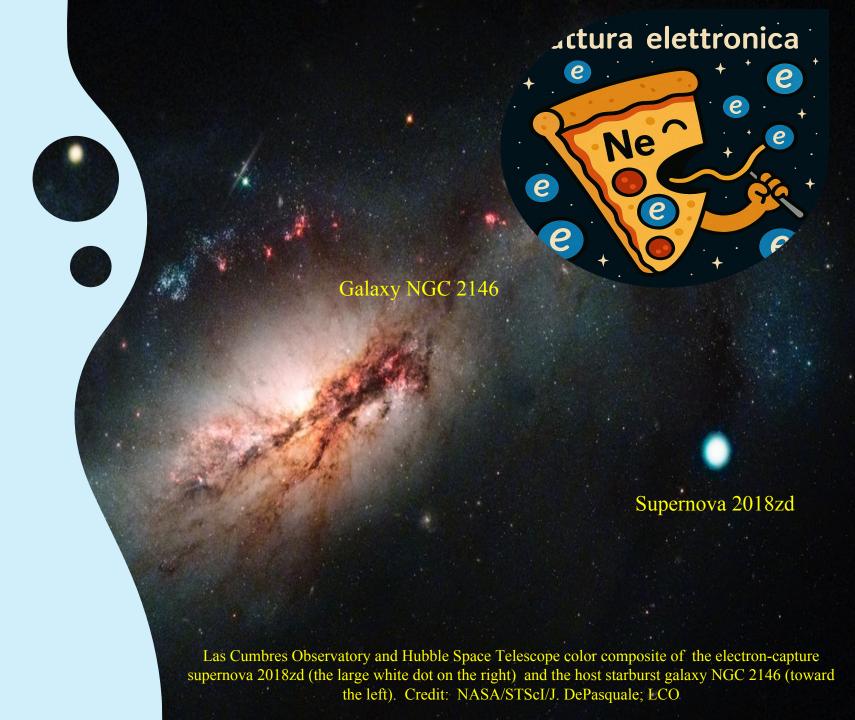
# Fermionic Dark Matter and Electron-Capture Supernovae

Pathways to Low-Mass Neutron Stars

Vishal Parmar

INFN Pisa, Italy

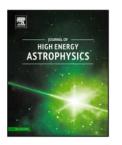
Nov 19, 2025 SEMINARIO SCIENTIFICO, INFN Pisa



Contents lists available at ScienceDirect

#### Journal of High Energy Astrophysics

journal homepage: www.elsevier.com/locate/jheap



Triggering electron capture supernovae: Dark matter effects in degenerate white-dwarf-like cores of super-asymptotic giant branch stars

Vishal Parmar <sup>a,©</sup>,\*, Domenico Scordino <sup>b,©</sup>, Ignazio Bombaci <sup>a,b,©</sup>

# Based

Home / Astronomy & Space / Astronomy









OCTOBER 22, 2025 FEATURE





#### Exploring how dark matter alters electroncapture supernovae and the birth of neutron stars

by Ingrid Fadelli, Phys.org



INFN, Sezione di Pisa, Largo B. Pontecorvo 3, I-56127, Pisa, Italy

Dipartimento di Fisica "E. Fermi", Università di Pisa, Largo B. Pontecorvo 3, I-56127, Pisa, Italy

# **Outline**

#### Introduction

- Nuclear Matter
- Neutron Star Anatomy
- Equation of State: Uncertainties and Implications

Electron capture supernovae (ECSNe): Discovery and Implication

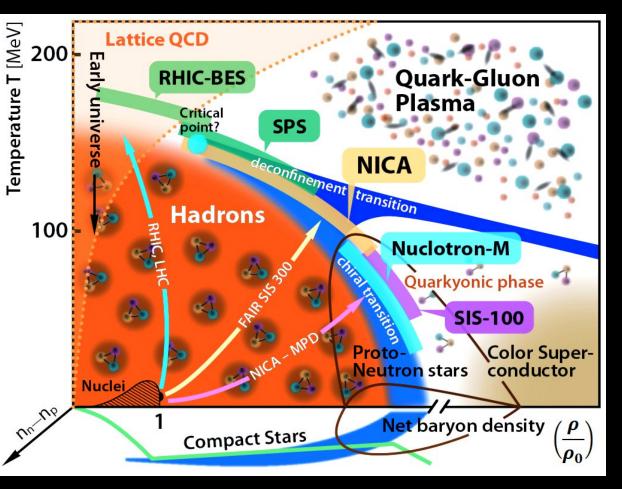
EoS for Dark Matter, Neutron star and White dwarf

Dark matter effect on Neutron star formed from ECSN

Implication and Future Prospective

# Nuclear Matter: Nuclei and Neutron Star

Nuclear Matter Phase Diagram



Neutron Star: A Physics Goldmine

Astrophysics: What is the minimum mass of a black hole?

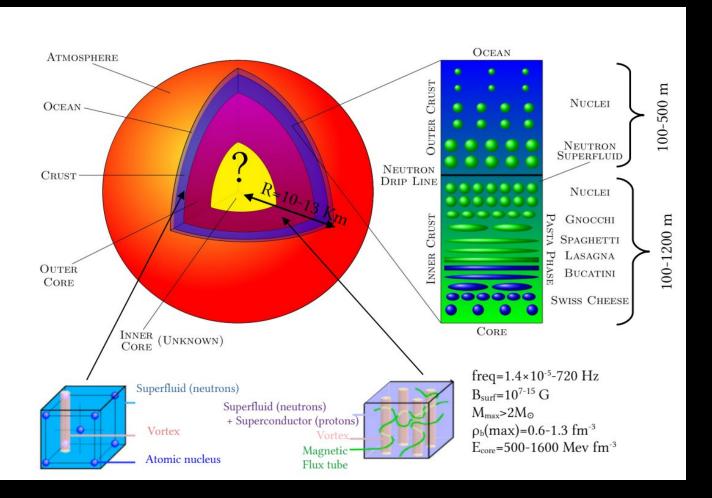
Nuclear Physics: What are the limits of nuclear existence and the EOS of nuclear matter?

Particle Physics: Dense quark matter and quark-gluon plasma?

General Relativity: Rapidly rotating neutrons stars as a source of gravitational waves?

Condensed-Matter Physics: Signatures for the liquid to crystalline state transition?

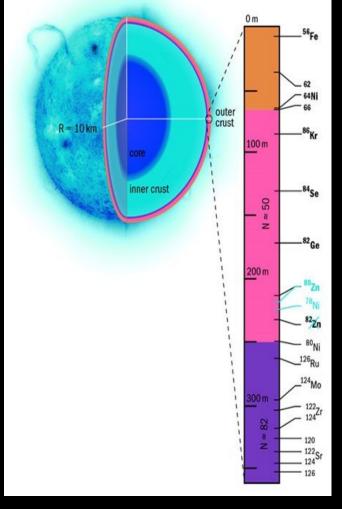
# **Anatomy of Neutron Star**

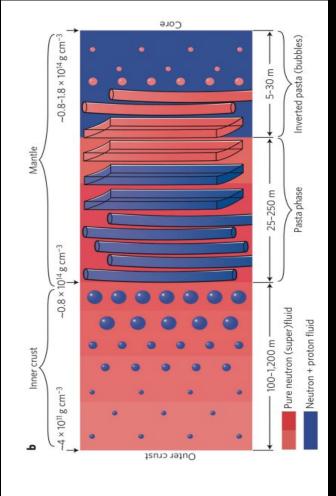


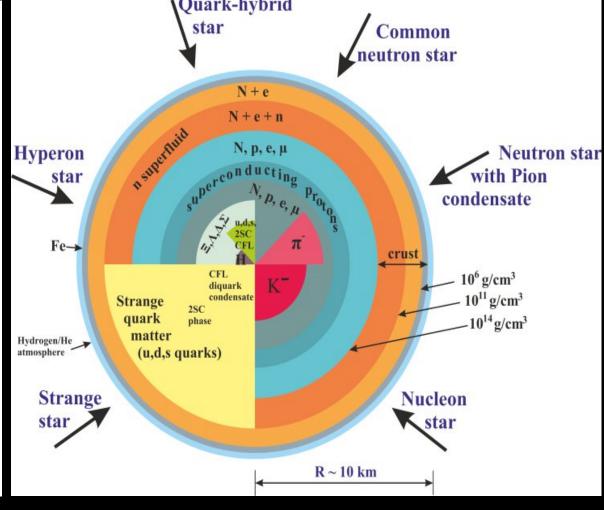
#### From Crust to Core

- Outer Crust:  $10^{-10}\rho_0 < \rho < 10^{-3}\rho_0$  "Coulomb Crystal" of progressively more neutron-rich nuclei.
- Inner Crust  $10^{-3}\rho_0 < \rho < 10^{-1}\rho_0$  "Nuclear Pasta" xotic shapes immersed in a neutron vapor.
- Outer/Inner Core  $10^{-1}\rho_0 < \rho < 10\rho_0$  "Fermi Liquid" of uniform neutron-rich atter ("Exotic Phases?").

Annual Review of Nuclear and Particle Science 56, 327–374 (2006)







#### **Outer Crust**

A solid Coulomb lattice of fully ionized atomic nuclei submerged in electron gas.

