



Results from the LHC Run I and perspectives for Run II

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CERN

OUTLINE

- The LHC at CERN
- Highlights from Run I
 - The discovery of the Higgs boson
 - Searches for new heavy resonances
 - Searches for Dark Matter in cascade (a.k.a. SUSY)
 - Searches for Dark Matter direct production (a.k.a. monojet, mono...)
- Perspective for Run II
- Conclusions

THE LARGE HADRON COLLIDER

27 Km tunnel
filled with
superconductive 8.3T
magnets kept @ ~3 K
Designed for 14 TeV
pp collisions
So far operated at
7TeV and 8 TeV



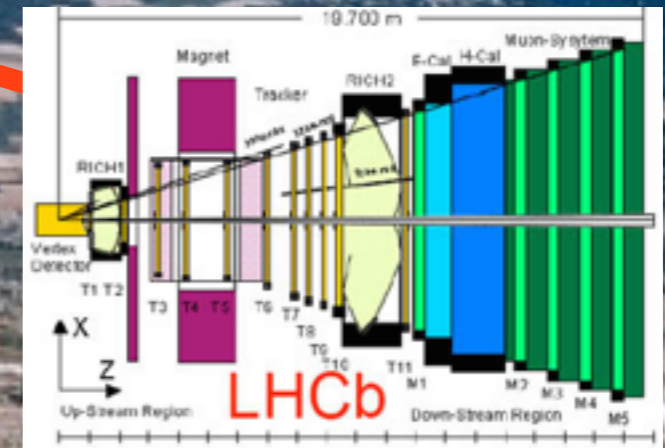
Switzerland

THE LHC EXPERIMENTS

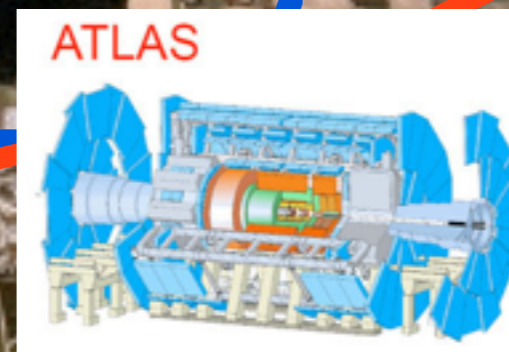
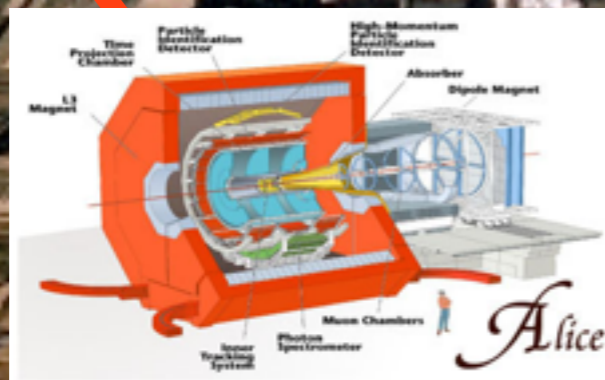
Multipurpose
(high-pT, HI, b physics)



dedicated to b physics



Dedicated to
Heavy Ions



Multipurpose
(high-pT, HI, b physics)

ATLAS & CMS PHYSICS GOAL

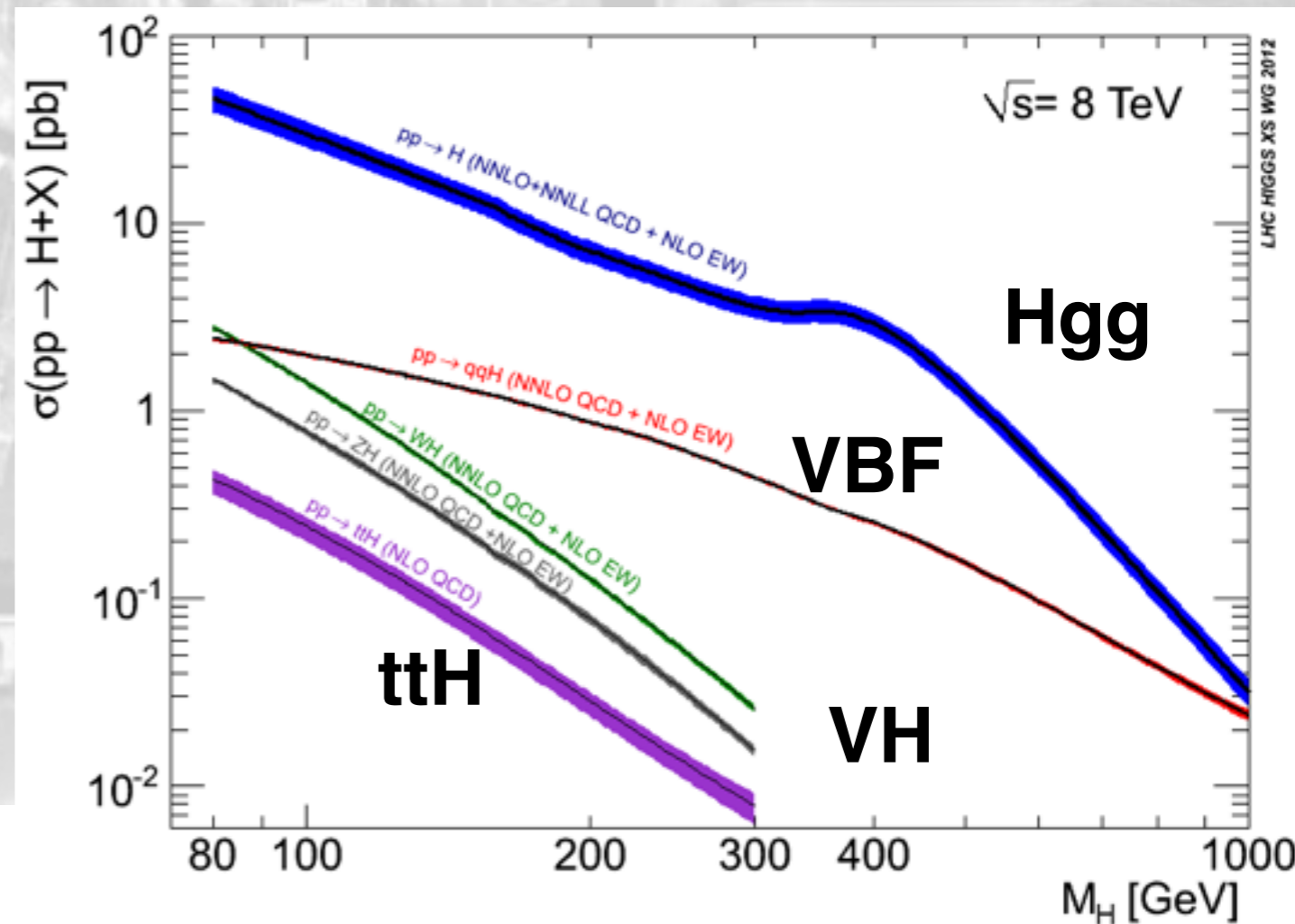
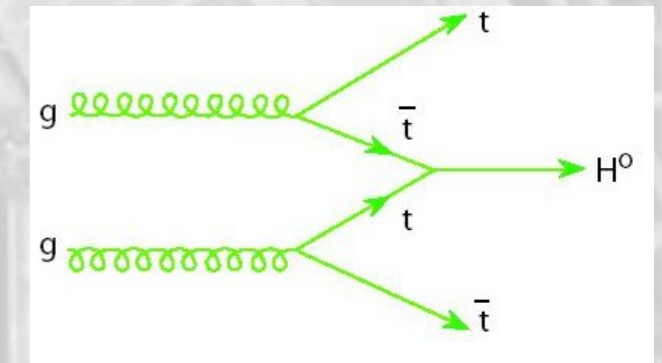
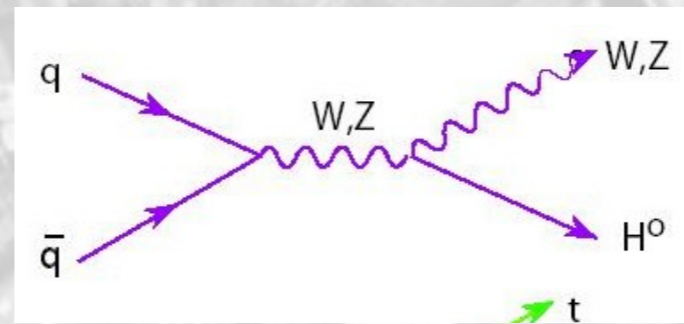
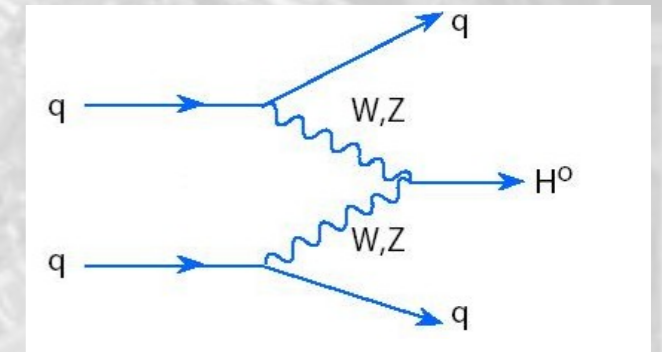
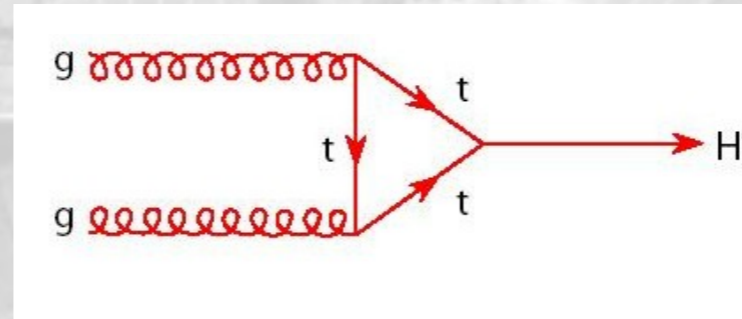
- Search for the Higgs boson
- Fully characterise EW symmetry breaking
- Explore the TeV scale
- Test SM with precision measurements (perturbative QCD, parton density functions, ...)
- Improve precision on SM parameters (e.g., masses of W, Z, and top)
- ...



Highlights from Run I: discovery of the Higgs boson

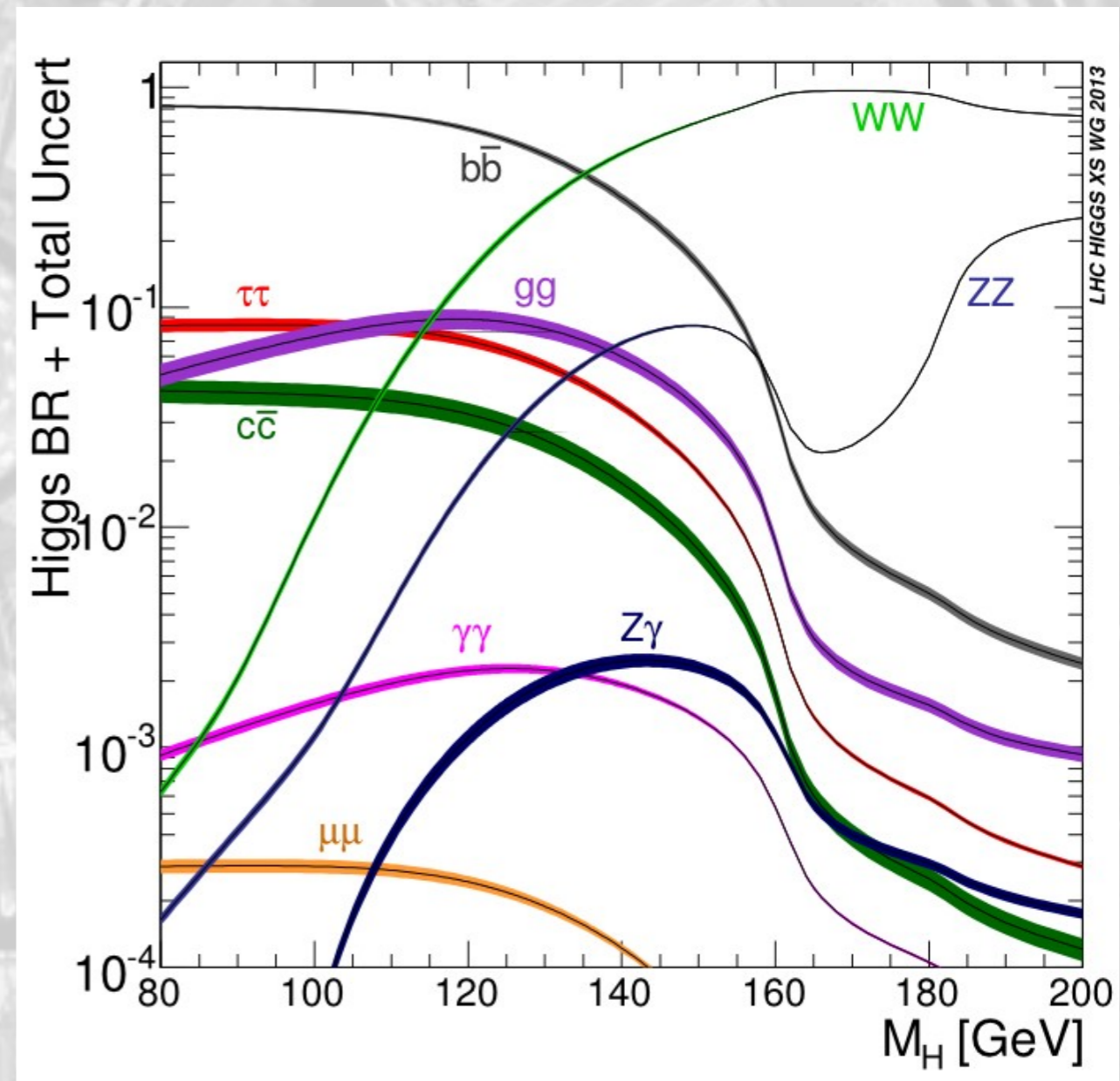
HIGGS BOSON PRODUCTION

- There are mainly four production mechanisms
 - gluon-gluon fusion (Hgg)
 - vector-boson fusion (VBF)
 - in association with vector boson (VH)
 - top-top fusion (ttH)
- While gg is dominant, all the mechanisms have been considered
 - redundancy
 - favourable S vs B for problematic channels

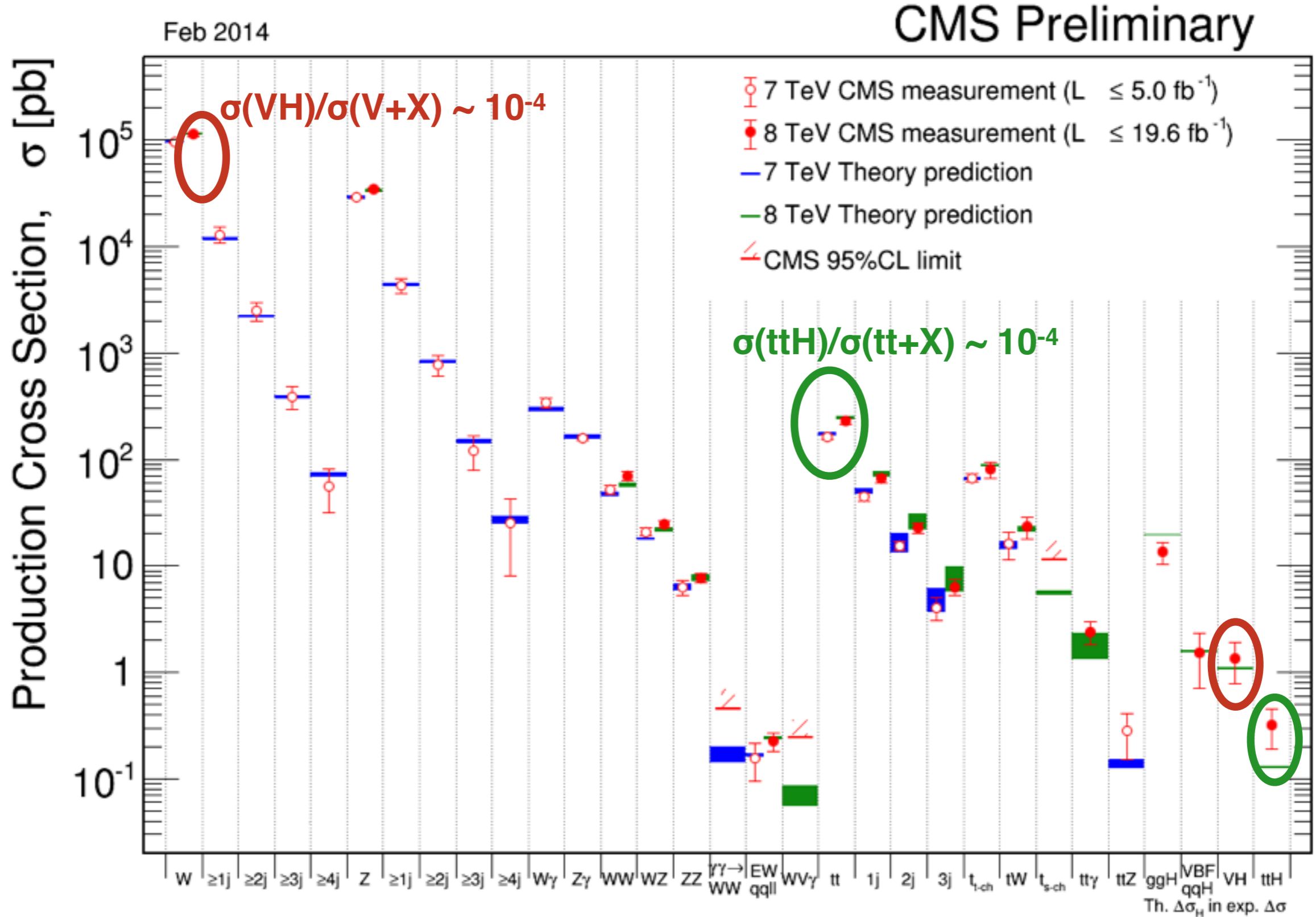


HIGGS BOSON DECAY

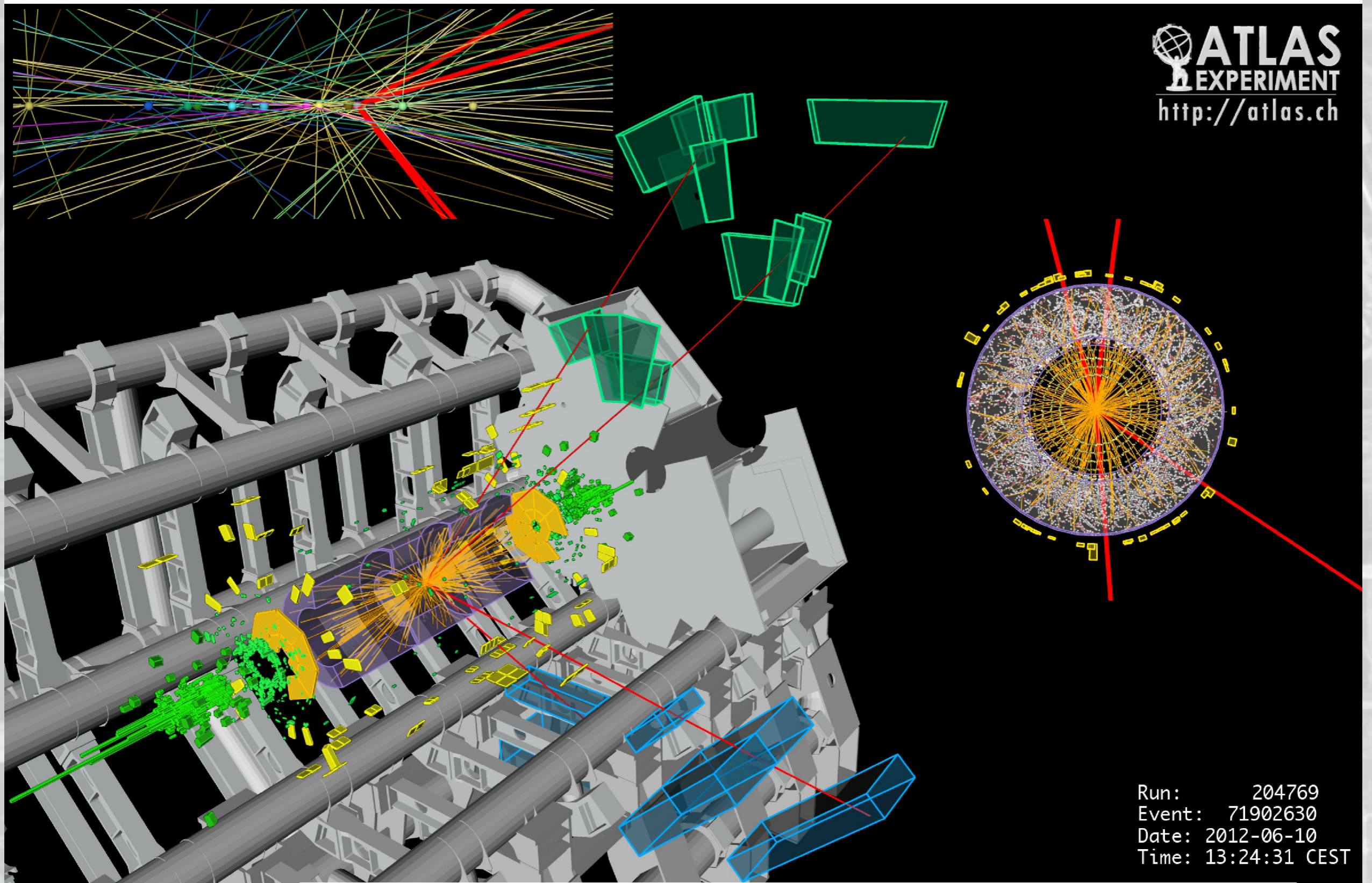
- The Higgs phenomenology is driven by the allowed decay modes
- Depending on the Higgs mass, this can result in a more or less rich set of possibilities
- Particularly for $m_H \sim 130$ GeV there are several possibilities (i.e., we ended up in a lucky spot)
 - Four leptons (the golden channel)
 - Diphoton
 - WW
 - $\tau\tau$ & bb : large background but needed to prove coupling to fermions



BACKGROUND IS CHALLENGING



$$H \rightarrow ZZ^{(*)} \rightarrow 4\ell$$



$$H \rightarrow ZZ^{(*)} \rightarrow 4\ell$$

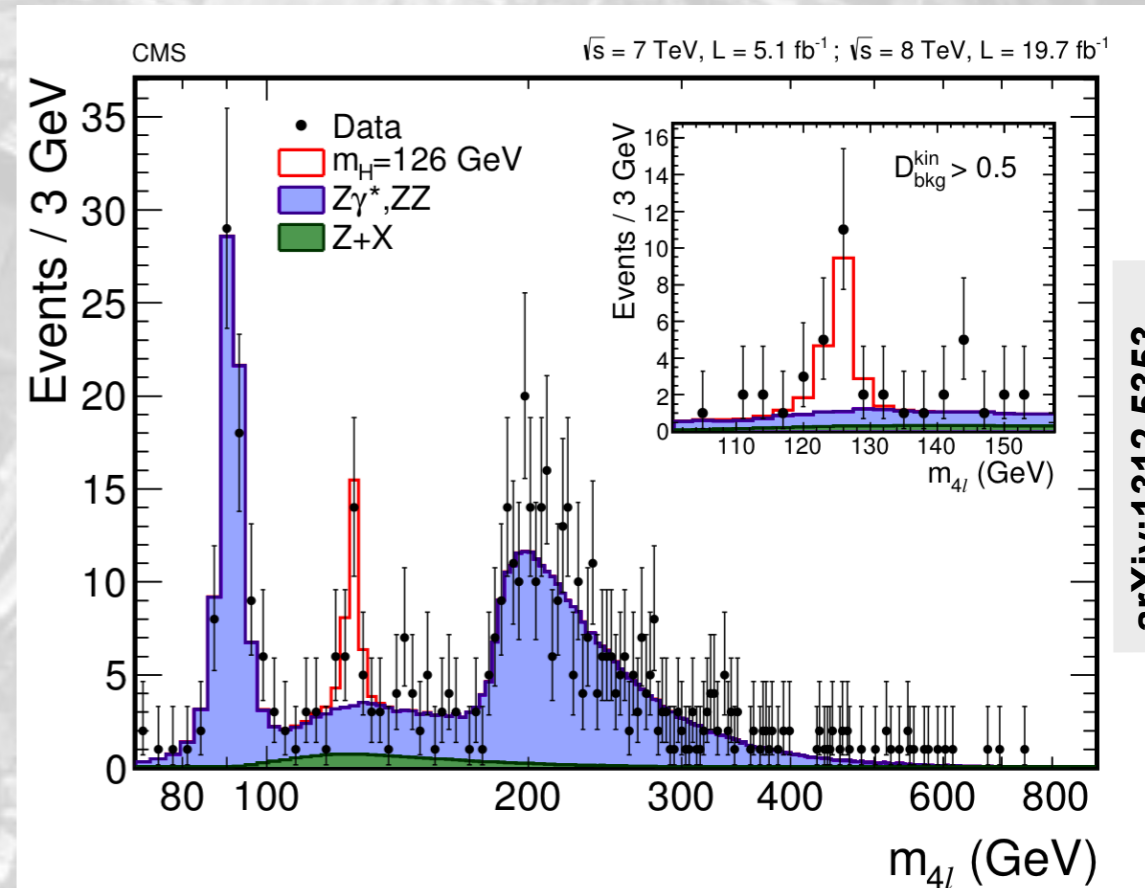
Require 4ℓ ($\ell=e,\mu$) isolated (not inside a jet)

For bkg, this is expensive to produce (need W or Z)

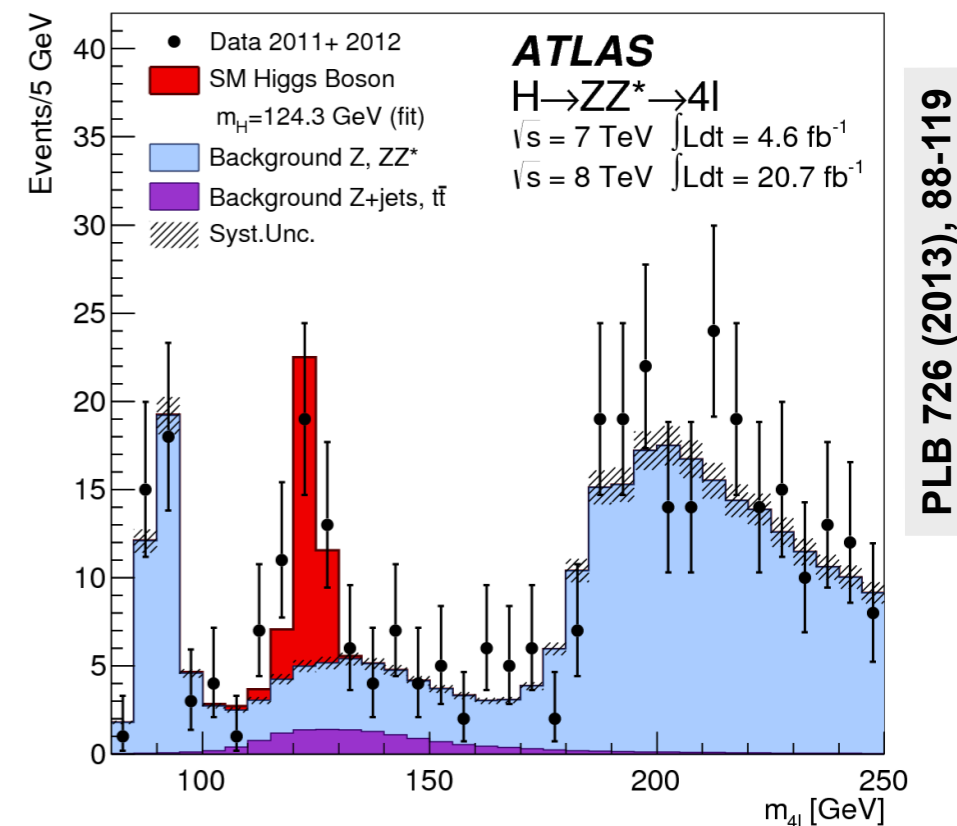
For signal, there is a (smaller) suppression due to $Z \rightarrow \ell\ell$ branching ratio

This is a signal difficult to miss

If we did not know about the Higgs boson, we would have found it in this channel (but most likely with a less-tuned analysis, i.e. we did not have a discovery by now, w/o the underlying model)

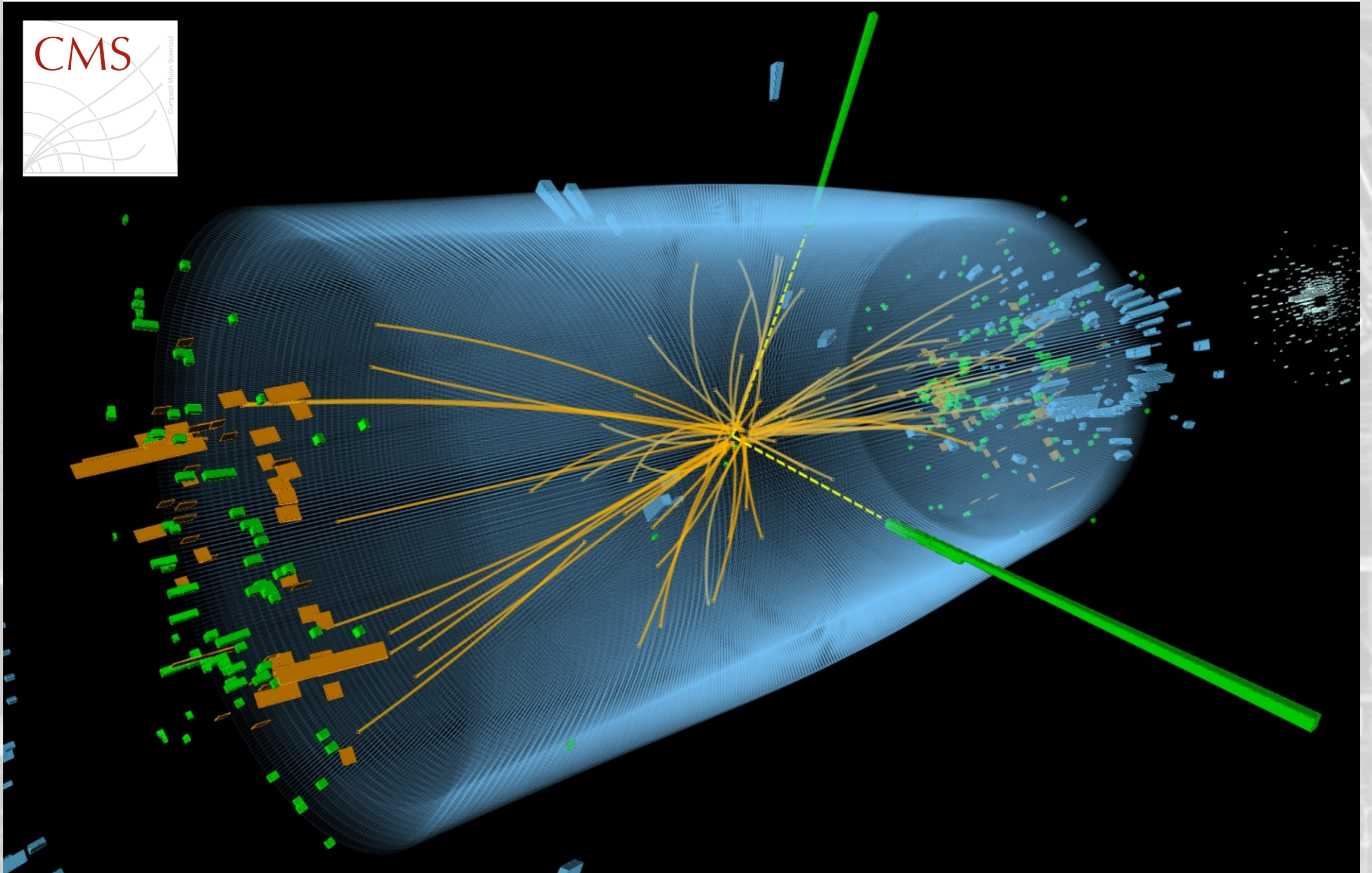


arXiv:1312.5353



PLB 726 (2013), 88-119

$$H \rightarrow \gamma\gamma$$

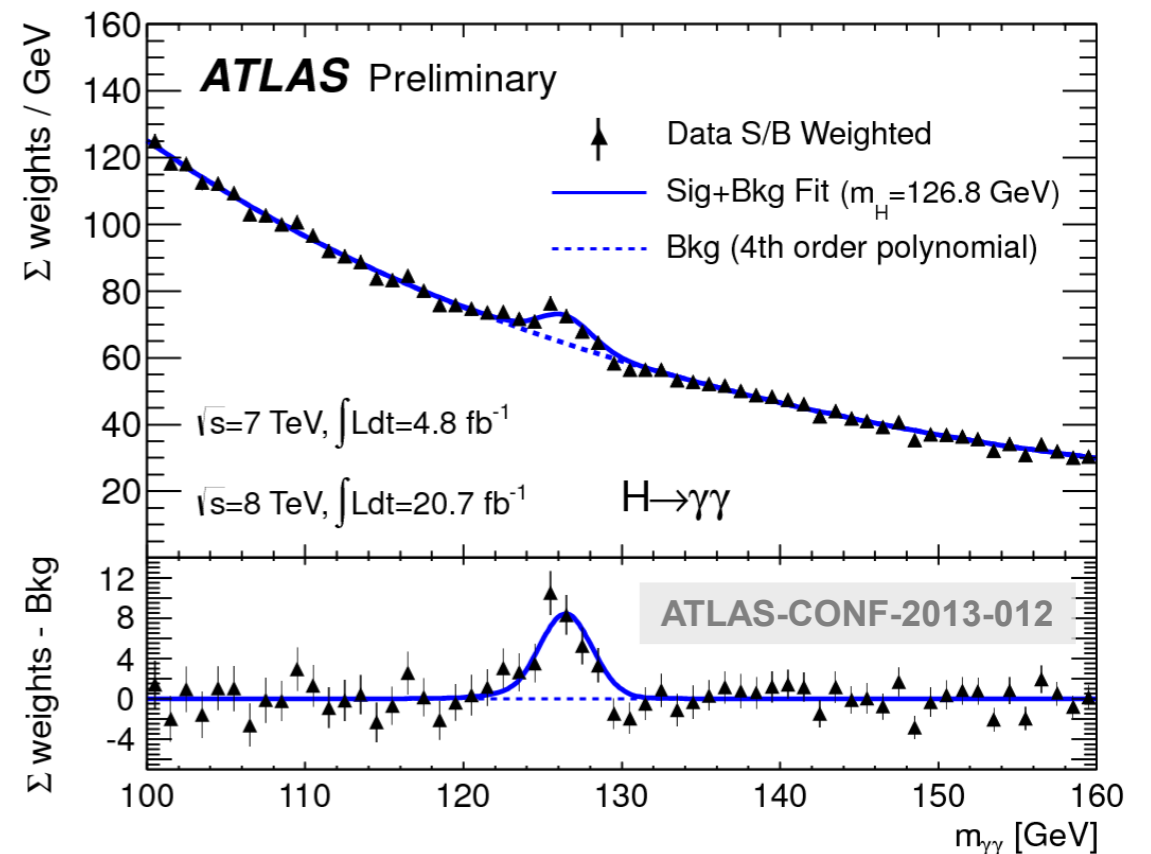
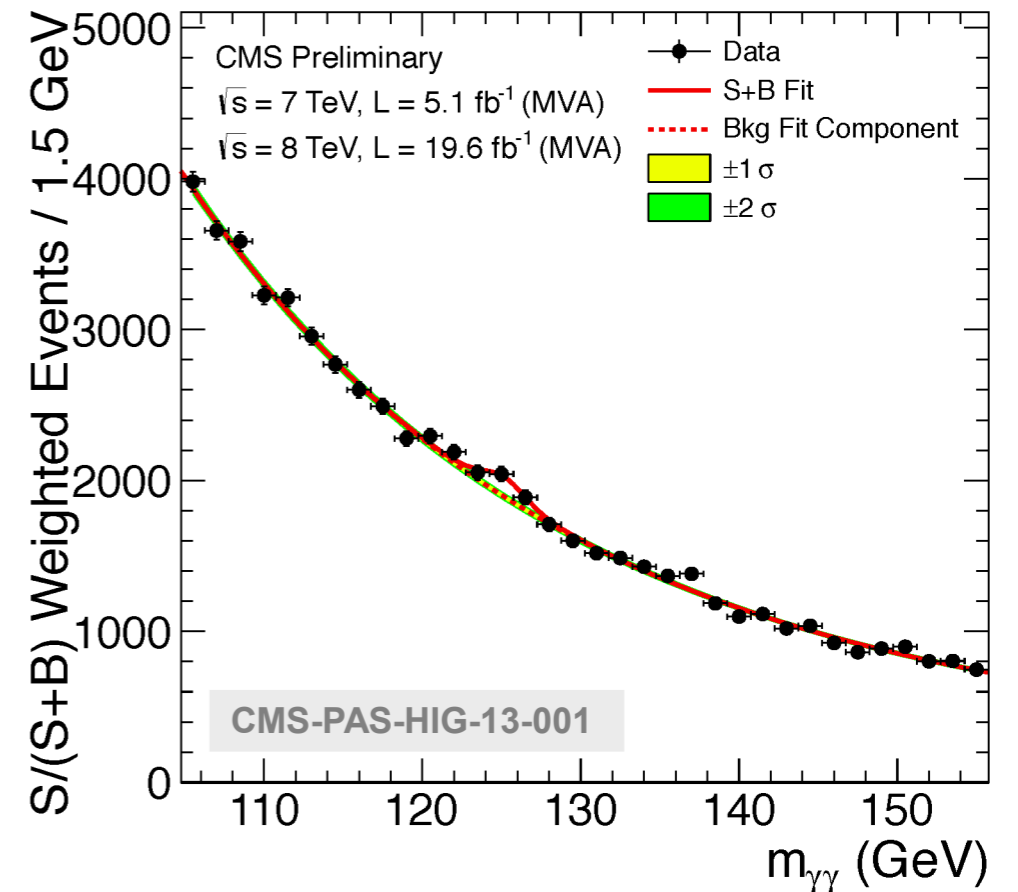


