### Electroweak Measurements

#### Vitaliano Ciulli Università e INFN Firenze Andrea Di Simone Universität Freiburg



#### Outline

W/Z physics • W charge asymmetry Z A<sub>FR</sub> asymmetry weak mixing angle VBF Z production Di-bosons cross section measurements limits on aTGC Conclusions

# W/Z bosons, introduction

Amongst the processes accessible at the LHC, the production and decay of W and Z bosons are of paramount importance

- experimentally, their leptonic decays present very clean signatures
- theoretically, we have very advanced tools at our disposal
  - NLO generators (integrated to PS in a consistent way)
  - NNLO predictions for cross sections (inclusive and differential)

In addition to the physics measurements they allow to perform, they are also a fundamental tool to understand the performance of the detectors

Tag&Probe (not covered in this talk)

# W/Z as probes

- In addition to being interesting *per se*, the study of the properties of the W and Z bosons gives insight on several parameters of the SM
- Both experiments are performing extensive studies
- Highlights given here include
  - W charge asymmetry
    - > ZA<sub>FB</sub>
      - weak mixing angle

W charge asymmetry sensitive to valence quark composition

Extrapolation to common fiducial volume allows direct comparison of ATLAS, CMS and LHCb









#### They arise from the parity violation of EW interactions

- As opposed to, for example, parity-conserving photon exchange
- The Z production+decay cross section gains a term proportional to the cosine of the scattering angle
- Coefficient depends on the left and right couplings, and vanishes if they are identical

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\cos\theta} = \frac{4\pi\alpha^2}{3s} \left[ \frac{3}{8}A(1+\cos^2\theta) + B\cos\theta \right]$$

This allows to define three non-vanishing observables:

- Longitudinal polarization asymmetry, ALR
- Unpolarized FB asymmetry
  - Polarized FB asymmetry



Z A<sub>FB</sub>

- An ambiguity in the definition of the scattering angle is present when the transverse momentum of the lepton pair in the lab frame is not negligible
  - Use a reference frame (Collins-Soper) which resolves this ambiguity by using a symmetric axis wrt the incoming partons

 $p_i^{\pm} = \frac{1}{\sqrt{2}}(E_i \pm p_{\mathrm{z},i})$ 



- In pp colliders, one extra complication arises from the fact that one does not know which beam the quark belonged to
- Result is that A<sub>FB</sub> is diluted
- Dilution less important if one limits the measurement to lepton pairs with high rapidity
  - In this case, one of the partons had high x, i.e. it was most likely a quark
    - assume that the direction of flight of the Z coincides with the direction of the quark
  - Note that this has a rather big effect on the measurement (see next slide)

# Z A<sub>FB</sub>, raw distributions



- Raw distributions allow to appreciate the effect of dilution
  - CMS: binning the measurement in bins of dilepton rapidity
  - ATLAS: using forward calorimetry to identify electrons up to |η|<4.9</li>

### Measuring the weak mixing angle

# $Z A_{\mbox{\scriptsize FB}}$ gives direct insight on the V and A couplings

- It is sensitive to the effective weak mixing angle
- Strategies for wma extraction are different
  between the two experiments
  - CMS: start from theory prediction of differential cross section, "fold in" all known effects, unbinned likelihood fit to observed cos(θ\*) distribution

muon final state only

- ATLAS: template fits on raw A<sub>FB</sub> spectra using MC samples generated with different values of wma
  - muon, electron and forward electron final states separately statistical combination of the results

#### Weak mixing angle: results

 $\sin^2 \theta_{\rm eff} = 0.2287 \pm 0.0020 \,(\text{stat.}) \pm 0.0025 \,(\text{syst.}).$ 

 $\sin^2 \theta_{W \text{ combined}}^{\text{eff}} = 0.2297 \pm 0.0004(\text{stat.}) \pm 0.0009(\text{syst.})$ 

CMS: muon channel @ 1.1/fb Phys. Rev. D 84, 112002 (2011)

ATLAS: e + forward-e + muon @ 4.7/fb ATLAS-CONF-2013-043

#### Even with only (part of) 2011 data, the result is *dominated by systematics*

- Mainly PDF, followed by detector-related effects (electron scale/resolution, allignment)
- ATLAS statistical precision is better than what one would expect from sqrt(N)
- The forward electrons provide a more precise measurement, even with smaller statistics