#### **Inner Crust**

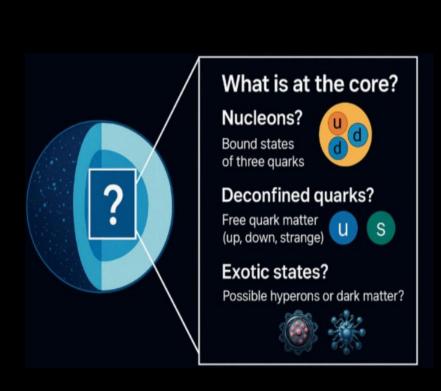
Neutrons "drip out of nuclei"

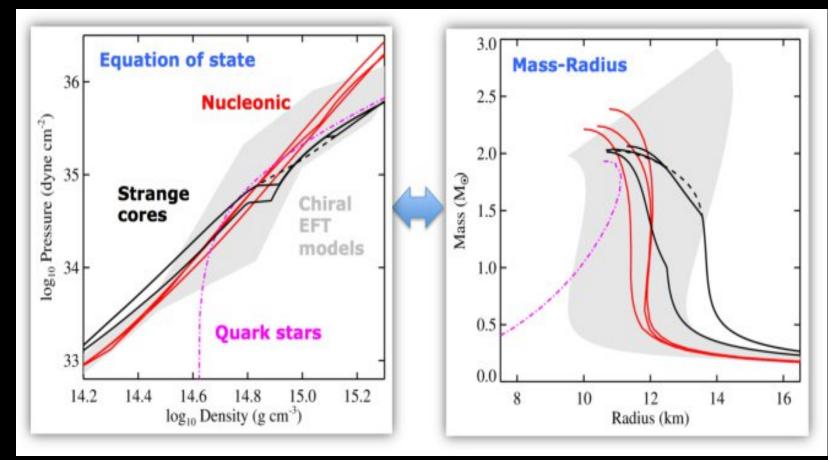
#### Core

Uniform neutron-rich matter in chemical equilibrium.

#### EoS and Mass-Radius of star

arXiv.1501.00042

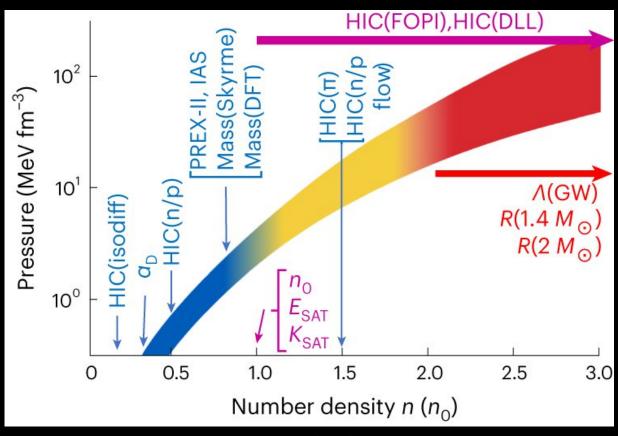




EoS must span 10-11 orders of magnitude in baryon density.

EoS needs joint analyses of nuclear theory, nuclear experiment and astrophysical observations

#### Constraints on the Ns EoS



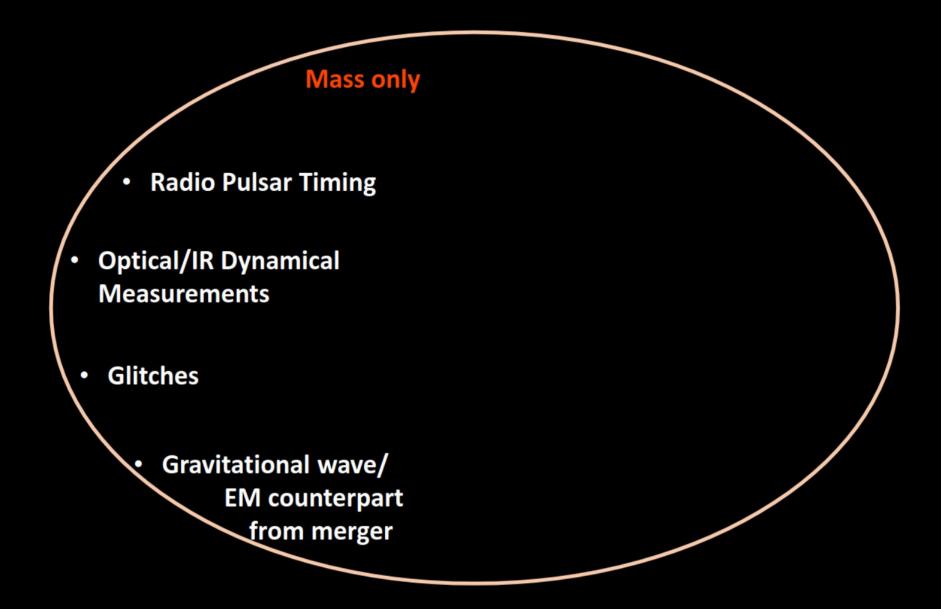
PSR J0952-0607 Romani et al. APJL 2022 NICER J0740+6620 J0348+0432 Dittmann et al. APJ 2024 Antoniadis et al. Scince 2013 NICER J0030+0451 GW 170817 disfavored Miller et al. APIL 201: PSR J0437-4715 Bauswein et al. APJL 2017 Choudhury et al  $M [M_{\odot}]$ GW 170817 Abbott et al. PRL 2019 GW 170817 disfavored Annala et al. PRL 2018 0.5 HESS J1731-347 Doroshenko et al. Nature Astronomy 2022 10 12 14 R [km]

Nature Astronomy volume 8, pages 328–336 (2024)

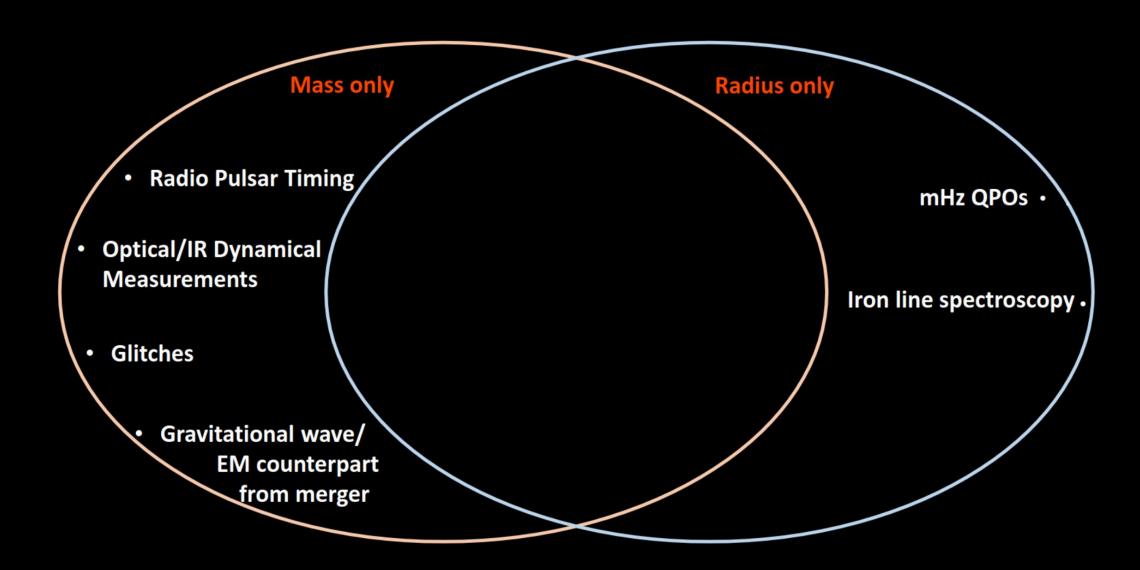
*Universe* **2025**, 11(5), 141

- A given NS EoS should at least satisfy <sup>2M⊙</sup> mass star.
- Ns mass and radius provide most significant information on the EoS.

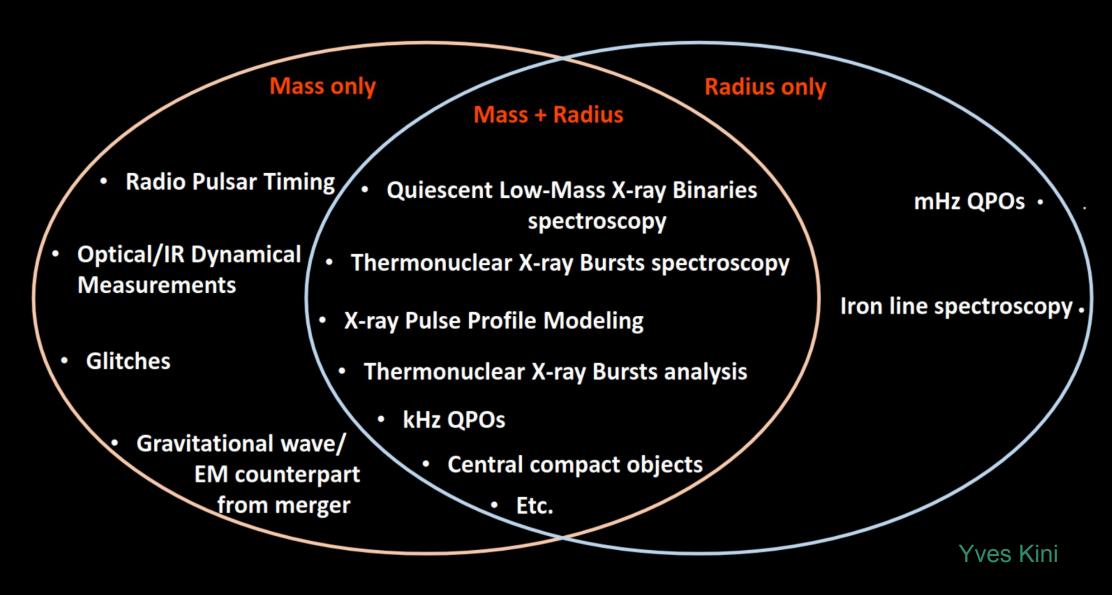
### Techniques to constrain the mass and/or Radius

