

Heavy ion collisions phenomenology overview

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Hadron 07
INFN Laboratori Nazionale di Frascati, October 2007

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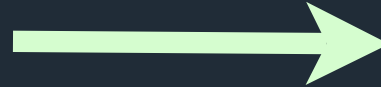
Fundamental Interactions
Searches – Higgs, SUSY,
extra-dimensions...



Increase
energy density

simple systems

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Collective properties
of the fundamental
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Increase
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“less simple” systems

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Increase
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“less simple” systems

What are the building blocks and how they organize?

- Spectroscopy
- Heavy-ion collisions

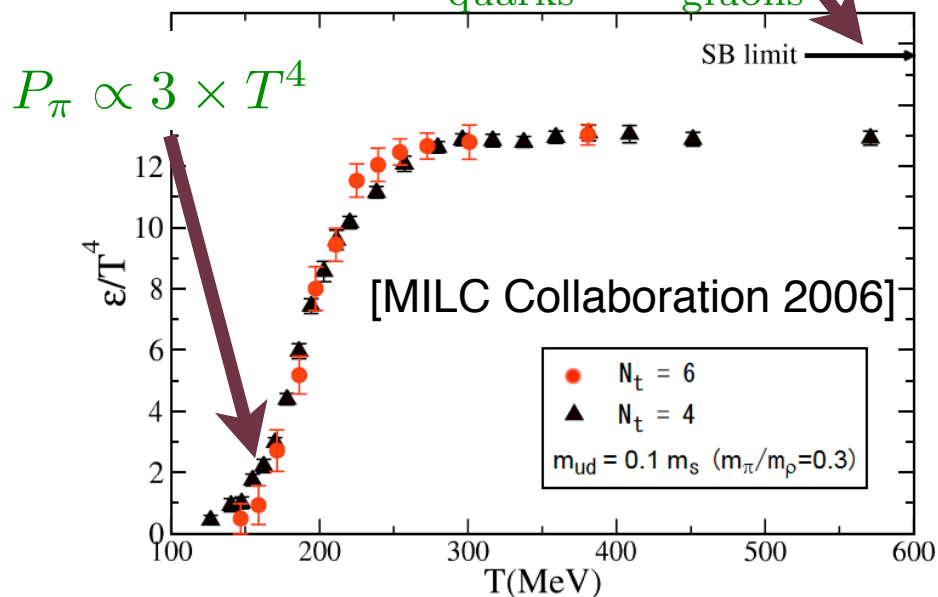
QCD has a rich dynamical structure

- ⇒ Two broken symmetries in the QCD vacuum
 - ↘ confinement
 - ↘ chiral symmetry is broken
- ⇒ Restored at high-temperatures ← asymptotic freedom

Equation of state

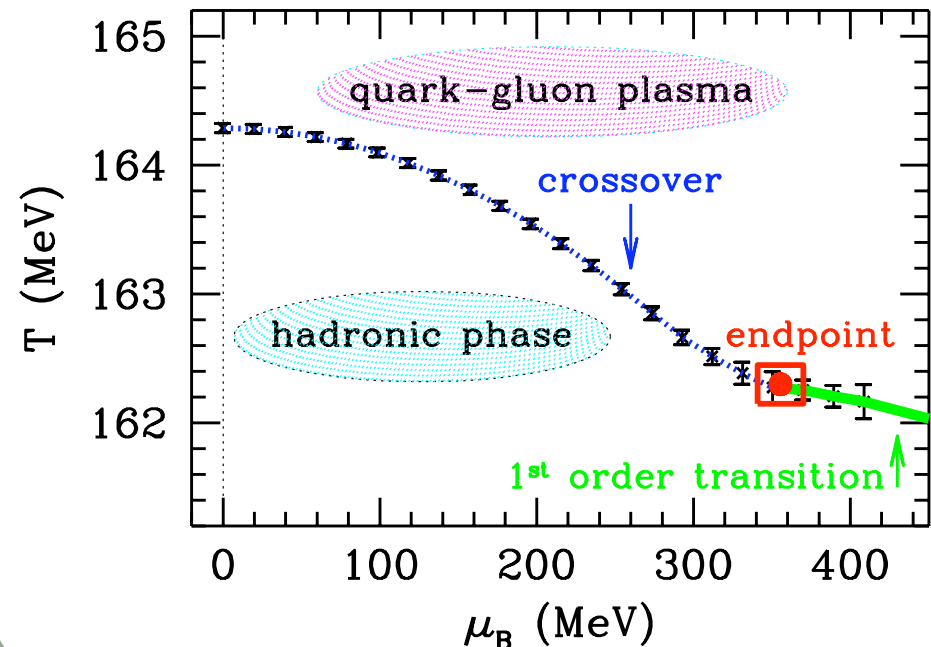
$$P_{QGP} \propto (\underbrace{2 \times 2 \times 3}_{\text{quarks}} + \underbrace{2 \times 8}_{\text{gluons}}) \times T^4$$

$$P_\pi \propto 3 \times T^4$$



Phase diagram

[Fodor, et al. 2004]



What do we expect to learn?

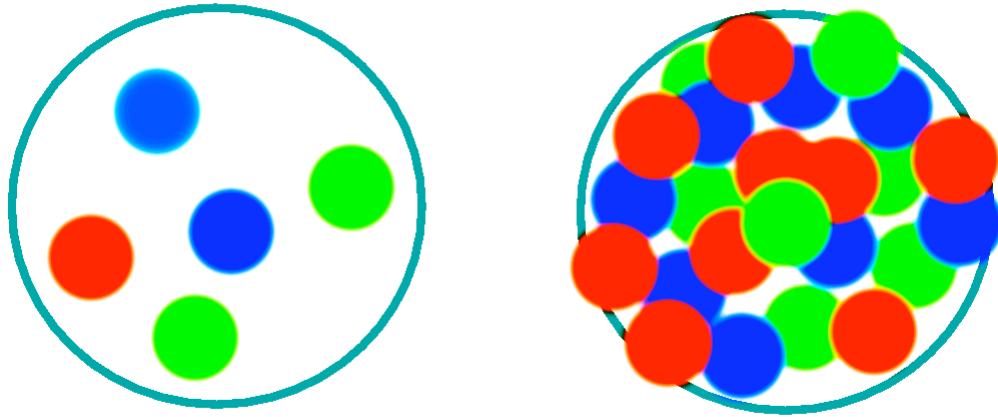
Specific questions in heavy-ion collisions

- ⇒ What is the initial state of the system and how is it produced?
 - ⇒ What is the structure of the colliding objects?
 - ⇒ What is the asymptotic limit of QCD?
- ⇒ What is the mechanism of thermalization?
 - ⇒ How is thermal equilibrium reached?
 - ⇒ What is the temperature of the created system?
- ⇒ What are the properties of the produced medium?
 - ⇒ How to measured them? – signals
 - ⇒ What is the relation with lattice QCD?

***(high-density)
Initial state***

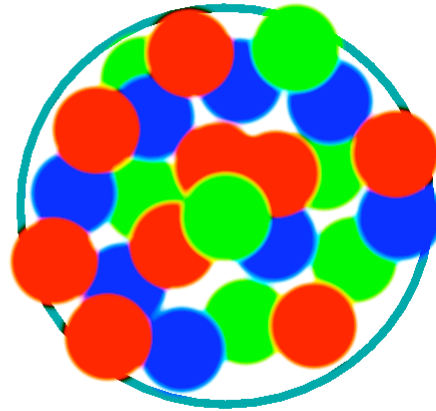
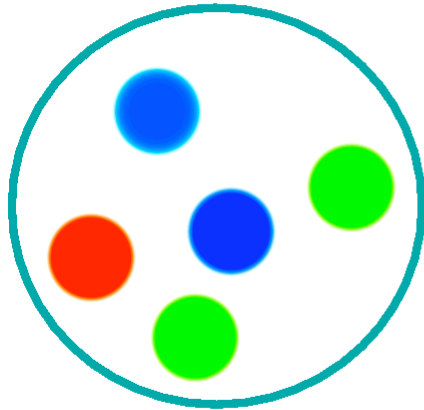
Saturation of partonic densities: picture

Saturation scale when interaction probability becomes $\mathcal{O}(1)$



Saturation of partonic densities: picture

Saturation scale when interaction probability becomes $\mathcal{O}(1)$



transverse area of the gluon

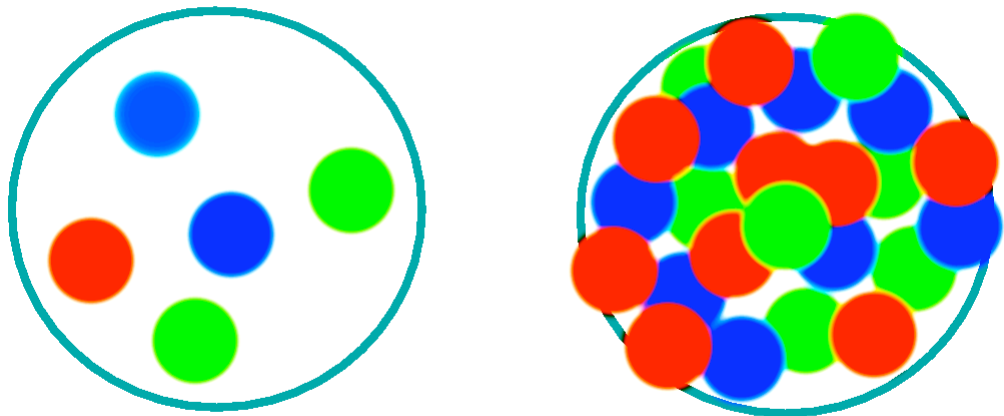
$$\alpha_s \frac{1}{Q_{\text{sat}}^2} A N_g(x, Q_{\text{sat}}^2) \sim \pi R_A^2$$

transverse area of the nucleus

$$R_A \sim A^{1/3}$$

Saturation of partonic densities: picture

Saturation scale when interaction probability becomes $\mathcal{O}(1)$



increasing energy (decreasing x)

$$N_g \sim \frac{1}{x^\lambda} \implies Q_{\text{sat}}^2 \sim \frac{A^{1/3}}{x^\lambda}$$

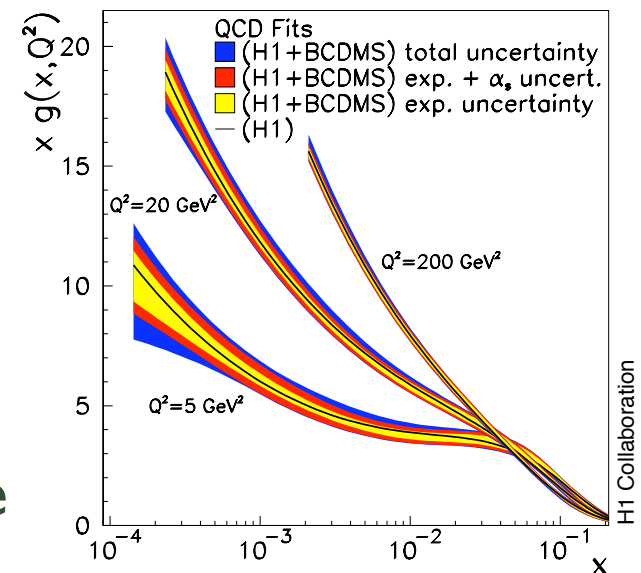
Strong fields and large occupation numbers.
Semiclassical approach possible:
Color Glass Condensate

transverse area of the gluon

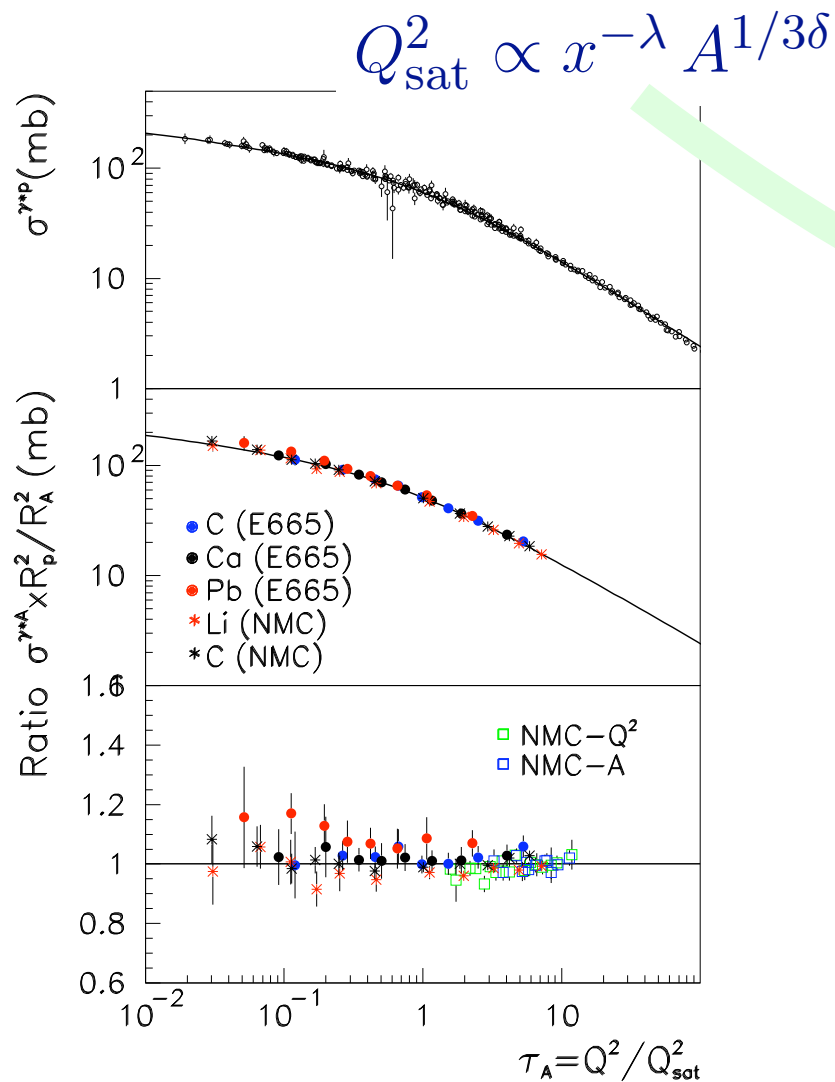
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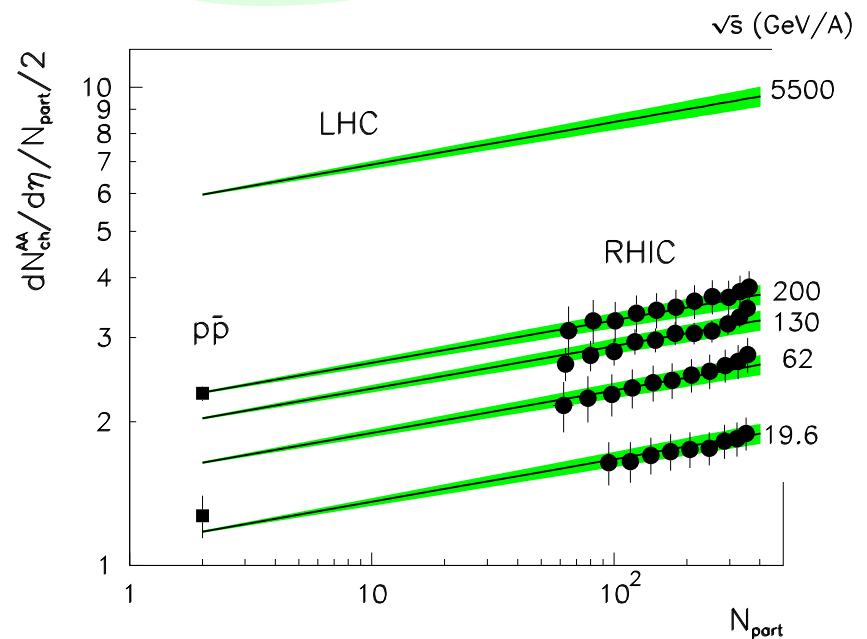
$$R_A \sim A^{1/3}$$