## EW Z production

- Electroweak production of Z boson, involving three diagrams
  - VBF, bremsstrahlung, multiperipheral
- Important benchmark to understand selection of forward jets and performance of additional veto on central jets
  - crucial for Higgs VBF analyses

• Signal is a Z boson plus two forward, well separated jets

 analysis requires two good-quality, high-p<sub>T</sub>, isolated, SFOS leptons, within 20GeV from nominal Z mass

• two highest  $p_{T}$  jets within  $|\eta| < 4.7$  are used as tagging jets Main background is DY Z+2jets

# EW Z production

- Cross section extracted through template fit on two distributions
- invariant mass of the tagging jets
- neural network output (BDT)
- Neural network yields better precision on the signal fraction
- Main systematic is JES+JER

   second largest is background modeling

  Observed cross section is in good agreement with NLO expectation (166fb)

 $\sigma^{\text{EWK}}_{\text{meas, }\mu\mu+\text{ee}} = 154 \pm 24(\text{stat}) \pm 46(\text{exp.syst.}) \pm 27(\text{th.syst}) \pm 3(\text{lumi}) \text{ fb}$ 

√s = 7 TeV

 $L = 5.0 \text{ fb}^{-1}$ 



## EW Z production

Measure detailed topology of selected events. E.g. average number of jets above 40GeV and phi separation of the two tagging jets, as a function of the eta separation of the tagging jets

m<sub>jj</sub> and NN output distributions, used for signal cross section measurement. Free parameters in the template fit are the normalizations of the DY background and of the signal



#### Dibosons

- Diboson production provides stringent tests of the electroweak sector of the SM
  - o deviations from the prediction may indicate New Physics
- In addition, these processes are background to many other channels
  - e.g. Higgs decays to ZZ
- Main backgrounds are W/Z+jets, ttbar
- Uncertainties vary considerably depending on the specific analysis
  - ZZ analyses are limited by statistical uncertainty, main systematic is lepton ID and reco
  - systematics are larger in channels including a W or a photon, with main contributions coming from ID and reco, and background estimate

### WW/WZ x-sec in leptonic decays



WZ and WW x-sec at **8 TeV** are systematically limited respectively by background and jet veto efficiency





#### CMS WW result slightly above expectations but still consistent

	√s	lumi	Measured $\sigma_{total}$ [pb]	MCFM NLO [pb]
ATLAS	7 TeV	4.7 fb <sup>-1</sup>	51.9 ± 2.0 (stat) ± 3.9 (sys) ± 2.0 (lumi)	<b>44.7</b> <sup>+ 2.1</sup> - 1.9
CMS	7 TeV	4.9 fb <sup>-1</sup>	52.4 ± 2.0 (stat) ± 4.5 (sys) ± 1.2 (lumi)	47.0 ± 2.0
CHS	8 TeV	3.5 fb <sup>-1</sup>	$69.9 \pm 2.8$ (stat) $\pm 5.6$ (sys) $\pm 3.1$ (lumi)	57.3 + 2.4

#### ZZ x-sec in leptonic decays

ZZ events, and their fully leptonic decays, offer very clean final state

#### Results are still statistically limited

	√s	lumi	Measured $\sigma_{total}$ [pb]	MCFM NLO [pb]	Z mass window in fidutial space
ATLAS	7 TeV	4.7 fb <sup>-1</sup>	7.2 ± 1.4 (stat) ± 0.8 (sys) ± 0.4 (lumi)	6.5 ± 0.3	66 < m⊪ <116 GeV
CMS	7 TeV	4.9 fb <sup>-1</sup>	6.2 ± 2.4 (stat) ± 1.1 (sys) ± 1.0 (lumi)	6.3 ± 0.4	60 < m <sub>ll</sub> < 120 GeV
ATLAS	8 TeV	20.0 fb <sup>-1</sup>	7.1 $\pm$ 0.4 (stat) $\pm$ 0.3 (sys) $\pm$ 0.2 (lumi)	7.2 ± 0.3	66 < m⊪ < 116GeV
CMS	8 TeV	5.3 fb <sup>-1</sup>	8.4 ± 1.0 (stat) ± 0.7 (sys) ± 0.4 (lumi)	7.7 ± 0.4	60 < m⊫ < 120 GeV







# WW(WZ) x-sec in semileptonic decays



 $\begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \int L \, dt = 4.7 \, \text{fb}^{-1} \\ \hline 1 \, \dot{s} = 7 \, \text{TeV} \end{array} \\ \begin{array}{c} W \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \end{array}$  \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets \end{array} \\ \begin{array}{c} \text{W} \rightarrow I(e, \mu)v + 2 \, jets