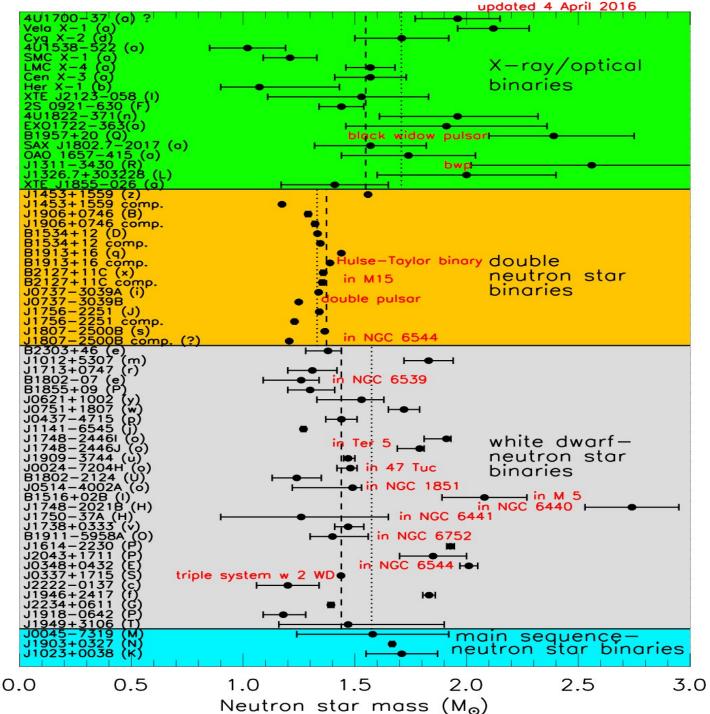


#### Techniques to constrain the mass and/or Radius



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DNSs are no more compatible with 1.4 M anymore.

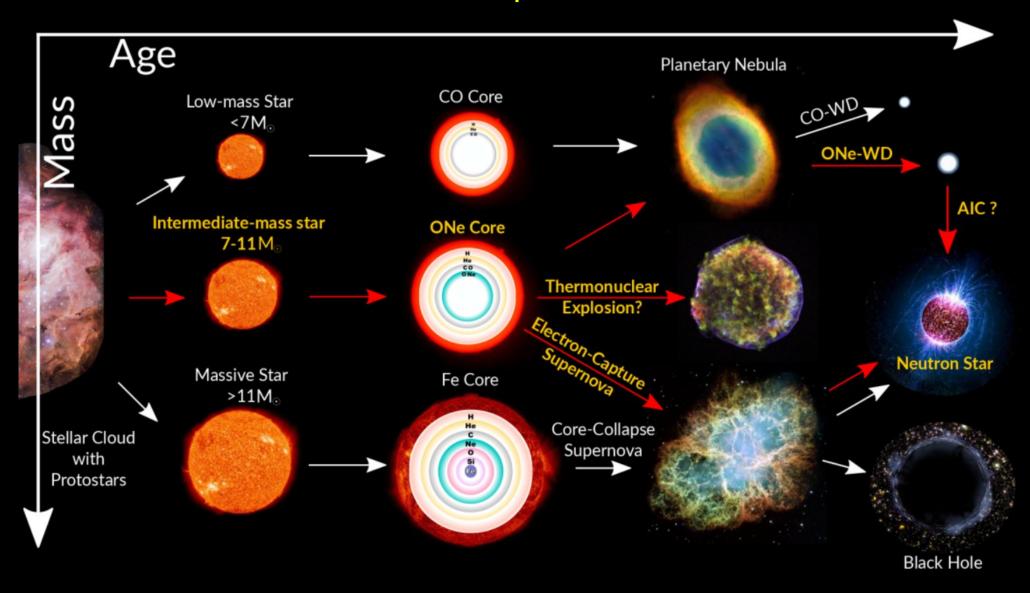
#### Some Notable Cases < 1.25 M☉

- 4U 1538--52  $(0.87 \pm 0.07 \text{ M}\odot)$
- SMC X-1  $(1.04 \pm 0.09 \text{ M}\odot)$
- Her X-1  $(1.07 \pm 0.36 \text{ M}\odot)$
- PSR J0453+1559  $(1.174 \pm 0.004 \text{ M}\odot)$
- PSR J1756--2251  $(1.230 \pm 0.007 \text{ M}\odot)$
- PSR J0737--3039B  $(1.249 \pm 0.0007)$  M $\odot$  )
- PSR J1946+2052 ~ 1.2 M<sub>O</sub>

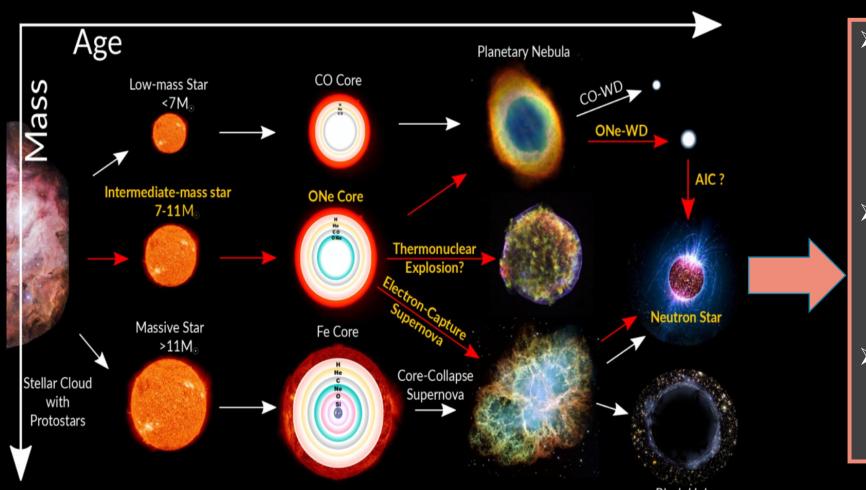
What is the origin of such low mass NSs?

https://stellarcollapse.org/index.php/nsmasses.html

#### Stellar Evolution: lifetime and death depends on mass

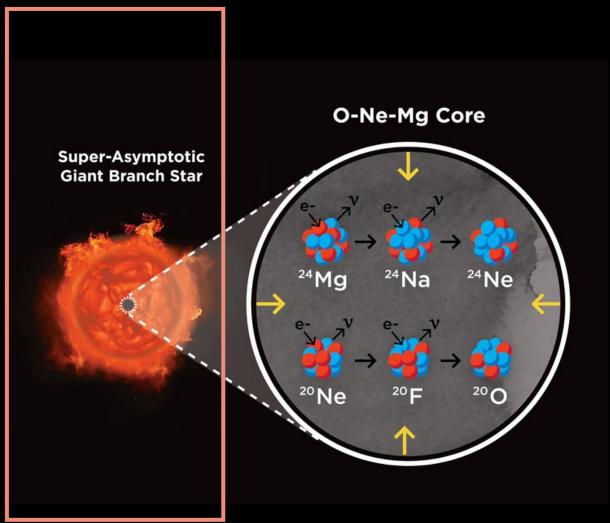


#### Stellar Evolution: lifetime and death depends on mass



- Nomoto and others developed detailed models of electron-capture supernovae, exploring how O–Ne–Mg cores collapse differently from iron-core supernovae.
- For decades, these events remained only theoretical since no unambiguous observational signature was found.
- ➤ Only recently, discoveries such as SN 2018zd provided the first clear evidence of an ECSN, confirming the long-standing predictions.

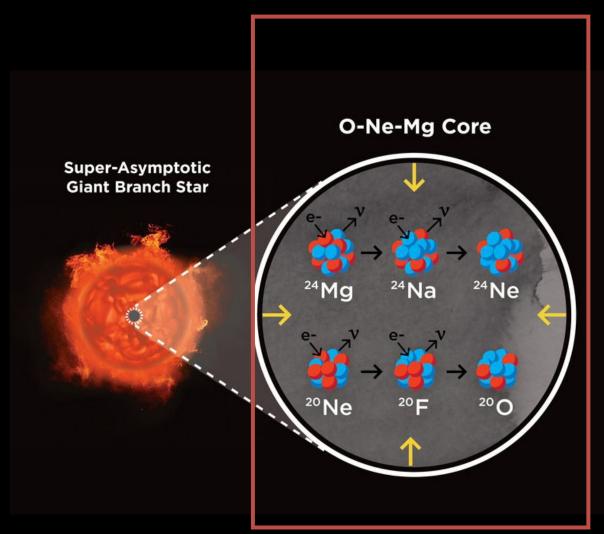
#### Electron Capture Supernovae



A super-asymptotic giant branch star is the end state of stars in a mass range of around 8-10 solar masses, whose core is pressure supported by electrons (e-)

Artist impressions of a super-asymptotic giant branch star (left) and its core (right) made up of oxygen (O), neon (Ne), and magnesium (Mg). Credit: S. Wilkinson; Las Cumbres Observatory

