

**Light Cluster Emission in Heavy Ion
Collisions and the Symmetry Energy**

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Organization of the following program points:

- 1. Comments on Light Cluster Emission**
- 2. Open discussion on aspects of clustering in heavy ion collisions**
- 3. Concluding remarks**

Outline for 1.:

- density and momentum dependence of the symmetry potential**
- pre-equilibrium emission of nucleons and light clusters as a probe of the symmetry potential**
- results from Xe+Sn (CT) and Sn+Sn (MSU)**
- question of light cluster dynamics (→ following discussion)**

with

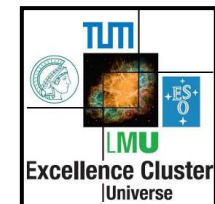
Malgorzata Zielinska-Pfabe, Piotr Decowski⁺ (Smith Coll., USA),

Maria Colonna (LNS Catania),

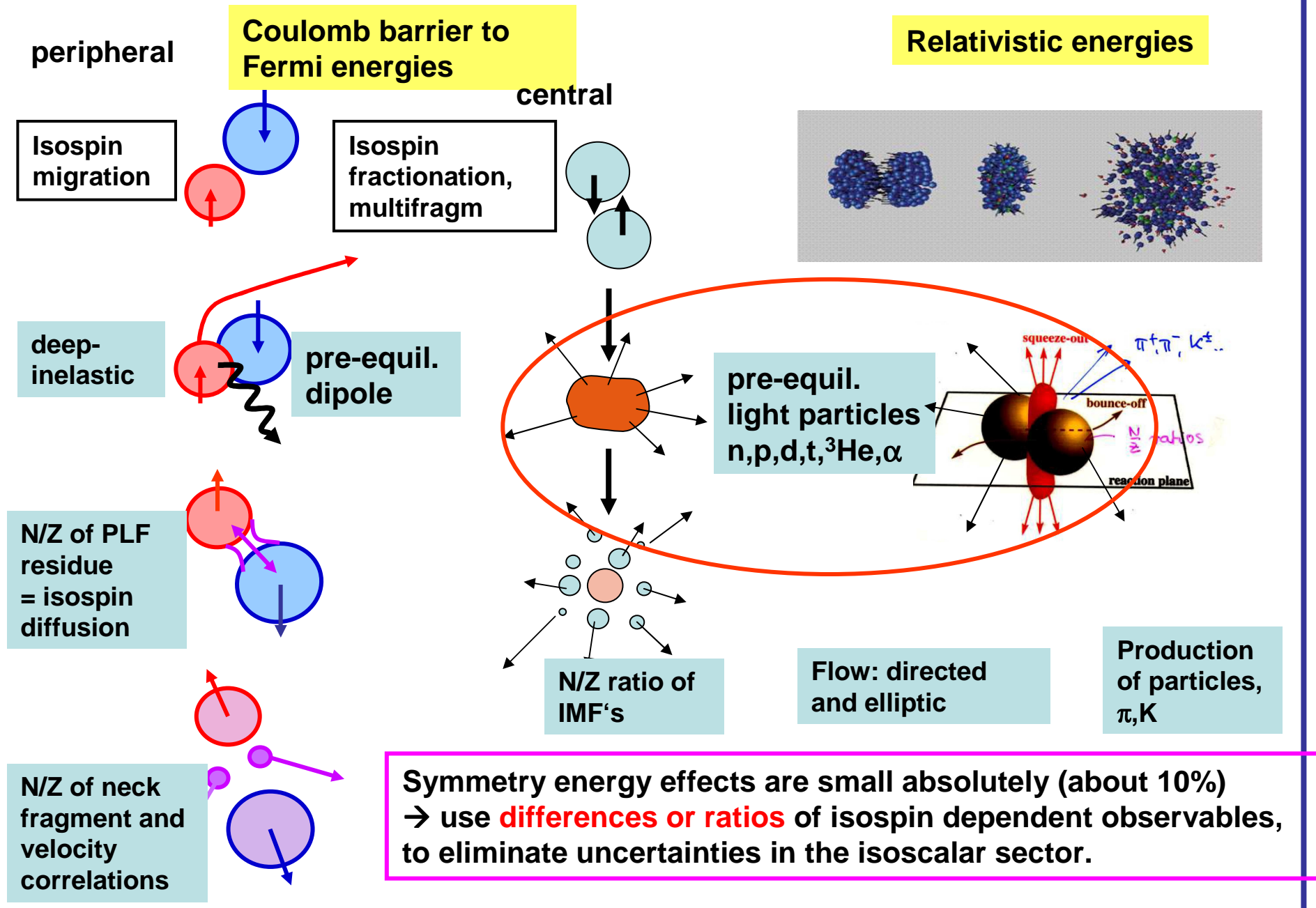
with encouragement from Remi Bougault (LPC Caen), Abdou

Chbihi (GANIL, Caen)

⁺deceased



Isospin sensitive Observables in Heavy Ion Collisions



Symmetry energy and symmetry potentials (effective masses)

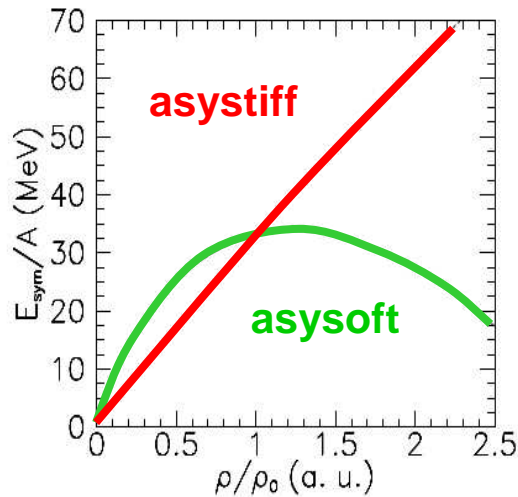
Momentum dependence of symmetry potential (isoscalar and isovector)

$$\varepsilon(\rho[f(\rho, k)]; \delta) \rightarrow U(\rho, k; \delta) = \underbrace{\frac{\partial \varepsilon(\rho, \delta)}{\partial f(\rho, k)}}_{U_\tau(\rho, k)} = U_0(\rho, k) + U_{sym}(\rho, k)(\tau\delta) + \dots$$

