Some news about GPDs & hard exclusive processes Dieter Müller

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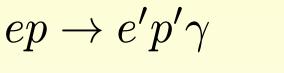
- From quasi-real to deeply virtual Compton scattering
- Higher twist contributions to DVCS
- One-to-one maps among GPD representations
- Data description/predicitions (partially includes DVMP)
- Unifying the partonic picture

Conclusions

A. Belitsky, Y. Ji;
V. Braun, A. Manashov; B. Pirnay,
M. Polyakov, K. Semenov-Tian-Shansky; O. Teryaev
K. Kumerički (KK), E. Aschenauer, S. Firzo, M. Murray, ...
K. Passek-Kumerički (KP-K), T. Lautenschlager, A. Schäfer; M. Meskauskas
D.S. Hwang

GPD related hard exclusive processes

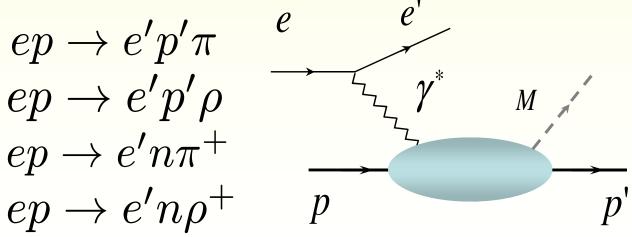
• Deeply virtual Compton scattering (clean probe)



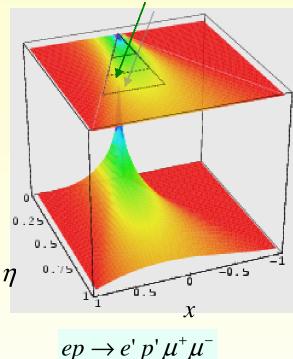
$$ep \rightarrow e'p'\mu^+\mu^-$$

 $\gamma p \rightarrow p' e^- e^+ \qquad p$ factorization proof for transversal cross sections [Collins Freund (99)]

• Deeply virtual meson production (flavor filter)



scanned area of the surface as a functions of lepton energy



twist-two observables:

longitudinal cross sections

transverse target spin asymmetries

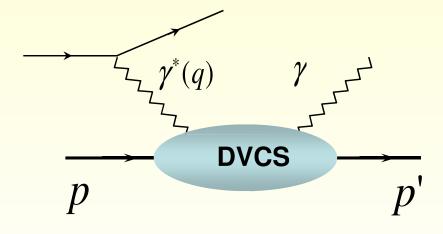
• etc.

factorization proof for longitudinal cross sections [Collins, Frankfurt, Strikman (96)]

GPDs embed non-perturbative physics

GPDs appear in various hard exclusive processes,

e.g., hard electroproduction of photons (DVCS)



[DM et. al (91/94) Radyushkin (96) **Ji (96)**]

$\mathcal{Q}^2 > 1 \text{GeV}^2$ GPD

 $t = \Delta^2 - \text{fix}$

 $\mathcal{F}(\xi, \mathcal{Q}^2, t) = \int_{-1}^1 dx \ C(x, \xi, \alpha_s(\mu), \mathcal{Q}/\mu) F(x, \xi, t, \mu) + O(\frac{1}{\mathcal{O}^2})$

CFF Compton form factor

observable

hard scattering part

perturbation theory

(our conventions/microscope)

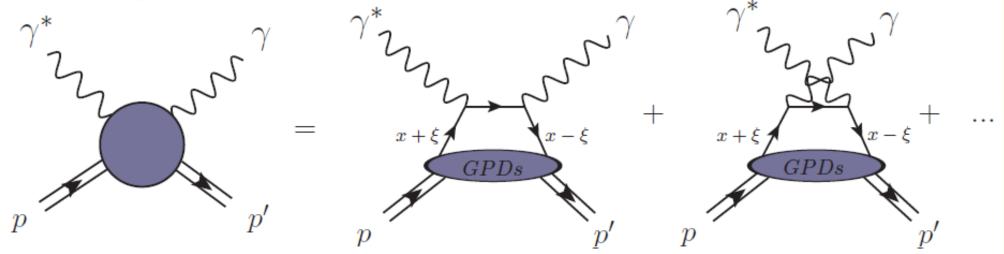
GPD

universal (conventional)

higher twist

depends on approximation

Higher twist contributions to DVCS



$$T_{\mu\nu} = i \int d^4x \, \mathrm{e}^{\frac{i}{2}(q_1+q_2)\cdot x} \langle p_2 | T \left\{ j_\mu(x/2) j_\nu(-x/2) \right\} | p_1 \rangle$$

