

Jet physics in ATLAS

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QCD, Heavy Flavours and Higgs physics

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Jets in the LHC era

At the Large Hadron Collider (LHC), jet production is the dominant high transverse-momentum (p_T) process.

It gives the first glimpse of physics at the TeV scale.

Jet cross sections and properties are key observables in high-energy particle physics.

Measured in e^+e^- , ep , $p\bar{p}$, and pp colliders, and in γp and $\gamma\gamma$ collisions.

- Measurements of the strong coupling constant.
- Information about the structure of the proton and photon.
- Tools for understanding the strong interaction
- Tools for searching for physics beyond the Standard Model.

ATLAS Detector overview

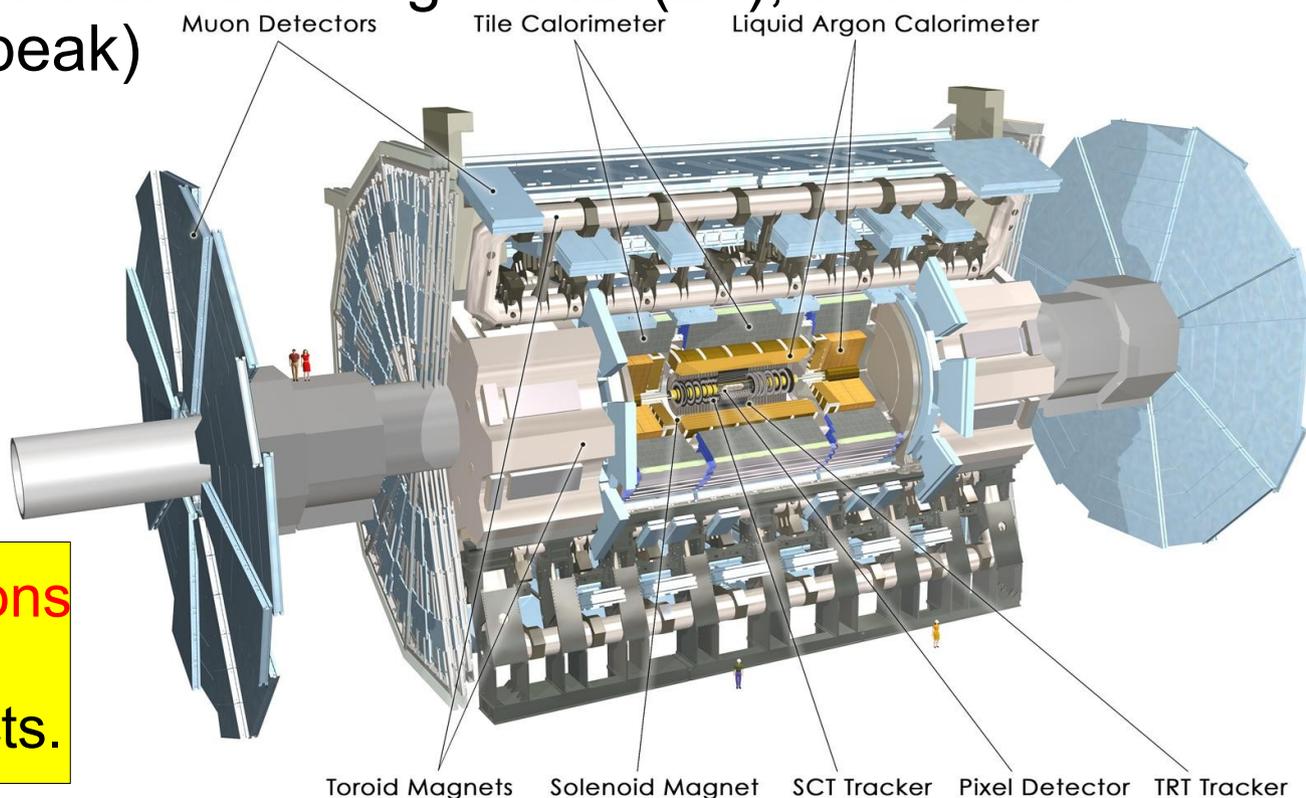
Magnetic field: **one solenoid** surrounding the ID (2T), **one toroid** (muon spectrometer - 4T peak)

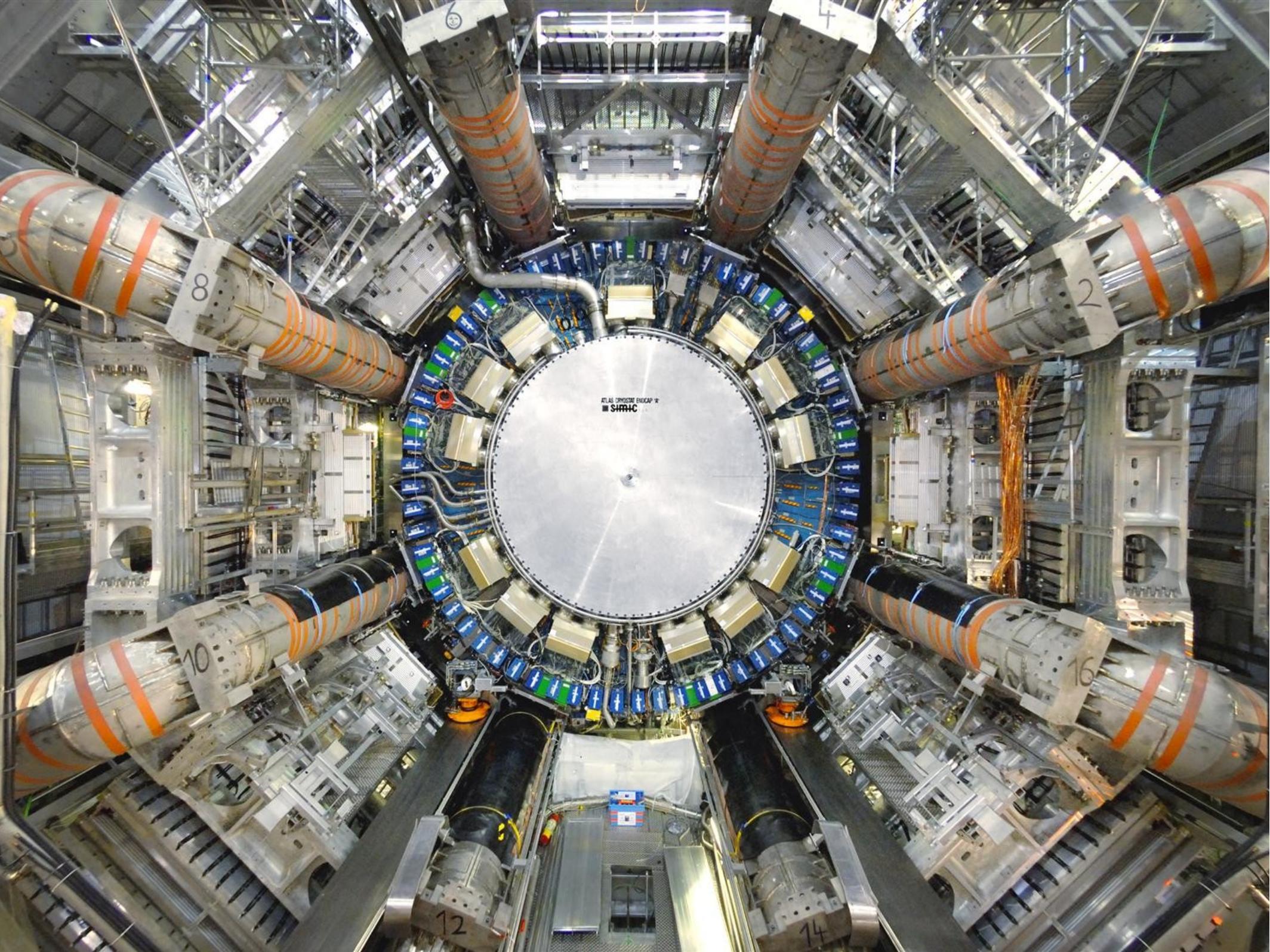
ID made up of **three different detectors** (Pixel, SCT, TRT):
High resolution tracking in $|\eta| < 2.5$

EM calorimeter - **two sections** covering up to $|\eta| \approx 3.2$.
High resolution on e/ γ objects.

HAD calorimeter - **3 sections** covering up to $|\eta| \approx 5$
Good containment, good resolution for jet measurement

Muon system (**4 different technologies**) covering up to $|\eta| = 2.7$
High precision muon momentum measurement (also standalone)





ATLAS EXPERIMENT ENDCAP #
SIRIAC

8

2

10

16

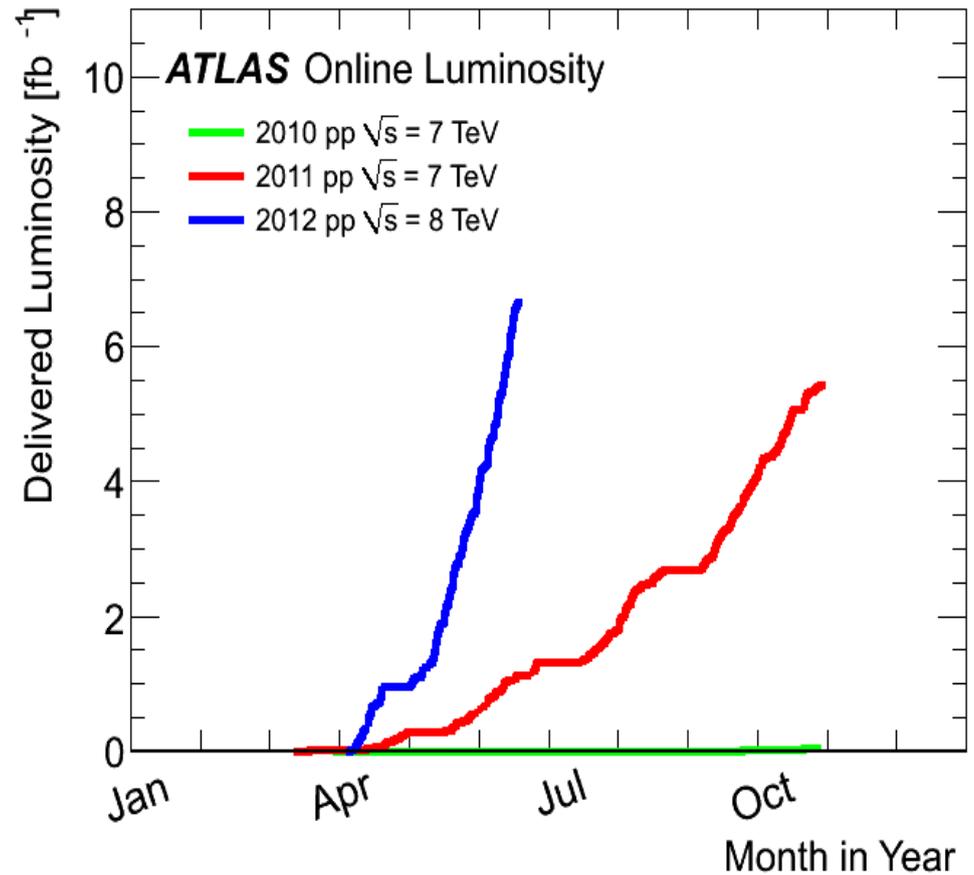
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14

Data Taking

Successful operation for the LHC machine and for the ATLAS experiment.

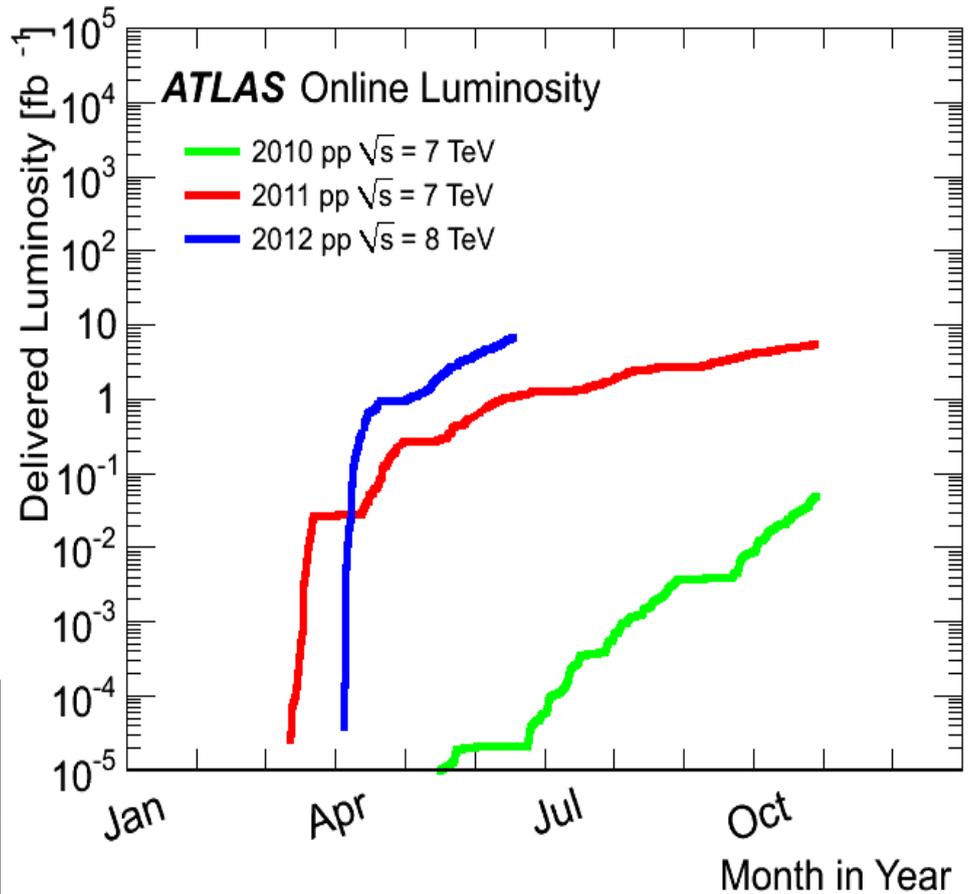
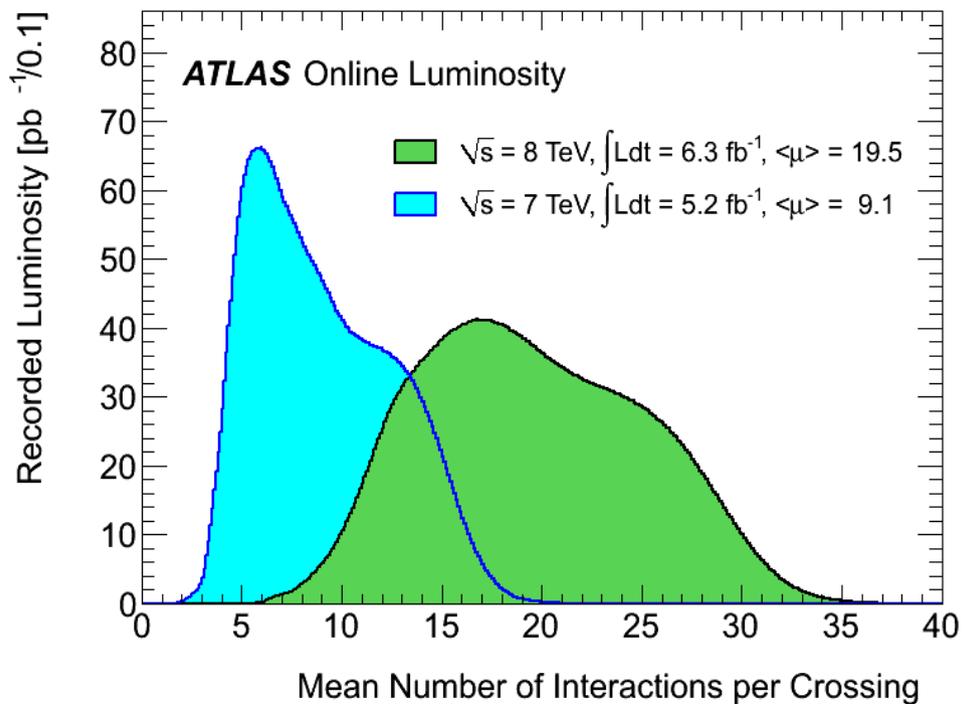
~40 pb⁻¹ in 2010 (sqrt(s)=7TeV)
~5 fb⁻¹ in 2011 (sqrt(s)=7TeV)
~ 6 fb⁻¹ june 2012 (sqrt(s)=8TeV)



Data Taking

Successful operation for the LHC machine and for the ATLAS experiment.

$\sim 40 \text{ pb}^{-1}$ in 2010 ($\sqrt{s}=7\text{TeV}$)
 $\sim 5 \text{ fb}^{-1}$ in 2011 ($\sqrt{s}=7\text{TeV}$)
 $\sim 6 \text{ fb}^{-1}$ june 2012 ($\sqrt{s}=8\text{TeV}$)



Change in pile-up conditions:

In 2012 ~ 9 interactions per bunch crossing.

Outlook

In this talk:

→ Jet reconstruction and performance

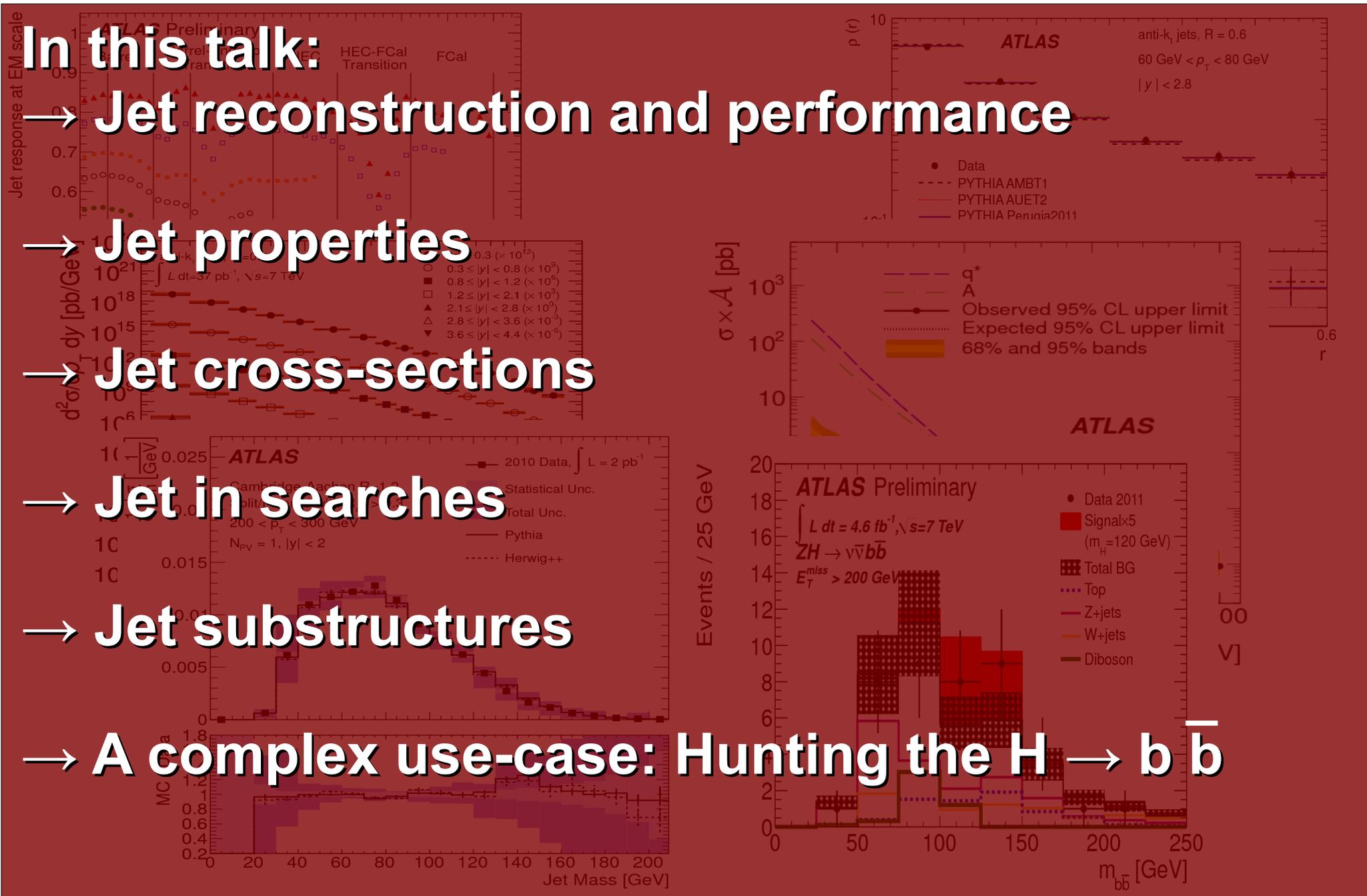
→ Jet properties

→ Jet cross-sections

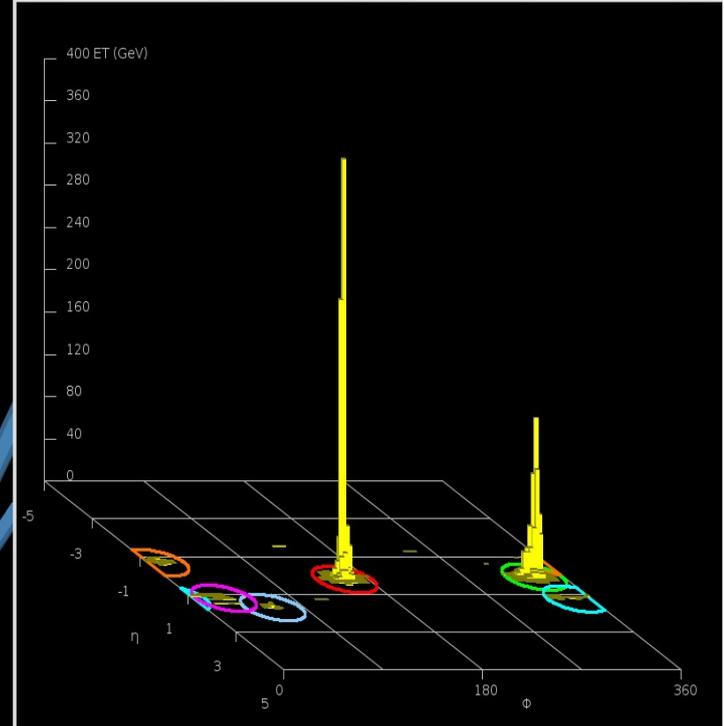
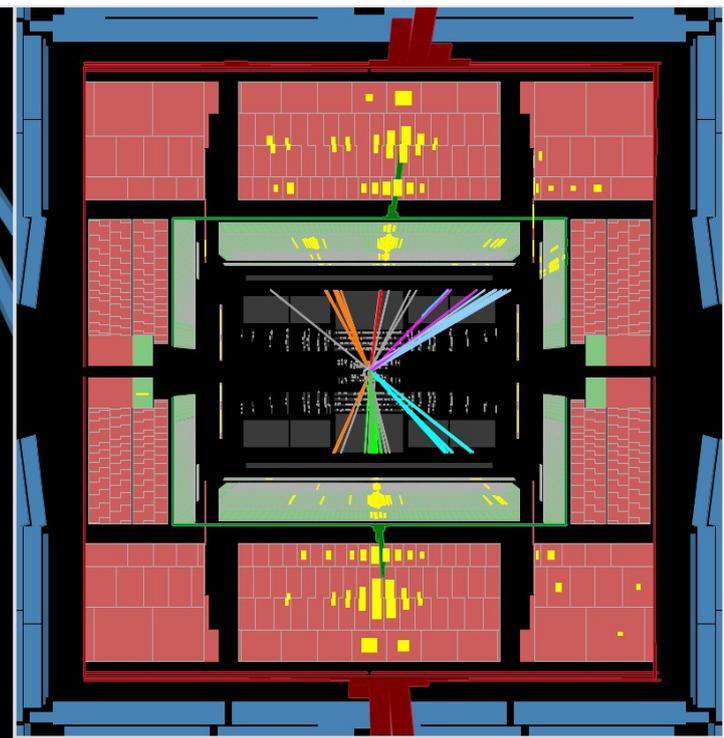
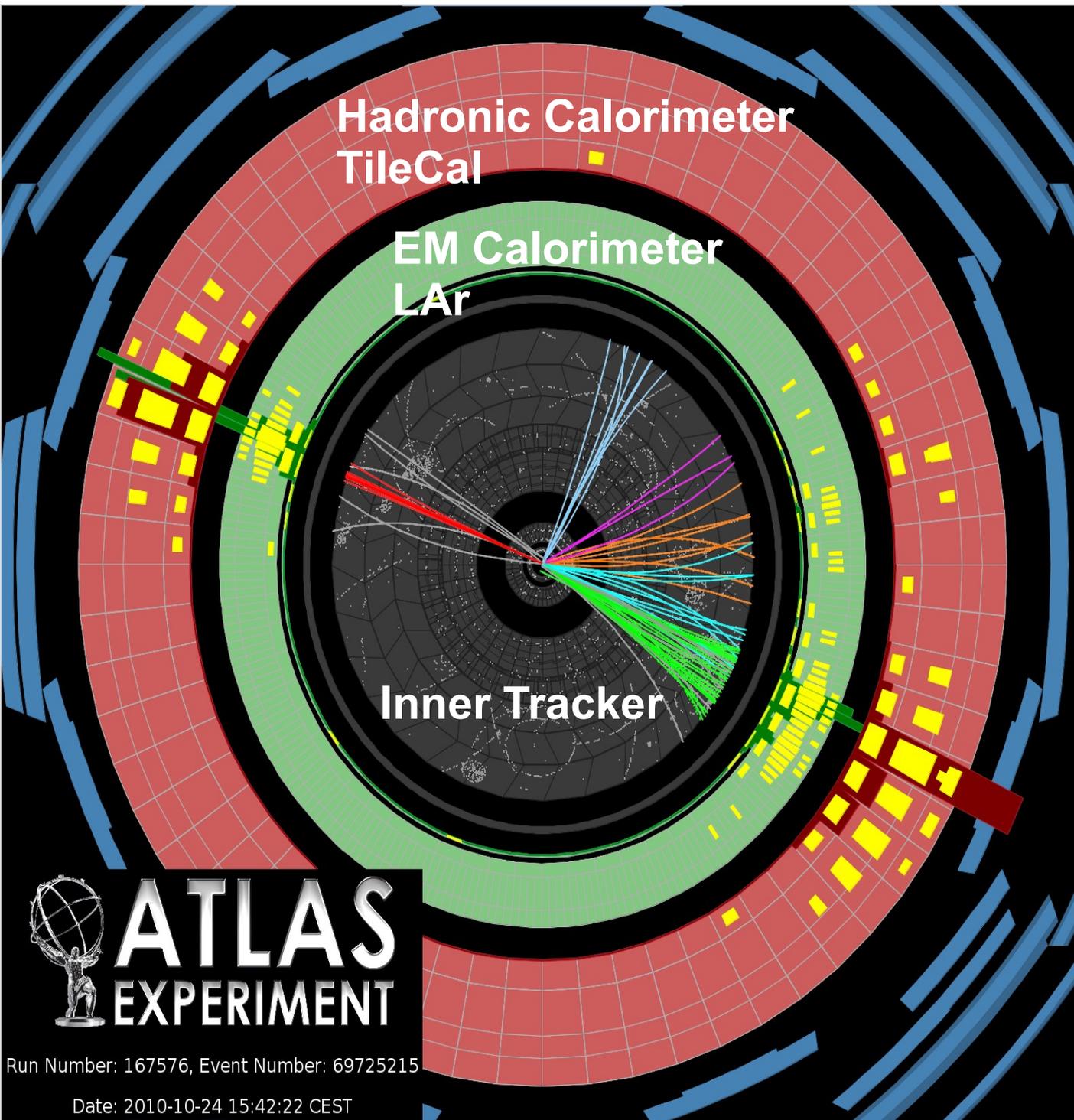
→ Jet in searches

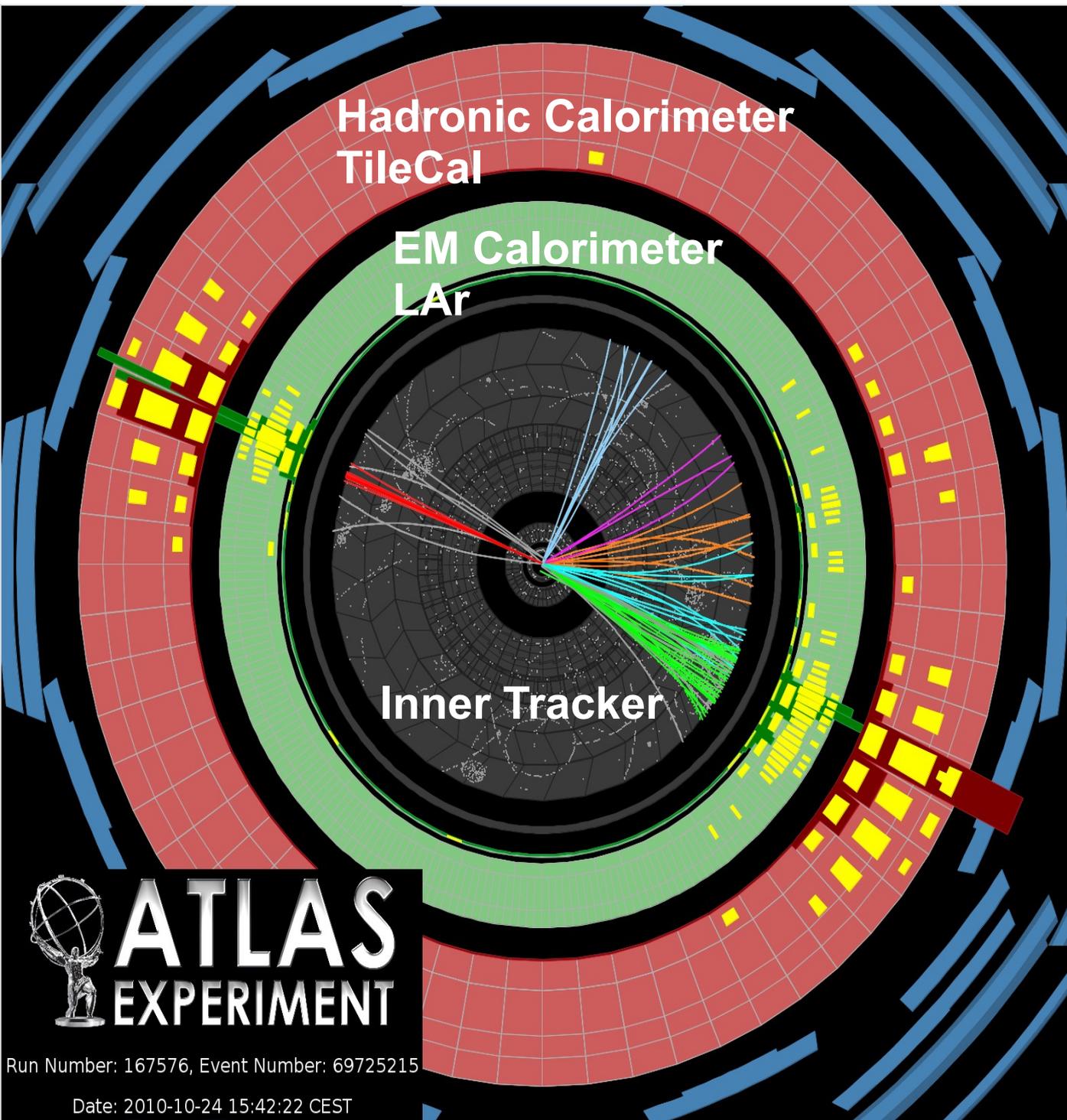
→ Jet substructures

→ A complex use-case: Hunting the $H \rightarrow b\bar{b}$

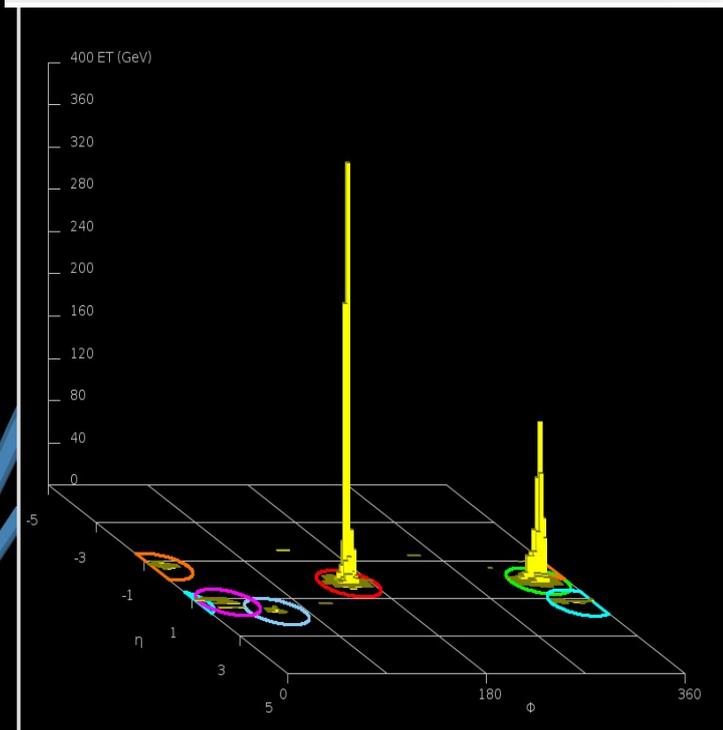
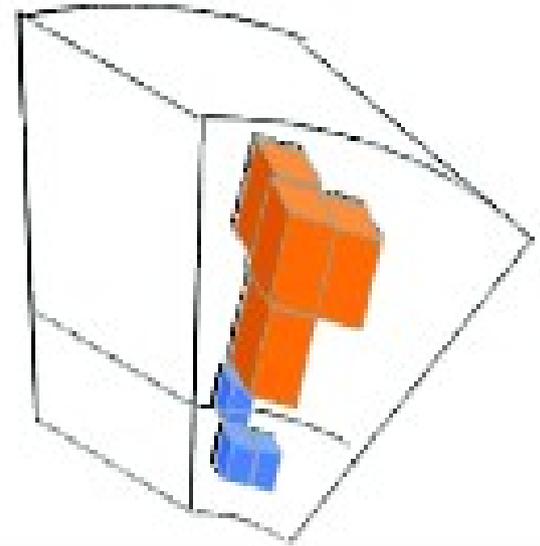


Jet reconstruction and performance





Topological clusters



Jet Reconstruction

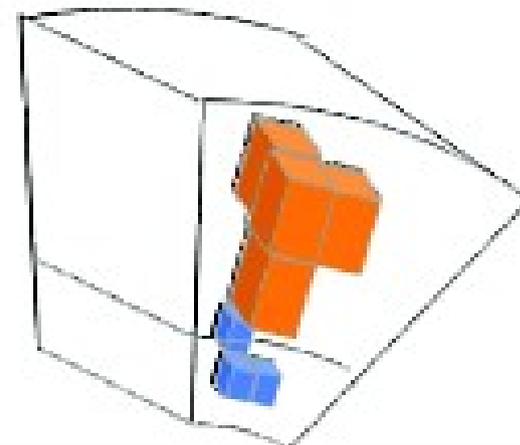
Topological clusters

Inputs:

3D Calorimetric clusters:

find local cell energy maxima and
cluster neighboring cells

Pro: noise suppression



Jet Algorithm: JHEP 0804 (2008) 063 [0802.1189]

The Anti- K_T (infrared safe) algorithm has been taken as the default jet algorithm.

$$d_{ij} = \min(k_{ti}^{2p}, k_{tj}^{2p}) \Delta R_{ij}^2 / R^2$$

The Anti- K_T is a sequential recombination jet algorithms with $p = -1$,

(K_T , $p=1$) which behaves like an idealized cone algorithm.

