



# Quarkonium production in pp, p-Pb, Pb-Pb collisions with ALICE at LHC

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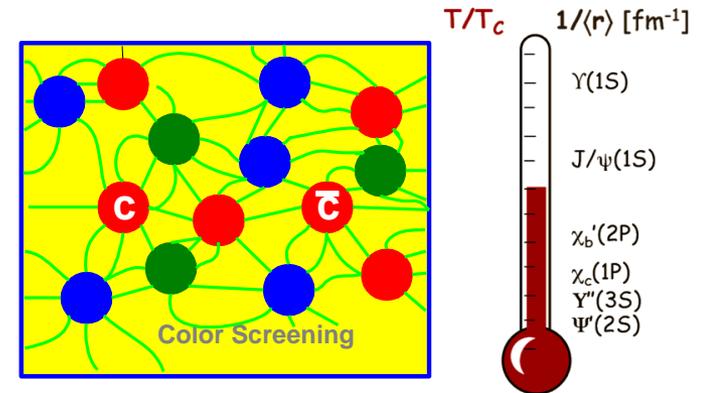
**QCD@Work - 7<sup>th</sup> International Workshop on QCD – Theory and Experiment**

# Outline

- Physics motivation
- ALICE setup
- ALICE results
  - Quarkonium in pp collisions at  $\sqrt{s}=2.76$  TeV
  - Pb-Pb measurements at  $\sqrt{s_{NN}} = 2.76$  TeV
    - nuclear modification factor
  - p-Pb measurements at  $\sqrt{s_{NN}} = 5.02$  TeV
    - cold nuclear matter effects
- Conclusion and outlook

# Probing the QGP : Quarkonium suppression and enhancement

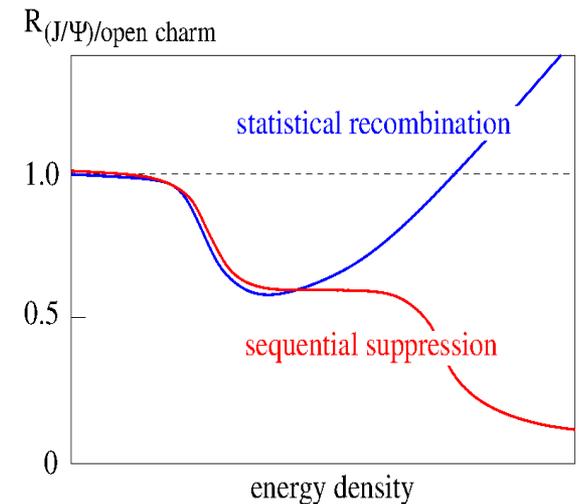
- Quarkonium ( $J/\psi$ ) suppression [Matsui,Satz; PLB 178 (1986) 416]
  - Color screening of strong interactions in QGP
  - Higher states are more easily dissociated, sequential suppression of resonances



- $J/\psi$  enhancement [Braun-Munzinger, Stachel; PLB 490(2000) 196]
  - The  $c\bar{c}$  production increases strongly with energy

In most central collisions [0-10%]	RHIC 200 GeV	LHC 2.76 TeV
$N_{c\bar{c}bar}/\text{event}$	13	115
$N_{b\bar{b}bar}/\text{event}$	0.1	3

- A (re)combination of  $c\bar{c}$  pairs to produce quarkonia may take place during the QGP stage or at the phase boundary



Digal, Satz,Vogt; PRC 85, (2012) 034906

# Different sources of medium effects

- Nuclear modification factor  $R_{AA}$ :  
Ratio of the quarkonium yield in AA ( $Y_{AA}$ ) with respect to the pp one, scaled by the overlap factor  $T_{AA}$  (from Glauber model)

