

GPD analyses of kaons leptonproduction

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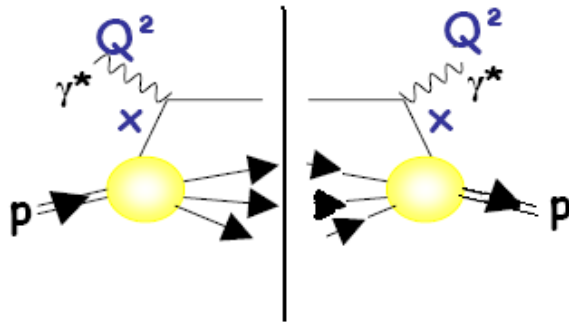
In collaboration with P. Kroll, Wuppertal

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- Handbag factorization .
- GPDs and amplitudes structure.
- Kaons leptonproduction .
- Polarized GPDs and transversity effects.
- Results on kaons production.
- Conclusion.

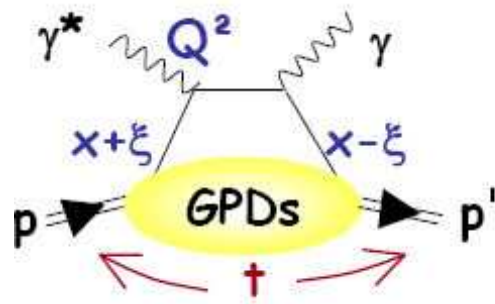
DIS and DVCD

- Deep Inelastic scattering



Cross section -
expressed in terms of
ordinary parton
distributions $q(x)$

- Deeply Virtual Compton Scattering



Amplitude - proportional to
Generalized Parton
Distributions
GPDs $H(x, \xi, t)$

GPDs – extensive information about hadron structure.

- Ordinary parton distribution connected with GPDs

$$H^g(x, 0, 0) = xg(x)$$

.....

- Hadron Form factors –are the GPDs moment

$$\int dx H(x, \xi, t) = F(t)$$

.....

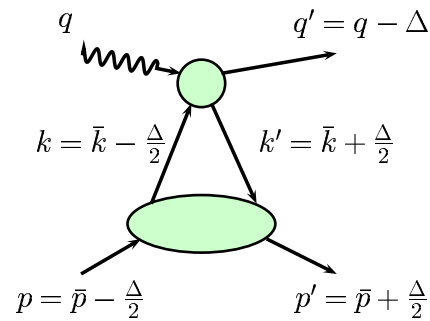
- Information on the parton angular momenta from Ji sum rules

$$\int x dx (H^q(x, \xi, 0) + E^q(x, \xi, 0)) = 2J^q$$

Handbag factorization of Mesons production amplitude

- Large Q^2 - factorization into a hard meson photoproduction off partons, and GPDs. (LL)

Radyushkin, Collins, Frankfurt Strikman



$$k = ((\bar{x} + \xi)p_+, \dots)$$

$$k' = ((\bar{x} - \xi)p_+, \dots)$$

$$k \neq k'; \quad \xi \sim \frac{x_B}{2 - x_B}$$

$L \rightarrow L$ transition - predominant. Other amplitudes are suppressed as powers $1/Q$

The process of meson production

- ϕ production (gluon&strange sea)
- ρ, ω production (gluon&sea&valence quarks)
- Pseudoscalar mesons- polarized distributions

The handbag model typically is valid at the range of large $Q^2 > 3\text{GeV}^2$ and low $x_B \leq 0.1$.

Modelling the GPDs

The double distributions for GPDs **Radyushkin '99** .

$$H_i(\bar{x}, \xi, t) = \int_{-1}^1 d\beta \int_{-1+|\beta|}^{1-|\beta|} d\alpha \delta(\beta + \xi \alpha - \bar{x}) f_i(\beta, \alpha, t) \quad (1)$$

simple for the double distributions.

$$f_i(\beta, \alpha, t) = h_i(\beta, t) \frac{\Gamma(2n_i + 2)}{2^{2n_i+1} \Gamma^2(n_i + 1)} \frac{[(1 - |\beta|)^2 - \alpha^2]^{n_i}}{(1 - |\beta|)^{2n_i+1}}, \quad (2)$$

★ $h_{val}^q(\beta, 0) = q_{val}(|\beta|) \Theta(\beta)$ –valence contribution (n=1).

PDF t -dependence —Regge parameterization. Regge form: $\alpha_i(t) = \alpha_i(0) + \alpha' t$

$$h(\beta, t) = N e^{b_0 t} \beta^{-\alpha(t)} (1 - \beta)^n \quad (3)$$

★ Amplitudes in terms of GPDs.

