

# Bjorken- $x$ dependences of light-quark sea

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2<sup>nd</sup> Workshop on “Probing Strangeness in  
Hard Processes”

LNF, Frascati, November 11-13, 2013

# Outline

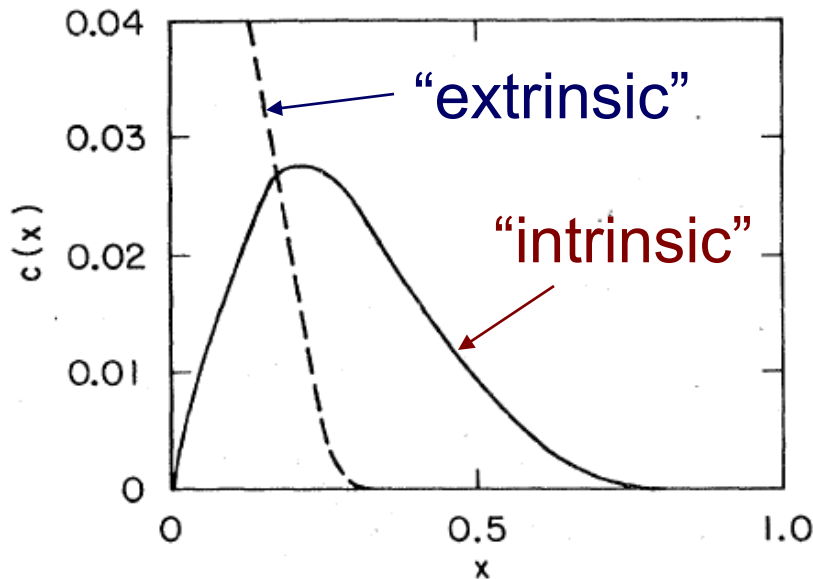
- Extraction of “intrinsic”  $\bar{u}$ ,  $\bar{d}$ , and  $\bar{s}$  sea in the nucleons from Drell-Yan and semi-inclusive DIS experiments
- Separation of “connected sea” from “disconnected sea” for  $\bar{u}(x) + \bar{d}(x)$
- Bjorken- $x$  and  $Q^2$  - dependences of strange quark distributions

# Search for the “intrinsic” quark sea

In 1980, Brodsky, Hoyer, Peterson, Sakai (BHPS) suggested the existence of “intrinsic” charm

$$|p\rangle = P_{3q} |uud\rangle + P_{5q} |uudQ\bar{Q}\rangle + \dots$$

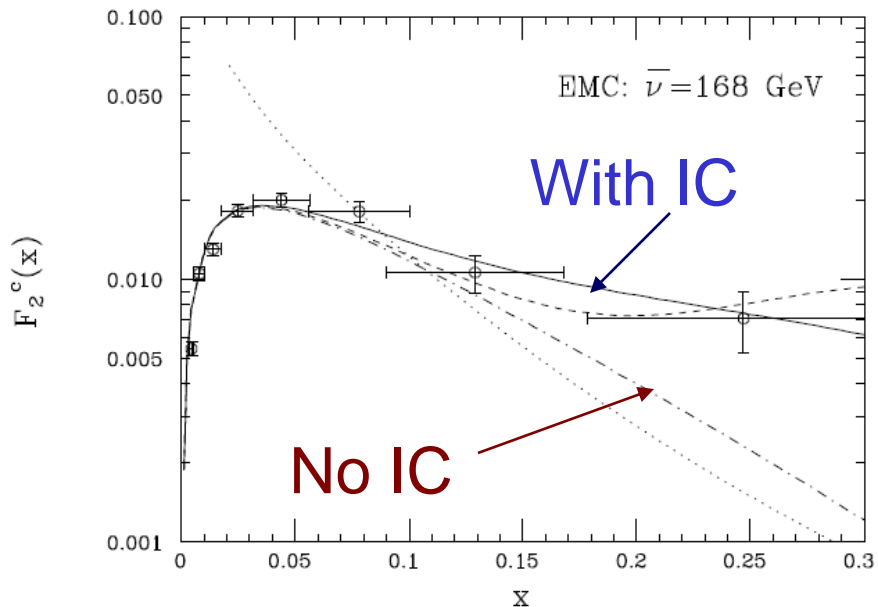
The “intrinsic”-charm from  $|uudc\bar{c}\rangle$  is “valence”-like and peak at large  $x$  unlike the “extrinsic” sea ( $g \rightarrow c\bar{c}$ )



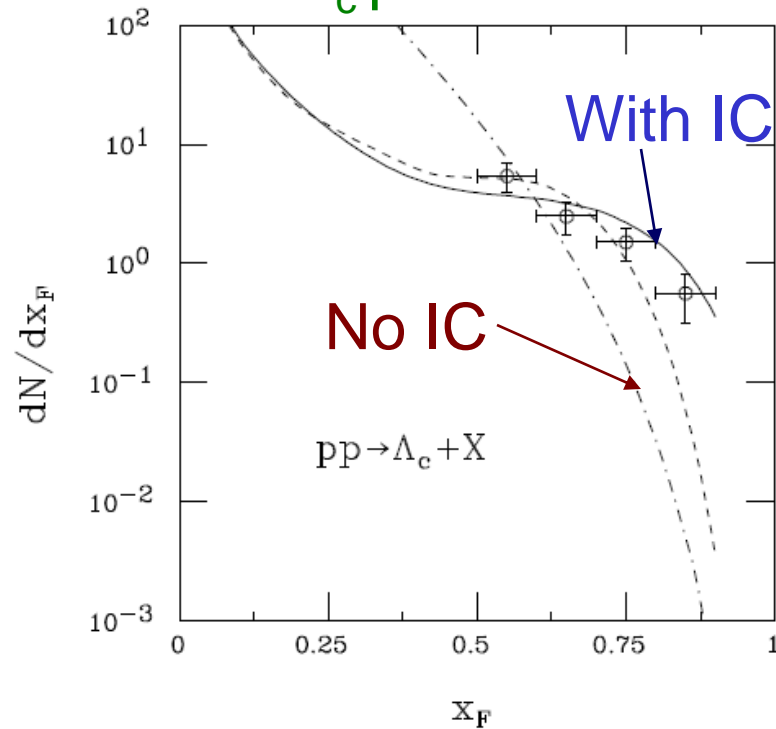
The  $|uudc\bar{c}\rangle$  intrinsic-charm can lead to large contribution to charm production at large  $x$

# “Evidence” for the “intrinsic” charm (IC)

DIS data



$\Lambda_c$  production



Gunion and Vogt (hep-ph/9706252)

“Evidence” appears to be rather weak

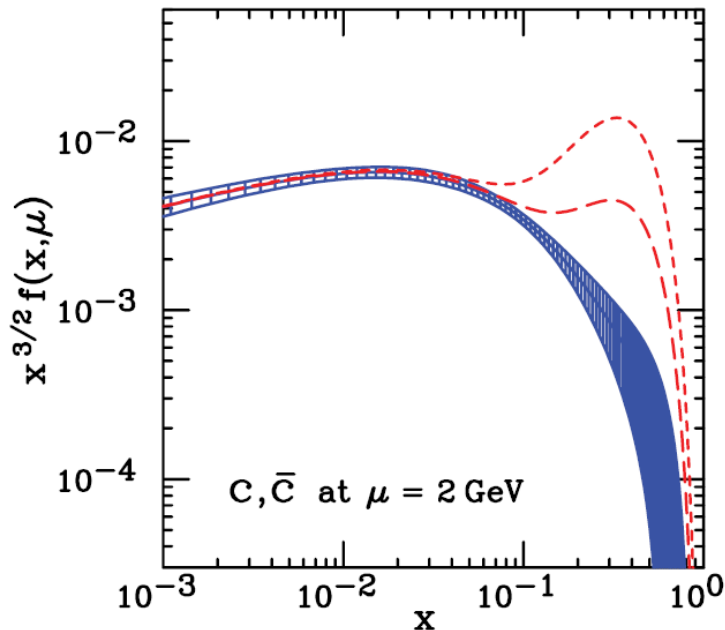
(subject to the uncertainties of charmed-quark parametrization in the PDF)

# A global fit by CTEQ to extract intrinsic-charm

PHYSICAL REVIEW D 75, 054029 (2007)

## Charm parton content of the nucleon

J. Pumplin,<sup>1,\*</sup> H. L. Lai,<sup>1,2,3</sup> and W. K. Tung<sup>1,2</sup>



Blue band corresponds to CTEQ6 best fit, including uncertainty

Red curves include intrinsic charm of 1% and 3% ( $\chi^2$  changes only slightly)

We find that the range of IC is constrained to be from zero (no IC) to a level 2–3 times larger than previous model estimates. The behaviors of typical charm distributions within this range are described, and their implications for hadron collider phenomenology are briefly discussed.

