

Partonic description of soft high-energy pp interactions

- Alternative s- and t-channel definitions of diffraction
- Partonic description of “soft” high-energy pp interactions, including diffraction, in terms of QCD/hard/BFKL-like Pomeron
- “SHRiMPS” Monte Carlo

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Diffraction 2012
Lanzarote, Sept.10-15

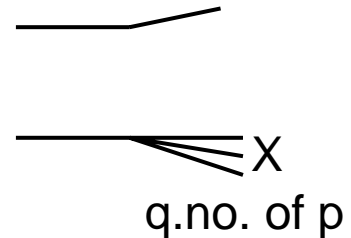
No unique definition of diffraction

1. Diffraction is elastic (or quasi-elastic) scattering caused, via **s-channel** unitarity, by the absorption of components of the wave functions of the incoming particles

e.g. $pp \rightarrow pp$,

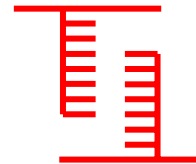
$pp \rightarrow pX$ (single proton dissociation, SD),

$pp \rightarrow XX$ (both protons dissociate, DD)



Good for quasi-elastic proc.

– but not high-mass dissociation



2. A diffractive process is characterized by a large rapidity gap (LRG), which is caused by **t-channel** “Pomeron” exchange. (or, to be more precise, by the exchange corresponding to the rightmost singularity in the complex angular momentum plane with vacuum quantum numbers).

Only good for very LRG events – otherwise

Reggeon/fluctuation contaminations

Elastic amp. $T_{el}(s,b)$

bare amp. $\Omega/2 = \text{---}$

$$\text{Im } T_{el} = \text{---} \text{---} = 1 - e^{-\Omega/2} = \sum_{n=1}^{\infty} \text{---} \text{---} \Omega/2$$

(s-ch unitarity)

proton dissociation ?

Low-mass diffractive dissociation

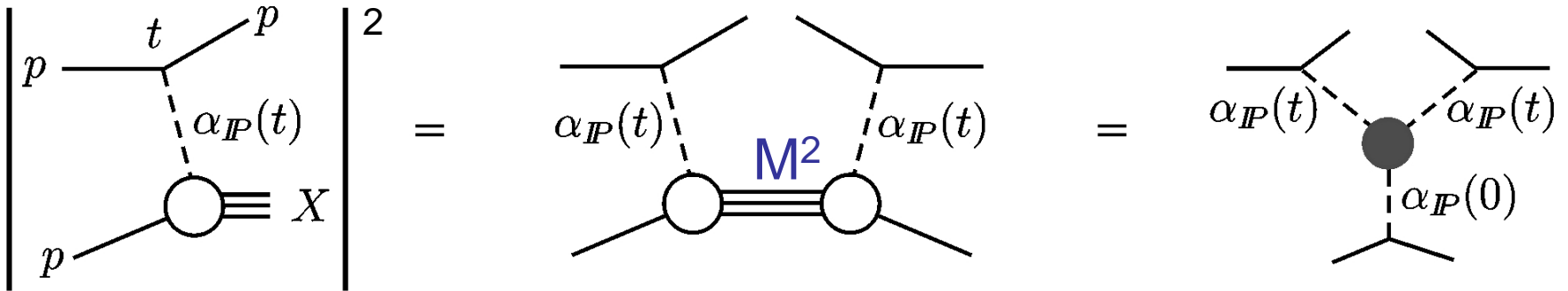
p^*
 → multichannel eikonal

introduce diff^{ve} estates ϕ_i, ϕ_k (comb^{ns} of p, p^*, \dots) which **only** undergo “elastic” scattering (Good-Walker)

$$\text{Im } T_{ik} = \text{---}^i \text{---}^k = 1 - e^{-\Omega_{ik}/2} = \sum \text{---} \text{---} \Omega_{ik}/2$$

what about high-mass diffractive dissociation ?