The proton non-flip amplitude is associated with F GPDs.

$$\mathcal{M}_{\mu'+,\mu+} \propto \int_{-1}^1 d\bar{x} H^a(\bar{x}, \xi, t) F_{\mu',\mu}^a(\bar{x}, \xi)$$

k_{\perp}/Q^2 and Sudakov corrections are taken into account.

$$H^a(x, 0, 0) = h^a(x), \quad H^g(x, 0, 0) = xg(x)$$

Quark (valence, sea), gluon PDFs are determined from CTEQ6 parameterization

★ Spin-flip contribution. Effects of E GPDs.

$$\mathcal{M}_{\mu'-,\mu+} \propto \frac{\sqrt{-t}}{2m} \int_{-1}^1 d\bar{x} E^a(\bar{x}, \xi, t) F_{\mu',\mu}^a(\bar{x}, \xi)$$

E parameters- from Pauli form factor. M. Diehl, ...,P.Kroll

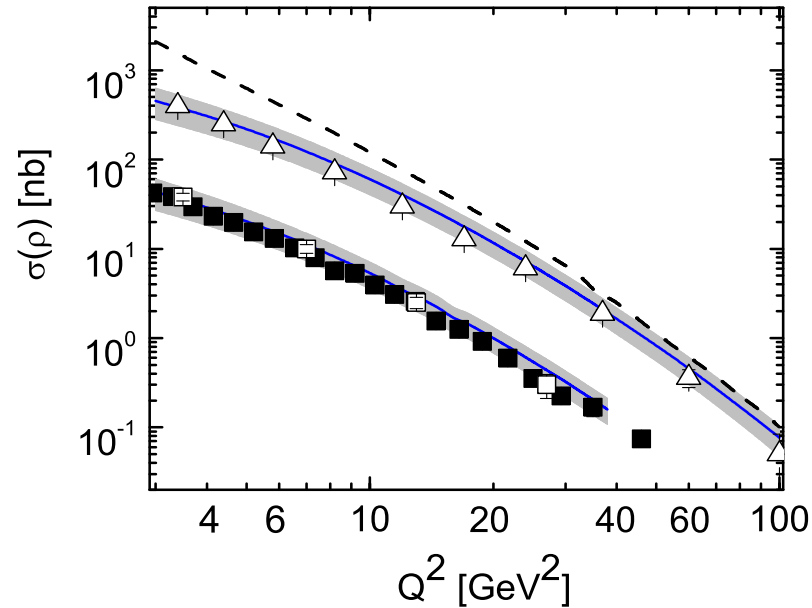
Standard connection with ordinary distribution :

$$E^a(x, 0, 0) = e^a(x)$$

Double distribution model is used to construct all GPDs.

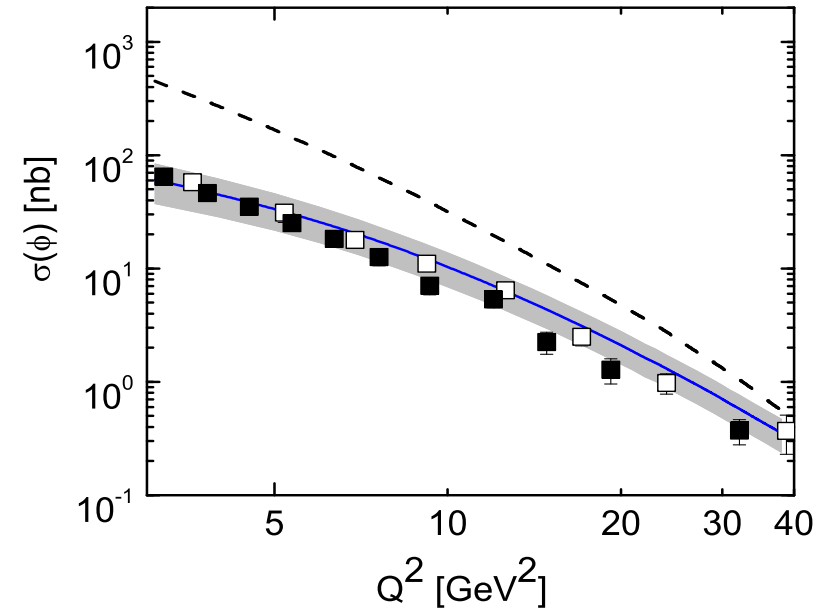
Cross sections of VM production

Q^2 dependence of cross sections of ρ and ϕ production at $W = 75\text{GeV}$. H1 and ZEUS data.



Cross sections of ρ production with errors from uncertainty in parton distributions at $W = 75\text{GeV}/10$ and $W = 90\text{GeV}$. Dashed line leading twist results.

