

Soft interactions in Herwig++

Manuel Bähr

University of Karlsruhe

First international workshop on
Multiple Partonic Interactions at the LHC,
Perugia, October 2008

Herwig++

- Main focus of this talk: Underlying event and MPI
- However, also other exciting new features will be available with the imminent next release (Herwig++ 2.3)
 - single W and Z , Higgs in gluon fusion and WH , ZH associated Higgs production at NLO QCD in the POWHEG scheme
 - 3 body BSM decays and off-shell effects
 - Sophisticated baryon decays to match the degree of precision in the meson decays
 - A lot of new things in the MPI model ...

Herwig++

- Main focus of this talk: Underlying event and MPI
- However, also other exciting new features will be available with the imminent next release (Herwig++ 2.3)
 - single W and Z , Higgs in gluon fusion and WH , ZH associated Higgs production at NLO QCD in the POWHEG scheme
 - 3 body BSM decays and off-shell effects
 - Sophisticated baryon decays to match the degree of precision in the meson decays
 - A lot of new things in the MPI model ...

Herwig++

- Main focus of this talk: Underlying event and MPI
- However, also other exciting new features will be available with the imminent next release (Herwig++ 2.3)
 - single W and Z , Higgs in gluon fusion and WH , ZH associated Higgs production at NLO QCD in the POWHEG scheme
 - 3 body BSM decays and off-shell effects
 - Sophisticated baryon decays to match the degree of precision in the meson decays
 - A lot of new things in the MPI model ...

Herwig++

- Main focus of this talk: Underlying event and MPI
- However, also other exciting new features will be available with the imminent next release (Herwig++ 2.3)
 - single W and Z , Higgs in gluon fusion and WH , ZH associated Higgs production at NLO QCD in the POWHEG scheme
 - 3 body BSM decays and off-shell effects
 - Sophisticated baryon decays to match the degree of precision in the meson decays
 - A lot of new things in the MPI model ...

Herwig++

- Main focus of this talk: Underlying event and MPI
- However, also other exciting new features will be available with the imminent next release (Herwig++ 2.3)
 - single W and Z , Higgs in gluon fusion and WH , ZH associated Higgs production at NLO QCD in the POWHEG scheme
 - 3 body BSM decays and off-shell effects
 - Sophisticated baryon decays to match the degree of precision in the meson decays
 - A lot of new things in the MPI model ...

Herwig++

- Main focus of this talk: Underlying event and MPI
- However, also other exciting new features will be available with the imminent next release (Herwig++ 2.3)
 - single W and Z , Higgs in gluon fusion and WH , ZH associated Higgs production at NLO QCD in the POWHEG scheme
 - 3 body BSM decays and off-shell effects
 - Sophisticated baryon decays to match the degree of precision in the meson decays
 - A lot of new things in the MPI model ...

Status semi-hard MPI

- Fully working model included from Herwig++ 2.1 onwards. It allows for the simulation of semi-hard ($p_t \geq p_t^{\min}$) multiple partonic interactions (MPI) to describe the underlying event (UE). Same functionality and physics than JIMMY.
- Detailed description including tuning results:
[arXiv:0803.3633, MB, S. Gieseke and M. H. Seymour]
- Main focus of this talk: Extension to soft ($p_t < p_t^{\min}$) scatters partly written up in [arXiv:0806.2949, MB, J. M. Butterworth and M. H. Seymour] ; all in [arXiv:??, MB, PhD thesis]

Outline

- 1 Introduction
 - Eikonal model basics
 - Soft part of the eikonal
- 2 Analytic constraints
 - Simple model
 - Hot-Spot model
- 3 Final states
 - Tevatron Run1
- 4 Conclusions

Outline

- 1 Introduction
 - Eikonal model basics
 - Soft part of the eikonal
- 2 Analytic constraints
 - Simple model
 - Hot-Spot model
- 3 Final states
 - Tevatron Run1
- 4 Conclusions

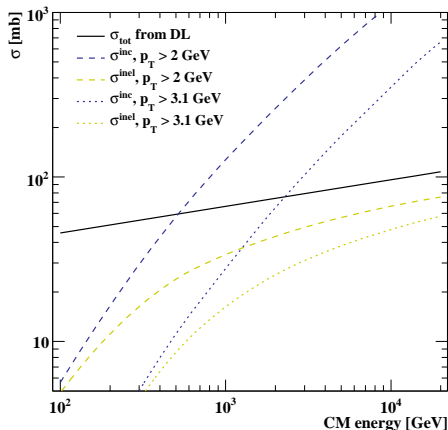
Model basics

based on [Butterworth, Forshaw, Seymour '96]

- Starting point: The inclusive cross section $pp \rightarrow jj$ with $p_t > p_t^{\min}$ ($\sigma_{\text{hard}}^{\text{inc}}$) may exceed σ_{tot}

$$\sigma_{\text{hard}}^{\text{inc}} = \int dx_1 dx_2 dp_t^2 f(x_1) f(x_2) \frac{d\hat{\sigma}}{dp_t^2}$$
- Source: Proliferation of low x partons, which increases the probability of more than one partonic collision: $\sigma^{\text{inc}} = \bar{n} \cdot \sigma_{\text{inel}}$
- Unitarization of $\sigma_{\text{hard}}^{\text{inc}}$:

$$a(\mathbf{b}, s) = \frac{1}{2i} \left[e^{-\chi(\mathbf{b}, s)} - 1 \right] \rightarrow \sigma_{\text{inel}} = \int d^2\mathbf{b} \left[1 - e^{-2\chi(\mathbf{b}, s)} \right]$$



Semi-hard part of the eikonal

- Impact parameter dependence is capable of describing the pedestal effect
- Distribution of additional scatters can be derived in field theory: $P_m(\mathbf{b}, s) = \frac{\bar{n}(\mathbf{b}, s)^m}{m!} e^{-\bar{n}(\mathbf{b}, s)}$
- Average multiplicity at fixed impact parameter

$$\bar{n}(\mathbf{b}, s) = \sum_{ij} \frac{1}{1 + \delta_{ij}} \int d^2\mathbf{b}_2 dx_1 dx_2 \int_{p_t^{\min}} dp_t^2 \frac{d\hat{\sigma}_{ij}}{dp_t^2} D_{i/A}(x_1, p_t^2, |\mathbf{b}_2|) D_{j/B}(x_2, p_t^2, |\mathbf{b} - \mathbf{b}_2|)$$

