

# Implications of an R-scan for Charm Physics

Johann H. Kühn



## I. CHARM and BOTTOM MASSES

- $m_Q$  from Sum Rules
- Experimental Analysis:  $m_c$

## II. Resonant Production of $\chi_{c1}$ and $\chi_{c2}$

- Concept
- Phenomenology
- Luminosity Requirements

## III. SUMMARY

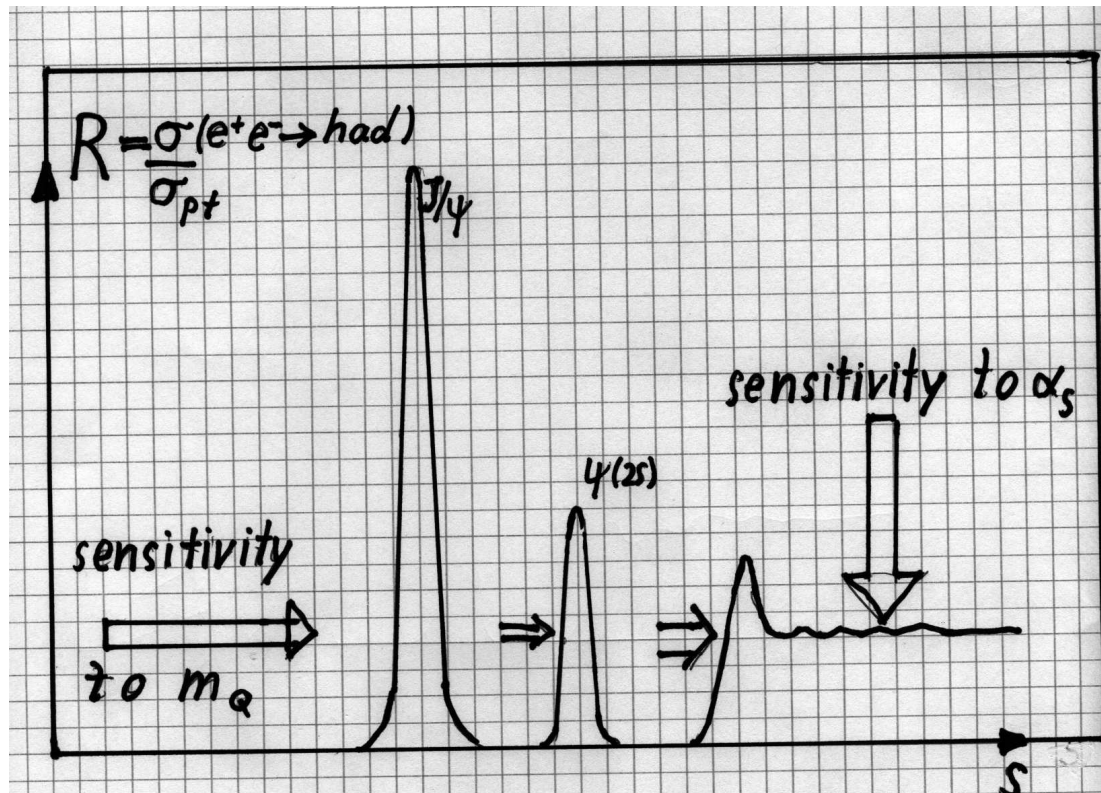
# **I. CHARM and BOTTOM MASSES**

**in collaboration with**

**K. Chetyrkin, Y. Kiyo, A. Maier, P. Maierhöfer, P. Marquard,  
A. Smirnov, M. Steinhauser, C. Sturm**

# I. 1. $m_Q$ from SVZ Sum Rules, Moments and Tadpoles

## Main Idea (SVZ)



Sensitivity to  $m_Q$  from location of  $Q\bar{Q}$  threshold.

Some definitions:

$$\left(-q^2 g_{\mu\nu} + q_\mu q_\nu\right) \Pi(q^2) \equiv i \int dx e^{iqx} \langle T j_\mu(x) j_\nu(0) \rangle$$

with the electromagnetic current  $j_\mu$ .

$$R(s) = 12\pi \text{Im} \left[ \Pi(q^2 = s + i\epsilon) \right]$$

Taylor expansion:  $\Pi_Q(q^2) = Q_Q^2 \frac{3}{16\pi^2} \sum_{n \geq 0} \bar{C}_n z^n$

with  $z = q^2 / (4m_Q^2)$  and  $m_Q = m_Q(\mu)$  the  $\overline{\text{MS}}$  mass.

$$\bar{C}_n = \bar{C}_n^{(0)} + \frac{\alpha_s}{\pi} \bar{C}_n^{(1)} + \left(\frac{\alpha_s}{\pi}\right)^2 \bar{C}_n^{(2)} + \left(\frac{\alpha_s}{\pi}\right)^3 \bar{C}_n^{(3)} + \dots$$

generic form

$$\begin{aligned}\bar{C}_n = & \bar{C}_n^{(0)} \\ & + \frac{\alpha_s}{\pi} \left( \bar{C}_n^{(10)} + \bar{C}_n^{(11)} l_{m_c} \right) \\ & + \left( \frac{\alpha_s}{\pi} \right)^2 \left( \bar{C}_n^{(20)} + \bar{C}_n^{(21)} l_{m_c} + \bar{C}_n^{(22)} l_{m_c}^2 \right) \\ & + \left( \frac{\alpha_s}{\pi} \right)^3 \left( \bar{C}_n^{(30)} + \bar{C}_n^{(31)} l_{m_c} + \bar{C}_n^{(32)} l_{m_c}^2 + \bar{C}_n^{(33)} l_{m_c}^3 \right) \\ & + \dots\end{aligned}$$

with  $\alpha_s = \alpha_s(\mu)$ ,  $l_{m_c} = \ln \left( \frac{m_c^2}{\mu^2} \right)$

