

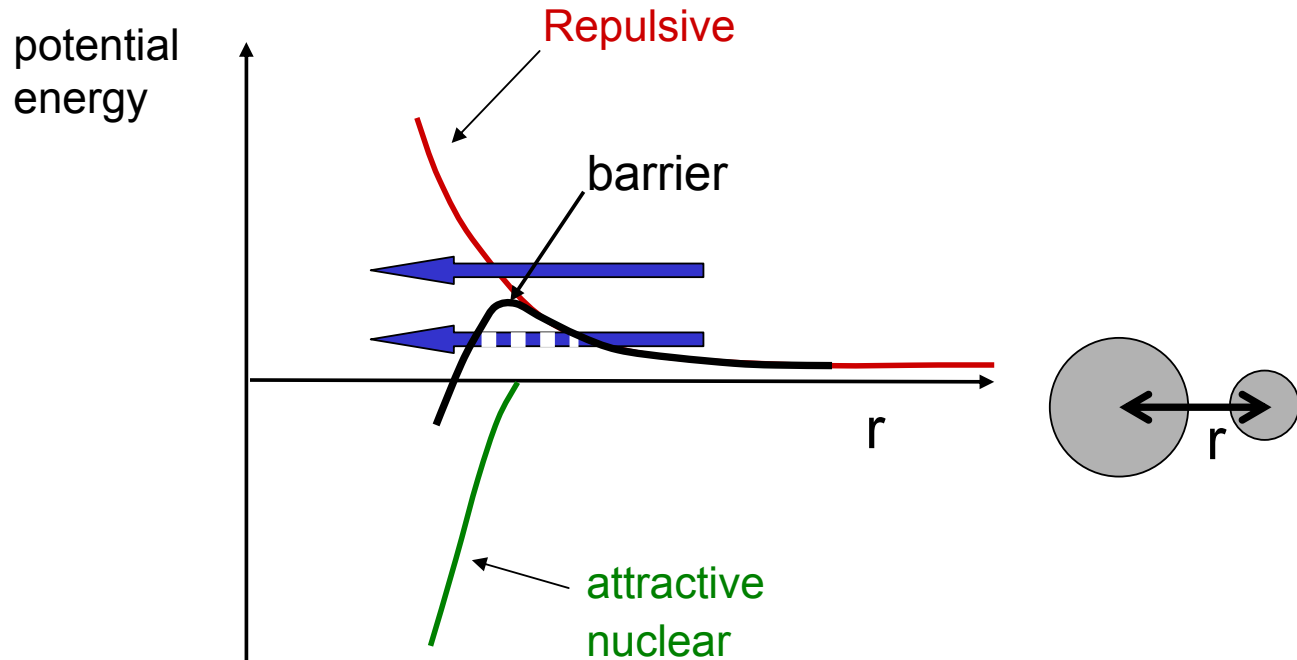
Reazioni tra ioni pesanti: dalle collisioni quasi-elastiche alla fusione completa

A.M.Stefanini

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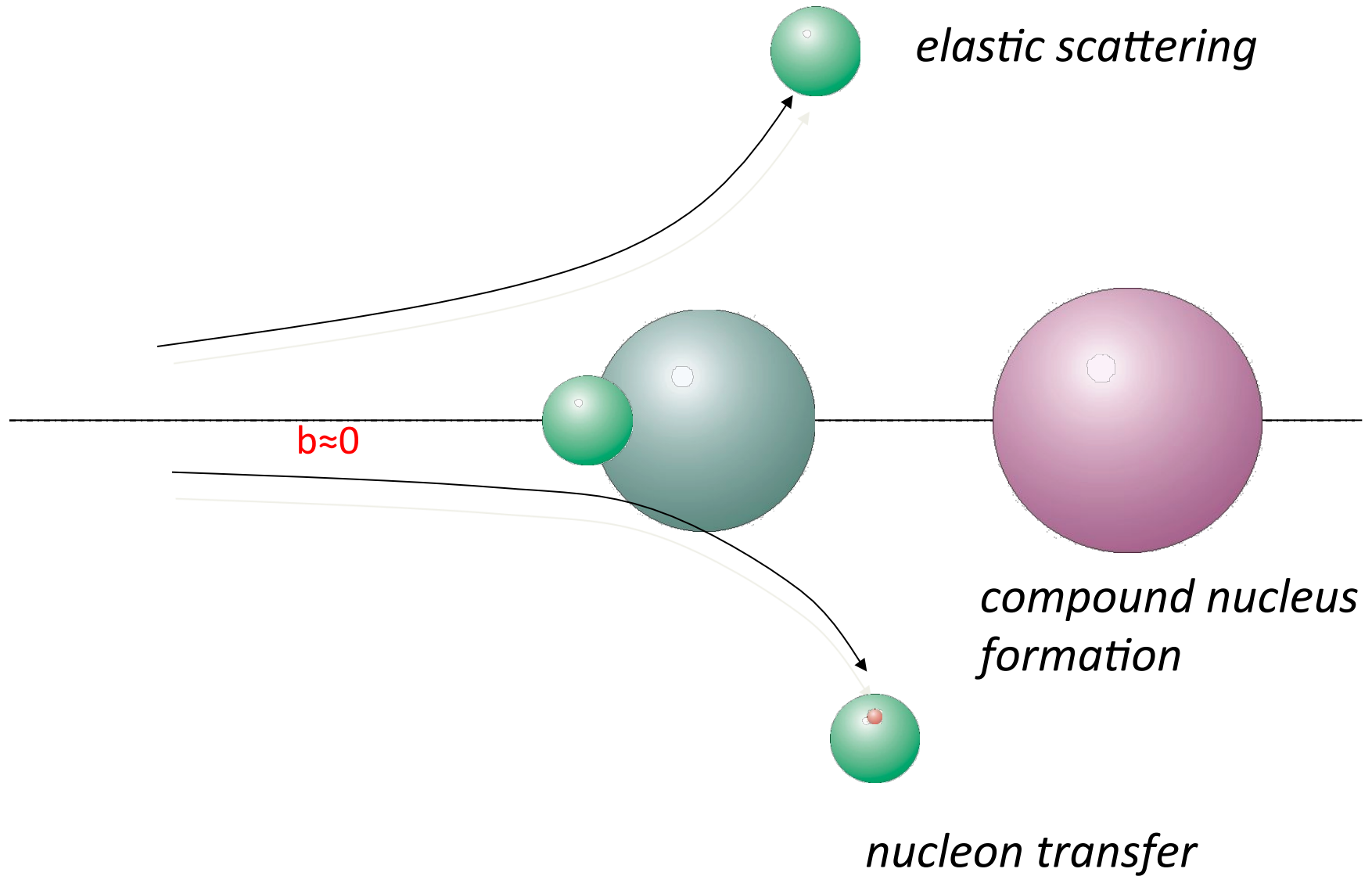


Attractive Nuclear Potential vs. Repulsive Electrostatic Potential



- Nuclear potential – a construct to account for the many body interactions of protons and neutrons
- Barrier passing or tunnelling → capture (fusion)

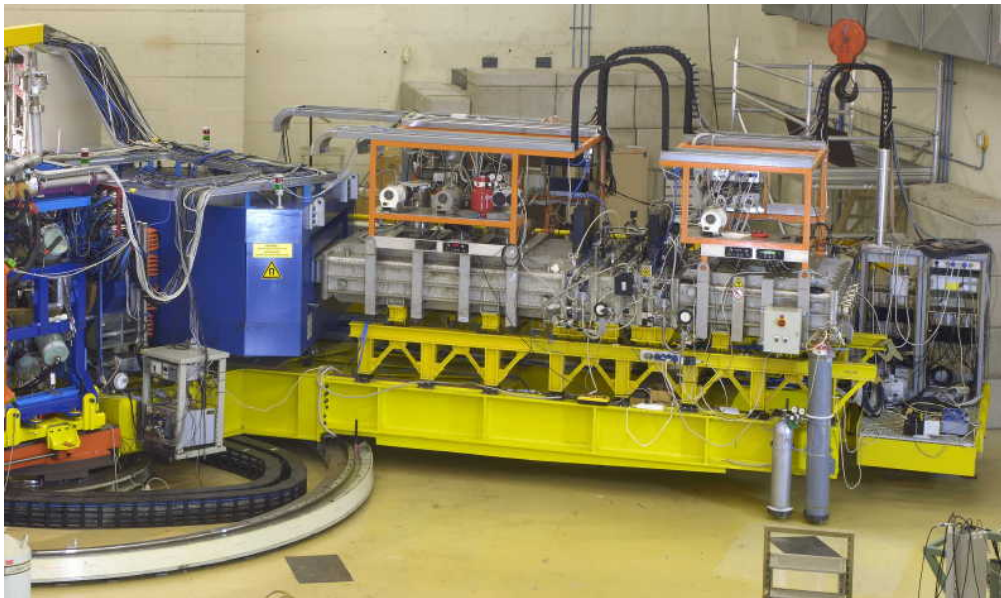
Heavy Ion Reactions @ the Coulomb Barrier



Two lines of research

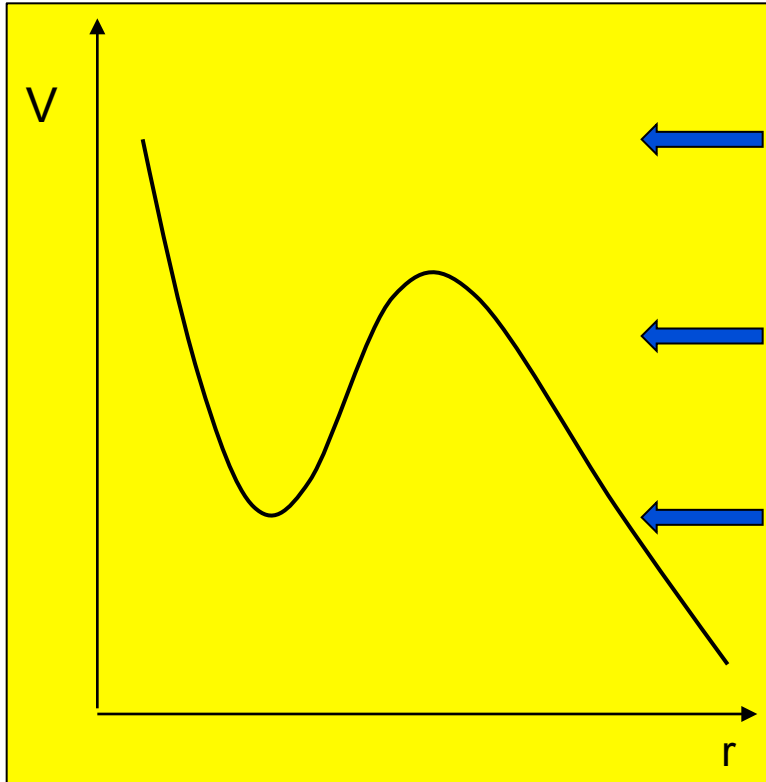


Heavy-ion fusion reactions
electrostatic beam separator
PISOLO



Nucleon transfer reactions
magnetic spectrometer
PRISMA

Nucleon transfer reactions: a smooth transition from deep-inelastic to quasi-elastic processes



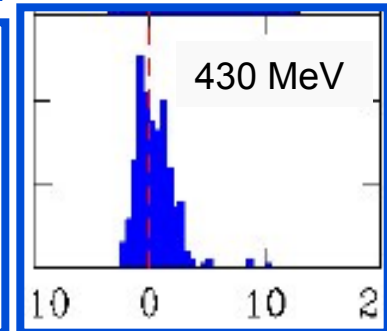
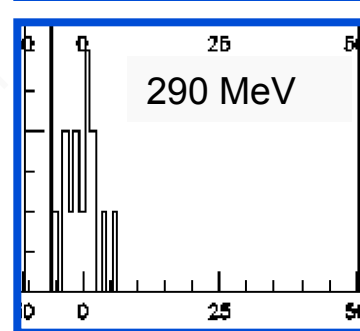
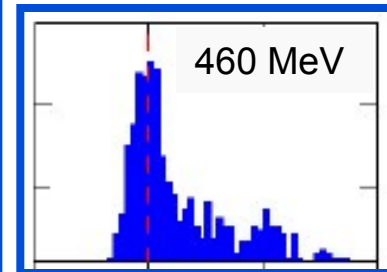
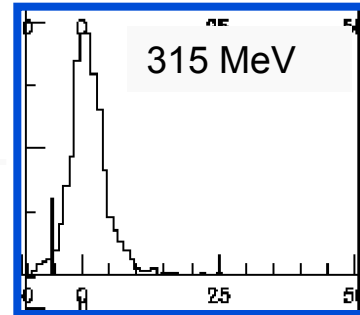
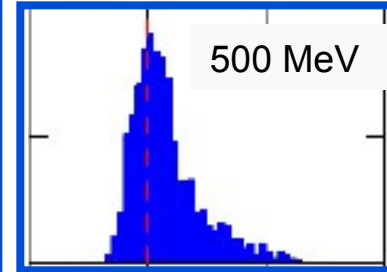
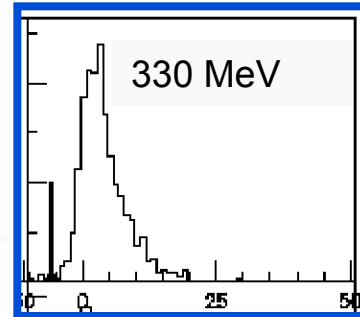
$E > E_b$

