

Channeling 2014

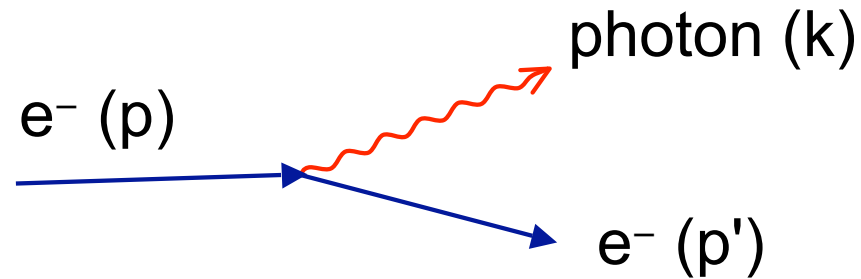
Oct. 5-10, **Capri**, Italy

Similarity between
Synchrotron radiation
and photons escaping from
a bent optical fiber

X. Artru and C. Ray
IPN-Lyon, France

Synchrotron case

The reaction :



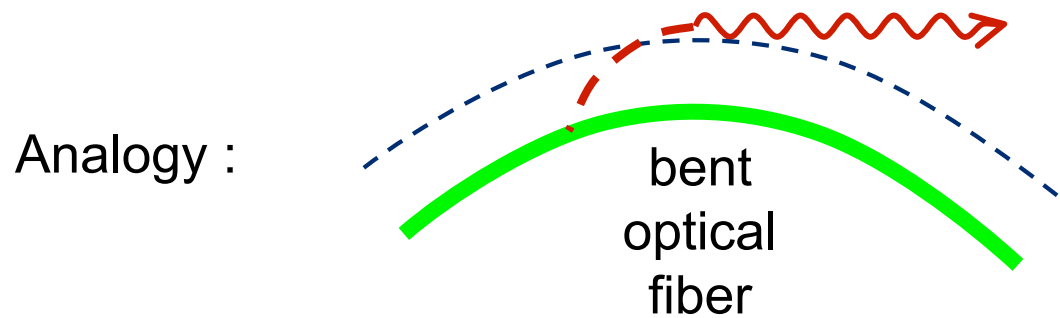
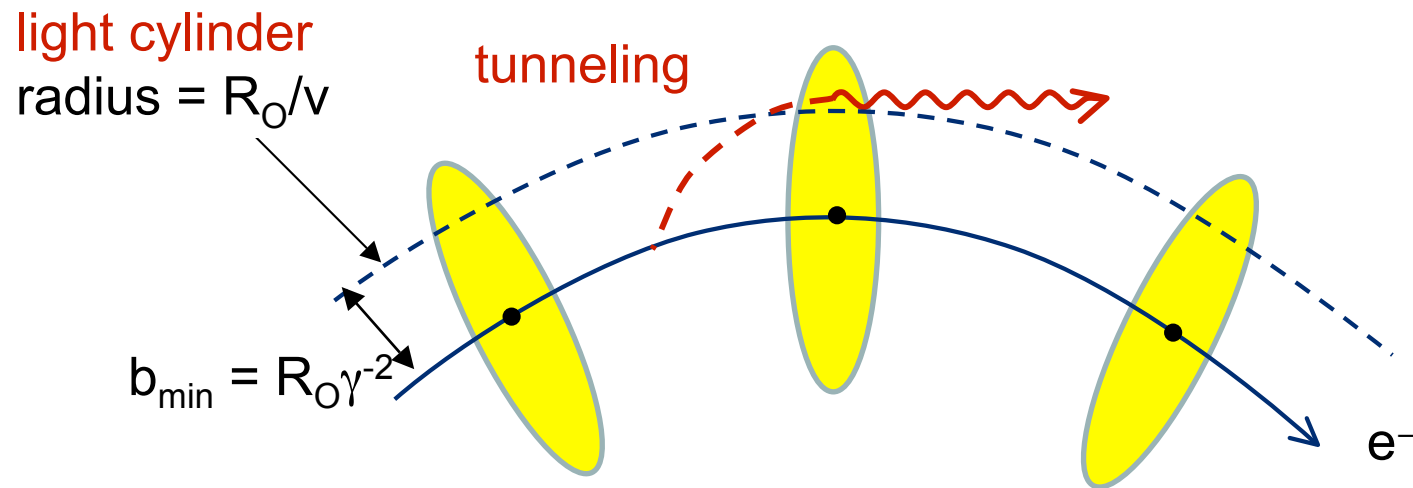
is kinematically forbidden