$\bar{C}_n^{(ij)}$  = pure numbers

for  $j \geq 1$  from RG

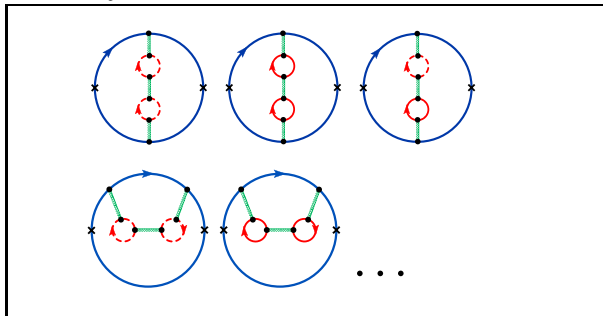
for  $j = 0$  : calculation

# Analysis in N<sup>3</sup>LO

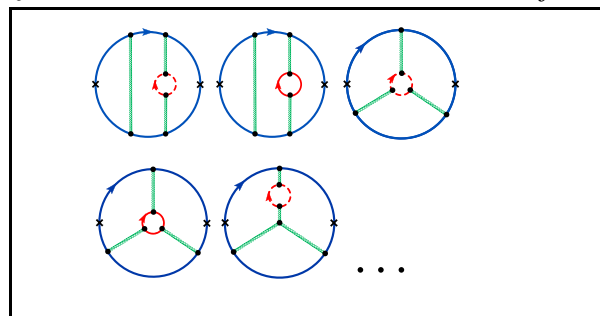
Algebraic reduction to 13 master integrals (Laporta algorithm);

numerical and analytical evaluation of master integrals

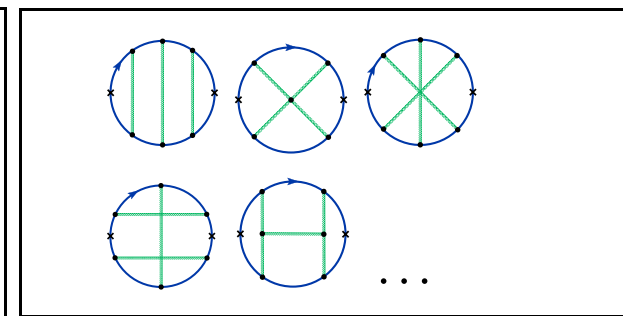
$n_f^2$ -contributions



$n_f^1$ -contributions



$n_f^0$ -contributions



 : heavy quarks,  : light quarks,

$n_f$ : number of active quarks

⇒ About **700 Feynman-diagrams**

$\bar{C}_0$  and  $\bar{C}_1$  in order  $\alpha_s^3$  (four loops!) (2006)

Reduction to master integrals

(Chetyrkin, JK, Sturm; Boughezal, Czakon, Schutzmeier)

$\bar{C}_2$  and  $\bar{C}_3$  (2008)

(Maier, Maierhöfer, Marquard, A. Smirnov)

All master integrals known analytically and double checked.

(Schröder + Vuorinen, Chetyrkin et al., Schröder + Steinhauser,

Laporta, Broadhurst, Kniehl et al.)

$\bar{C}_4 - \bar{C}_{10}$ : extension to higher moments by Padé method, using analytic information from low energy ( $q^2 = 0$ ), threshold ( $q^2 = 4m^2$ ), high energy ( $q^2 = -\infty$ ) (Kiyoyama, Maier, Maierhöfer, Marquard, 2009)



## Relation to measurements

$$\mathcal{M}_n^{\text{th}} \equiv \frac{12\pi^2}{n!} \left( \frac{d}{dq^2} \right)^n \Pi_c(q^2) \Big|_{q^2=0} = \frac{9}{4} Q_c^2 \left( \frac{1}{4m_c^2} \right)^n \bar{C}_n$$

Perturbation theory:  $\bar{C}_n$  is function of  $\alpha_s$  and  $\ln \frac{m_c^2}{\mu^2}$

dispersion relation:

$$\Pi_c(q^2) = \frac{q^2}{12\pi^2} \int ds \frac{R_c(s)}{s(s-q^2)} + \text{subtraction}$$

$$\Rightarrow \mathcal{M}_n^{\text{exp}} = \int \frac{ds}{s^{n+1}} R_c(s)$$

constraint:  $\mathcal{M}_n^{\text{exp}} = \mathcal{M}_n^{\text{th}}$

$$\Rightarrow m_c = \frac{1}{2} \left( \frac{9}{4} Q_c^2 \bar{C}_n / \mathcal{M}_n^{\text{exp}} \right)^{1/(2n)}$$

## qualitative considerations

$$\mathcal{M}_n = \int_{\text{threshold}} \frac{ds}{s^{n+1}} R_c(s) \sim \text{dimension } (m_c)^{-2n}$$

