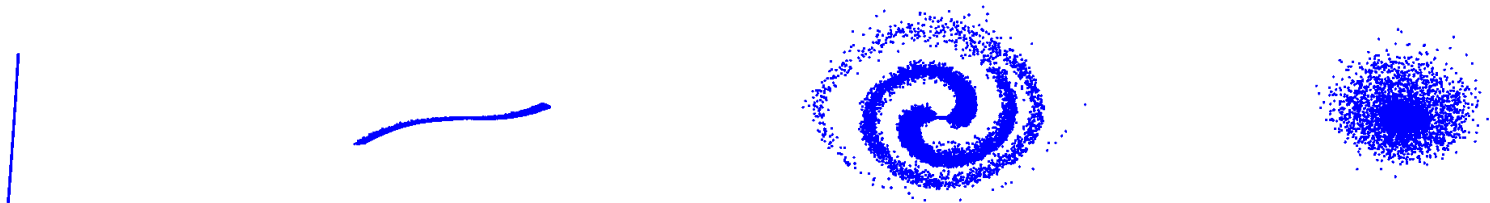


# Study of LWFA as Injectors for Synchrotron Light Sources

Steffen Hillenbrand,  
Ralph Assmann, Oliver Jansen, Vitali Judin, Anke-Susanne Müller, Alexander Pukhov



# Outline



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- Motivation
- Laser Wake Field Acceleration Simulations
- Transfer Line
- Behaviour in Synchrotron
- Summary

# Outline



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## ■ Motivation

- LWFA Simulations
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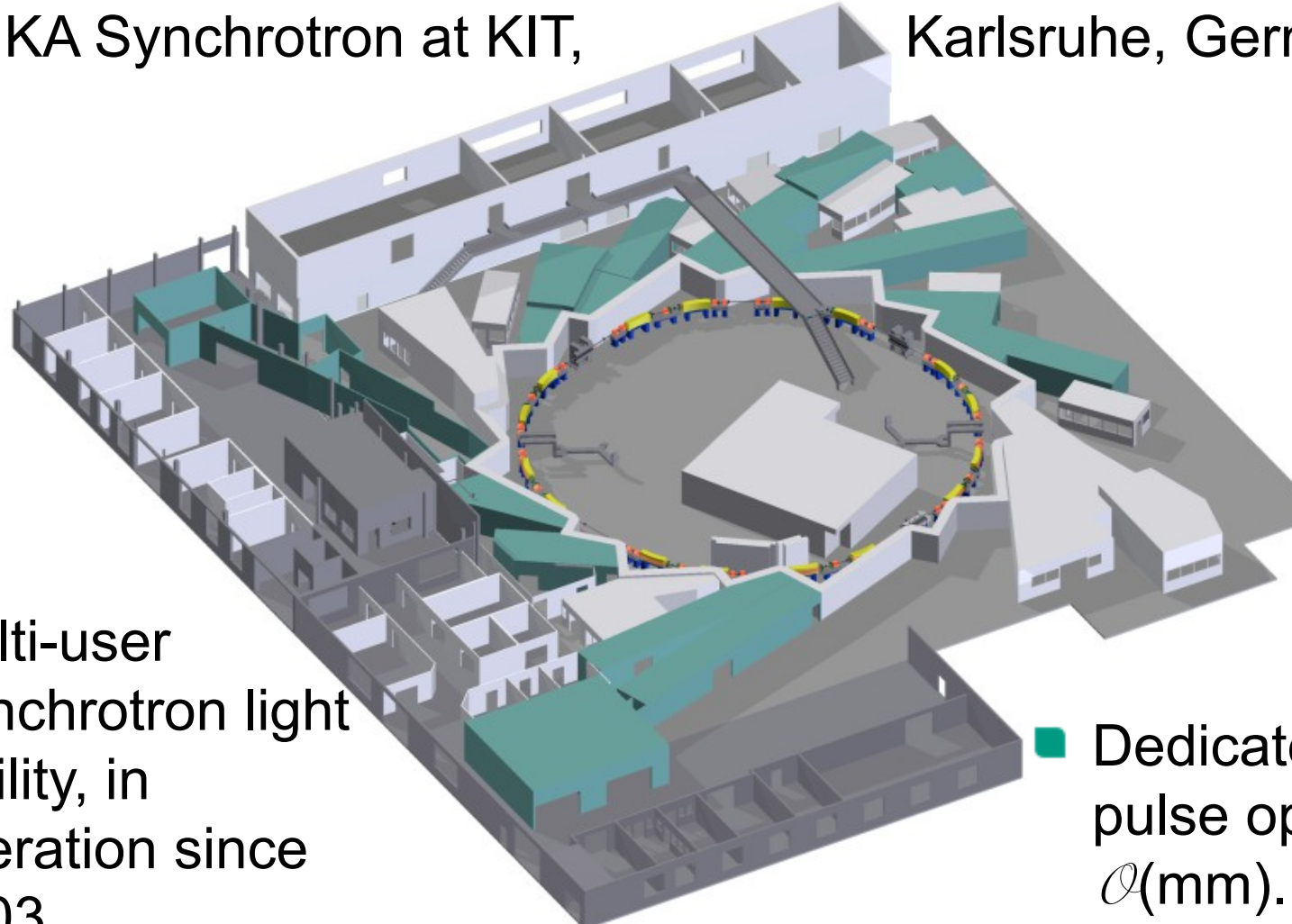
# The ANKA Synchrotron