Cambridge/Aachen ($p=0$) and K_T ($p=1$) algorithms used for jet-substructure studies.

C/A algorithm has a distance which depends only on angular distance.

Calorimetric calibration

Electromagnetic (EM) scale:
Baseline cluster calibration, established using test beam with e and μ in the calorimeters. Good estimate of the energy deposited by γ and e . 60-70% estimate of the energy deposited for hadrons and jets

Hadronic Calibration. Why?:

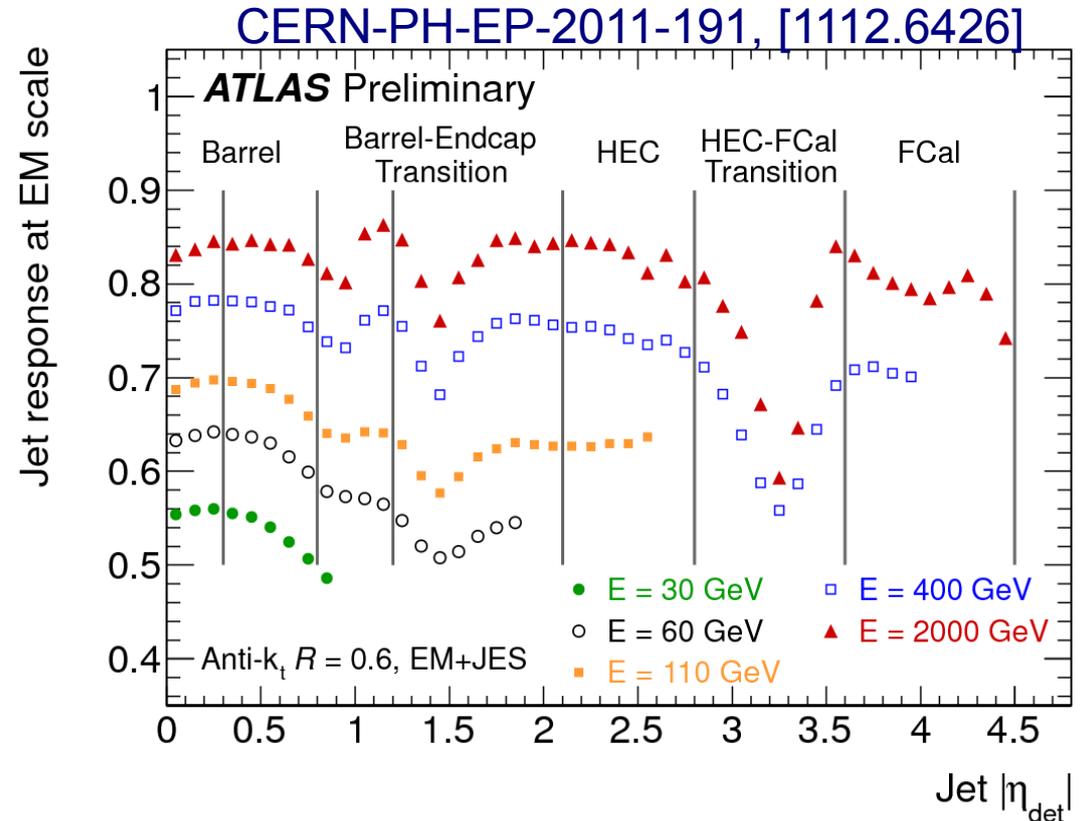
In the ATLAS Calorimeters,

- Response to hadrons lower than response to electrons.
- Energy losses in inactive regions of the detector.
- Out of cone effects.

Hadronic Jet Calibration driven by MC description (**EM+JES**).

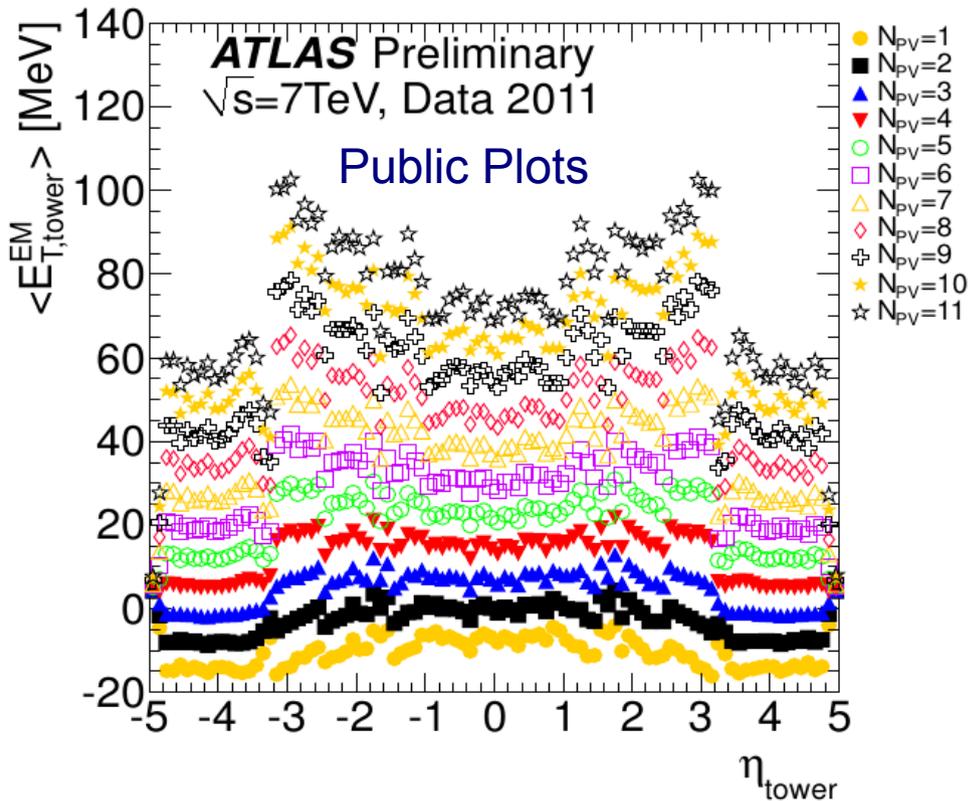
1) Pile-up subtraction accounting for in-time and out-of-time pile-up

2) Correction factor: $p_T^{\text{Calibrated}} = C(E^{\text{EM}}, \eta) p_T^{\text{EM}}$

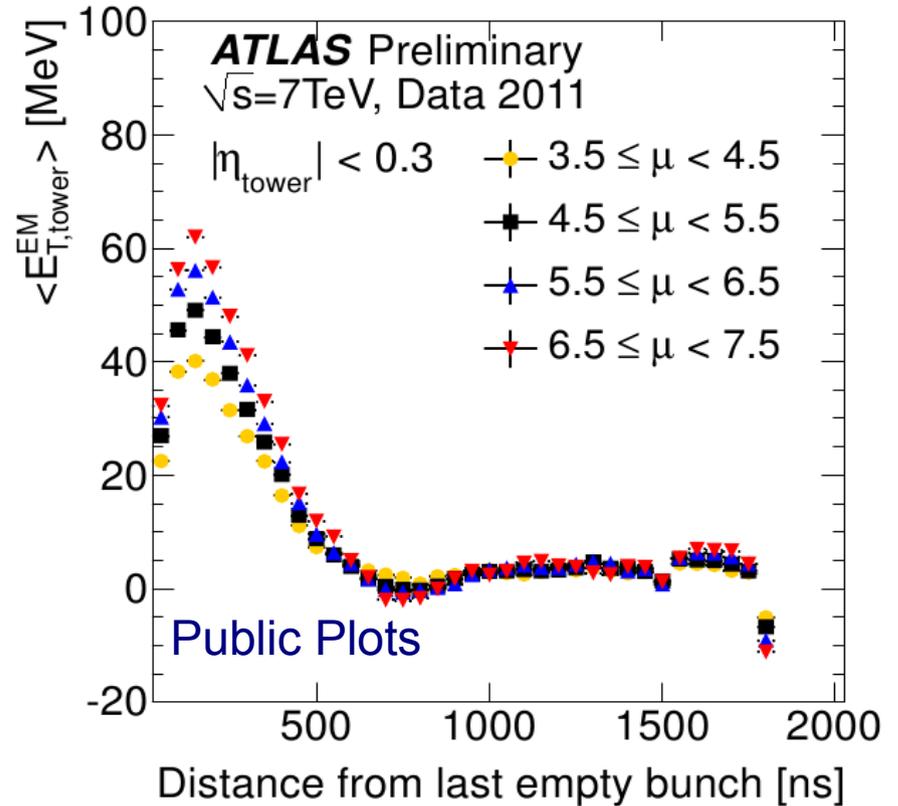


Pile-up correction

In-time pile-up (2010-2011-2012):
multiple interactions in same
bunch crossing
additional soft diffuse radiation



Out-of-time pile-up (2011-2012):
overlapping signal from collisions in
other bunch crossings
affects calorimeter energy
reconstruction



Checks on the EM scale simulation

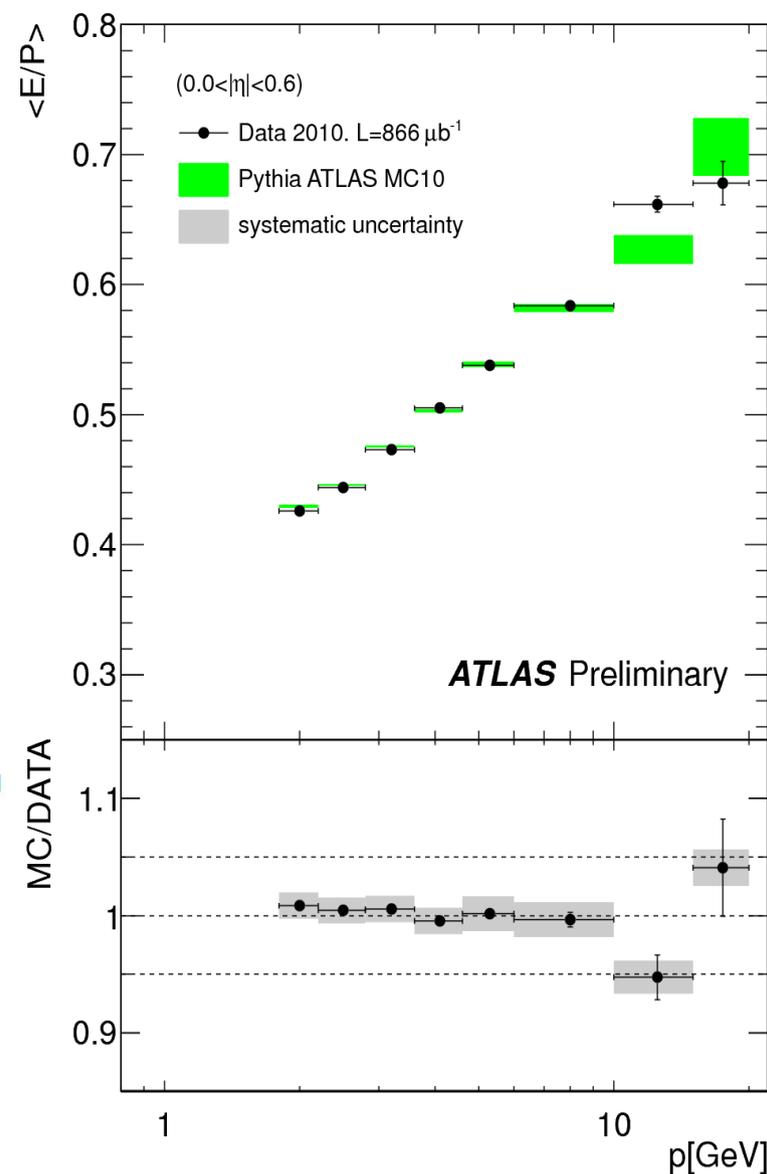
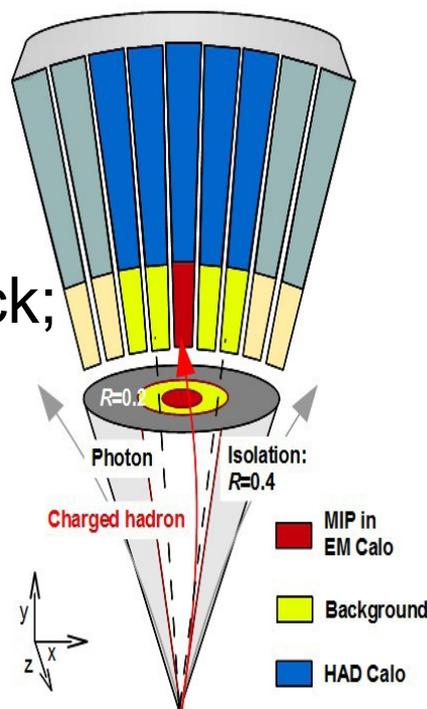
The simulation has been validated using test-beam and collision data.

Checks on Collisions data:

- Select **isolated tracks**;
- Collect **the energy in the calorimeter** around the track;
- Compare to MC.

$\langle E/P \rangle$ measured in

- $|\eta| < 2.3$
- $2 \text{ GeV} < p < 20 \text{ GeV}$



The calorimeter response to isolated hadrons shows agreement between Data and MC at the **5%** level for most of the calorimeter.

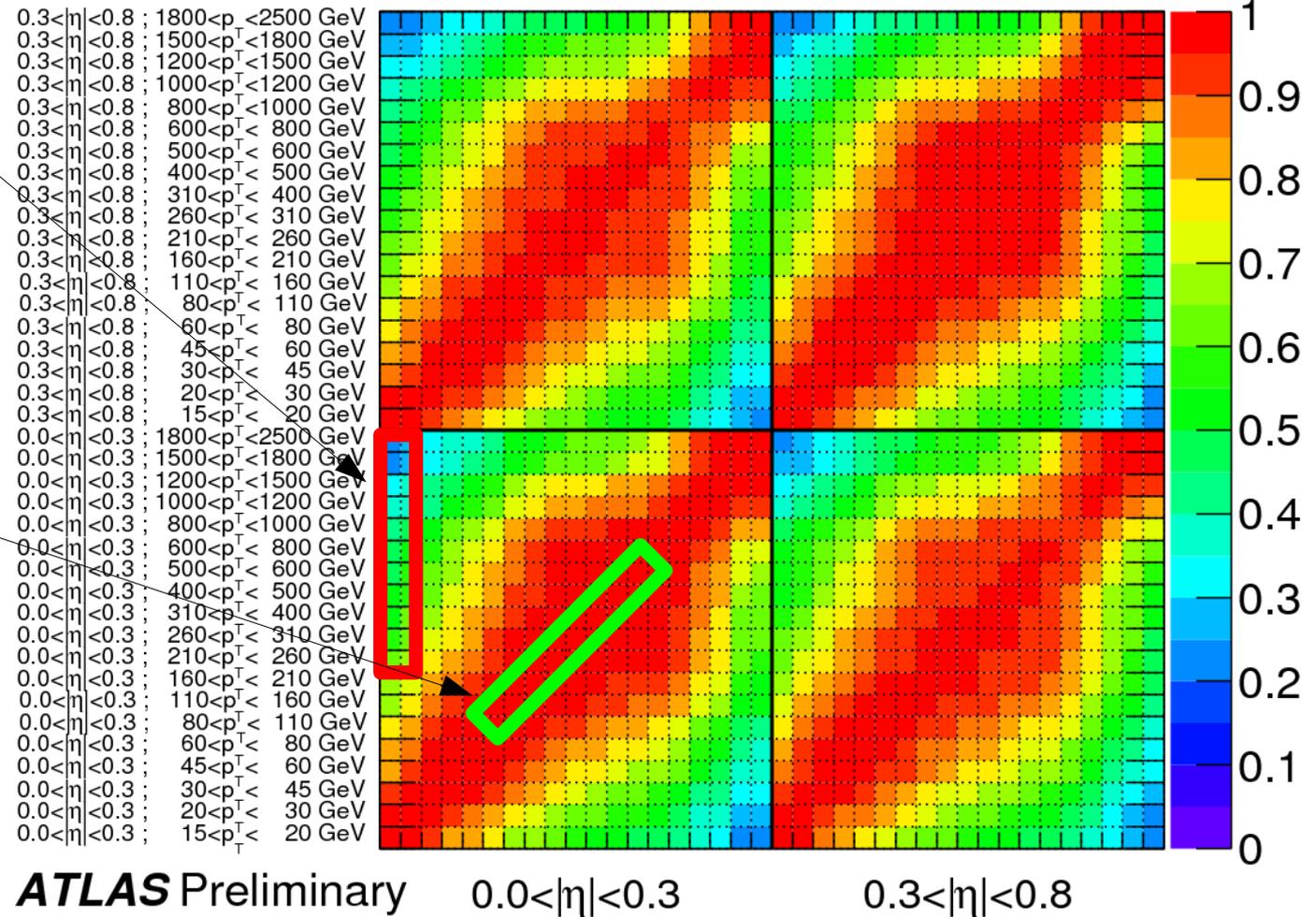
CERN-PH-EP-2012-005 [1203.1302]

Correlations on the calorimeter response

Propagate **single isolated hadron uncertainties** to jets to obtain estimate of calorimeter JES

Long Range correlations

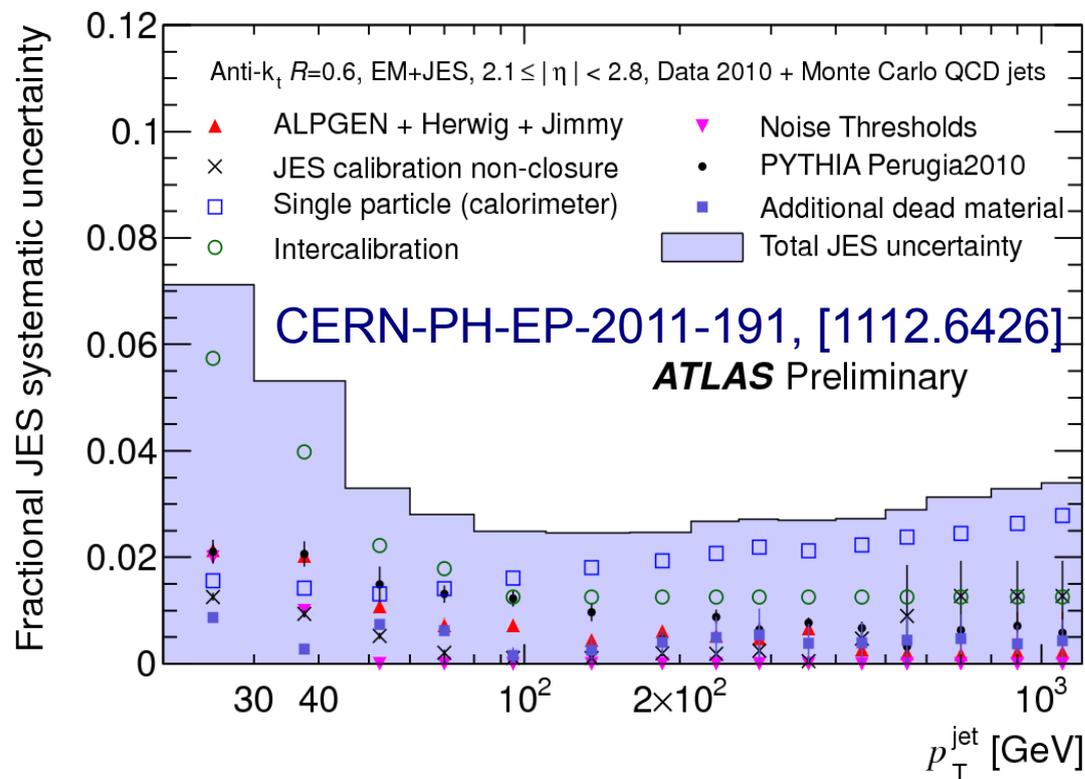
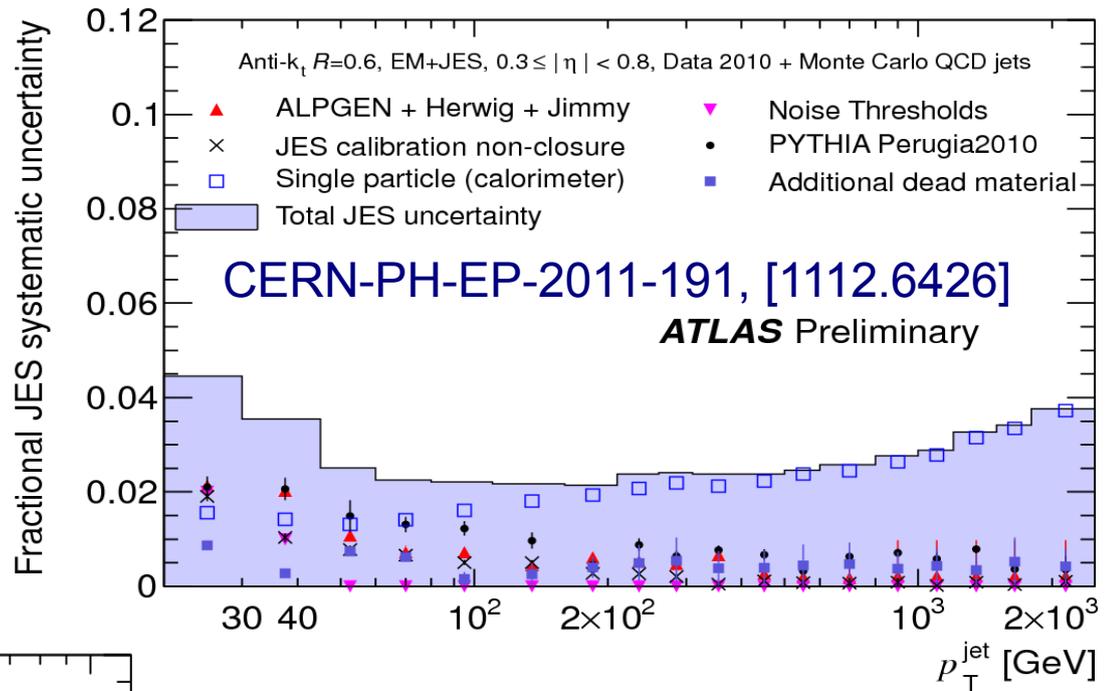
Strong diagonal correlations



CERN-PH-EP-2012-005 [1203.1302]

Jet Energy Scale Uncertainty

Uncertainties:
 ~2% for central jets with
 $p_T \sim 100$ GeV



Dominant:

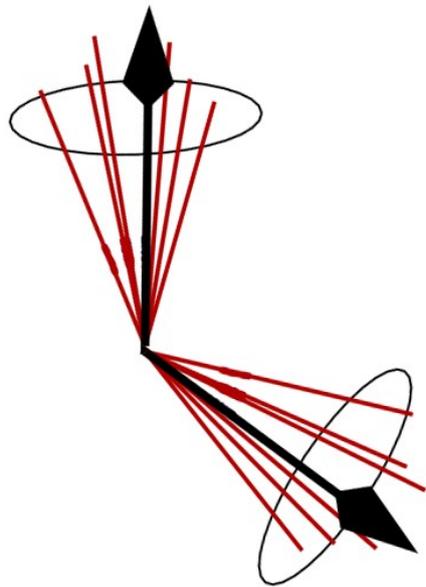
- Single particle response
- Noise description
- Dead Material
- η intercalibration

Smaller:

- Hadronization
- Underlying Event
- Parton Shower
- Pileup

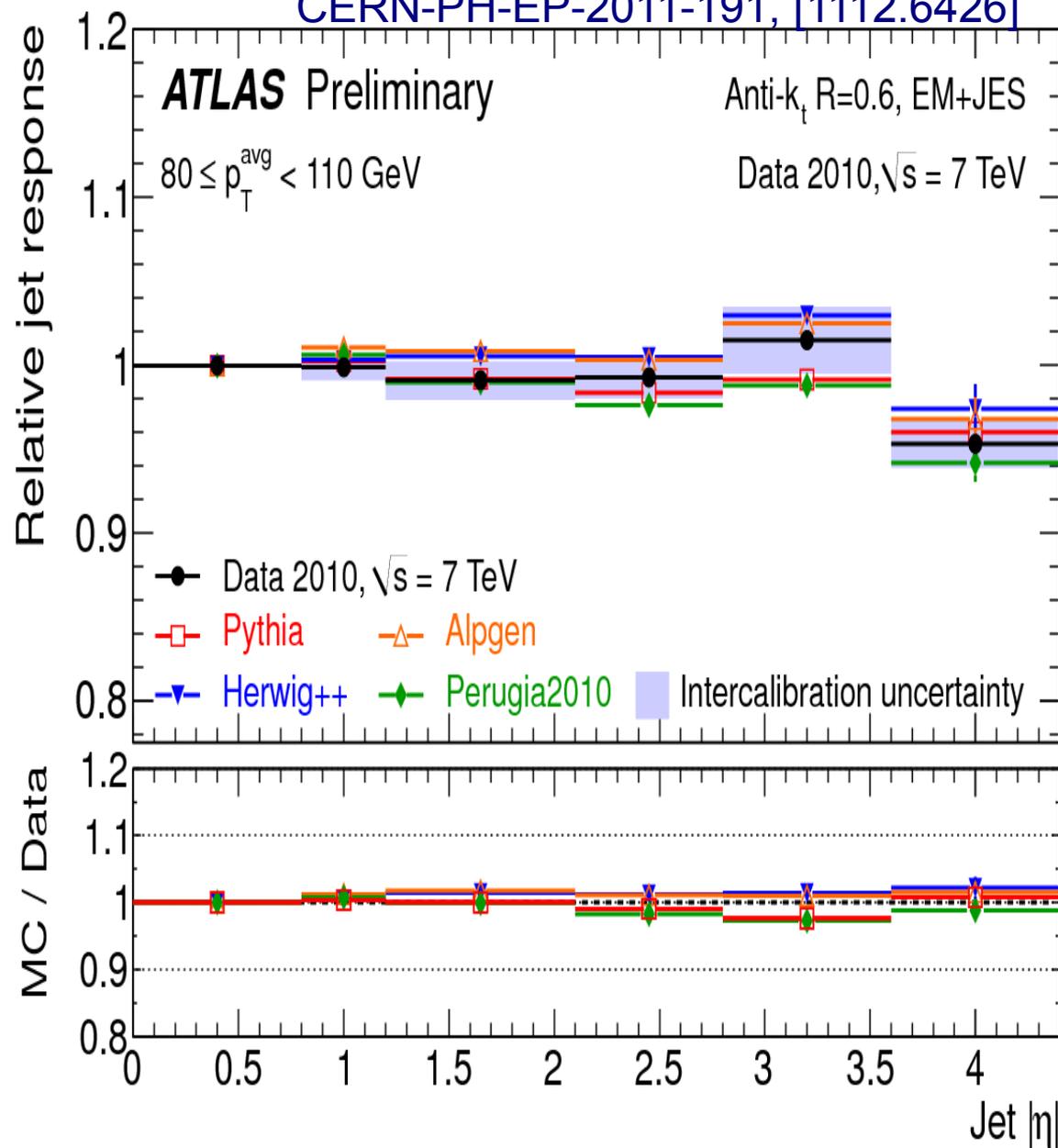
Jet Calibration VS. η

Reference region



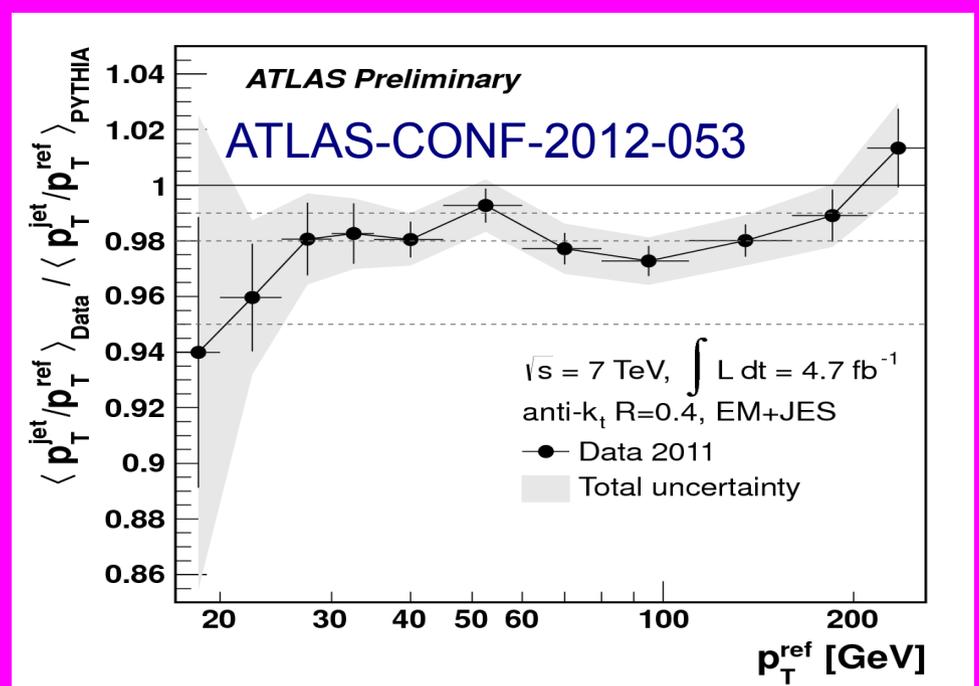
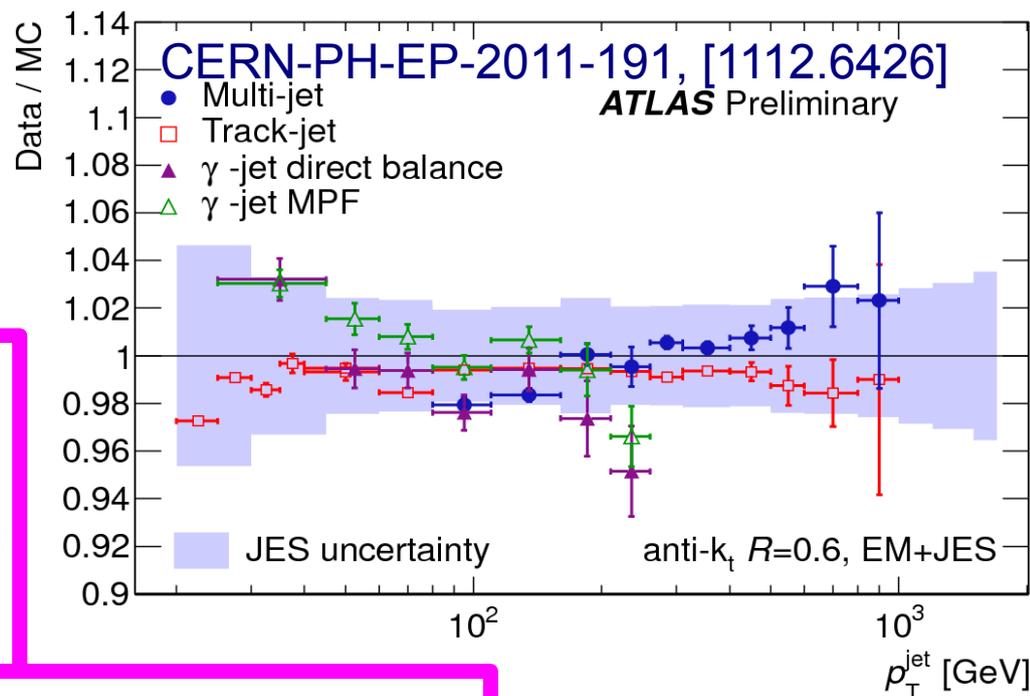
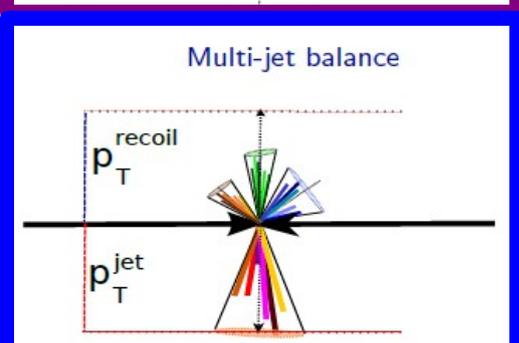
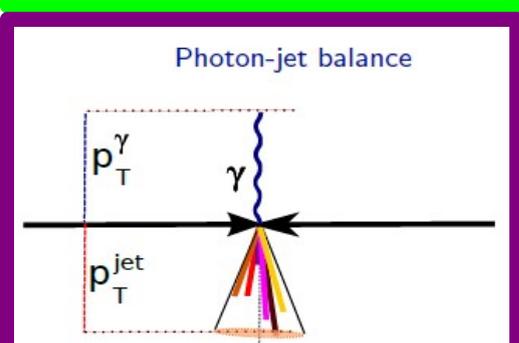
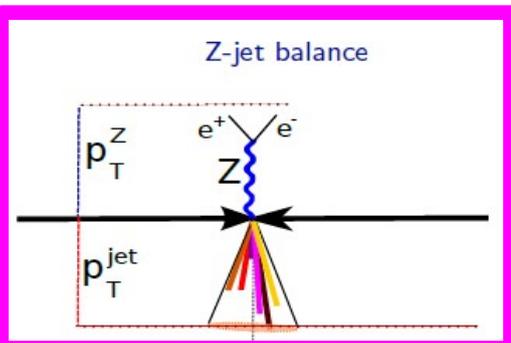
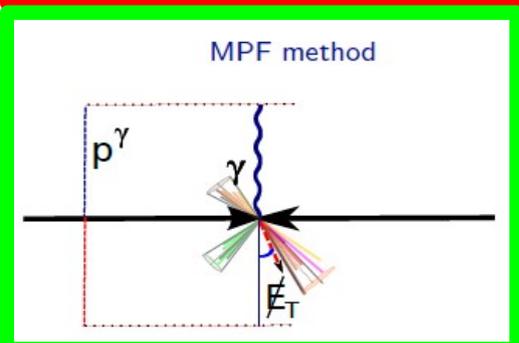
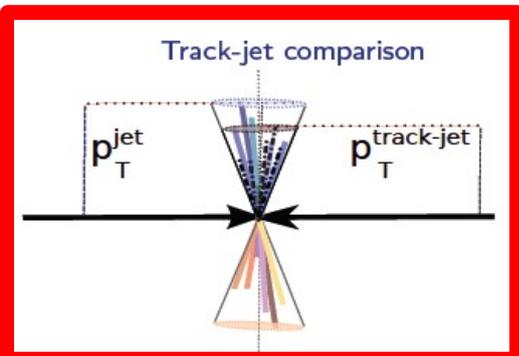
Small (~2%)
deviation at high
pseudo-rapidity

CERN-PH-EP-2011-191, [1112.6426]



Checks on the jet energy scale uncertainty

C. Doglioni



Constrain of the JES by using these measurements:

ONGOING
(public soon).

