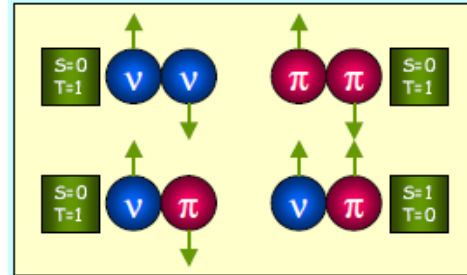
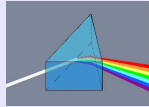


Probing nucleon-nucleon correlations in heavy ion transfer reactions

Suzana Szilner
Ruđer Bošković Institute
Zagreb, Croatia



PRISMA collaboration



Catania



Zagreb

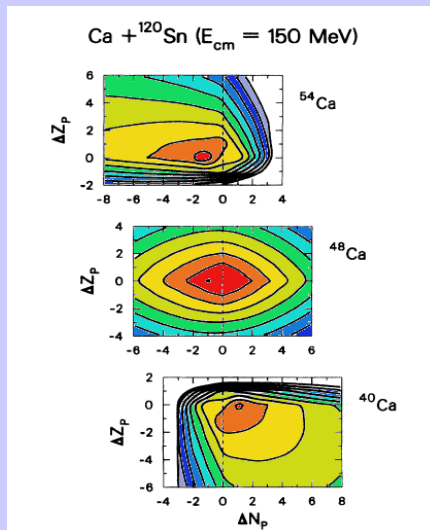


Nucleus Nucleus 2015

21-26 June 2015 Dipartimento di Fisica ed Astronomia, Università di Catania
Europe/Rome timezone

Transfer reactions among heavy ions

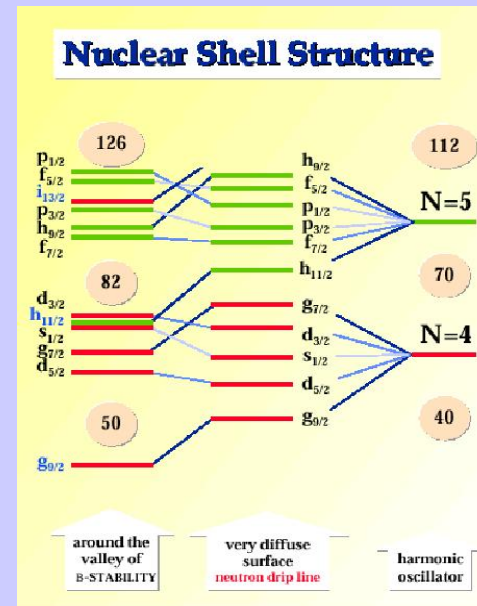
- ✓ A study of multiple particle transfers.
- ✓ The transition from quasi elastic to deep inelastic processes.
- ✓ A tool for the population of neutron rich nuclei.



C.H. Dasso et al., PRL 73 (1994) 1907

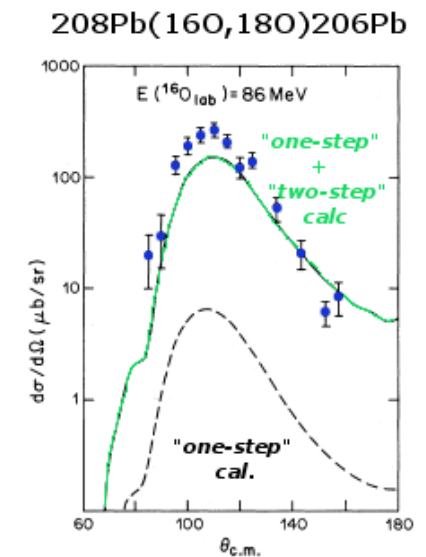
A study of properties near shell closure:

- ✓ single particle states
- ✓ coupling of the particle/hole to the collective boson



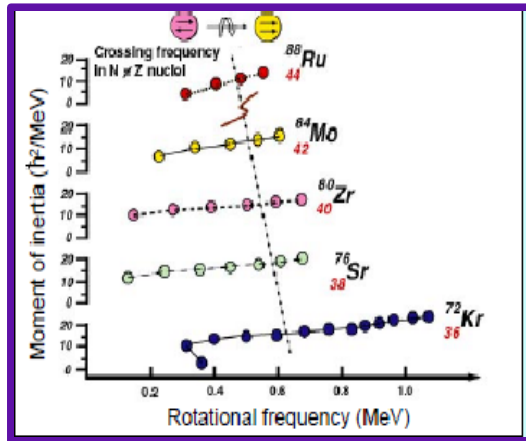
A study of the residual interaction (correlations):

- ✓ two particle transfer (absolute values)
- ✓ population of specific states (pairing vibration/rotation)

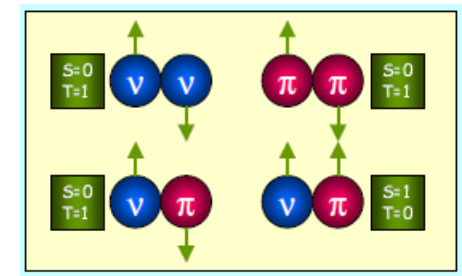
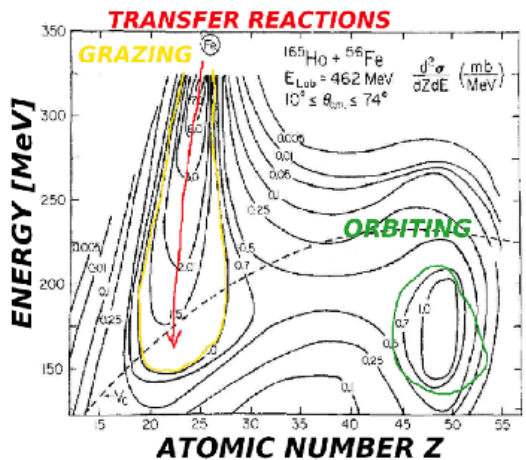


B.F. Bayman et al., PRC 26 (1982) 1509

Probing correlations



Higher rotational frequency for “pair” break in N=Z



How the correlations that go beyond a mean field description can be probed (static and dynamics properties and effects)?

- Binding energies: the ground states → description in terms of superfluid condensates, in which the pairs of nucleons form the Cooper pairs
- Significantly different behavior at medium to high spins of rotational bands
- Enhanced probability to add or remove a nucleon-nucleon pair.

What will be a signature in the heavy ion transfer reactions?

HI **advantages**: test of correlation properties in multi-neutron and multi-proton transfer processes via simultaneous comparison of observables for $\pm nn/\pm pp/\pm np$ pairs

HI **drawbacks**: difficult experimental conditions (A,Z,Q-value resolutions, total efficiency) and difficult theoretical treatment (complex structure of the two interacting ions, QE and DIC processes, multistep processes, many open channels and CC effects)

Magnetic spectrometers for transfer reaction studies

Q3D, split-pole

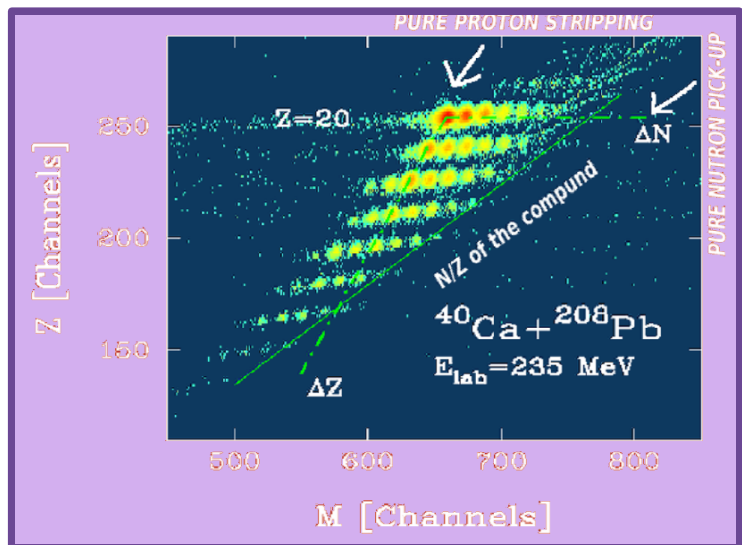
1980

- ✓ excited states populated in light ion transfer reactions (en. resolution ~few tenths keV)
- ✓ distribution of atomic charge states (magnetic elements of different complexity to focus momenta at the focal plane)

TOF spectrometers

1990

- ✓ focus ions of different atomic charge states to a (small) focal plane
- ✓ good A and Z resolution, and detection efficiency, large energy dynamic range of transfer products



S.Szilner et al, PRC 71 (2005) 044610

Large solid angle spectrometers

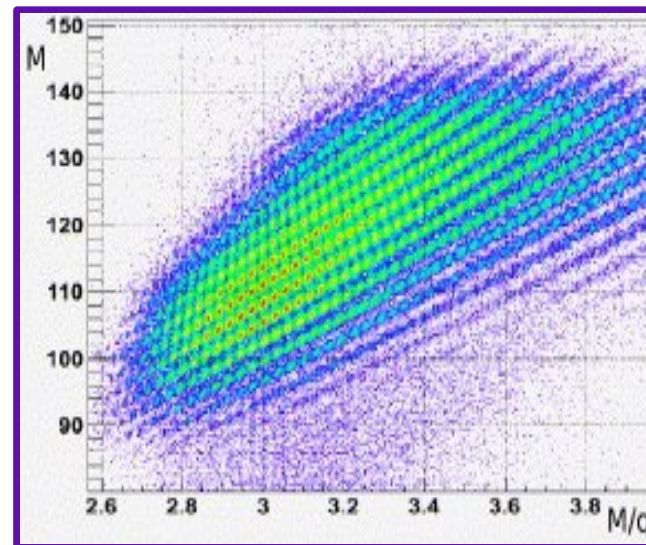
2000

PRISMA, VAMOS, MAGNEX

coupling to large γ arrays

CLARA, EXOGAM, AGATA

- ✓ simple magnetic elements and “complex” detector systems
- ✓ good A and Z resolution, detection efficiency



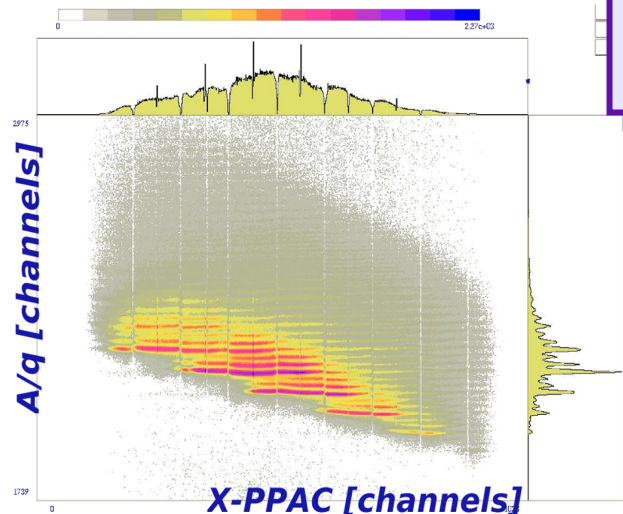
S. Pullanhiotan et al, NIM A 593 (2008) 343

