

Efficient turbulent amplification of magnetic field driven by dynamo effect at supernova remnant shocks

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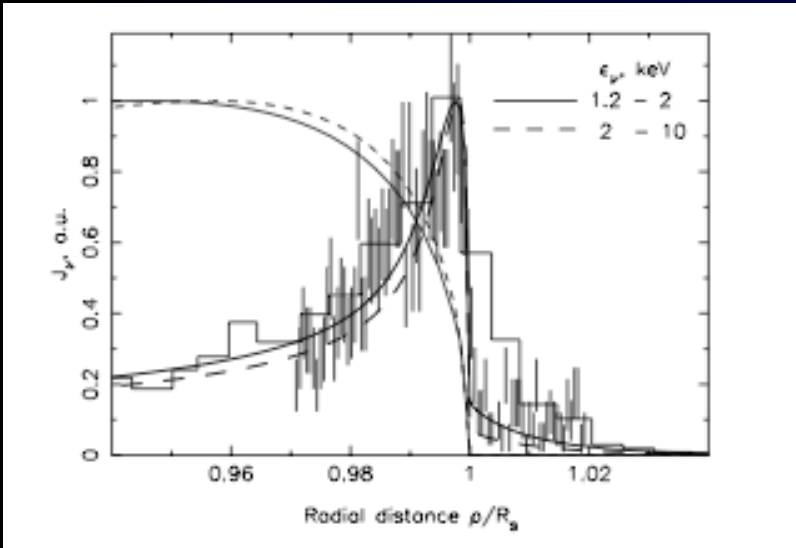
Outline

Multiwavelength evidence of amplified magnetic field

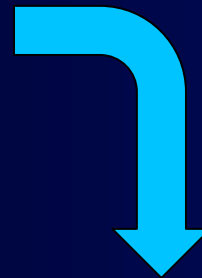
MHD approach to 2D rippled shocks

Secular evolution of magnetic energy at the shock

Evidence of large B in supernova remnant shocks

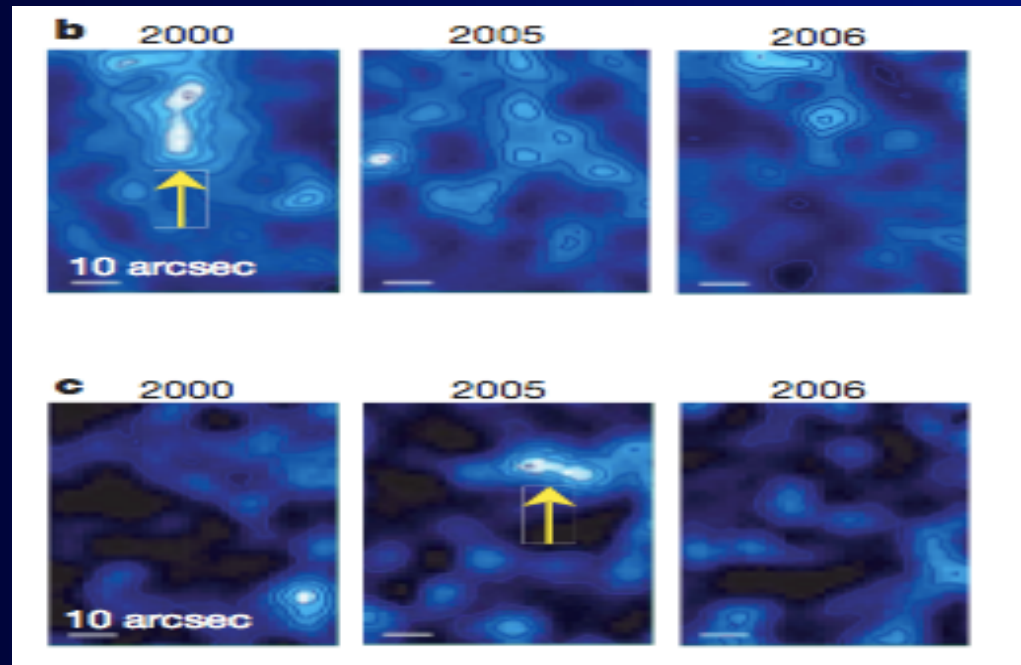
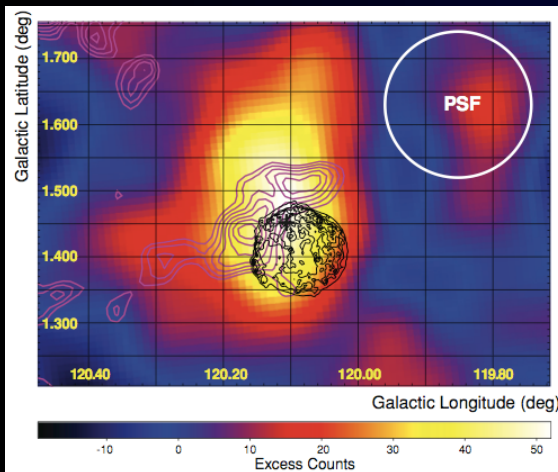


