

Improvements in the Jet Energy Resolution for $H \rightarrow b\bar{b}$ at D0

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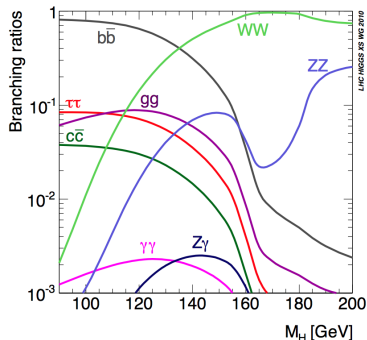


- Improvements to the Jet Energy Resolution
- Propagation to Higgs analysis: $WH \rightarrow \ell \nu b \bar{b}$
- Summary

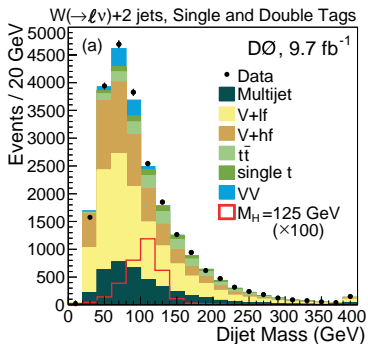




- Dominant decay mode for Higgs boson with $m_H = 125$ GeV is $H \rightarrow b\bar{b}$.



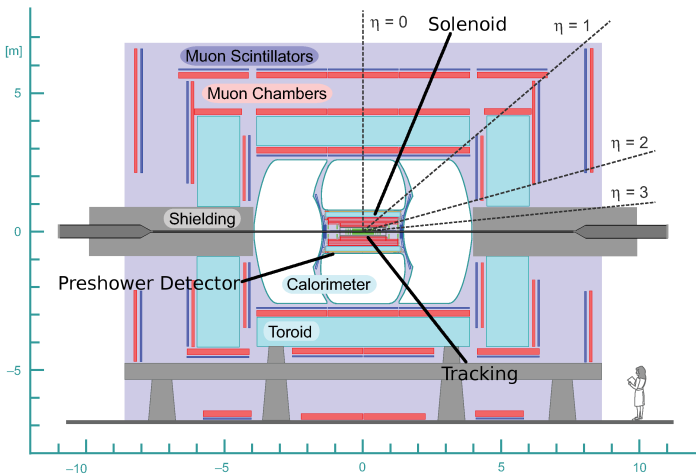
- Most of the sensitivity in $H \rightarrow b\bar{b}$ analyses comes from the dijet mass.



- Phys. Rev. D 88, 052008 (2013)

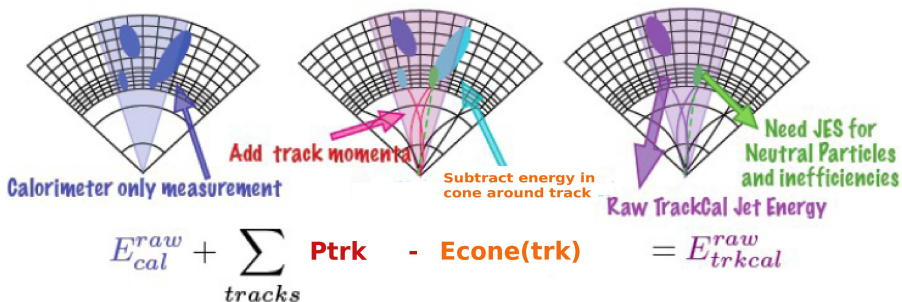


The D0 Detector





- At low energies, tracker has better resolution than calorimeter, so we can try to use the track momentum measurement instead of the calorimeter energy measurement.
- $E_{trkcal\ jet} = \text{raw jet } E + \sum_{trks}(P_{trk} - E_{cone}(trk))$

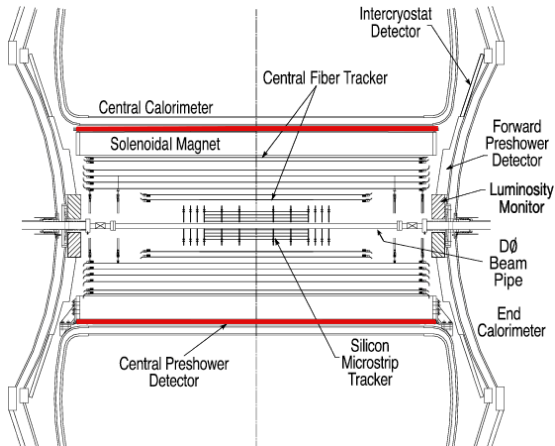




Central Preshower Detector Energy



- Jets are reconstructed using only calorimeter information
- Can add in energy from central preshower (CPS) detector
 - This corrects for energy lost in solenoid and upstream material.