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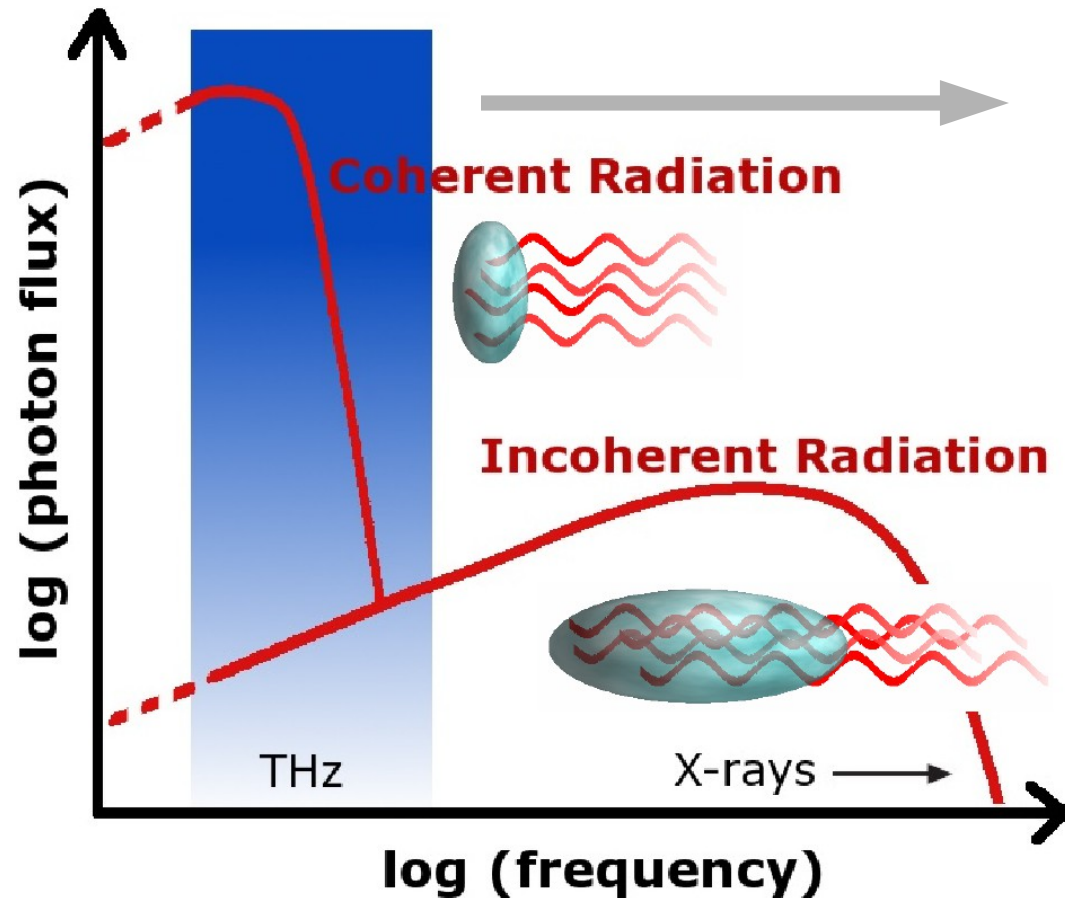
- Exemplary Simulations have been carried out using the ANKA Synchrotron at KIT, Karlsruhe, Germany.



- Multi-user Synchrotron light facility, in operation since 2003.
- Dedicated short pulse operation,  $\mathcal{O}(\text{mm})$ .

