

# Innovative X- $\gamma$ ray sources based on LASER-produced plasmas

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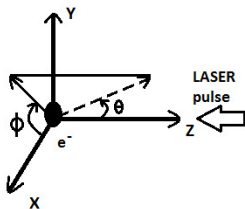
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# Overview

- 1 Theory of the TB radiation
  - Electron dynamics
  - Thomson Backscattering radiation
- 2 ENEA TB source
  - ABC + LINAC parameters
  - ABC + 5 MeV LINAC
  - ABC + 10 MeV LINAC

# Interaction



**Figure:** Geometrical configuration for the scattering between the LASER pulse and the electron bunch[1].

$$\xi = k_i x^i = \omega_0 \left( t + \frac{z}{c} \right)$$

$$\vec{a} = \hat{y} a_0 e^{-\frac{\xi^2}{2\omega_0^2 \tau^2}} \cos \xi \quad ; \quad a_0 \sim 8.5 \times 10^{-10} \sqrt{I_0 \lambda_0^2}$$

# Electron dynamics

$\omega_0 \tau \gg 1 \leftarrow$  *Long Pulse*

$$x = x_0 \xi$$

$$y = y_0 \xi + y_1 a_0 e^{-\frac{\xi^2}{2\omega_0^2 \tau^2}} \sin \xi$$

$$z = z_0 \xi + z_1 a_0 e^{-\frac{\xi^2}{2\omega_0^2 \tau^2}} \sin \xi + \frac{z_2}{2} a_0^2 e^{-\frac{\xi^2}{\omega_0^2 \tau^2}} \left( \xi - \frac{1}{2} \sin 2\xi \right)$$

## Radiation

Two variables generalized Bessel functions[2]

$$\frac{d^2 N}{d\Omega d\omega} = \frac{\alpha \omega}{(2\pi\omega_0)^2} \left| \sum_{n,l=-\infty}^{n,l=+\infty} \int d\xi J_n\left(\rho_1 \frac{\omega}{\omega_0}\right) J_{n-2l}\left(\rho_2 \frac{\omega}{\omega_0}\right) (\widetilde{\beta}_\theta, \widetilde{\beta}_\phi) e^{i(\rho_0 \frac{\omega}{\omega_0} - n)\xi} \right|^2$$

$$\rho_0 = 1 - \sin\theta \frac{\omega_0}{c} (x_0 \cos\phi + y_0 \sin\phi) - \frac{\omega_0}{c} (1 + \cos\theta) \left( z_0 + \frac{a_0^2}{2} z_2 \right)$$

$$\rho_1(\xi) = \frac{\omega_0}{c} a_0 e^{-\frac{\xi^2}{2\omega_0^2 \tau^2}} (y_1 \sin\phi \sin\theta + z_1 (1 + \cos\theta))$$

$$\rho_2(\xi) = -\frac{\omega_0}{c} \frac{a_0^2 e^{-\frac{\xi^2}{\omega_0^2 \tau^2}}}{4} z_2 (1 + \cos\theta)$$

# Linear Thomson Backscattering Radiation

$$a_0 \ll 1$$

$$\frac{d^2N}{d\Omega d\omega} = \frac{\alpha \omega T^2}{64\gamma_0^2 \pi^2} a_0^2 e^{\frac{-T^2}{4\tau^2}} \left(1 - 4 \frac{(\gamma_0 \theta \sin\phi - \gamma_0 \theta_e \sin\phi_e)^2}{(1 + (\gamma_0 \chi)^2)^2}\right) \text{sinc}^2\left(\frac{T}{2}(\rho_0 \omega - \omega_0)\right)$$

$$\text{Fundamental emission frequency} \rightarrow \omega_f = \frac{4\gamma_0^2 \omega_0}{1 + \gamma_0^2 \chi^2}$$

$$\chi = \sqrt{\theta^2 + \theta_e^2 - 2\theta_e \theta \cos(\phi - \phi_e)}$$

## Radiation from an electron bunch

$$\star \frac{d^2 N}{d\Omega d\omega} = \frac{F N_e \alpha}{4\pi^2} \omega a_0^2 e^{\frac{-T^2}{4\tau^2}} T^2 \int d\theta_e d\phi_e d\gamma_e \Theta(\theta_e) \Phi(\phi_e) \Gamma(\gamma_e) \left\{ \left( 1 - 4 \frac{(\gamma_e \theta \sin\phi - \gamma_e \theta_e \sin\phi_e)^2}{(1 + (\gamma_e \chi)^2)^2} \right) \text{sinc}\left(\frac{T}{2}(\rho_0 \omega - \omega_0)\right) \right\}$$

$$\star N_{ph}^{tot} = F N_e \frac{1}{2} \alpha \omega_0 T a_0^2 e^{\frac{-T^2}{4\tau^2}} \psi^2 \frac{1 + \psi^2 + \frac{2}{3}\psi^4 - \frac{1}{4}\gamma_0^2 (\theta_e^{div})^2 (1 + \psi^2 + \frac{1}{3}\psi^4)}{(1 + \psi^2)^3} \quad [3]$$

