

# Structure of Mass Gap Between Two Spin Multiplets

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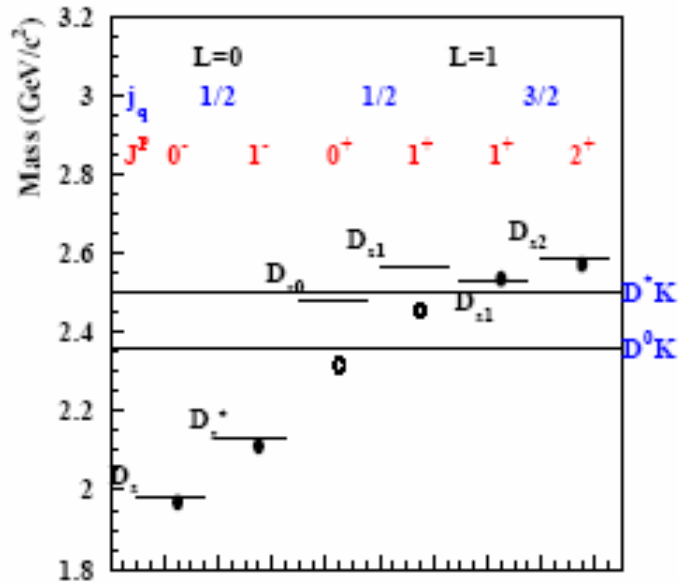
with T. Morii (Kobe U.) and K. Sudoh (KEK)

## Based on

- **Structure of mass gap between two spin multiplets,**  
hep-ph/0710.0325
- **Radial Excitations of Heavy-Light System,**  
Eur. Phys. J. C31(2007) 701
- **New Heavy-Light Mesons  $Qq\bar{q}$ ,**  
Prog. Theor. Phys. 51 (2007) 1077 (hep-ph/0605019).
- **$0^+$  and  $1^+$  States of  $B$  and  $B_s$  Mesons,**  
Phys. Lett. B606 (2005) 329 (hep-ph/0411034).
- **Spectroscopy of heavy mesons expanded in  $1/m_Q$ ,**  
Phys. Rev. D56 (1997) 5646.

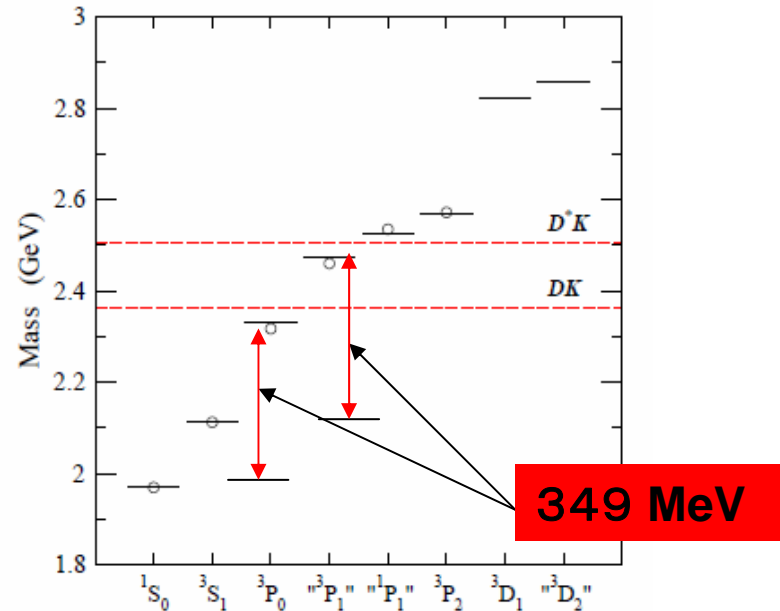
# Motivation1

- Successful prediction/reproduction of  $D_s$  mass spectra using our semi-relativistic potential model
  - Lowering  $0^+$  and  $1^+$  of  $D_0^*$  (2308) and  $D_1'$  (2427) compared with other potential models



prediction by conventional potential model (Godfrey & Isgur, PRD43, 1679 (1991))