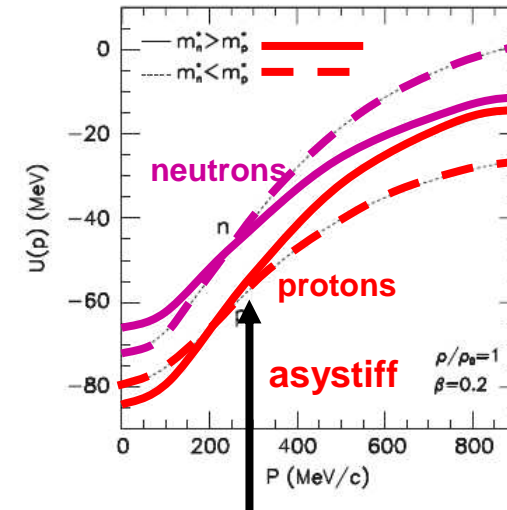
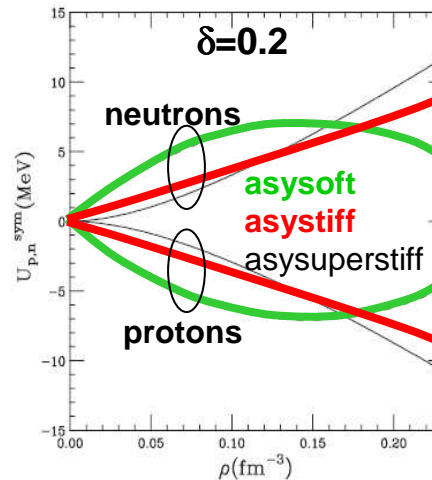
Effective proton-neutron mass splitting

$$\frac{m^*_\tau}{m} = \left(1 + \frac{m}{\hbar^2 k} \frac{\partial U_\tau}{\partial k} \right)^{-1}$$

Symmetry energy



Neutron/proton potentials as a fct of density for $p=p_F$ (left) momentum for $p=p_0$ (right)



crossing of potentials

→ leads to crossing of emission spectra

Symmetry energy and symmetry potentials (effective masses)

B.A. Li, X. Han, PLB727, 276 (2013)

establish a relation $\{E_{\text{sym}}, L\} \leftrightarrow \{U_{\text{sym}}, m_{n,p}^*\}$ using the Hugenholtz-Van Hove theorem

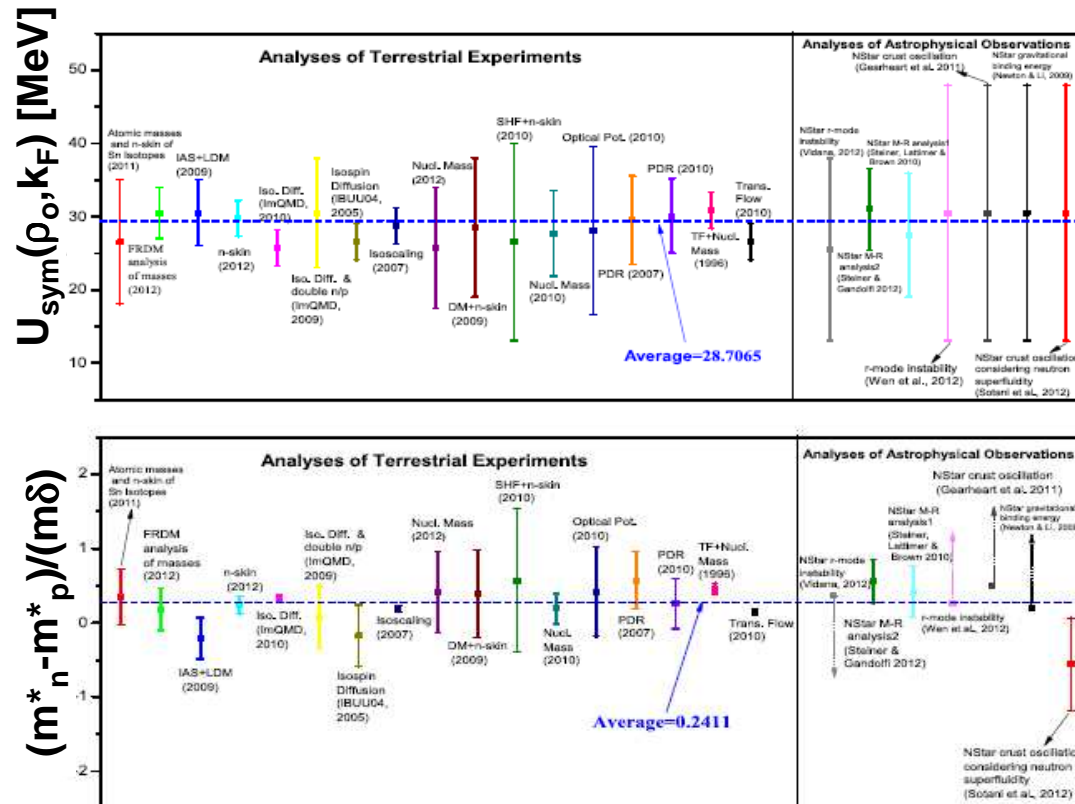
$$E_{\text{sym}}(\rho) = \frac{1}{3} \frac{\hbar^2 k_F^2}{2m_0^*} + \frac{1}{2} U_{\text{sym}}(\rho, k_F),$$

$$L(\rho) = \frac{2}{3} \frac{\hbar^2 k_F^2}{2m_0^*} + \frac{3}{2} U_{\text{sym}}(\rho, k_F) + \frac{\partial U_{\text{sym}}}{\partial k} \Big|_{k_F k_F},$$

$$U_{\text{sym}}(\rho_0, k_F) = 2 \left[E_{\text{sym}}(\rho_0) - \frac{1}{3} \frac{m}{m_0^*} E_F(\rho_0) \right],$$

$$\frac{dU_{\text{sym}}}{dk} \Big|_{k_F}(\rho_0) = \left[L(\rho_0) - 3E_{\text{sym}}(\rho_0) + \frac{1}{3} \frac{m}{m_0^*} E_F(\rho_0) \right] / k_F,$$