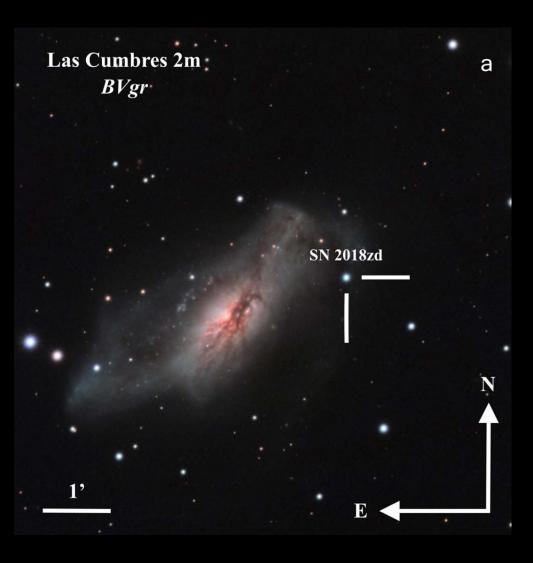
#### Electron Capture Supernovae



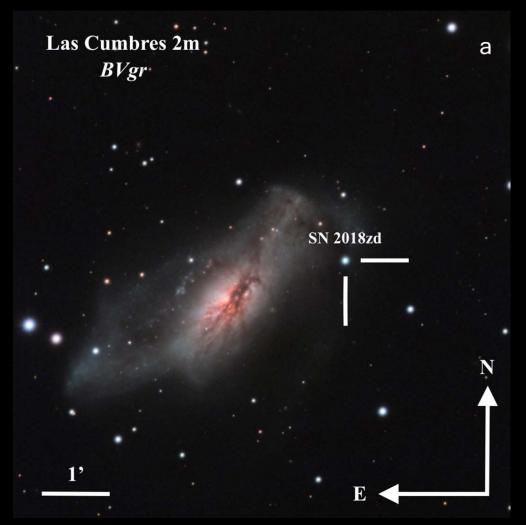
Artist impressions of a super-asymptotic giant branch star (left) and its core (right) made up of oxygen (O), neon (Ne), and magnesium (Mg). Credit: S. Wilkinson; Las Cumbres Observatory

A super-asymptotic giant branch star is the end state of stars in a mass range of around 8-10 solar masses, whose core is pressure supported by electrons (e-)

When the core becomes dense enough, neon and magnesium start to eat up electrons (so called electron-capture reactions), reducing the core pressure and inducing a core-collapse supernova



• SN 2018zd, the first strong electron-capture supernova candidate, was discovered on March 2, 2018 by Japanese amateur astronomer Koichi Itagaki using a 0.35-m telescope at his Yamagata observatory while comparing fresh CCD images of NGC 2146 with earlier references.



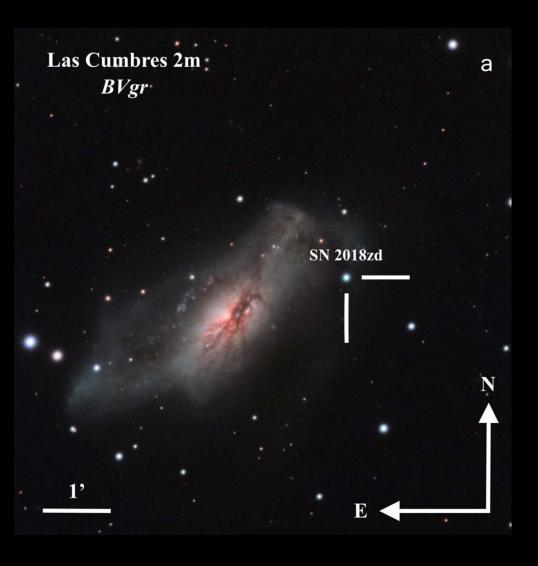
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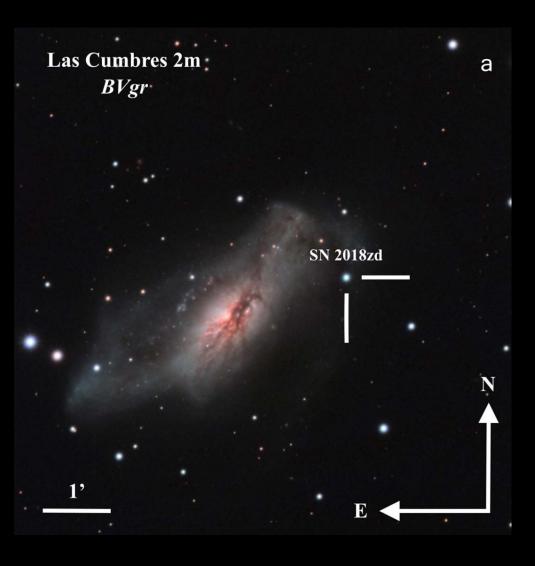


Hiramatsu, D., *et al.* The electron-capture origin of supernova 2018zd. *Nat Astron* **5**, 903–910 (2021)

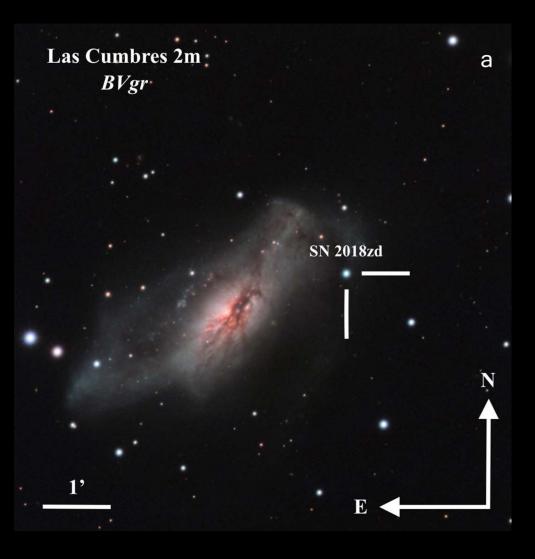
Koichi Itagaki's homespun observatory in Yamagata is one of three he maintains across Japan in hopes of spotting supernovae before anyone else. Credit: Science.



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- Major observatories (Keck, HST, Swift, Gemini, etc.) launched rapid photometric and spectroscopic campaigns, confirming unusual light-curve and spectral features.



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- An international team analyzed progenitor detections, nucleosynthetic signatures, and ejecta properties, confirming SN 2018zd as the first strong case of an electron-capture supernova.



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#### What it means for the Crab Nebula



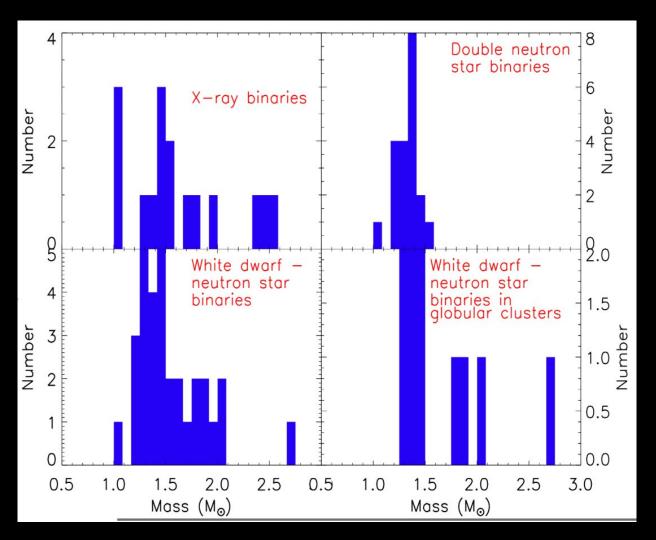
- In 1054 A.D., a supernova called SN 1054 occurred in the Milky Way. According to Chinese and Japanese records, it was so bright that it could be seen in the daytime for 23 days and at night for nearly two years.
- The resulting remnant the Crab Nebula has been studied in great detail. It was previously the best candidate for an electron-capture supernova, but this was uncertain partly because the explosion happened nearly 1,000 years ago. The new result increases the confidence that the historic SN 1054 was an electron-capture supernova.
- Interaction between the ejecta and pre-supernova mass loss explains the brightness, supporting the Crab's origin as an electron-capture supernova.

Credit: NASA, ESA, NRAO/AUI/NSF and G. Dubner (University of Buenos Aires)

# However

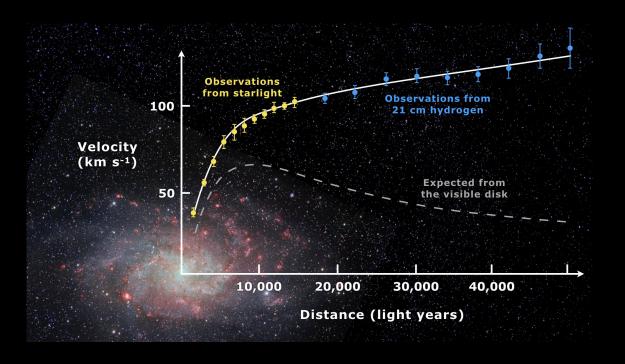
- O-Ne White Dwarfs: Theory predicts electron-capture supernovae (ECSN) from O-Ne WDs with a progenitor WD mass ≈ 1.35 M☉, collapsing to a neutron star of ≈ 1.25 M☉ (after neutrino losses).
- Expected Mass Floor: This sets a "canonical minimum" NS mass around 1.25 M☉.
- Observed Outliers: Yet, some pulsars are significantly lighter:

• Possible Explanations:???