2007/10/09

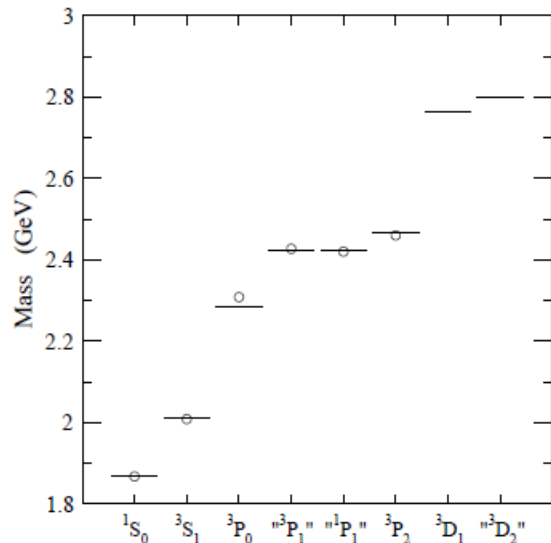


prediction by our semi-relativistic potential model (Prog. Theor. Phys.117 (2007) 1077)

Hadron07

# Other Mass Spectra

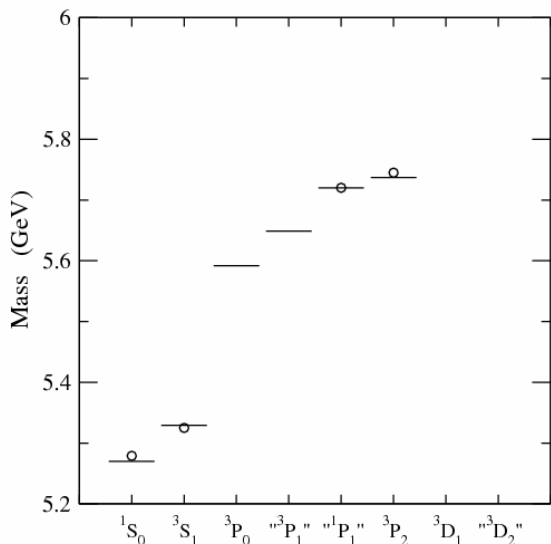
**D**



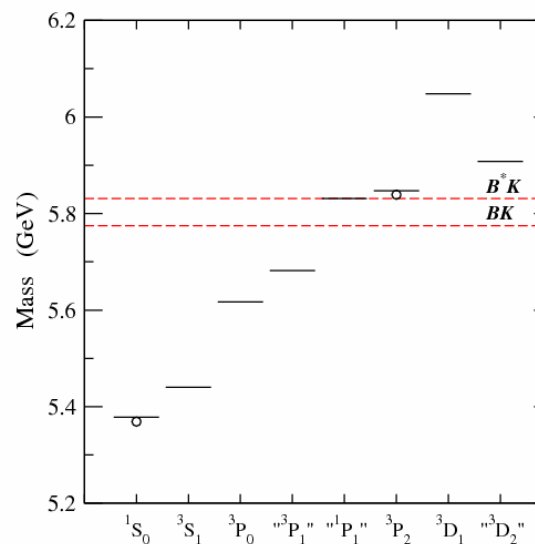
Successful reproduction of the following spectra

- $D_0^*$ (2308) and  $D_1'$ (2427) by Belle
- $D_{s0}$ (2860) and  $D_s^*$ (2715) by BaBar & Belle (n=2;  $0^+$  and  $1^-$  states of  $D_s$ )
- $B_1$ (5720) and  $B_2^*$ (5745) by D0 (1+ and 2+ states of B)
- $B_{s2}^*$ (5839) by D0 ( $2^+$  state of  $B_s$ )

**B**



**$B_s$**



# Motivation2

- Successful prediction of the mass gap using an effective Lagrangian approach (Bardeen & Hill)
  - $\Delta M(m_c) = 338 \text{ MeV}$
  - Later corrected to the experimentally observed value by B+Eichten+H using **S(3)**
  - **Wonderful agreement** with the experiments (349 MeV) !
  - Is this prediction?
  - Is this describing the underlying physics?
  - Is an effective theory constructed from a four-fermi theory a **true theory**?
  - Their prediction has no light quark flavor dependency.
  - BEH tried to predict the **mass gap for D as 255 MeV** but the experiment gives more than 400 MeV.  
 **$D(0^+) - D(0^-) = 2308 - 1867 = 441 \text{ MeV}$**   
 **$D(1^+) - D(1^-) = 2427 - 2008 = 419 \text{ MeV}$**
  - Hence there is no way to distinguish their prediction with that for heavy D, B, or Bs mesons
- If finding the mass gap eq. depending on the light quark mass, there might be a hint to get to the **true** underlying physics.