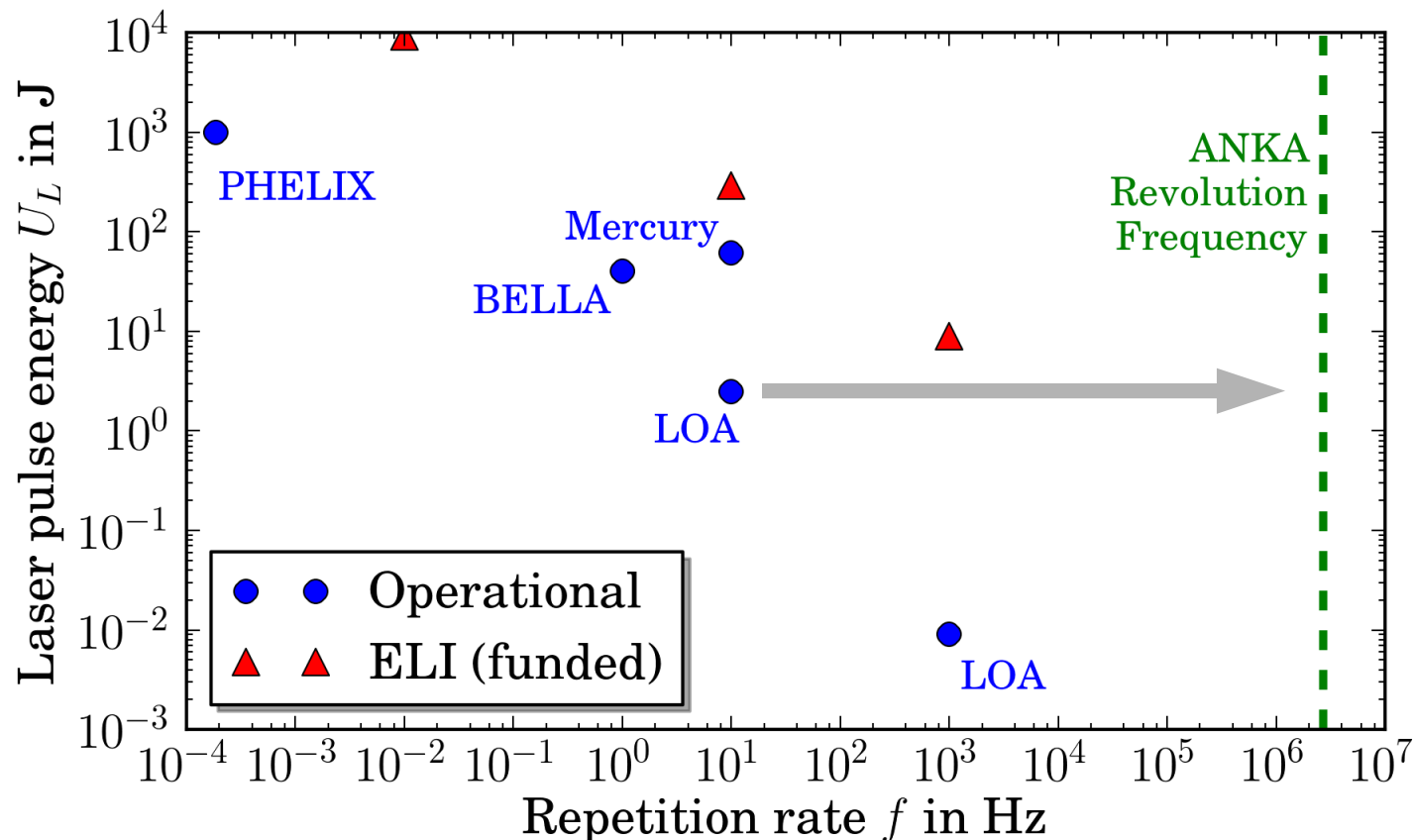
# Radiation Spectrum

- For bunch lengths shorter than the emitted radiation wavelength, emission becomes coherent.
- THz region currently difficult to access with Synchrotrons.
- LWFA inherently deliver ultra short pulses.



# Repetition Frequency

- Storing and “re-using” LWFA bunches in a Synchrotron would allow for much higher statistics for experiments.



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# LWFA Simulation I / II

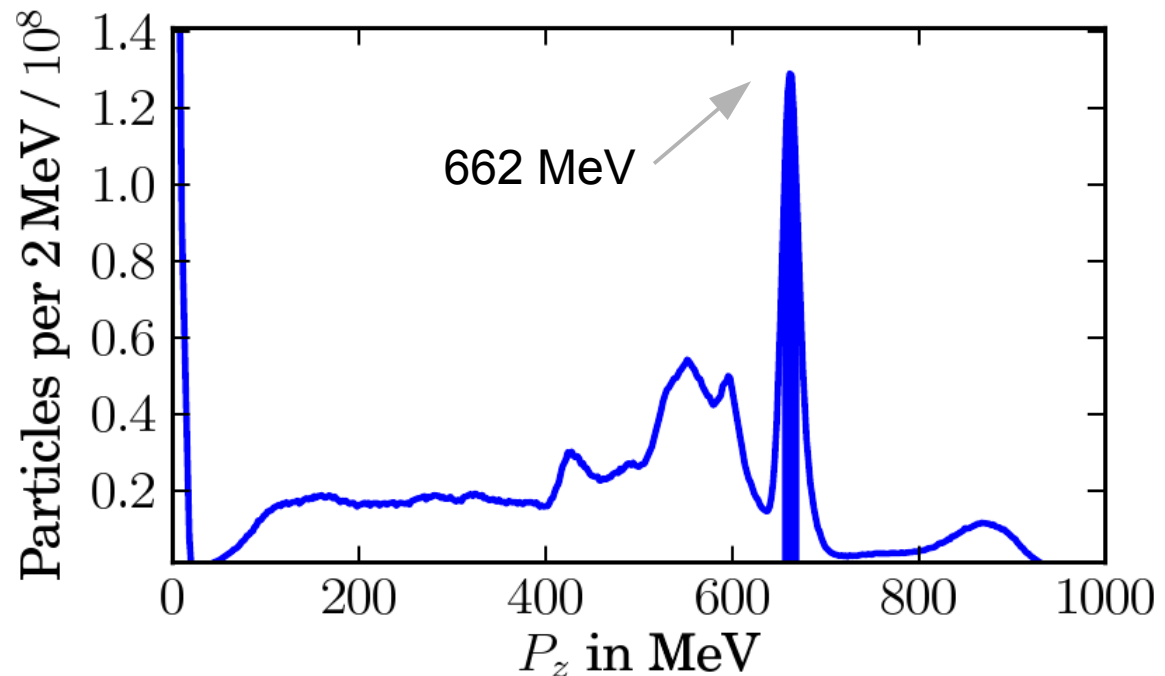
O. Jansen, A. Pukhov, HHUD Duesseldorf  
Using the 3D-PIC code VLPL [1]



- Energy spectrum at time of laser depletion is shown to the right.
- Energy acceptance of the ANKA Synchrotron ( $p_0 \pm 1\%$ ) is indicated by the solid area.

Table 1: VLPL Input Parameters

Plasma density $n_0$ in $\text{cm}^{-3}$	$9 \times 10^{18}$
Laser wavelength $\lambda_L$ in nm	800
RMS Laser pulse duration $\tau_L$ in fs	13
RMS Laser radius at focus $r_L$ in nm	720
Laser pulse energy $E_L$ in J	4.6



[1] Alexander Pukhov.  
Three-dimensional  
electromagnetic relativistic  
particle-in-cell code VLPL  
J. of Plasma Physics,  
61(3):425, 1999.



# LWFA Simulation II / II

O. Jansen, A. Pukhov, HHUD Duesseldorf  
Using the 3D-PIC code VLPL



■ All particles outside the ANKA energy acceptance have been discarded.