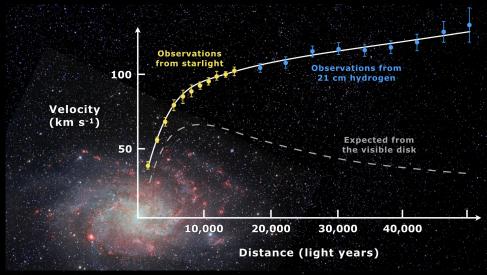


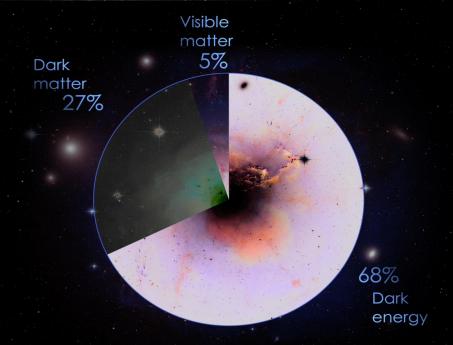
https://stellarcollapse.org

## **DARK Matter**

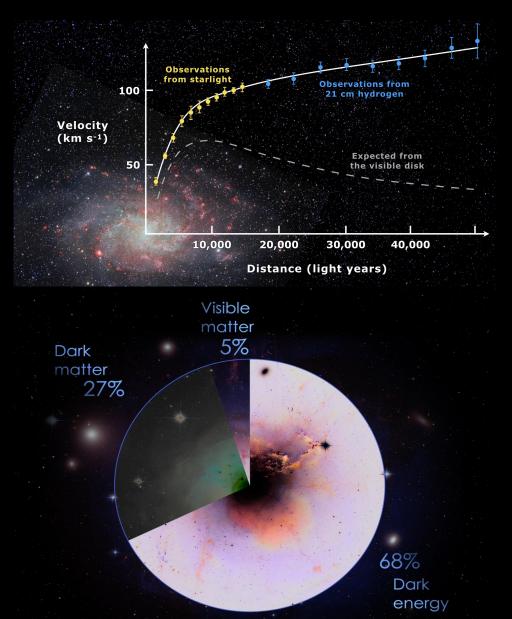


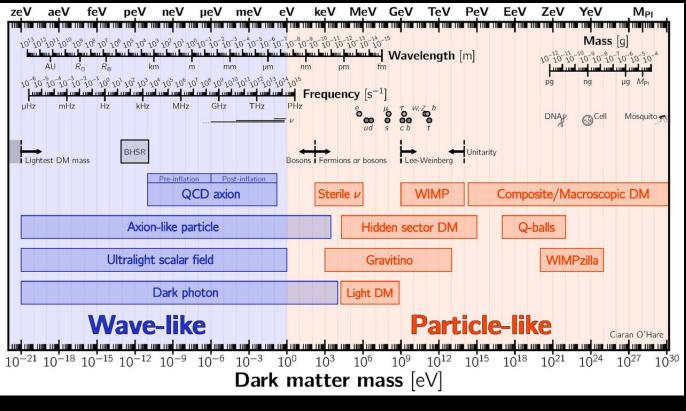
# **DARK Matter**





#### **DARK Matter**





Neutron stars and white dwarfs act as natural dark-matter detectors: capture, heating, cooling, and structural shifts let us infer DM particle properties indirectly.

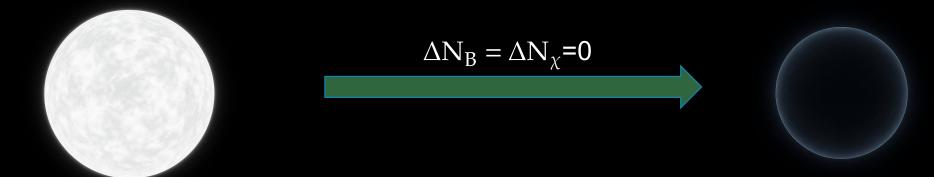
- > What the Mass of the NS from the ECSN without DM? Can DM explain the low mass NSs?
- We map the Ne WD at the EC threshold with the final Remnant NS in the presence of DM. No dynamical simulation.
- > Assume no accretion or loss and therefore, conserving Baryon number and Dark Matter number.

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Ne WD at the Electron capture threshold

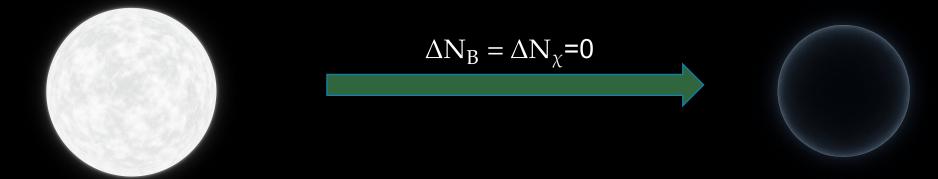
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Ne WD at the Electron capture threshold

Final Remnant NS

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Ne WD at the Electron capture threshold

Final Remnant NS

What we need? 1- WD EOS, 2- NS EOS (composition dependent), 3- DM model

In this work, we consider the case of unmagnetized and zero temperature (T = 0) stellar matter, corresponding to white dwarfs where the temperature is lower then the crystallization temperature. The total energy density of this system can be written as

$$\varepsilon = n_X M_N (A,Z) + \varepsilon_e + \varepsilon_L$$

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Energy density contribution due to ions (i.e. nuclei, since we consider fully ionized

atoms) and  $n_{\chi}$  is the number density of ions which is related to the number

density of electrons n<sub>e</sub> by the charge neutrality condition.

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$$\varepsilon = n_X M_N(A,Z) + \varepsilon_e + \varepsilon_L$$

$$\varepsilon_{e} = \frac{m_{e}c^{2}}{8\pi^{2}\lambda_{e}^{3}} \left[ x_{e} \left( 1 + 2x_{e}^{2} \right) \sqrt{1 + x_{e}^{2}} - \ln\left( x_{e} + \sqrt{1 + x_{e}^{2}} \right) \right]$$
 is the energy

density of electrons.

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$$\varepsilon = n_X M_N (A,Z) + \varepsilon_e + \varepsilon_L$$

 $\varepsilon_{\rm L} = {\rm Ce} \, 2 \, {\rm n_e}^{4/3} \, {\rm f} \, {\rm (Z)}$  is the Coulomb lattice energy, considering point like ions. C is the lattice constant.

#### Electron capture instability

• Inside the WD, at sufficiently high mass-densities, electron capture by a nucleus  ${}_{Z}^{A}X$  can become energetically favorable,

$$\begin{array}{ccc} A & & & A \\ Z X + e^- \rightarrow & & A \\ Z - 1 Y + v_e \end{array}$$

• Within the star, the pressure \$P(r)\$ must vary continuously. The Gibbs free energy per nucleon is given by:

$$g = \frac{\varepsilon + P}{n}$$

• The pressure (\$P^\*\$) at which the onset of electron capture, takes place can then be determined by the condition that the Gibbs free energy per baryon remains unchanged before and after the electron capture reaction

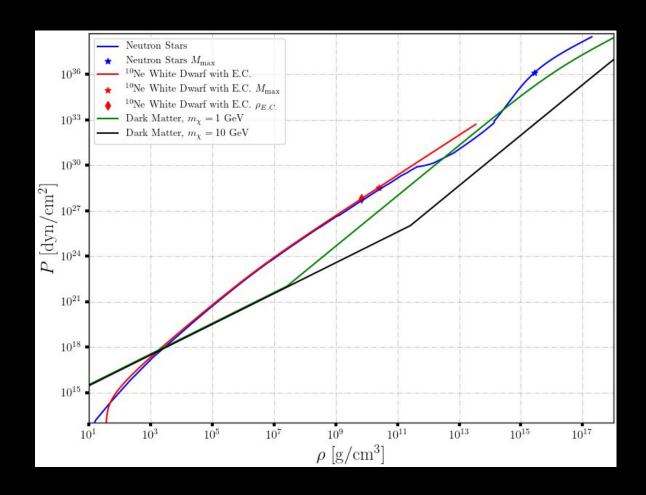
$$g(A,Z,P^*) = g(A,Z-\Delta Z,P^*)$$

#### NS and DM Model

- To model the ordinary matter (OM) fluid of a DM admixed neutron star (DANS) we use a new microscopic EOS using the Brueckner-Hartree-Fock (BHF) quantum many-body approach. (Bombaci, I., Logoteta, D., 2018).
- For DM, we consider non-selfannihilating ideal fermionic DM.
- The DM component is modeled as a cold, degenerate Fermi gas, a widely adopted approximation in both compact-object and galactic contexts.