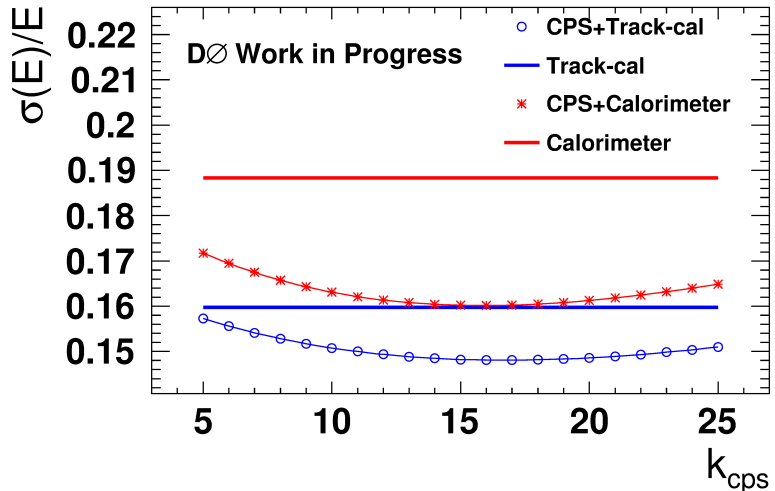


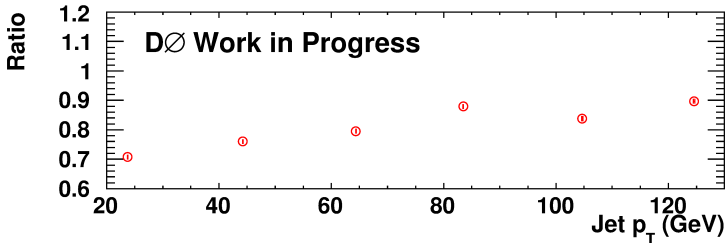
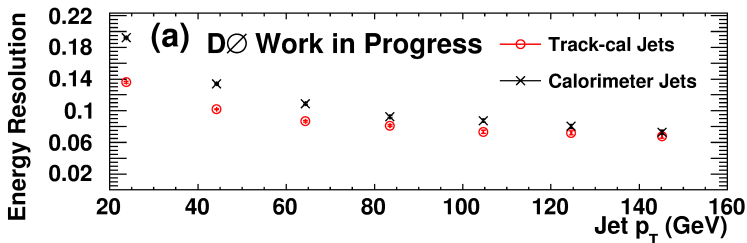


Adding Central Preshower Energy to TrackCal Jets

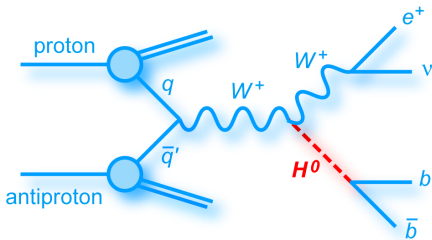


- Define our corrected energy as $E = E_{trkCal} + K_{CPS} * E_{CPS}$.
- K_{CPS} is chosen to minimize resolution ($K_{CPS} = 18$).

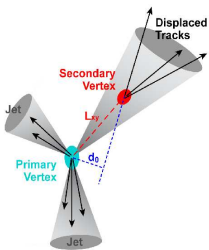




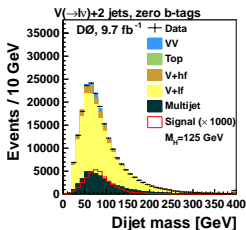
(a) $\sigma(E)/E, |\eta_{jet}| < 0.4$



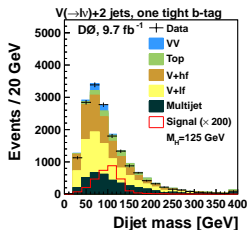
- One charged lepton (e or μ) with $p_T > 15$ GeV
- $\cancel{E}_T > 15$ (e), 20 (μ) GeV
- Jet $p_T > 20$ GeV, $|\eta| < 2.5$.



