ATLAS results on diffraction and forward physics







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4 analyses: 1)	Vetoing one side of detector
2)	Soft rapidity gaps
3)	Jet veto
4)	$\sigma_{tot}, \sigma_{el}, \sigma_{inel}$ with ALFA

The ATLAS Detector



ATLAS Forward detectors



Gaps in the detector

Typical Single-Diffraction event: activity on one side (proton remnants + possibly hard scatter), the other part of detector empty.



Diffraction enhanced MinBias events

"Measurement of the inelastic pp cross section at \sqrt{s} = 7 TeV with the ATLAS detector"

Nature Comm 2 (2011) 1926

L= 20 μb^{-1} , single fill in March 2010, <µ> ~ 0.01

Event selection:

1) Veto activity in MBTS on one side of IP 2) $N_{trk} \ge 1$ ($p_T > 0.5$ GeV, $|\eta| < 2.5$)

Study:

Ratio of diffraction enhanced to inclusive events
 Inelastic cross section

Diffraction enhanced MinBias events





- syst.unc. from background, MBTS response, material

 $f_D = 26.9^{+2.5}_{-1.0}$ - syst.unc. from varying DD/SD event ratio between 0.0 and 1.0 using default DL model and still consistent with data $R_{SS} \pm 1\sigma$

MC generators based on triple Pomeron exchange:



Pythia 8 +DL with IP intercept $\alpha(t) = \alpha(0) + \alpha' t$ $\frac{d\sigma_{SD}}{d\xi} \sim \frac{(1+\xi)}{\xi^{1+\varepsilon}}, \ \varepsilon = \alpha(0) - 1$

> Measured $\sigma_{inel} = 60.33 \pm 2.10$ mb for $\xi > 5 \times 10^{-6}$. Extrapolating to full ξ -range using DL MC gives $\sigma_{inel}(\xi > m_p^2/s) = 69.4 \pm 2.4 \pm 6.9$ (extr.) mb

"Rapidity gap cross section in pp interactions at \sqrt{s} = 7 TeV"

Eur. Phys. J. C72 (2012) 1926

L= 7.1 μb^{-1} , taken in March 2010, μ < 0.005

Event selection:

- 1) Only good tracks with $p_T > 200 \text{ MeV}$
- 2) Only calorimeter cells with E > noise ($\sim 5\sigma_{noise}$)
- 3) Look for empty η -rings starting from the edge of calorimeter ($|\eta| < 4.9$)

Study:

1) Gap spectrum

- 2) Predictions of various MC models
- 3) Diffraction dynamics

Definitions of gaps

Detector level:

Detector divided into η -rings of size 0.1

in the calorimeter ($|\eta| < 4.9$)

Ring is considered empty if

- No track with p_T > 200 MeV ($|\eta|$ <2.5)
- No calo cell with E > 4.5-5.5 $\sigma_{noise}(\eta)$ ($|\eta|$ <4.9)



Phase space divided into η -rings of size 0.1 Ring is considered empty if

- No stable particle with p_T > 200 MeV

Observable: Forward gap $\Delta \eta_F$

- The largest consecutive set of empty rings starting from the edge of the calorimeter





General properties of the gap spectrum



The effect of $p_{T,min}$ of clusters and tracks



- □ Too large a $p_{T,min}$ produces artificial gaps
- **\square** $p_{T,min}$ = 200 MeV was found to be optimum:
- Tracks are within the acceptance of inner tracker
- Selection efficiency of calorimeter clusters > 50%
- Above noise threshold in any part of calorimeter

□ Data show an exponential decrease at low $\Delta \eta_F$ (dominance of ND) and plateau at large $\Delta \eta_F$ (dominance of diffr. processes) – as expected from previous HERA and Tevatron results □ This general feature is followed by

all MC models.