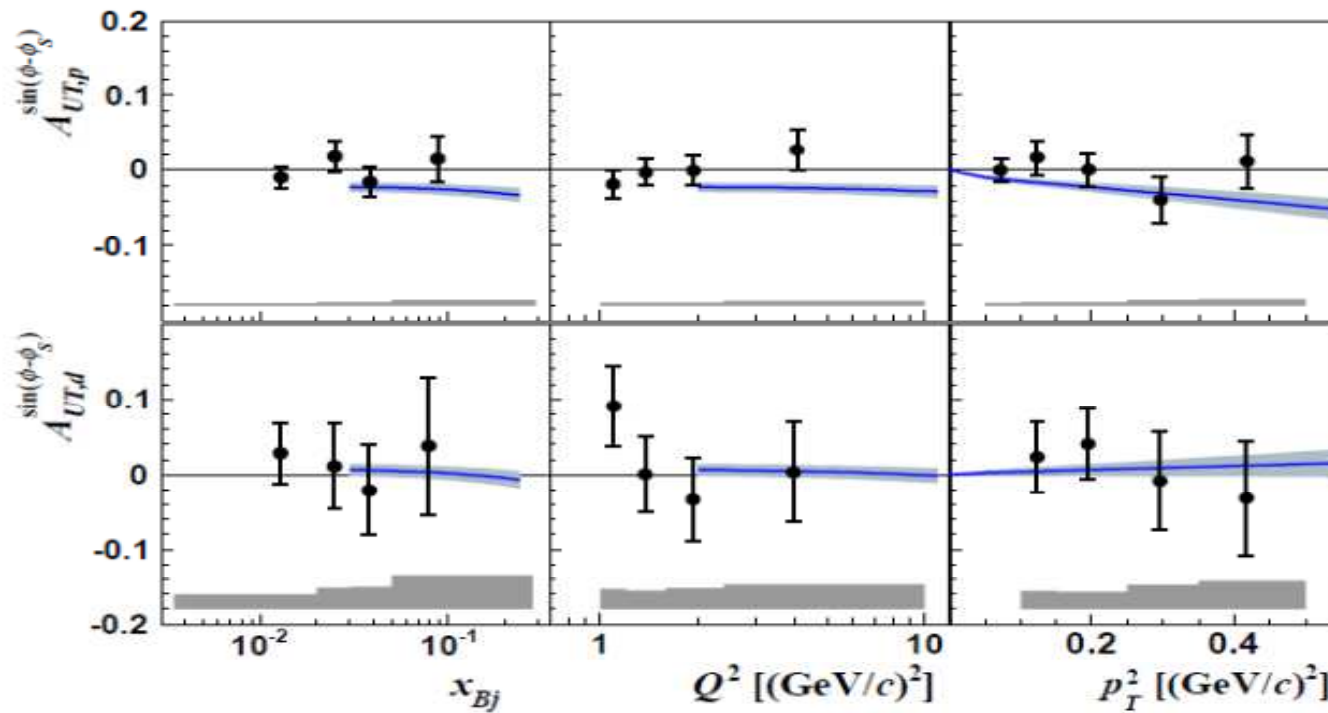
★ Power corrections $\sim k_{\perp}^2/Q^2$ in propagators are important at low Q^2 —1/10 suppression at $Q^2 \sim 3\text{GeV}^2$



Cross sections of ϕ production with errors from uncertainty in parton distributions at $W = 75\text{GeV}$. Dashed line leading twist results.

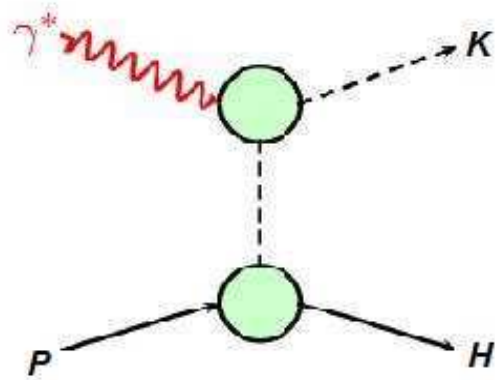
A_{UT} asymmetries in VM production.

$$A_{UT} = \propto \frac{\text{Im} \langle E \rangle^* \langle H \rangle}{|\langle H \rangle|^2}$$



Our results on $A_{UT}^{\sin(\phi - \phi_s)}$ asymmetry for ρ production on the proton and deuteron together with COMPASS data.

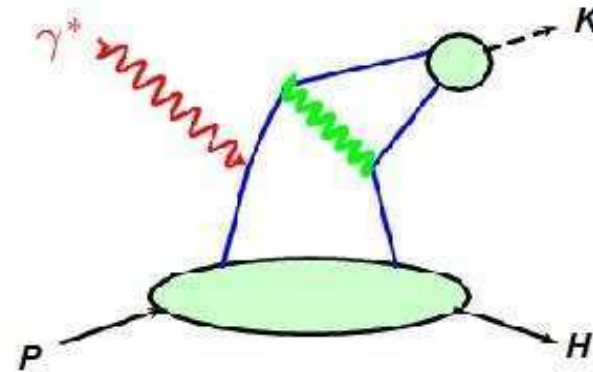
Pseudoscalar K meson production.