**No conclusive evidence for intrinsic-charm**

# Search for the lighter “intrinsic” quark sea

$$|p\rangle = P_{3q} |uud\rangle + P_{5q} |uudQ\bar{Q}\rangle + \dots$$

No conclusive experimental evidence  
for intrinsic-charm so far

Are there experimental evidences for the intrinsic

$|uudu\bar{u}\rangle$ ,  $|uudd\bar{d}\rangle$ ,  $|uuds\bar{s}\rangle$  5-quark states ?

$$P_{5q} \sim 1/m_Q^2$$

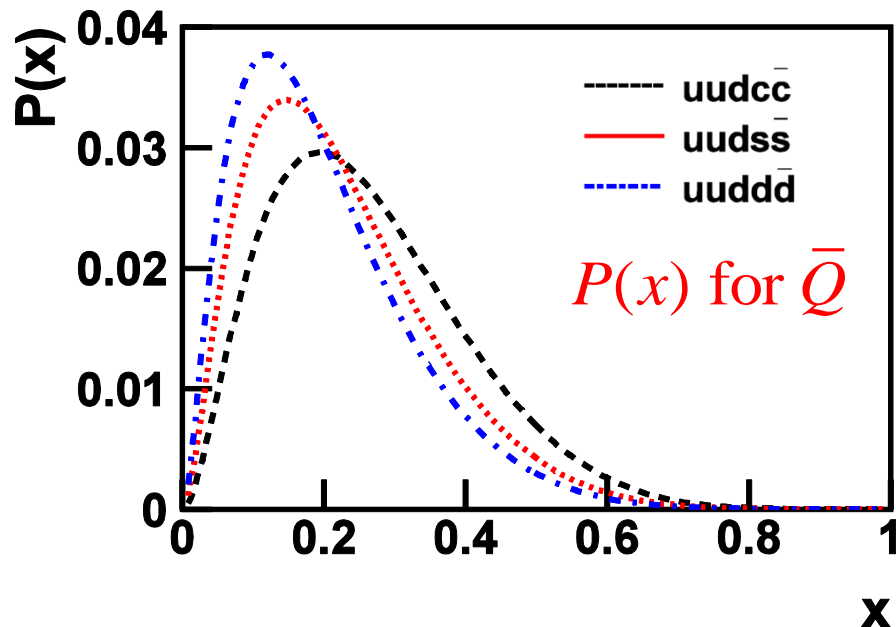
The 5-quark states for lighter  
quarks have larger probabilities!

# $x$ -distribution for “intrinsic” light-quark sea

$$|p\rangle = P_{3q} |uud\rangle + P_{5q} |uudQ\bar{Q}\rangle + \dots$$

Brodsky et al. (BHPS) give the following probability for quark  $i$  (mass  $m_i$ ) to carry momentum  $x_i$

$$P(x_1, \dots, x_5) = N_5 \delta(1 - \sum_{i=1}^5 x_i) [m_p^2 - \sum_{i=1}^5 \frac{m_i^2}{x_i}]^{-2}$$



In the limit of large mass for quark  $Q$  (charm):

$$P(x_5) = \frac{1}{2} \tilde{N}_5 x_5^2 [(1-x_5)(1+10x_5+x_5^2) - 2x_5(1+x_5)\ln(1/x_5)]$$

One can calculate  $P(x)$  for antiquark  $\bar{Q}$  ( $\bar{c}, \bar{s}, \bar{d}$ ) numerically

# How to separate the “intrinsic sea” from the “extrinsic sea”?

- Select experimental observables which have no contributions from the “extrinsic sea”
- “Intrinsic sea” and “extrinsic sea” are expected to have different  $x$ -distributions
  - Intrinsic sea is “valence-like” and is more abundant at larger  $x$
  - Extrinsic sea is more abundant at smaller  $x$



# How to separate the “intrinsic sea” from the “extrinsic sea”?

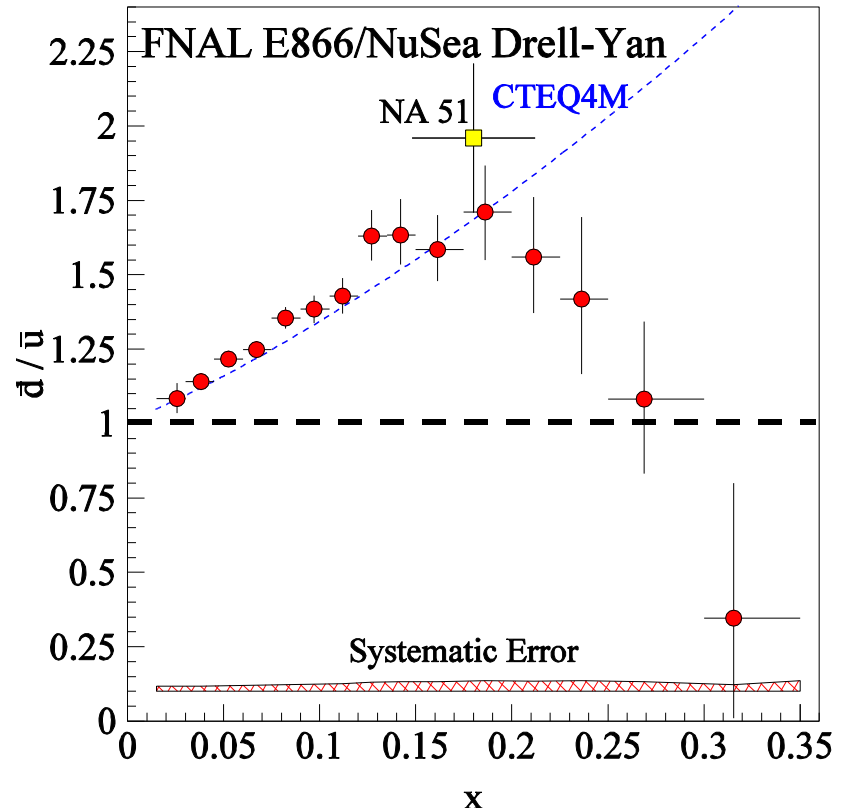
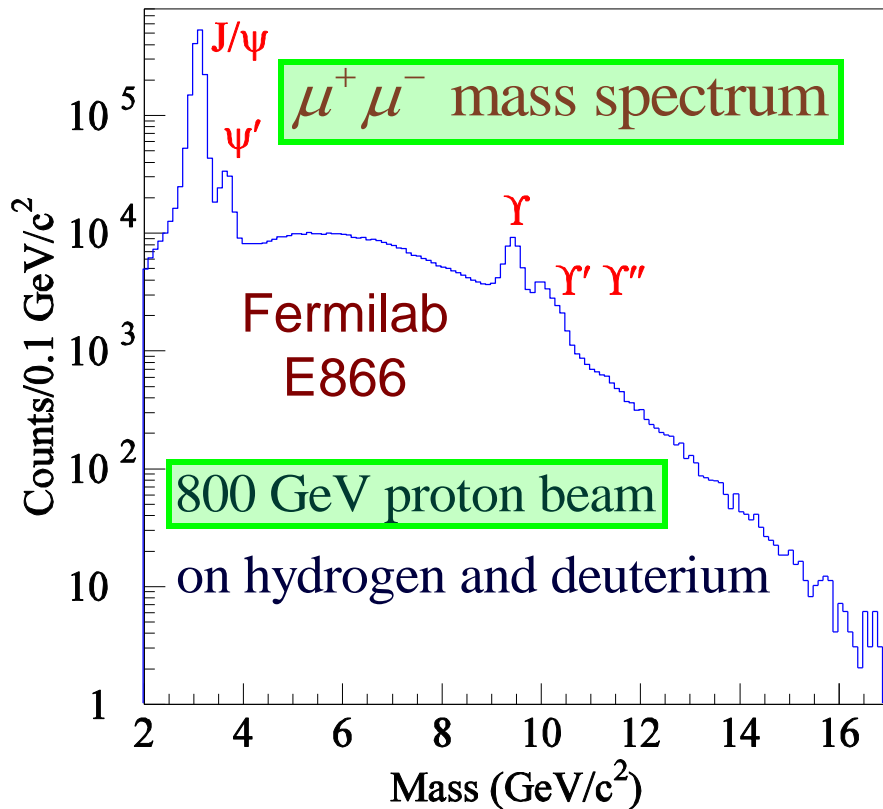
- Select experimental observables which have no contributions from the “extrinsic sea”

$\bar{d} - \bar{u}$  has no contribution from extrinsic sea ( $g \rightarrow \bar{q}q$ )  
and is sensitive to "intrinsic sea" only



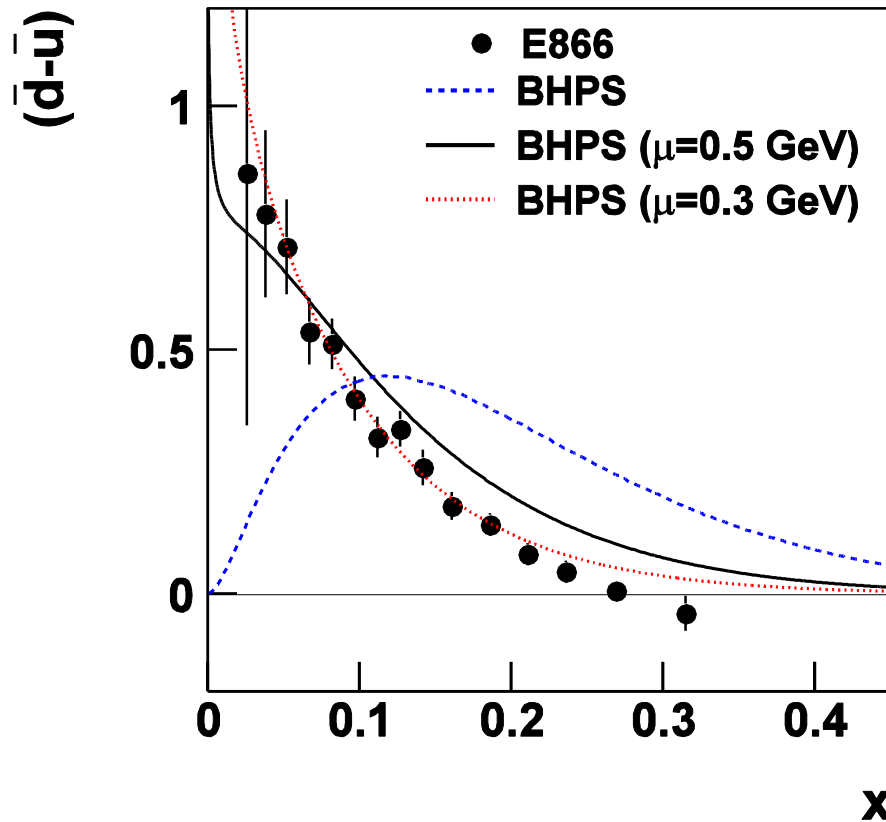
# $\bar{d} / \bar{u}$ flavor asymmetry from Drell-Yan

$$\left( \frac{d^2\sigma}{dx_1 dx_2} \right)_{D.Y.} = \frac{4\pi\alpha^2}{9sx_1x_2} \sum_a e_a^2 [q_a(x_1)\bar{q}_a(x_2) + \bar{q}_a(x_1)q_a(x_2)]$$



at  $x_1 > x_2$  : Drell-Yan:  $\sigma^{pd} / 2\sigma^{pp} \sim \frac{1}{2} (1 + \bar{d}(x_2) / \bar{u}(x_2))$