as a ***local*** process between ***classical particles***,

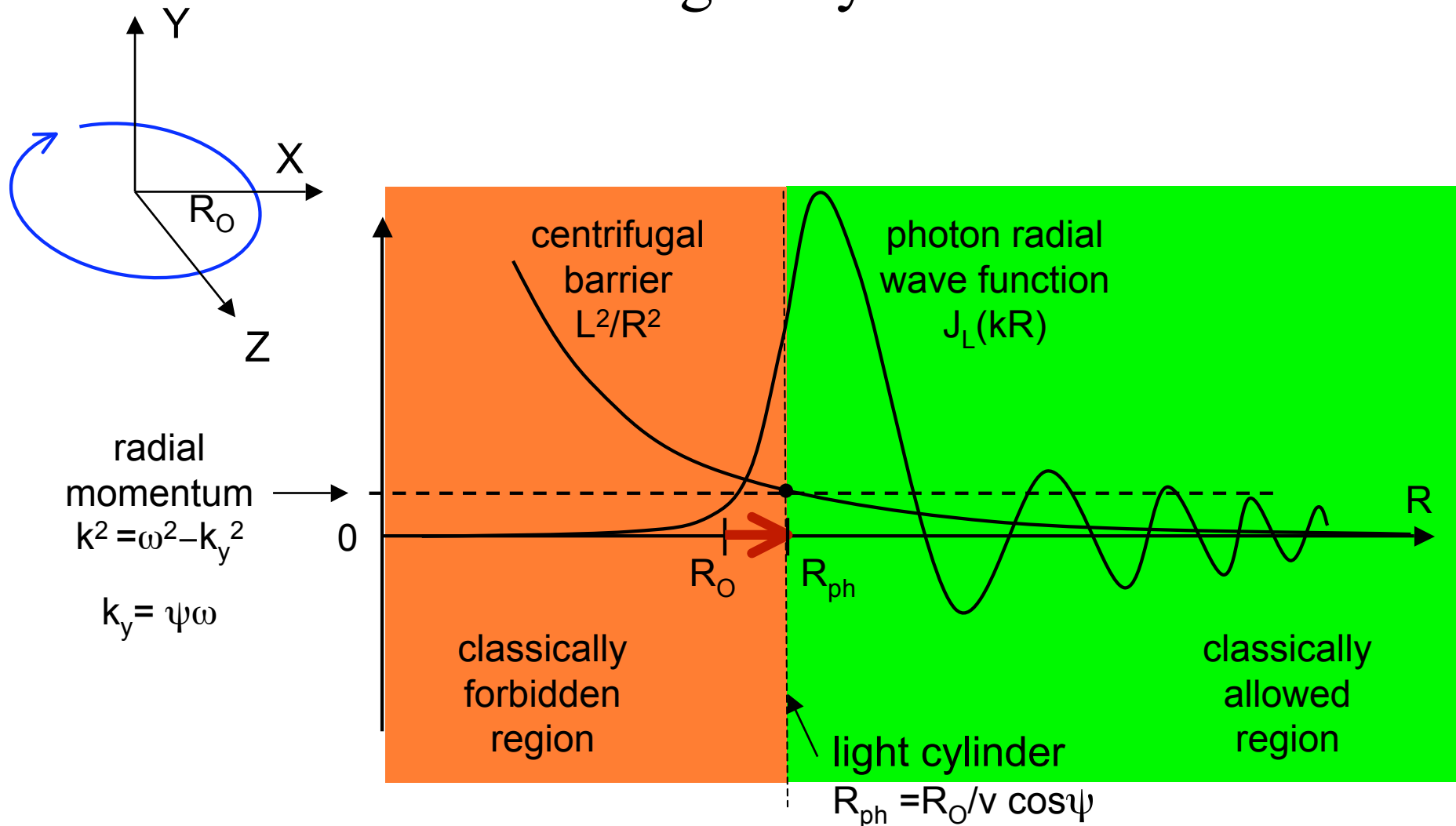
but it can occur via a quantum **tunnelling effect**.