Semi-hard part of the eikonal

- Impact parameter dependence is capable of describing the pedestal effect
- Distribution of additional scatters can be derived in field theory: $P_m(\mathbf{b}, s) = \frac{\bar{n}(\mathbf{b}, s)^m}{m!} e^{-\bar{n}(\mathbf{b}, s)}$
- Average multiplicity at fixed impact parameter

$$\bar{n}(\mathbf{b}, s) = \sum_{ij} \frac{1}{1 + \delta_{ij}} \int d^2\mathbf{b}_2 dx_1 dx_2 \int_{p_t^{\min}} dp_t^2$$

$$\frac{d\hat{\sigma}_{ij}}{dp_t^2} f_{i/A}(x_1, p_t^2) f_{j/B}(x_2, p_t^2) G_A(|\mathbf{b}_2|) G_B(|\mathbf{b} - \mathbf{b}_2|)$$

Semi-hard part of the eikonal

- Impact parameter dependence is capable of describing the pedestal effect
- Distribution of additional scatters can be derived in field theory: $P_m(\mathbf{b}, s) = \frac{\bar{n}(\mathbf{b}, s)^m}{m!} e^{-\bar{n}(\mathbf{b}, s)}$
- Average multiplicity at fixed impact parameter

$$\bar{n}(\mathbf{b}, s) = A(\mathbf{b}) \cdot \sigma_{\text{hard}}^{\text{inc}}(s; p_t^{\text{min}})$$

Semi-hard part of the eikonal

- Impact parameter dependence is capable of describing the pedestal effect
- Distribution of additional scatters can be derived in field theory: $P_m(\mathbf{b}, s) = \frac{\bar{n}(\mathbf{b}, s)^m}{m!} e^{-\bar{n}(\mathbf{b}, s)}$
- Average multiplicity at fixed impact parameter

$$\bar{n}(\mathbf{b}, s) = A(\mathbf{b}) \cdot \sigma_{\text{hard}}^{\text{inc}}(s; p_t^{\text{min}})$$

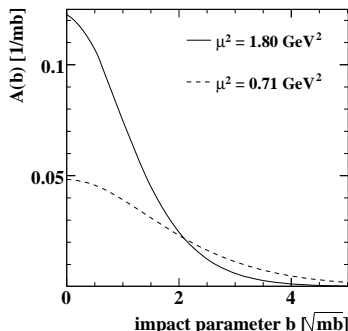
- Determination of χ : $\sigma_{\text{inel}} = \int d^2\mathbf{b} [1 - e^{-2\chi(\mathbf{b}, s)}]$

$$\sigma_{\text{inel}} = \int d^2\mathbf{b} \sum_{m=1}^{\infty} P_m(\mathbf{b}, s) = \int d^2\mathbf{b} (1 - e^{-\bar{n}(\mathbf{b}, s)})$$

$$\rightarrow \chi_{\text{QCD}}(\mathbf{b}, s) = \frac{1}{2}\bar{n}(\mathbf{b}, s) = \frac{1}{2}A(\mathbf{b})\sigma_{\text{hard}}^{\text{inc}}$$

Overlap function

- $A(b)$ is the overlap function of the two colliding particles
- Convolution of individual spatial parton distributions: $G_h(\mathbf{b})$
- Individual distributions are proportional to EM form-factors. But μ is not fixed to the radius measured there (0.71 GeV^2)!



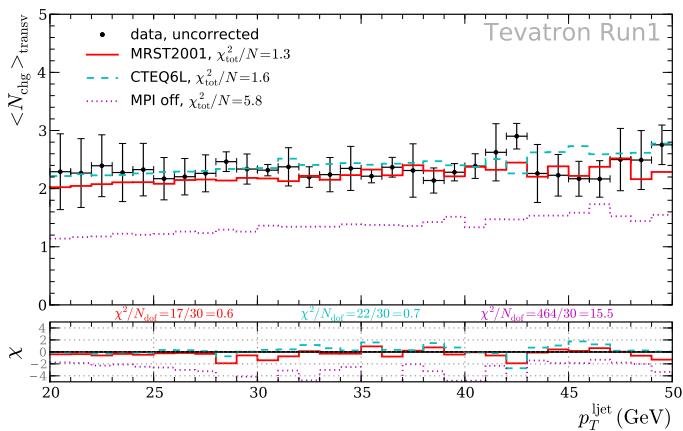
$$A(b = |\mathbf{b}|) = \int d^2\mathbf{b}' G_{h_1}(|\mathbf{b}'|) G_{h_2}(|\mathbf{b} - \mathbf{b}'|)$$

$$G_{\bar{p}}(\mathbf{b}) = G_p(\mathbf{b}) = \int \frac{d^2\mathbf{k}}{(2\pi)^2} \frac{e^{i\mathbf{k}\cdot\mathbf{b}}}{(1 + \mathbf{k}^2/\mu^2)^2}$$

→ main parameters: p_t^{\min}, μ^2

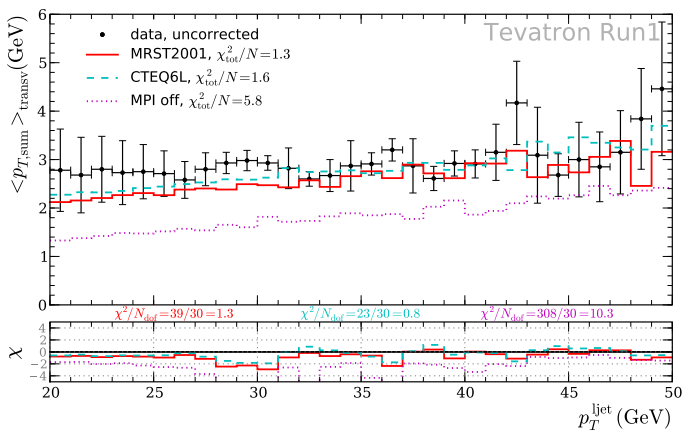
Results

- semi-hard MPI model describes the CDF Run1 UE data very well (overall $\chi^2 = 1.3$)
- However, soft scatters are intrinsically missing



