Signatures of the Inert Doublet Dark Matter Model

Alejandro Ibarra

Technische Universität München





Based on works with

Camilo Garcia-Cely, JCAP 1309 (2013) 025. Camilo Garcia-Cely, Michael Gustafsson, in preparation

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Could be composed by scalar doublets

The Inert Doublet Model

Deshpande, Ma '78

Barbieri, Hall, Rychkov '06 Lopez Honorez, Nezri, Oliver, Tytgat '06 Cirelli, Fornengo, Strumia '06 Majumdar, Ghosal '06 Gustafsson, Lundström, Bergström, Edsjö '07 Agrawal, Dolle, Krenke '08 Andreas, Tytgat, Swillens '09 Nezri, Tytgat, Vertongen '09 Arina, Ling, Tytgat '09 Hambye, Ling, Lopez Honorez, Rocher '09 Lopez Honorez, Yaguna '10 Klasen, Yaguna, Ruiz-Alvarez '13 AI, Garcia Cely '13

Simple extension of the Standard Model accounting for the dark matter of the Universe.

Assumptions:

- 1) Introduce a scalar field, η , with identical gauge quantum numbers as the Standard Model Higgs boson.
- 2) Postulate that the vacuum possess an exact \mathbb{Z}_2 symmetry, under which η is odd, while all the Standard Model particles are even: $\begin{array}{c} \eta \rightarrow -\eta \\ SM \rightarrow SM \end{array}$

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Implications

- 1) The \mathbb{Z}_2 symmetry ensures that η contains an absolutely stable particle \rightarrow dark matter candidate.
- 2) The exotic doublet η does not interacts at tree level with any of the Standard Model fermions. The doublet η is "inert".
- 3) The inert doublet η , communicates with the observable sector via gauge interactions and via the Higgs portal, e.g. $(\Phi^{\dagger}\Phi) (\eta^{\dagger}\eta)$

 \rightarrow thermal dark matter production possible

Lagrangian of the scalar sector $\mathcal{L} = \mathcal{L}_{\Phi} + \mathcal{L}_{\eta} - V_{\text{int}}$ $\mathcal{L}_{\Phi} = (D_{\mu}\Phi)^{\dagger}(D^{\mu}\Phi) - m_{\Phi}^2(\Phi^{\dagger}\Phi) - \lambda_1(\Phi^{\dagger}\Phi)^2$ $\mathcal{L}_{\eta} = (D_{\mu}\eta)^{\dagger} (D^{\mu}\eta) - m_{\eta}^2 (\eta^{\dagger}\eta) - \lambda_2 (\eta^{\dagger}\eta)^2$ $V_{\rm int} = \lambda_3(\Phi^{\dagger}\Phi)(\eta^{\dagger}\eta) + \lambda_4(\Phi^{\dagger}\eta)(\eta^{\dagger}\Phi) + \frac{1}{2} \Big(\lambda_5(\Phi^{\dagger}\eta)(\Phi^{\dagger}\eta) + \text{h.c.}\Big)$

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After the electroweak symmetry breaking

$$\left< \Phi \right> = \begin{pmatrix} 0 \\ \left< \Phi^0 \right> \end{pmatrix} \ ,$$

Physical scalar fields: h, H^0, A^0, H^{\pm} :

$$\Phi = \begin{pmatrix} 0 \\ \langle \Phi^0 \rangle + h \end{pmatrix} , \qquad \eta = \begin{pmatrix} H^+ \\ \frac{1}{\sqrt{2}} \left(H^0 + iA^0 \right) \end{pmatrix}$$



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Dark matter candidates

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After the electroweak symmetry breaking

$$\left< \Phi \right> = \begin{pmatrix} 0 \\ \left< \Phi^0 \right> \end{pmatrix} \ ,$$

 $\langle \eta \rangle = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$ Preserved

Physical scalar fields: h, H^0, A^0, H^{\pm} :

$$\Phi = \begin{pmatrix} 0 \\ \langle \Phi^0 \rangle + h \end{pmatrix} , \qquad \eta = \begin{pmatrix} 1 \\ \frac{1}{\sqrt{2}} (H^0 + iA^0) \end{pmatrix}$$

Mass splittings between the exotic scalars and their interactions with *h* determined by the quartic couplings λ_3 , λ_4 , λ_5 .

Mass splittings

$$M_{H^+}^2 - M_{H^0}^2 \propto (\lambda_4 + \lambda_5) \qquad M_{A^0}^2 - M_{H^0}^2 \propto \lambda_5$$



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Interactions



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Interactions



Rich phenomenology. Parameter space constrained by

- Perturbativity
- Vacuum stability
- Collider searches
- Electroweak precision tests
- Relic dark matter abundance
- Direct dark matter searches
- Indirect dark matter searches



WIMP dark matter production



WIMP dark matter production



Relic abundance of DM particles

$$\Omega h^2 \simeq \frac{3 \times 10^{-27} \, \mathrm{cm}^3 \, \mathrm{s}^{-1}}{\langle \sigma v \rangle}$$

Correct relic density,
$$\Omega h^2 = 0.1$$
, if

$$\langle \sigma v \rangle \simeq 3 \times 10^{-26} \,\mathrm{cm}^3 \,\mathrm{s}^{-1}$$











 Ωh^2

Irrelevant if $m_{A^0} - m_{H^0} \gtrsim 100 \text{ keV}$

If allowed, typically leads to a too large scattering rate.

$$\sigma(H^0 p) \simeq 4 \times 10^{-44} \lambda_{H^0}^2 \left(\frac{1 \,\mathrm{TeV}}{M_{H^0}}\right)^2 \,\mathrm{cm}^2$$

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The decays of the W, Z and Higgs bosons produce
a flux of antiprotons, electrons/positrons, neutrinos and gamma-rays that could be detected at the Earth.