High-mass diffractive dissociation



described by
triple-Pomeron
diagram, plus
screening corrections

Elastic amp. $T_{el}(s,b)$

bare amp. $\Omega/2 = \overline{\quad}$

$$\text{Im } T_{el} = \overline{\text{oval}} = 1 - e^{-\Omega/2} = \sum_{n=1}^{\infty} \overline{\text{bars}} \Omega/2 \quad (-20\%)$$

(s-ch unitarity)

Low-mass diffractive dissociation

p^*
 \rightarrow multichannel eikonal

introduce diff^{ve} estates ϕ_i, ϕ_k (comb^{ns} of p, p^*, \dots) which **only** undergo “elastic” scattering (Good-Walker)

$$\text{Im } T_{ik} = \overline{\text{oval}_{ik}} = 1 - e^{-\Omega_{ik}/2} = \sum \overline{\text{bars}} \Omega_{ik}/2 \quad (-40\%)$$

include high-mass diffractive dissociation

(SD -80%)

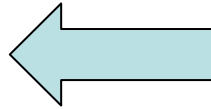
$$\Omega_{ik} = \overline{\text{bars}}_{ik} + \overline{\text{Y-shape}}_{ik} \} M + \overline{\text{Y-shape}}_{ik} + \dots + \overline{\text{Y-shape}}_{ik} + \dots$$

High-energy pp interactions

soft

hard

Reggeon Field Theory
with phenomenological
soft Pomeron



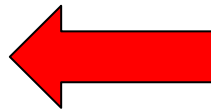
pQCD
partonic approach

smooth transition using
QCD / “BFKL” / hard Pomeron

There exists only one Pomeron, which makes
a smooth transition from the hard to the soft regime