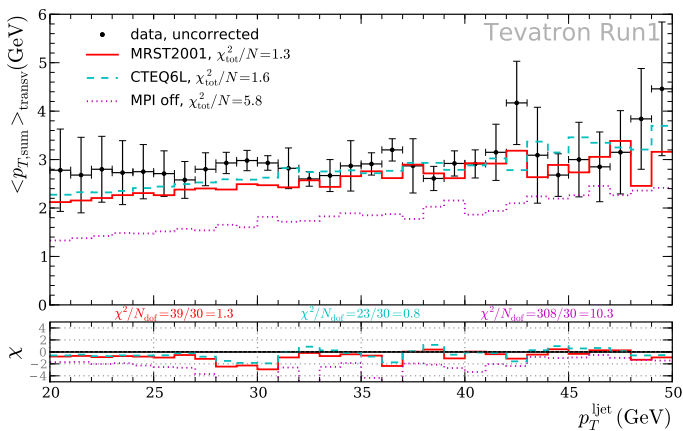
Results

- semi-hard MPI model describes the CDF Run1 UE data very well (overall $\chi^2 = 1.3$)
- However, soft scatters are intrinsically missing



Results

- semi-hard MPI model describes the CDF Run1 UE data very well (overall $\chi^2 = 1.3$)
- However, soft scatters are intrinsically missing



Outline

- 1 Introduction
 - Eikonal model basics
 - Soft part of the eikonal
- 2 Analytic constraints
 - Simple model
 - Hot-Spot model
- 3 Final states
 - Tevatron Run1
- 4 Conclusions

$\chi_{\text{soft}}(\mathbf{b}, s)$

- When we use $\chi_{\text{QCD}}(\mathbf{b}, s)$, we mean

$$\chi_{\text{tot}}(\mathbf{b}, s) = \chi_{\text{QCD}}(\mathbf{b}, s) + \chi_{\text{soft}}(\mathbf{b}, s)$$

- Assume similar structure for $\chi_{\text{soft}}(\mathbf{b}, s)$ and $\chi_{\text{QCD}}(\mathbf{b}, s)$:

$$\chi_{\text{soft}}(\mathbf{b}, s) = \frac{1}{2} A_{\text{soft}}(\mathbf{b}) \sigma_{\text{soft}}^{\text{inc}},$$

where $\sigma_{\text{soft}}^{\text{inc}}$ is a free parameter of our model.

- Choose simplest Ansatz for A_{soft} , $A_{\text{soft}}(b) \equiv A(b; \mu)$, which then leads to:

$$\chi_{\text{tot}}(\mathbf{b}, s) = \frac{A(b, \mu)}{2} (\sigma_{\text{hard}}^{\text{inc}} + \sigma_{\text{soft}}^{\text{inc}})$$

Determination of $\sigma_{\text{soft}}^{\text{inc}}$

- Exploit the connection to the total cross section within the eikonal model:

$$\sigma_{\text{tot}}(s) = 2 \int d^2\mathbf{b} \left[1 - e^{-\chi_{\text{tot}}(\mathbf{b},s)} \right],$$

- Use experimental total cross section values to fix $\sigma_{\text{soft}}^{\text{inc}}$ accordingly (for fixed CM energy).
- For $\sqrt{s} > 2$ TeV: Get σ_{tot} from the Regge fit by Donnachie & Landshoff
 - ① DL '92 [arXiv:hep-ph/9209205, D & L]
 - ② DL '92 fixed to the CDF σ_{tot} measurement
 - ③ DL '04 [arXiv:hep-ph/0402081, D & L]

Determination of $\sigma_{\text{soft}}^{\text{inc}}$

- Exploit the connection to the total cross section within the eikonal model:

$$\sigma_{\text{tot}}(s) = 2 \int d^2\mathbf{b} \left[1 - e^{-A(b;\mu)/2(\sigma_{\text{hard}}^{\text{inc}} + \sigma_{\text{soft}}^{\text{inc}})} \right],$$

- Use experimental total cross section values to fix $\sigma_{\text{soft}}^{\text{inc}}$ accordingly (for fixed CM energy).
- For $\sqrt{s} > 2$ TeV: Get σ_{tot} from the Regge fit by Donnachie & Landshoff
 - ① DL '92 [arXiv:hep-ph/9209205, D & L]
 - ② DL '92 fixed to the CDF σ_{tot} measurement
 - ③ DL '04 [arXiv:hep-ph/0402081, D & L]

Determination of $\sigma_{\text{soft}}^{\text{inc}}$

- Exploit the connection to the total cross section within the eikonal model:

$$\sigma_{\text{tot}}(s) = 2 \int d^2\mathbf{b} \left[1 - e^{-A(b;\mu)/2(\sigma_{\text{hard}}^{\text{inc}} + \sigma_{\text{soft}}^{\text{inc}})} \right],$$

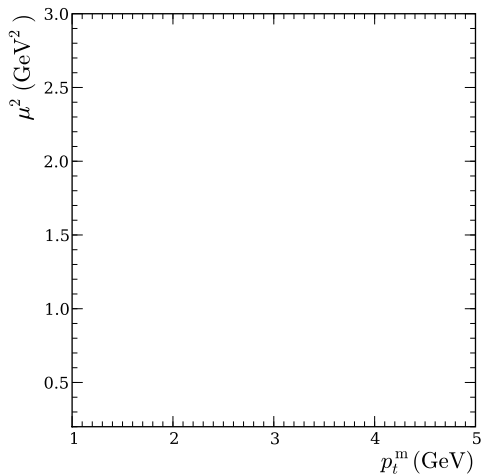
- Use experimental total cross section values to fix $\sigma_{\text{soft}}^{\text{inc}}$ accordingly (for fixed CM energy).
- For $\sqrt{s} > 2$ TeV: Get σ_{tot} from the Regge fit by Donnachie & Landshoff
 - 1 DL '92 [arXiv:hep-ph/9209205, D & L]
 - 2 DL '92 fixed to the CDF σ_{tot} measurement
 - 3 DL '04 [arXiv:hep-ph/0402081, D & L]

Outline

- 1 Introduction
 - Eikonal model basics
 - Soft part of the eikonal
- 2 Analytic constraints
 - Simple model
 - Hot-Spot model
- 3 Final states
 - Tevatron Run1
- 4 Conclusions

Tevatron parameter space

- Find constraints on the parameter choices
- Require $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb, while describing the total cross section
- Require elastic t -slope, $b_{\text{el}}(s) = \left[\frac{d}{dt} \left(\ln \frac{d\sigma_{\text{el}}}{dt} \right) \right]_{t=0}$, to be correctly described
$$b_{\text{el}}(s) = \int d^2b \frac{b^2}{\sigma_{\text{tot}}} \left[1 - e^{-\chi_{\text{tot}}(b,s)} \right]$$
- Final state tune of **semi-hard MPI** (approximate constraint for MRST2001)



