

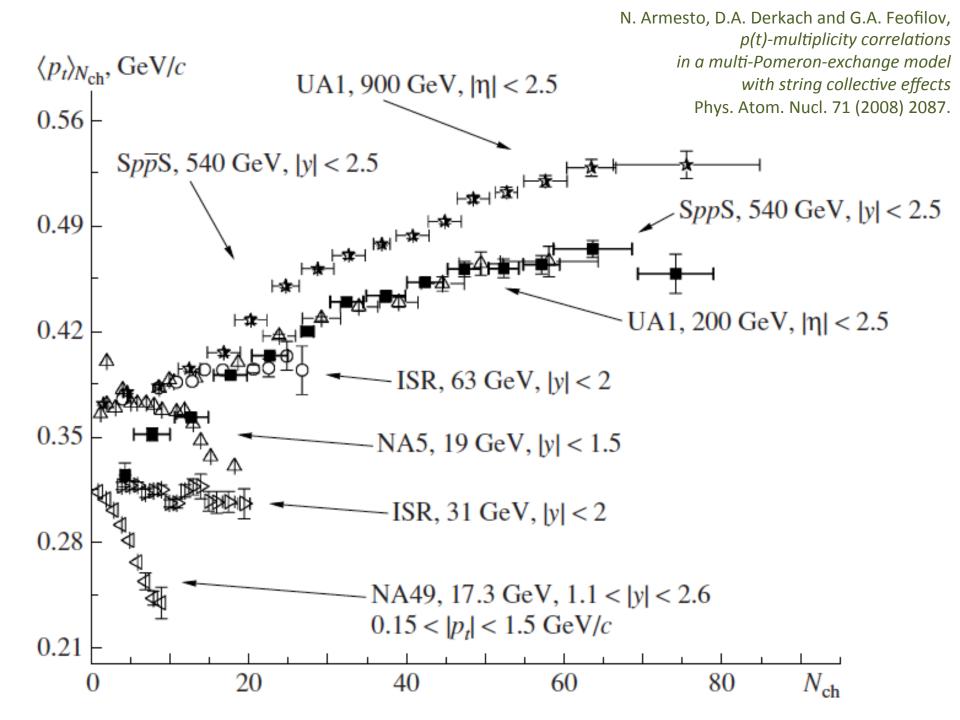
### Why mean p<sub>T</sub> is interesting

Michal Praszalowicz

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•  $\langle p_T \rangle (N_{ch})$  – correlations are sensitive to the fine details of dynamics





- $\langle p_T \rangle (N_{ch})$  correlations are sensitive to the fine details of dynamics
- complete change of behavior from small to large energies

#### MULTIPLICITY DEPENDENCE OF $p_t$ SPECTRUM AS A POSSIBLE SIGNAL FOR A PHASE TRANSITION IN HADRONIC COLLISIONS

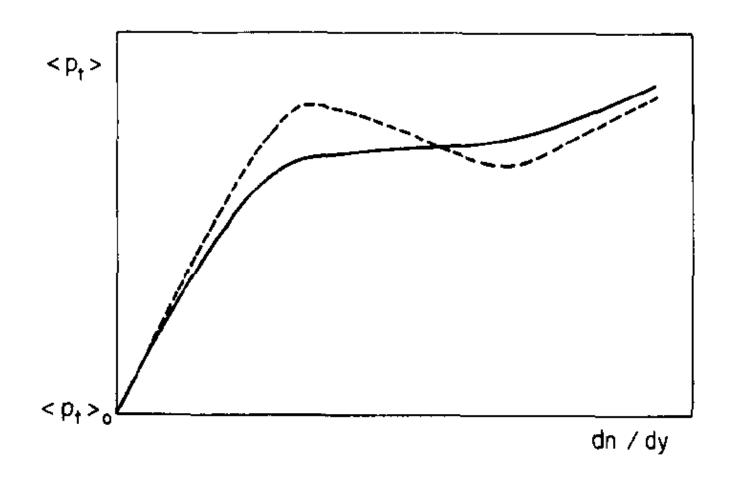


L. VAN HOVE

CERN, Geneva, Switzerland

Received 25 August 1982

Phys. Lett. B 118 (1982) 139





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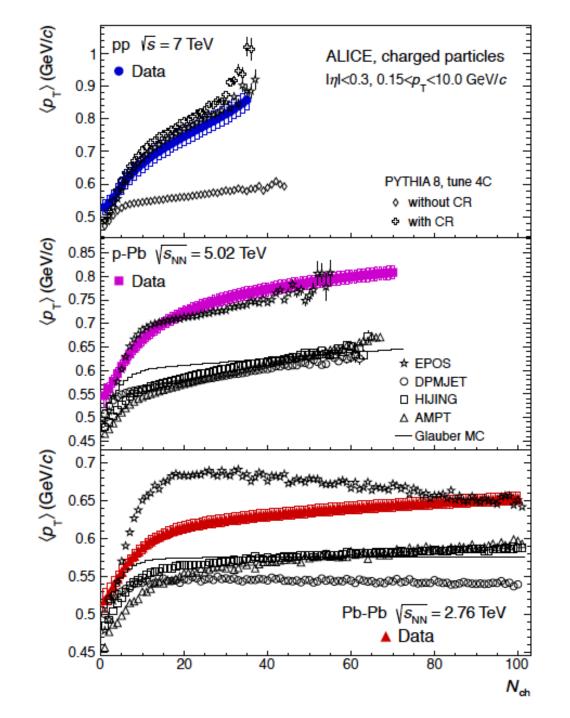