$H \rightarrow \gamma\gamma$

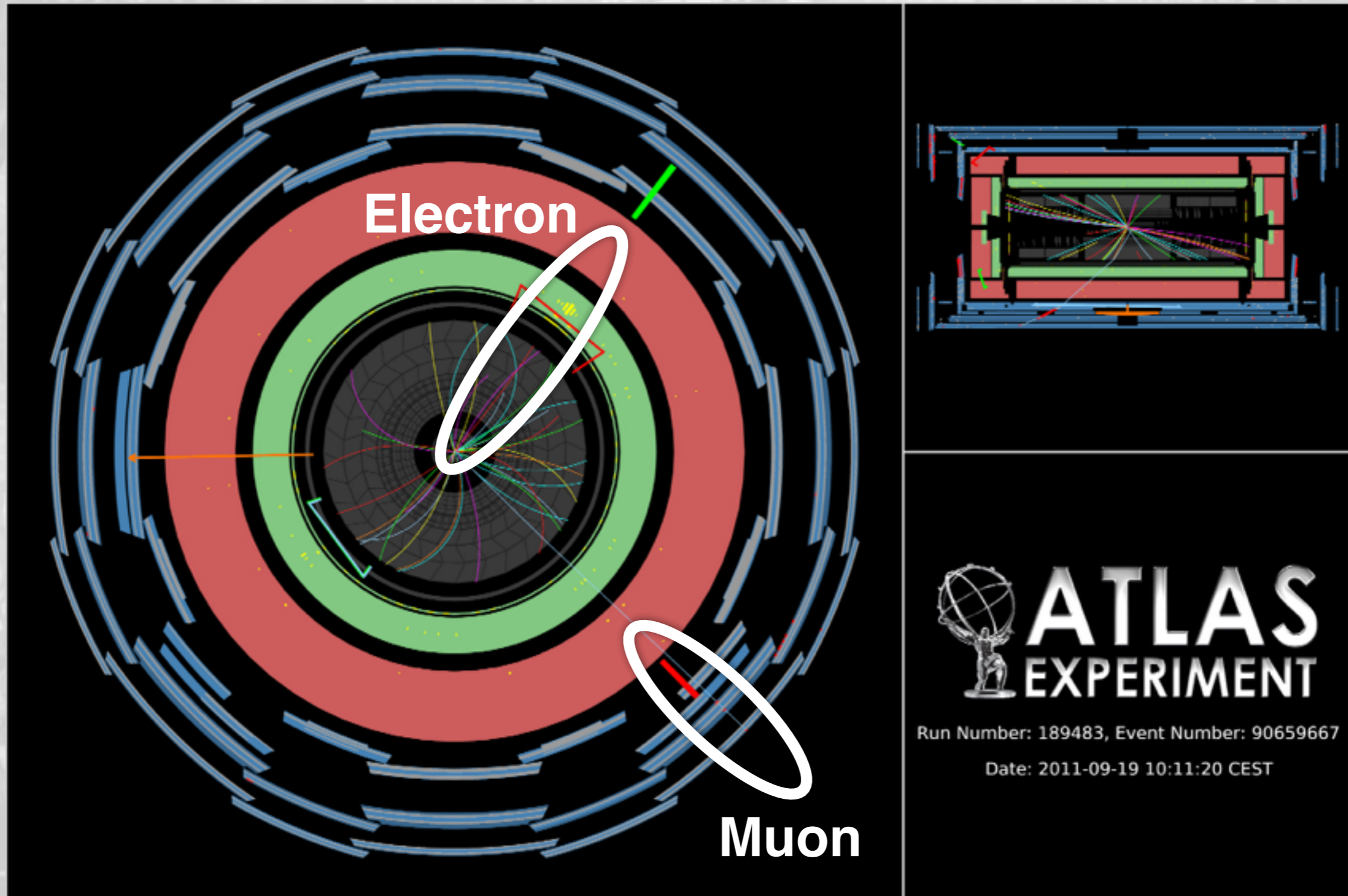
- Select events with two photons
- Classify events according to the “goodness” of the photon reconstruction (high purity vs. high efficiency)
- Perform di-photon invariant mass fit in each category

- signal hypothesis from signal MC
- bkg from analytic function (tested on MC and QCD data control samples)

Combine the categories

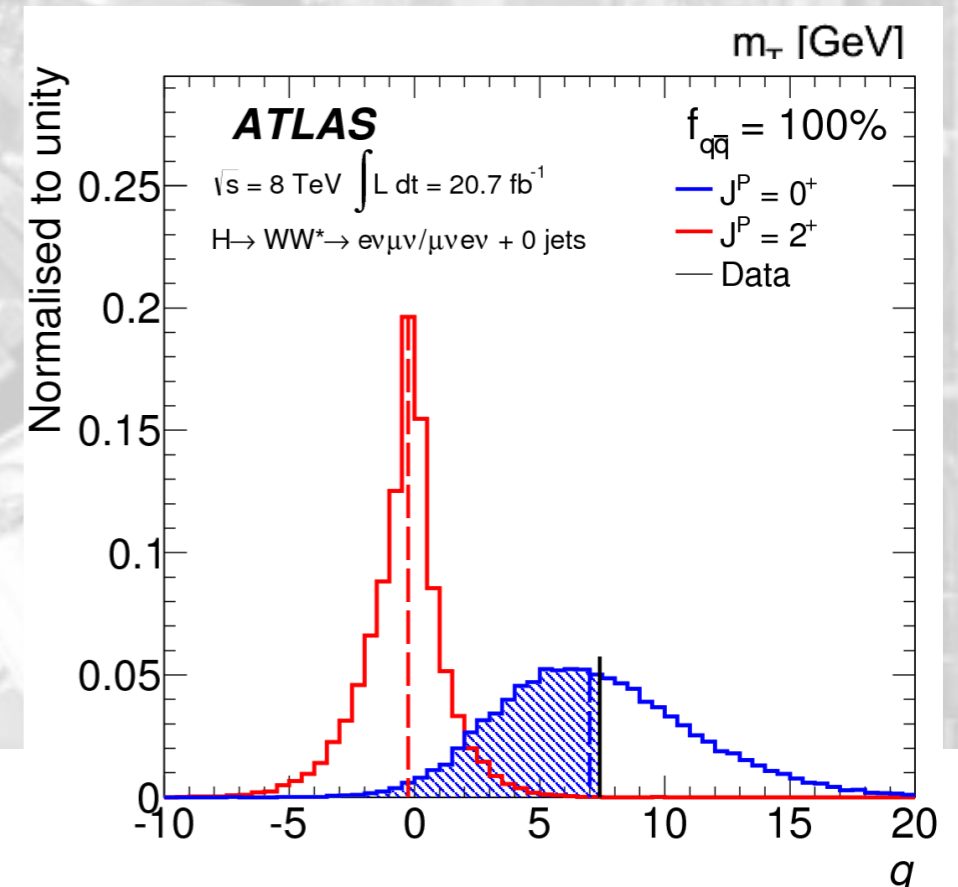
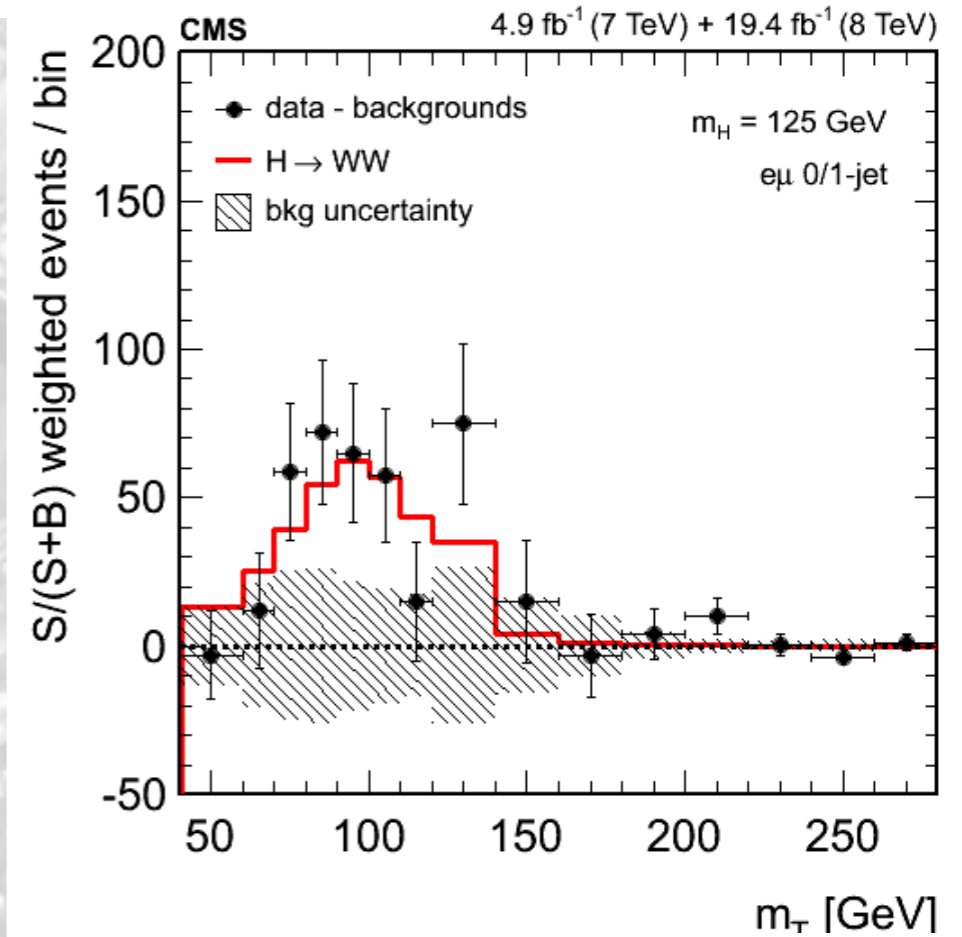


$$H \rightarrow WW^{(*)} \rightarrow 2\ell 2V$$



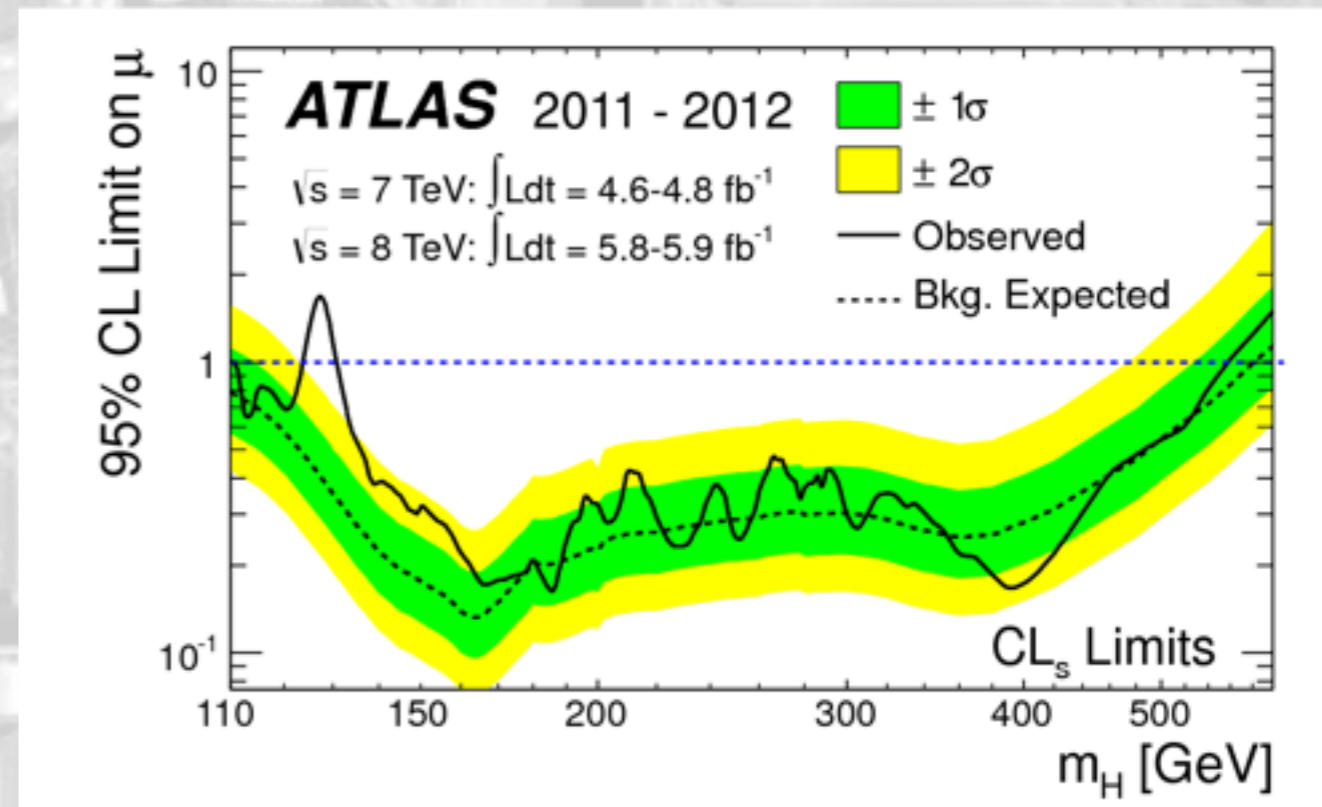
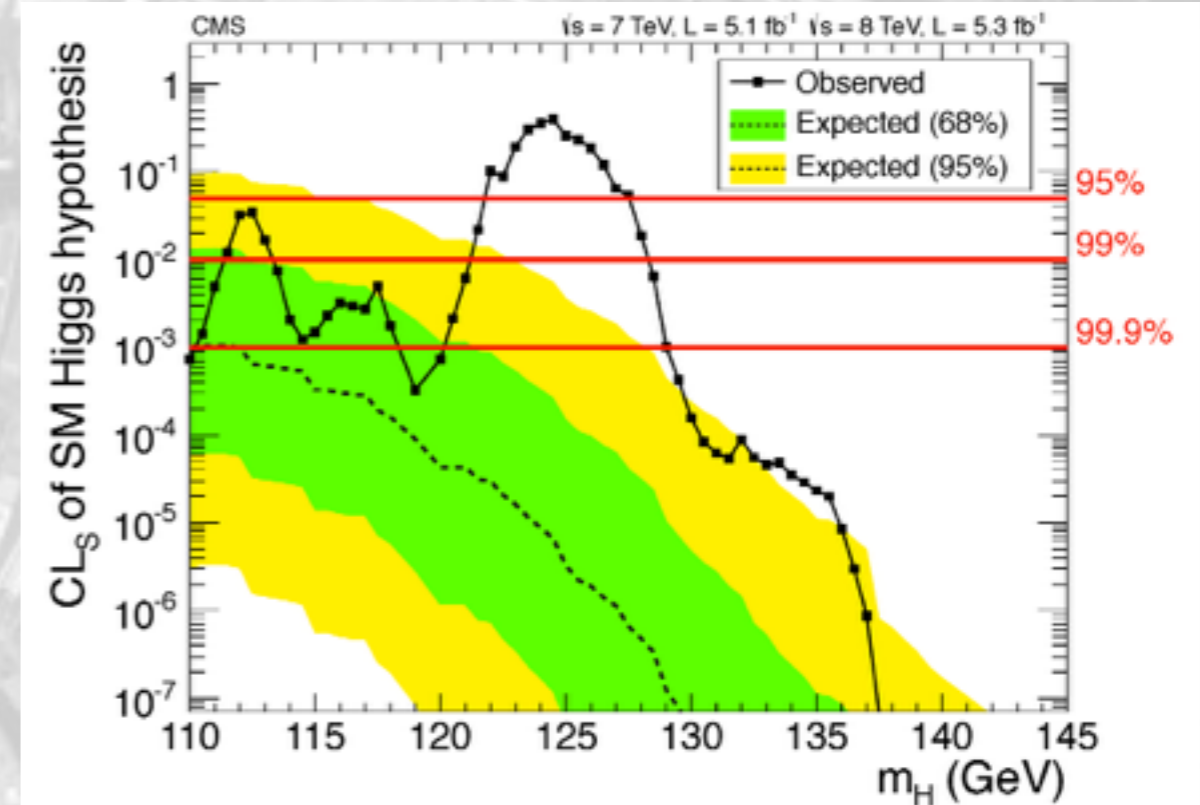
$H \rightarrow WW^{(*)} \rightarrow 2\ell 2V$

- Select events with 2ℓ ($\ell=e,\mu$) and missing transverse energy (unobserved neutrinos should balance observed ℓ s)
- Open kinematic in the final state
- Conservation of transverse momentum allows to guess the Higgs mass (with worse resolution)
- We can “see” a signal, but we cannot measure the resonance mass with good accuracy
- Or we can use Machine Learning to



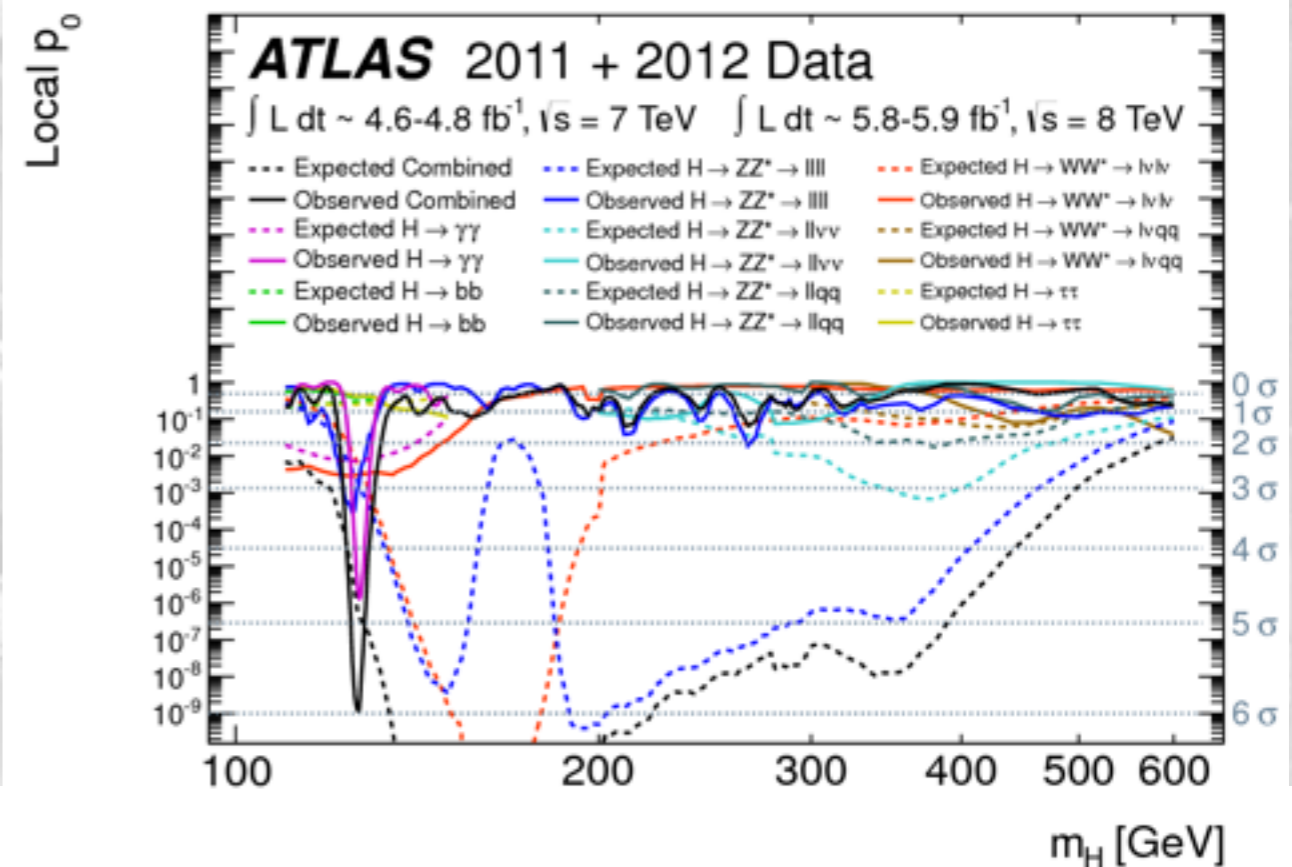
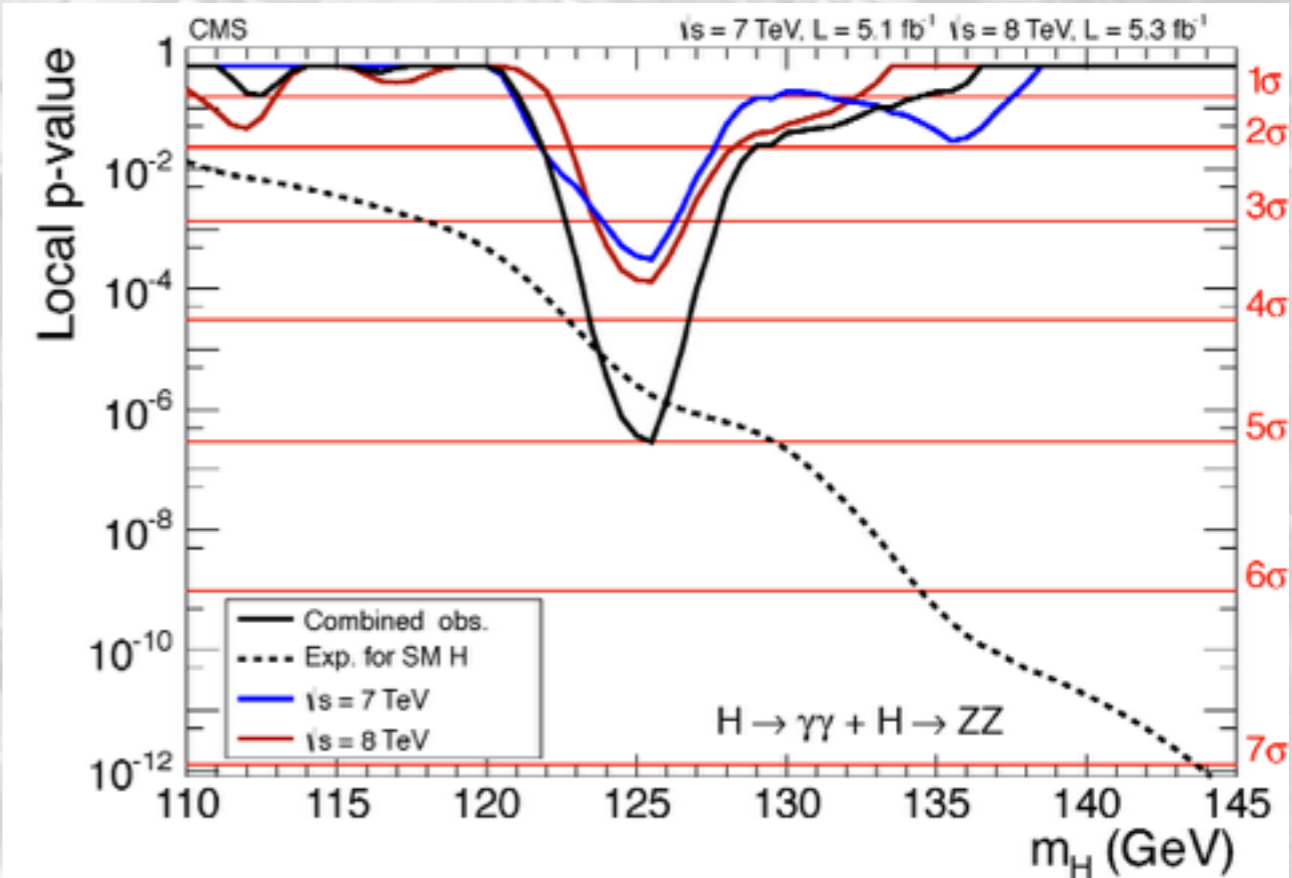
TRYING TO EXCLUDE A SIGNAL

- When you don't know if you have a signal, you first try to exclude it
- If the signal is there, your limit will be poor (and worse than expectation)
- If it is much worse, you might have discovered a signal...
- ... or you might have discovered that your analysis is terrible
- these plots are not the right plots to establish the presence of a signal



ESTABLISHING A DISCOVERY

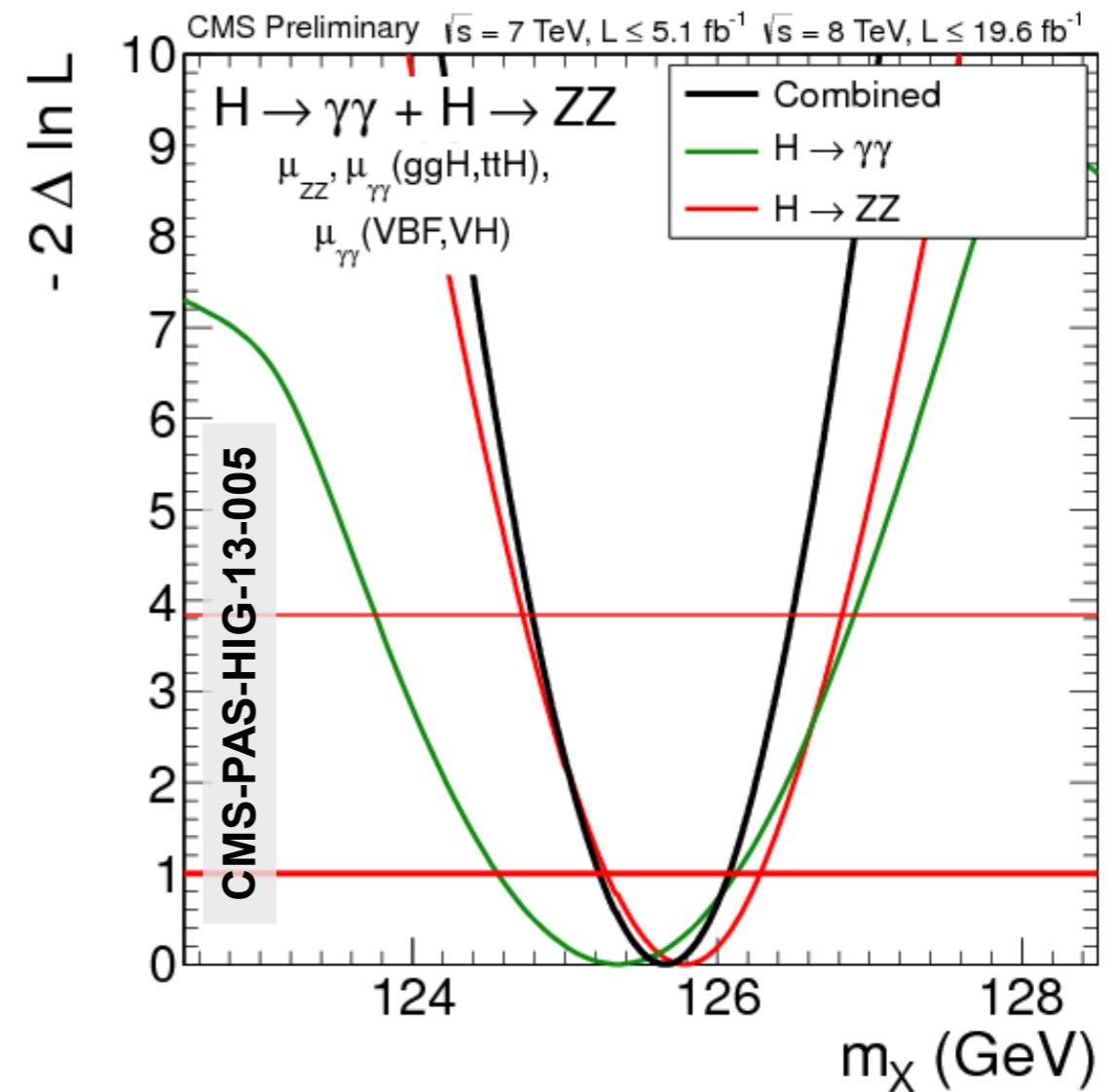
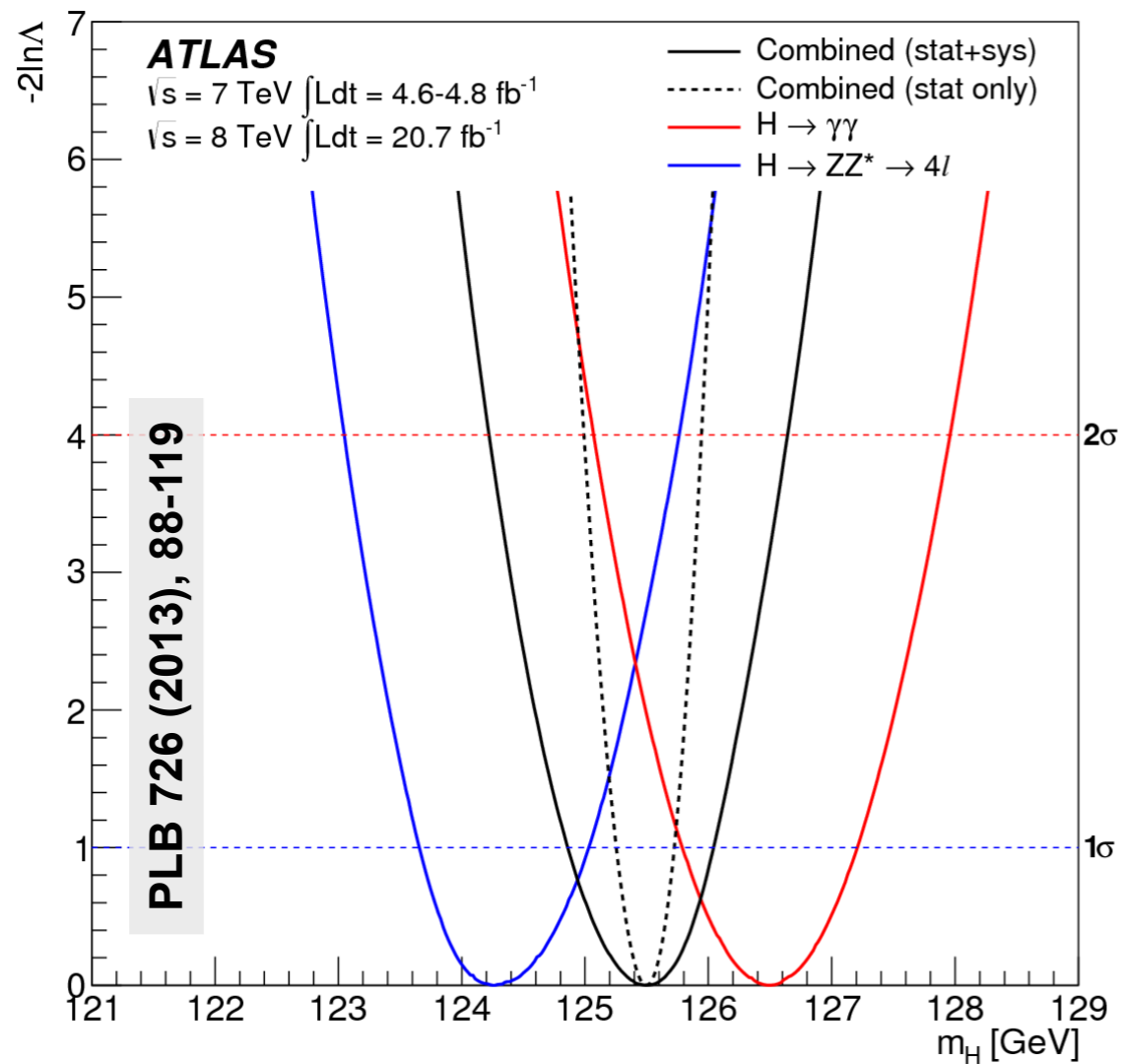
- To claim a discovery, you need to exclude the possibility that your background could mimic a signal
- To do so, you measure (with toy experiments) the probability that a bkg-only sample gives a result as signal-like as what you see on data
- The signal is stringer than the conventional 5σ threshold so...