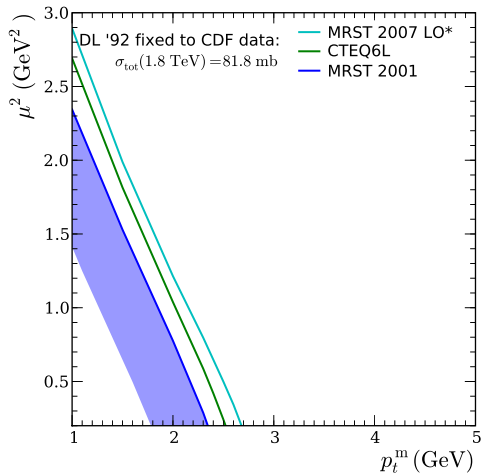
Tevatron parameter space

- Find constraints on the parameter choices
- Require $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb, while describing the total cross section

- Require elastic t -slope, $b_{\text{el}}(s) = \left[\frac{d}{dt} \left(\ln \frac{d\sigma_{\text{el}}}{dt} \right) \right]_{t=0}$, to be correctly described

$$b_{\text{el}}(s) = \int d^2b \frac{b^2}{\sigma_{\text{tot}}} \left[1 - e^{-\chi_{\text{tot}}(b,s)} \right]$$

- Final state tune of **semi-hard MPI** (approximate constraint for MRST2001)



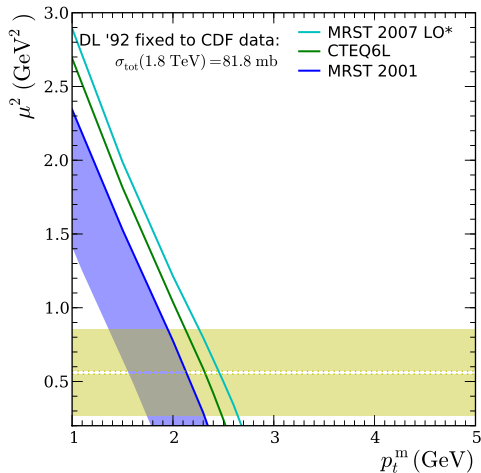
Tevatron parameter space

- Find constraints on the parameter choices
- Require $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb, while describing the total cross section

- Require elastic t -slope, $b_{\text{el}}(s) = \left[\frac{d}{dt} \left(\ln \frac{d\sigma_{\text{el}}}{dt} \right) \right]_{t=0}$, to be correctly described

$$b_{\text{el}}(s) = \int d^2\mathbf{b} \frac{b^2}{\sigma_{\text{tot}}} \left[1 - e^{-\chi_{\text{tot}}(\mathbf{b}, s)} \right]$$

- Final state tune of semi-hard MPI (approximate constraint for MRST2001)



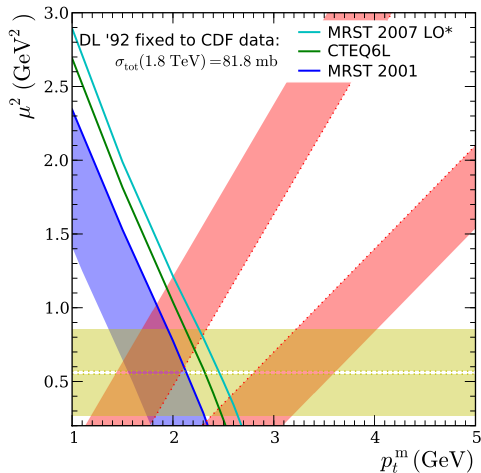
Tevatron parameter space

- Find constraints on the parameter choices
- Require $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb, while describing the total cross section

- Require elastic t -slope, $b_{\text{el}}(s) = \left[\frac{d}{dt} \left(\ln \frac{d\sigma_{\text{el}}}{dt} \right) \right]_{t=0}$, to be correctly described

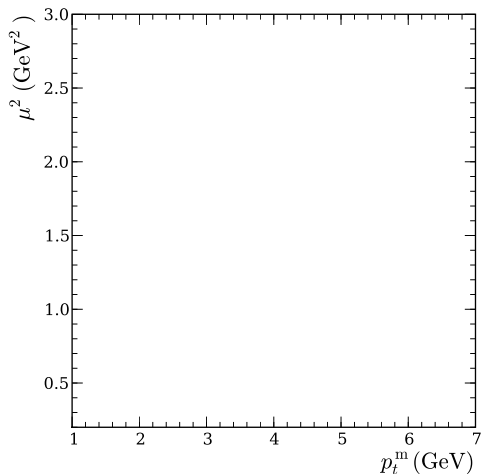
$$b_{\text{el}}(s) = \int d^2\mathbf{b} \frac{b^2}{\sigma_{\text{tot}}} \left[1 - e^{-\chi_{\text{tot}}(\mathbf{b}, s)} \right]$$

- Final state tune of **semi-hard MPI** (approximate constraint for MRST2001)



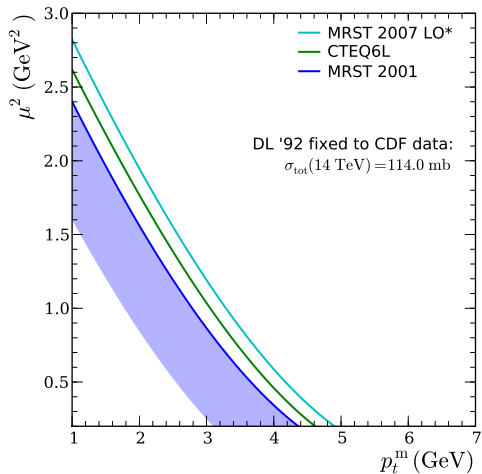
LHC parameter space

- What to expect from measurements at 14 TeV?
- $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb. σ_{tot} from Regge fit
- Require $\bar{n}_{\text{hard}} < 10$
- Require elastic t -slope to be correctly described.
Get range of possible measurements from DL '92 vs. e.g. [arXiv:0710.2494, Khoze, Martin, Ryskin] , [arXiv:0708.1506, Gotsman, Levin, Maor] predictions