# Comparison between the $\bar{d}(x) - \bar{u}(x)$ data with the intrinsic 5- $q$ model



The data are in good agreement with the 5- $q$  model after evolution from the initial scale  $\mu$  to  $Q^2=54 \text{ GeV}^2$

The difference in the two 5-quark components can also be determined

$$P_5^{uudd\bar{d}} - P_5^{uudu\bar{u}} = 0.118$$

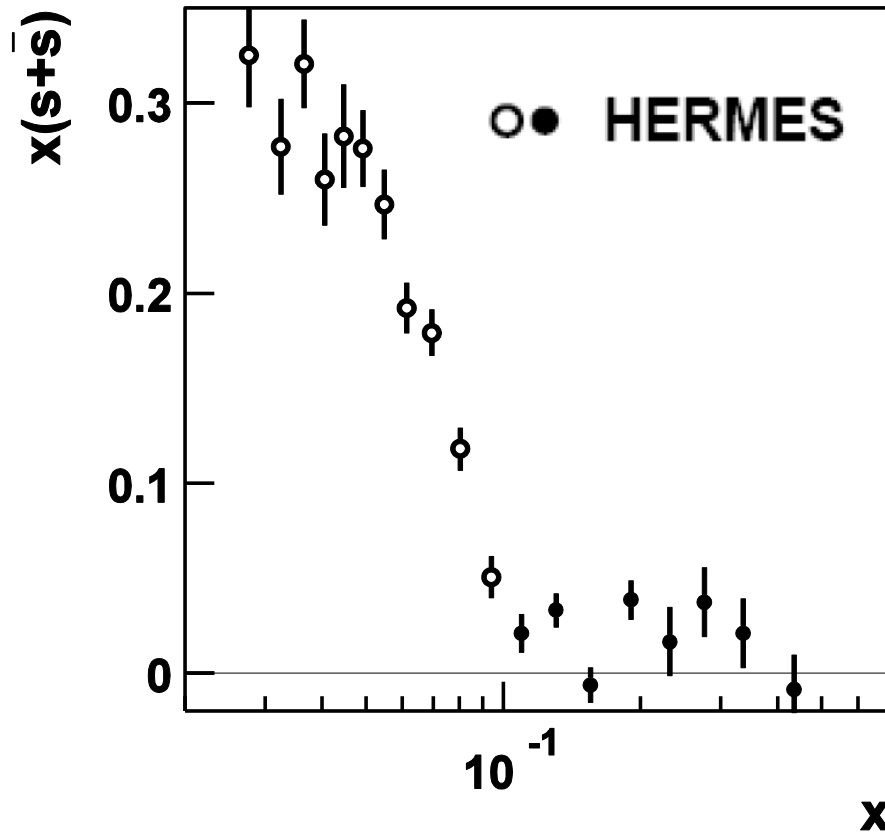
(W. Chang and JCP , PRL 106, 252002 (2011))

# How to separate the “intrinsic sea” from the “extrinsic sea”?

- “Intrinsic sea” and “extrinsic sea” are expected to have different  $x$ -distributions
  - Intrinsic sea is “valence-like” and is more abundant at larger  $x$
  - Extrinsic sea is more abundant at smaller  $x$

An example is the  $s(x) + \bar{s}(x)$  distribution

# Comparison between the $s(x) + \bar{s}(x)$ data with the intrinsic 5- $q$ model

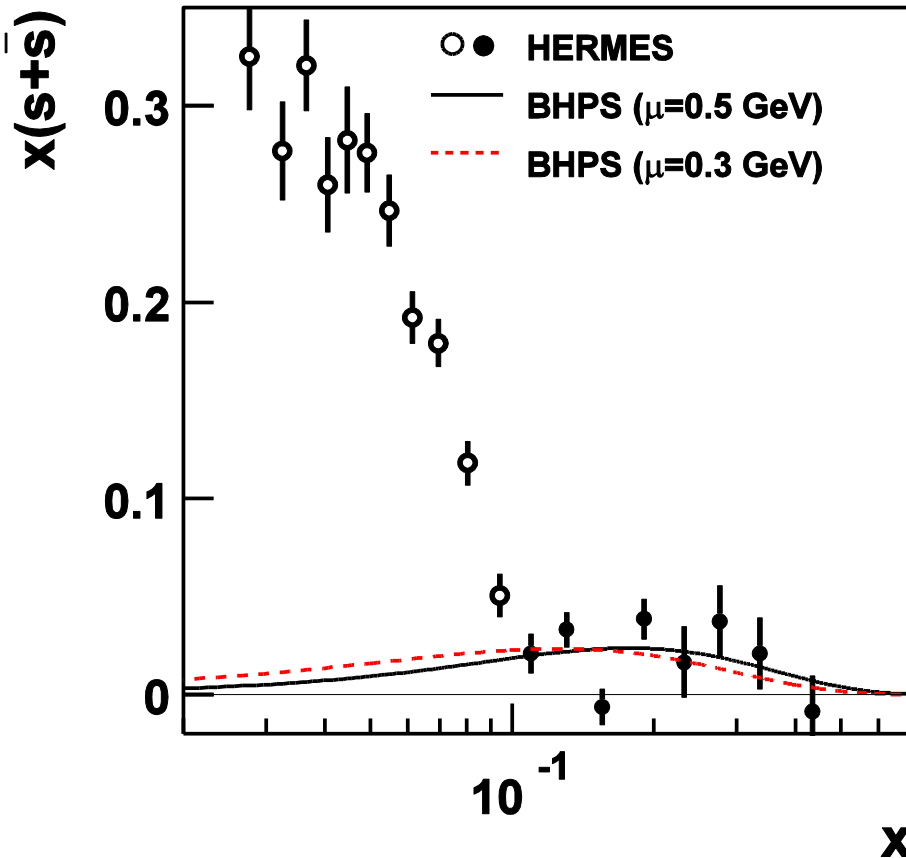


$s(x) + \bar{s}(x)$  from HERMES kaon  
SIDIS data at  $\langle Q^2 \rangle = 2.5 \text{ GeV}^2$

The data appear to consist  
of two different components  
(intrinsic and extrinsic?)

HERMES collaboration, Phys. Lett.  
B666, 446 (2008)

# Comparison between the $s(x) + \bar{s}(x)$ data with the intrinsic 5- $q$ model



$s(x) + \bar{s}(x)$  from HERMES kaon  
SIDIS data at  $\langle Q^2 \rangle = 2.5 \text{ GeV}^2$

Assume  $x > 0.1$  data are dominated  
by intrinsic sea (and  $x < 0.1$  are  
from QCD sea)

This allows the extraction of the  
intrinsic sea for strange quarks

(W. Chang and JCP, PL B704, 197(2011))

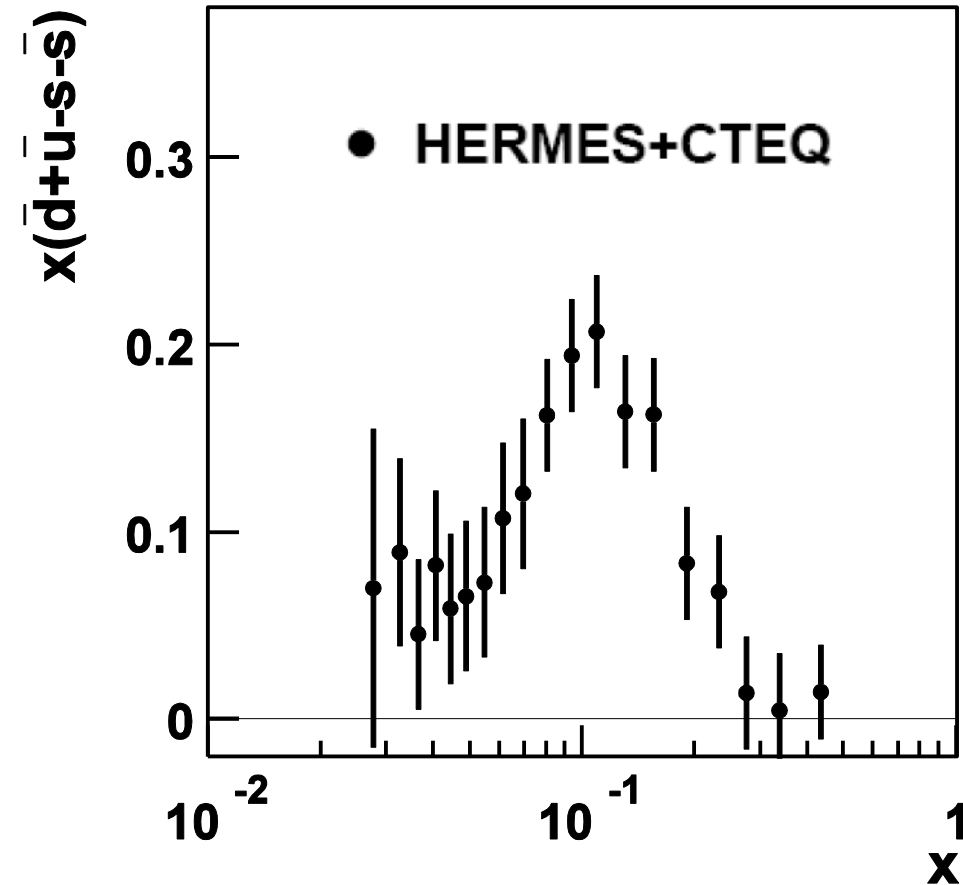
$$P_5^{uud\bar{s}} = 0.024$$

# How to separate the “intrinsic sea” from the “extrinsic sea”?

- Select experimental observables which have no contributions from the “extrinsic sea”