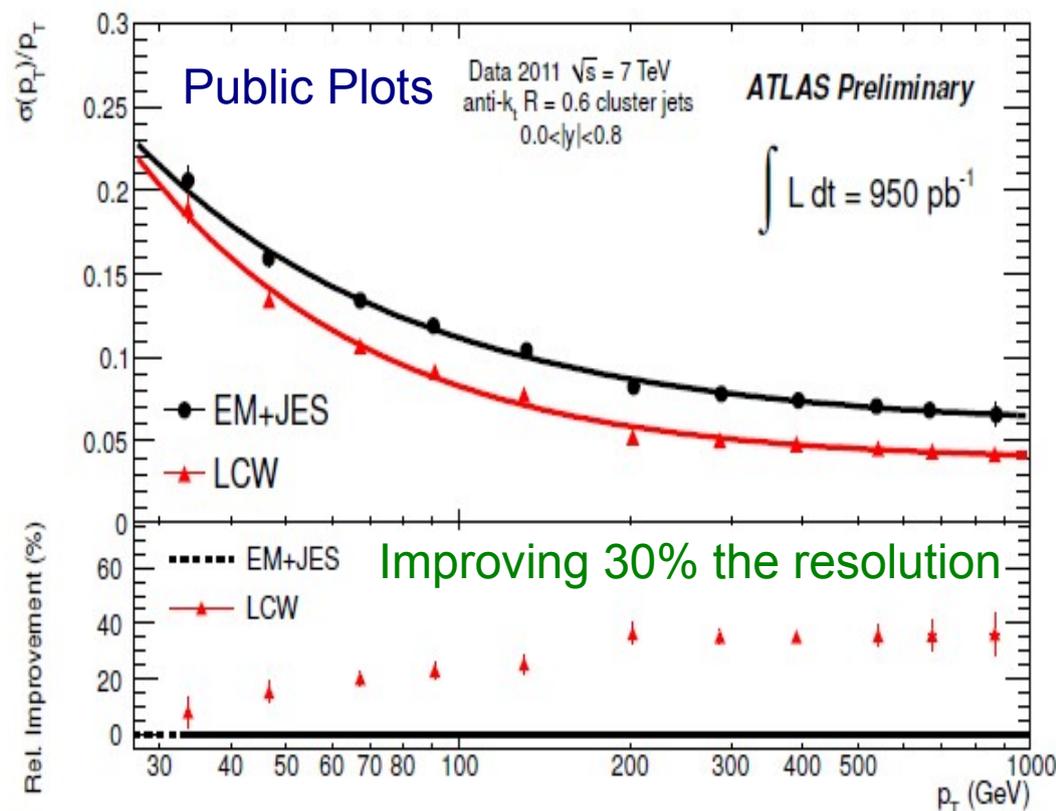
Alternative hadronic calibrations

Local Cluster Weighting (LCW) :

Factorized corrections derived from cluster properties in single pion MC, independent from jet context.

Pro:

- Improving the jet energy resolution
- Calibration does not depend on jet context. More flexibility for substructure studies.
- Coherent hadronic calibration for the clusters which do not belong to any jet (used in the Missing Et)



So why did we use EM+JES for the first analyses?

→ *EM+JES allow direct and fast estimate of uncertainties.*

LCW already used for some measurements in 2011-2012

Heavy flavor jets

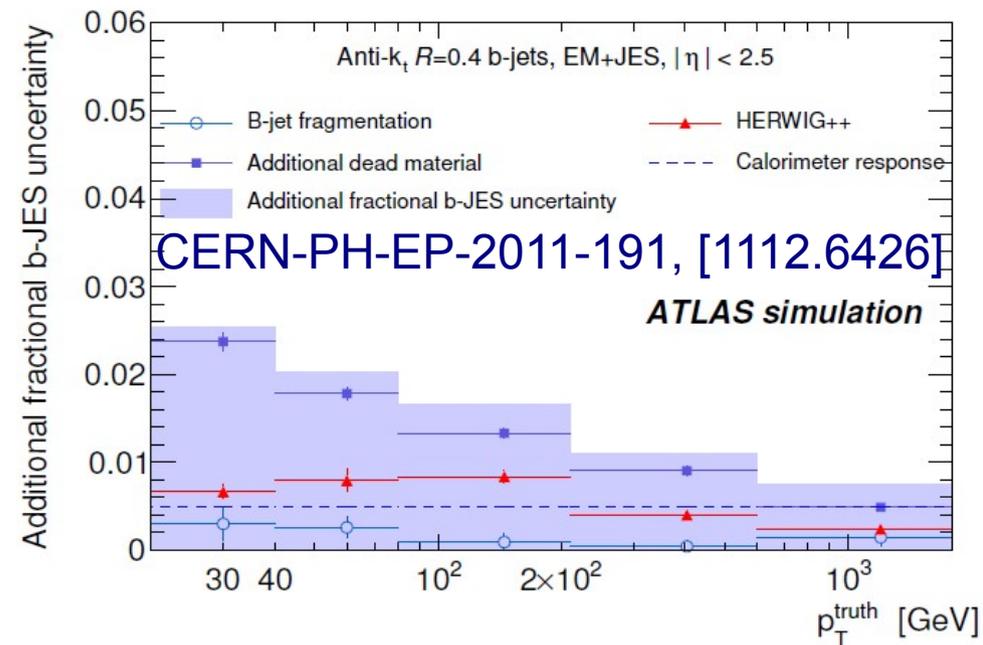
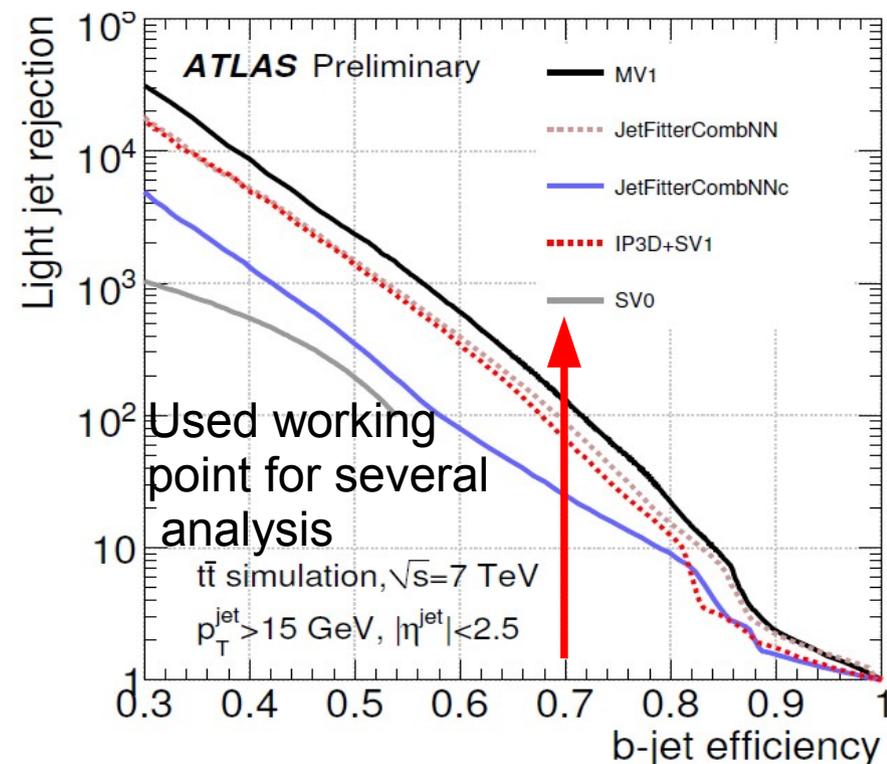
Heavy b-jets identified exploiting the long lifetime of b-hadrons.

Several tagging algorithms combined to improve the performance (MV1).

Calibration done with the standard EM+JES (or LCW).

Ad-hoc estimate of the systematic uncertainty on the JES.

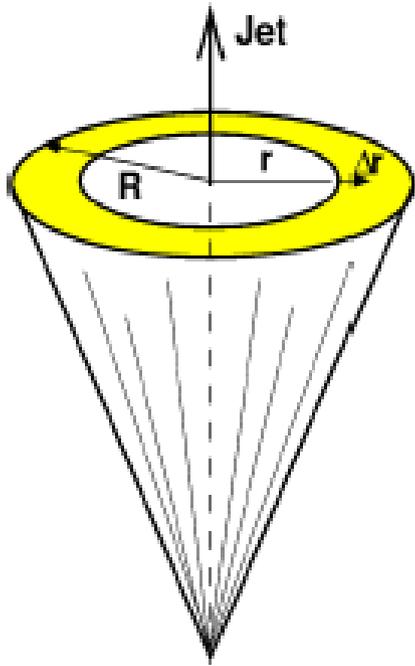
ATLAS-CONF-2012-043



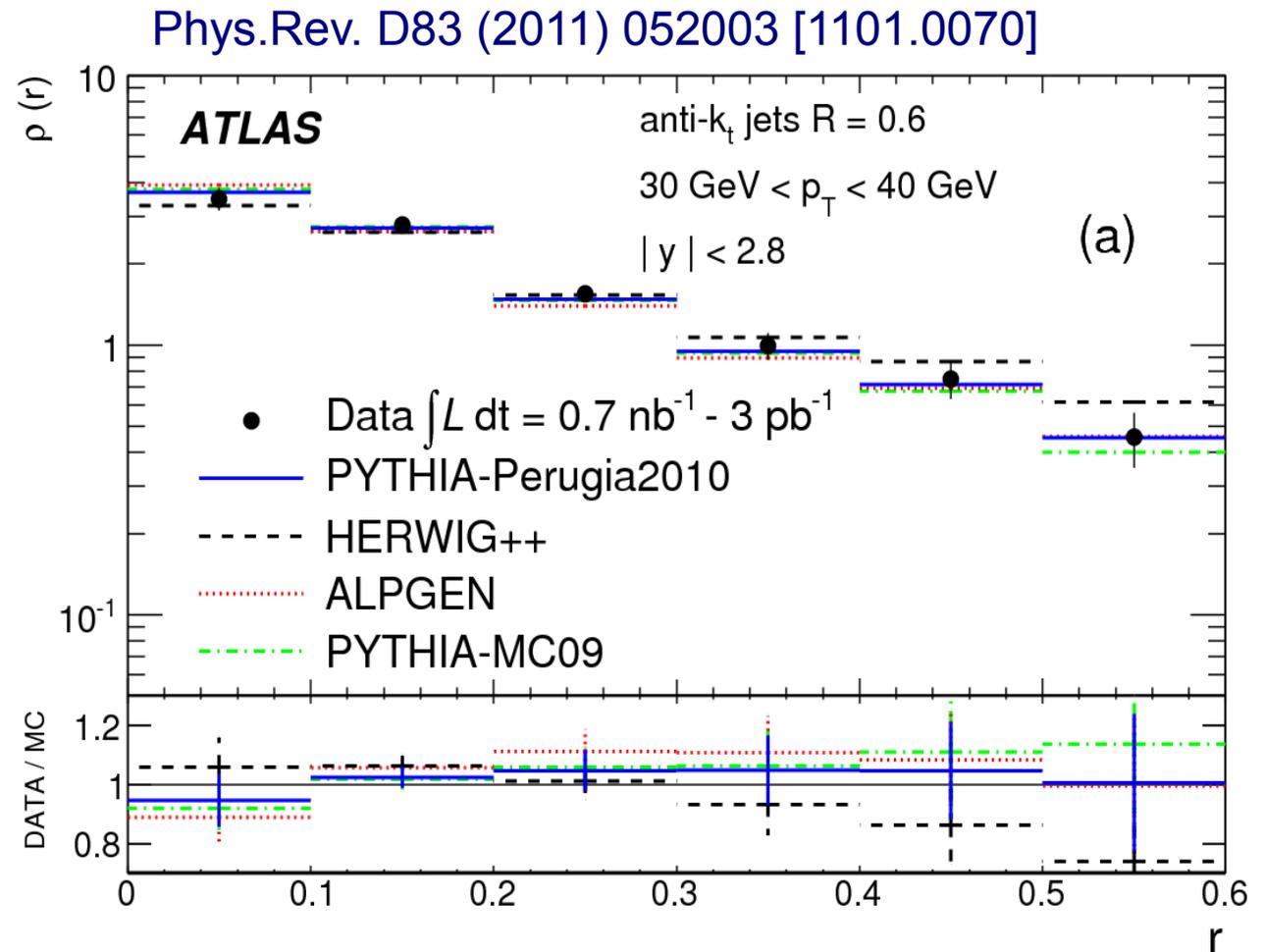
Jet properties: From the constituents to the topologies

Jet Shapes

Energy flow around the jet core

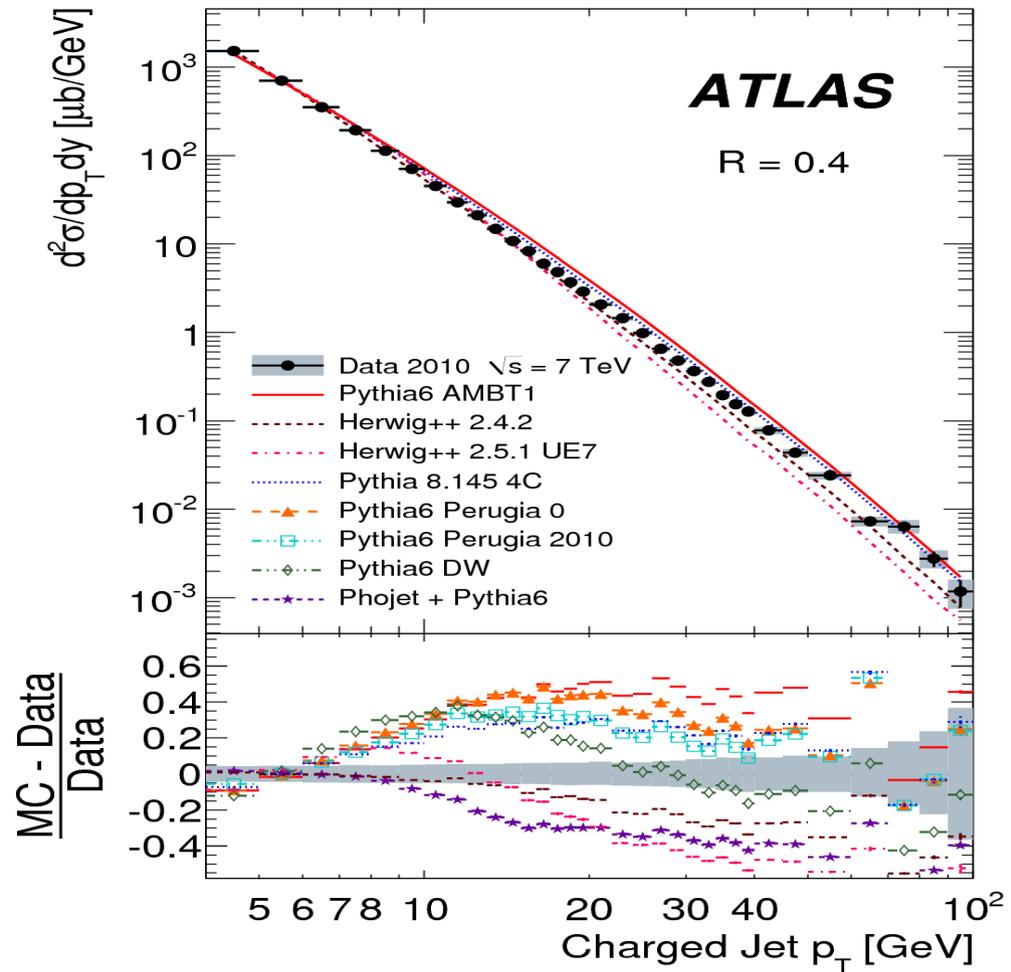
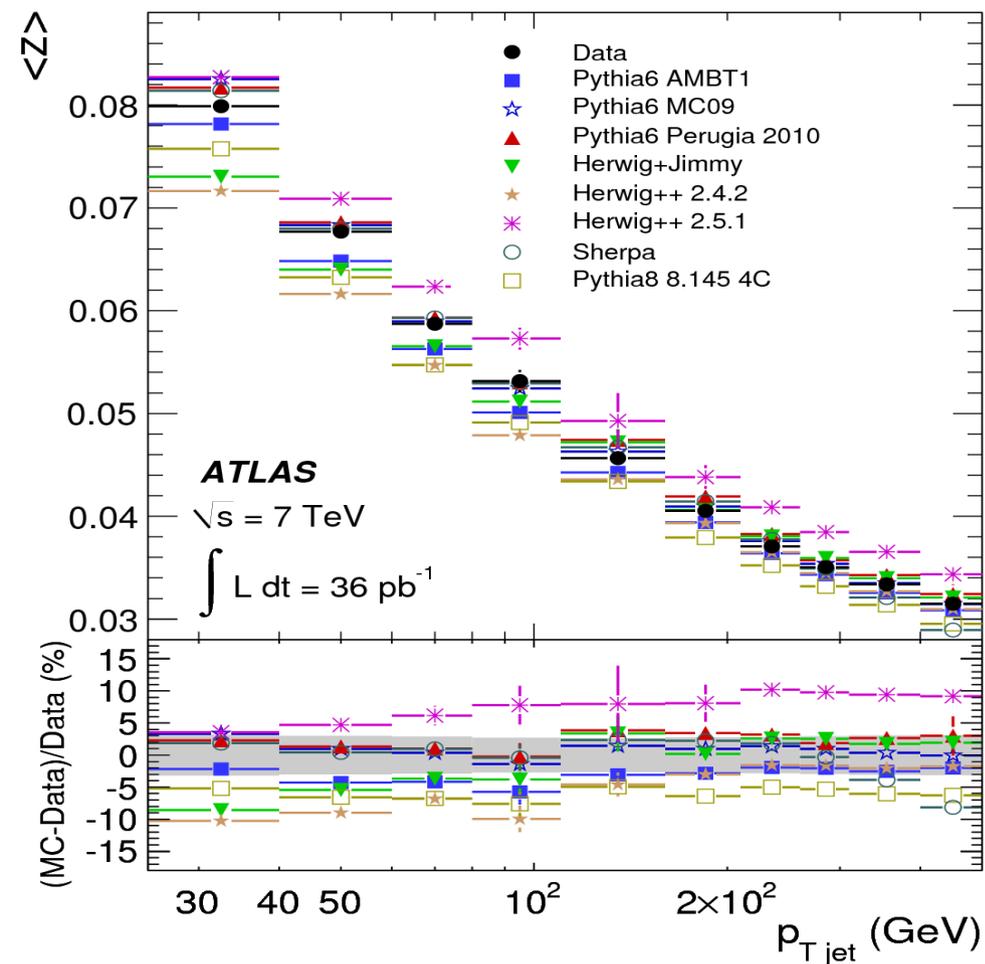


Sensitive to parton-shower, underlying event and fragmentation details.



Tracks in jets: Fragmentation

The tracks are a useful input to the jet clustering to study the jet fragmentation in charged particles and to improve the fragmentation models used in the MC simulations

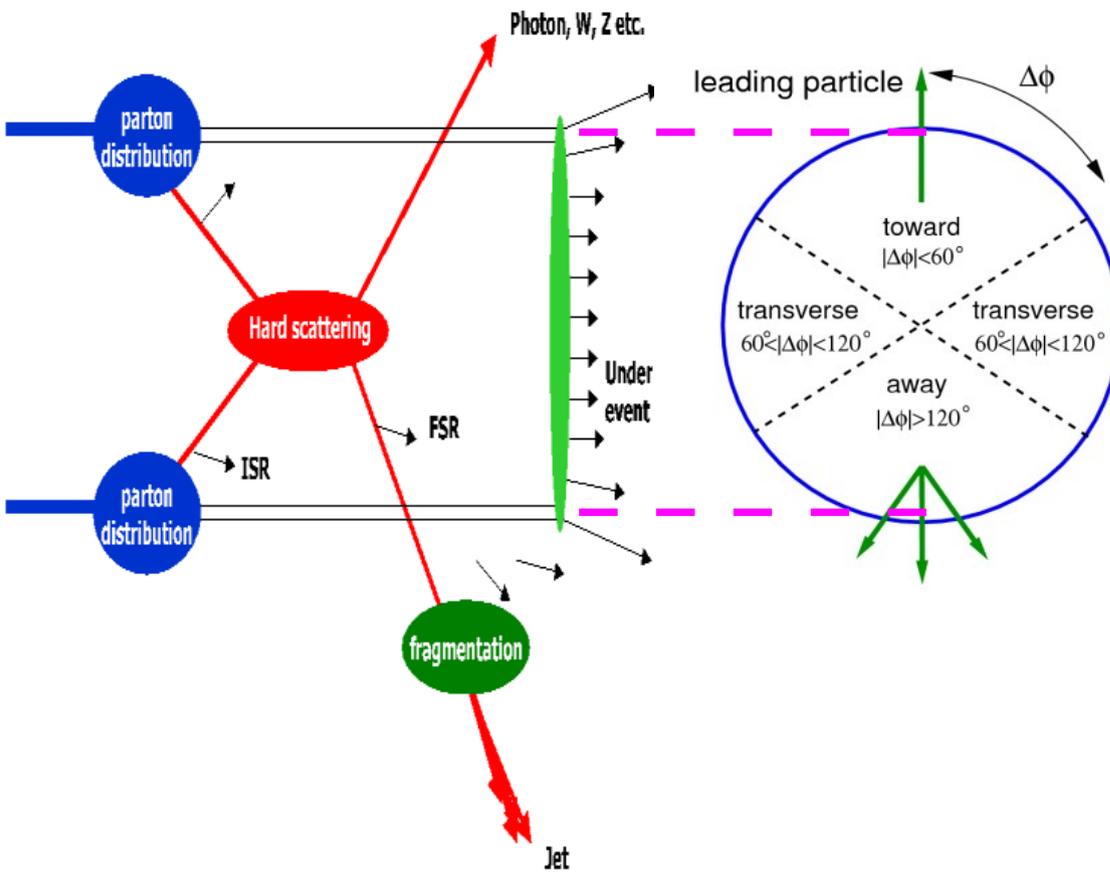


$$Z_{\text{track}} = p_T^{\text{track}} / p_T^{\text{Jet}}$$

is the relative p_T contribution to the jet.

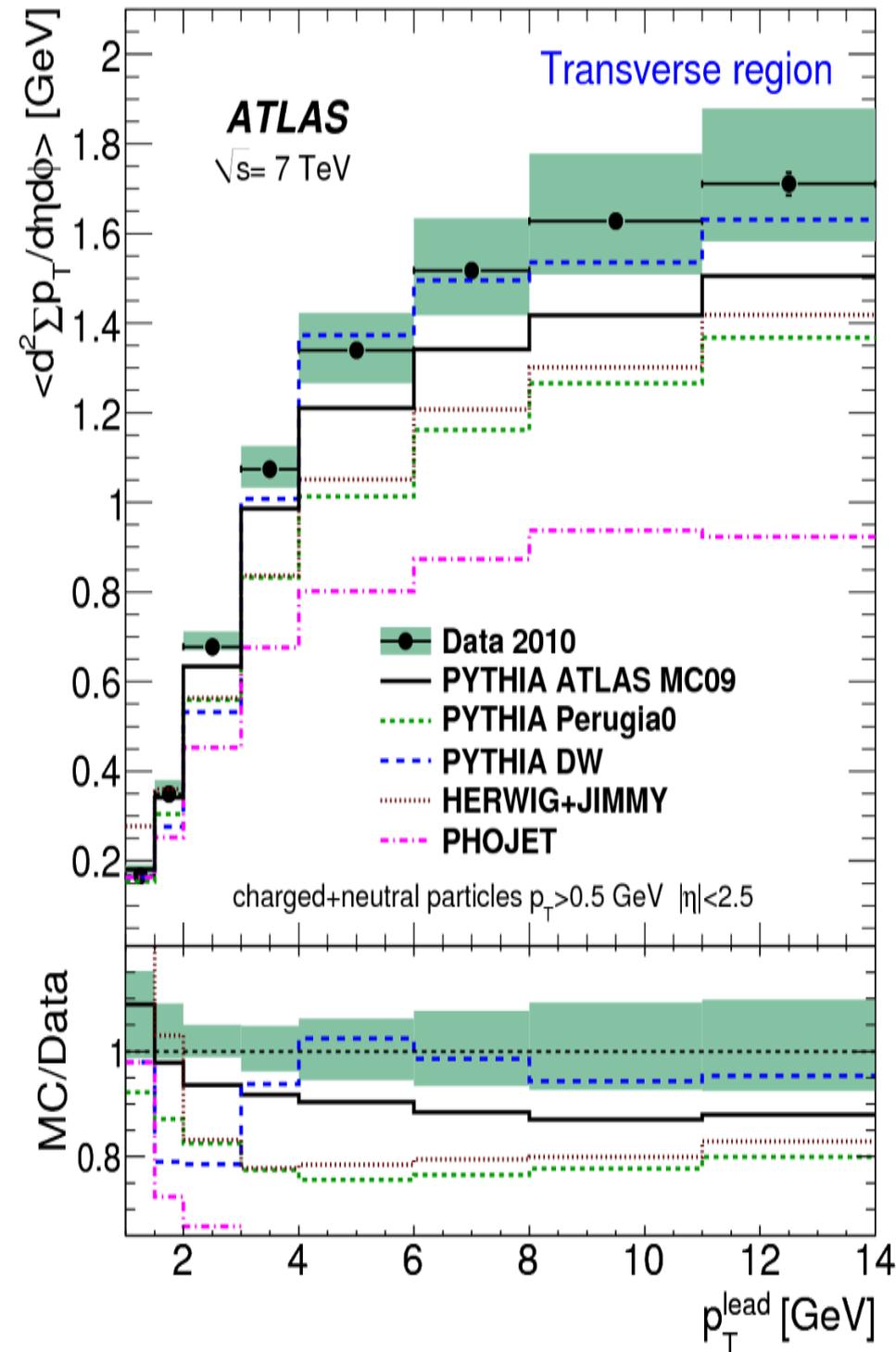
Eur.Phys.J.C 71 (2011) 1795 [1109.5816]

Underlying event

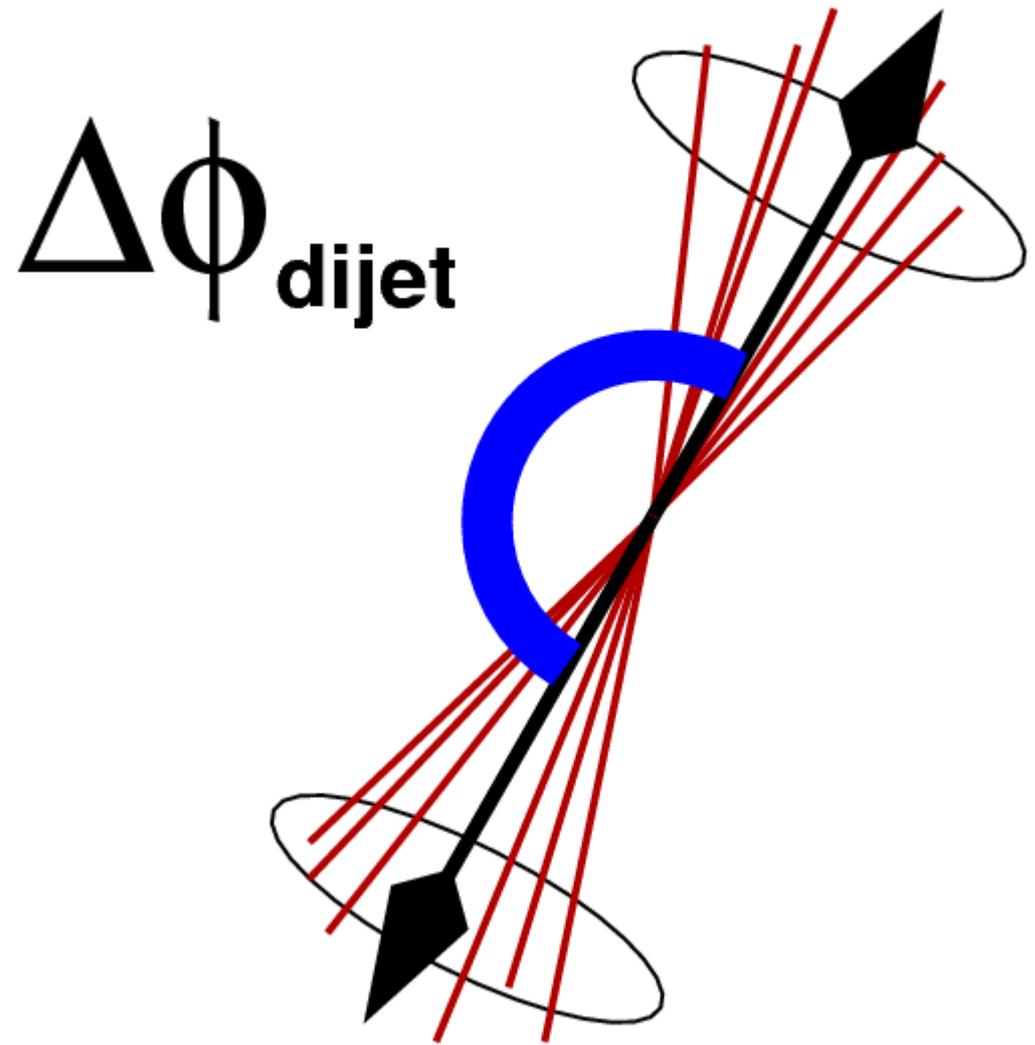


Measuring the properties of the underlying event is crucial to properly describe in the MC jet vetos, low p_T jets, isolations,...

Phys. Rev. D83 (2011) 112001 [1012.0791]



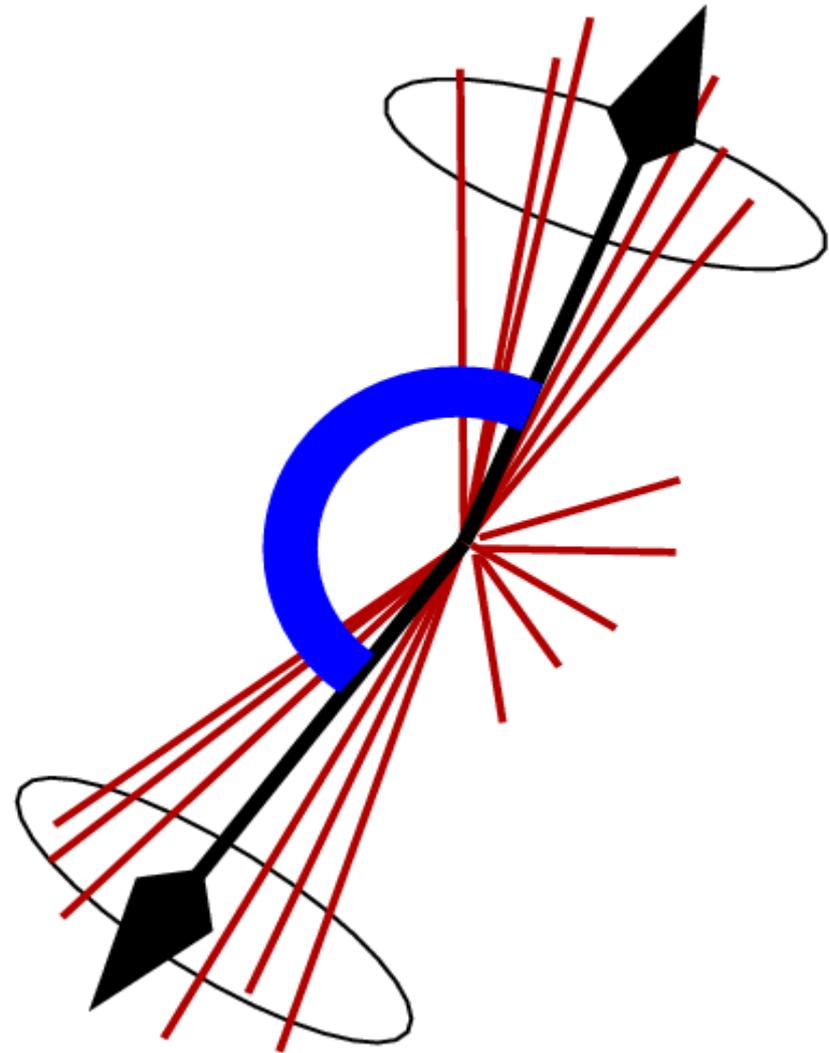
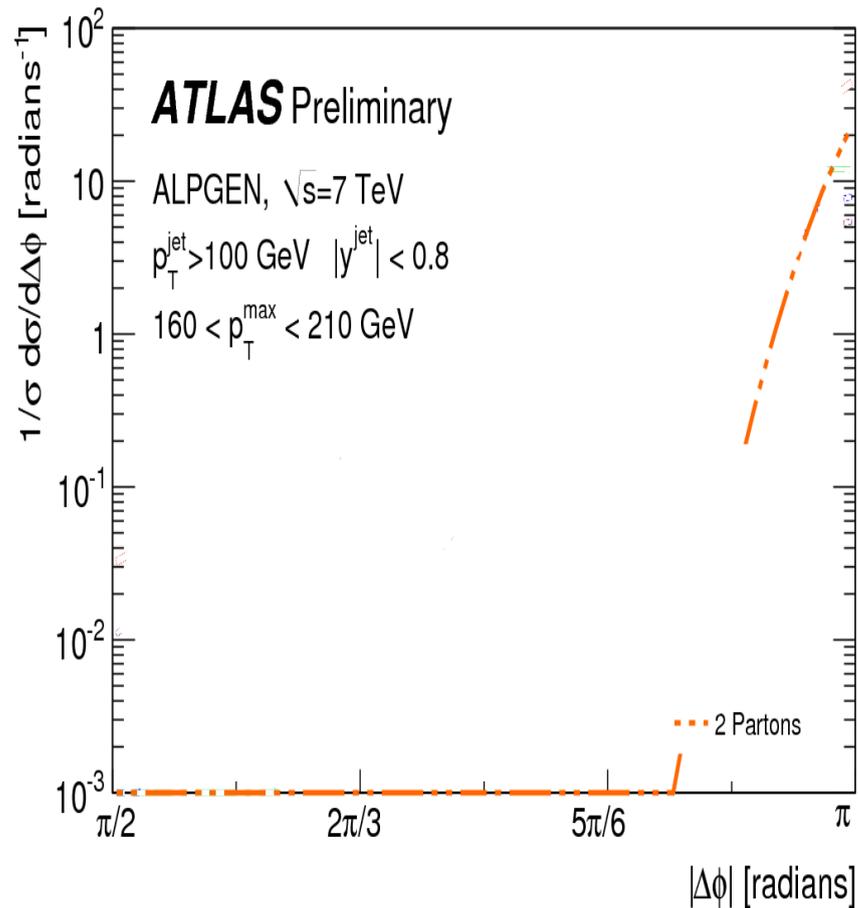
Azimuthal decorrelation



Dijet production at leading order (LO) results in two jets with equal p_{T} and correlated azimuthal angles $\Delta\phi=\pi$.