- Details are not described by any of the models.
- Data lie in between the DL slope and slope of the other models





Details about composition and components in PYTHIA 6, PYTHIA 8 and PHOJET

- □ All models overshoot the data (i.e. overshoot σ_{inel}) □ PYTHIA 6, 8: unexpectedly large DD
- □ PHOJET: much smaller DD but contribution of CD

Jet veto

1) "Measurement of dijet production with a veto on additional central jet activity in pp collisions at $\sqrt{s} = 7$ TeV using the ATLAS detector" JHEP 1109 (2011) 053 L= 36 pb^{-1} , all 2010 data

A new study: 2) "Measurement of jet vetoes and azimuthal decorrelations in dijet events produced in pp collisions at $\sqrt{s} = 7$ TeV using the ATLAS detector" arXiv:1407.5756 [hep-ex] L= 38 pb^{-1} (all 2010 data) + L= 4.5 fb^{-1} (all 2011 data)

i) optimised 2-jet trigger procedure to reach Δy of up to 8 on 2010 data (low pile-up) ii) add 2011 data to extend the high-pt region up to 1.5 TeV

Event selection:

Jet anti-kt: R=0.6, $|\eta_{jet}| < 4.4$, $p_{T,jet1} > 60 \text{ GeV}$, $p_{T,jet2} > 50 \text{ GeV}$ Veto scale $Q_0 = p_T > 20$ (30) GeV for 2010 (2011) data

Define Gap fraction = $\sigma_{jj}(Q_0)/\sigma_{jj}$, σ_{jj} = inclusive dijet x-section,

 $\sigma_{ij}(Q_0)$ = dijet cross section with no jets of $p_T > Q_0$ between η_{jet1} and η_{jet2}

Results:

- 1) Gap fraction vs. Δy for $Q_0 = 20$ GeV
- 2) Gap fraction vs. $\overline{p_T}$ for $Q_0 = 30$ GeV
- 3) Gap fraction vs. Q_0 for slices of y
- 4) < N_{jets} > vs. Δy and $\overline{p_T}$
- 5) Cross sections vs. Δy and $\Delta \phi$
- 6) $<\cos\Delta\phi>$ and $<\cos2\Delta\phi>$ vs. Δ y and $\overline{p_T}$

Jet veto

Goal = study higher-order QCD effects in the dijet production

These are most visible when in the dijet system (2 highest Et jets):

- 1) jet separation (Δy) is large
- 2) jet veto between jets (with scale Q_0) is applied
- 3) azimuthal decorrelation between jets $(\pi \Delta \varphi)$ is large

Resummation of HO QCD effects: - in $ln(Q^2)$: DGLAP approach - in ln(1/x) : BFKL approach

DGLAP effects if p_T/Q_0 large

BFKL effects if ∆y large

Gap fraction sensitive to

Color-Singlet-Exchange if $p_T/Q_0 \& \Delta y$ large

Azim. decorrelations sensitive to both DGLAP and BFKL effects



Theoretical predictions:

Powheg: NLO dijet (DGLAP form.) interfaced with Pythia8 or Herwig to resum soft and collinear emissions using parton shower approximation

HEJ: leading-log calc. of multi-jet (at least 2 hard jets) production in Mueller-Navelet limit – interfaced to ARIADNE (color-dipole cascades producing soft and collinear radiation in the gap between jets → contr. from low-x BFKL-like log. terms)



Jet veto



Exponential decrease:
exchange of color in the t-channel
Plateau:

a) steeply falling PDFs at large x
disabling additional parton radiation
[at extreme Δy or p_T gap fr. must reach unity]
b) Contribution of color-singlet
exchange (CSE) processes
[cross section for CSE rises with Δy]
> All models fall at high Δy

➢ Powheg underestimates, HEJ overestimates azim. decorrelation
 ➢ With increasing ∆y and p_T, the jet veto makes the dijet system more back-to-back

• Best agreement with Powheg+Py8 except for low- p_T , high Δy region • Interfacing ARIADNE to HEJ improves significantly description of dates

Total and elastic pp cross section

"Measurement of the total cross section from elastic scattering in pp collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector"

submitted to Nucl. Phys. B, arXiv: 1408.5778 [hep-ex]

L= 78.7 μb^{-1} , special run in 2011, $<\mu_{elastic}> \sim 0.01$

Event selection:

- 1) Dedicated run in 2011 (together with Totem), $\beta^* = 90$ m
- 2) ATLAS subdetector: Forward Proton detector ALFA
- 3) Geometrical, fiducial and back-to-back signatures to select elastics & reduce background

Results:

- 1) Differential elastic cross section as a function of t
- 2) Total pp cross section via optical theorem: $\sigma_{tot} \propto 4\pi \cdot \text{Im}(f_{el})_{t \to 0}$

4) Slope of the elastic cross section at small t

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ALFA detector + Measurement principle



Vertical Roman Pots :

- Thin windows at bottom and front: 0.2mm and 0.5mm
- Main detectors (MD), Overlap detectors (for alignment), Trigger detectors
- MD = 10 staggered layers, 1 layer=64 single-cladded square (0.5mm) scintil. fibres in u- and v-directions
- Spatial resolution ~ 30 μm
- Reconstruction efficiency ~ 90%
- Distance of the closest approach in the run: $6.5\sigma_y \sim 5.8$ mm

To measure $d\sigma/dt$:

- 1) Measure track position y
- 2) Calculate scattering angle θ_y^* using the beam optics, i.e. using transport matrix elements:

$$\begin{pmatrix} y \\ \theta_y \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} y^* \\ \theta_y^* \end{pmatrix}$$

In the simple case (high beta*, phase advance 90 deg., parallel-to-point focusing):



Experimental procedure (1)

Signal selection & Background rejection



Event selection:

- L1 elastic trigger
- Data quality & track quality
- Geometrical & fiducial cuts:
- acceptance & signal selection

& background rejection

Final sample: ~ 805k selected events

(~ 1k Pile-up events)

Data-driven Background estimates:



Experimental procedure (2)





Independently measured Luminosity (3 methods: BCM, LUCID and Inner Det.): $L = 78.7 \pm 1.9 \ \mu b^{-1}$ (2.3% uncertainty) **Event reconstruction efficiency** (from data, separately for each arm):

 $\epsilon_{reco,1}$ = 89.74 ± 0.04(stat) ± 0.61(syst) % $\epsilon_{reco,2}$ = 88.00 ± 0.05(stat) ± 0.92(syst) %

[Trigger plates placed behind MD in arm1, but in front of MD in arm2]



Correction procedure (separately for each arm):

$$\left(\frac{d\sigma}{dt}\right)_{i} = \frac{1}{t_{i}} \cdot \frac{\mathbf{M}^{-1}[N_{i} - B_{i}]}{A_{i} \cdot \varepsilon^{reco} \cdot \varepsilon^{trig} \cdot \varepsilon^{DAQ} \cdot \mathbf{L}_{int}}$$

- A = Acceptance (t)
- M = Unfolding procedure (symbolic)
- N = Number of selected events
- B = Number of estimated background events
- e^{reco} = reconstruction efficiency
- ϵ^{trig} = trigger efficiency
- ϵ^{DAQ} = dead-time correction
- L_{int} = integrated luminosity

Main system. uncertainties:

- Luminosity
- Nominal beam energy

Fitting elastic cr.section $\rightarrow \sigma_{tot}, \sigma_{el}, \sigma_{inel}, B$ -slope



- > Extraction of σ_{el} : assume Nuclear term only and B(t)=const: Integrating over full t-range $\sigma_{el} = 24.00 \pm 0.19$ (stat.) ± 0.57 (syst.) mb [Observed in -t $\epsilon < 0.0025$, 0.38> GeV²: $\sigma_{el} = 21.66 \pm 0.02$ (stat.) ± 0.58 (syst.) mb (90% of the total σ_{el})]
- $> \sigma_{inel} = \sigma_{tot} \sigma_{el} \longrightarrow \sigma_{inel} = 71.34 \pm 0.36 \text{ (stat.)} \pm 0.83 \text{ (syst.) mb}$

Comparison with previous measurements





The same run in 2011, Lumi-dependent method:

ATLAS: $\sigma_{tot} = 95.4 \pm 1.4 \text{ mb}$ (Lumi unc=2.3%) TOTEM: $\sigma_{tot} = 98.6 \pm 2.2 \text{ mb}$ (Lumi unc=4%) \rightarrow Difference = 1.3 σ

ATLAS value ~2σ below COMPETE fit, but closer to predictions by Block & Halzen, KMR, Soffer.