$\bar{d} + \bar{u} - s - \bar{s}$  has no contribution from extrinsic sea ( $g \rightarrow \bar{q}q$ )  
and is sensitive to "intrinsic sea" only

# Comparison between the $\bar{u}(x) + \bar{d}(x) - s(x) - \bar{s}(x)$ data with the intrinsic $5-q$ model



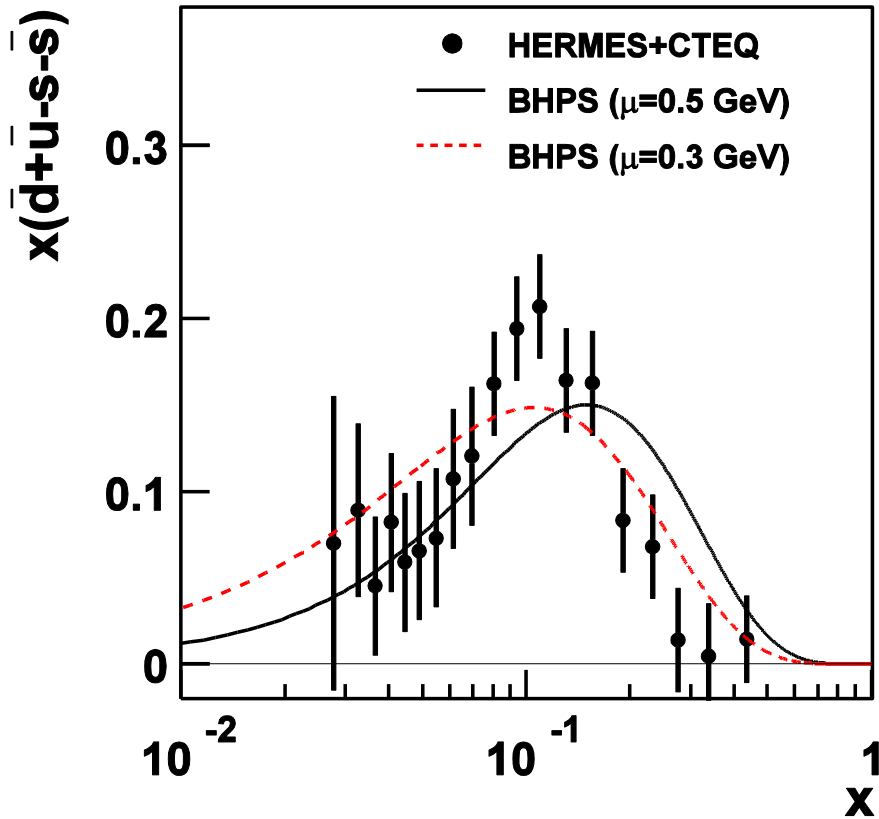
$\bar{d}(x) + \bar{u}(x)$  from CTEQ6.6  
 $s(x) + \bar{s}(x)$  from HERMES

$\bar{u} + \bar{d} - s - \bar{s}$  has  
no contribution  
from extrinsic sea

A valence-like  $x$ -distribution is observed



# Comparison between the $\bar{u}(x) + \bar{d}(x) - s(x) - \bar{s}(x)$ data with the intrinsic 5- $q$ model



$\bar{d}(x) + \bar{u}(x)$  from CTEQ6.6  
 $s(x) + \bar{s}(x)$  from HERMES

$$\bar{u} + \bar{d} - s - \bar{s}$$

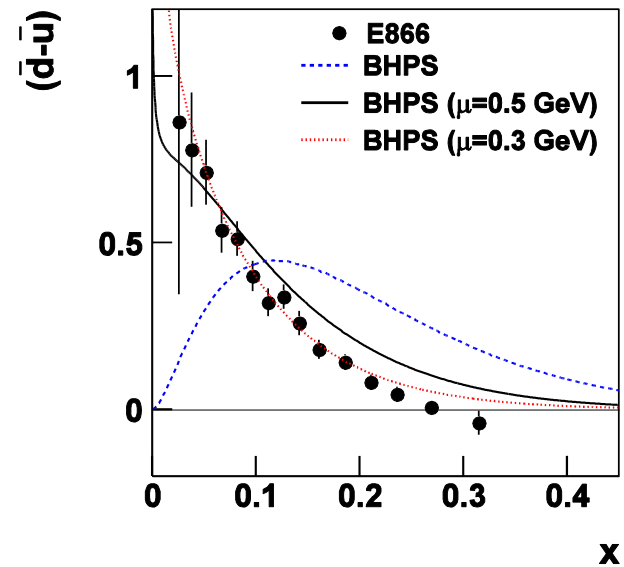
$$\sim P_5^{uudu\bar{u}} + P_5^{uudd\bar{d}} - 2P_5^{uuds\bar{s}}$$

(not sensitive to extrinsic sea)

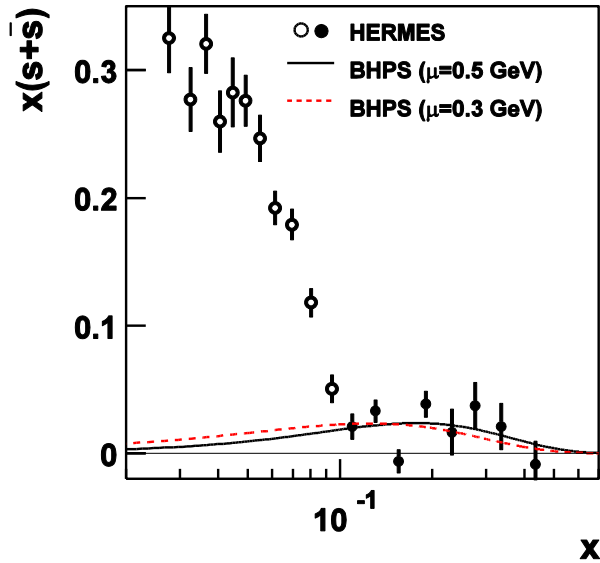
(W. Chang and JCP, PL B704, 197(2011))

$$P_5^{uudu\bar{u}} + P_5^{uudd\bar{d}} - 2P_5^{uuds\bar{s}} = 0.314$$

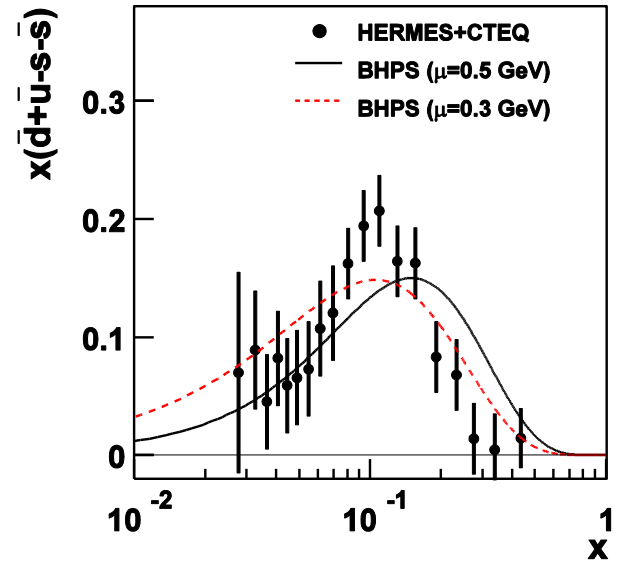
# Extraction of the various five-quark components for light quarks



$$P_5^{uudd\bar{d}} - P_5^{uud\bar{u}\bar{u}} = 0.118$$



$$P_5^{uud\bar{s}} = 0.024$$



$$P_5^{uud\bar{u}} + P_5^{uudd\bar{d}} - 2P_5^{uud\bar{s}} = 0.314$$

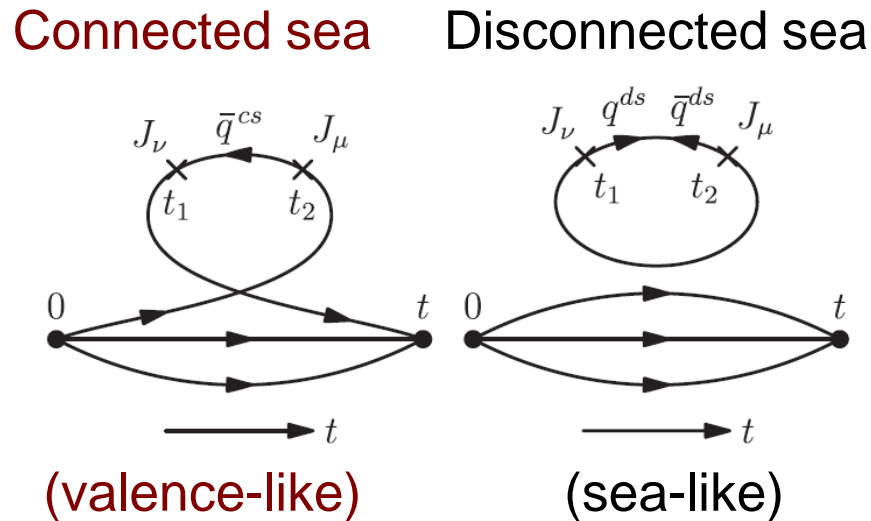
$$P_5^{uudd\bar{d}} = 0.240; \quad P_5^{uud\bar{u}\bar{u}} = 0.122; \quad P_5^{uud\bar{s}} = 0.024$$

# Future Possibilities

- Search for intrinsic charm and beauty at RHIC and LHC.
- Intrinsic gluons in the nucleons (Hoyer and Roy)?
- Spin-dependent observables of intrinsic sea?
- Global fits including intrinsic u, d, s sea?
- Intrinsic sea for hyperons and mesons?
- Connection between intrinsic sea and lattice QCD formalism?