*Transversal filling factor*  $\rightarrow F \equiv \frac{\int dr_e R(r_e) e^{\frac{-2r_e^2}{w_0^2}}}{\int dr_e R(r_e)} = \frac{w_0^2}{w_0^2 + 4w_e^2}$

$$\star \psi = \gamma_0 \theta$$

## ABC + LINAC @ ENEA Research Center of Frascati

$$\omega_0 = \frac{2\pi c}{\lambda} = \frac{2\pi c}{1.054\mu m} \sim 1.79 \times 10^{15} \text{ rad/s}$$

$$\tau = 3 \text{ ns}$$

$$w_0 \sim 100 \mu m$$

$$E_L \sim 20 \text{ J}$$

$$I_0 = 2.12 \times 10^{13} \text{ W/cm}^2 \longrightarrow a_0 \sim 0.004$$

$$\epsilon_n \sim 2 \text{ mm mrad} \longrightarrow w_e \sim 2 \text{ mm}, \theta_e^{\text{max}} < 1 \text{ mrad}$$

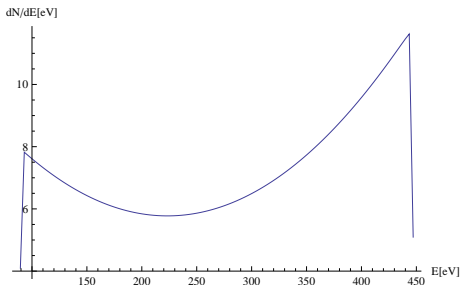
$$I_e \sim 6 \text{ A} \longrightarrow N_e \sim 5 \times 10^8$$

$$\tau_e = 15 \text{ ps}$$

$$\gamma_0 = 10, 20$$

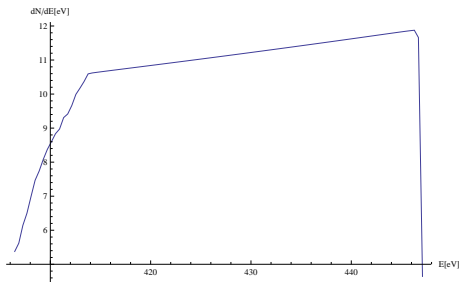


# Emission Spectrum, maximum acceptance



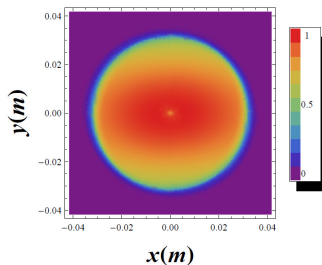
Thomson spectrum obtained by fixing the maximum acceptance semi-aperture  $\theta_{max} \sim 1/\gamma_0$  ( $\gamma_0 = 10$ ).

# Emission Spectrum, 10% monochromaticity



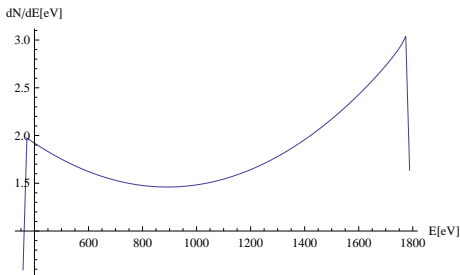
Thomson spectrum, relative to the case with  $\gamma_0 = 10$ , obtained by fixing the acceptance semi-aperture  $\theta_{max} \sim 1/(10\sqrt{10})$  in such a way to detect just radiation near 0.45 keV (on-axis) with a 10% monochromaticity degree.

# Radiation spot



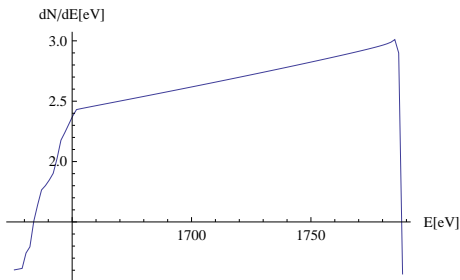
Thomson radiation spot on a screen at  $1\text{ m}$ , relative to the case with  $\gamma_0 = 10$ ; red stands for  $\sim 1.4 \times 10^2$  photons/cm<sup>2</sup>.

# Emission Spectrum, maximum acceptance



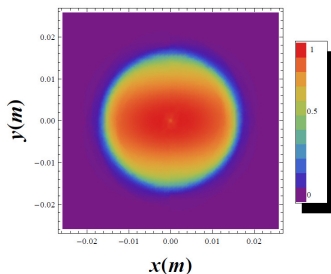
Thomson spectrum obtained by fixing the maximum acceptance semi-aperture  $\theta_{max} \sim 1/\gamma_0$  ( $\gamma_0 = 20$ ). The number of particles randomly generated for the calculation is  $\sim 10000$ .

# Emission Spectrum, 10% monochromaticity



Thomson spectrum, relative to the case with  $\gamma_0 = 20$ , obtained by fixing the acceptance semi-aperture  $\theta_{max} \sim 1/(20\sqrt{20})$  in such a way to detect just radiation near 1.8 keV (on-axis) with a 10% monochromaticity degree.

# Radiation spot



Thomson radiation spot on a screen at  $1\text{ m}$ , relative to the case with  $\gamma_0 = 20$ ; red stands for  $\sim 5.6 \times 10^2$  photons/cm<sup>2</sup>.

# References



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High Brilliance X- $\gamma$  ray sources based on LASER-matter interaction a high intensities

*Master degree thesis, D. Giulietti Supervisor, Physics Department of Pisa University*

Thank you for your time