Impact parameter point of view

Synchrotron radiation ~ **slow leakage** of the Coulomb field



Wave tunnelling in Synchrotron Radiation



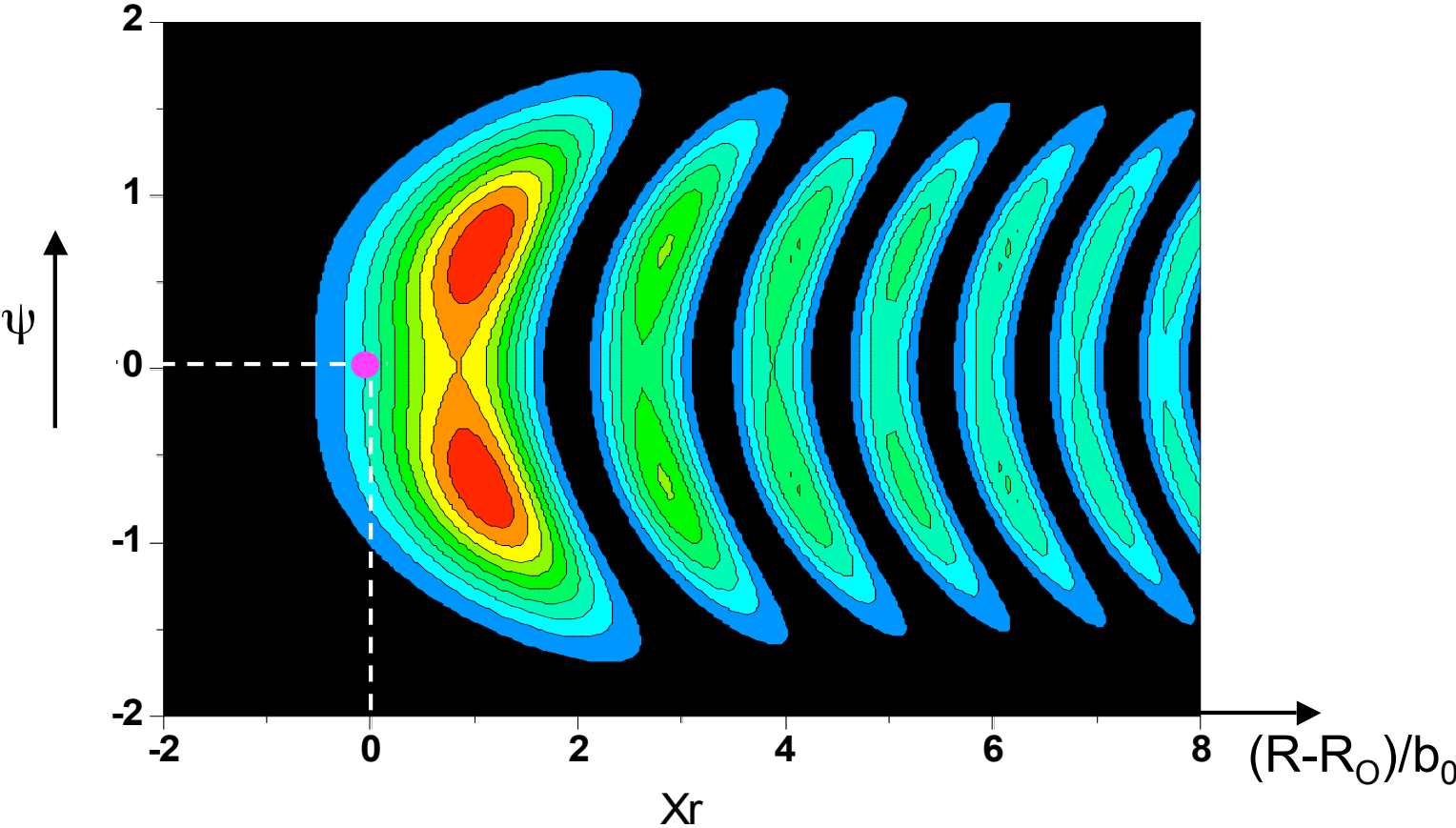
$$J_L(kR) \sim \text{Airy function } (R_{ph} - R)$$