## Connected-Sea Partons

Keh-Fei Liu,<sup>1</sup> Wen-Chen Chang,<sup>2</sup> Hai-Yang Cheng,<sup>2</sup> and Jen-Chieh Peng<sup>3</sup>

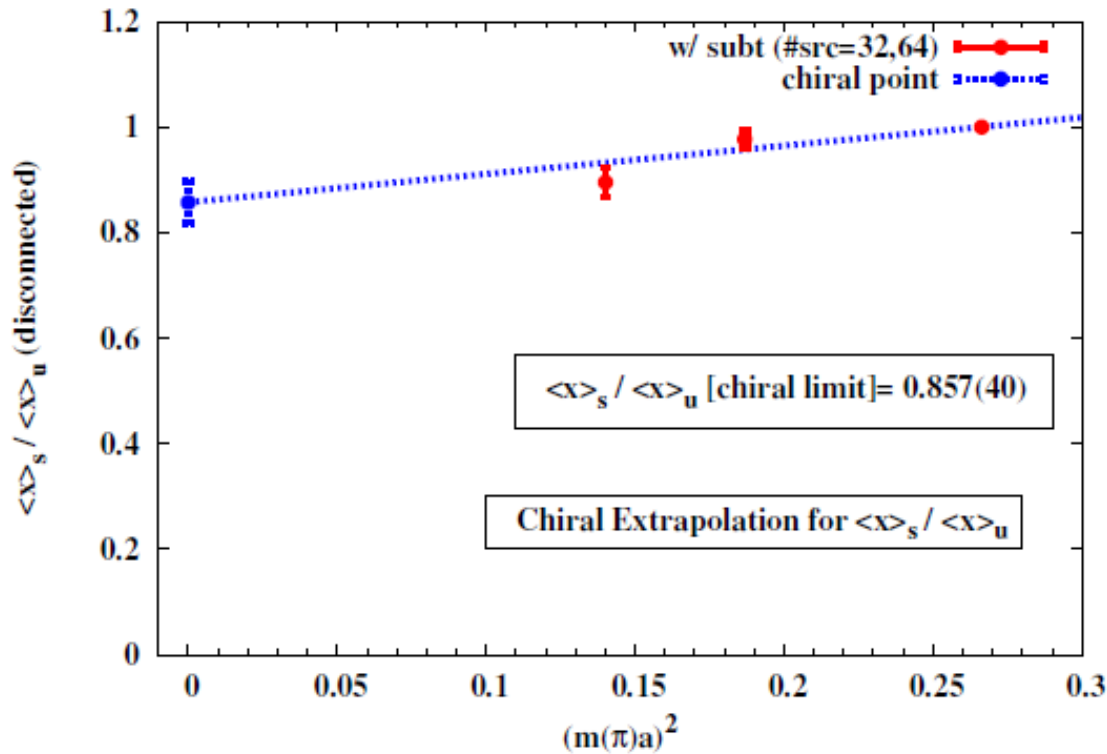


Two sources of sea:  
Connected sea (CS) and  
Disconnected sea (DS)

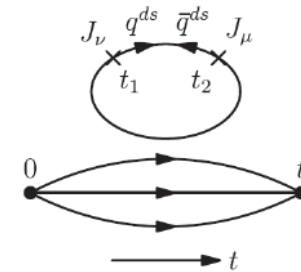
CS and DS have  
different Bjorken- $x$  and  
flavor dependences

- $x$  – dependence: at small  $x$ , CS  $\sim x^{-1/2}$ ; DS  $\sim x^{-1}$
- Flavor dependence:  $\bar{u}$  and  $\bar{d}$  have both CS and DS;  $\bar{s}$  is entirely DS

# Can one separate the “connected sea” from the “disconnected sea” for $\bar{u} + \bar{d}$ ?



Disconnected sea



$$R = \frac{\langle x \rangle_{s+\bar{s}}}{\langle x \rangle_{u+\bar{u}}} = 0.857(40)$$

for disconnected sea

(Doi et al., Pos lattice 2008, 163.)

Lattice QCD shows that disconnected sea is roughly SU(3)-flavor independent

Can one separate the “connected sea” from the “disconnected sea” for  $\bar{u} + \bar{d}$  ?

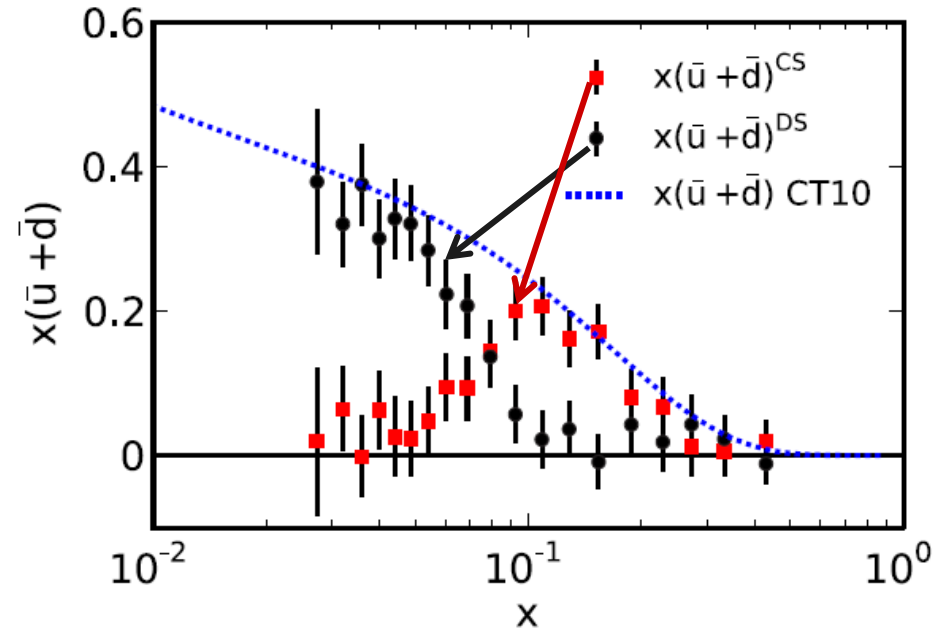
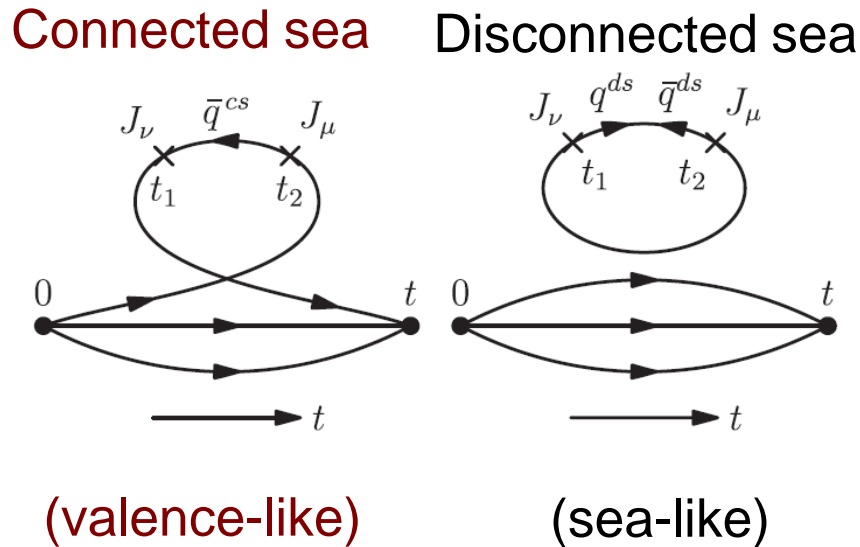
A) Lattice QCD shows that disconnected sea is roughly SU(3)-flavor independent

$$R = \frac{\langle x \rangle_{s+\bar{s}}}{\langle x \rangle_{u+\bar{u}}} = 0.857(40) \text{ for disconnected sea}$$

$$\text{B) } [\bar{u}(x) + \bar{d}(x)]_{\text{disconnected sea}} = \frac{1}{R} [s(x) + \bar{s}(x)]$$

$$\text{C) } [\bar{u}(x) + \bar{d}(x)]_{\text{connected sea}} = [\bar{u}(x) + \bar{d}(x)]_{\text{PDF}} - [\bar{u}(x) + \bar{d}(x)]_{\text{disconnected sea}}$$