Look at events with one W decaying leptonically and a second boson (W or Z) hadronically

The main systematics are JES and W+jets background

Both ATLAS and CMS constraint normalization and shape of background in the fit. CMS also constraints the JES to W mass

Measurements performed at 7 *TeV* only

Results are in agreement with SM expectations





σ(WW + WZ)	measured (pb)	expected (pb)
ATLAS	72 ± 9 (stat.) ± 15 (syst.) ± 13 (MC stat.)	63.4 ± 2.6
СМЗ	68.9±8.7(stat.)±9.7(syst.)± 1.5 (lum.)	65.6 ± 2.2





#### Measurements performed at **7 TeV**

Look at events with one W or one Z decaying leptonically and an isolated photon





The final state with the Z decaying into neutrinos and an isolated photon is also considered



# **W**γ/**Z**γ





V. Ciulli (INFN FI), A. Di Simone (Uni Freiburg)





Results are in good agreement with generator predictions MCFM slightly lower for W $\gamma$  results both in ATLAS and CMS



Coupling	Parameters	Channel
wwγ	λγ, Δκγ	₩₩,₩Y
wwz	$\lambda$ Z, Δκz, Δg <sub>1</sub> <sup>Z</sup>	WW,WZ
ZZγ	h₃ <sup>Z</sup> , h₄ <sup>Z</sup>	Zγ
Ζγγ	h₃ <sup>ץ</sup> , h₄ <sup>γ</sup>	Zγ
ZZZ	$f_4^Z, f_5^Z$	ZZ
ΖγΖ	f4 <sup>Y</sup> , f5 <sup>Y</sup>	ZZ

- Triple gauge couplings in SM:
  - Charged triple gauge couplings (WWZ, WW) allowed
  - Neutral triple gauge couplings (ZZZ, ZZγ) forbidden
- aTGCs modify total production rate as well as event kinematics.
- Maximum likelihood fit is done, leaving one (or two) of the aTGC parameter free.

Form factor used in old results:  $\Delta g(s) = \Delta g/(1+s/\Lambda^2)^2$ for recent we have  $\Lambda = \infty$   $Z\gamma \rightarrow II\gamma$  and  $Z\gamma \rightarrow vv\gamma$  combined to improve limits on neutral aTGC CMS Limits are tighter due to 400 GeV  $p_{\tau}$  threshold of last bin in the

 $vv\gamma$  channel, which drives the limit

**Cross section** measurements described in the previous slides can be interpreted in terms of potential contributions from aTGCs

**Results for** charged aTGC are similar for ATLAS and **CMS** CMS also published limits with  $WW \rightarrow Ivij$  $(\Delta g_1^Z = 1)$ 

Feb 2013			
			ATLAS Limits HI CMS Limits HI CDF Limit HI
۳	<b>⊢</b> −−−−	Zγ	-0.015 - 0.016 4.6 fb <sup>-1</sup>
n <sub>3</sub>	н	Zγ	-0.003 - 0.003 5.0 fb <sup>-1</sup>
	⊢I	Zγ	-0.022 - 0.020 5.1 fb <sup>-1</sup>
ьZ	<b>⊢−−−−</b> 1	Zγ	-0.013 - 0.014 4.6 fb <sup>-1</sup>
n <sub>3</sub>	н	Zγ	-0.003 - 0.003 5.0 fb <sup>-1</sup>
	⊢I	Zγ	-0.020 - 0.021 5.1 fb <sup>-1</sup>
h <sup>Y</sup> 100	⊢I	Zγ	-0.009 - 0.009 4.6 fb <sup>-1</sup>
n <sub>4</sub> x100	Н	Zγ	-0.001 - 0.001 5.0 fb <sup>-1</sup>
hZ. 100	<u>н</u>	Zγ	-0.009 - 0.009 4.6 fb <sup>-1</sup>
$n_4 \times 100$	Н	Zγ	-0.001 - 0.001 5.0 fb <sup>-1</sup>
-0.5	0	0.5	1 1.5 x10
			aTGC Limits @95% C.L

