# FastSim dE/dx parameters

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# Introduction

- The goal is to choose suitable dE/dx parameters in FastSim to correctly match the behaviour from BAD 1500.
- We cannot match both dE/dx resolution and K/pi separation, so we focus on K/pi separation.
- Also we wish to choose suitable dE/dx parameters to implement "effective cluster counting" without modifying code.

# BaBar

- <u>BaBar NIM Paper</u>: 7.5% dE/dx resolution
- <u>Sasha Telnov's BAD 1500:</u> It's complicated, but around 7.5%
- More importantly, K/pi separation (at 3-4 GeV):



## K/pi separation in FastSim



#### Fixing K/pi separation to BAD 1500

- Recall that the dE/dx resolution is determined by a  $\sigma_0 = p_1 \left(\frac{\mu_0}{C}\right)^{p_2} L^{p_3}$ 3-parameter formula:
- We keep p2 the same, then tune p1 and p3 to match the BAD 1500 results at theta = 0.5 and 1.0 in the momentum range 3-4 GeV

 $p_1' \approx 0.00140$  $p_3' \approx -0.565$ 

#### K/pi separation with new parameters



## Angular dependence

The two points are better matched than before, and the overall shape seems to agree.



#### dE/dx resolution with new parameters

dEdx Mean, RMS, Resolution Revised Theta 0.5 Pions



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## Proposal

- Change the p1 dE/dx parameter to 0.00140 (from 0.00154)
- Change the p3 dE/dx parameter to -0.565 (from -0.34)
- These choices can be fine-tuned with more precise comparison to BAD 1500 data (here the values were taken from graphs with a ruler)

# Cluster Counting

It would be difficult to implement microscopic simulation of cluster counting in FastSim, so instead we would use "parametrized" cluster counting. A certain amount of dE/dx would correspond to a cluster

The uncertainty in counting follows simple Poisson statistics, reflected in a choice of p2 = 1/2

Cluster-counting efficiency is reflected in the choice of p1. Expected efficiency is around 60%



From Garfield in He:Isobutane 90:10

## Cluster counting parameters

From Poisson statistics, we get p2 = 0.5, p3 = -0.5

By fixing the ratio of ionization to clusters at the minimally-ionizing point, we determine p1 to be 5.81e-4

This choice of p1 would correctly model the dE/dx resolution, however the more important quantity is the K/pi separation.

We have seen that in FastSim, in order to get the correct K/pi separation, the dE/dx resolution is overstated by a factor of ~2 using the standard parameters, and slightly less using my revised parameters shown on slide 7.

Thus for cluster counting, we would multiply the p1 parameter by the same factor of 2, in order to obtain the correct K/pi separation behavior.

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## Proposal for cluster counting

- Change p1 to 0.00116 (from 0.00154), using a factor of 2 for the inflated dE/dx resolution.
- Change p2 to 1/2 (from 1).
- Change p3 to -1/2 (from -0.34).
- Note that the exact value of p1 can be tuned with better estimates of Nmin, better estimate of counting efficiency, and more precise comparison with non-cluster counting FastSim.

#### K/pi separation with cluster counting



#### dE/dx resolution with cluster counting



## Cluster counting evaluation

The cluster counting configuration can be used on its own, but the main purpose is to have a large production made of  $B \rightarrow K \nu \bar{\nu}$ with full backgrounds. The semileptonic analysis of this process can reveal the quantitative gains to be had from a clustercounting drift chamber.

Analysis code already exists for this, both from previous FPID studies and developed myself.

A more detailed proposal is in preparation.

# Backup Slides

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# SuperB DCH in FastSim

dE/dx measurements for each cell hit by a track are drawn from a Gaussian distribution with mean given by the Bethe formula and standard deviation given by:

A fraction of the individual measurements are used to form a truncated mean, which is then fed to the normal reconstruction code.

$$\mu_0 = \left[\frac{dE}{dx}\right]$$

Bethe formula result normalized to material density, per unit length

$$\sigma_0 = p_1 \left(\frac{\mu_0}{C}\right)^{p_2} L^{p_3}$$

Track segment length in cm (cell\_thickness/sin(theta) at large momentum Bethe formula result for

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# dE/dx resolution in FastSim

For a track crossing M layers, with a T truncated mean fraction, you get k = MT samples. Model a straight uniform track, then the truncated mean will also be a Gaussian variable with the same mean, and reduced standard deviation

The resolution we use is the RMS/mean of the truncated mean measurements:

Given that the second parameter is identically 1, this reduces to a constant function of the track segment length:

> Putting in the numerical factors (k = 28 = 40\*0.7, L = 1.375 cm)