$$\sim \exp\left\{ - (\omega/\omega_c) (1 + \gamma^2\psi^2)^{3/2} / 3 \right\} \leftarrow \text{tunnelling factor}$$

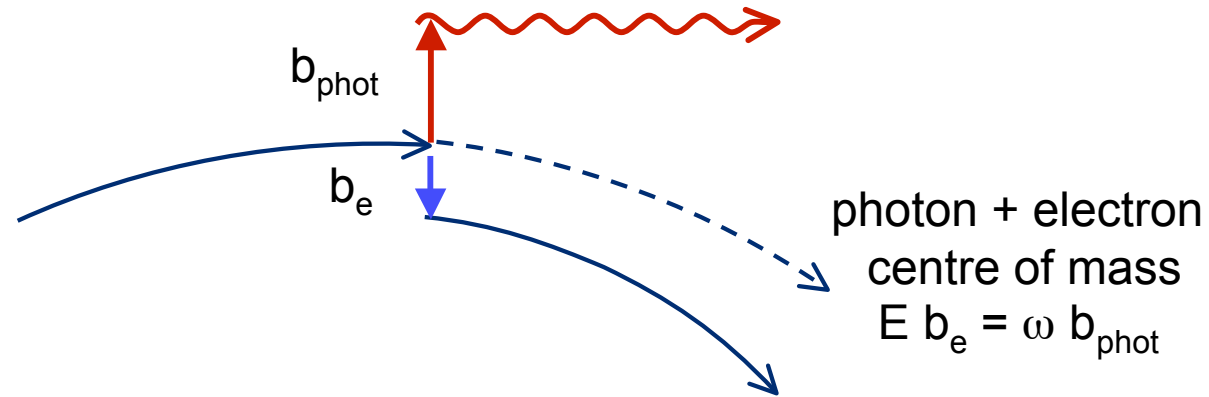
$$\omega_c = \gamma^3/R_0$$

Horizontal Profile of Synchrotron Radiation

[X. A. and C. Ray (2006)]

