Hydrodynamics of quantum critical points

Pisa University September 24, 2013

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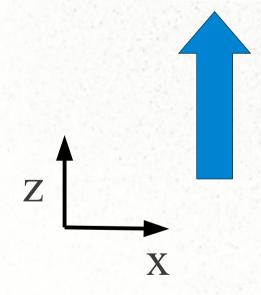
Bom Soo Kim, Yaron Oz

Quantum Critical Points

Example: Ising model

One-dimensional spin chain

$$H_I = -J \sum_{i} \hat{\sigma}_i^z \hat{\sigma}_{i+1}^z$$









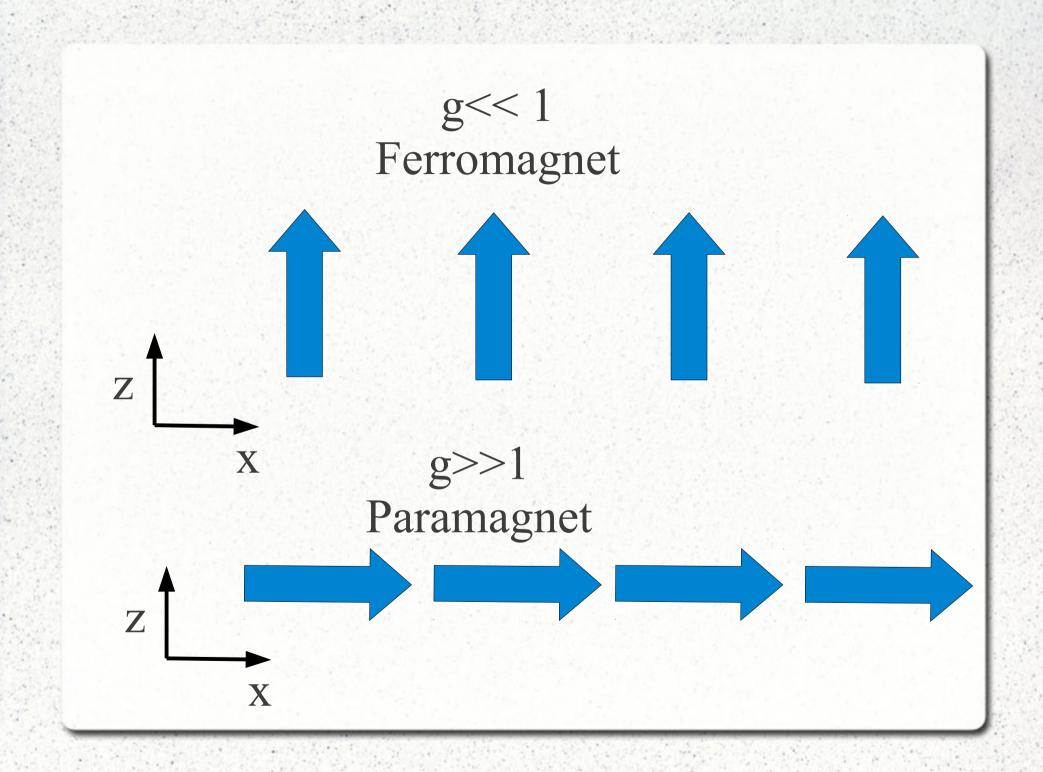
Ferromagnet

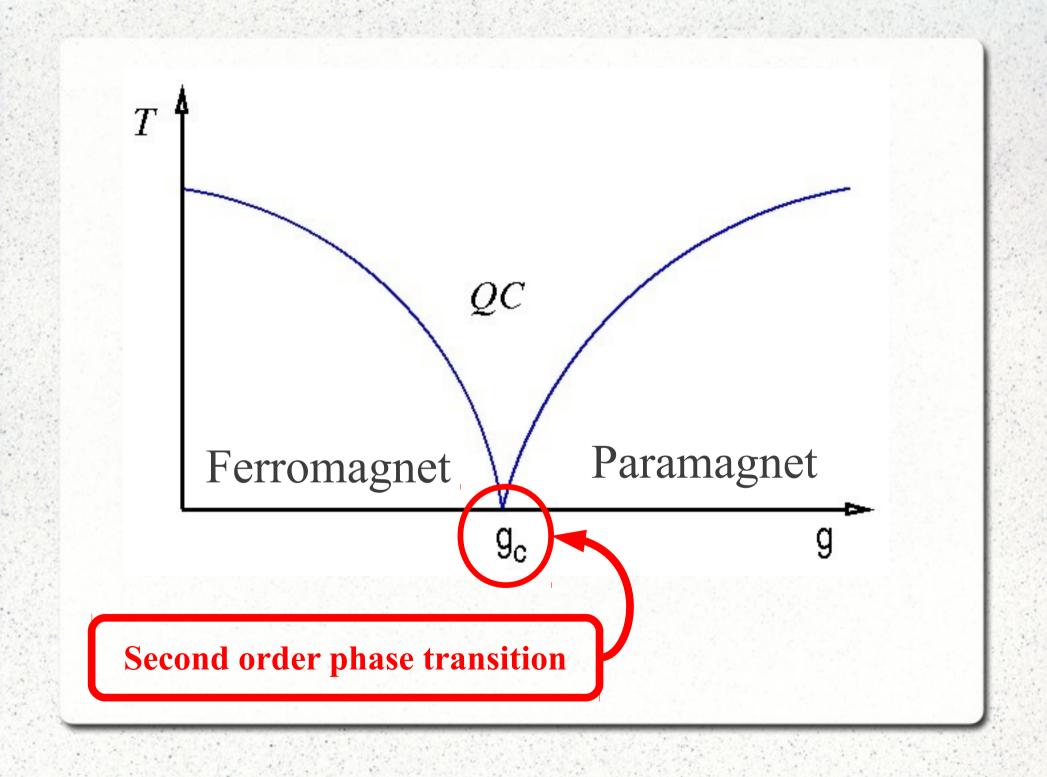