The PRISMA spectrometer

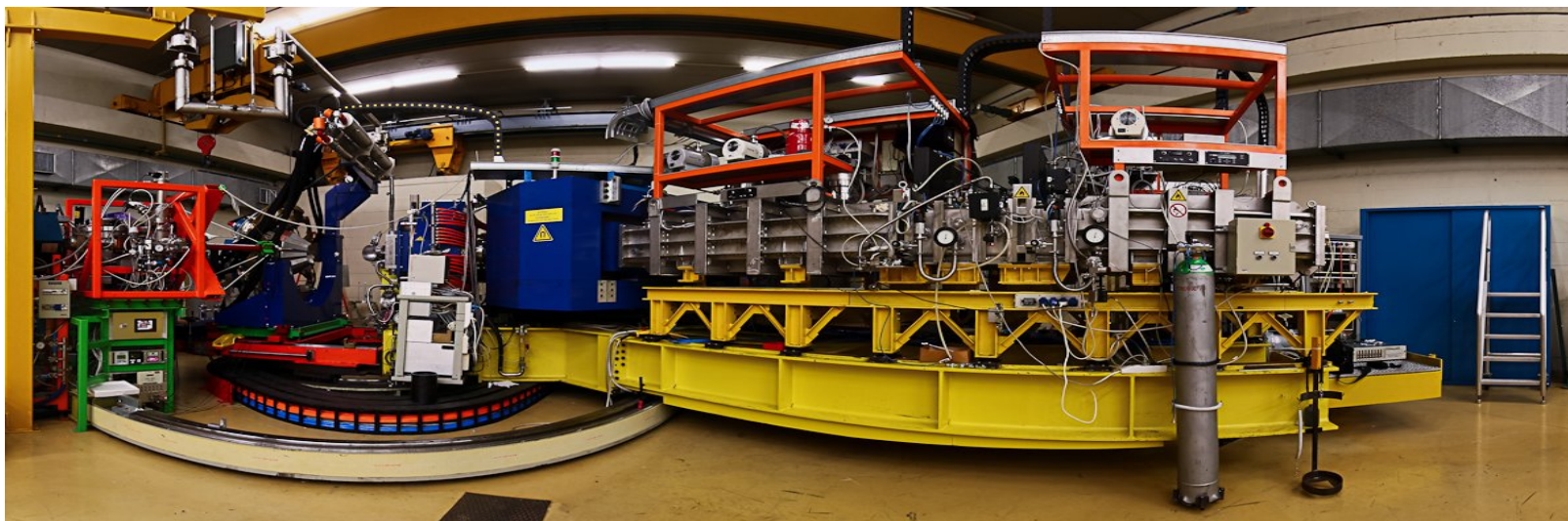
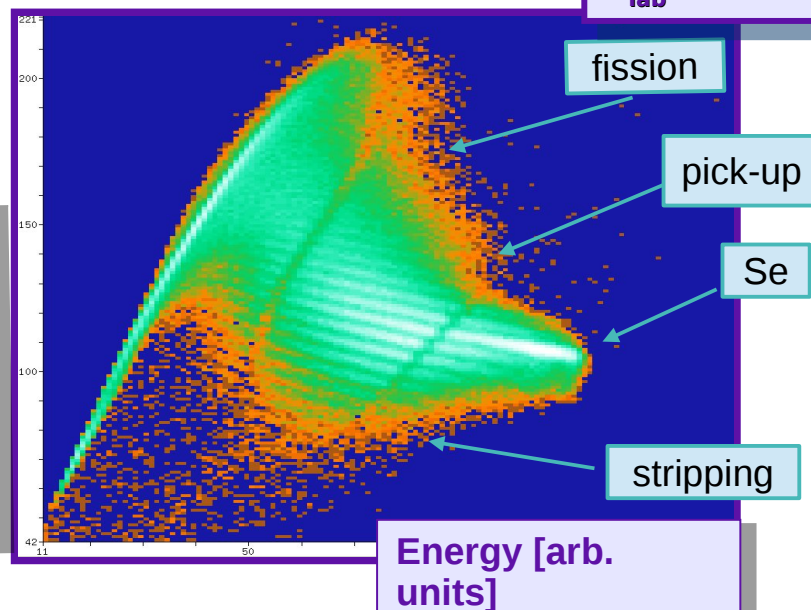
reconstruction of the ion trajectory inside (simple) magnetic elements, ray tracing procedure: position sensitive detectors of large area

$^{82}\text{Se} + ^{238}\text{U}$
 $E_{\text{lab}} = 505 \text{ MeV}$
 $\theta_{\text{lab}} = 64^\circ$

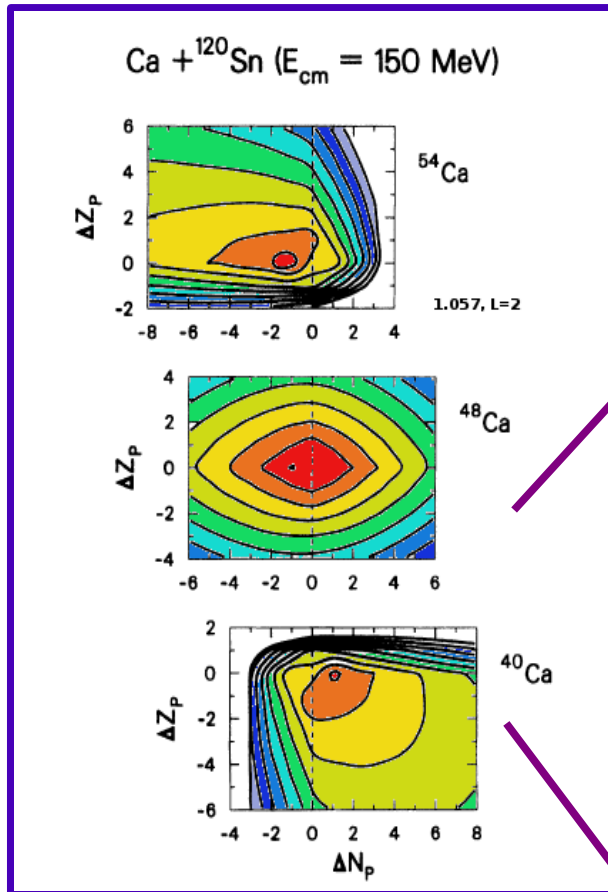
$^{90}\text{Zr} + ^{208}\text{Pb}$
 $E_{\text{lab}} = 560 \text{ MeV}$
 $\theta_{\text{lab}} = 61^\circ$



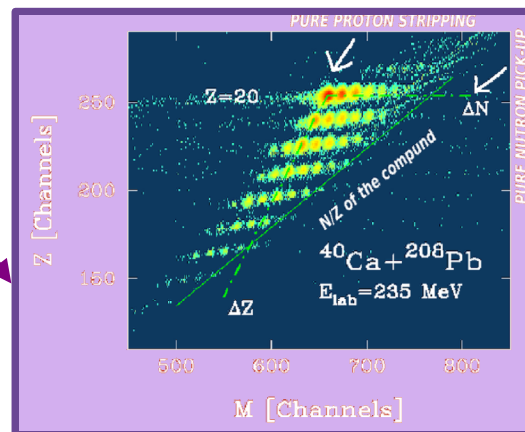
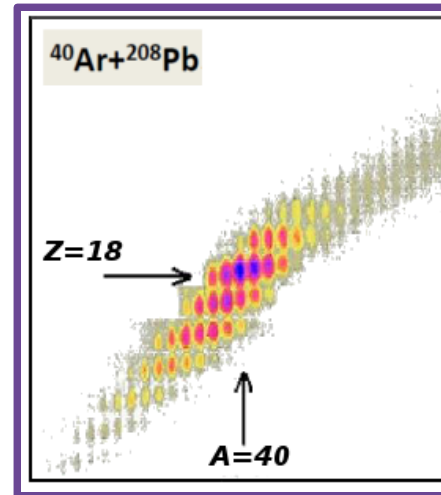
ΔE [arb. units]



Properties of transfer reactions at the Coulomb barrier

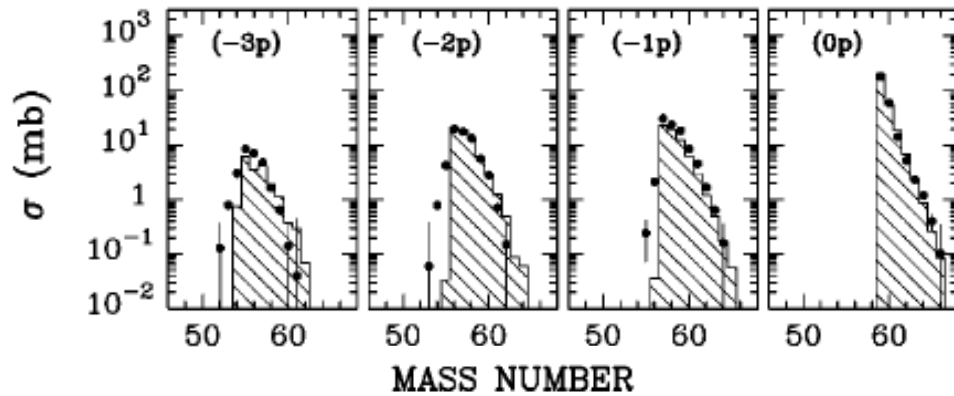


C. H. Dasso, G. Pollaro, and A. Winther, PRL 73 (1994) 1907
L. Corradi, G. Pollaro and S. Szilner, J. Phys. G 36 (2009) 113101 (Special Topic)

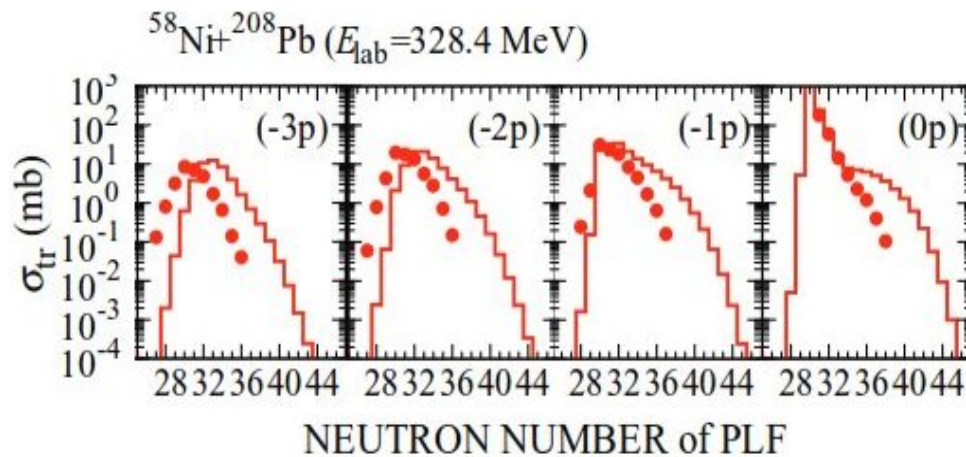


- ✓ The transfer process is governed by **optimum Q-value** and nuclear structure properties.
- ✓ Nuclei are located on the left side of the charge equilibration line --> dominance of a **direct mechanism**
- ✓ For the massive proton transfer channels → the isotopic distributions drift towards lower masses (neutron **evaporation**)
- ✓ **GRAZING model**: calculates the evolution of the reaction by taking into account:
 - ✓ the **relative motion** (nuclear + Coulomb field), and
 - ✓ the **intrinsic degrees of freedom of projectile and target** (surface vibration and the **single-nucleon transfer** channels).
- ✓ The multinucleon transfers are described via a **multistep** mechanism.
- ✓ The model takes into account the effect of neutron evaporation.

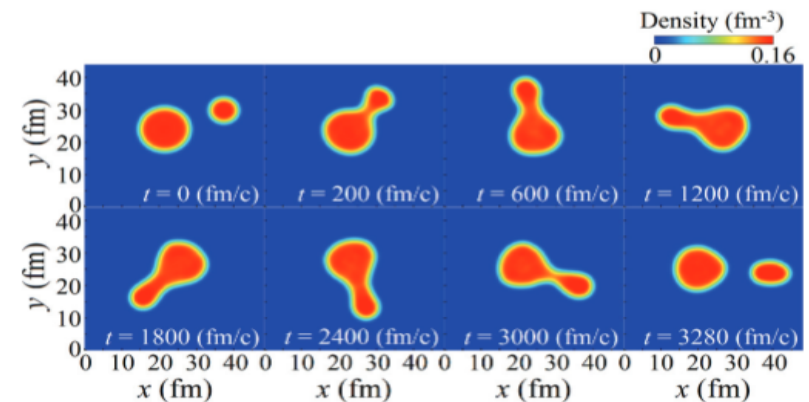
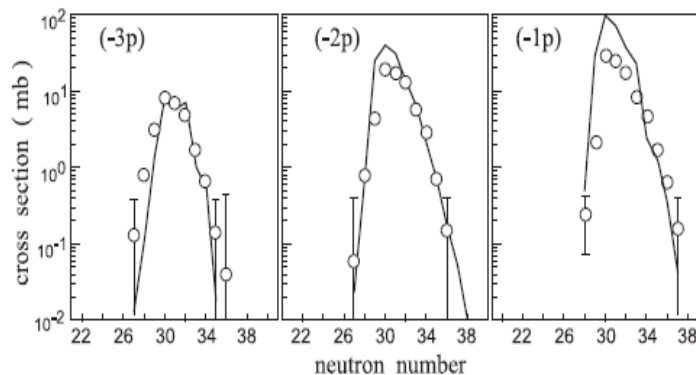
Multinucleon transfer reactions : experiment vs. theory



EXP: $^{58}\text{Ni}+^{208}\text{Pb}$,
L. Corradi et al., PRC 66 (2002) 024606
GRAZING or CWKB, G. Pollaro



