

## Evidence for new resonances

# in the combined analysis of recent hyperon photoproduction data 

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- PWA group in HISKP:
A.Anisovich, E.Klempt, V.Nikonov, A.Sarantsev, U.Thoma (http://pwa.hiskp.uni-bonn.de)
- Approach and recent results were presented in the talk of A.Anisovich
- The main task: search for new baryon resonances
- Polarization data are sensitive to weak signals
- Double polarization is necessary for a complete experiment
- CB ELSA double polarization data taking started (talks of U.Thoma and R.Beck)
$N(2070) D_{15}$ was discovered in combined fit of $\eta$ and $\pi^{0}$ photoproduction.


The check can be made using double polarization data.


The $P_{11}(1840)$ and $P_{13}(1900)$ states


A relatively narrow resonance is needed in 1850-1900 MeV region
$\gamma p \rightarrow \Lambda K^{+}$: dashed $P_{13}$, dotted $S_{11}$, dash-dotted $K^{*}$ exchange.
$\gamma p \rightarrow \Sigma K^{+}$: dashed $P_{13}$, dash-dotted $P_{11}$, dotted $K$ exchange.

The main proof is $K^{0} \Sigma^{+}$


dashed $P_{11}$ (1840)
$P_{11}$ gives the best $\chi^{2}$
$P_{13}$ gives the second best $\chi^{2}$
Could be both, but with pure data it is not a proof.


$\Lambda K^{+}$(left) and $\Sigma K^{+}$(right) recoil polarization (CLAS)
Dotted line is the fit without $P_{11}(1840)$

New $C_{x}, C_{z}$ data:
Evidence for the $N(1900) P_{13}$
No model could describe them.
In our approach we also had no good description of $C_{x}, C_{z}$ with old set of resonances.
Systematic discrepancies were observed.


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Refit of new data with old resonances. Maximal freedom for non-resonance contributions. Description of other data is much worse than before.
We tried to add new resonances.
As the firsr step, further resonances were introduced as Breit-Wigner amplitudes and different quantum numbers were tested.
The best \(\chi^{2}\) was obtained by introducing the second \(P_{13}\) state.
Solution 1: \(\quad 1885 \pm 15 \mathrm{MeV}\) mass and \(180 \pm 25 \mathrm{MeV}\) width, with \(\Delta \chi^{2}=1540\).
Solution 2: \(\quad 1975 \pm 15 \mathbf{M e V}\) mass and \(200 \pm 20 \mathbf{M e V}\) width.
Replacing:
\(S_{11} \quad \Delta \chi^{2}=950\).
\(D_{15} \Delta \chi^{2}=970\).
\(P_{11} \quad \Delta \chi^{2}=205\).
\(F_{15} \quad \Delta \chi^{2}\) small.
\(P_{33} \quad \Delta \chi^{2}\) smaller by a factor 2 than for a \(P_{13}\).
\(F_{17}, G_{17}\) did not improve the fit.
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In a final step, the $P_{13}$ was parameterized as 3 -pole 8 -channel K-matrix with $\pi N, \eta N, \Delta(1232) \pi$ ( $P$ and $F$-waves), $N \sigma, D_{13}(1520) \pi$ ( $S$-wave), $K \Lambda$ and $K \Sigma$ channels. This resulted in the fit solutions 1 and 2 which both are compatible with B-W fits.
In addition, both solutions are compatible now with elastic $\pi N$ scattering.
From the fit, properties of resonances in the $P_{13}$-wave were derived.
The lowest-mass pole is identified with the established $N(1720) P_{13}$,
the second pole with the badly known $N(1900) P_{13}$.
The third pole is introduced at about 2200 MeV .
It improves the quality of the fit in the high-mass region but its quantum numbers cannot be deduced safely from the present data base.



$C_{x}$ (full circles) and $C_{z}$ (open circles) for $\gamma \mathrm{p} \rightarrow \Lambda \mathrm{K}^{+}$. The solid and dashed curves are results of our fit obtained with solution 1 (left) and solution 2 (right) for $C_{x}$ and $C_{z}$.


$C_{x}$ (full circles) and $C_{z}$ (open circles) for $\gamma \mathrm{p} \rightarrow \Sigma \mathrm{K}^{+}$. The solid and dashed curves are results of our fit obtained with solution 1 (left) and solution 2 (right) for $C_{x}$ and $C_{z}$.

## $\sigma_{t o t}\left(\gamma p \rightarrow K^{0} \Sigma^{+}\right)$from CB-ELSA



Red line - $P_{13}(1900)$
Blue line - $P_{11}(1860)$ (improved $P$ in $K \Lambda$ and $K \Sigma$ data)


The total cross section for $\gamma p \rightarrow \Lambda K^{+}$for solution 1 (a) and solution 2 (b). The solid curves are the results of our fits, dashed lines are the $P_{13}$ contribution, dotted lines are the $S_{11}$ contribution and dash-dotted lines are the contribution from $K^{*}$ exchange.