CERN-PH-EP-2013-111 July 2, 2013

## Multiplicity dependence of the average transverse momentum

The ALICE Collaboration





- $\langle p_T \rangle (N_{ch})$  correlations are sensitive to the fine details of dynamics
- complete change of behavior from small to large energies
- possible sign of phase transition
- difficult to describe by untuned MonteCarlos
- saturation included in EPOS does a good job



$$\frac{dN}{dyd^2p_{\mathrm{T}}} = S_{\perp}\mathcal{F}(\tau) \quad \tau = \frac{p_{\mathrm{T}}^2}{Q_s^2(x)} \qquad Q_s(x) = Q_0 \left(\frac{x_0}{x}\right)^{\lambda/2}$$



$$\frac{dN}{dyd^2p_{\mathrm{T}}} = S_{\perp}\mathcal{F}(\tau) \quad \tau = \frac{p_{\mathrm{T}}^2}{Q_s^2(x)} \qquad dp_{\mathrm{T}}^2 = \frac{2}{2+\lambda}\bar{Q}_s^2(W)\,\tau^{-\lambda/(2+\lambda)}d\tau$$

$$\bar{Q}_s(W) = Q_0\left(\frac{W}{Q_0}\right)^{\lambda/(2+\lambda)}$$



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$$rac{dN}{dy} = S_{\perp} \int \mathcal{F}( au) d^2p_{\mathrm{T}} = S_{\perp} \bar{Q}_s^2(W) \int \mathcal{F}( au) \dots d au = rac{1}{\kappa} S_{\perp} \bar{Q}_s^2(W)$$



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$$\frac{dN}{dy} = \frac{1}{\kappa} S_{\perp} \bar{Q}_s^2(W) \ \rightarrow \ \bar{Q}_s^2(W) = \frac{\kappa}{S_{\perp}} \frac{dN}{dy}$$



for  $y \sim 0$  (central rapidity) *i.e.* for  $x_1 \sim x_2 = x$  and for symmetric systems

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saturation scale = gluon density per transverse area Michal Praszalowicz



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parton – hadron duality: power-like growth of particle multiplicity