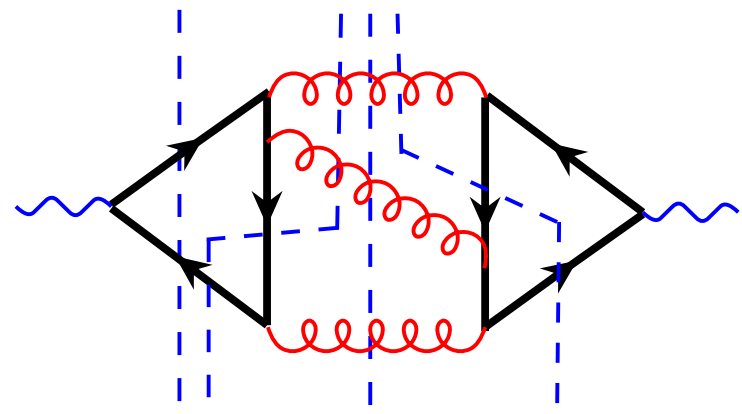
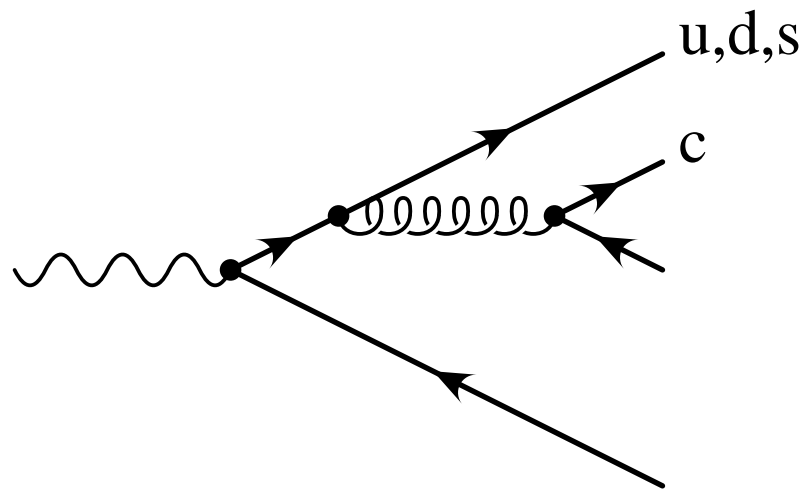
- depends moderately on  $\alpha_s$ !
- $\Pi(q^2)$  is an analytic function with branch cut from  $(2m_D)^2$  to  $\infty$ .
- averaging over resonances reduces influence of long distances (binding effects).
- $\Pi(q^2 = 0)$  and its derivatives at  $q^2 = 0$  are short distance quantities.  
 $\Rightarrow$  pQCD is applicable.

## I. 2. Experimental Analysis: $m_c$

$$\mathcal{M}_n^{\text{exp}} \equiv \int \frac{ds}{s^{n+1}} R_{\text{charm}}(s)$$

$$\Rightarrow m_c = \frac{1}{2} \left( \frac{9}{4} Q_c^2 \bar{C}_n / \mathcal{M}_n^{\text{exp}} \right)^{1/(2n)}$$

- $\bar{C}_n$  from calculation
- $\mathcal{M}_n^{\text{exp}}$  from experiment



## Ingredients (charm)

### experiment:

- $\Gamma_e(J/\psi, \psi')$  from BES & CLEO & BABAR (PDG)
- $\psi(3770)$  and  $R(s)$  from BES
- $\alpha_s = 0.1187 \pm 0.0020$

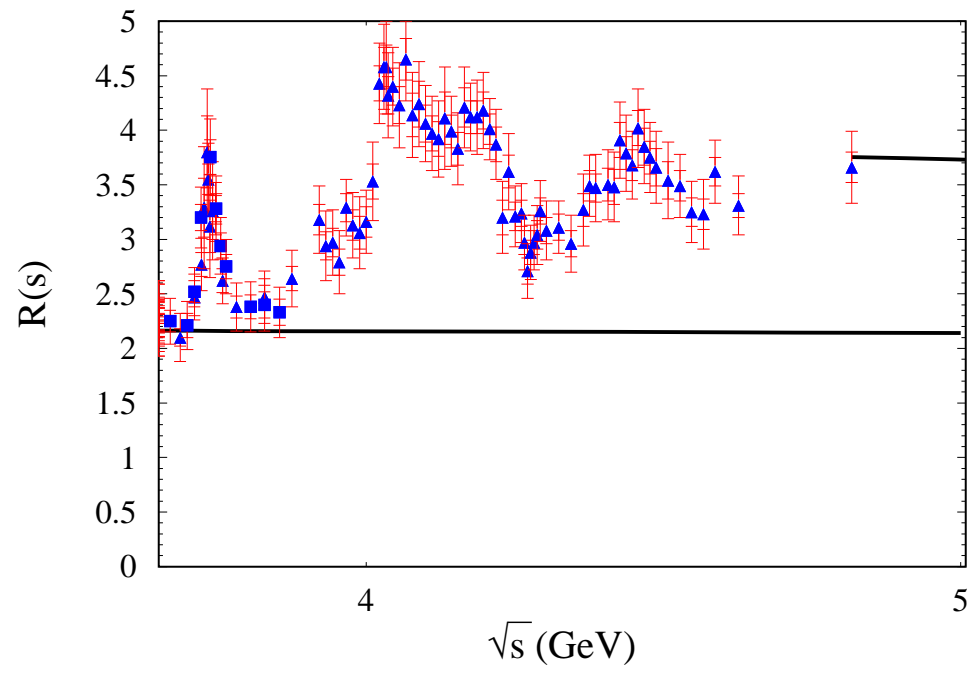
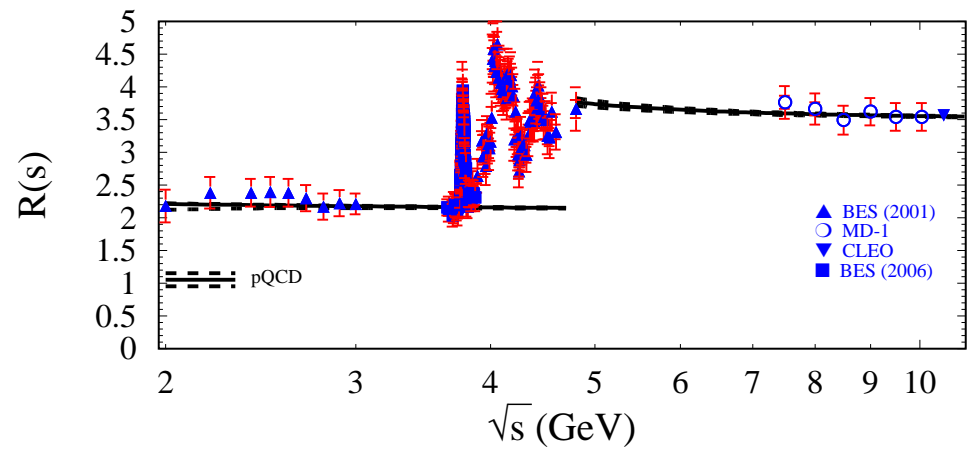
### theory:

- N<sup>3</sup>LO for  $n = 1, 2, 3, 4$
- include condensates

$$\delta\mathcal{M}_n^{\text{np}} = \frac{12\pi^2 Q_c^2}{(4m_c^2)^{(n+2)}} \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle a_n$$

(including NLO-terms)

- estimate of non-perturbative terms  
(oscillations, based on [Shifman](#))
- careful extrapolation of  $R_{uds}$
- careful definition of  $R_c$



Contributions from

- narrow resonances:  $R = \frac{9\Pi M_R \Gamma_e}{\alpha^2(s)} \delta(s - M_R^2)$  (PDG)
- threshold region ( $2 m_D - 4.8 \text{ GeV}$ ) (BESS)
- perturbative continuum ( $E \geq 4.8 \text{ GeV}$ ) (Theory)

$n$	$\mathcal{M}_n^{\text{res}}$ $\times 10^{(n-1)}$	$\mathcal{M}_n^{\text{thresh}}$ $\times 10^{(n-1)}$	$\mathcal{M}_n^{\text{cont}}$ $\times 10^{(n-1)}$	$\mathcal{M}_n^{\text{exp}}$ $\times 10^{(n-1)}$	$\mathcal{M}_n^{\text{np}}$ $\times 10^{(n-1)}$
1	0.1201(25)	0.0318(15)	0.0646(11)	0.2166(31)	-0.0001(2)
2	0.1176(25)	0.0178(8)	0.0144(3)	0.1497(27)	0.0000(0)
3	0.1169(26)	0.0101(5)	0.0042(1)	0.1312(27)	0.0007(14)
4	0.1177(27)	0.0058(3)	0.0014(0)	0.1249(27)	0.0027(54)

Different relative importance of resonances vs. continuum for  $n = 1, 2, 3, 4$ .

## Results ( $m_c$ )

PRD 80: (2009) 074010

$n$	$m_c(3\text{ GeV})$	exp	$\alpha_s$	$\mu$	np	total
1	986	9	9	2	1	13
2	976	6	14	5	0	16
3	978	5	15	7	2	17
4	1004	3	9	31	7	33

Remarkable consistency between  $n = 1, 2, 3, 4$

and stability ( $O(\alpha_s^2)$  vs.  $O(\alpha_s^3)$ );

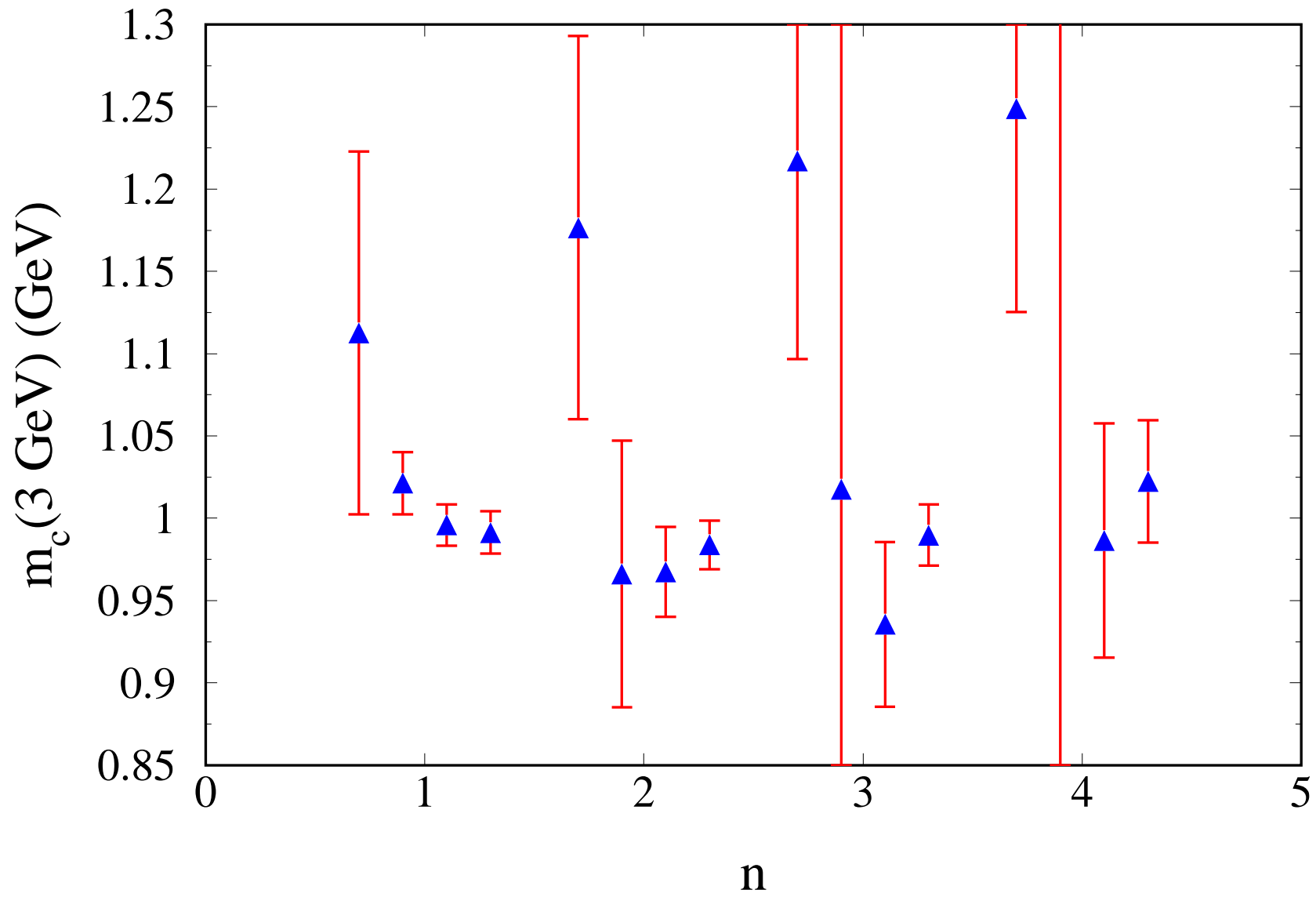
preferred scale:  $\mu = 3\text{ GeV}$ ,

conversion to  $m_c(m_c)$ :