$$\alpha_P^{\text{eff}} \sim 1.08 + 0.25 t$$

up to Tevatron energies



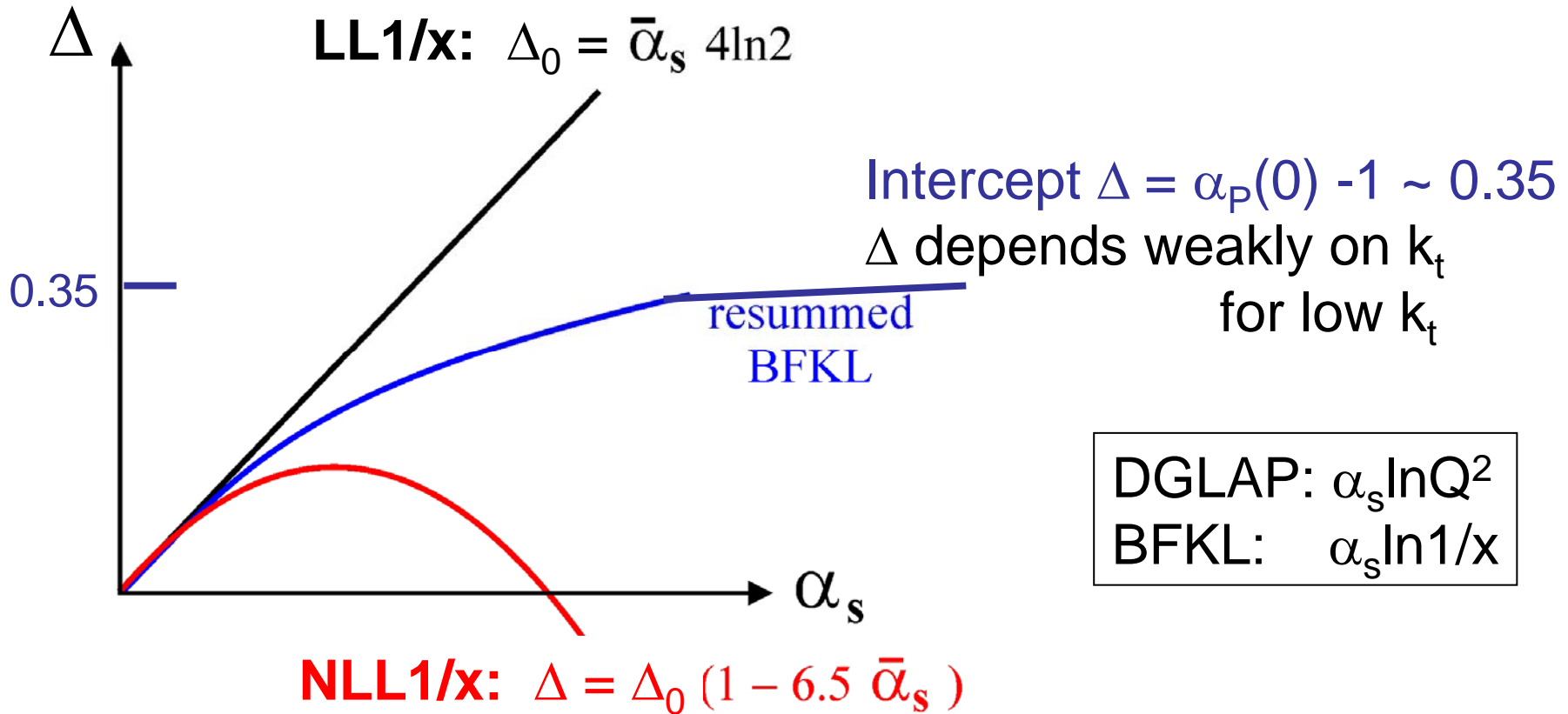
with absorptive
(multi-Pomeron) effects

$$\alpha_P^{\text{bare}} \sim 1.35 + 0 t$$

small

BFKL stabilized

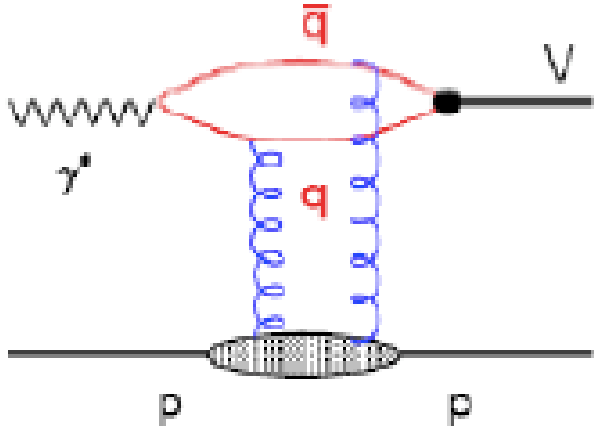
$$\Delta = \alpha_P(0) - 1$$



Small-size “BFKL” Pomeron is natural object to continue from “hard” to “soft” domain

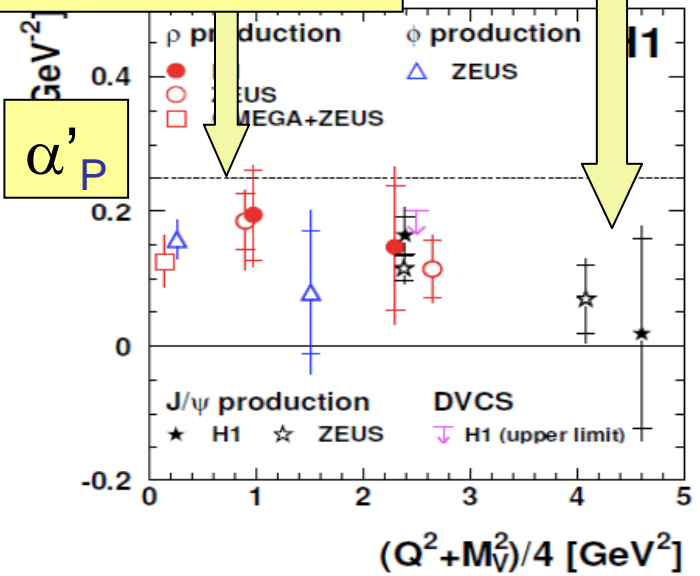
Vector meson prodⁿ at HERA
 ~ bare QCD Pom. at high Q^2
 ~ no absorption