Sharp X-ray edges
(Berezhko et al. 2003)



Time variability of synchrotron
emission from filaments
~0.03 pc
(Uchiyama 2007)

..and energetic particles?



Acciari et al. 2011

Real shocks are corrugated

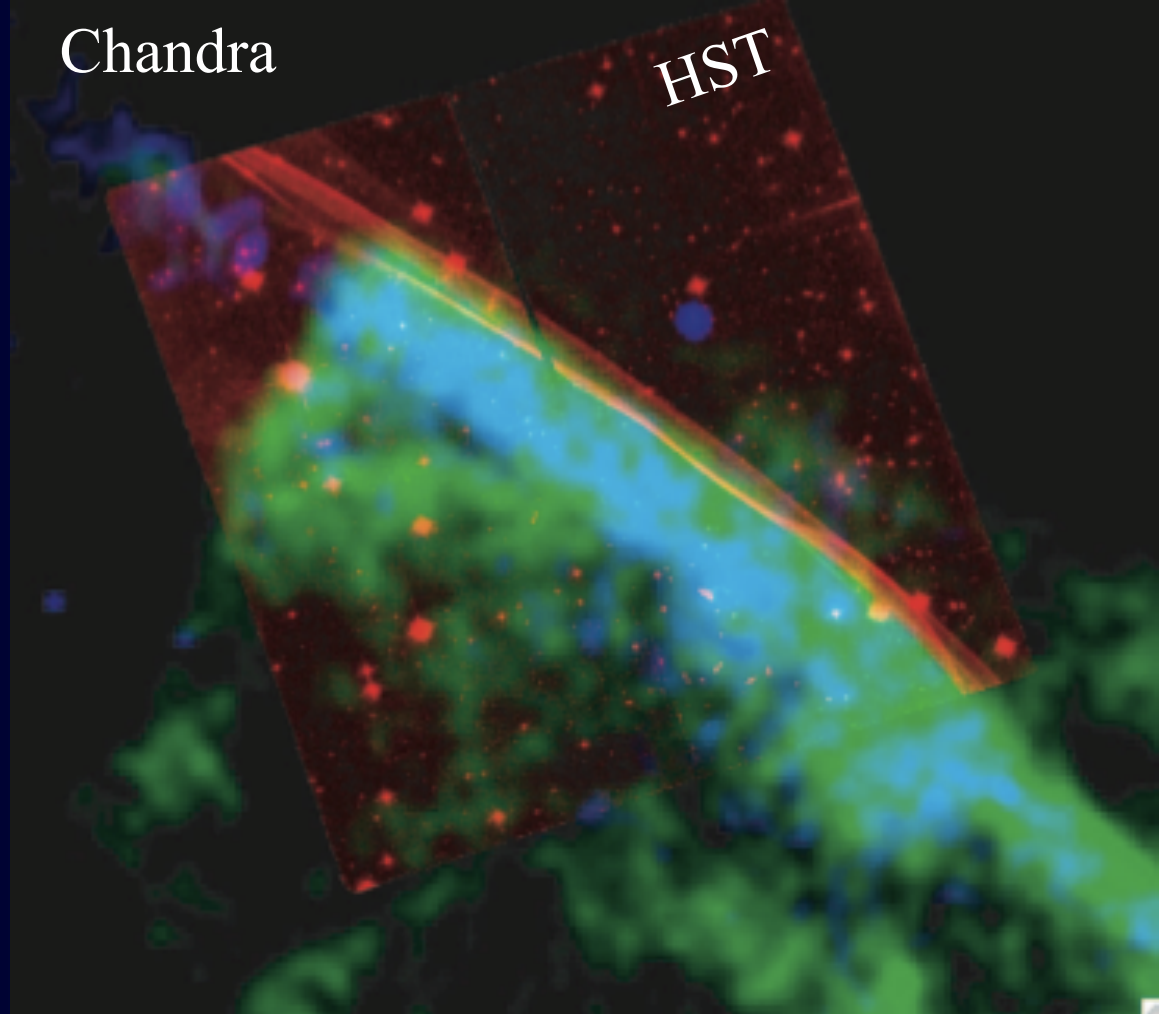