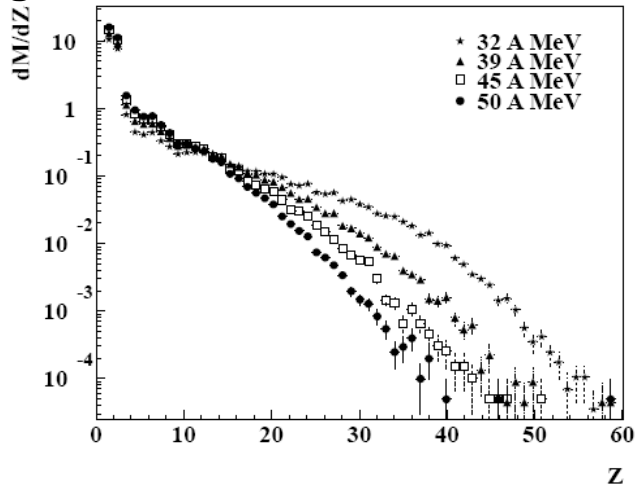
Use existing determinations in the literature for $\{E_{\text{sym}}, L\}$ to obtain values for $\{U_{\text{sym}}, m_{n,p}^*\}$



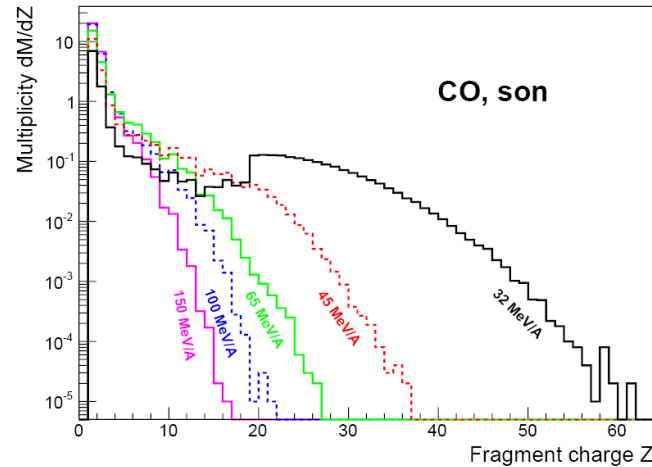
- On the average $m_n^* > m_p^*$,
 - mainly determined by nuclear structure and astrophysics,
 - however, error bars are large.

Study of Light Fragment Emission: $^{136,124}\text{Xe} + ^{124,112}\text{Sn}$, $E = 32, \dots, 150 \text{ A MeV}$

fragment distributions at different energies: exp (left), these calculations (right)



S. Hudan, et al., PRC67, 064613



H.H. Wolter, et al., EPJ Web of Conf. &&, 03097 (2014)

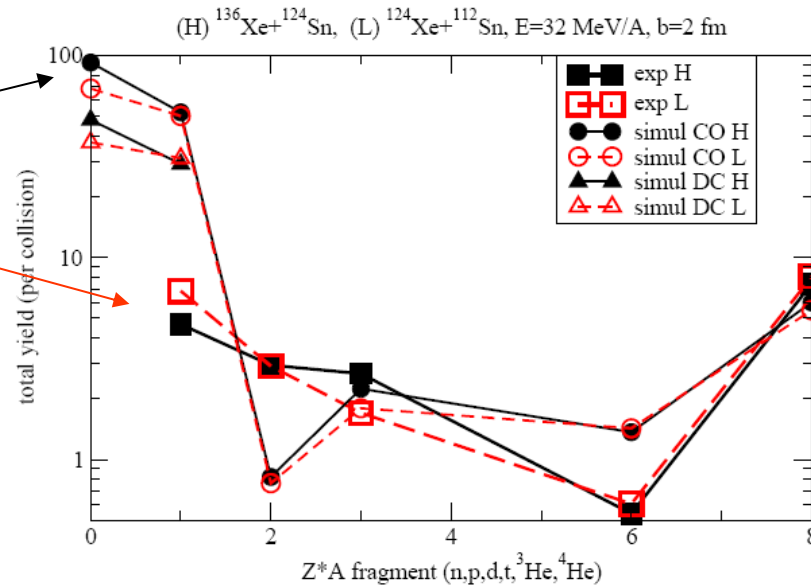
Cluster recognition: (both a-posteriori)

DC: „density cut“ coordinate space coalescence with $\rho_c > \rho_0/8$

CO: „coalescence“ in coordinate and momentum space, $r_0=1.6\text{fm}$, $p_0=1.3\text{fm}^{-1}$

overpredict nucleons with CO

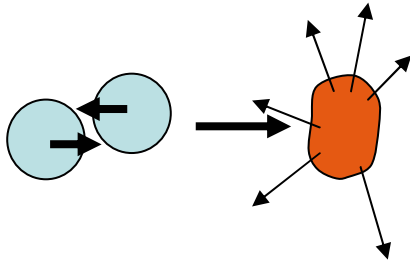
data R. Bougault, A. Chbihi (Ganil, prelim, IWM11)