Time Dependent Hartree-Fock theory
K.Sekizawa, K.Yabana, PRC 88 (2013) 014614



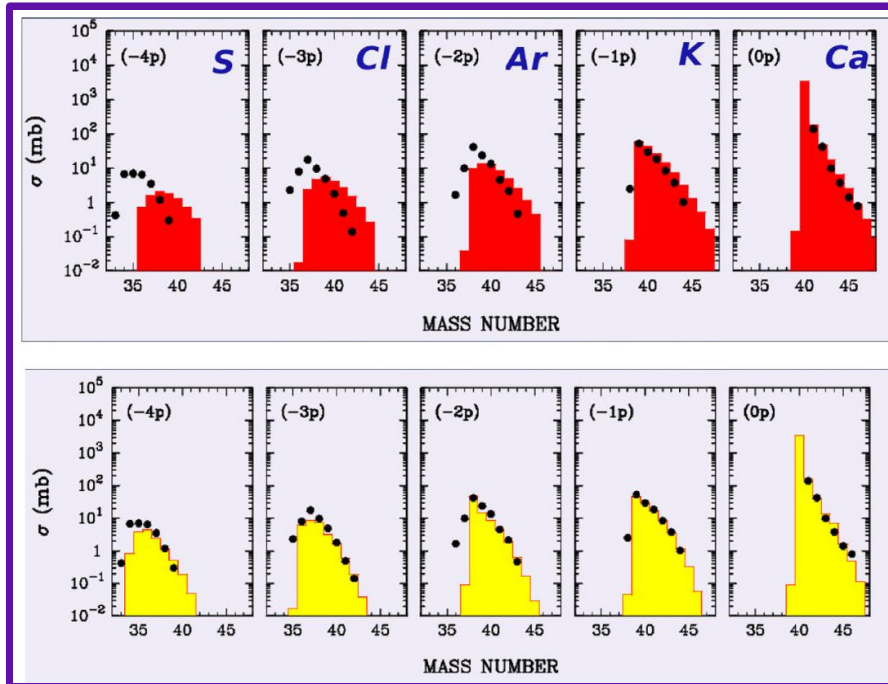
Langevin-type dynamical equations of motion
V. Zagrebaev, W. Greiner PRL 101 (2008) 122701

Total cross sections - nucleon-nucleon correlations

$^{40}\text{Ca} + ^{208}\text{Pb}$

CWKB cal.
successive
transfer

+
simultaneous
transfer
(microscopic
pair form
factor)
+
neutron
evaporation



- ✓ The shape of the yield distribution reflects the **optimum Q-value**
- ✓ The theory describes well **(0p)** and **(-1p)** but underestimates **(-2p)**
- ✓ The contribution of **a direct pair mode** (“macroscopic”) both for neutrons and protons has been added in the calculations.
- ✓ The same strength of the form factor for neutrons and protons.
- ✓ The **pair mode** alters little the cross section for neutron transfer but is **essential for the proton** transfer.

$$F_P(r) = \beta_P \frac{\partial V(r)}{\partial A} \simeq \left(\frac{\beta_P R}{3A} \right) \frac{\partial V(r)}{\partial r},$$

- ✓ **The residual interaction: components responsible for the couplings (phonon - single particle), and for nucleon - nucleon correlations**
- ✓ **Inclusive data - difficult to have a clear signature of pair mode**

L. Corradi et al., PRC 66 (2002) 024606

S. Szilner et al., PRC 71 (2005) 044610

L. Corradi, G. Pollaro and S. Szilner, J. Phys. G 36 (2009) 113101 (Special Topic)

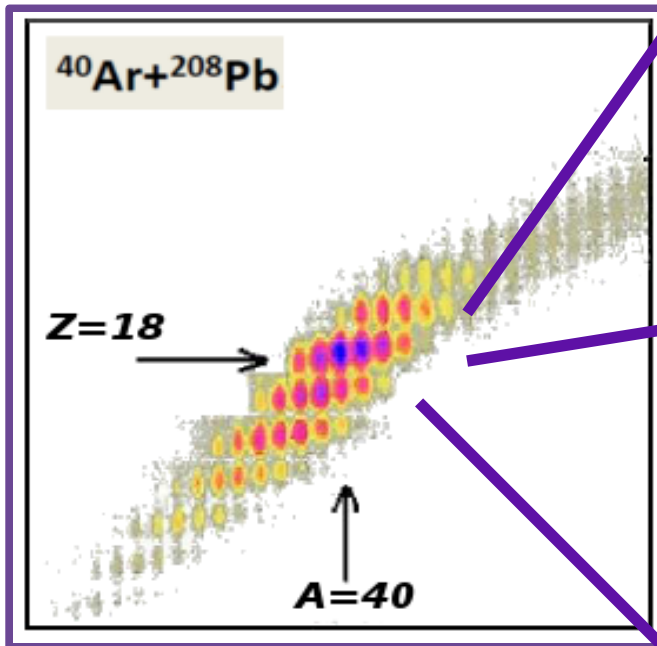
“coupling to large γ arrays”

exp: $^{40}\text{Ar}+^{208}\text{Pb}$

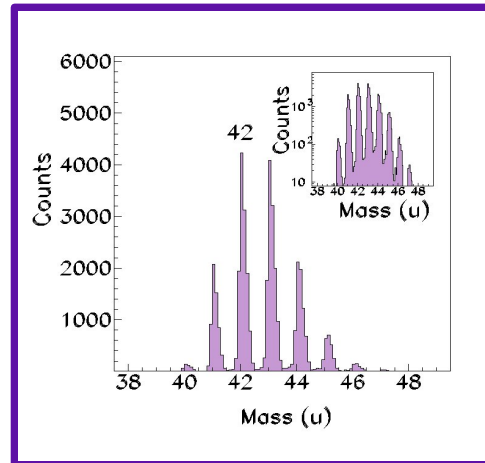
$\theta_{\text{lab}} = 46^\circ$

54° (grazing)

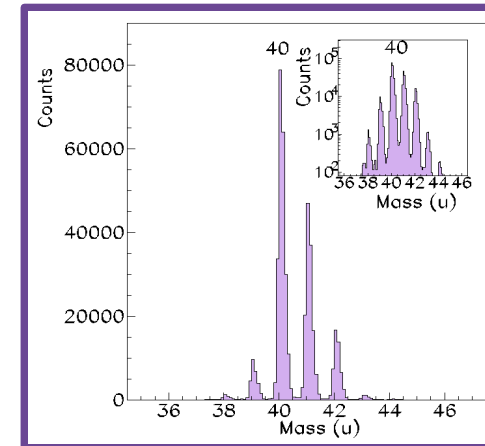
59°



T. Mijatović, Ph.D. Thesis, 2015
Diff. and total cross sections

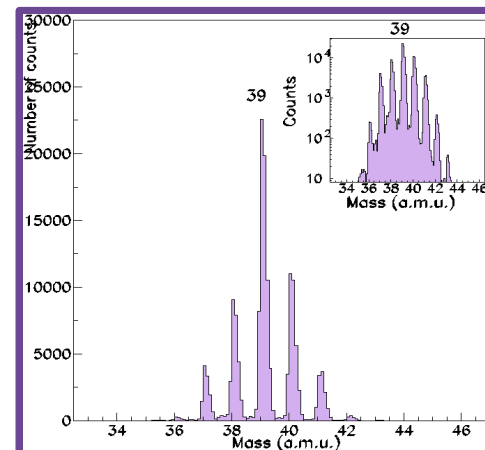


Z=19: K isotopes
A=40-46
(+1p±xn) transfer channels



Z=18: Ar isotopes
A=38-43
Neutron transfer channels

S. Szilner et al,
PRC 84 (2011) 014325

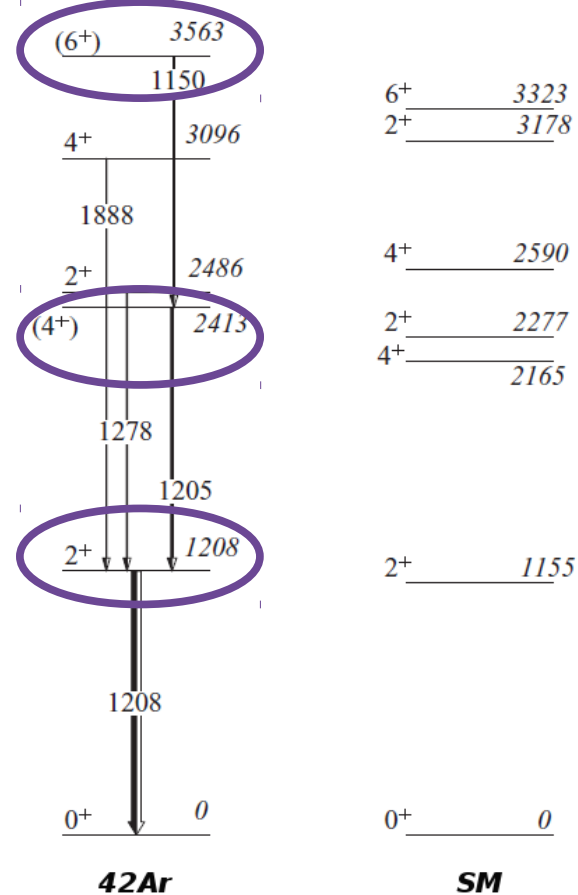
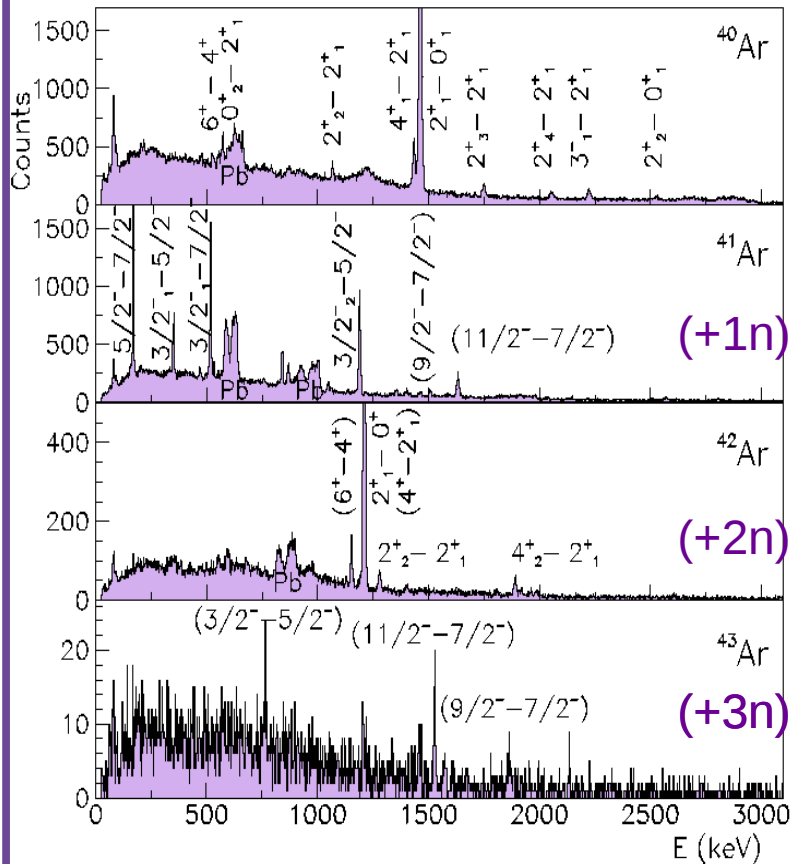
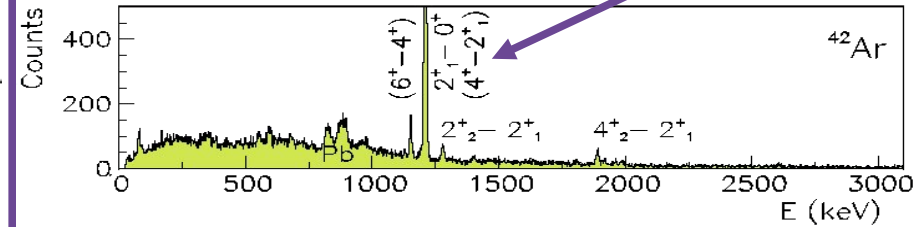
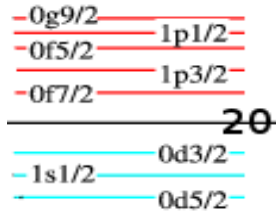