Geometric scaling and data



$$\frac{1}{N_{\text{part}}} \left. \frac{dN^{AA}}{d\eta} \right|_{\eta \sim 0} = N_0 \sqrt{s}^\lambda N_{\text{part}}^{\frac{1-\delta}{3\delta}}$$



Stasto, Golec-Biernat, Kwiecinski 2001

Armesto, Salgado, Wiedemann 2004

Kharzeev, Levin, McLerran, Nardi 2000...

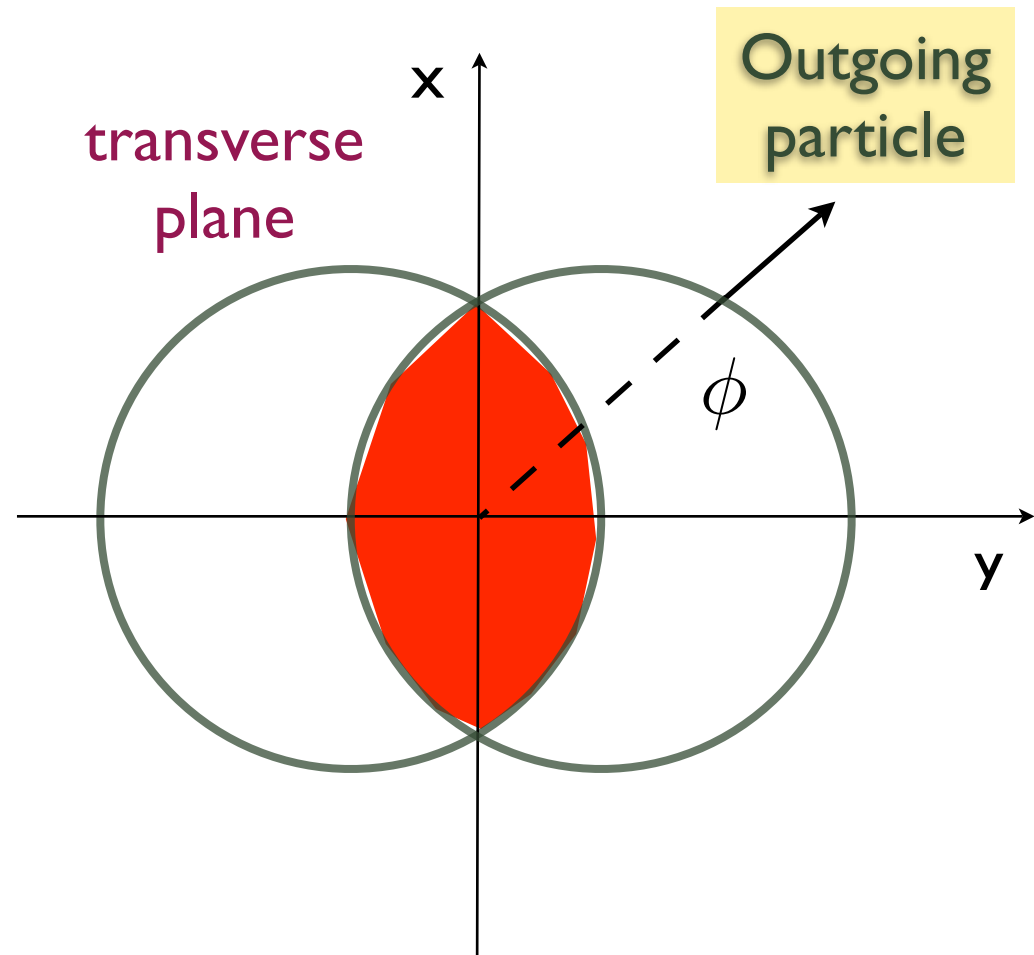
Armesto, Salgado, Wiedemann 2004

Checks of hydrodynamical evolution (thermalization)

The essential measurement for hydro

⇒ Recall the Euler equation

$$\frac{d\beta}{dt} = -\frac{c^2}{\epsilon + P} \nabla P$$



The essential measurement for hydro

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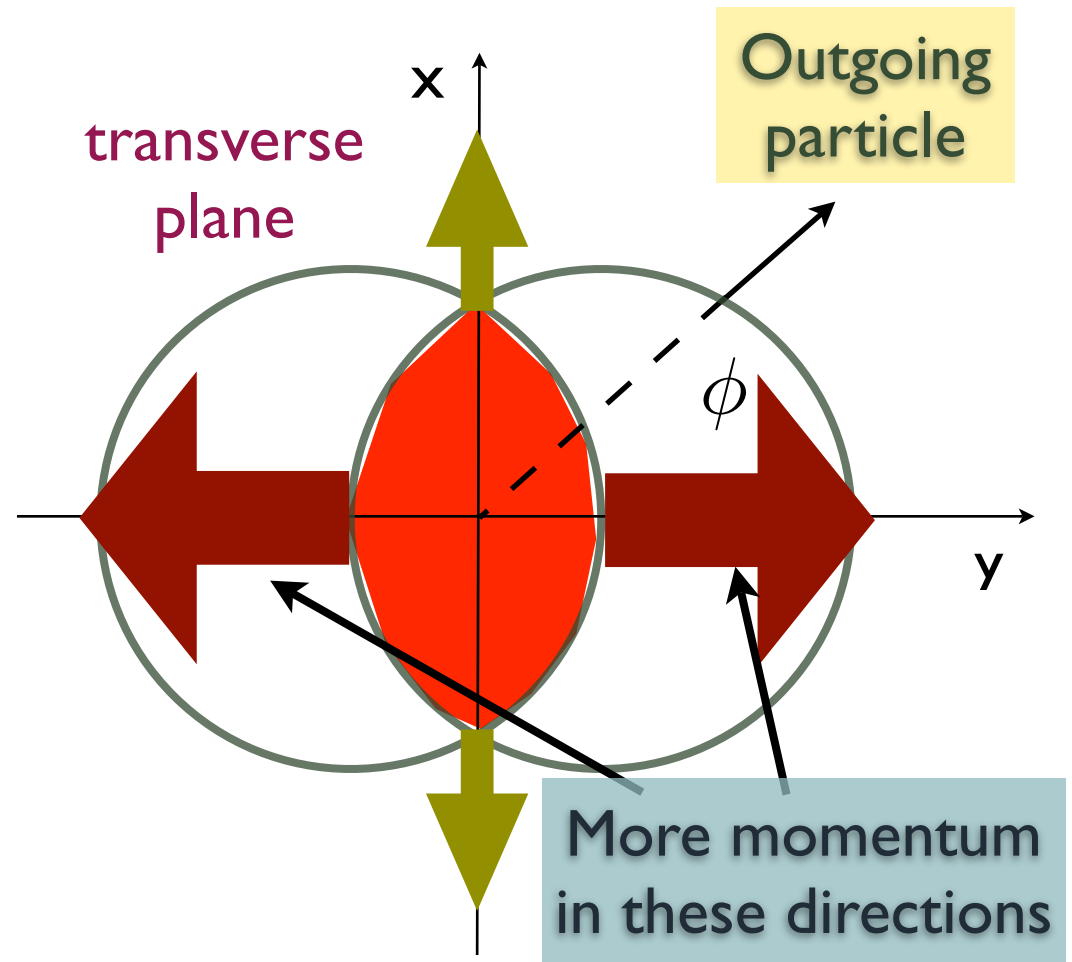
$$\frac{d\beta}{dt} = -\frac{c^2}{\epsilon + P} \nabla P$$

$$\epsilon = 3P \implies \nabla_x P < \nabla_y P$$

⇒ Elliptic flow normally measured by the second term in the Fourier expansion

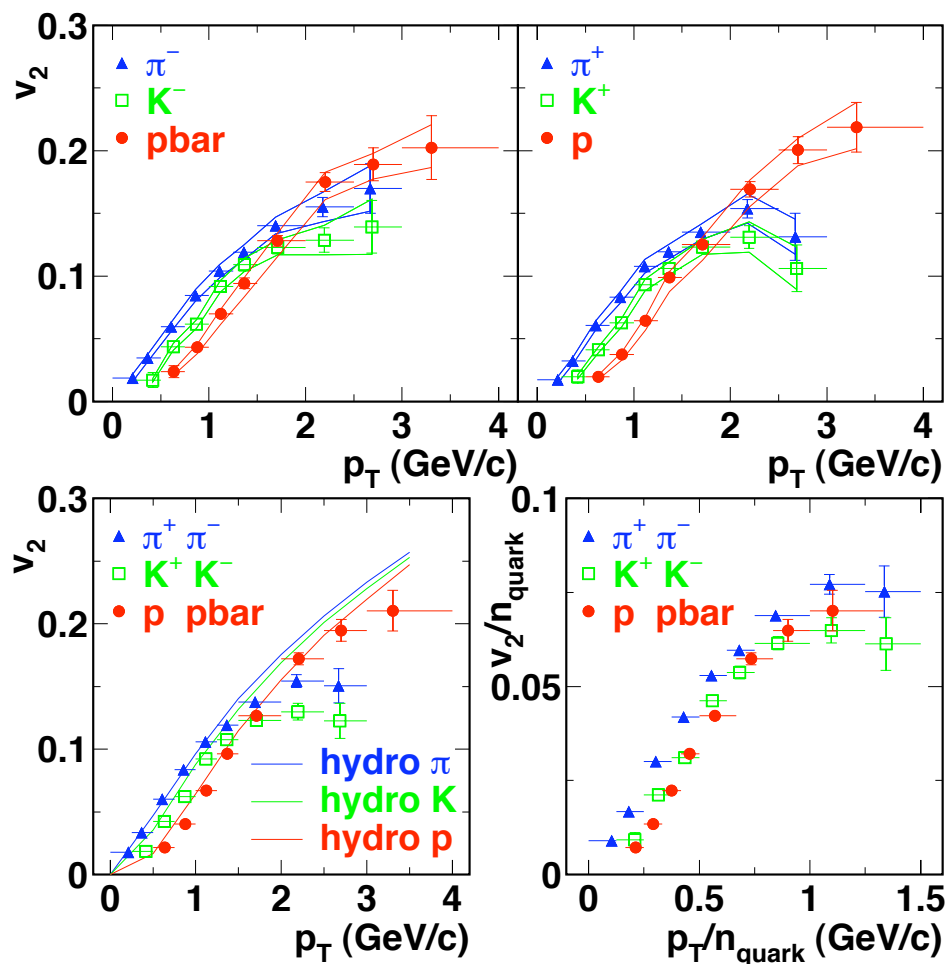
$$\frac{dN}{d\phi} \propto 1 + 2v_2 \cos(2\phi)$$

Initial conditions at thermalization time need to be given (ex. CGC)



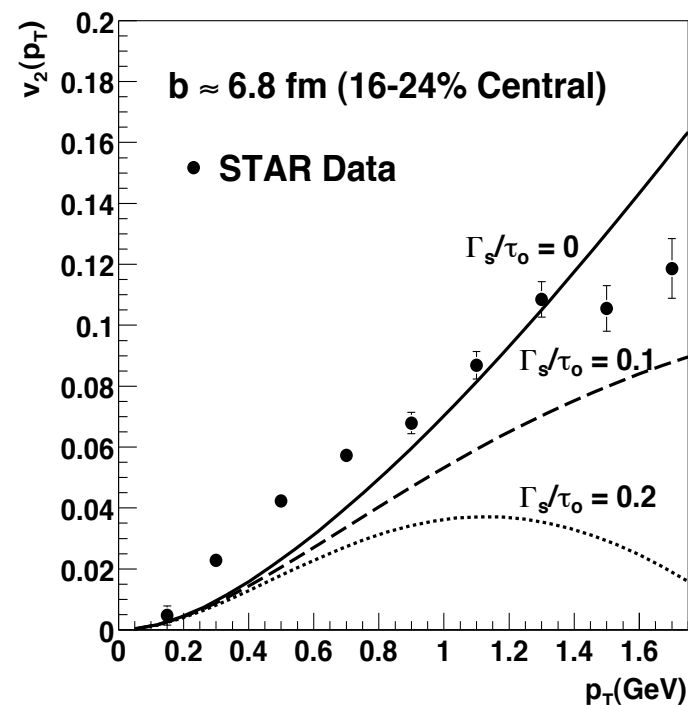
One of the first measurements at RHIC

Large elliptic flow compatible with ideal hydrodynamics



[data: STAR]