saturation scale = gluon density

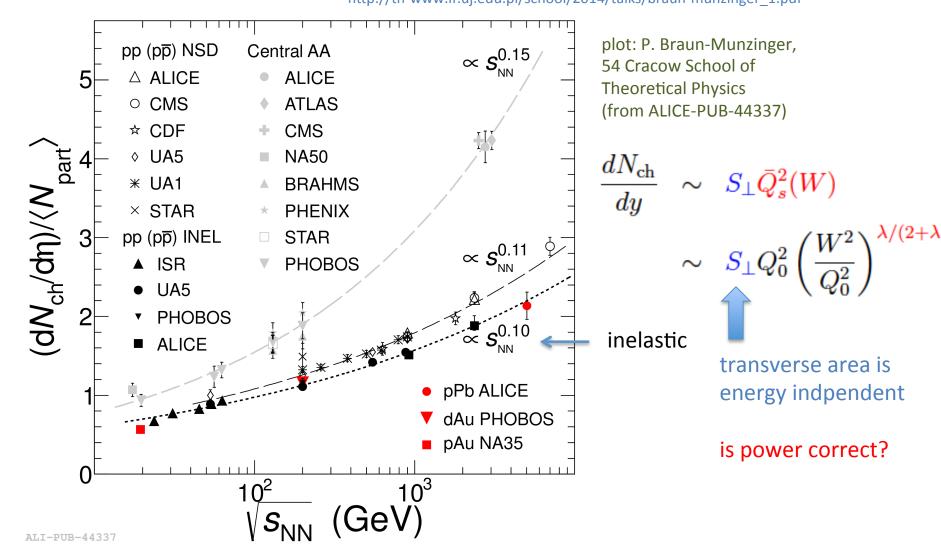
per transverse area

Michal Praszalowicz

## The state of the s

#### Power-like growth of multiplicity

http://th-www.if.uj.edu.pl/school/2014/talks/braun-munzinger\_1.pdf

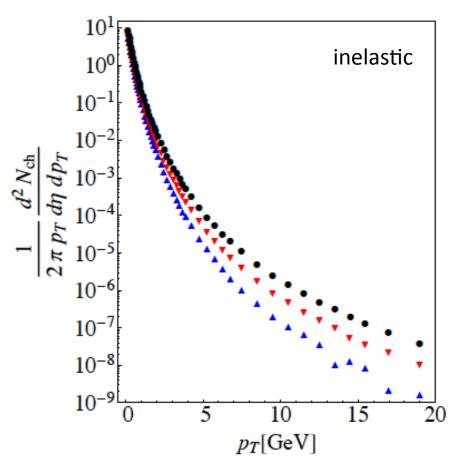


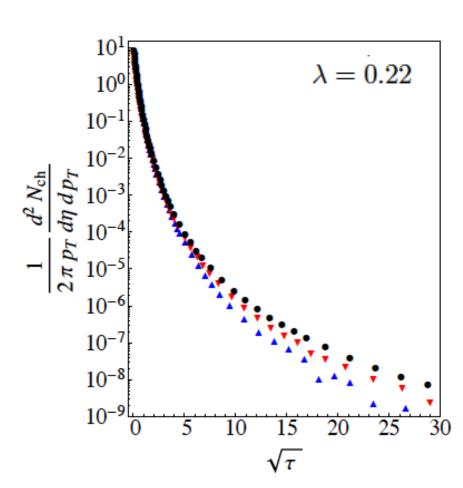


#### Determination of lambda

$$\frac{dN_{\rm ch}}{dyd^2p_{\rm T}} = S_{\perp} \mathcal{F}(\tau) \qquad \tau = \frac{p_{\rm T}^2}{Q_{\rm sat}^2(p_{\rm T}/\sqrt{s})} = \frac{p_{\rm T}^2}{1\,{\rm GeV}^2} \left(\frac{p_{\rm T}}{\sqrt{s}\times 10^{-3}}\right)^{\lambda}$$

ALICE 1307.1093 [nucl-ex], Eur.Phys.J C73 (2013) 2662



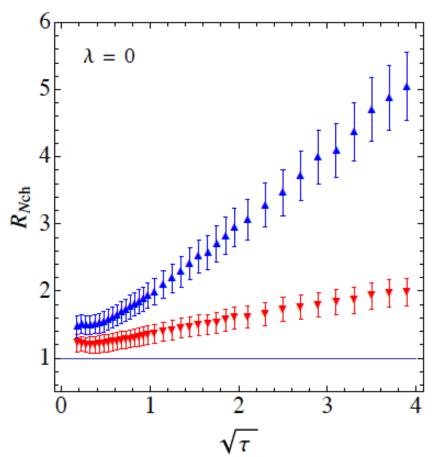


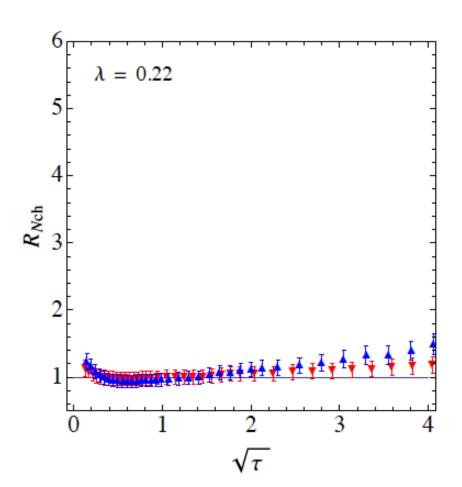


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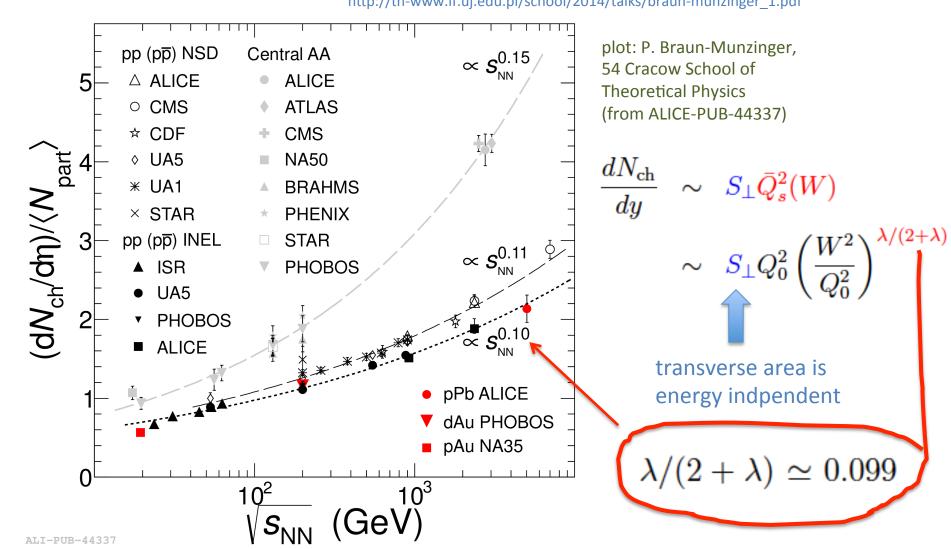
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#### Power-like growth of multiplicity