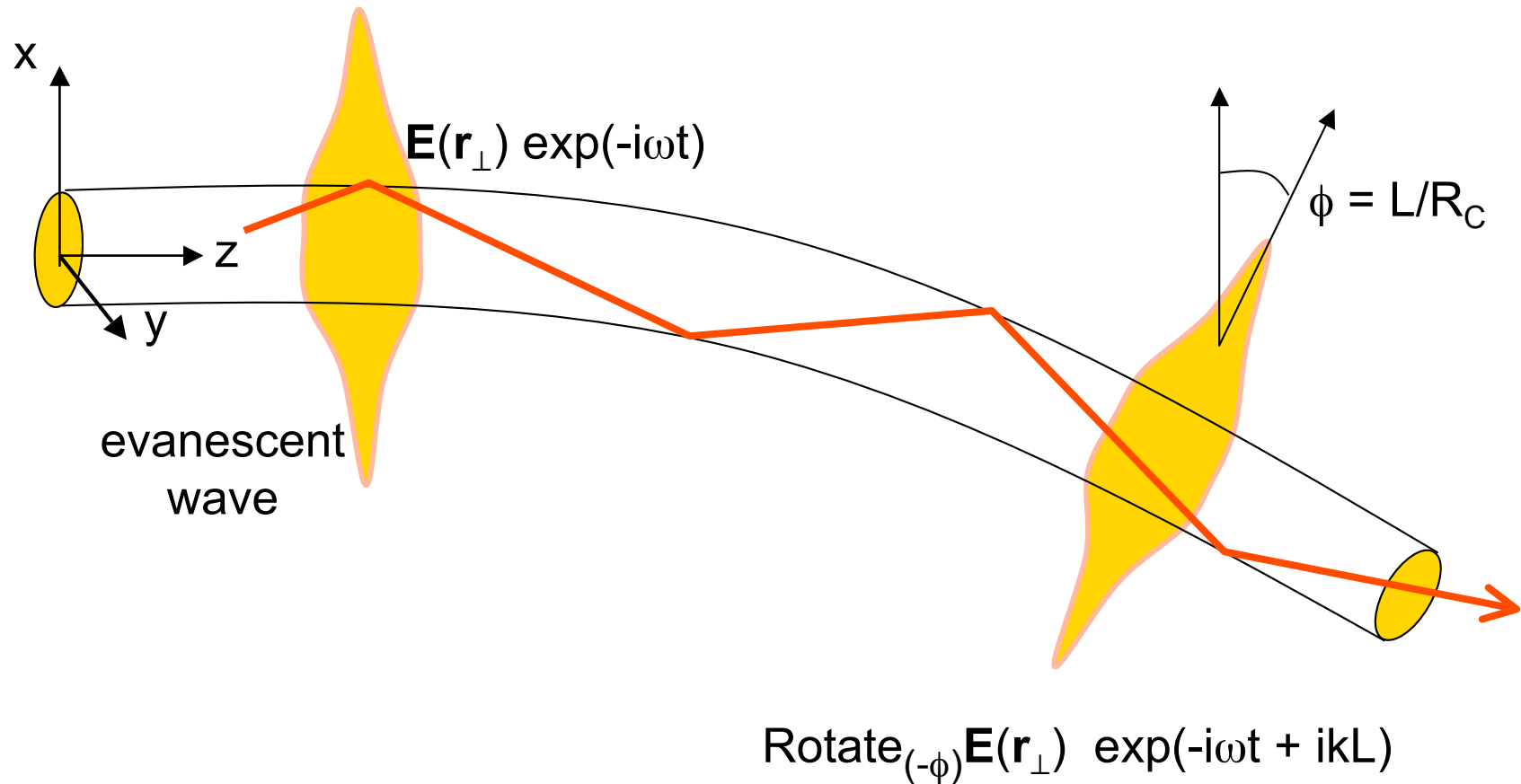


Side slipping

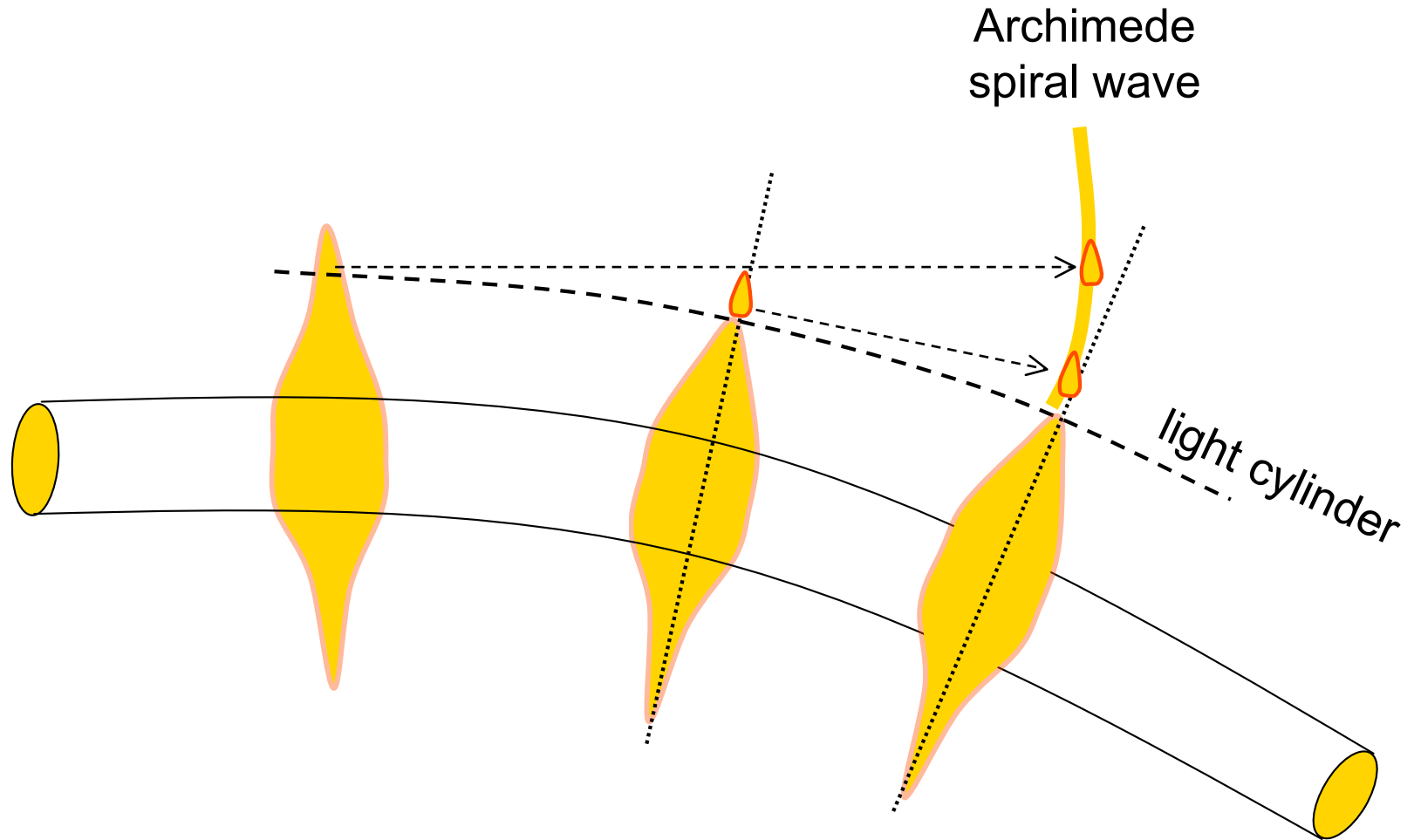


- $b_e \sim \lambda_C = 400 \text{ fm}$
- Side-slipping is responsible for the d^3X/dt^3 term of the Abraham-Lorentz equation