SNR 1006

Chandra

HST

shock ripples $\sim 10^{17}$ cm

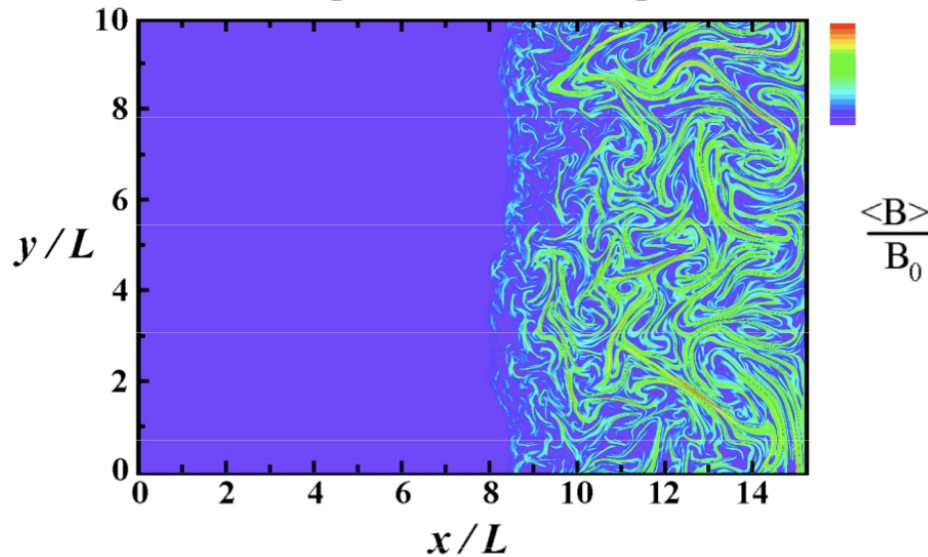
Smaller for SNR
closer to galactic plane



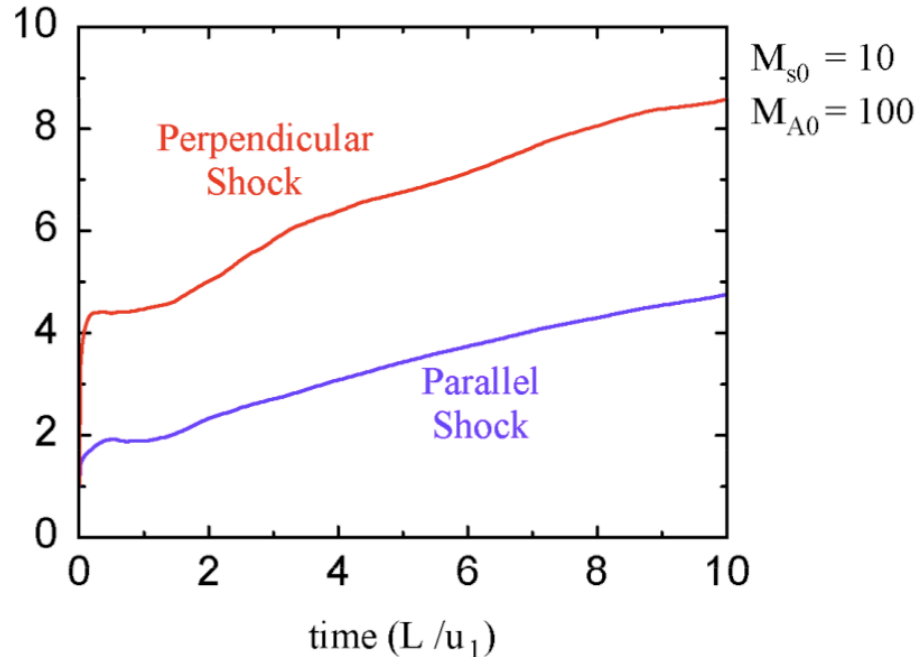
Raymond et al. (2007)