H: $^{136}\text{Xe} + ^{124}\text{Sn}$
L: $^{124}\text{Xe} + ^{112}\text{Sn}$

general pattern ok, but differences in detail

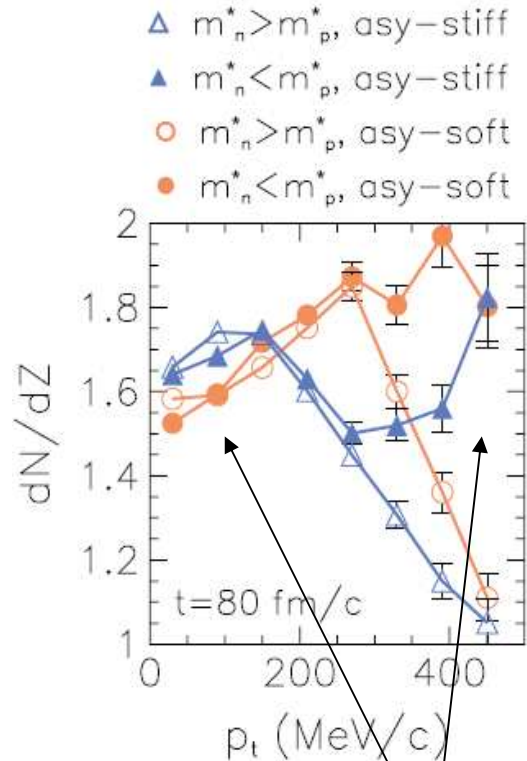
Ratios of emitted pre-equilibrium particles in central collisions



Early emitted neutrons and protons reflect difference in potentials in expanded source, esp. ratio $Y(n)/Y(p)$.

Previous studies:

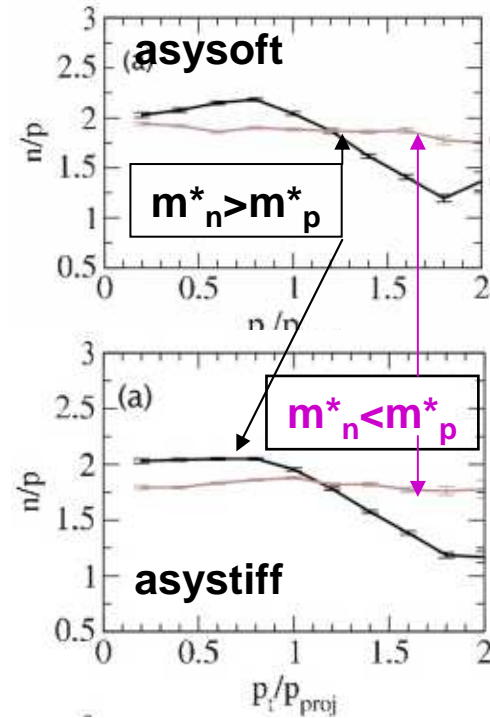
$^{132}\text{Sn} + ^{124}\text{Sn}$, 100 A MeV



splitting dep. on stiffness of symm energy resp. effect, mass

(J. Rizzo et al., PRC72,064609 (2005))

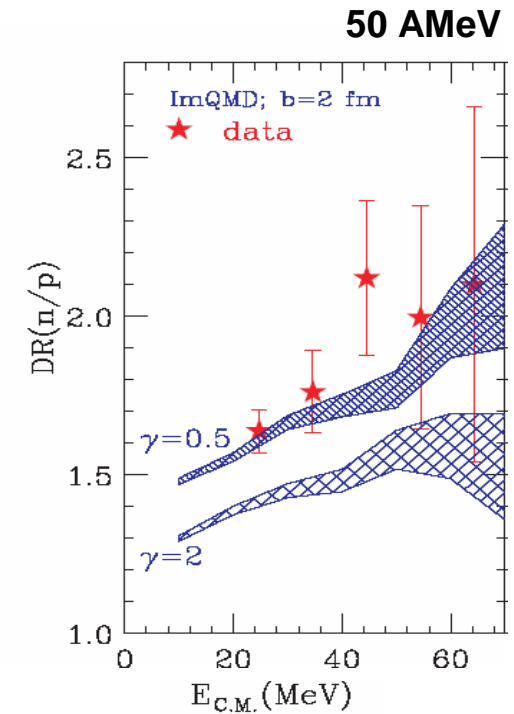
$^{197}\text{Au} + ^{197}\text{Au}$, 600 A MeV



essentially only dep. on eff. mass

V.Giordano, et al., PRC 81(2010))

Double Ratios $\frac{^{124}\text{Sn} + ^{124}\text{Sn}}{^{112}\text{Sn} + ^{112}\text{Sn}}$



no effective mass depend.

M. Famiano, et al., PRL 97,052701 (06)
M.B.Tsang, et al., PRL 102, 122701 (09)

Light Fragment Emission: $^{136,124}\text{Xe} + ^{124,112}\text{Sn}$, $E = 32, \dots, 150$ A MeV ,