$E \sim E_b$

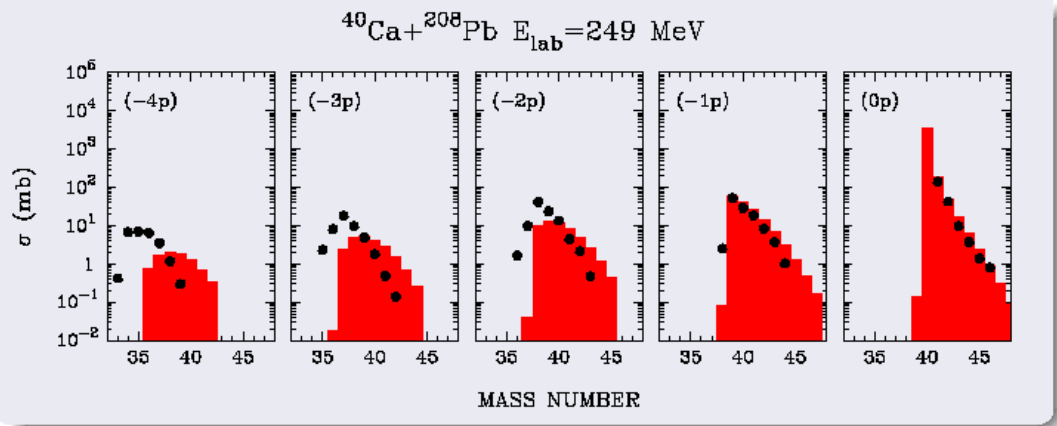
$E < E_b$

$^{96}\text{Zr}(^{40}\text{Ca}, ^{42}\text{Ca})$
 $Q_{gs} = +5.6 \text{ MeV}$

$^{116}\text{Sn}(^{60}\text{Ni}, ^{62}\text{Ni})$
 $Q_{gs} = +1.3 \text{ MeV}$



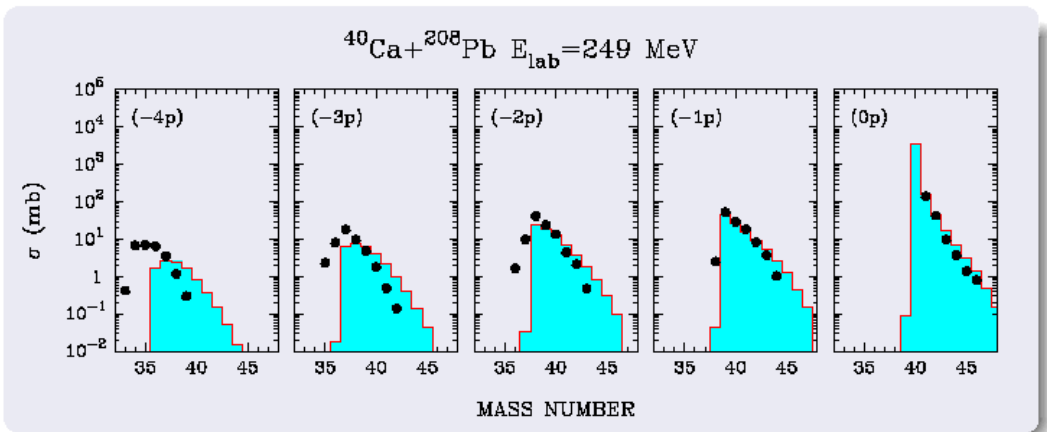
Multi-nucleon transfer: experiment vs. theory



1pt

independent
particle transfer

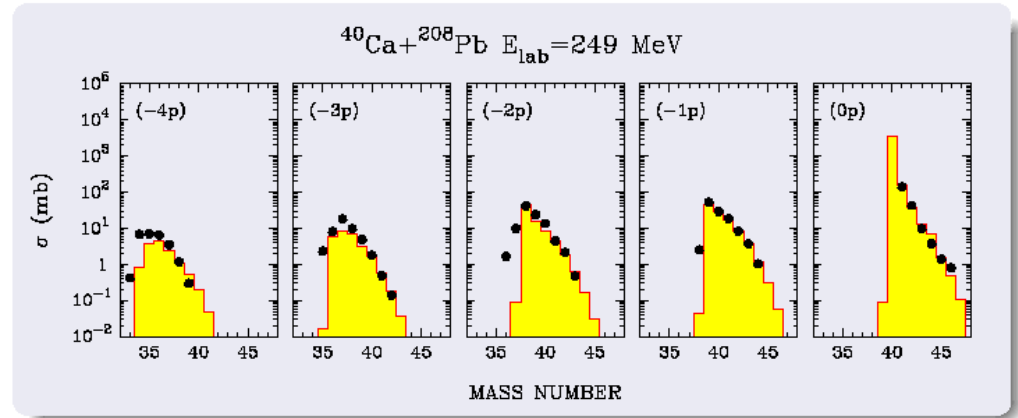
data : LNL
theory :
GRAZING code
and CWKB



+ pair transfer modes

1pt+2pt

$$F_{fi}^{pair}(r) = \beta_p \frac{\partial V^{opt}(r)}{\partial A}$$



+Evap.

L.Corradi, G.Pollarolo,
S.Szilner
J.Phys. G36(2009)113101
(Topical Review)

Sub-barrier nucleon transfer reactions

Few reaction channels are open



uncertainties related to the nuclear potential are reduced

Q-value distributions are narrow



probe nucleon correlations close to the g.s., interplay between single and multiple particle transfer

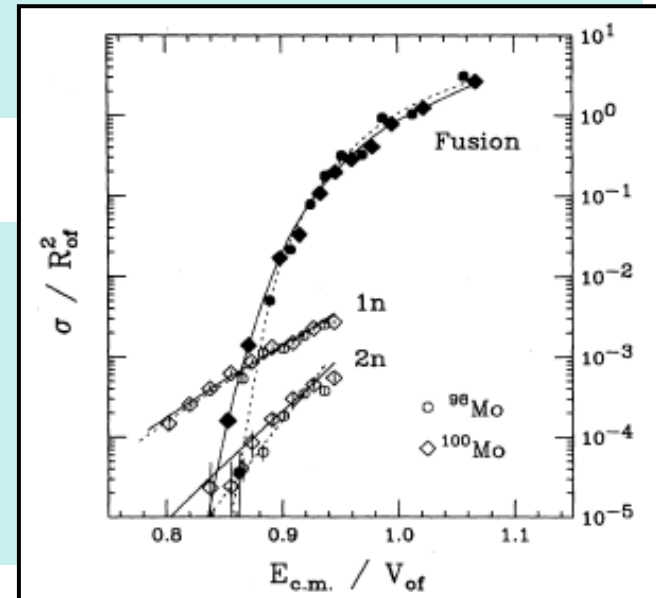
Tunneling between the colliding nuclei



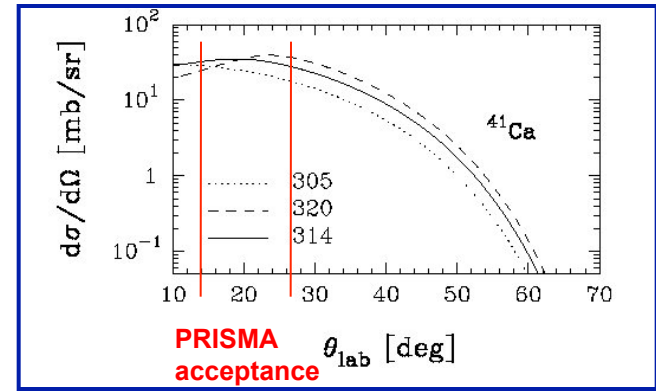
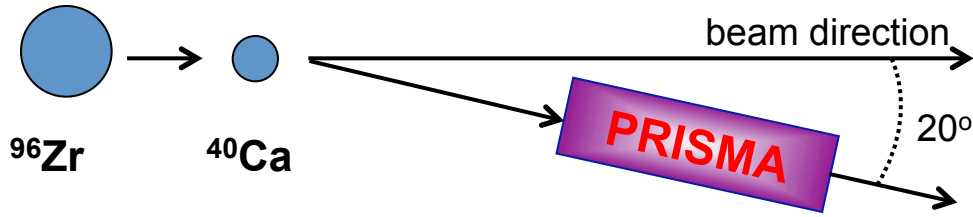
transfer and fusion are measured in an overlapping range of energy and angular momentum

but ... experimental difficulties !

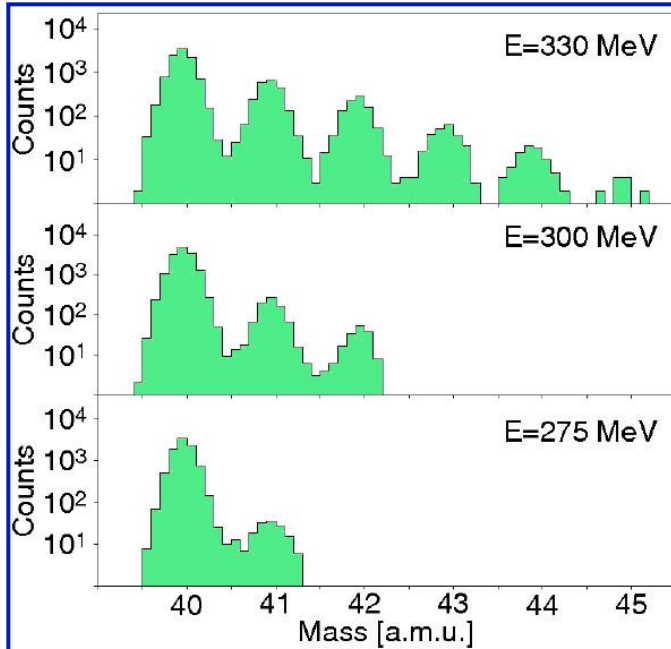
- Angular distributions are backward-peaked and projectile-like particles have low kinetic energies
- Cross sections are very small (hi efficiency needed)
- Complete identification of reaction products in A, Z and Q is difficult



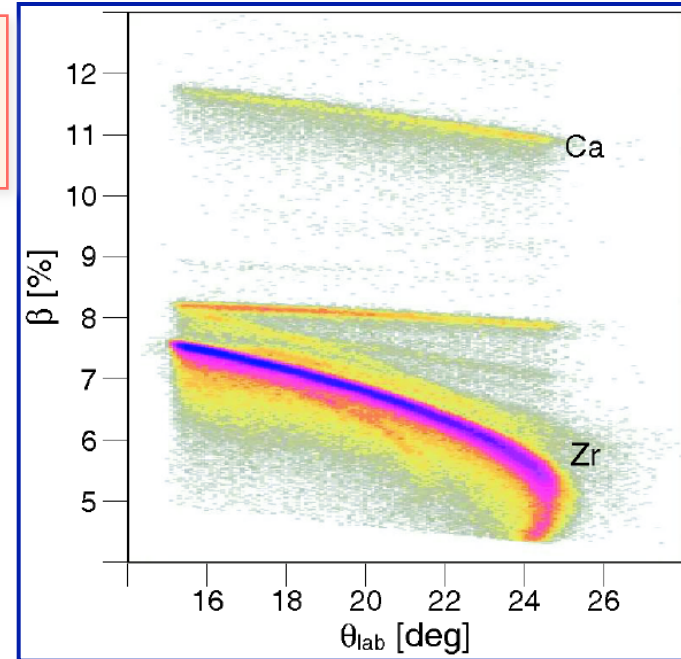
Study of the reaction $^{96}\text{Zr} + ^{40}\text{Ca}$ using the inverse kinematics