ESTABLISHING A DISCOVERY



THE HIGGS BOSON MASS



ATLAS

125.5 \pm 0.2 (stat) $^{+0.5}_{-0.6}$ (syst) GeV

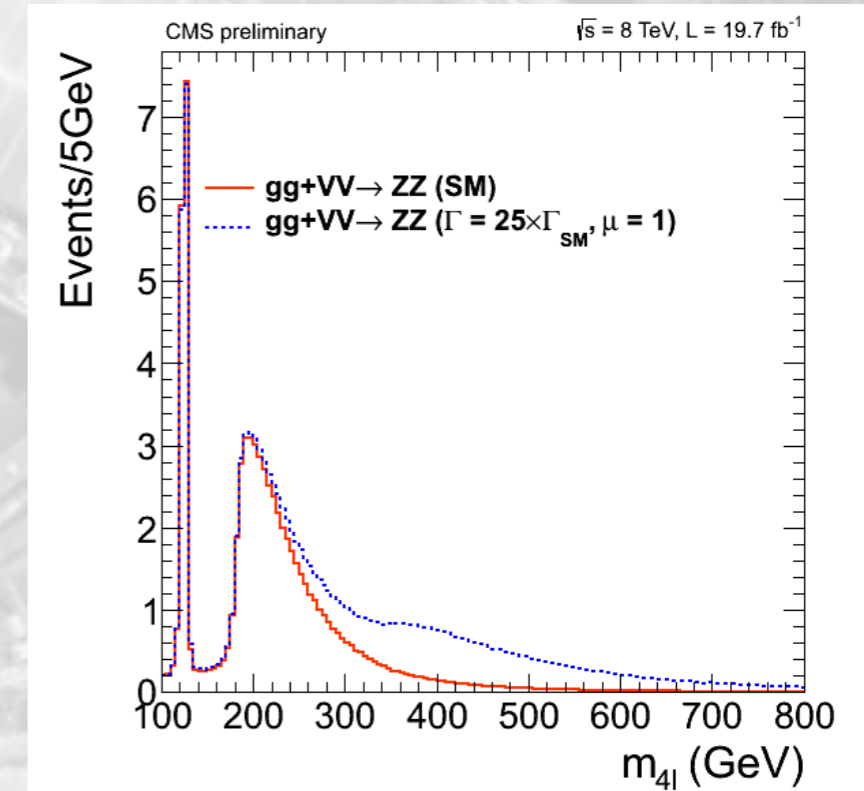
CMS (new ZZ(4l) not used)

125.7 \pm 0.3 (stat) \pm 0.3 (syst) GeV

THE HIGGS BOSON WIDTH

- The Higgs width (~ 4 MeV) was considered too small for the LHC to measure it (expected precision ~ 5 GeV)

	$4l$	$2l2\nu$	Combined
Expected 95% CL limit, r	11.5	10.7	8.5
Observed 95% CL limit, r	6.6	6.4	4.2
Observed 95% CL limit, Γ_H (MeV)	27.4	26.6	17.4

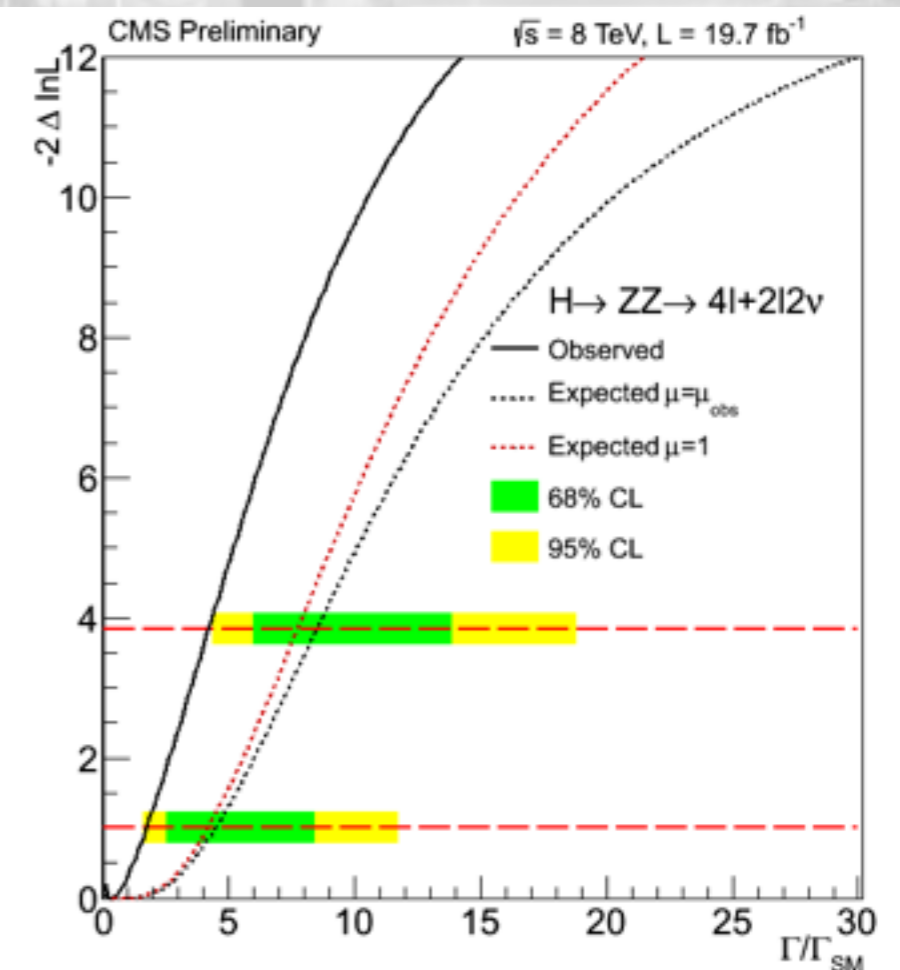


- It was pointed out how to exploit interference effects with ZZ events

Caola & Melnikov

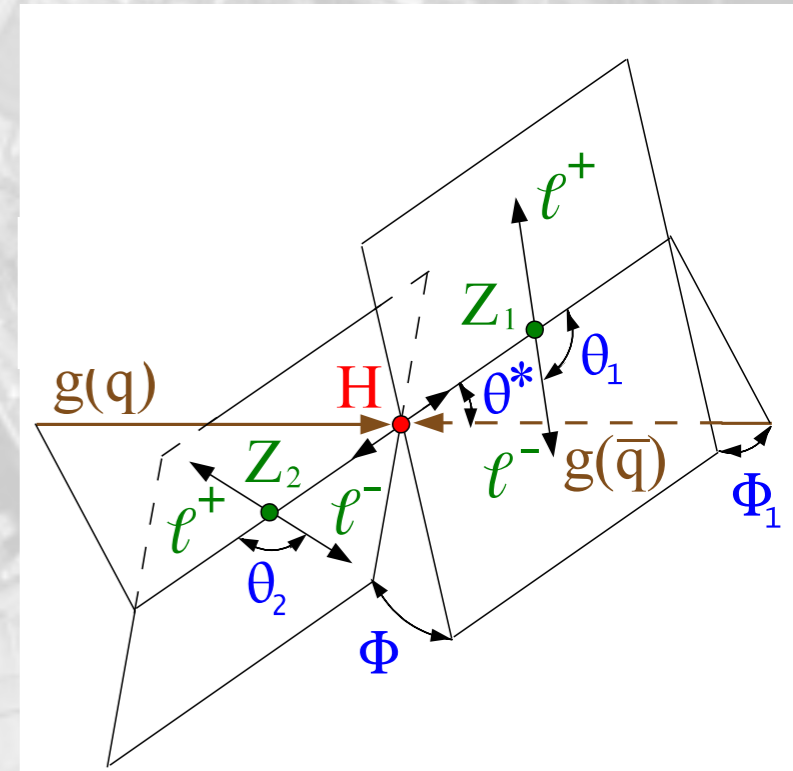
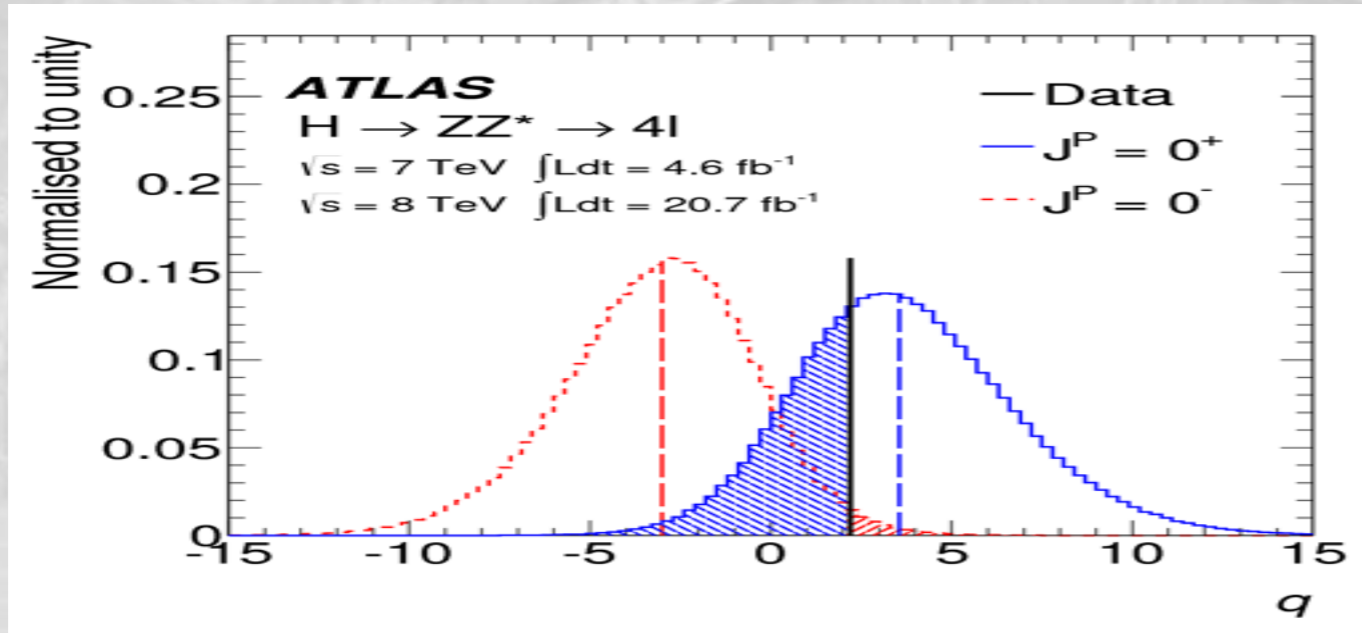
<http://arxiv.org/abs/1307.4935>

- Limit obtained ~ 20 MeV (not there yet, but not that far)

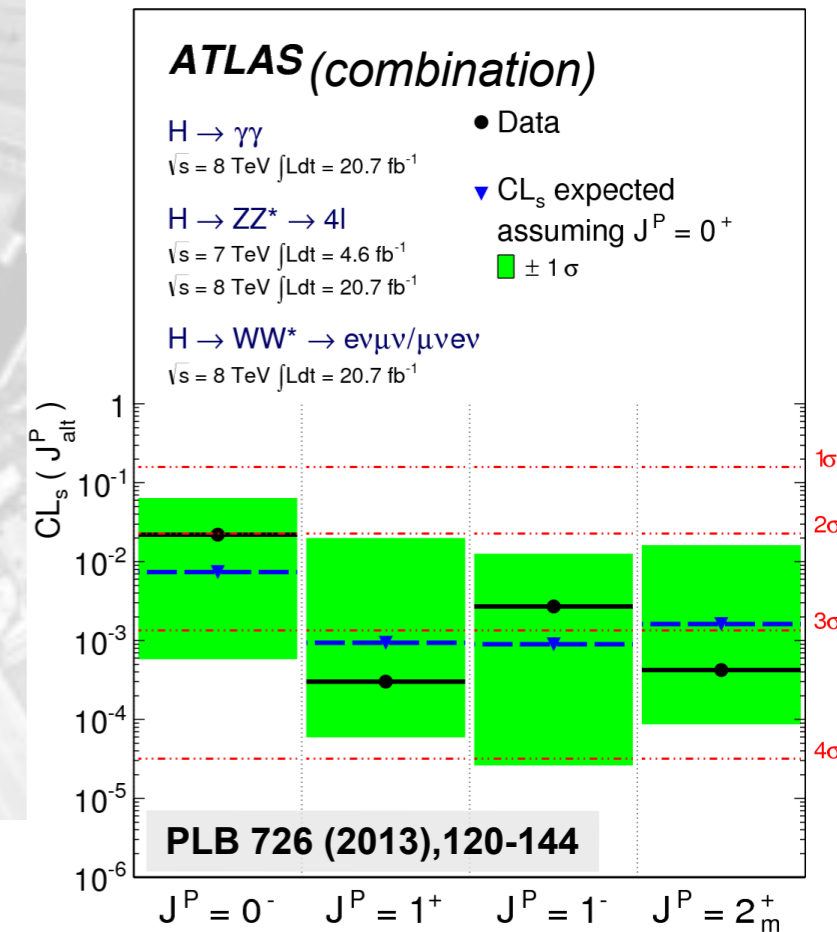
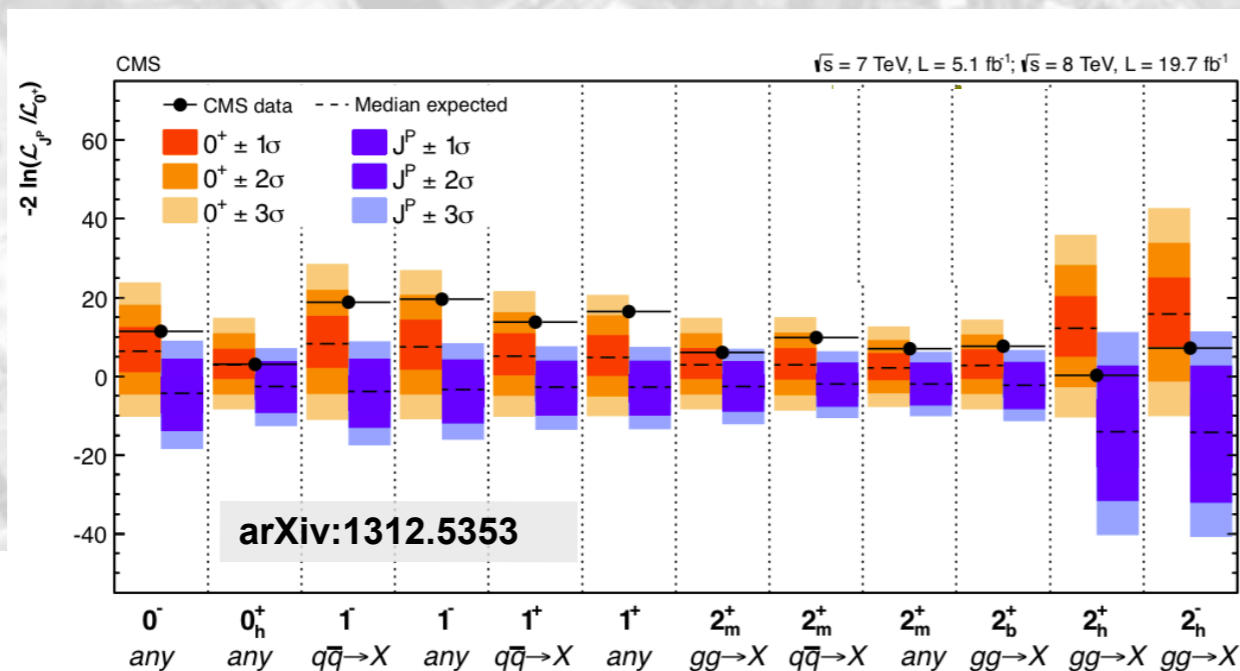


HIGGS PROPERTIES

- Quantum numbers can be measured from angular distribution of the H decay products

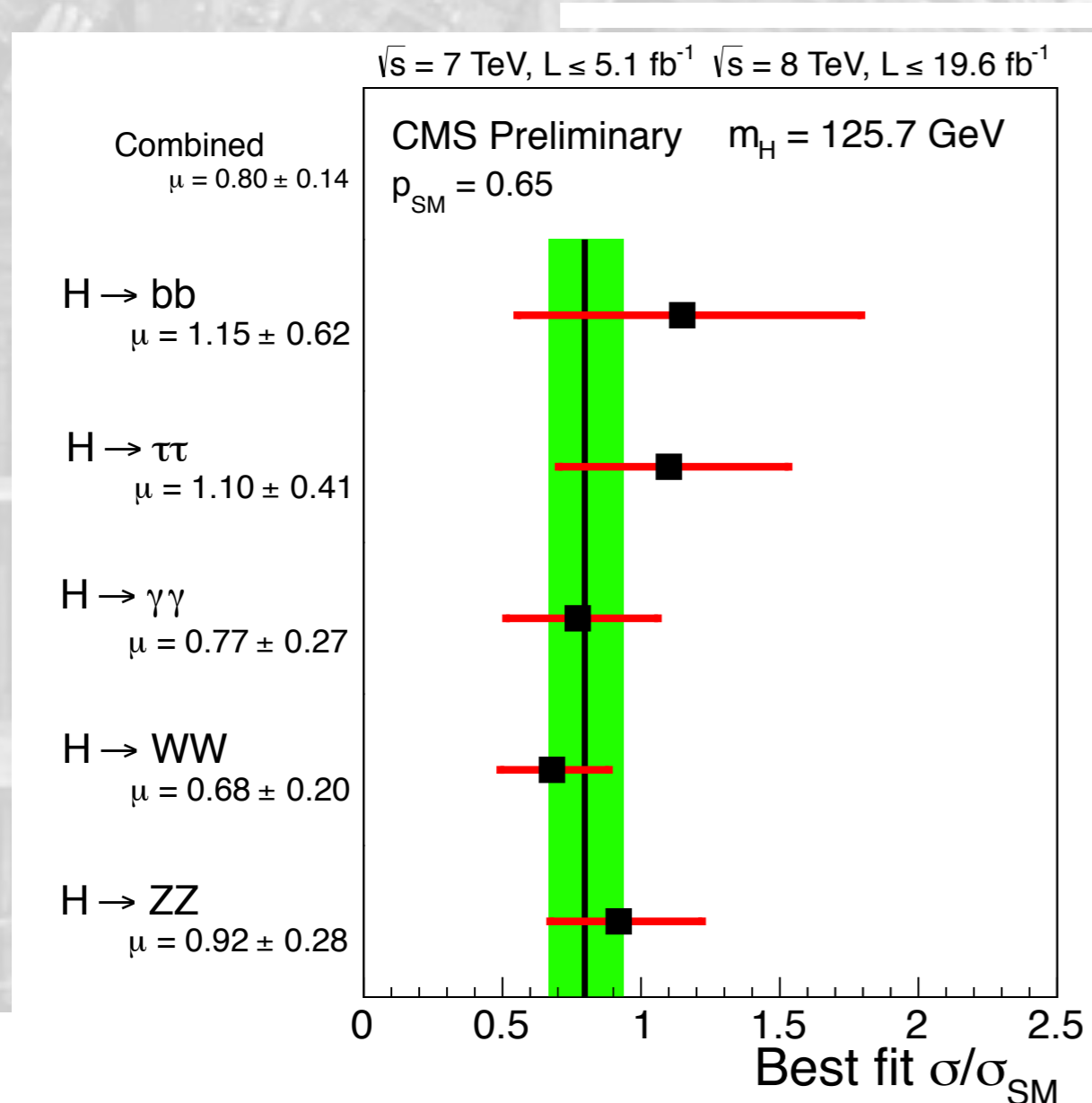


- Different channels probe different Higgs properties (e.g., spin, CP, etc)

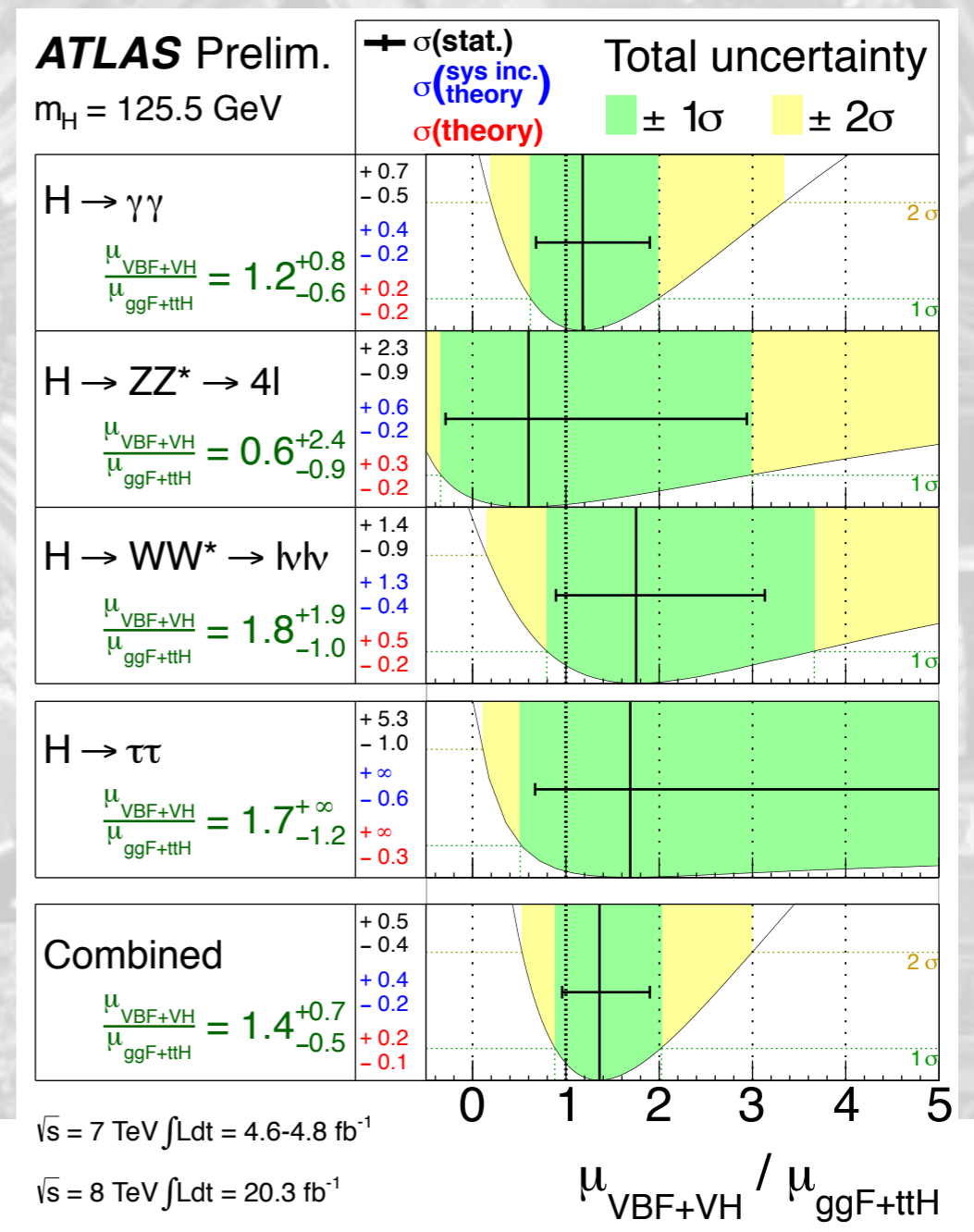


COUPLING FIT

- Established coupling to fermions and vector bosons
- Couplings scale as expected in the SM
- Deviations are possible (within errors) but cannot be of O(1) w/o introducing tension between different channels



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COUPLING FIT

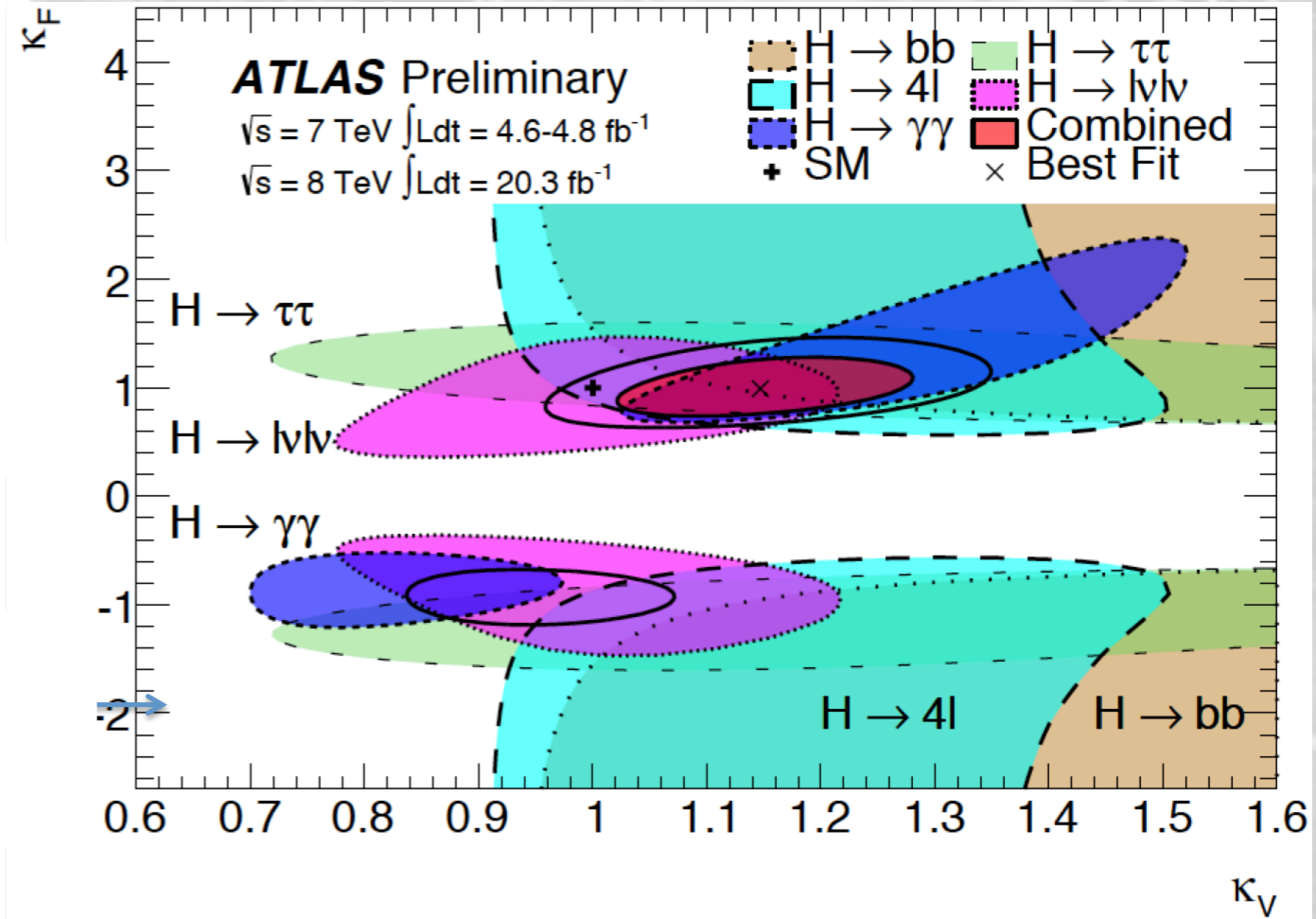
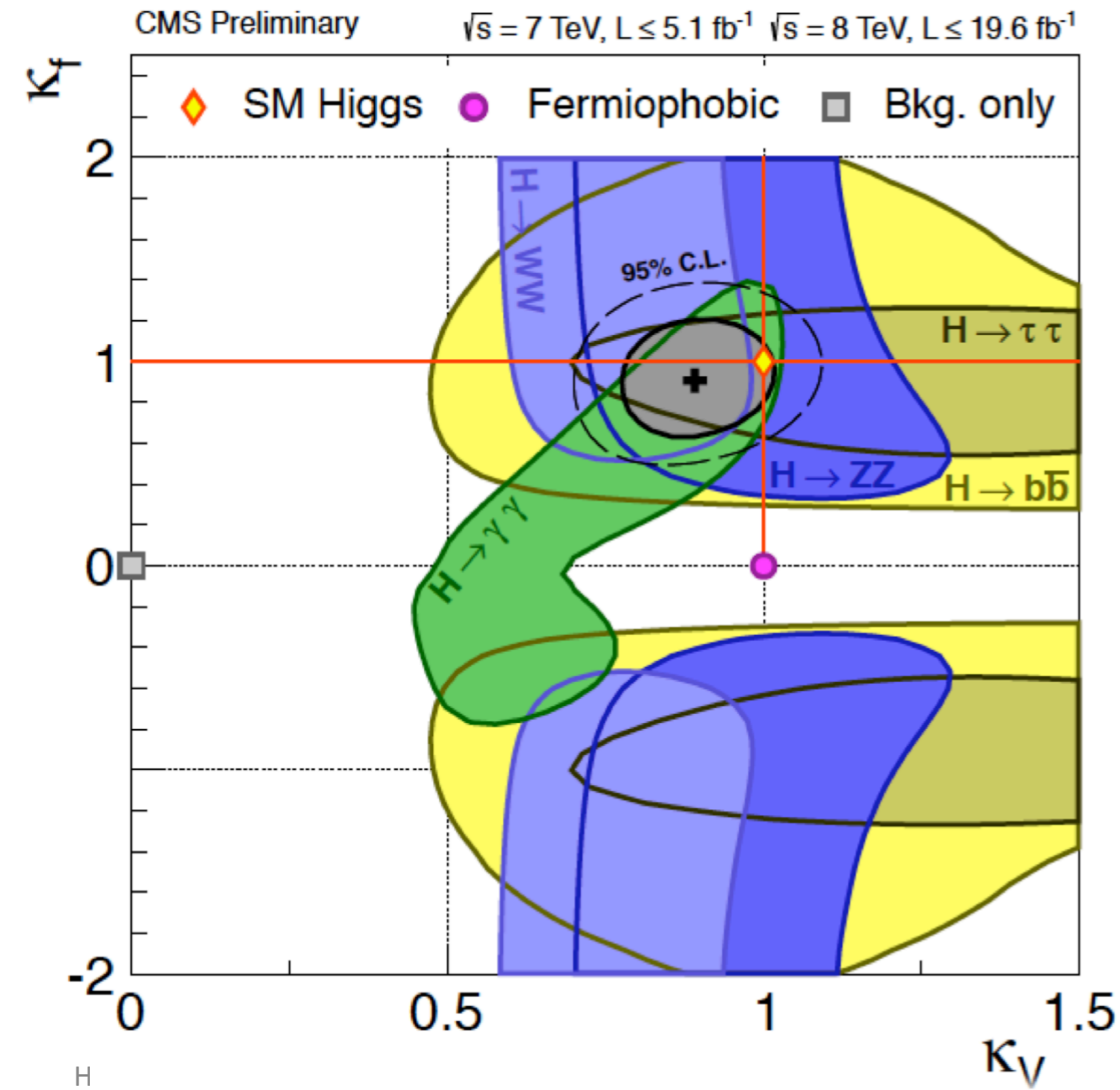
- Generalize couplings introducing multiplicative factors
- One channel can be function of more than one parameter

$$\kappa_g^2(\kappa_b, \kappa_t) = \frac{\kappa_t^2 \cdot \sigma_{ggH}^{tt} + \kappa_b^2 \cdot \sigma_{ggH}^{bb} + \kappa_t \kappa_b \cdot \sigma_{ggH}^{tb}}{\sigma_{ggH}^{tt} + \sigma_{ggH}^{bb} + \sigma_{ggH}^{tb}}$$

- Simplified analysis assuming universal fermion (κ_f) and vector (κ_v) deviations

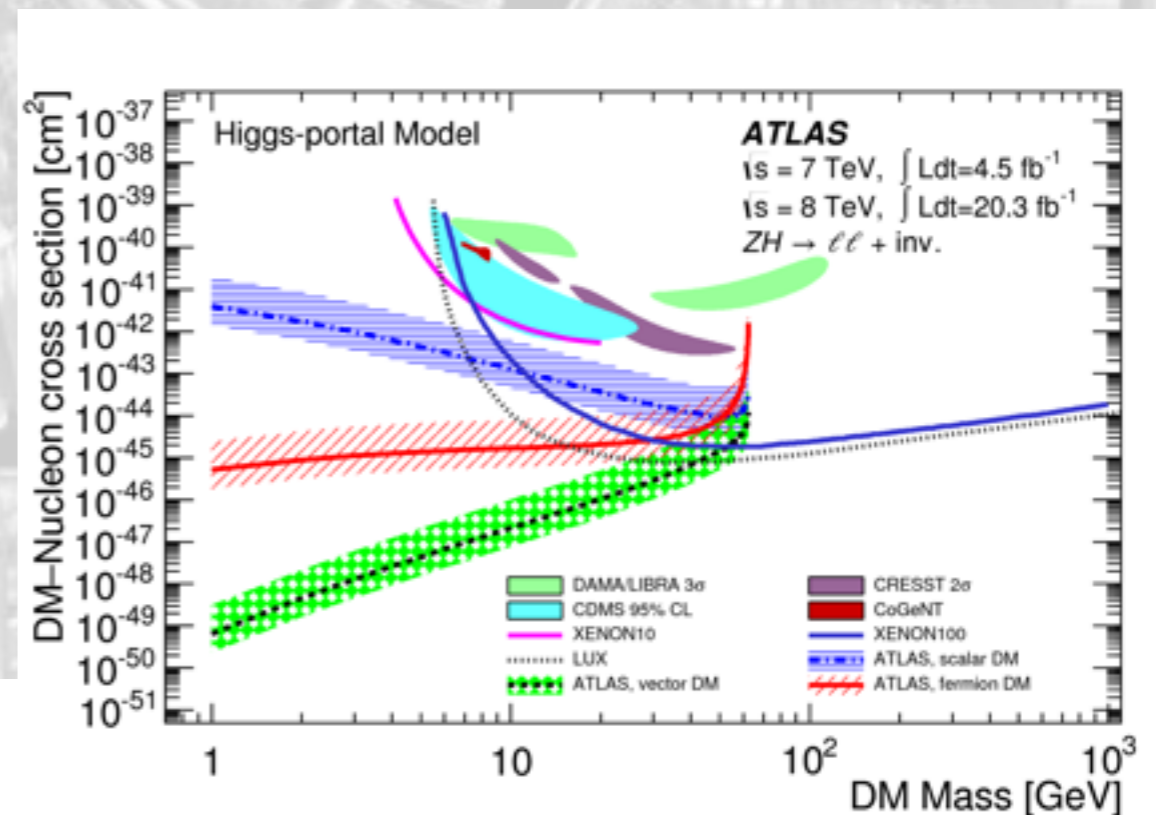
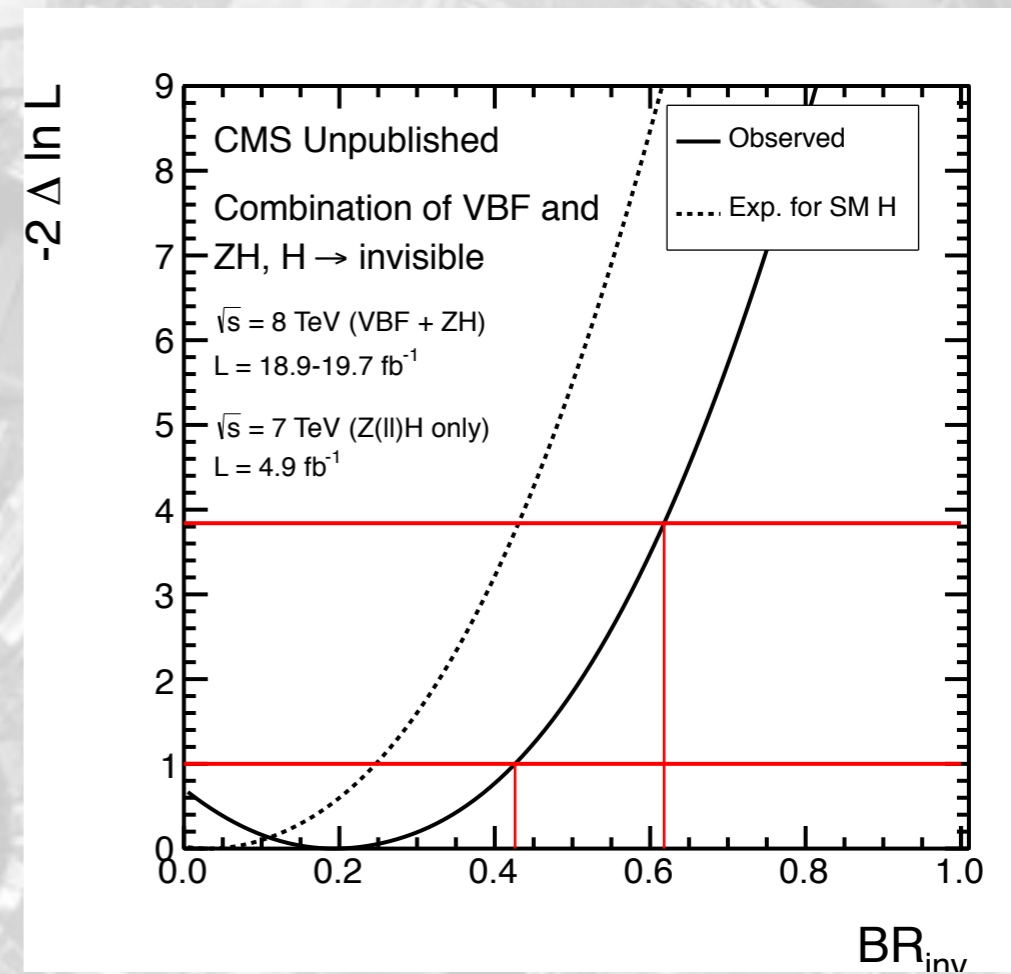
Analysis	Prod.	Decay	Analysis	Prod.	Decay
H	κ	κ	H	κ	κ
H	κ	κ	H	κ	κ
H	κ	κ	H	κ	κ
H	κ	κ			
H	κ	κ			

COUPLING FIT



HIGGS AND NEW PHYSICS

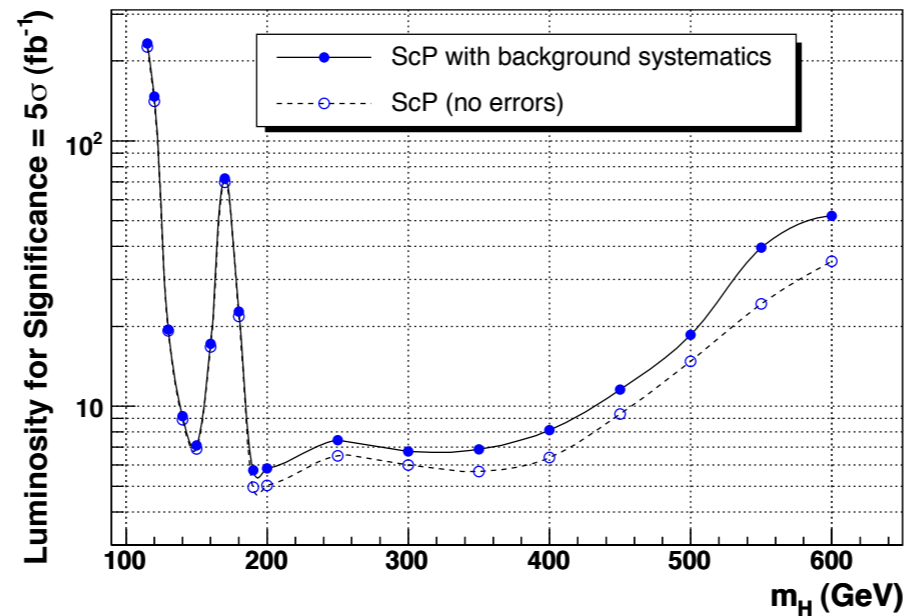
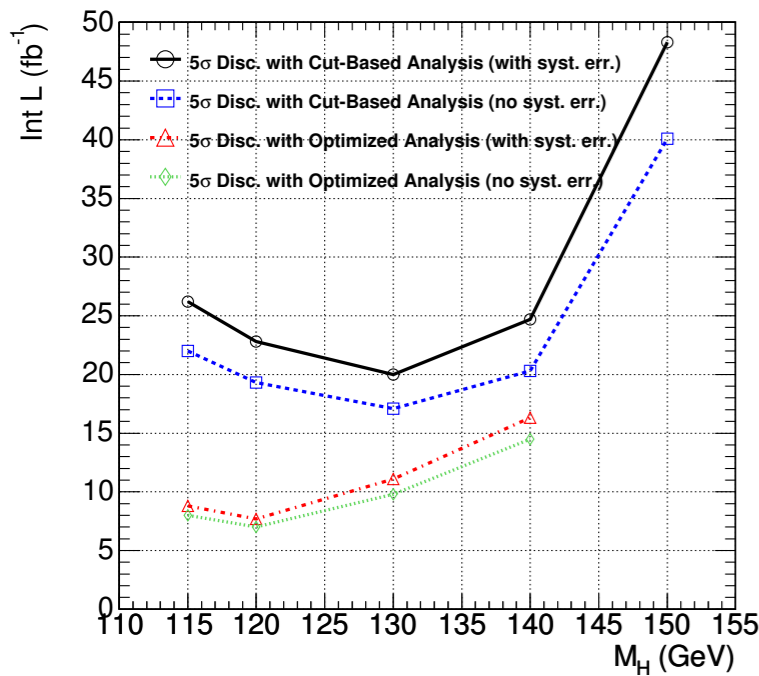
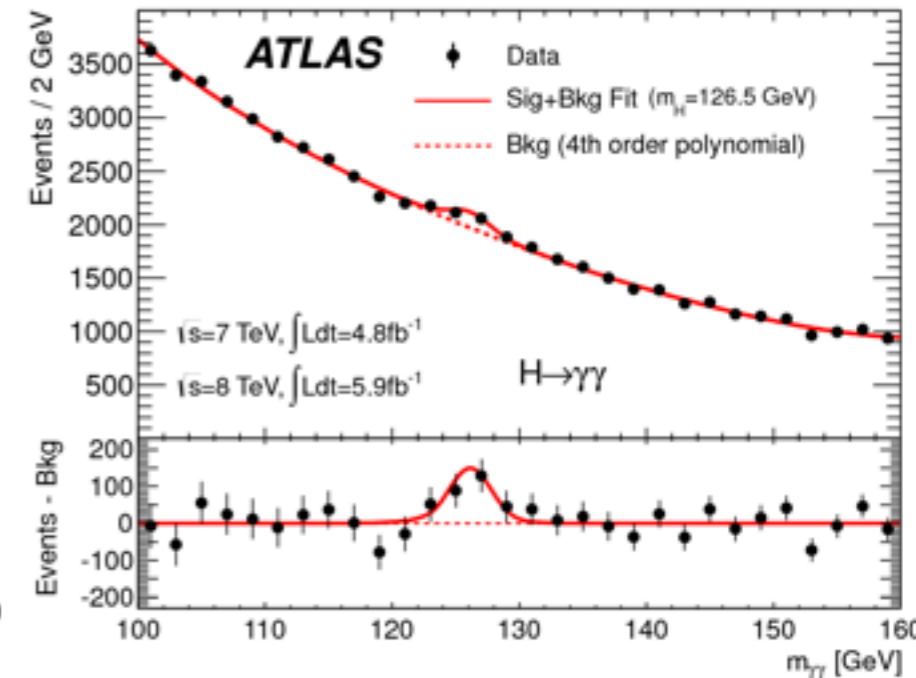
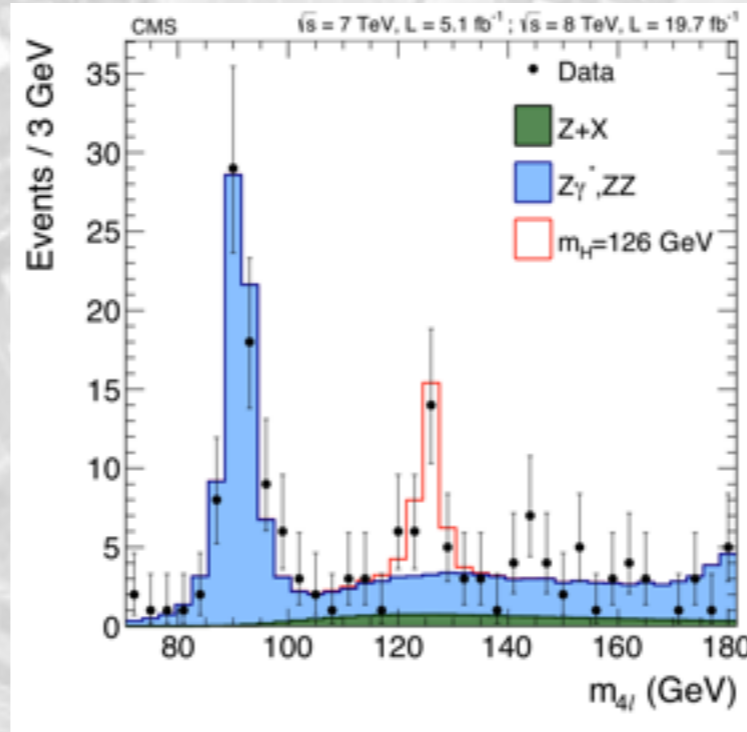
- The Higgs could give us access to new particles in the decay
- These particles could be invisible, and probed with W/Z +invisible searches
- Beside constraining the BR to invisible, these searches have implication on DM models
- Any conclusion, on the other hand, is model dependente (i.e., mind the assumptions)



THIS WAS FAST!!!