• collinear factorization approach (calculating Feynman diagrams on partonic level)

operator product expansion (in terms of light-ray operators)

$$Tj_{\mu}(x/2)j_{\nu}(-x/2) \stackrel{\text{LO}}{=} \frac{S_{\mu\nu\alpha\beta}ix^{\alpha}}{(x^{2}-i\epsilon)^{2}} \left[\overline{\psi}(x/2)\gamma^{\beta}\psi(-x/2) - \overline{\psi}(-x/2)\gamma^{\beta}\psi(x/2)\right] \\ + \frac{i\epsilon_{\mu\nu\alpha\beta}ix^{\alpha}}{(x^{2}-i\epsilon)^{2}} \left[\overline{\psi}(x/2)\gamma^{\beta}\gamma^{5}\psi(-x/2) + \overline{\psi}(-x/2)\gamma^{\beta}\gamma^{5}\psi(x/2)\right]$$
• expansion in leading $1/x^{2}$ singularities is easily done by projection on the

• expansion in leading 1/x² singularities is easily done by projection on the light cone $n_{\mu} \sim q_{\mu} + ...$ and $n_{\mu}^* \sim P_{\mu} + ...$ or $n_{\mu} = q_{2\mu}$ and $n_{\mu}^* = q_{1\mu} + ... q_{2\mu}$ with $q_{\mu} = (q_{1\mu} + q_{2\mu})/2$ and $P_{\mu} = p_{1\mu} + p_{2\mu}$

$$T_{\mu\nu} \stackrel{\text{LO}}{=} -g_{\mu\nu}^{\perp} \sum_{q} \int_{-1}^{1} dx \left[\frac{e_{q}^{2}}{\xi - x - i\epsilon} - \frac{e_{q}^{2}}{\xi + x - i\epsilon} \right] q(x,\xi,t,\mathcal{Q}^{2}|s_{1},s_{2}) \\ -i\epsilon_{\mu\nu}^{\perp} \sum_{q} \int_{-1}^{1} dx \left[\frac{e_{q}^{2}}{\xi - x - i\epsilon} + \frac{e_{q}^{2}}{\xi + x - i\epsilon} \right] \tilde{q}(x,\xi,t,\mathcal{Q}^{2}|s_{1},s_{2})$$

consequences of 1/Q truncation and restriction to leading order in pQCD

- DVCS tensor structure depends on the choice of *n*
- scaling variable $\xi \sim x_B/(2-x_B)$ depends on the choice of *n*
- gauge and translation invariance holds only to leading power accuracy
- DVCS tensor structure is not complete

to overcome these problems one should go

- to twist-3 accuracy, yields 4 other GPDs (LT photon helicity flips) [done, 2000]
- to NLO, yields 4 gluon transversity GPDs (TT photon helicity flips) [done, 2000]
- twist-4 accuracy pushes ambiguity to the 1/Q⁴ level [Braun, Manashov 2011] but yields new parton correlation functions, however, no new structures

What is the problem in calculating higher twist contributions?

`correct' decomposition of $\partial^{\mu} \mathcal{O}_{\mu\mu_{1}\mu_{2}}... \stackrel{\text{tree}}{=} 0 \text{ and } \partial^{2} \mathcal{O}_{\mu\mu_{1}\mu_{2}}...$ guiding principle conformal covariance (same evolution as twist-2 operators)

Conformal PWE of GPDs

• GPD support is a consequence of Poincaré covariance (polynomiality)

$$H_j(\eta, t, \mu^2) = \int_{-1}^{1} dx \, c_j(x, \eta) H(x, \eta, t, \mu^2) \,, \qquad c_j(x, \eta) = \eta^j C_j^{3/2}(x/\eta)$$

• conformal moments evolve autonomously (to LO and beyond in a special scheme)

$$\mu \frac{d}{d\mu} H_j(\eta, t, \mu^2) = -\frac{\alpha_s(\mu)}{2\pi} \gamma_j^{(0)} H_j(\eta, t, \mu^2)$$