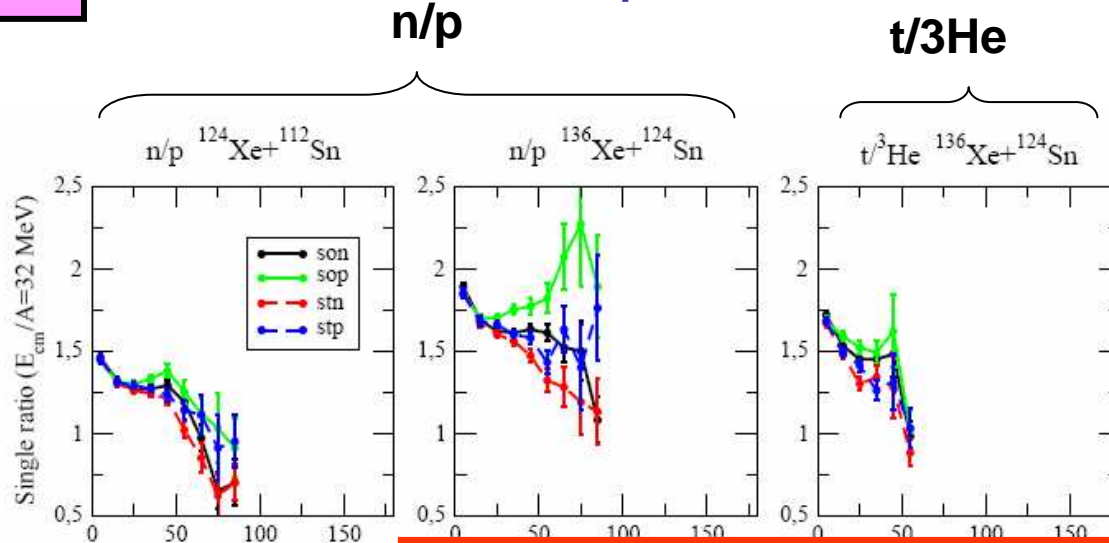
Single ratios

effects smaller for heavier clusters,
but have to improve statistics

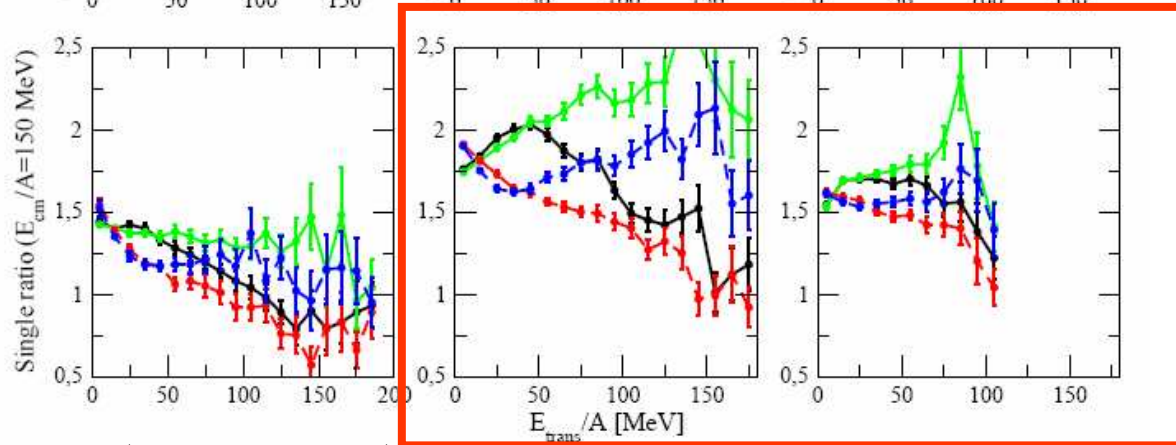
Variation
of density dep. and
effective mass splitting
with isoscalar potential
held fixed
(Bombaci-Gale-
Bertsch-DasGupta
type)

son: asysoft, $m_n^* > m_p^*$
 stn: asystiff, $m_n^* > m_p^*$
 sop: asysoft, $m_n^* < m_p^*$
 stp: asystiff, $m_n^* < m_p^*$

E=32 A MeV



E=150 A MeV



look in more detail
on next page

neutron poor
 $^{124}\text{Xe} + ^{112}\text{Sn}$

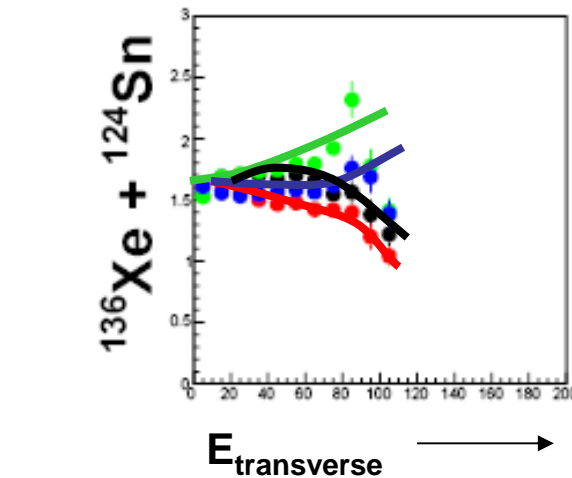
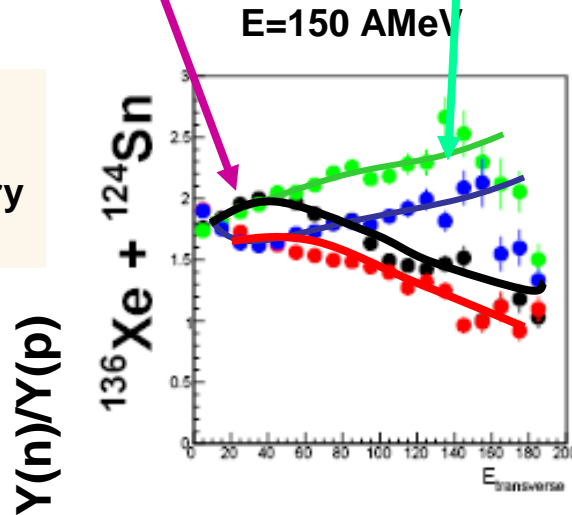
neutron rich
 $^{136}\text{Xe} + ^{124}\text{Sn}$

effects smaller for neutron poor

Asy-EOS dominates for slow particles; asysoft has larger repulsion at lower densities