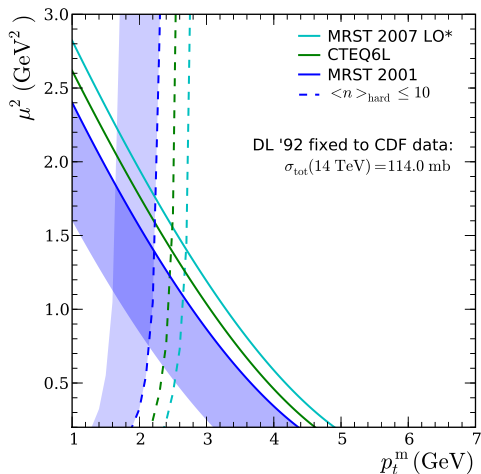
LHC parameter space

- What to expect from measurements at 14 TeV?
- $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb. σ_{tot} from Regge fit
- Require $\bar{n}_{\text{hard}} < 10$
- Require elastic t -slope to be correctly described.
Get range of possible measurements from DL '92 vs. e.g.
[arXiv:0710.2494, Khoze, Martin, Ryskin] ,
[arXiv:0708.1506, Gotsman, Levin, Maor] predictions



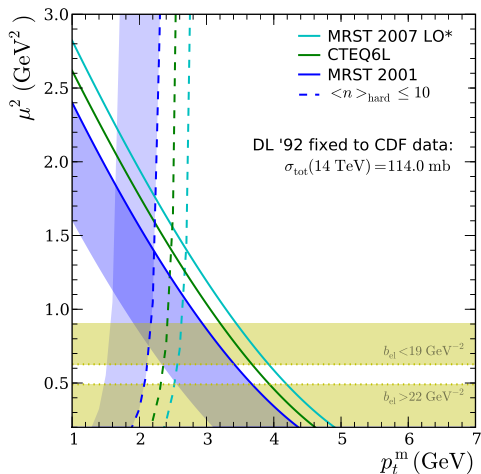
LHC parameter space

- What to expect from measurements at 14 TeV?
- $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb. σ_{tot} from Regge fit
- Require $\bar{n}_{\text{hard}} < 10$
- Require elastic t -slope to be correctly described.
Get range of possible measurements from DL '92 vs. e.g.
[arXiv:0710.2494, Khoze, Martin, Ryskin] ,
[arXiv:0708.1506, Gotsman, Levin, Maor] predictions



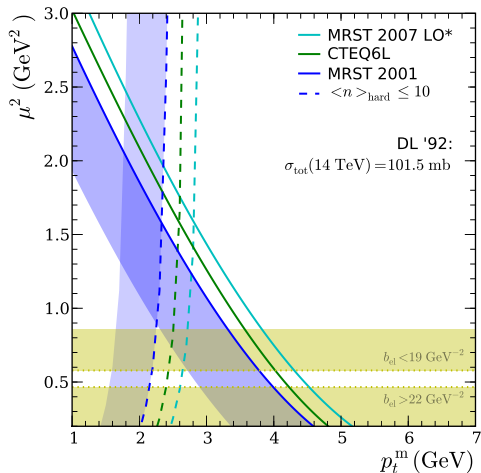
LHC parameter space

- What to expect from measurements at 14 TeV?
- $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb. σ_{tot} from Regge fit
- Require $\bar{n}_{\text{hard}} < 10$
- Require elastic t -slope to be correctly described.
Get range of possible measurements from DL '92 vs. e.g.
[arXiv:0710.2494, Khoze, Martin, Ryskin] ,
[arXiv:0708.1506, Gotsman, Levin, Maor] predictions



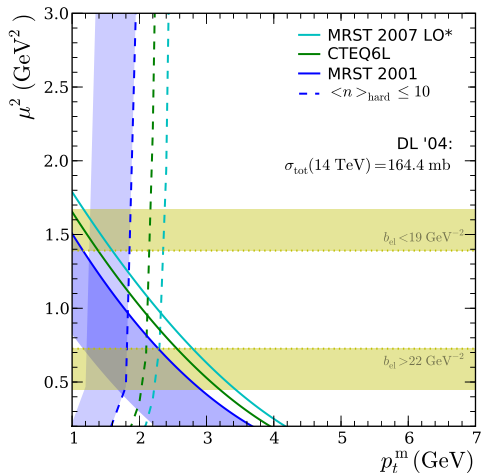
LHC parameter space

- What to expect from measurements at 14 TeV?
- $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb. σ_{tot} from Regge fit
- Require $\bar{n}_{\text{hard}} < 10$
- Require elastic t -slope to be correctly described.
Get range of possible measurements from DL '92 vs. e.g.
[arXiv:0710.2494, Khoze, Martin, Ryskin] ,
[arXiv:0708.1506, Gotsman, Levin, Maor] predictions



LHC parameter space

- What to expect from measurements at 14 TeV?
- $\sigma_{\text{soft}}^{\text{inc}} > 0$ mb. σ_{tot} from Regge fit
- Require $\bar{n}_{\text{hard}} < 10$
- Require elastic t -slope to be correctly described.
Get range of possible measurements from DL '92 vs. e.g.
[arXiv:0710.2494, Khoze, Martin, Ryskin] ,
[arXiv:0708.1506, Gotsman, Levin, Maor] predictions



Observations

- $\sigma_{\text{soft}}^{\text{inc}}$ rises artificially fast (expect $\sim s^{0.08}$)
- Forced to have energy dependent parameters (would like to have the choice, i.e. let measurements decide)
- The measurement of the elastic t -slope fixes μ^2 at Tevatron:

$$\mu^2 = 0.56 \pm 0.01 \text{ GeV}^2$$

$\sigma_{\text{eff}} = (\int d^2\mathbf{b} A^2(b))^{-1}$ as measured by CDF in the $\gamma + 3j$ channel suggests:

$$\mu^2 = 3.0 \pm 0.5 \text{ GeV}^2$$

→ Relax the constraint of identical overlap functions:

$$A_{\text{soft}}(b) = A(b, \mu_{\text{soft}})$$

If $\mu > \mu_{\text{soft}}$: Hot Spots

Outline

- 1 Introduction
 - Eikonal model basics
 - Soft part of the eikonal
- 2 Analytic constraints
 - Simple model
 - Hot-Spot model
- 3 Final states
 - Tevatron Run1
- 4 Conclusions

Fixing parameters

- Basis:

$$\chi_{\text{tot}}(\mathbf{b}, s) = \frac{1}{2} \left[A(b, \mu) \sigma_{\text{hard}}^{\text{inc}}(s; p_t^{\text{min}}) + A(b, \mu_{\text{soft}}) \sigma_{\text{soft}}^{\text{inc}} \right]$$