• inverse relation is given as series of (mathematical) generalized distributions:

$$H(x,\eta,t) = \sum_{j=0}^{\infty} (-1)^{j} p_{j}(x,\eta) H_{j}(\eta,t) , \ p_{j}(x,\eta) \propto \theta(|x| \le \eta) \frac{\eta^{2} - x^{2}}{\eta^{j+3}} C_{j}^{3/2}(-x/\eta)$$

- various ways of resummation were proposed:
- smearing method [Radyushkin (97); Geyer, Belitsky, DM., Niedermeier, Schäfer (97/99)]
- mapping to a kind of forward PDFs [A. Shuvaev (99), J. Noritzsch (00)]
- `dual' parameterization [M. Polyakov, A. Shuvaev (02), Polyakov (07), Semenov-Tian-Shansky]
- based on conformal light-ray operators [Balitsky, Braun (89); Kivel, Mankewicz (99)]
- Mellin-Barnes integral [DM, Schäfer (05); A. Manashov, M. Kirch, A. Schäfer (05)]

Purely mathematical problems/exercises:

- show that GPD representations are in one-to-one correspondence
- give inverse transformation (in principle done for Mellin-Barnes integral)
- support properties of Shuvaev's 'forward-like GPDs' is not known (i.e. Shuvaev's claim is wrong -- known since more than one decade that)

New: [DM, M. Polyakov, K. Semenov-Tian-Shansky (??)]

i. it is explicitly shown that `dual parameterization' and Mellin-Barnes integral + SO(3) PWE are equivalent

$$\begin{split} H(x \ge -\eta, \eta, t)) \ &= \ \sum_{\nu=0}^{\infty} \int_{0}^{1} dy K_{2\nu}(x, \eta | y) y^{2\nu} Q_{2\nu}(y, t) + `D - term', \\ H(x \ge -\eta, \eta, t) \ &= \ \sum_{\nu=0}^{\infty} \frac{1}{2i} \int_{c-i\infty}^{c+i\infty} dj \, \frac{\eta^{2\nu} p_{j+2\nu}(x, \eta)}{\sin(\pi[j+1])} \, H_{j+2\nu, j+1}(t) \, \hat{d}_{00}^{j+1}(\eta) + `D - term' \end{split}$$

(state-of-art formalism (NLO/NNLO, twist-4) is worked out in MB representation)

ii. (numerical) map of double distributions to conformal moments by means of Appell's F_4 function

Strategies to analyze DVCS data

(ad hoc) modeling: VGG code [Goeke et. al (01) based on Radyushkin's DDA] BMK model [Belitsky, DM, Kirchner (01) based on RDDA] `aligned jet' model [Freund, McDermott, Strikman (02)] Goloskokov/Kroll (05) based on RDDA (pinned down by DVMP)
`dual' model [Polyakov,Shuvaev 02;Guzey,Teckentrup 06;Polyakov 07] " -- " [KMP-K (07) in MBs-representation] polynomials [Belitsky et al. (98), Liuti et. al (07), Moutarde (09)]

dynamical models: not applied [Radyushkin et.al (02); Tiburzi et.al (04); Hwang DM (07)]... (respecting Lorentz symmetry)

flexible models:any representation by including unconstrained degrees of freedom(for fitting)KMP-K (07/08) for H1/ZEUS in MBs-integral-representation