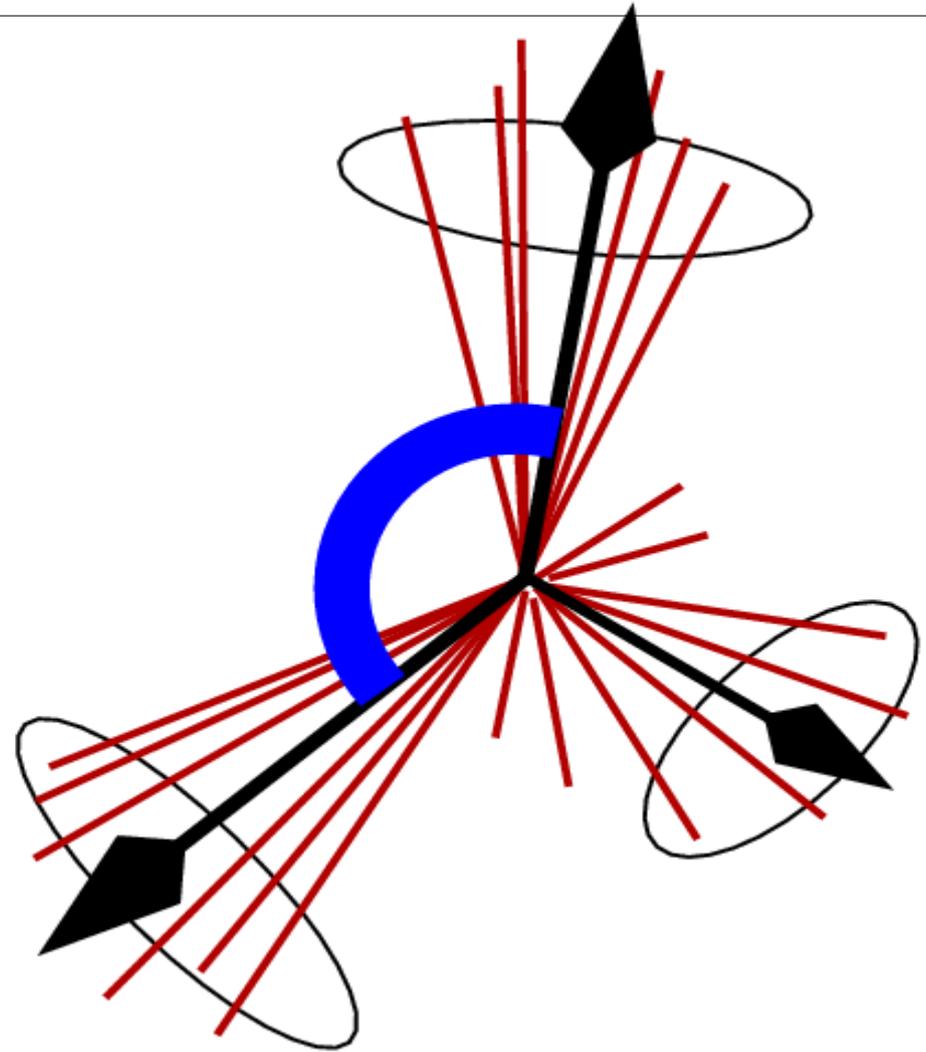
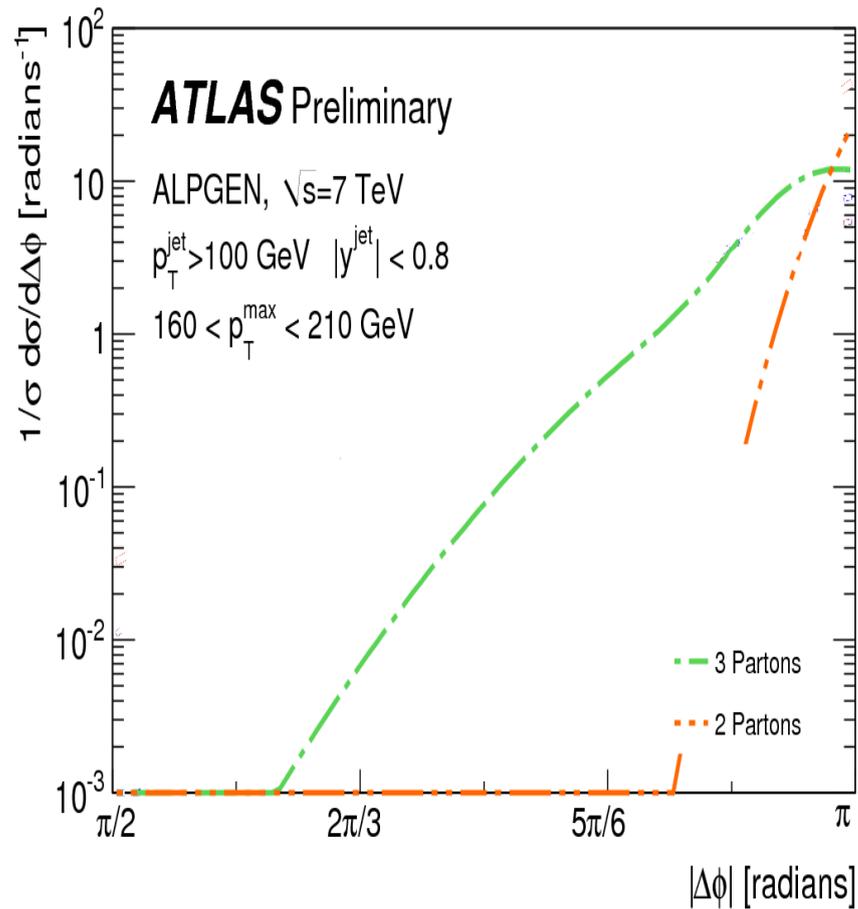
Phys.Rev.Lett. 106 (2011) 172002 [1102.2696]

Azimuthal decorrelation

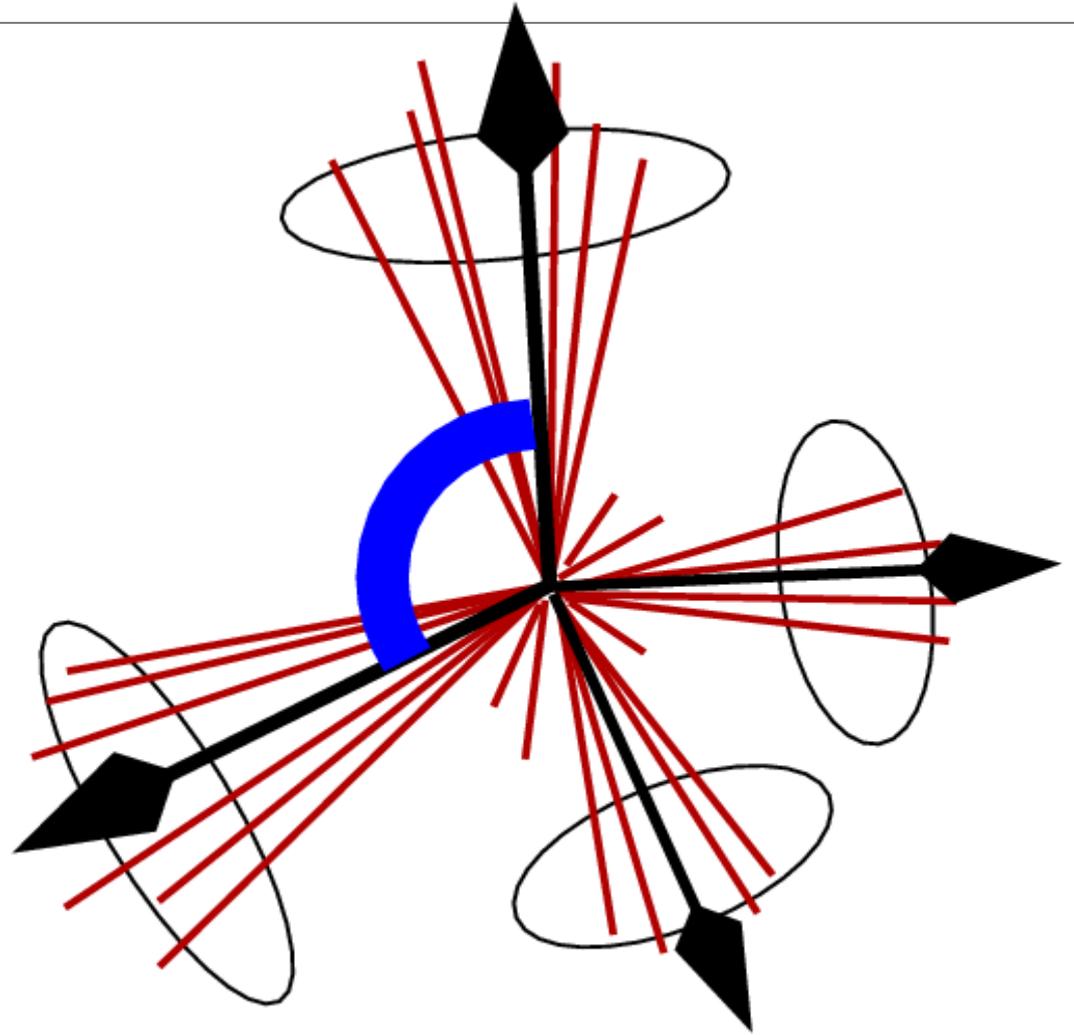
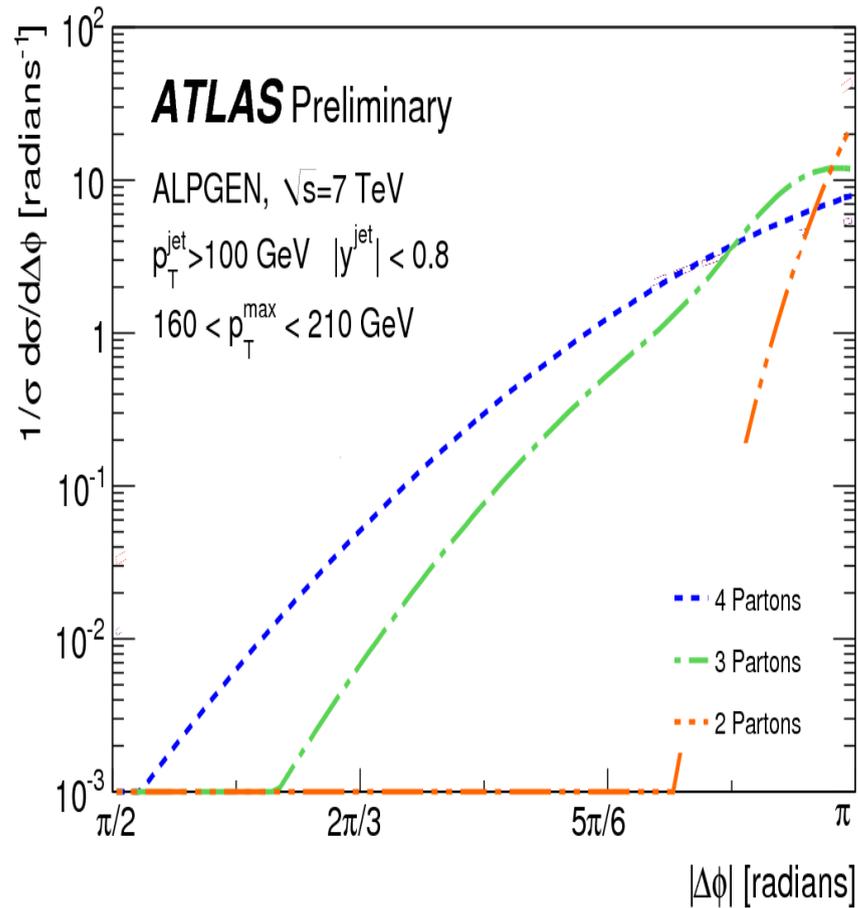


Soft radiation in dijet events starts to produce a decorrelation in $\Delta\phi$.

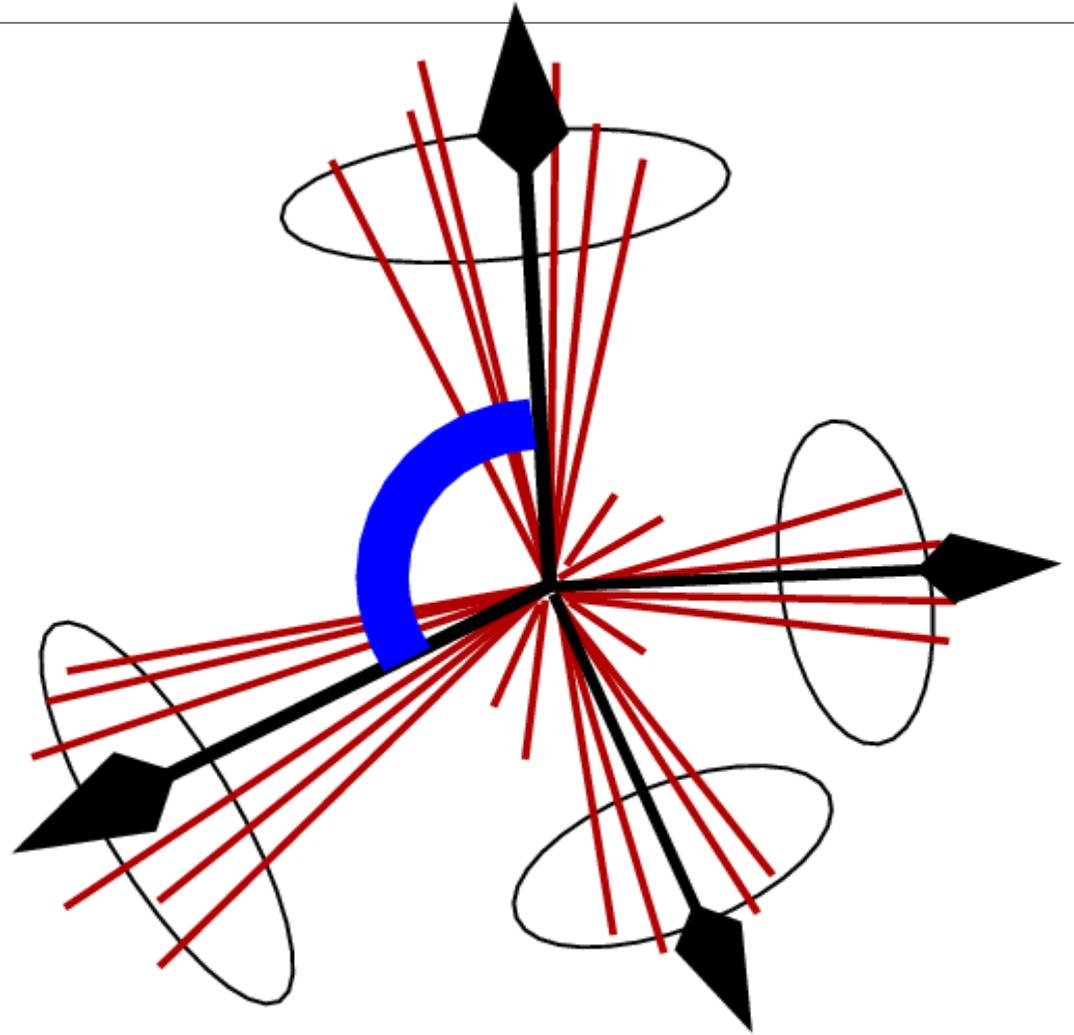
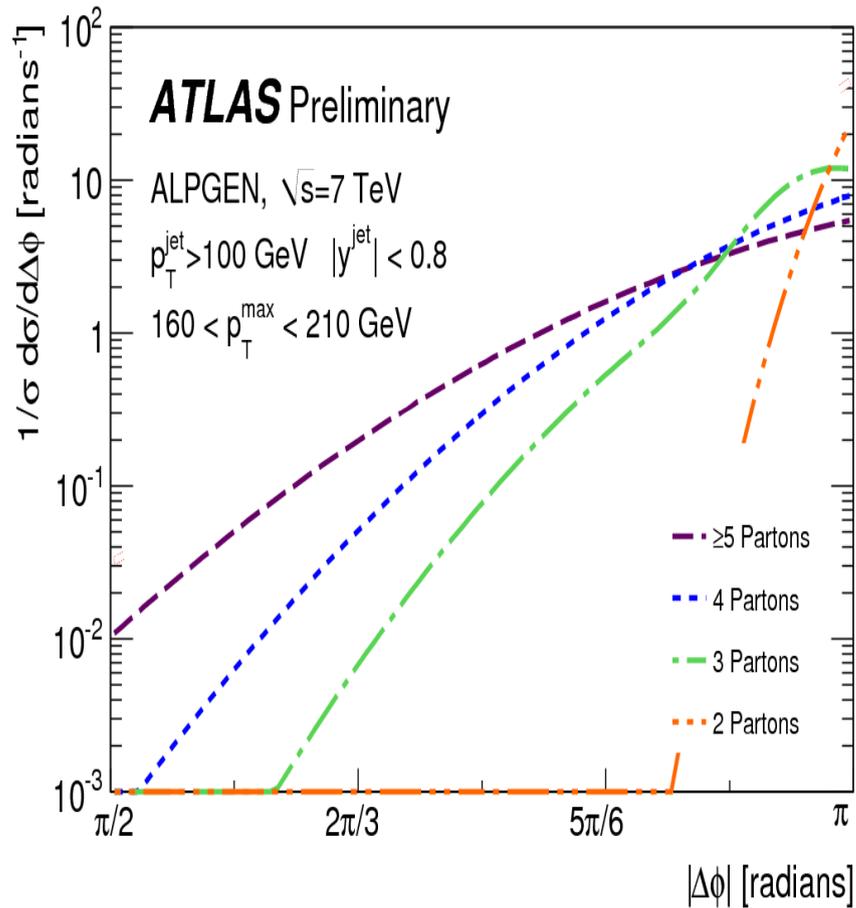
Azimuthal decorrelation



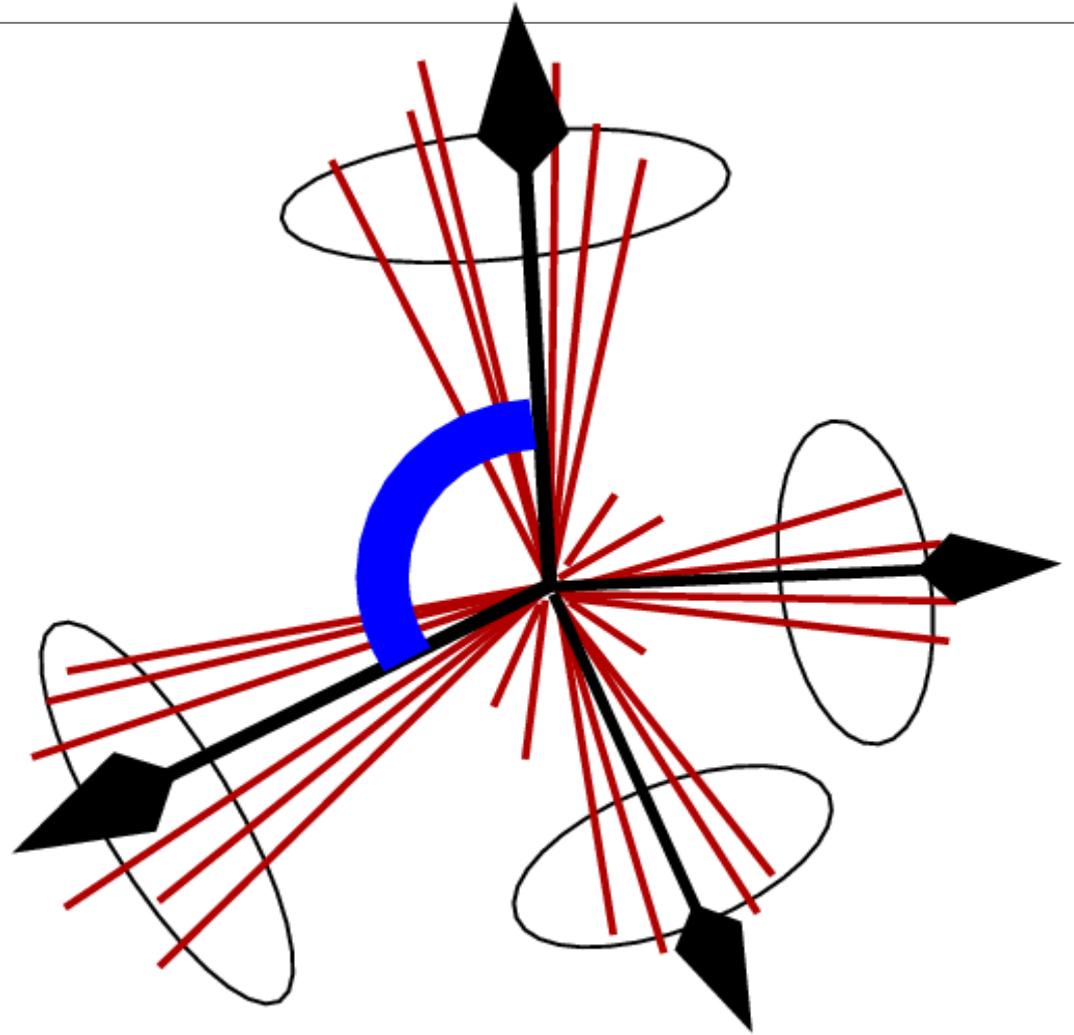
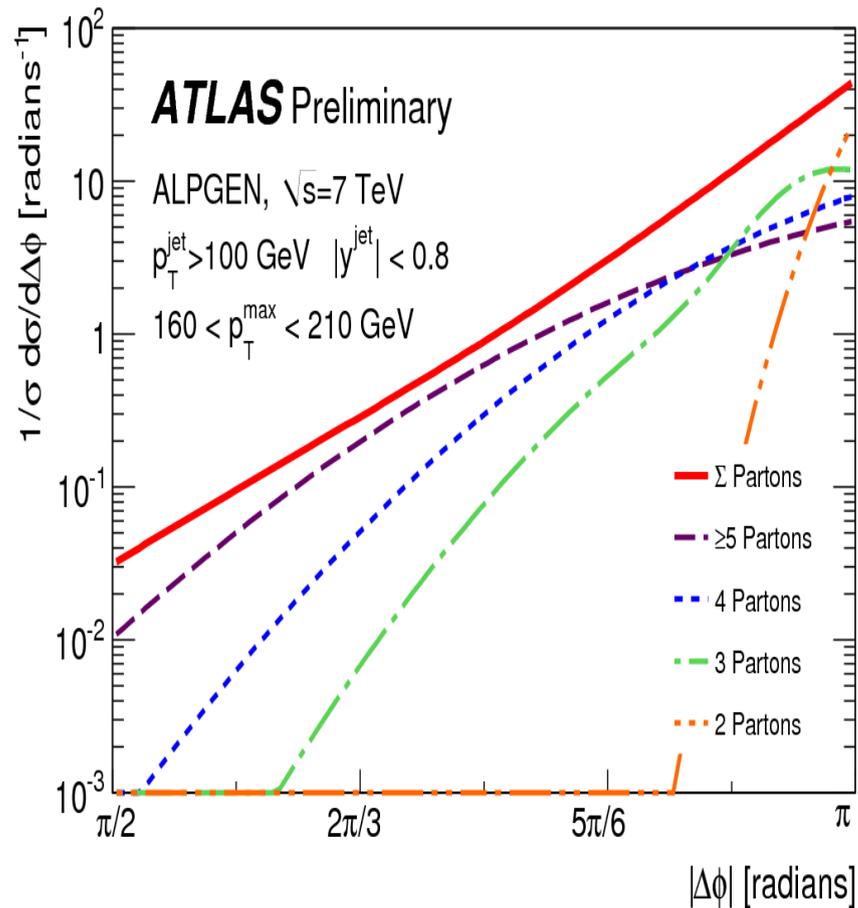
Azimuthal decorrelation



Azimuthal decorrelation



Azimuthal decorrelation



The azimuthal decorrelation $\Delta\phi$ is a test higher-order perturbative QCD (pQCD) calculations without requiring the reconstruction of additional jets and a way to examine the transition between soft and hard QCD processes with a single observable.

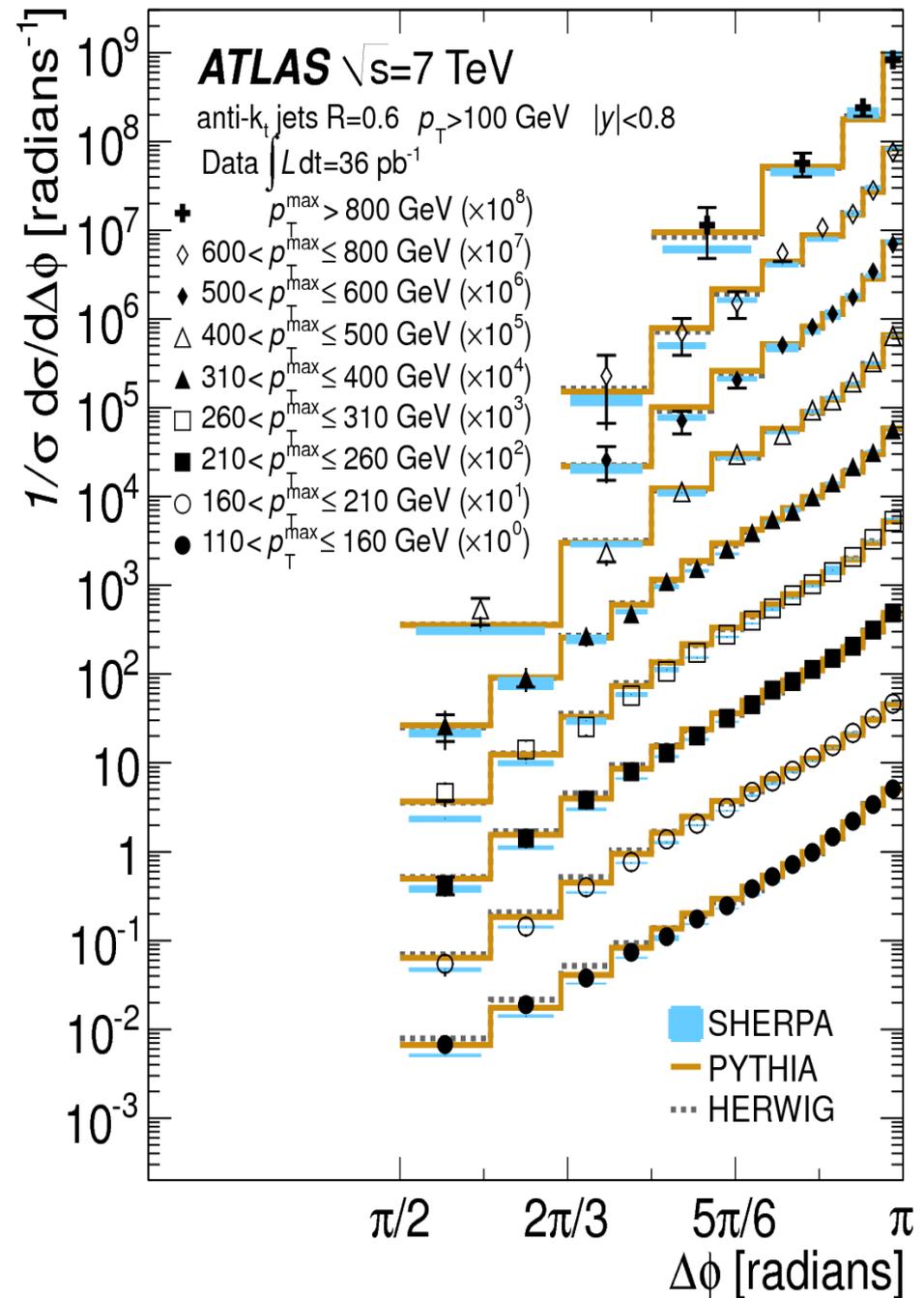
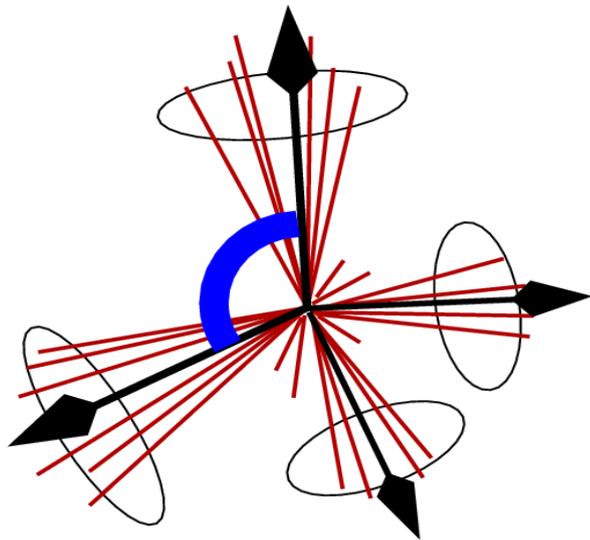
Phys.Rev.Lett. 106 (2011) 172002 [1102.2696]

Azimuthal decorrelation

Monte Carlos show good agreement with data.

Similar agreement for the NLO calculation (NLOJET++)

Phys.Rev.Lett. 106 (2011) 172002 [1102.2696]



Cross Sections

Cross Section: Inclusive Single Jet

Measurements of inclusive cross-sections are important verifications of perturbative QCD and probes of new physics (e.g. quark compositeness, etc.).

Cross Sections:

Inclusive single-jet double-differ. cross-sections as a function of p_T and y

Transverse momentum: $p_T > 20 \text{ GeV}$ $\frac{d^2\sigma}{dP_{T,jet} d|y|}$ Rapidity: $|y| < 4.4$

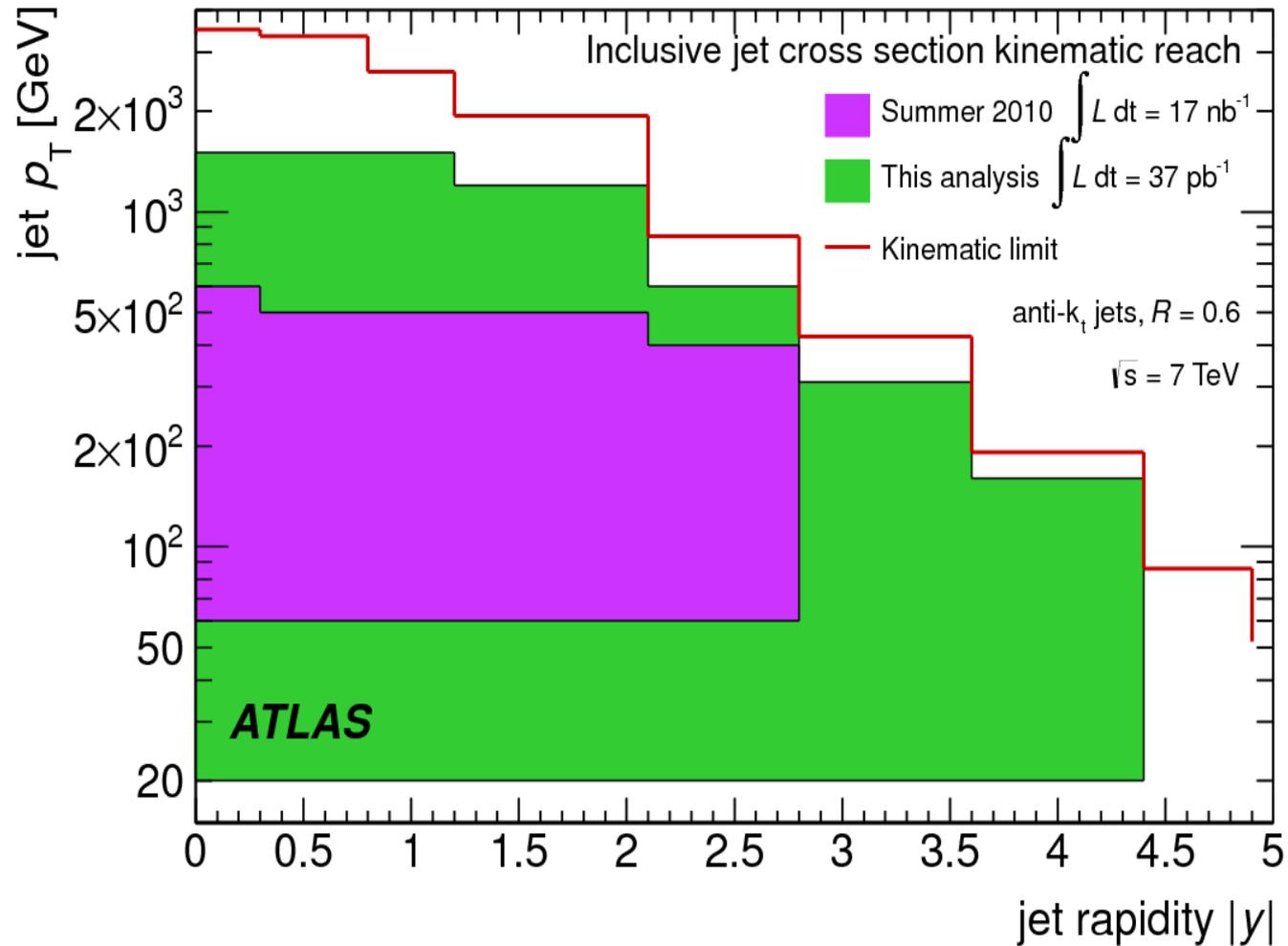
Jet Algorithm: Anti- K_T jets with $R=0.4$ and $R=0.6$