## Connected-Sea Partons

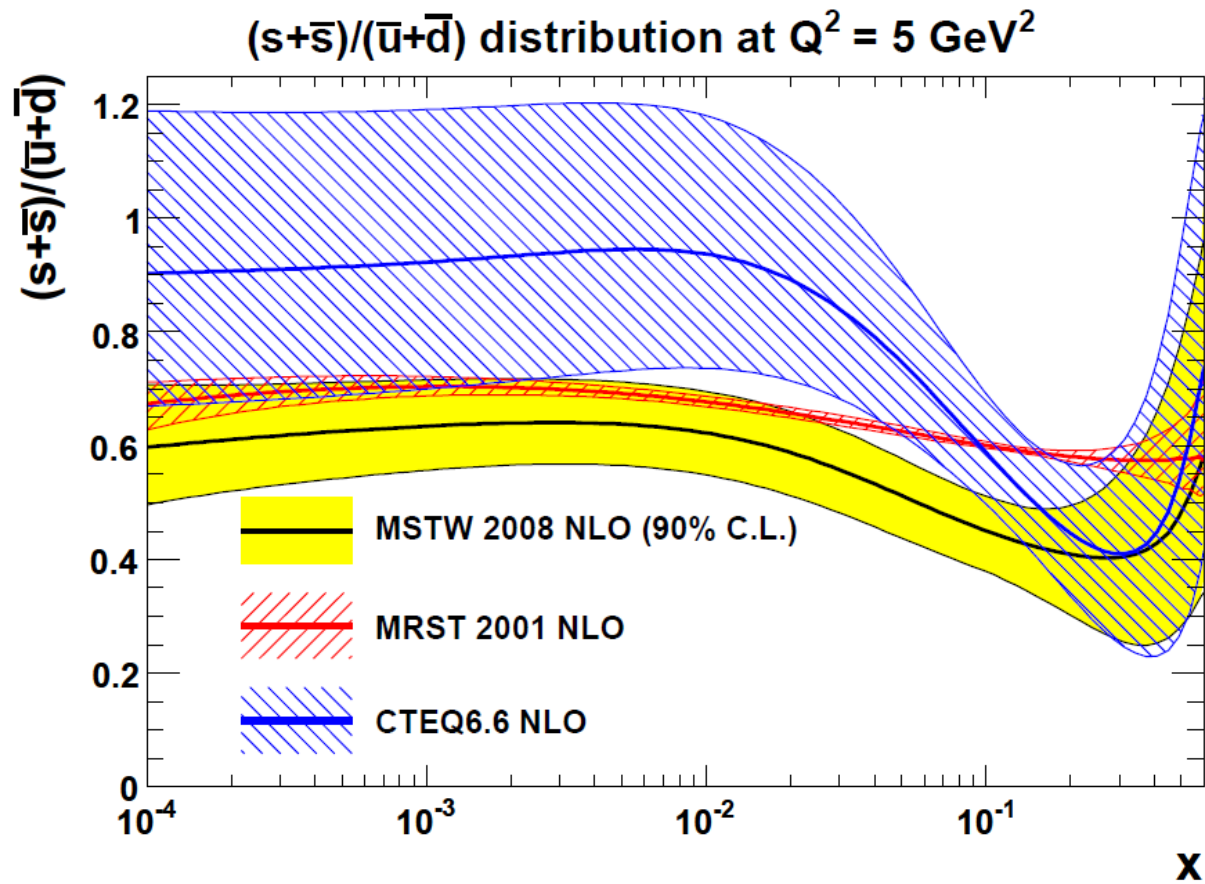
Keh-Fei Liu,<sup>1</sup> Wen-Chen Chang,<sup>2</sup> Hai-Yang Cheng,<sup>2</sup> and Jen-Chieh Peng<sup>3</sup>

- Using input from lattice QCD, one can separate the connected sea from the disconnected sea for  $\bar{u}(x) + \bar{d}(x)$
- For  $\bar{u} + \bar{d}$  at  $Q^2 = 2.5 \text{ GeV}^2$ , momenta carried by CS and DS are roughly equal

$$\text{Is } s(x) + \bar{s}(x) = \bar{u}(x) + \bar{d}(x) ?$$

Expectation:

$s$  and  $\bar{s}$  are suppressed relative to  $\bar{u}$  and  $\bar{d}$  due to larger  $s$ -quark mass





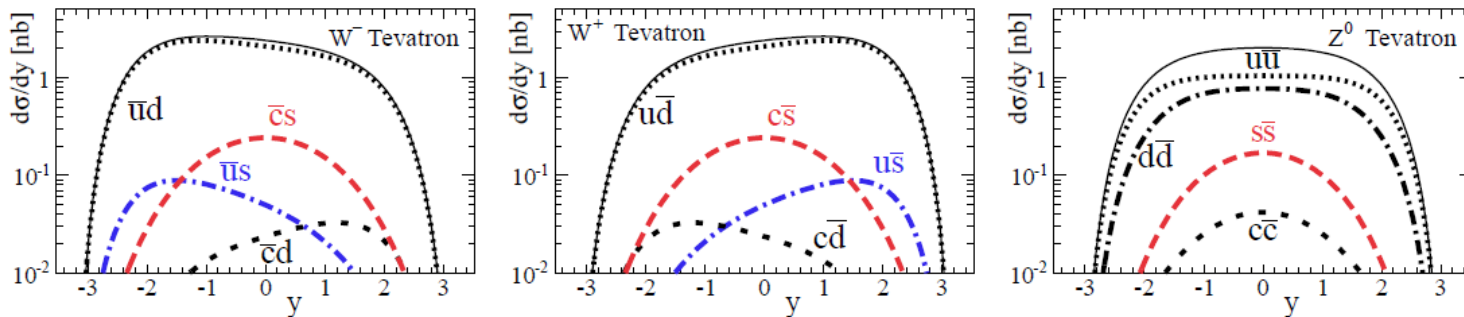
# Strange sea from inclusive W/Z production

Inclusive W / Z production at Tevatron/LHC

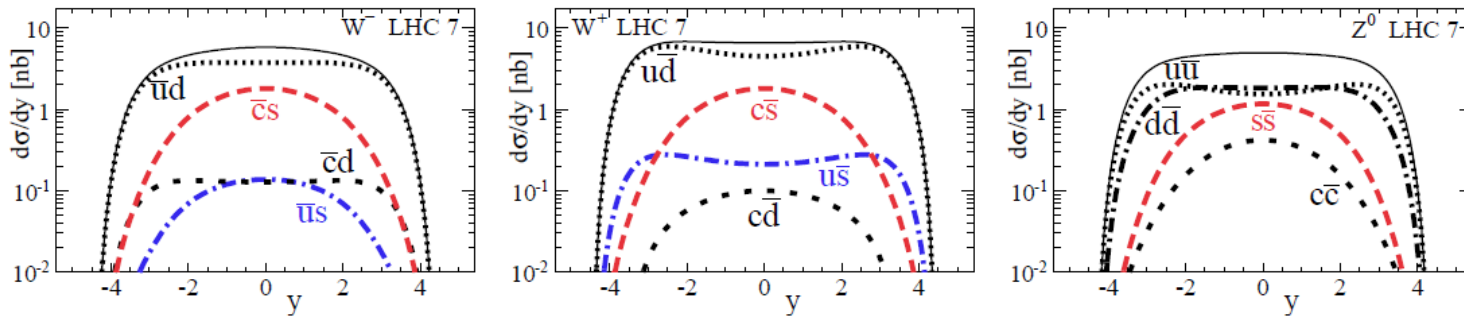
$$W^+ : (u \text{ or } c) + (\bar{d} \text{ or } \bar{s}) \rightarrow W^+$$

$$W^- : (\bar{u} \text{ or } \bar{c}) + (d \text{ or } s) \rightarrow W^-$$

$$Z^0 : s + \bar{s} \rightarrow Z^0$$

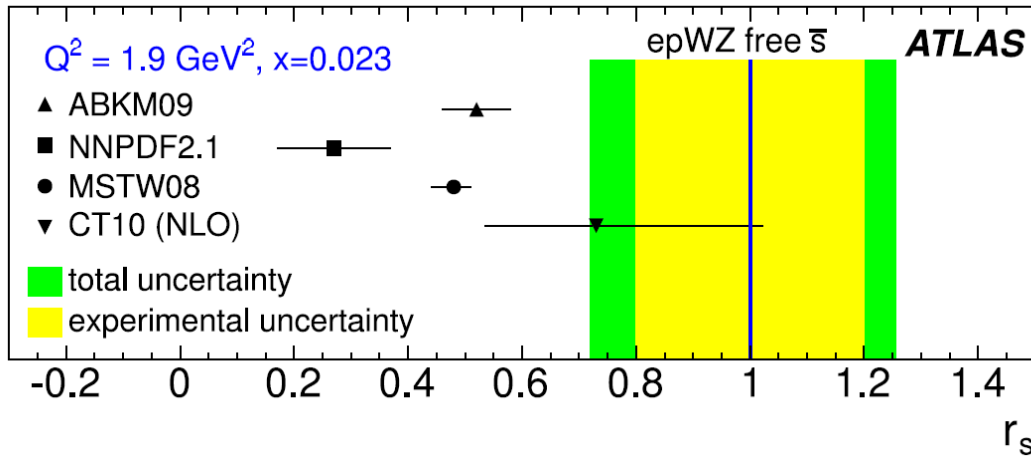
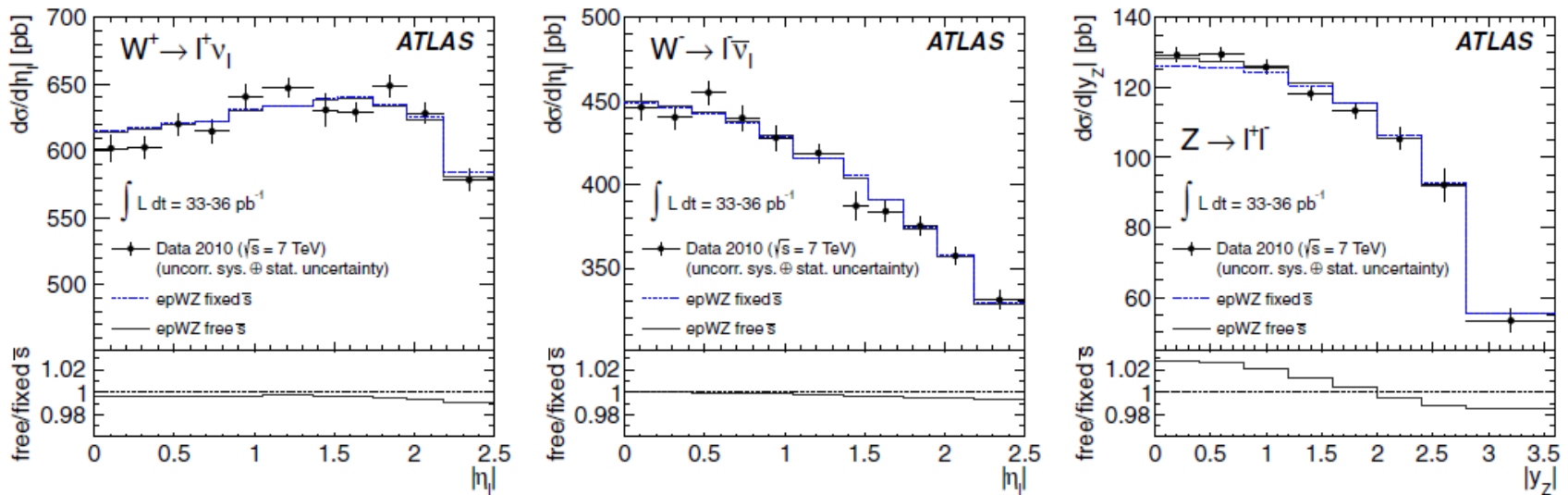


(a)  $d\sigma/dy$  for  $W^-$  (left),  $W^+$  (middle),  $Z^0$  (right) boson production at the Tevatron.