- $m_c(3\text{ GeV}) = 986 \pm 13\text{ MeV}$

- $m_c(m_c) = 1279 \pm 13\text{ MeV}$



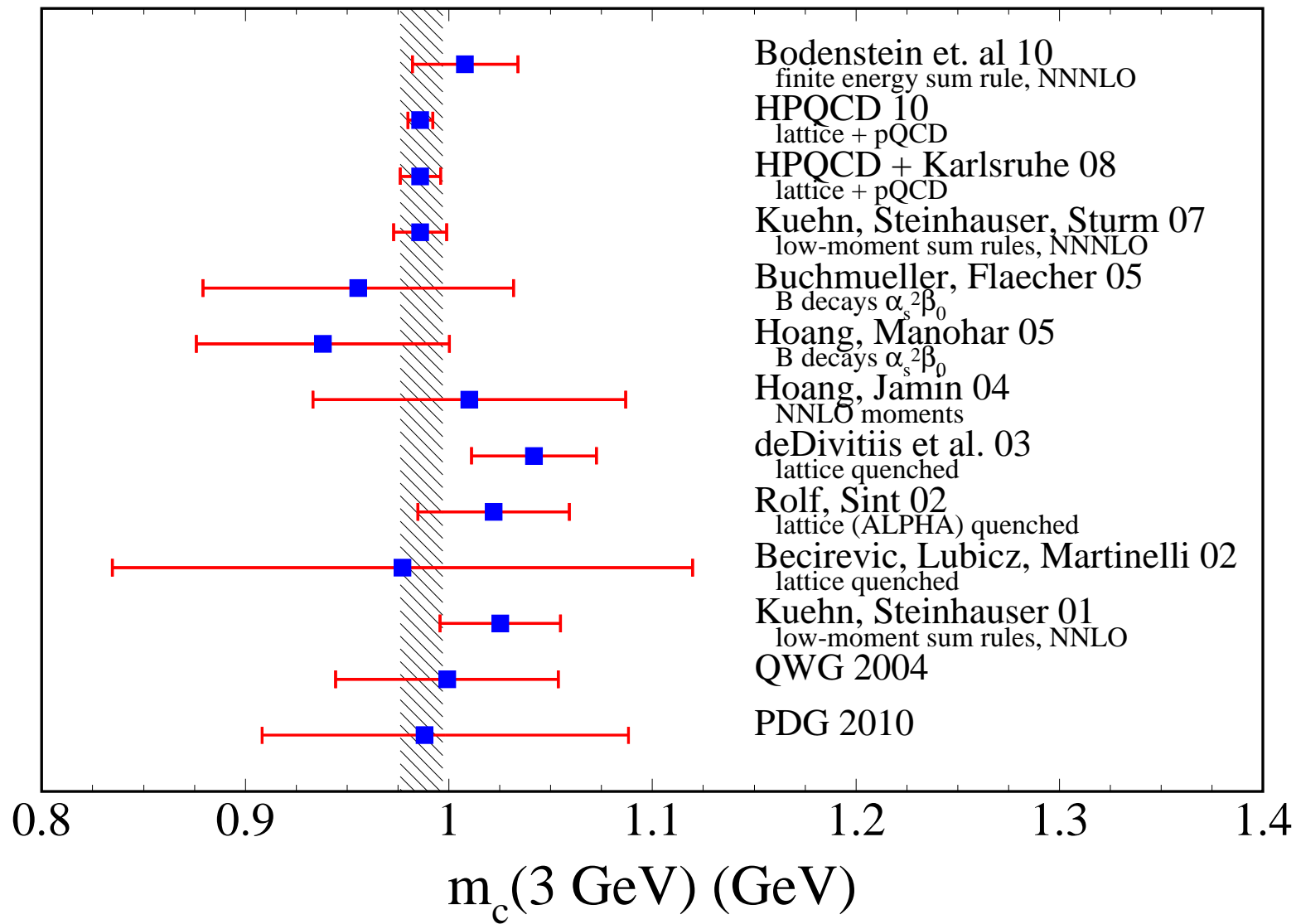


## Perturbative stability

$$\begin{aligned}
 m_c &= \frac{1}{2} \left( \frac{9Q_c^2 \bar{C}_n^{\text{Born}}}{4 \mathcal{M}_n^{\text{exp}}} \right)^{\frac{1}{2n}} (1 + r_n^{(1)} \alpha_s + r_n^{(2)} \alpha_s^2 + r_n^{(3)} \alpha_s^3) \\
 &\propto 1 - \begin{pmatrix} 0.328 \\ 0.524 \\ 0.618 \\ 0.662 \end{pmatrix} \alpha_s - \begin{pmatrix} 0.306 \\ 0.409 \\ 0.510 \\ 0.575 \end{pmatrix} \alpha_s^2 - \begin{pmatrix} 0.262 \\ 0.230 \\ 0.299 \\ 0.396 \end{pmatrix} \alpha_s^3, \tag{1}
 \end{aligned}$$

error from next order  $\leq r_n^{\text{max}} \alpha_s^4 < 2$  to 3 permille

(smaller than  $\mu$ -variation!)



## Potential improvements in analysis

1) Combined fit to lowest three moments

⇒ optimal usage of experimental information,

⇒ requires dedicated analysis of correlation.

2) More “aggressive” choice for  $\delta\alpha_s$ :

$$\alpha_s = 0.1189 \pm 0.002 \quad \Rightarrow \quad \alpha_s = 0.1184 \pm 0.0007 \text{ (PDG 2012)}$$

3) Theory error from perturbation series (2 – 3 permille)

instead of  $\mu$ -variation

⇒  $\delta m_c$  below 10 MeV

## Potential experimental improvements

1.) narrow resonances ( $J/\Psi, \Psi'$ ) dominate:

$$\Gamma_e(1S) = 5.55 \pm 0.14 \pm 0.02 \text{ keV}; \quad \Gamma_e(2S) = 2.35 \pm 0.04 \text{ keV}$$

improvement? correlations?