Z=20

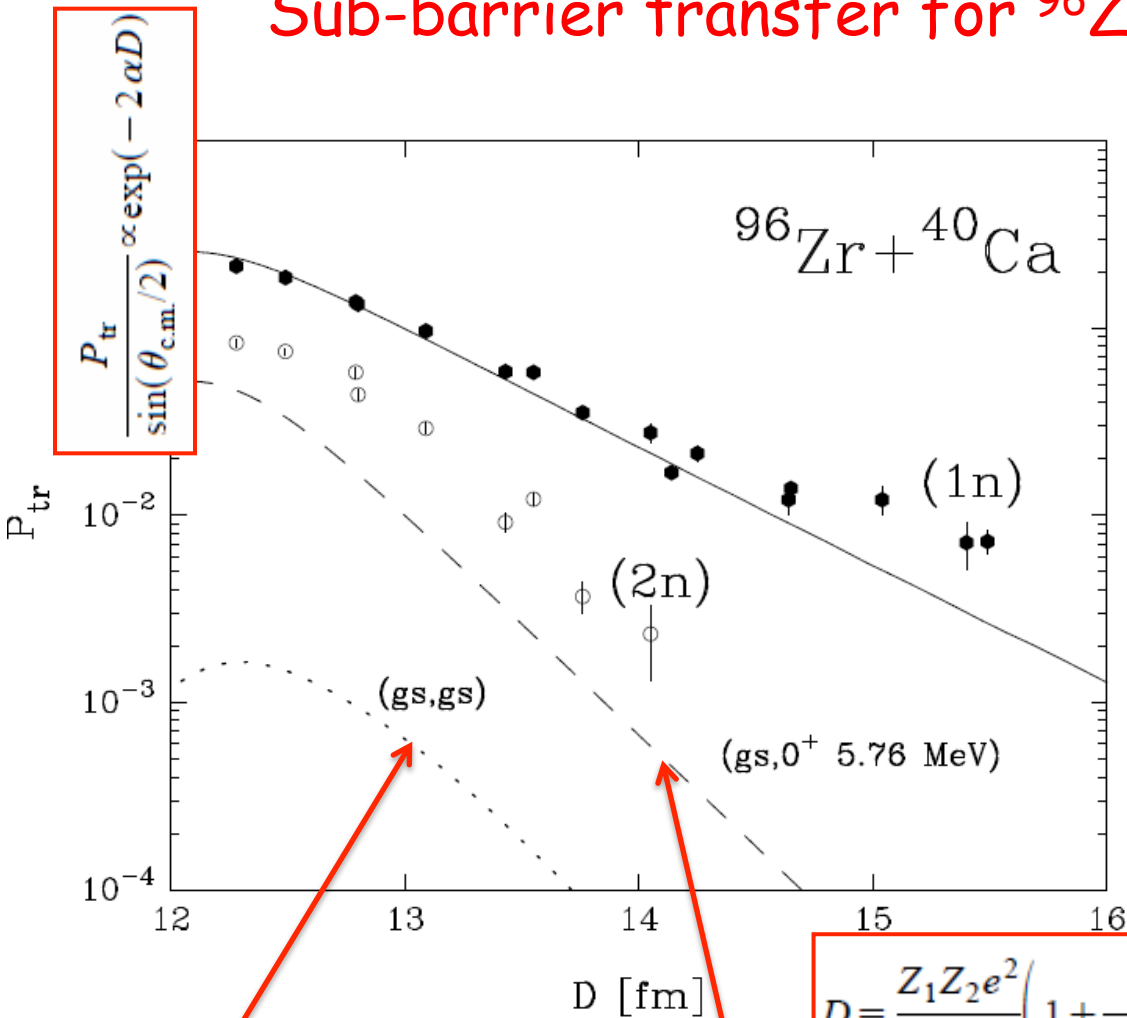


Excitation functions
 $E_{\text{beam}} = 330 \text{ MeV} - 275 \text{ MeV}$
6 MeV step



Nucleon transfer channels have been measured down to 25 % below the Coulomb barrier

Sub-barrier transfer for $^{96}\text{Zr} + ^{40}\text{Ca}$



$$\frac{P_{\text{tr}}}{\sin(\theta_{\text{c.m.}}/2)} \propto \exp(-2\alpha D)$$

+1n well reproduced by theory in slope and absolute value
+2n : theory gives the correct slope, but underestimates the cross section

Is this due to the contribution of other excited states ?

(absorption is reproduced by theory)

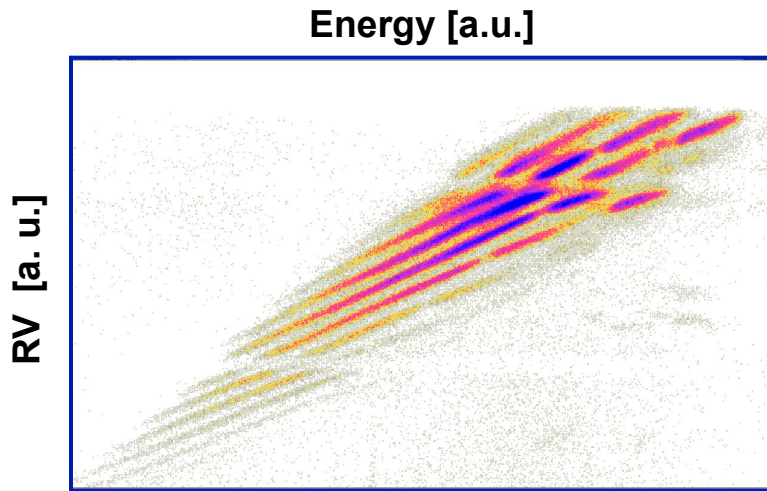
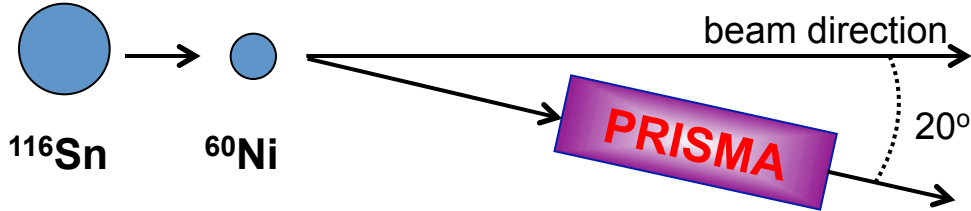
2n transfer to g.s.

2n transfer to 0^+ (5.76 MeV)

$$D = \frac{Z_1 Z_2 e^2}{2E_{\text{c.m.}}} \left(1 + \frac{1}{\sin(\theta_{\text{c.m.}}/2)} \right)$$

$$\alpha = \sqrt{\frac{2\mu B}{\hbar^2}}$$

The reaction $^{116}\text{Sn} + ^{60}\text{Ni}$ using the inverse kinematics

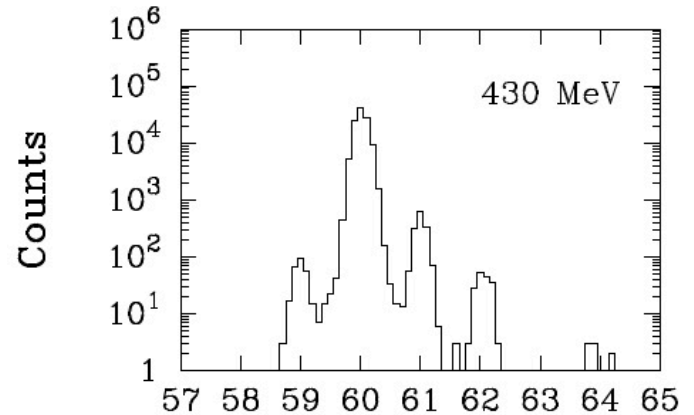
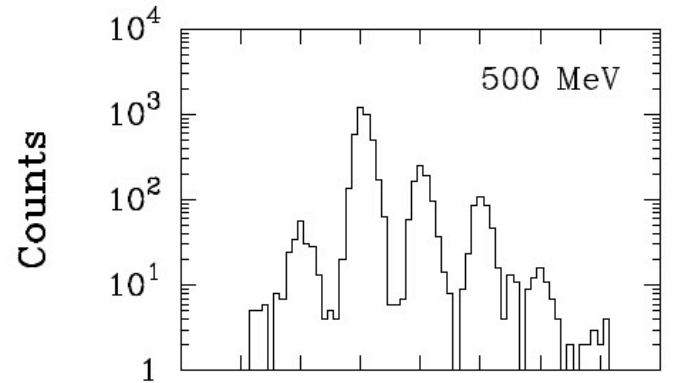


Ground state Q-values

	+1n	+2n	+3n	+4n
$^{96}\text{Zr} + ^{40}\text{Ca}$	+ 0.51	+ 5.53	+ 5.24	+ 9.64
$^{116}\text{Sn} + ^{60}\text{Ni}$	- 1.74	+ 1.31	- 2.15	- 0.24

Excitation functions

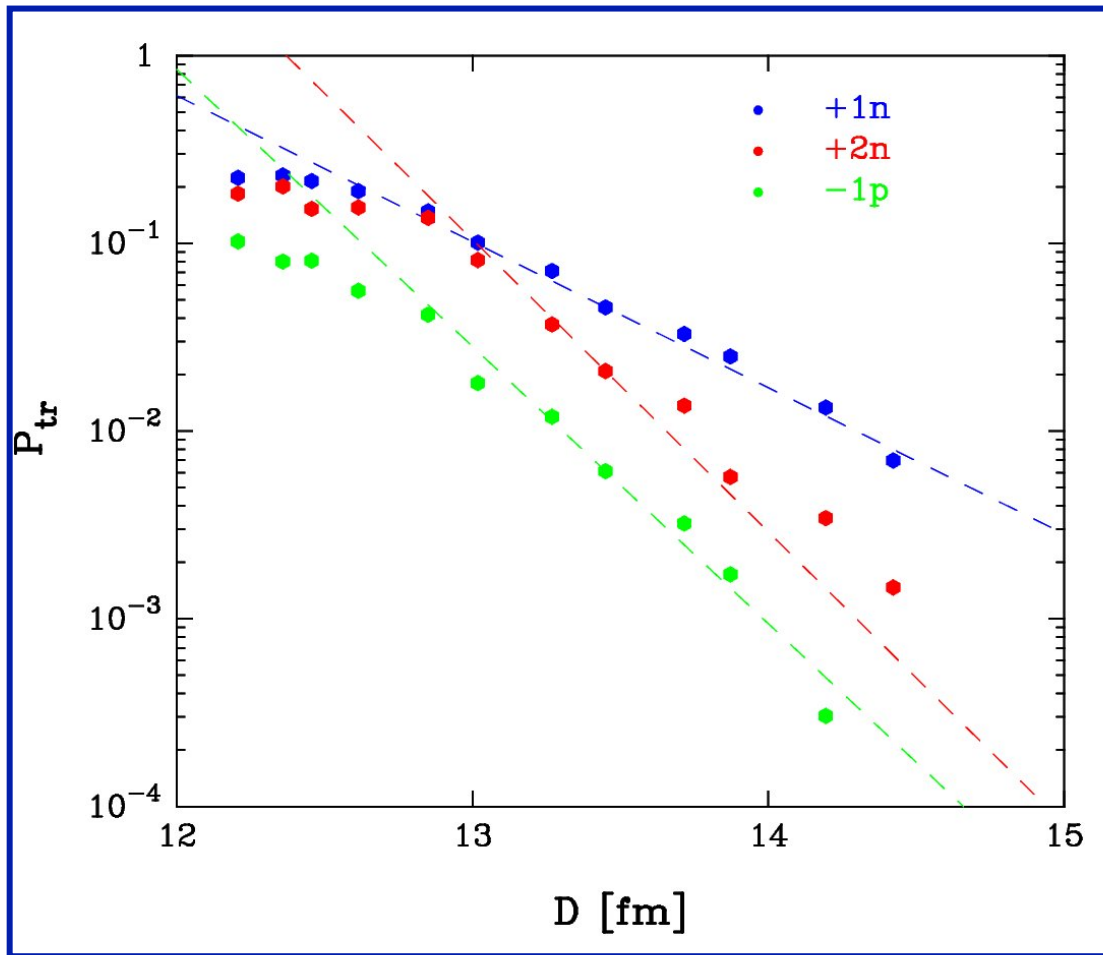
$E_{\text{beam}} = 410 \text{ MeV} - 500 \text{ MeV}$
($D \sim 12.3 \text{ to } 15.0 \text{ fm}$)



Mass [amu]

→ Most recent measurement for the case of $^{92}\text{Mo} + ^{54}\text{Fe}$: study of pair correlations populating $\pm(\text{nn})$, $\pm(\text{pp})$ and $\pm(\text{np})$ transfer channels