Z=17: Cl isotopes
A=37-42
(-1p±xn) transfer channels

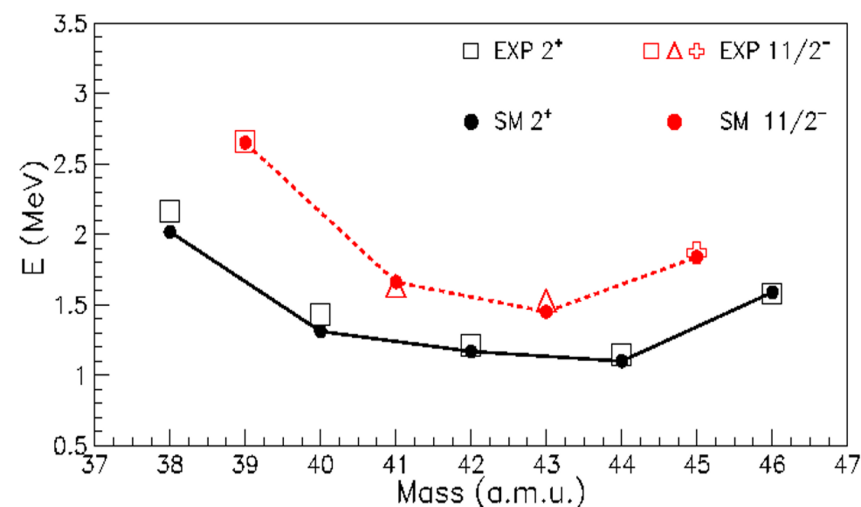
S. Szilner et al,
PRC 87 (2013) 054322

The character of states populated in transfer reactions

Ar isotopes: neutron transfer channels

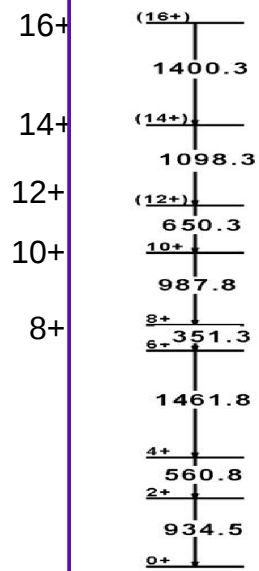


$0g_{9/2}$
 $0f_{5/2}$
 $0f_{7/2}$
 $1p_{1/2}$
 $1p_{3/2}$
20
 $0d_{3/2}$
 $1s_{1/2}$
 $0d_{5/2}$


$$2^+ \otimes \nu f_{7/2} \rightarrow 11/2^-, 9/2^-, \dots, 3/2^-$$


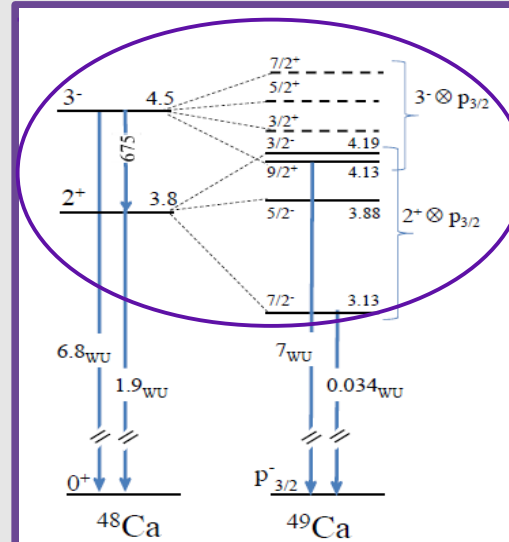
- ## A strong interplay between single-particle and collective degrees of freedom and the reaction dynamics

“yrast” states; “fermion-boson” coupling



$^{90}\text{Zr} + ^{208}\text{Pb}$ reaction:
 ^{92}Zr : a strong population of the “yrast” states (up to 16^+ at 7.4 MeV)

M. Varga Pajtler,
 Ph.D. Thesis,
 University of Zagreb, 2014.

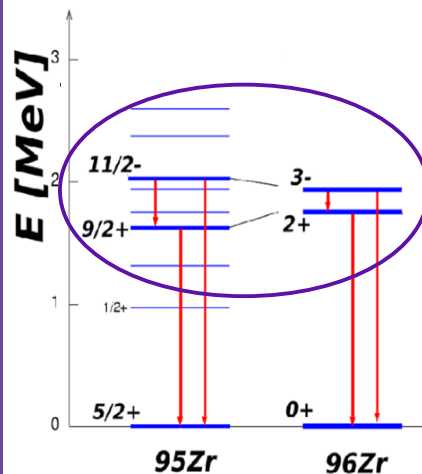


$^{48}\text{Ca} + ^{64}\text{Ni}$ reaction:
 ^{49}Ca : a strong population of “boson-fermion” multiplets

D. Montanari et al.,
 PLB 697, 288 (2011);
 PRC 85 (2012) 044301

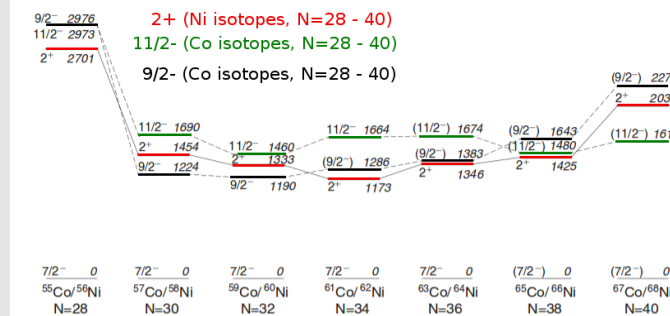
$(d_{3/2})^- \times 3^- \rightarrow 1/2^-, 3/2^-, \dots, 11/2^-$

$(d_{3/2})^- \times 2^+ \rightarrow 1/2^+, 3/2^+, \dots, 9/2^+$



$^{40}\text{Ca} + ^{96}\text{Zr}$ reaction:
 ^{95}Zr : a strong population of the $9/2^+$ and $11/2^-$ states (“boson-fermion” couplings)

S. Szilner et al, PRC 76 (2007) 024604



F. Recchia et al, PRC 85 (2012) 064305
Co isotopes

Fe: S. Lunardi et al, PRC 76 (2007) 034303
 Ar: D. Mengoni et al, PRC 82 (2010) 024308
 S. Bhattacharyya et al, PRL 101 (2008) 03501 ...

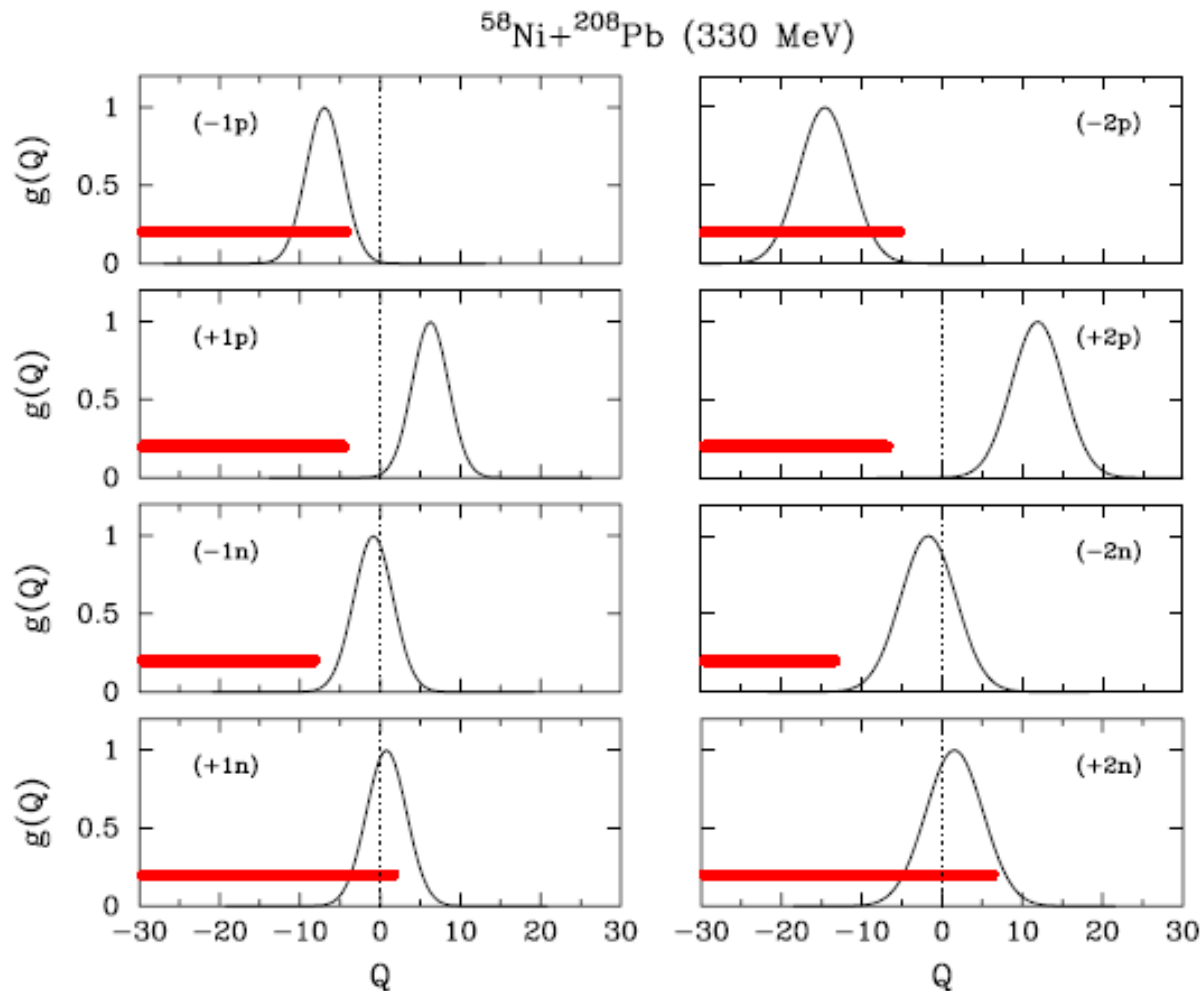
Properties of transfer reactions at the Coulomb barrier: optimum Q-value

cut-off function

$$g(Q) = \exp\left(-\frac{(Q - Q_{\text{opt}})^2}{\hbar^2 \ddot{r}_0 \kappa_{a_1'}}\right)$$

transfer probability

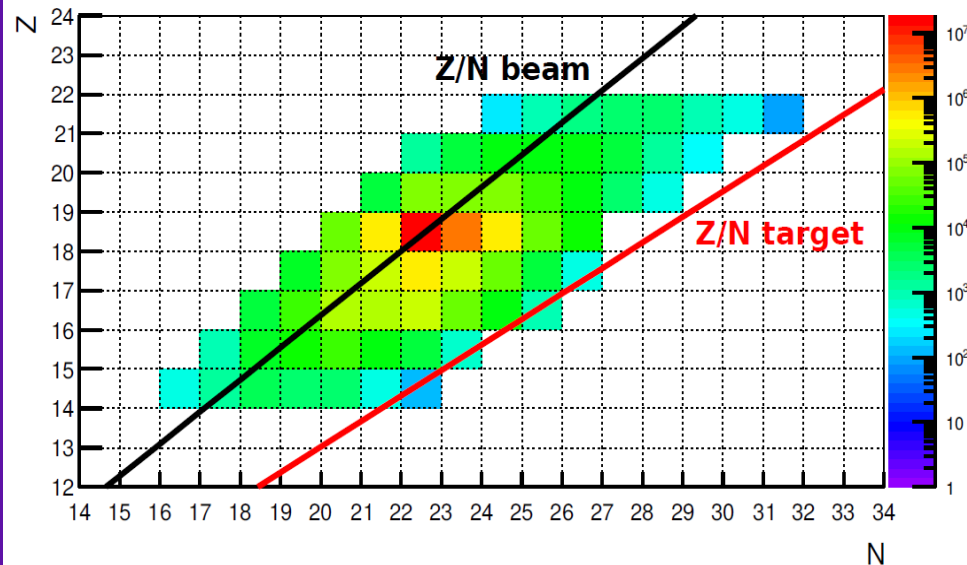
$$P_{\beta\alpha} = \sqrt{\frac{1}{16\pi\hbar^2 |\ddot{r}_0| \kappa_{a_1'}}} |f_{\beta\alpha}(0, r_0)|^2 g(Q_{\beta\alpha})$$



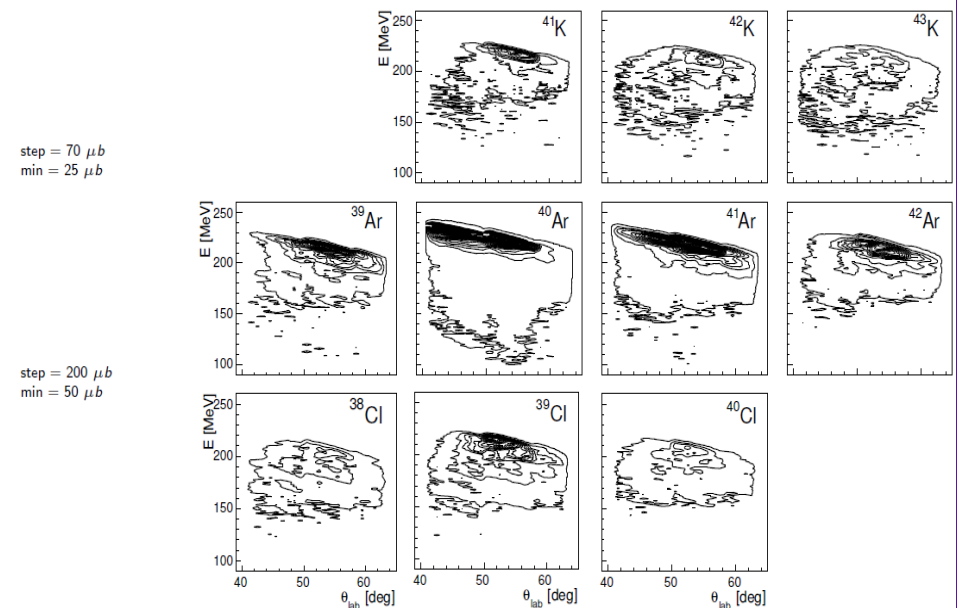
open reaction channels are those compatible with the optimum Q-value window (kinematical condition).