$$R_{AA} = \frac{Y_{AA}}{\langle T_{AA} \rangle \sigma_{pp}}$$

If yield scales with the number of binary collisions

$$\rightarrow R_{AA} = 1$$

and if,

$$\rightarrow R_{AA} \neq 1$$

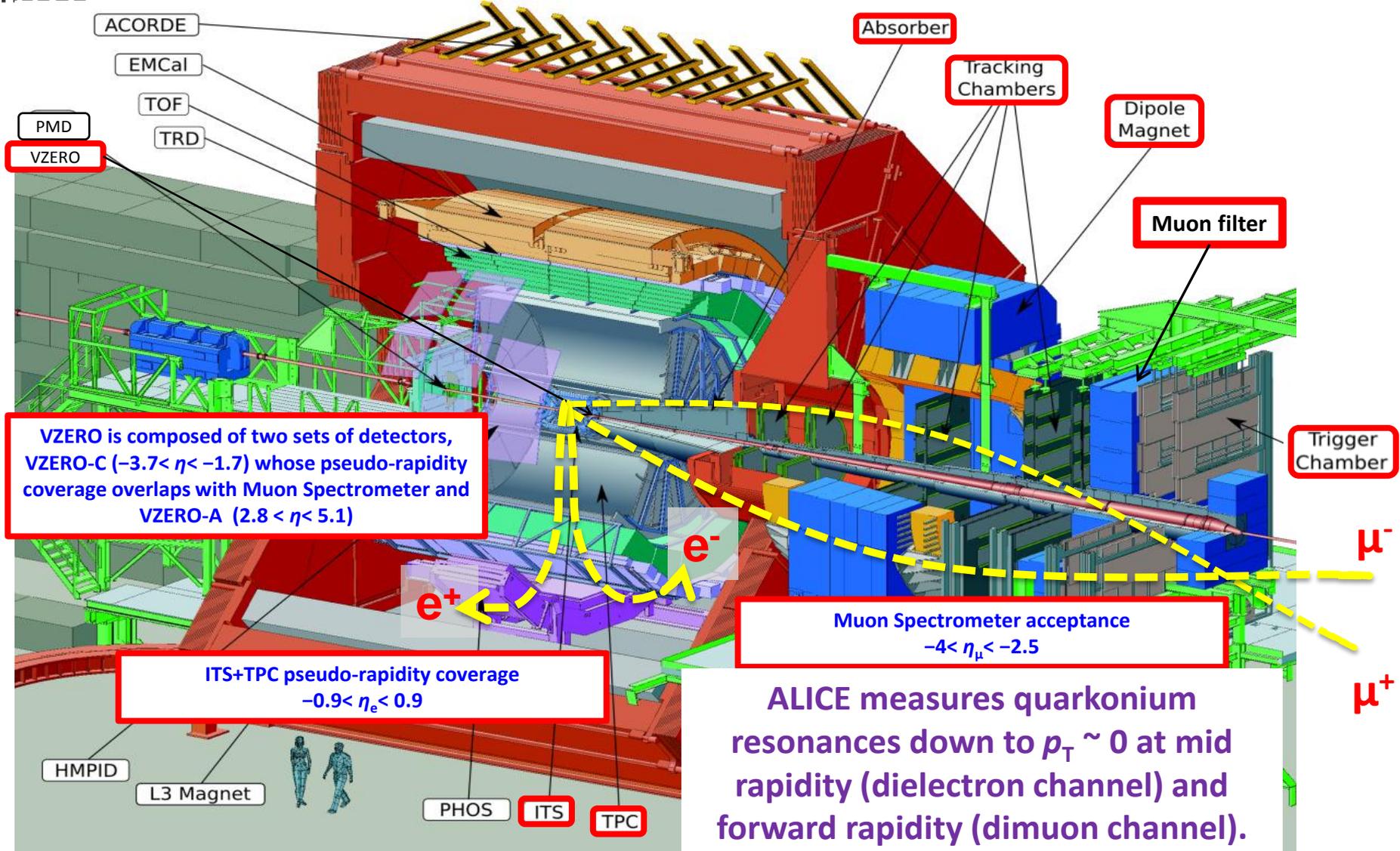
there are medium effects

## Hot Medium effects:

- Quarkonium suppression
- Enhancement due to recombination

## Cold Nuclear Matter effects (CNM):

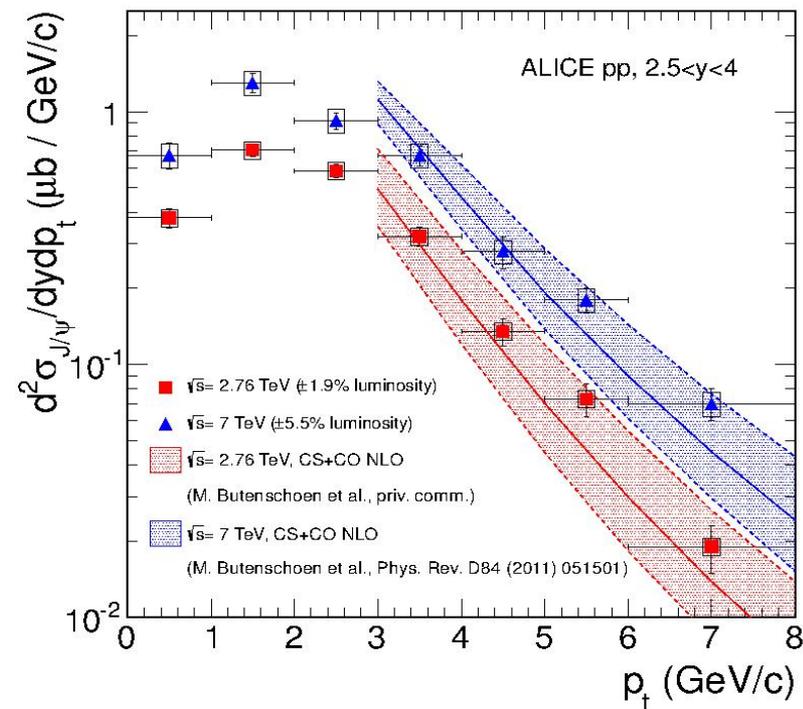
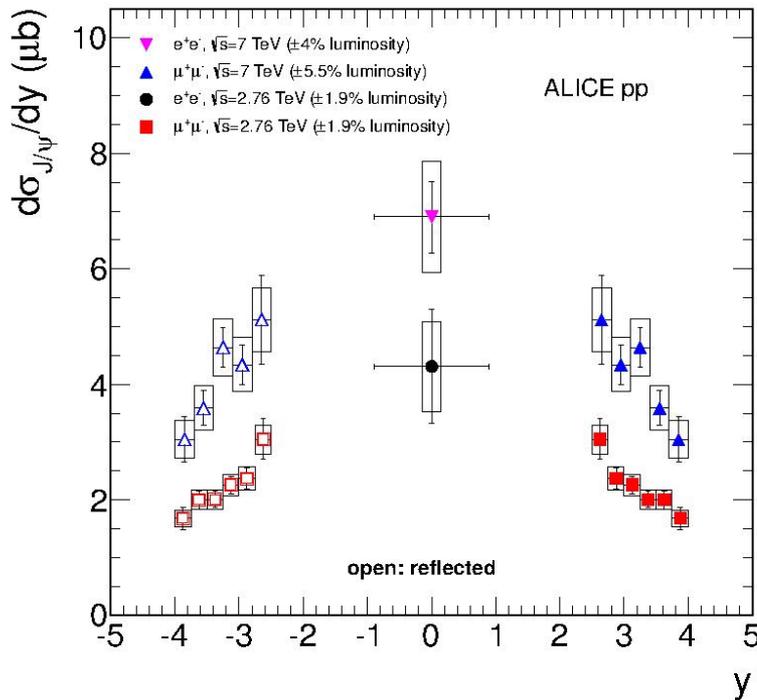
- Nuclear parton shadowing/gluon saturation
- Parton energy loss
- $c\bar{c}$  in medium break-up



Data taking at  $\sqrt{s}=2.76$  TeV essential to build the  $R_{AA}$  reference, result based on  $L_{int}^e=1.1 \text{ nb}^{-1}$  and  $L_{int}^\mu=19.9 \text{ nb}^{-1}$

$$\sigma_{J/\psi}(|y| < 0.9) = 7.75 \pm 1.78(\text{stat.}) \pm 1.39(\text{syst.}) + 1.16(\lambda_{HE} = 1) - 1.63(\lambda_{HE} = -1) \mu\text{b}$$

$$\sigma_{J/\psi}(2.5 < y < 4) = 3.34 \pm 0.13(\text{stat.}) \pm 0.27(\text{syst.}) + 0.53(\lambda_{CS} = 1) - 1.07(\lambda_{CS} = -1) \mu\text{b}.$$

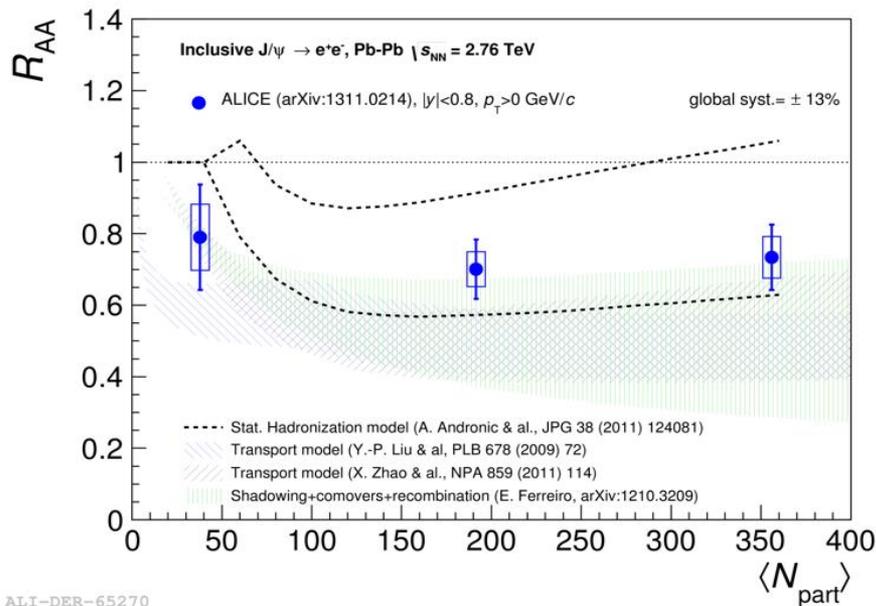


ALICE Coll., PLB 718 (2012) 295

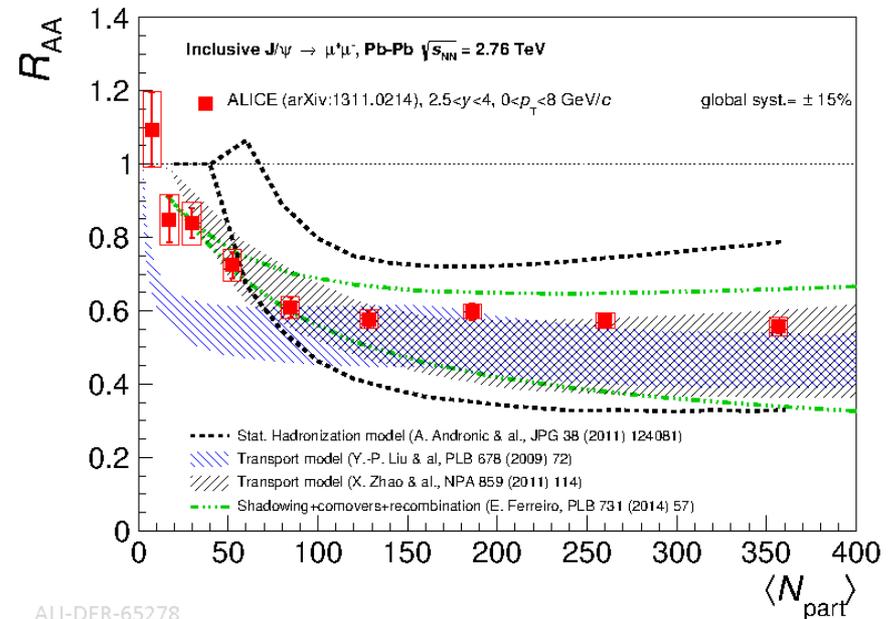
(Note the updated high precision results of pp at  $\sqrt{s}=7$  TeV is available at [arXiv:1403.3648](https://arxiv.org/abs/1403.3648) )

A reduction of the J/ψ yield wrt to pp collisions is observed at SPS ( $\sqrt{s}=17\text{GeV}$ ), RHIC ( $\sqrt{s}=200\text{GeV}$ ) and finally LHC ( $\sqrt{s}=2.76\text{TeV}$ )! [ALICE central and forward rapidity luminosities are  $23\mu\text{b}^{-1}$  and  $70\mu\text{b}^{-1}$ , respectively]

Julian Book, QM (2014)



