Schouten identities and the (diff. eq.s of the) two-loop sunrise graph

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Based on [arXiv:1311.3342], in collaboration with Ettore Remiddi

Why studying the two-loop massive sunrise?

- Phenomenologically: needed for any two-loop calculation in realistic theories with massive particles
 - 1. Electro-weak sector of SM
 - 2. Massive QCD (see $gg \rightarrow t\bar{t}$ at two loops!)
- ▶ Mathematically: it is the first (≈ *easiest*) diagram that escapes our understanding through *multiple-polylogs*.
 - 1. Functions: Elliptic functions, iterated integrals, concept of transcendental weight...?
 - 2. Differential Equation method: it is possible to find a canonical basis?

The two-loop massive sunrise graph

$$= \int \mathfrak{D}^{d} k \, \mathfrak{D}^{d} l \, \frac{(k \cdot p)^{n_{4}} \, (l \cdot p)^{n_{5}}}{(k^{2} + m_{1}^{2})^{n_{1}} (l^{2} + m_{2}^{2})^{n_{2}} ((k - l - p)^{2} + m_{3}^{2})^{n_{3}}}$$
$$= \int \mathfrak{D}^{d} k \, \mathfrak{D}^{d} l \, \frac{(k \cdot p)^{n_{4}} \, (l \cdot p)^{n_{5}}}{D_{1}^{n_{1}} D_{2}^{n_{2}} D_{3}^{n_{3}}}$$

How do we compute all these integrals (in dimensional regularisation !!) ?

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How do we compute all these integrals (in dimensional regularisation!!) ?

Standard way to proceed:

 Use Integration-by-parts id.s (IBPs) to reduce them to Master Integrals

$$\int \mathfrak{D}^d k \, \mathfrak{D}^d l \, \left(v_j^\mu \, \frac{\partial}{\partial \, k_i^\mu} \, \frac{(k \cdot p)^{n_4} \, (l \cdot p)^{n_5}}{D_1^{n_1} D_2^{n_2} D_3^{n_3}} \, \right) = 0$$

Use IBPs to derive differential equations (DE) in the external invariants for the MIs.

$$p^2 = p_\mu p^\mu \quad
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One finds for the Sunrise:

Sunrise graph reduced to 4 MIs

 $M_1(d; p^2), \quad M_2(d; p^2), \quad M_3(d; p^2), \quad M_4(d; p^2),$

Deriving DE we find 4 coupled differential equations:

$$\frac{\partial}{\partial \rho^2} \begin{pmatrix} M_1 \\ \dots \\ M_4 \end{pmatrix} = \begin{pmatrix} c_{11} & \dots & c_{14} \\ \dots & \dots & \dots \\ c_{41} & \dots & c_{44} \end{pmatrix} \begin{pmatrix} M_1 \\ \dots \\ M_4 \end{pmatrix} + \text{sub-topologies}$$

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How to solve a system of differential equations?

Different choices of the basis of MIs can simplify DE.

• Try to triangularise the system as $d \rightarrow 4$

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- ► This concept has been generalised by J.Henn's Canonical Basis.
- Find a basis $m_j(d; p^2)$ such that the system takes canonical form

$$\frac{\partial}{\partial p^2} \left(\begin{array}{cc} m_1 \\ \dots \\ m_n \end{array} \right) = (d-4) \left(\begin{array}{ccc} c_{11} & \dots & c_{1n} \\ \dots & \dots & \dots \\ c_{n1} & \dots & c_{nn} \end{array} \right) \left(\begin{array}{ccc} m_1 \\ \dots \\ m_n \end{array} \right)$$

Where now c_{ij} is function ONLY of $p^2!!$

 \rightarrow also now some criteria start to be known! [arXiv:1304.1806; arXiv:1401.2979; arXiv:1404.2922; arXiv:1404.4853]

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 \rightarrow 4 coupled DE in d = 4 and in d = 2 (and in *any even number of dimensions...*).

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$$\begin{split} \widetilde{S} &= \int \mathfrak{D}^d k \mathfrak{D}^d l \frac{1}{D_1 D_2 D_3}, \quad \widetilde{S}_1 = \int \mathfrak{D}^d k \mathfrak{D}^d l \frac{1}{D_1^2 D_2 D_3}, \\ \widetilde{S}_2 &= \int \mathfrak{D}^d k \mathfrak{D}^d l \frac{k \cdot p}{D_1 D_2 D_3}, \quad \widetilde{S}_3 = \int \mathfrak{D}^d k \mathfrak{D}^d l \frac{l \cdot p}{D_1 D_2 D_3} \end{split}$$

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Can we generalise this? \longrightarrow Schouten (Pseudo-)Identities!