- Instead of introducing a new parameter, we **fix** $\mu_{\text{soft}}, \sigma_{\text{soft}}^{\text{inc}}$ by requiring $\sigma_{\text{tot}}, b_{\text{el}}$ to be correctly described simultaneously

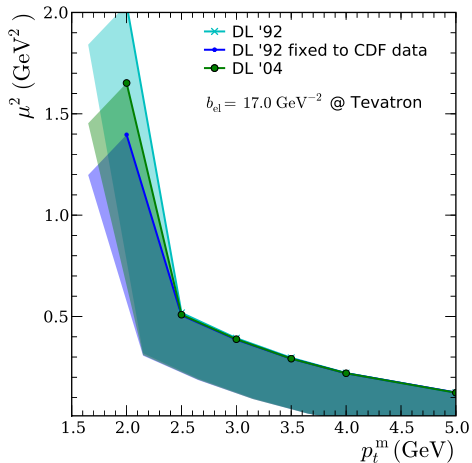
$$\sigma_{\text{tot}}(s)^{\text{measured}} - 2 \int d^2\mathbf{b} \left[1 - e^{-\chi_{\text{tot}}(\mathbf{b}, s)} \right] \stackrel{!}{=} 0$$

$$b_{\text{el}}(s)^{\text{measured}} - \int d^2\mathbf{b} \frac{b^2}{\sigma_{\text{tot}}} \left[1 - e^{-\chi_{\text{tot}}(\mathbf{b}, s)} \right] \stackrel{!}{=} 0$$

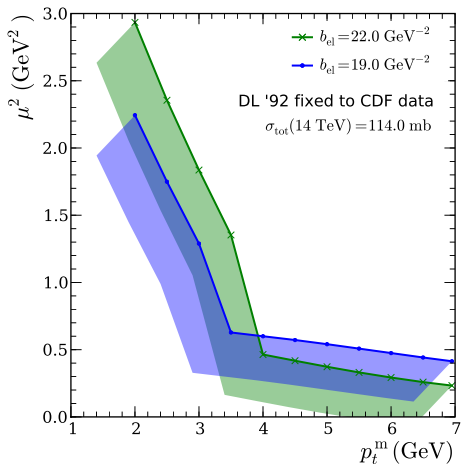
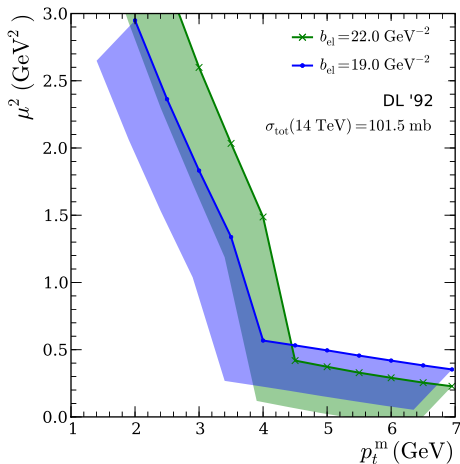
Tevatron parameter space

Only one constraint:

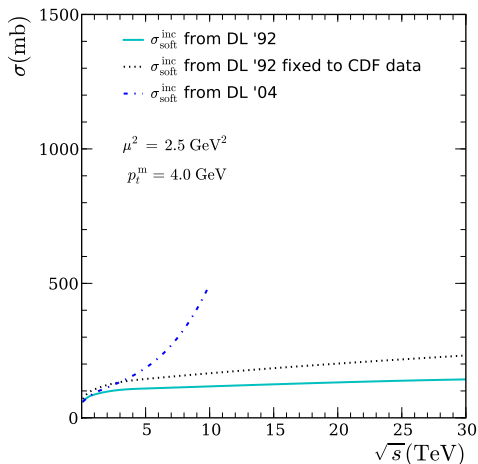
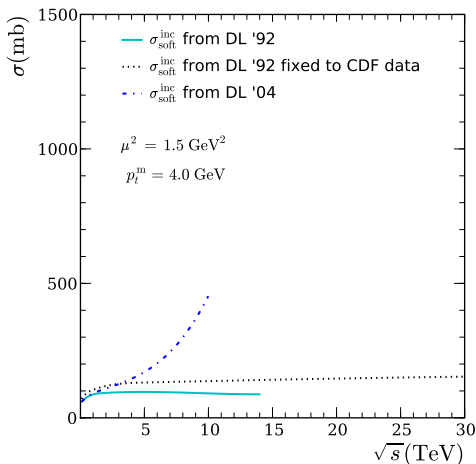
Correct total cross section and elastic slope



LHC parameter space

Same for LHC except for uncertainty in b_{el} and σ_{tot} 

Resulting soft cross section



Outline

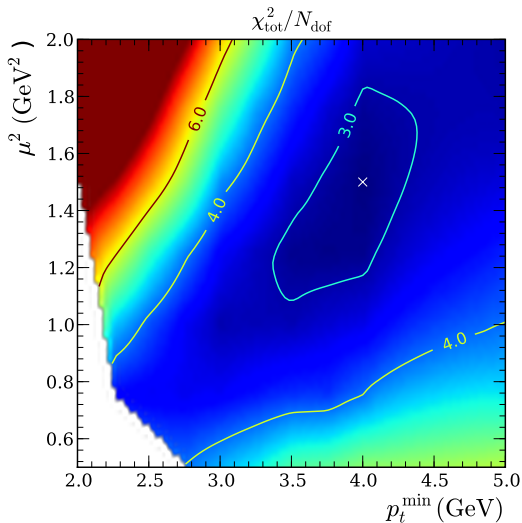
- 1 Introduction
 - Eikonal model basics
 - Soft part of the eikonal
- 2 Analytic constraints
 - Simple model
 - Hot-Spot model
- 3 Final states
 - Tevatron Run1
- 4 Conclusions

Intro

- So far: Indirect constraints from total and elastic cross sections
- Eikonal model is used to calculate the multiplicities of semi-hard and soft additional scatters in our MC implementation (which I have no time to talk about)
- Now use the results from that implementation to fit the remaining free parameters (p_t^{\min}, μ^2) to data