http://th-www.if.uj.edu.pl/school/2014/talks/braun-munzinger 1.pdf





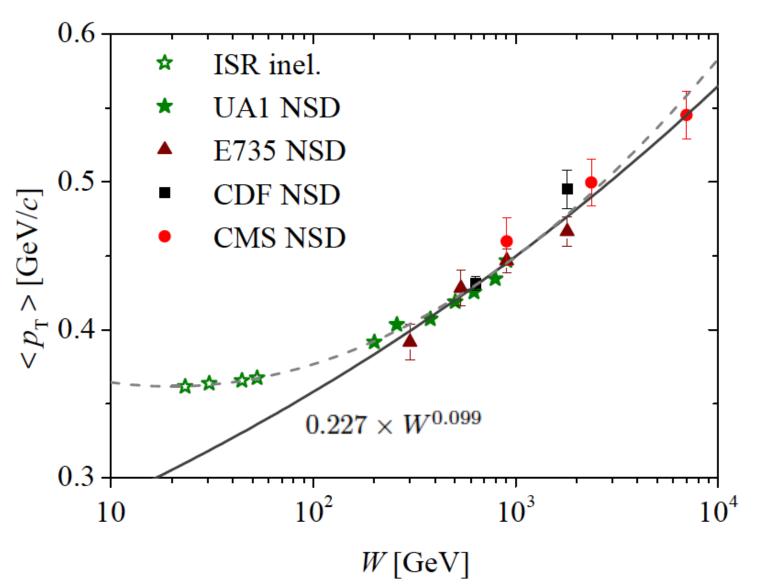
#### Average transverse momentum

$$rac{dN_{
m ch}}{dyd^2p_{
m T}} = S_{\perp} \, {\cal F}( au)$$

$$\langle p_{\rm T} \rangle = \frac{\int p_T \frac{dN_{\rm g}}{dy d^2 p_{\rm T}} d^2 p_T}{\int \frac{dN_{\rm g}}{dy d^2 p_{\rm T}} d^2 p_T} \sim \bar{Q}_{\rm s}(W) \sim Q_0 \left(\frac{W}{Q_0}\right)^{\lambda/(2+\lambda)}$$



#### Average transverse momentum





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- geometrical scaling predicts energy dependence of  $\langle p_T \rangle$



#### Mean $p_T$ as a function of $N_{ch}$

$$\langle p_{\rm T} \rangle \sim \bar{Q}_{\rm s}(W)$$



#### Mean $p_T$ as a function of $N_{ch}$

interaction radius

$$\langle p_{\rm T} \rangle \sim \bar{Q}_{\rm s}(W) \sim \sqrt{\frac{dN/dy}{S_{\perp}}} \sim \frac{1}{R} \sqrt{\frac{dN}{dy}}$$



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interaction radius

phenomenological formula:

$$\langle p_{\mathrm{T}} 
angle = rac{lpha}{\hbar} + rac{1}{R} \sqrt{rac{dN}{dy}}$$

nonperturbaitive coefficient

 $\alpha$ ,  $\beta$  do not depend on energy, do depend on particle species

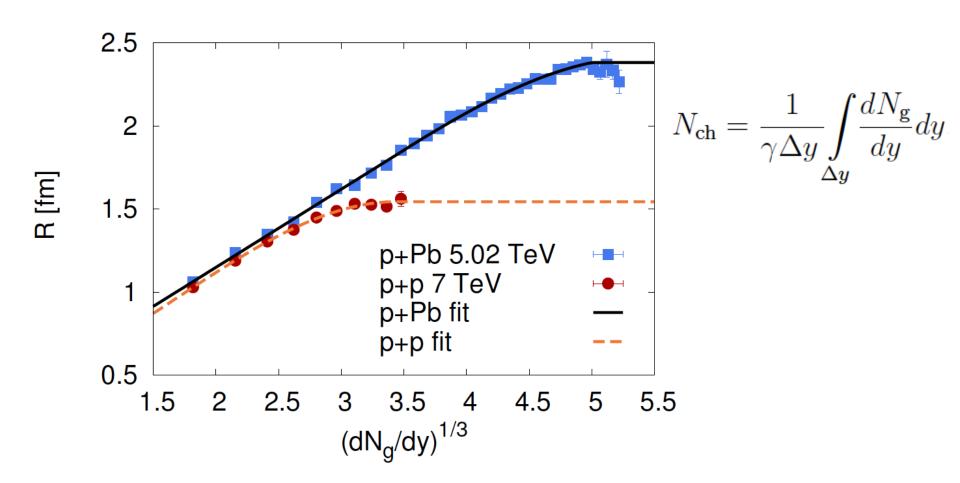


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#### Interaction radius

A. Bzdak, B. Schenke, P. Tribedy and R. Venugopalan, *Initial state geometry and the role of hydrodynamics in proton-proton, proton-nucleus and deuteron-nucleus collisions*, Phys. Rev. C 87 (2013) 064906, [arXiv:1304.3403 [nucl-th]].