# *D mesons*

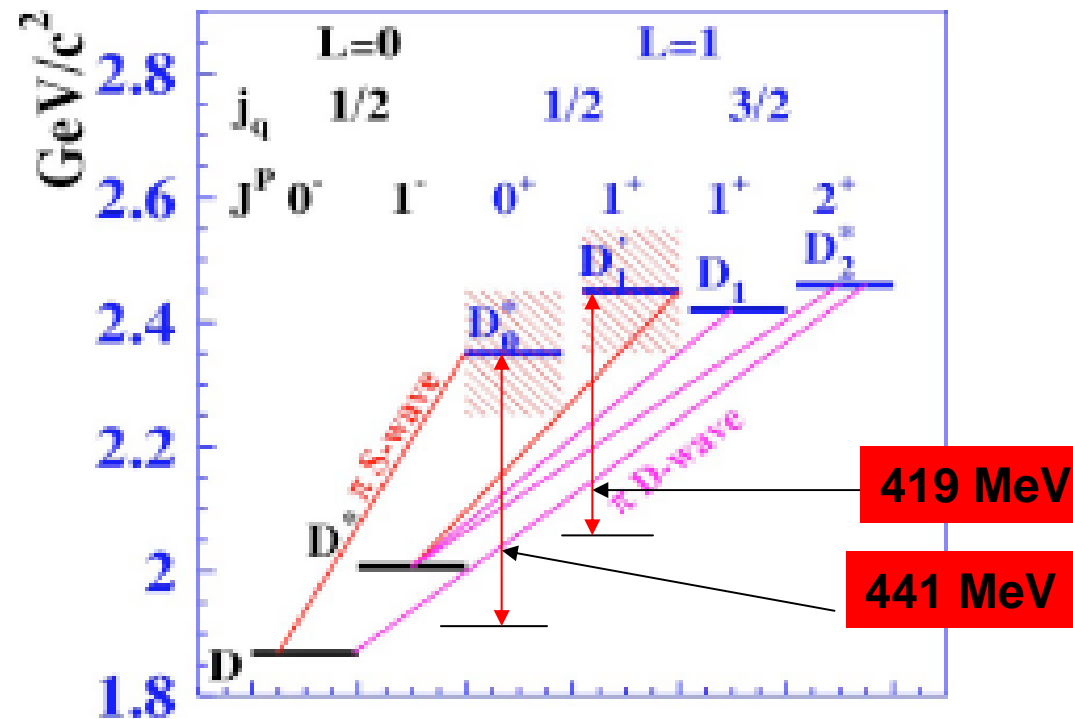


FIG. 1. Spectroscopy of *D*-meson excitations. The lines show possible single pion transitions.

# Our Model Hamiltonian for $Q\bar{q}$ System

- Start with an effective Hamiltonian for 2-body bound states

$$H = (\vec{\alpha}_q \cdot \vec{p}_q + \beta_q m_q) + (\vec{\alpha}_Q \cdot \vec{p}_Q + \beta_Q m_Q) + \beta_q \beta_Q S \\ + \left[ 1 - \frac{1}{2} \left\{ \vec{\alpha}_q \cdot \vec{\alpha}_Q + (\vec{\alpha}_q \cdot \vec{n})(\vec{\alpha}_Q \cdot \vec{n}) \right\} \right] V$$

$$S(r) = \frac{r}{a^2} + b, \quad V(r) = -\frac{4}{3} \frac{\alpha_s}{r}$$

Semi-relativistic approach to heavy meson

→ Foldy-Wouthuysen-Tani (FWT) transformation  
to the heavy quark ~1/m<sub>Q</sub> expansion

- Heavy Meson System (eigenvalue equation): mass of a heavy meson =  $m_Q + E^l$

$$H\psi_l = E^l\psi_l$$

$$H = H_{\text{FWT}} - m_Q = m_Q H_{-1} + H_0 + \frac{1}{m_Q} H_1 + \frac{1}{m_Q^2} H_2 + \dots$$

$$E^l = E_0^l + \frac{1}{m_Q} E_1^l + \frac{1}{m_Q^2} E_2^l + \dots$$

$$\psi_l = \psi_{l0} + \frac{1}{m_Q} \psi_{l1} + \frac{1}{m_Q^2} \psi_{l2} + \dots$$