Integrated Luminosity: $\sim 40 \text{ pb}^{-1}$

The cross section is corrected by the detector effects

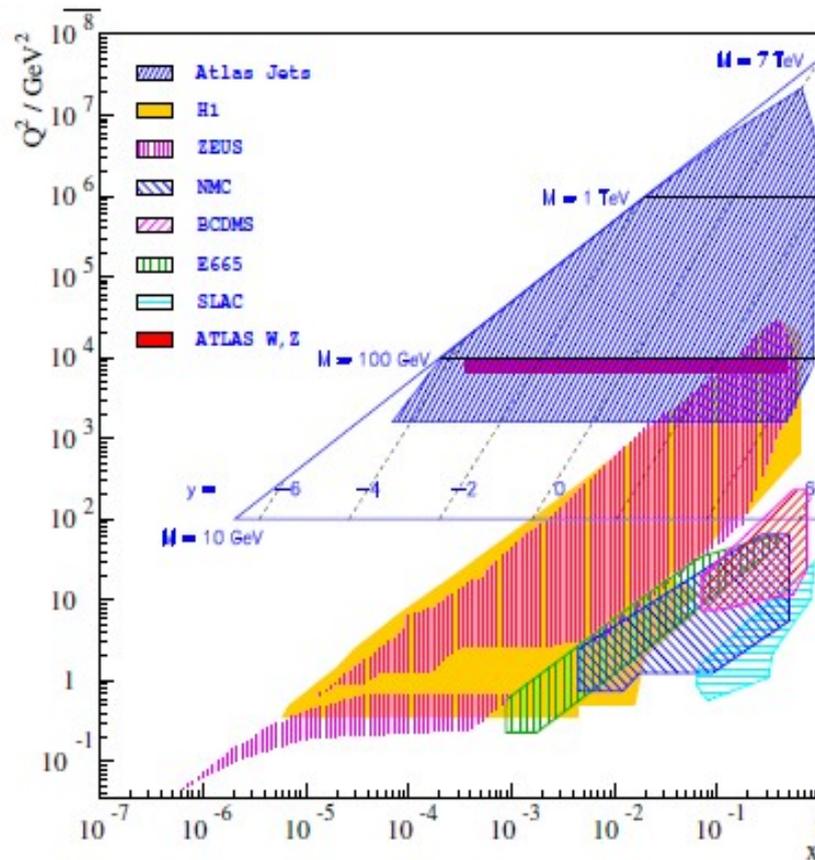
Cross Section: Inclusive Single Jet

Measurement done in a new regime of phase space



Cross Section: Inclusive Single Jet

Measurement done in a new regime of phase space

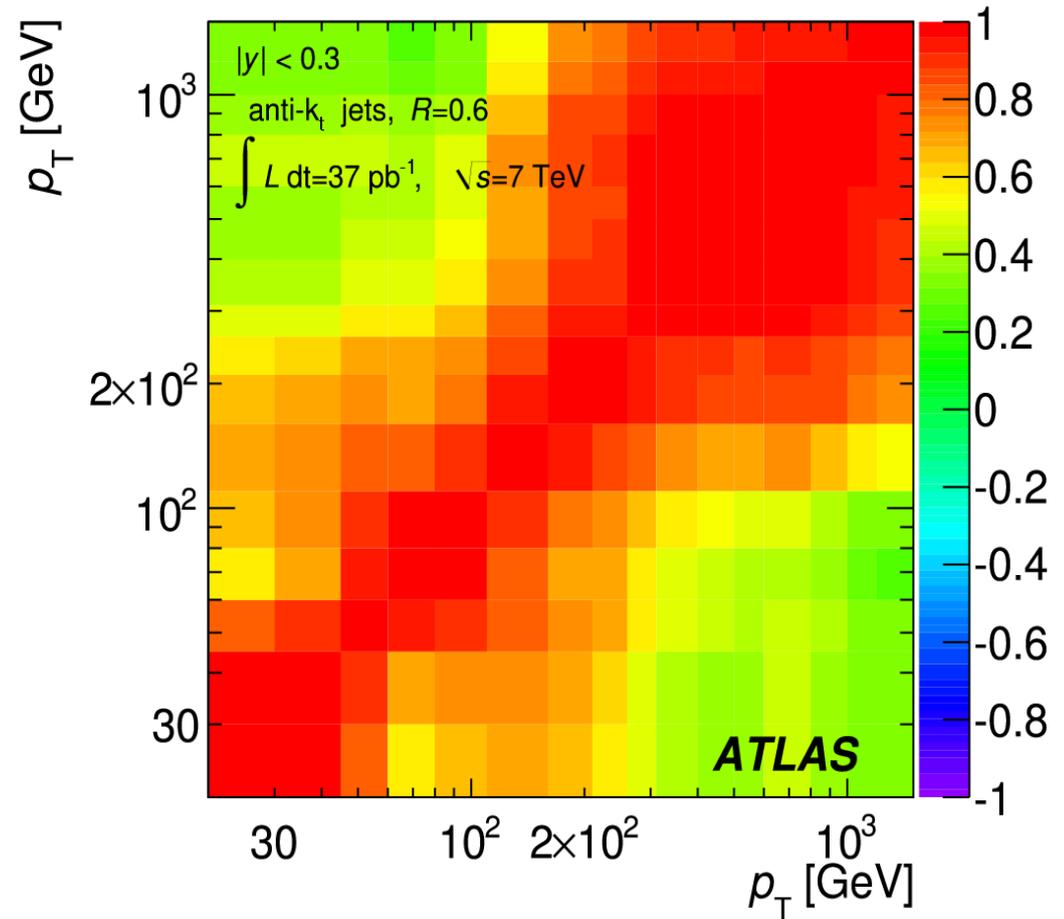
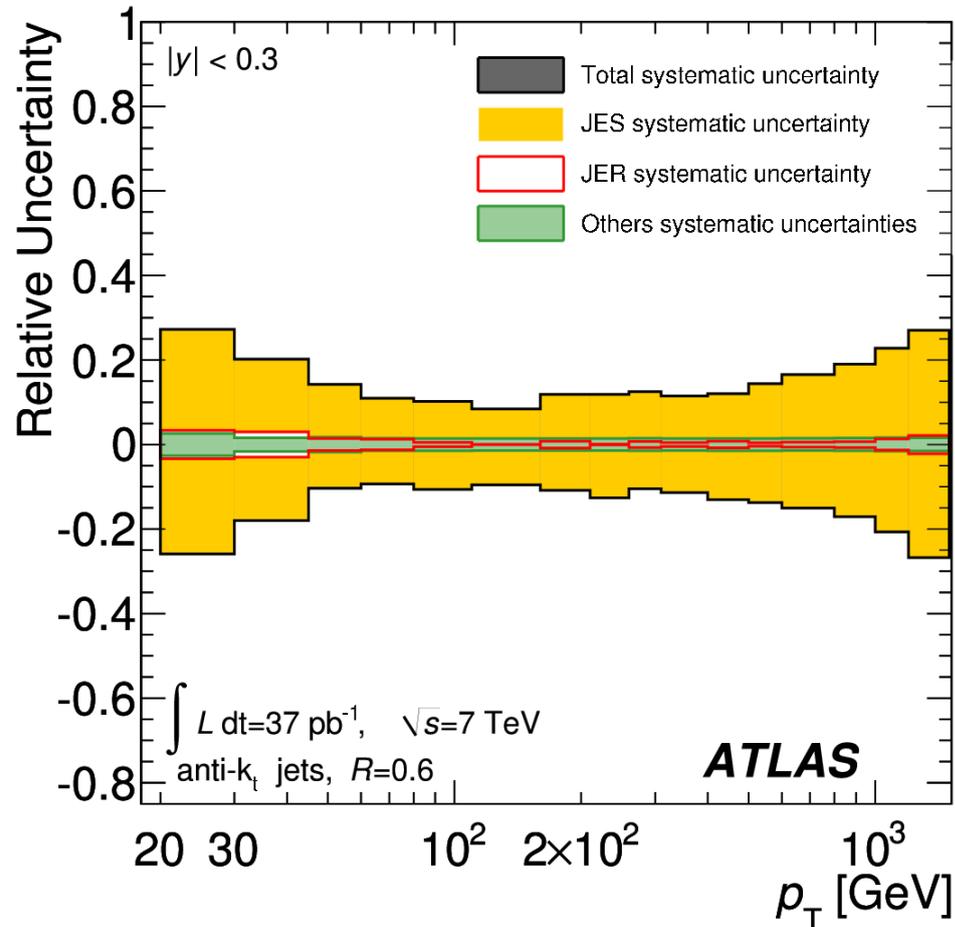


Already used for:
1) determination of α_s
2) constrain the gluon PDF

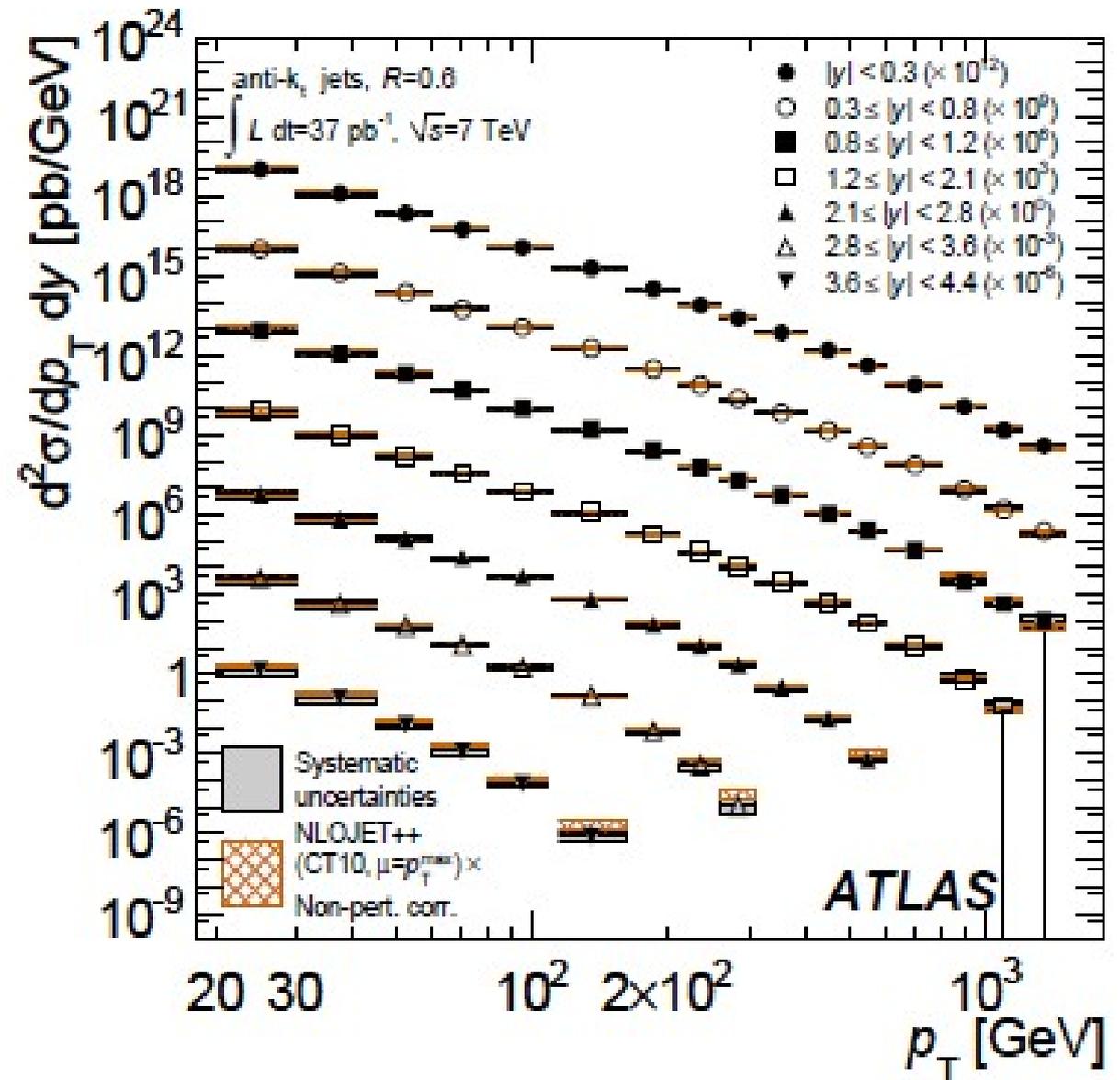
$$7 \cdot 10^{-5} < x < 0.9$$
$$Q^2 > 2 \cdot 10^7 \text{ GeV}^2$$

Cross Section: Inclusive Single Jet

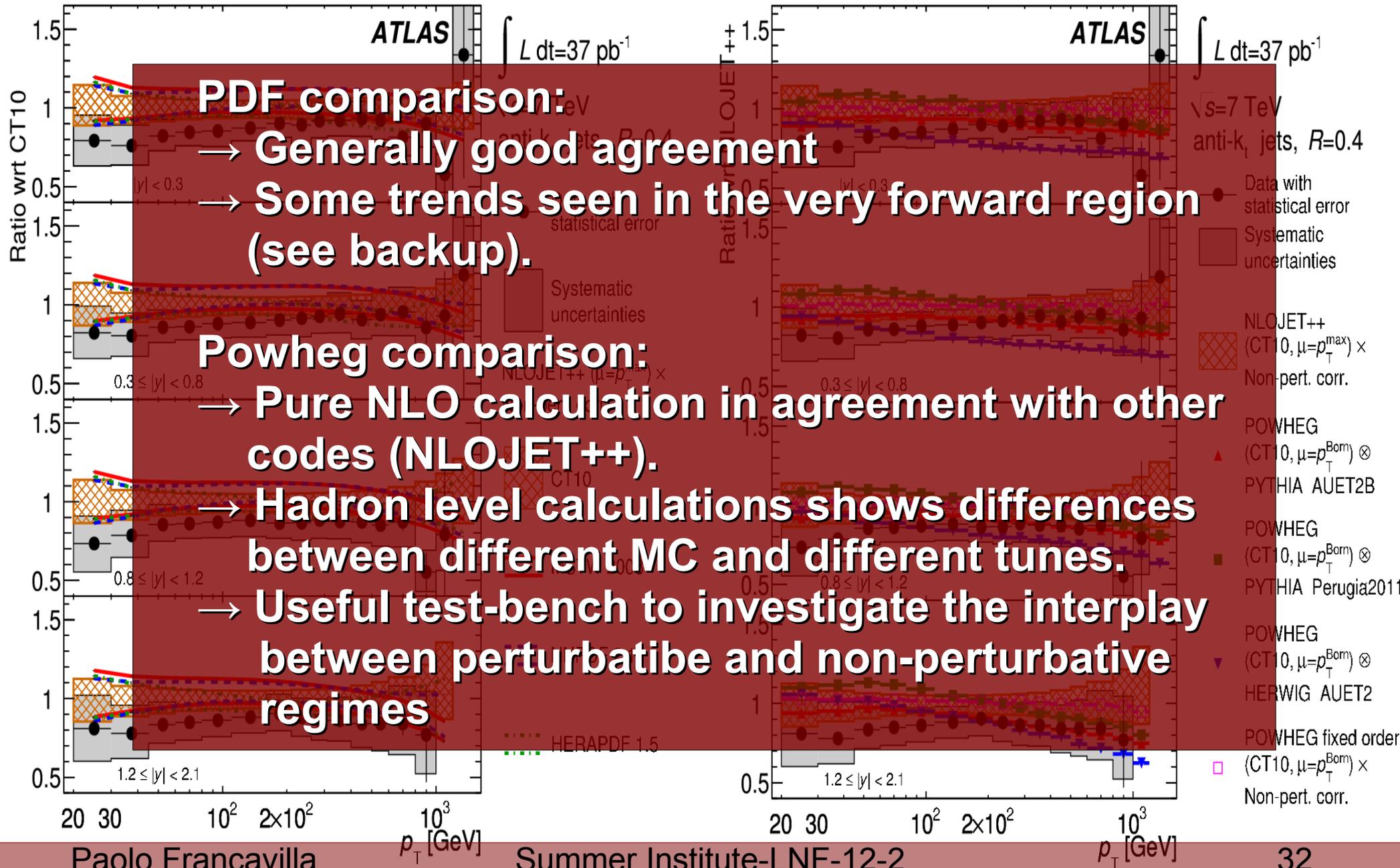
QCD fits gained from the complete study of the systematic correlations: ~ 90 independent components



Cross Section: Inclusive Single Jet



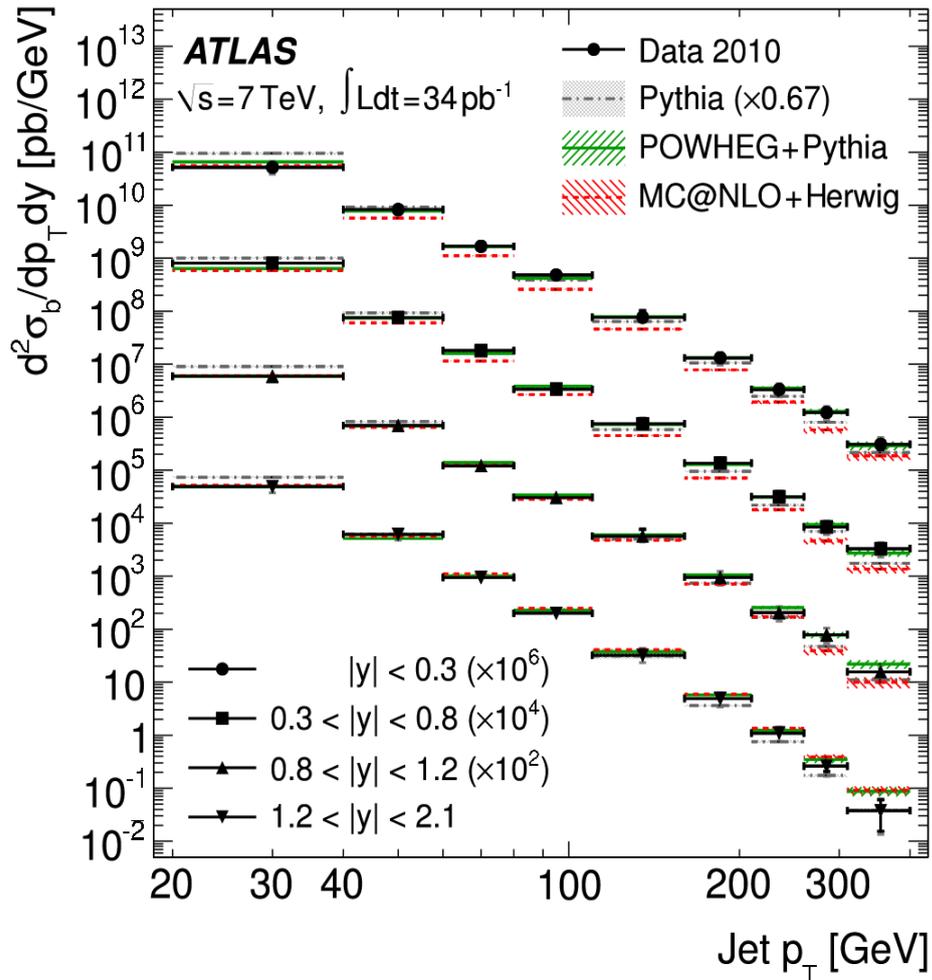
Cross Section: Inclusive Single Jet R=0.4



Cross Section with heavy flavour jets

Bottom

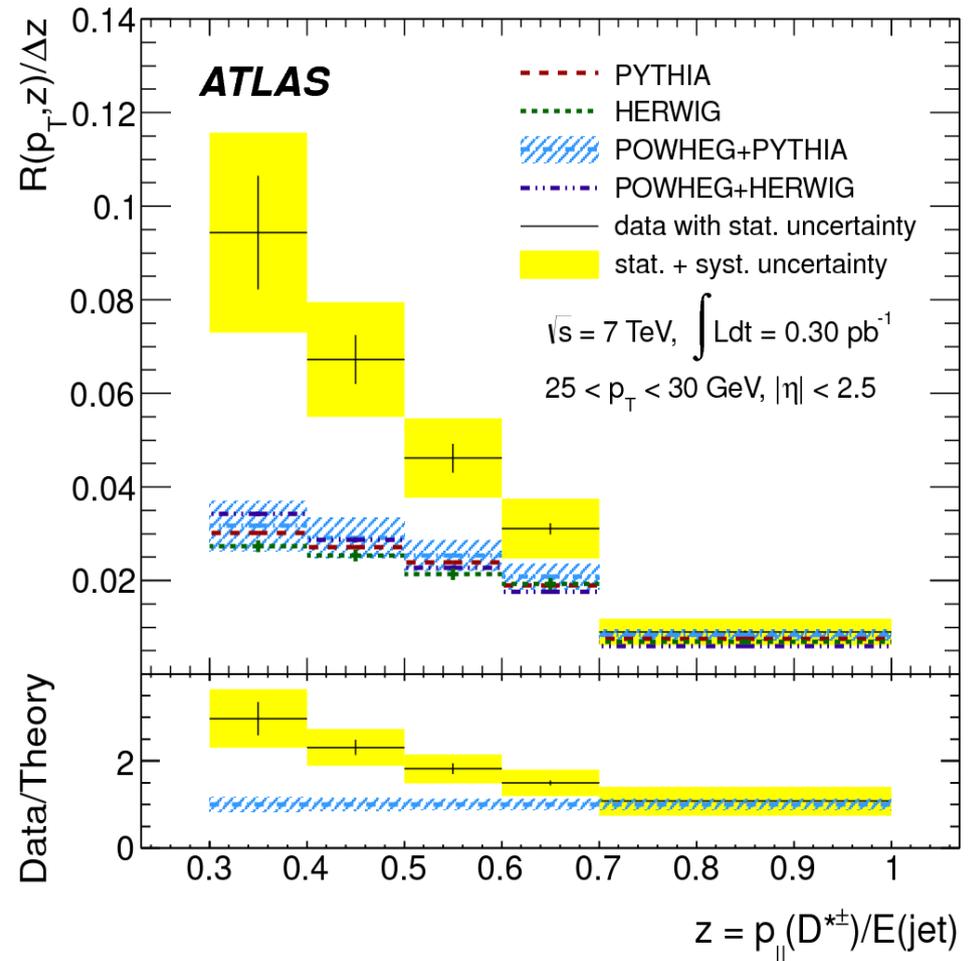
B-jets cross section.
General good agreement.



Eur.Phys.J.C 71 (2011) 1846 [1109.6833]

Charm

Production of jets with associated D^* .
MC fail to describe the D^* relative energy.



Phys. Rev. D85 (2012) 052005 [1112.4432]

Cross Section: Di-Jets

Cross Sections:

Di-jet cross-sections as a function of di-jet mass and angle.

$$d^2\sigma/dM_{1,2} dy^*$$

$M_{1,2}$ is invariant mass of first two leading jets with $P_{T,1} > 30$ GeV and $P_{T,2} > 20$ GeV

$$y^* = \frac{1}{2} \ln \left(\frac{1 + |\cos \theta^*|}{1 - |\cos \theta^*|} \right)$$

Y^* is the rapidity in the two-jet center of mass system.

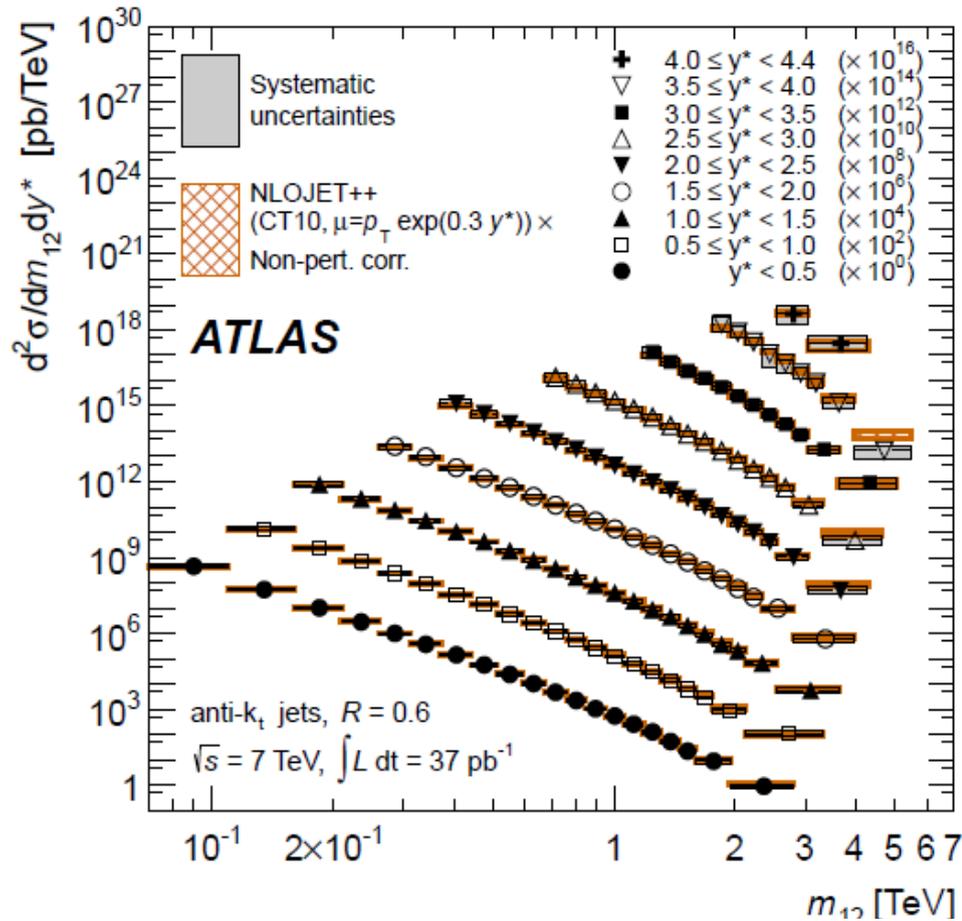
Jet Algorithm: Anti- K_T jets with $R=0.4$ and $R=0.6$

Integrated Luminosity: $\sim 40 \text{ pb}^{-1}$ (2010) and $\sim 5 \text{ fb}^{-1}$ (2011)

The cross section is corrected by the detector effects

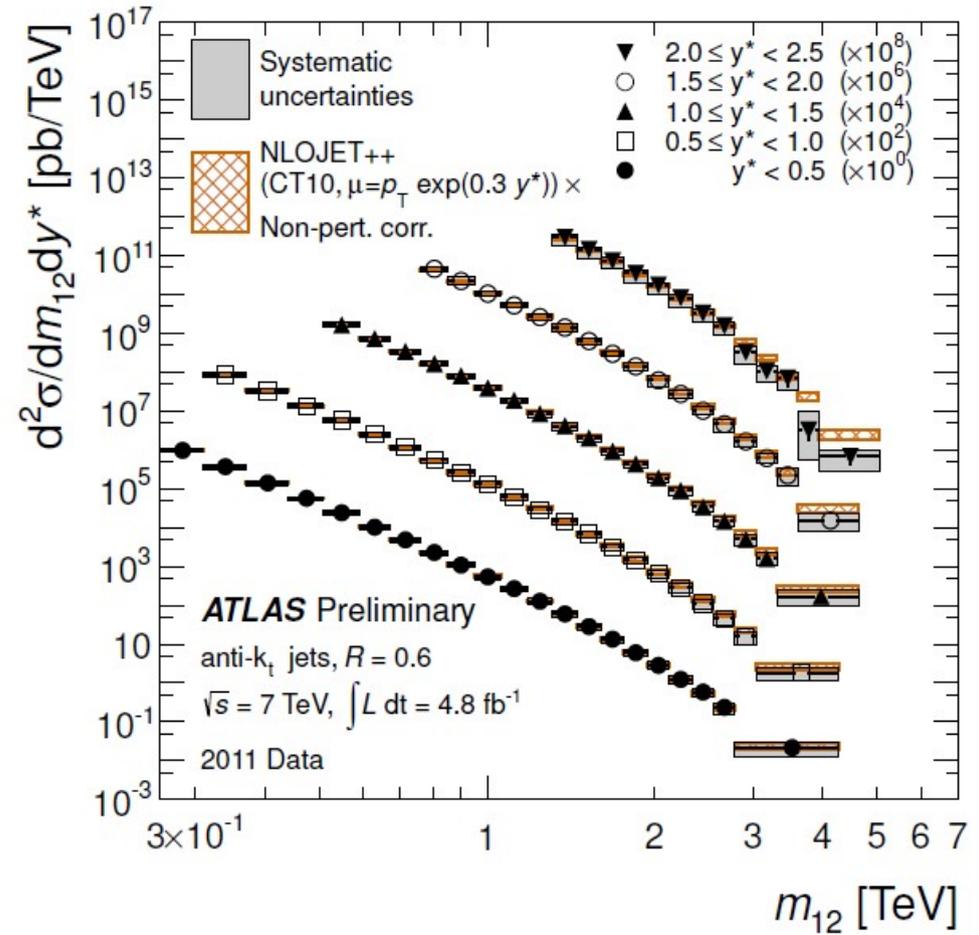
Cross Section: Di-Jets

2010



CERN-PH-EP-2011-192 [1112.6297]

2011



ATLAS-CONF-2012-021

Cross Section: Multi-Jets

A first step toward the measurement of complex QCD final states

- Important as a measurement in itself
(i.e. to extract the strong coupling constant)
- Fundamental to start the controls for the QCD background for searches.

Cross Sections:

Multi-Jet cross section:

Multi Jet rates

p_T spectrum for the 1st, 2nd, 3rd, 4th jet (ordered in p_T)

H_T distribution for different multiplicity

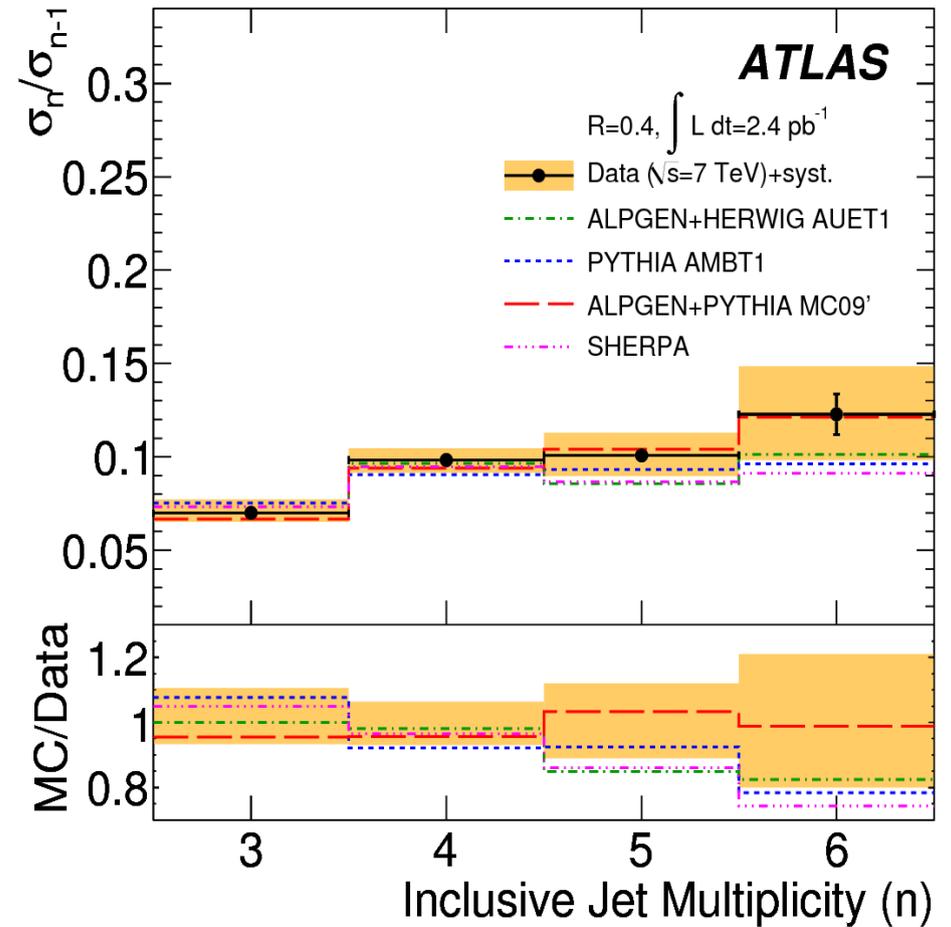
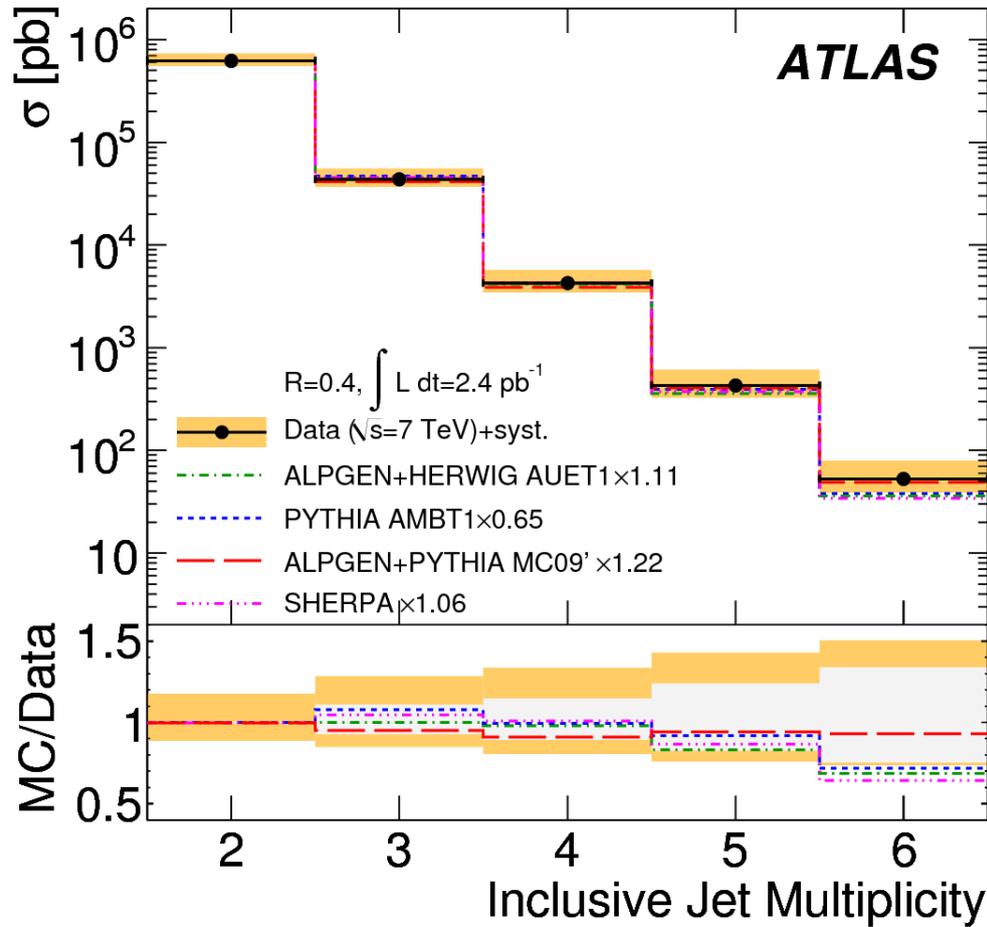
Cuts: leading jets: $p_T > 80$ GeV, subleading jets $p_T > 60$ GeV

Jet Algorithm: Anti- K_T jets with $R=0.4$

Integrated Luminosity: 2.4 pb^{-1}

The cross section is corrected by the detector effects

Cross Section: Multi-Jets



AlpGen+Pythia describes better the data.

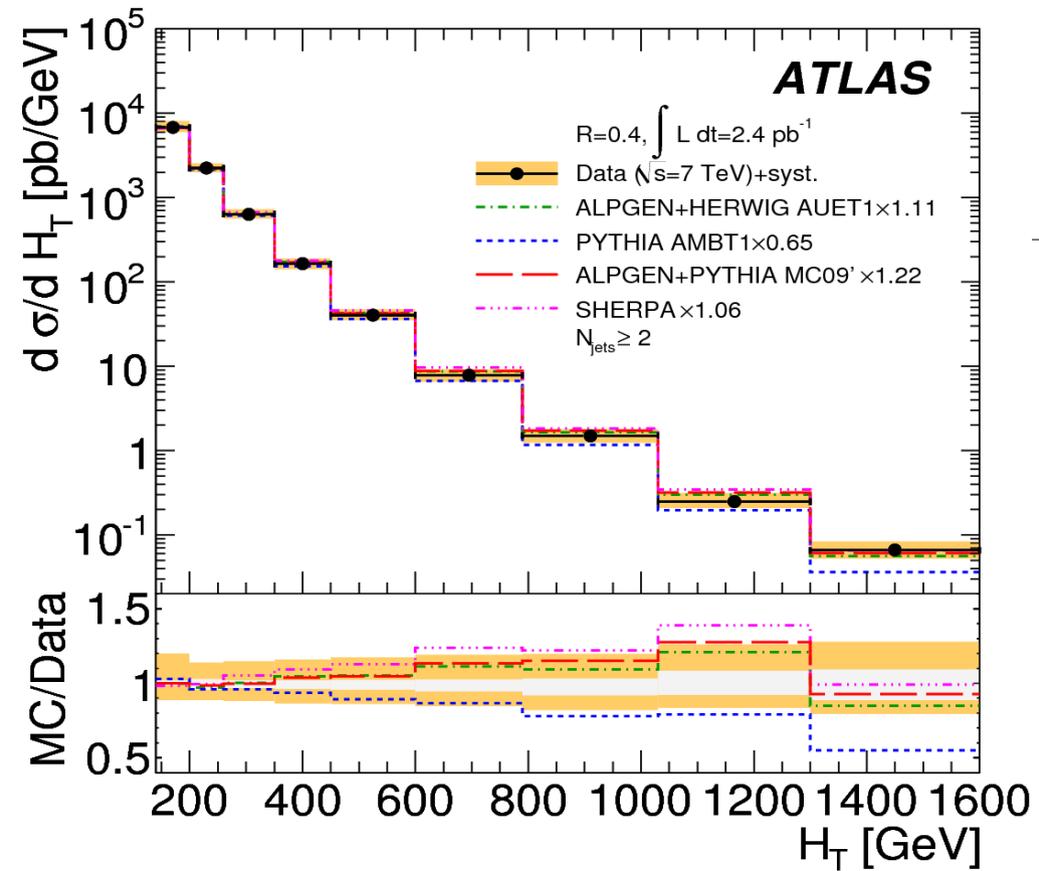
Pythia has a factor 0.65

Cross Section: Multi-Jets

Eur.Phys.J. C71 (2011) 1763 [1107.2092]

$H_T = \Sigma p_T$ of selected jets

Inclusive variable to describe
the events.