CFFs (real and imaginary parts) and GPD fits/predictions

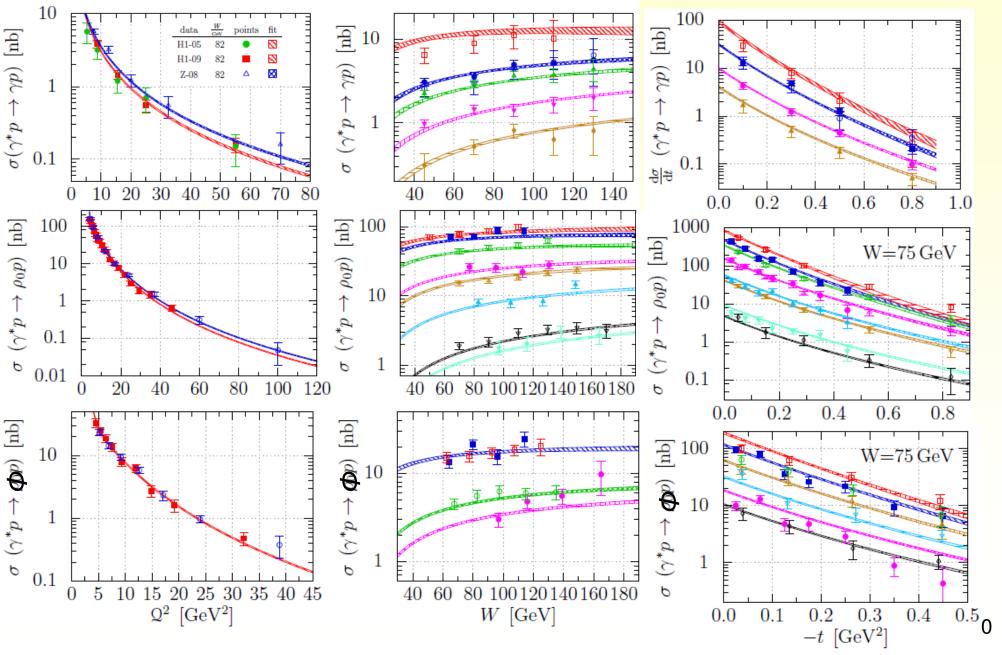
i. CFF extraction with formu	ulae (local) [BMK (01), HALL-A (06)] and [KK,DM, Murray]
least	square fits (local) [Guidal, Moutarde (08)]
neura	al networks – a start up <mark>[KMS (11)]</mark>
ii. `dispersion integral' fits	[KMP-K (08),KM (08)]
iii. flexible GPD model fits	[KM (08), AFKM (13), KMM (13), LSM (13)]
	/GG code, however also BMK01 (up to 2005)
•	Goloskokov/Kroll (07) model based on RDDA ⁸
[DVCS: Kroll, Moutarde, Sabatie (13)]

Status of theory	Mankiewicz et. al (97); Ji,Osborne (97/98);
<i>twist-two</i> DVCS coefficients at <i>next-to-leading</i> order	Pire, Szymanowski, Wagner (11); DM, Pire, Szymanowski
large-ξ: logarithmical enhancement valence region: weak evolution implies moderate effects K. F	Wagner (11)]
✓ anomalous dimensions and evolution kernels at <i>next-to</i>	-leading order
evolution effects can be called moderate, except for H/E at sn NLO analyses have to include NLO evolution	+ Fleund (01)]
✓ gluon transversity at <i>next-to-leading</i> order [Belitsky, DM	[DM (06); (00)] KMP-K, Schaefer (06)]
<i>next-to-next-to-leading</i> order for DVCS in a specific conform	nal subtraction scheme
 NLO → NNLO corrections can be called moderate w.r.t. LO → ✓ twist-three including quark-gluon-quark correlation at LO ✓ partially, twist-three sector at next-to-leading order [Kivel, Mathematication] 	NLO [Anikin,Teryaev, Pire (00); Polyakov et. al (00), Belitsky DM (00); Kivel et. al, Weiss, Radyushkin (00)] ankiewicz (03)]
? `target mass corrections' (not understood) [Belitsky DM (01)]	9
kinematical twist-four corrections [Braun, Manashov (11)]	