# Effective Hamiltonian Expanded in $1/m_Q$

$$(H_{\text{FWT}} - m_Q) \otimes \psi_{\text{FWT}} = \tilde{E} \psi_{\text{FWT}}$$

$$H_{\text{FWT}} - m_Q = H_{-1} + H_0 + H_1 + H_2$$

$$H_{-1} = -(1 + \beta_Q) m_Q$$

$$H_0 = \vec{\alpha}_q \cdot \vec{p} + \beta_q m_q - \beta_q \beta_Q S + \left\{ 1 + \frac{1}{2} \vec{\alpha}_q \cdot \vec{\alpha}_Q + (\vec{\alpha}_q \cdot \vec{n})(\vec{\alpha}_Q \cdot \vec{n}) \right\} V$$

$$H_1 = -\frac{1}{2m_Q} \beta_Q \vec{p}^2 + \frac{1}{m_Q} \beta_q \vec{\alpha}_Q \cdot \left( \vec{p} + \frac{1}{2} \vec{q} \right) S + \frac{1}{2m_Q} \vec{\gamma}_Q \cdot \vec{q} V$$

$$-\frac{1}{2m_Q} \left[ \beta_Q \left( \vec{p} + \frac{1}{2} \vec{q} \right) + i \vec{q} \times \beta_Q \vec{\Sigma}_Q \right] \cdot \left[ \vec{\alpha}_q + (\vec{\alpha}_q \cdot \vec{n}) \vec{n} \right] V$$

$$H_2 = \frac{1}{2m_Q} \beta_q \beta_Q \left( \vec{p} + \frac{1}{2} \vec{q} \right)^2 S - \frac{i}{4m_Q^2} \vec{q} \times \vec{p} \cdot \beta_q \beta_Q \vec{\Sigma}_Q S - \frac{1}{8m_Q^2} \vec{q}^2 V - \frac{i}{4m_Q^2} \vec{q} \times \vec{p} \cdot \vec{\Sigma}_Q V$$

$$-\frac{1}{8m_Q^2} \left\{ (\vec{p} + \vec{q})(\vec{\alpha}_Q \cdot \vec{p}) + \vec{p}[\vec{\alpha}_Q \cdot (\vec{p} + \vec{q})] + i \vec{q} \times \vec{p} \gamma_5 \right\} \cdot \left[ \vec{\alpha}_q + (\vec{\alpha}_q \cdot \vec{n}) \vec{n} \right] V$$

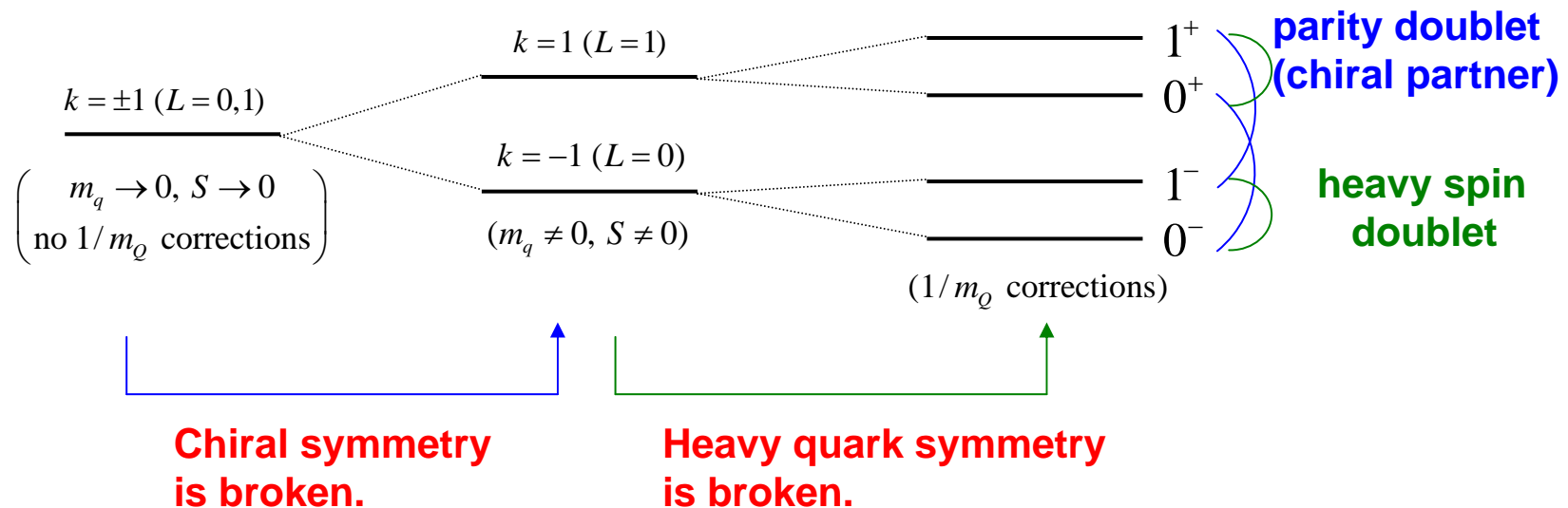
$$\because \vec{p} = \vec{p}_q = -\vec{p}_Q, \quad \vec{p}' = \vec{p}'_q = -\vec{p}'_Q, \quad \vec{q} = \vec{p}' - \vec{p}$$

spin-spin interaction

LS force

# Patterns of Symmetry Breaking

- mass gap between degenerate masses
  - In the heavy quark limit, masses for  $0^-$  and  $1^-$ , and also  $0^+$  and  $1^+$  are degenerate
  - In the chiral limit, all these four states are degenerate (S has mass dimension)