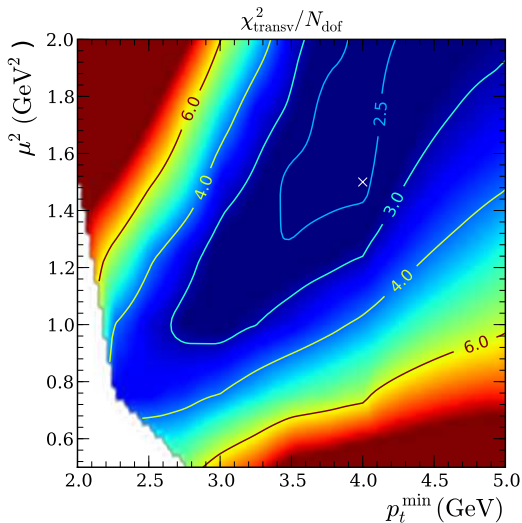
Parameter space at Tevatron

- χ^2 for Rick's Run1 Jet analysis for **all** regions
- only the transverse region

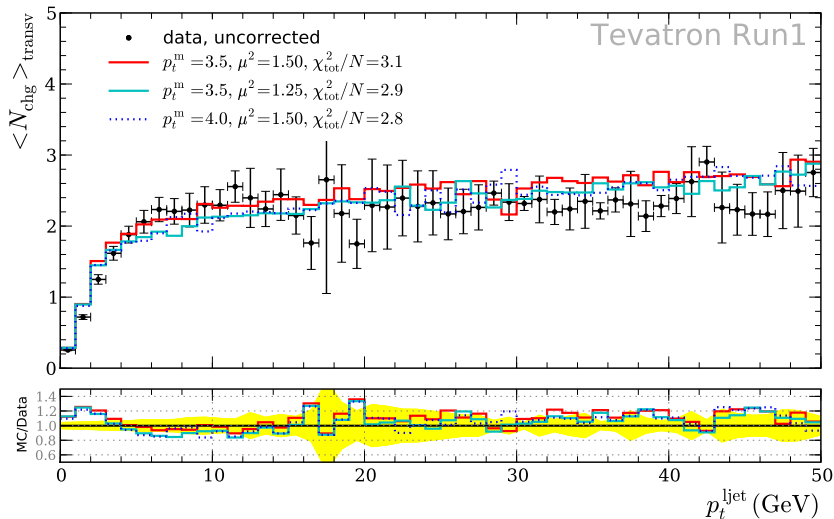


Parameter space at Tevatron

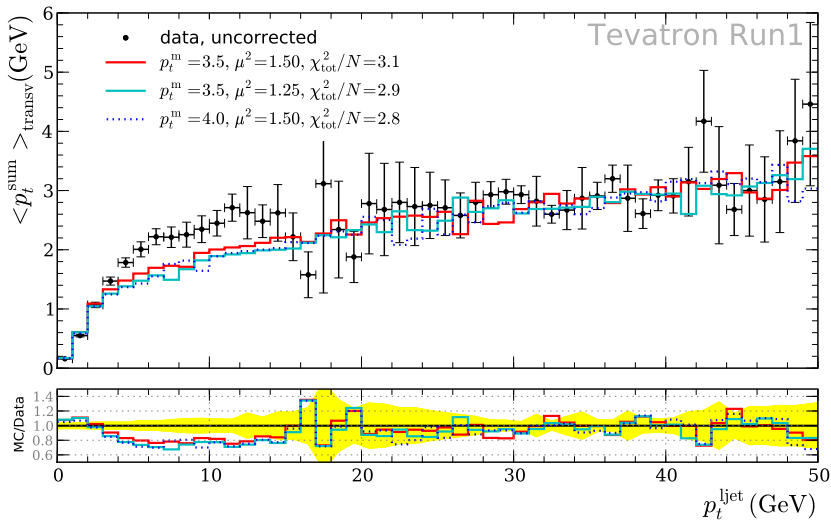
- χ^2 for Rick's Run1 Jet analysis for **all** regions
- only the transverse region



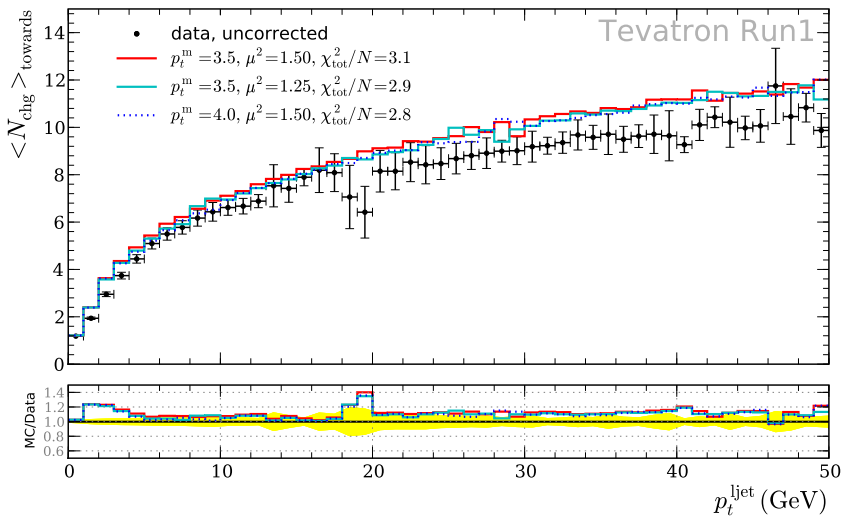
Detailed look at observables: Transverse Region



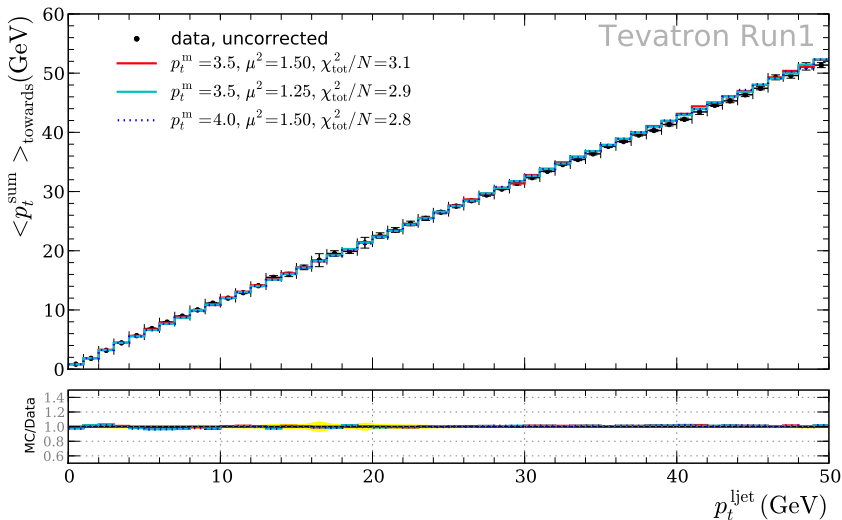
Detailed look at observables: Transverse Region