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Kaplan et al., 2009; Zurek, 2014, Logoteta et al., 2021

## Structure equations for DM admixed compact stars

Since the non-gravitational interaction between DM and ordinary matter is negligible compared to interaction which determine ordinary matter EOSs, it is possible to write TOV equation as

$$\frac{dP_j}{dr} = -G \frac{m_{\text{tot}}(r) \,\varepsilon_j(r)}{c^2 r^2} \left(1 + \frac{P_j(r)}{\varepsilon_j(r)}\right) \left(1 + \frac{4\pi r^3 P_{\text{tot}}(r)}{c^2 m_{\text{tot}}(r)}\right) \left(1 - \frac{2G m_{\text{tot}}(r)}{c^2 r}\right)^{-1}$$

and

$$\frac{dm_j(r)}{dr} = \frac{4\pi}{c^2} r^2 \varepsilon_j(r)$$

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and

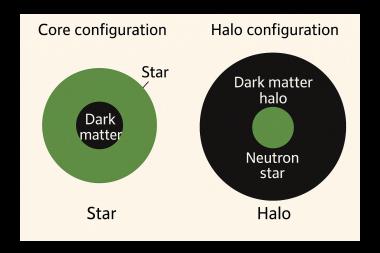
$$\frac{\mathrm{dm}_{j}(r)}{\mathrm{dr}} = \frac{4\pi}{c^{2}} r^{2} \varepsilon_{j} (r)$$

Where,

$$\triangleright$$
  $m_{tot}(r) = m_{OM}(r) + m_{DM}(r)$ 

$$P_{tot}(r) = P_{OM}(r) + P_{DM}(r)$$

- With boundary conditions  $m_j(0) = 0$ ,  $\varepsilon_j(0) = \varepsilon_{c,j}$
- The total gravitation mass is then given by  $M_j \equiv m_j (R_j) = \frac{4\pi}{c^2} \int_0^{R_j} r^2 \varepsilon_j(r) dr^{-1}$
- $\sim$  With  $M_{tot} = M_{OM} + M_{DM}$



In our scenario we assume that during the collapse of the DAWD-like core into the remnant DANS both the total baryon number and the total DM particle number are conserved. In other words, we assume that no accretion or ejection of either ordinary matter and DM components occurs during the stellar transition.

#### **DM** Number fraction

$$y_{\chi} = \frac{N_{D}}{N_{B} + N_{D}}$$

Where,

$$N_B = \int_0^{R_{OM}} n_B(r) dV$$
 
$$N_D = \int_0^{R_{DM}} n_D(r) dV$$
 
$$And dV = \left(1 - \frac{2Gm_{tot}(r)}{c^2r}\right)^{-1/2} 4\pi r^2 dr$$

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#### **DM** Mass fraction

The baryonic and dark mass of the DM admixed compact star (DACS) (DAWD or DANS) then follows as

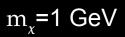
$$M_B = m_n N_B \text{ And } M_D = m_{\chi} N_D$$

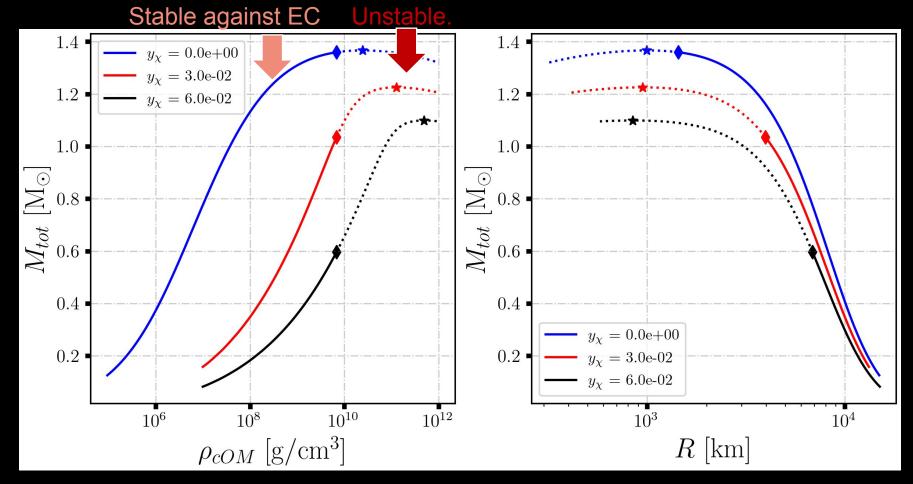
The mass fraction, generally used in literature as

$$f_{\chi} = \frac{M_{DM}}{M_{OM} + M_{DM}}$$

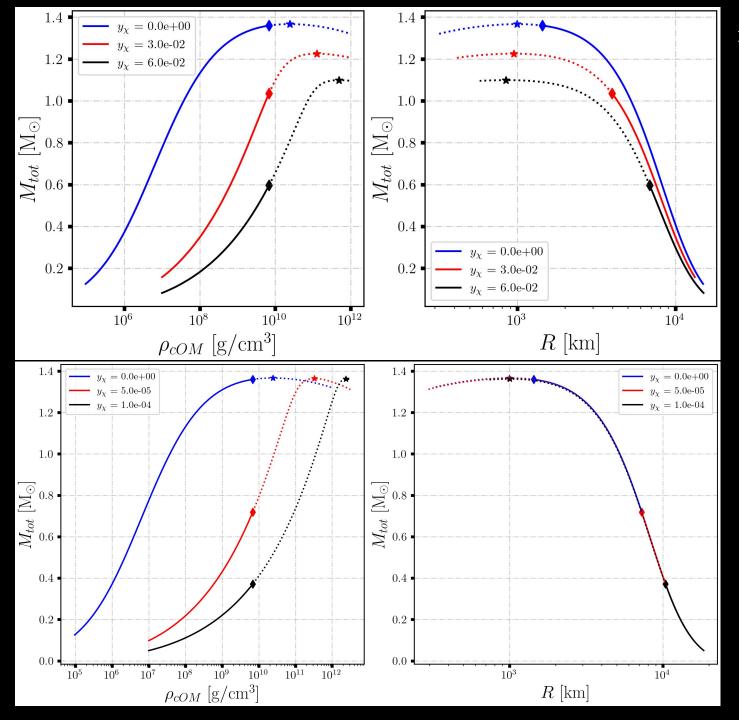
within our scenario will not be constant during the stellar process which forms the remnant DANS of an ECSN.

#### Effect of DM on degenerate white dwarfs





In the case of electron capture on 20 Ne,  $20 \text{Ne} + \text{e}^- \rightarrow 20 \text{F} + \nu_e \, \rho *$  (20 Ne) = 6.82 × 10^9 g/cm<sup>3</sup> Consequently,  $20 \text{F} + \text{e}^- \rightarrow \text{O}_{20} + \nu_e$  having  $\rho *$  (20 F) = 1.40 × 10^9 g/cm<sup>3</sup>, undergoes further EC



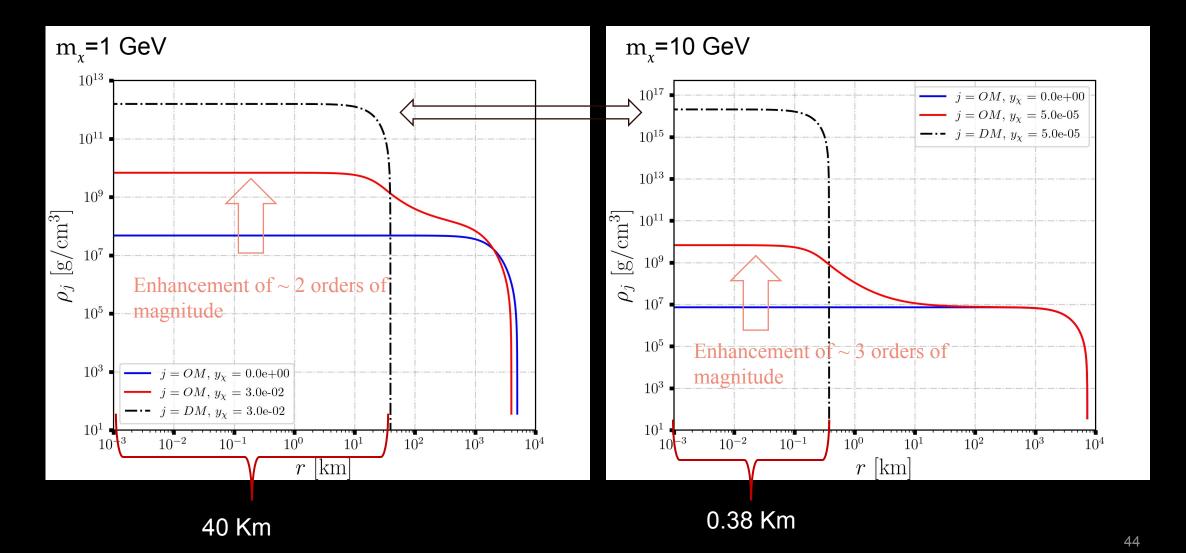
$$m_x$$
=1 GeV

Presence of DM in white dwarfs leads to a reduction in the Chandrasekhar mass, to a decrease in the corresponding ordinary matter radius, and an increase in the ordinary matter central density.

$$m_r = 10 \text{ GeV}$$

Even a tiny addition of DM leads to significant changes in the mass–central-density