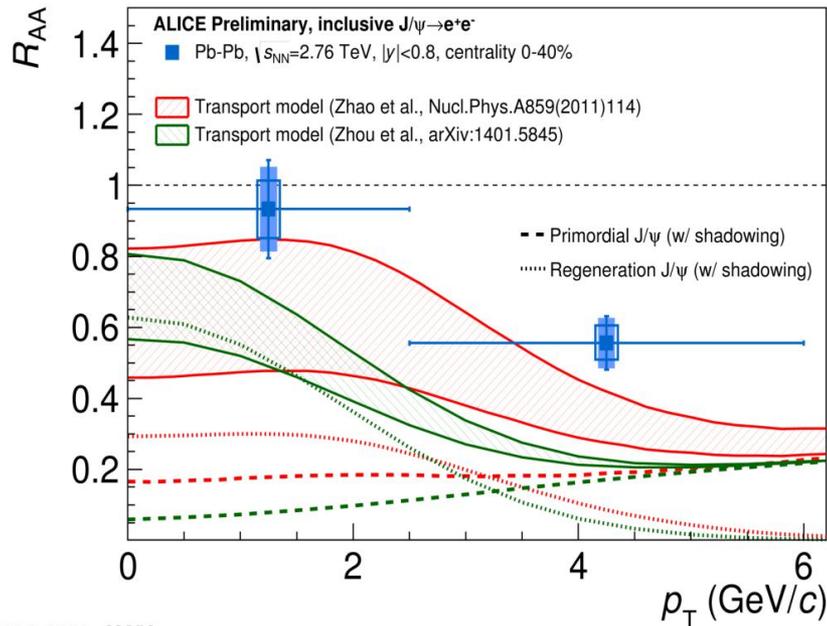
ALI-DER-65270



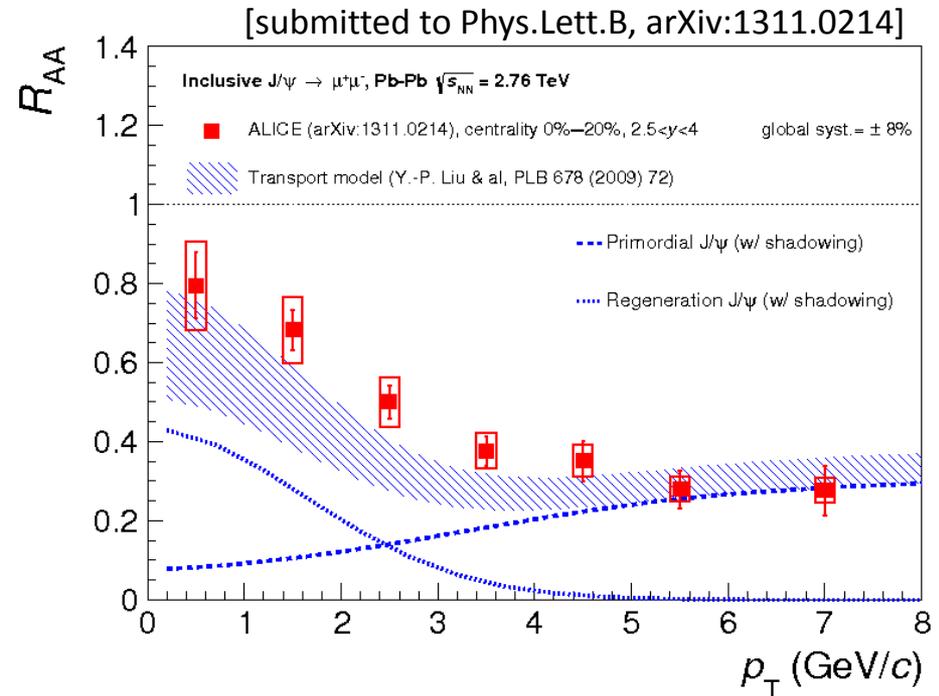
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Statistical hadronization and transport models which respectively feature a full and a partial J/ψ suppression of charm quarks at hadronization or in the QGP phase can describe the data. Note that both models take into account the recombination of  $c\bar{c}$  pair to J/ψ during the QGP stage or at the phase boundary.

Another signature of recombination is visible when the  $R_{AA}$  is plotted as function of  $p_T$ . For central events, the model prediction of  $\sim 50\%$  low- $p_T$  J/ψ production via recombination and no recombination at high- $p_T$  well reproduces the data in left plot for mid-rapidity and right plot for forward rapidity.

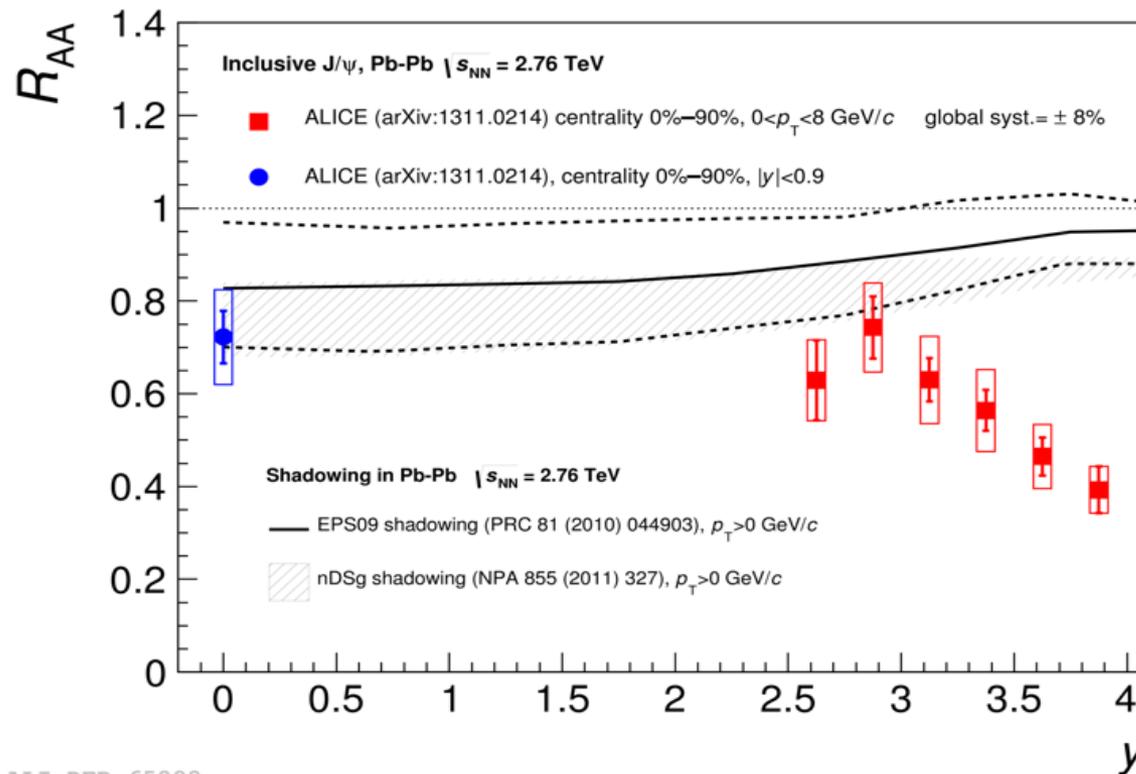


ALI-PREL-69850



# J/ψ $R_{AA}$ in Pb-Pb

Inclusive J/ψ measured also as a function of rapidity:  $R_{AA}$  decreases by 40% from  $y=2.5$  to  $y=4$



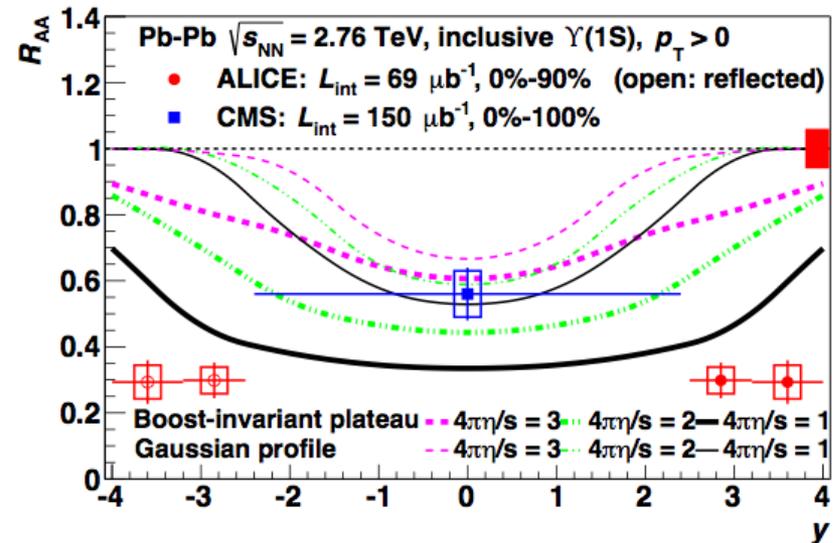
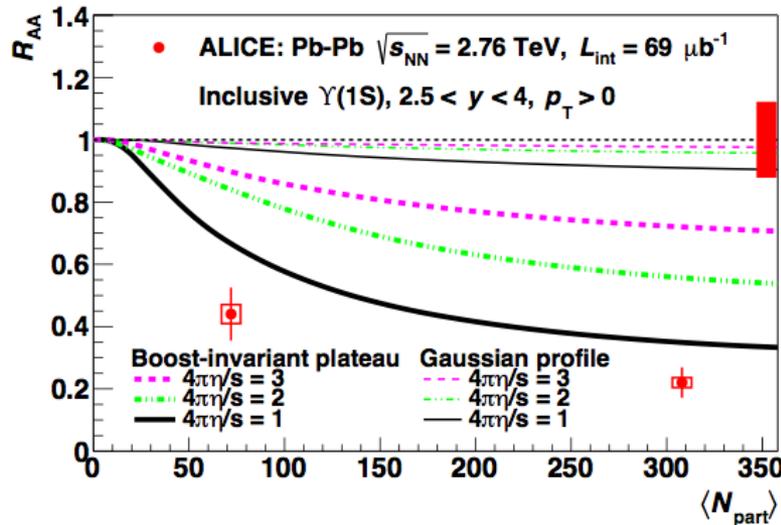
ALI-DER-65282