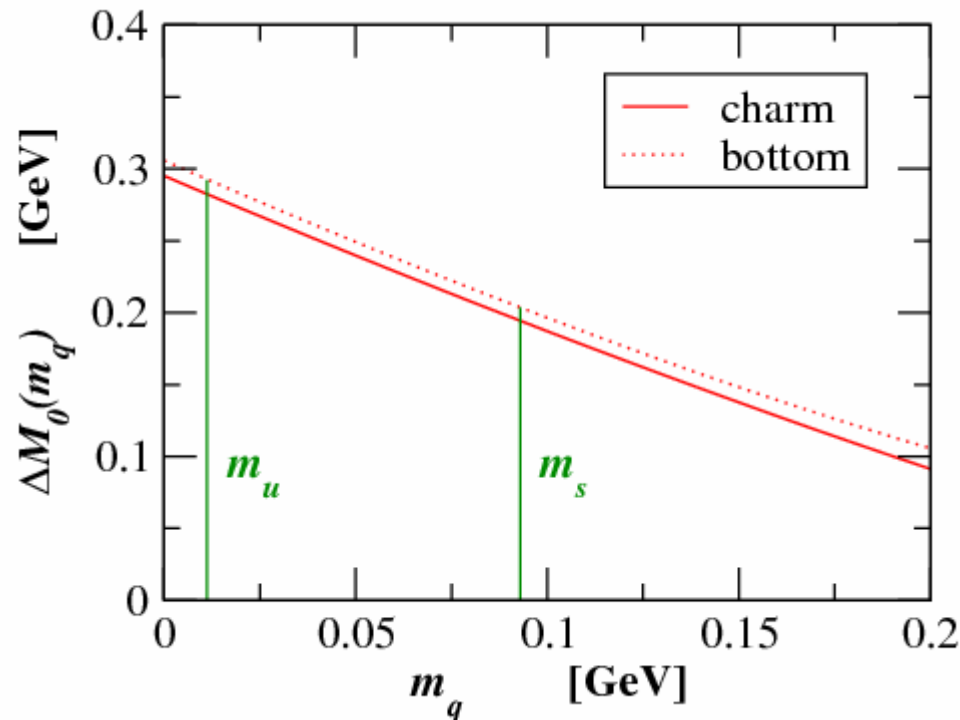




# Numerical Results of $\Delta M$

(without  $1/m_Q$ )

- mass gap between degenerate masses (figure at the lowest order, i.e., without  $1/m_Q$  corrections)
  - In the heavy quark limit, masses for  $0^-$  and  $1^-$ , and also  $0^+$  and  $1^+$  are degenerate
  - Light quark mass dependence of  $\Delta M(m_q) = M(0^+) - M(0^-) = M(1^+) - M(1^-)$



# Numerical Formula for $\Delta M$

- First derive numerical mass gap between two spin (degenerate) multiplets

$$\Delta M_0 = M(0^+) - M(0^-) = M(1^+) - M(1^-) \approx \Lambda_Q - m_q$$

$$\Lambda_Q = 300 \text{ MeV} \quad \text{and } m_q \text{ light quark mass}$$

- More precisely

$$\Delta M_0 = M(0^+) - M(0^-) = M(1^+) - M(1^-) = g_0 \Lambda_Q - g_1 m_q$$

$$g_0 = 0.984, \quad g_1 = 1.080 \quad \text{for } D/D_s$$

$$g_0 = 1.017, \quad g_1 = 1.089 \quad \text{for } B/B_s$$

- From these equations, the mass gap for a degenerate system is essentially given by (Note negative dependency on  $m_q$ )