The effect of viscosity



[Teaney 2003]

Two important results

- Initial time for the evolution is very small
- Viscosity (non-perfect fluid behavior) is small

Hard Probes

Hard probes in heavy-ion collisions

- ⇒ SPS $\sqrt{s} = 20$ GeV ($Q \sim 1$ GeV) → marginal access to HP
- ⇒ RHIC $\sqrt{s} = 200$ GeV ($Q \sim 10$ GeV) → access to HP
- ⇒ LHC $\sqrt{s} = 5500$ GeV ($Q \gtrsim 100$ GeV) → HP and QCD evolution

$$\sigma^{pp \rightarrow h} = f_p(x_1, Q^2) \otimes f_p(x_2, Q^2) \otimes \underbrace{\sigma(x_1, x_2, Q^2)}_{\text{RHIC}} \otimes D(z, Q^2) + \left(\frac{1}{Q^2}\right)^n$$

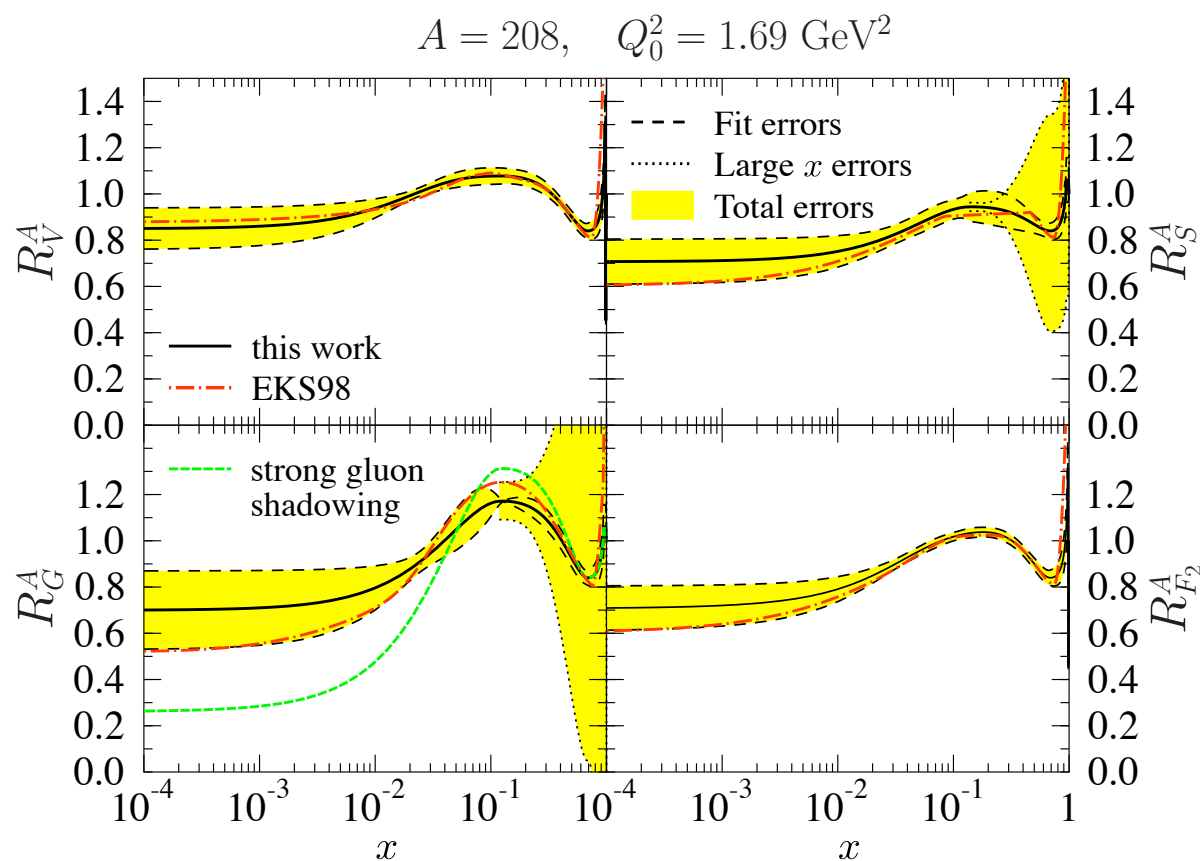
Diagram illustrating the kinematic access of different colliders to the terms in the cross-section formula:

- LHC** (red arrow) accesses the parton distribution functions $f_p(x_1, Q^2)$ and $f_p(x_2, Q^2)$.
- RHIC** (red arrow) accesses the hard cross-section $\sigma(x_1, x_2, Q^2)$.
- SPS** (red arrow) accesses the fragmentation function $D(z, Q^2)$.

- ⇒ The extension of the medium modifies the long-distance terms
 - ⇒ New evolution equations for $f_A(x, Q^2); D(z, Q^2)$
- ⇒ Kinematical access to evolution: large- Q^2 , small- x → LHC

Nuclear PDFs: uncertainties

PDFs determined by a global DGLAP fit

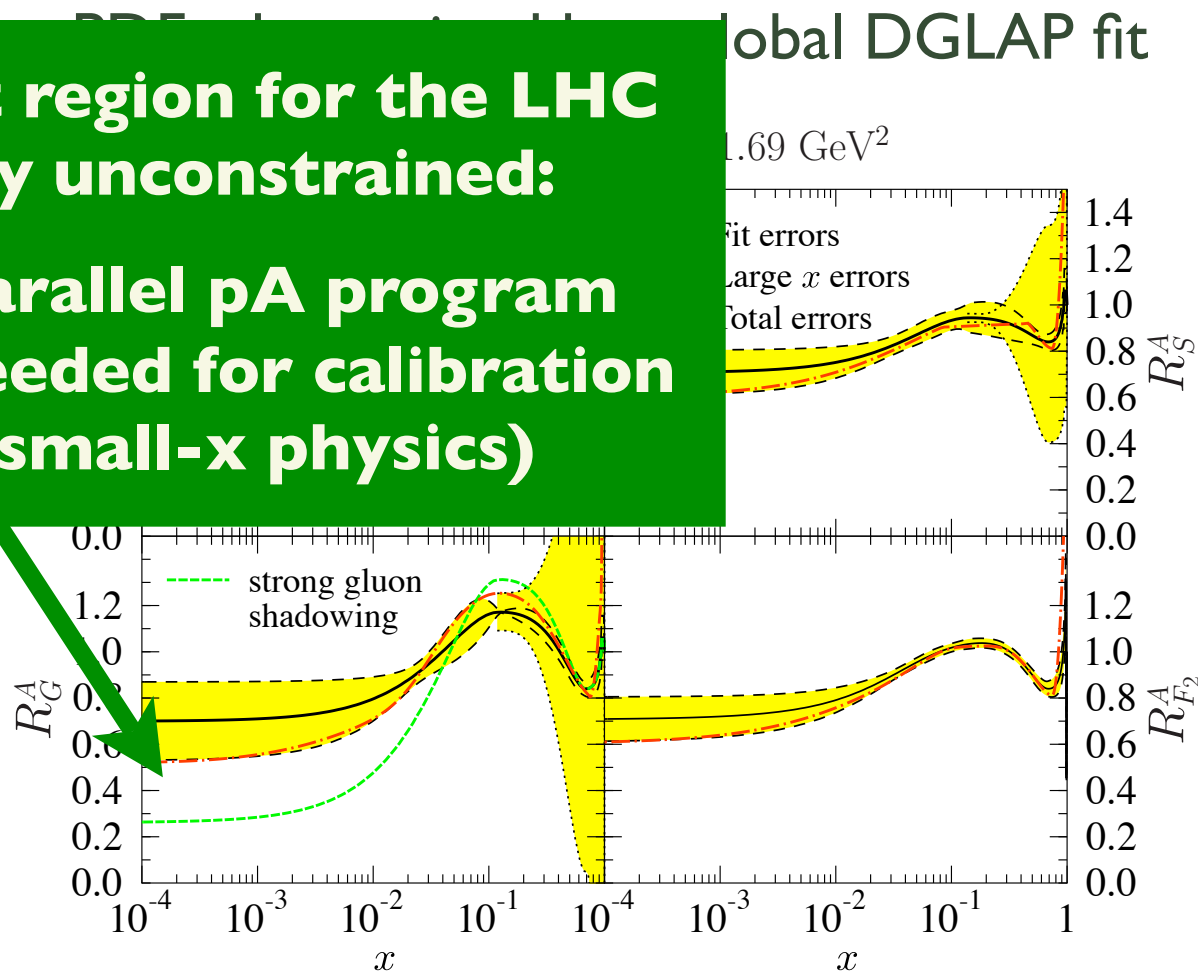


[Eskola, Kolhinen, Paukkunen, Salgado 2007]

[Other sets: De Florian, Sassot 2004; Frankfurt, Guzey, Strikman 2005; Hirai, Kumano, Nagai 2007]

Nuclear PDFs: uncertainties

**Relevant region for the LHC
largely unconstrained:**
→ **a parallel pA program
will be needed for calibration
(and small- x physics)**



[Eskola, Kolhinen, Paukkunen, Salgado 2007]

[Other sets: De Florian, Sassot 2004; Frankfurt, Guzey, Strikman 2005; Hirai, Kumano, Nagai 2007]

A conceptually simple example, J/Ψ suppression

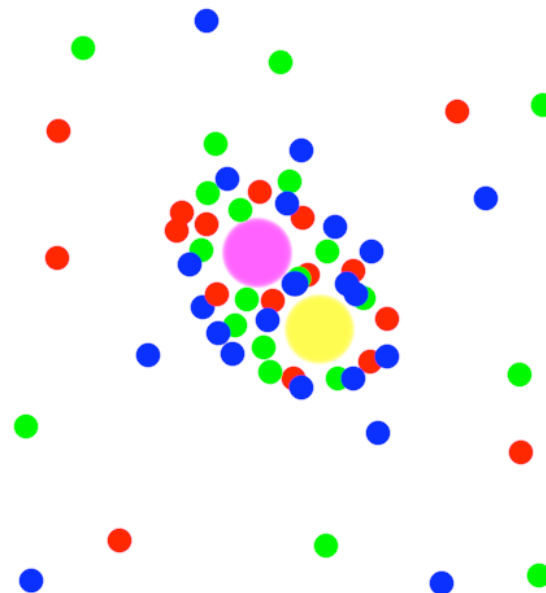
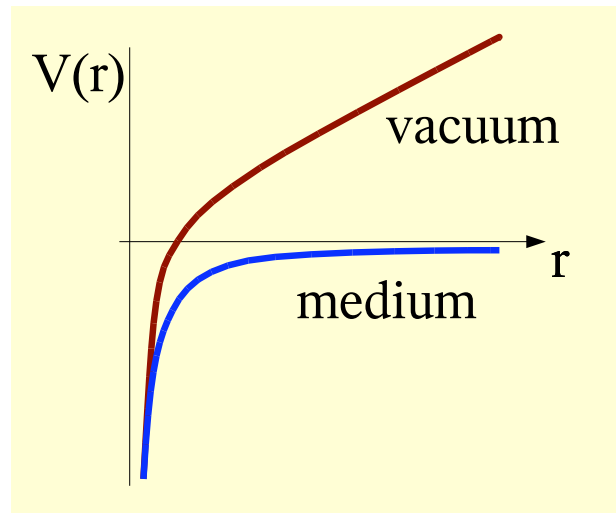
⇒ A J/Ψ is a $c\bar{c}$ bound state.

$$\sigma^{hh \rightarrow J/\Psi} = f_i(x_1, Q^2) \otimes f_j(x_2, Q^2) \otimes \sigma^{ij \rightarrow [c\bar{c}]}(x_1, x_2, Q^2) \langle \mathcal{O}([c\bar{c}] \rightarrow J/\Psi) \rangle$$

⇒ The potential is screened by the medium

→ The long-distance part is modified $\langle \mathcal{O}([c\bar{c}] \rightarrow J/\Psi) \rangle \rightarrow 0$

⇒ The J/Ψ production is suppressed [Matsui, Satz 1986]