[submitted to Phys.Lett.B, arXiv:1311.0214]

# $\Upsilon(1S)R_{AA}$ in Pb-Pb

A strong suppression have been observed in the inclusive measurement of  $\Upsilon(1S)$  state in heavy-ion collision at forward rapidity ( $2.5 < y < 4.0$ )

arXiv:1405.4493



CMS: PRL109 (2012) 222301

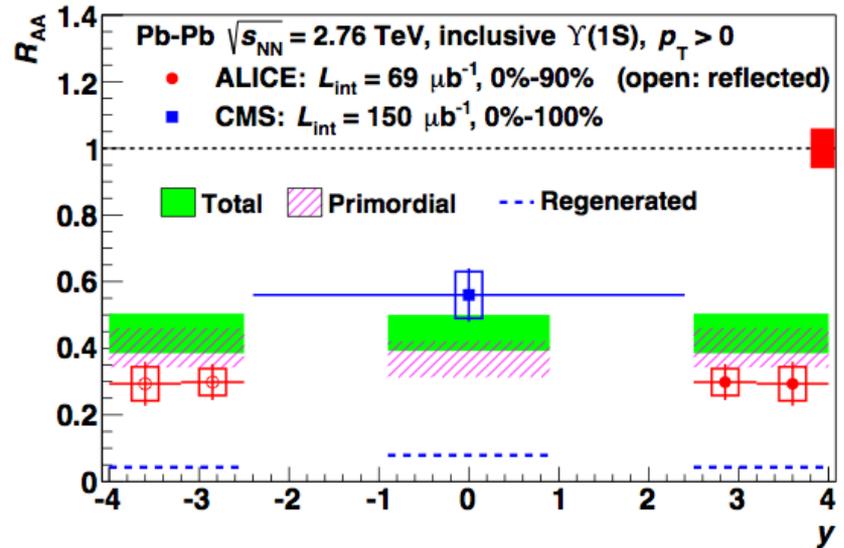
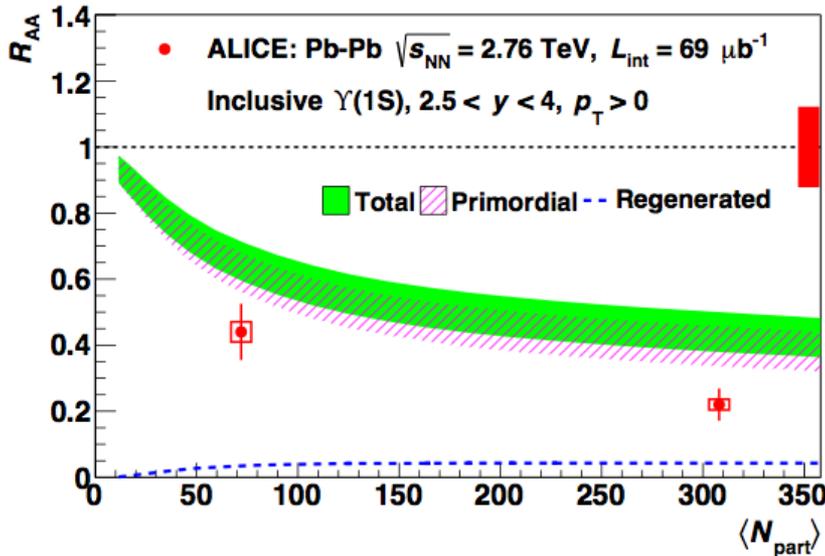
M. Strickland, [arXiv:1207.5327]

- Thermal suppression of bottomonium states
- Anisotropic hydro model
- Two temperature rapidity profiles: Boost invariant or Gaussian
- Three shear viscosities
- Feed down from higher mass states included
- No CNM included
- No regeneration included

In all cases the model underestimates the measured  $\Upsilon(1S)$  suppression at forward rapidity

# $\Upsilon(1S)R_{AA}$ in Pb-Pb

arXiv:1405.4493

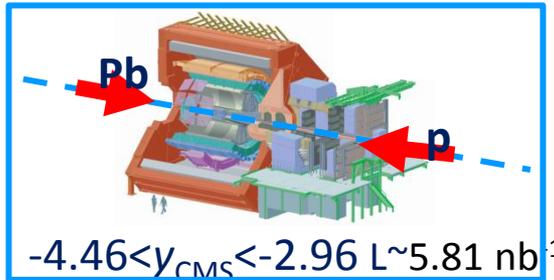
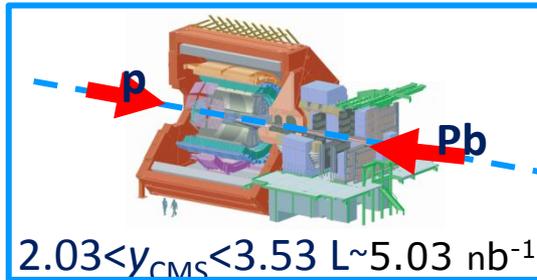


CMS: PRL109 (2012) 222301

A. Emerick et al., [EPJ A48 (2012) 72]

- Transport model
- Suppression of  $\Upsilon(1S)$  resonances by the QGP
  - Mainly of the higher mass states
- Small regeneration component included
- Feed down from higher mass states included
- CNM included via an “effective”  $\sigma_{ABS} = 0-2$  mb

Model does not reproduce the strong rapidity dependence of the  $R_{AA}$  and underestimates the  $\Upsilon(1S)$  suppression at forward rapidity



ALICE collected p-Pb/Pb-p data at  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$  at the beginning of 2013

$\Delta y = 0.465$  towards the direction of p-beam

**Nuclear modification factor  $R_{pA}$ :**

$$R_{pA} = \frac{Y_{pA}}{\langle T_{pA} \rangle \cdot \sigma_{pp}}$$

The full coverage of the ALICE muon spectrometer  $2.5 < y_{LAB} < 4$  can be exploited

**Modification factor  $Q_{pA}$ :**

$$Q_{pA} = \frac{Y_{pA}}{\langle T_{pA}^{mult} \rangle \cdot \sigma_{pp}}$$

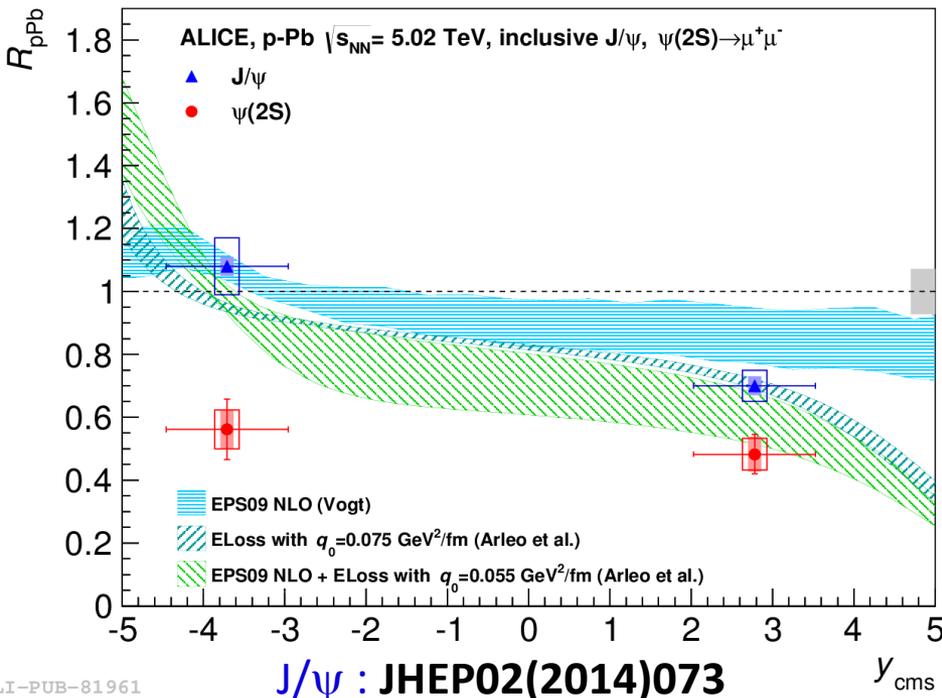
For centrality dependent studies  $Q_{pA}$  instead of  $R_{pA}$