Transfer probabilities for $^{116}\text{Sn} + ^{60}\text{Ni}$

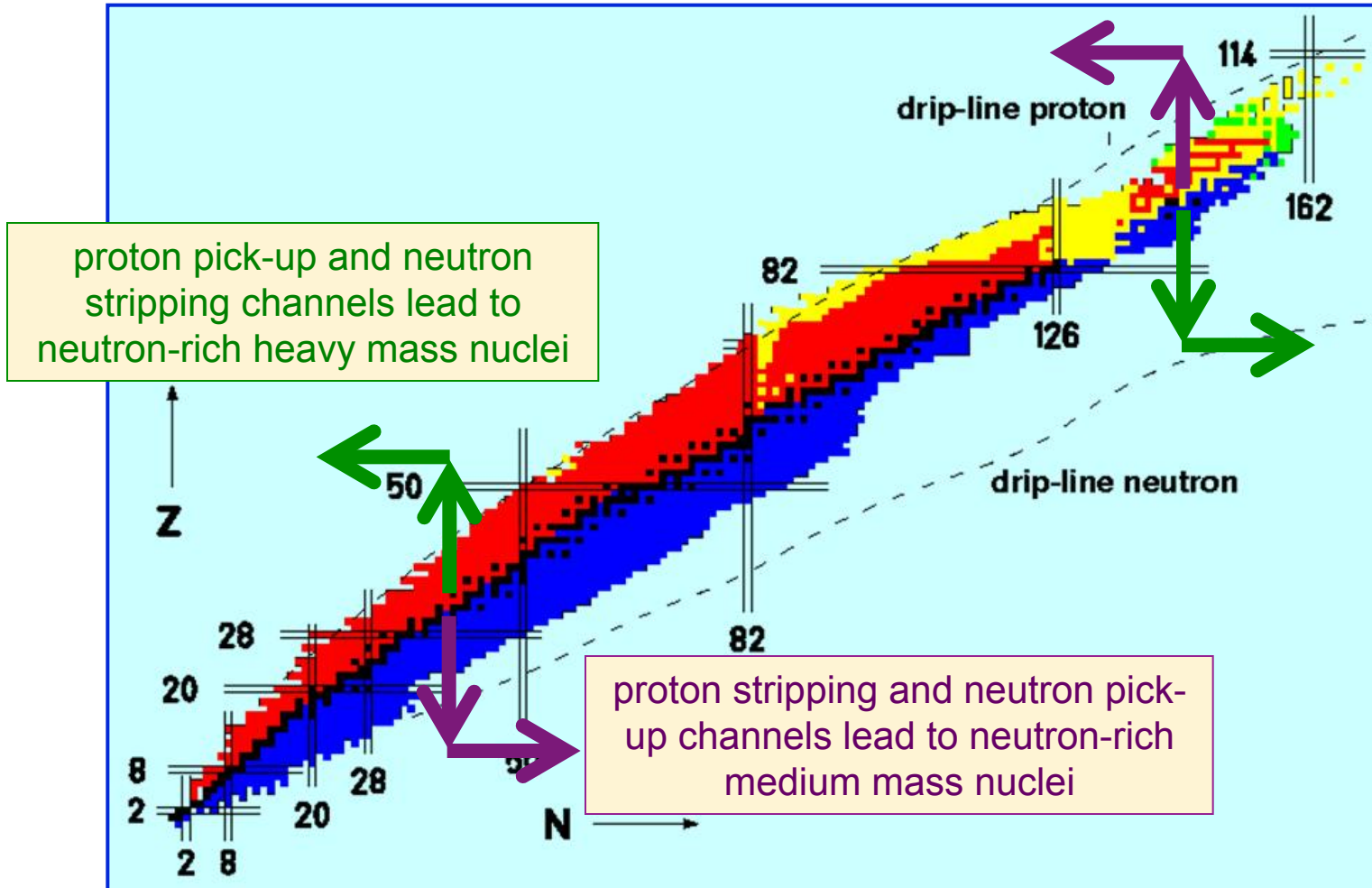


+1n and +2n slopes are in agreement with those expected from binding energies

An enhancement factor is needed for the +2n channel also in this system

Microscopic calculations in progress

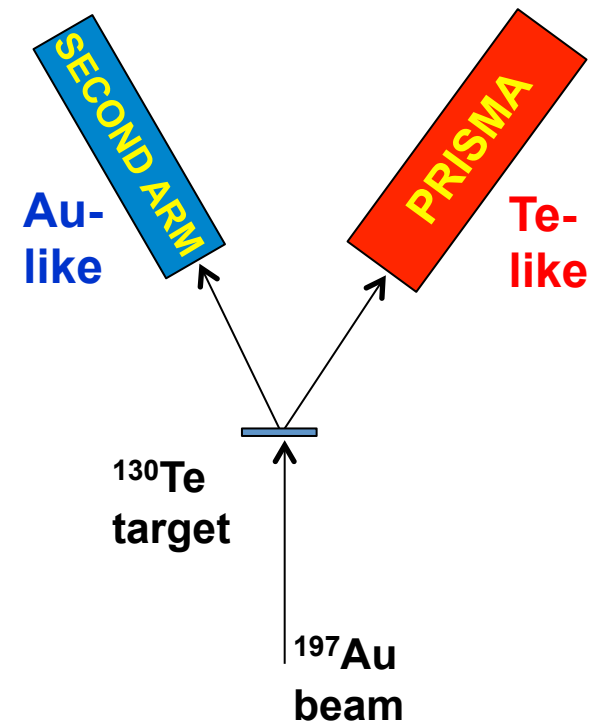
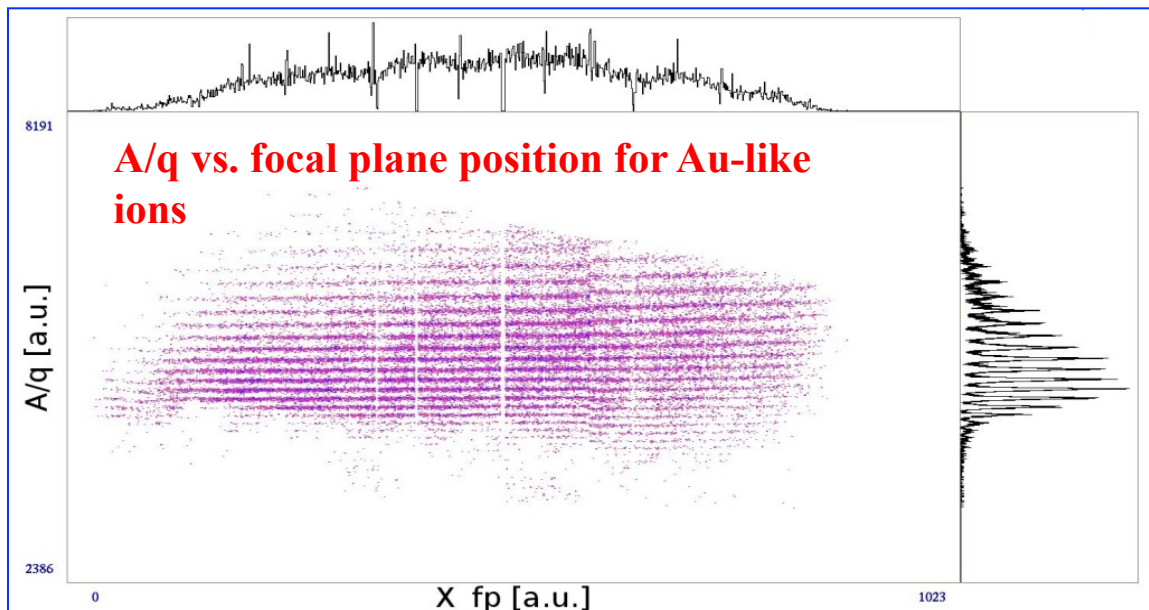
What about the heavy partner in multi-nucleon transfer reactions ?



Certain regions of the nuclear chart, like that below ^{208}Pb or in the actinides, can be hardly accessed by fragmentation or fission reactions, and multi-nucleon transfer represents a suitable mechanism to approach those neutron rich areas.

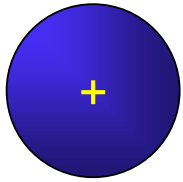
Neutron rich nuclei populated via multinucleon transfer reactions: preliminary results on the $^{197}\text{Au}+^{130}\text{Te}$ system

$E_L=1070$ MeV, $\theta_L=37^\circ$ and $E_L=1300$ MeV, $\theta_L=27^\circ$
Detected both Te- and Au-like transfer products

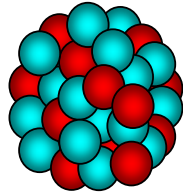


We obtained excellent A/q resolution.
The present problem is to get atomic charge states discrimination to uniquely identify final A

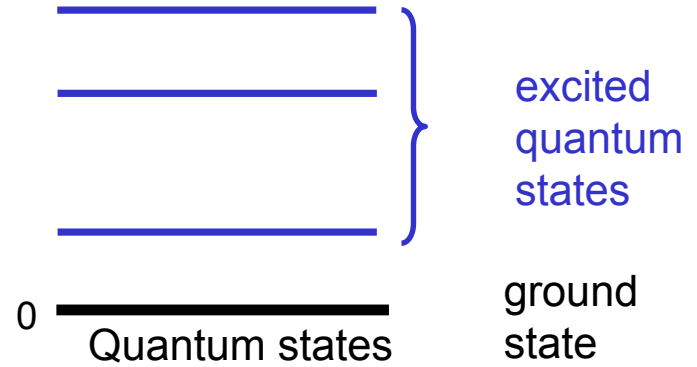
Sometimes the colliding nuclei may decide to proceed to fusion,
even at sub-barrier energies ...