Pole

and

handbag



contributions to Kaons production.

★ Meson pole (charge mesons) is essential mainly in $\mathcal{M}_{0+,0+}$ amplitude.

★ Contribute to $\mathcal{M}_{0-,0+}$, $\mathcal{M}_{0+,++}$, $\mathcal{M}_{0-,++}$ amplitudes too.

General case of Pseudoscalar meson production

$$\mathcal{M}_{0+,0+}^M \propto \sqrt{1-\xi^2} \left[\langle \tilde{H}^M \rangle - \frac{\xi^2}{1-\xi^2} \langle \tilde{E}_{n.p.}^M \rangle - \frac{\xi(m_{Ni} + M_{Nf})Q^2}{1-\xi^2} \frac{\rho_M}{t - m_M^2} \right]; \quad (4)$$

$$\mathcal{M}_{0-,0+}^M \propto \frac{\sqrt{-t'}}{(m_{Ni} + M_{Nf})} \left[\xi \langle \tilde{E}_{n.p.}^M \rangle + (m_{Ni} + M_{Nf})Q^2 \frac{\rho_M}{t - m_M^2} \right].$$

Masses: M - produced pseudoscalar meson, N^i -initial nucleon (proton), M_{Nf} -final nucleon (Λ , Σ)

$$\langle \tilde{F} \rangle = \sum_{\lambda} \int_{-1}^1 d\bar{x} \mathcal{H}_{0\lambda,0\lambda}(\bar{x}, \dots) \tilde{F}(\bar{x}, \xi, t), \quad .$$

The hard scattering amplitudes-transverse quark motion

$$H_{0\lambda,0\lambda}^a(\bar{x}, \xi) = \frac{8\pi\alpha_s(\mu_R)}{\sqrt{2N_c}} \int_0^1 d\tau \int \frac{d^2\mathbf{k}_{\perp}}{16\pi^3} \phi_{V\mu'}(\tau, k_{\perp}^2) f_{0\lambda,0\lambda}^a(\mathbf{k}_{\perp}, \bar{x}, \xi, \tau) D. \quad (5)$$

$$\phi_V(\mathbf{k}_{\perp}, \tau) = 8\pi^2 \sqrt{2N_c} f_V a_V^2 \exp \left[-a_V^2 \frac{\mathbf{k}_{\perp}^2}{\tau\bar{\tau}} \right]. \quad (6)$$

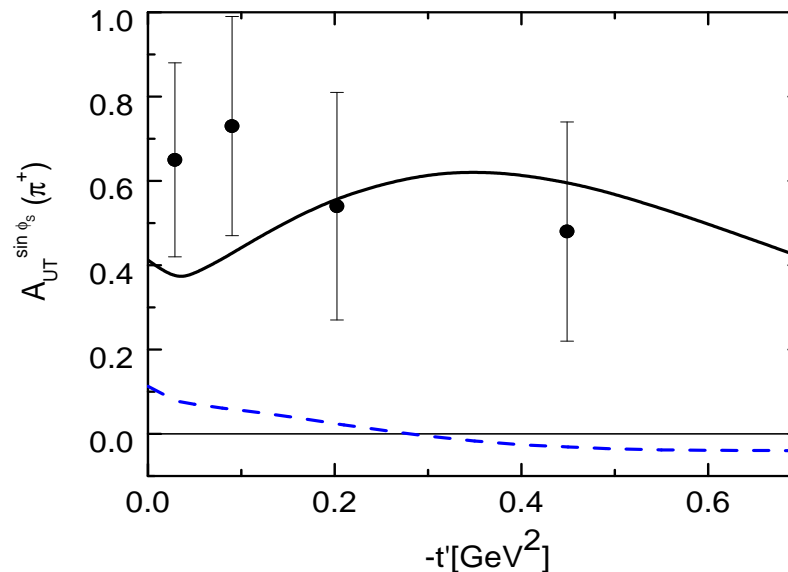
Meson pole contribution (charge meson production)

$$\rho_M = g_{MN^i Nf} F_{MN^i Nf}(t) F_M(Q^2) \quad (7)$$

Why leading twist effects is not enough at low Q^2 ?

At low Q^2 we have problems with understanding of some observables.

Example: $A_{UT}^{\sin(\phi_s)}$ asymmetry in π^+ production.



$$A_{UT}^{\sin(\phi_s)} \propto \text{Im}[M_{0-,++}^* M_{0+,0+}]$$

The handbag amplitude $M_{0-,++} \propto t'$. Small pole effect in $M_{0-,++}$ can not explain asymmetry. New not small contribution to $M_{0-,++}$ amplitude is needed.