Run I marked the first success of the LHC program

the Higgs boson was indeed found



Two things to keep in mind

It arrived earlier than expected

we knew what to search for, and this helped A LOT

Discovery Lumi for 14 TeV collisions (where S/B is more favorable)

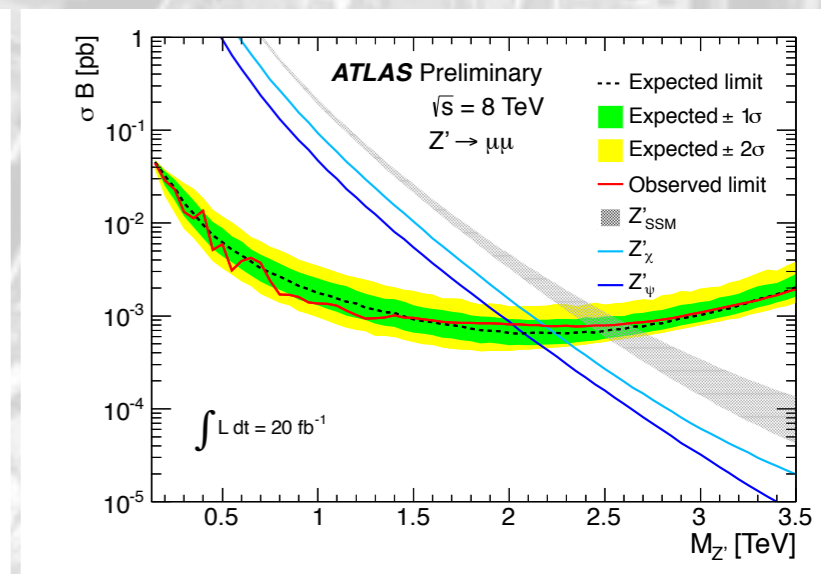
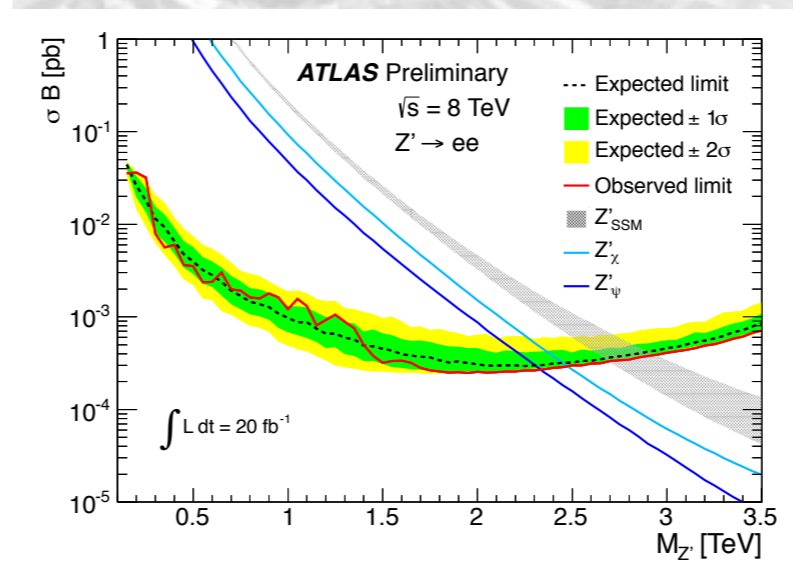
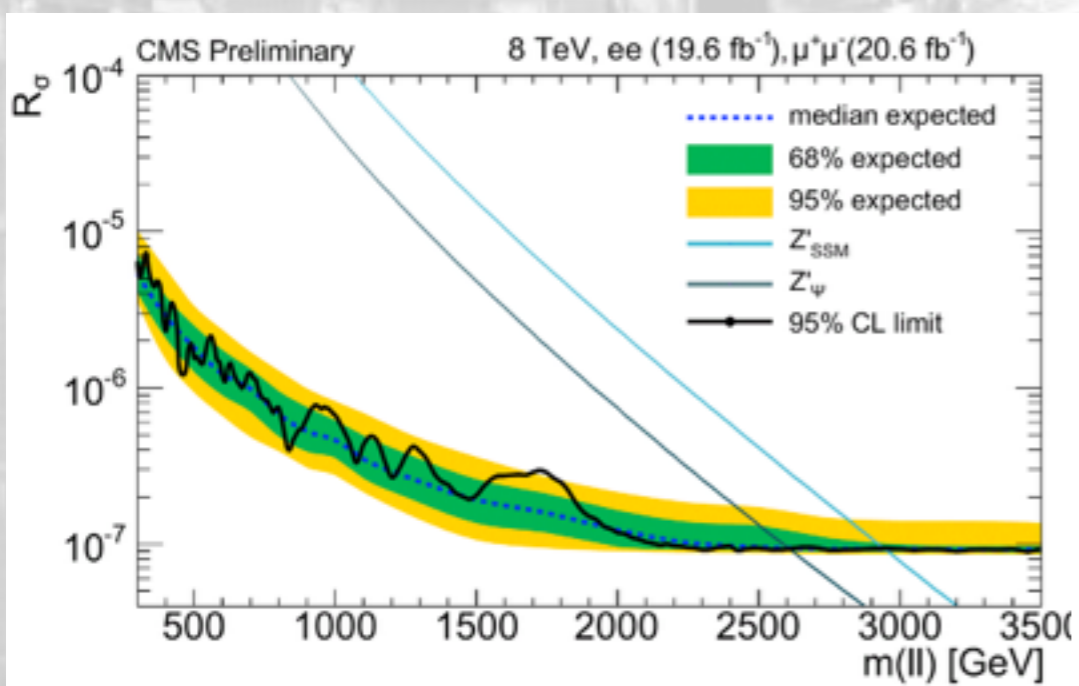
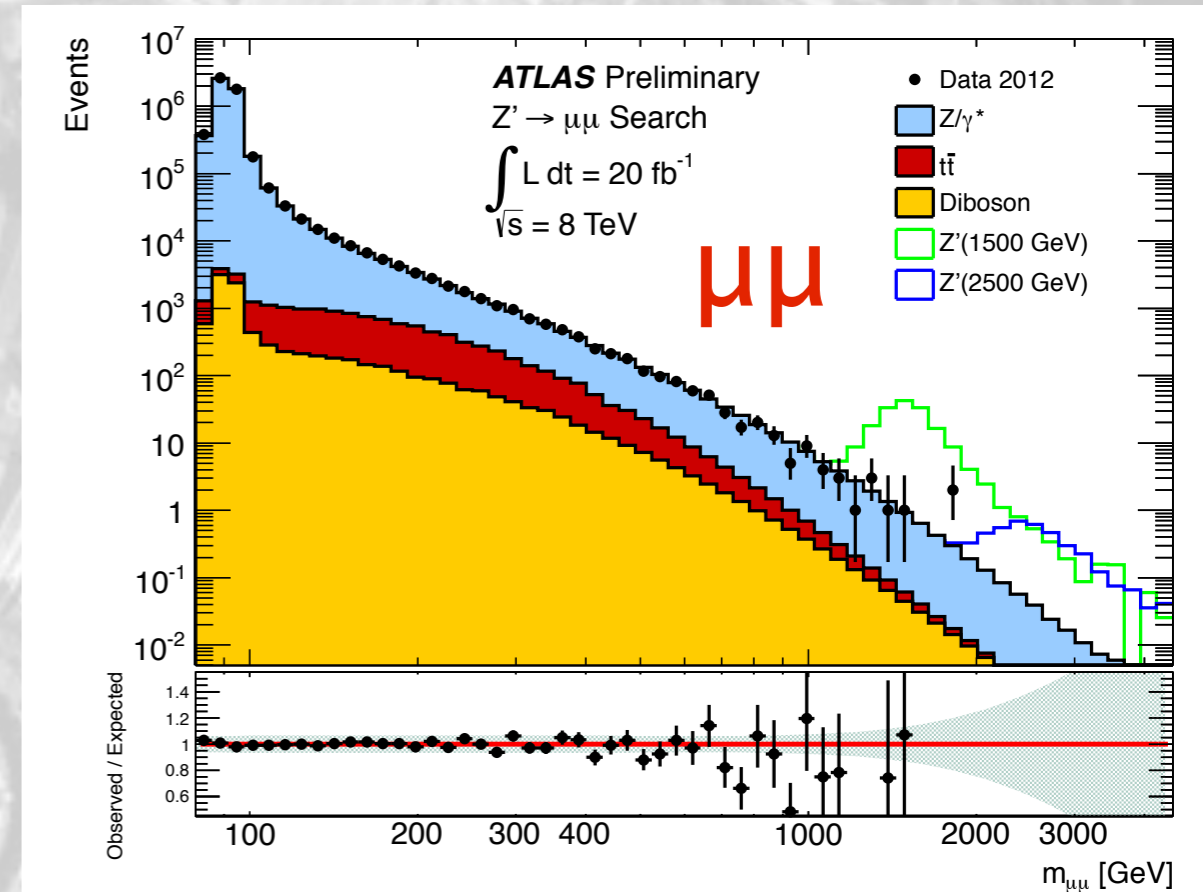
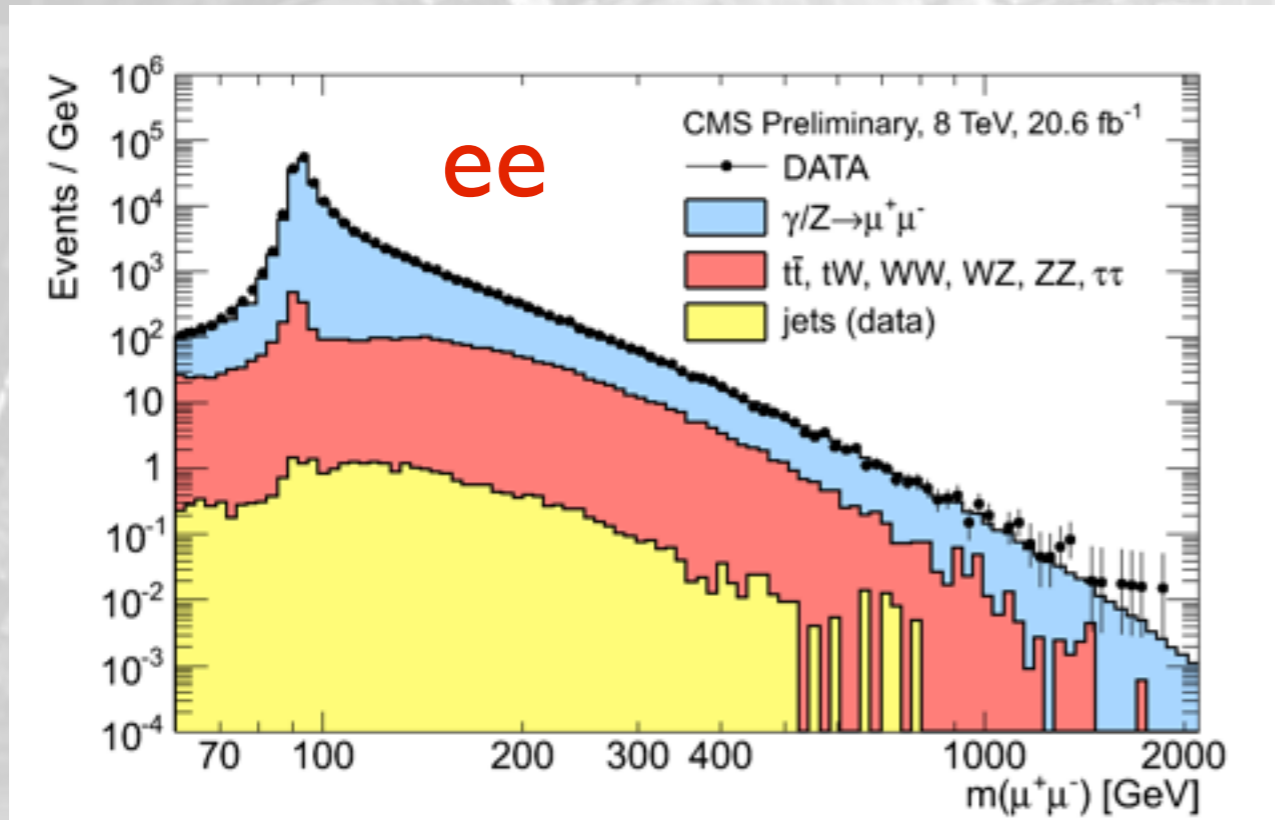


Highlights from Run I: search for heavy resonances

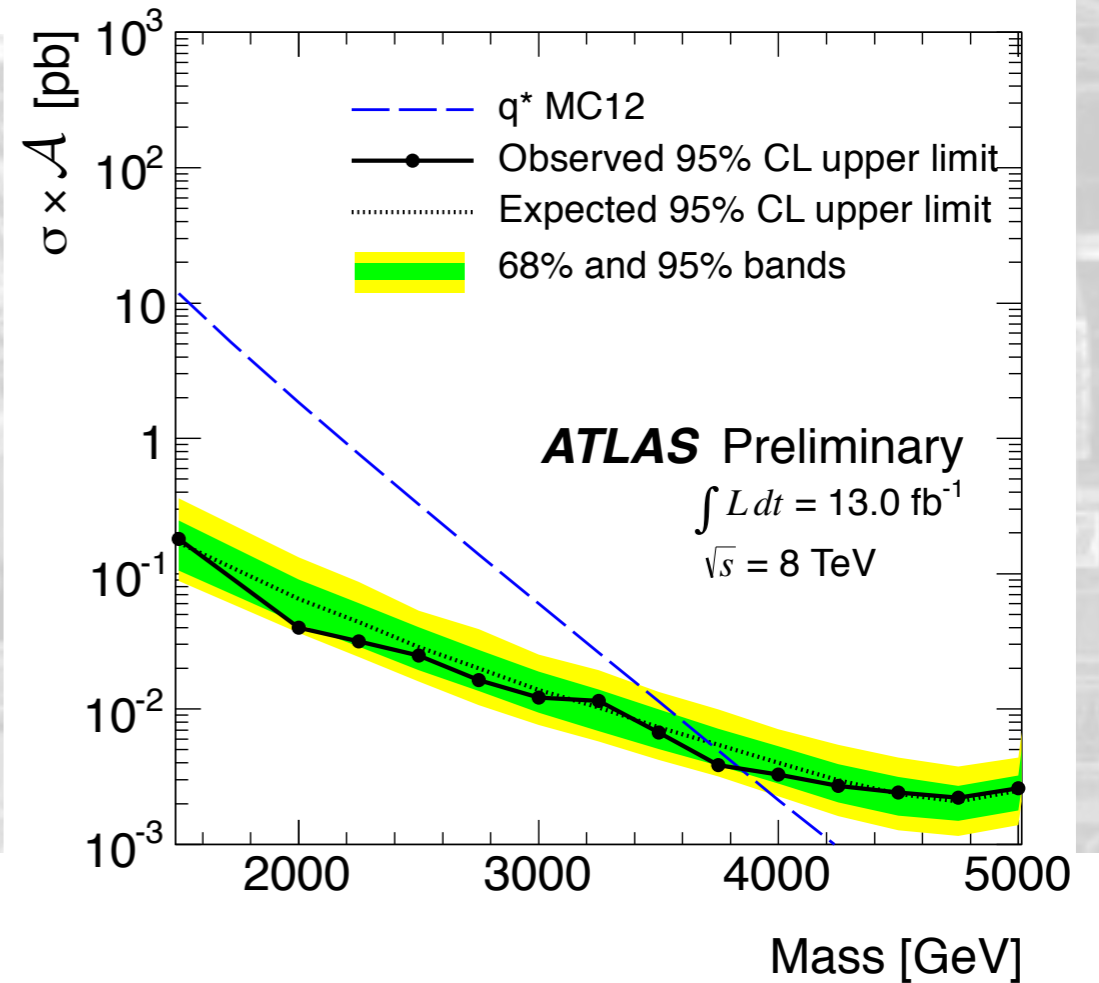
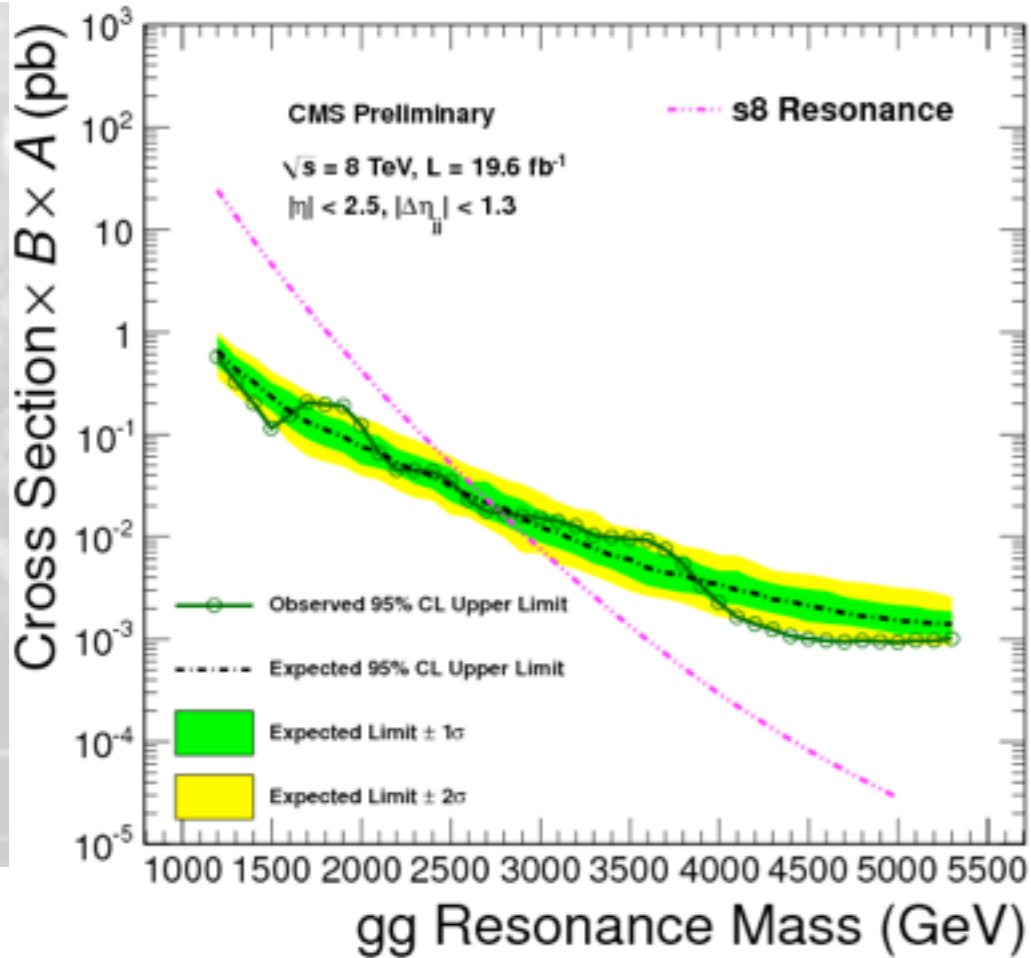
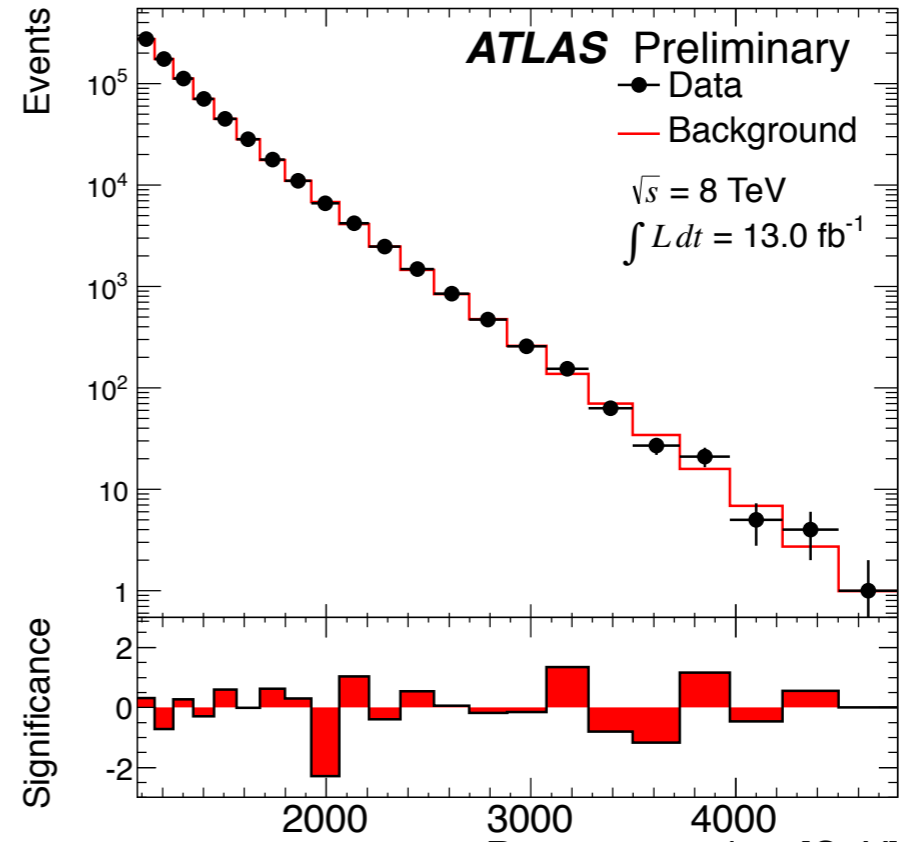
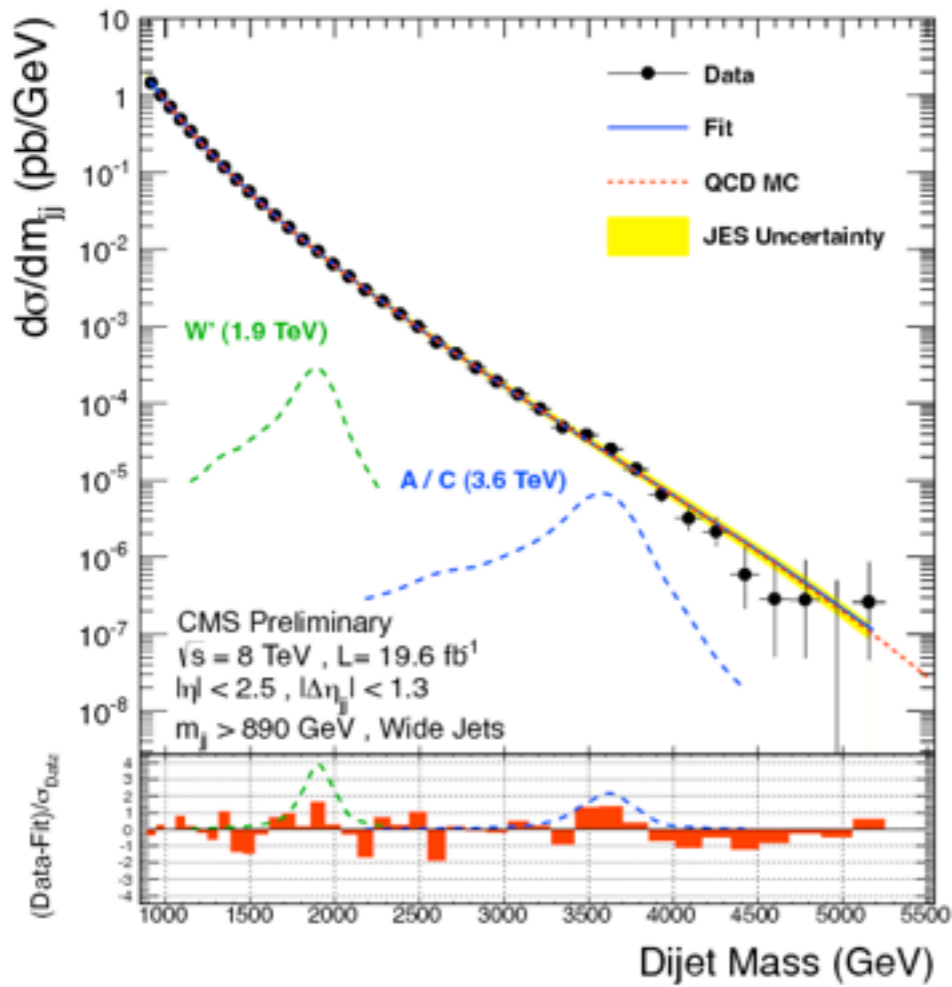
MOTIVATION

- Many models predict new resonances, coupling to fermions and/or vector bosons
 - New gauge interactions
 - Extra dimensions
 - Compositeness
- Analyses are generic (i.e., not tuned on a model) and quite simple
 - select two objects
 - look for a bump in diobject mass spectrum

DILEPTON SEARCH



DIJET SEARCH



FROM JETS TO BOOSTED OBJECTS

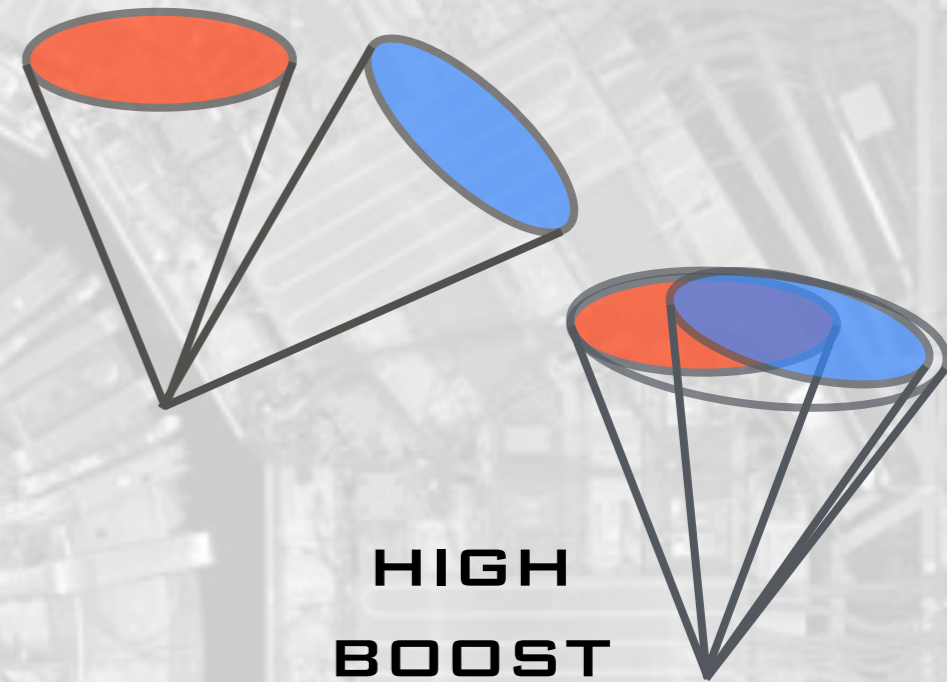
Jet substructure is now a common approach to look for boosted objects reconstructed as single jets

Several techniques implemented. But a cut on the “right” jet mass is still the strongest ingredient for bkg rejection in NP searches (at least for Ws and Zs)

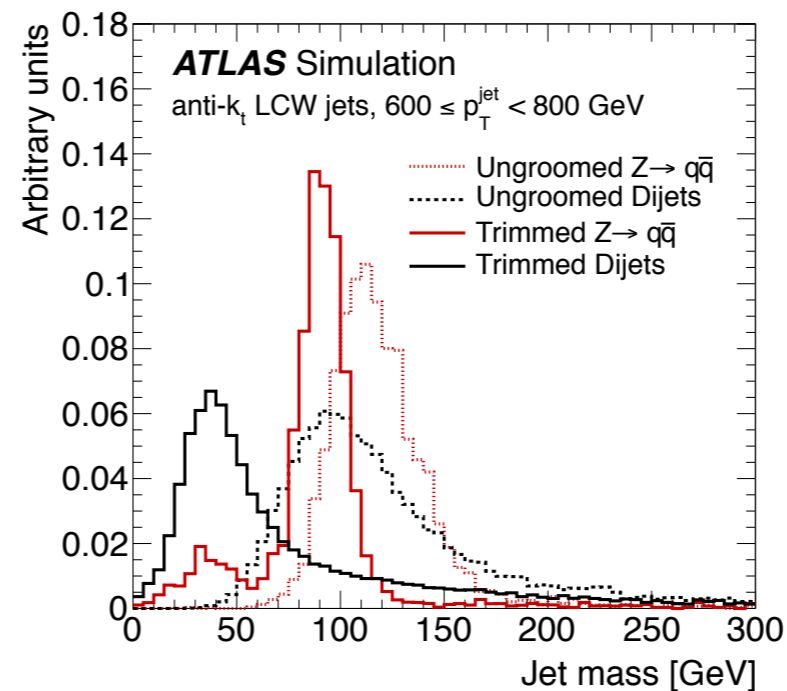
The right mass comes from grooming. We are still exploring possibilities

This is the Run II Jet R&D

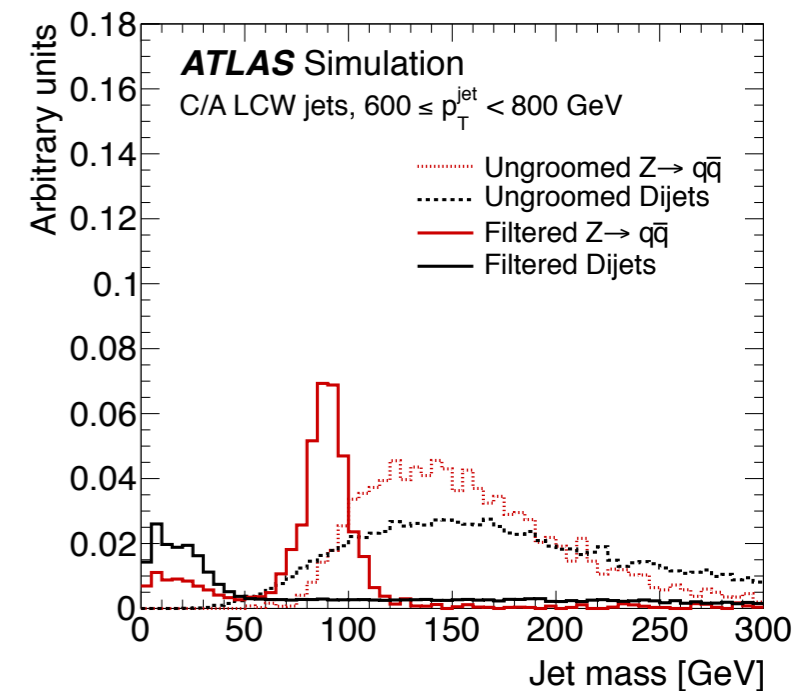
LOW BOOST



HIGH BOOST



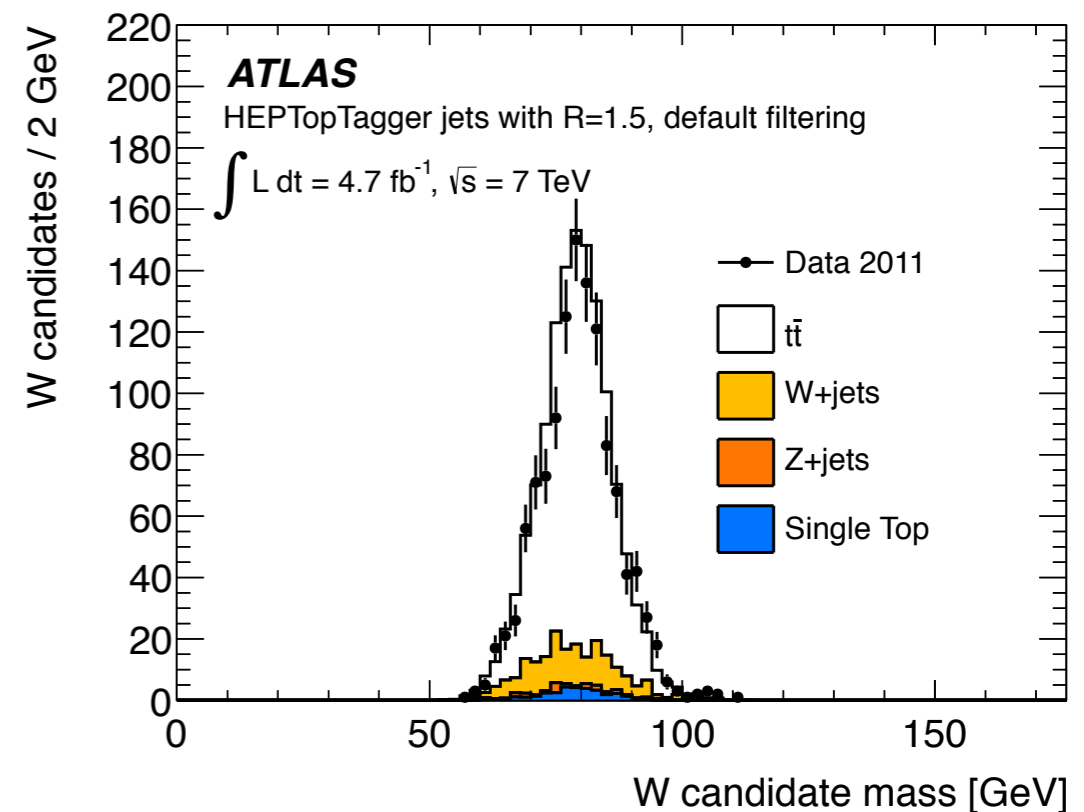
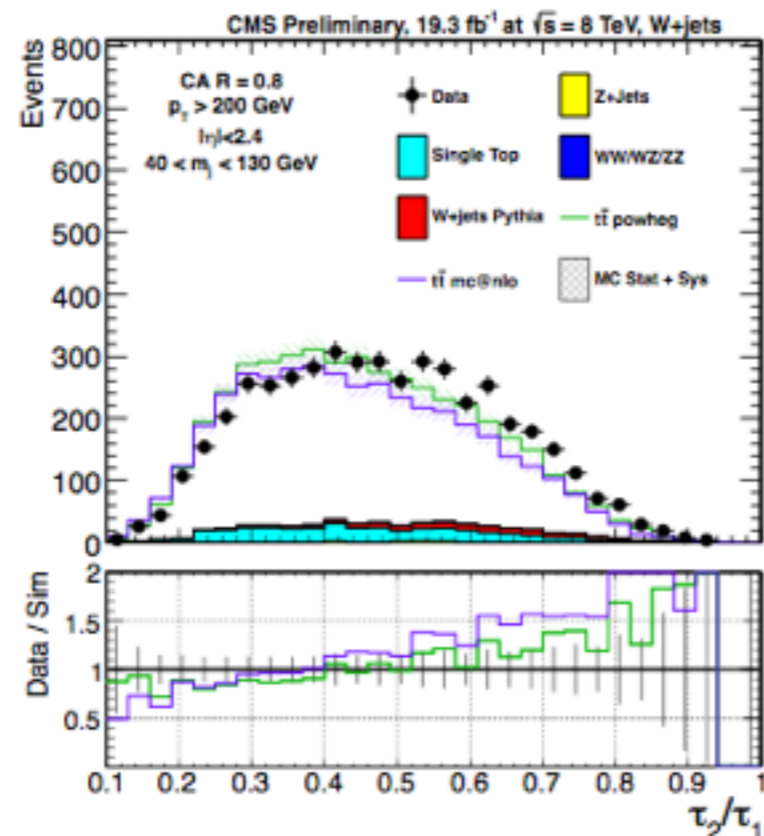
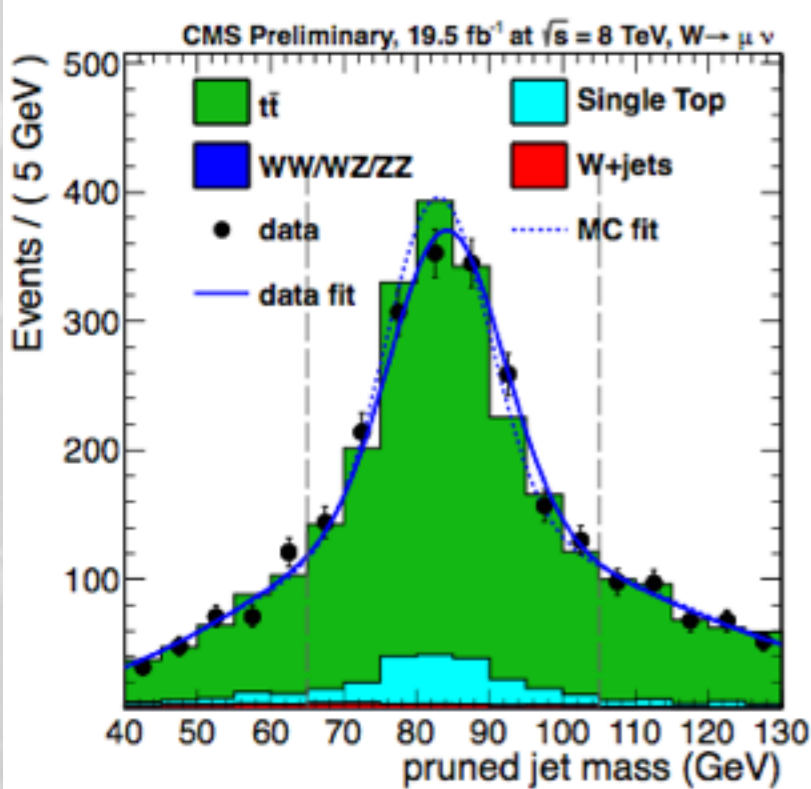
(a) anti- k_t , $R = 1.0$