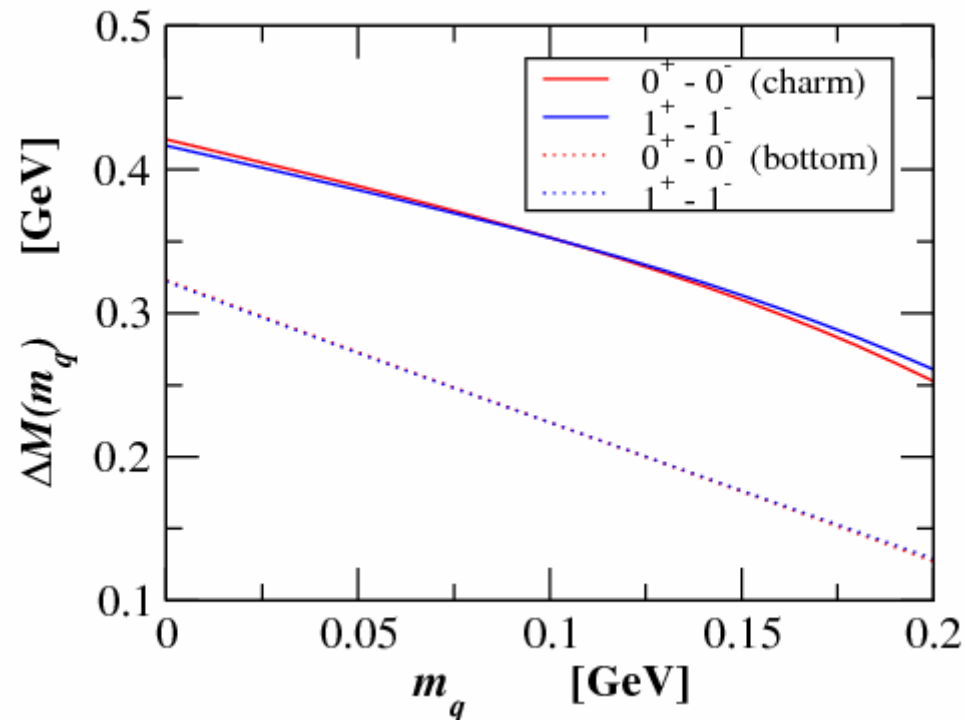
$$\Delta M_0 = \Lambda_Q - m_q$$

# Numerical Results of $\Delta M$

(first order)

- mass gap between masses of spin multiplets (figure at the next order, i.e., with  $1/m_Q$  corrections included)
  - Light quark mass dependence is given by the equation

$$\Delta M(m_q) = M(0^+) - M(0^-) \simeq M(1^+) - M(1^-) \simeq g_0 \Lambda_Q - g_1 m_q + \frac{c + d \cdot m_q}{m_Q}$$



# Prediction

- mass gap between masses of spin multiplets
  - Light quark mass dependence is given by the equation

$$\Delta M(m_q) = M(0^+) - M(0^-) \simeq M(1^+) - M(1^-) \simeq g_0 \Lambda_Q - g_1 m_q + \frac{c + d \cdot m_q}{m_Q}$$

- where

$$c = 1.28 \times 10^5 \text{ MeV}^2$$

$$d = 4.26 \times 10^2 \text{ MeV}$$

this predicts

$$B(0^+) - B(0^-) \approx B(1^+) - B(1^-) \approx 322 \text{ MeV}$$

$$B_s(0^+) - B_s(0^-) \approx B_s(1^+) - B_s(1^-) \approx 240 \text{ MeV}$$

# Numerical Values 1

TABLE I: Optimal values of parameters.

Parameters	$\alpha_a^c$	$\alpha_a^b$	$a$ (GeV $^{-1}$ )	$b$ (GeV)
	$0.261 \pm 0.001$	$0.393 \pm 0.003$	$1.939 \pm 0.002$	$0.0749 \pm 0.0020$
$m_{u,d}$ (GeV)		$m_s$ (GeV)	$m_c$ (GeV)	$m_b$ (GeV)
	$0.0112 \pm 0.0019$	$0.0929 \pm 0.0021$	$1.032 \pm 0.005$	$4.639 \pm 0.005$
# of data	# of parameter		total $\chi^2$ /d.o.f	
18	8		107.55	

TABLE II: Degenerate masses of model calculations and their mass gap between  $0^+(1^+)$  and  $0^-(1^-)$  for  $n = 1$ .

	$M_0(D)$	$M_0(D_s)$	$M_0(B)$	$M_0(B_s)$
$0^-/1^-$	1784	1900	5277	5394
$0^+/1^+$	2067	2095	5570	5598
$0^+(1^+) - 0^-(1^-)$	283	195	293	204

Decreasing as  
D  $\rightarrow$  Ds  
(u  $\rightarrow$  s)

TABLE III: Model calculations of the mass gap. Values in brackets are taken from the experiments. Units are MeV.