Example: Ising model

One-dimensional spin chain

$$H_I = -J \sum_{i} \left(g \hat{\sigma}_i^x \right) + \hat{\sigma}_i^z \hat{\sigma}_{i+1}^z \right)$$

Magnetic field





Energy gap:

$$\Delta \sim J|g - g_c|^{z\nu}$$

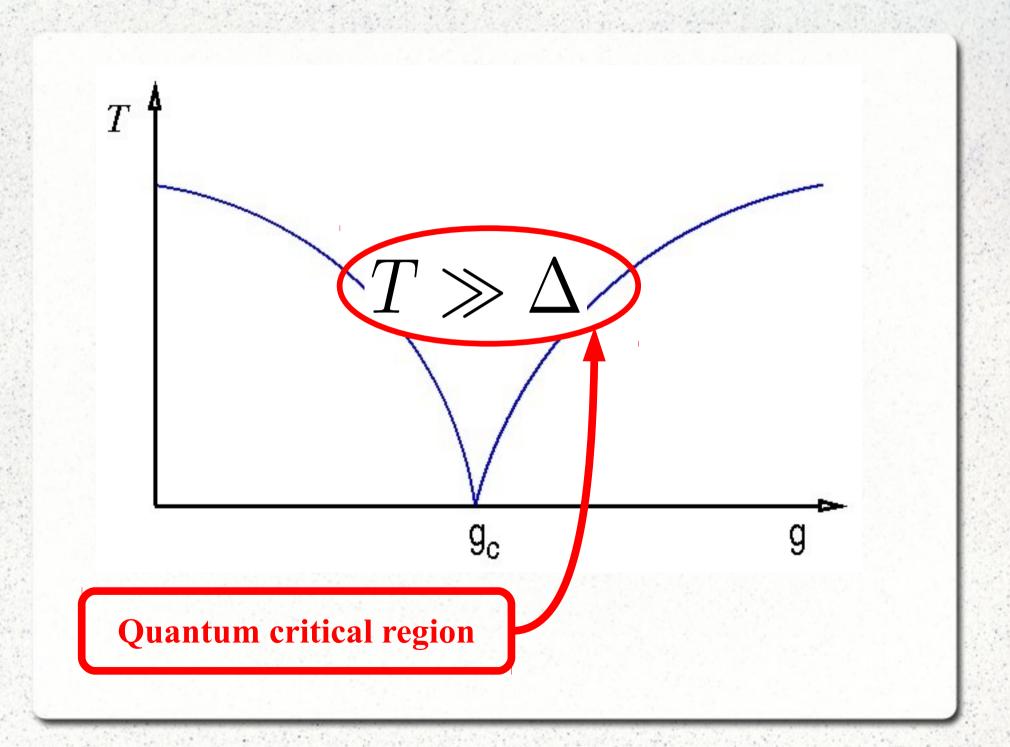
Correlation length:

$$\xi^{-1} \sim \Lambda |g - g_c|^{\nu}$$

Critical point:

$$\Delta \to 0 \quad \xi \to \infty$$

(Ising model: z = v = 1)