Not an elementary particle

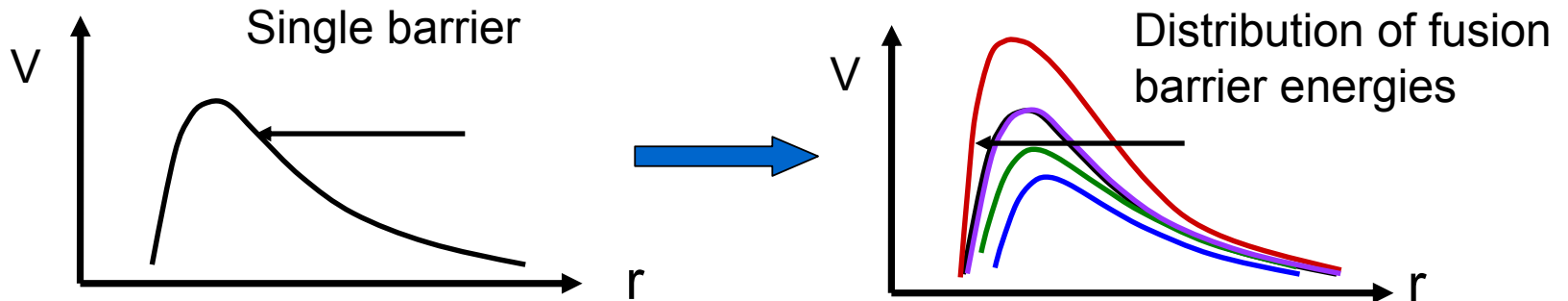


Many body quantum system



- Colliding nuclei in a superposition of quantum states

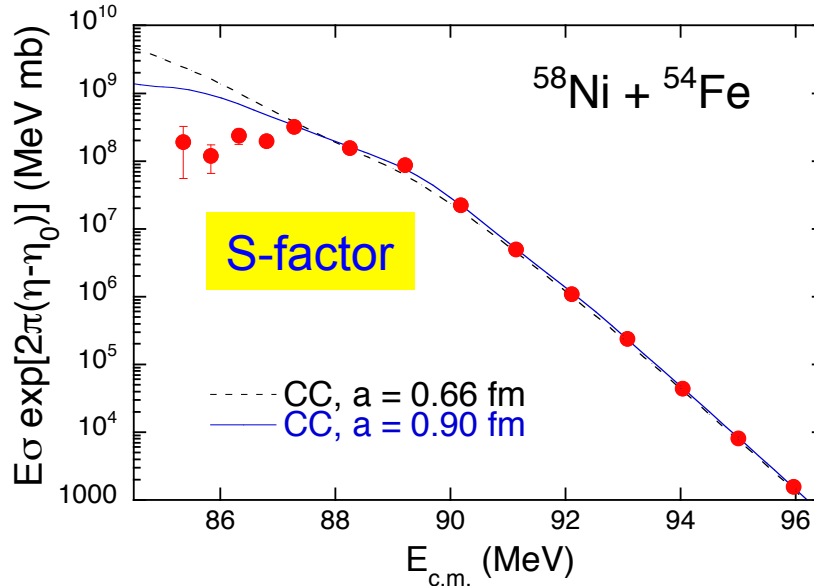
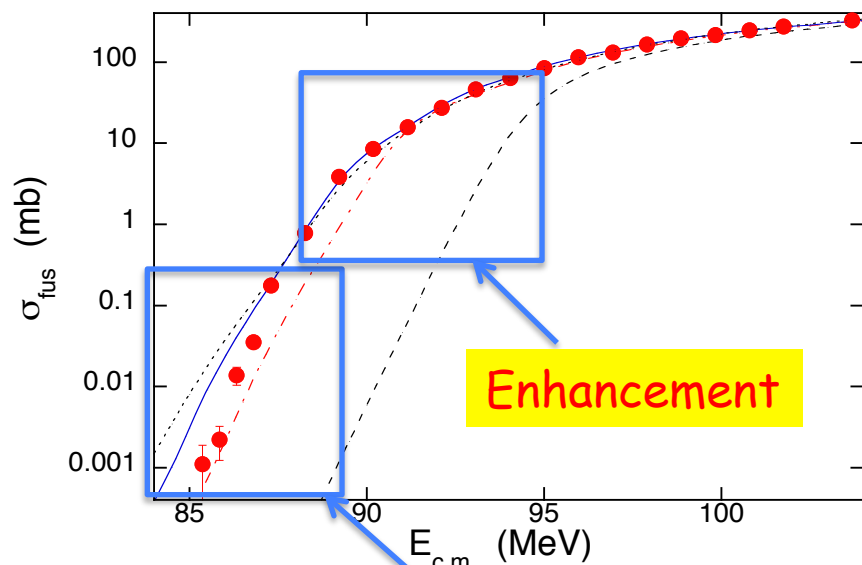
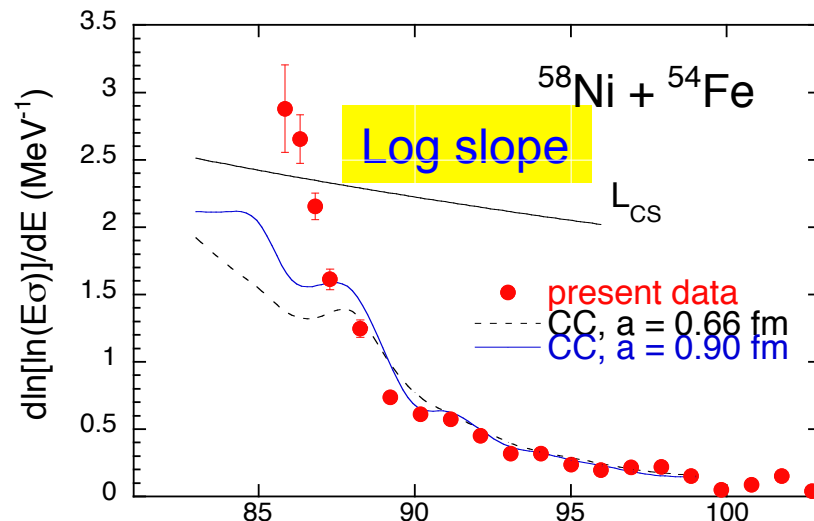
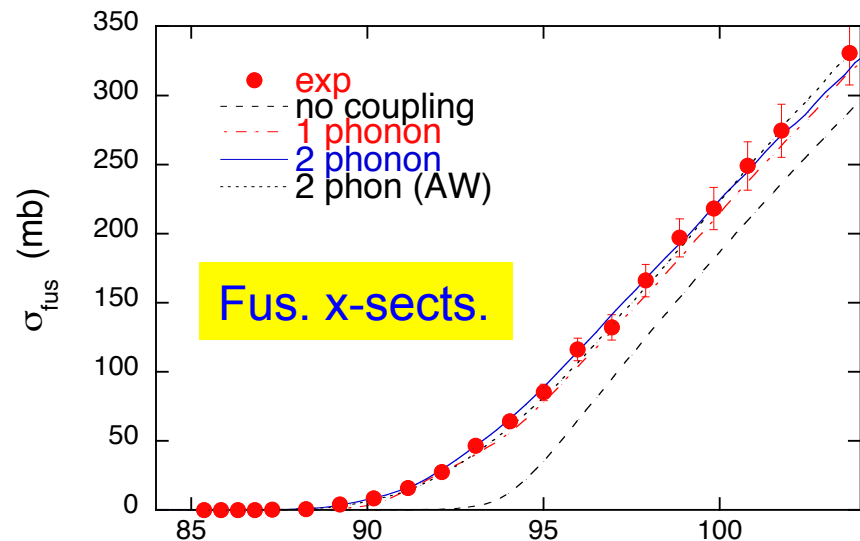
Dasso et al., NPA 405 (1983) 221



- Experimental barrier distribution $d^2(\sigma E)/dE^2$

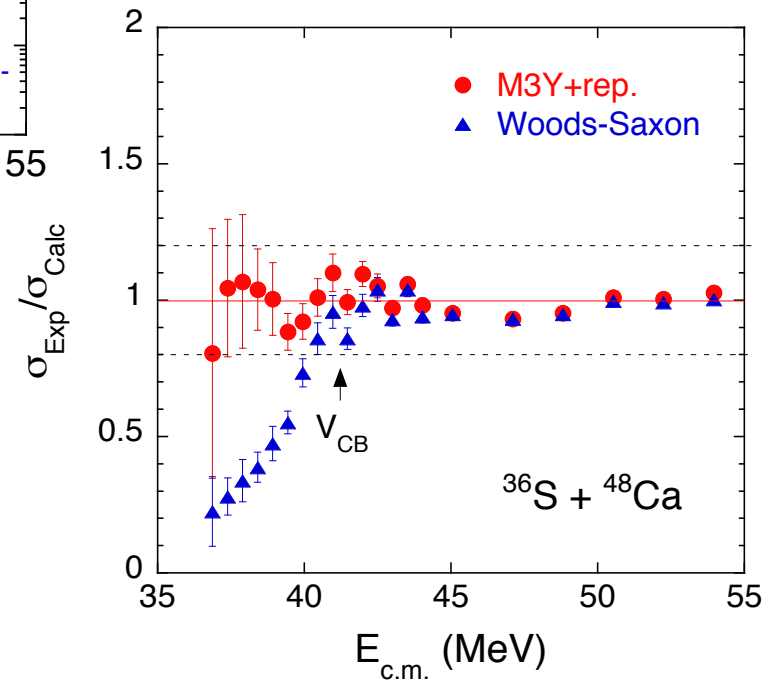
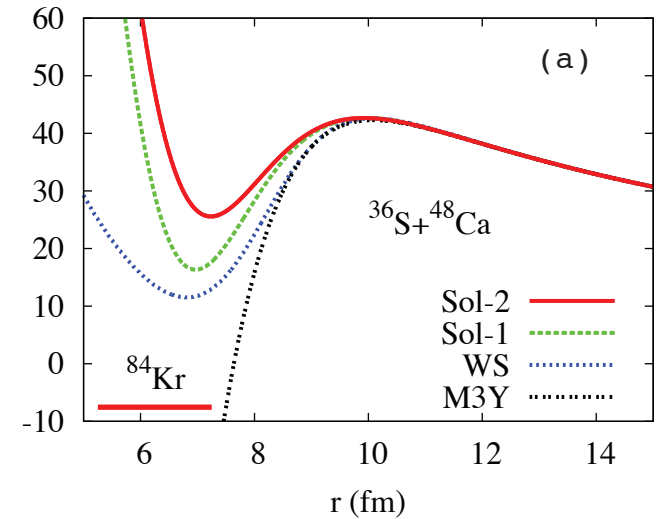
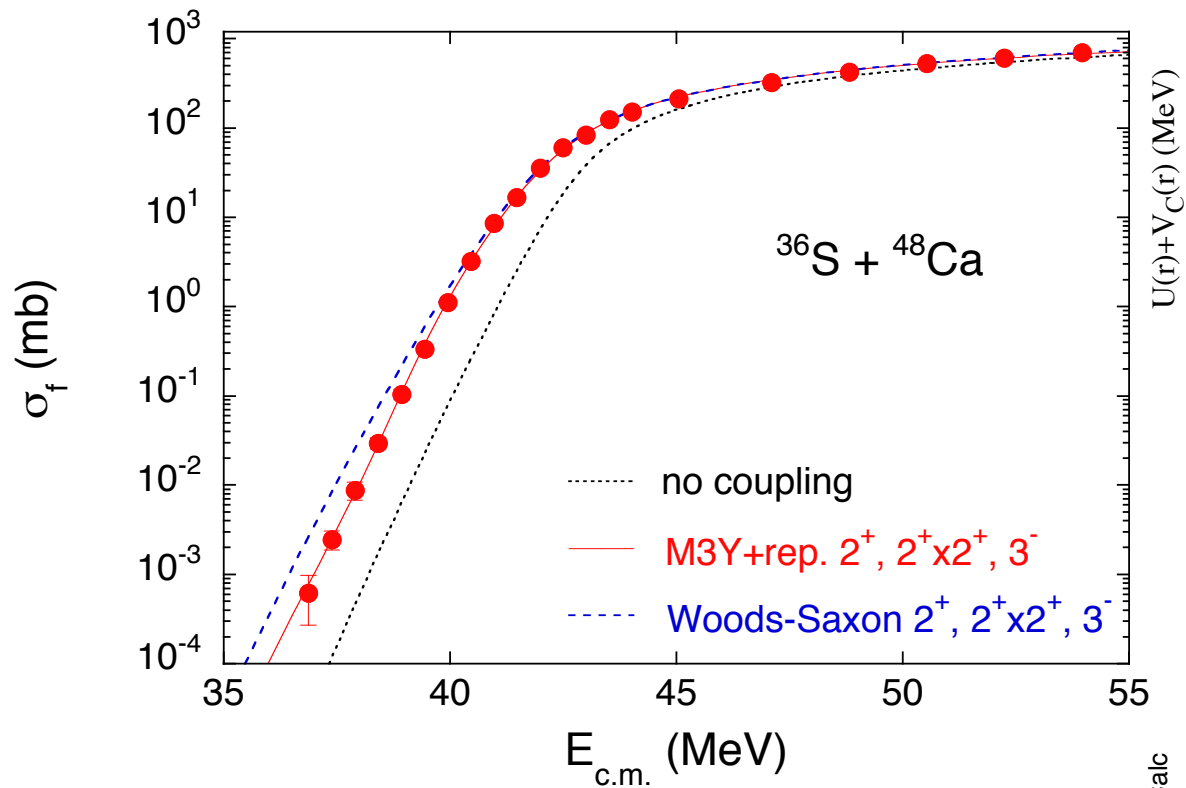
Rowley et al., PLB 254(1991)25

Fusion of stiff nuclei: the case of $^{58}\text{Ni} + ^{54}\text{Fe}$



Fusion hindrance !