A conceptually simple example, J/Ψ suppression

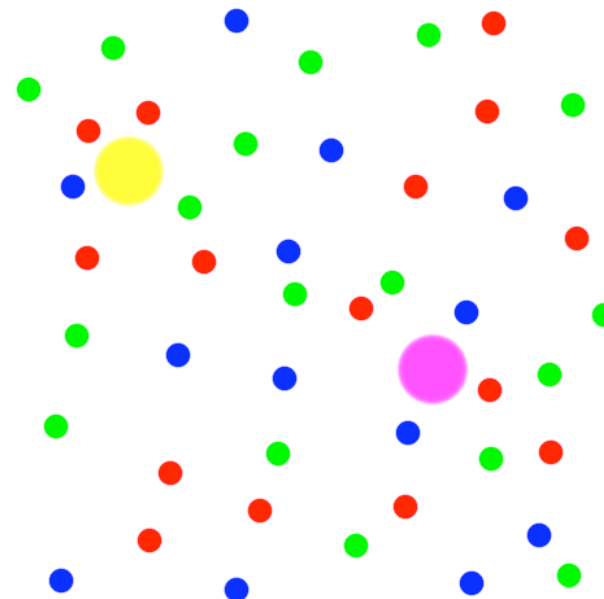
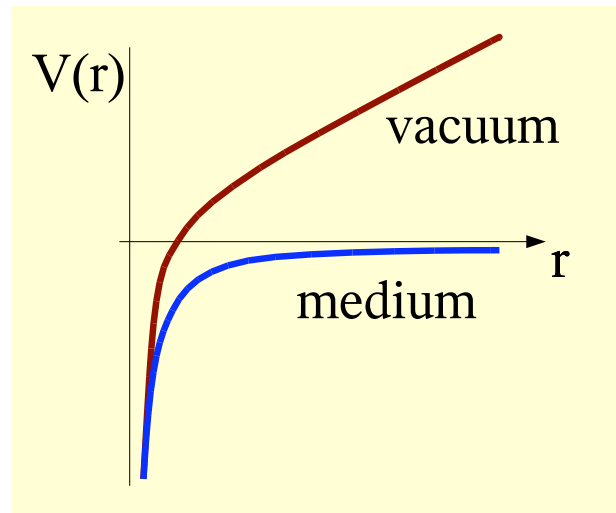
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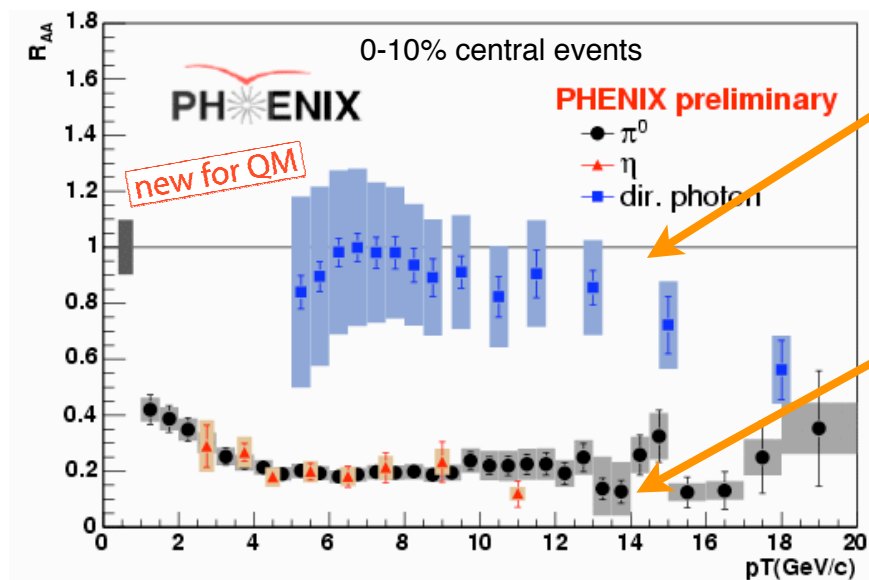
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Effects on high- p_t particles

$$R_{AA} = \frac{dN^{AA}/dp_t}{N_{\text{coll}} dN^{pp}/dp_t}$$



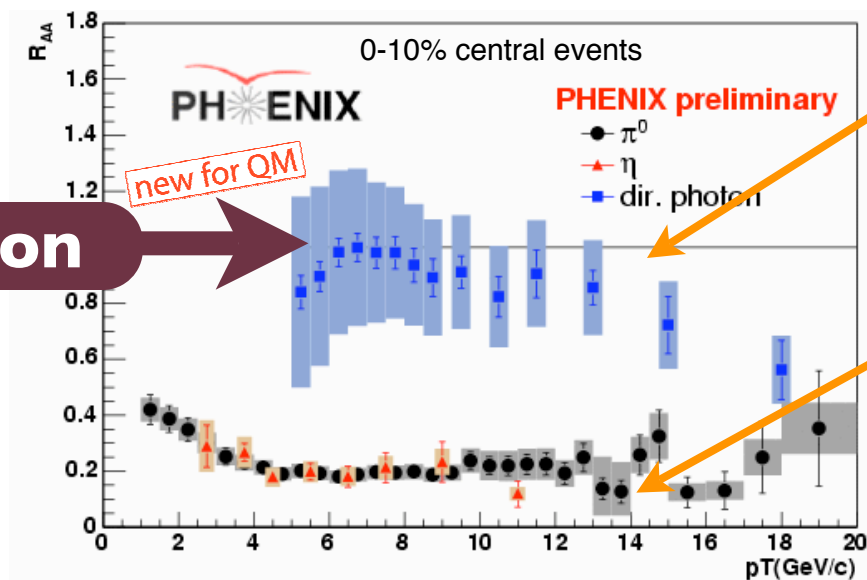
photons

mesons

Photons don't interact with the medium quarks and gluons do

Effects on high- p_t particles

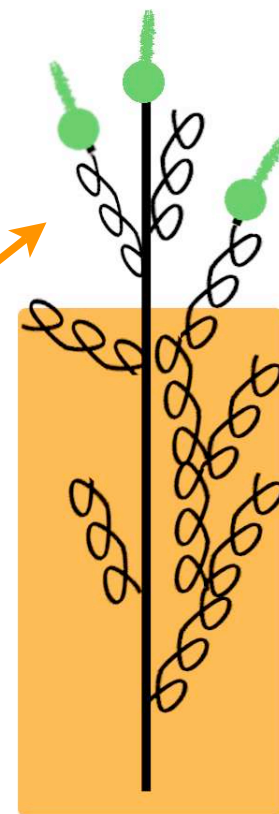
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Calibration

photons

mesons

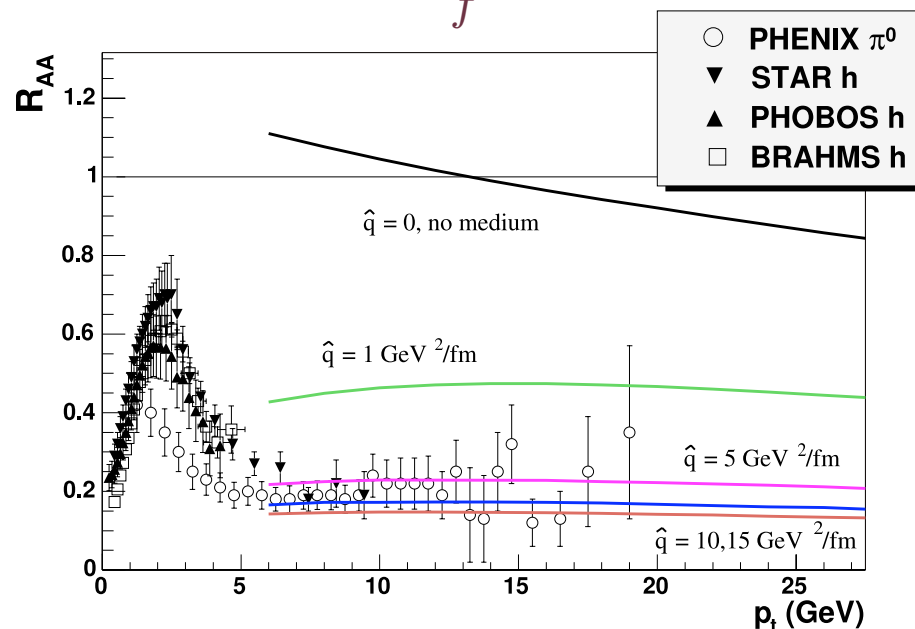


Photons don't interact with the medium quarks and gluons do

⇒ Well calibrated and abundant probes of the medium

The inclusive particle suppression

$$d\sigma_{\text{med}}^{AA \rightarrow h+X} = \sum_f d\sigma_{\text{vac}}^{AA \rightarrow h+X} \otimes P_f(\Delta E, L, \hat{q}) \otimes D_{f \rightarrow h}^{(\text{vac})}(z, \mu_F)$$



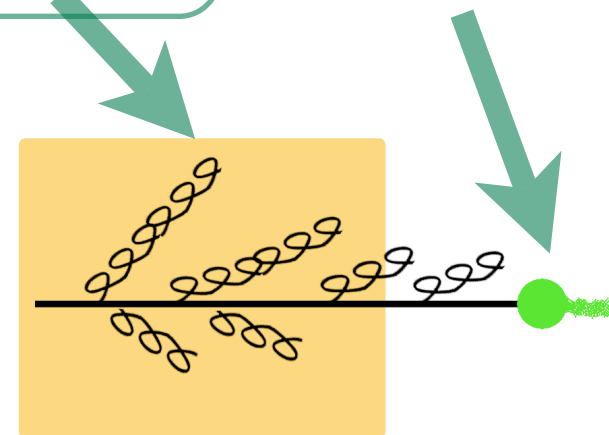
[Eskola, Honkanen, Salgado, Wiedemann 2005]

⇒ Data favors a large transport coefficient

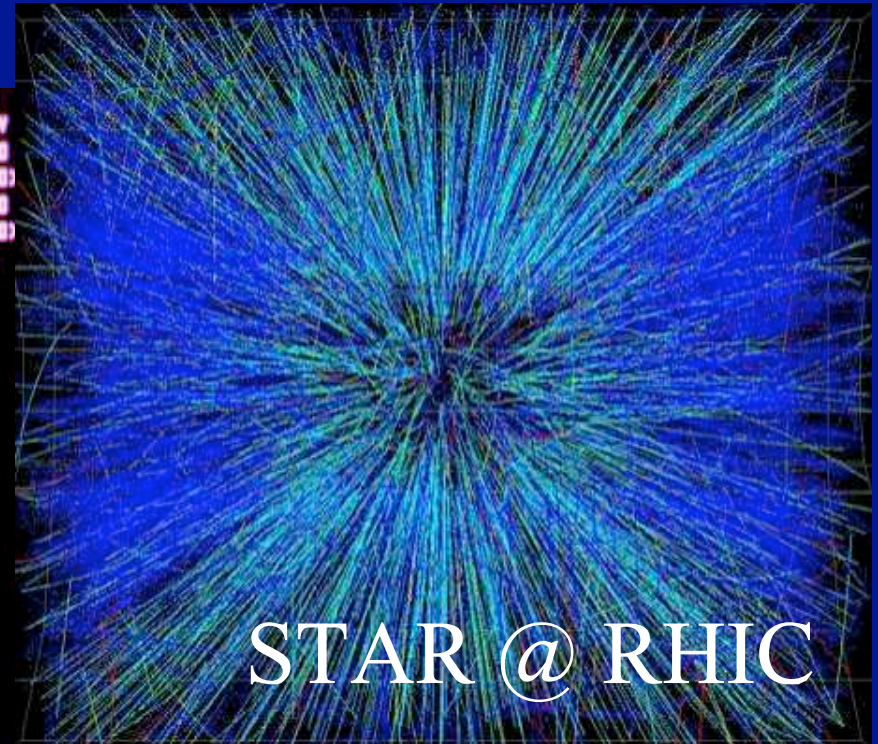
$$\hat{q} \simeq 5 \dots 15 \text{ GeV}^2/\text{fm}$$

⇒ Transport coefficient

$$\hat{q} \simeq \frac{\langle k_{\perp}^2 \rangle}{\lambda} \propto n(\xi) \sigma$$



Jets in HIC



Jets in HIC



Prediction: jet broadening

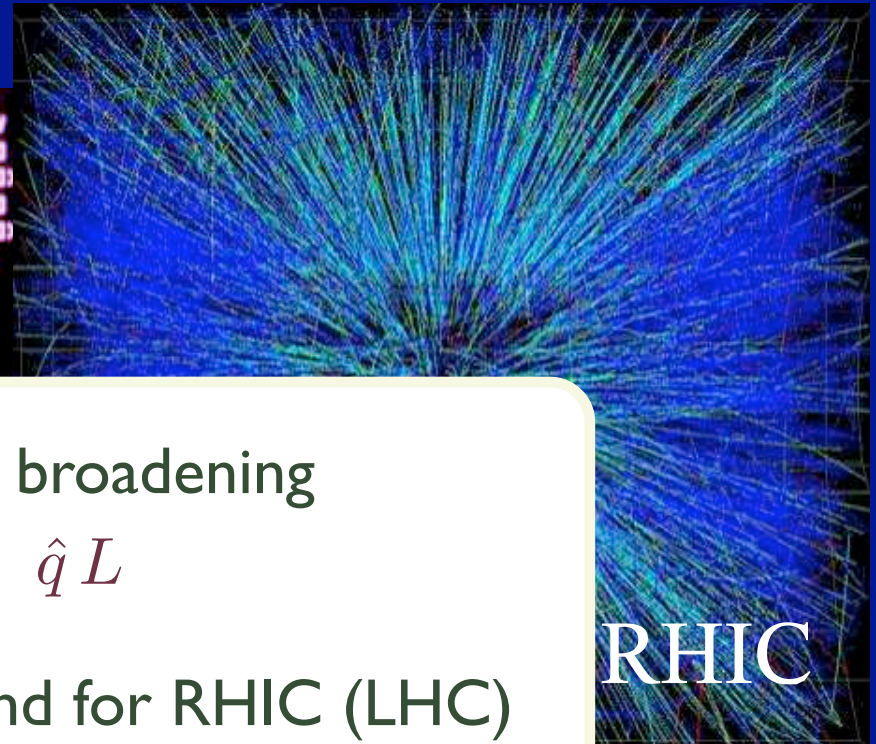
$$\langle k_t \rangle \sim \hat{q} L$$

Multiplicity background for RHIC (LHC)

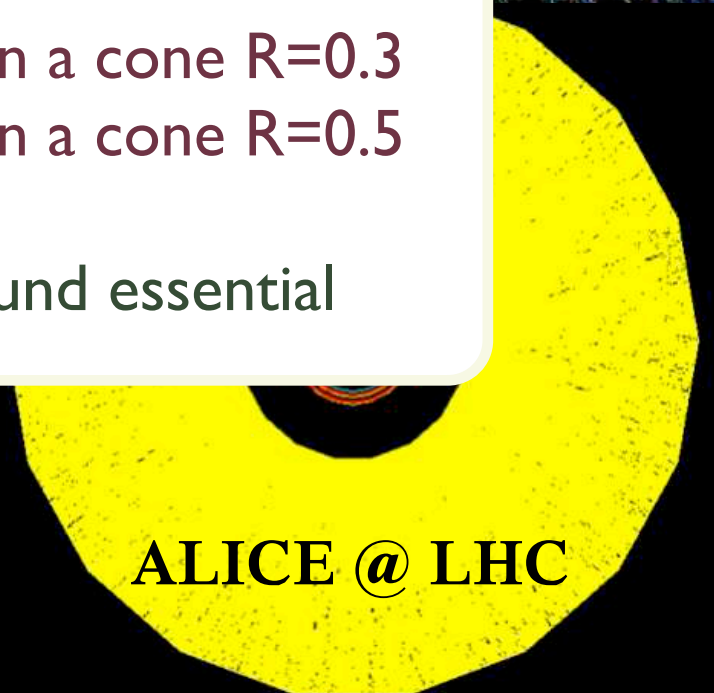
⇒ $E^{\text{bg}} \sim 20$ (100) GeV in a cone $R=0.3$