Shocks: numerical simulations

Magnetic Field Strength



Mean Downstream Field Strength



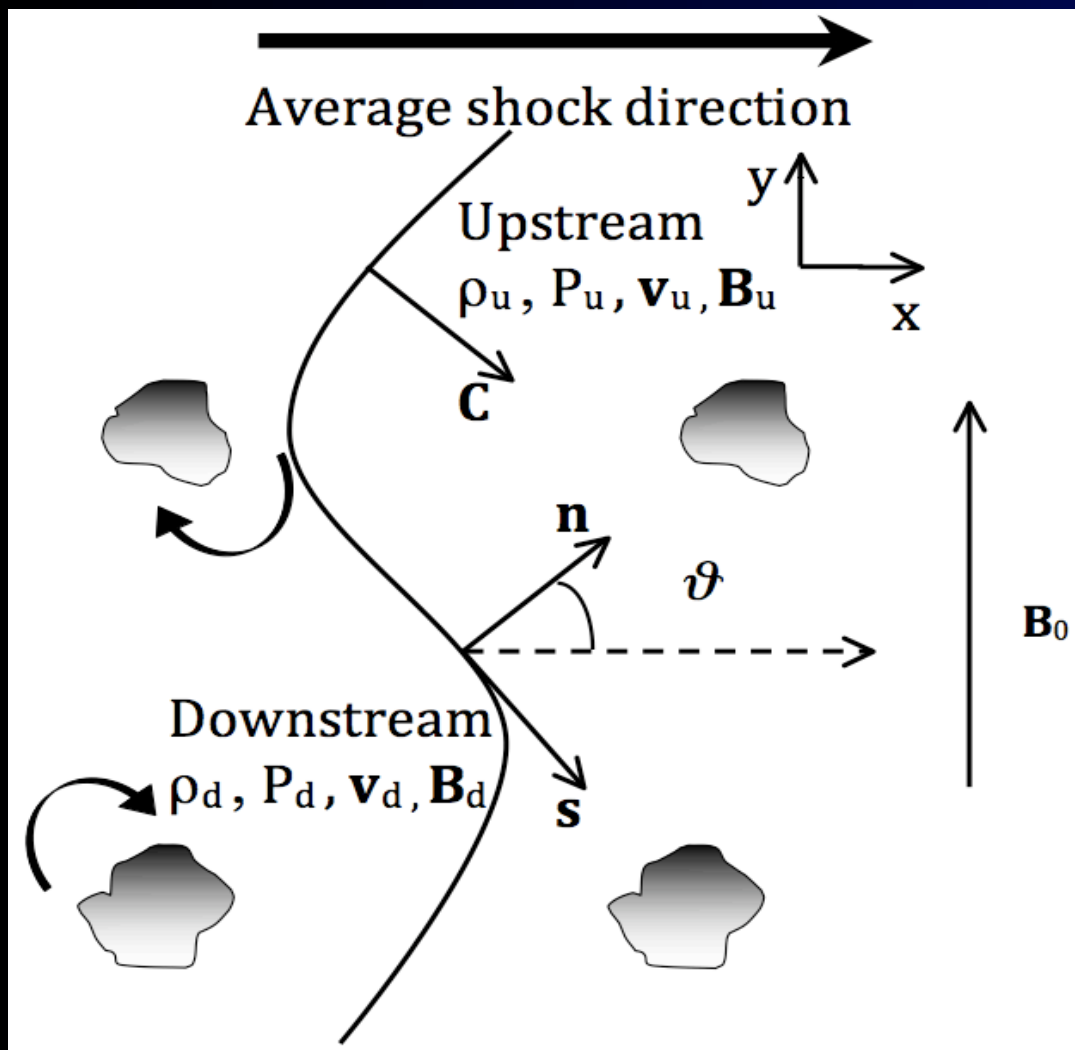
Significant amplification
of magnetic energy
(Giacalone Jokipii 2007)

Evidence of efficient magnetic field amplification downstream
of **gamma-ray burst** outflows (numerical simulations, Mizuno et al. 2011)

MHD analytical approach to rippled shock

2D shock, arbitrary Alfvén Mach number

Jump conditions can be applied locally at rippled shocks



Fraschetti (2013)