Q^2



$\alpha'_P(0) \sim 0.25$
 after absorption

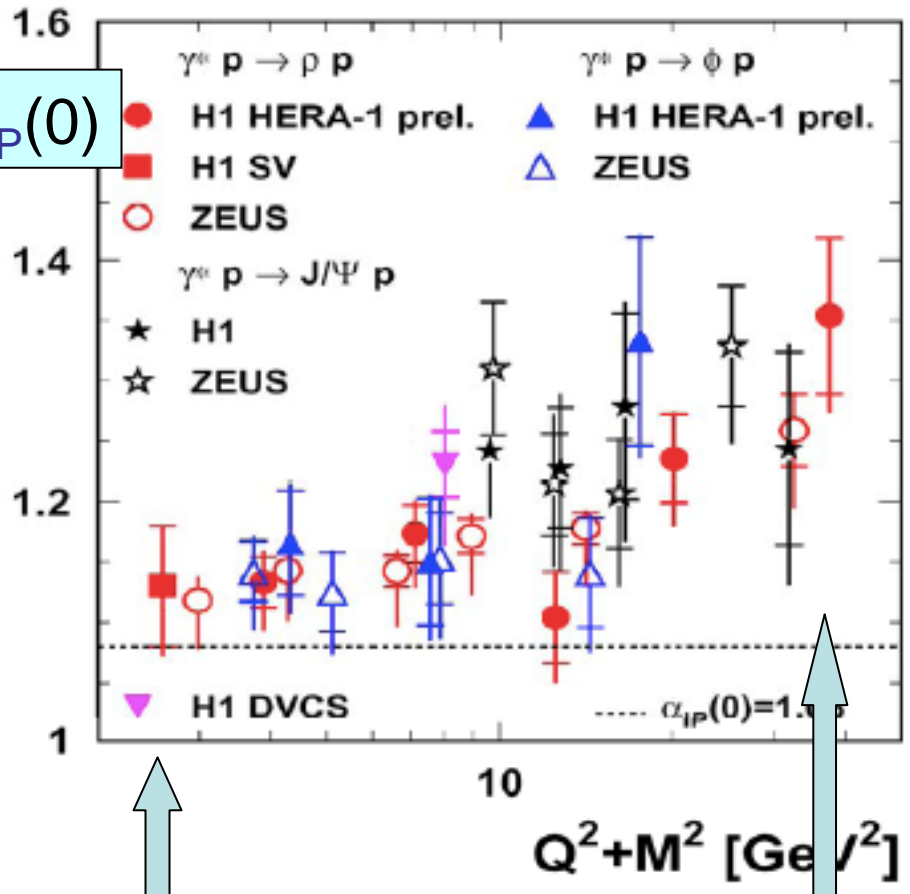
$\alpha'_P{}^{\text{bare}}(0) \sim 0$



α'_P

hard energy dependences

$\alpha_P(0)$



$\alpha_P(0) \sim 1.1$
 after absorption

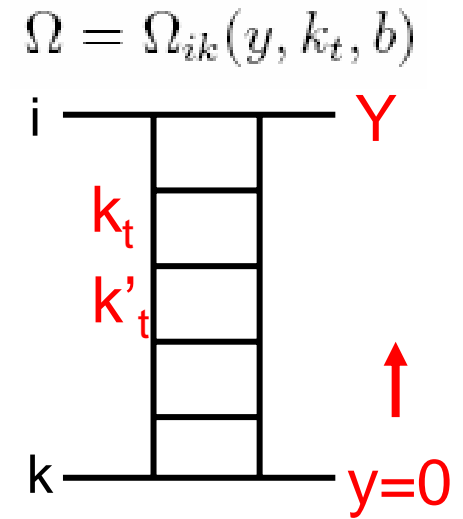
$\alpha_P{}^{\text{bare}}(0) \sim 1.35$

Partonic structure of “bare” Pomeron

BFKL evolⁿ in rapidity generates ladder

$$\frac{\partial \Omega(y, k_t)}{\partial y} = \bar{\alpha}_s \int d^2 k'_t K(k_t, k'_t) \Omega(y, k'_t)$$

- At each step k_t and b of parton can be changed – so, in principle, we have 3-variable integro-diff. eq. to solve Khoze, Martin, Ryskin
- **Inclusion of k_t crucial to match soft and hard domains. Moreover, embodies less screening at larger k_t .**
- KMR model uses simplified form of the kernel K with the main features of BFKL – diffusion in $\log k_t^2$, $\Delta = \alpha_P(0) - 1 \sim 0.35$
- b dependence during the evolution is prop' to the Pomeron slope α' , which is v.small ($\alpha' < 0.05 \text{ GeV}^{-2}$) -- so ignore. Only b dependence comes from the starting evolⁿ distribⁿ

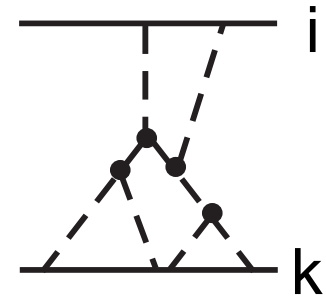


● Evolution gives \longrightarrow

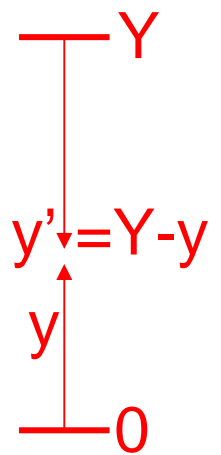
$$\Omega = \Omega_{ik}(y, k_t, b)$$

Multi-Pomeron contributions

Now include rescatt of intermediate partons with the “beam” i and “target” k