(b)  $d\sigma/dy$  for  $W^-$  (left),  $W^+$  (middle),  $Z^0$  (right) boson production at the LHC with  $\sqrt{S} = 7$  TeV.

# Strange sea from inclusive W/Z production



$$r_s = (s + \bar{s}) / 2\bar{d} = 1.00^{+0.09}_{-0.10}$$

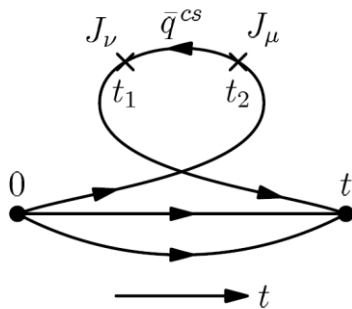
$$\text{at } x=0.013, Q^2 = M_Z^2$$

# Strange sea content is strongly $x$ dependent

- Perturbative sea at small  $x$  is roughly SU(3) symmetric
- Non-perturbative sea at larger  $x$  is SU(3) asymmetric

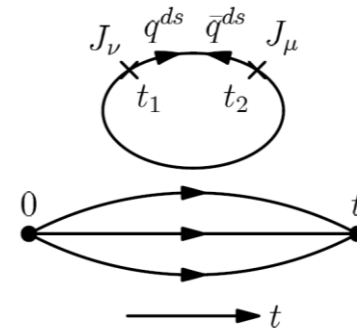
Can be well understood from Lattice QCD  
(PRL 109 (2012)252002)

Connected sea



Generate "valence-like"  
 $\bar{u}(x)$  and  $\bar{d}(x)$  (no  $\bar{s}(x)$ )  
at larger  $x$

Disconnected sea



Generate roughly symmetric  
 $s(x), \bar{s}(x), \bar{u}(x)$  and  $\bar{d}(x)$   
at small  $x$

# Strange sea content is strongly $Q^2$ dependent

$$\kappa(x, Q) = \frac{2s(x, Q)}{\bar{u}(x, Q) + \bar{d}(x, Q)}$$

Kusina et al., PRD 85 (2012)  
094028

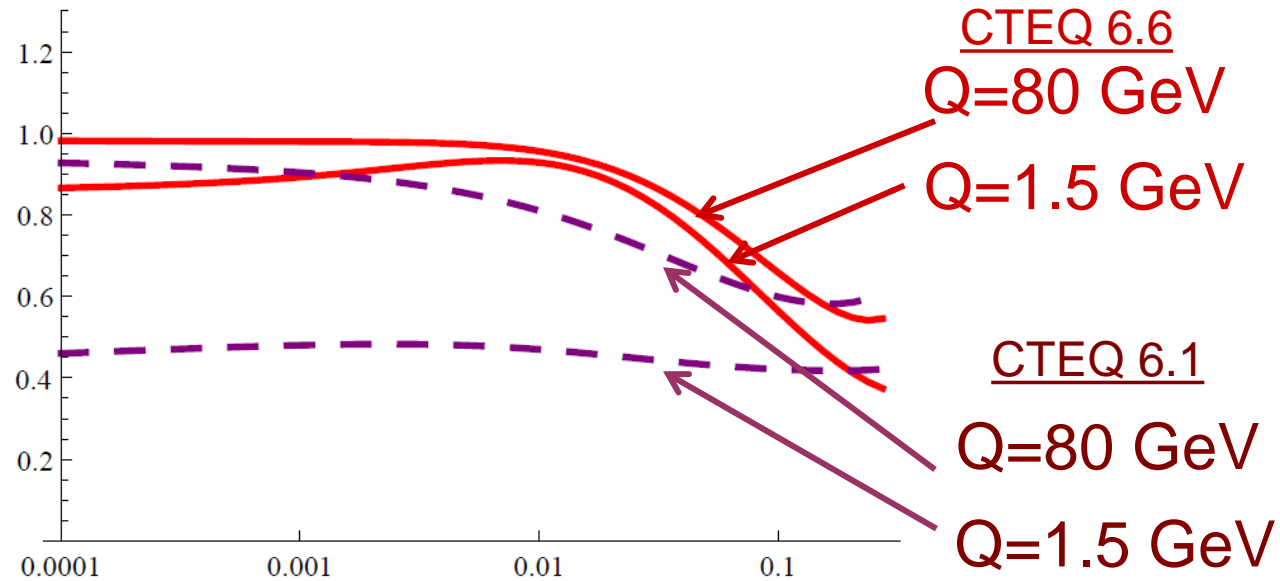


Figure 5:  $\kappa(x, Q)$  vs.  $x$  showing the evolution from low to high scales. The solid (red) lines are for CTEQ6.6, and the dashed (purple) lines are for CTEQ6.1. The lower

W/Z productions are sensitive to  $s(x), \bar{s}(x)$  at very large  $Q^2$  scale ( $Q^2 = M_{W/Z}^2$ ), dominated by perturbative roughly SU(3) symmetric sea!

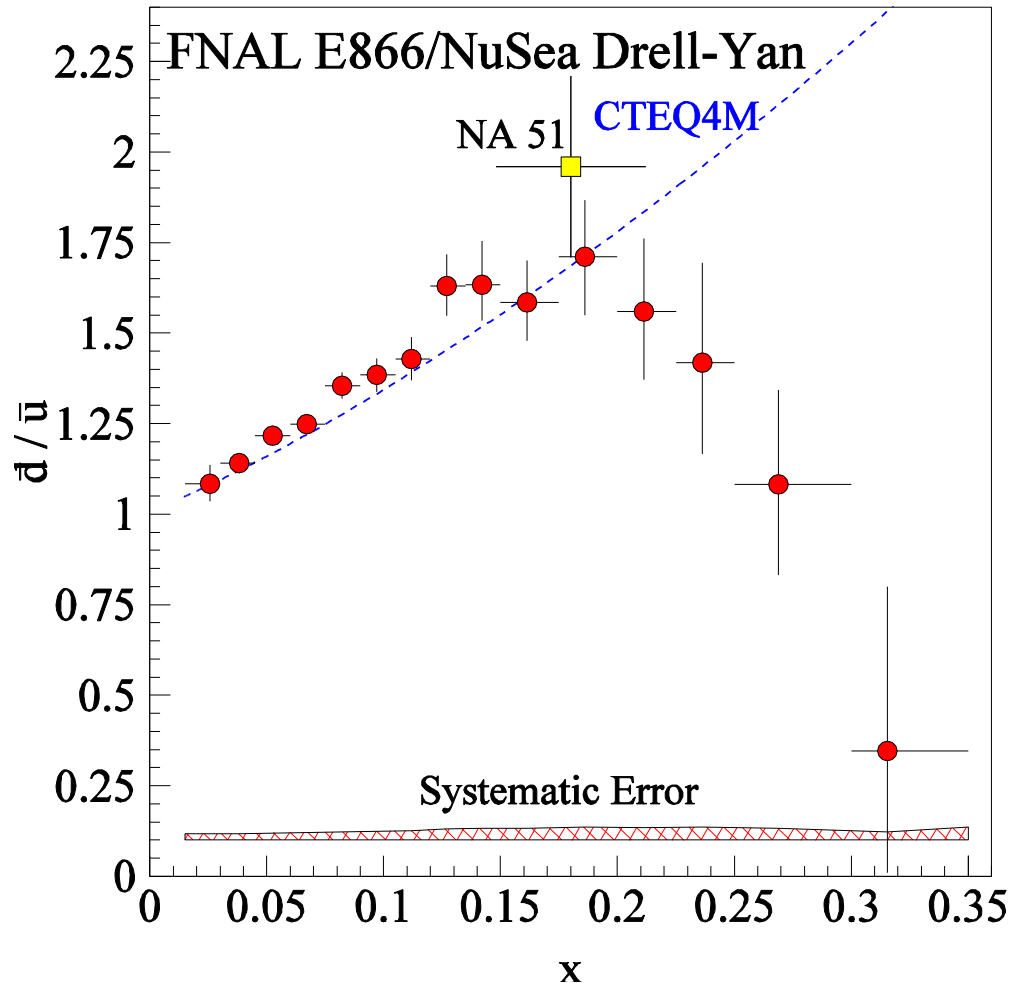
Measurements at low  $Q^2$  are very important

# Conclusions

- Evidences for the existence of "intrinsic" light-quark seas ( $\bar{u}, \bar{d}, \bar{s}$ ) in the nucleons.
- Clear evidence for intrinsic charm remains to be found.
- The flavor structures of the nucleon sea and their Bjorken- $x$  dependence provide strong constraints on theoretical models.
- The concept of connected and disconnected seas in Lattice QCD offers useful insights on the flavor- and  $x$ -dependences of the sea.
- Ongoing and future Drell-Yan and SIDIS experiments will provide crucial new information.