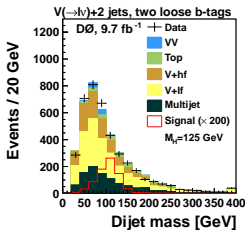
(a) B Hadron Decay



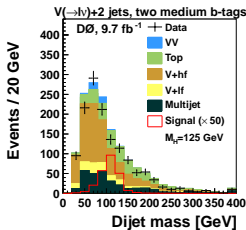
(b) 0 b-tags



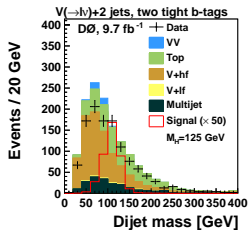
(c) 1 tight b-tag



(d) 2 loose b-tags



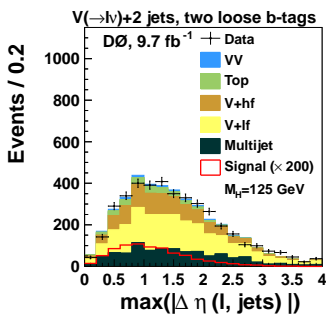
(e) 2 medium b-tags



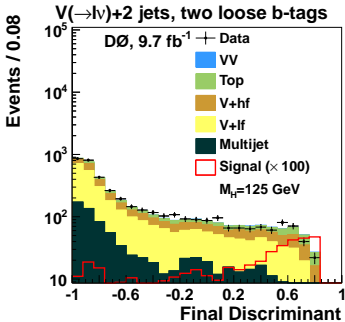
(f) 2 tight b-tags



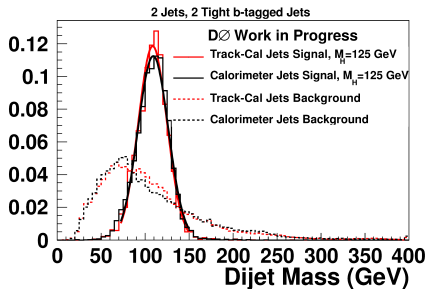
- Use a boosted decision tree to distinguish between signal and background
- Separate boosted decision tree trained for each jet, b-tag and lepton multiplicity to distinguish signal from total background



(a) Max $|\Delta\eta(\ell, jets)|$

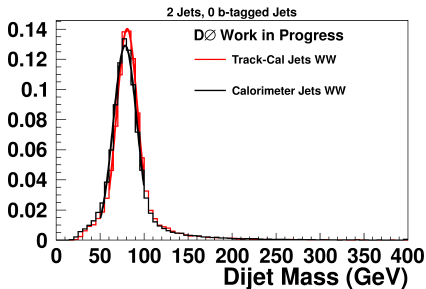


(b) BDT Output



(a) 2 jets, 2 tight b-tags

- Track-cal jets width: 16.4
- Calorimeter jets width: 17.6



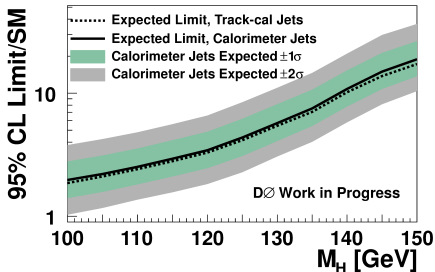
(b) 2 jets, 0 b-tags

- Track-cal jets width: 12.0
- Calorimeter jets width: 12.9

7% improvement in dijet mass resolution for $H \rightarrow b\bar{b}$ and $W \rightarrow jj$ events!

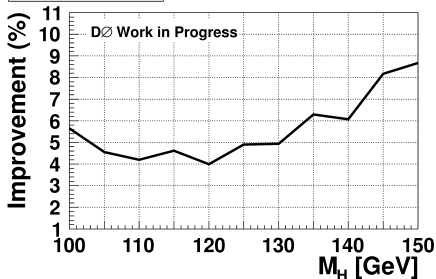


Comparison of Expected limits: $WH \rightarrow l\nu bb$



(a) Expected 95% C.L. limits

Percent Improvement



(b) Percent improvement

~5% improvement in expected limits!



- Implemented a correction to the jet energy for charged particles within a jet
- Propagated this correction to D0 WH analysis
 - Improved jet energy resolution by $\sim 20\%$,
 - dijet mass resolution by $\sim 7\%$,
 - and expected upper limit on $WH \rightarrow \ell\nu b\bar{b}$ cross section by $\sim 5\%$.