■ Properties of the remaining bunch are given in Tab. 2

Table 2: LWFA  $e^-$ -Beam Parameters

Central energy $p_0$ in MeV	662	
Applied energy cut in MeV	655 - 669	
RMS energy spread $\delta$ in %	0.5	←
Bunch charge $q$ in pC	160	←
Number of particles $N$	$1.0 \times 10^9$	
Geometric emittance $\epsilon_{geo}$ in $m \times rad$	$1.8 \cdot 10^{-8}$	
Normalized emittance $\epsilon_{norm}$ in $m \times rad$	$2.3 \cdot 10^{-5}$	
RMS Bunch length $\sigma_z$ in $\mu m$	1.1	
RMS Bunch length $\sigma_z$ in fs	3.7	
RMS Bunch radius $\sigma_r$ in $\mu m$	1.6	←
RMS Divergence in rad	0.01	←
Twiss $\alpha_x = \alpha_y$	0.0	
Twiss $\beta_x = \beta_y$ in m	$1.4 \times 10^{-4}$	

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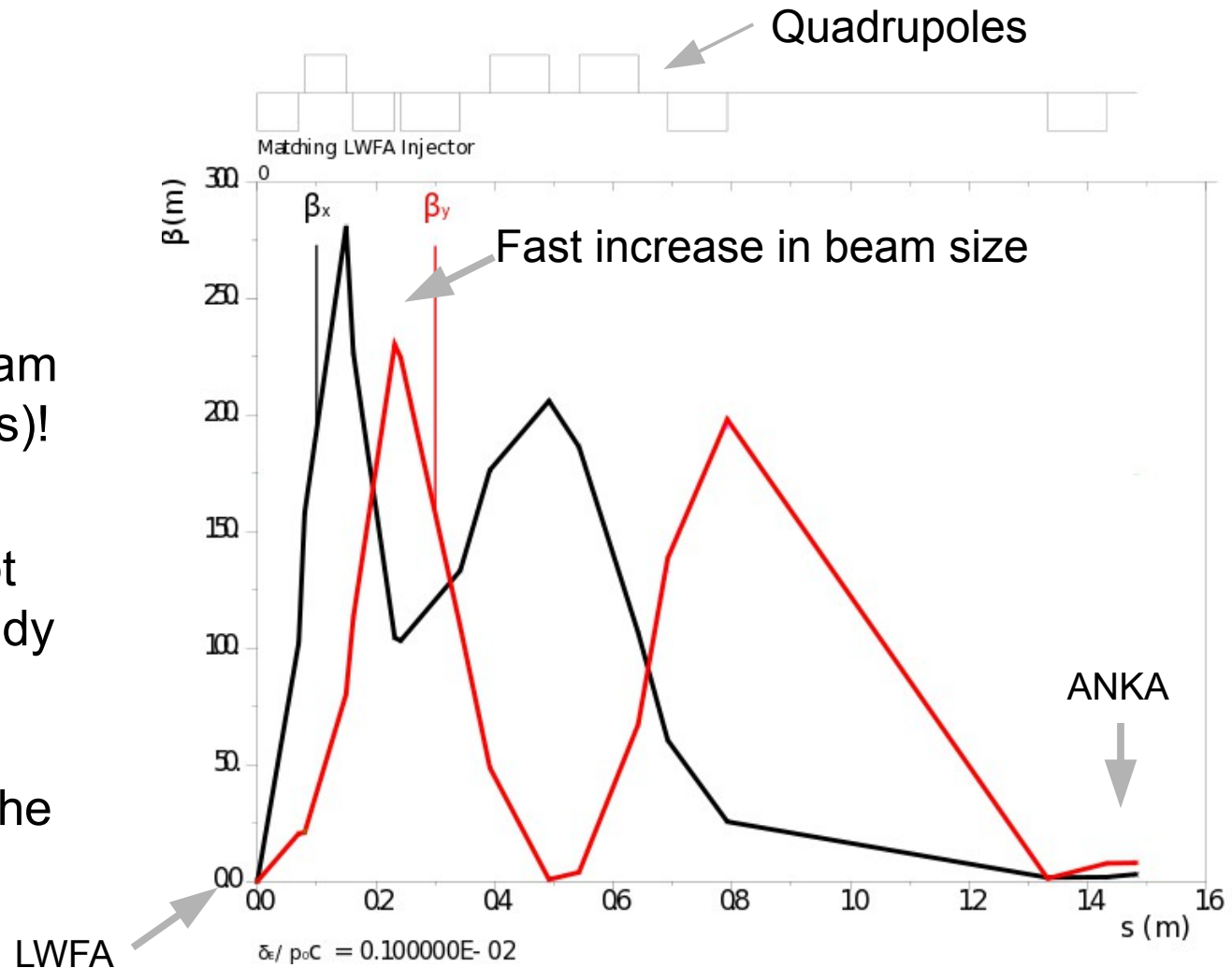
# Transfer Line



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- Exemplary transfer line.
- Note the initial explosion in beam size ( $\beta$ -functions)!
- Chromaticity not corrected, already resulting in a significant lengthening of the bunch.





- Higher order multipoles could be used to (at least partially) compensate the chromaticity, mitigating the bunch lengthening (Cf. eg. [2]).
- If not achromatic, bunch already lengthens to a few mm in the transfer line (factor of  $\sim 5000$ )!
- LWFA simulations have been performed until laser depletion. Density drop at plasma exit should increase beam size and reduce divergence, mitigating the constraints on the transfer line.