[Phys Lett B 727 (2013) 371–380]

In LHC there were no data for pp collision at  $\sqrt{s} = 5.02 \text{ TeV}$ , therefore the reference cross section  $\sigma_{pp}$  is obtained by means of an interpolation procedure based on ALICE data at  $\sqrt{s_{NN}} = 7$  and  $2.76 \text{ TeV}$

# J/ψ and ψ(2S) $R_{pA}$ vs $y$

Theoretical models including initial state effects as shadowing and parton energy loss predict the same suppression (within <5%) for J/ψ and ψ(2S). However, while they are qualitatively in agreement with the J/ψ results they cannot reproduce the ψ(2S) suppression.



J/ψ : JHEP02(2014)073

ψ(2S) : arXiv : 1405.3796

possible if formation time  
( $\tau_f \sim 0.05-0.15$  fm/c) < crossing time ( $\tau_c$ )

forward-y:

$$\tau_c \sim 10^{-4} \text{ fm/c}$$

backward-y:

$$\tau_c \sim 7 \cdot 10^{-2} \text{ fm/c}$$

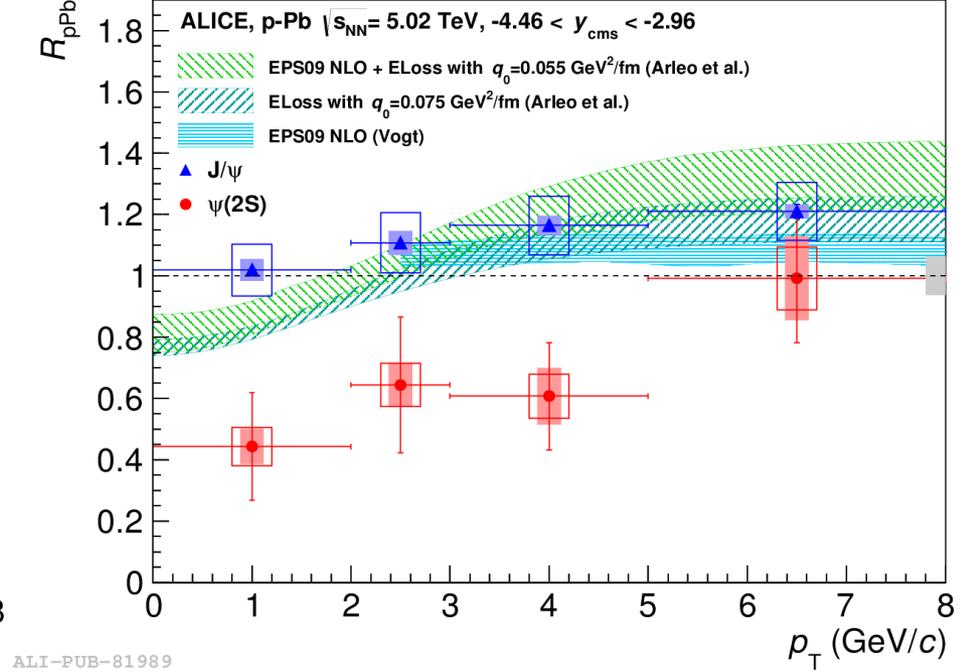
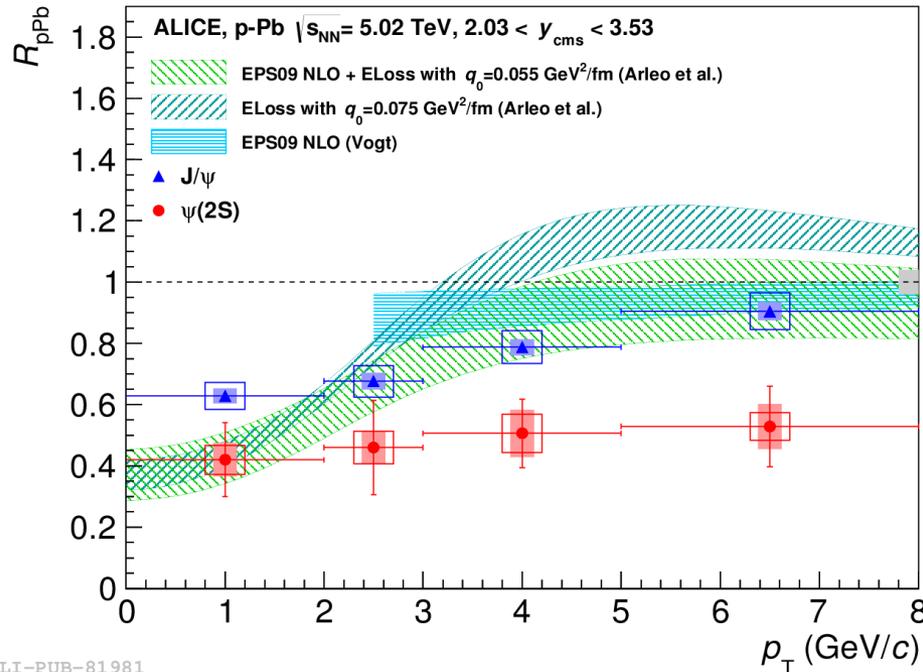
$$\tau_c = \frac{\langle L \rangle}{(\beta_z \gamma)}$$

D. McGlinchey, A. Frawley and R. Vogt, PRC 87,054910 (2013)

➔ break-up effects excluded at forward-y

➔ at backward-y, since  $\tau_f \sim \tau_c$ , break-up in CNM can hardly explain the very strong difference between J/ψ and ψ(2S) suppressions