(b) C/A, $R = 1.2$

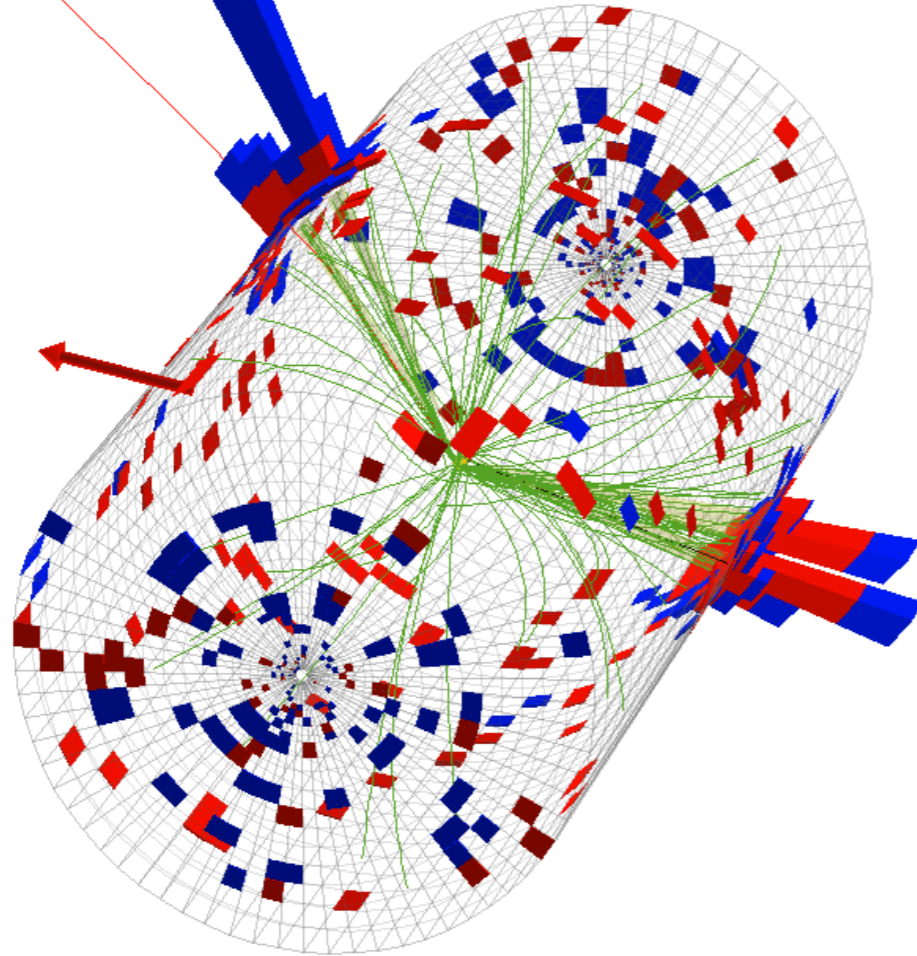
VALIDATION WITH DATA

- BOOSTED $t\bar{t}$ RECONSTRUCTED AS ONE b-JET + 1 LEPTON RECOILING AGAINST ONE bJET AND ONE JET (THE W CANDIDATE)
- PEAK IN THE JET MASS: WE ARE SEEING BOOSTED WS
- STUDY SUBSTRUCTURE VARIABLES DATA VS MC
- MC GET SUBSTRUCTURE QUITE RIGHT ($\sim 5\%$ SYSTEMATIC ON PREDICTED EFFICIENCY)

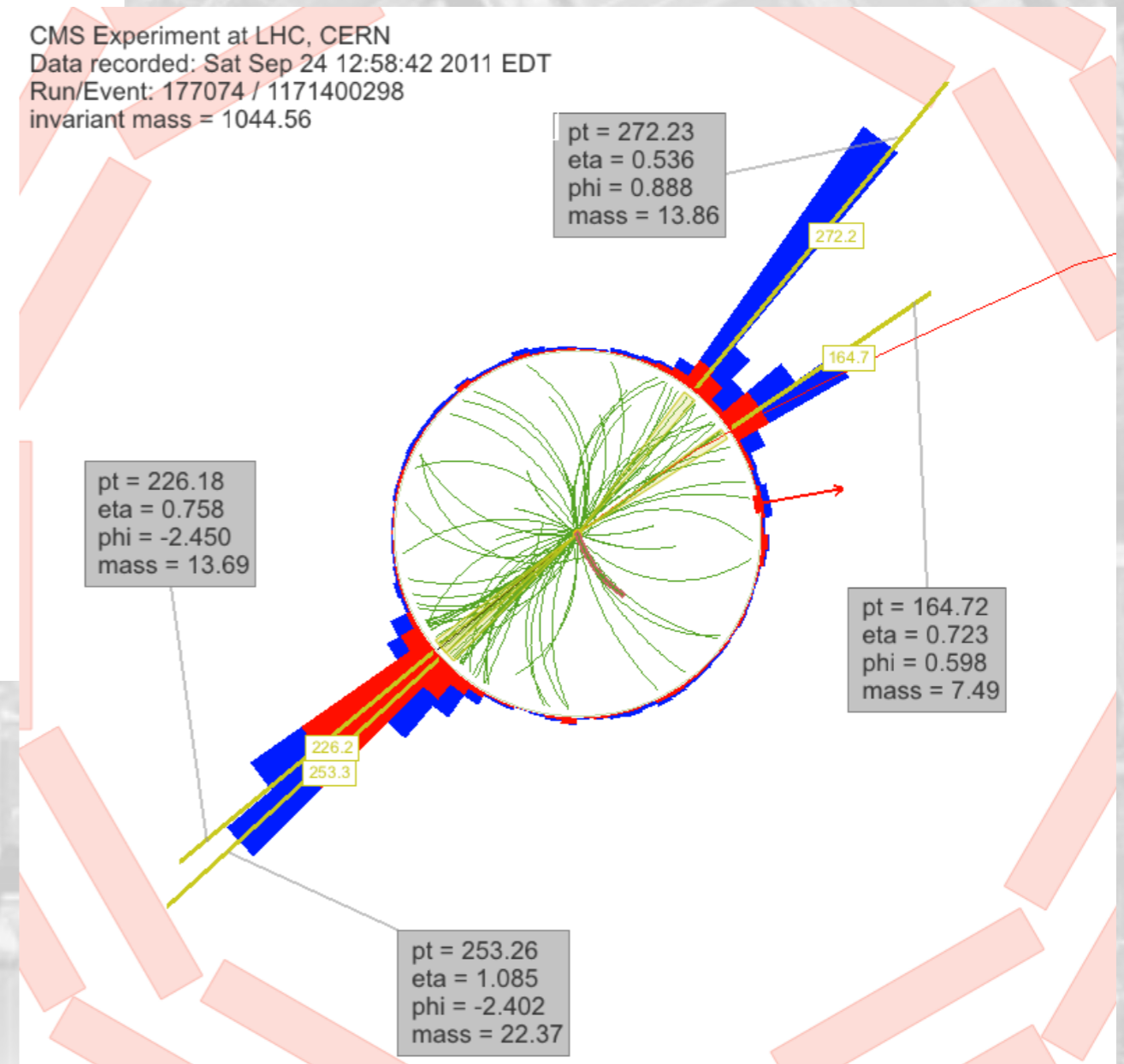


A DOUBLE-TAG EVENT

CMS Experiment at LHC, CERN
Data recorded: Sat Sep 24 12:58:42 2011 EDT
Run/Event: 177074 / 1171400298
invariant mass = 1044.56

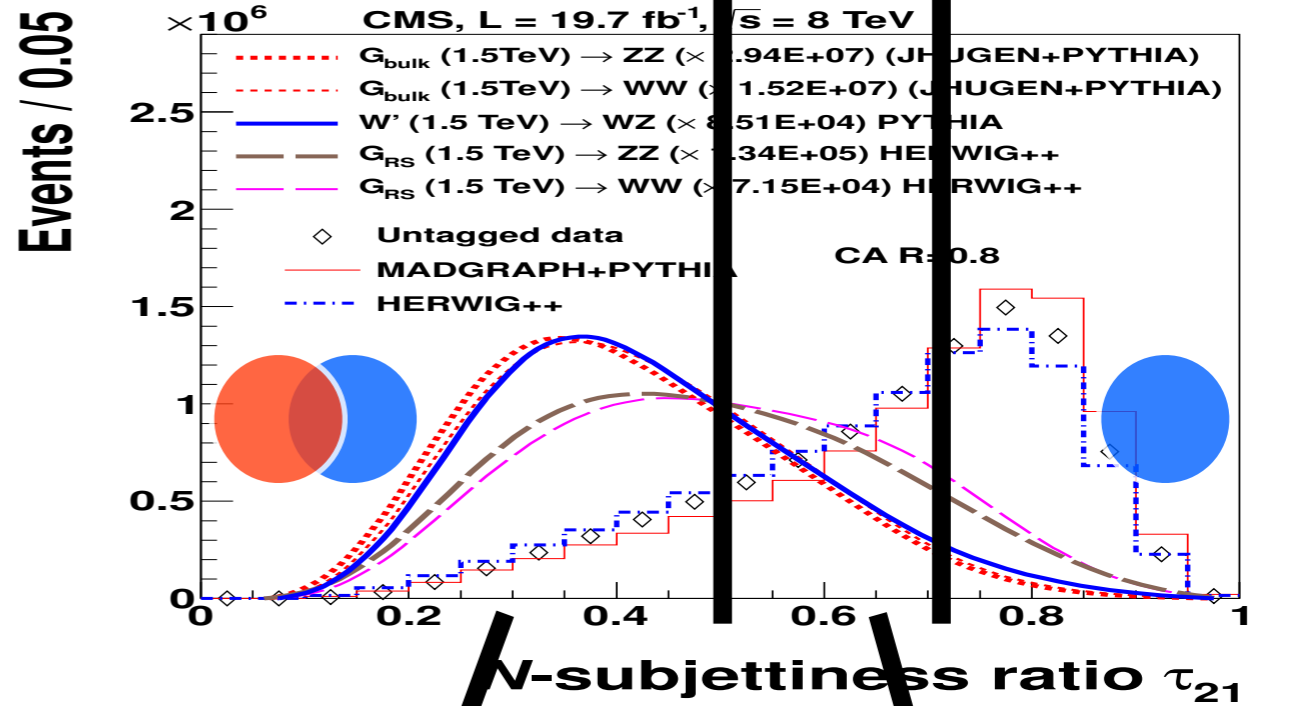


CMS Experiment at LHC, CERN
Data recorded: Sat Sep 24 12:58:42 2011 EDT
Run/Event: 177074 / 1171400298
invariant mass = 1044.56



ZZ/ZW/WW SEARCHES

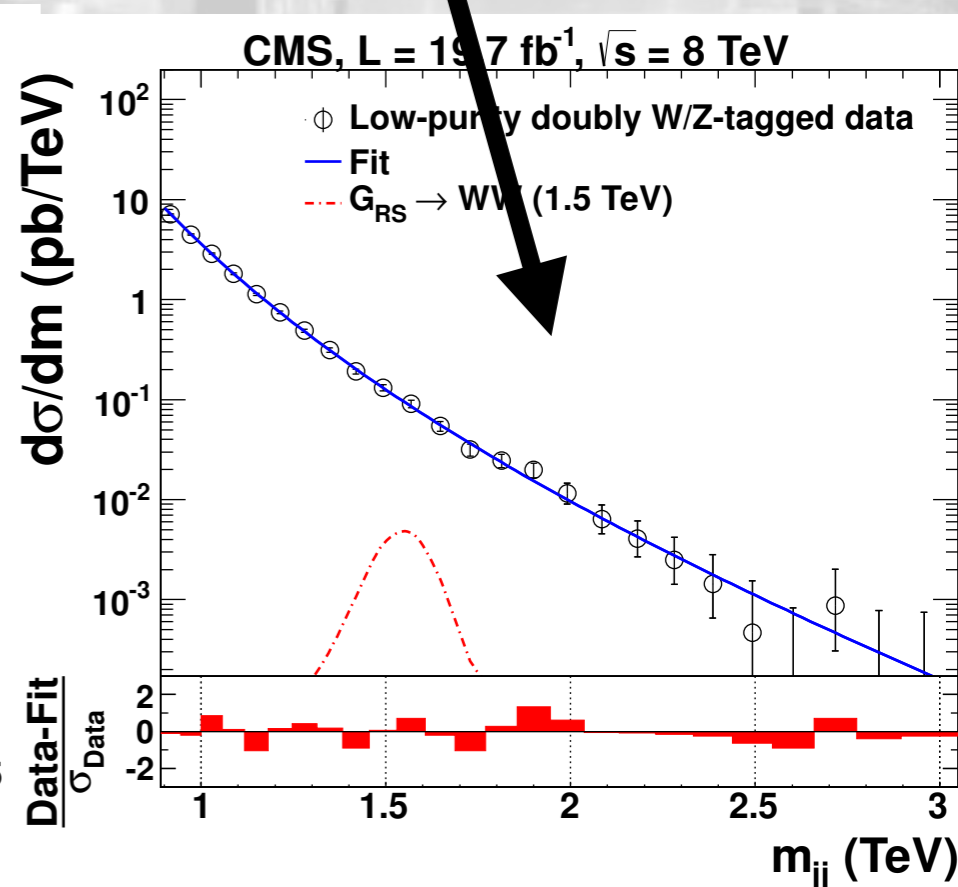
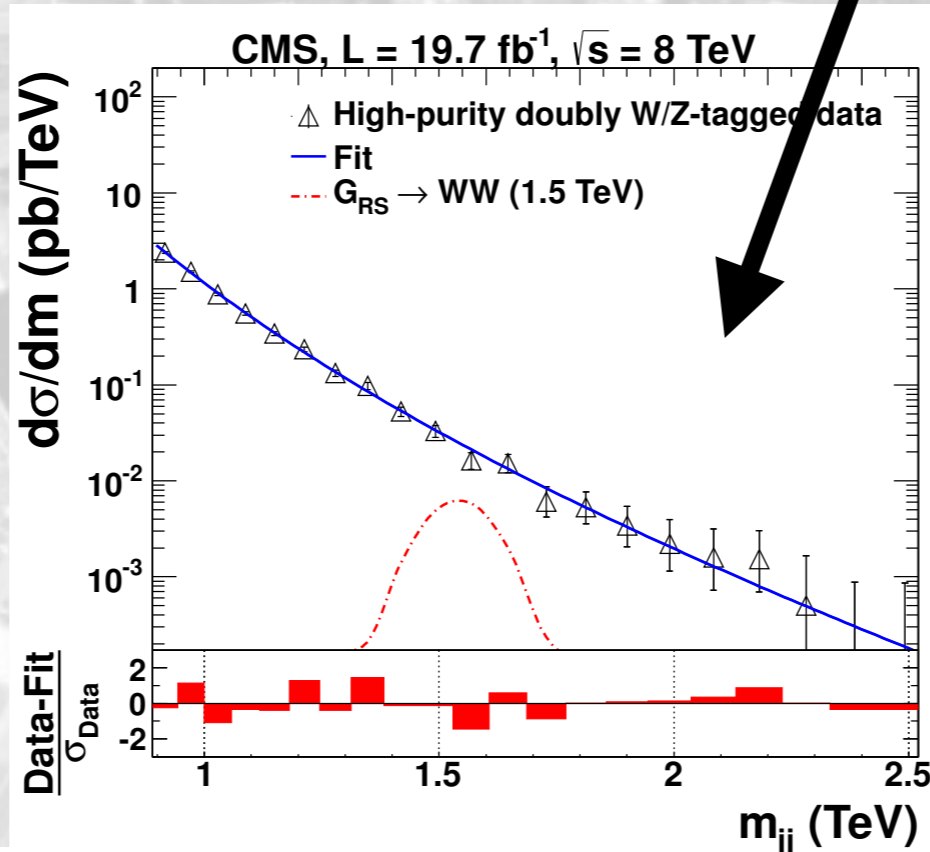
- Technique applied to **jj (WW/WZ/ZZ)**, $j\ell\nu$ (WW/WZ) and $j\ell\ell$ (WZ/ZZ) final states



High-purity

Low-purity

- Background determined with shape fit (as Hgg) or with absolute prediction from jet sideband



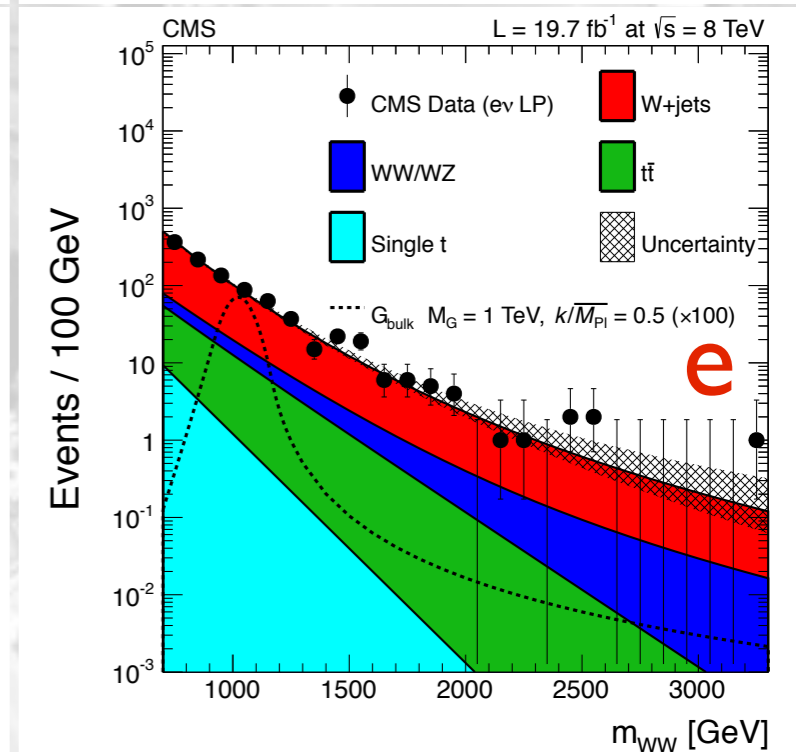
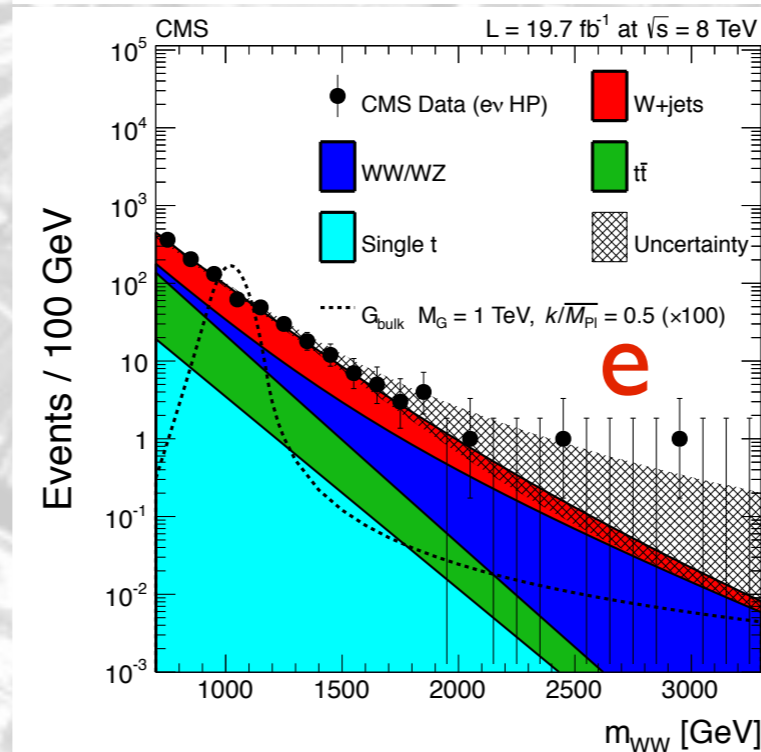
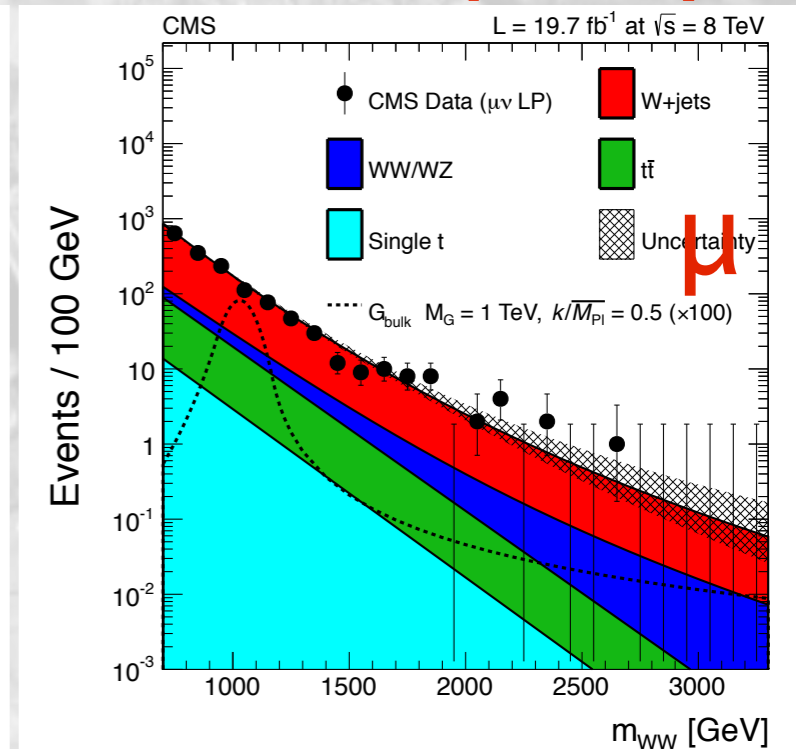
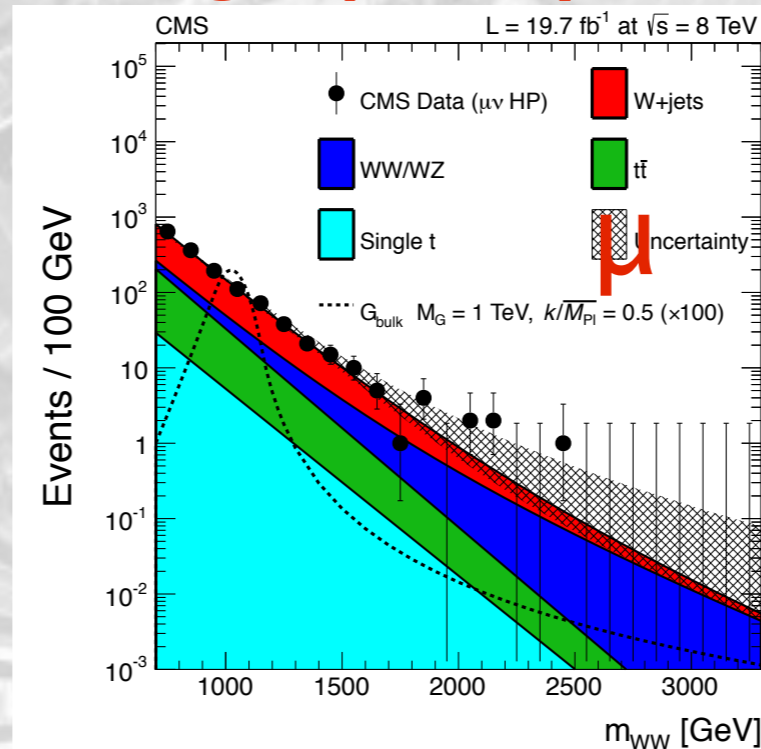
ZZ/ZW/WW SEARCHES

- Technique applied to jj (WW/WZ/ZZ), $j\ell\nu$ (WW/WZ) and $j\ell\ell$ (WZ/ZZ) final states

- Background determined with shape fit (as Hgg) or with absolute prediction from jet sideband

high-purity

low-purity

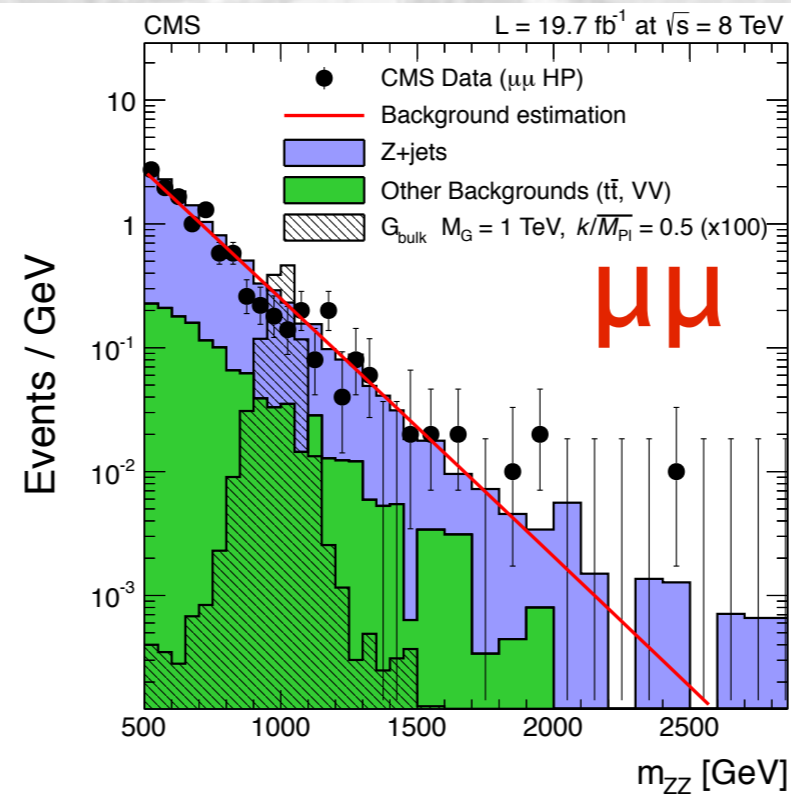


ZZ/ZW/WW SEARCHES

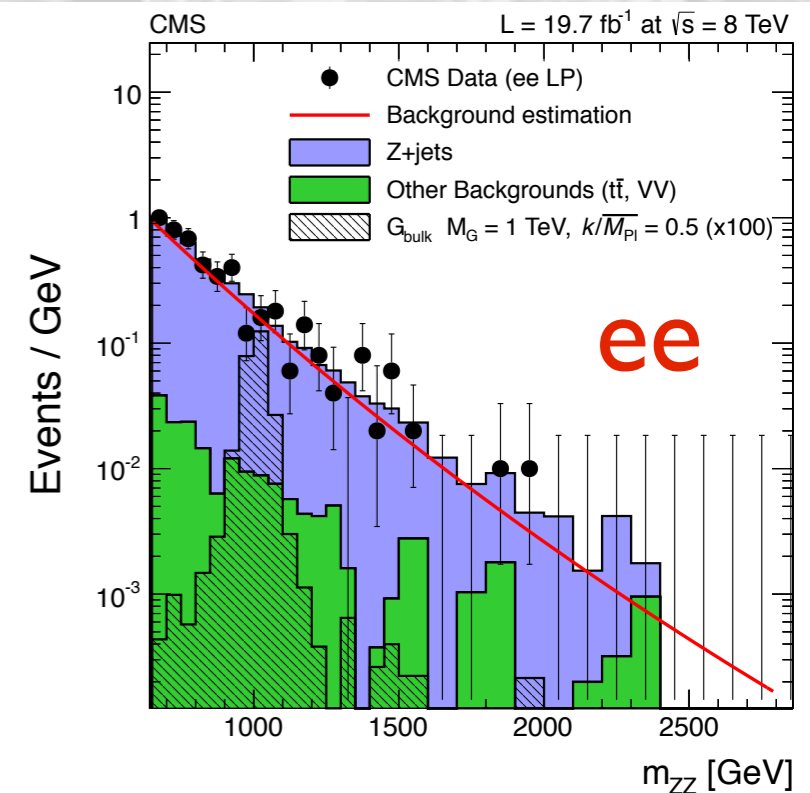
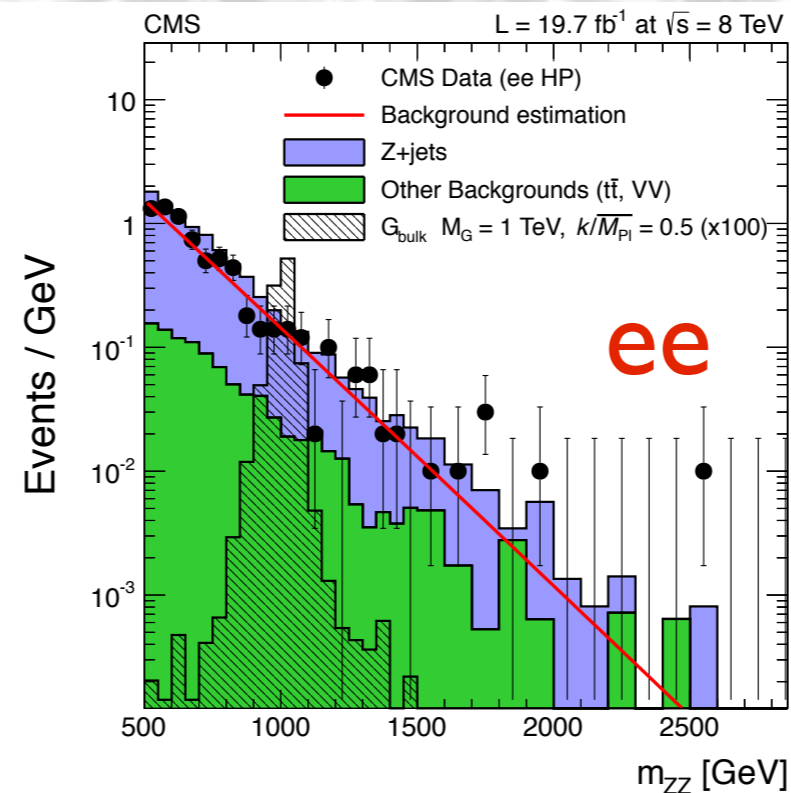
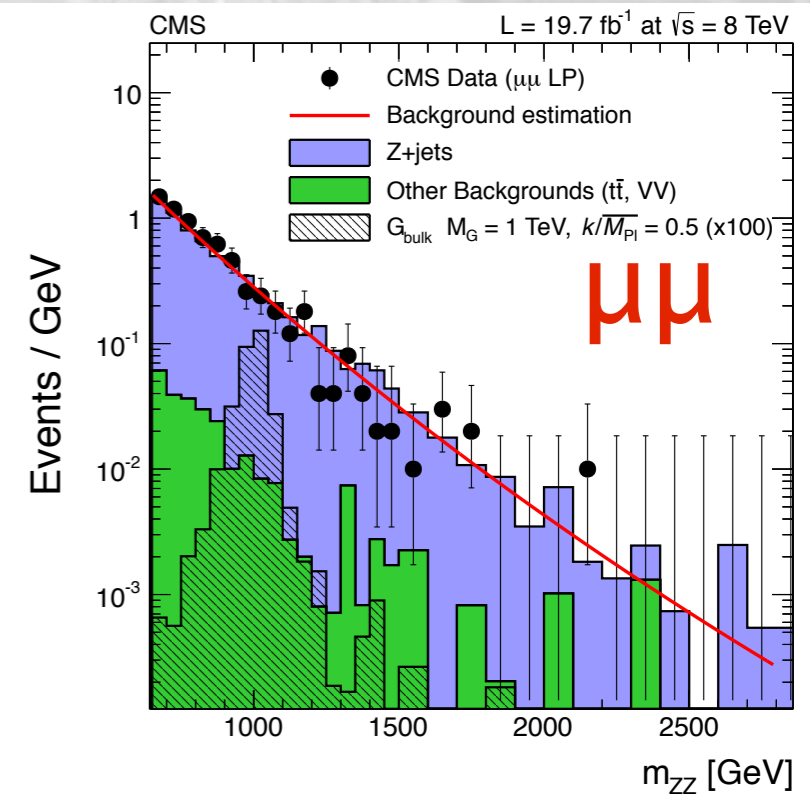
- Technique applied to jj (WW/WZ/ZZ), $j\ell\nu$ (WW/WZ) and $j\ell\ell$ (WZ/ZZ) final states

- Background determined with shape fit (as Hgg) or with absolute prediction from jet sideband