# Radial mass-density profiles at the same total gravitational mass



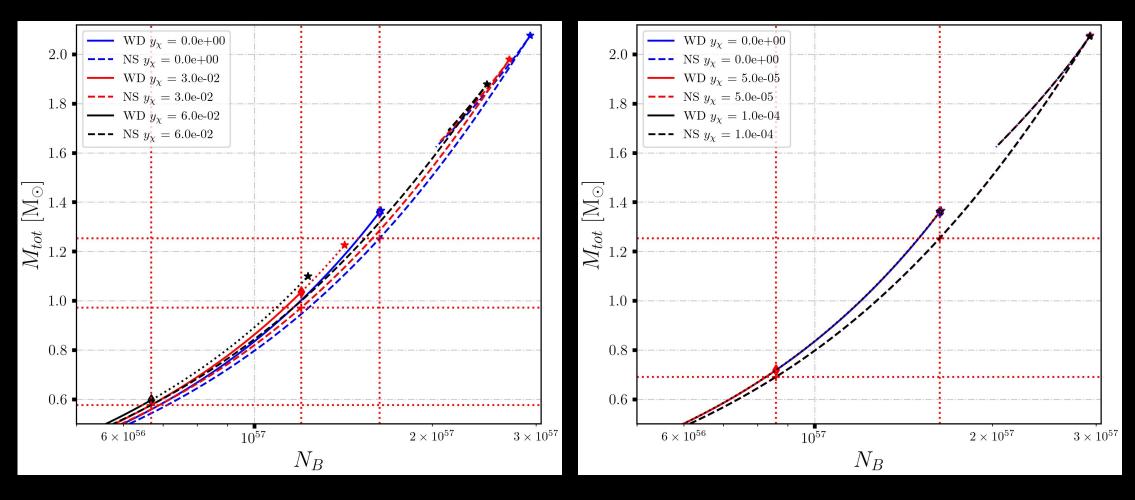
## To Quantify

			λ			
$\mathcal{Y}_{\chi}$	$M_{ch}~[{ m M}_{\odot}]$	$R_{ch}[\mathrm{km}]$	$\rho_{c_{OM}}$ [g/cm <sup>3</sup> ]	$\rho^*$ [g/cm <sup>3</sup> ]	$M^*$ [M $_{\odot}$ ]	$R^*[km]$
No DM		Increases wit	h DM mass and f	raction Decrea	ises with DM	I mass and fraction
0	1.366	998.92	$2.4201 \times 10^{10}$	$6.82 \times 10^9$	1.359	1440.7
$m_{\chi} = 1 \text{ Ge}$	V					
$3 \times 10^{-2}$	1.225	956.05	$1.2328 \times 10^{11}$	$6.82 \times 10^9$	1.0353	3970.2
$6 \times 10^{-2}$	1.028	837.10	$4.9770 \times 10^{11}$	$6.82 \times 10^9$	0.5964	6865.6
$m_{\chi} = 10 \text{ G}$	eV					
$5 \times 10^{-5}$	1.3634	1035.8	$3.1257 \times 10^{11}$	$6.82 \times 10^9$	0.7184	7272.6
$1 \times 10^{-4}$	1.3616	1005.13	$2.3919 \times 10^{12}$	$6.82 \times 10^9$	0.3707	10400.0

The EC threshold density remain same, due to no interaction of NM and OM. However, the gravitational interaction between OM and DM, leads to a redistribution of OM within the WD due to the influence of the DM. As a result, the EC process that triggers the ECSN is initiated at a lower critical stellar mass compared to white dwarfs without DM.

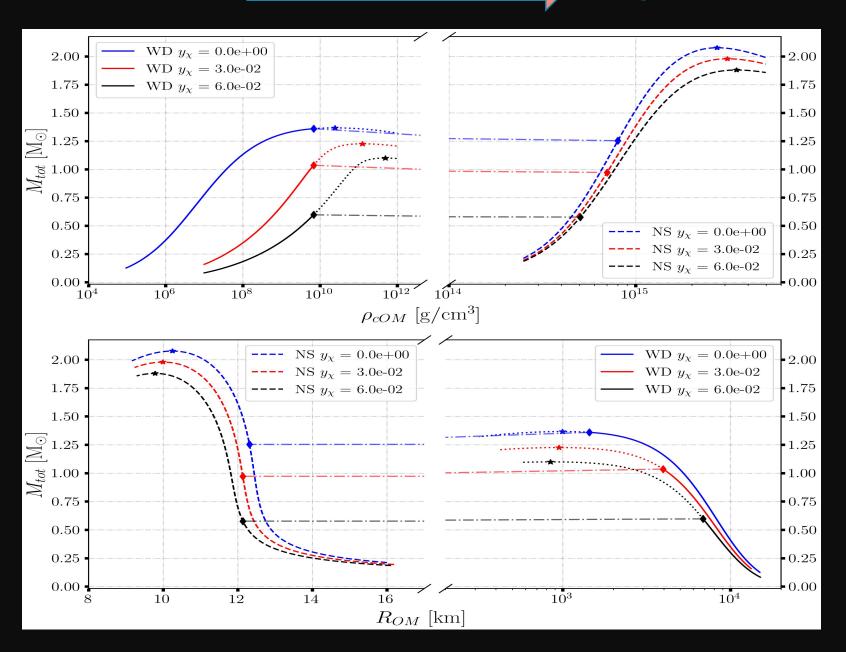
# WD to NS transition: Conserving Baryon number and Dark matter number

Bombaci and Datta, 2000





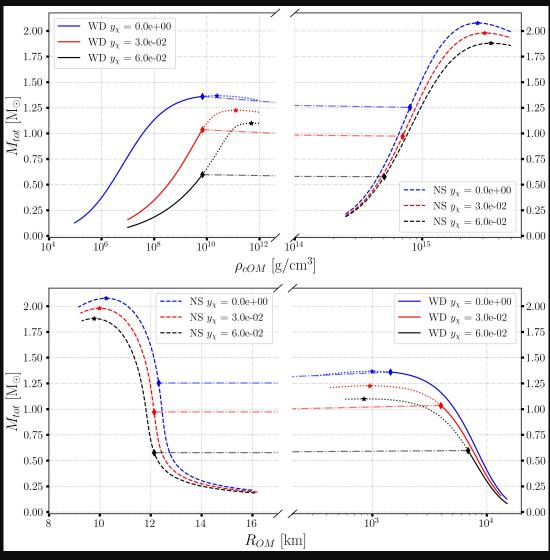
#### NS

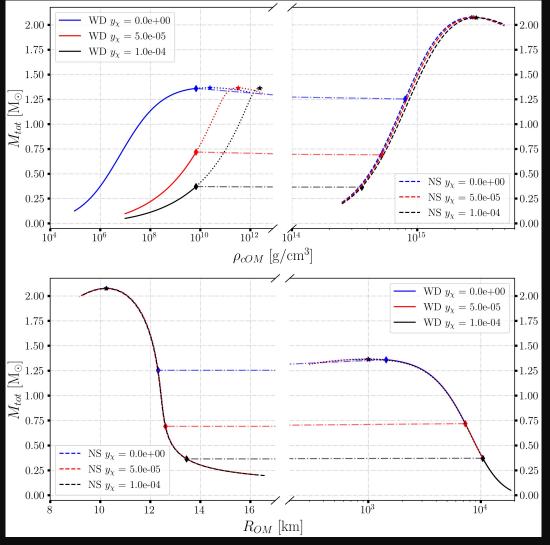


m<sub>x</sub>=1 GeV

WD to NS transition in the presence of DM

### WD to NS transition in the presence of DM





WD NS

$y_{\chi}$	$M_{ m OM}^* \ [{ m M}_{\odot}]$	R* <sub>OM</sub> [km]	$M_{ m DM}^* \ [{ m M}_{\odot}]$	R* <sub>DM</sub> [km]	<i>M</i> * [M <sub>☉</sub> ]	$M_{ m OM}^{ m NS}$ $[{ m M}_{\odot}]$	R <sub>OM</sub> [km]	$M_{ m DM}^{ m NS}$ $[{ m M}_{\odot}]$	R <sub>DM</sub> [km]	$M_{ m tot}^{ m NS}$ $[{ m M}_{\odot}]$	$E^{conv}$ [×10 <sup>53</sup> erg]
No DM											
0	1.359	1440.7			1.359	1.253	12.317			1.253	1.903
$m_{\chi} = 1 \text{ Ge}$	V										
$ 3 \times 10^{-2} \\ 6 \times 10^{-2} $	1.0020 0.5582	3970.193 6865.600	0.0333 0.0382	40.7306 38.6725	1.0353 0.5964	0.9385 0.5379	12.1383 12.1418	0.0329 0.0378	5.7617 6.6232	0.971 0.576	1.144 0.357
$m_{\chi} = 10 \text{ Ge}$	eV										
$   \begin{array}{r}     5 \times 10^{-5} \\     1 \times 10^{-4}   \end{array} $	0.7180 0.3703	7272.626 10399.74	$3.854 \times 10^{-4}$ $3.975 \times 10^{-4}$	0.3883 0.3837	0.7184 0.3707	0.6905 0.3643	12.612 13.453	$3.853 \times 10^{-4}$ $3.970 \times 10^{-4}$	0.3617 0.3666	0.692 0.365	0.472 0.102

WD •



# With No DM, ECSN can not describe NS PSR J0453+1559 with 1.17 M

$y_{\chi}$	$M_{ m OM}^*$ $[{ m M}_{\odot}]$	R* <sub>OM</sub> [km]	$M^*_{ m DM} \ [{ m M}_\odot]$	R* <sub>DM</sub> [km]	<i>M</i> * [M <sub>☉</sub> ]	$M_{ m OM}^{ m NS}$ $[{ m M}_{\odot}]$	R <sub>OM</sub> [km]	$M_{ m DM}^{ m NS} \ [{ m M}_{\odot}]$		R <sub>DM</sub> [km]	$M_{ m tot}^{ m NS}$ $[{ m M}_{\odot}]$	$E^{conv}$ [×10 <sup>53</sup> erg]
No DM												
0	1.359	1440.7			1.359	1.253	12.317		<u></u>		1.253	1.903
$m_{\chi} = 1 \text{ GeV}$	V											
$3 \times 10^{-2}$ $6 \times 10^{-2}$	1.0020 0.5582	3970.193 6865.600	0.0333 0.0382	40.7306 38.6725	1.0353 0.5964	0.9385 0.5379	12.1383 12.1418	0.0329 0.0378		5.7617 6.6232	0.971 0.576	1.144 0.357
$m_{\chi} = 10 \text{ Ge}$	eV											
$5 \times 10^{-5}$ $1 \times 10^{-4}$	0.7180 0.3703	7272.626 10399.74	$3.854 \times 10^{-4}$ $3.975 \times 10^{-4}$	0.3883 0.3837	0.7184 0.3707	0.6905 0.3643	12.612 13.453	3.853 × 3.970 ×		0.3617 0.3666	0.692 0.365	0.472 0.102