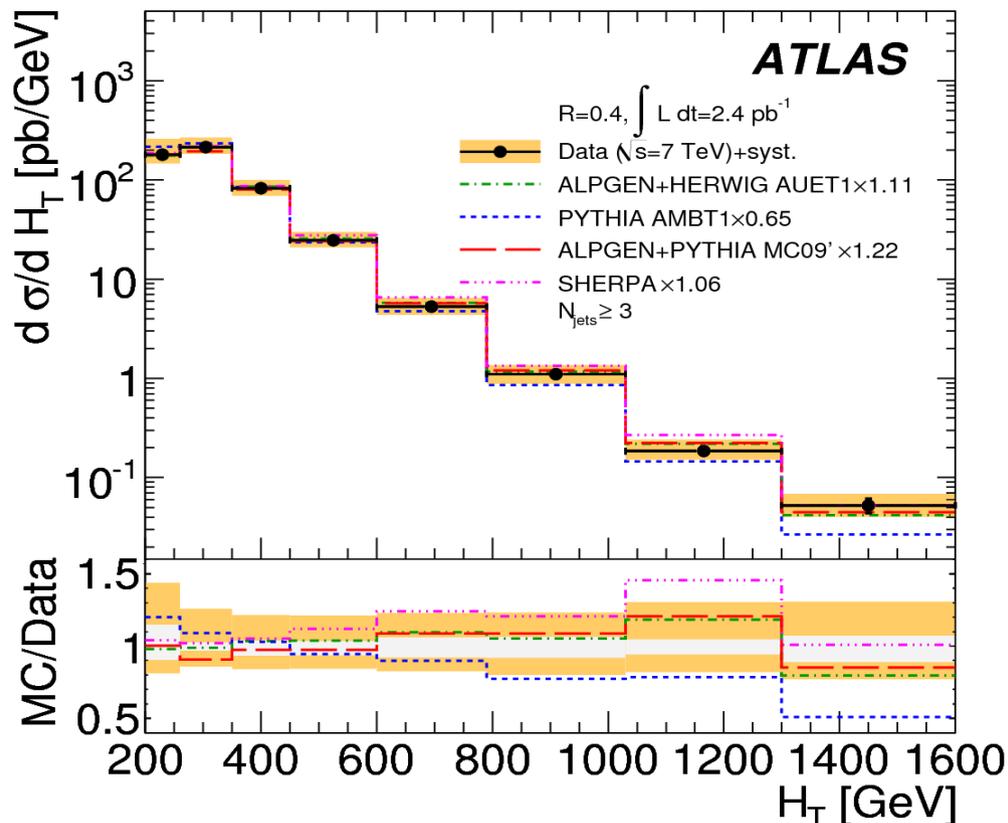
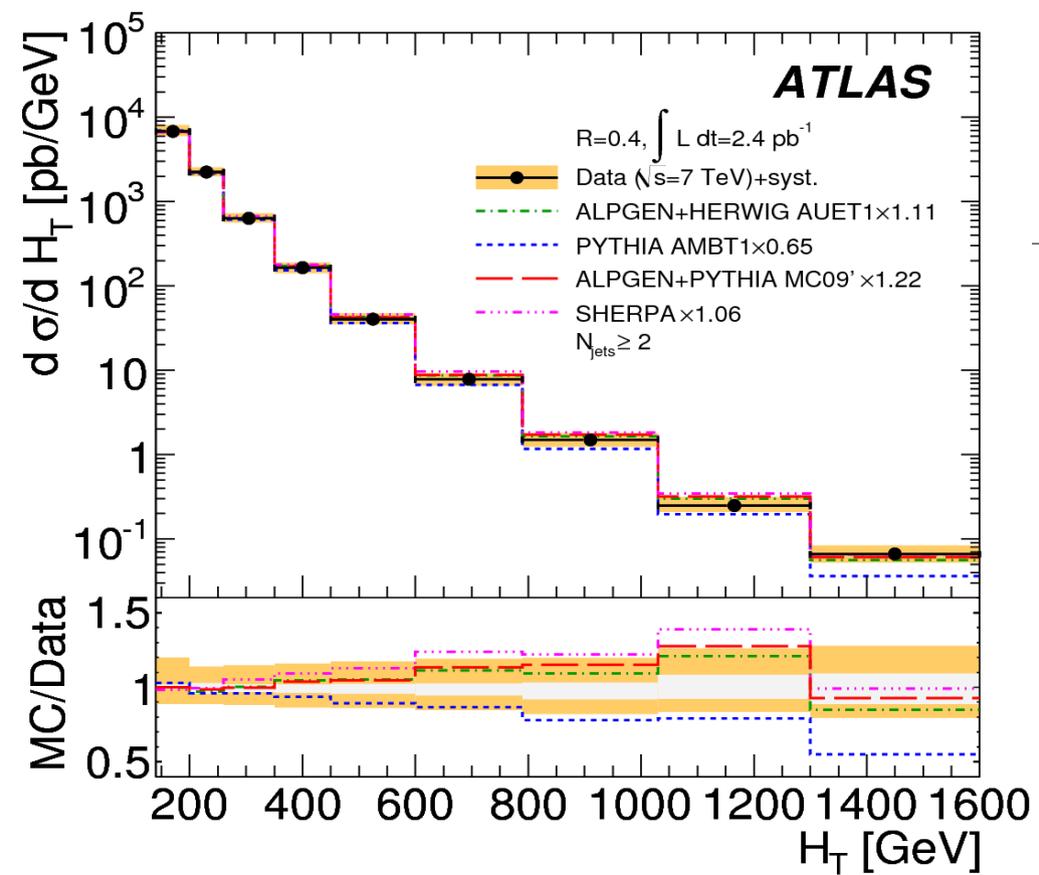


Cross Section: Multi-Jets

Eur.Phys.J. C71 (2011) 1763 [1107.2092]

$H_T = \Sigma p_T$ of selected jets

Inclusive variable to describe the events.

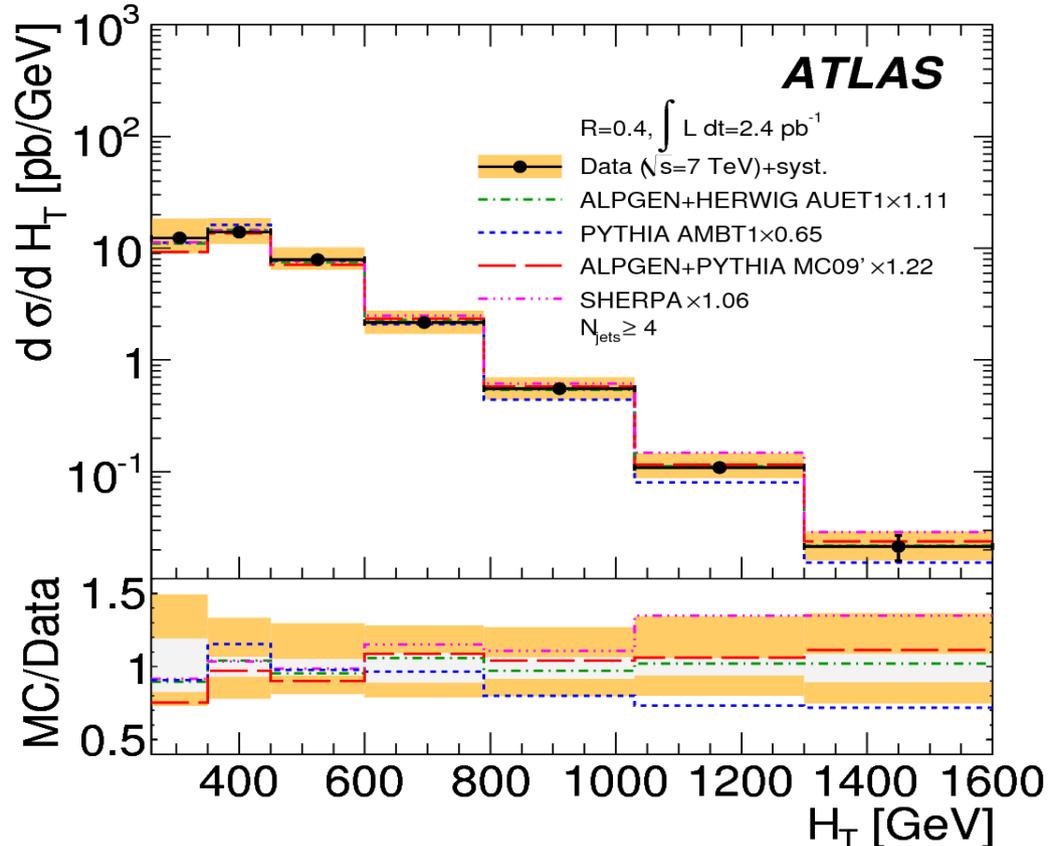
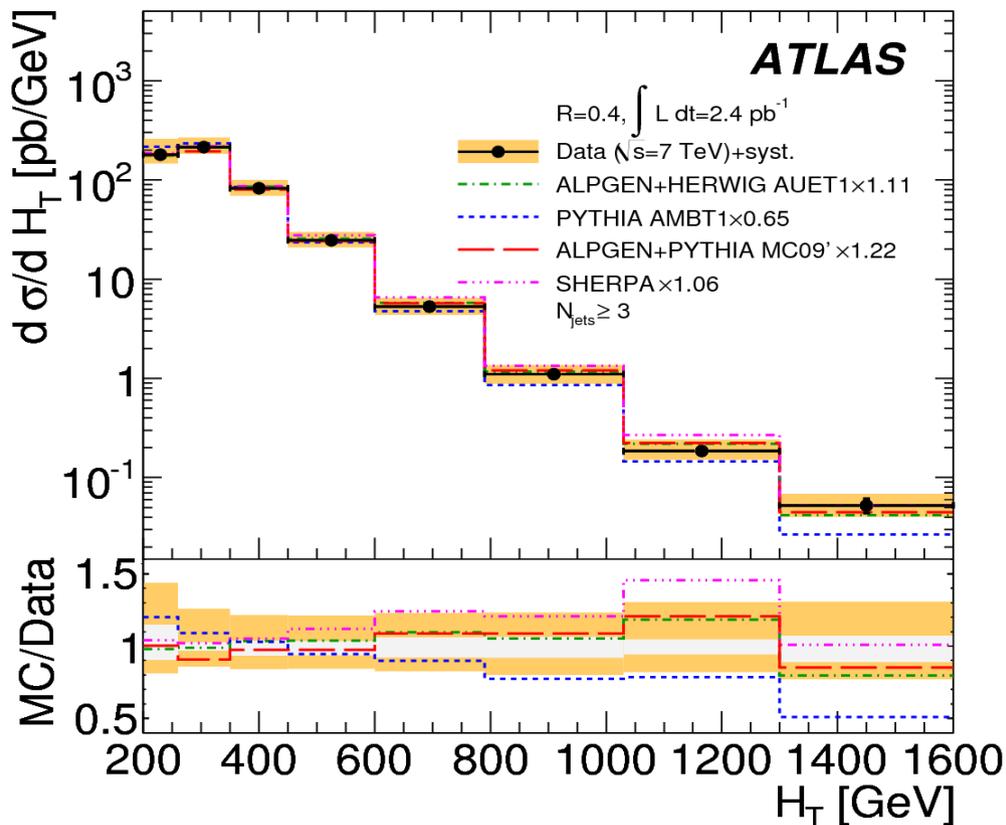
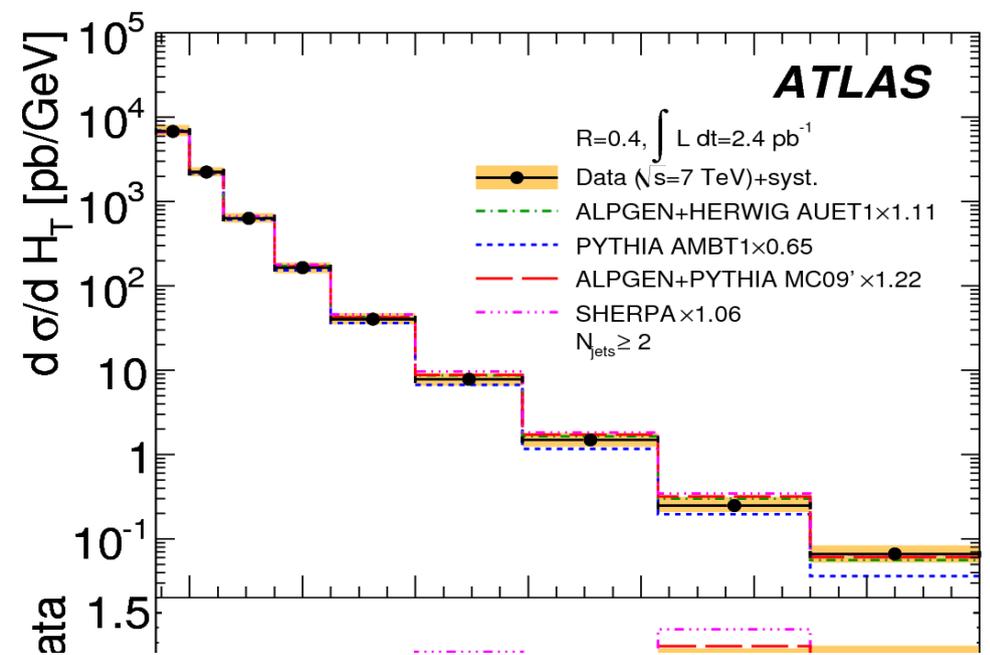


Cross Section: Multi-Jets

Eur.Phys.J. C71 (2011) 1763 [1107.2092]

$H_T = \Sigma p_T$ of selected jets

Inclusive variable to describe the events.



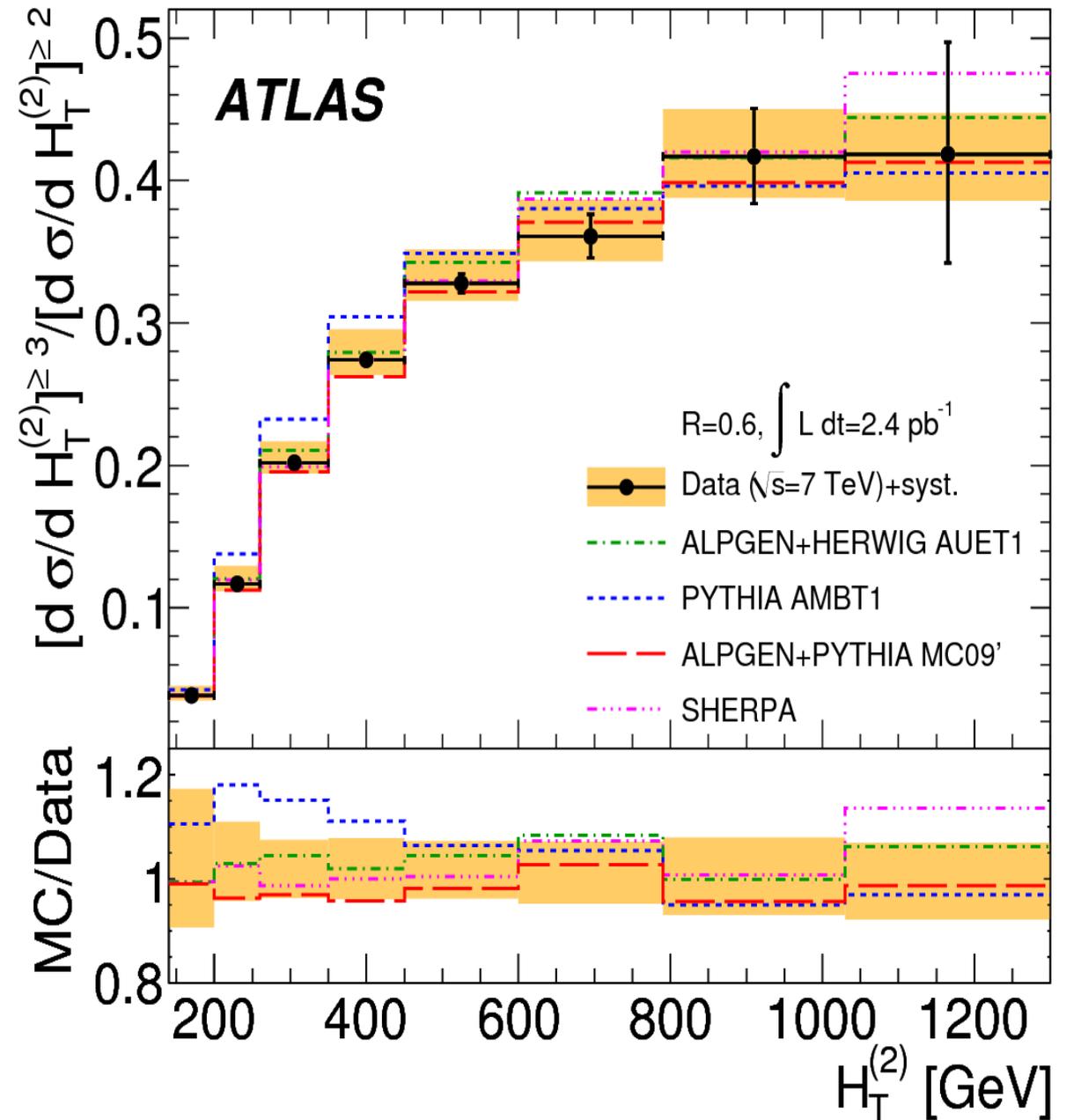
Cross Section: Multi-Jets

By making the ratio,
part of the systematics
cancel out.

Useful as an input for
the strong coupling
constant evaluation

(rough indication
of the scaling violation).

Eur.Phys.J. C71 (2011) 1763 [1107.2092]



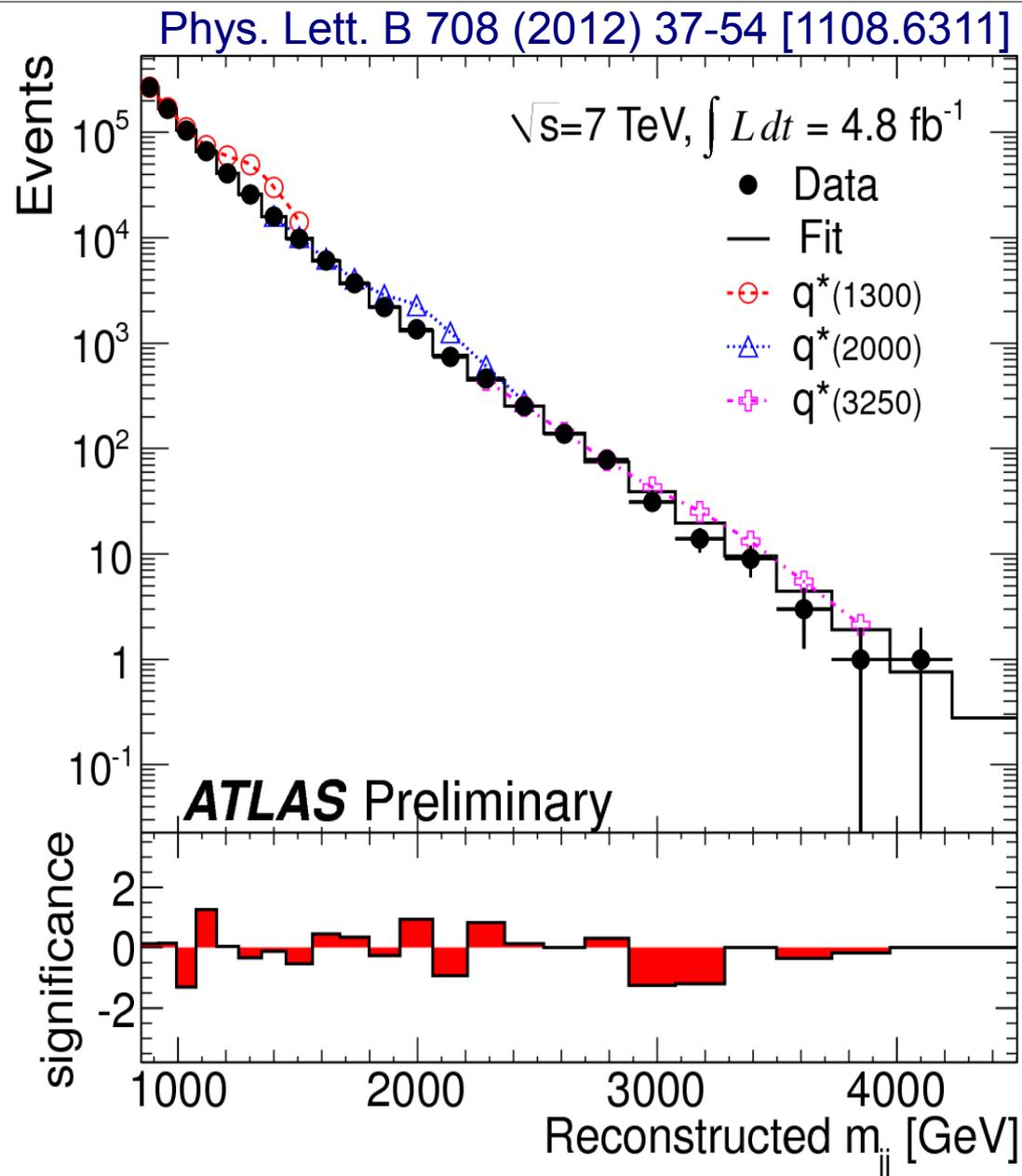
Jets in searches

Exclusions: Di-jet Mass

By using the di-jet measurements, limits on new physics can be studied.

Search for bumps in the di-jet spectrum.

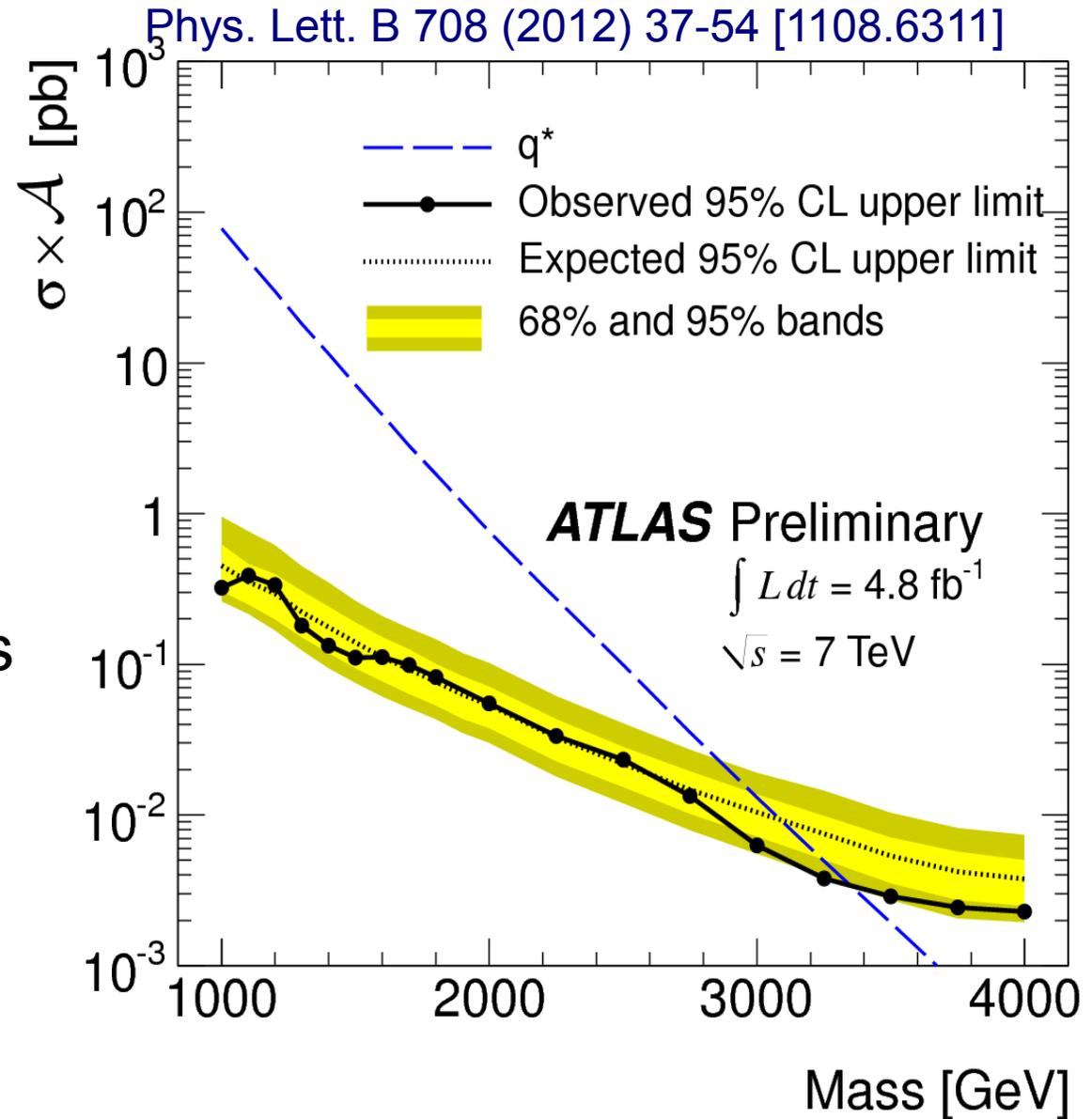
The fluctuations are not statistically significant.



Exclusions: Di-jet Mass

Assuming a narrow di-jet resonance, the di-jet mass measurement can be used to exclude a certain cross section for the production of a resonance at a certain mass

This result can be used to exclude regions in the plane masses/couplings for effective theories.



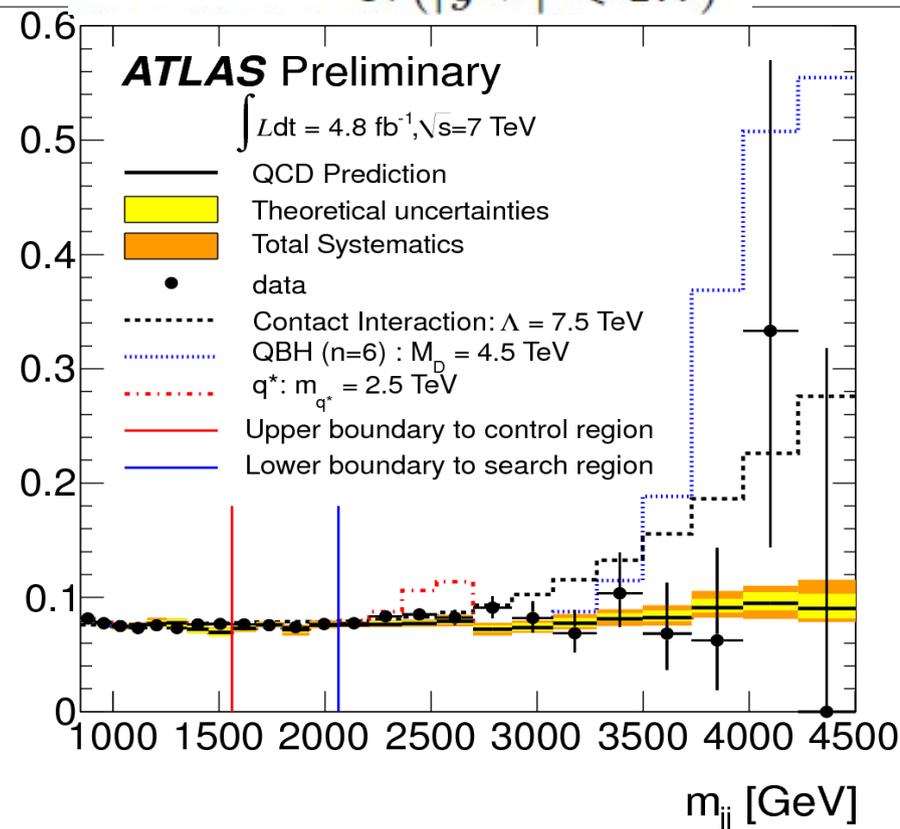
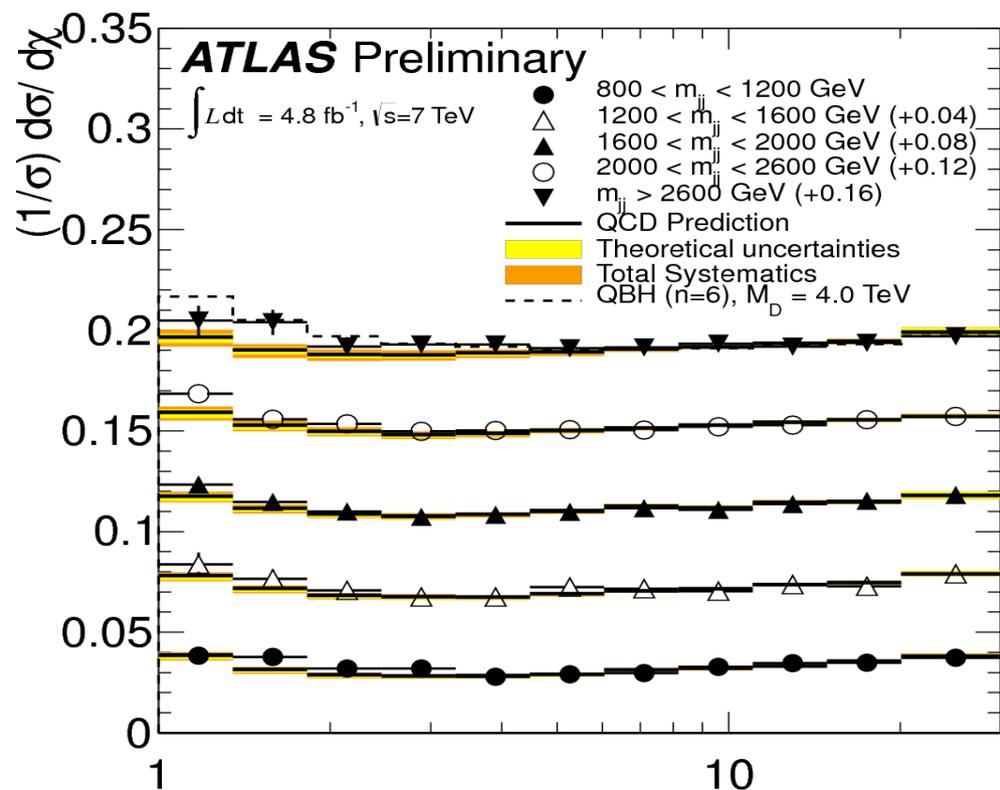
Exclusions: Di-jet Angles

Even the angular distribution for jets F_χ have an important role to constrain models of new physics:

i.e. contact interaction

$$\chi = \exp(|y_1 - y_2|)$$

$$F_\chi(m_{jj}) = \frac{N(|y^*| < 0.6)}{N(|y^*| < 1.7)}$$

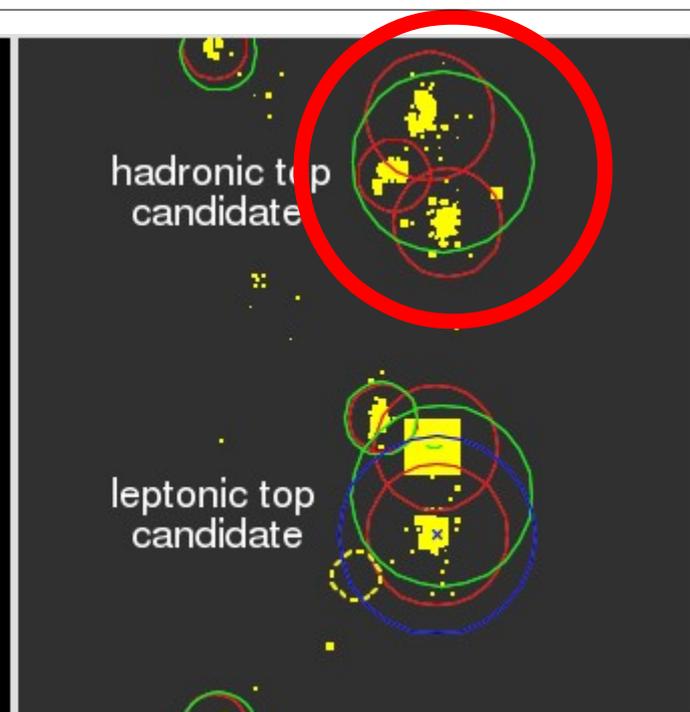
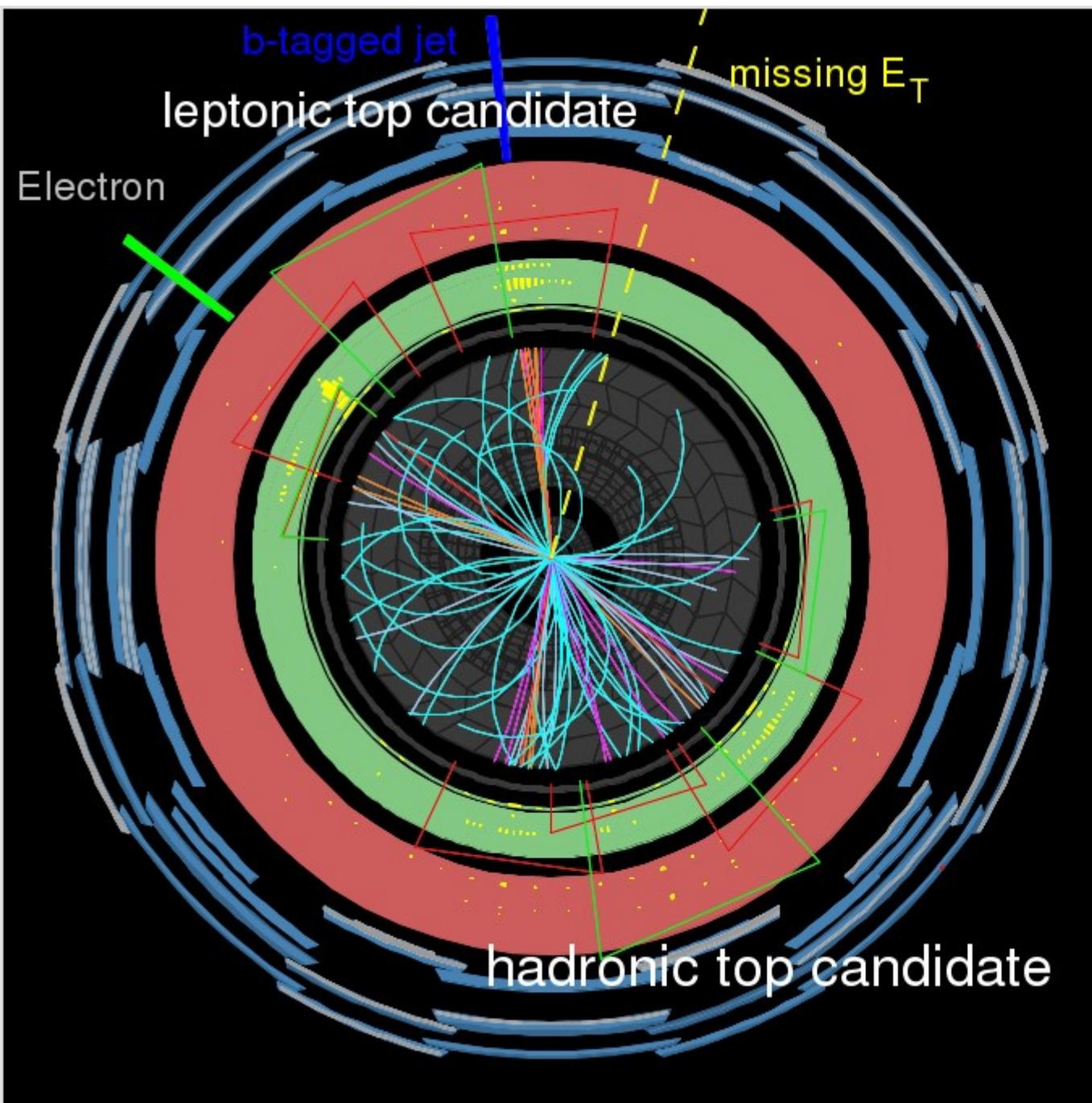


Model, and Analysis Strategy	95% C.L. Limits (TeV)	
	Expected	Observed
Excited quark, mass of q^*		
Resonance in m_{jj}	3.09	3.35
Resonance in $F_\chi(m_{jj})$	2.97	2.58
Colour octet scalar, mass of s_8		
Resonance in m_{jj}	1.94	1.94
Quantum Black Hole for $n = 6, M_D$		
$F_\chi(m_{jj})$	4.14	4.11
11-bin χ , $m_{jj} > 2.6 \text{ TeV}$	4.23	3.96
Contact interaction, Λ , destructive interference		
$F_\chi(m_{jj})$	8.2	7.6
11-bin χ , $m_{jj} > 2.6 \text{ TeV}$	8.7	7.8

Phys. Lett. B 708 (2012) 37-54 [1108.6311] χ

Jets substructure

Not just a crowded event



 **ATLAS**
EXPERIMENT

Run Number: 166658, Event Number: 34533931

Date: 2010-10-11 23:57:42 CEST

Observation of heavy particles in 1 fat-jet

t-tbar events are good candidates to look for heavy particles reconstructed in one fat-jet.

top \rightarrow Wb \rightarrow qqb: Three local sub-jets.