Antiprotons

The decays of the W, Z and Higgs bosons producea flux of antiprotons, electrons/positrons, neutrinos and gamma-rays that could be detected at the Earth.

Positrons

Positron fraction well fitted by:

$$\Phi_{e^+}^{\mathrm{IS}}(E) = \underbrace{C_{e^+} E^{-\gamma_{e^+}}}_{\substack{\mathrm{Secondary}\\ \mathrm{production}}} + \underbrace{C_s E^{-\gamma_s} \exp(-E/E_s)}_{\substack{\mathrm{Sources}}}$$

and corrected for solar modulation effects. \rightarrow Limits on dark matter properties

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Positrons

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Gamma rays: diffuse flux

The decays of the W, Z and Higgs bosons producea flux of antiprotons, electrons/positrons, neutrinos and gamma-rays that could be detected at the Earth.

Gamma rays: dwarf galaxies

However, astrophysical processes could produce an antiproton flux or a featureless gamma-ray flux that could mimic the annihilations $H^0 H^0 \rightarrow W^+ W^-$

How to discover dark matter in indirect dark matter searches?

Smoking gun for dark matter: no (known) astrophysical process can produce a sharp feature in the gamma-ray energy spectrum

Three gamma-ray spectral features have been identified:

Gamma ray line

Internal bremsstrahlung

Smoking gun for dark matter: no (known) astrophysical process can produce a sharp feature in the gamma-ray energy spectrum

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Gamma ray line

Gamma ray box

Internal bremsstrahlung

The three of them arise in the inert doublet model!

Gamma-ray lines

+ symmetry related and non-W contributions

Gustafsson

Gamma-ray lines

Gamma-ray lines

Gamma-ray lines

Gamma lines in the high-mass regime? In progress...

AI, Garcia-Cely, Gustafsson

Gamma-ray boxes

Gamma-ray boxes

Gamma-ray boxes

Internal Bremsstrahlung

$$H^0 \longrightarrow W^-$$

 $H^+ \longrightarrow \gamma$
 H^+
 $H^0 \longrightarrow W^+$

Internal Bremsstrahlung

$$\Delta(p) \propto \left((p_{H^0} - p_W)^2 - M_{H^+}^2 \right)^{-1} \\ \approx (M_{H^0}^2 + M_W^2 - M_{H^+}^2 - 2M_{H^0} E_W)^{-1}$$

In the case $M_{H^0} \approx M_{H^+} \gg M_W$ the scalar propagator gets enhanced when E_W is small.

Internal Bremsstrahlung

 $H^0 \longrightarrow W^ H^+ \longrightarrow \gamma$ $H^+ H^+$ $H^0 \longrightarrow W^+$

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The condition $M_{H^0} \approx M_{H^+} \gg M_W$ is automatically fulfilled in the high mass window of the inert doublet dark matter model.

Enhancement of the amplitude (and the rate) when E_{γ} is close to the kinematic end-point.

Internal Bremsstrahlung

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Internal Bremsstrahlung

Especially for large dark matter masses and λ_3 small, the spectral feature is very prominent.

Internal Bremsstrahlung: benchmark points

BMP	$M_{H^0}(\text{GeV})$	$M_{H^+}(\text{GeV})$	$M_{A^0}(\text{GeV})$	λ_3	λ_4	λ_5	λ_{H^0}	λ_{A^0}
1	559.99	561.85	560.67	-0.02	-0.06	-0.01	-0.05	-0.03
2	983.75	993.60	991.79	0.17	-0.38	-0.26	-0.23	0.03
3	2088.3	2090.99	2100.26	0.39	0.46	-0.83	0.01	0.83
4	3596.67	3597.99	3609.7	-0.07	1.24	-1.55	-0.19	1.36
5	4212.49	4213.09	4225.29	-0.58	1.61	-1.78	-0.37	1.41
6	5382.08	5382.21	5392.59	1.08	1.82	-1.87	0.51	2.38

AI, Garcia Cely arXiv:1306.4681

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Internal Bremsstrahlung: H.E.S.S. limits

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Future gamma-ray telescopes might explore the region where spectral features could be observed.

Internal Bremsstrahlung + lines

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Internal Bremsstrahlung + lines

Conclusions

The inert doublet model is a simple extension of the Standard Model which accounts for the dark matter of the Universe.

It offers a rich phenomenology and predicts potentially observable effects in direct and indirect dark matter searches, collider searches and electroweak precision tests.

Large portions of the parameter space will be explored in the near future by direct dark matter search experiments (LUX, XENON1T).

It generically predicts gamma-ray spectral features which, if observed, would constitute a strong hint for dark matter annihilations in the Milky Way center.