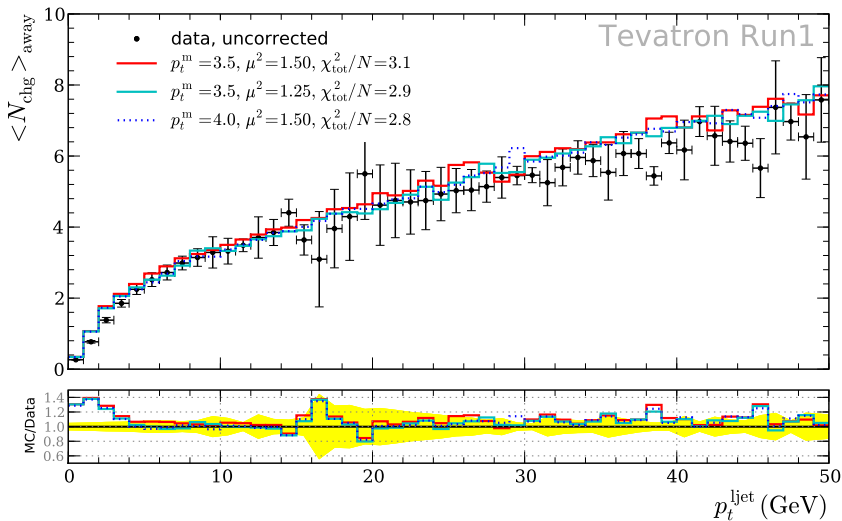
Detailed look at observables: Towards Region



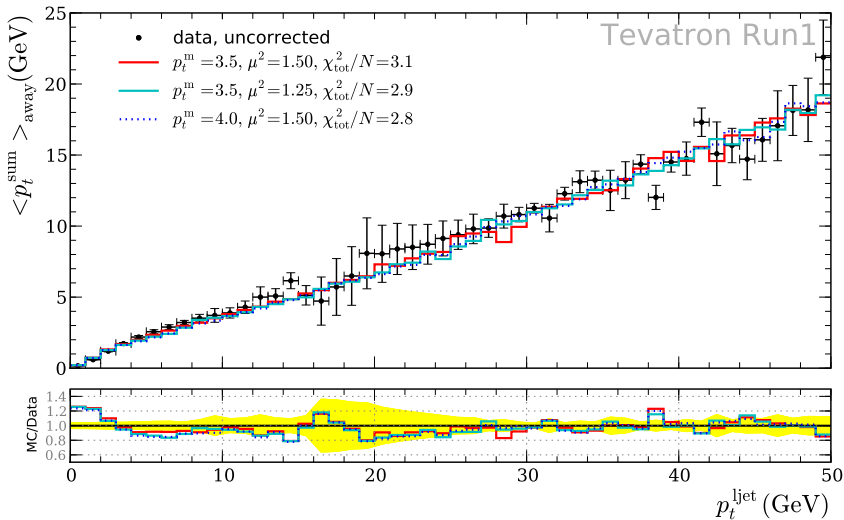
Detailed look at observables: Towards Region



Detailed look at observables: Away Region



Detailed look at observables: Away Region



Things to do

So far: Implemented the new model and tuned to CDF's Run1 data

- 1 Look at Run2 MinBias data (see talk by Niccolo Moggi)
- 2 Look at UE in DY (see talks by Rick Field and Hendrik Hoeth)
- 3 Look at Older MinBias data
- 4 Look at New LHC data
- 5 Look into possible energy dependence of spatial parton distributions (see talks by Ted Rogers and Rohini Godbole)
- 6 Consistently incorporate diffractive interactions
- 7 Allow the possibility of colour reconnections

Conclusions

- We extended the existing model to soft partonic interactions.
First time MinBias is available in Herwig ever!
- The ue activity is directly coupled to the total and elastic cross section
 - Large impact of first measurements of these quantities at LHC
 - Extrapolation to larger energies constraint by predictions for these quantities
- **Hot-Spot model** to describe inconsistency between σ_{eff} and b_{el}
- Fully exclusive simulation of **multiple hard scatterings** + low p_t jets possible (again no time to talk about it), e.g. $\gamma j + jj$, like sign W 's, several b-jet pairs (... you name it) **with arbitrary and independent cuts.**

Conclusions

- We extended the existing model to soft partonic interactions.
First time MinBias is available in Herwig ever!
- The ue activity is directly coupled to the total and elastic cross section
 - Large impact of first measurements of these quantities at LHC
 - Extrapolation to larger energies constraint by predictions for these quantities
- **Hot-Spot model** to describe inconsistency between σ_{eff} and b_{el}
- Fully exclusive simulation of **multiple hard scatterings** + low p_t jets possible (again no time to talk about it), e.g. $\gamma j + jj$, like sign W 's, several b-jet pairs (... you name it) **with arbitrary and independent cuts**.

Conclusions

- We extended the existing model to soft partonic interactions.
First time MinBias is available in Herwig ever!
- The $u\bar{e}$ activity is directly coupled to the total and elastic cross section
 - Large impact of first measurements of these quantities at LHC
 - Extrapolation to larger energies constraint by predictions for these quantities
- **Hot-Spot model** to describe inconsistency between σ_{eff} and b_{el}
- Fully exclusive simulation of **multiple hard scatterings** + low p_t jets possible (again no time to talk about it), e.g. $\gamma j + jj$, like sign W 's, several b-jet pairs (... you name it) **with arbitrary and independent cuts.**

Conclusions

- We extended the existing model to soft partonic interactions.
First time MinBias is available in Herwig ever!
- The ue activity is directly coupled to the total and elastic cross section
 - Large impact of first measurements of these quantities at LHC
 - Extrapolation to larger energies constraint by predictions for these quantities
- **Hot-Spot model** to describe inconsistency between σ_{eff} and b_{el}
- Fully exclusive simulation of **multiple hard scatterings** + low p_t jets possible (again no time to talk about it), e.g. $\gamma j + jj$, like sign W 's, several b-jet pairs (... you name it) **with arbitrary and independent cuts**.