# J/ψ and ψ(2S) $R_{pA}$ vs $p_T$



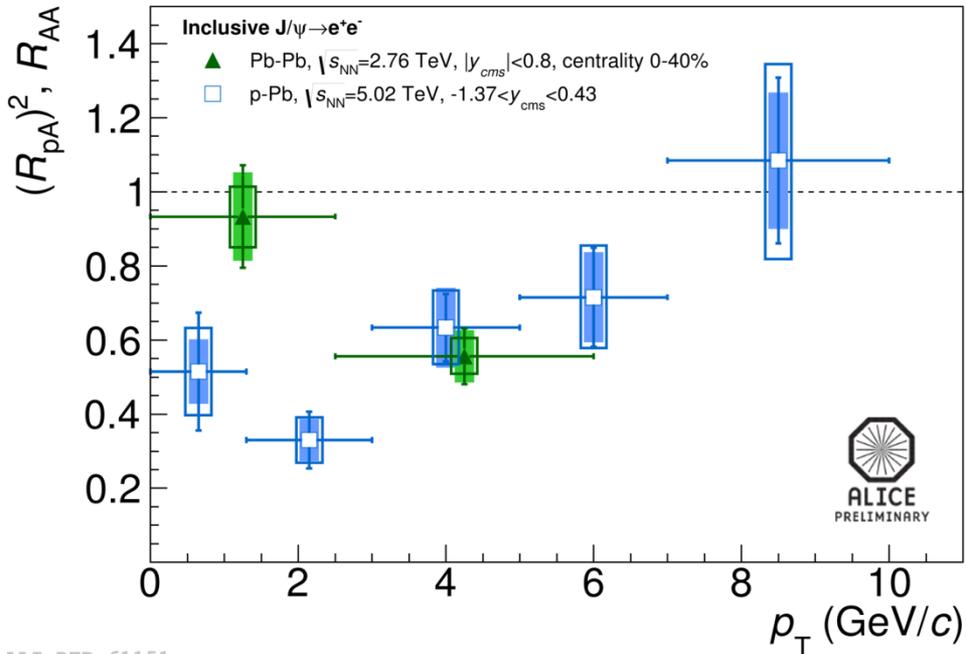
ALI-PUB-81981

ALI-PUB-81989

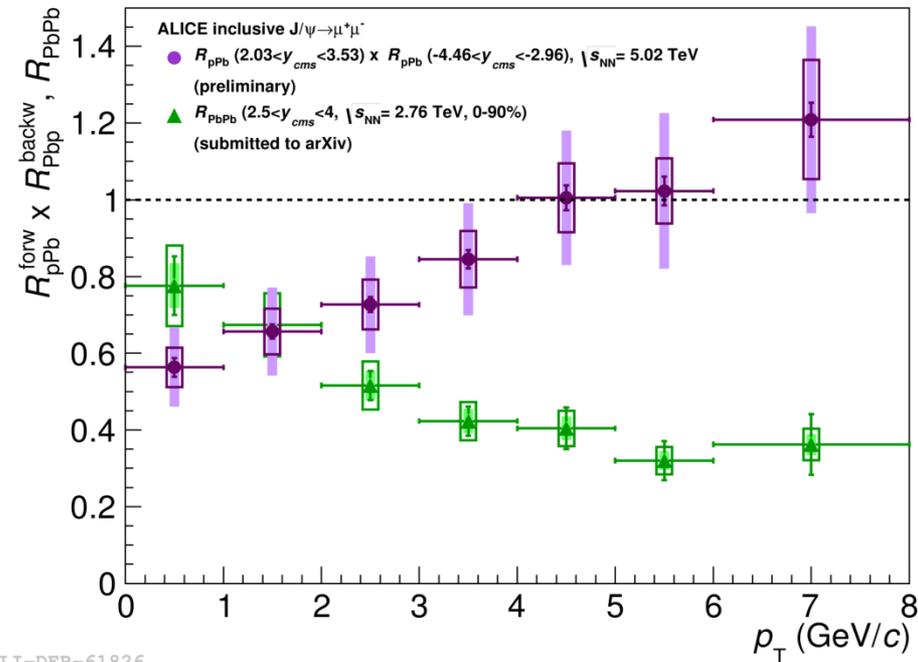
As already observed for the  $p_T$ -integrated results, ψ(2S) is more suppressed than the J/ψ

Theoretical models are in fair agreement with the J/ψ, but clearly overestimate the ψ(2S) results

# J/ψ $R_{pA}$ and $R_{AA}$ in mid and forward rapidity

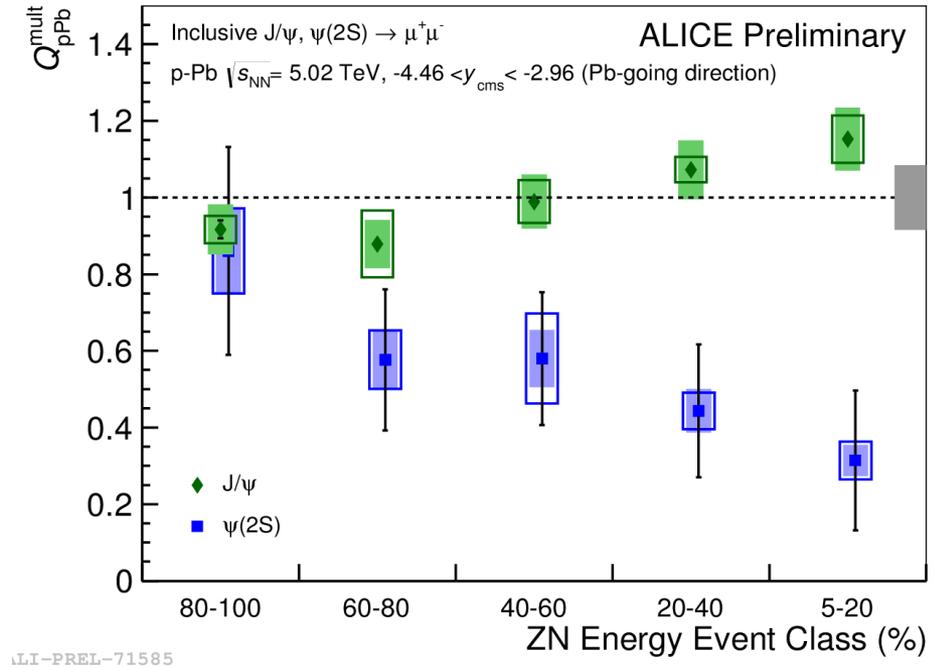
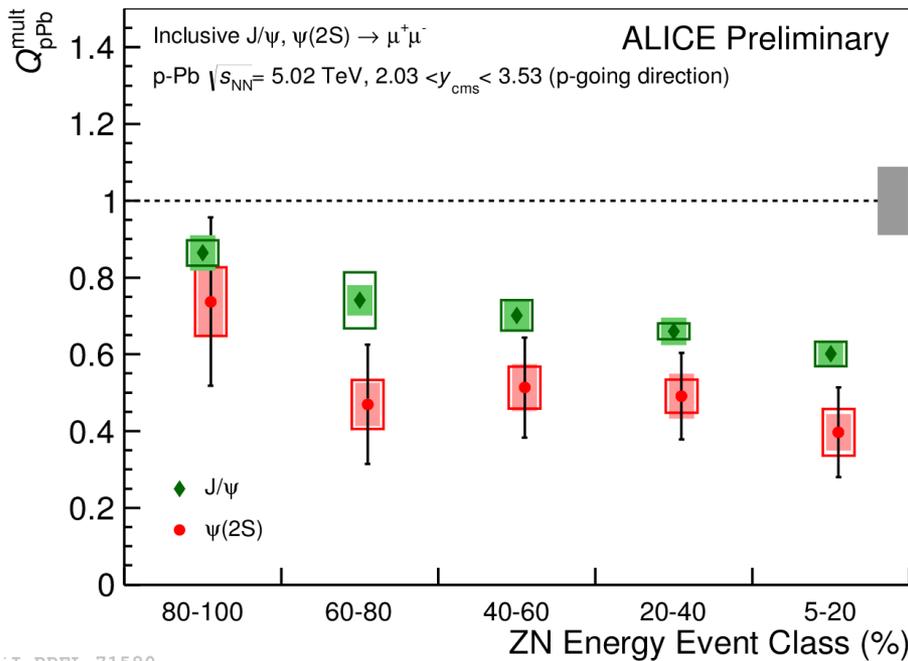


ALI-DER-61151



ALI-DER-61826

The CNM effects in AA is extrapolated using the pA results. The  $R_{AA}$  is compared to  $R_{pA} \times R_{Ap}$  ( $R_{pA}^2$  at midrapidity) one can see that CNM effects are not enough to describe the high  $p_T$  suppression in AA and that at low  $p_T$  there is an enhancement which may be a hint for recombination



forward-y: J/ψ and ψ(2S) show a similar decreasing pattern vs event activity

backward-y: the J/ψ and ψ(2S) behaviour is different, with the ψ(2S) significantly more suppressed for larger event activity classes