Other Examples

- Other spin chain models
- O(N) (quantum rotor) models
- Hubbard models
- Fermi liquids
- Bose-Einstein condensates

Experimental relevance

- Insulators with magnetic properties
- Strange metals:
 - Heavy fermion compounds
 - Materials with high Tc superconductivity

Experiments can be made only at finite temperature

What is the signature of the quantum critical point?

Equilibration/mean free time

Quantum critical:

$$au \sim \hbar/k_B T$$

Non-critical:

$$au_{.} \sim (\hbar/k_B T) e^{|\Delta|/k_B T}$$

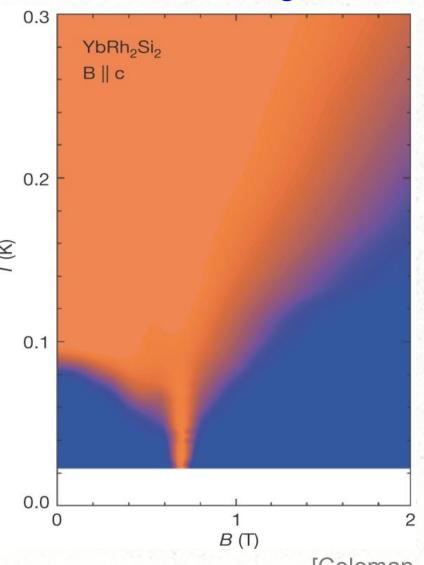
Large separation of time scales: classical relaxation (collisions of quasiparticles)

Equilibration/mean free time

Resistivity in critical region:

$$\rho_{DC} = \frac{m_*}{4\pi e^2 n} \tau^{-1} \sim \frac{m_*}{4\pi e^2 n} \frac{k_B T}{\hbar}$$





Orange:
Resistivity Linear
in Temperature

[Coleman, Schofield Nature 433, 226]

Quantum Critical Points

- Continuum limit: quantum field theory
- Scaling symmetry at the critical point:

$$t \to \lambda^z t, \quad x^i \to \lambda x^i$$

▶ Free field example: "Lifshitz theory"

$$S = \int dt \, d^dx \, \left[\frac{1}{2} (\partial_t \phi)^2 - \frac{\kappa}{2z} ((\partial_i^2)^{z/2} \phi)^2 \right]$$

Connections to High Energy Physics

- There are now many examples of gravity duals with scale invariance
- Canonical example: conformal field theories i.e. relativistic quantum critical points
- Time scales in strongly coupled gauge theories with gravity duals

$$\tau \sim \frac{\eta}{Ts} \sim \hbar/k_B T$$

Hydrodynamic description very successful!

From gauge/gravity duality models, we expect a universal hydrodynamic description of scale-invariant theories

This was already argued by condensed matter physicists, but not developed in the same way as for relativistic theories

[Sachdev & Ye]

We initiate the formulation of hydrodynamics of quantum critical points at finite temperature

Hydrodynamics in QCP

Graphene: [Fritz, Schmalian, Muller, Sachdev]

z=1 (relativistic fermions)

Hydrodynamic Drude model

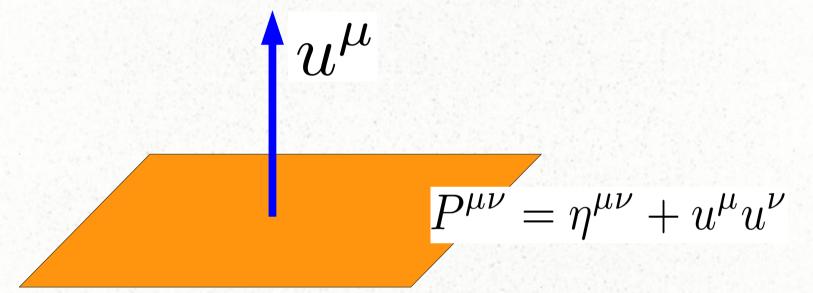
Fermions at unitarity: [Cao et al.; Son & Wingate]

z=2 (non-relativistic conformal invariance)