Fusion excitation function of $^{36}\text{S} + ^{48}\text{Ca}$



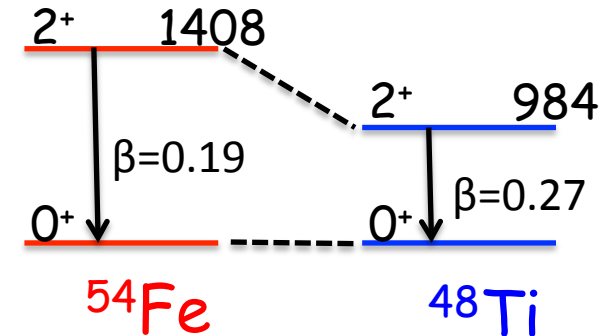
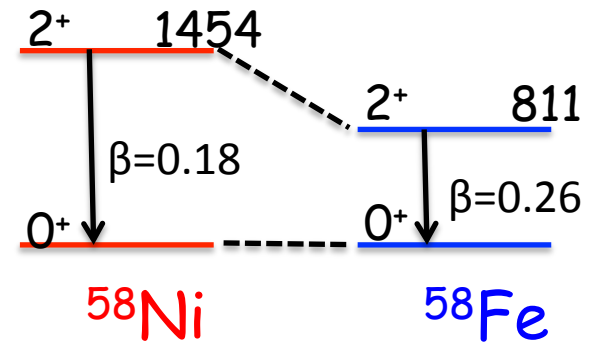
Influence of low-energy nuclear structure on hindrance

stiff nuclei-less fusion-hindrance near the barrier



Ni52 80 ms 0+	Ni53 45 ms (7/2-)	Ni54 600 ms 0+	Ni55 212.1 ms 7/2-	Ni56 5.9 d 0+	Ni57 35.60 h 3/2-	Ni58 0+ 68.077	Ni59 7.6E+4 y 3/2-	Ni60 0+
Co51 100 ms [7/2-]	Co52 18 ms [1+,6+]	Co53 240 ms (7/2-)*	Co54 193.23 ms 0+*	Co55 17.53 h 7/2-	Co56 77.27 d 4+	Co57 271.79 d 7/2-	Co58 70.82 d 2+*	Co59 7/2-*
Fe50 500 ms 0+	Fe51 305 ms (5/2-)	Fe52 8.275 h 0+*	Fe53 8.51 m 7/2-*	Fe54 0+ 5.8	Fe55 2.73 y 3/2-	Fe56 0+ 91.72	Fe57 1/2- 2.2	Fe58 0+ 0.28*
Mn49 384 ms 5/2-	Mn50 283.0 ms 0+*	Mn51 46.2 m 5/2-*	Mn52 5.591 d 6+*	Mn53 3.74E+6 y 7/2-	Mn54 312.3 d 3+*	Mn55 5/2-*	Mn56 2.5785 h 3+*	Mn57 87/2 s 5/2-*
Cr48 21.56 h 0+	Cr49 42.3 m 5/2-	Cr50 1.8E+17 y 0+ ECEC 4.345	Cr51 27.702 d 7/2-	Cr52 0+ 83.789	Cr53 3/2- 9.501	Cr54 0+ 2.365*	Cr55 3.497 m 3/2-	Cr56 5.94 m 0+ 3/2-
V47 32.6 m 3/2-	V48 15.9735 d 4+*	V49 330 d 7/2-	V50 1.4E+17 y 6+ EC,β- 0.250	V51 7/2- 99.750	V52 3.75 m 3+ β-	V53 1.61 m 7/2- β-	V54 49.8 s 3+ β-	V55 6.54 s (7/2-) β-
Ti46 0+ 8.0	Ti47 5/2- 7.3	Ti48 0+ 73.8	Ti49 7/2- 5.5	Ti50 0+ 5.4	Ti51 5.76 m 3/2- β-	Ti52 1.7 m 0+ β-	Ti53 32.7 s (3/2-) β-	Ti54 30 s 0+ β-
Sc45 7/2-*	Sc46 83.79 d 4+*	Sc47 3.345 d 7/2-	Sc48 43.67 h 6+	Sc49 57.2 m 7/2-	Sc50 102.5 s 5+*	Sc51 12.4 s (7/2-)	Sc52 8.2 s 3+	Sc53 150 ms [7/2-]

Z=28

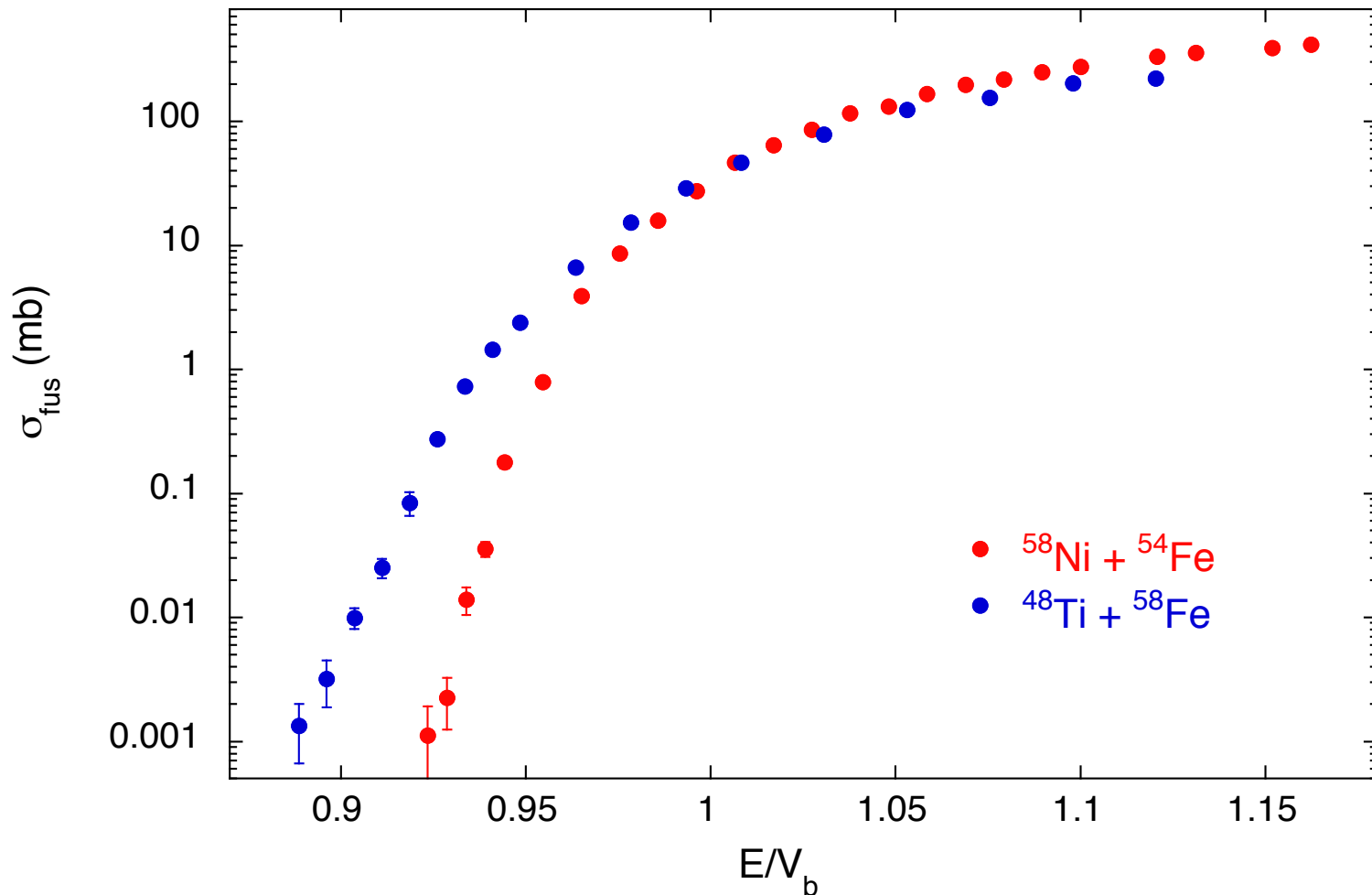


N=28

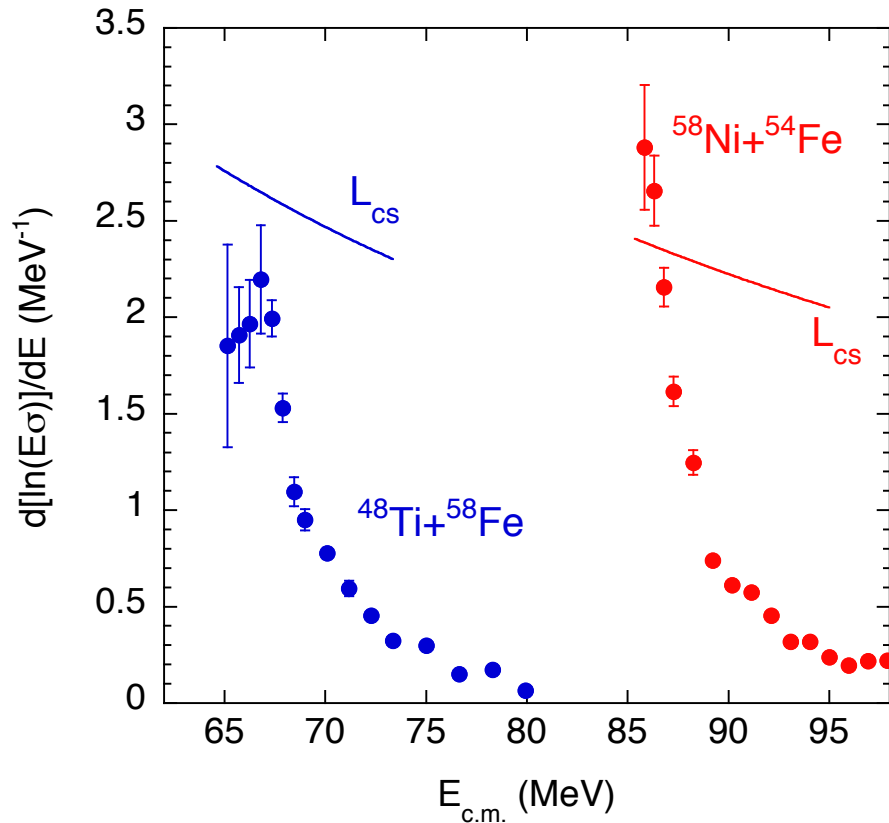
$^{48}\text{Ti} + ^{58}\text{Fe}$

soft nuclei-more fusion-hindrance far below the barrier

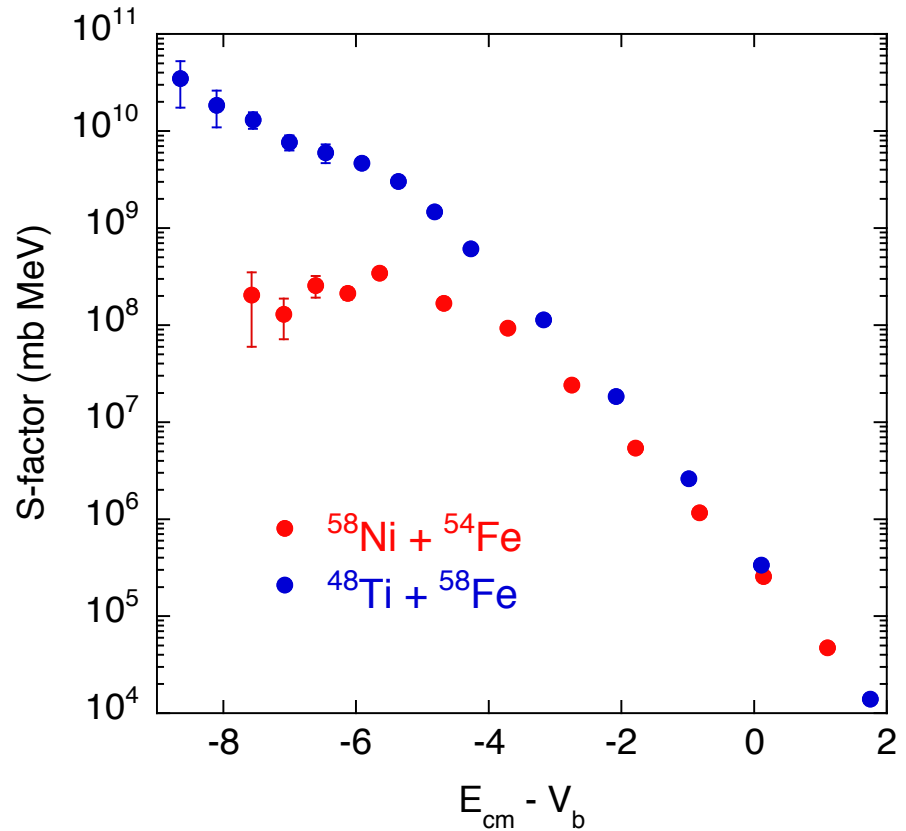
Fusion excitation functions of $^{48}\text{Ti} + ^{58}\text{Fe}$ and $^{58}\text{Ni} + ^{54}\text{Fe}$, plotted vs. the energy relative to the nominal barrier. V_b is the barrier height produced by the Akyüz-Winther potential.