Mass gap ( $n = 1$ )	$\Delta M(D)$	$\Delta M(D_s)$	$\Delta M(B)$	$\Delta M(B_s)$
$0^+ - 0^-$	414 (441)	358 (348)	322	239
$1^+ - 1^-$	410 (419)	357 (348)	320	242

( $n = 2$ )	$\Delta M(D)$	$\Delta M(D_s)$	$\Delta M(B)$	$\Delta M(B_s)$
$0^+ - 0^-$	308	274	206	160
$1^+ - 1^-$	350	327	216	171

# Numerical Values2

TABLE IV:  $D/D_s$  meson mass spectra for both the calculated and experimentally observed ones. Units are MeV.

$^{2s+1}L_J(J^P)$	$M_{\text{calc}}(D)$	$M_{\text{obs}}(D)$	$M_{\text{calc}}(D_s)$	$M_{\text{obs}}(D_s)$
$^1S_0(0^-)$	1869	1867	1967	1969
$^3S_1(1^-)$	2011	2008	2110	2112
$^3P_0(0^+)$	2283	2308	2325	2317
$^3P_1(1^+)$	2421	2427	2467	2460

TABLE V:  $B/B_s$  meson mass spectra for both the calculated and experimentally observed ones. Units are MeV.

$^{2s+1}L_J(J^P)$	$M_{\text{calc}}(B)$	$M_{\text{obs}}(B)$	$M_{\text{calc}}(B_s)$	$M_{\text{obs}}(B_s)$
$^1S_0(0^-)$	5270	5279	5378	5369
$^3S_1(1^-)$	5329	5325	5440	—
$^3P_0(0^+)$	5592	—	5617	—
$^3P_1(1^+)$	5649	—	5682	—

recent experiments find

$$B(1^+) = 5720(5720), B(2^+) = 5737(5745)$$

$$B_s(2^+) = 5847(5839) \text{ MeV}$$

numbers in brackets are from expt.

# Explanation due to Effective Theory

- Lets try to explain the light quark dependency of  $\Delta M$  by using the chiral effective theory with heavy quark symmetry

(Bardeen et al., PRD68 (2003) 054024)

- mass gap is given by

Parity doublets :  $\mathcal{H} = (0^-, 1^-)$ ,  $\mathcal{H}' = (0^+, 1^+)$

$$L_{mass} = \frac{g_\pi}{4} \left[ \text{Tr}(\bar{\mathcal{H}}' \tilde{\sigma} \mathcal{H}') - \text{Tr}(\bar{\mathcal{H}} \tilde{\sigma} \mathcal{H}) \right]$$

$$\langle \tilde{\sigma} \rangle = f_\pi I_3 \quad \tilde{\sigma} = \sigma \oplus m_q$$

**Plus** sign instead of **minus** sign

- Goldberger-Treiman relation**

$$\Delta M(m_q) = M_X(1^+) - M_X(1^-) = M_X(0^+) - M_X(0^-) = g_\pi f_\pi$$

$$\Delta M(m_c) = 349, \Delta M(\infty) = 338 \text{ for } D_s$$

349 MeV is an input and cannot be identified to be DsJ

# Discussions and Summary

- Global Flavor SU(3) Recovery (Dmitrasinovic, PRL 94 (2005) 162002)
  - He tried to explain by tetraquark.

$$M(D(0^+)) = 2308, \quad M(D_s(0^+)) = 2317$$

$$M(D(1^+)) = 2427, \quad M(D_s(1^+)) = 2460$$

- This can be nicely explained by our mass gap expression, i.e., monotonously decreasing behavior of the mass gap between two spin multiplets.
- Mass gap of Heavy Baryons
  - $QQ = 3^* \rightarrow QQq = 3^* \times 3$  (like a heavy meson)
  - May apply our formula. For instance,

$$\Delta M = M'(ccq) - M(ccq) = \Lambda_Q - m_q + \frac{c + d \cdot m_q}{m_{Q1} + m_{Q2}}$$

- We have derived an empirical formula for the mass gap between two spin multiplets, i.e.,  $(0^-, 1^-)$  and  $(0^+, 1^+)$ , which is given above.
  - May give us a hint to find the underlying (true?) physics.