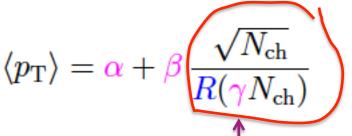


#### Scaling of mean $p_T$

$$\langle p_{
m T} 
angle = lpha + eta \, rac{\sqrt{N_{
m ch}}}{R(\gamma N_{
m ch})}$$
 parton-hadron duality  $\uparrow$ 



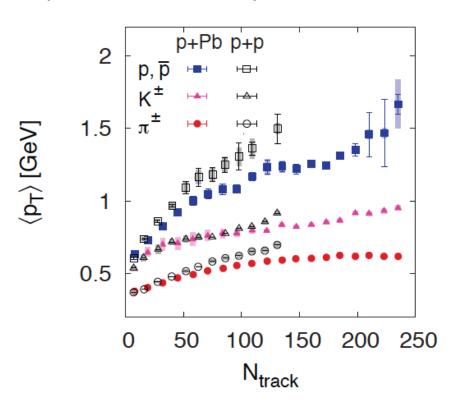
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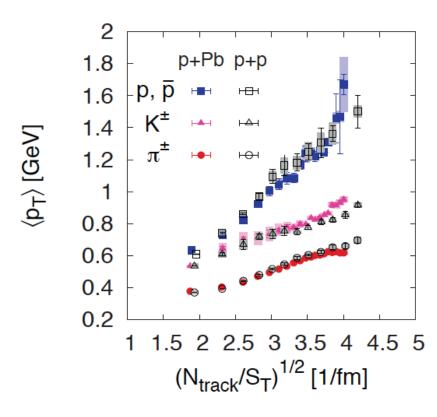


scaling variable

parton-hadron duality

CMS Collaboration, Eur. Phys. J. C72 (2013) 2164, C74 (2014) 2847





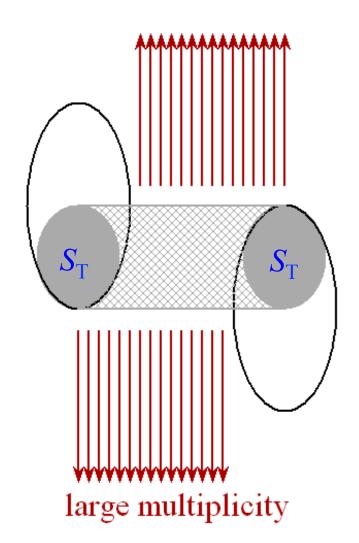


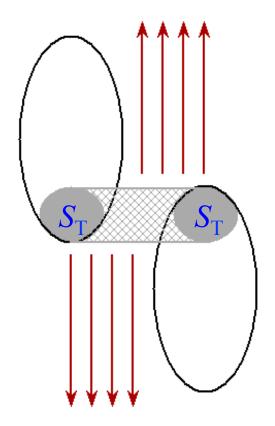
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#### Interaction radius

Transverse size and expansion time (longitudinal size) are proportional for fixed multiplicity





similar effect in multipomeron model, where string tension is growing with multiplicity

M. A. Braun, C. Pajares Phys. Lett. B 287, 154(1992) Nucl. Phys. B 390, 542, 559 (1993)

N. Armesto, D.A. Derkach, G.A. Feofilov Phys. of At. Nuclei 71, 2087 (2008)

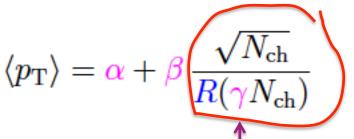
small multiplicity



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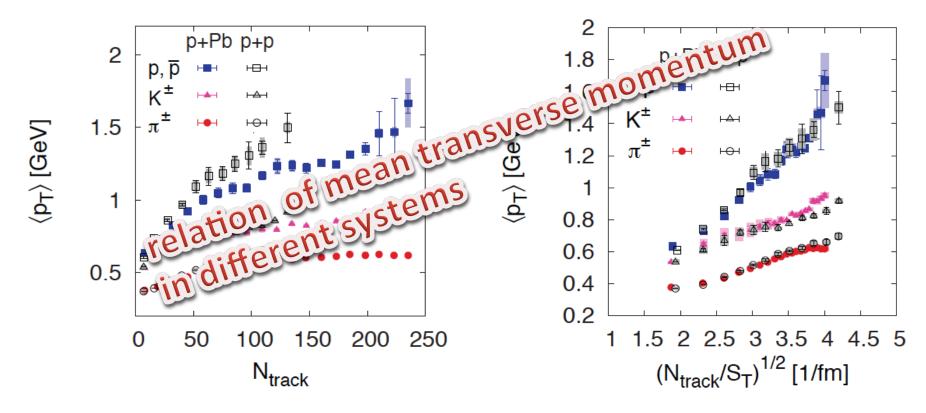
#### Scaling of mean $p_T$