Two useful candles:

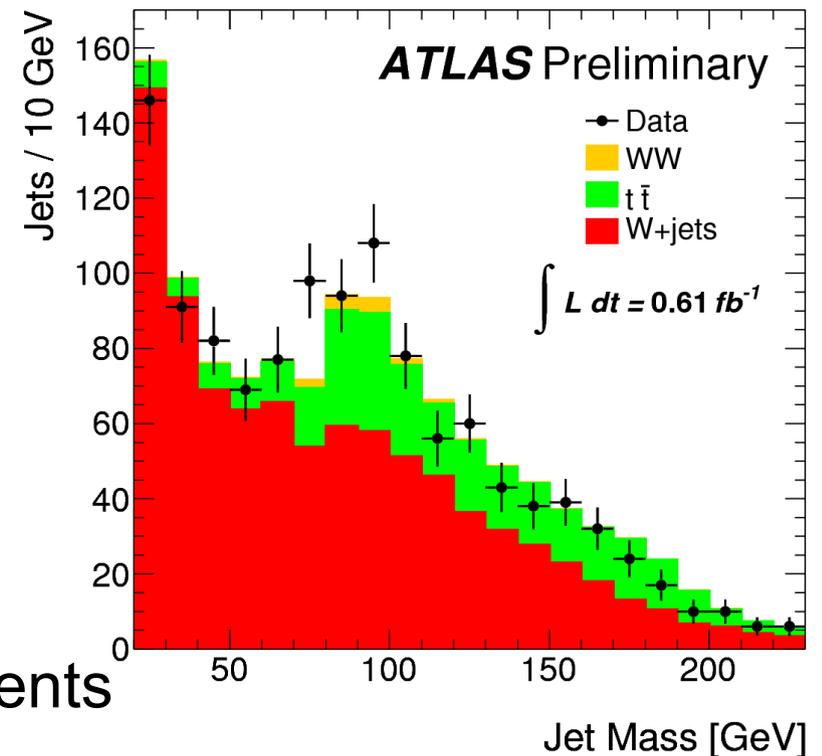
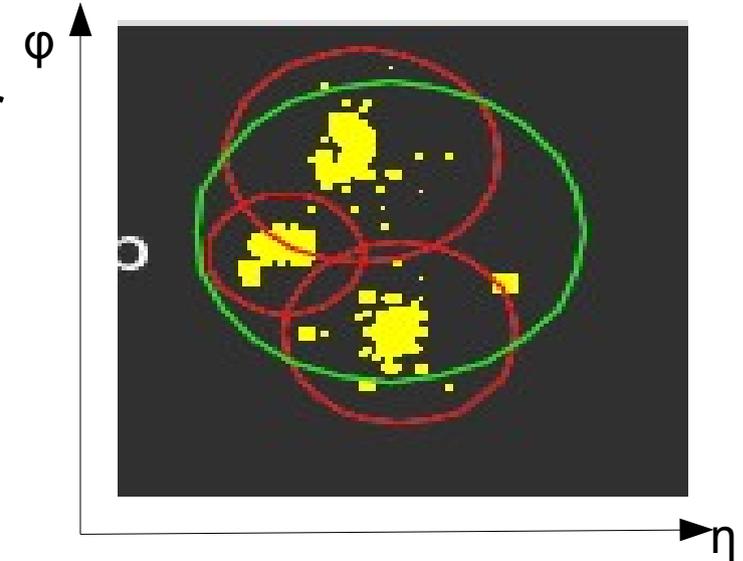
1) FatJet(qqb): The fat-jet integrates in its catchment area the three quarks.

We expect to have a mass for this jet \sim 170-180 GeV

2) FatJet(qq): The fat jet integrates in its catchment area the two q from W decay.

We expect to have a mass for this jet \sim 70-90 GeV

Seen in MET+lepton+jets($p_T > 180$ GeV) events



ATLAS-CONF-2011-103

The jet mass as observable

Different aspects need to be under control to use of the jet mass to look for particles reconstructed as a single jet.

The jet mass is a complex QCD observable:

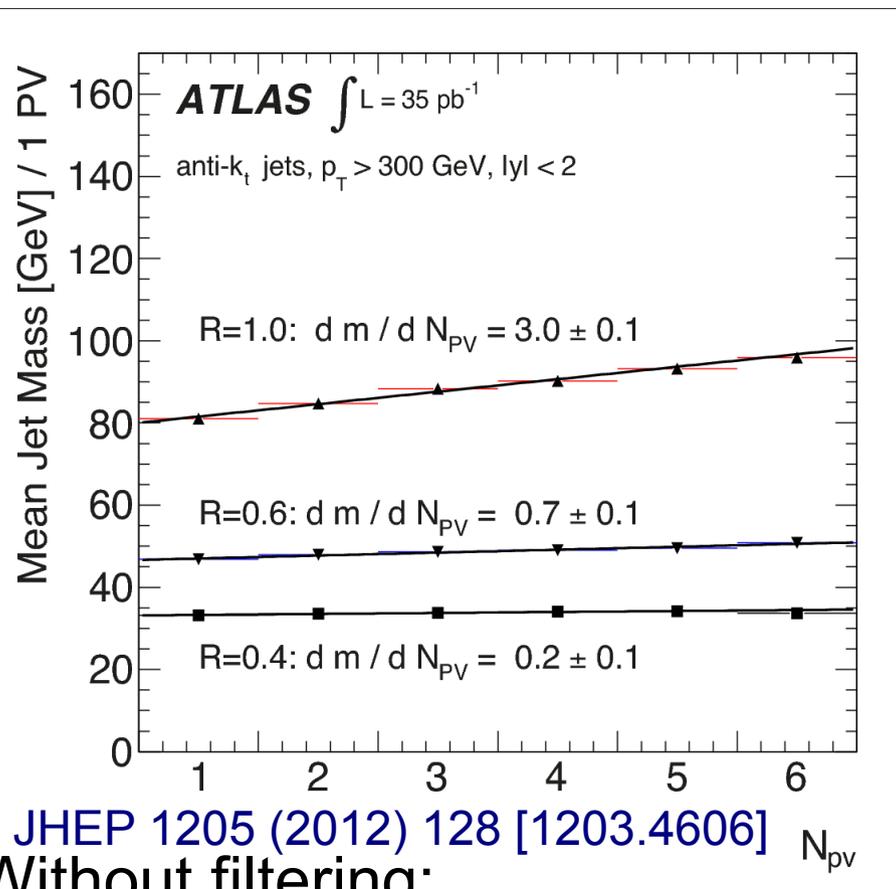
- In Monte Carlo, non trivial dependence on Parton shower+ non perturbative effects (fragmentation+UE)
- At the detector level, the capability to distinguish local deposition depends on the calorimeter granularity and on the size of the hadronic shower in the detector.
- The particles produced in extra partonic interaction (pile-up) can drastically change the mass.

In the last years different techniques have been developed to overcome these limitations, and some of them have been tested on data

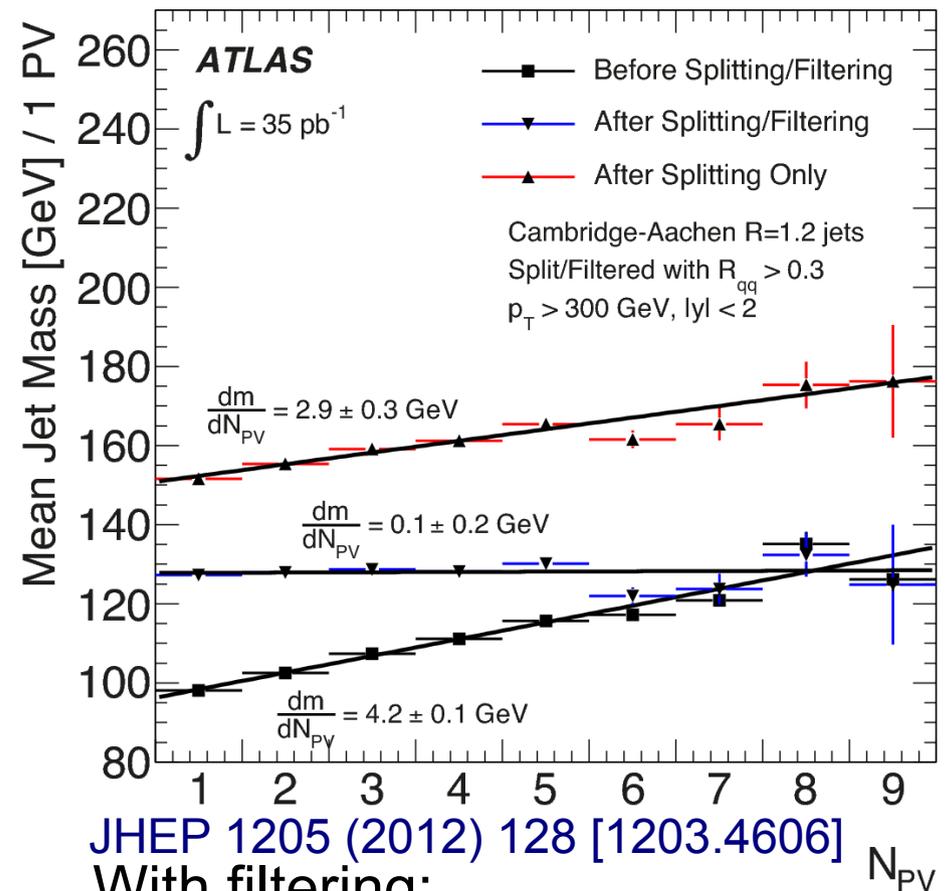
The jet mass in presence of pile-up

Jet filtering: filter away UE and pile-up contamination while retaining hard perturbative radiation from the heavy particle decay (i.e. Higgs).

Phys.Rev.Lett.100:242001,2008 [0802.2470]



Without filtering:
 Impact of pile-up increase with R



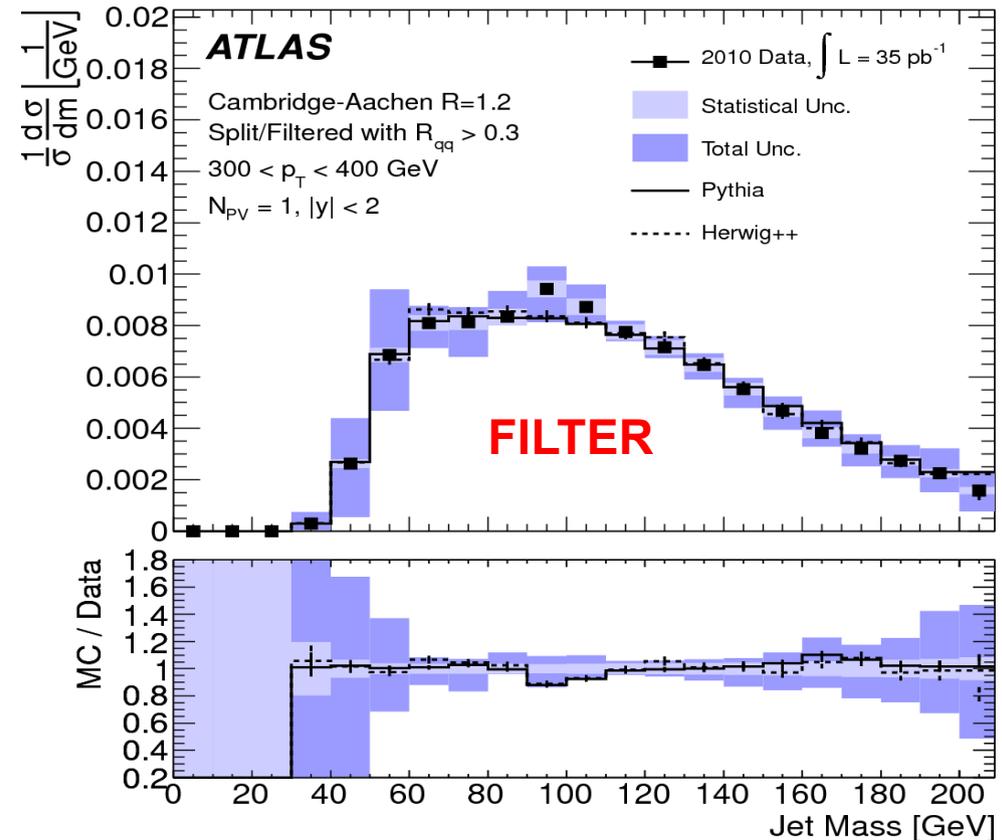
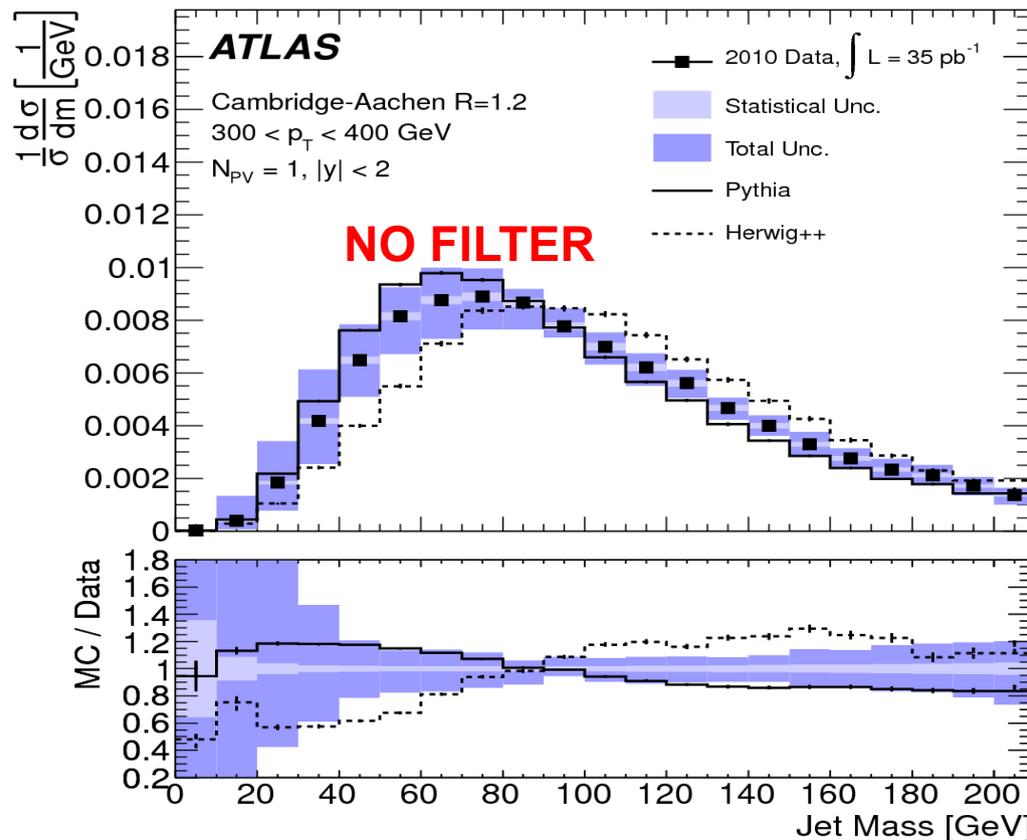
With filtering:
 Stable mass even for big R

Measuring the QCD jet mass

Mass of single jet:

JHEP 1205 (2012) 128 [1203.4606]

- Validation of the Monte Carlo description;
- Validation of the filtering techniques on data;
- Information on parton-shower properties:

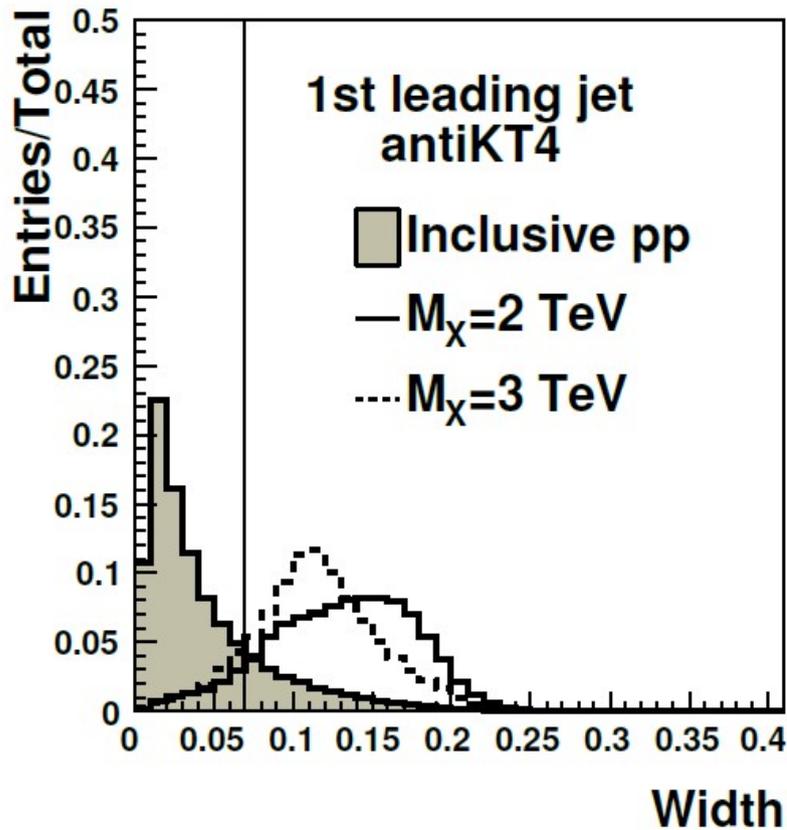


Filtered mass reduces dependence on soft physics
Better agreement data-MC(s)

Jet width

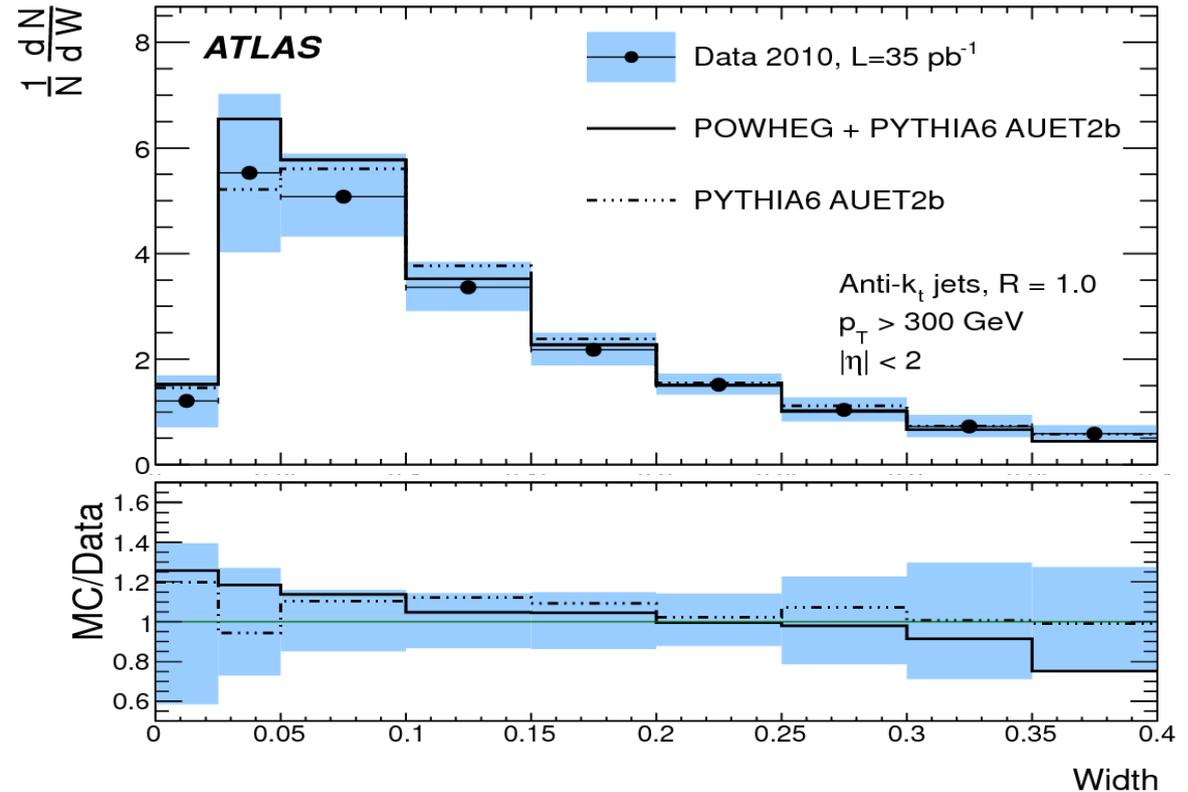
The shape of the jet is giving extra information on its properties, and it can be used to discriminate resonances from QCD jets. Several variables measured in .

Expectation



ANL-HEP-PR-10-2 [1002.3982]

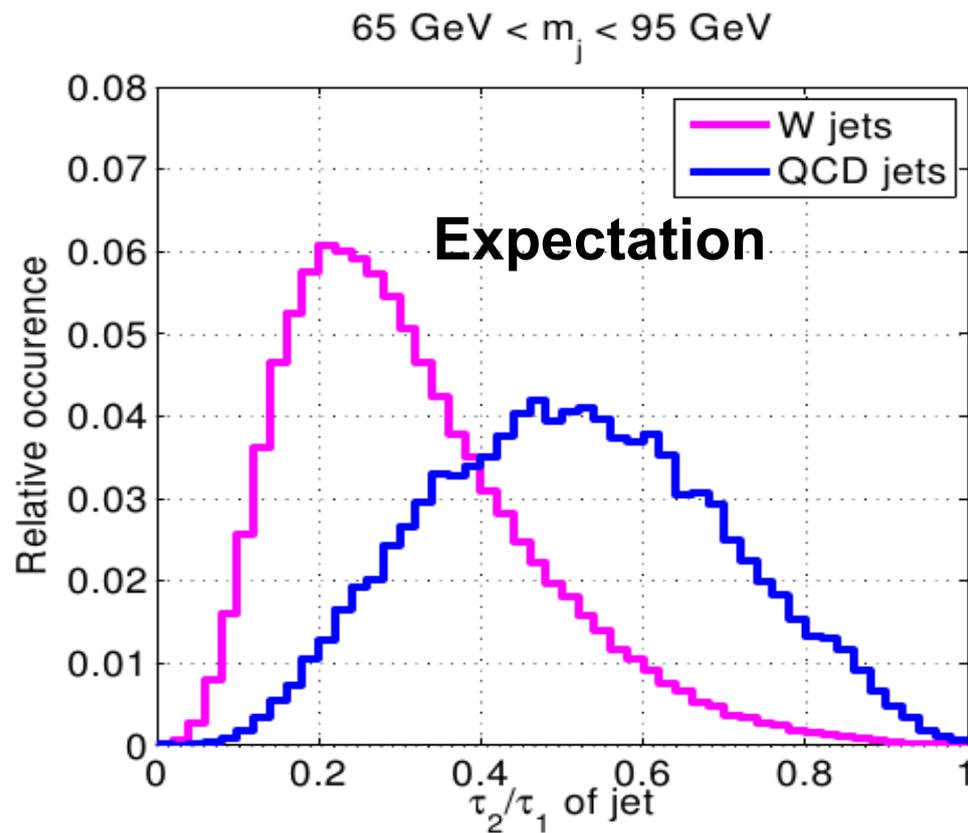
Measurement in QCD jets



ATLAS-CONF-2012-044

Sub-jetness

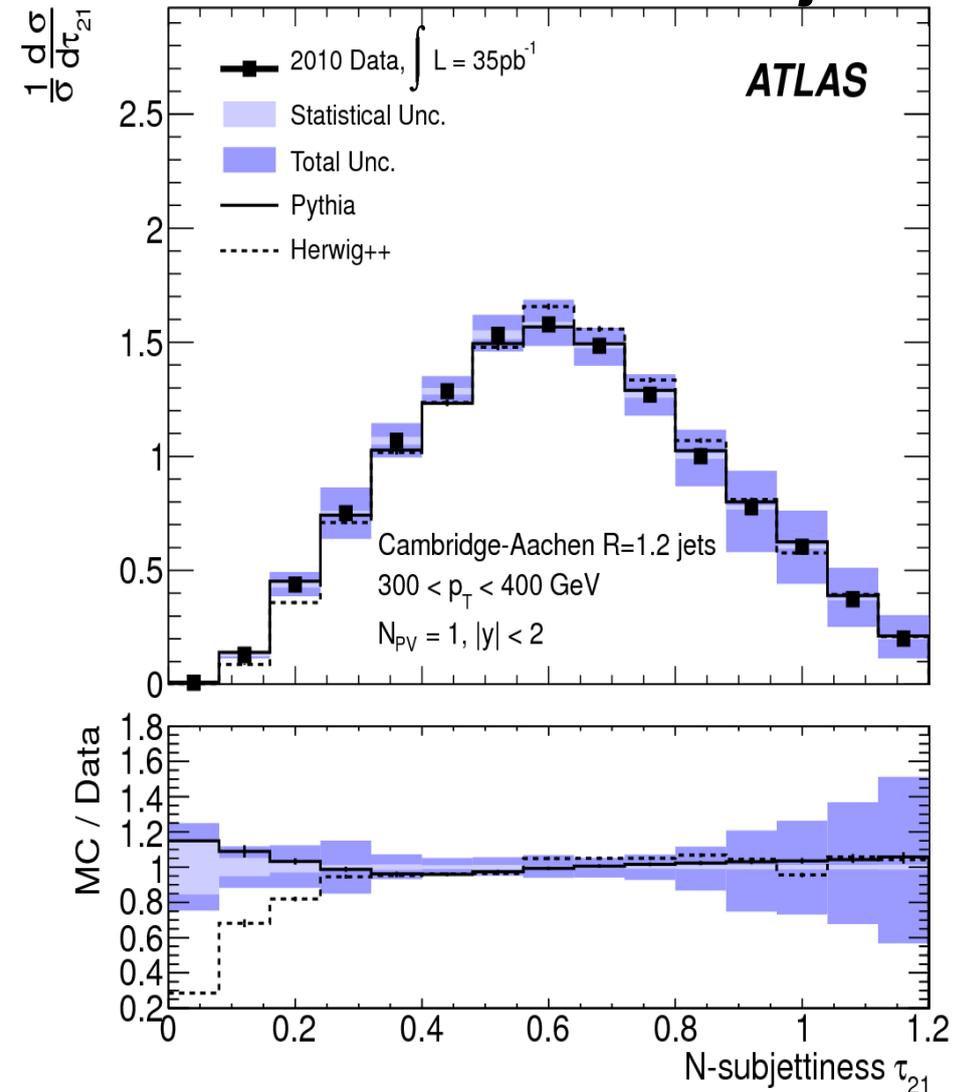
τ_n sub-jetness: Testing the hypothesis that a jet is composed by n -subjets, by using its constituents.



JHEP 1103:015,2011 [1011.2268]

Current MC describes quite well this variable.

Measurement in QCD jets



JHEP 1205 (2012) 128 [1203.4606]

A complex use-case:

$H \rightarrow b \bar{b}$

Introduction

$H \rightarrow b\bar{b}$ is the dominant decay mode for low masses.

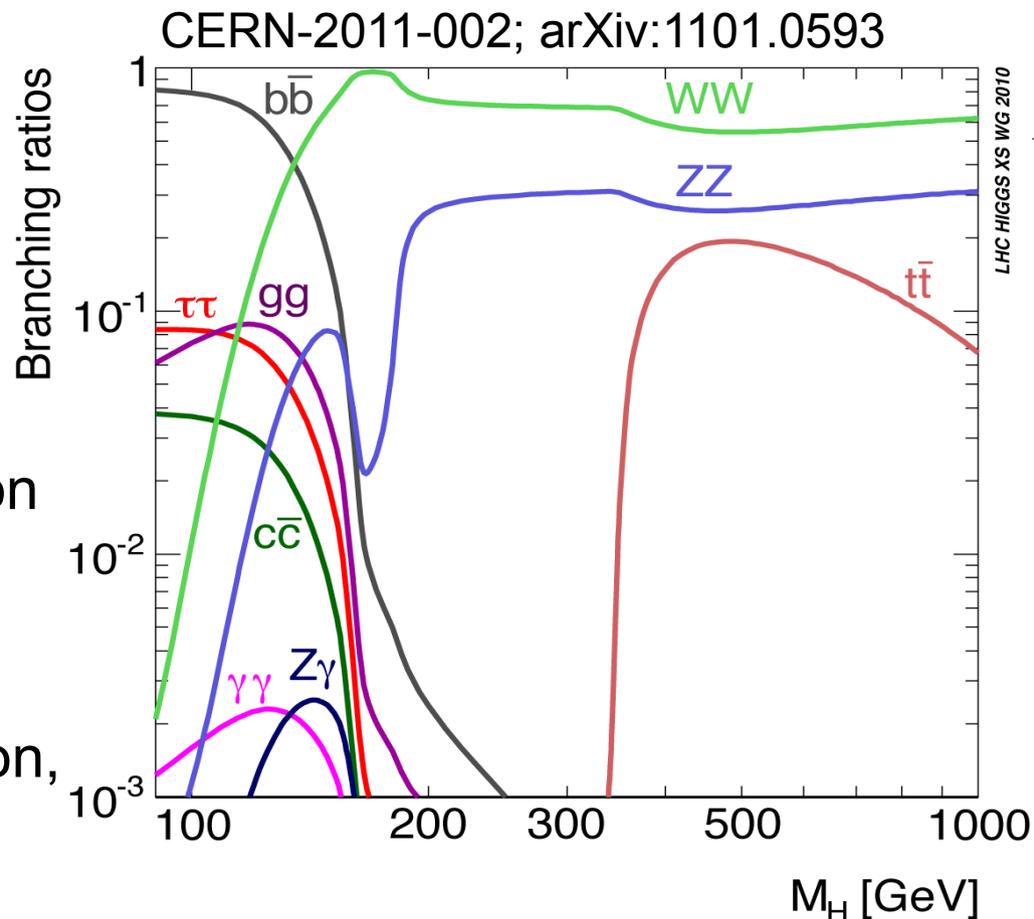
To reduce the QCD background, one can study the associated production with vector bosons VH :

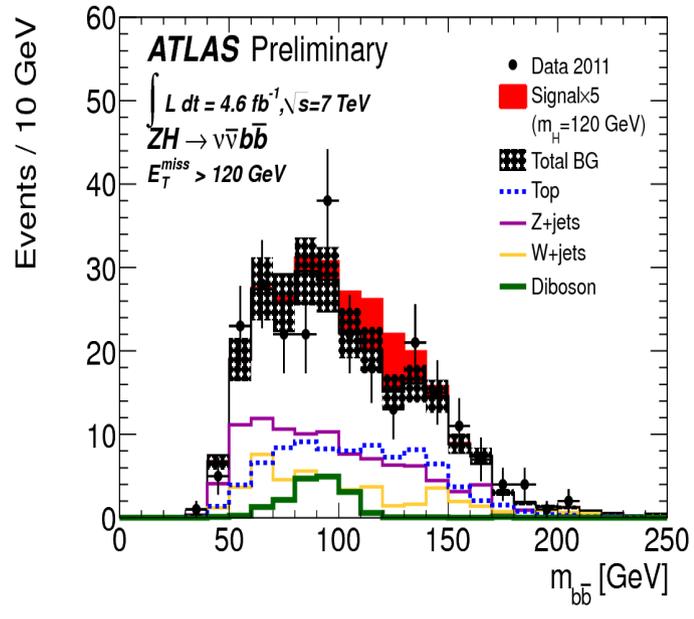
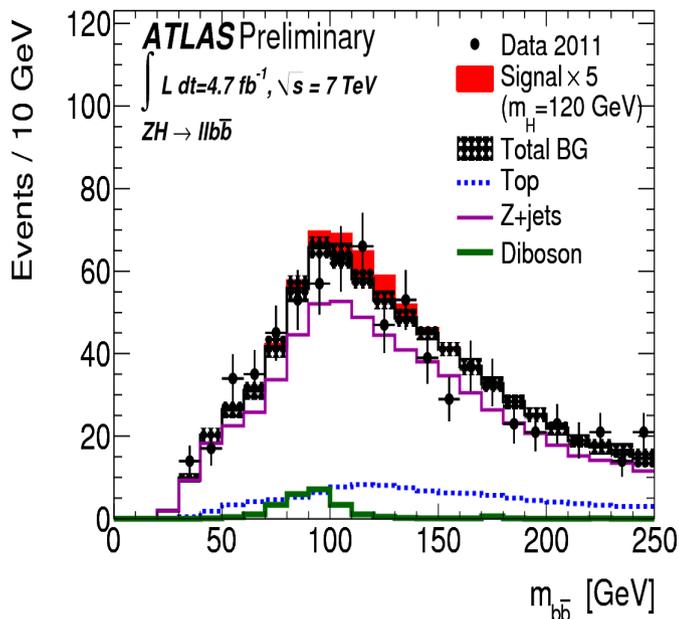
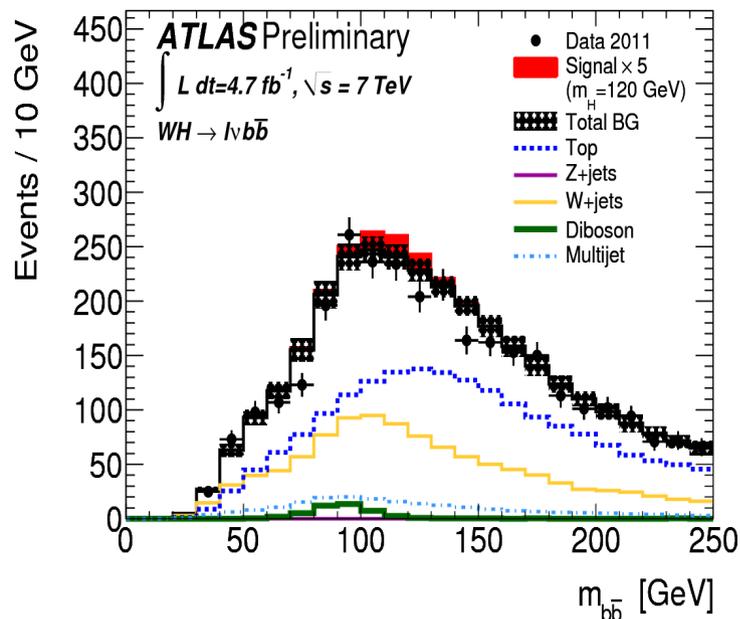
WH and ZH .