Thank you for listening!



Backup





$$E_{track-cal} = E_{raw}^{calo} + \sum_{tracks} (P_{track} - E_{track}^{\Delta R=0.15}) + 18E_{CPS}, \text{ if } |\eta_{jet}| < 1.2,$$

$$E_{track-cal} = E_{raw}^{calo} + 13E_{CPS}, \text{ if } 1.2 < |\eta_{jet}| < 1.6,$$

$$E_{track-cal} = E_{raw}^{calo}, \text{ if } |\eta_{jet}| > 1.6.$$



- Need to choose a cone size so that we subtract the energy coming from the track.

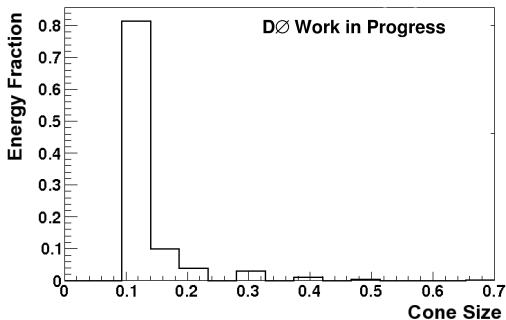


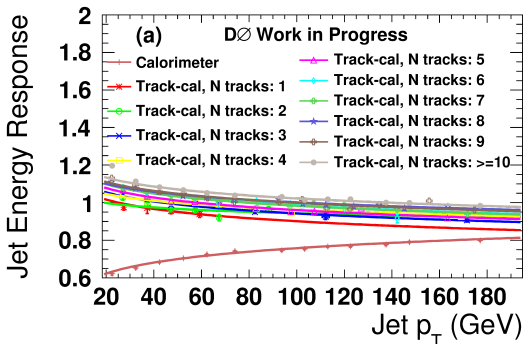
Figure: Fraction of energy deposited vs cone size



- Track selection was optimized to select high resolution tracks, and cut down on fake tracks
 - —track Z - PV Z— < 0.5 cm
 - Radial DCA (PV vs. Track) < 0.05 cm
 - Fake Track Killer quality > 0
 - Nsmt hits > 0
 - significance of curvature > 5
 - $$\text{Sig curv} = \frac{P_{Ttrk}}{\sqrt{\sigma_{trk} P_T}}$$



- Calculate response as $p_T^j / p_T^{j \text{ truth}}$ in jet p_T bins using γ +jet events.
- Plot mean as a function of true jet p_T to obtain a response function.



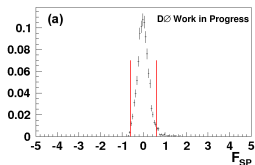
$$(a) R = p_T^j / p_T^{j \text{ truth}}, |\eta_{jet}| < 0.4$$



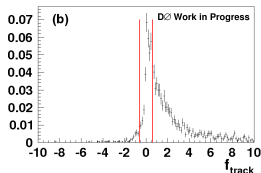
Some things to be careful about



- Need to be careful not to subtract out too much energy (for example if a neutral particle deposits energy near the track)
 - Did a study in single pion events to have some sort of metric of when too much energy is too much:
 - Considered $\sigma = (E_{cone} - P_{trk} * R(P_{trk})) / P_{trk}$,
 - $R(P_{trk}) =$ single pion response.
 - σ parameterized as a function of $P_{trk} * R(P_{trk})$, and jet η .
- Only apply algorithm if $|E_{cone} - P_{trk} * R(P_{trk})| / P_{trk} < 2\sigma$