$$\left\{ \begin{array}{l} \text{evolve up from } y=0 \\ \frac{\partial \Omega_k(y)}{\partial y} = \bar{\alpha}_s \int d^2 k'_t \exp(-\lambda(\Omega_k(y) + \Omega_i(y'))/2) K(k_t, k'_t) \Omega_k(y) \\ \text{evolve down from } y'=Y-y=0 \\ \frac{\partial \Omega_i(y')}{\partial y'} = \bar{\alpha}_s \int d^2 k'_t \exp(-\lambda(\Omega_i(y') + \Omega_k(y))/2) K(k_t, k'_t) \Omega_i(y') \end{array} \right.$$



where $\lambda\Omega_{i,k}$ reflects the different opacity of protons felt by intermediate parton, rather the proton-proton opacity $\Omega_{i,k}$ $\lambda \sim 0.3$

solve iteratively for $\Omega_{ik}(y, k_t, b)$ inclusion of k_t crucial

Note: data prefer $\exp(-\lambda\Omega) \rightarrow [1 - \exp(-\lambda\Omega)] / \lambda\Omega$
 Form is consistent with generalisation of AGK cutting rules

In principle, knowledge of $\Omega_{ik}(y, k_t, b)$ allows the description of all soft, semi-hard pp high-energy data:

σ_{tot} , $d\sigma_{\text{el}}/dt$, $d\sigma_{\text{SD}}/dtdM^2$, DD, DPE...

LRG survival factors S^2

PDFs and diffractive PDFs at low x and low scales

Indeed, such a model can describe the main features of all the data, in a semi-quantitative way, with just a few physically motivated parameters:

Gotsman, Levin, Maor have similar multi-Pomeron model as KMR, except that they do not include the k_T dependence (the internal structure) of the Pomeron.