ω downstream at distance smaller than local curvature radius

$$\mathbf{v} = (v_x, v_y, 0)$$

$$\boldsymbol{\omega} = \nabla \times \mathbf{v} = (0, 0, \omega_z)$$

Local orthonormal frame:

$$(x, y) \longrightarrow (n, s)$$

Prasad (2001)

To first order in θ
(large curvature radius)

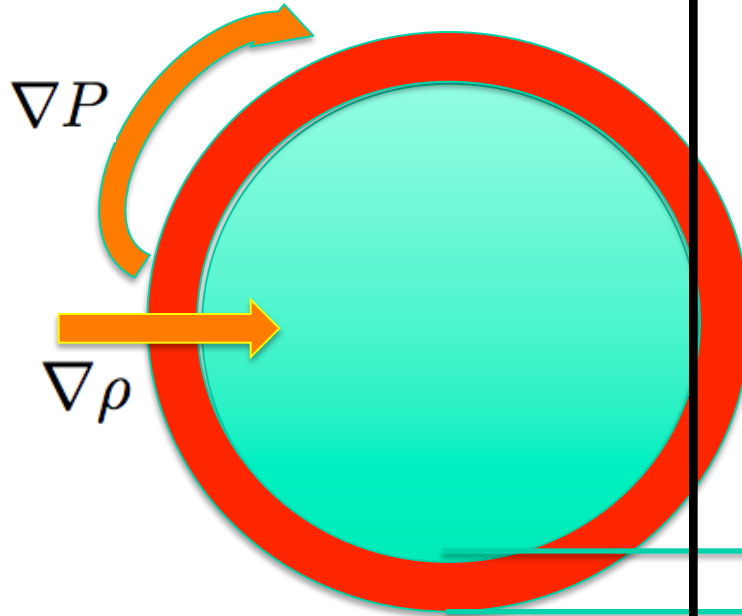
Perpendicular seed field
upstream

Average shock direction of motion



up

down



ℓ_F =Field length

Shock

$$\frac{\nabla \rho \times \nabla P}{\rho^2} \text{ non-zero}$$

Vorticity downstream

$$\delta\omega_z = \frac{r-1}{r} \left[\left(\frac{C_r}{\rho} \right)_u \partial_s \rho + \partial_s C_r \right] - \frac{B_n \delta B_s}{4\pi \rho C_r} \partial_s \vartheta$$

r: compression
 C_r : shock speed

$$M_A = C_r \sqrt{4\pi \rho} / B_0$$

Fraschetti (2013)

Energy deposited in vortical motion grows with shock speed
 Shear or power spectrum

Finite curvature radius (zero for planar shock)

Turbulent field backreaction
 Strongly rippled \rightarrow higher amplified field
 Large B_0 makes resistance to field lines tangling

growth rate

backreaction

Small-scale dynamo


At each scale, the growth of magnetic field depends on vorticity
 No assumption on the magnetic power spectrum

Exponential growth

$$\frac{d\varepsilon}{dt} = 2\beta\varepsilon$$

$$\varepsilon = B^2 / 8\pi\rho$$

Kulsrud, 2005



$$\frac{d\varepsilon}{dt} = 2(\tau^{-1} - \alpha\varepsilon)\varepsilon$$

Growth time-scale

$$\tau \sim \frac{\tau}{r-1} \frac{1}{C_r} \frac{R_c \ell_F}{R_c + \ell_F} \quad \tau \sim \ell_F / C_r$$