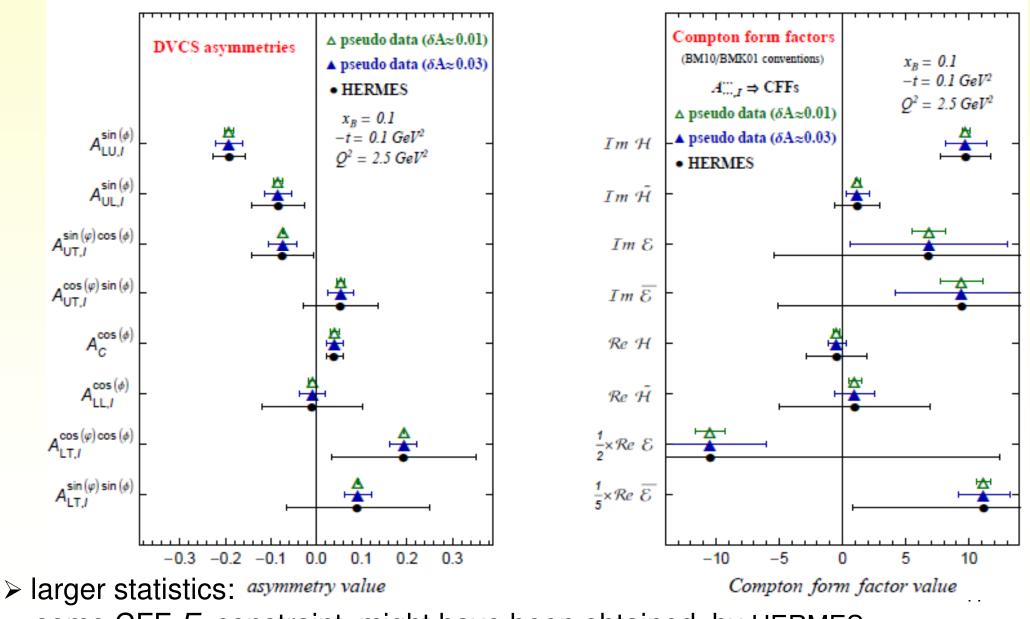
[Belitsky, DM (97);

DIS+DVCS+DVMP phenomenology at small-x_B (H1,ZEUS)

works somehow without DIS at LO [T. Lautenschlager, DM, A. Schäfer (13)] works at NLO ($Q^2 > 4 \text{ GeV}^2$), done with Bayes theorem (probability distribution function)



- > a complete measurement allows in principle to pin down all CFFs KK, DM, Murray (13)
- missing information in incomplete measurements can be filled with noise (Guidal's philosophy: use noise together with hypotheses and model constraints, our results are compatible)



some CFF E constraint might have been obtained by HERMES

KM models are available at WWW

http://calculon.phy.hr/gpd/ — binary code for cross sections

% xs.exe

xs.exe ModelID Charge Polarization Ee Ep xB Q2 t phi

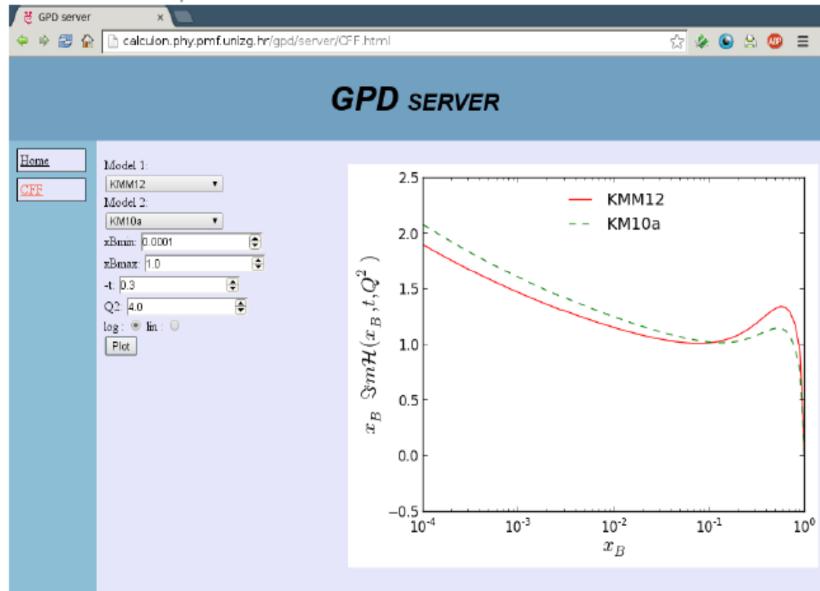
returns cross section (in nb) for scattering of lepton of energy Ee on unpolarized proton of energy Ep. Charge=-1 is for electron.

```
ModelID is one of
0 debug, always returns 42,
1 KM09a - arXiv:0904.0458 fit without Hall A,
2 KM09b - arXiv:0904.0458 fit with Hall A,
3 KM10 - preliminary hybrid fit with LO sea evolution, from Trento presentation,
4 KM10a - preliminary hybrid fit with LO sea evolution, without Hall A data
5 KM10b - preliminary hybrid fit with LO sea evolution, with Hall A data
xB Q2 t phi -- usual kinematics (phi is in Trento convention)
% xs.exe 1 -1 1 27.6 0.938 0.111 3. -0.3 0
```