---

[2] C. Widmann, Design of a Dispersive Beam Transport Line for the JETI Laser Wakefield Accelerator, IPAC11

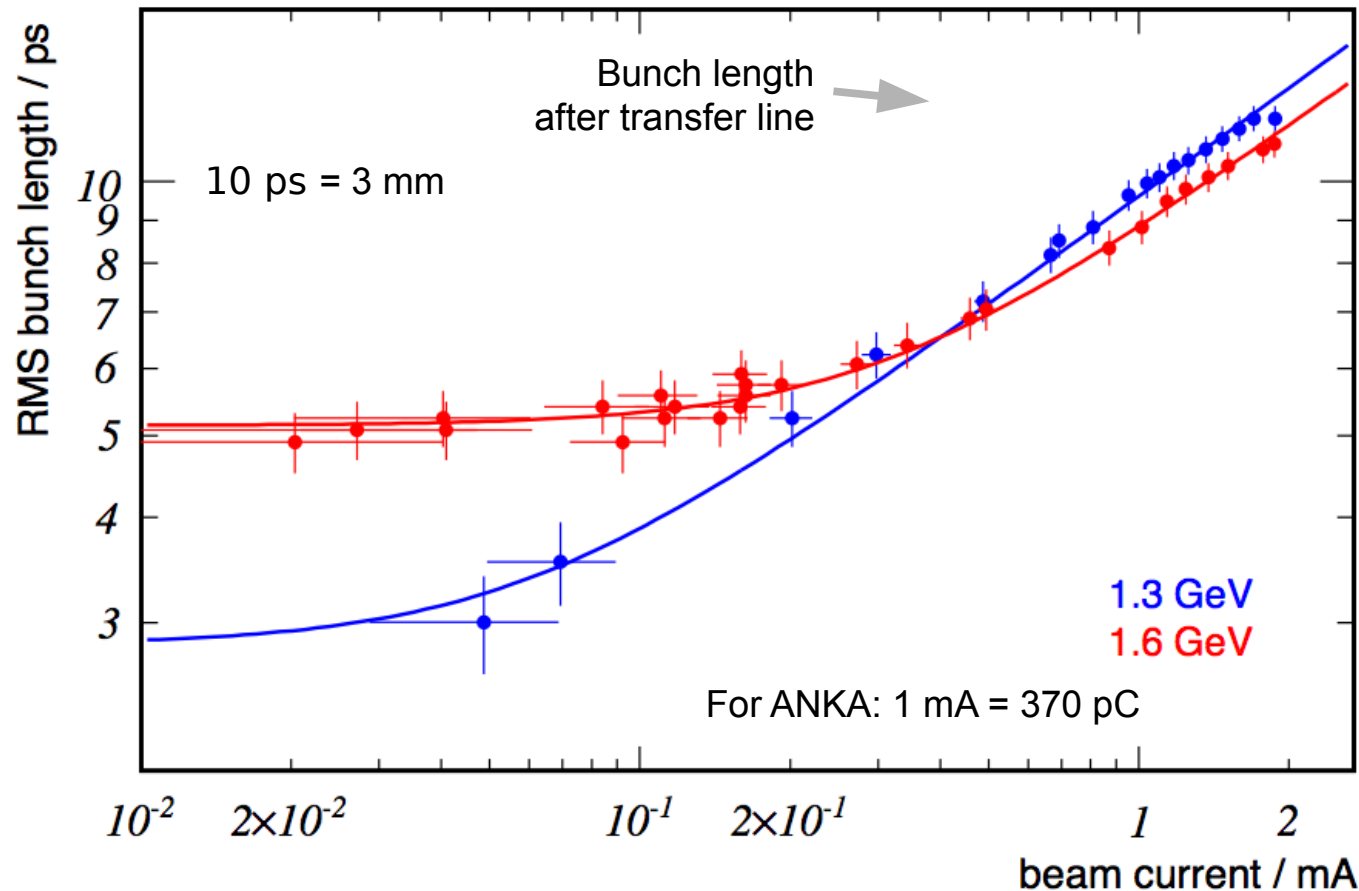
# Comparison ANKA low- $\alpha$



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- Cf. *equilibrium* bunch length measurements at ANKA [3]



- [3] N. Hiller et al., Status of Bunch Deformation and Lengthening Studies at the ANKA Storage Ring, IPAC11

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# ANKA Lattice Parameter



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## ■ ANKA injection optic, scaled to 662 MeV central energy:

Central energy $p_0$ in MeV	662	
Cavity voltage in kV	200	
Cavity frequency in MHz	499	
Circumference in m	110.4	←
Revolution time in ns	368	
Momentum compaction factor	0.008	←
Natural RMS energy spread	$2.4 \times 10^{-4}$	
Natural geometric emittance in $\text{m} \times \text{rad}$	$6.8 \cdot 10^{-9}$	← Comparable to LWFA bunch
Radiation energy damping time in ms	79	
Linear energy acceptance in %	1.1	
Synchrotron tune in kHz	22.7	
Synchrotron tune in turns	119.5	
Bunch length in mm	4.0	
Bunch length in ps	13.4	

# Reminder: Damping Time



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- For electron storage rings, the emittance is given as the equilibrium between radiative damping and quantum excitation.
- Exact damping times depend on lattice structure and beam energy.
- For ANKA at 662 MeV, the (longitudinal) energy damping time is  $\sim 80$  ms, corresponding to  $\sim 200.000$  revolutions.
- Naively, one could expect that lengthening process happens on this time scale.