Elliptic flow with low viscosity

Hydrodynamics

Time-like Killing vector: defines rest frame of the fluid



$$\eta_{\mu\nu}u^{\mu}u^{\nu} = -1$$
 $u^{\mu} = (1, \beta^i)/\sqrt{1 - \beta^2}$
 $\beta^i = \frac{v^i}{c}$

Time derivatives:

$$\partial_t \phi = u^\mu \partial_\mu \phi$$

Space derivatives:

$$\nabla^2 \phi = P^{\mu\nu} \partial_\mu \partial_\nu \phi$$

Example: z=2 Lifshitz scalar

$$\mathcal{L} = \frac{1}{2} (u^{\mu} \partial_{\mu} \phi)^2 - \frac{\kappa}{4} (P^{\mu\nu} \partial_{\mu} \partial_{\nu} \phi)^2$$

Symmetry generators

Time and space translations:

$$P^{\parallel} = u^{\mu} \partial_{\mu}, \quad P^{\perp}_{\mu} = P_{\mu}^{\ \nu} \partial_{\nu}$$

Anisotropic dilatations:

$$D = zx^{\mu}u_{\mu}P^{\parallel} - x^{\mu}P_{\mu}^{\perp}$$

$$[D, P^{\parallel}] = zP^{\parallel} , \quad [D, P_{\mu}^{\perp}] = P_{\mu}^{\perp}$$

Ward identities

Conservation equations: $\partial_{\mu}T^{\mu\nu}=0$

"Trace" of energy-momentum tensor:

$$zT_{\mu\nu}u^{\mu}u^{\nu} - T_{\mu\nu}P^{\mu\nu} = 0$$

Lorentz symmetry is broken:

$$T^{\mu\nu} \neq T^{\nu\mu}$$

Constitutive relations

Ideal energy-momentum tensor:

$$T^{\mu\nu} = (\varepsilon + p)u^{\mu}u^{\nu} + p\eta^{\mu\nu}$$

Scale symmetry and equation of state:

$$zT_{\mu\nu}u^{\mu}u^{\nu} - T_{\mu\nu}P^{\mu\nu} = 0$$

$$z\varepsilon = dp$$

Temperature dependence

Scale symmetry:
$$z\varepsilon=dp$$

Thermodynamic relations:

$$\varepsilon + p = Ts$$

$$s = \frac{\partial p}{\partial T}$$

Temperature dependence:

$$\varepsilon \sim p \sim T^{\frac{z+d}{z}}$$

Constitutive relations

Energy-momentum tensor:

$$T^{\mu\nu} = (\varepsilon + p)u^{\mu}u^{\nu} + p\eta^{\mu\nu} + \eta^{\mu\nu} + \eta^{\mu\nu}$$

$$T^{\mu\nu} \neq T^{\nu\mu}$$

Landau frame condition:

$$T^{\mu\nu}u_{\nu} = -\varepsilon u^{\mu}$$

Second law of thermodynamics

$$\partial_t S = \int d^d x \, \partial_t s \ge 0$$

Local form:
$$\partial_{\mu}s^{\mu} \geq 0$$

$$\partial_{\mu}T^{\mu\nu}u_{\nu}=0$$

Entropy current: $s^{\mu} = su^{\mu} + \cdots$

Symmetric terms:

$$\pi_S^{(\mu\nu)} = -\eta P^{\mu\alpha} P^{\nu\beta} \Delta_{\alpha\beta} - \frac{\zeta}{d} P^{\mu\nu} \partial_{\alpha} u^{\alpha}$$

$$\Delta_{\alpha\beta} = 2\partial_{(\alpha}u_{\beta)} - \frac{2}{d}P_{\alpha\beta}(\partial_{\sigma}u^{\sigma})$$

$$\partial_{\mu}s^{\mu} \geq 0 \implies \eta \geq 0 \quad \zeta \geq 0$$

Anti-symmetric terms: $\partial_{\mu}s^{\mu}\geq 0$

$$\partial_{\mu}s^{\mu} \geq 0$$

$$\pi_A^{[\mu\nu]} = -\alpha^{\mu\nu\alpha\beta} (\partial_{[\alpha} u_{\beta]} - u_{[\alpha} u^{\rho} \partial_{\rho} u_{\beta]})$$

$$\tau_{\mu\nu}\alpha^{\mu\nu\sigma\rho}\tau_{\sigma\rho} \ge 0$$

Rotational invariance:
$$\alpha^{0i0j} = \alpha \delta^{ij} \geq 0$$

Temperature dependence fixed by scaling:

$$\eta \sim \zeta \sim \alpha \sim T^{\frac{d}{z}}$$

New transport coefficient:

- Dissipation due to non-inertial motion of the fluid
- Distinguishes Lifshitz from relativistic theories

Application: Drude model of strange metal

Non-relativistic limit

$$c \to \infty$$

$$\varepsilon = \rho c^2 - \rho \frac{v^2}{2} + U$$

Ideal equations of motion: zU=dp

$$zU = dp$$

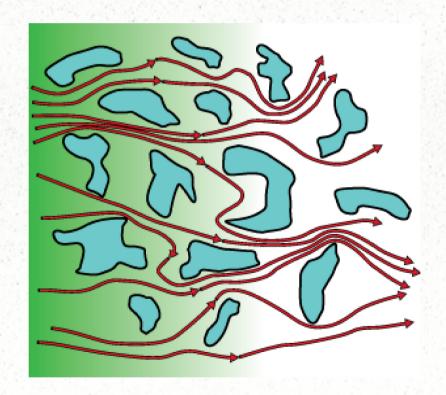
$$\partial_t \rho + \partial_i (\rho v^i) = 0$$

$$\partial_t U + \partial_i (U v^i) + p \partial_i v^i = 0$$

$$\partial_t (\rho v^i) + \partial_i (\rho v^j v^i) + \partial^i p = 0$$

Drude model: electron fluid moving though medium

- Force by external electric field
- Drag force



Ideal equations of motion:

$$\partial_t \rho + \partial_i (\rho v^i) = 0$$

$$\partial_t U + \partial_i \left(U v^i \right) + p \partial_i v^i = \lambda \rho v^2$$

$$\partial_t(\rho v^i) + \partial_j(\rho v^j v^i) + \partial^i p$$

$$= \rho E^i - \lambda \rho v^i$$

Scaling dimensions:

$$[T] = z$$

$$[v^i] = z - 1$$

$$[p] = [U] = z + d$$

$$[\rho] = d + 2 - z$$

$$[\lambda] = z \longrightarrow \lambda \sim k_B T/\hbar$$

Conductivity:

$$J^i = \rho v^i \simeq \frac{\rho}{\lambda} E^i$$

Resistivity linear in T independent of z and d

$$\partial_t U + \partial_i (Uv^i) + p\partial_i v^i$$

$$= \frac{\eta}{2} \sigma^{ij} \sigma_{ij} + \frac{\zeta}{d} (\partial_i v^i)^2 + \frac{\alpha}{2} (V_A^i)^2$$

Shear:
$$\sigma_{ij} = 2\partial_{(i}v_{j)} - (2/d)\delta_{ij}\partial_k v^k$$

$$V_A^i = 2D_t v^i + \omega^{ij} v_j$$

Linear acceleration:

$$D_t \equiv \partial_t + v^i \partial_i$$

Coriolis acceleration:

$$\omega_{ij} = 2\partial_{[i}v_{j]}$$

Corrections to the conductivity

$$\boldsymbol{\sigma}_{xx}(E_x) = \frac{\rho}{\lambda} \left[1 + \frac{1}{\rho \lambda E_x} \left(\eta \partial_y^2 E_x + \left(\frac{\alpha}{6\lambda^2} \partial_y^2 E_x^3 \right) \right) \right]$$

Estimate for $E_x = E_0 y/L$

$$\frac{\boldsymbol{\sigma}_{xx} - \boldsymbol{\sigma}_{xx}^{0}}{\boldsymbol{\sigma}_{xx}^{0}} \sim 10^{-11} \left(\frac{m_e}{m_*}\right)^2 \left(\frac{T}{K}\right)^{-3} \frac{\alpha/\rho}{\text{sec}} (\partial E_0)^2$$

$$(\partial E_0) = \frac{E_0/L}{N \, \text{m}^{-1} \, \text{C}^{-1}}$$

Future directions

- Transport coefficients beyond first order
- Additional conserved currents
- Superfluids with Lifshitz scaling
- Fluids in a curved space: Weyl anomaly?
- Anomalous currents

Grazie!