This window has its origin in the matching of the orbits before and after the transfer process.

The $^{40}\text{Ar} + ^{208}\text{Pb}$ system

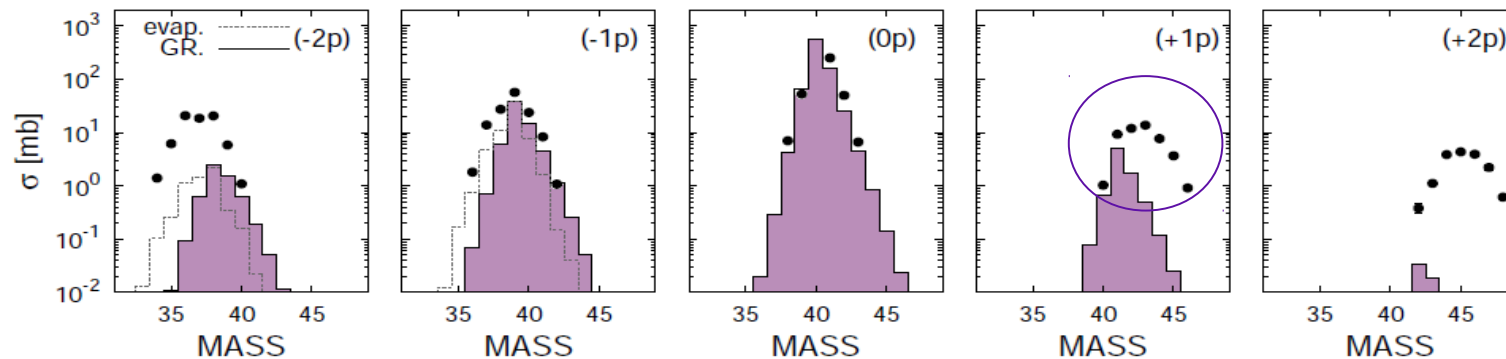


Z- M distribution



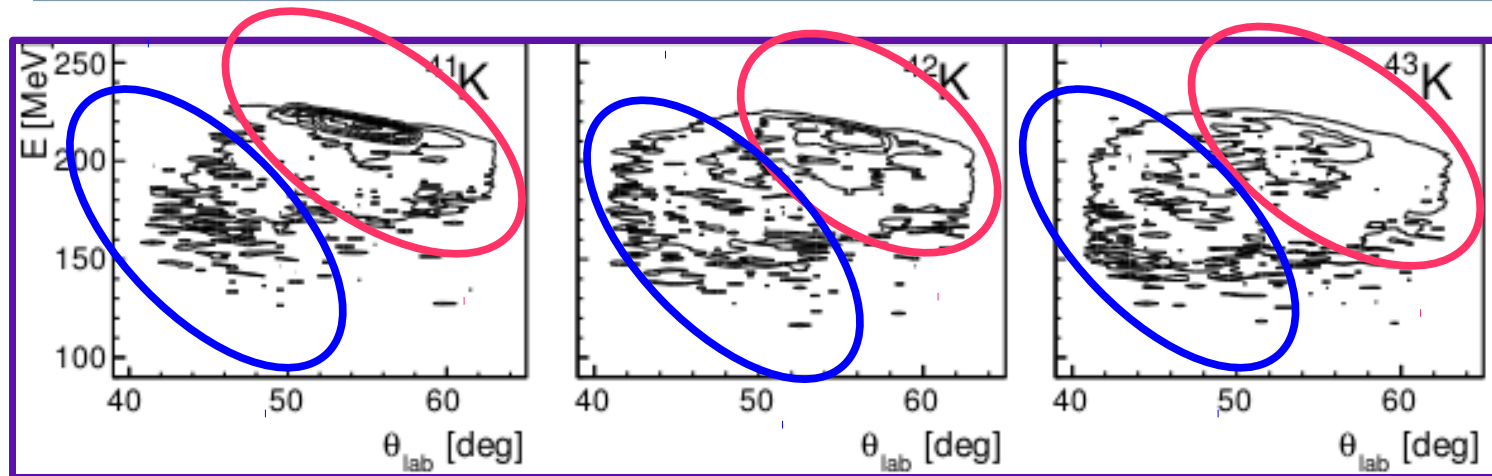
Wilczynski plots

$^{40}\text{Ar} + ^{208}\text{Pb}$ ($E_{\text{LAB}} = 260$ MeV)



EXP vs.
GRAZING:
the (+np)
channel –
factor 5

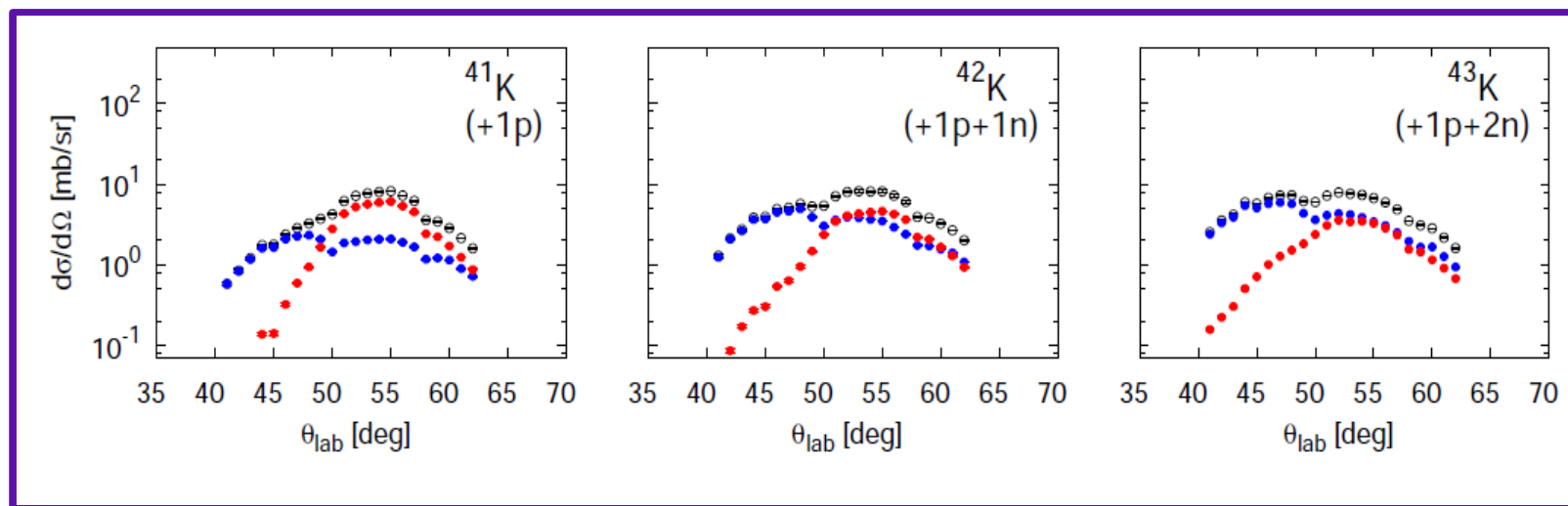
The $^{40}\text{Ar}+^{208}\text{Pb}$ system



Quasi-elastic
processes

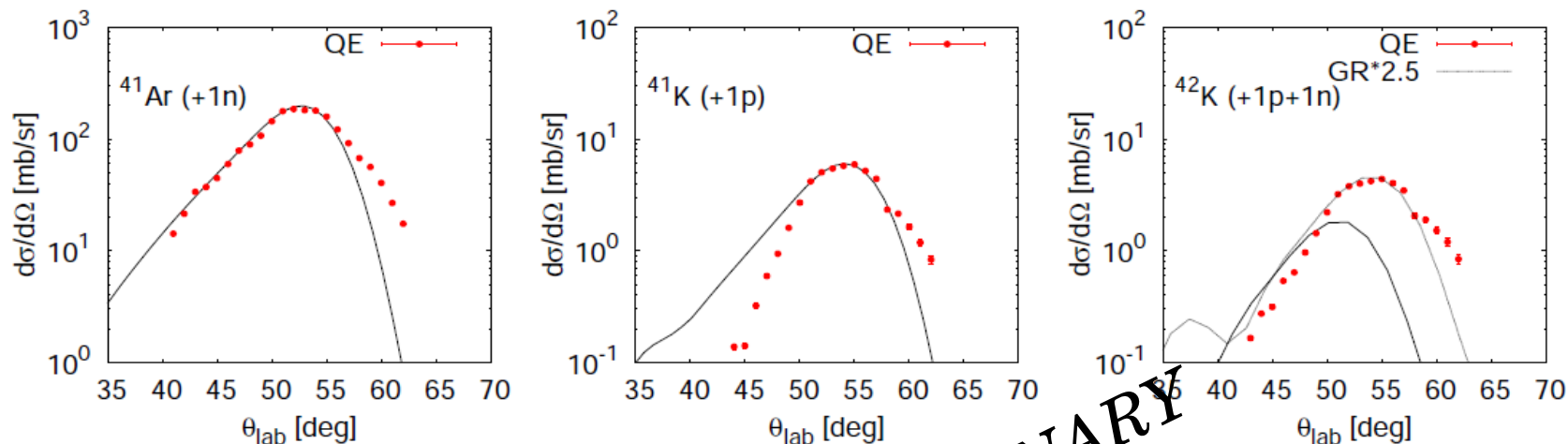
Deep-inelastic
processes

Wilczynski plots, (+1p+xn) channels

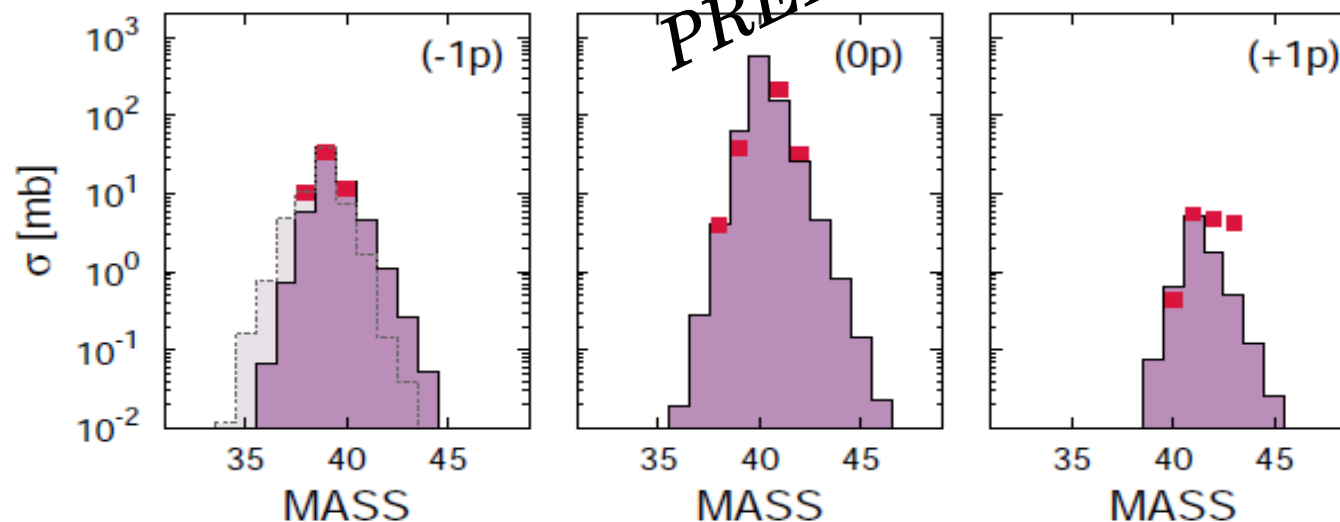


Angular distributions (integrated over energy: TOTAL, QE, DIC)

