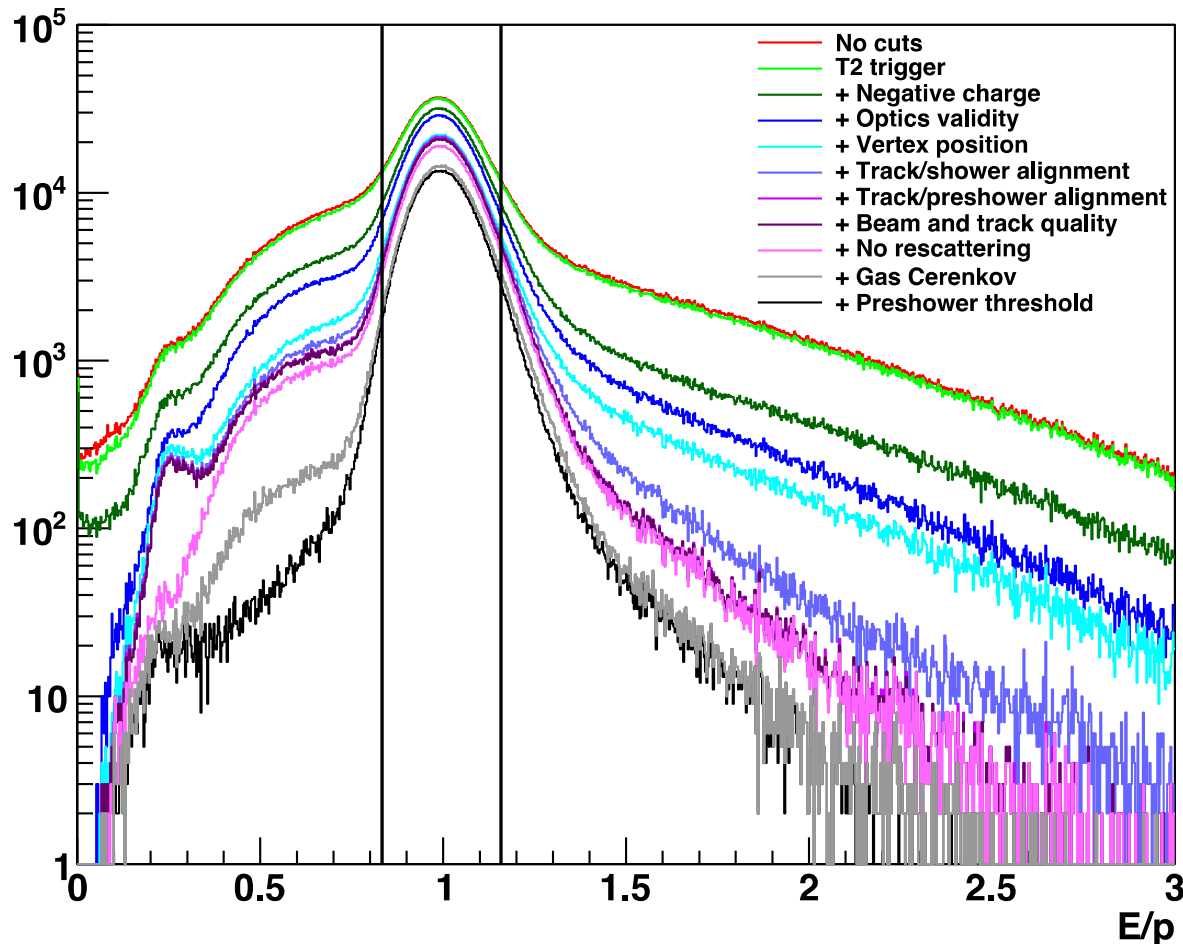
# e<sup>+</sup>/e<sup>-</sup> Ratios extracted from data (E06-014)

## e<sup>+</sup>/e<sup>-</sup> Comparison



Courtesy of D.Fløy, M. Posik

# Data reduction (E06-014)



- BigBite data shown
- Negligible pion contamination
- Errors associated with momentum reconstruction (tracking) form highest backgrounds

Courtesy of D. Parno

# E12-06-121 Updated Kinematics

Table 3: Kinematic bins and expected rates for the SHMS. The uncertainties for  $A_{\parallel}$  and  $A_{\perp}$  are *statistical* only.

SHMS Setting	$E'_{bin}$ [GeV]	$Q^2$ [GeV <sup>2</sup> ]	x	W [GeV]	$e^-$ rate [Hz]	$\pi^-$ rate [Hz]	$t_{\parallel}$ [hrs]	$t_{\perp}$ [hrs]	$\Delta A_{\parallel}$ [ $\cdot 10^{-4}$ ]	$\Delta A_{\perp}$ [ $\cdot 10^{-4}$ ]
$\theta_0 = 11^\circ$  $E'_{cent} = 7.5$ GeV	7.112	2.875	0.394	2.305	1058	11	12	113	2.0	0.5
	7.709	3.116	0.504	1.988	708	3.1	12	113	2.3	0.7
	8.304	3.357	0.663	1.610	259	0.83	12	113	3.7	0.1
	8.900	3.597	0.912	1.109	2.7	0.21	12	113	36	10
$\theta_0 = 13.3^\circ$  $E'_{cent} = 7.0$ GeV	6.647	3.922	0.480	2.267	268	3.1	12	113	3.5	1.0
	7.203	4.250	0.596	1.941	139	0.8	12	113	4.8	1.5
	7.758	4.578	0.752	1.548	31.6	0.16	12	113	10	3.1
	8.314	4.906	0.972	1.012	0.10	0.033	12	113	173	55
$\theta_0 = 15.5^\circ$  $E'_{cent} = 6.3$ GeV	5.997	4.798	0.511	2.342	96	1.9	12	113	5.7	1.8
	6.496	5.197	0.614	2.037	49	0.47	12	113	7.8	2.5
	6.995	5.597	0.744	1.677	13.5	0.11	12	113	15	4.7
	7.494	5.996	0.911	1.215	0.29	0.025	12	113	98	33
$\theta_0 = 18.0^\circ$  $E'_{cent} = 5.6$ GeV	5.348	5.756	0.542	2.397	35	1.1	12	113	9.5	3.1
	5.790	6.235	0.637	2.106	17	0.25	12	113	13	4.4
	6.233	6.711	0.749	1.769	5.1	0.05	12	113	24	8.1
	6.675	7.187	0.885	1.350	0.38	0.01	12	113	87	30

Table 4: Expected rates for the three HMS settings. The uncertainties for  $A_{\parallel}$  and  $A_{\perp}$  are *statistical* only.

$\theta_0$ [°]	$E'_{cent}$ [GeV]	$Q^2$ [GeV <sup>2</sup> ]	x	W [GeV]	$e^-$ rate [Hz]	$\pi^-$ rate [Hz]	$t_{\parallel}$ [hrs]	$t_{\perp}$ [hrs]	$\Delta A_{\parallel}$ [ $\cdot 10^{-4}$ ]	$\Delta A_{\perp}$ [ $\cdot 10^{-4}$ ]
13.5	4.305	2.617	0.208	3.293	954	765	8	117	2.0	0.6
16.4	5.088	4.555	0.410	2.727	218	15	12	113	3.9	1.2
20.0	4.000	5.31	0.404	2.951	76	66	10	115	6.0	1.8
25.0	2.500	5.15	0.323	3.417	20	84	13	112	10.7	3.1

# Updated Kinematics