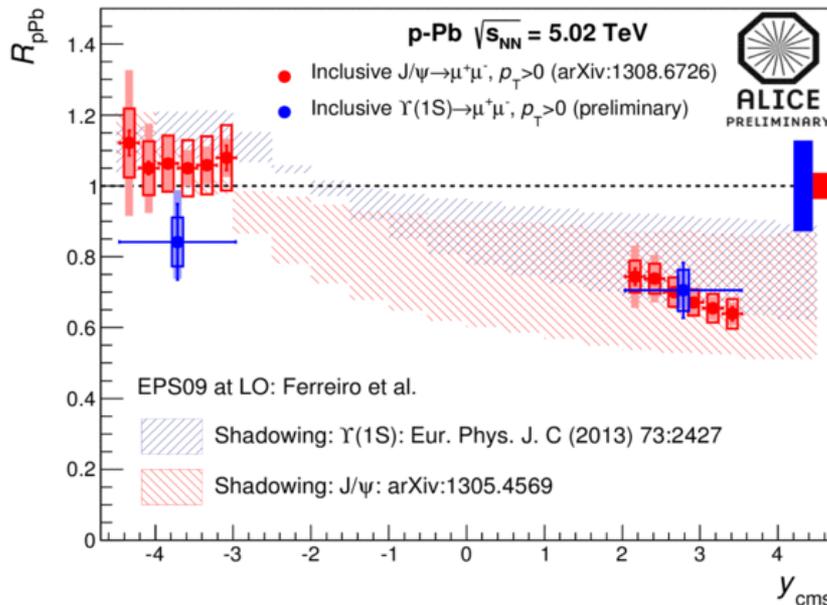


ALICE

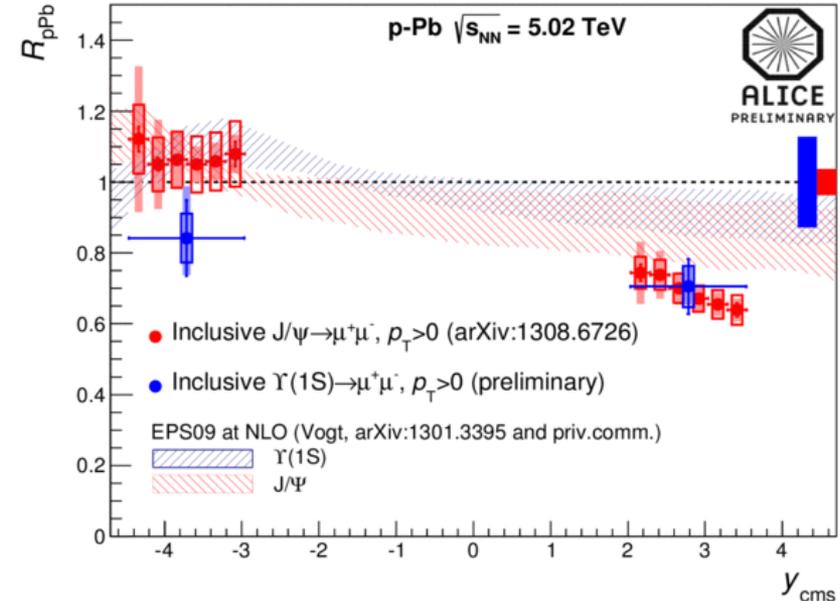
# $R_{pA}$ of $\Upsilon(1S)$



- Comparison with ALICE J/ $\psi$   $R_{pPb}$ 
  - Forward: similar suppression
  - Backward: slightly lower  $R_{pPb}$  for  $\Upsilon(1S)$ , but compatible within uncertainties



ALI-DER-58992

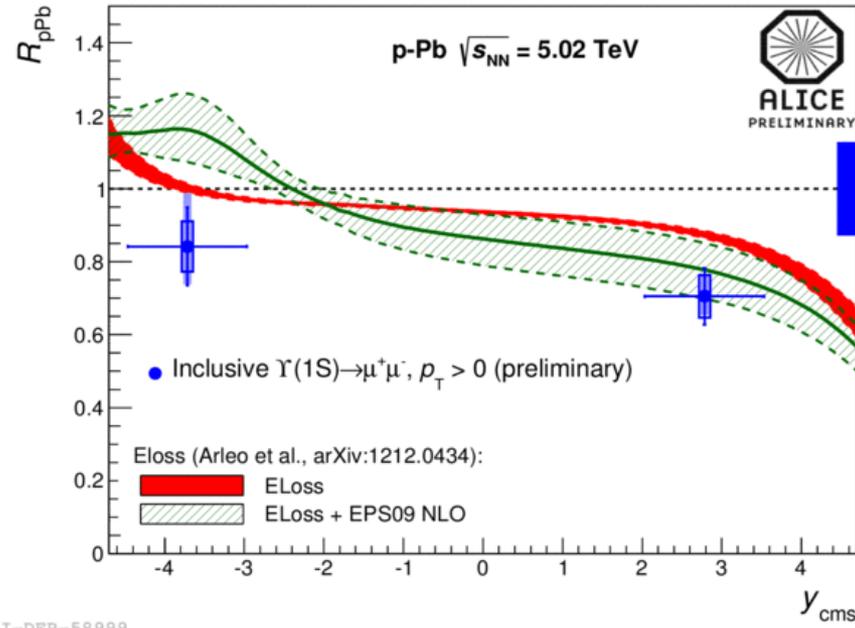


ALI-DER-58992

- Ferreiro et al. [EPJC 73 (2013) 2427]
  - 2 $\rightarrow$ 2 production model at LO
  - EPS09 shadowing parameterization at LO
  - Fair agreement with measured  $R_{pPb}$ 
    - Although slightly overestimates it in the antishadowing region

- Vogt [arXiv:1301.3395]
  - CEM production model at NLO
  - EPS09 shadowing parameterization at NLO
  - Fair agreement with measured  $R_{pPb}$  within uncertainties
    - Although slightly overestimates it

# $R_{pA}$ of $\Upsilon(1S)$



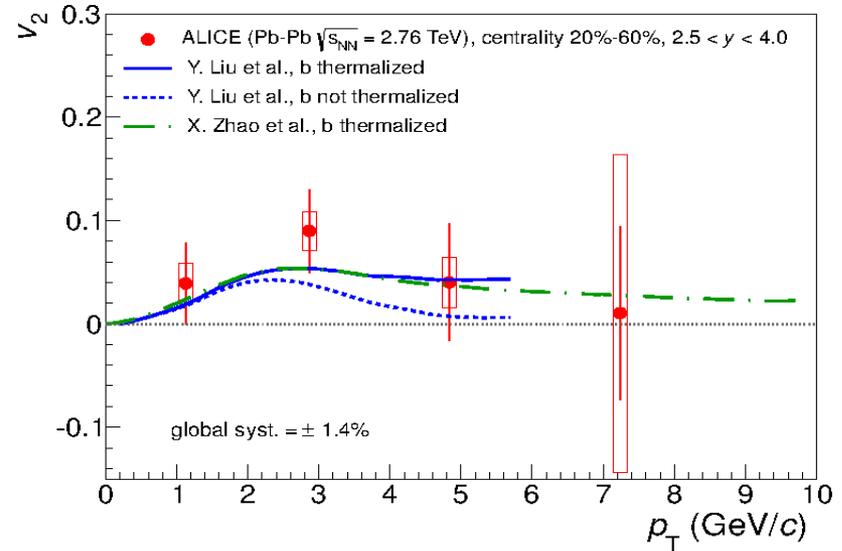
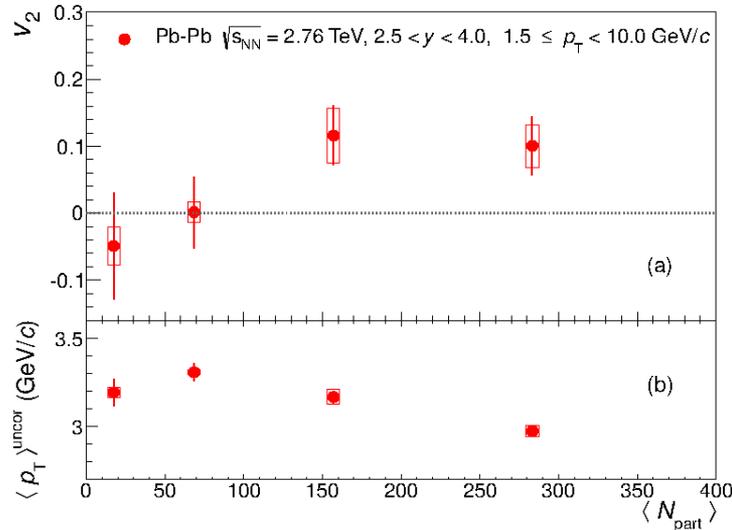
- Arleo et al. [JHEP 1303 (2013) 122]
  - Model including a contribution from coherent parton energy loss
  - With or without shadowing (EPS09)
  - Forward: Better agreement with ELoss and shadowing
  - Backward: Better agreement with ELoss only

- ALICE findings in Pb-Pb collisions:
  - Two main mechanisms at play
    - Suppression by color screening or dissociation in QGP
    - Re-generation (for charmonium only!) at high  $\sqrt{s}$  can qualitatively explain the main features of the results
  - $R_{AA}$  exhibits a weak centrality dependence at all  $y$  for charmonium, however  $\Upsilon(1S)$  shows centrality dependence.
  - Less suppression at low  $p_T$  with respect to high  $p_T$ , with stronger  $p_T$  dependence for central events
  - Stronger suppression when rapidity increases for  $J/\psi$ , whereas there is no rapidity dependence for  $\Upsilon(1S)$ .
- ALICE p-Pb results :
  - $R_{pA}$  result shows an increasing suppression of the  $J/\psi$  and  $\Upsilon(1S)$  yield towards forward  $y$
  - pure nuclear shadowing and/or energy loss seem to reasonably describe the data, indicating that final state absorption for the  $J/\psi$  may indeed be negligible at LHC energies
  - $\psi(2S)$  is significantly more suppressed than the  $J/\psi$  in both  $y$  regions
    - A similar  $p_T$  dependence of  $\psi(2S)$  suppression, with respect to the  $J/\psi$ , is observed within uncertainties at forward rapidity
    - at backward- $y$ ,  $\psi(2S)$  suppression shows an increase, with event activity, stronger than  $J/\psi$
    - initial state nuclear effects (shadowing, energy loss) alone cannot account for the modification of the  $\psi(2S)$  yields
    - final state effects, as the resonance break-up with interactions with CNM, are unlikely, because of short cc pair crossing time. Other final states effects as the cc pair interaction with the hadronic medium created in p-Pb collisions should be considered