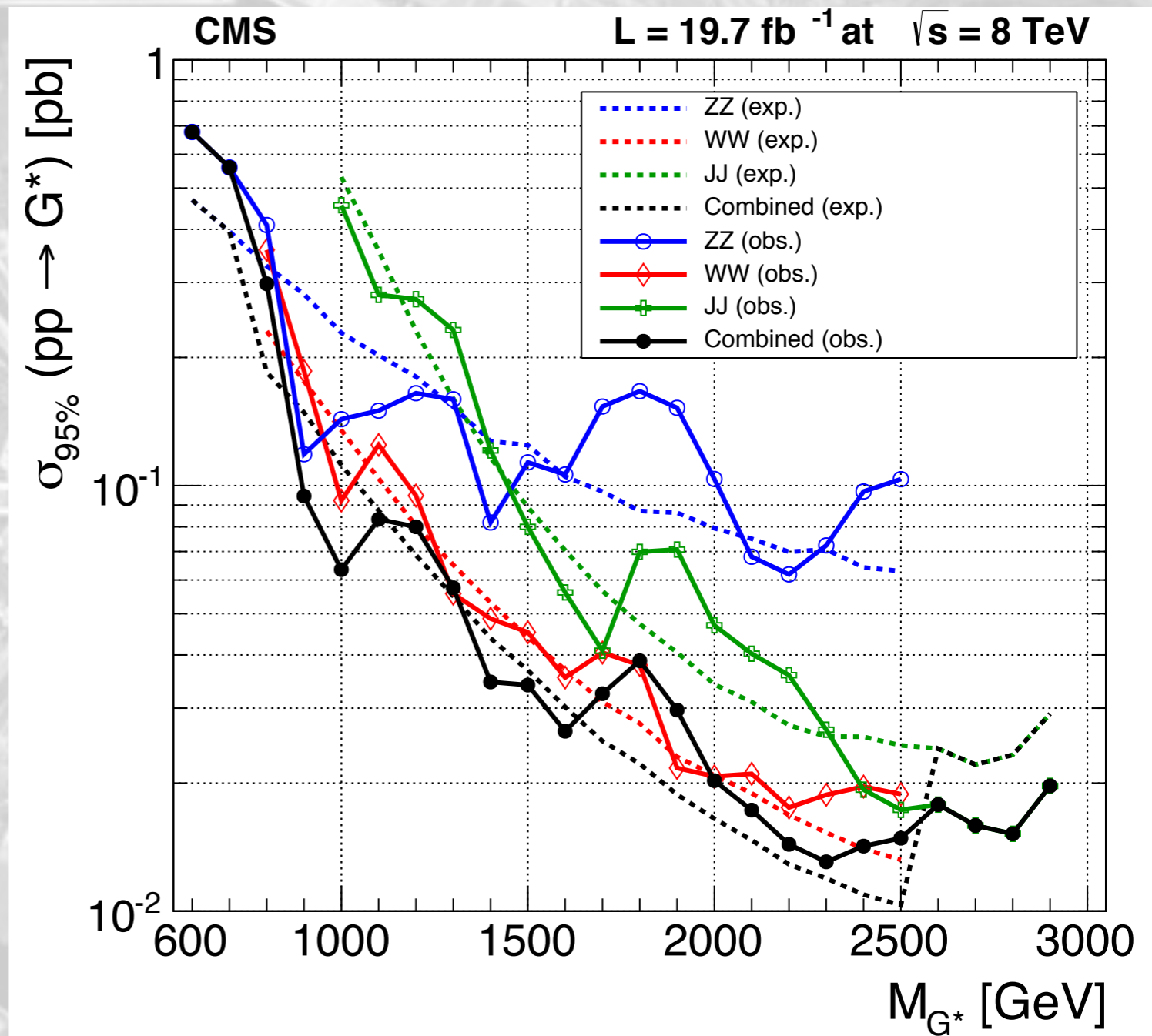
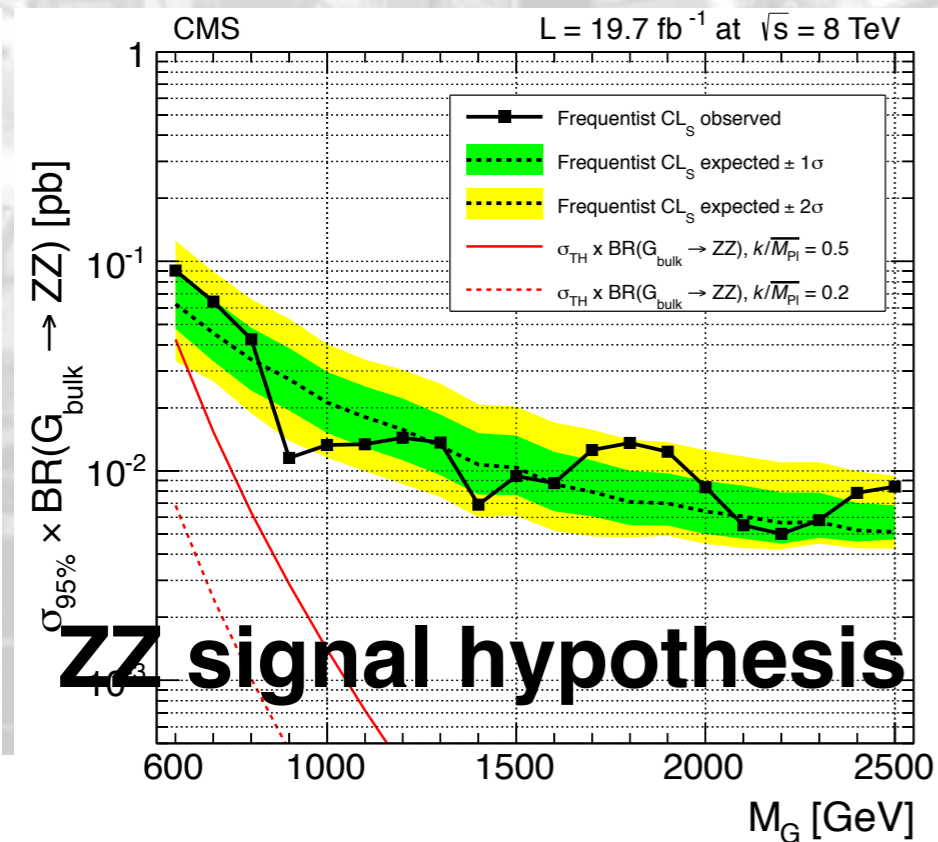
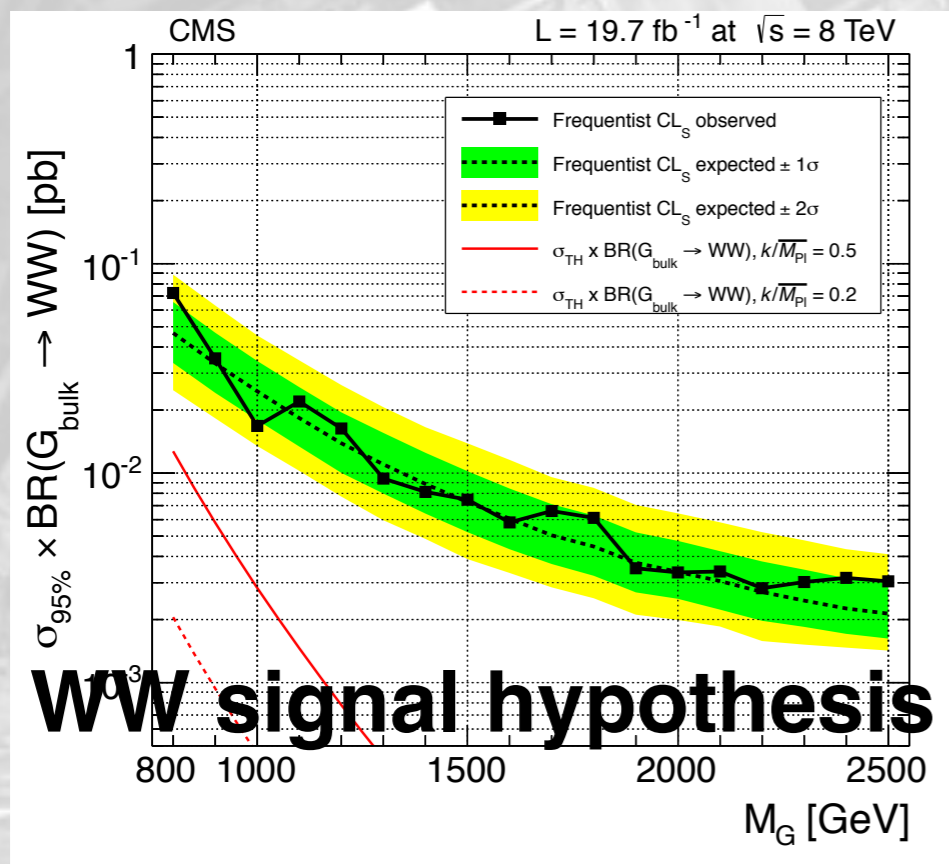
high-purity



low-purity



INTERPRETATION

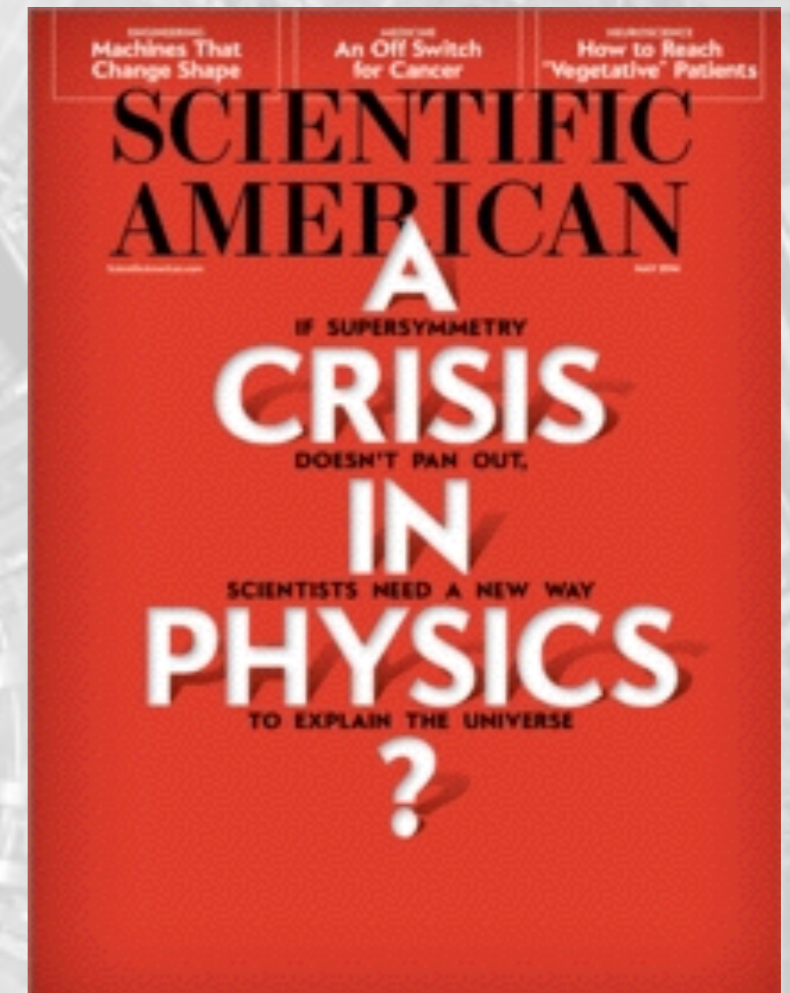
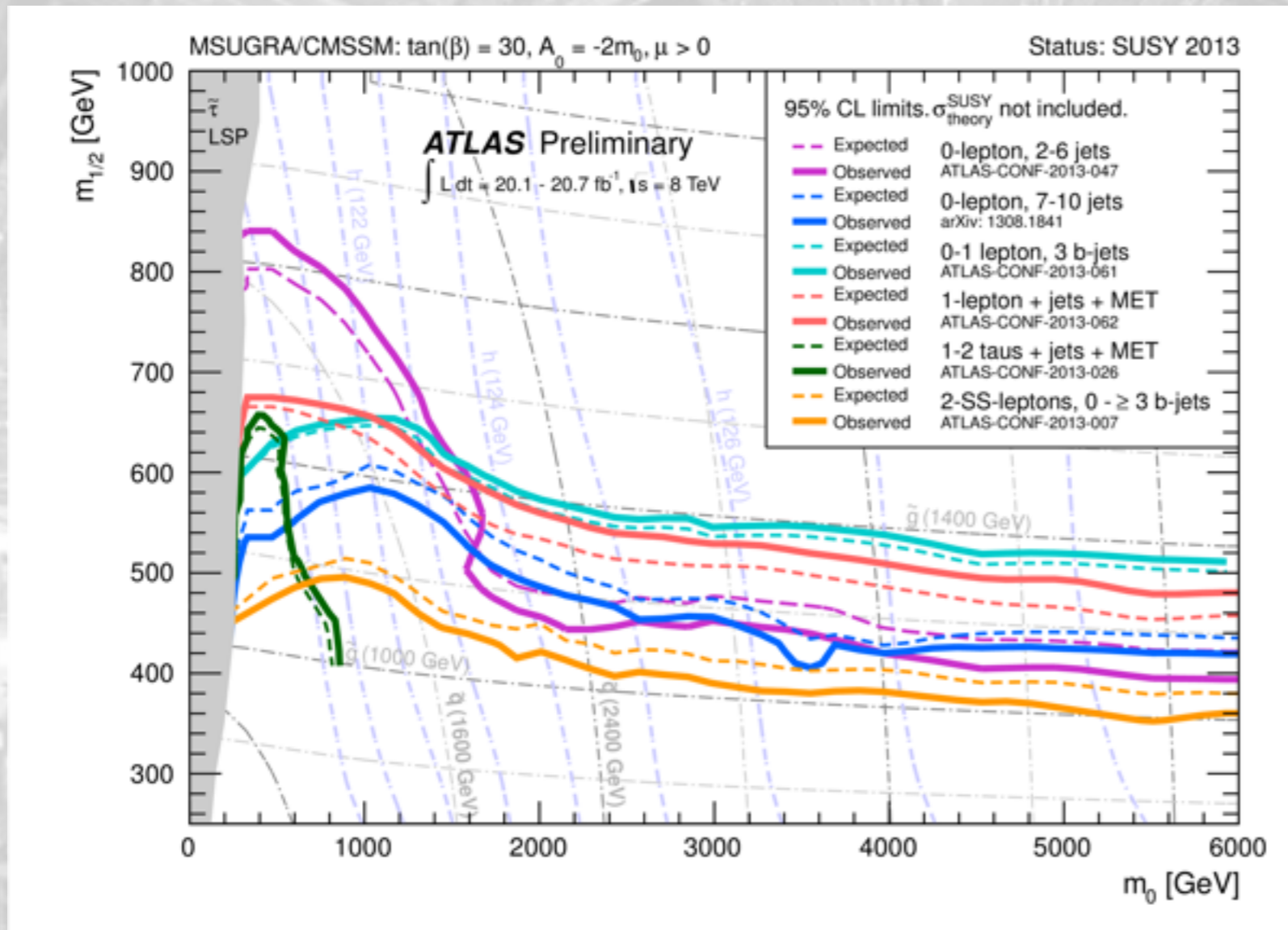




**Highlights from Run I:
search for DM in cascade
(aka SUSY searches)**

THERE WAS NOTHING BEHIND THE CORNER

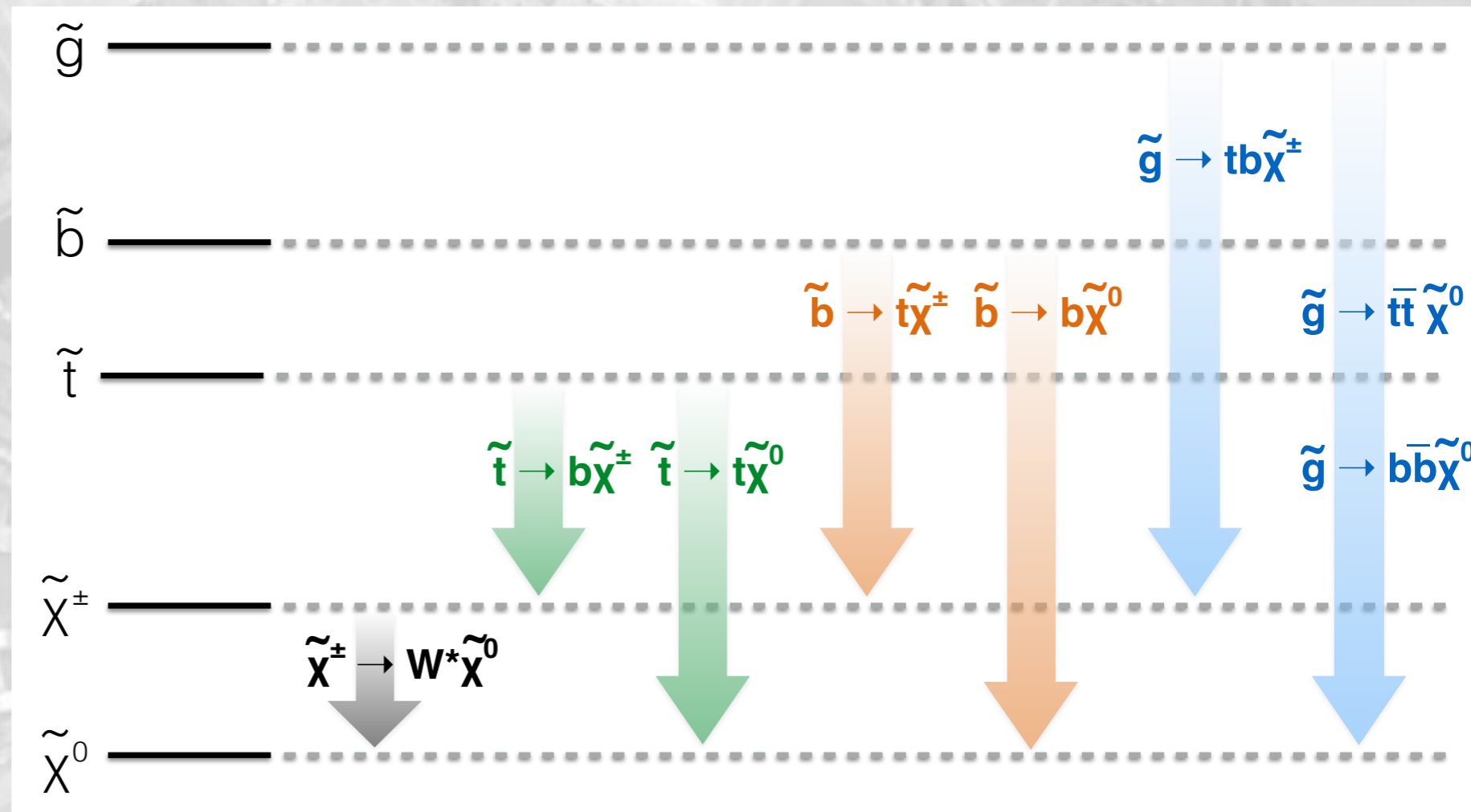
- this was somehow expected (e.g. EW precision, flavour)
- Now we know it as a fact



Run I forced us to look for “new” paradigms

NATURAL SUSY

- Multijet final states with many b quarks (4t, 3t1b, 2t2b, 1t3b, 4b) from gluinos
- High-pT leptons from W decays
- Same/opposite charge lepton pairs, with same or different flavour



- Rather than focusing on one signature, we decided to design an inclusive search
- Rather than focusing on the tail of some kinematic distribution, we decided to use a loose selection: more signal, but also more background

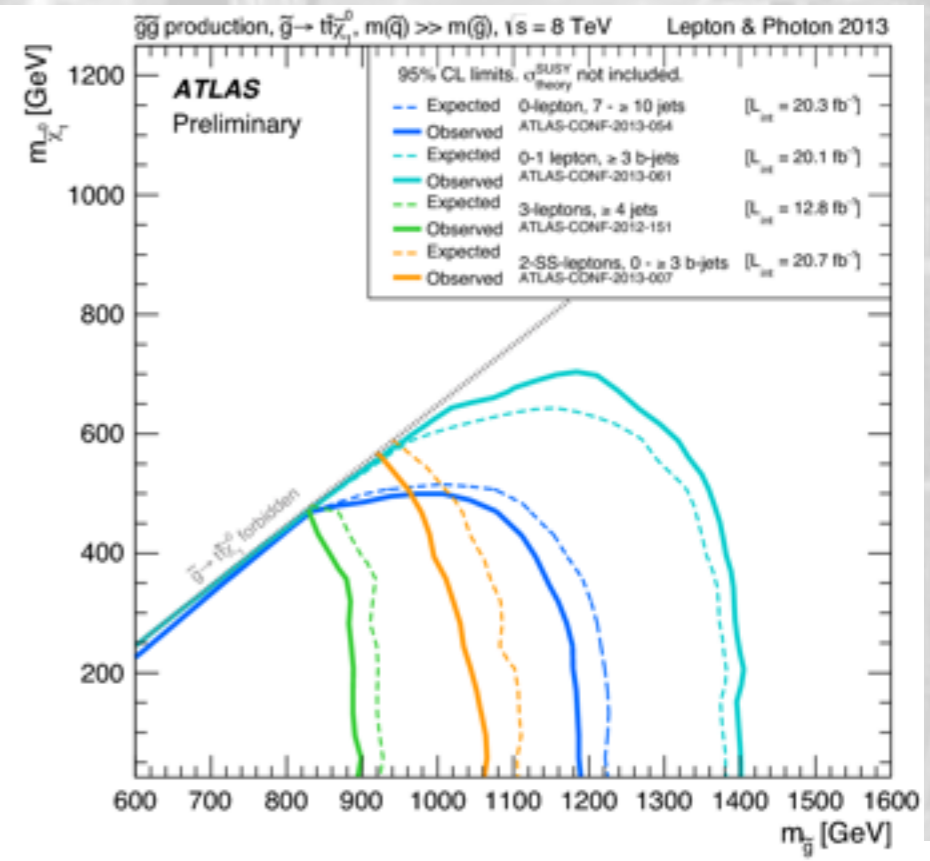
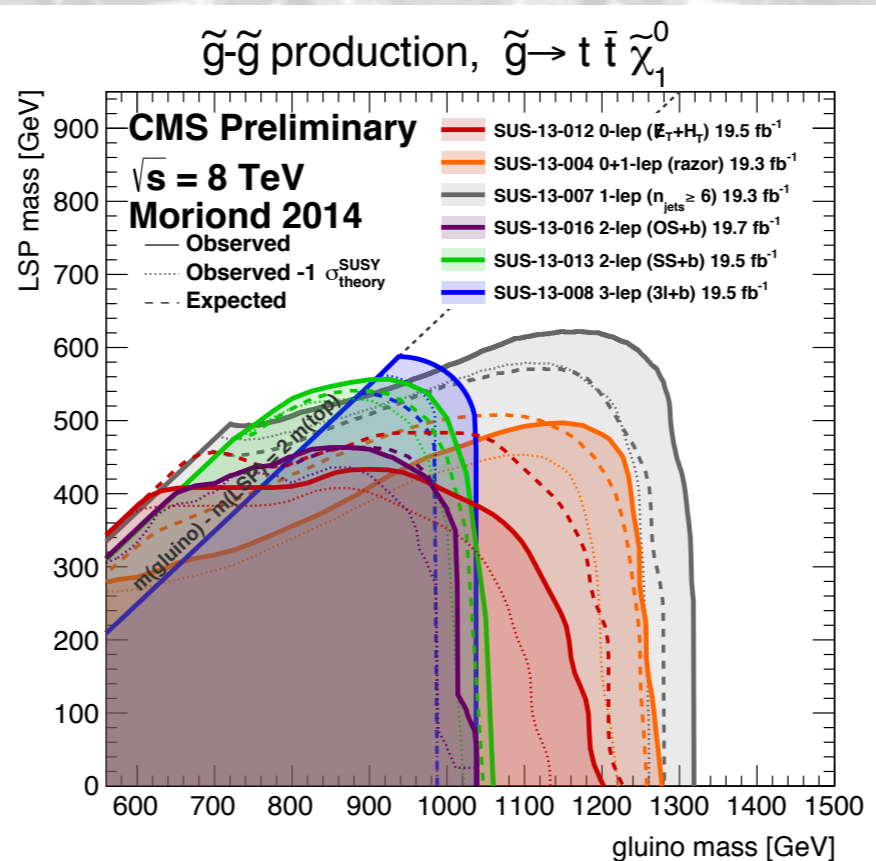
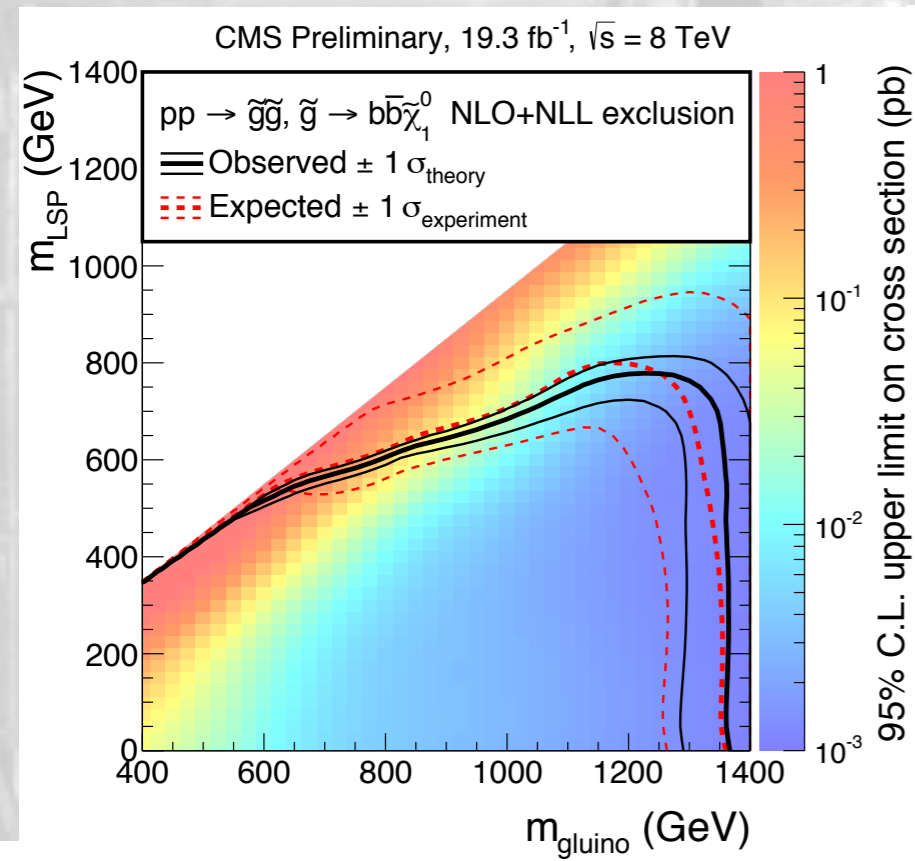
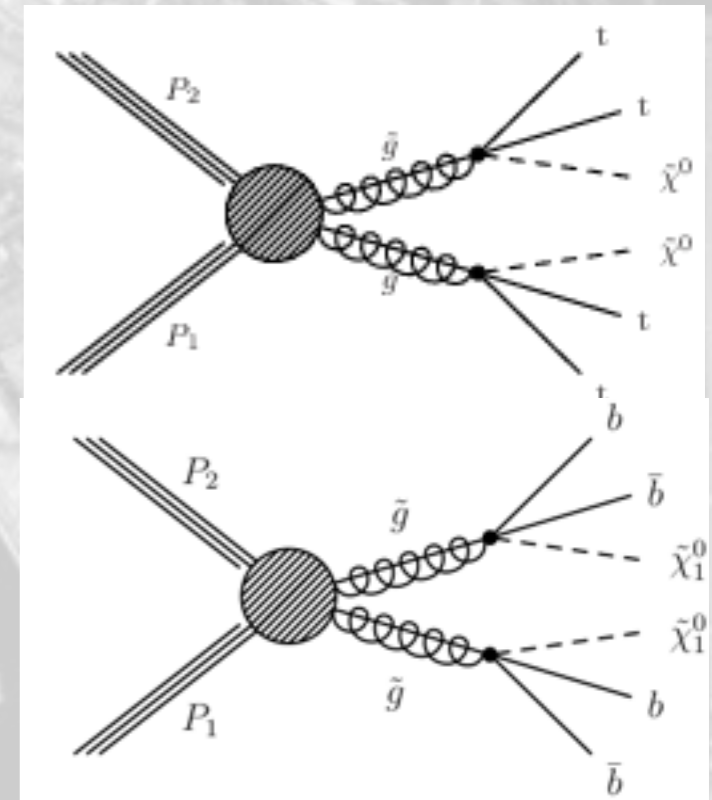
SEARCH FOR GLUINOS

Gluino is a model killer if kinematically accessible

- large xsec
- rich final states

Loose sensitivity if gluino (LSP) heavier than ~ 1300 (~ 600) GeV

As long as this is not the case, squarks produced in cascade are also largely excluded (limits robust vs intermediate squark masses)

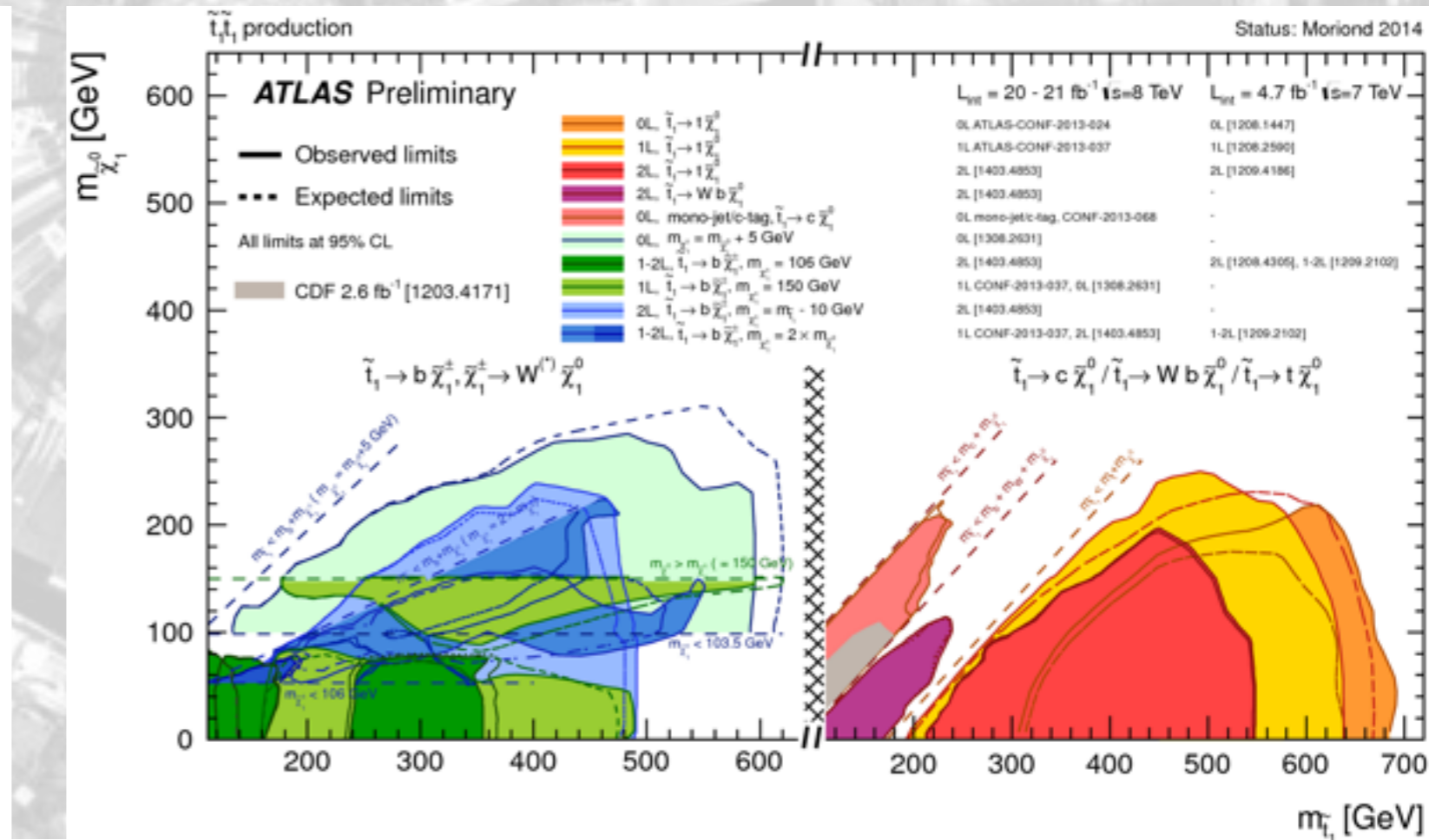
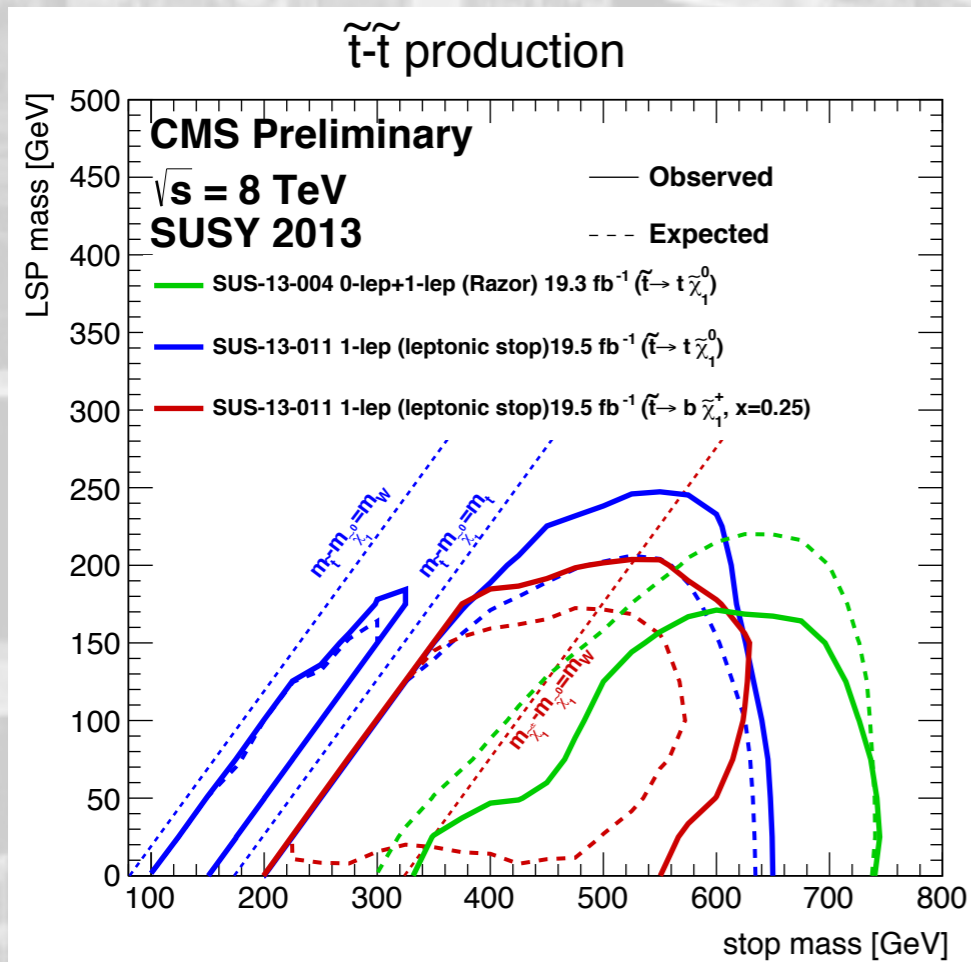
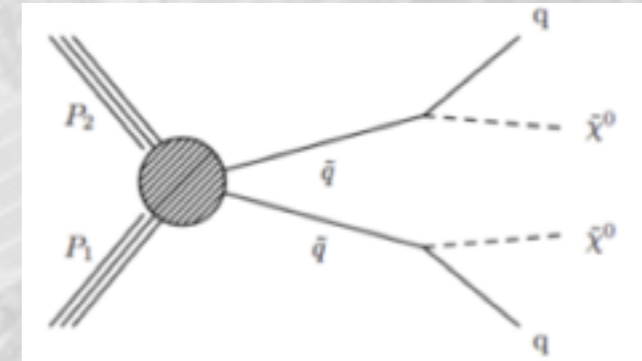