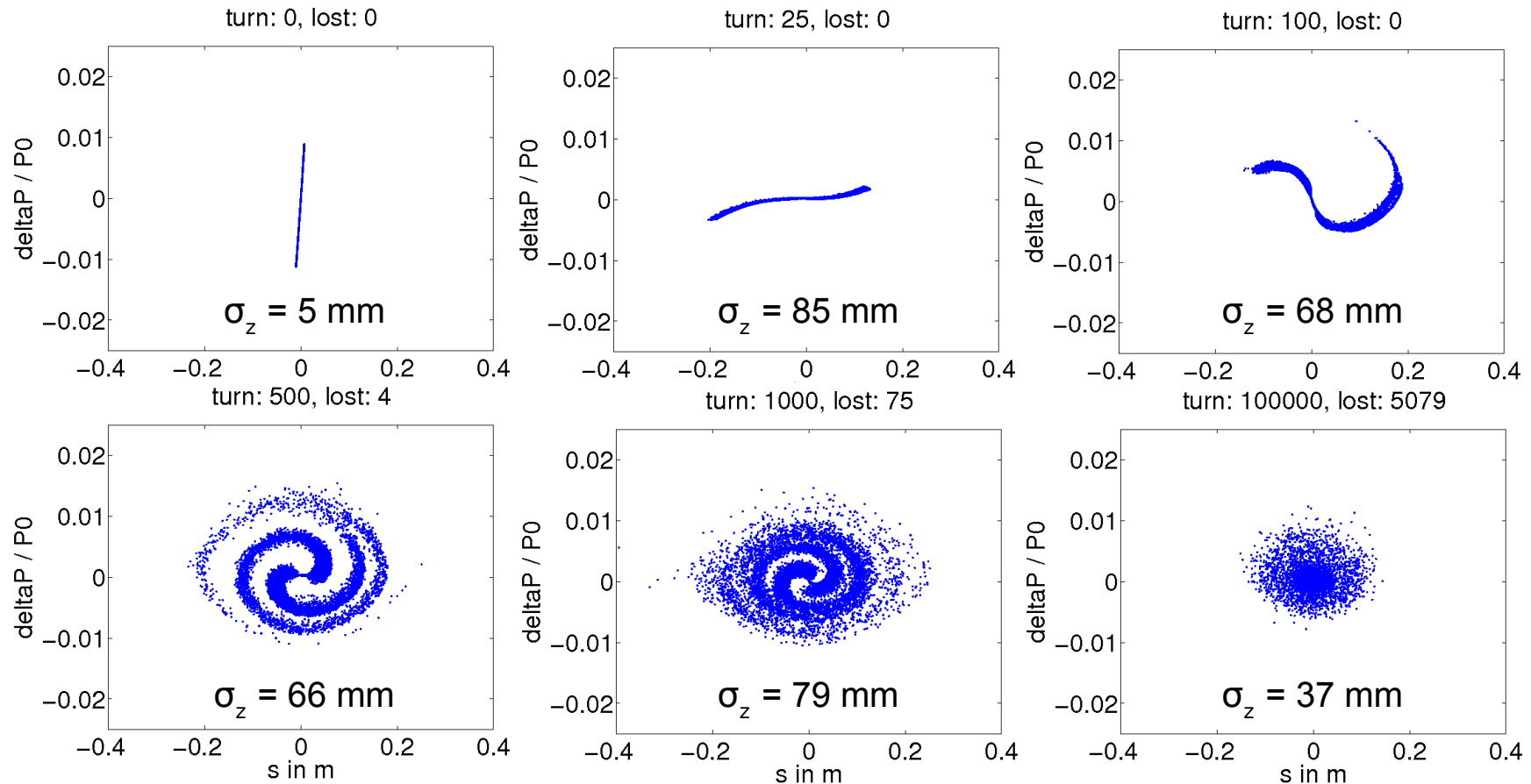


# Longitudinal Results I / II

Using AT [4] with 10.000 simulated particles



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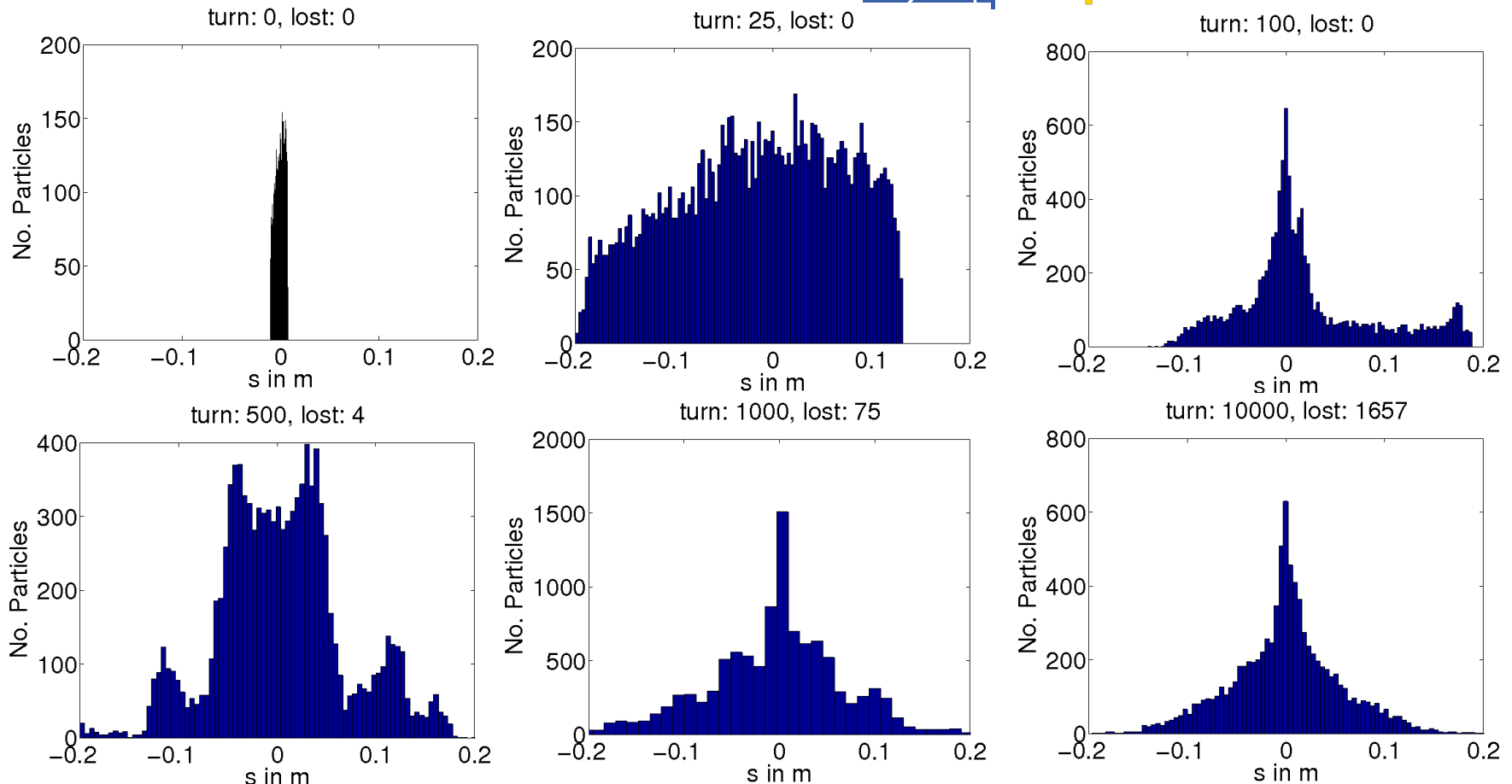
■ Bunch length increases rapidly (independent of initial length)!