The field in a bent fiber



Escape of the field

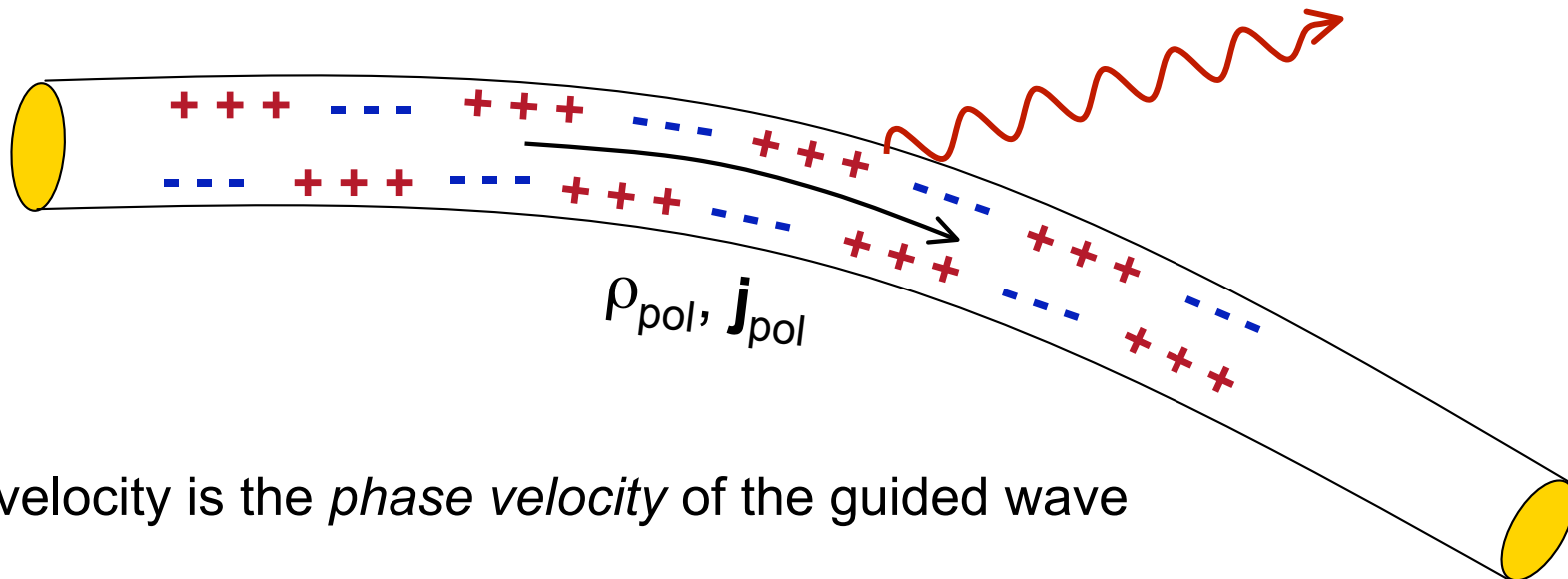


Mathematical treatment

Volume Current Method [M. Kuznetsov and H. Haus ; White] : the wave inside the fiber produces a polarization 4-current

$$\mathbf{j}_{\text{pol}}^{\mu} = (\rho_{\text{pol}}, \mathbf{j}_{\text{pol}}) ; \quad \rho_{\text{pol}} = \nabla \cdot \mathbf{E} ; \quad \mathbf{j}_{\text{pol}} = \nabla \times \mathbf{B} - \partial_t \mathbf{E}$$

One assumes that ρ_{pol} and \mathbf{j}_{pol} are *nearly the same as in a straight fiber*. This 4-current follows the bends of the fiber \Rightarrow it radiates a kind of Synchrotron Radiation.



The velocity is the *phase velocity* of the guided wave

Analytical results - momentum space

For a narrow fiber, $v_{\text{phase}} = \omega/k$ is close to 1. The escaping light is collimated as SR at $\gamma \gg 1$

Notations: $\theta_0 = (kR_0)^{-1/3}$; $\xi = (1/v - \cos\psi) / \theta_0^2$; $A(\xi) = 2^{4/3} \pi \text{Ai}(2^{1/3} \xi)$

Synchrotron Radiation :

$$\mathbf{E}(\mathbf{k}) = \mathbf{e} \theta_0 R_0 \left\{ -i\theta_0 A'(\xi) \mathbf{e}_{\text{horiz}} + \psi A(\xi) \mathbf{e}_{\text{vert}} \right\}$$


Optical Fiber : - replace \mathbf{e} by \mathbf{f}_z

- add 2 other terms : $-\theta_0 R_0 A(\xi) [f_x \mathbf{e}_{\text{horiz}} + f_y \mathbf{e}_{\text{vert}}]$


$\mathbf{f} = 2\text{-dim Fourier transform of } \mathbf{j}_{\text{pol}}(x,y) \text{ with argument } (0, k_y)$