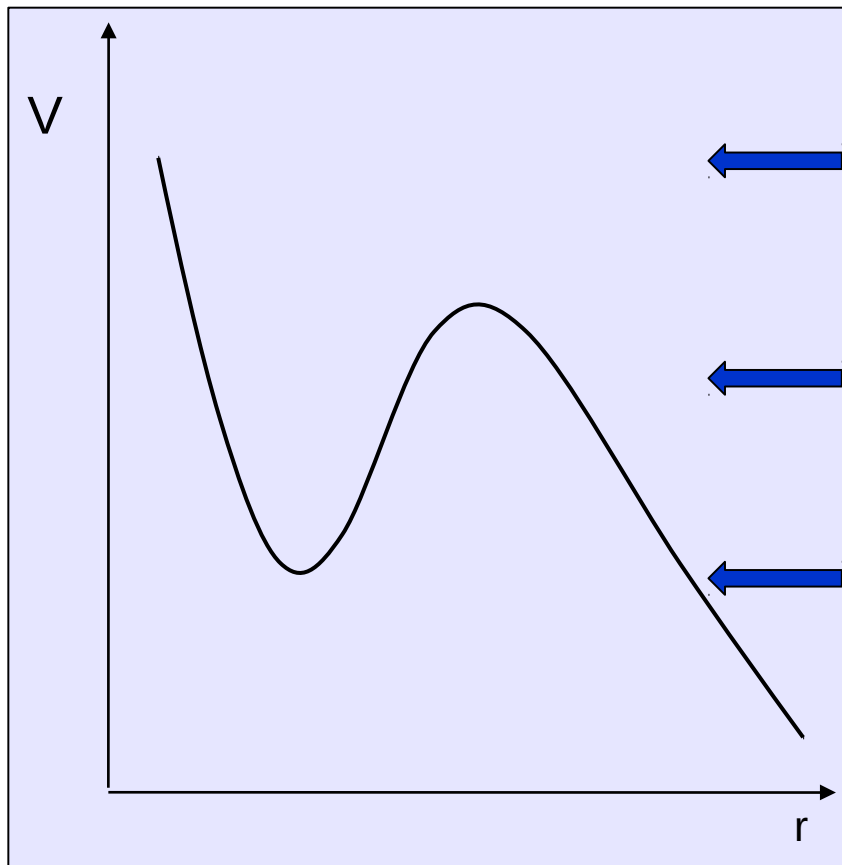
The $^{40}\text{Ar}+^{208}\text{Pb}$ system : QE ang. distribution and total cross sections



EXP vs.
GRAZING
(quasi-
elastic):
the (+np)
channel –
factor 2.5



A smooth transition between QE and DIC processes



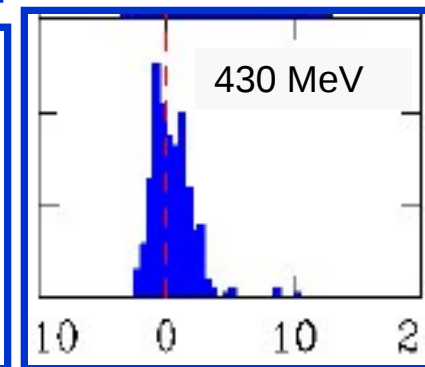
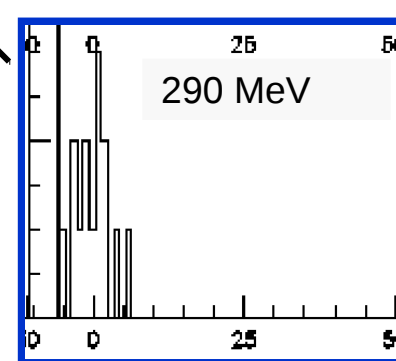
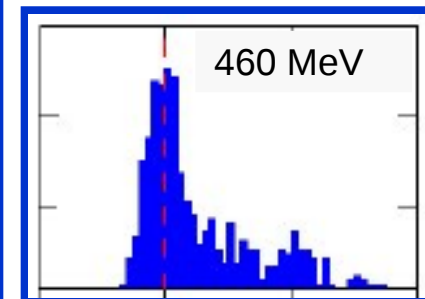
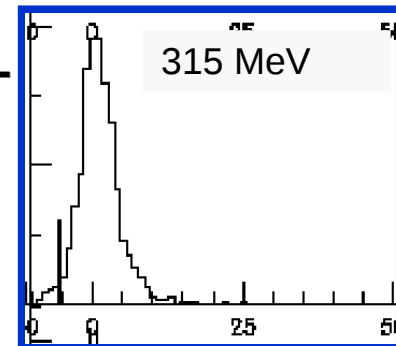
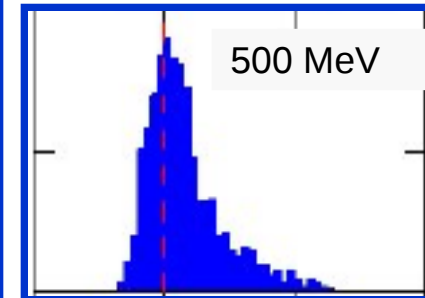
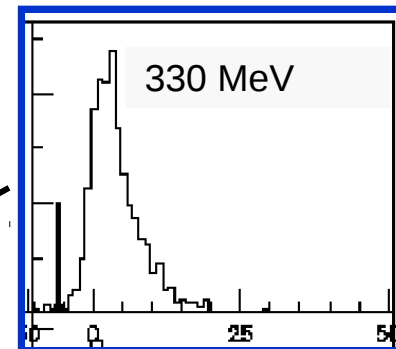
$E > E_b$

$E \sim E_b$

$E < E_b$

$^{96}\text{Zr}(^{40}\text{Ca}, ^{42}\text{Ca})$
 $Q_{\text{gs}} = +5.6 \text{ MeV}$

$^{116}\text{Sn}(^{60}\text{Ni}, ^{62}\text{Ni})$
 $Q_{\text{gs}} = +1.3 \text{ MeV}$



Below the barrier Q-values gets very narrow and without DIC components:

- 1) $E > E_b$, large number of open channels, DIC components
- 2) $E < E_b$, narrow Q-value distributions: no evaporation effects

L. Corradi et. al., PRC 84 (2011) 034603

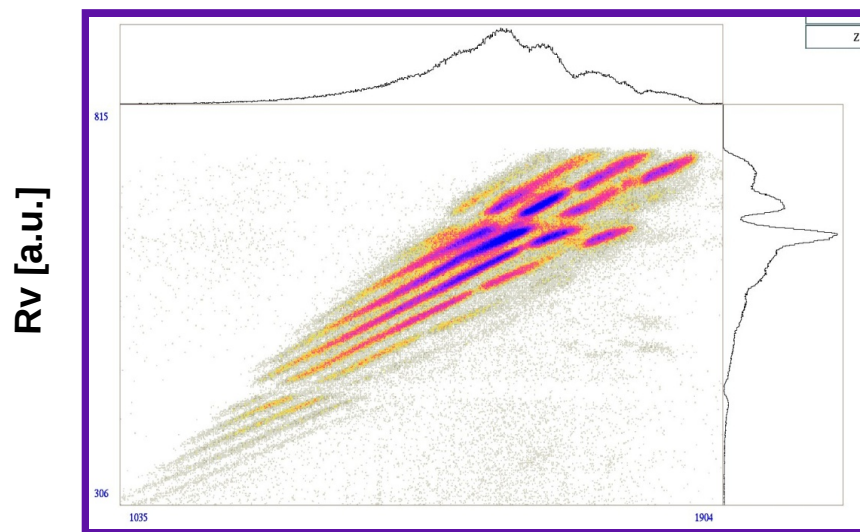
D. Montanari et. al., PRL 113 (2014) 052501

$^{60}\text{Ni} + ^{116}\text{Sn}$: detection of (light) target like ions in inverse kinematics with PRISMA

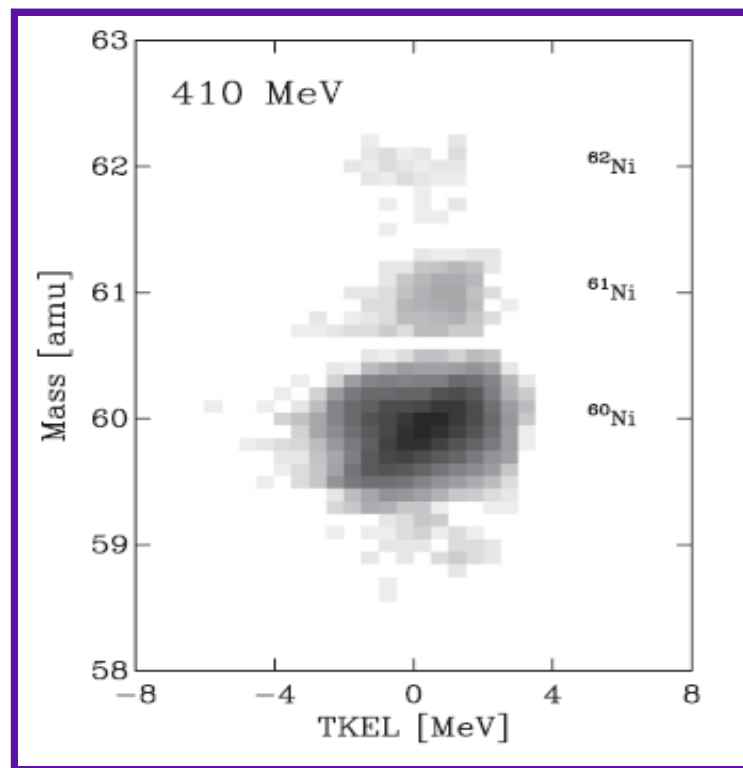
excitation function:

$$E_{\text{beam}} = 410 \text{ MeV} - 500 \text{ MeV}$$

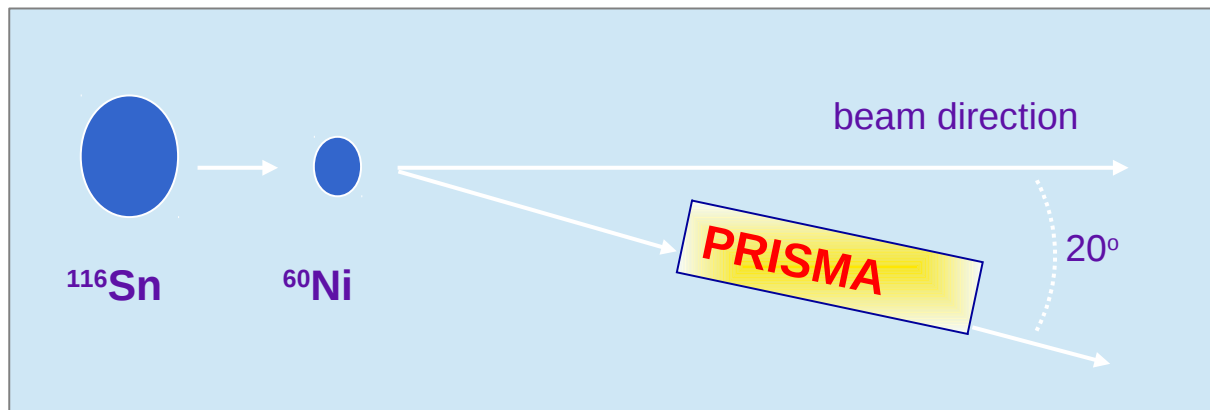
($D \sim 12.3$ to 15.0 fm)



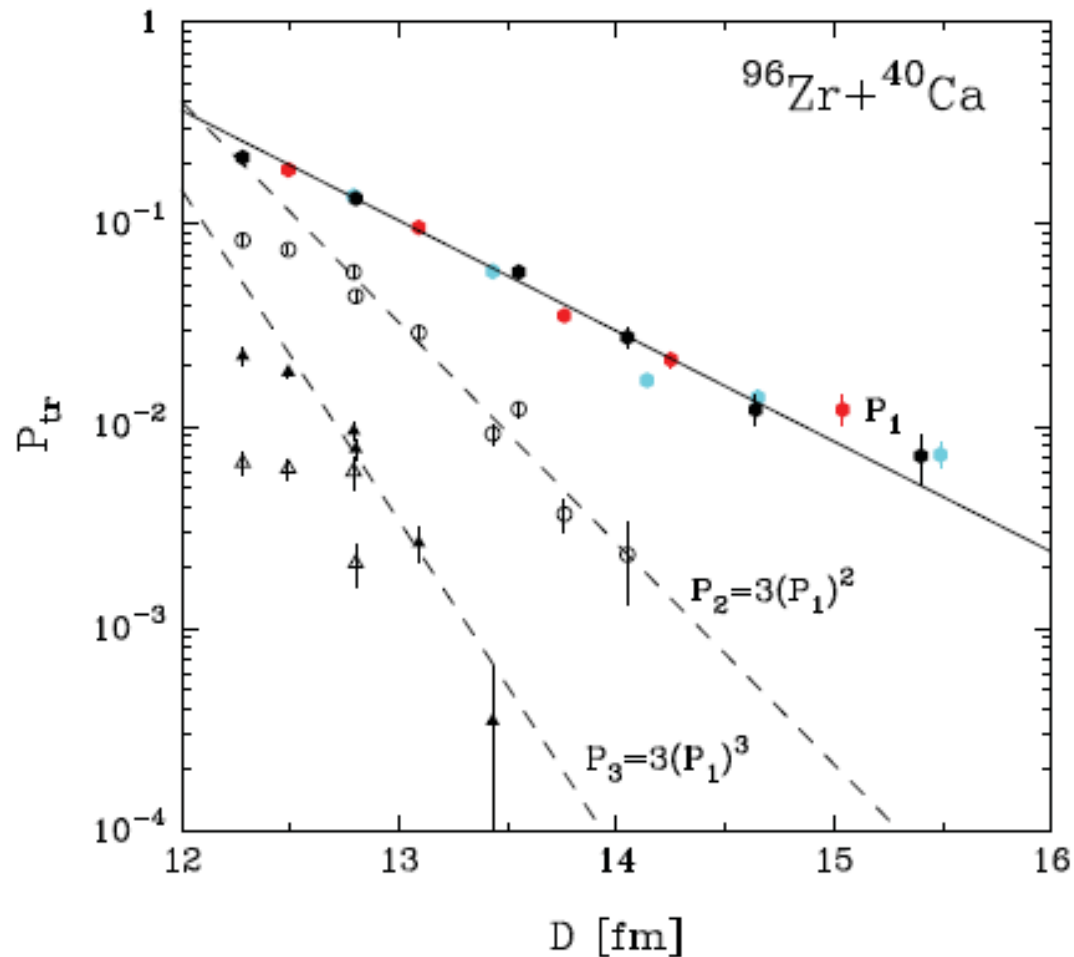
Energy [a.u.]