SEARCH FOR SQUARKS

Both inclusive and exclusive searches

sbottom and stop excluded up to $m \sim 750$ for large mass splits

reduced sensitivity for heavy LSP or for split $\sim m_{\text{top}}$

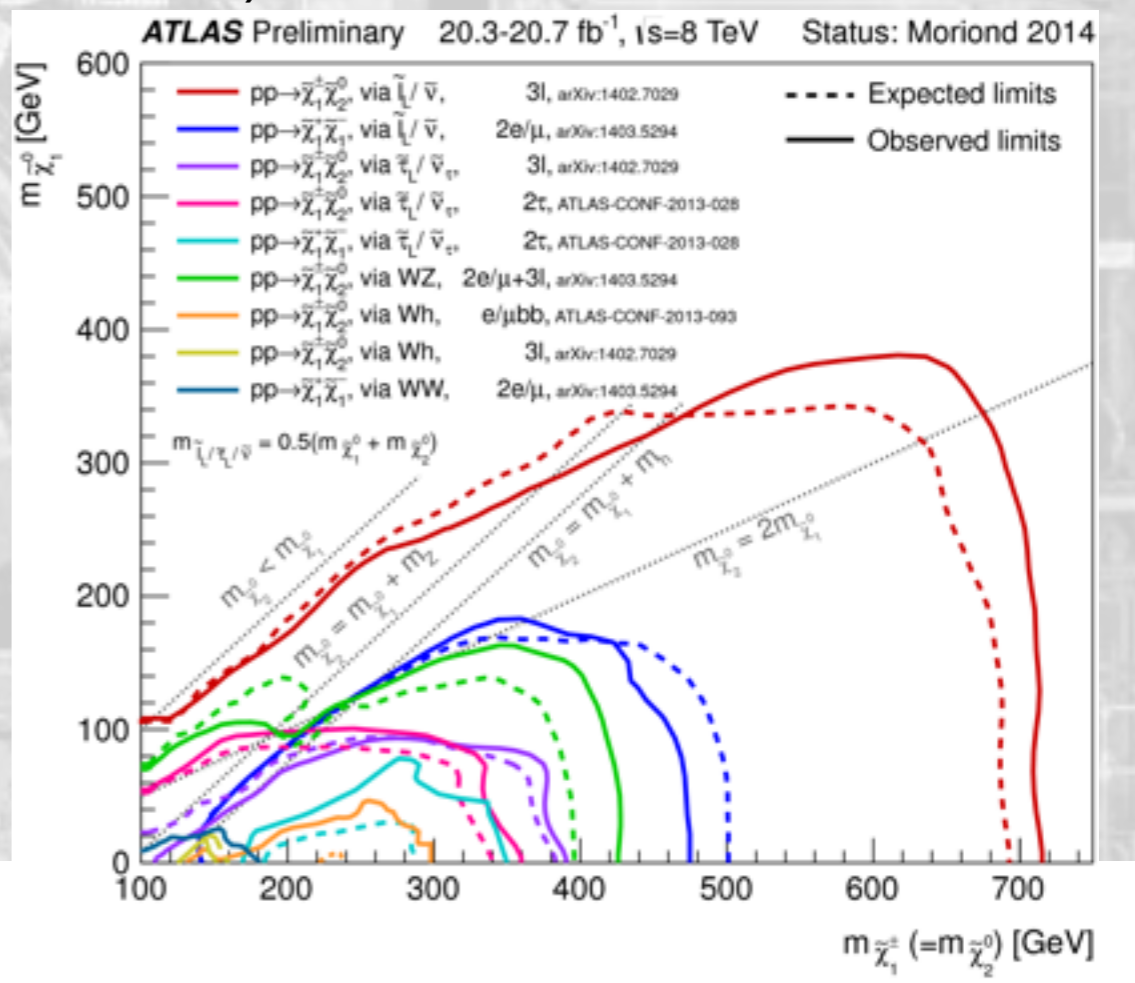
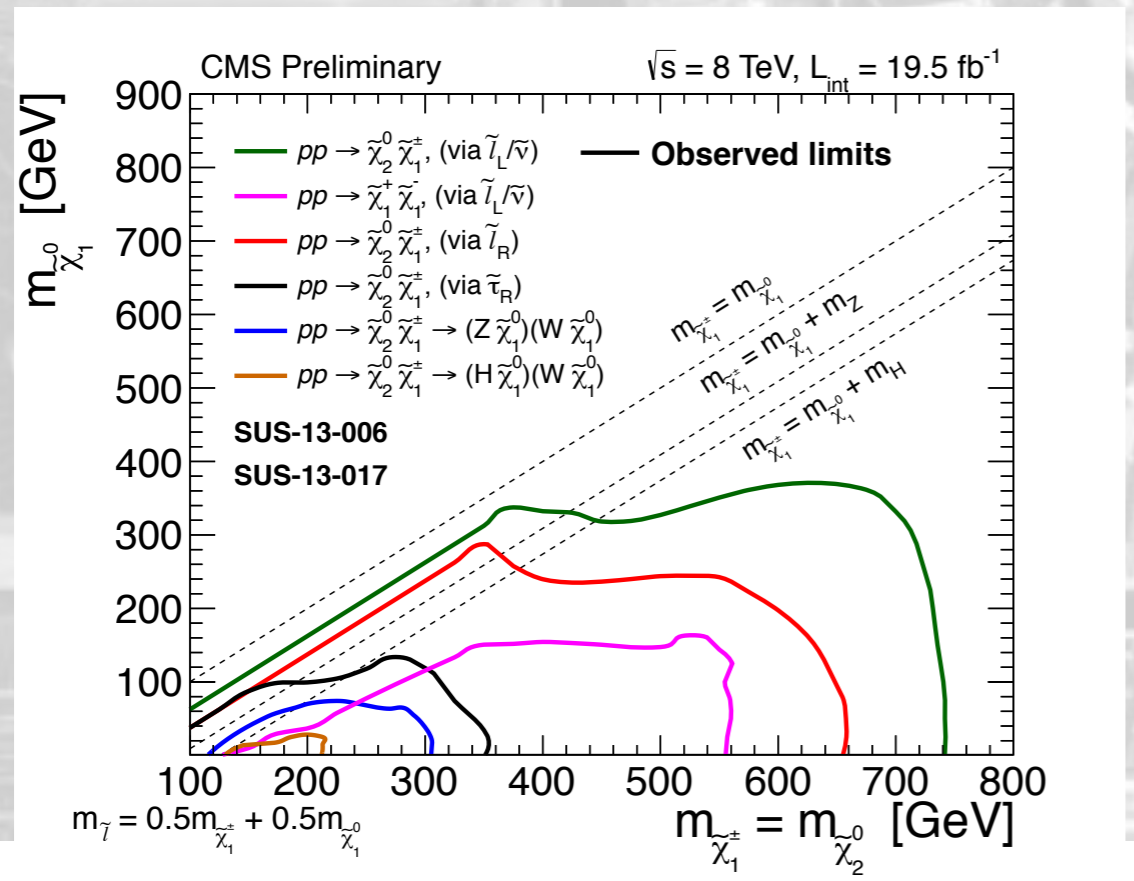
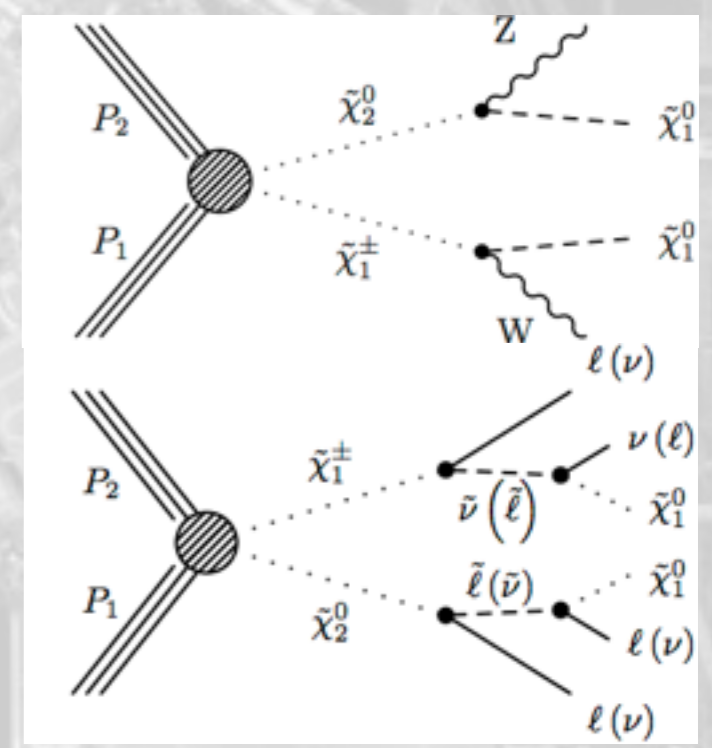


SEARCH FOR EWKINOS

Search for Ewkinos direct production more challenging


- smaller xsec
- limited kinematic handles (particularly at small splits)

Use combination of several final states, including Higgs bosons (now that we know the mass)



A FEW REMARKS

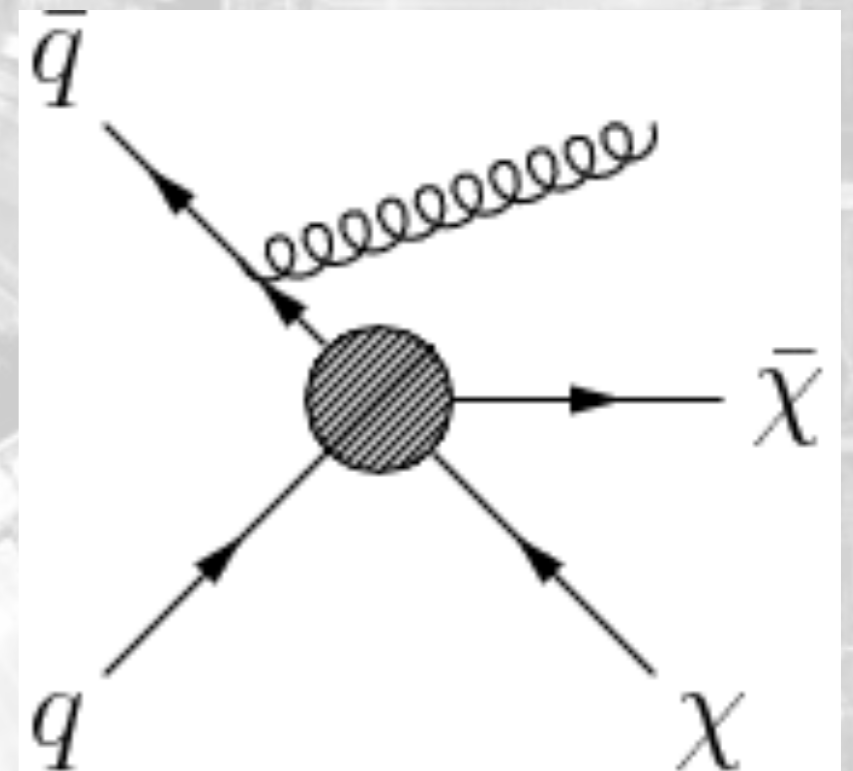
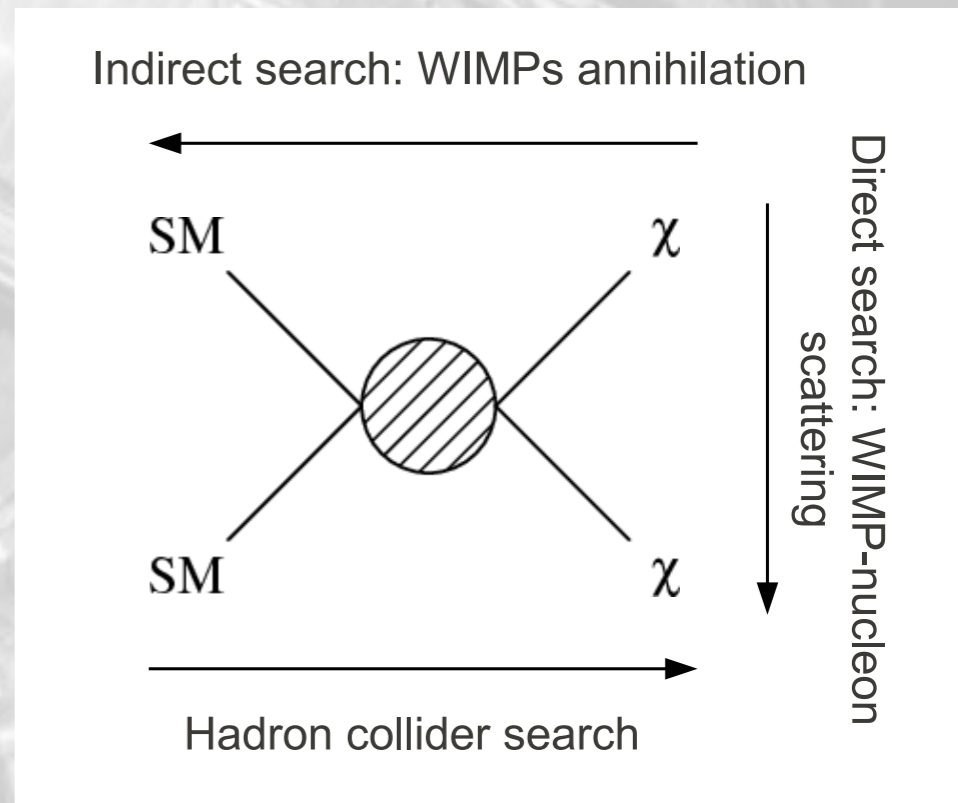
- Simplified models are very useful to design searches
- But they can be misleading if the results are generalised
- Typically
 - 100% BR is assumed for a given decay mode
 - A string constrain with this assumption can vanish for generic BRs
 - Two limits with 100% BR assumption can be contradicting and difficult to combine
 - The cross section is computing assuming other sparticles are decoupled (ignored many t-channel diagrams). This can underestimate/overestimate the cross section
- The SMS limits are ballpark right, but they cannot be taken literally
- In other words, we did not see any SUSY but the limits could be much more relaxed than what Simplified Models imply



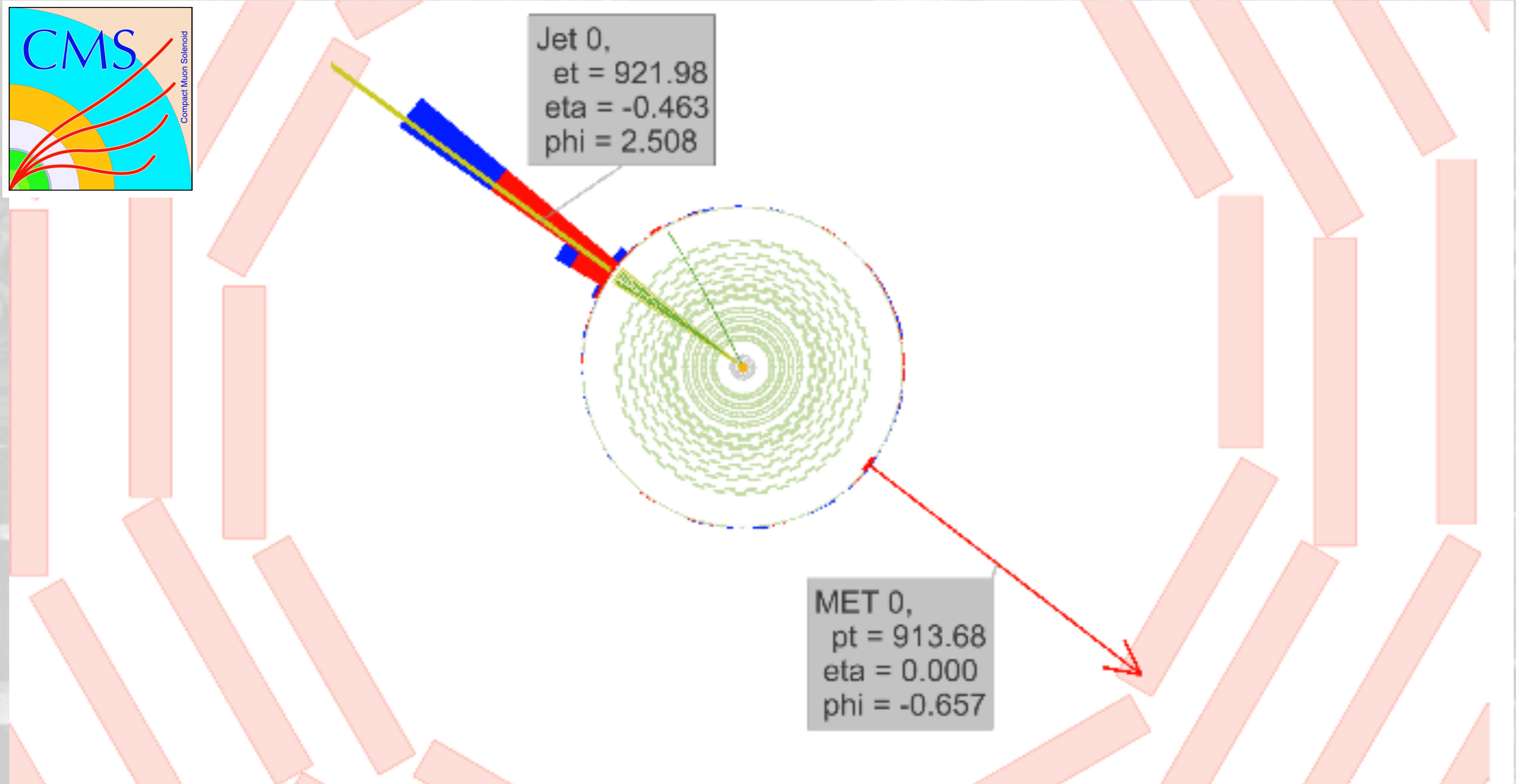
**Highlights from Run I:
search for DM production
(aka monojet)**

INVISIBLE AND MONOJET

- DM can be produced at the LHC with a similar process than scattering in underground experiments
- But DM is invisible with our detectors, so these events don't even pass the trigger
- We then exploit Initial State Radiation to look for 2DM+jet production
- These events look like a jet recoiling against nothing
- The main SM background (e.g., $Z(\text{nn})$ +jets) can be studied with data (e.g., $Z(\text{mm})$ +jets)

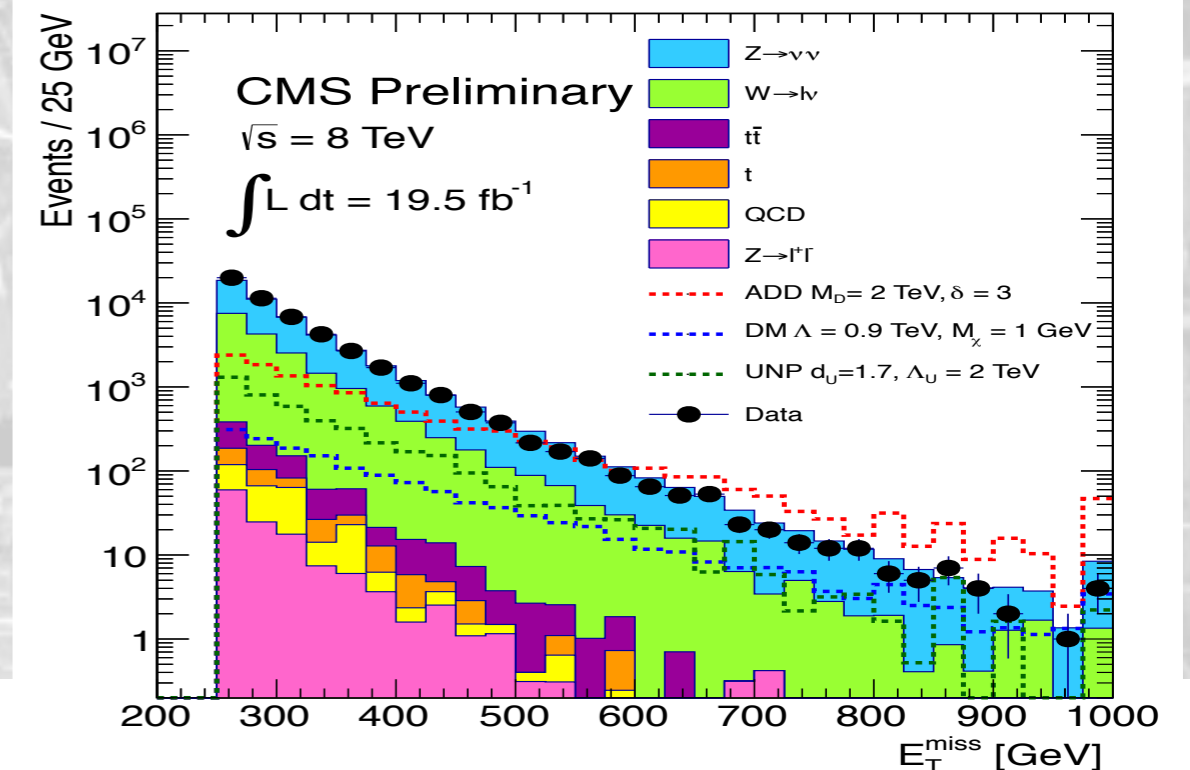
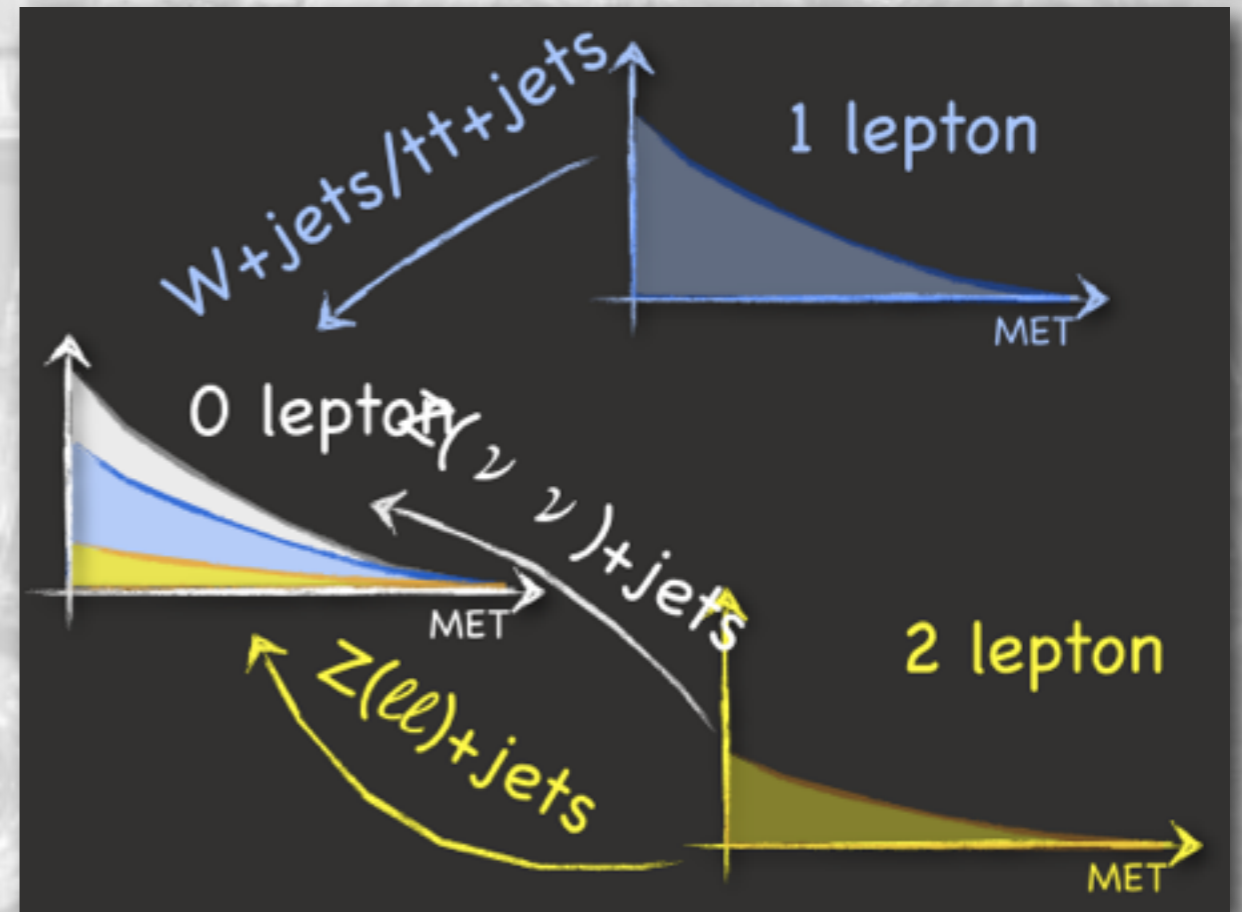


INVISIBLE AND MONOJET



SEARCH STRATEGY

- Reject the QCD background by kinematic cuts (large missing E_T , large jet p_T)
- Left with W/Z +jets and $t\bar{t}$ production
- Measure the background in control samples (1 lepton, 2 lepton, etc)
- Use the MC to scale the observed data yield to the 0lepton sample:
 $(1-\epsilon)/\epsilon$
- Compare prediction to observation

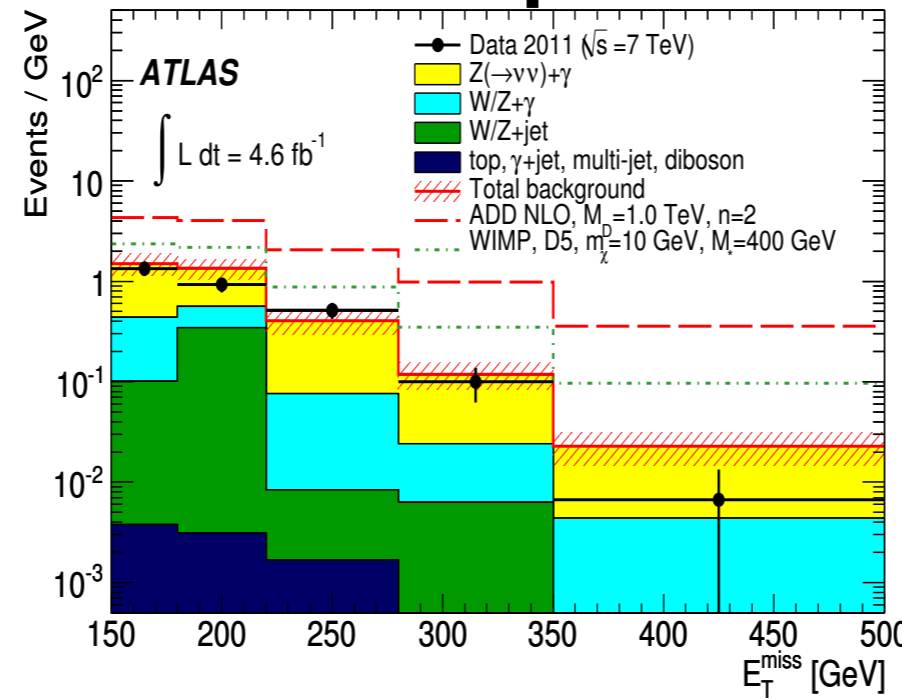


MORE THAN MONOJET

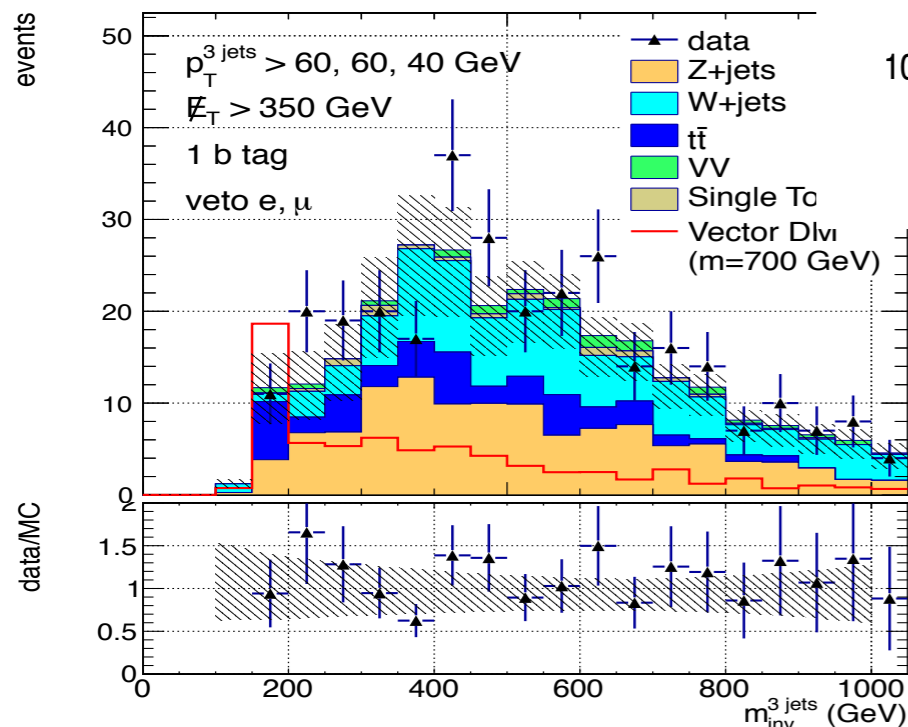
More rare objects than a jet could be radiated (i.e., less background)

Depending on the object, one could be sensitive to different models (e.g. DM coupling more to 3rd generation quarks, etc)

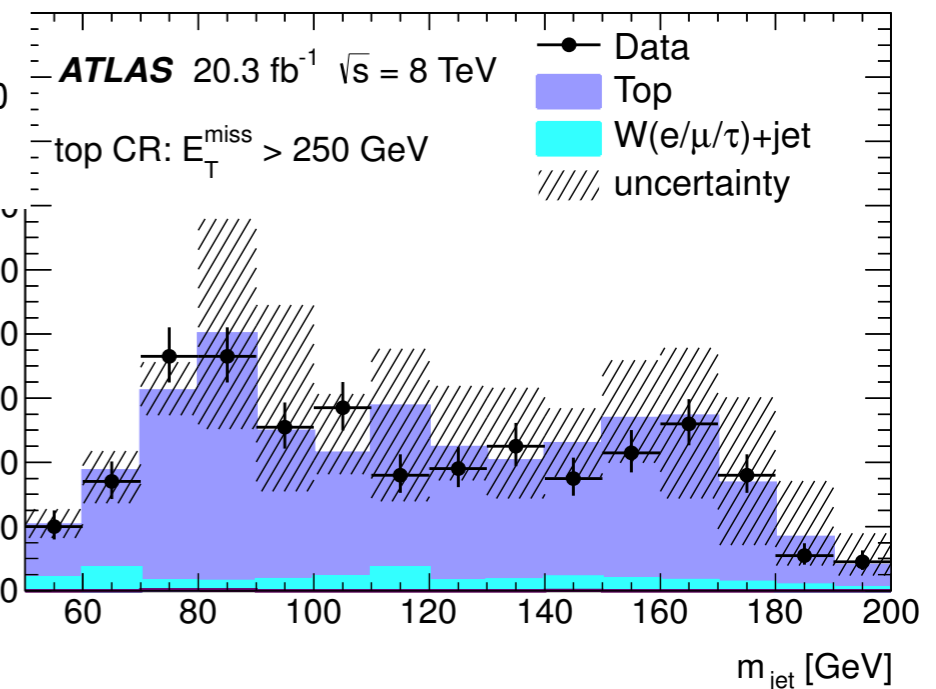
Monophoton



Monotop



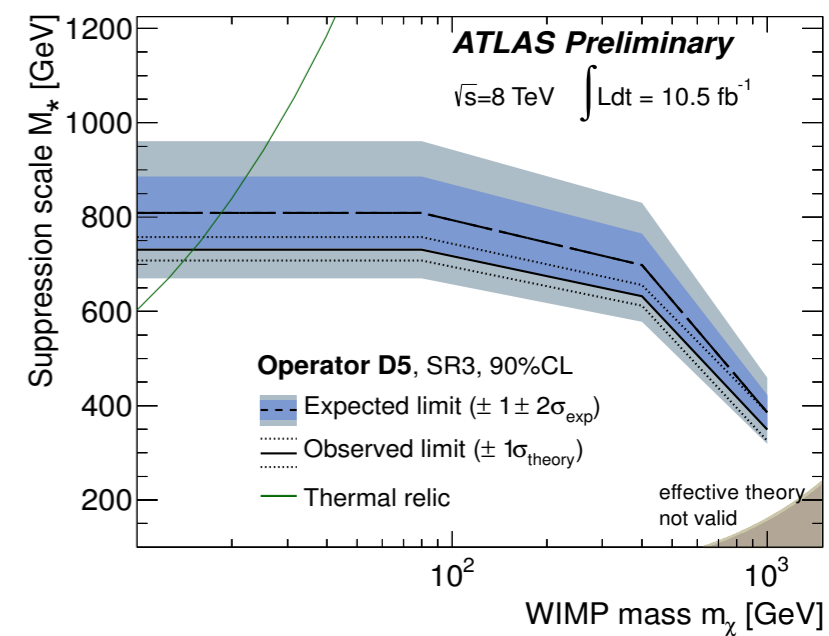
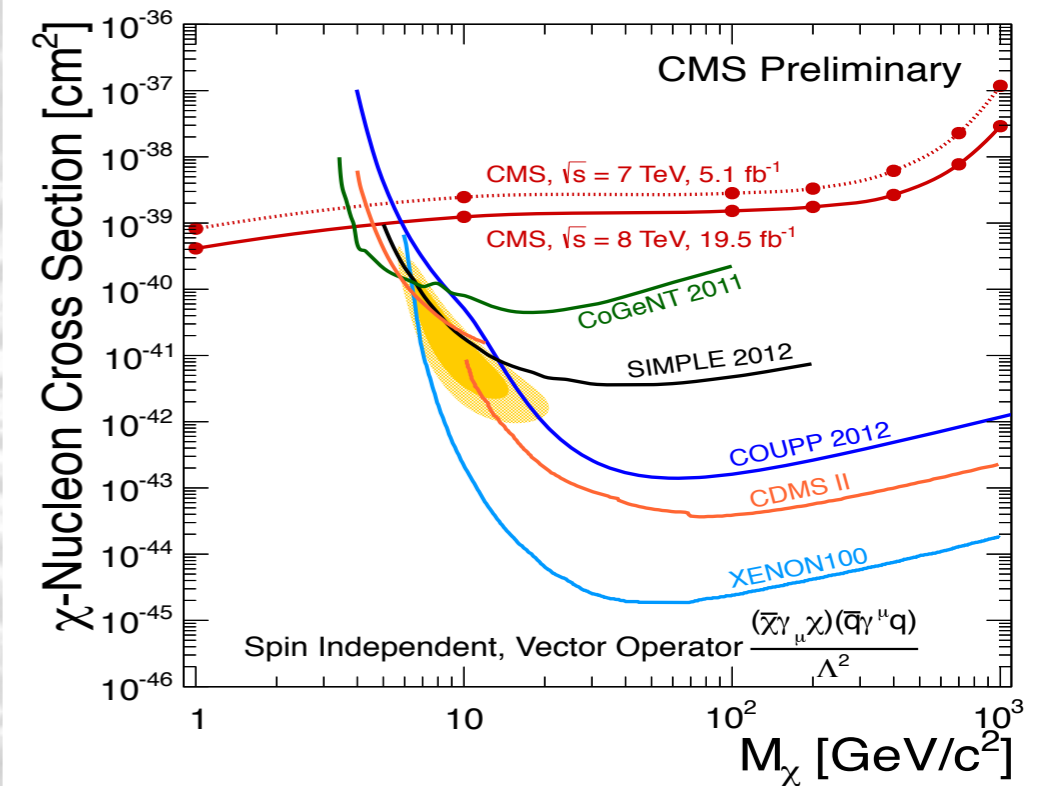
MonoZ/W



INTERPRETATION

- So far, result interpreted in terms of EFT, integrating out some heavy mediator
- Different kinds of mediator imply different operators
- This allows to project our results on the same plane as the underground experiment...
- ... but is an EFT with mediator masses < 1 TeV and collision energy 8 TeV reasonable?
- Work ongoing to adopt more specific models (e.g. SUSY simplified Models)

Name	Initial state	Type	Operator
D1	qq	scalar	$\frac{m_q}{M_*^3} \bar{\chi} \chi \bar{q} q$
D5	qq	vector	$\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$
D8	qq	axial-vector	$\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$
D9	qq	tensor	$\frac{1}{M_*^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$
D11	gg	scalar	$\frac{1}{4M_*^3} \bar{\chi} \chi \alpha_s (G_{\mu\nu}^a)^2$

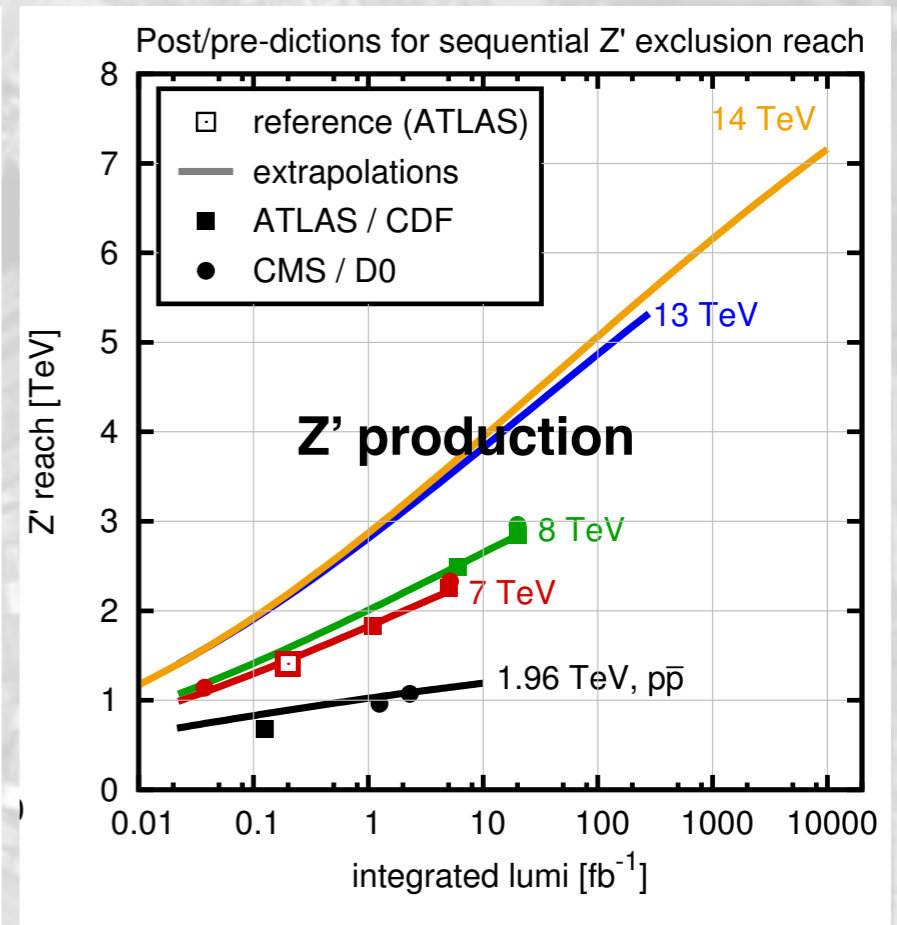
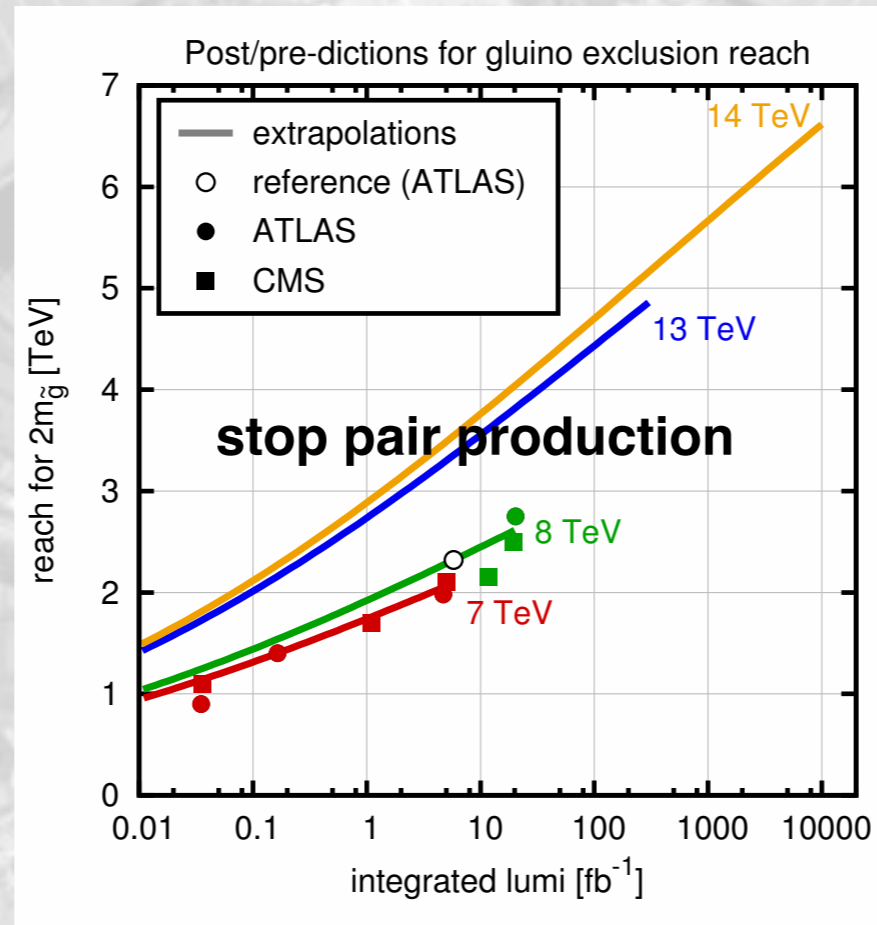
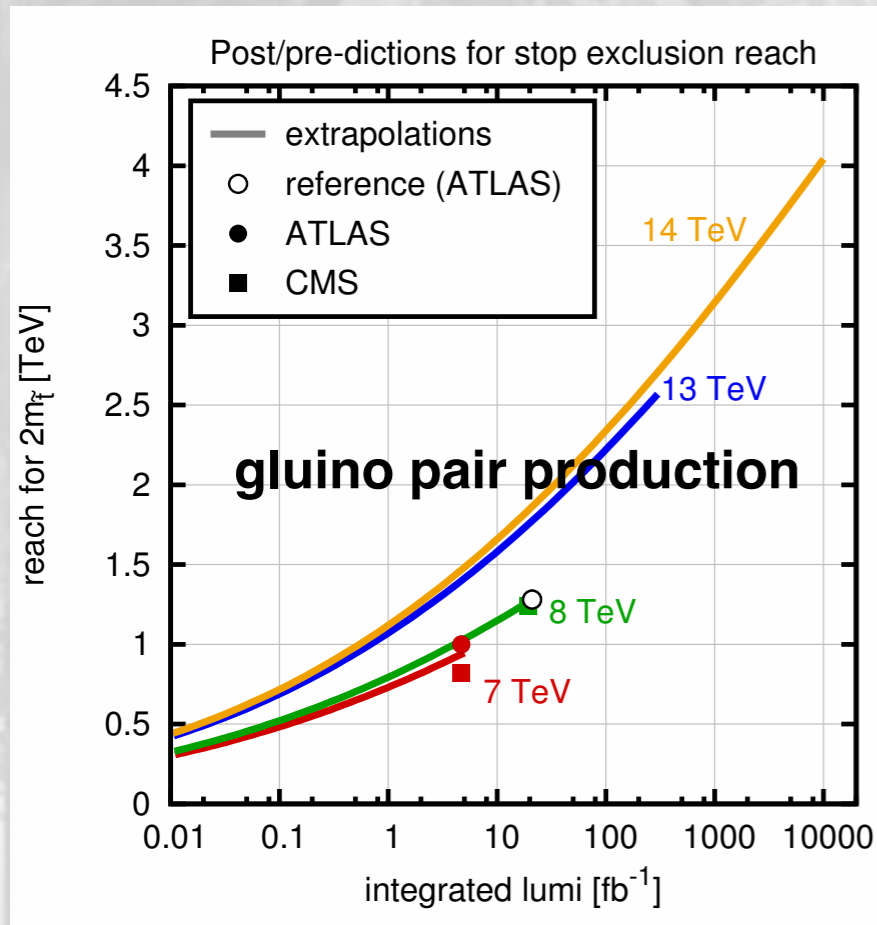




Perspectives for Run II

WHY GOING TO HIGHER ENERGY?