Logarithmic derivative of the fusion excitation function of $^{48}\text{Ti}+^{58}\text{Fe}$ and $^{58}\text{Ni}+^{54}\text{Fe}$, and comparison of the S-factors for the two systems

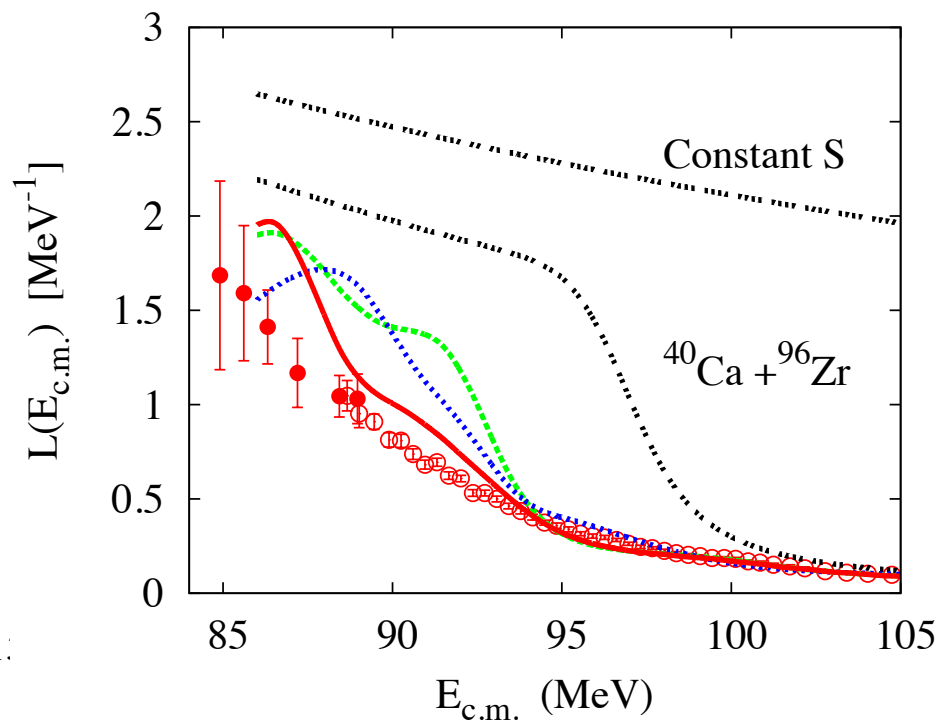
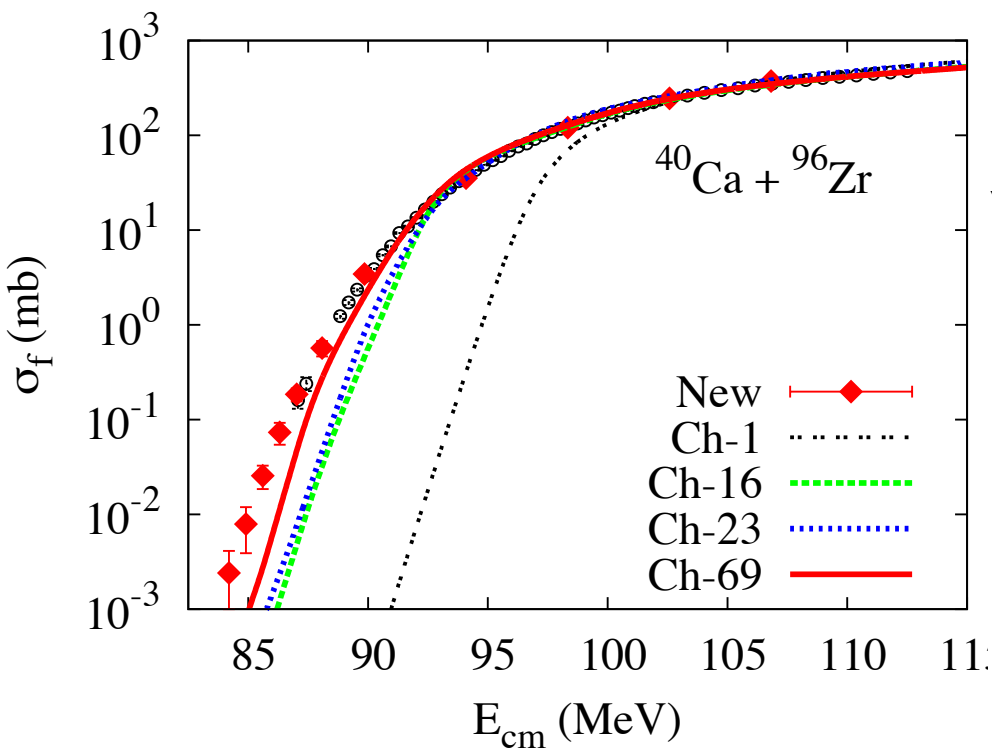


The slope of $^{48}\text{Ti}+^{58}\text{Fe}$ saturates below the barrier, while it keeps increasing for $^{58}\text{Ni}+^{54}\text{Fe}$



A clear maximum of the S-factor develops for $^{58}\text{Ni}+^{54}\text{Fe}$, but no maximum is observed for $^{48}\text{Ti}+^{58}\text{Fe}$

Couplings to transfer channels and hindrance in $^{40}\text{Ca} + ^{96}\text{Zr}$



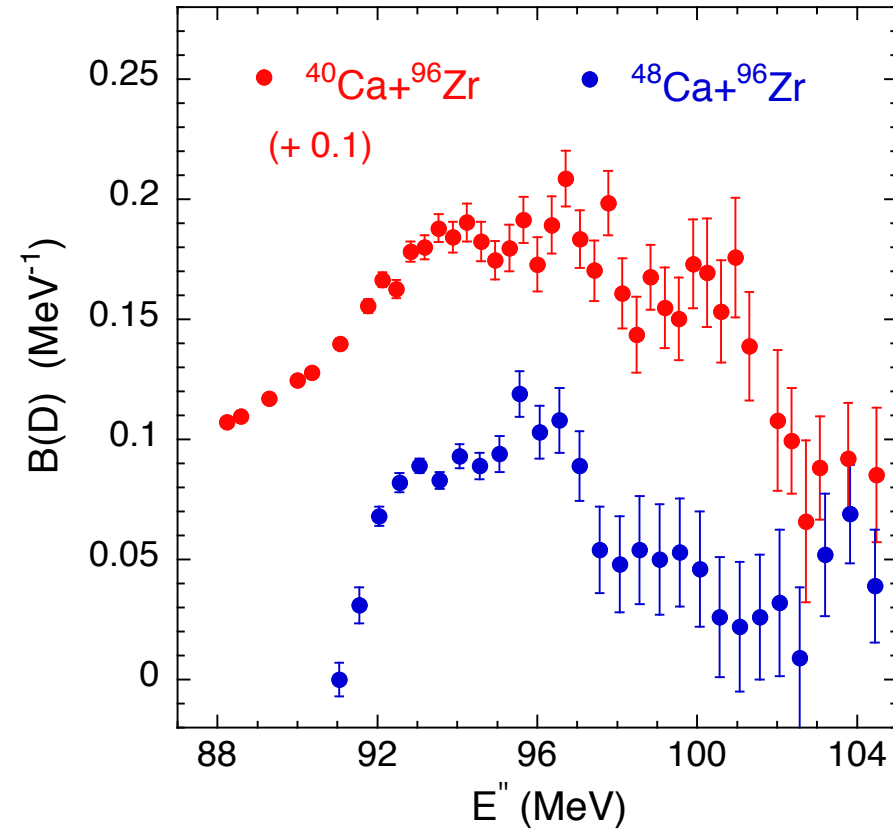
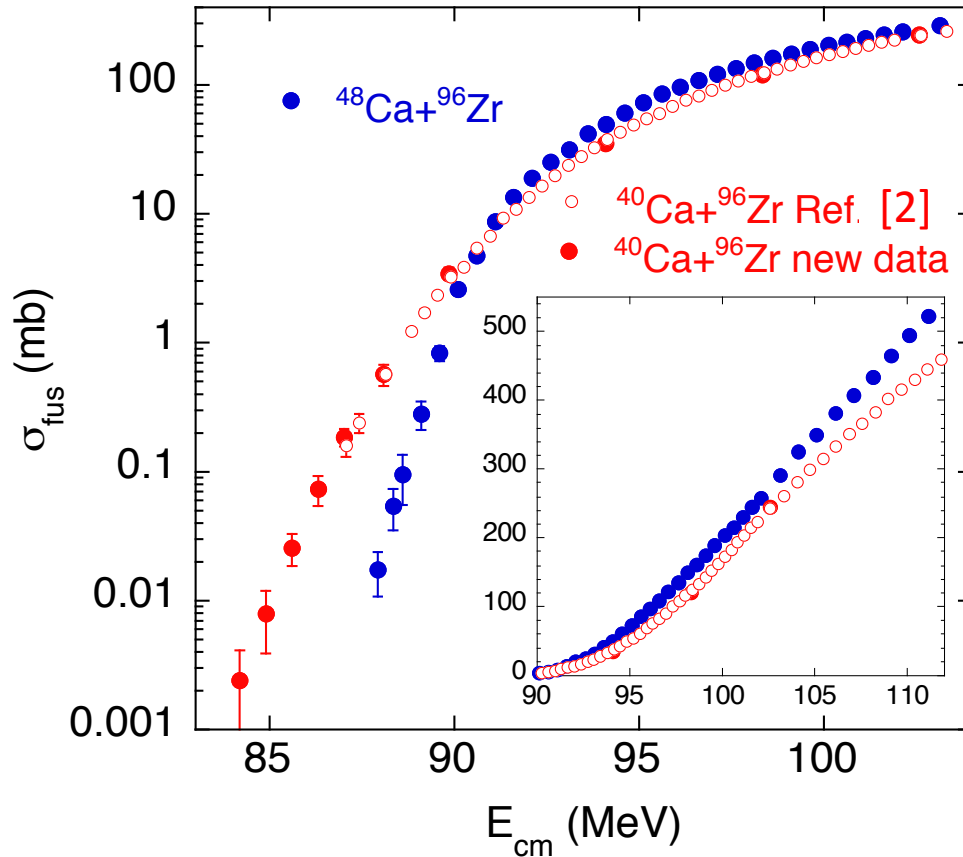
The cross sections are strongly underestimated below a few mb.
 The low-lying $2+$ and $3-$ states (Ch-23) produce strong effects.
 Couplings to $Q > 0$ one- and two-nucleon transfer channels (Ch-69) bring further significant enhancements, even at the level of a few μb , but no indication of hindrance appears yet.

Comparing the two cases of $^{40,48}\text{Ca} + ^{96}\text{Zr}$

sub-barrier slopes



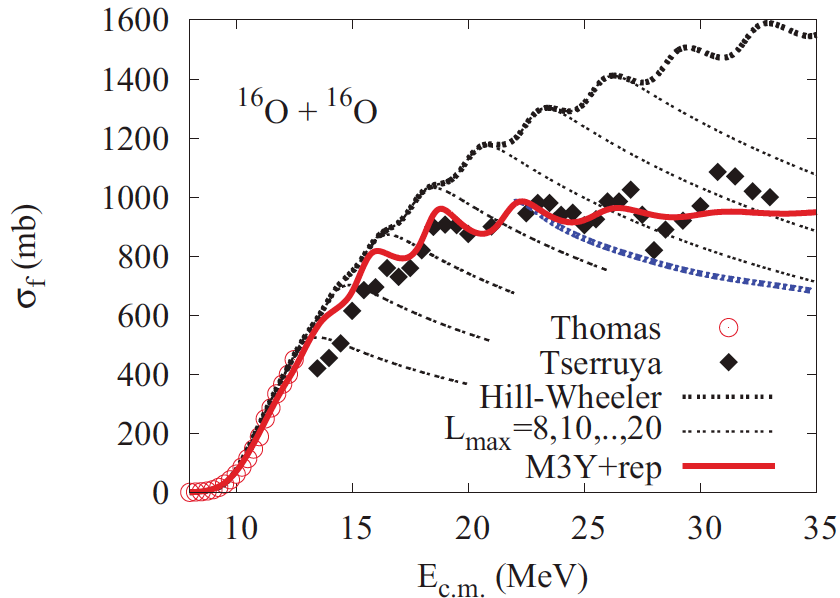
barrier distributions



The cross section for $^{48}\text{Ca} + ^{96}\text{Zr}$ decreases very sharply below the barrier and, indeed, this system shows hindrance. The decrease of the excitation function for $^{40}\text{Ca} + ^{96}\text{Zr}$ is by far slower.