#### NS PSR J0453+1559 with 1.17 M WD $M_{ m OM}^{ m NS}$ $M_{ m DM}^{ m NS}$ $R_{\rm OM}^{\rm NS}$ $R_{ m DM}^{ m NS}$ $M_{ m tot}^{ m NS}$ $M_{\rm OM}^*$ $M_{ m DM}^*$ $E^{conv}$ $R_{\mathrm{OM}}^*$ $R_{ m DM}^*$ $M^*$ $y_{\chi}$ $[\times 10^{53} \text{ erg}]$ $[\mathrm{M}_{\odot}]$ $[\mathrm{M}_{\odot}]$ $[\mathrm{M}_{\odot}]$ $[\mathrm{M}_{\odot}]$ $[\mathrm{M}_{\odot}]$ [km] [km] [km] $[\mathrm{M}_{\odot}]$ [km] No DM 1.253 1.253 1.359 1440.7 1.359 12.317 1.903 $m_{\chi} = 1 \text{ GeV}$ $3 \times 10^{-2}$ 0.0333 0.0329 1.0020 3970.193 40.7306 1.0353 0.9385 12.1383 5.7617 0.971 1.144 $6 \times 10^{-2}$ 0.5582 6865.600 0.0382 38.6725 0.5964 0.5379 12.1418 0.0378 6.6232 0.576 0.357 $m_{\chi} = 10 \text{ GeV}$ $5 \times 10^{-5}$ $3.854 \times 10^{-4}$ $3.853 \times 10^{-4}$ 0.7180 7272.626 0.3883 0.7184 0.6905 12.612 0.3617 0.692 0.472 $1 \times 10^{-4}$ $3.975 \times 10^{-4}$ 0.3703 0.3837 13.453 $3.970 \times 10^{-4}$ 0.365 10399.74 0.3707 0.3643 0.3666 0.102

With No DM, ECSN can not describe

WD

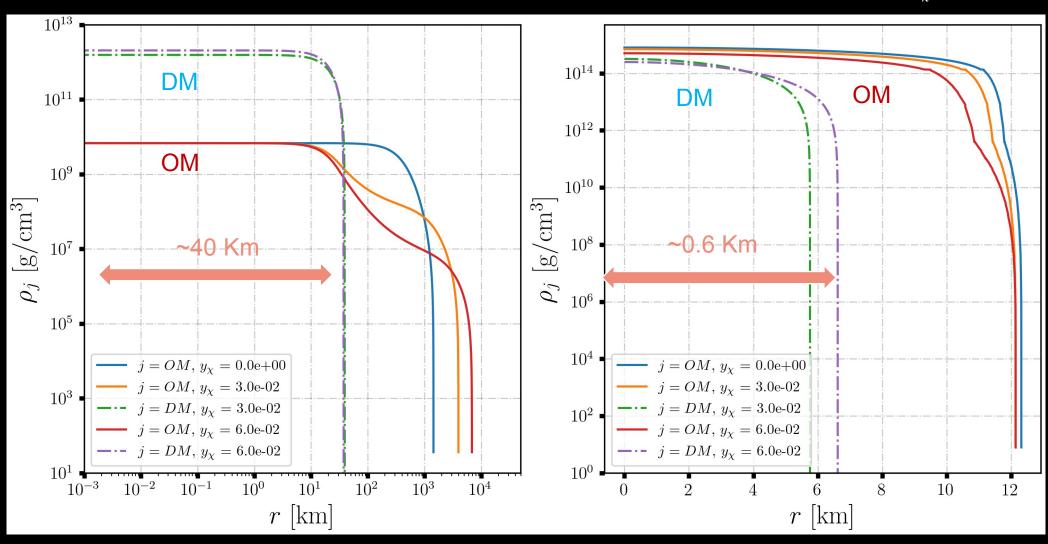
# With No DM, ECSN can not describe NS PSR J0453+1559 with 1.17 M

$y_{\chi}$	$M_{ m OM}^*$ $[{ m M}_{\odot}]$	R* [km]	$M^*_{ m DM} \ [{ m M}_\odot]$	R* <sub>DM</sub> [km]	<i>M</i> * [M <sub>☉</sub> ]	$M_{ m OM}^{ m NS}$ $[{ m M}_{\odot}]$	R <sub>OM</sub> [km]	$M_{ m DM}^{ m NS}$ $[{ m M}_{\odot}]$		$R_{ m DM}^{ m NS}$ [km]	$M_{ m tot}^{ m NS}$ $[{ m M}_{\odot}]$	$E^{conv}$ [×10 <sup>53</sup> erg	g]
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$6 \times 10^{-2}$	0.5582	6865.600	0.0382	38.6725	0.5964	0.5379	12.1418	0.0378		6.6232	0.576	0.357	
$m_{\chi} = 10 \text{ G}$	eV												
$5 \times 10^{-5}$	0.7180	7272.626	$3.854 \times 10^{-4}$	0.3883	0.7184	0.6905	12.612	3.853 ×			0.692	0.472	
$1 \times 10^{-4}$	0.3703	10399.74	$3.975 \times 10^{-4}$	0.3837	0.3707	0.3643	13.453	3.970 ×	$10^{-4}$	0.3666	0.365	0.102	L

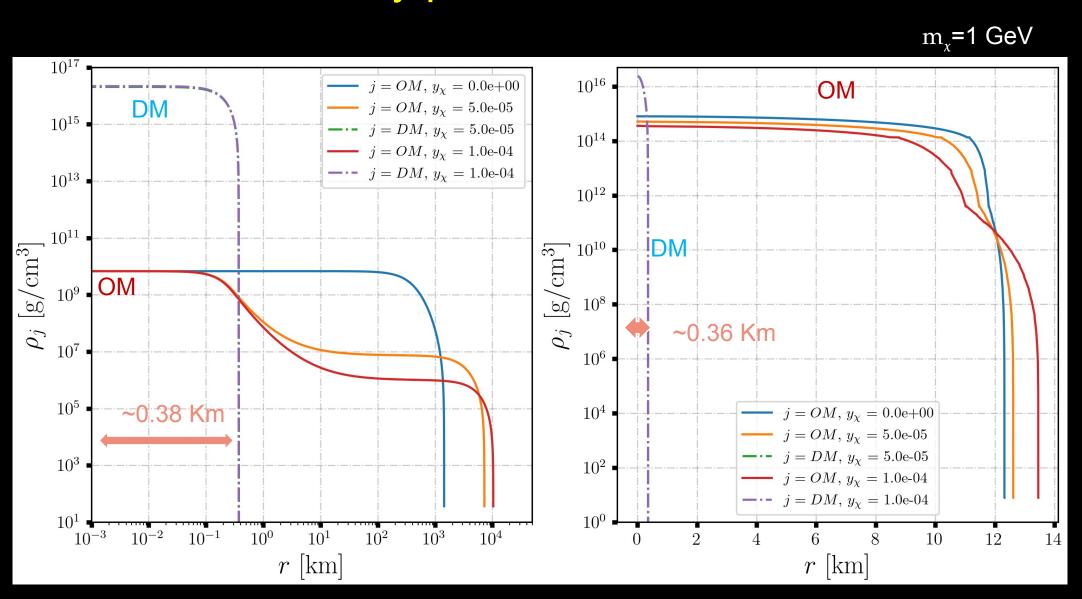
DM can reduce the NS mass considerably from ECSN. DM can results in further reduction in the energy released

## Mass-density profiles for WD and NS

 $m_x$ =1 GeV



# Mass-density profiles for WD and NS



## Conclusion

- ➤ Recent observations have established electron-capture supernovae as a real, distinct pathway to neutron-star formation, validating decades of theoretical work on collapsing O–Ne–Mg cores.
- ➤ Incorporating dark matter in the progenitor/core raises or lowers the electron-capture threshold and, at fixed baryon number, predicts systematically lower gravitational NS masses; this is directly testable with the emerging population of low NS mass systems
- A synthesis of EC-WD microphysics with multimessenger data (NICER radii, GW tidal deformabilities, precision pulsar timing) provides a coherent inference framework that can tightens constraints on the cold dense-matter EOS and crust properties.

# Future Implications and Prospective

- Can we identify candidate neutron stars that may have formed through electron-capture supernovae using our current understanding of nuclear theory and observational constraints?
- For such neutron stars, what can we infer about their internal structure, including possible signatures of dark-matter admixture or deviations from standard nuclear microphysics?
- Given a confirmed ECSN candidate, to what extent can we constrain the mass and fractional abundance of the dark-matter component responsible for modifying the collapse threshold?

# THANK YOU





#### Collaborators

Prof. Ignazio Bombaci (INFN Pisa & UNIPI)

&

Domenico Scordino (Former UNIPI student)