Vanilla New Physics facing the Higgs boson

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July 4th, celebration day



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Next step: what's the shape of the Higgs?

Some interesting hints that something may be going on:



Consistent with SM at 95% C.L., but...

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Fits to σ/σ_{SM} and BR/BR_{SM}

Many fits roughly in agreement

(Corbett et al., Giardino et al., Espinosa et al., Plehn et al, Montull et al.)



Corbett et al., 2012

Everything compatible with SM, but best fit points hint to:



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Key quantities:

where

$$R_{\gamma\gamma} = \frac{\Gamma(h \to \gamma\gamma)}{\Gamma(h \to \gamma\gamma)_{SM}} \simeq \left| 1 + \sum_{i} \delta_{C}^{i} + \sum_{j} \delta_{NC}^{j} \right|^{2}$$

$$R_{GG} = \frac{\sigma(GG \to h)}{\sigma(GG \to h)_{SM}} \simeq \left| 1 - 9.7 \sum_{i} \frac{t_{c}^{i}}{N_{c}^{i} Q_{i}^{2}} \delta_{C}^{i} \right|^{2}$$

$$\delta_{C,NC}^{i} = \frac{N_{c}^{i} Q_{i}^{2} \frac{2g_{hii}}{m_{i}} A_{i}^{i} \left(\frac{m_{h}^{2}}{4m_{i}^{2}}\right)}{A_{W} + A_{t}}$$
Loop function

$$\delta_{C,NC}^{i} = \frac{N_{c}^{i} Q_{i}^{2} \frac{2g_{hii}}{m_{i}} A_{i}^{i} \left(\frac{m_{h}^{2}}{4m_{i}^{2}}\right)}{0.021}$$

Since

$$\mathcal{A}_W + \mathcal{A}_t \simeq -0.075$$

if we want to increase $h \rightarrow \gamma \gamma$ we need $\delta_{C,NC} < 0$. Possibilities:



with

$$A_{vector}\simeq -5A_{fermion}, \quad A_{fermion}\simeq (2-4)A_{scalar}$$

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A closer look to $2g_{hii}/m_i$

Ellis et al. 76, Shifman et al. 79:



Conclusion:

eigenvalues repulsion with vev-dependent mixing can do the job

Example:
$$M = \begin{pmatrix} m_1 & cv \\ cv & m_2 \end{pmatrix}$$

$$\Rightarrow \lambda_{1,2}^{2} = \frac{m_{1}^{2} + m_{2}^{2} + 2c^{2}v^{2} \mp (m_{1} + m_{2})\sqrt{(m_{1} - m_{2})^{2} + 4c^{2}v^{2}}}{2}$$

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SM + (uncolored/colored) fermions

A chiral fourth generation?

Mixing with SM fermions?

Does not decouple from Higgs system, production \sim 9 times larger than in SM

Possible, but better be small so to avoid problems with precision measurements and flavor

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In this talk



Analysis of effects on



Models considered

$$\begin{array}{c|cccc} SU(2)_{L} & U(1)_{Y} & SU(3)_{c} \\ \hline 2+1 & 3+2 \\ \hline D_{L,R}=2 & T_{L,R}=3 & \hat{y}=y-\frac{1}{2} & \textbf{1}, \textbf{3}, \textbf{6}, \textbf{8} \\ S_{L,R}=1 & D_{L,R}=2 & y & \textbf{1}, \textbf{3}, \textbf{6}, \textbf{8} \end{array}$$

with lagrangian

$$-\mathcal{L}^{2+1} = c \,\overline{D}_R \,HS_L + c \,\overline{D}_L HS_R + m_1 \overline{D}_R D_L + m_2 \overline{S}_R S_L + \text{h.c.}$$
$$-\mathcal{L}^{3+2} = c \overline{D}_R \,T_L H + c \overline{D}_L \,T_R H + m_1 \,\overline{T}_L T_R + m_2 \,\overline{D}_L D_R + \text{h.c.}$$

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Models considered

SU(2)_LU(1)_Y
$$SU(3)_c$$
 $2+1$ $3+2$ $D_{L,R} = 2$ $T_{L,R} = 3$ $\hat{y} = y - \frac{1}{2}$ **1**, **3**, **6**, **8** $S_{L,R} = 1$ $D_{L,R} = 2$ y **1**, **3**, **6**, **8**

Particle content:

	Q coupled to h	Q not coupled to h
2 + 1	(<i>y</i> , <i>y</i>)	y-1
	(y+1/2, y+1/2)	
3 + 2	+	y - 3/2
	(y-1/2, y-1/2)	

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2+1 Uncolored



3+2 Uncolored



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Constraints from oblique parameters: 2 + 1 case