By requiring a boost for the vector boson, the signal/background increase.

In this way, the $b\bar{b}$ system start to be in a collimated topology. One of the channels where analysis with jet substructure techniques can be really useful.

Several experimental and theoretical aspects enters in the analysis. They need to be dominated to get the best sensitivity on this channel.





Systematic component	B-tagging	Jets and MET	Background normalization	Theory
	1-10%	2-15%	3-5%	1-20%

Different backgrounds need to be understood.

Experimental systematic uncertainties, and background normalization can be constrained by measuring control region + side bands properties.

Theory uncertainty is mostly dominated by the $V p_T$ uncertainty and the m_{bb} shape variation from different MC predictions.

Questions for “classic” and “sub-jet” analysis

- 1) how well do we know the truth properties of b-jets:
→ B jet shapes measurement (i.e. in tt events?)
- 2) The b-bbar system can emit an extra gluon. Is this properly described for the signal and for the background in the Monte Carlo?
- 3) How well do we simulate the boosted b-bar system:
We need to model several background, mostly with Monte Carlo.
Different generators are available, and they could be in disagreement.
We need to define region where these properties for the boosted b-bar system can be measured and constrained.
- 4) How well the B hadrons properties in jets are reproduced by MC?
(we have seen that the D^* are not properly simulated).
How much is this affecting the final result?

Different challenges on these channel. Stay tuned!!

Conclusions

Jet reconstruction and performances:

→ Good understanding of the jet energy scale and resolution in 2010 and 2011 ATLAS data.

Jet properties:

→ Several measurements used to tune the MC simulations, and to understand the non perturbative effects, such as UE and fragmentation.

Cross section measurements:

→ Good agreement of data and pQCD.

First steps towards precision measurements (i.e. inclusive jet cross section).

Searches with jets:

→ No evidence of new phenomena in inclusive jet final state → limit setting.

Jet substructure:

→ inclusive QCD substructure measurements:
first milestone for identification of boosted objects.

Improvements in the description of more exclusive final state will “boost” our searches with jet.

BACKUP

ATLAS Calorimeter System

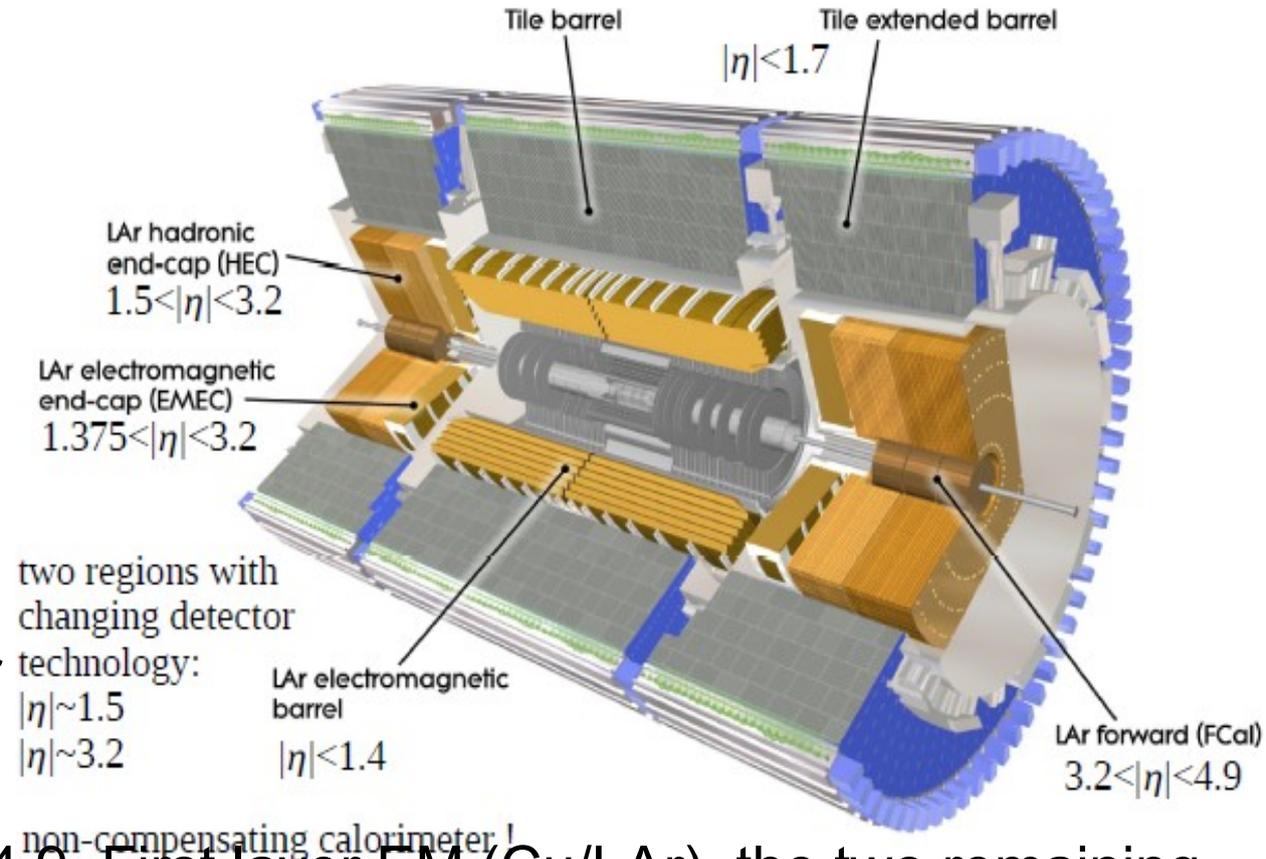
EM LAr: $|\eta| < 3$ - Pb/LAr calorimeter, high resolution for e/ γ objects. e/h ~ 1.7

Central hadronic calorimeter (**TileCal**): $|\eta| < 1.7$: Fe(82%), scintillator (18%) - e/h = 1.36

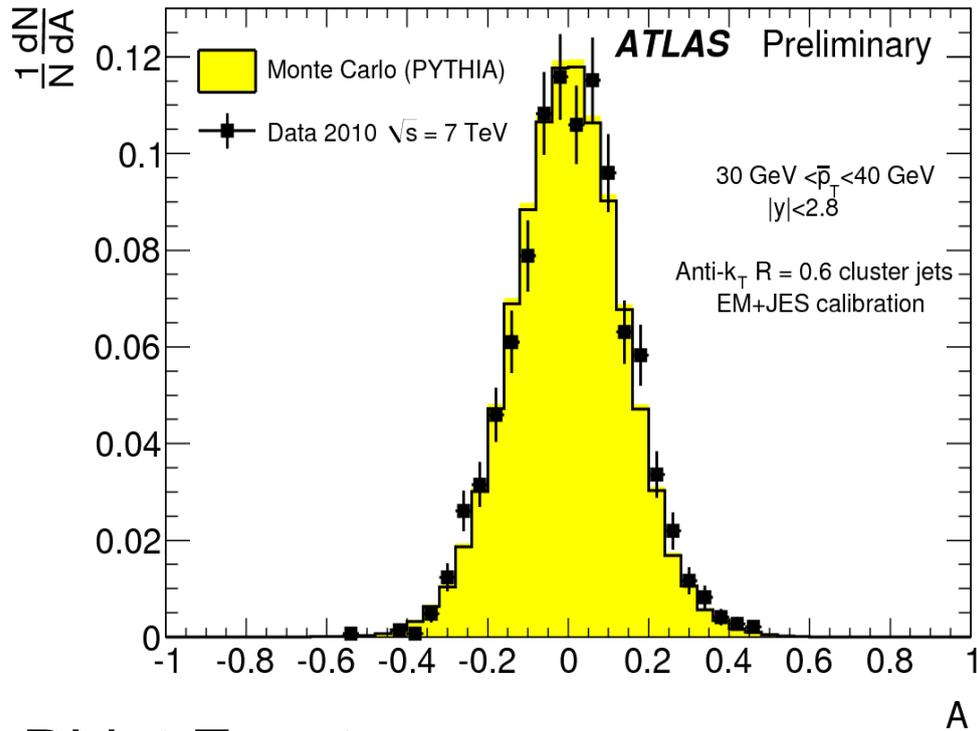
End Cap Hadronic Calorimeter (**HEC**): $1.7 < |\eta| < 3.2$ - Cu/LAr

Forward calorimeter: $3 < |\eta| < 4.9$. First layer EM (Cu/LAr), the two remaining layers HAD.

Highly hermetical ($|\eta| < 5$), non compensating calorimeters.



Jet Resolution



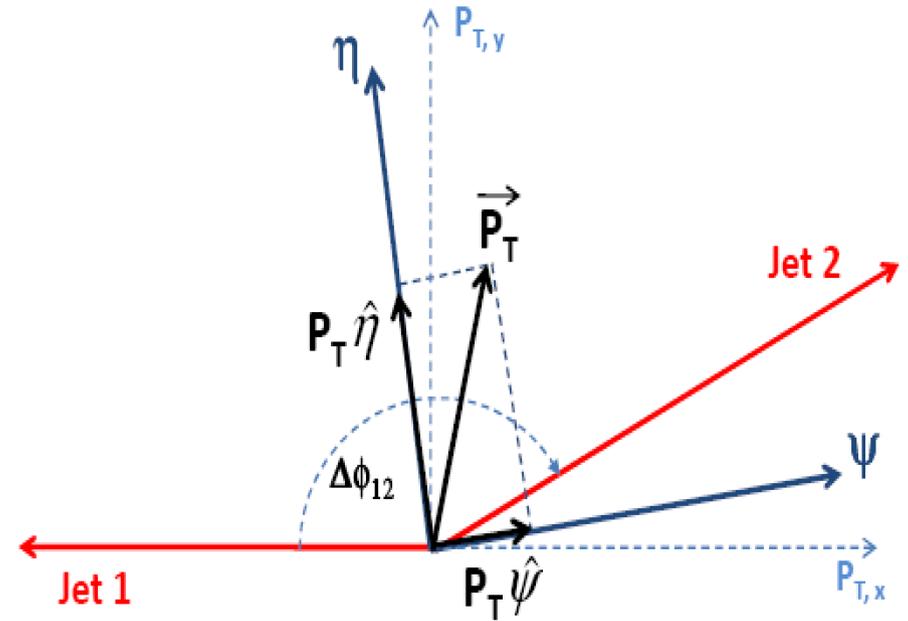
Di-jet Events

$\Delta\phi > 2.8$ and $p_T[3^{\text{rd}} \text{ jet}] < 10 \text{ GeV}$

$$A = (p_T[1] - p_T[2]) / \langle p_T \rangle$$

$$\sigma(p_T)/p_T = \sqrt{2} \sigma_A$$

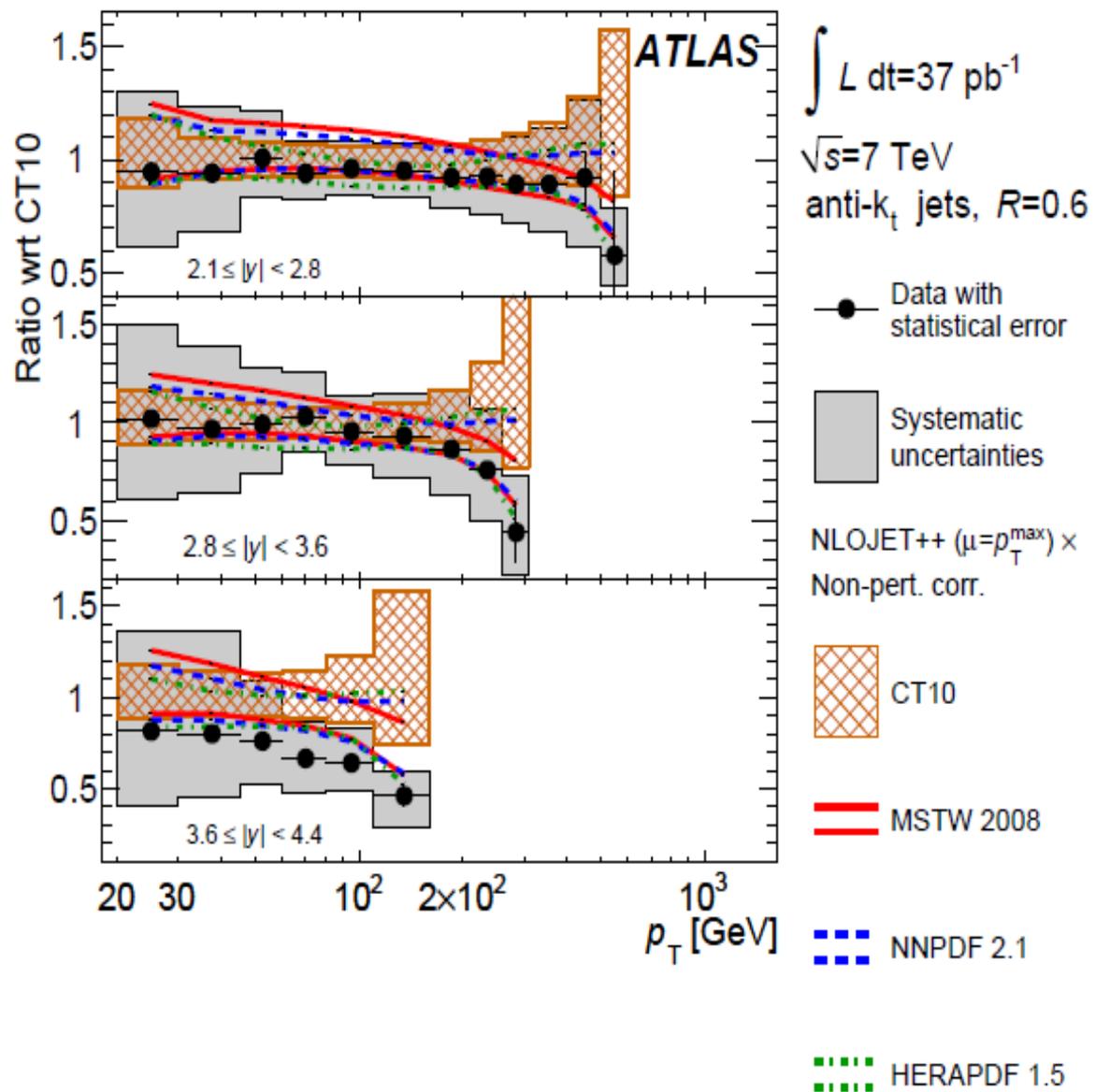
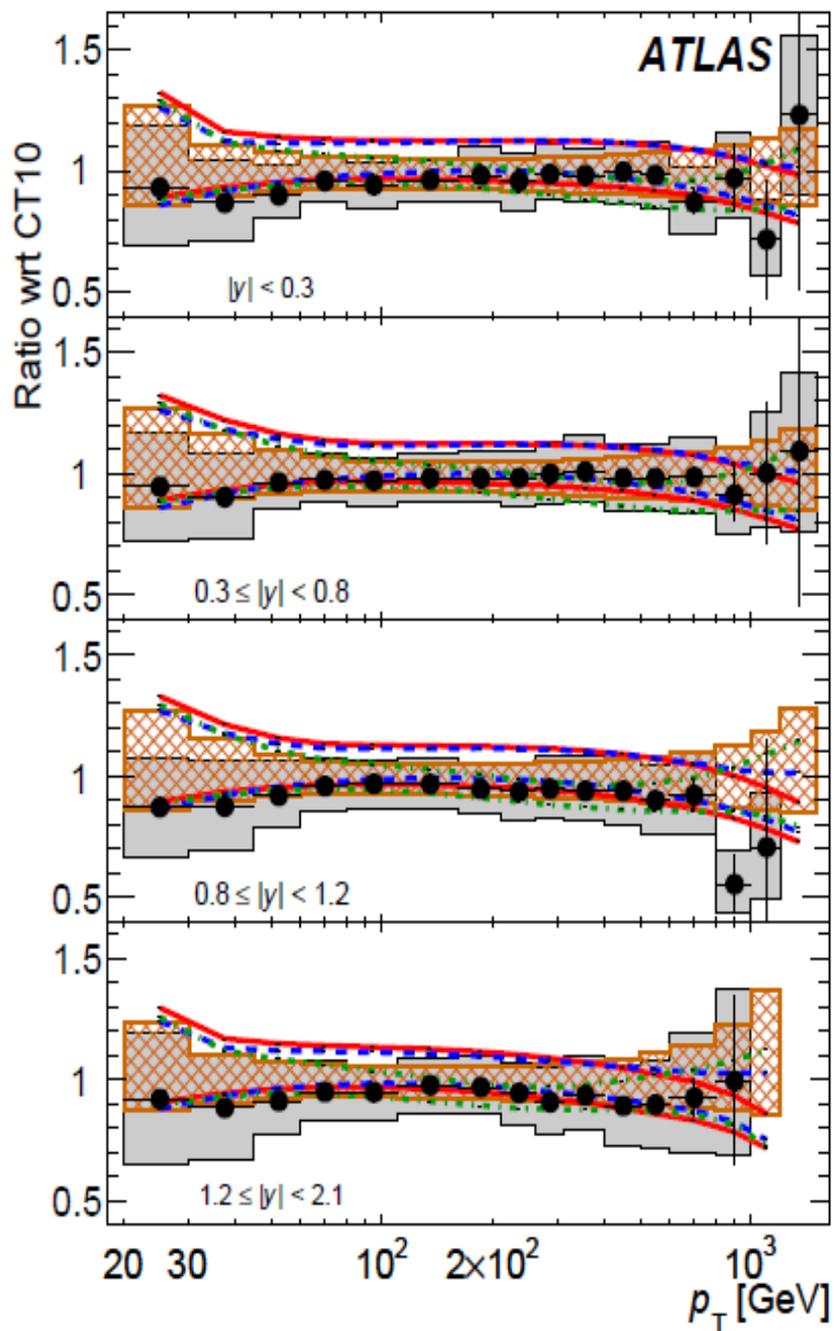
Estimate of the unbalance due to soft radiation.

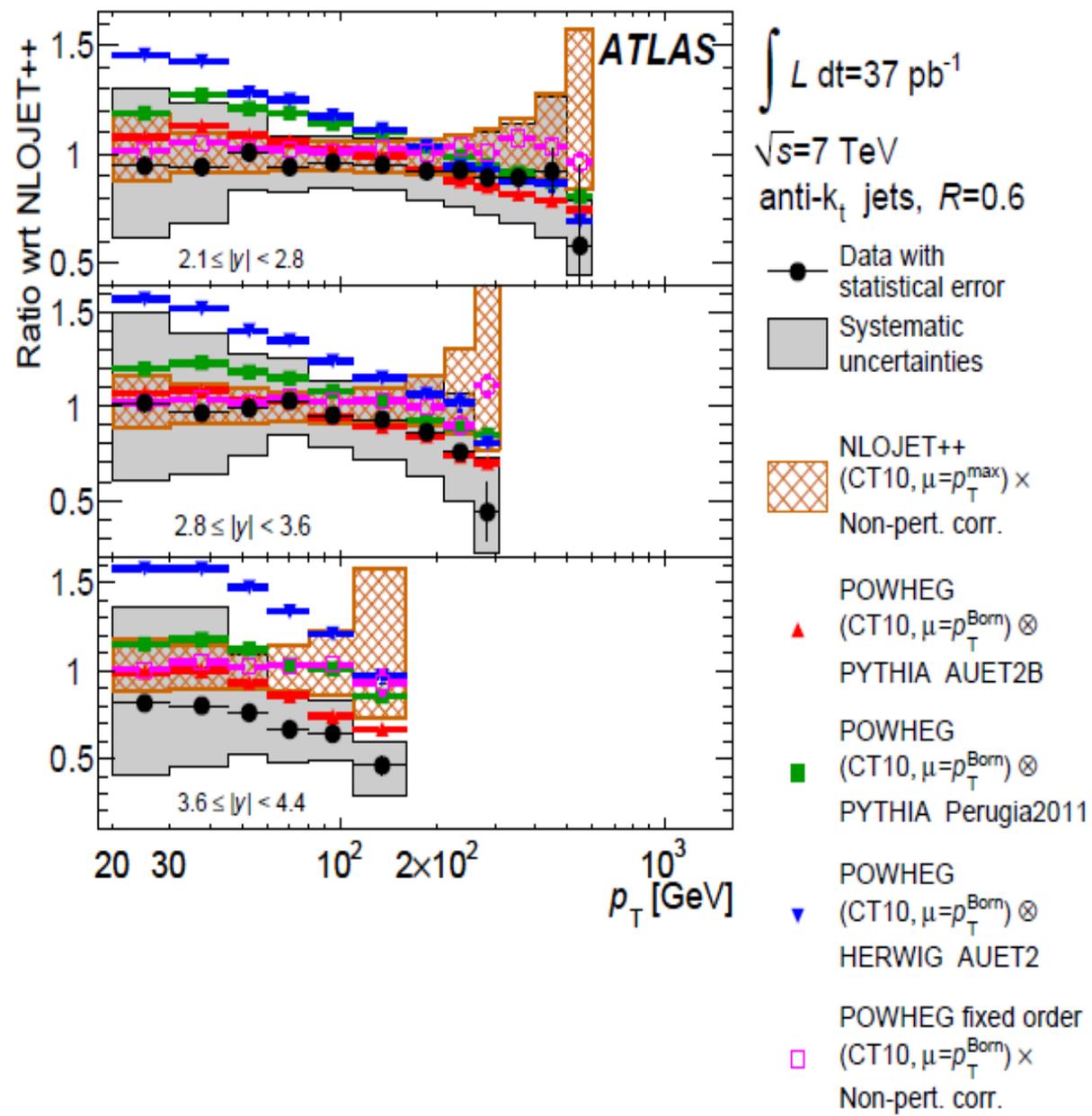
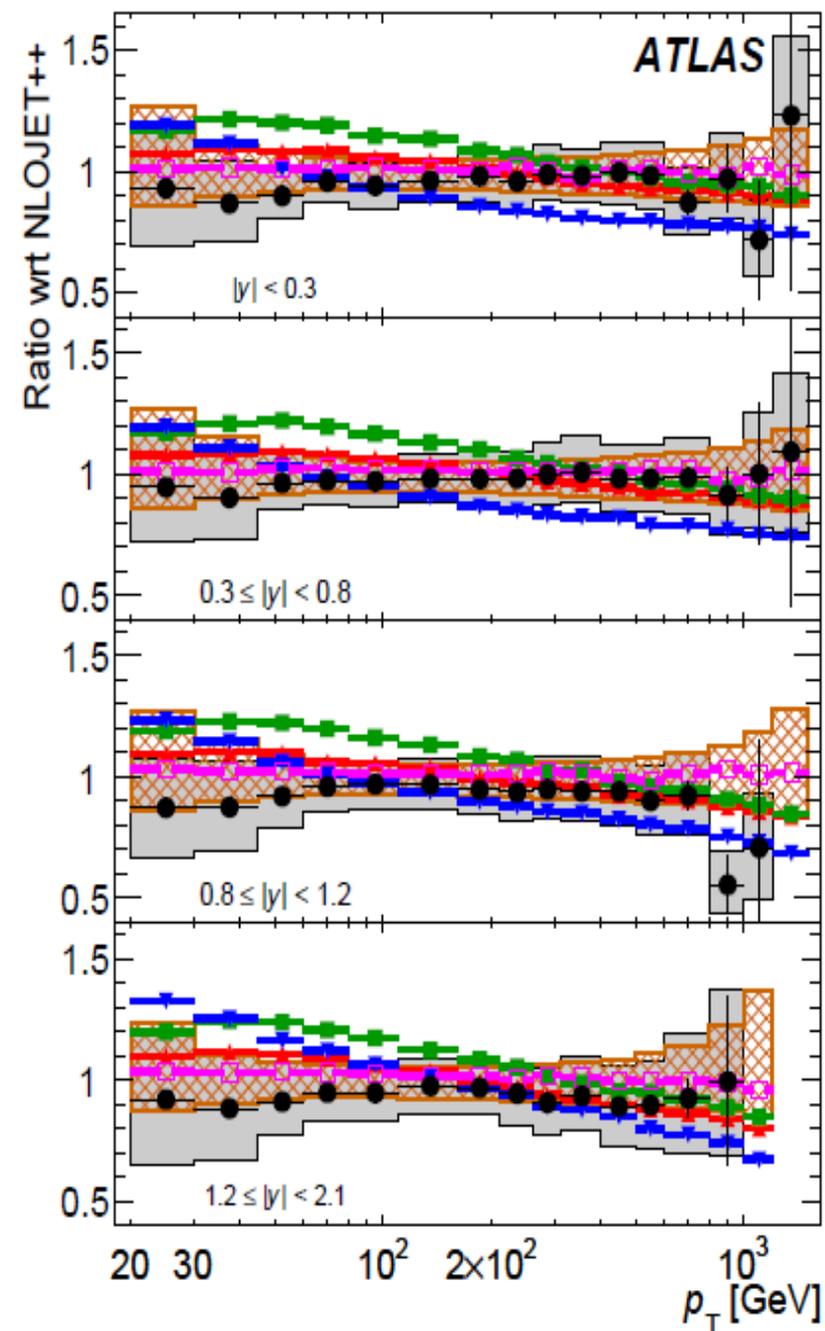


Particle Level: $\sigma_\psi \sim \sigma_\eta \neq 0$ Radiation

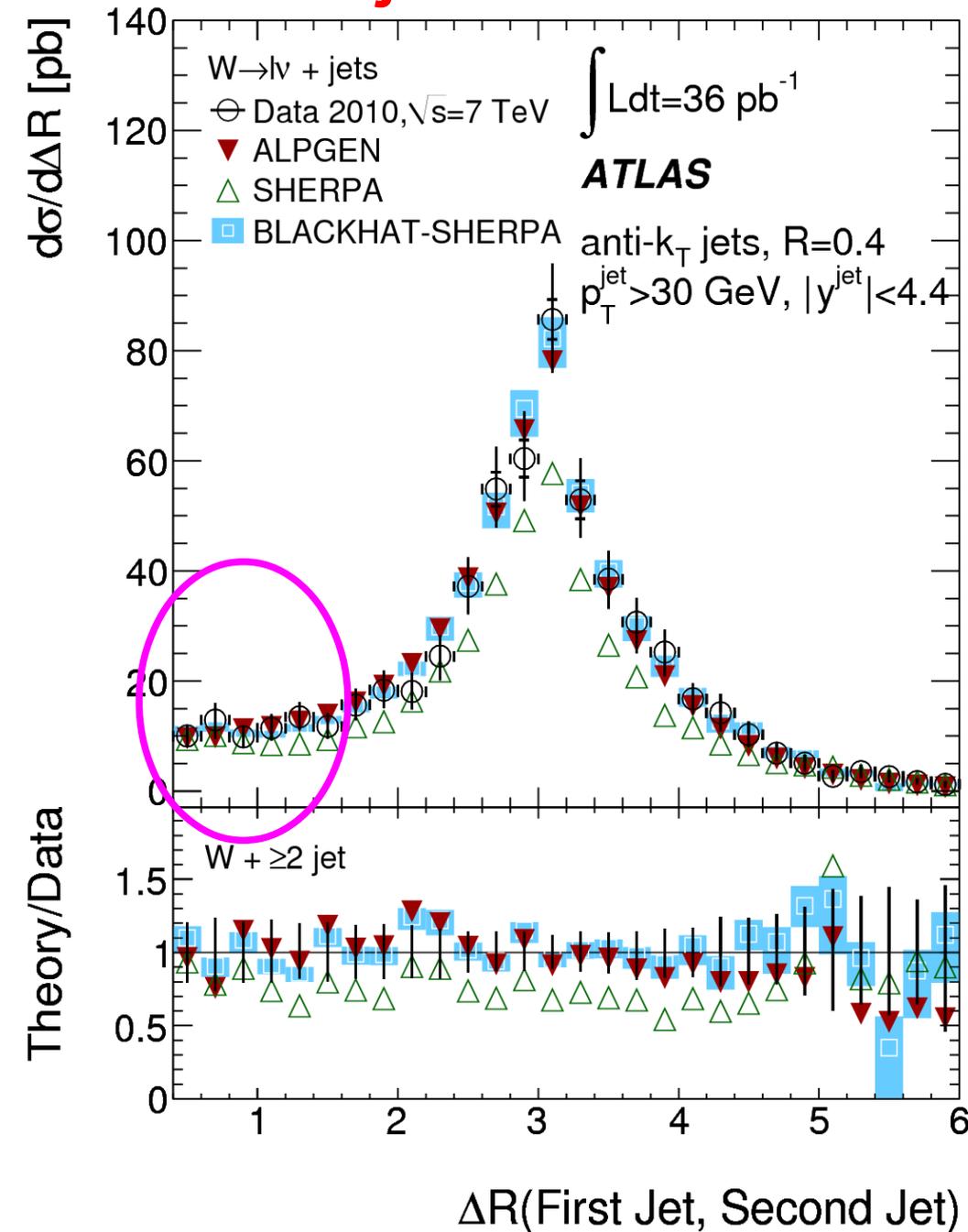
Detector Level:

$$\sigma(p_T)/p_T = \frac{\sqrt{(\sigma_\psi^2 - \sigma_\eta^2)}}{\sqrt{2} \langle p_T \rangle |\cos(\Delta\phi_{12})|}$$

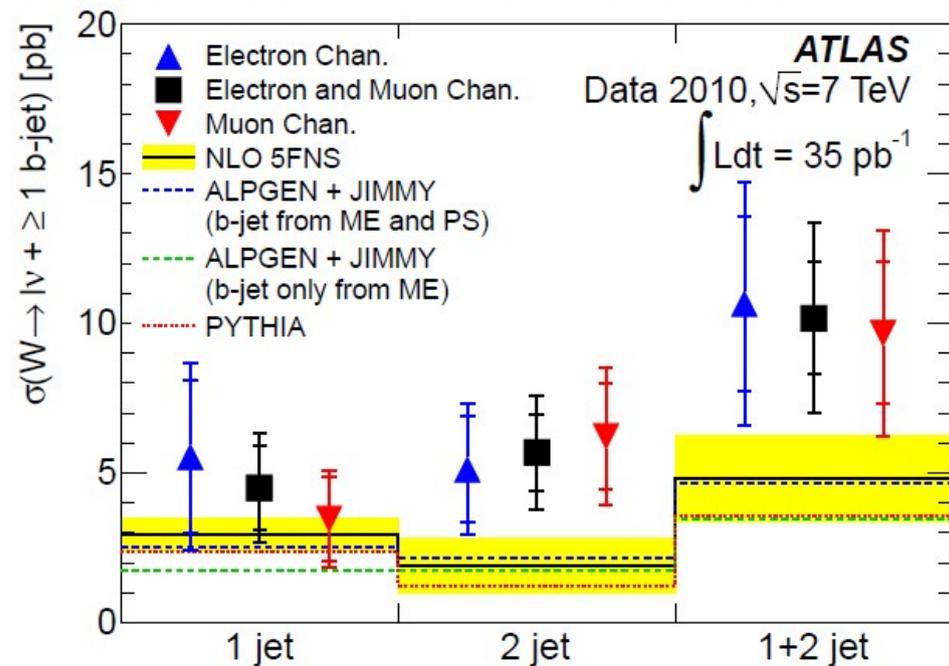




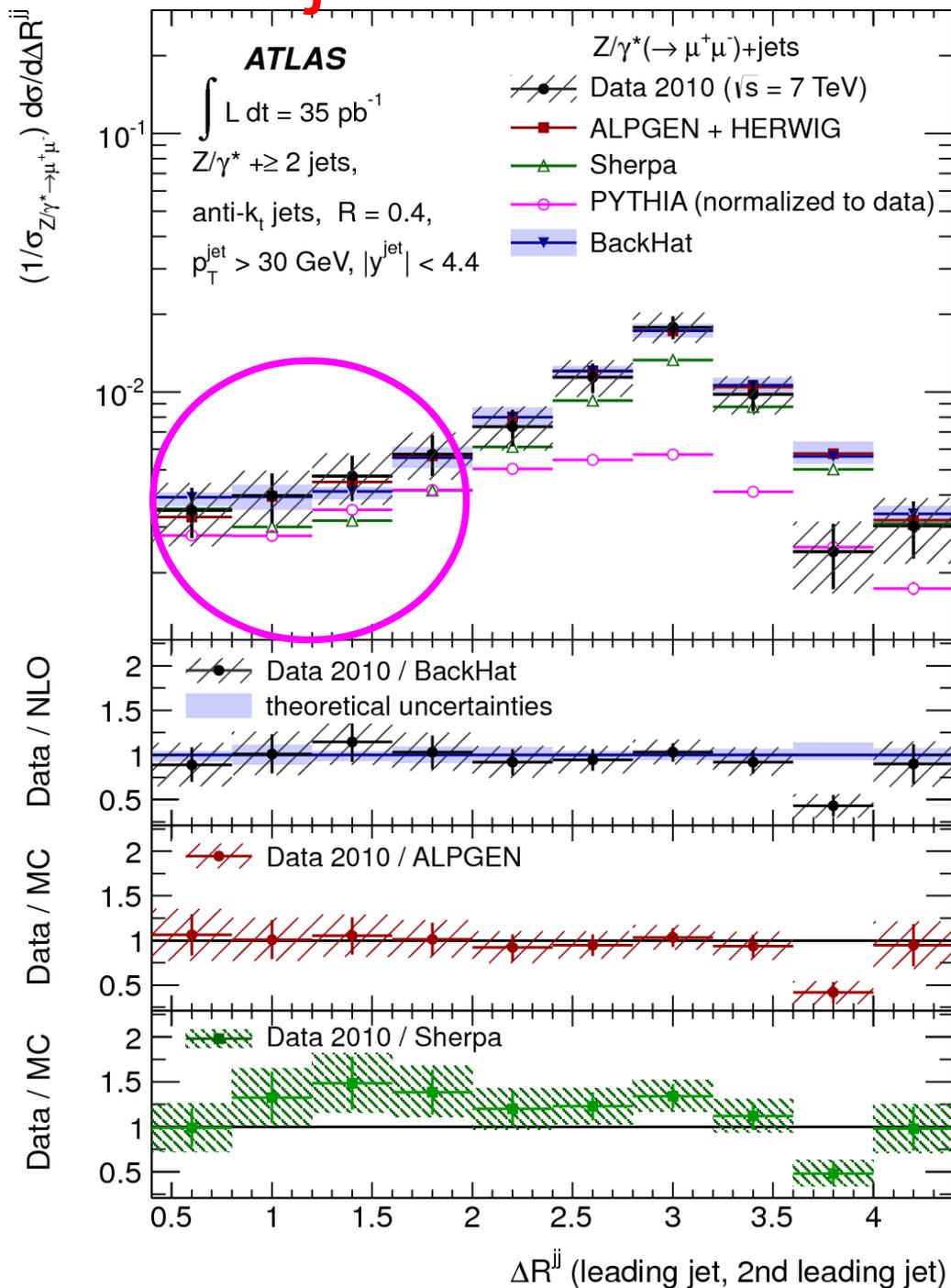
W+j cross section



Wb cross section



Z+j cross section



Zb cross section

