Probing NP with tau LFV and EDM measurements

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Workshop on τ -charm at high luminosity. La Biodola, 27–31/05/2013

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Still, lots of work needed!!

Standard Model

All Observed Flavour Changing Neutral Currents can be accomodated in Yukawa couplings:

$$\mathcal{L}_Y = H \bar{Q}_i Y_{ij}^d d_j + H^* \bar{Q}_i Y_{ij}^u u_j$$

Only masses and CKM mixings, $V_{
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a) what is the origin of the Yukawa structures??
 b) why is there a CP-violating phase in CKM??

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New Physics

New flavour structures generically present \Rightarrow measure of new observables provides new information on flavour origin...

SUSY Flavour (and CP) problems

Soft masses fixed by $m_{3/2}$. $O(m_{3/2})$ elements in soft matrices.

$$\Rightarrow$$
 Severe FCNC problem !!!

CP broken, we can expect all complex paramaters have

O(1) phases. \implies Too large EDMs !!!

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SM Flavour and CP

Fermion masses fixed by M_W . If O(1) elements in Yukawa matrices and O(1) phases

Impossible reproduce masses, mixings

and CP observables !!!

FLAVOURED NEW PHYSICS

2 Higgs Doublet Models

- Four possible Yukawa matrices. \Rightarrow Large FCNC.
- Discrete symmetry (type I, type II) to forbid FCNC. \Rightarrow No connection with structure of flavour matrices.
- Alignment of Yukawa matrices \Rightarrow ad hoc requirement, no connection with struct. of flavour matrices.
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Supersymmetry

• Five sfermion mass matrices and Three trilinear matrices \Rightarrow Lots of new observables to understand flavour.

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EW production: Direct limits on chargino and slepton masses \rightarrow EW production chargino/neutralino and slepton. $M\gtrsim 300$ GeV

ATLAS-CONF-2013-049, CMS PAS SUS-12-022



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LHC-Flavour fact. Feedback:

- 1. LHC measures squark, gluino, neutralino mass scale
- 2. With this information Flavour factories determine $\tan\beta$ and mixings
- 3. These parameters help reanalyze LHC data
- 4. ...

MI CONSTRAINTS

 $\tan \beta = 10, m_{\tilde{l}} = 400 \text{ GeV}, M_2 = 150 \text{ GeV}.$

δ_{13}^l	$\tau \to e\gamma$	au ightarrow e e e	$ au o e \mu \mu$
LL	0.15	_	-
RR	-	-	-
LR/RL	0.04	0.5	-
δ_{23}^l	$ au o \mu \gamma$	$ au o \mu \mu \mu$	$ au o \mu e e$
LL	0.12	-	-
RR	_	-	-
LR/RL	0.03	-	0.5

$$\begin{split} &\mathsf{BR}(\tau \to e \,\gamma) < 1.2 \ \times \ 10^{-8}, \, \mathsf{BR}(\tau \to e \,\mu \,\mu) \ < 2. \ \times \ 10^{-7}, \\ &\mathsf{BR}(\tau \to \mu \,\gamma) \ < 4.8 \ \times \ 10^{-8}, \, \mathsf{BR}(\tau \to \mu \,\mu \,\mu) \ < 2 \ \times \ 10^{-7} \end{split}$$

τ -CHARM PHYSICS

- 3rd generation (Yukawa) couplings larger, \Rightarrow sizeable effects possible
- Importance of polarization \rightarrow Decay of a polarized τ . Access to different (LFV) couplings.
 - \Rightarrow LFV decays: $\tau \rightarrow \mu \gamma$, $\tau \rightarrow lll$.
 - \Rightarrow CP violation. τ Electric Dipole Moment (and also AMM)

FLAVOUR SYMMETRIES IN SUSY

- Very different elements in Yukawas: $y_t \simeq 1$, $y_u \simeq 10^{-5}$
- Expect couplings in a "fundamental" theory $\mathcal{O}(1)$
- Small couplings generated at higher order or function of small vevs.
- Froggatt-Nielsen mechanism and flavour symmetry to understand small Yukawa elements. Example: $U(1)_{fl}$



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Symmetry Breaking.

Unbroken symmetry applies
 both to fermion and sfermions.

Diagonal soft masses allowed

by symmetry.

• Nonuniversality in soft terms proportional to symm. breaking.

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We can <u>relate</u> the structure in Yukawa matrices to

the nonuniversality in Soft Breaking masses !!!

Symmetric texture

• Non-Abelian flavour symmetries.

$$Y^{d,e} = \begin{pmatrix} 0 & 1.5 \varepsilon^3 & 0.4 \varepsilon^3 \\ 1.5 \varepsilon^3 & \Sigma \varepsilon^2 & 1.3 \Sigma \varepsilon^2 \\ 0.4 \varepsilon^3 & 1.3 \Sigma \varepsilon^2 & 1 \end{pmatrix} y_b$$

• Universal sfermion masses in in unbroken limit:

in unbroken limit:

$$\mathcal{L}_{m^2} = m_0^2 \Phi^{\dagger} \Phi = m_0^2 (\phi_1 \phi_2 \phi_3)^* \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}$$
• After symmetry breaking:

$$M_{\tilde{D}_R,\tilde{E}_L}^2 \simeq \begin{pmatrix} 1+\bar{\varepsilon}^3 & \bar{\varepsilon}^3 & 0\\ \bar{\varepsilon}^3 & 1+\bar{\varepsilon}^2 & \bar{\varepsilon}^2\\ 0 & \bar{\varepsilon}^2 & 1+\bar{\varepsilon} \end{pmatrix} m_0^2$$