Status report on “SHRiMPS” Monte Carlo

Seek MC that describes all aspects of minimum bias
-- total, differential elastic Xsections, diffraction, jet prod...—
in a unified framework; capable of modelling exclusive
final states.

Incorporate the KMR model in SHERPA MC framework

Krauss, Hoeth, Zapp + KMR

KMR model is based on bare QCD Pomeron, with
absorptive multi-Pomeron rescattering corrections →

“SHRiMPS” MC

= **S**oft-**H**ard **R**eactions **i**nvolving **M**ulti-**P**omeron **S**catt.

Special properties of “SHRiMPS” Monte Carlo

- Based on partonic model of Pomeron, which enables **BFKL-like** structure to be continued into **soft** domain, increasingly subject to absorptive corrections
- Stronger absorption of low k_T partons **automatically** gives effective infrared **cutoff** k_{\min} which increases with collider energy
(Existing general purpose **DGLAP-based** MCs have external parameter giving an energy dependent cutoff.
“BFKL-like diffusion in $\ln k_T$ + absorption of low k_T ”
can be approximately mimicked by DGLAP)
- Consistently includes **low-mass diffraction**, via 2-channel eikonal.

Special properties of “SHRiMPS” MC continued

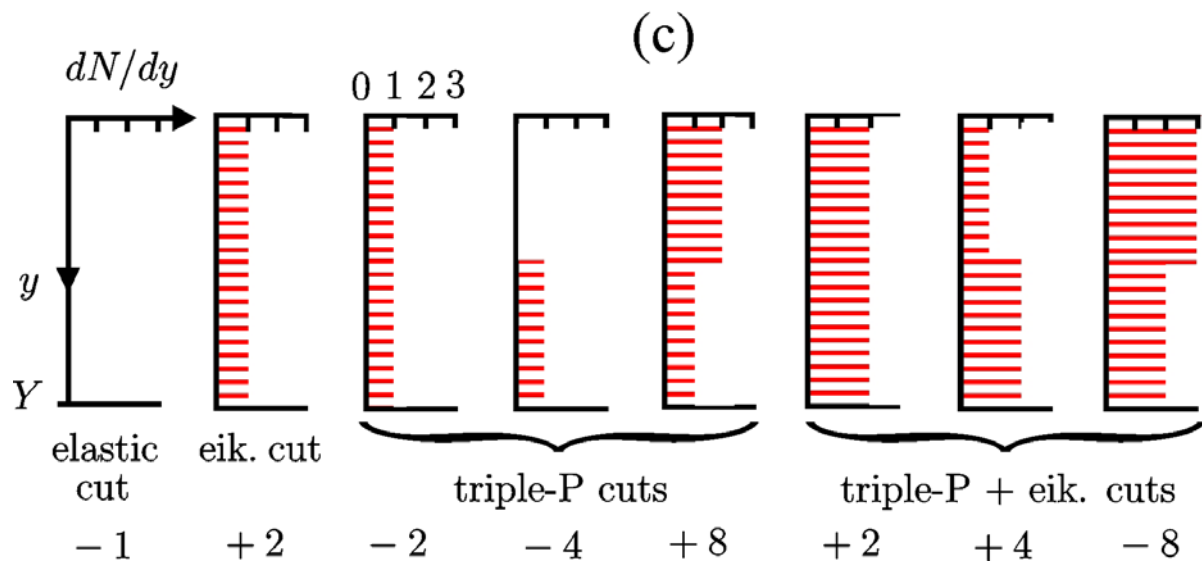
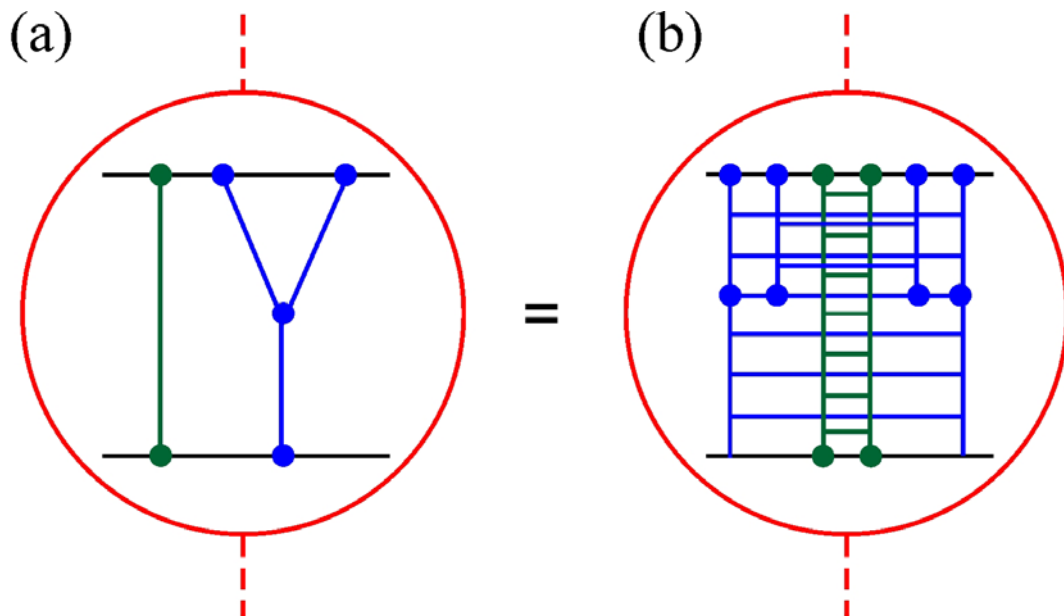
- Consistent inclusion of (absorptive) multi-Pomeron effects.