0.18584386497251

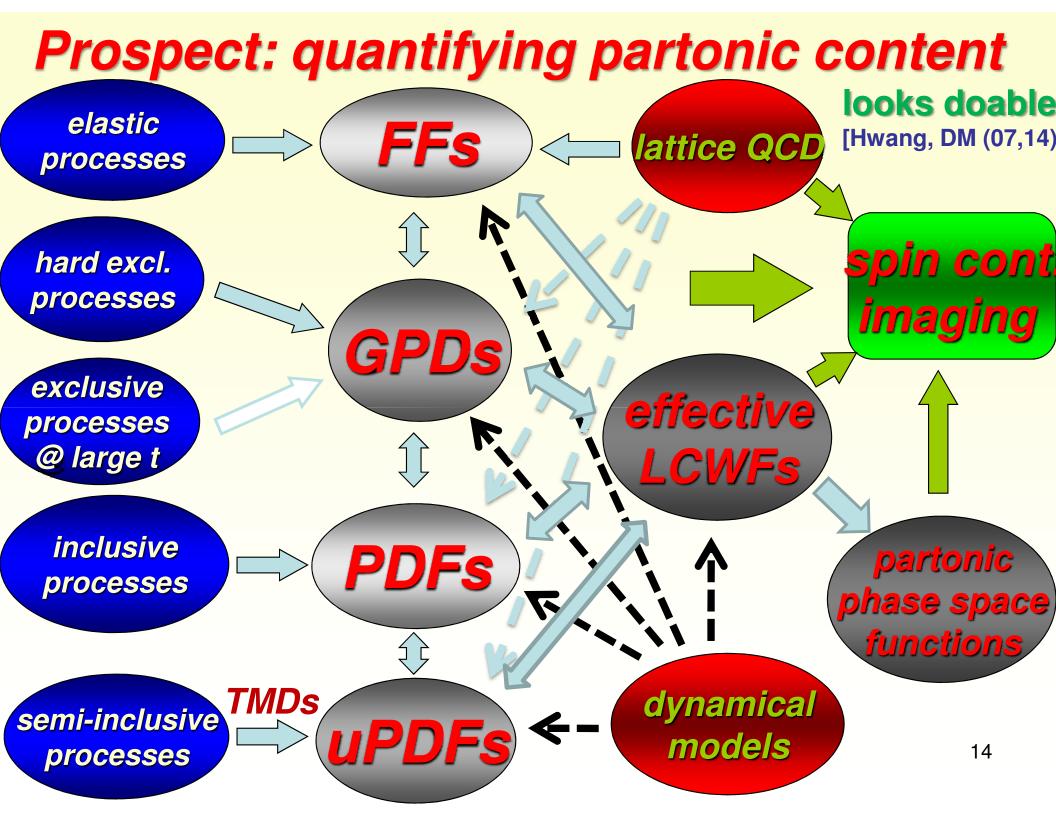
GPD page and server

Durham-like CFF/GPD server page



Ξ

• Do we need "Les Houches Accord" CFF/GPD interface?



Summary GPDs are intricate and (thus) a promising tool

- > to reveal the transverse distribution of partons (to some extend done at small x_B)
- > to address the spin content of the nucleon (not possible at present in pheno.)
- > providing a bridge to LCWFs & non-perturbative methods (e.g., lattice)
- modeling in terms of effective LCWFs is doable (requires efforts)

first decade of hard exclusive leptoproduction measurements

- CFFs have their own interest, bridging low and high virtuality regimes
- should be straightforward to improve global (flexible) model fits to DVCS
- DVCS and DVMP data are describable in global NLO fits at small x
- moving on: to NLO, kinematical twist, full GPD models, DVCS+DVMP+...
- covering the kinematical region between HERA (COMPASS) experiments within a high luminosity machine and dedicated detectors is needed to quantify exclusive and inclusive QCD phenomena: handle on GPD E & 3D
- some kind of education is desired before one can enter GPD phenomenology
- theory development is needed to address phenomenological goals