⇒ $E^{\text{bg}} \sim 50$ (250) GeV in a cone $R=0.5$

Control over background essential

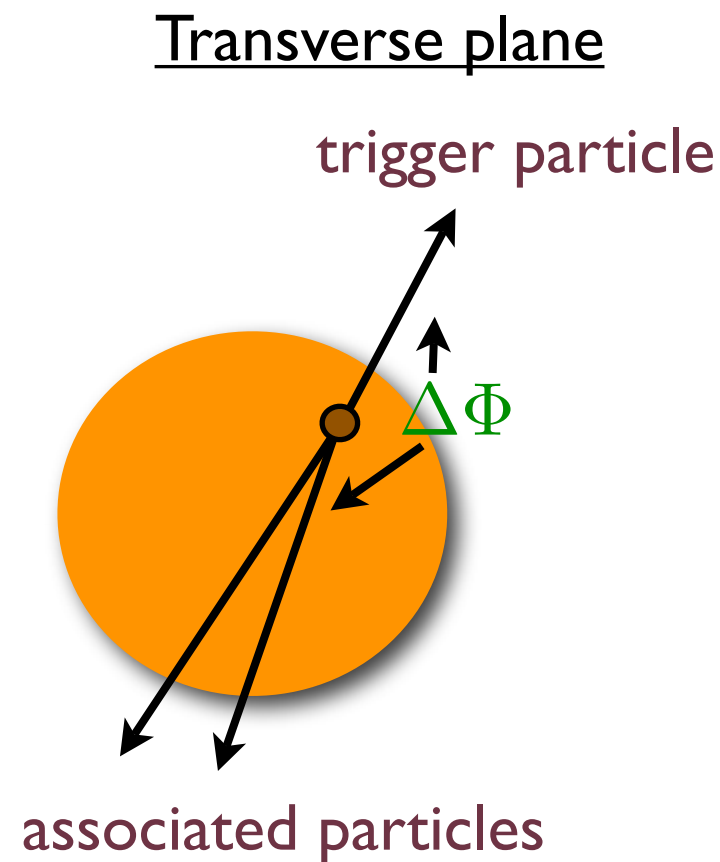
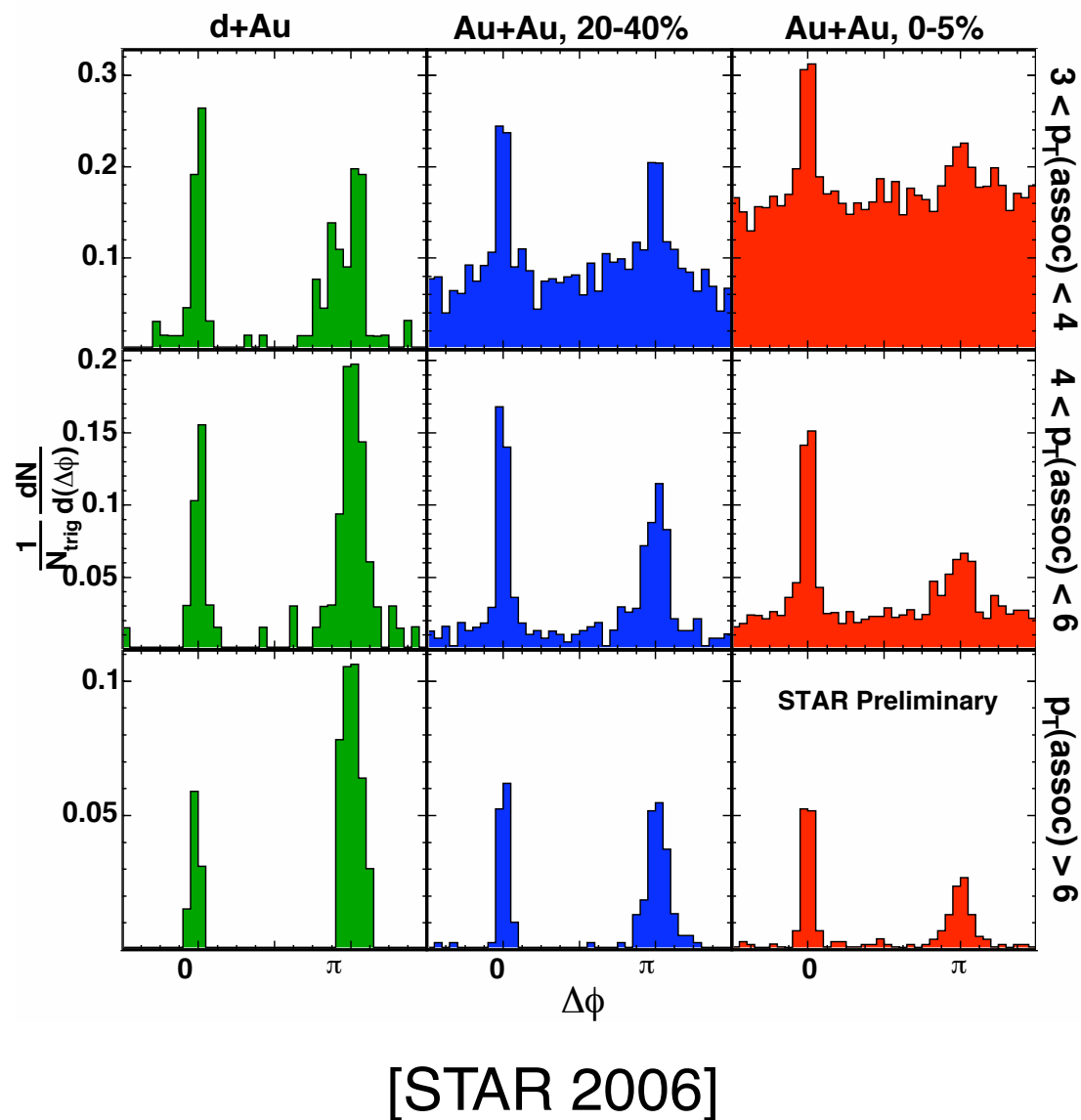


RHIC

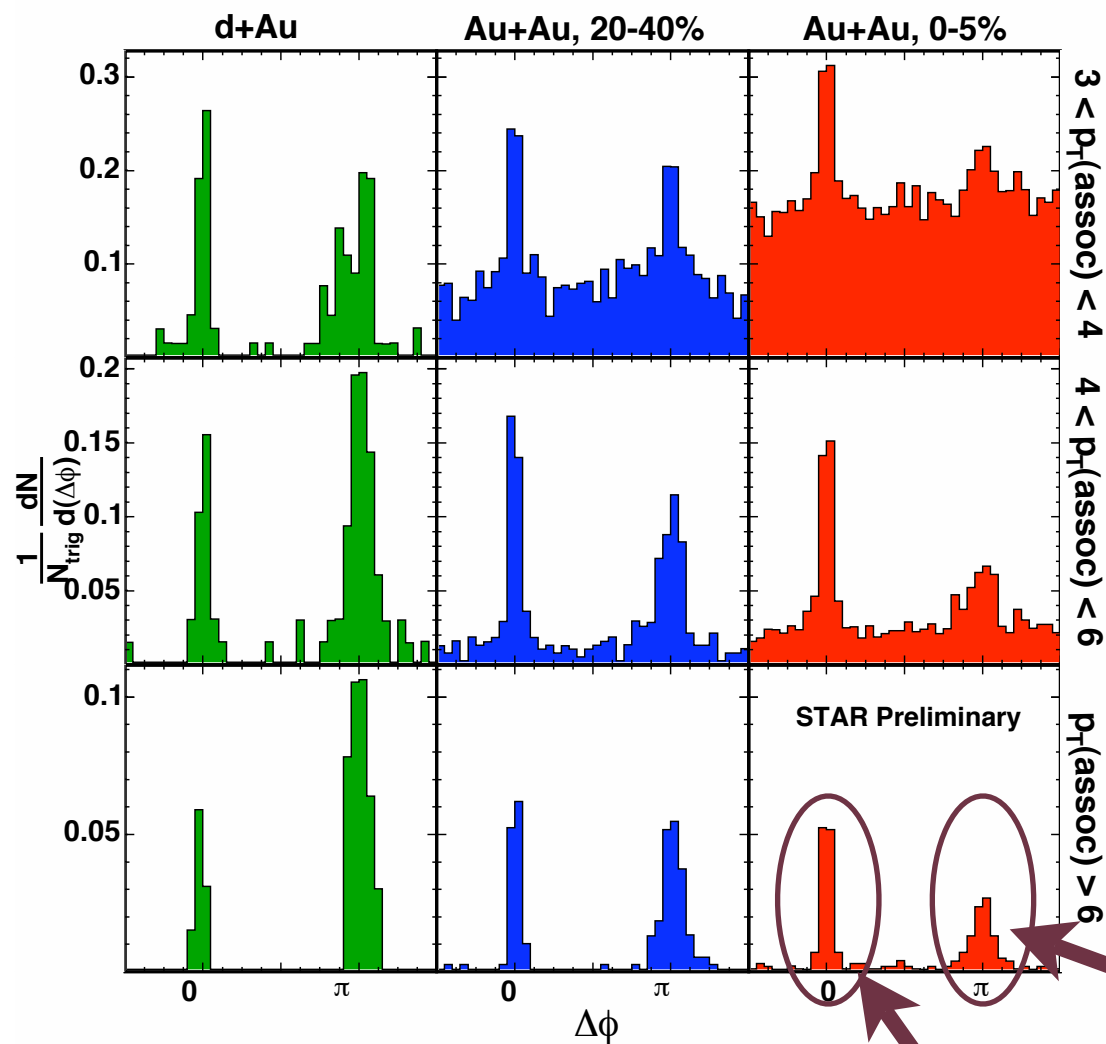


ALICE @ LHC

RHIC: two particle correlations

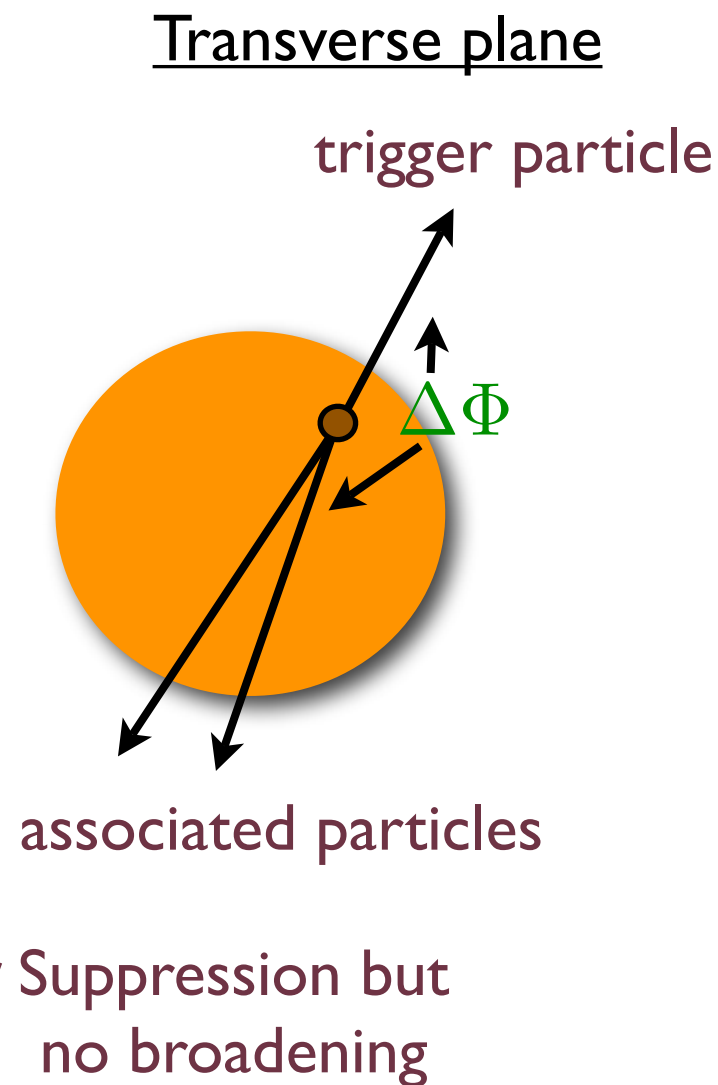


RHIC: two particle correlations



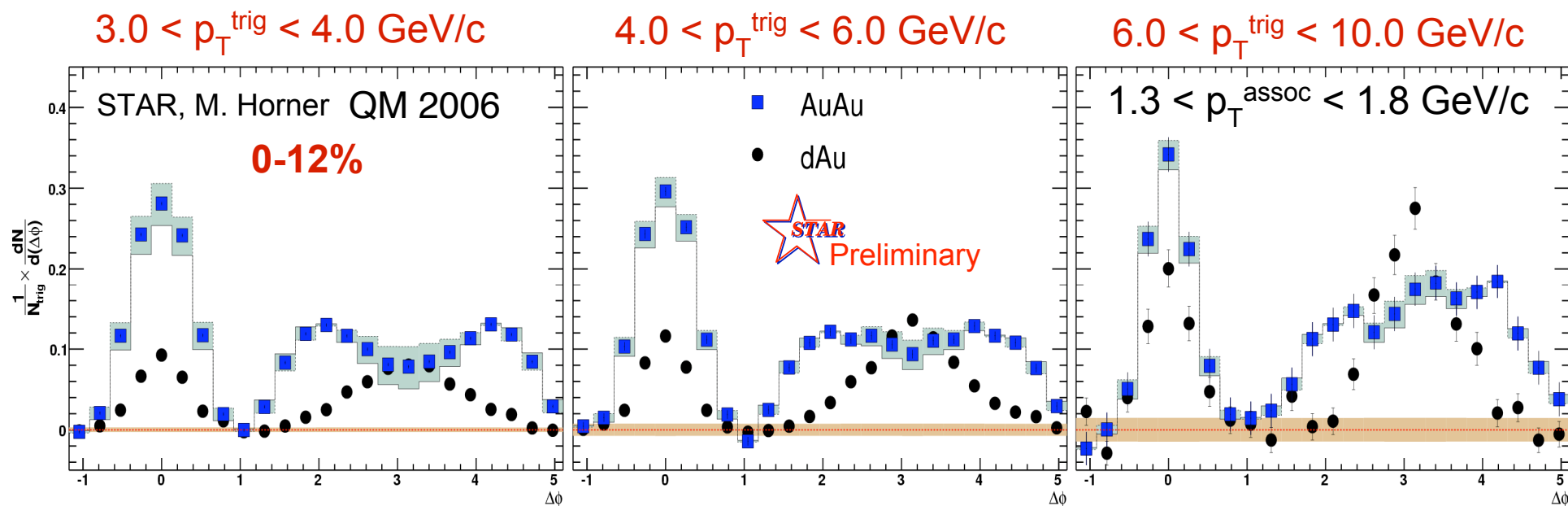
[STAR 2006]