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Constraints from oblique parameters: 3 + 2 case



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$$R_{\gamma\gamma} = \frac{\Gamma(h \to \gamma\gamma)}{\Gamma(h \to \gamma\gamma)_{SM}} = |1 + \sum_{i} \delta_{C}^{i} + \sum_{j} \delta_{NC}^{j}|^{2}$$
$$R_{GG} = \frac{\sigma(GG \to h)}{\sigma(GG \to h)_{SM}} = |1 - 9.7\sum_{i} \frac{t_{c}^{i}}{N_{c}^{i}Q_{i}^{2}} \delta_{C}^{i}|^{2}$$

where



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General expectation



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In principle,

SM+ colored vector-like fermions

should be able to accommodate <u>both</u> $GG \rightarrow h \& h \rightarrow \gamma \gamma$.

Again: consider SM + (2 + 1) or (3 + 2) colored vector-like fermions:

effect on S, T and U?

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R_{GG} depends on *t_c* = tr(*T^AT^B*) = 1/2 (3), 5/2 (6), 3 (8) *R_{γγ}* = |1 + δ_c|² with δ_c ∝ *N_c S*, *T*, *U* ∝ *N_c*

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2+1 with colored **3** and **8**



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3+2 with colored **3**



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3+2 with colored **8**



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2+1 - lightest color triplet pair production

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A model building disclaimer

RGE for the Higgs quartic:
$$\frac{d\lambda}{d \log \mu} = -\frac{N_c}{8\pi^2}c^4$$

 λ driven negative at a scale $\Lambda \sim 3 m_{light}$

Arkani-Hamed et al 2012, Reece 2012

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new fields should enter at this scale to postpone vacuum instability (i.e. new scalars, but their effect on $R_{\gamma\gamma}$ and R_{GG} should be small)

In our case, in most of parameter space $m_{heavy} \lesssim (2-2.5) m_{light}$

- ▶ We have a new particle, it looks like the SM Higgs boson
- Some interesting hints that something may be going on in the loops (but right now everything consistent with SM at 95% C.L.)
- Addition of new particles to the SM:
 - Scalars: (relatively) small effect, mixing with Higgs
 - Fermions: mixed, unmixed (this talk)
 - Vectors: ?
 - Spin 2: Urbano 2012

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General message:

- Uncolored Fermions:
 - \blacktriangleright Increasing $R_{\gamma\gamma}$ up to $\gtrsim 2$ requires large electric charges $Q\gtrsim 2$
 - ► A modest increase (up to ~ 20%) not too problematic even with Q ~ 1 and light ~ 200 - 300 GeV particles
- Colored Fermions:
 - $R_{GG} \sim 0.5$ simpler with octects, but more difficult to pass *S*, *T* and *U*
 - *R*_{γγ} approximately as in the uncolored case (but more difficult to pass EWPM according to the number of colors)

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BACK UP

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New Gfitter result

$$\begin{split} \alpha(M_Z^2) \, S^{\mathsf{NP}} &= \quad \frac{4 s_W^2 c_W^2}{M_Z^2} \left[\Pi_{ZZ}^{\mathsf{NP}}(M_Z^2) - \Pi_{ZZ}^{\mathsf{NP}}(0) - \Pi_{\gamma\gamma}^{\mathsf{NP}}(M_Z^2) - \frac{c_W^2 - s_W^2}{c_W s_W} \Pi_{\gamma Z}^{\mathsf{NP}}(M_Z^2) \right] \,, \\ \alpha(M_Z^2) \, T^{\mathsf{NP}} &= \quad \frac{\Pi_{WW}^{\mathsf{NP}}(0)}{M_W^2} - \frac{\Pi_{ZZ}^{\mathsf{NP}}(0)}{M_Z^2} \,, \\ \alpha(M_Z^2) \, U^{\mathsf{NP}} &= \quad 4 s_W^2 \left[\frac{\Pi_{WW}^{\mathsf{NP}}(M_W^2) - \Pi_{WW}^{\mathsf{NP}}(0)}{M_W^2} - c_W^2 \left(\frac{\Pi_{ZZ}^{\mathsf{NP}}(M_Z^2) - \Pi_{ZZ}^{\mathsf{NP}}(0)}{M_Z^2} \right) \right. \\ &\left. - 2 s_W c_W \, \frac{\Pi_{\gamma Z}^{\mathsf{NP}}(M_Z^2)}{M_Z^2} - s_W^2 \, \frac{\Pi_{\gamma\gamma}^{\mathsf{NP}}(M_Z^2)}{M_Z^2} \right] \,, \end{split}$$

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