 $\mu = \mu_0$ 

$$\sigma = \frac{\sigma_0}{\sqrt{k}}$$
$$\Gamma = p_1 \frac{\mu_0^{p_2 - 1}}{\sqrt{k}C^{p_2}} L^{p_3}$$

$$\Gamma = \frac{p_1 L^{p_3}}{\sqrt{kC}}$$

 $\Gamma = 0.161$ 

## Confirmation in simulation



#### Fixing K/pi separation to BAD 1500

Separation is inversely proportional to the width of the dE/dx distribution:

So we choose two new parameters to match the BaBar separation at the two  $\frac{1}{p'_1} \left(\frac{\sin \theta_1}{L}\right)^{p'_3} = \frac{1}{A} \frac{1}{p_1} \left(\frac{\sin \theta_1}{L}\right)^{p_3}$ 

The two equations uniquely determine the new parameters.

$$p_1' = Ap_1 \left(\frac{\sin\theta_1}{L}\right)^{p_3' - p_3} \approx 0.00140 \qquad p_3' = \frac{\log\frac{B}{A}}{\log\frac{\sin\theta_1}{\sin\theta_2}} + p_3 \approx -0.565$$

This dependence on L is unfortunate, as geometry changes, the parameter must be updated. Hereon I use 1.375 cm from the non-updated geometry in FastSim. 20

$$\Delta_{FS}(\theta) \propto \frac{1}{p_1} \left(\frac{\sin\theta}{L}\right)^{p_3}$$

$$\frac{1}{p_1'} \left(\frac{\sin\theta_2}{L}\right)^{p_3'} = \frac{1}{B} \frac{1}{p_1} \left(\frac{\sin\theta_2}{L}\right)^{p_3}$$

#### Deriving Parameters for Cluster Counting

$\mu_c = \begin{bmatrix} \frac{dE}{dx} \end{bmatrix}$ (1) the mean ionization per unit length normalized by gas density, given by the Bethe formula.	$C$ $p_1$ $\Box$	ionization per unit length of a minimally-ionizing particle parameter with dimension
$\mu_{c_{cell}}=\mu_{c}L$ (2) mean ionization in a cell of size L	$p_2$	of ionization density* $cm^{-p_3}$ dimensionless
$\sigma_c = p_1 \left( rac{\mu_c}{C}  ight)^{p_2} L^{p_3}$ (3) standard deviation of the io	$p_3$ nizatio	parameters n per unit length
$\sigma_{c_{cell}} = \sigma_c L$ (4) standard deviation of the total ionization in a cell of size L.	lpha	average ionization corresponding to one cluster
$\mu_{N_{cell}} = lpha \epsilon \mu_{c_{cell}}$ (5) mean number of clusters counted in a cell, proportional to the total ionization in that cell.	$\epsilon$	cluster-counting efficiency
$\sigma_{N_{cell}} = \sqrt{\mu_{N_{cell}}} \qquad \mbox{(6) standard deviation of the number of clusters counted in a cell from Poisson statistics.}$	$N_{min}$	number of clusters created per cm by a minimally-ionizing particle
$\sigma_{N_{cell}} = \alpha \epsilon \sigma_{c_{cell}} \qquad (7) \text{ also the standard deviation of the number of clusters counted in a cell, following from equation 5}$	Eq substit	uate (6) and (7), ute (2),(3), and (5)
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#### Deriving Parameters for Cluster Counting (2)

$$\begin{split} \sqrt{\alpha \epsilon \mu_c L} &= \alpha \epsilon L p_1 \left(\frac{\mu_c}{C}\right)^{p_2} L^{p_3} \\ \end{split}$$
This holds for all cell sizes L and mean ionization densities, so we immediately get:
$$p_2 &= \frac{1}{2} \\ p_3 &= -\frac{1}{2} \\ \end{cases}$$
The rest of the parameters are globbed up into p1:
$$p_1 &= \frac{C^{p_2}}{\sqrt{\alpha \epsilon}} = \sqrt{\frac{C}{\alpha \epsilon}} \\ N_{min} \text{ is the minimum number of clusters created, so} \\ N_{min} \alpha &= \frac{N_{min}}{C} \\ \end{pmatrix} p_1 = \frac{C}{\sqrt{N_{min} \epsilon}}$$

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#### Deriving Parameters for Cluster Counting (3)

Numerical values for some constants (in appropriate units):

$$C = 1.622 \times 10^{-3}$$
  

$$\epsilon \approx 0.6$$
  

$$N_{min} \approx 13$$
  

$$p_1 = 5.808 \times 10^{-4}$$

This choice of p1 would correctly model the dE/dx resolution, however the more important quantity is the K/pi separation.

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