(a) Single Pions, 5 GeV



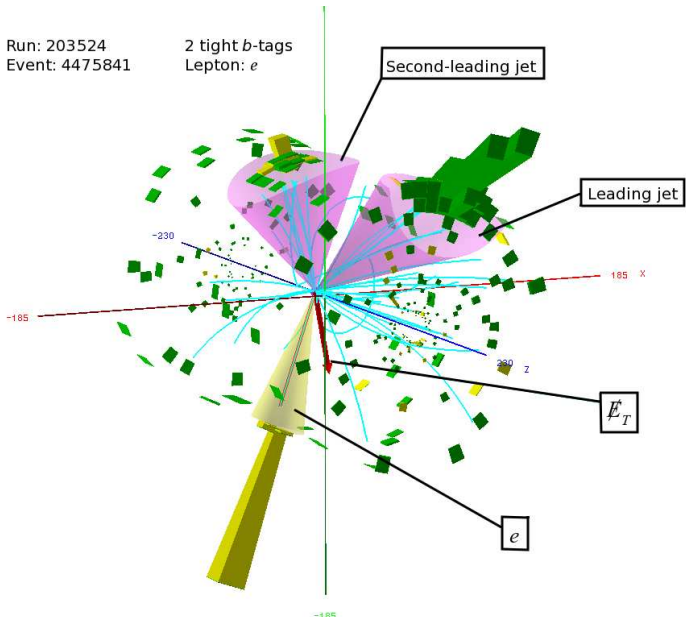
(b) photon+jet, $4.5 < P_{trk} < 5$ GeV





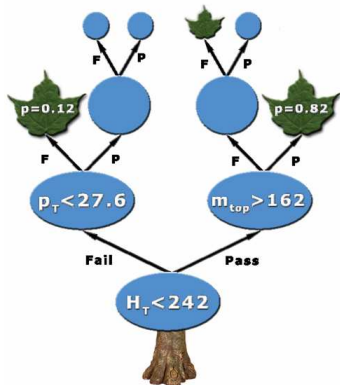
Run: 203524
Event: 4475841

2 tight b -tags
Lepton: e





- Start with a list of well modelled variables that have some separation between signal and at least one background.
- Select best variable, and best cut on that variable, to maximally separate signal and background.
- Repeat for signal and background separated samples
- Keep repeating until a state of pure background or signal is reached, or until a minimum number of events is reached.





- Boosting retrains tree to correct for events that were misclassified in the previous tree.
- This builds up a forest of trees which and returns a weighted response:
- $w_i = \frac{1}{\sum_{j=1}^{N_{trees}} \alpha_j} \log(\alpha_i) p_i,$
 - α is the tree weight, p is the node purity.



MVA _{MJ} Input Variables	Description
η^ν	Pseudorapidity of the missing E_T vector
E_T^{Sig}	E_T significance, a measure of the consistency of the observed E_T with respect to zero E_T , accounting for the uncertainty in the calorimeter objects that contribute to E_T
$\Delta\eta(\ell, \nu)$	Separation in η between the lepton and the reconstructed neutrino, $ \eta^\ell - \eta^\nu $
$T_{W \rightarrow \ell\nu}$	Twist of the $\ell\nu$ system = $\arctan(\Delta\phi(\ell, \nu)/\Delta\eta(\ell, \nu))$
$\cos\theta(\ell)_{\ell\nu CM}$	Cosine of the angle between the charged lepton and the proton beam axis in the CM of $\ell\nu$ system
$V(j_{12})$	Velocity of the dijet system
m^{Asym}	Mass asymmetry between $\ell\nu$ system and the dijet system: $(M_{\ell\nu} - m_{j_{12}})/(M_{\ell\nu} + m_{j_{12}})$
C	Centrality is $(\sum_i p_T^i)/(\sum_i \vec{p}_i)$, where i runs over ℓ and all jets
E_T	Missing transverse energy
p_T^{VIS}	Magnitude of the vector sum of the \vec{p}_T of the visible particles
$\max \Delta\eta(\ell, \{j_1 \text{ or } j_2\}) $	Maximum $\Delta\eta$ between the charged lepton and the leading or second leading jet