Effective mass dominates for fast particles; smaller eff. mass favors emission

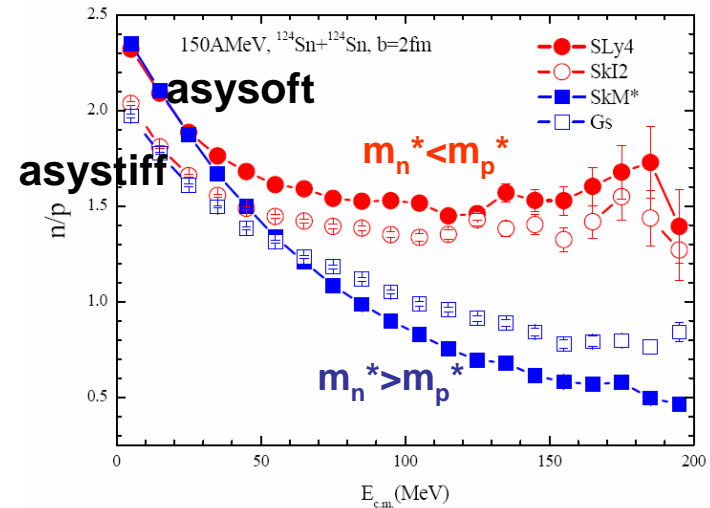
Possibility to separate density and momentum dependence of symmetry energy



Effect also exists for light clusters (easier to measure) but somewhat reduced

son: asysoft, $m_n^* > m_p^*$
 stn: asystiff, $m_n^* > m_p^*$
 sop: asysoft, $m_n^* < m_p^*$
 stp: asystiff, $m_n^* < m_p^*$

similar conclusions for Sn+Sn collisions (MSU)

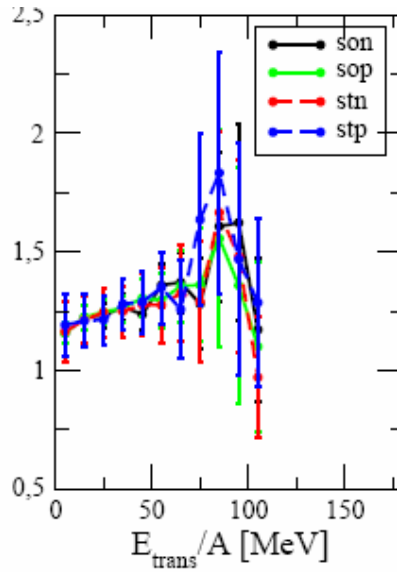
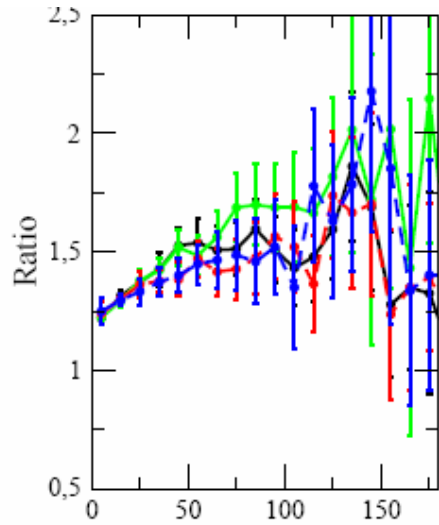


Y. Zhang, M.B.Tsang, et al., PLB 732, 186 (2014)

Double Ratios

$$\frac{n/p_{136\text{Xe}+124\text{Sn}}}{n/p_{124\text{Xe}+112\text{Sn}}}$$

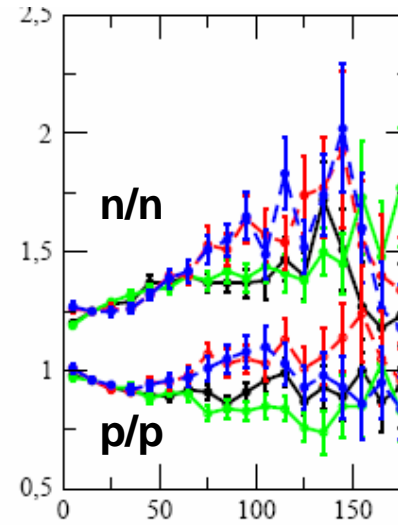
$$\frac{t/{}^3\text{He}_{136\text{Xe}+124\text{Sn}}}{t/{}^3\text{He}_{124\text{Xe}+112\text{Sn}}}$$



sensitivity much smaller

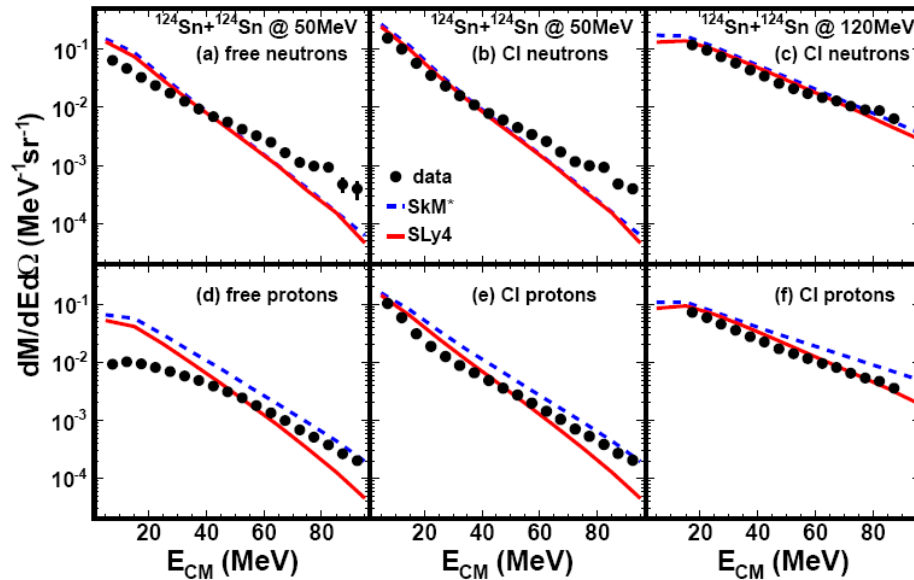
Singe Ratios

$$\frac{n_{136\text{Xe}+124\text{Sn}}}{n_{124\text{Xe}+112\text{Sn}}}, \text{ resp. }, \frac{p_{136\text{Xe}+124\text{Sn}}}{p_{124\text{Xe}+112\text{Sn}}}$$