[4] A.Terebilo, Accelerator Toolbox for MATLAB, SLAC-PUB-9732, 2001

# Longitudinal Results II / II



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- Density profiles could still lead to interesting radiation properties.
- Has to be investigated further.

# Momentum Compaction



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- The momentum compaction factor is defined as:

$$\alpha_c = 1/L \times \oint [D(s)/\rho(s)] ds$$

- It gives the path length difference  $\Delta L$  per revolution for a particle with an energy deviation  $\Delta p$  via the relation

$$\alpha_c \frac{\Delta p}{p_0} = \frac{\Delta L}{L}$$

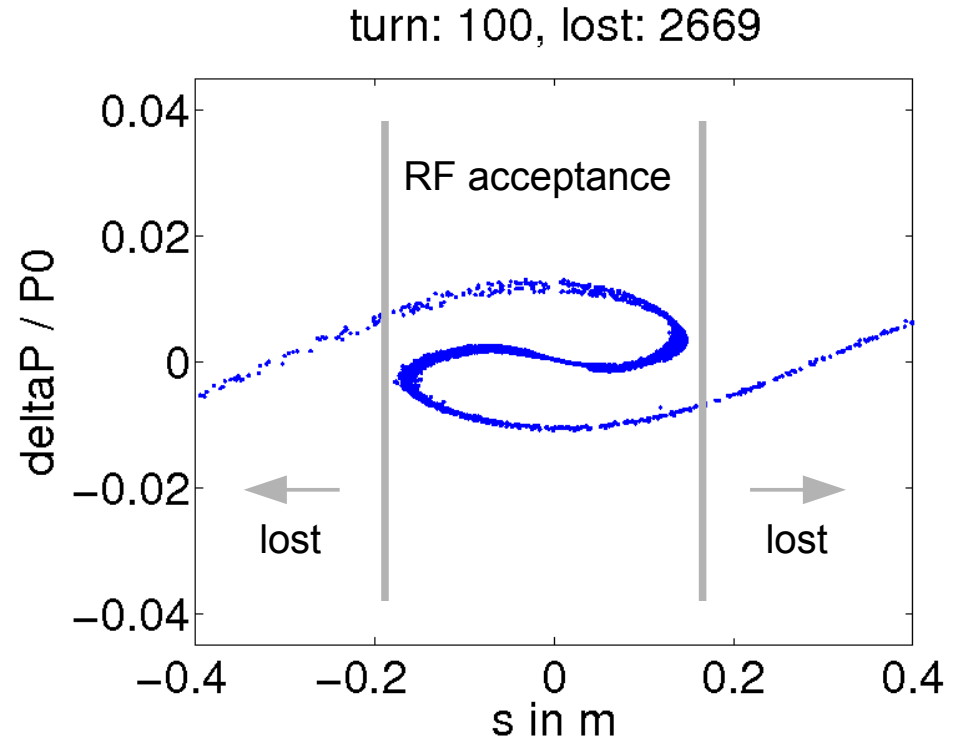
- For ANKA,  $10^{-4} < \alpha_c < 10^{-2}$ .
- The observed lengthening is consistent with analytical estimates.

- Operating at a lower  $\alpha_c$  and applying a stricter energy cut could slow the lengthening process by a factor  $\sim 1000$ .
- Assuming lengthening in the transfer line can be suppressed,  
the bunch length would still reach the length customary for Synchrotrons with dedicated low- $\alpha_c$  operation within a few 100 turns.
- Space charge and CSR (Coherent Synchrotron Radiation) effects should even accelerate this process.

# Comments II / II



- Only longitudinal losses considered so far, see example to the right for  $p_0 \pm 3\%$ .
- An increase of energy acceptance to a few percent seems feasible.