			.	ATLAS Limits CMS Limits D0 Limit LEP Limit	
Δκ	H	WW		-0.043 - 0.043	4.6 fb <sup>-1</sup>
ΔĸZ	H	WV		-0.043 - 0.033	5.0 fb <sup>-1</sup>
	<b>⊢</b> ●-	LEP Co	mbination	-0.074 - 0.051	0.7 fb <sup>-1</sup>
λ	$\vdash$	WW		-0.062 - 0.059	4.6 fb <sup>-1</sup>
Λz	H	WW		-0.048 - 0.048	4.9 fb <sup>-1</sup>
	$\vdash$	WZ		-0.046 - 0.047	4.6 fb <sup>-1</sup>
	H	WV		-0.038 - 0.030	5.0 fb <sup>-1</sup>
	ю	D0 Con	nbination	-0.036 - 0.044	8.6 fb <sup>-1</sup>
	H	LEP Co	mbination	-0.059 - 0.017	0.7 fb <sup>-1</sup>
۸az	$\vdash$	WW		-0.039 - 0.052	4.6 fb <sup>-1</sup>
∆9 <sub>1</sub>	<b>⊢−−−</b> 1	ww		-0.095 - 0.095	4.9 fb <sup>-1</sup>
	$\vdash$	WZ		-0.057 - 0.093	4.6 fb <sup>-1</sup>
	$\vdash \circ \dashv$	D0 Con	nbination	-0.034 - 0.084	8.6 fb <sup>-1</sup>
	H	LEP Co	mbination	-0.054 - 0.021	0.7 fb <sup>-1</sup>
0.5	<u> </u>		1	1 5	
-0.5	U	0.5		1.5	





aTGCs

### neutral aTGC from ZZ

			ATLAS Limits
εŶ		ZZ	-0.015 - 0.015 4.6 fb <sup>-</sup>
$\mathbf{f}_4$	<b>  </b>	ZZ	-0.013 - 0.015 5.0 fb <sup>-</sup>
۶Z		ZZ	-0.013 - 0.013 4.6 fb <sup>-</sup>
1 <sub>4</sub>	<b>⊢−−−−</b> 1	ZZ	-0.011 - 0.012 5.0 fb <sup>-</sup>
¢۲	<b>  </b>	ZZ	-0.016 - 0.015 4.6 fb
1 <sub>5</sub>	<b>⊢−−−−−</b> 1	ZZ	-0.014 - 0.014 5.0 fb
۶Z		ZZ	-0.013 - 0.013 4.6 fb
1 <sub>5</sub>	<b>⊢−−−−</b> 1	ZZ	-0.012 - 0.012 5.0 fb
0.5		0.5	
-0.5	U	0.5	I I.O XIU









#### **Conclusions**

# LHC data provides a wealth of information concening the SM

- In particular its EW sector
- Physics of the W and Z bosons plays a crucial role in the exploitation of this huge potential
  - both single- and di-boson production
- With the present statistics and with the experimental systematics under control, it becomes possible to use these analyses to probe other areas of the SM
  - PDFs
  - weak mixing angle

#### Discussion

#### wma systematics

- Differences in the experimental method are reflected in the systematics
  - precision on weak mixing angle in ATLAS is largely driven by the forward-electrons channel
    - wrt central electrons, it only has about 1/3 of the total events, but the statistical error is 2/3
    - drawback is that detector-related systematics are larger (e. g. electron energy scale)
  - CMS systematics due to FSR much larger than for ATLAS
    - could this be due to using only muons?
- In general, the measurement is dominated by systematics
  - more importantly, it is dominated by *theory* systematics
    - PDF
    - QCD/EW NLO

#### wma systematics

**ATLAS** 

	CC electrons	CF electrons	Muons	Combined
Uncertainty source	$(10^{-4})$	$(10^{-4})$	$(10^{-4})$	$(10^{-4})$
PDF	9	5	9	7
MC statistics	9	5	9	4
Electron energy scale	4	6	_	4
Electron energy smearing	4	5	_	3
Muon energy scale	_	—	5	2
Higher-order corrections	3	1	3	2
Other sources	1	1	2	2

source	correction	uncertainty
PDF	-	±0.0013
FSR	-	$\pm 0.0011$
LO model (EWK)	-	$\pm 0.0002$
LO model (QCD)	+0.0012	$\pm 0.0012$
resolution and alignment	+0.0007	$\pm 0.0013$
efficiency and acceptance	-	$\pm 0.0003$
background	-	$\pm 0.0001$
total	+0.0019	±0.0025

CMS

#### wma systematics

#### Not clear how to improve the measurement in the future

- will the "analytical folding" procedure à la CMS be still doable with improved statistical precision?
- template fits à la ATLAS require a generator capable of changing the effective wma value without altering the masses of the bosons
- Both measurements rely heavily on LO generator (Pythia)
  - but NLO QCD has a large effect, so one should probably move to NLO generators
  - with improved statistics, NLO EW will also become relevant
  - how to define the quark-lepton angle in an NLO world?