2.) threshold region: improvement at BESS?

subtraction of  $u\bar{u}, d\bar{d}, s\bar{s} \Rightarrow$  precise reference point below charm threshold needed

3.) continuum above 4.8 GeV:

missing data substituted by theory  $\Rightarrow$  small error

excellent agreement with data

improved calibration at  $\sim 4.5$  GeV desirable

$\rightarrow$  crucial information over wide range

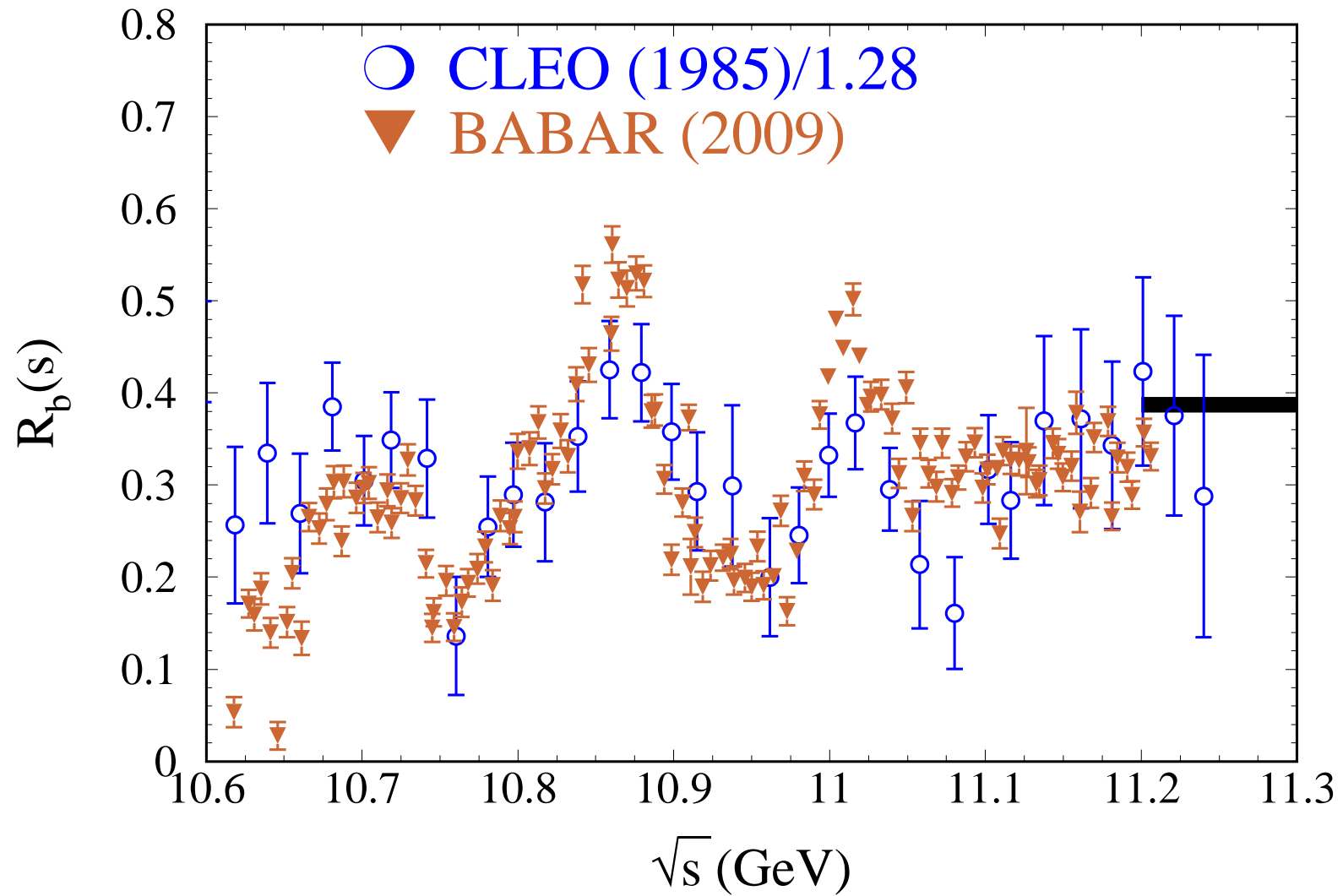
In total  $\delta m_c \approx 7$  MeV may be within reach

## Experimental Ingredients for $m_b$

Contributions from

- narrow resonances ( $\Upsilon(1S) - \Upsilon(4S)$ ) (PDG)
- threshold region (10.618 GeV – 11.2 GeV) (BABAR 2009)
- perturbative continuum ( $E \geq 11.2$  GeV) (Theory)
- different relative importance of resonances vs. continuum for  $n = 1, 2, 3, 4$

$n$	$\mathcal{M}_n^{\text{res},(1S-4S)}$ $\times 10^{(2n+1)}$	$\mathcal{M}_n^{\text{thresh}}$ $\times 10^{(2n+1)}$	$\mathcal{M}_n^{\text{cont}}$ $\times 10^{(2n+1)}$	$\mathcal{M}_n^{\text{exp}}$ $\times 10^{(2n+1)}$
1	1.394(23)	0.287(12)	2.911(18)	4.592(31)
2	1.459(23)	0.240(10)	1.173(11)	2.872(28)
3	1.538(24)	0.200(8)	0.624(7)	2.362(26)
4	1.630(25)	0.168(7)	0.372(5)	2.170(26)



BELLE?