scaling variable

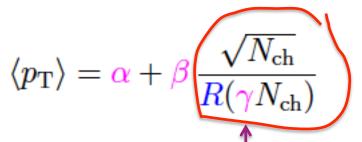
parton-hadron duality

CMS Collaboration, Eur. Phys. J. C72 (2013) 2164, C74 (2014) 2847





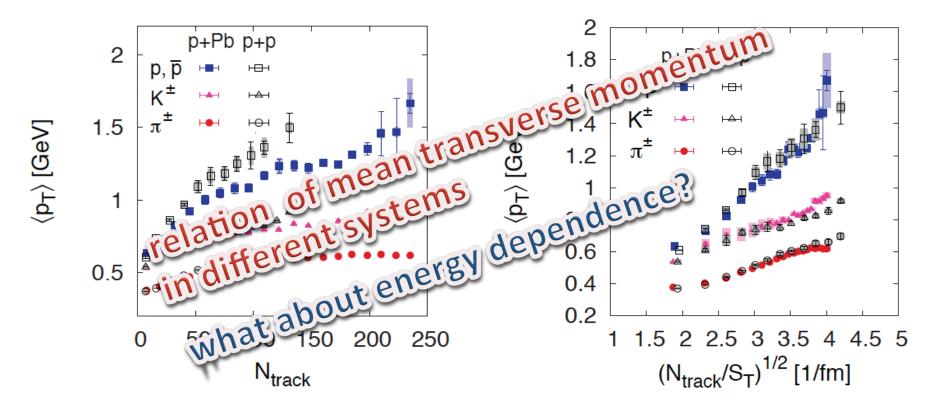
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scaling variable

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CMS Collaboration, Eur. Phys. J. C72 (2013) 2164, C74 (2014) 2847



## Energy dependence of mean $p_T$ - apparent paradox?



$$rac{dN_{
m ch}}{dy} \sim S_{\perp} \bar{Q}_s^2(W)$$

$$\sim S_{\perp} Q_0^2 \left(\frac{W^2}{Q_0^2}\right)^{\lambda/(2+\lambda)}$$

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If one *fixes* multiplicity and *then* changes energy, transverse area has to change accordingly

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energy indpendent



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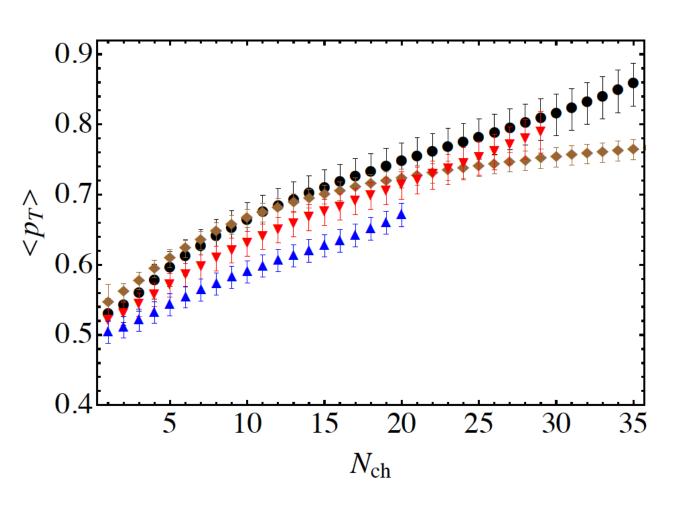
new scaling variable

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#### Mean $p_T$ scaling

ALICE Collaboration, Phys. Lett. B727 (2013) 371 [arXiv:1307.1094 [nucl-ex]]