The total cross section for $\gamma p \rightarrow \Sigma K^{+}$for solution 1 (a) and solution 2 (b). The solid curves are the results of our fits, dashed lines are the $P_{13}$ contribution, dash-dotted lines are the $P_{11}$ contribution and dotted lines are the contribution from $K$ exchange.

$\gamma p \rightarrow \Lambda K^{+}$(left) and $\gamma p \rightarrow \Sigma K^{+}$(right). Only energy points where $C_{x}$ and $C_{z}$ were measured are shown. The solution 1 is shown as solid line and solution 2 (hardly visible since overlapping) as a dashed line.


Real (a) and imaginary (b) part of the $P_{13}$ elastic scattering amplitude solution 1 and solution 2

Polarization variables are related as: $C_{x}^{2}+C_{z}^{2} \leq \operatorname{Min}\left(\left(1-\Sigma^{2}\right),\left(1-P^{2}\right)\right)$


$\gamma \mathrm{p} \rightarrow \mathrm{K}^{+} \Lambda$ (left) and $\gamma \mathrm{p} \rightarrow \mathrm{K}^{+} \Sigma^{0}$ (right) from CLAS (open circle) and GRAAL (black circle).
The solid and dashed curves are solution 1 and 2 respectively.


The beam asymmetries as a function of $W$ for $\gamma \mathrm{p} \rightarrow \mathrm{K}^{+} \Lambda$ (left) and $\gamma \mathrm{p} \rightarrow \mathrm{K}^{+} \Sigma$ (right). The solid and dashed curves are the result of our fit obtained with solution 1 and 2 , respectively.

The masses, widths are given in MeV, the branching ratios in \% and helicity couplings in $\mathbf{1 0}^{-3} \mathrm{GeV}^{-1 / 2}$. The helicity couplings and phases were calculated as residues in the pole position.

|  | Solution 1 |  | Solution 2 |  |
| :--- | :---: | :---: | :---: | :---: |
| $M_{\text {pole }}$ | $1640 \pm 80$ | $1870 \pm 15$ | $1630 \pm 60$ | $1960 \pm 15$ |
| $\Gamma_{\text {tot }}^{\text {ole }}$ | $480 \pm 60$ | $170 \pm 30$ | $440 \pm 60$ | $195 \pm 25$ |
| $M_{B W}$ | $1800 \pm 100$ | $1885 \pm 15$ | $1780 \pm 80$ | $1975 \pm 15$ |
| $\Gamma_{\text {tot }}^{B W}$ | $700 \pm 100$ | $180 \pm 25$ | $680 \pm 80$ | $200 \pm 25$ |
| $A_{1 / 2}$ | $140 \pm 80$ | $-(15 \pm 15)$ | $160 \pm 40$ | $-(18 \pm 8)$ |
| $\varphi_{1 / 2}$ | $-(10 \pm 15)^{\circ}$ | - | $(10 \pm 15)^{\circ}$ | $(40 \pm 15)^{\circ}$ |
| $A_{3 / 2}$ | $150 \pm 80$ | $-(40 \pm 15)$ | $70 \pm 30$ | $-(35 \pm 12)$ |
| $\varphi_{3 / 2}$ | $-(40 \pm 30)^{\circ}$ | $-(20 \pm 15)^{\circ}$ | $(0 \pm 20)^{\circ}$ | $-(40 \pm 15)^{\circ}$ |
| $\operatorname{Br}_{N \pi}$ | $8 \pm 4$ | $5 \pm 3$ | $11 \pm 4$ | $6 \pm 3$ |
| $\operatorname{Br}_{N \eta}$ | $13 \pm 4$ | $21 \pm 8$ | $5 \pm 2$ | $15 \pm 3$ |
| $\operatorname{Br}_{\Delta \pi(P)}$ | $48 \pm 10$ | $3 \pm 2$ | $28 \pm 6$ | $7 \pm 2$ |
| $\operatorname{Br}_{\Delta \pi(F)}$ | $2 \pm 2$ | $4 \pm 3$ | $11 \pm 4$ | $21 \pm 5$ |
| $\operatorname{Br}_{K \Lambda}$ | $15 \pm 6$ | $10 \pm 5$ | $5 \pm 2$ | $12 \pm 3$ |
| $\operatorname{Br}_{K \Sigma}$ | $<1$ | $20 \pm 8$ | $<1$ | $8 \pm 2$ |
| $\operatorname{Br}_{D 13} \pi$ | $10 \pm 6$ | $8 \pm 3$ | $38 \pm 6$ | $5 \pm 3$ |
| $\operatorname{Br}_{N \sigma}$ | $4 \pm 2$ | $30 \pm 12$ | $2 \pm 2$ | $26 \pm 8$ |
|  |  |  |  |  |



Predictions for T polarization.
Could be possible to choose between three solutions.

## Conclusions

- New $P_{13}$ state has found fitting new CLAS $C_{x}, C_{z}$ data
- A qualitatively good description of all fitted observables was obtained.
- The analysis of the first double polarization data on hyperon photoproduction reveals an exciting result.
- Systematic measurements with further single and double polarization observables - as being planned and carried out at several laboratories - are urgently needed.