$n$	$m_b(10\text{ GeV})$	exp	$\alpha_s$	$\mu$	total	$m_b(m_b)$
1	3597	14	7	2	16	4151
2	3610	10	12	3	16	4163
3	3619	8	14	6	18	4172
4	3631	6	15	20	26	4183

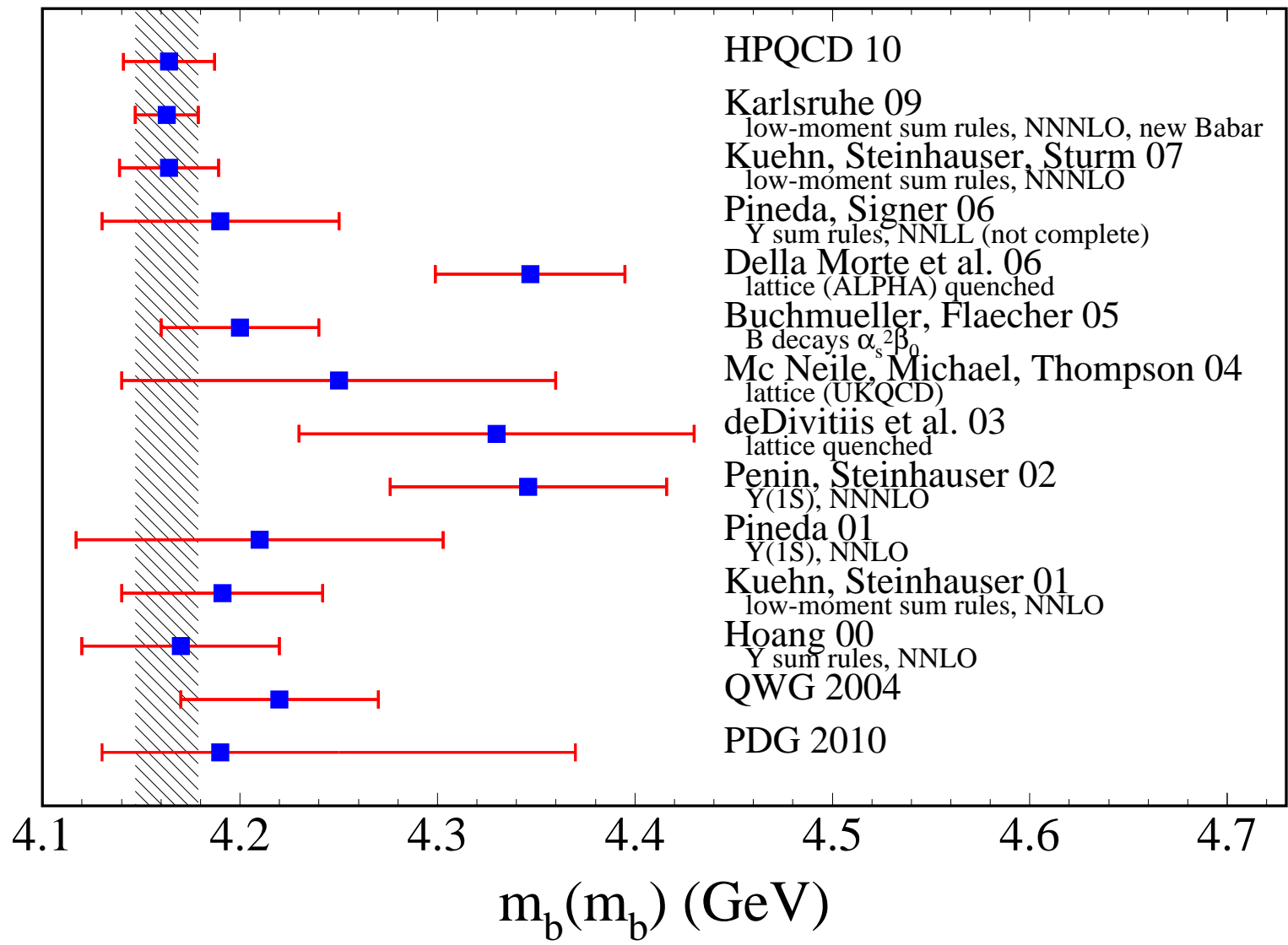
Consistency ( $n = 1, 2, 3, 4$ ) and stability ( $O(\alpha_s^2)$  vs.  $O(\alpha_s^3)$ );

- $m_b(10\text{ GeV}) = 3610 \pm 16\text{ MeV}$
- $m_b(m_b) = 4163 \pm 16\text{ MeV}$

potential improvements:

- $\delta\alpha_s$
- combined analysis of 3 lowest moments
- improved data





## $\alpha_s$ -dependence

$$m_c(3 \text{ GeV}) = \left( 986 - \frac{\alpha_s - 0.1189}{0.002} \cdot 9 \pm 10 \right) \text{ MeV}$$

$$m_b(10 \text{ GeV}) = \left( 3610 - \frac{\alpha_s - 0.1189}{0.002} \cdot 12 \pm 11 \right) \text{ MeV}$$

$$m_b(m_b) = \left( 4163 - \frac{\alpha_s - 0.1189}{0.002} \cdot 12 \pm 11 \right) \text{ MeV}$$

$$m_b(M_Z) = \left( 2835 - \frac{\alpha_s - 0.1189}{0.002} \cdot 27 \pm 8 \right) \text{ MeV}$$

$$m_b(161.8 \text{ GeV}) = \left( 2703 - \frac{\alpha_s - 0.1189}{0.002} \cdot 28 \pm 8 \right) \text{ MeV}$$

$$m_c(3 \text{ GeV}) = 986(13) \text{ MeV}$$

$$m_b(10 \text{ GeV}) = 3610(16) \text{ MeV}$$
$$m_b(m_b) = 4163(16) \text{ MeV}$$