- ▶ Valid only in **integer** number of dimensions → *pseudo-identities*!
- in d = n ∈ N dimensions only n vectors can be linearly independent.
- 1. d = 1 dimension: a^{μ}, b^{μ} cannot be independent $\rightarrow a^{\mu}b^{\nu}\epsilon_{\mu\nu} = 0$.
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Squaring relations above we obtain Schouten polynomials:

► For Example take in 3 dimensions

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Squaring it we get

$$\rightarrow P_3(a, b, c) = a^2 b^2 c^2 - a^2 (b \cdot c)^2 - b^2 (a \cdot c)^2 - c^2 (a \cdot b)^2 + 2(a \cdot b)(b \cdot c)(a \cdot c)$$

(Gram determinants of the *n* vectors!)

The polynomial can now be considered as *d* dimensional! (*d* never appears on r.h.s!)

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Sunrise depends on **3** vectors $k, l, p \rightarrow$ Schouten in 2 dimensions!!

\rightarrow "natural" number of dimensions for the Sunrise is 2!

The Sunrise satisfies a **second-order differential equation** in *d* = 2 !! [arXiv:1112.4360] by S.Müller-Stach, S.Weinzierl, R.Zayadeh [arXiv:1302.7004] by L.Adams, C.Bogner, S.Weinzierl

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Note that as $k, l \to \infty$ then $P_3(d; k, l, p) \approx k^2 l^2$

Consider now:

$$Z(d; n_1, n_2, n_3) = \int \mathfrak{D}^d k \mathfrak{D}^d l \frac{P_3(d; k, l, p)}{D_1^{n_1} D_2^{n_2} D_3^{n_3}}$$

Power-counting in d = 2

 $Z(d \rightarrow 2; 1, 1, 1)$ UV div, $Z(d \rightarrow 2; 2, 1, 1)$ UV div, (with permutations) $Z(d \rightarrow 2; 2, 2, 1)$, $Z(d \rightarrow 2; 2, 1, 2)$, $Z(d \rightarrow 2; 1, 2, 2)$, UV finite!

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For all Zs that are <u>finite</u> we then also have at once:

 $Z(d \rightarrow 2; n_1, n_2, n_3) \rightarrow 0$.

We can the reduce them to the 4 MIs

$$Z(d; n_1, n_2, n_3) = \sum_{i=0}^{3} C_i(d; p^2) S_i(d; p^2) + \text{sub-topologies}$$

• This gives in principle a **relation** among the MIs in d = 2!

$$0 = \sum_{i=0}^{3} C_i(d \rightarrow 2; p^2) S_i(d \rightarrow 2; p^2) + \text{sub-topologies}\Big|_{d \rightarrow 2}$$

Question: There are infinite converging Zs! Which ones give non-trivial information (if any...)?

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$$0 = \sum_{i=0}^{3} C_i(d \rightarrow 2; p^2) S_i(d \rightarrow 2; p^2) + \text{sub-topologies}\Big|_{d \rightarrow 2}$$

Question: There are infinite converging Zs! Which ones give <u>non-trivial</u> information (if any...)?

Apparently only those with the minimal powers required for convergence!

For all Zs that are <u>finite</u> we then also have at once:

 $Z(d \rightarrow 2; n_1, n_2, n_3) \rightarrow 0$.

We can the reduce them to the 4 MIs

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Some examples of Zs:

• If the powers are too high **no information** as $d \rightarrow 2$

$$Z(d; 2, 2, 2) = -\frac{(d-1)(d-2)}{4} \times \left[(d-3)S(d; p^2) + m_1^2 S_1(d; p^2) + m_2^2 S_2(d; p^2) + m_3^2 S_3(d; p^2) \right]$$

► On the contrary:

$$\begin{aligned} &Z_2(d;p^2) = Z(d;2,1,2,p^2) \\ &= \frac{(d-1)}{12} \left[-(d-2)p^2 + (d-3)(m_1^2 - 2m_2^2 + m_3^2) \right] S(d;p^2) \\ &+ \frac{(d-1)}{12} (p^2 + m_1^2 - 3m_2^2 + 3m_3^2) \ m_1^2 S_1(d,p^2) \\ &- \frac{(d-1)}{6} (p^2 + m_2^2) \ m_2^2 S_2(d;p^2) \\ &+ \frac{(d-1)}{12} (p^2 + 3m_1^2 - 3m_2^2 + m_3^2) \ m_3^2 S_3(d;p^2) \\ &+ \frac{(d-1)(d-2)}{24} \left[T(d;m_1,m_2) - 2 \ T(d;m_1,m_3) + T(d;m_2,m_3) \right] \end{aligned}$$