Calculation of $M_{0-,++}$ – transversity effects.

$M_{\mu'\nu',\mu\nu} \propto \sqrt{-t'}^{|\mu-\nu-\mu'+\nu'|}$ from angular momentum conservation.

$M_{0-,++} \propto \sqrt{-t'}^0 \propto \text{const}$ but handbag amplitude $\propto t'$

$M_{0-,++}$ -is determined by twist 3 contribution $\rightarrow \text{const}$.

Transversity GPDs (H_T, E_T, \dots) contribute

$$\mathcal{M}_{0-, \mu+}^{twist-3} \propto \int_{-1}^1 d\bar{x} \mathcal{H}_{0-, \mu+}^{twist-3}(\bar{x}, \dots) [H_T + \dots O(\xi^2 E_T)].$$

We calculate twist-3 amplitude and use twist-3 meson wave function.

Double distribution model

$$H_T^a(x, 0, 0) = \delta^a(x)$$

transversity PDFs –from azimuthal asymmetry in semi-inclusive DIS (Anselmino model)

$$\delta^a(x) = C N_T^a x^{1/2} (1-x) [q_a(x) + \Delta q_a(x)],$$

Estimation of $M_{0+,++}$ – transversity effects.

Amplitude is important in some asymmetries and cross section σ_T, σ_{TT} e.g.

$$\mathcal{M}_{0+,\mu+}^{twist-3} \propto \frac{\sqrt{-t'}}{4m} \int_{-1}^1 d\bar{x} \mathcal{H}_{0-,\mu+}^{twist-3}(\bar{x}, \dots) \bar{E}_T.$$

Similar calculation of twist-3 amplitude as for H_T

$$\bar{E}_T(\beta, 0, 0) = e_T(\beta); \quad e_T(\beta, t) = N e^{b_0 t} \beta^{-\alpha(t)} (1 - \beta)^n \quad (8)$$

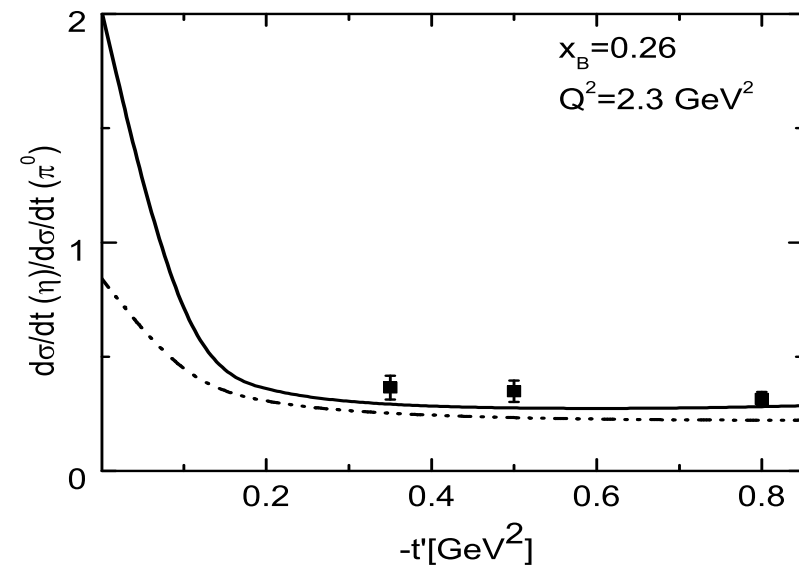
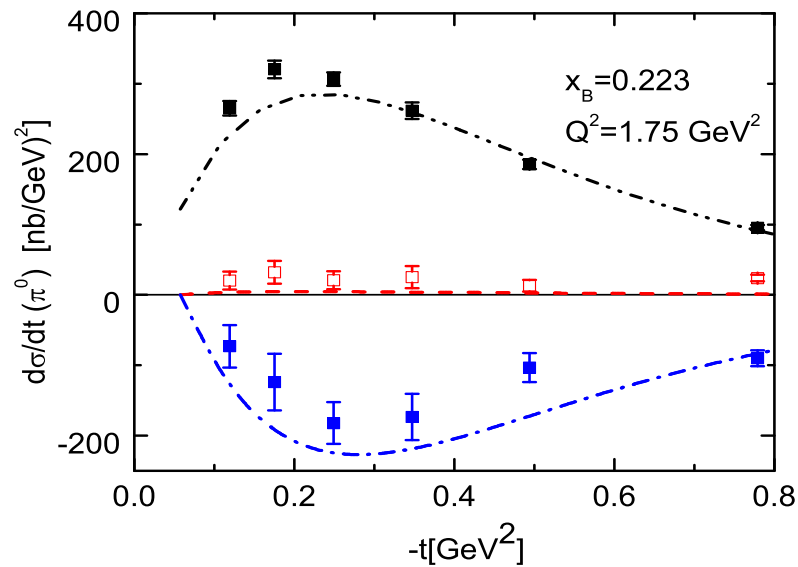
Double distribution model for \bar{E}_T

Parameters are taken from the lattice results for the moments of E_T

Moments for u and d are large and have the same sign and **not very different each other**

$$\star \text{ Enhancement for } \pi^0: \bar{E}_T^0 = 2/3 \bar{E}_T^u + 1/3 \bar{E}_T^d$$

Transversity effects at CLAS.