Unchanged



Removing the cut-off at RHIC

[Similar results for PHENIX and also SPS (Ceres)]



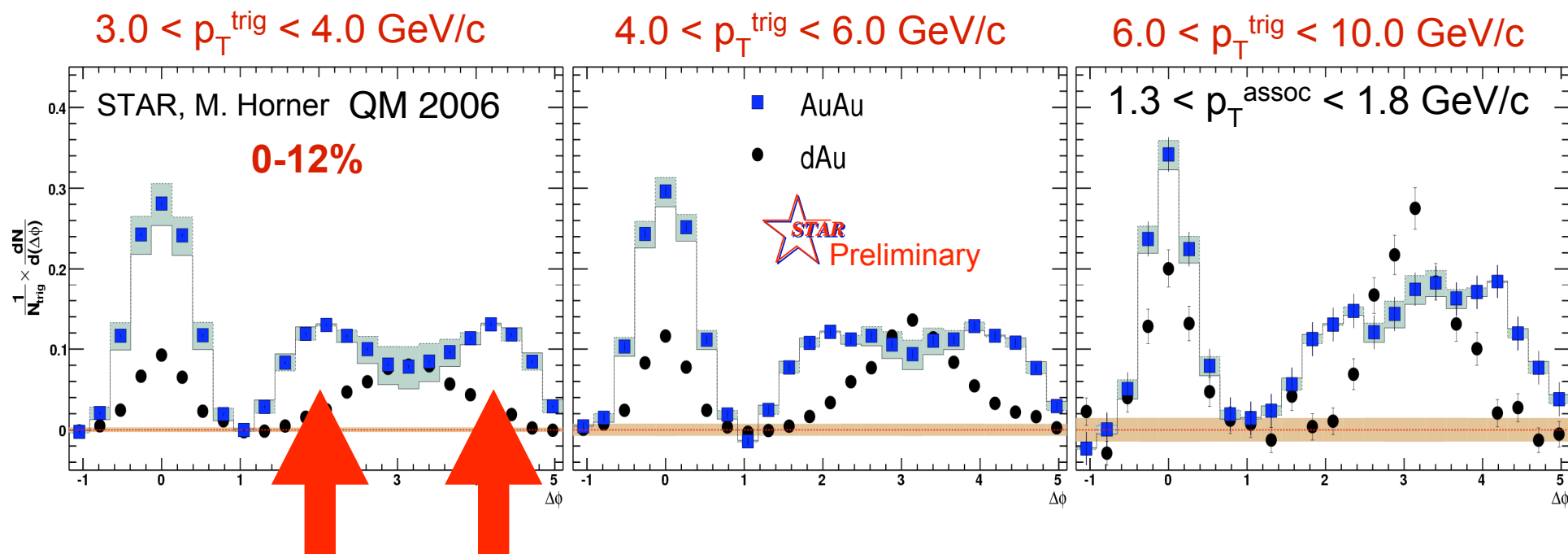
⇒ Nontrivial angular dependences in the away side

⇒ Large broadening

⇒ Two-peaks when $p_t^{\text{trigg}} \sim p_t^{\text{assoc}}$

Removing the cut-off at RHIC

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⇒ Nontrivial angular dependences in the away side

⇒ Large broadening

⇒ Two-peaks when $p_t^{\text{trigg}} \sim p_t^{\text{assoc}}$

Interpretations: two opposite assumptions

All lost energy is transferred to the medium

⇒ Hydro evolution

[Satarov, Stoecker, Mishustin 2005; Casalderrey-Solana, Shuryak, Teney 2005....]

⇒ Colored wakes

[Muller, Ruppert, Renk 2006, Chakraborty, Mustafa, Thoma 2006...]



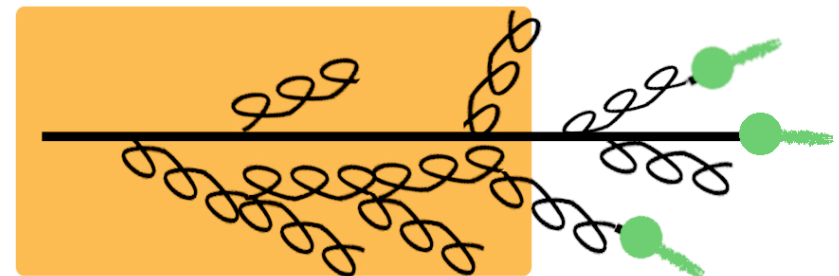
Negligible energy transferred: Energy loss by radiation

⇒ Medium-induced gluon radiation

[Polosa, Salgado 2007]

⇒ Cherenkov radiation

[Dremin 1979, 2005; Koch, Majumder, Wang 2006]



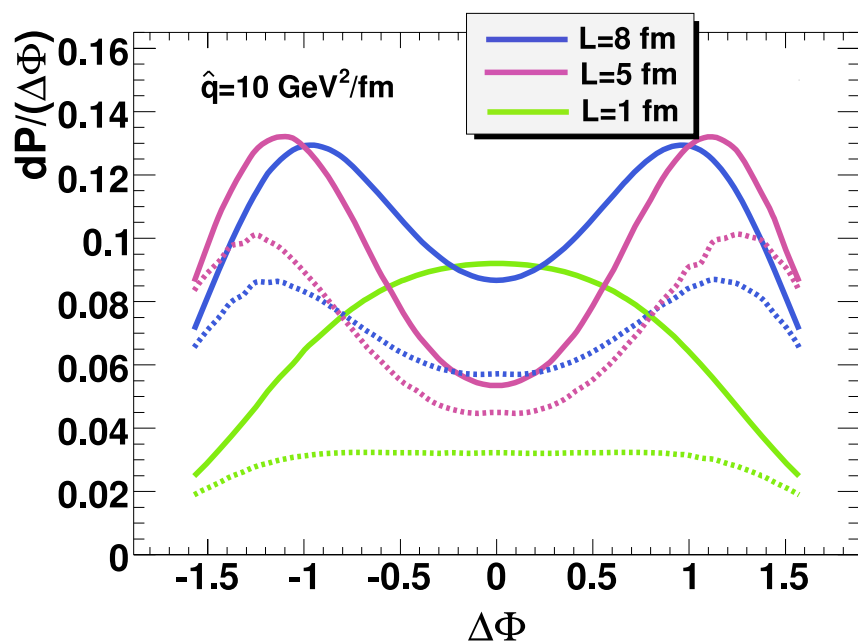
Parton shower for opaque media

⇒ When $p_t \sim \omega \lesssim \hat{q}^{1/3} \sim 3 \text{ GeV} \rightarrow$ Large \hat{q} needed

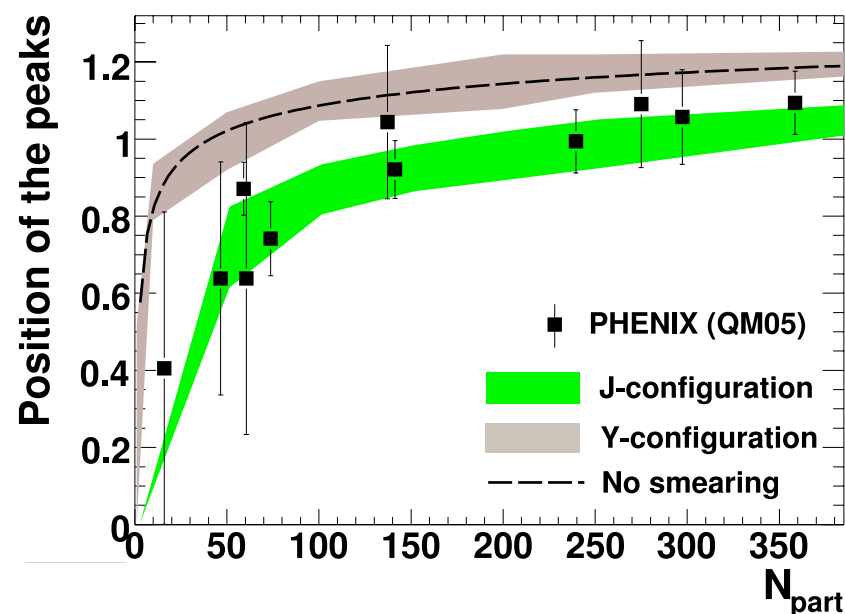
⇒ totally coherent limit and large angle radiation

⇒ Probability of only one splitting has non-trivial angular dependence

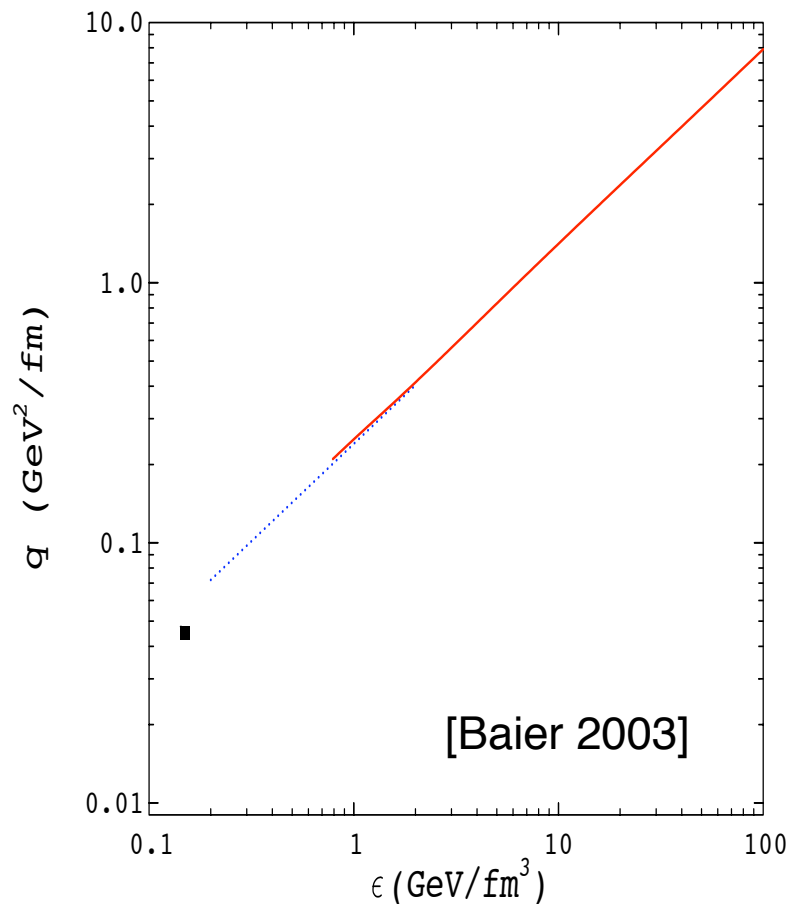
$$d\mathcal{P} = dz d\theta \frac{\alpha_s C_R}{16\pi^2} E L \sin \theta \cos \theta \exp \left\{ -\frac{\alpha_s C_R}{16\pi} E L \cos^2 \theta \right\}$$



[Polosa, Salgado 2007]



Interpretation of the value of \hat{q}



⇒ Transport coefficient for an ideal quark-gluon gas

$$\hat{q}_{\text{ideal gas}} \simeq \frac{72}{\pi} \xi(3) \alpha_s^2 T^3 \simeq 2\epsilon^{3/4}$$

[Baier and Schiff 2006]

⇒ Fits to the data

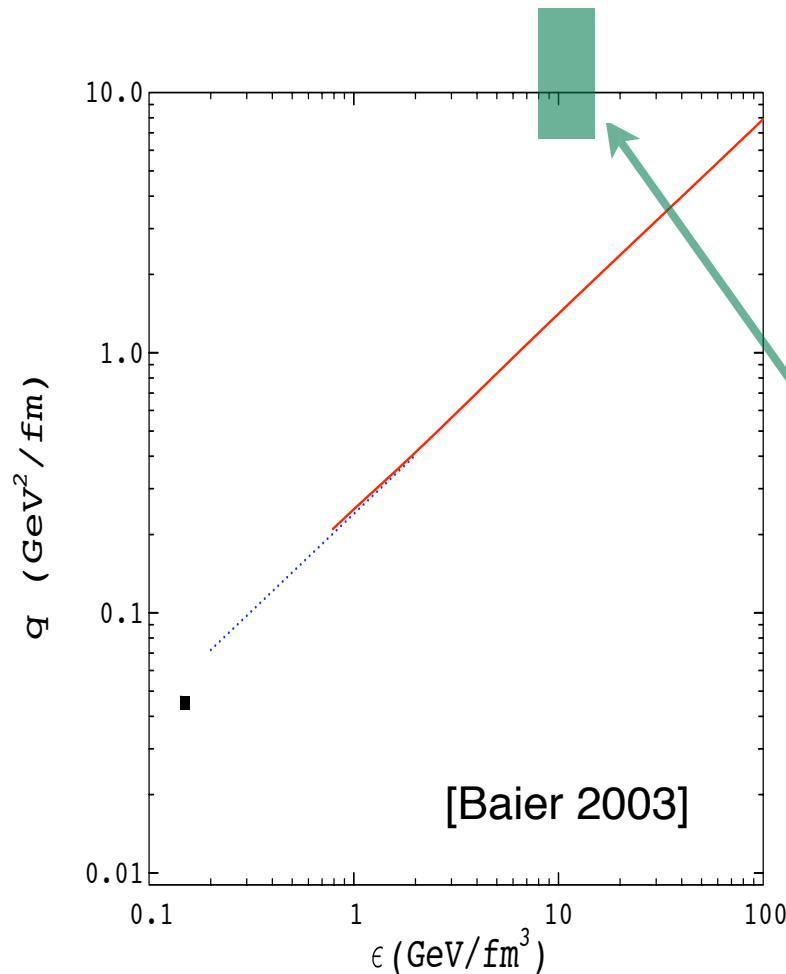
$$\hat{q} > 5 \hat{q}_{\text{ideal gas}} \quad [\text{Eskola et al. 2004}]$$

$$\hat{q} \simeq 4.2 \hat{q}_{\text{ideal gas}} \quad [\text{Renk et al. 2007}]$$

⇒ Geometry plays a crucial role

⇒ Model of the medium? sQGP?