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$$+ \frac{(d-1)}{12} (p^{2} + m_{1}^{2} - 3m_{2}^{2} + 3m_{3}^{2}) m_{1}^{2}S_{1}(d, p^{2})$$

$$- \frac{(d-1)}{6} (p^{2} + m_{2}^{2}) m_{2}^{2}S_{2}(d; p^{2})$$

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$$+ \frac{(d-1)(d-2)}{24} \left[T(d; m_{1}, m_{2}) - 2 T(d; m_{1}, m_{3}) + T(d; m_{2}, m_{3}) \right]$$

Which in d = 2 becomes:

$$Z_{2}(2; p^{2}) = 0 = -\frac{1}{12}(m_{1}^{2} - 2m_{2}^{2} + m_{3}^{2})S(2; p^{2})$$

$$+ \frac{1}{12}(p^{2} + m_{1}^{2} - 3m_{2}^{2} + 3m_{3}^{2}) m_{1}^{2}S_{1}(2, p^{2})$$

$$- \frac{1}{6}(p^{2} + m_{2}^{2}) m_{2}^{2}S_{2}(2; p^{2})$$

$$+ \frac{1}{12}(p^{2} + 3m_{1}^{2} - 3m_{2}^{2} + m_{3}^{2}) m_{3}^{2}S_{3}(2; p^{2})$$

$$+ \frac{1}{96} \ln \frac{m_{2}^{2}}{m_{1}m_{3}}$$

In the very same way we can obtain similar relations for:

$$Z_1(d; p^2) = Z(d; 1, 2, 2), \qquad Z_3(d; p^2) = Z(d; 2, 2, 1),$$

• One **two** of them are really independent in d = 2:

$$Z_1(d;p^2) + Z_2(d;p^2) + Z_3(d;p^2) = -rac{(d-1)(d-2)}{4}p^2 \ S(d;p^2) \ ,$$

▶ We can choose 2 of them as **new MIs**! For example choose as basis:

$$S(d; p^2), \quad S_1(d; p^2), \quad Z_2(d; p^2), \quad Z_3(d; p^2).$$

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Differential equations in the new basis are in **block form** in d = 2

$$\frac{d}{d p^2} \begin{pmatrix} S \\ S_1 \end{pmatrix} = \begin{pmatrix} c_{11} & c_{12} & c_{13} & c_{14} \\ c_{21} & c_{22} & c_{23} & c_{24} \end{pmatrix} \begin{pmatrix} S \\ S_1 \\ Z_2 \\ Z_3 \end{pmatrix} + \text{sub-topologies}$$
$$\frac{d}{d p^2} \begin{pmatrix} Z_2 \\ Z_3 \end{pmatrix} = (d-2) \begin{pmatrix} c_{31} & c_{32} & c_{33} & c_{34} \\ c_{41} & c_{42} & c_{43} & c_{44} \end{pmatrix} \begin{pmatrix} S \\ S_1 \\ Z_2 \\ Z_3 \end{pmatrix} + \text{sub-topologies}$$

The second system **decouples** completely in d = 2!

 \rightarrow order by order in (d-2) we can derive a second order differential equation satisfied by $S(d; p^2)$. [see Weinzierl et al., 2011/2013]

▶ Let us express Z_1 , Z_2 , Z_3 in terms of the \widetilde{S}_j !

$$Z_1 = - \frac{(d-1)(d-2)}{4} \widetilde{S}_2$$

$$Z_2 = + \frac{(d-1)(d-2)}{4} \widetilde{S}_3$$

$$Z_3 = -\frac{(d-1)(d-2)}{4} \left[p^2 \widetilde{S}_0 - \widetilde{S}_2 + \widetilde{S}_3 \right].$$

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Conclusions and outlook

- Schouten identities are a new class of (pseudo-)identities among MIs!
- ▶ Schouten identities can be used to select a "good" basis of MIs as $d \approx n \in \mathbb{N}$
- They show that in d = 2 the Sunrise has only 2 MIs (the other two are linearly dependent!)

 \rightarrow **Decoupling** of the system in <u>*d*</u> = <u>*n*</u> seems to be due to "hidden relations" between the MIs in <u>*d*</u> = <u>*n*</u> ???!!!

- 1. Can this be used in more involved cases?
- 2. What happens with IR divergences?

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Thanks!

Back-up slide

What about d = 4?

- Schouten Polynomial has also the role of shifting the dimensions $d \rightarrow d + 2$ [R. Lee]
- Implies that:

$$Z_1(d; p^2) = \int \mathfrak{D}^d k \mathfrak{D}^d l \frac{P_3(d; k, l, p)}{D_1^1 D_2^2 D_3^2} \approx \int \frac{\mathfrak{D}^{d+2} k \mathfrak{D}^{d+2} l}{D_1^1 D_2^2 D_3^2}$$

and the same way for the others...

Which implies that

$$\int rac{\mathfrak{D}^d k \, \mathfrak{D}^d l}{D_1^1 \, D_2^2 \, D_3^2}$$
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