# Backup Slides

# Does $\bar{d} / \bar{u}$ drop below 1 at large $x$ ?



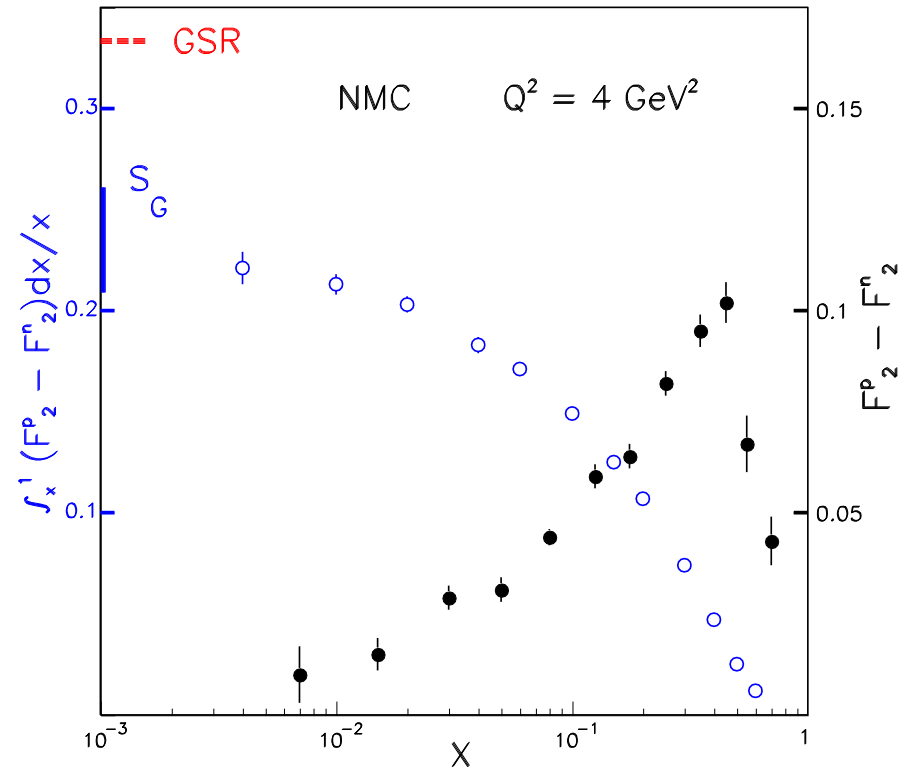
No existing models can explain sign-change

for  $\bar{d}(x) - \bar{u}(x)$  at any value of  $x$

# Revisit the NMC measurement of the Gottfried Sum rule

## The Gottfried Sum Rule

$$\begin{aligned}
 S_G &= \int_0^1 [(F_2^p(x) - F_2^n(x)) / x] dx \\
 &= \frac{1}{3} + \frac{2}{3} \int_0^1 (\bar{u}_p(x) - \bar{d}_p(x)) dx \\
 &= \frac{1}{3} \quad (\text{if } \bar{u}_p = \bar{d}_p)
 \end{aligned}$$



New Muon Collaboration (NMC) obtains

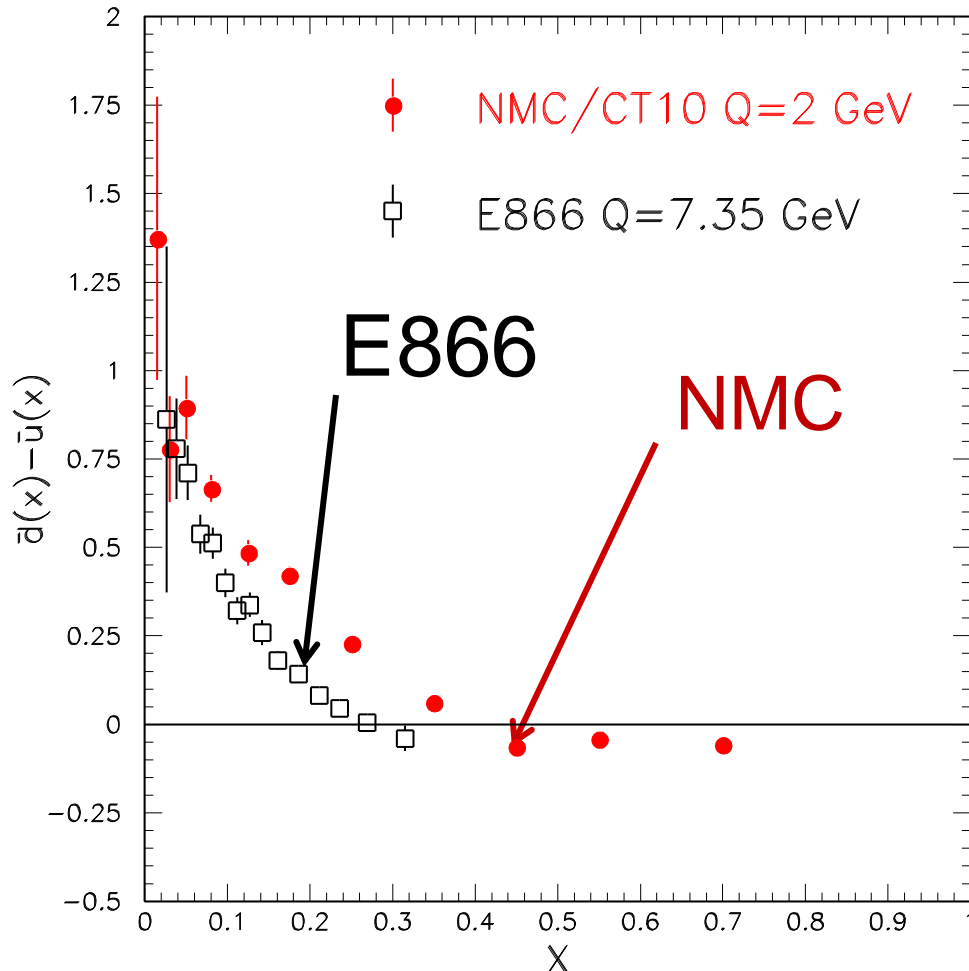
$$S_G = 0.235 \pm 0.026$$

( Significantly lower than 1/3 ! )  $\Rightarrow \bar{d} \neq \bar{u}$  ?



# Extracting $\bar{d}(x) - \bar{u}(x)$ from the NMC data

$$\bar{d}(x) - \bar{u}(x) = [u_V(x) - d_V(x)]_{CT10} / 2 - 3/2 * [F_2^p(x) / x - F_2^n(x) / x]_{NMC}$$

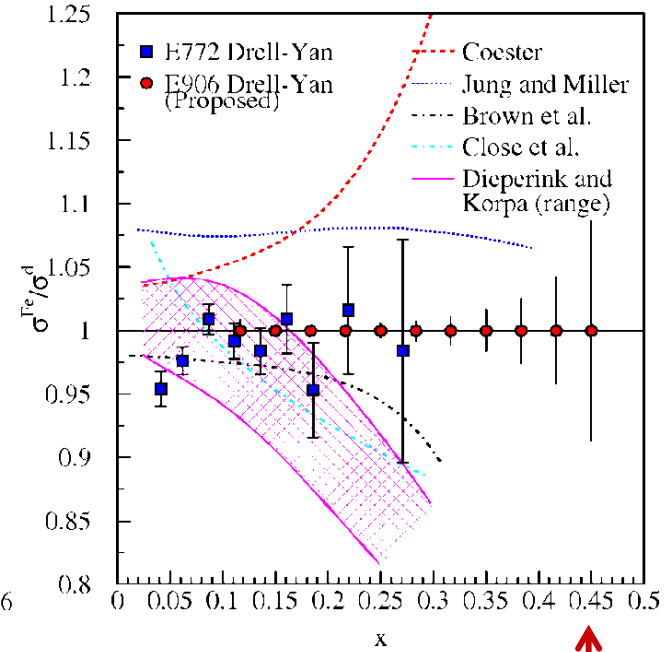
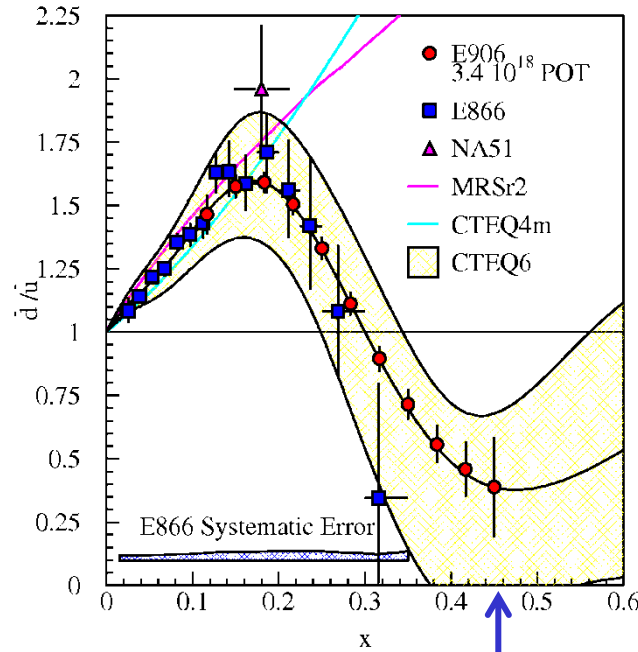


The NMC data, together with the recent PDF, already suggest that  $\bar{d}(x) - \bar{u}(x) < 0$  at large  $x$  !

(JCP, W.C. Chen, K.F. Liu)

# Drell-Yan Experiment at Fermilab

SeaQuest Experiment ( Unpolarized Drell-Yan using 120 GeV proton beam)



Main goals: 1) Measure  $\bar{d} / \bar{u}$  flavor asymmetry up to  $x \sim 0.45$   
 2) Measure EMC effect of antiquarks up to  $x \sim 0.45$

- Commission run took place in February – April 2012
- 2-year production run expected in 2013-2015