Predictions for π^0 production at CLAS energy range together with CLAS data. Full line- $\sigma_T + \epsilon\sigma_L$, red dashed line- σ_{LT} , blue dashed-dotted- σ_{TT}

η/π^0 production ratio at CLAS energy range together with preliminary data.

Transversity contributions are essential in the cross section of pseudoscalar meson production. Model predictions for π^0 cross sections and η/π^0 ratio were confirmed later by CLAS experiment.

Kaon production. Coupling constants and GPDs

- Wave function -non symmetric over $\tau \rightarrow \bar{\tau}$. Reason-quarks have different masses.
- Coupling constants-SU(3) predictions

$$g_{K^+p\Lambda} \sim -13.3; \quad g_{K^+p\Sigma^0} \sim -3.5;$$

Pole contribution is larger for $K^+p\Lambda$ channel .

- Form Factors: $F_{K^+}(Q^2) \sim 0.9 F_{\pi^+}(Q^2)$ in agreement with CLEO data

GPDs in kaon production.

Proton- hyperon transition GPDs contracted with the help of SU(3) flavor symmetry.

$$\star \gamma p \rightarrow K^+\Lambda : \quad \tilde{F}_{p \rightarrow \Lambda} \sim -\frac{1}{\sqrt{6}}[2\tilde{F}^u - \tilde{F}^d - \tilde{F}^s]$$

$$\star \gamma p \rightarrow K^+\Sigma^0 : \quad \tilde{F}_{p \rightarrow \Sigma^0} \sim -\frac{1}{\sqrt{2}}[\tilde{F}^d - \tilde{F}^s]$$

$$\star \gamma p \rightarrow K^0\Sigma^+ : \quad \tilde{F}_{p \rightarrow \Sigma^+} \sim -[\tilde{F}^d - \tilde{F}^s]$$

Polarized distributions contribute.

Transversity effects at kaons production.

★ Parameterization the same as for π^0 production .

Large H_T contributions for Kaons especially in $K^+\Lambda$ channel.

$$H_T(p \rightarrow \Lambda) \sim -\frac{1}{\sqrt{6}}[2H_T^u - H_T^d]$$

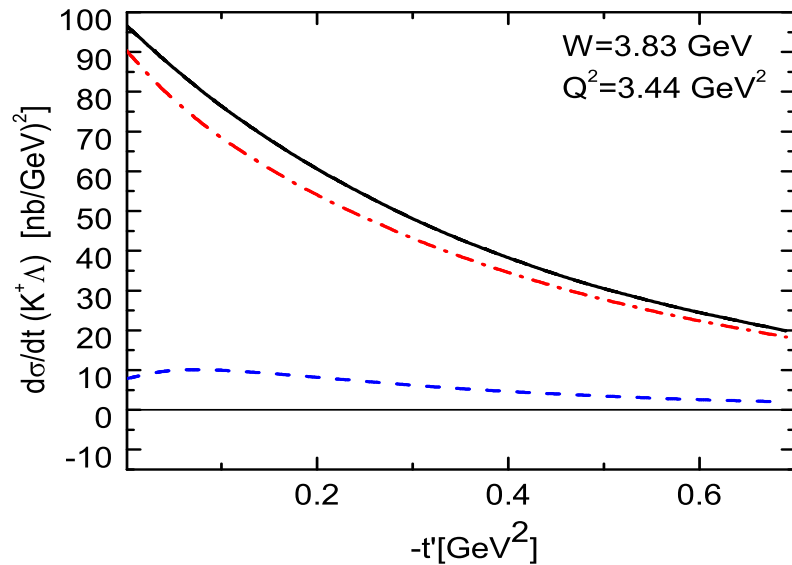
$$H_T(p \rightarrow \Sigma^0) \sim -\frac{1}{\sqrt{2}}[H_T^d]$$