# Thank You

# J/ψ elliptic flow in Pb-Pb

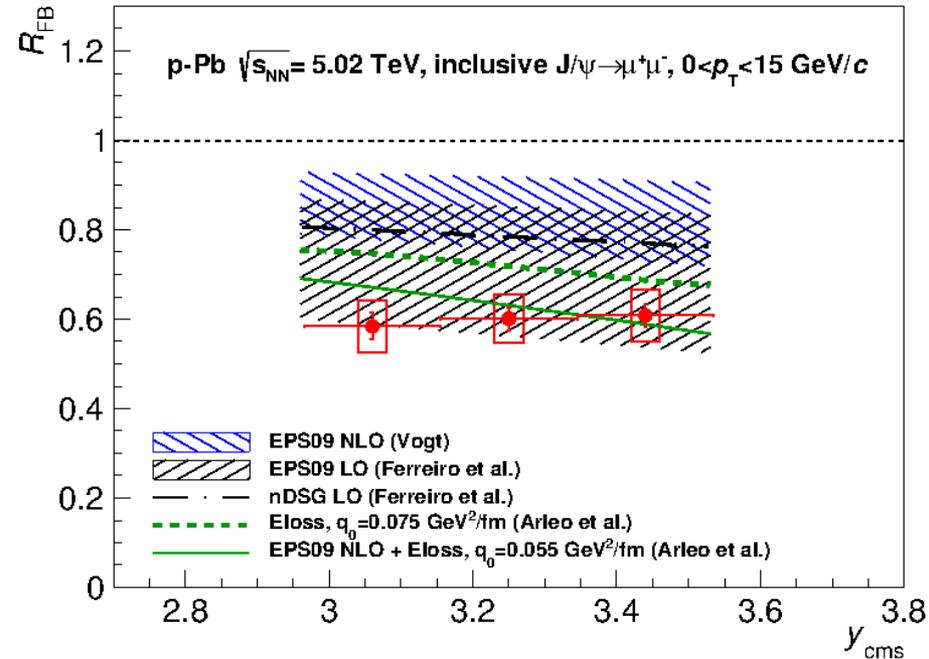
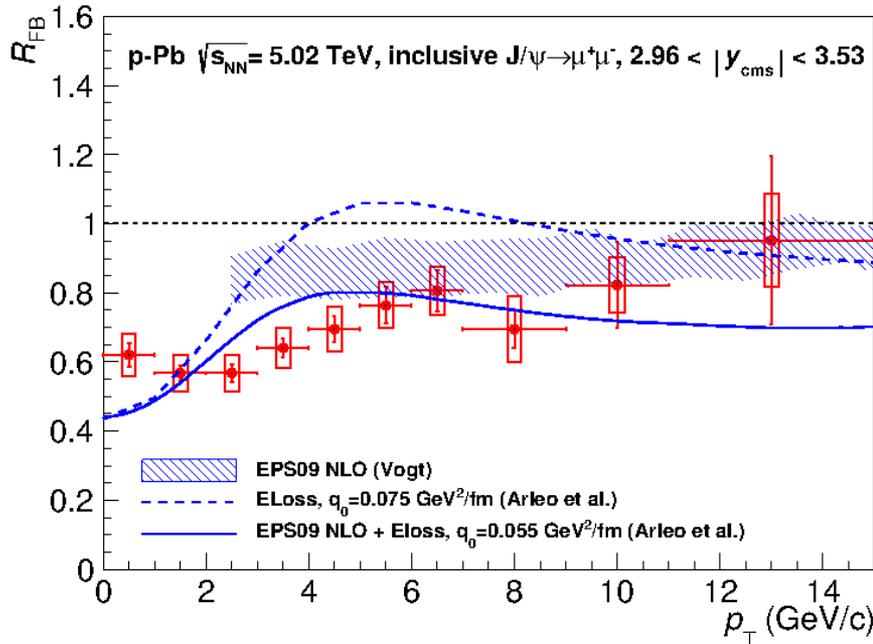
ALICE coll, Phys. Rev. Lett. 111, 162301 (2013)



- Hint of non-zero  $v_2$
- Complement  $R_{AA}$  results: a significant  $v_2$  and less suppression with respect to RHIC, SPS are indications for an observation of (re)combination from charm quarks in the QGP phase
- In qualitative agreement with transport models with 50% regeneration

# J/ψ $R_{FB}$ in p-Pb

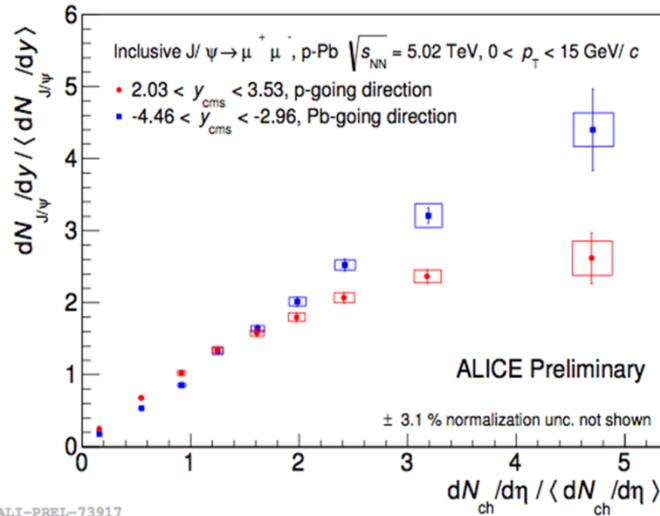
The  $R_{FB}$  rapidity dependence has also been investigated



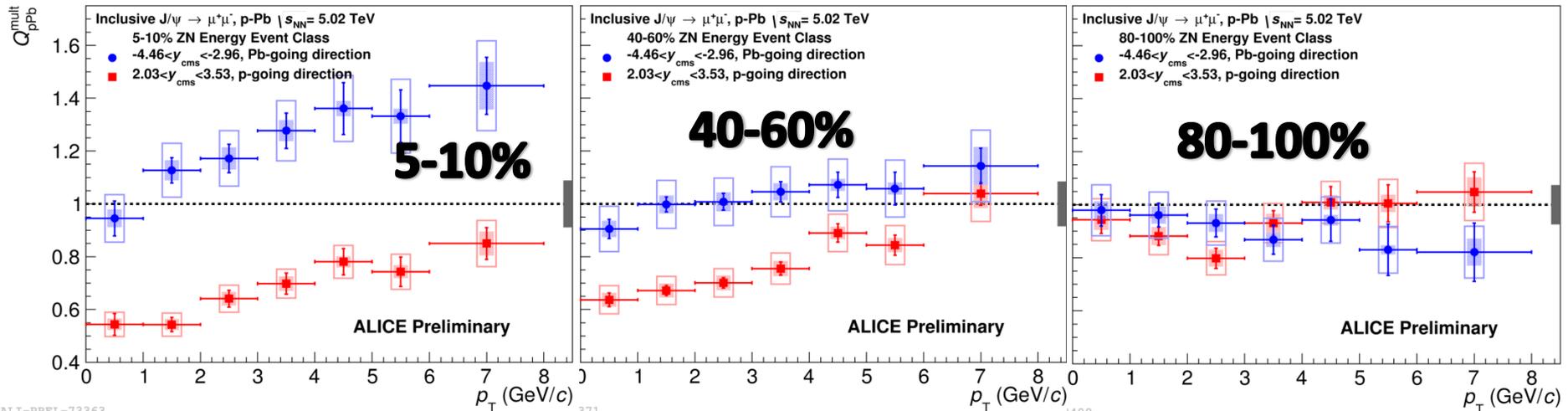
ALICE Coll. JHEP02(2014)073

- Comparison with theoretical models confirms previous observations done on the  $y$ -integrated results
- Calculations including both shadowing and energy loss seem consistent with the data

# J/ψ as function of multiplicity in p-Pb



ALI-PREL-73917



ALI-PREL-73363

371

1409