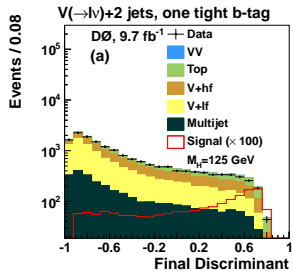
Final MVA Input Variables



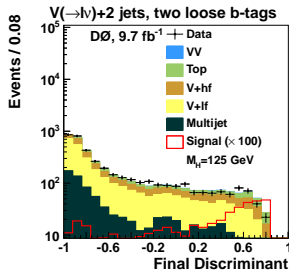
Variable	Description	2T	2M	2L	1T
$MVA_{MJ}(WH)$	Output of the multivariate discriminant trained to distinguish $WH \rightarrow \ell\nu b\bar{b}$ from the MJ background	1	1		
m_{b1b2}	The dijet invariant mass	2	4	3	1
$p_T^W / (p_T^\ell + \cancel{E}_T)$	The transverse momentum of the reconstructed W divided by the sum of the transverse momentum of the lepton and the missing transverse energy	3	6	4	2
b_{ID}^{j12}	Averaged b -jet identification output for the highest energy b -tagged jets	4	13	1	4
$\cos(\chi^*)$	$\chi^* = \angle(\ell, \text{the direction of the W boson spin})$ in $\ell\nu$ system CM frame [?]	5	3		
$\max \Delta\eta(\ell, \{j_1 \text{ or } j_2\}) $	Maximum $\Delta\eta$ between the charged lepton and the leading or second leading jet	6	11	2	3
$q^\ell \times \eta^\ell$	Product of the lepton charge and its pseudorapidity	7	2	6	6
$\Delta\mathcal{R}(\ell, j_1)$	$\Delta\mathcal{R}$ between the charged lepton and the leading jet	8	5		
$\min[SIG(j_{12}, \{j_1 \text{ or } j_2\})]$	Minimum SIG of the leading or second leading jet defined as $p_T^{\min}(j_1, j_2)\Delta\mathcal{R}(j_1, j_2) / \sum_{i=1}^2 p_T^{j_i}$ with respect to the dijet system.	9	15	9	5
	Based on the pull variables described in Ref. [?]				
$q^\ell \times \eta^{j1}$	Product of the lepton charge and pseudorapidity of the leading jet	10	7	11	9
$\mathcal{V}(j_{12})$	Velocity of the dijet system	11	12	7	11
$\cos(\theta^*)$	$\theta^* = \angle(W, \text{incoming } u\text{-type quark})$ in the Higgs CM frame [?]	12	10		
$m_{\ell\nu j_2}$	Invariant mass of the system consisting of the charged lepton, reconstructed neutrino, and second leading jet	13	16	12	13
m_T^{j12}	Transverse mass of the leading and second leading jets	14	14		
\mathcal{C}	Centrality is $(\sum_i p_T^i) / (\sum_i \vec{p}_i)$, where i runs over ℓ and all jets	15	8	8	10
$\sum (p_T)^{VIS}$	Scalar sum of the p_T of the visible particle	16	9		
m^{Asym}	Mass asymmetry between $\ell\nu$ system and the dijet system: $(M_{\ell\nu} - m_{j_{12}}) / (M_{\ell\nu} + m_{j_{12}})$			5	8
\mathcal{A}	Aplanarity is $3\lambda_3/2$ where λ_3 is the smallest eigenvalue of the normalized momentum tensor $S^{\alpha\beta} = (\sum_i p_i^\alpha p_i^\beta) / (\sum_i \vec{p}_i ^2)$, where $\alpha, \beta = 1, 2, 3$ correspond to the x, y, z momentum components, and i runs over selected objects. Without arguments, it is calculated for all visible objects			10	12
p_T^{j2}	p_T of the second leading jet			13	7



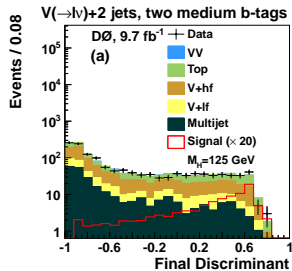
MVA Output Distributions



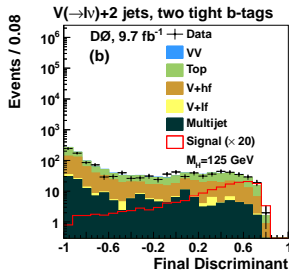
(a) 1 tight b-tag



(b) 2 loose b-tags



(c) 2 medium b-tags





- Consider two hypotheses: Background only (H_0), and Signal + Background (H_1)
- Assume data comes randomly from a Poisson distribution, and generate pseudo data for two hypotheses
- Use Log-Likelihood Ratio (LLR) test statistic to evaluate statistical significance.
 - LLR is the -2 times the log of the ratio of likelihoods for the test and null hypotheses:
 - $LLR = -2\log\left(\frac{LH(H_1)}{LH(H_0)}\right)$.
 - Most powerful test statistic for distinguishing between two hypotheses.
 - Calculated for each set of pseudo-data for background only and signal + background hypotheses.

$$LLR = \sum_{i=0}^{N_c} \sum_{j=0}^{N_{bins}} s_{ij} - d_{ij} \ln \left(1 + \frac{s_{ij}}{b_{ij}} \right) \quad (1)$$