$$H_T(p \rightarrow \Sigma^+) \sim -[H_T^d]$$

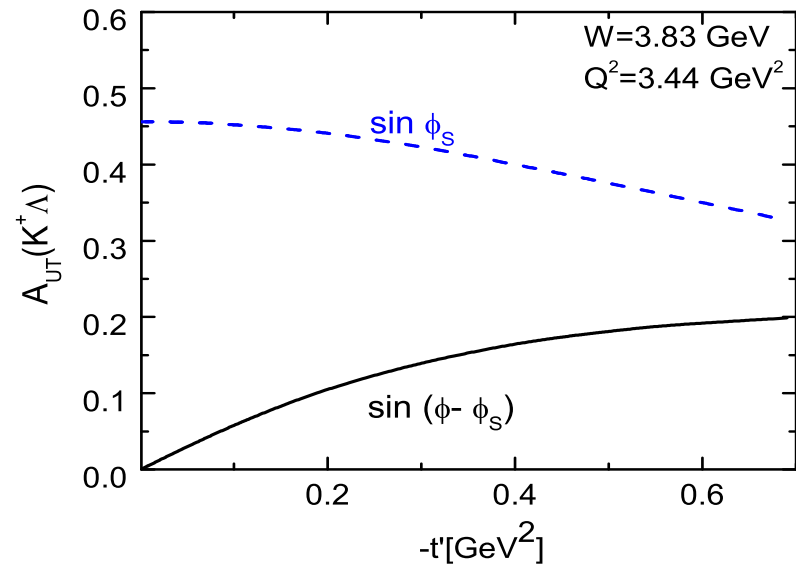
\bar{E}_T effects. $H_T \rightarrow \bar{E}_T$

In all processes of Kaons production we have and large \bar{E}_T contribution.

$\gamma p \rightarrow K^+ \Lambda$ reaction.



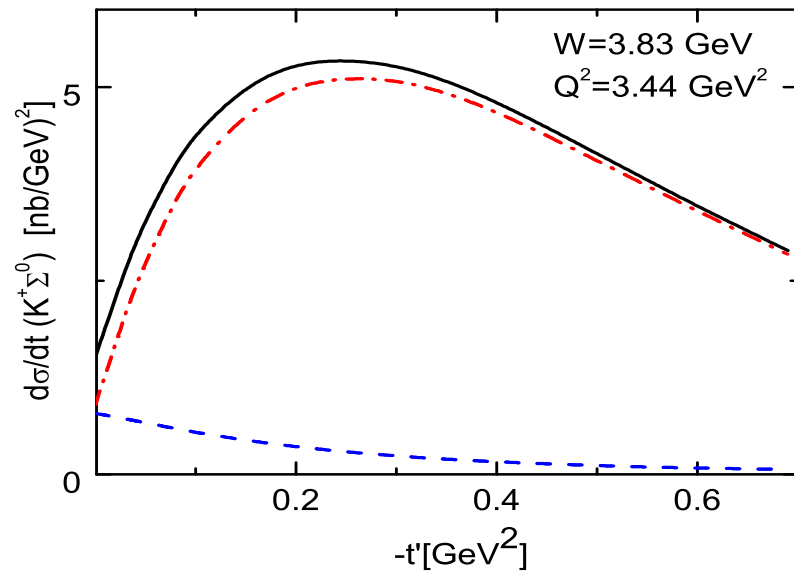
Cross section of $K^+ \Lambda$ production.
Dashed line- σ_L , dashed-dotted - σ_T



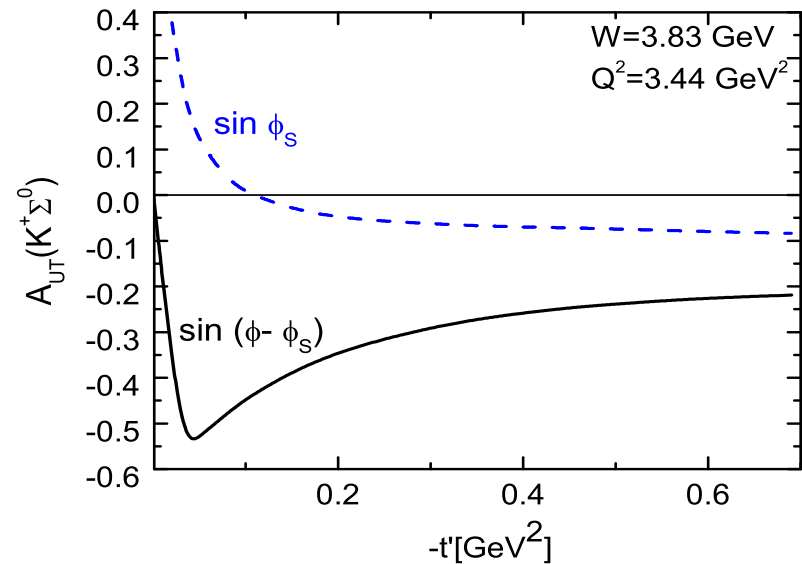
$\sin(\phi - \phi_s)$ and $\sin(\phi_s)$ moments of transverse target asymmetry of $K^+ \Lambda$ production.

Pole contribution is essential here- H_T is large no dip in σ_T . σ_T - is large with respect to the leading twist σ_L . Twist-3 transversity effects predominate at low Q^2 .