#### W+jets background systematicys on WW measurement





The shape of the W+jets bkg is taken from MC with ME+PS (LO)

- uncertainty is modeled changing ren/fac scale and matching parameters
- scale variations in LO MC are known to be very large compared to data
- problem is not the rate but shape and exclusivity of the selection
- uncertainty will be smaller using NLO ME+PS generators?



#### Results on aTGC are still dominated by statistical errors

- current constraints on charged aTGCs < 10%
  - already improved over LEP on several parameters
- expected to reach few % at 8 TeV

#### Perspectives on aQGC

- can be measured from
  - tri-bosons final states
  - Vector Boson Scattering
- SM cross-sections are ~2 order of magnitude lower than dibosons
- Theoretical interpretation of anomalies is much more difficult
  - quartic couplings expected also from aTGC
  - $\circ$  unitarity is violated quite soon  $\rightarrow$  add form factor

#### $\gamma\gamma \rightarrow WW$

### Search for eµ events with no additional tracks aQGC enhance x-sec at high p<sub>T</sub>(eµ)

- results on 5 fb<sup>-1</sup> of 7 TeV data
  - $\sim$  2 events with p<sub>T</sub>(eµ) > 30 GeV compatible with SM



no events with  $p_T$ (eµ) > 100 GeV

Limits on aQGC with  $\Lambda$ =500 GeV about 2 order of magnitude better than LEP



Without form factor the limit is much smaller, but dominated by the high energy  $\gamma\gamma$  region ~ 1 TeV that is above the unitarity bound



#### SM and the wma

- SM only needs three input parameters
  - All other observables can be predicted from these
- The most common prescription uses the Z mass, the fine structure constant and the Fermi constant
  - In this context, wma and the W mass can be calculated (given some value of m<sub>t</sub> and m<sub>H</sub>)
- At tree level, all is nice and simple
  - The problem arises when incorporating higher orders
  - Depending on where one decides to "hide" HO contributions in the formulas, one ends up having different predictions for the calculated quantities, with different dependence on m<sub>t</sub>
- A striking consequence of this is that the PDG lists not one, but *FIVE* different possible values of the weak mixing angle

#### SM and the wma

- The most relevant normalization scheme for Z-A<sub>FB</sub> measurements is the Effective one, since it is directly related to the coupling of Z to fermions
- Bulk of EW corrections is absorbed into effective couplings
- One then defines an *effective* weak mixing angle in such a way that the new couplings are proportional to the tree-level ones
  - The predictions for the asymmetries stay formally identical to the tree-level expressions, modulo using the effective

anale				
5	Scheme	Notation	Value	
	On-shell	$s_W^2$	0.2233	
	NOV	$s_{M_Z}^2$	0.2311	
	$\overline{\mathrm{MS}}$	$\widehat{s}_Z^2$	0.2313	
	$\overline{\mathrm{MS}}$ ND	$\widehat{s}_{ ext{ND}}^2$	0.2315	
	Effective angle	$\overline{s}_{f}^{2}$	0.2316	



#### **CMS** Detector

SILICON TRACKER Pixels (100 x 150 µm<sup>2</sup>) ~1m<sup>2</sup> ~66M channels Microstrips (80-180µm) ~200m<sup>2</sup> ~9.6M channels

#### CRYSTAL ELECTROMAGNETIC CALORIMETER (ECAL) ~76k scintillating PbWO<sub>4</sub> crystals

#### PRESHOWER Silicon strips ~16m<sup>2</sup> ~137k channels

STEEL RETURN YOKE ~13000 tonnes

SUPERCONDUCTING SOLENOID Niobium-titanium coll carrying ~18000 A

Total weight Overall diameter Overall length Magnetic field : 14000 tonnes : 15.0 m : 28.7 m : 3.8 T HADRON CALORIMETER (HCAL) Brass + plastic scintillator ~7k channels

#### CALORIMETER Steel + quartz fibres ~2k channels

FORWARD

MUON CHAMBERS

Barrel: 250 Drift Tube & 480 Resistive Plate Chambers Endcaps: 473 Cathode Strip & 432 Resistive Plate Chambers