Asymmetric texture

• Abelian flavour symmetries.

$$Y^{d,e} = \begin{pmatrix} \varepsilon^4 & \varepsilon^3 & \varepsilon^3 \\ \varepsilon^3 & \varepsilon^2 & \varepsilon^2 \\ \varepsilon & 1 & 1 \end{pmatrix} y_b$$

• In principle nonuniversal masses in unbroken symmetry:

$$\mathcal{L}_{m^2} = m_1^2 \ \phi_1^* \phi_1 + m_2^2 \ \phi_2^* \phi_2 + m_3^2 \ \phi_3^* \phi_3$$

• After symmetry breaking:

$$M_{\tilde{D}_R,\tilde{E}_L}^2 \simeq \begin{pmatrix} 1 & \bar{\varepsilon} & \bar{\varepsilon} \\ \bar{\varepsilon} & c & b \\ \bar{\varepsilon} & b & a \end{pmatrix} m_0^2$$

LEPTON FLAVOUR VIOLATION

Off-diagonal entries in slepton masses generate *LFV* processes:

$$\mathsf{BR}(l_i \to l_j \gamma) \simeq \frac{3\pi \alpha_2^3}{G_F^2} \left| \frac{(\delta_{\mathbf{L}}^e)_{ij}}{m_{\tilde{l}}^2} \frac{\mu M_2 \tan \beta}{(M_2^2 - \mu^2)} F_{2L}(a_2, b) \right|^2$$

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Brown(light):Present (fut.) $\mu \rightarrow e\gamma$, Green (light): Present (fut.) $\tau \rightarrow \mu\gamma$.

FLAVOURED EDMS

– SUSY EDMs in presence of flavour-blind phases ($\varphi_{\mu}, \varphi_{A}$) directly proportional to lepton masses,

$$d_{\chi^+}^l \simeq \frac{-\alpha \ e \ m_l \tan \beta}{4\pi \sin^2 \theta_W} \ \frac{\text{Im}[M_2 \ \mu]}{m_{\tilde{\nu}_e}^2} \ \frac{A(r_1) - A(r_2)}{m_{\chi_1^+}^2 - m_{\chi_2^+}^2}$$

- Still, if $\varphi_{\mu}, \varphi_{A} = 0$, contributions to EDMs from offdiagonal elements in sfermion masses:

 $d_e \propto (\delta^e_{LL})_{1i} (\delta^e_{LR})_{i1} f_1 + (\delta^e_{LR})_{1i} (\delta^e_{RR})_{i1} f_2 + (\delta^e_{LL})_{1i} (\delta^e_{LR})_{ij} (\delta^e_{RR})_{j1} f_3$

- In 2HdM d_l proportional to three masses and mixings

 $d_l \propto \; m_l \; m_i^2 \; |K_{li}|^2 \; f$

Three leptonic EDMs must be measured independently to discriminate the source!!!



From light to dark: $d_e = 1 \times 10^{-30}, 5 \times 10^{-29}, 1 \times 10^{-29}$ e cm

DECAY OF POLARIZED τ

- τ polarized, $P_{z,x}$, in the presence of initial beam polarization, P_e :

$$P_{z'}^{(-)}(\theta, P_e) = -\frac{8G_F s}{4\sqrt{2}\pi\alpha} \operatorname{Re}\left\{\frac{g_V^l - Q_b g_V^c \Upsilon(s)}{1 + Q_c^2 \Upsilon(s)}\right\} \times \left(\frac{g_A^\tau}{p^0} \frac{|\vec{p}|}{p^0} + 2 g_A^e \frac{\cos\theta}{1 + \cos^2\theta}\right) + P_e \frac{\cos\theta}{1 + \cos^2\theta}$$

– Large τ polarization can be used to measure (or constrain) different MI through angular distrib.

$$A_{L}^{ij} = \frac{\alpha_2}{4\pi} \frac{\delta_{\text{LL}}^{ij}}{m_{\tilde{l}}^2} \left(A(x_a) + B(x_a) \tan \beta \right) + \frac{\alpha_1}{4\pi} \frac{\delta_{\text{RL}}^{ij}}{m_{\tilde{l}}^2} \left(\frac{M_1}{m_{l_i}} \right) C(x_a),$$

$$A_{R}^{ij} = \frac{\alpha_1}{4\pi} \left[\frac{\delta_{\text{RR}}^{ij}}{m_{\tilde{l}}^2} \left(A'(x_a) + \tan \beta B'(x_a) \right) + \frac{\delta_{\text{LR}}^{ij}}{m_{\tilde{l}}^2} \left(\frac{M_1}{m_{l_i}} \right) C'(x_a) \right],$$

•

Conclusions

au-Charm factory necessary to solve the flavour problem.

• New colorLR flavour structures provide valuable information on origin of flavour

- LFV processes, $\tau \to \mu \gamma$, $\tau \to lll$, constrain 3rd generation couplings.
- Ratios of leptonic EDMs depend on flavour structures and new physics model.
- Beam polarization: new tool to be explored.
- LFV and EDMs can explore large areas of flavour MSSM.