excellent channel
separation at
 $D \sim 15$ fm



Experimental transfer probabilities



P_{tr} slope

$$P_{tr} \propto e^{-2\alpha D} \quad \alpha = \sqrt{\frac{2mB}{\hbar^2}}$$

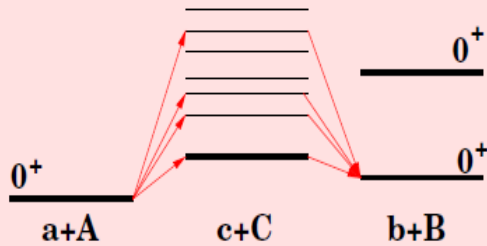
$B \rightarrow$ binding energy

slopes of P_{tr} vs D are as expected from the binding energies (tail of the formfactor)

a bare phenomenological analysis shows an “enhanced” pair transfer, $P_{2n} \sim 3(P_{1n})^2$ and $P_{3n} \sim P_{1n}(P_{2n}) \sim 3(P_{1n})^3$

$^{60}\text{Ni} + ^{116}\text{Sn}$: two particle transfer (semiclassical theory, microscopic calculations, 2nd order Born app.)

$$c_\beta(\ell) = c_\beta^{(1)} + c_\beta^{ort} + c_\beta^{succ}$$

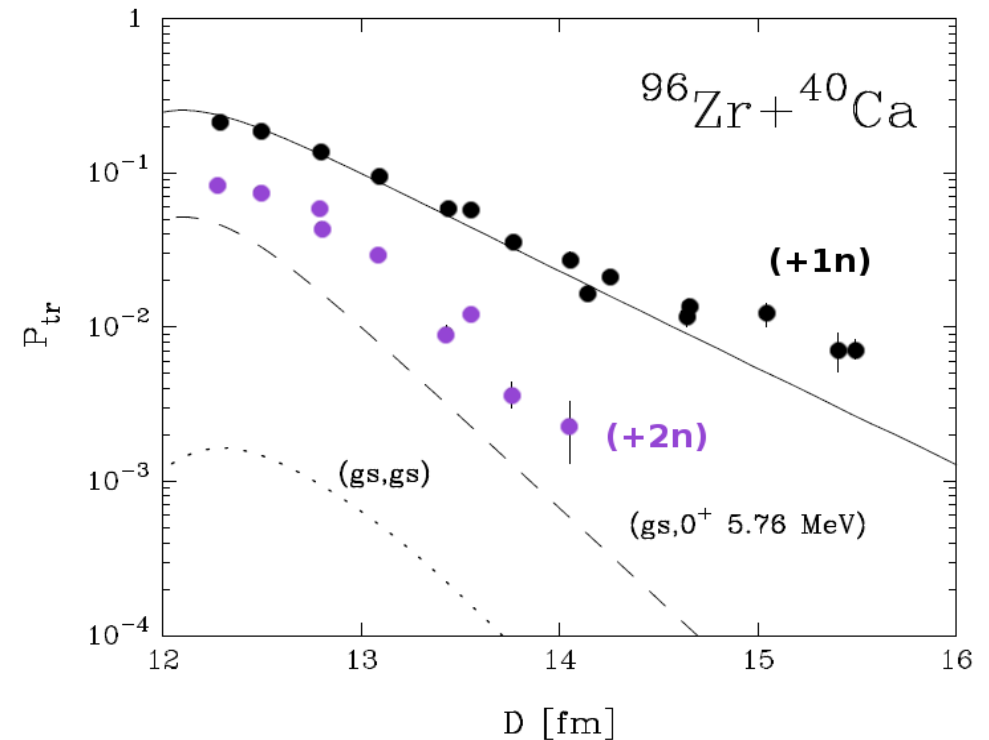
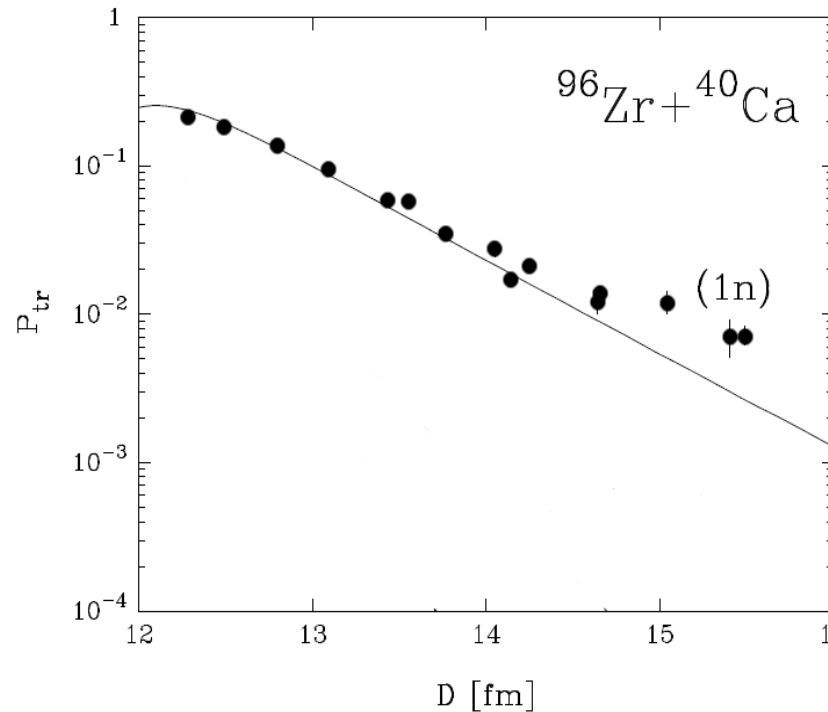


	nlj
^{116}Sn	$1g_{9/2}$
	$2d_{5/2}$
	$1g_{7/2}$
	$3s_{1/2}$
	$2d_{3/2}$
	$1h_{11/2}$
	$2f_{7/2}$
	$3p_{3/2}$
^{60}Ni	$2p_{3/2}$
	$2p_{1/2}$
	$1f_{5/2}$
	$1g_{9/2}$

3 terms : simultaneous, orthogonal and successive
only the successive term contributes to the transfer amplitude
Only 0+ to 0+ transition can be reduced to a simple expression

$$\begin{aligned}
 (c_\beta)_{succ} = & \frac{1}{\hbar^2} \sum_{a_1, a'_1} B^{(A)}(a_1 a_1; 0) B^{(a)}(a'_1 a'_1; 0) 2 \frac{(-1)^{j_1 + j'_1}}{\sqrt{(2j_1 + 1)} \sqrt{(2j'_1 + 1)}} \sum_{m_1 m'_1} (-1)^{m_1 + m'_1} \\
 & \times \int_{-\infty}^{+\infty} dt f_{m_1 m'_1}(\mathcal{R}) e^{i[(E_\beta - E_\gamma)t + \delta_{\beta\gamma}(t) + \hbar(m'_1 - m_1)\Phi(t)]/\hbar} \\
 & \times \int_{-\infty}^t dt f_{-m_1 - m'_1}(\mathcal{R}) e^{i[(E_\gamma - E_\alpha)t + \delta_{\gamma\alpha}(t) - \hbar(m'_1 - m_1)\Phi(t)]/\hbar} .
 \end{aligned}$$

Comparison between experimental and theoretical transfer probabilities



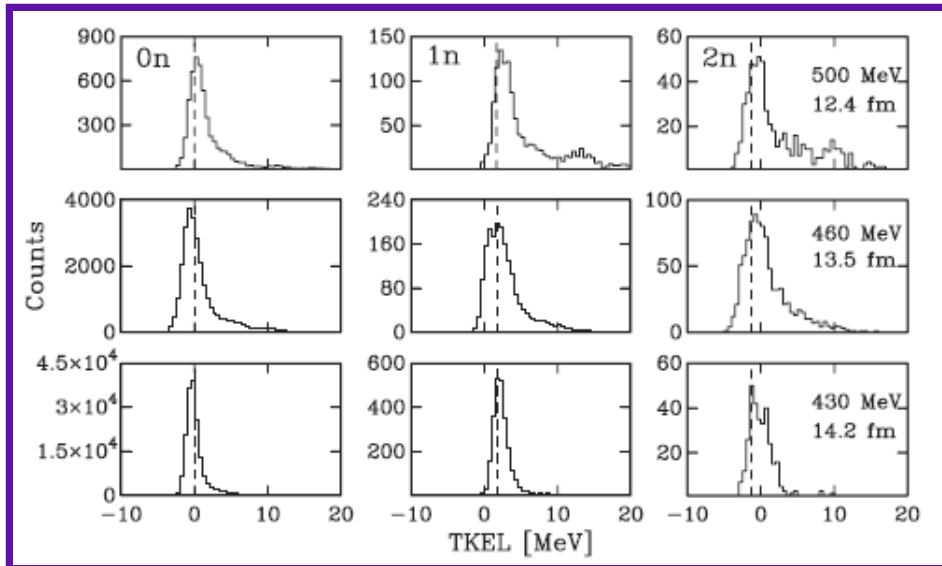
-to obtain P_{tr} : summed over all possible transitions that can be constructed from the single particle states in projectile and target
 - the set of single particle states covers a full shell below the Fermi level for ^{96}Zr and a full shell above for ^{40}Ca

Two particle transfer (semiclassical theory, microscopic calc.)
 3 terms : simultaneous, orthogonal and successive (only the successive term contributes to the transfer amplitude)

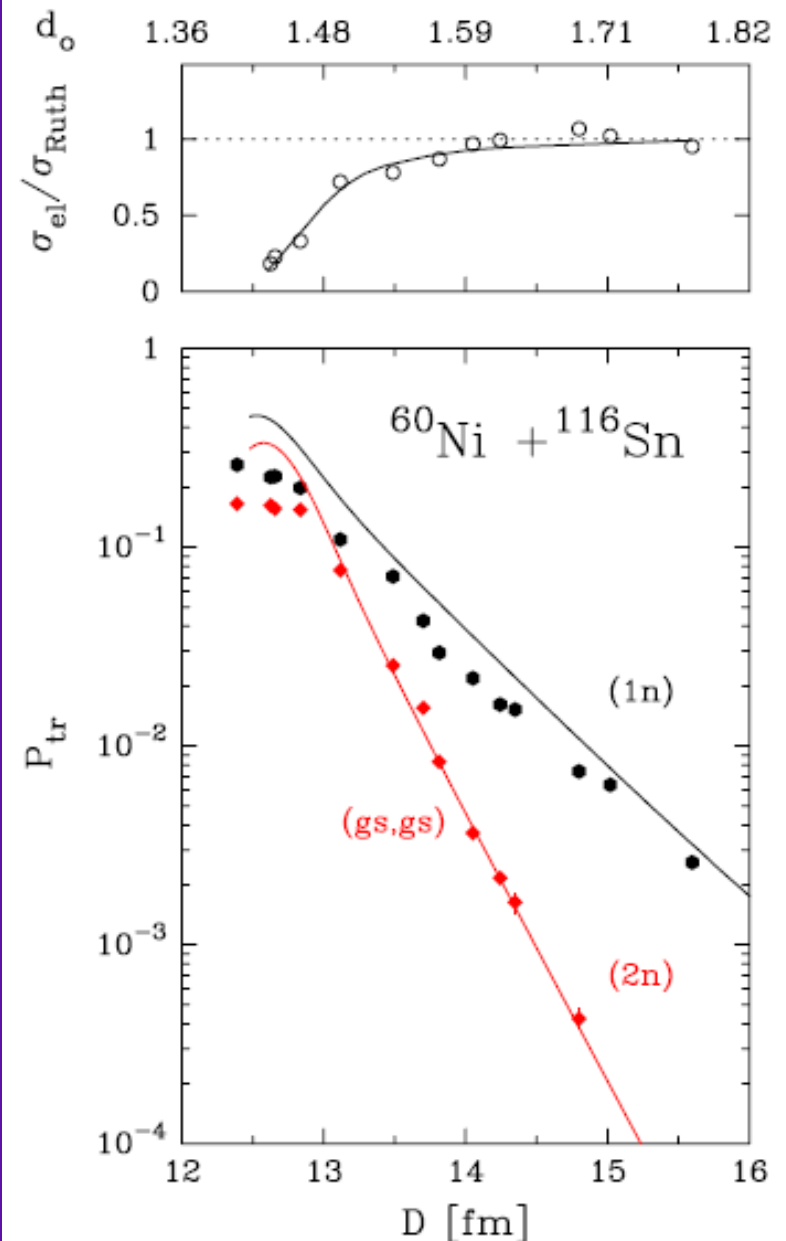
$^{60}\text{Ni} + ^{116}\text{Sn}$: neutron pair transfer far below the Coulomb barrier