Back-reaction

$$\alpha \sim \vartheta / (R_c C_r)$$

Exact solution

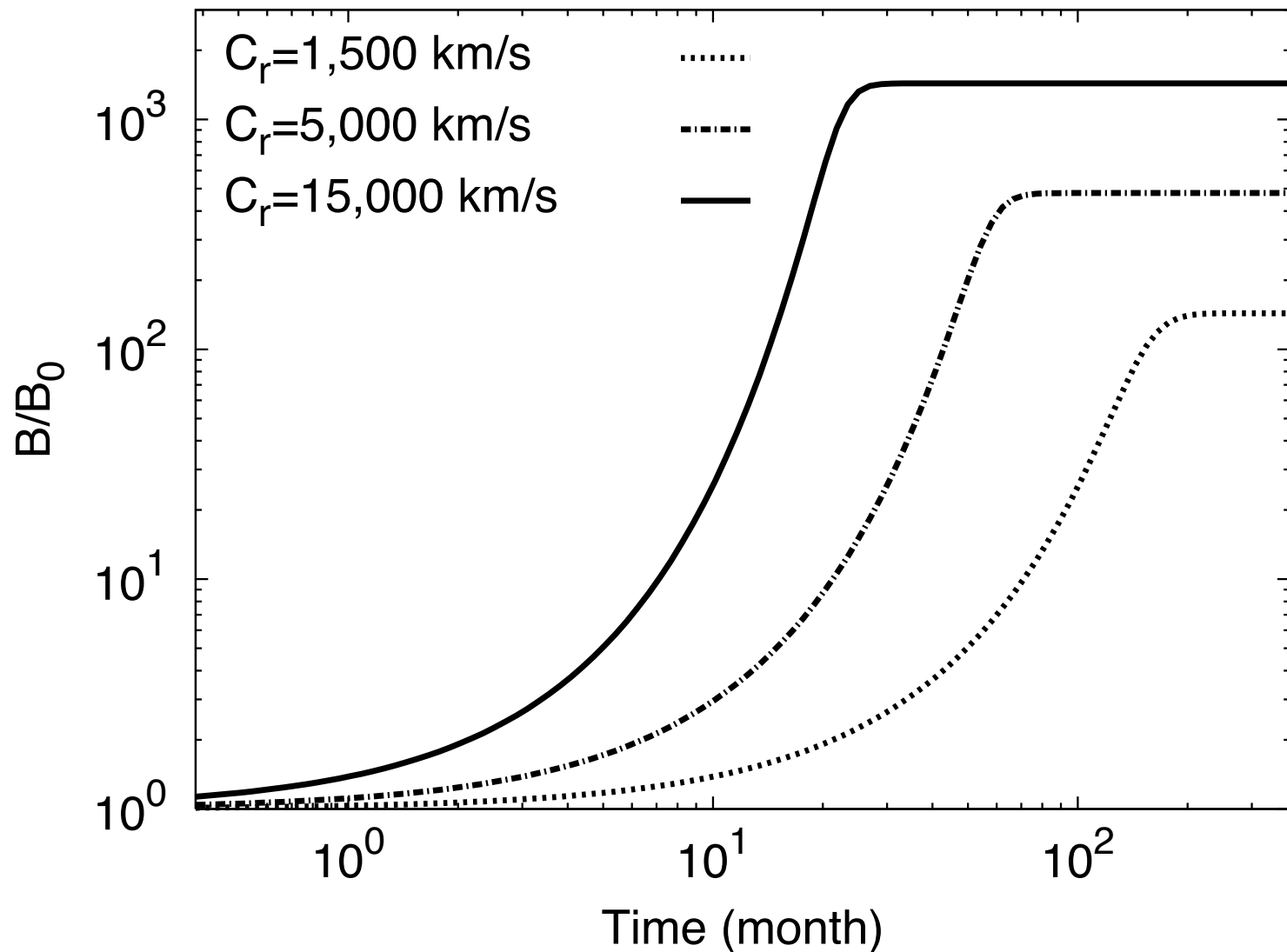
$$\left(\frac{B}{B_0}\right)^2(t) = \frac{e^{2t/\tau}}{1 - \alpha\tau(1 - e^{2t/\tau})v_A^2/2}$$

Fraschetti (2013)

Saturation

$$t \gg \tau$$

$$B/B_0 \sim M_A$$



Ripples scale $= 10^{17}$ cm, rapid growth
(Raymond et al. 2007)

Bohm diffusion?

Energetic electrons

$$D_B(E) \simeq (3.3 \times 10^{23} \text{ cm}^2 \text{ s}^{-1}) \times E_{\text{TeV}} \times B_{100}^{-1}$$

$$\tau_{\text{acc}} \simeq D(E)/V_{\text{sh}}^2 \quad (\text{energy losses})$$

Field growth

$$\tau \sim \ell_F / C_r$$

Summary

Turbulent vortical motions are an efficient tracker of strong magnetic field in astrophysical shocks

Real shocks in nature are rippled but MHD jump conditions can still be applied locally

Magnetic field amplification with saturation as short as a few months agrees with observations (X, optical)

Prediction of turbulence evolution on secular scale