Improvement by factor 2 conceivable

# II. Resonant Production of $\chi_{c1}$ and $\chi_{c2}$

## Concept

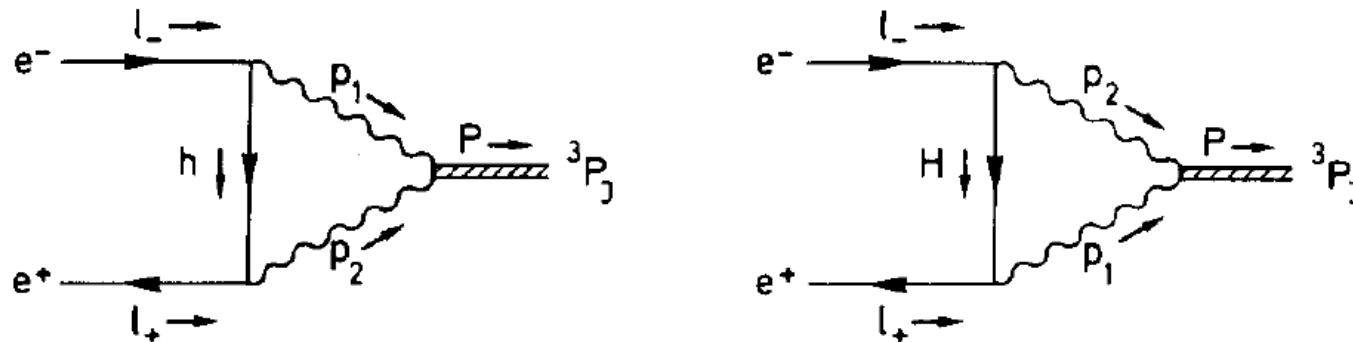
direct coupling of  $\chi_{cJ}$  to  $e^+e^-$

$$\mathcal{A}(e^+e^- \rightarrow \chi_0) = O(m_e/M_\chi) \approx 0$$

$$\mathcal{A}(e^+e^- \rightarrow \chi_1) = g_1 \bar{v} \gamma_5 \not{\epsilon} u$$

$$\mathcal{A}(e^+e^- \rightarrow \chi_2) = g_2 \bar{v} \gamma^\mu u \epsilon_{\mu\nu} (l_+^\nu - l_-^\nu) / M_\chi$$

through two virtual photons



quarkonium: short distance calculation

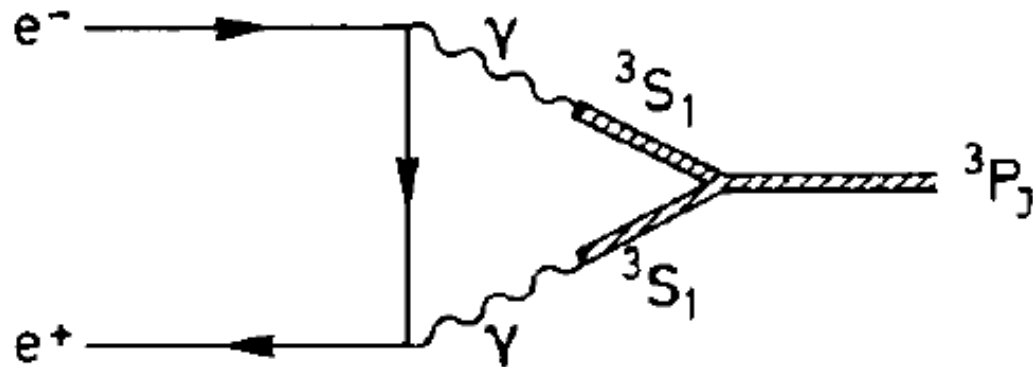
$$g_1 = -\frac{\alpha^2}{M^2} 32 a \ln \frac{2b}{M}$$

$$g_2 = +\frac{\alpha^2}{M^2} 32 a \left[ \sqrt{2} \ln \frac{2b}{M} + \frac{\sqrt{2}}{3} (i\pi + \ln 2 - 1) \right]$$

with  $M \equiv M_\chi$ ,  $a = \sqrt{\frac{3}{4\pi m_c}} \Phi'(0)$ ,

$b = \text{“binding energy”} \rightarrow b = \pm 0.5 \text{ GeV}$

vector dominance:



Results:  $\Gamma(\chi_J \rightarrow e^+e^-)$  in eV

J	1	2
quarkonium		
$b = +0.5$ GeV	0.023	0.013
$b = -0.5$ GeV	0.17	0.27
vector dominance	0.46	0.014
unitarity limit	0.044	0.0023

consider  $J = 1$  and use  $\Gamma = 0.1$  eV

## Strategy:

R-scan around  $M(\chi_1) = 3.511 \text{ GeV}$

A) all final states

$$\Rightarrow R_{\text{peak}} = \frac{9\pi}{2\alpha\sqrt{2\pi}} \frac{\Gamma_1}{\Delta M} c_{\text{rad}} \approx 0.05 \frac{\Gamma_1[\text{eV}]}{\Delta M[\text{MeV}]}$$

$\Delta M = \text{energy spread} = 1.1 \text{ MeV}; \quad c_{\text{rad}} \approx 0.5$

$$R_{\text{peak}} = 0.005 \text{ above background of } R = 2$$

( $\Rightarrow 100 \text{ pb}^{-1}$  per point with 5 points  $\Rightarrow 3$  sigma)

B) selected final state  $\chi_1 \rightarrow \gamma J/\psi (\rightarrow \mu^+ \mu^-, e^+ e^-)$

(Br  $\approx 0.35 \cdot 0.12 \approx 0.054$ )

background from  $e^+ e^- \rightarrow \gamma J/\psi$  (radiative return)

similar requirement on luminosity

## SUMMARY

precise charm and bottom quark masses can be further improved through:

- improved  $\Gamma(J/\psi \rightarrow e^+ e^-)$  etc.
- improved  $R_{\text{charm}}$  and  $R_{\text{bottom}}$
- improved  $\alpha_s$
- combined analysis of lowest three moments

direct resonant production of  $\chi_{1c}$  is possible

- through a dedicated scan around 3.511 GeV
- perhaps through radiative return  
→ improved insight in charmonium dynamics