ATLAS: $\sigma_{el} = 24.0 \pm 0.6$ mb (Lumi unc=2.3%) Totem: $\sigma_{el} = 25.4 \pm 1.1$ mb (Lumi unc=4%) \rightarrow Difference = 1.1 σ

Summary

Four ATLAS analyses presented (two of them very fresh):

□ 10%-fraction of diffraction-enhanced to inclusive events measured in minimum bias events with veto on one side of the detector.

 \Box $\sigma_{inel} = 69.4 \pm 2.4 \pm 6.9$ (extr.) measured by extrapolating to $\xi < 5 \times 10^{-6}$ using DL model

- ♦ Forward gap distribution measured for different $p_{T,min}$ cuts (200, 400, 600, 800 MeV)
- Plateau observed for large gaps suggests presence of diffractive processes
- None of the models able to describe the details of the gap spectrum
- * σ_{inel} measured for various ξ values. Models not able to describe the ξ-evolution
- ✤ Pomeron intercept in DL model measured from fit to spectrum at large gaps. It is compatible with the default DL value ($\alpha_{IP}(0) = 1.085$).

> Jet Gap fraction and azimuthal decorrelation measurements combined into a single analysis that extends the kinematic reach and tests models based on either DGLAP or BFKL dynamics.

- First ATLAS measurement of σ_{tot} using fit to t-dependence of elastic cross section and independently measured luminosity. The value 95.4 ± 1.4 mb in agreement with Totem result.
- Measured total σ_{el} and σ_{inel} , and B-slope also agree with Totem.
- More data from 2012 data at $\sqrt{s}=8$ TeV ($\beta^* = 90$ m and 1km) in the pipeline.

BACKUP SLIDES

Diffraction needs very low Pile-up



Pile-up = soft particles sitting on top of the hard-scale event, influencing efficiencies of various finding algorithms (Primary vertex, triggers, jets and other usual objects).

Diffractive signature: in the absence of forward proton detectors we have to rely on gaps. Gaps are easily seen only when Pile-up is low \rightarrow concentrate on data 2010.

Special behaviour of HERWIG++ 2.5.1 UE7-2 which describes ATLAS MinBias data



□ HERWIG++ 2.5.1 MinBias UE7-2 does not contain diffractive processes, yet it produces very large gaps

- Different models of hadronization:
- No Empty Evts = excludes soft events
- No CR = excludes Color Reconnection
- □ HERWIG++ 2.5.1 MinBias UE7-2 fails to describe the shape for any of the tunings

Estimate of Pomeron intercept $\alpha_{IP}(0)$ in Donnachie-Landshoff model



 $(\alpha_{IP}(0) \text{ governs the slope at large } \Delta \eta_F$: Flatness $\rightarrow \alpha_{IP}(0) = 1$, Increase $\rightarrow \alpha_{IP}(0) > 1$)

 $\alpha_{I\!P}(t=0) = 1.058 \pm 0.003(\text{stat})^{+0.034}_{-0.039}(\text{syst}) \leftarrow \text{due to large model uncert. still}$ compatible with default DL $\alpha_{IP}(0)$

Cross section of ~ 1mb per unit of rapidity (see KMR, arXiv:1102.2844)

Best fit:

Estimate of the inelastic cross section



 \Box Inelastic cross section obtained for $\xi > \xi_{cut}$ from MC thanks to strong correlation $\Delta \eta_F \sim -\ln \xi$

- □ Small corrections (~2%) applied to include particles with p_T < 200 MeV and to account for hadronization fluctuations
- Total uncertainty dominated by uncertainty on luminosity
- **Comparison with Totem result and with previous ATLAS analysis ("Measurement of \sigma_{inel} (7 TeV)**
- Models fail to describe the evolution from low ξ (Totem) to large ξ (ATLAS) but RMK better in describing the slope (thanks to adding IP IP IR terms to IP IP IP terms)