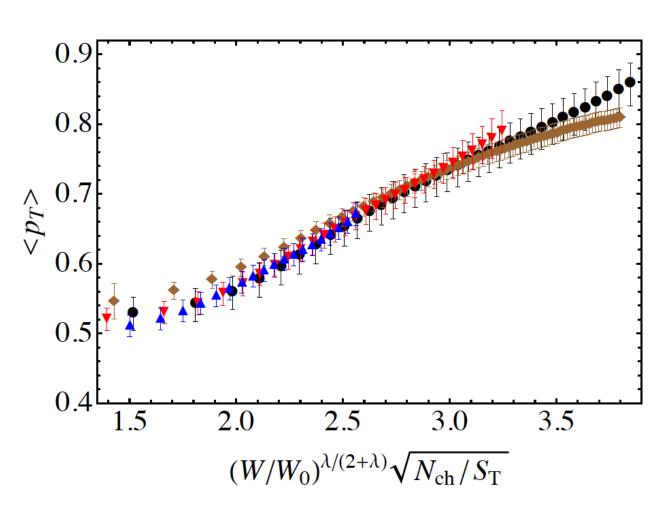


pp 7 TeV pp 2.76 TeV pPb 5.02 TeV pp 0.9 TeV



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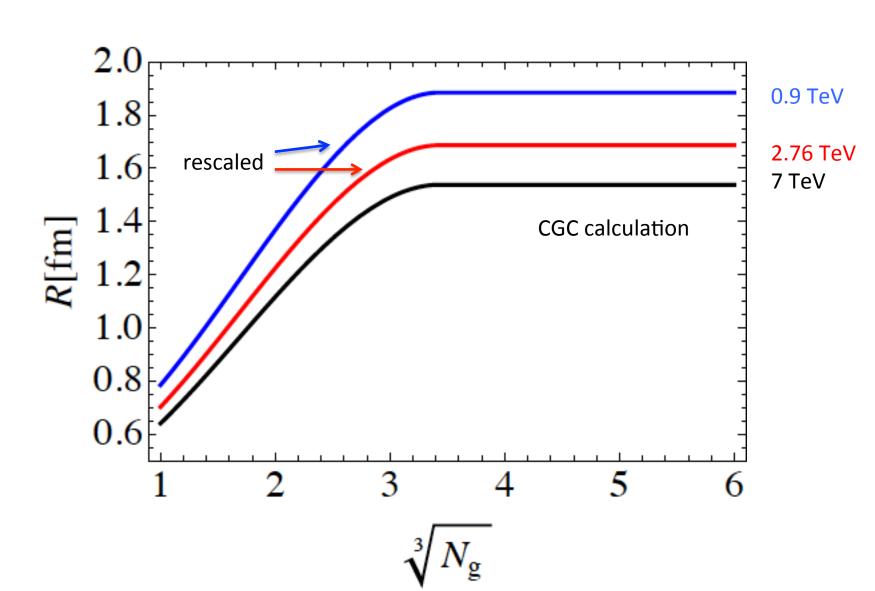
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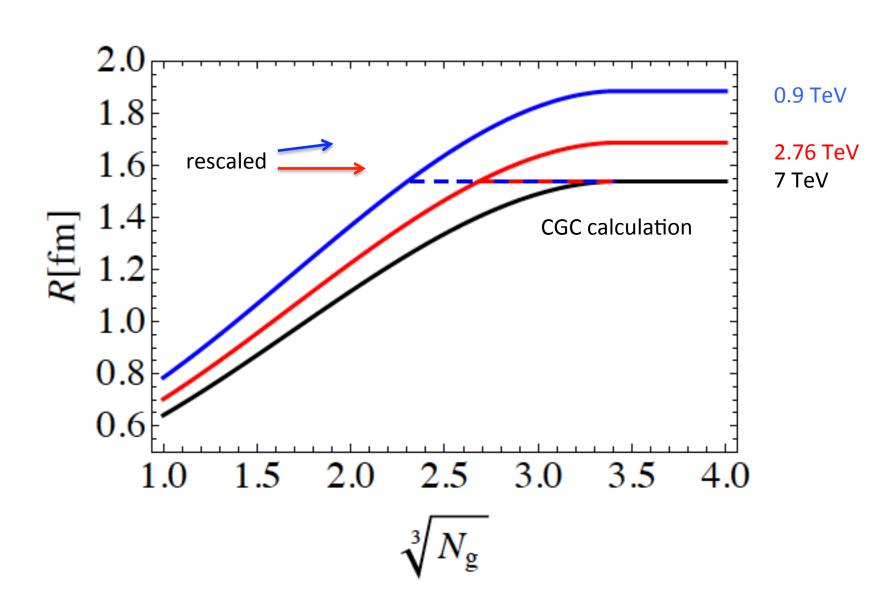


#### Effective R dependence on W





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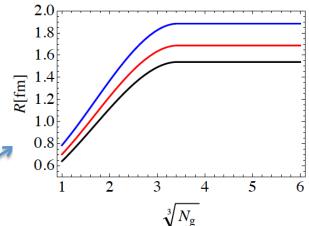




## Mean $p_T$ vs. interaction radius R

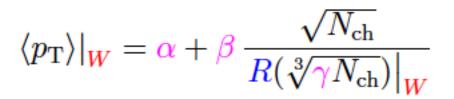
$$\langle p_{\mathrm{T}} 
angle |_{W} = \alpha + \beta \left. \frac{\sqrt{N_{\mathrm{ch}}}}{R(\sqrt[3]{\gamma N_{\mathrm{ch}}}) \right|_{W}}$$

fit  $\alpha$ ,  $\beta$  and  $\gamma$  to 7 TeV and then use R from Fig.

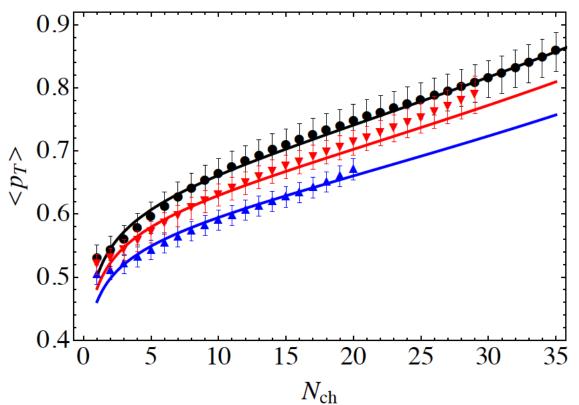


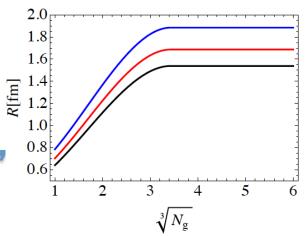


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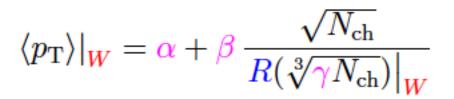