Analytical results - impact parameter space

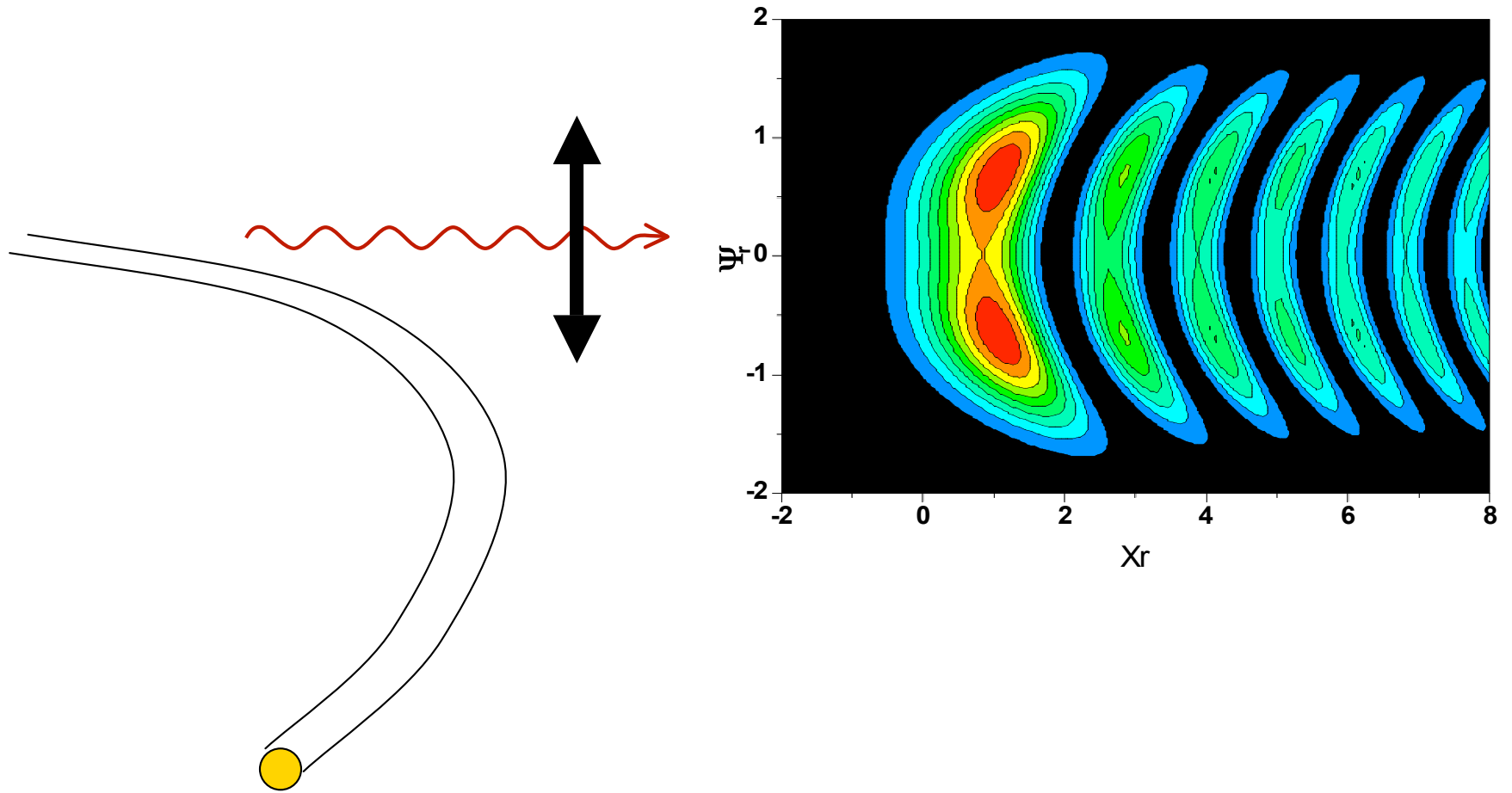
Synchrotron Radiation


$$dW/(dx, d\psi, d\omega) = (\omega\theta_0^2/2\pi) A^2(\xi) dW/(d\phi, d\psi, d\omega)$$

Optical Fiber


$$dW/(dx, d\psi, dt) = (\omega\theta_0^2/2\pi) A^2(\xi) dW/(d\phi, d\psi, dt)$$

Expected from a optical fiber experiment



Conclusion

- Bent fibers can be used to understand and test properties of Synchrotron Radiation.
- They have at least a pedagogical interest.

Thank you for your attention !