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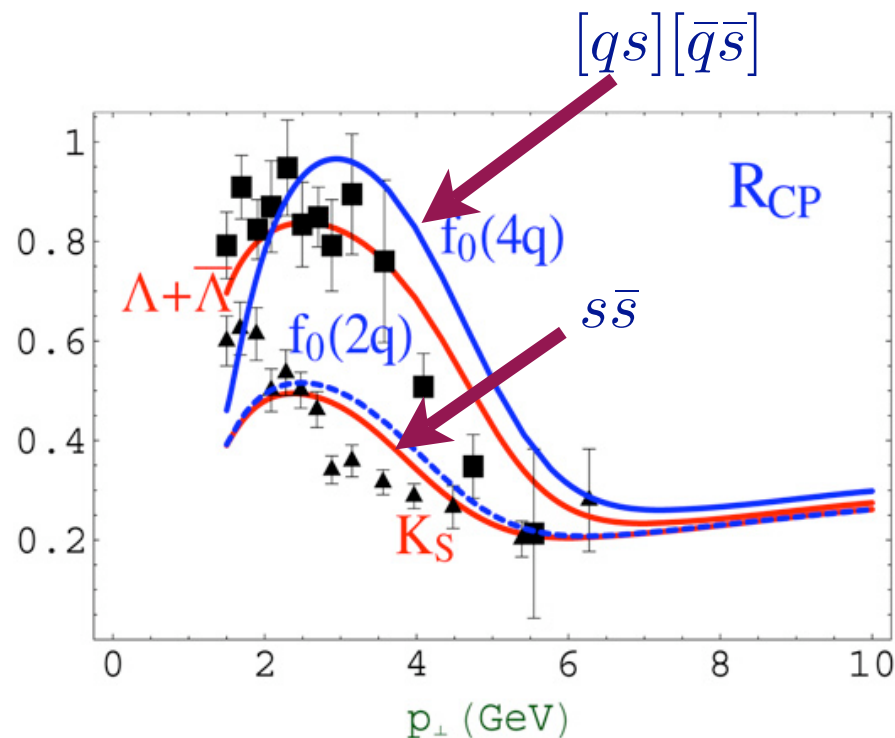
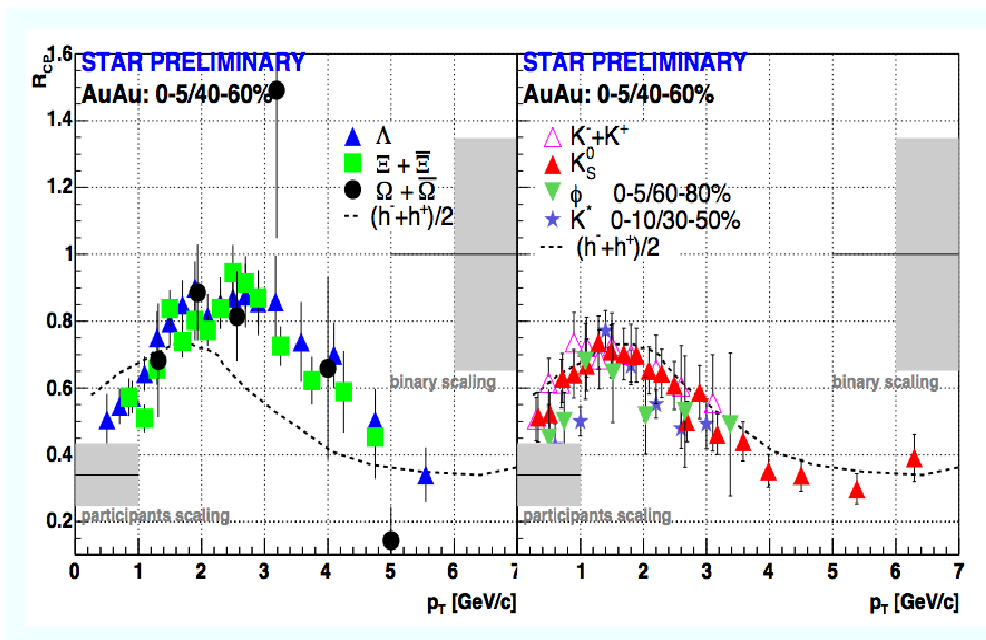
What is the order of magnitude of the NLO correction?

What are the degrees of freedom?

- 👁 Above $3T_c$ quark and gluon quasiparticles (pert. QCD)
- 👁 Below $3T_c$ (where experiments are)
 - Possible existence of bound states (lattice)
 - Large Interaction cross sections (viscosity, jet quenching)
 - Still unclear \rightarrow several proposals
 - Ideal fluid formed by colored bound states, diquarks...?
 - Strongly coupled Quark Gluon Plasma (sQGP)?

Exotic meson spectroscopy and HIC

Baryon/meson hierarchy of suppressions observed at RHIC

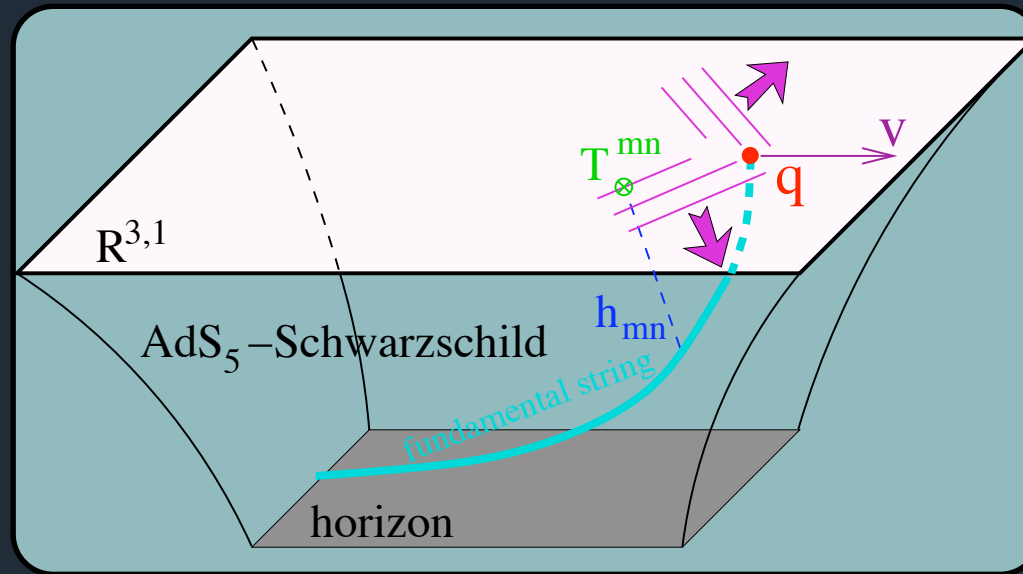


A nice possibility for a clean measurement of the exotic structure

$$f_0(980) = s\bar{s} \quad \text{OR} \quad [qs][\bar{q}\bar{s}] ?$$

[Maiani, Polosa, Riquer, Salgado (2007)]

[Nonaka, Muller, Asakawa, Bass, Fries 2004; Chen, Ko, Liu, Nielsen, Greco, Lee, Liu 2004, 2007]



The String Theory connection...

The observables

⇒ Applied to the jet quenching parameter:

$$\langle W^A(\mathcal{C}) \rangle \simeq \exp \left[-\frac{1}{4\sqrt{2}} \hat{q} r^2 L_- \right]$$

$\hat{q} = 4.5, 10.6, 20.7 \text{ GeV}^2/\text{fm}$
 $T = 300, 400, 500 \text{ MeV}$

[Liu, Rajagopalan, Wiedemann; Armesto, Edelstein, Mas...2006]

⇒ The viscosity-to-entropy ratio

$$\frac{\eta}{s} = \frac{1}{4\pi}$$

$\eta \propto \text{area of horizon}$
 $s \propto \text{area of horizon}$

Universal lower bound?

[Kovtun, Son, Starinets 2003]

⇒ The hydrodynamic behavior

→ Bjorken hydrodynamics (and more)

[Janik, Peschanski 2006; Kovchegov, Taliotis 2007...]

⇒ Shock waves; heavy quark energy loss; bound states....

[Gubser; Herzog, Karch, Kovtun, Kozcaz, Yaffe; Casalderrey-Solana, Teaney.... 2006]

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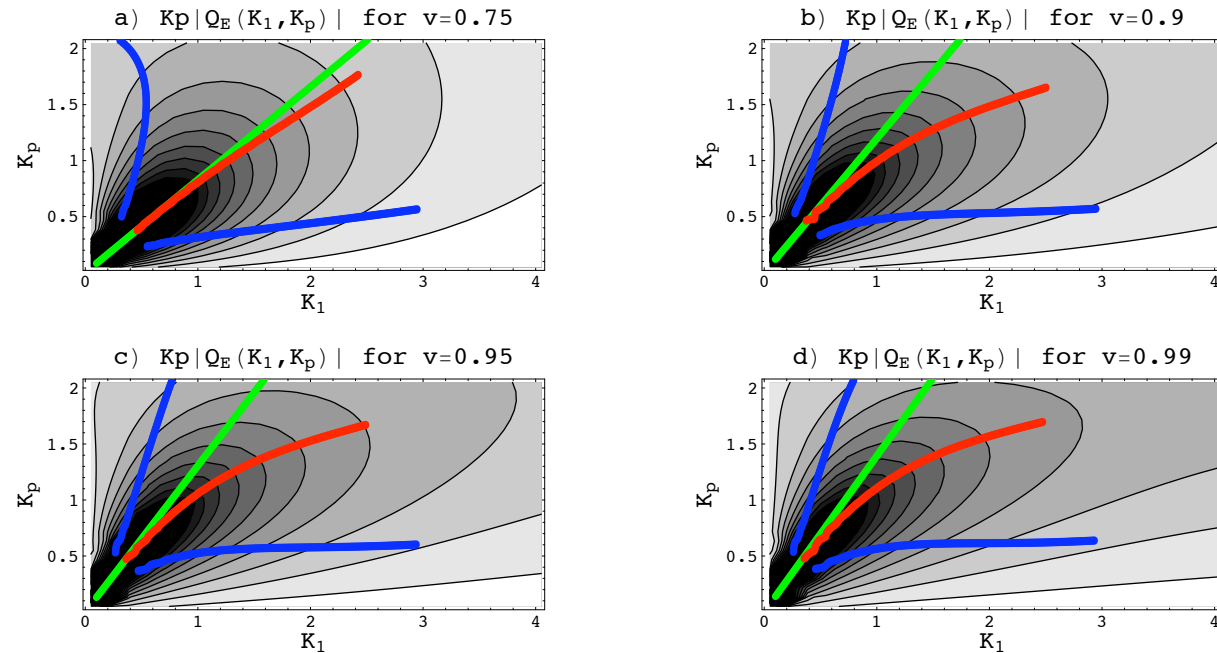
[Janik, Peschanski 2006; Kovchegov, Taliotis 2007...]

⇒ Shock waves; heavy quark energy loss; bound states....