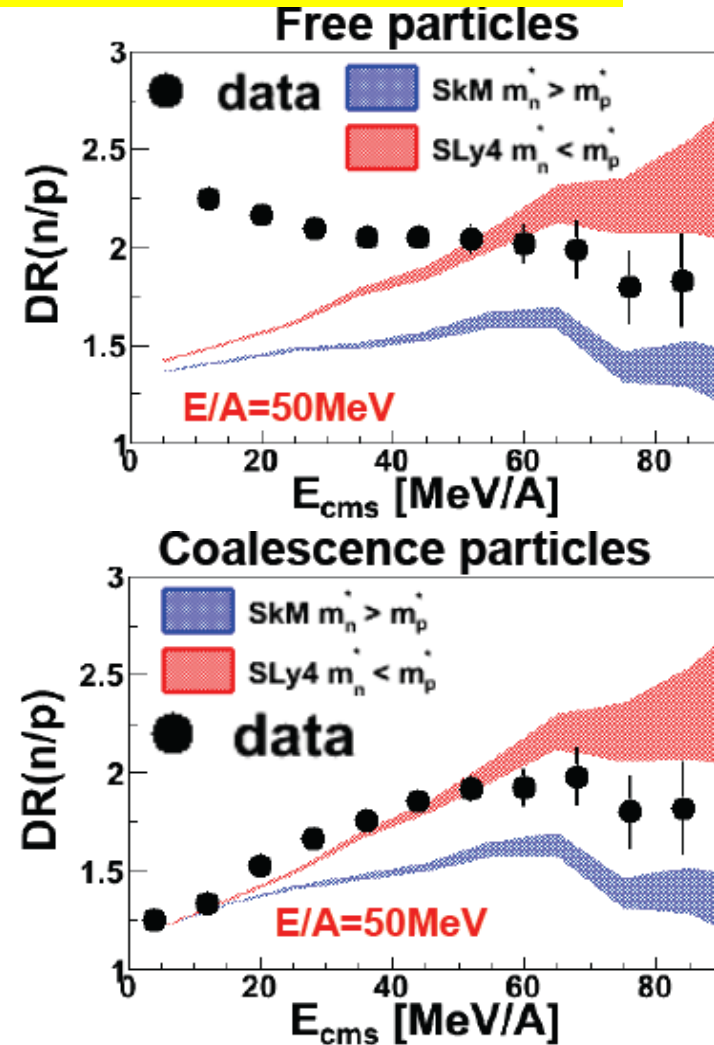
sensitivity only to asy-stiffness

Comparison with data: problem of light cluster description in transport approaches (recent work from MSU)



coalescence invariant spectra
agree better with experiment:
A poor man's substitute for not
treating light clusters properly in
the simulation

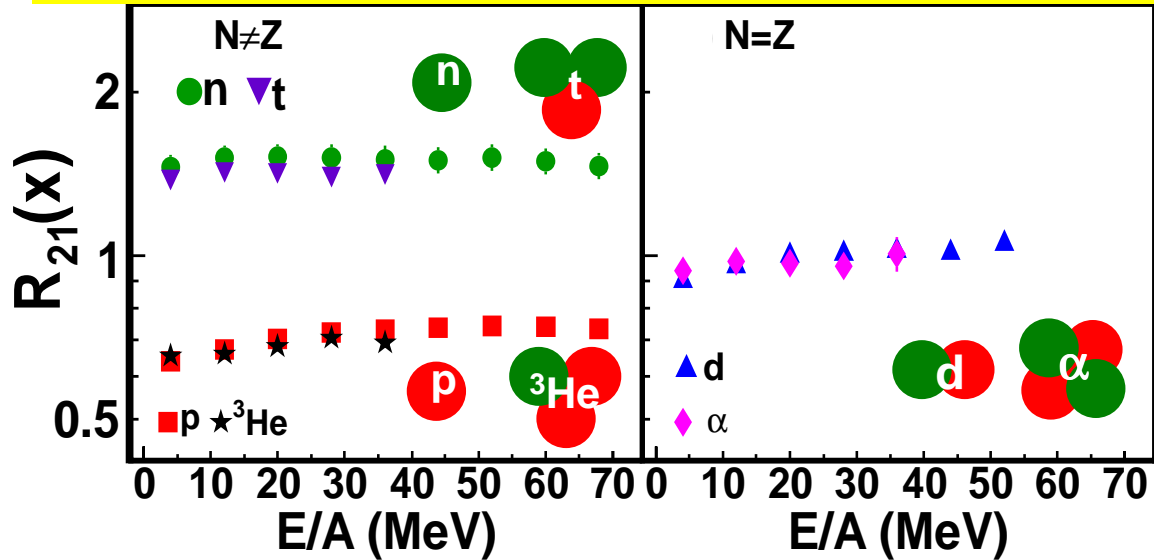
favors slightly $m_n^* < m_p^*$ (in contrast to
other analyses, see above)
→ more work required!



Y.X. Zhang, M.B.Tsang, et al., PLB 732, 186 (2014)
D.D.S.Coupland, arXiv 1406.4546

An interesting note:

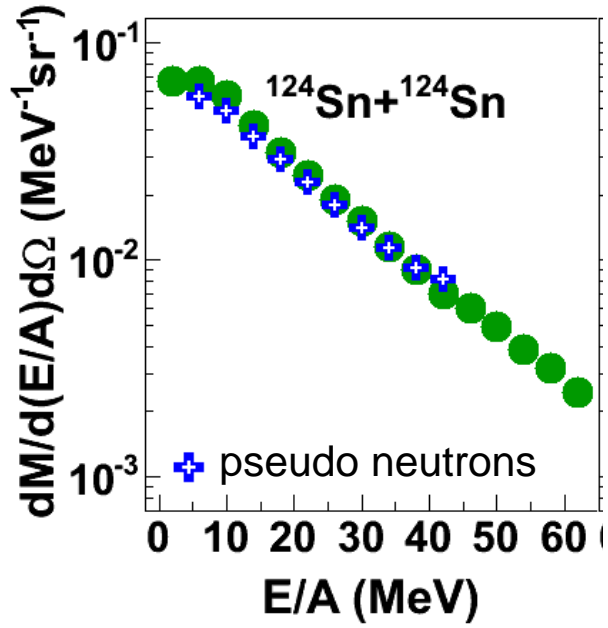
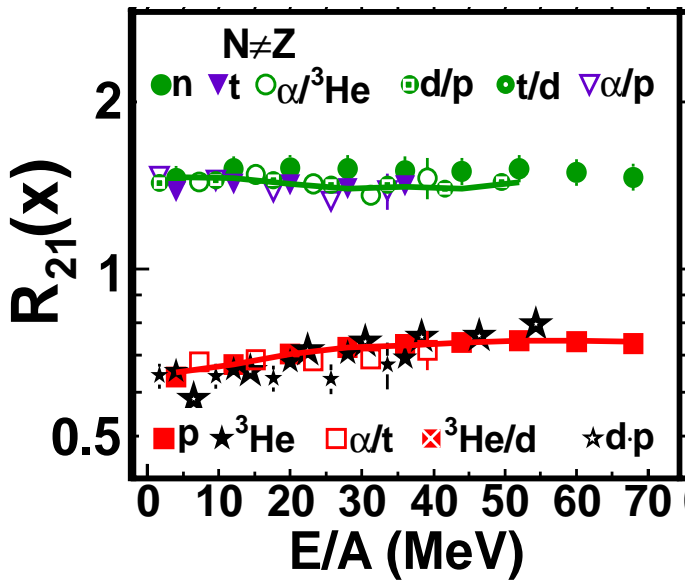
„Chemical potential scaling“ Z.Chajecski, et al., arXiv 1402.5216



$$R_{21}(N, Z) = \frac{dM_2(N, Z)}{dM_1(N, Z)}$$

expect from chemical potential dependence without correlation contributions

$$R_{21}(N, Z) = \exp\left[\frac{(N\Delta\mu_n + Z\Delta\mu_p)}{T}\right]$$



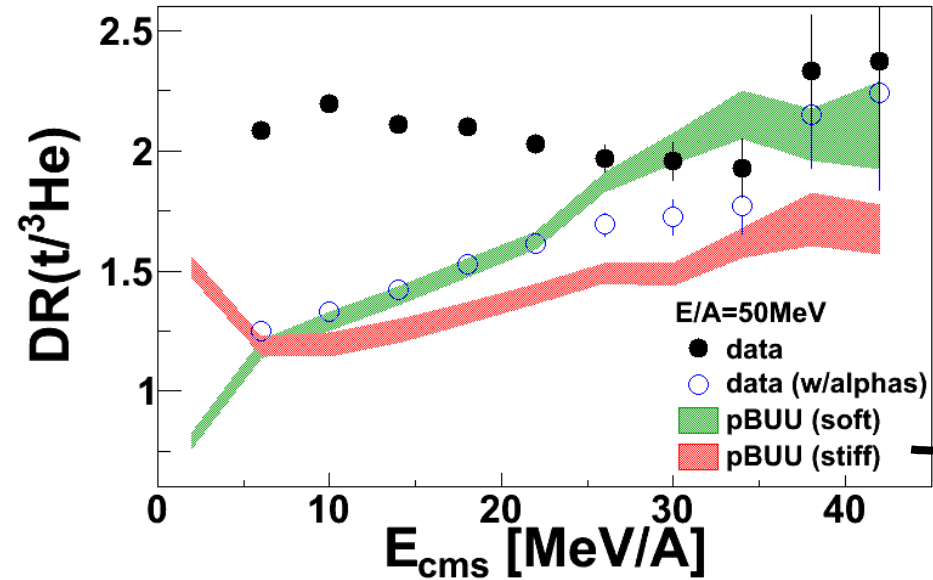
$$\Rightarrow Y(n) = \frac{Y(t)}{Y(^3\text{He})} Y(p)$$

“Pseudo neutron yields”

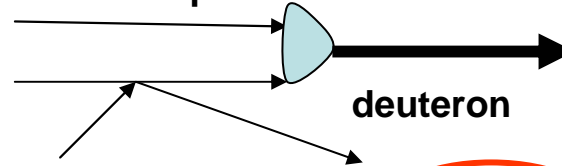
correlation effects seem to be small for light cluster yields (?)

Interdependence of light cluster yields

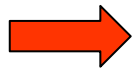
W.Lynch, INPC, Florence, 2013:
 $^{124,112}\text{Sn} + ^{124,112}\text{Sn}$, 50 A MeV;
 Z. Chajecski, NuSYM13



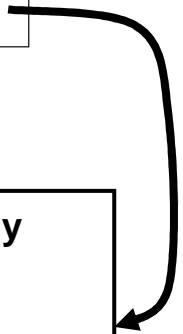
LC's as explicit degrees of freedom coupled to nucleon distribution functions by 3-body collisions of type $NNN \rightarrow ND$



P. Danielewicz and Q. Pan, PRC 46 (1992): (d, t, ^3He , **but no α !**) coupled transport eqs.



cluster yields depend on each other



Conclusions:

- pre-equilibrium nucleon and light cluster emission directly probe density and momentum dependence of symmetry potential
- the single ratio spectra may allow to separate this density and momentum dependence.
The energy range around 100 to 200 MeV is most sensitive. At higher energy sensitive mainly to momentum dependence
- more asymmetric systems and the n/p ratio are most sensitive. The t/³He ratio carries a similar information
- more precise data recently available, but not yet precise enough to determine effective masses (at some stress with other determinations)
- touches question of dynamical light cluster description in HIC
(→ see following discussion)