Medium-light systems and oscillations

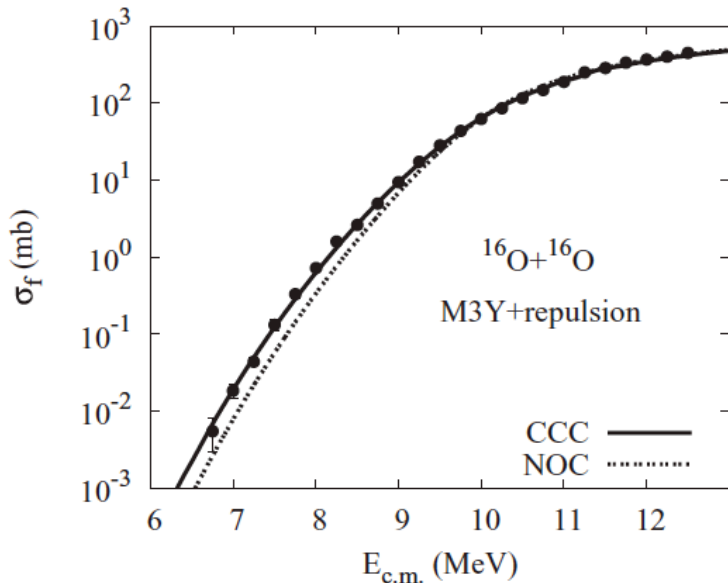
Structures in fusion excitation function of light systems



- Oscillatory structures were recently interpreted due to the penetration of successive centrifugal barriers

$$V_B(L) = V_{CB} + \frac{\hbar^2 L(L+1)}{2\mu R_{CB}^2}$$

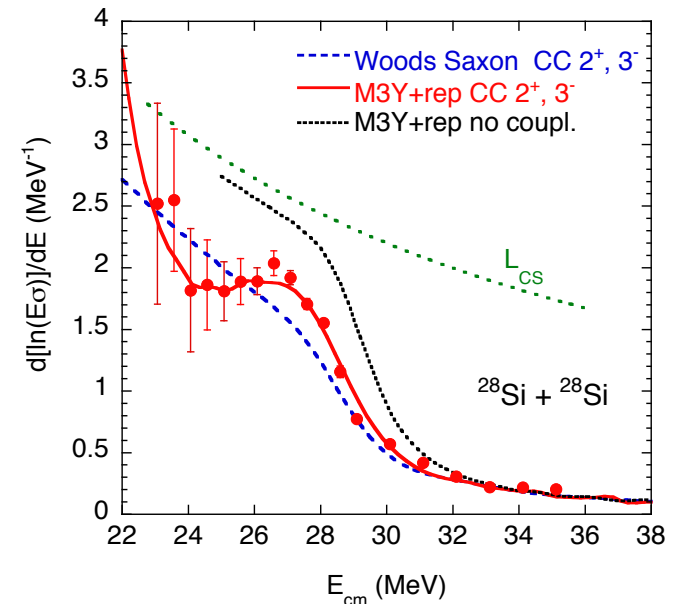
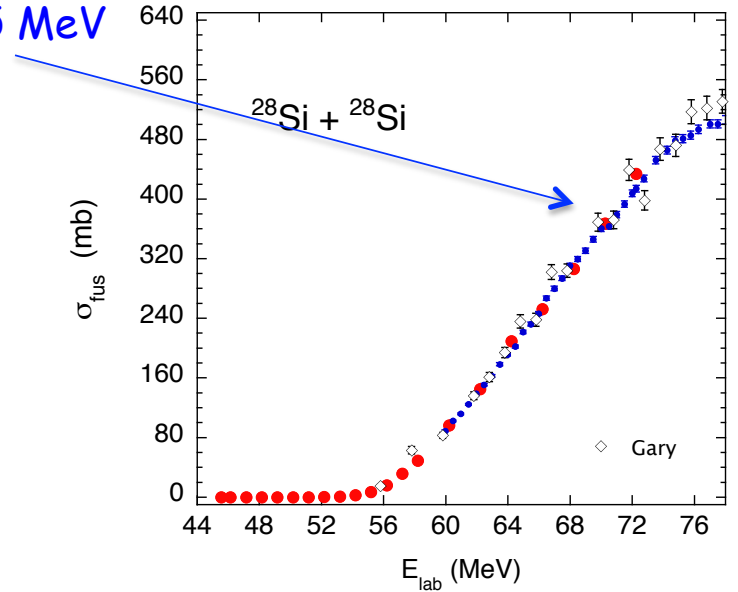
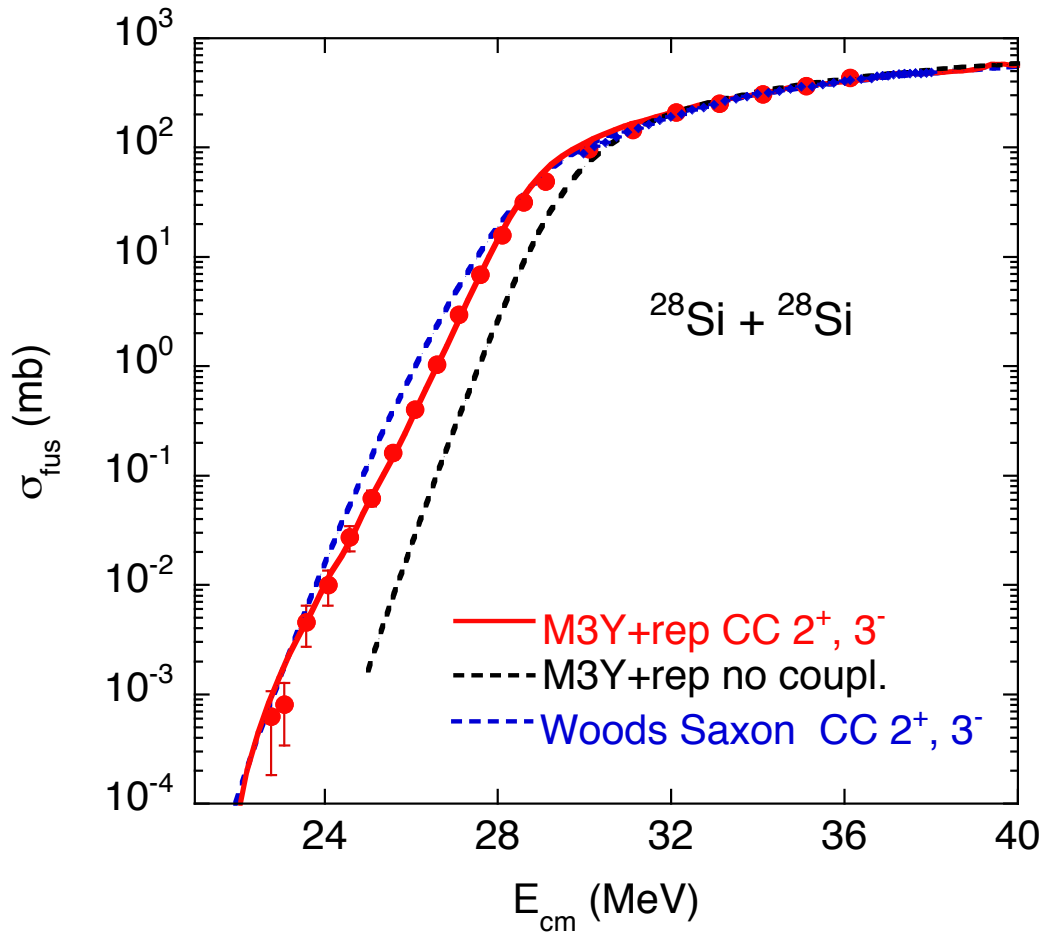
- A shallow ion-ion potential (M3Y+rep) is needed to fit the data above (and below) the barrier.
- $\rightarrow\rightarrow$ Look for above-barrier oscillations in heavier systems, where sub-barrier hindrance is stronger and better established.
- Clearly observing such oscillations would put strong constraints on the ion-ion potential in a wide energy range.



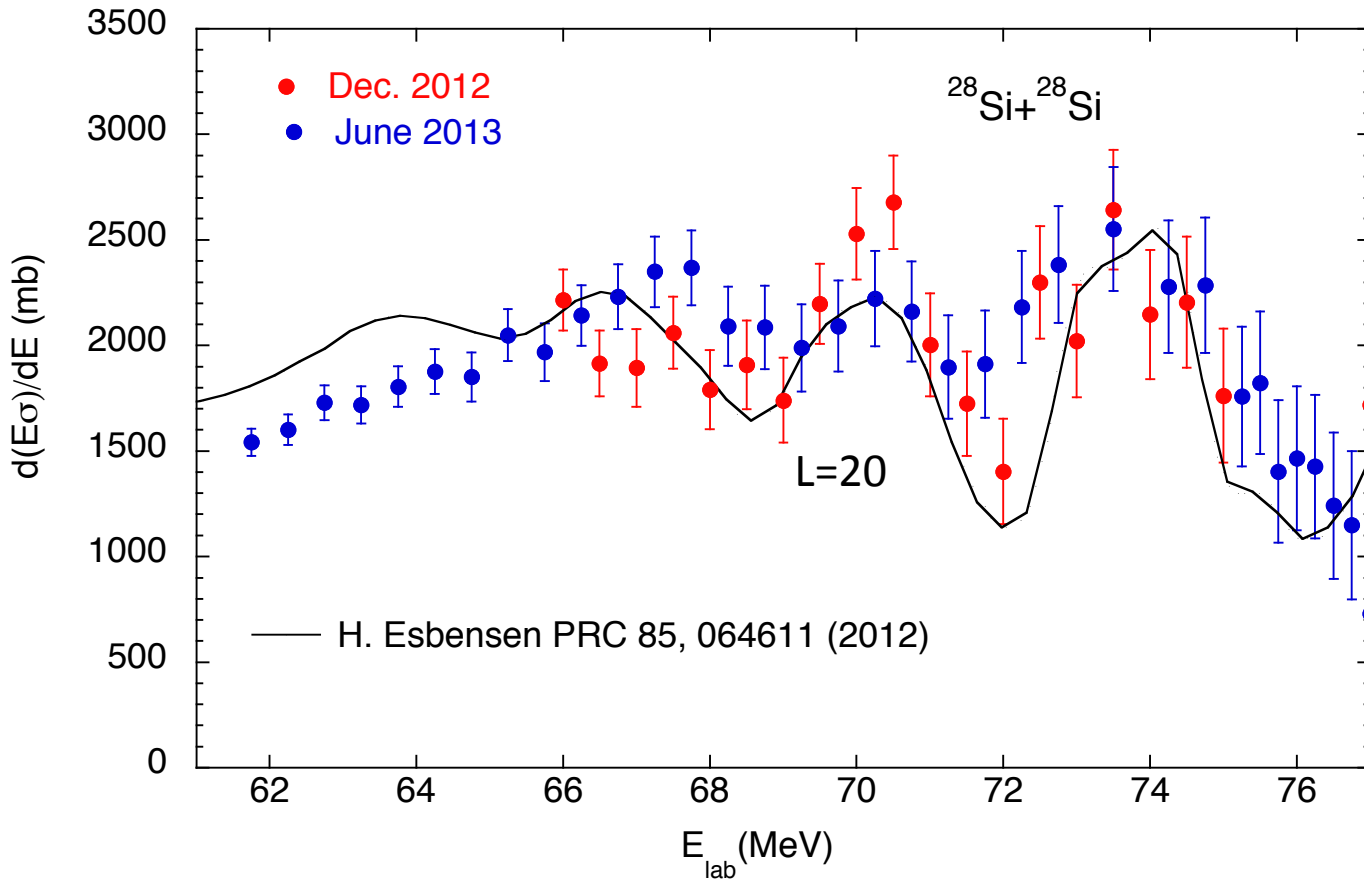
H.Esbensen, PRC77, 054608 (2008); PRC85, 064611 (2012)
C.Y.Wong, PRC86, 064603 (2012)

The case of $^{28}\text{Si} + ^{28}\text{Si}$: new measurement of the fusion excitation function

(the blue points were measured with $\Delta E_{\text{lab}} = 0.5 \text{ MeV}$ and 1% statistical error)



We take the energy-weighted derivative of the exc. function, and we compare with the result of a recent CC calculation



The energy difference between successive barriers:

$$\Delta V_B = V_B(L + 1) - V_B(L) \approx \frac{\hbar^2 2(L + 1)}{2\mu R_{CB}^2}$$

(twice as that for a symmetric system like $^{28}\text{Si} + ^{28}\text{Si}$, i.e. $\Delta V_{B(\text{c.m.})} \approx 1.52 \text{ MeV}$)

Summary

- Few- and multi-nucleon transfer reactions are a suitable tool for studies both of nuclear structure and reaction dynamics
- Suitable conditions for investigating pair correlations in heavy ion collisions are offered by sub-barrier transfer reactions
- Interesting results have been obtained for the systems $^{96}\text{Zr} + ^{40}\text{Ca}$ and $^{116}\text{Sn} + ^{60}\text{Ni}$ using the inverse kinematics and the PRISMA spectrometer
- Heavy-ion fusion below the barrier continues displaying interesting, and sometimes unexpected, features, such as the hindrance phenomenon
- The low-lying collective structure of the colliding nuclei strongly influences the threshold of hindrance, through *CC* effects
- A shallow potential (M3Y+Rep) allows reproducing the x-sects. far below the barrier; oscillations have been observed above the barrier for $^{28}\text{Si} + ^{28}\text{Si}$, that are described in the same theoretical frame

Our collaboration in recent experiments

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