The experimental transfer probabilities are well reproduced, for the **first** time with **heavy ion** reactions, in **absolute values** and in **slope** by **microscopic** calculations which incorporate nucleon-nucleon **correlations**

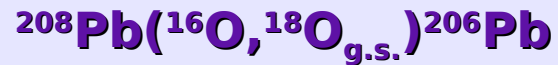
Transfer strength very close to the g.s. to g.s. transitions



D. Montanari, L. Corradi, S. Szilner, G. Pollaro et. al.,
Phys. Rev. Lett. 113 (2014) 052501



Absolute cross sections for one and two-nucleon transfer reactions

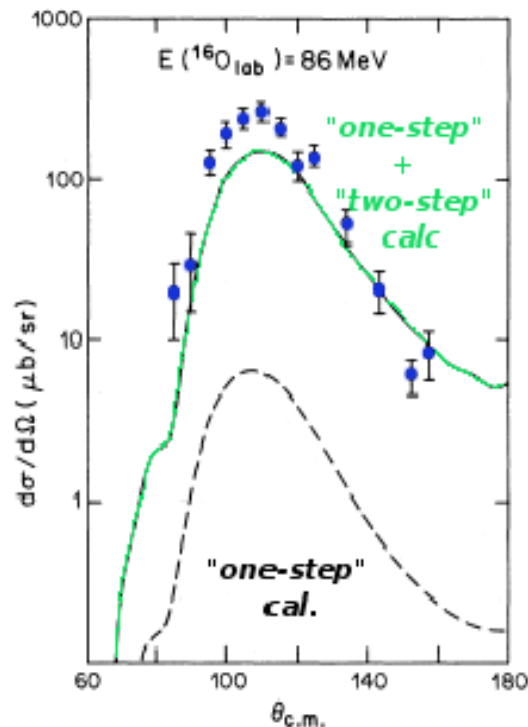


— successive+simultaneous
 - - - simultaneous

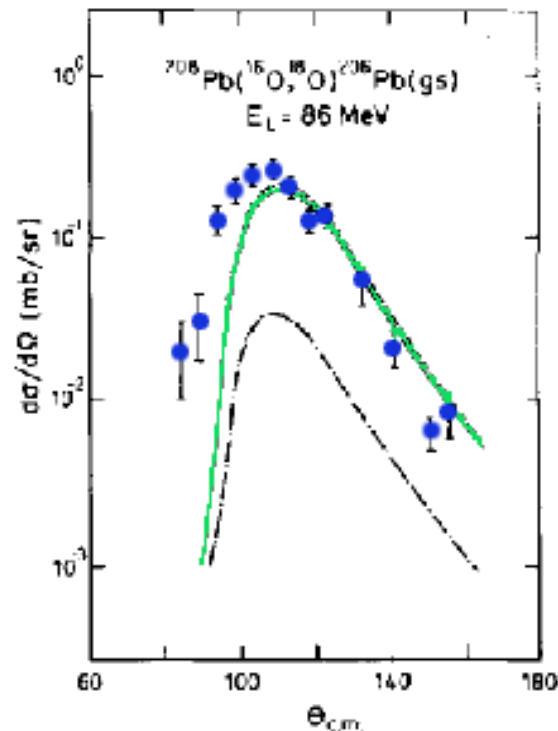
full quanto-mechanical

semi-classical

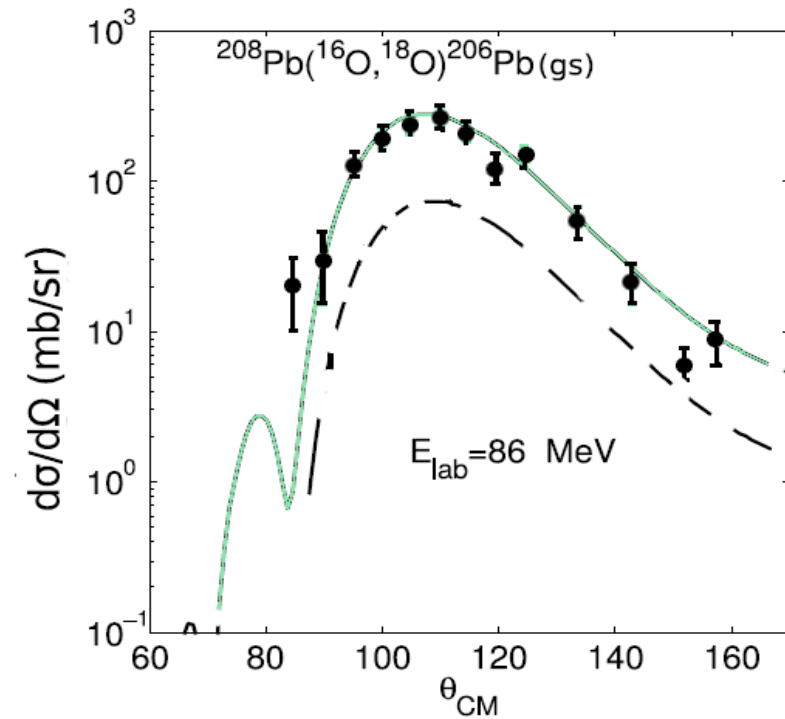
full quanto-mechanical



B.F.Bayman and J.Chen,
 Phys. Rev. C 26 (1982) 1509



E.Maglione, G.Pollarolo, A.Vitturi,
 R.A.Brogia and A.Winther
 Phys. Lett. B 162 (1985) 59



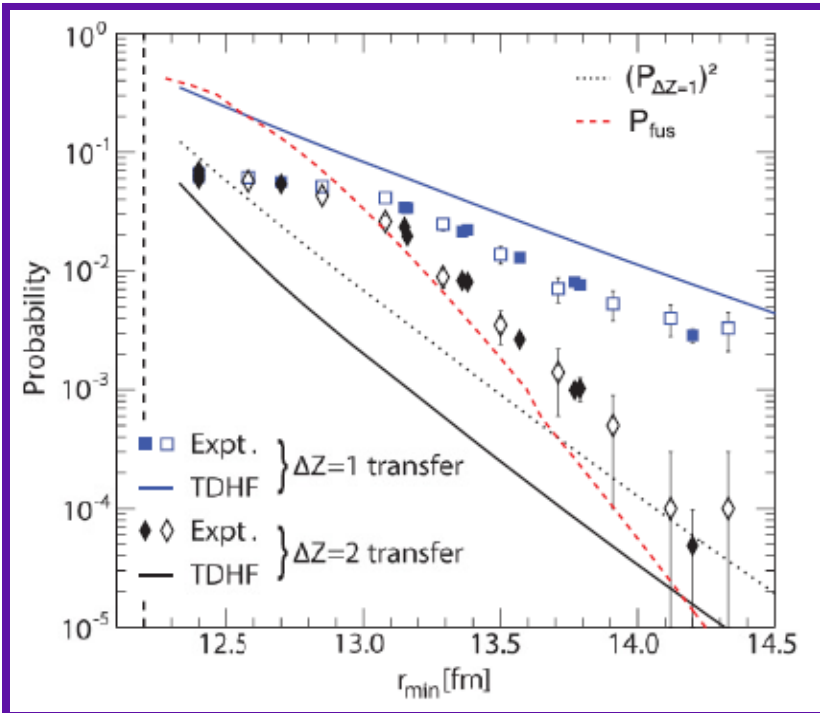
G.Potel, A.Idini, F.Barranco, E.Vigezzi
 and R.A.Brogia, Rep. Prog. Phys. 76
 (2013) 106301;

G. Potel et al, PRL 105 (2010) 172502

Sub-barrier transfer : TDHF or TDHF+BCS

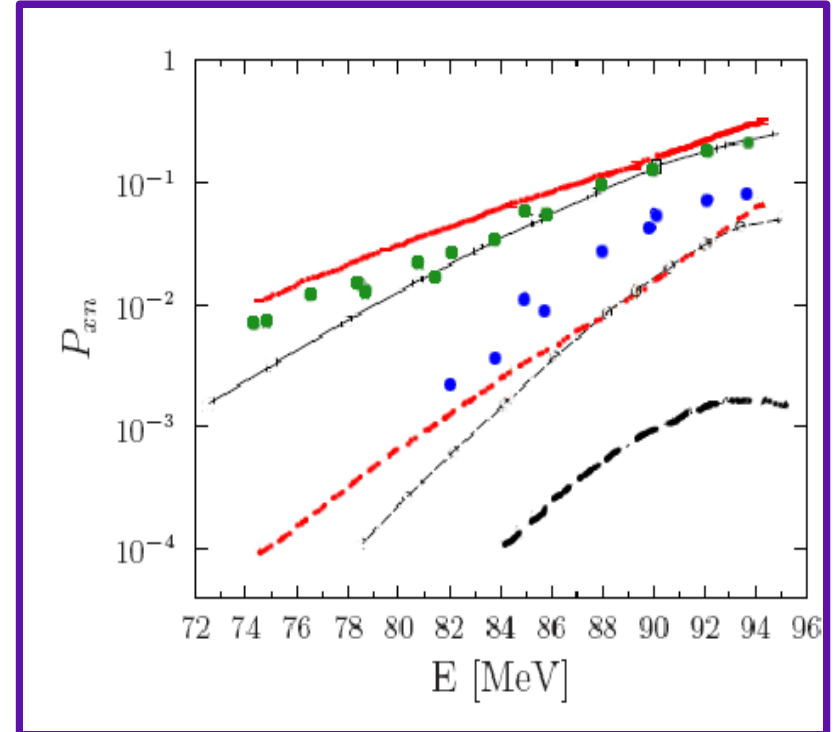
$^{16}\text{O} + ^{208}\text{Pb}$

—— (2n) TDHF
 - - - - (1n) TDHF



$^{40}\text{Ca} + ^{96}\text{Zr}$

—— (2n) TDHF+BCS
 - - - - (1n) TDHF+BCS



L.Corradi et.al., Phys. Rev. C 84 (2011) 034603
 EXP (1n) and (2n);
 (1n) c.c.; (2n) (g.s. \rightarrow g.s.)
 (g.s. \rightarrow 0+ at ~6MeV)

C.Simenel, PRL105(2010)192701

M.Evers et al, PRC84(2011)054614

G.Scamps et al., EPJ Web Conf. 86 (2015) 00042

Summary

- ✓The comparison between **data and theory**: **elementary modes** of the complex mechanism can be probed.
- ✓“large” **spectrometers coupled to** “large” **gamma arrays** are powerful tools to study the fine details of such processes.
- ✓The total, differential cross sections and individual state yield distribution reflect a strong interplay between **single-particle and collective degrees** of freedom and the **reaction dynamics**: transfer reactions are sensitive to the transferred angular momentum (good matching), the nuclear matrix element contains the spectroscopic information (both of the projectile and target).
- ✓The importance of components responsible for the couplings (phonon-single particle), and for particle **correlations** (residual interaction)
- ✓**Sub-barrier transfer** reaction measurement (nuclei interact at large distances): good probe for pair **correlations**

OUTLOOK:

gamma-particle coincidences
proton transfer channels at large D
proton rich nuclei (np correlations)
neutron rich nuclei (density dependent forces)
very heavy systems
microscopic calculations for high multipolarity states

*L. Corradi, T. Mijatović, E. Fioretto, A. Gadea, D. Montanari, A.M. Stefanini, J.J. Valiente-Dobon, E. Farnea, G. Montagnoli, C.A. Ur, S. Lunardi, C. Michelagnoli, N. Marginean, F. Haas, S. Courtin, D. Lebhertz, A. Goasduff, M.-D. Salsac
D. Jelavić Malenica, N. Soić
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Kokopelli: links distant
and diverse
communities together.



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