$\gamma p \rightarrow K^+ \Sigma^0$ reaction.



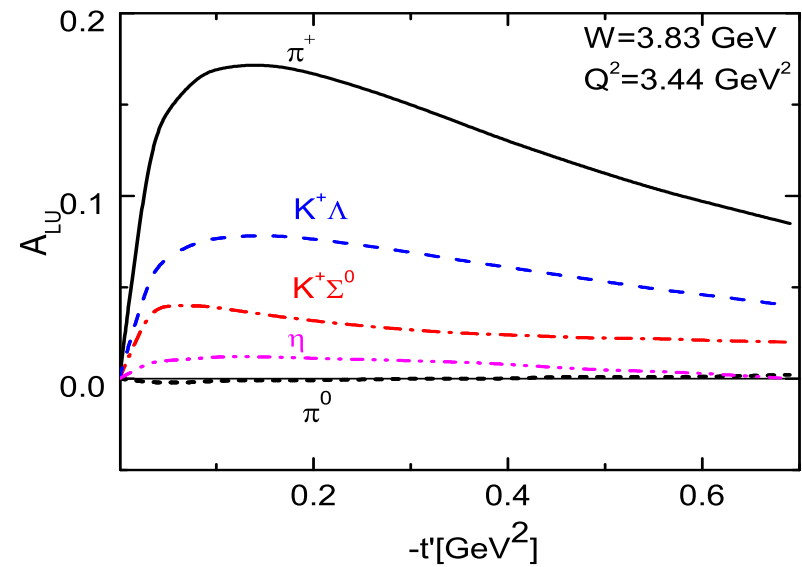
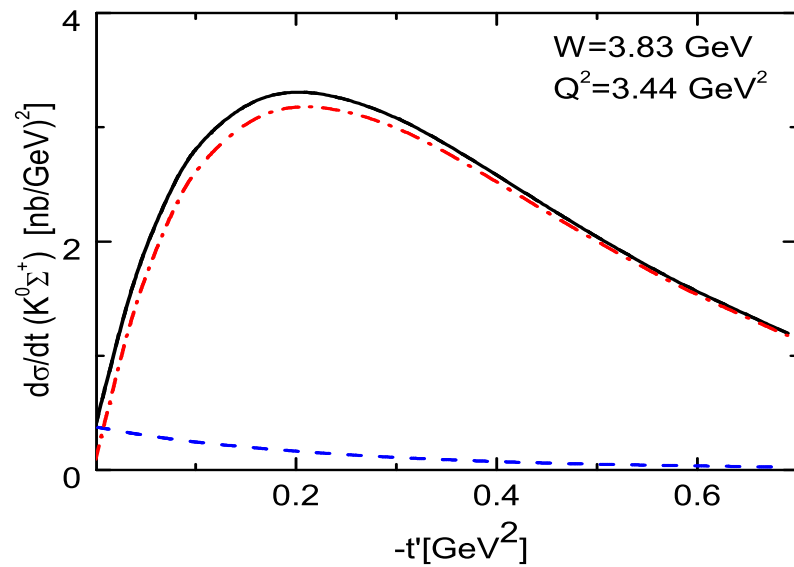
Cross section of $K^+ \Sigma^0$ production.



$\sin(\phi - \phi_s)$ and $\sin(\phi_s)$ moments of transverse target asymmetry of $K^+ \Sigma^0$ production.

Pole contribution is much smaller here (small $K^+ p \Sigma^0$ coupling constant) - dip in σ_T . σ_T is large with respect to σ_L - E_T effects. We predict large transversity effects in $K^0 \Sigma^+$ channel.

$\gamma p \rightarrow K^0 \Sigma^+$ reaction and A_{LU} asymmetry for meson channels.



Cross section of $K^0 \Sigma^+$ production.
 No pole contribution -sizeable dip in σ_T observed. σ_T is large- E_T effects with respect to σ_L . Large transversity effects in this reaction too.

The beam spin asymmetry for various pseudoscalar-meson channels

Conclusion

- Polarized GPDs are essential pseudoscalar K mesons production.
- GPDs are calculated using PDF on the bases of DD representation.
- At low Q^2 H_T effects are mostly essential in $K^+\Lambda$ channel.
- E_T effects are predominated in $K^+\Sigma^0$ and $K^0\Sigma^+$ channels . Dip in cross section at small momentum transfer is predicted.
- Moments of A_{UT} asymmetry are expected to be not small.
- Transversity H_T and E_T contributions are twist-3 effects. They decrease with Q^2 growing. At high Q^2 the leading twist σ_L will predominate.
- Future JLAB12 and COMPASS results will should throw the light on importance of transversity effects in pseudoscalar mesons production at low Q^2 .

Thank You!