Max. energy deviation	100 %	10 %	5 %	3 %	0.25 %
Bunch charge $q$ in pC	24000	370	270	240	61
Number of particles $N$	$1.5 \cdot 10^{11}$	$2.3 \cdot 10^9$	$1.7 \cdot 10^9$	$1.5 \cdot 10^9$	$3.8 \cdot 10^8$

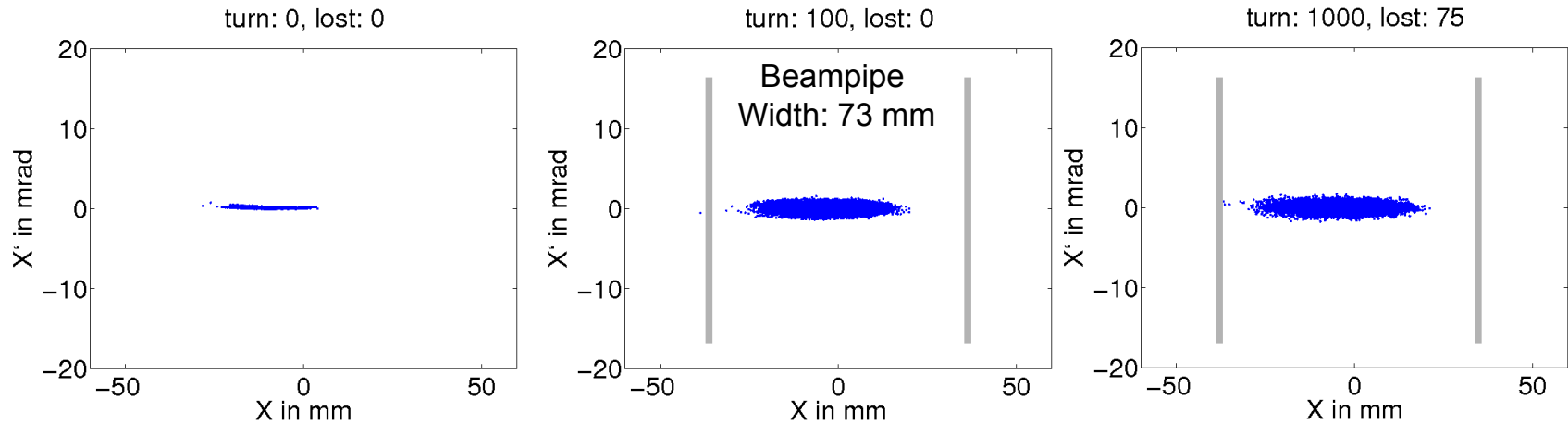
# Transverse Results



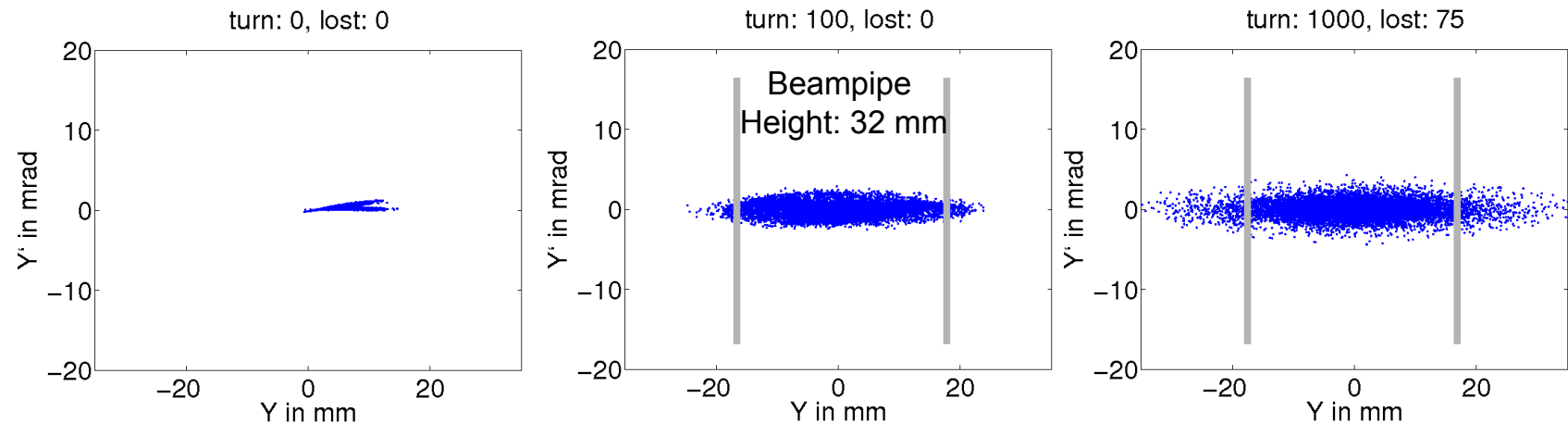
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Horizontal:



Vertical:



- Losses to the beam pipe have not yet been considered.
- They would soon become a major contribution.

# Outline



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# Summary



- The possibility to inject LWFA generated bunches into a Synchrotron has been studied.
- Preserving their ultra short length seems challenging.
- Evolution of bunch density profile could still lead to interesting radiation properties, but needs to be investigated further.



**Fine**



# Thank you for your attention!

For more information, please see:

Steffen Hillenbrand *et al.*,

"Study of Laser Wakefield Accelerators as Injectors for Synchrotron Light Sources",  
WEPEA012, Proceedings of IPAC'13



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# Backup slides

# Transfer Line I / II



- Initially round beam has to be matched to flat ANKA lattice parameters (Tab. 3).
- Note the 5 orders of magnitude difference in the  $\beta$ -functions (beam size)!
- The matching has been performed using MAD-X [2].
- The large initial divergence made the use of pulses quadrupole magnets [3] necessary.

Table 3: ANKA Twiss Parameters at Injection Point

Horizontal beta $\beta_x$ in m	16.6
Vertical beta $\beta_y$ in m	6.5
Horizontal alpha $\alpha_x$	-0.03
Vertical alpha $\alpha_y$	-0.07



[2] W. Herr and F. Schmidt. A MAD-X primer. CERN-AB-2004-027-ABP, 2004.

[3] M. Winkler *et al.*, Development and test of iron-free quadrupole lenses with high magnetic flux densities, Nucl. Inst. B, 2003 DOI: 10.1016/S0168-583X(02)02120-1