# Backups

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- Have time, then show these

# Intuitive Explanation of $\Delta M$

$m_q$  dependence of the mass gap:

$$\begin{aligned}\Delta M &= M^{1(0^+)} - M^{-1(0^-)} = M^{1(1^+)} - M^{-1(1^-)} \\ &= \int \frac{d^3x}{4\pi r^2} \left\{ \Phi_1^\dagger(r) \begin{pmatrix} m_q + S + V & -\partial_r + \frac{1}{r} \\ \partial_r + \frac{1}{r} & -m_q - S + V \end{pmatrix} \Phi_1(r) - \Phi_{-1}^\dagger(r) \begin{pmatrix} m_q + S + V & -\partial_r - \frac{1}{r} \\ \partial_r - \frac{1}{r} & -m_q - S + V \end{pmatrix} \Phi_{-1}(r) \right\} \\ &= \int dr \left[ \Phi_1^\dagger(r) K_1 \Phi_1(r) - \Phi_{-1}^\dagger(r) K_{-1} \Phi_{-1}(r) \right] + m_q \int dr \left[ \Phi_1^\dagger(r) \beta \Phi_1(r) - \Phi_{-1}^\dagger(r) \beta \Phi_{-1}(r) \right].\end{aligned}$$

where

$$\Phi_k(r) = \begin{pmatrix} u_k(r) \\ v_k(r) \end{pmatrix}, \quad K_k = \begin{pmatrix} S(r) + V(r) & -\partial_r + \frac{k}{r} \\ \partial_r + \frac{k}{r} & -S(r) + V(r) \end{pmatrix}$$

$k=-1(+1)$  corresponds to  $L=0(L=1)$  states, then  $k=+1$  is more relativistic than  $k=-1$ , which means a lower component  $v_1(r)$  is larger than  $v_{-1}(r)$ . Hence  $(u_1)^2 - (v_1)^2 = \Phi_1^\dagger(r) \beta \Phi_1(r)$  becomes smaller than  $\Phi_{-1}^\dagger(r) \beta \Phi_{-1}(r)$ .

Thus, the coefficient of  $m_q$  becomes negative.

TABLE II: Degenerate masses of model calculations and their mass gap between  $0^+(1^+)$  and  $0^-(1^-)$  for  $n = 1$ .

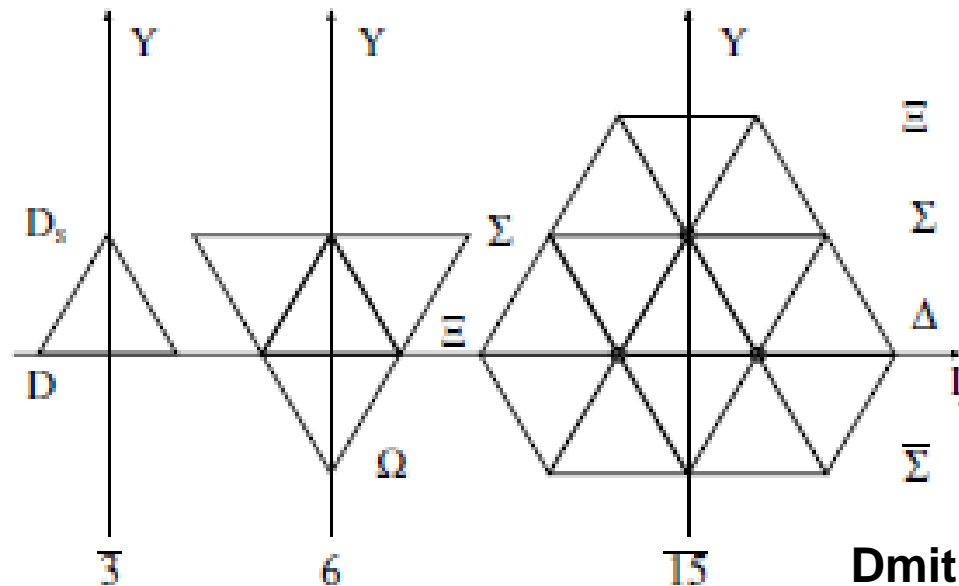
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$0^+(1^+) - 0^-(1^-)$	283	195	293	204

# Tetra-quark model

$$|D^0 \subset \bar{\mathbf{3}}_A\rangle = \frac{1}{2} |c(s(\bar{u}\bar{s} - \bar{s}\bar{u}) - d(\bar{d}\bar{u} - \bar{u}\bar{d}))\rangle,$$

$$|D^+ \subset \bar{\mathbf{3}}_A\rangle = \frac{1}{2} |c(s(\bar{d}\bar{s} - \bar{s}\bar{d}) - u(\bar{d}\bar{u} - \bar{u}\bar{d}))\rangle,$$

$$|D_s^+ \subset \bar{\mathbf{3}}_A\rangle = \frac{1}{2} |c(u(\bar{u}\bar{s} - \bar{s}\bar{u}) - d(\bar{d}\bar{s} - \bar{s}\bar{d}))\rangle,$$



Dmitrasinovic, PRL,  
94 (2005) 162002