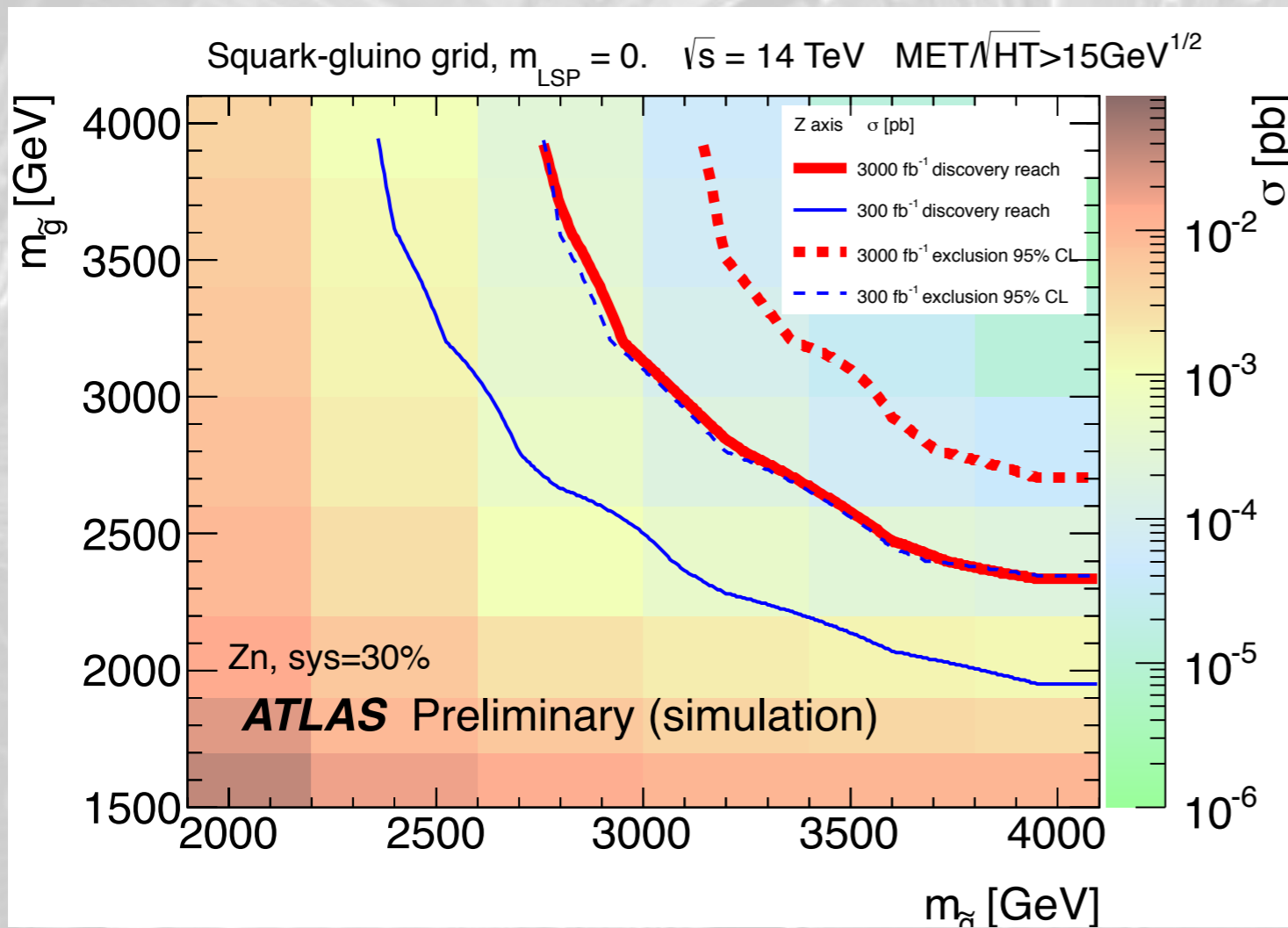
- Going higher in energy we gain faster than integrating more data



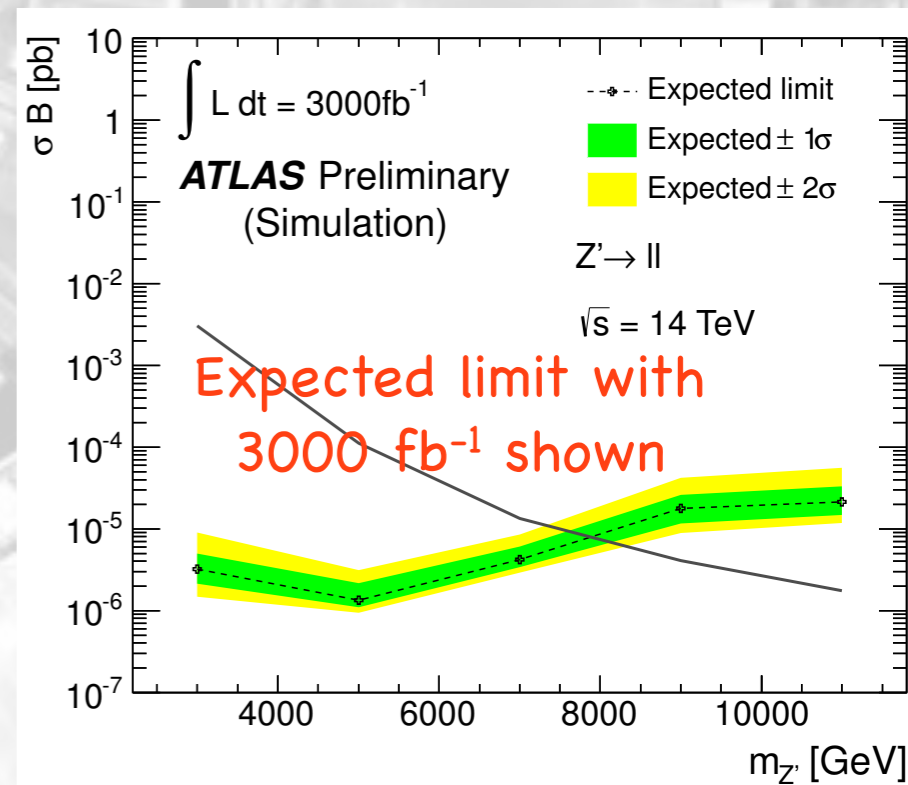
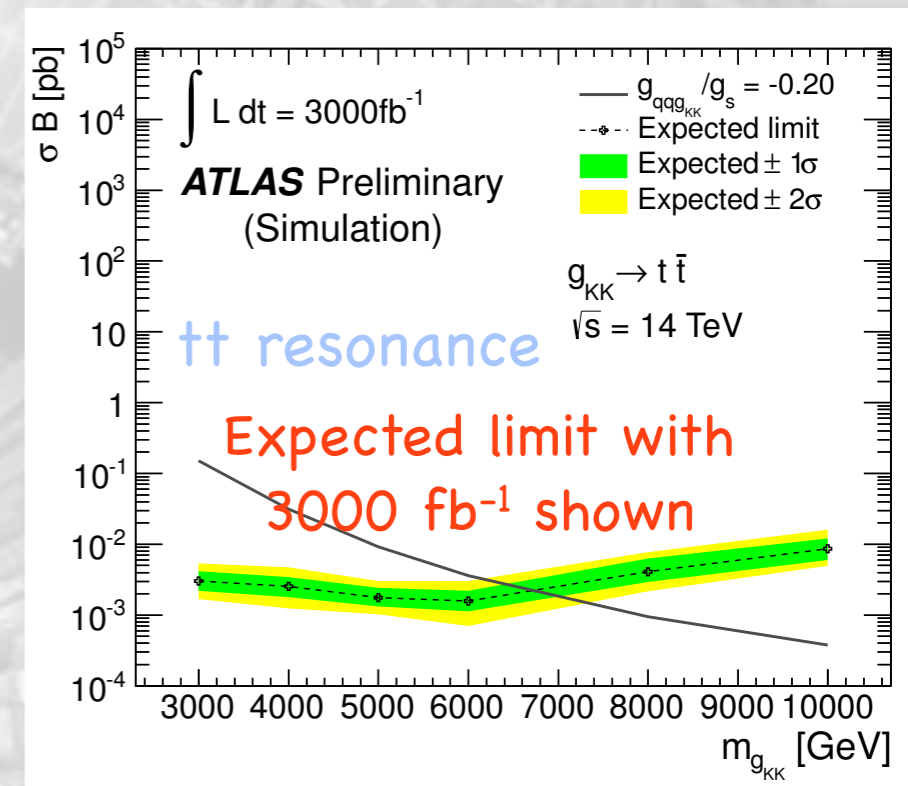
By G. Salam & A. Weiler (<http://cern.ch/collider-reach/>)

- As a thumb rule, increase energy by X is like increasing lumi by X²
- And there are things which could be just too heavy to be produced at low energy

PERSPECTIVES FOR 13 TEV

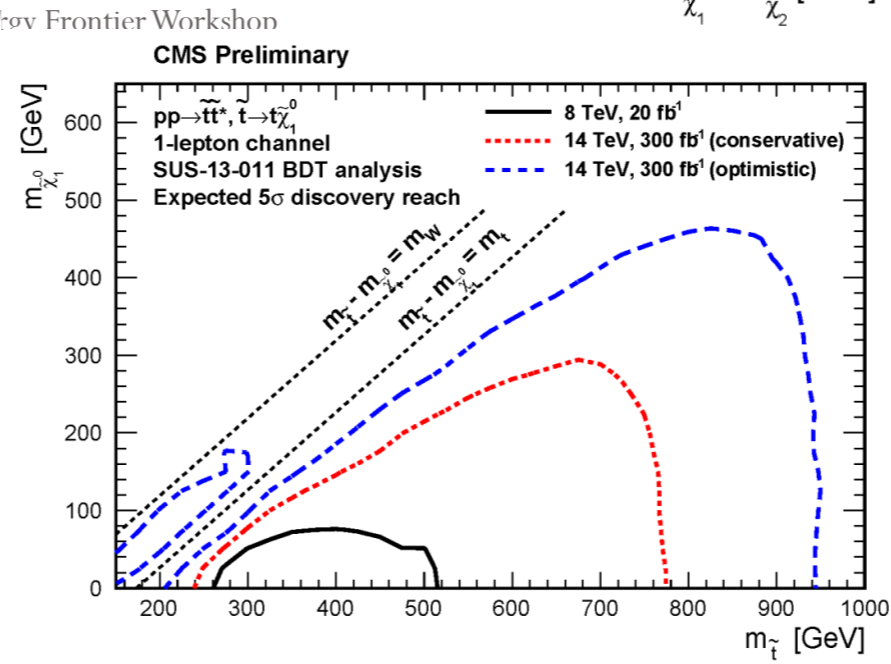
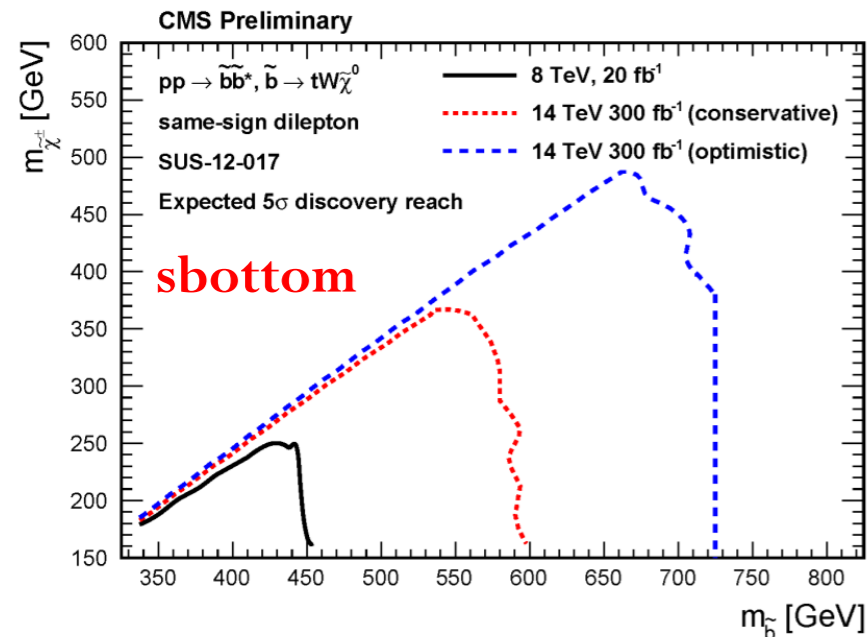
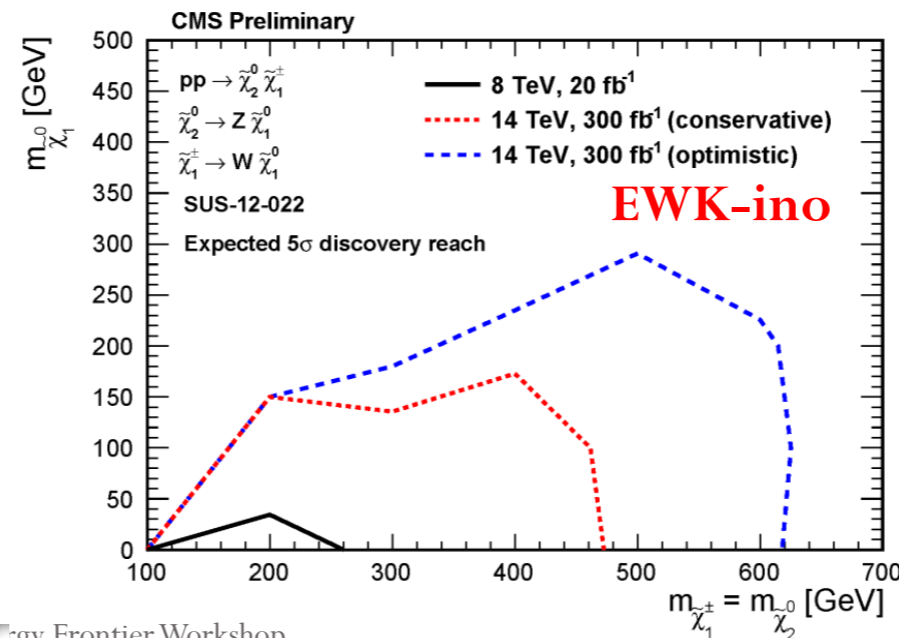
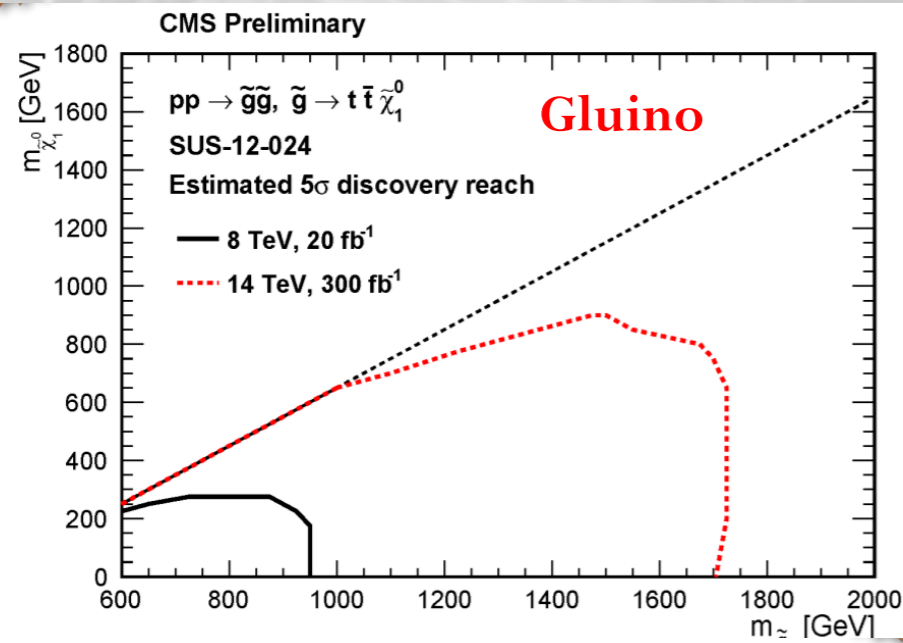


Similar extrapolations from CMS



PERSPECTIVES FOR 13 TeV

- Extrapolated with pessimistic (same systematics as now) and optimistic (scale systematics with luminosity) models
- The true value should be in the middle
- 5σ discovery reach shown



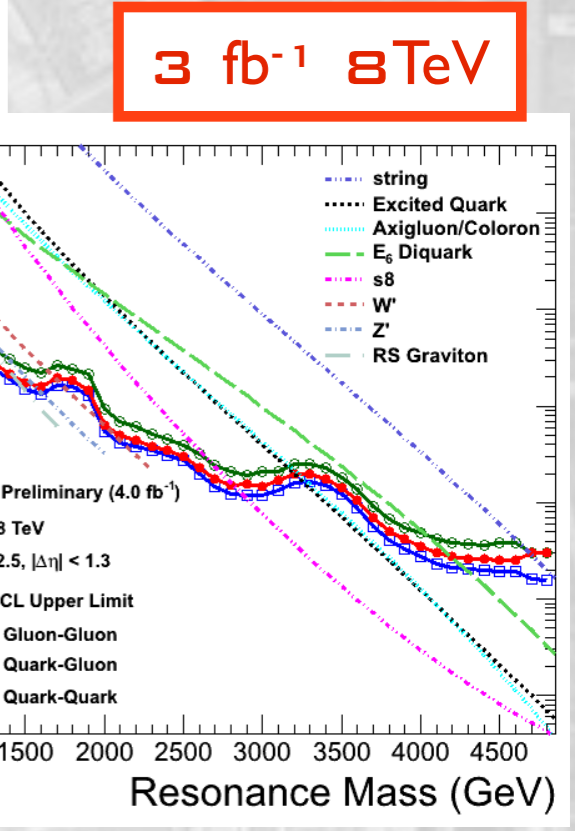
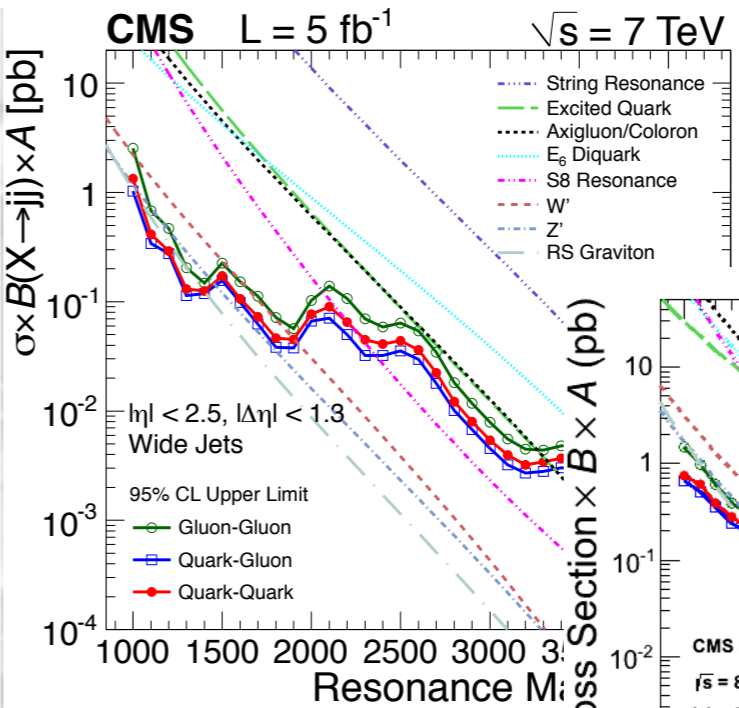
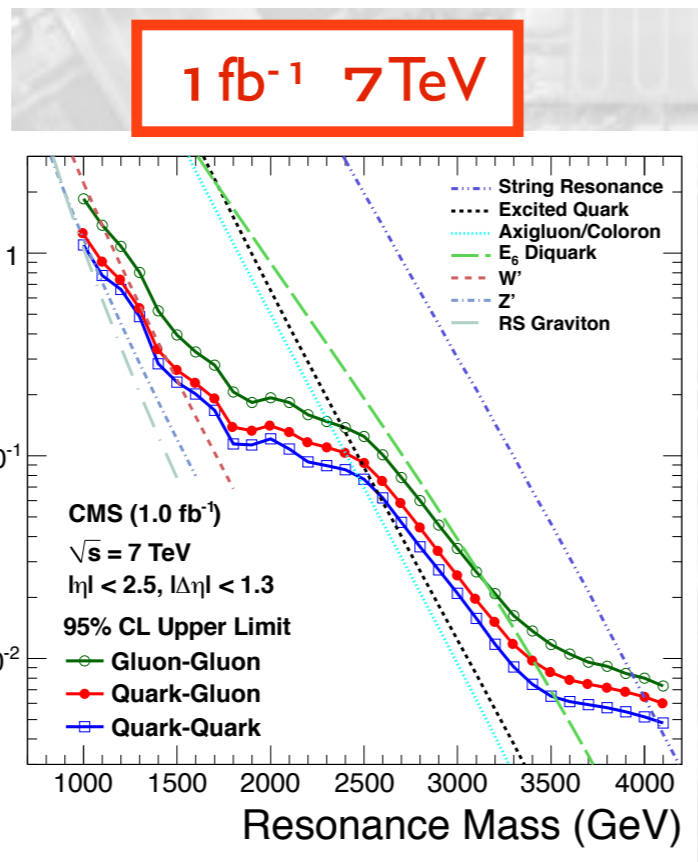
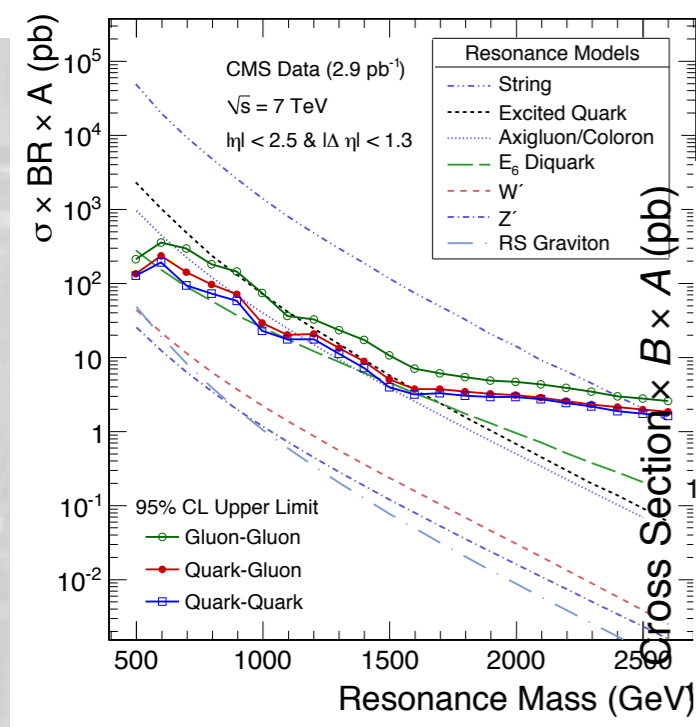
5σ discovery reach

- **Guino: up to 1.7 TeV**
- **Sbottom: ~600 – 700 GeV**
- **EWK-ino: ~500 – 600 GeV**

**SIMILAR
 EXTRAPOLATIONS
 FROM ATLAS**

3 pb⁻¹ 7 TeV

LOW MASS & SMALL XSEC



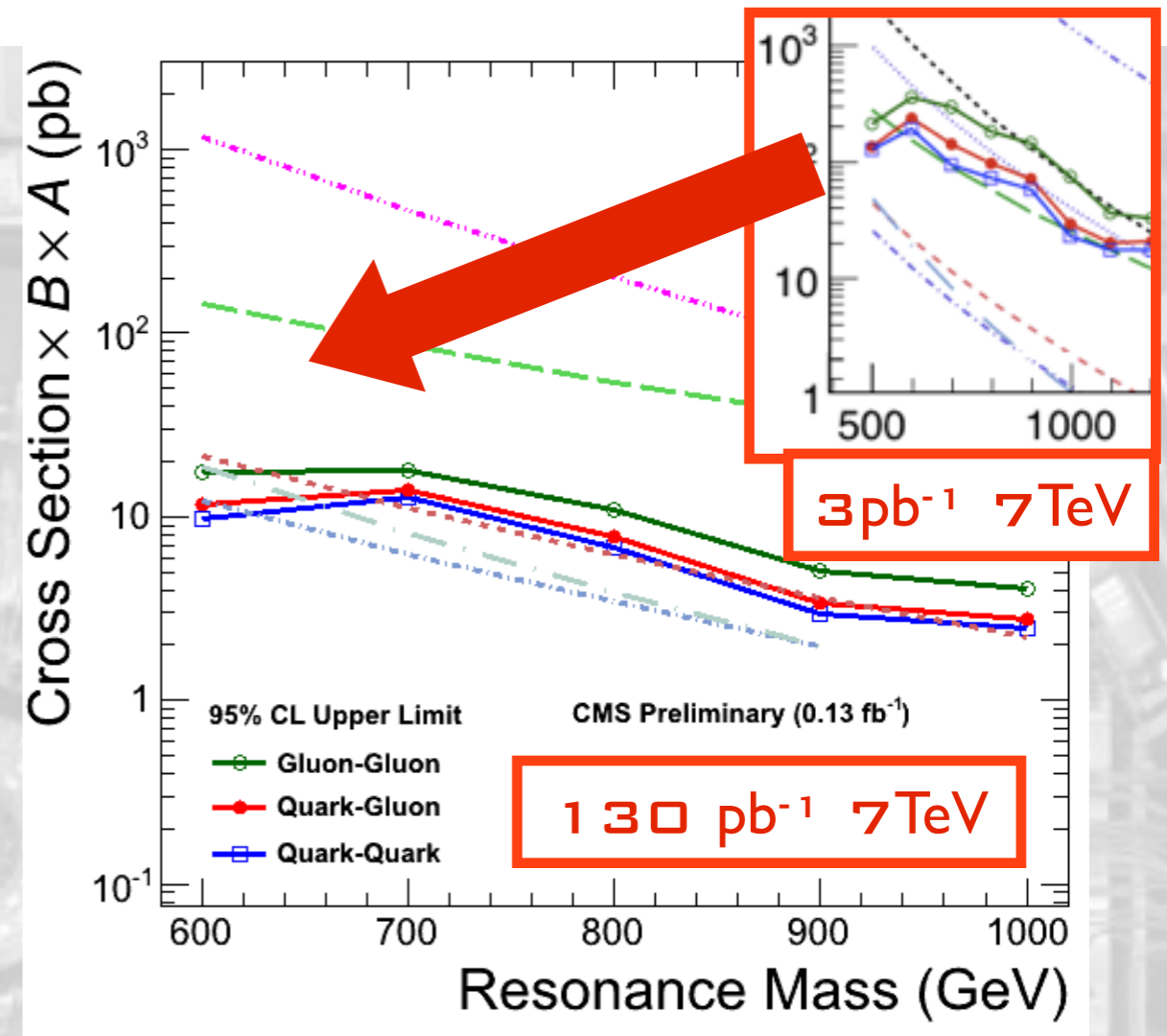
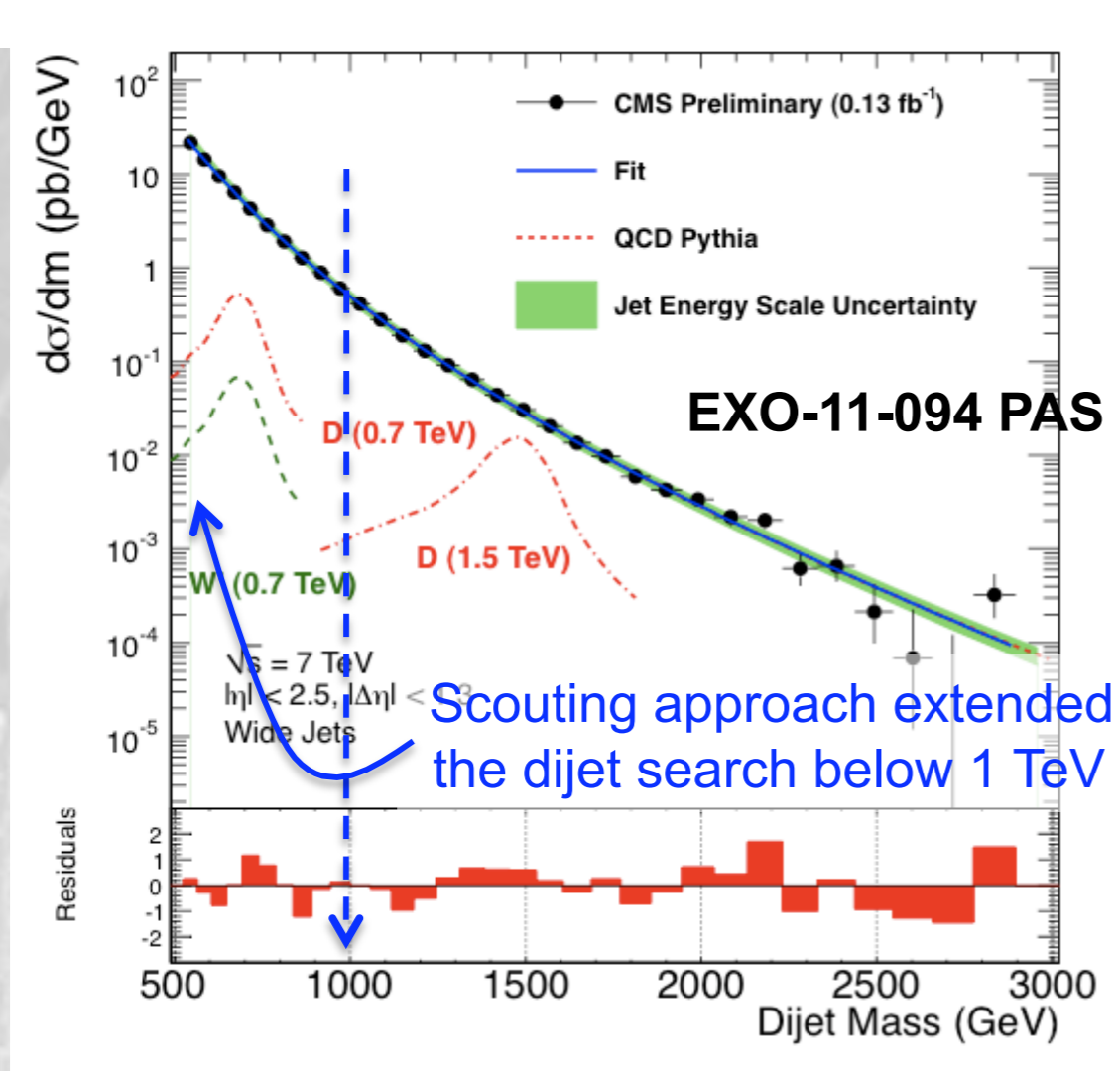
5 fb⁻¹ 7 TeV

E	L	MINMASS	σ MINMASS
7	0.003	400	~ 200
7	1	900	~ 1
7	5	1000	~ 1
8	3	1000	~ 0.6

- UNEXPLORED TERRITORY LEFT BEHIND WHEN THE LUMINOSITY INCREASED

- TRIGGER IMPROVEMENTS COMPENSATED THE (SLOWER) RATE INCREASE AFTER

THE DIJET DATA SCOUTING



- 16 HOUR RUN AT THE END OF 2011 RUN (7TeV)
- COLLECTED ~4 TIMES THE STATISTICS WE HAD IN 2010 (35 pb⁻¹) WITH EQUIVALENT TRIGGER
- IMPROVED THE LIMIT PUBLISHED IN 2010 BY ONE ORDER OF MAGNITUDE
- 18 fb⁻¹ RESULTS@8TeV TO BE RELEASED SOON

CONCLUSIONS

- The first LHC run was a big success, per se and in perspective
- We achieved more than expected (e.g., early Higgs discovery) despite the lower energy
- The achievement is a consequence of progresses in experimental techniques (e.g., jet tagging, kinematic variables, pileup suppression, bandwidth extension with scouting and delayed reprocessing...)
- We know what to do. Let's start the accelerator again, and let's see what comes out