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Shock waves and AdS/CFT

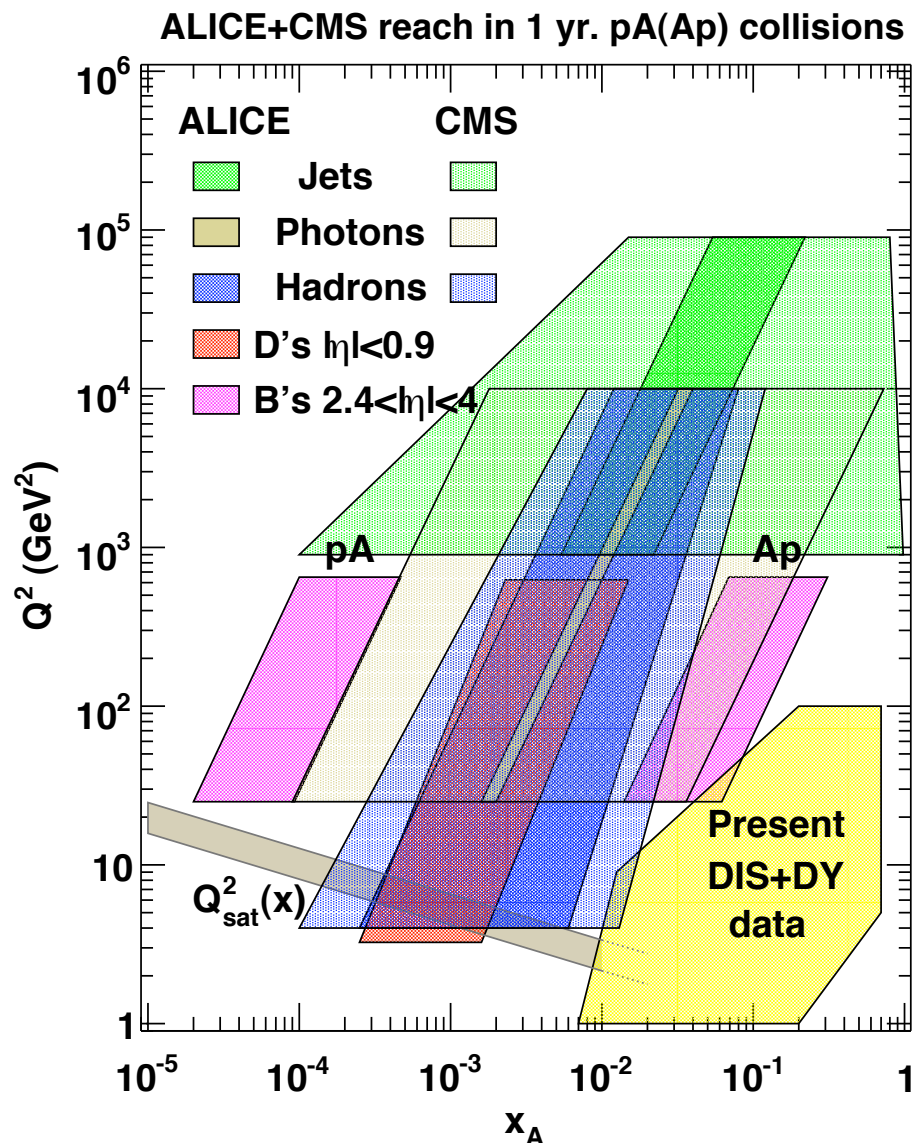
$T^{\mu\nu}$ computed for a quark moving with constant velocity in a medium



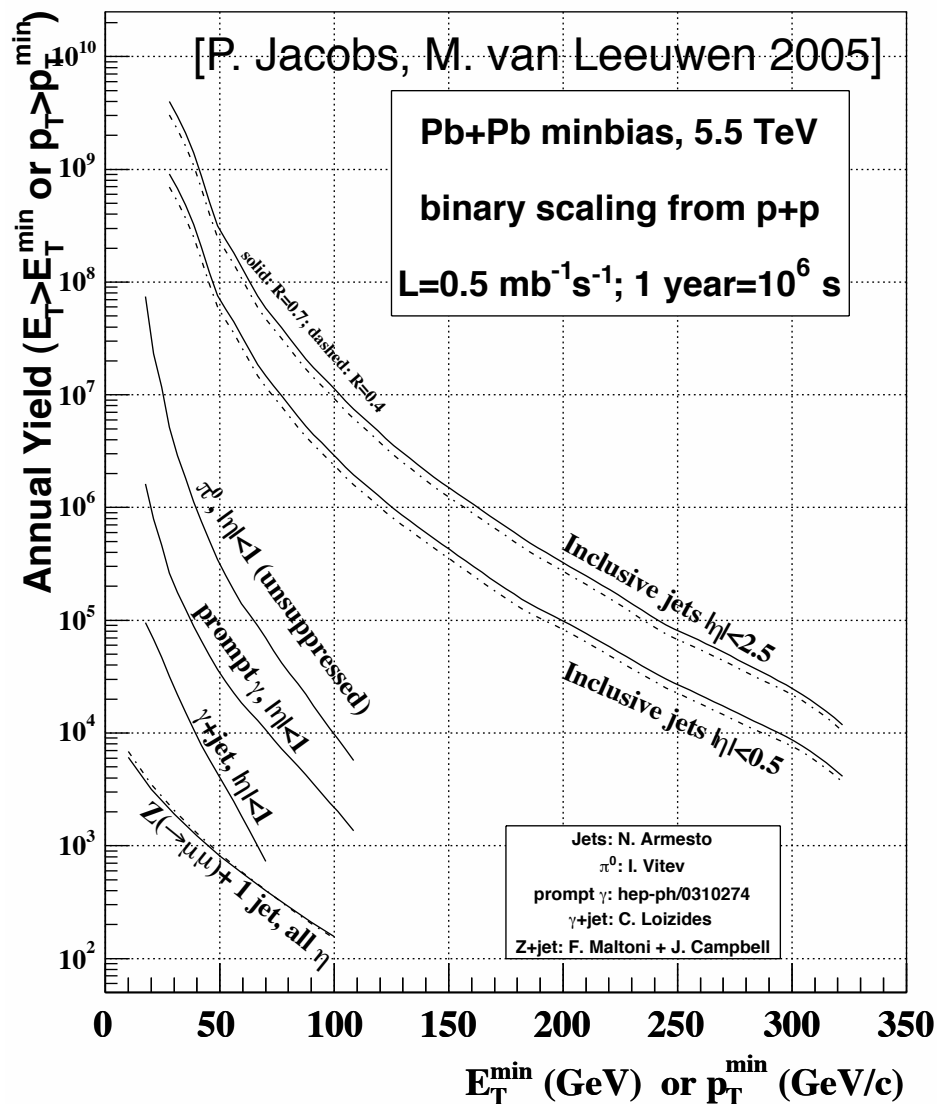
Yet, despite the potential stumbling blocks, it is exciting to see a simple type IIB string theory construction approaching quantitative comparisons with a data-rich experimental field.

[Friess, Gubser, Michalogiorgakis, Pufu hep-th/0607022]

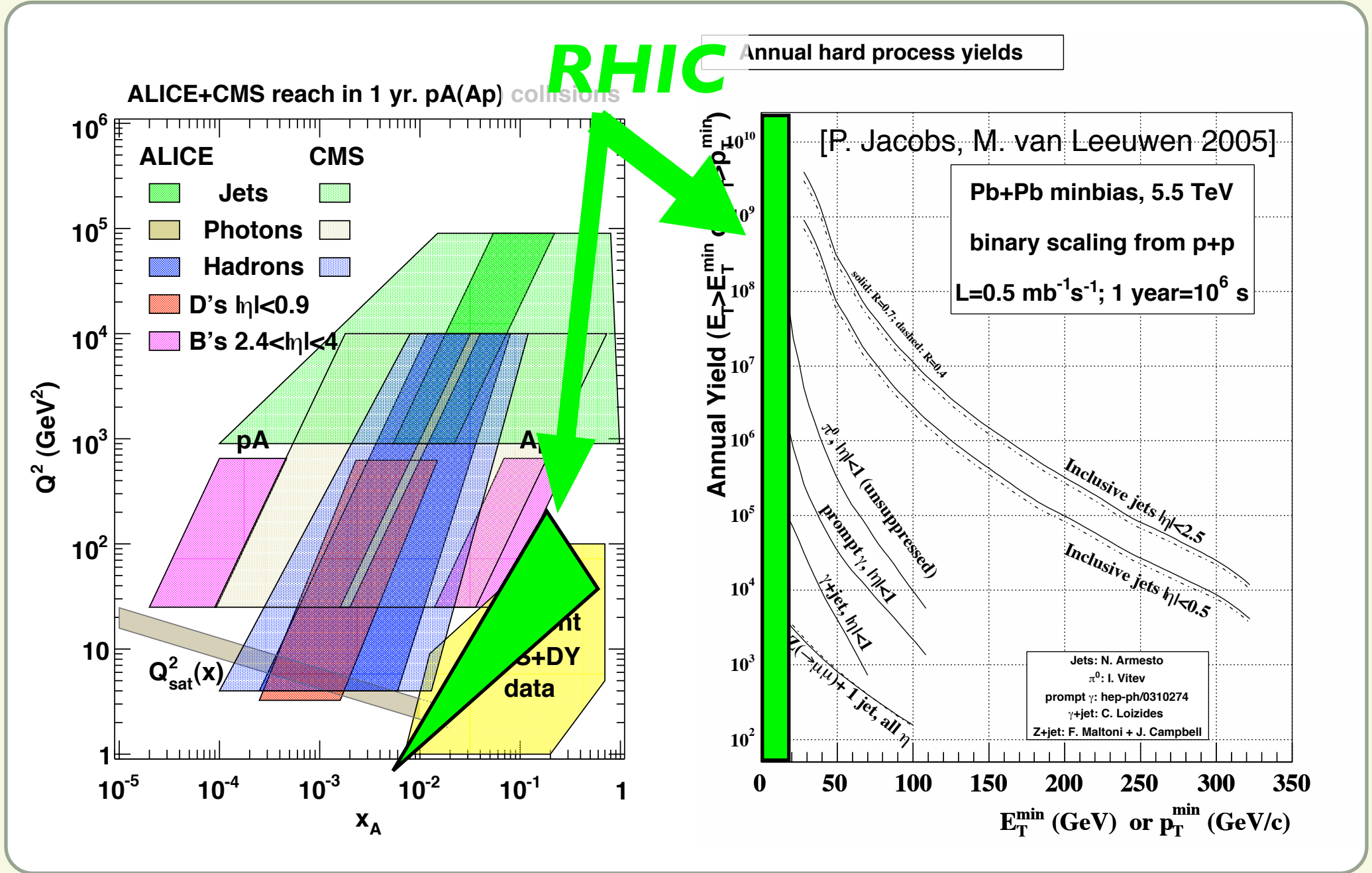
New regimes at the LHC



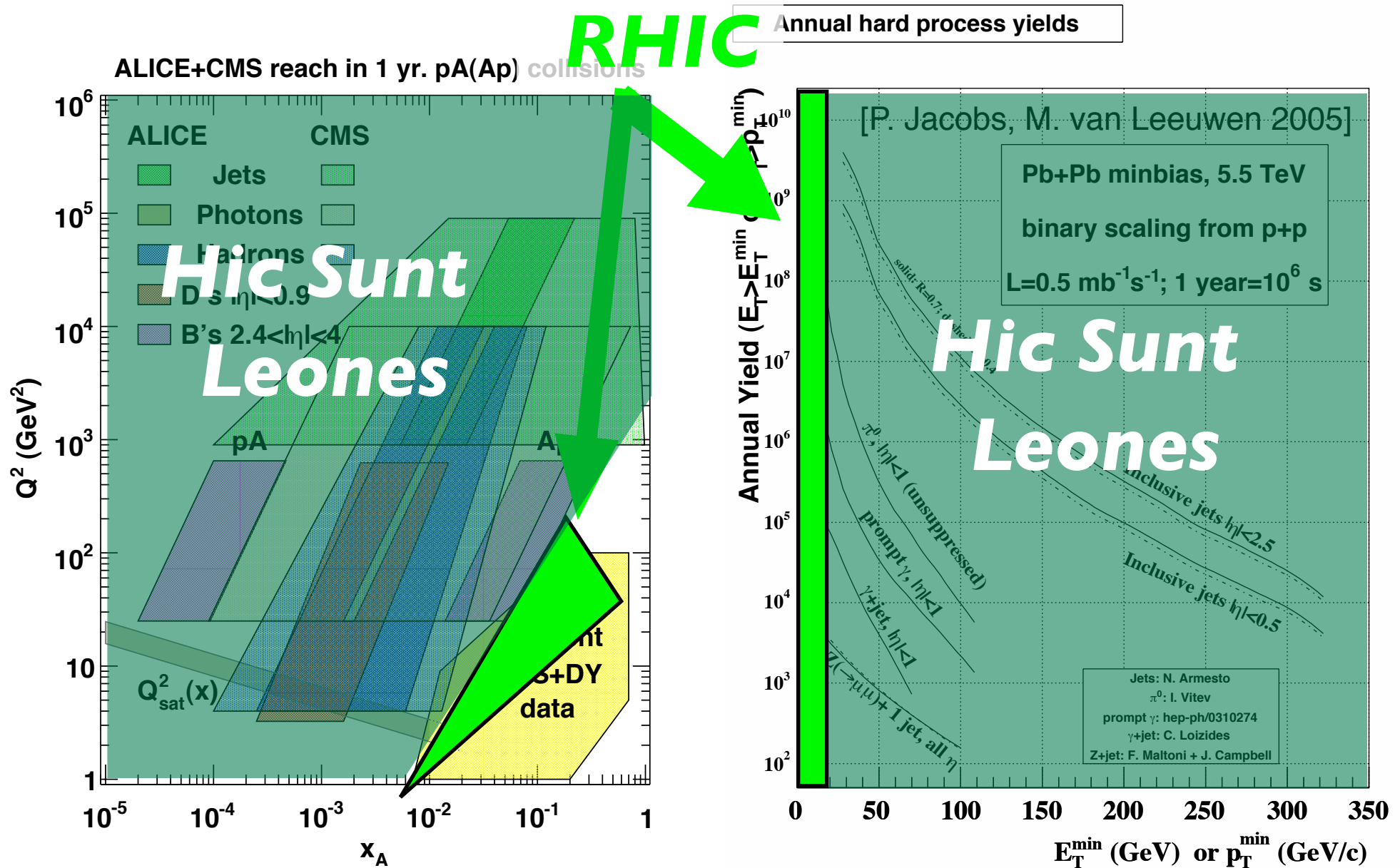
Annual hard process yields



New regimes at the LHC



New regimes at the LHC



Summary

- ⇒ Developing paradigm from experimental data+theory
 - ↘ Early thermalization and ideal fluid behavior
 - ↘ Very dense medium
 - ↘ Strongly coupled quark-gluon plasma → perfect liquid ?
 - ⇒ What are the relevant degrees of freedom (building blocks)?
 - ↘ Not a free gas of quarks and gluons in present experiments
 - ↘ Quasiparticles? bound (colored) states?
 - ⇒ Different fields are contributing to these developments
 - ↘ String-theory computations (attempt to) face experimental data
- ⇒ LHC: new regimes of QCD where in-medium evolution dominates