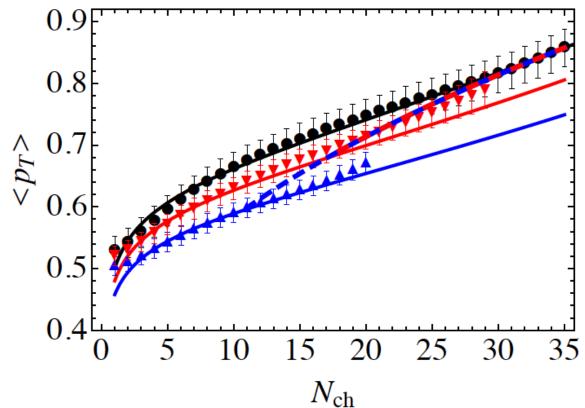
pp 7 TeV pp 2.76 TeV pp 0.9 TeV

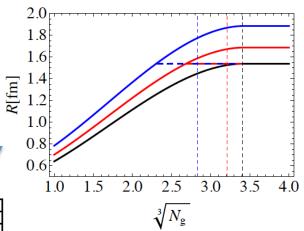


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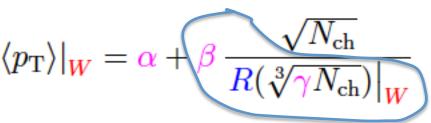


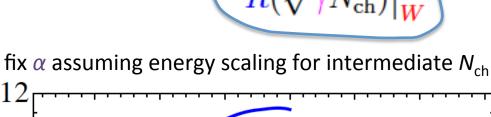
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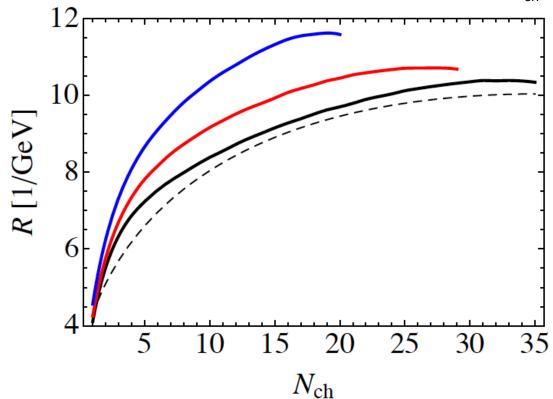


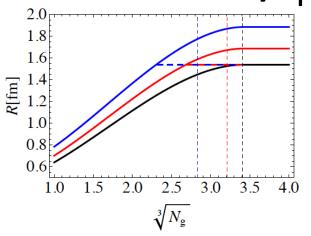
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- interaction radii cannot be too large, new saturation seen in data?

#### Interaction radius R from mean $p_{T}$



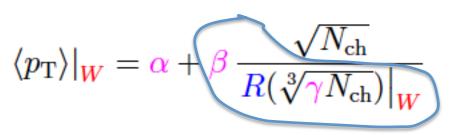




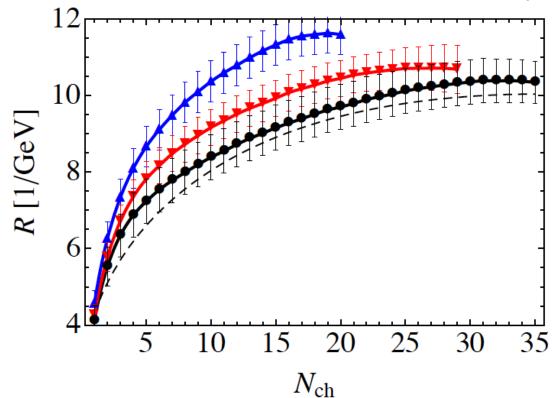


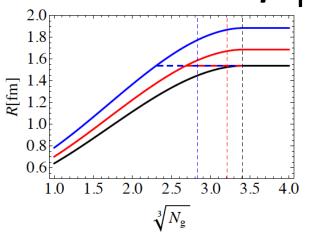
extracted interaction radii at 2.76, 0.9 and 7 TeV

#### Interaction radius R from mean $p_{T}$



fix  $\alpha$  assuming energy scaling for intermediate  $N_{\rm ch}$ 





extracted interaction radii at 2.76, 0.9 and 7 TeV



- $\langle p_T \rangle (N_{ch})$  correlations are sensitive to the fine details of dynamics
- complete change of behavior from small to large energies
- possible sign of phase transition
- difficult to describe by untuned MonteCarlos
- saturation included in EPOS does a good job
- geometrical scaling predicts energy dependence of  $\langle p_T \rangle$
- sensitivity to the interaction radius dependence on  $N_{\rm ch}$
- radii calculated from CGC describe well  $\langle p_{\top} \rangle$  in different systems
- longitudinal vs. transverse sizes one scale problem
- induced energy dependence of  $S_T$  for fixed  $N_{ch}$
- scaling of  $\langle p_T \rangle (N_{ch})$  induced by energy dependence of  $Q_{sat}$
- interaction radii cannot be too large, new saturation seen in data?
- $< p_T > (N_{ch})$  is sensitive to the space-time characteristics of the interaction volume, at <u>large multiplicities</u> universal behavior expected



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