A multi-Pomeron diagram **simultaneously** describes several different processes depending on which Pomeron ladders are cut

- (i) multiparticle production results from “cut” ladders
- (ii) processes with rapidity gaps (no cut ladders in gap)

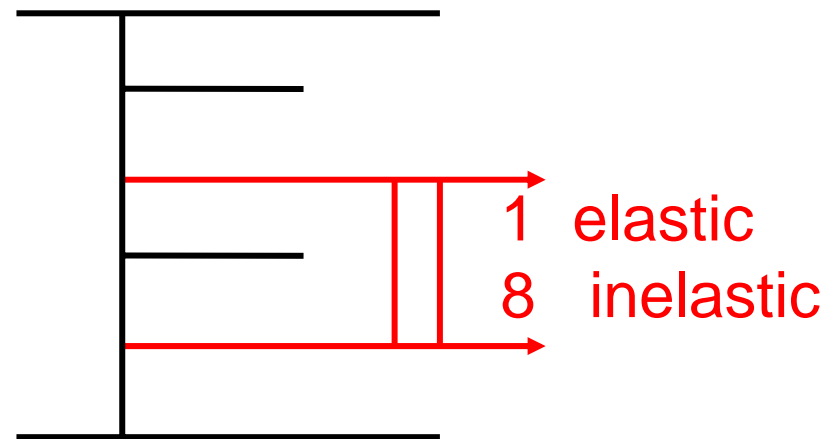
example →

Example: 8 ways to cut Pomeron ladders in this diagram



Special properties of “SHRiMPS” MC continued

- Consistent inclusion of (absorptive) multi-Pomeron effects. A multi-Pomeron diagram **simultaneously** describes several different processes depending on which Pomeron ladders are cut
 - multiparticle production results from “cut” ladders
 - processes with rapidity gaps (no cut ladders in gap)
- Moreover, account not only for multiple interactions of incoming particles, but also the possibility of **additional** (elastic and inelastic rescattering) **interactions of new partons produced by previous ladder.**



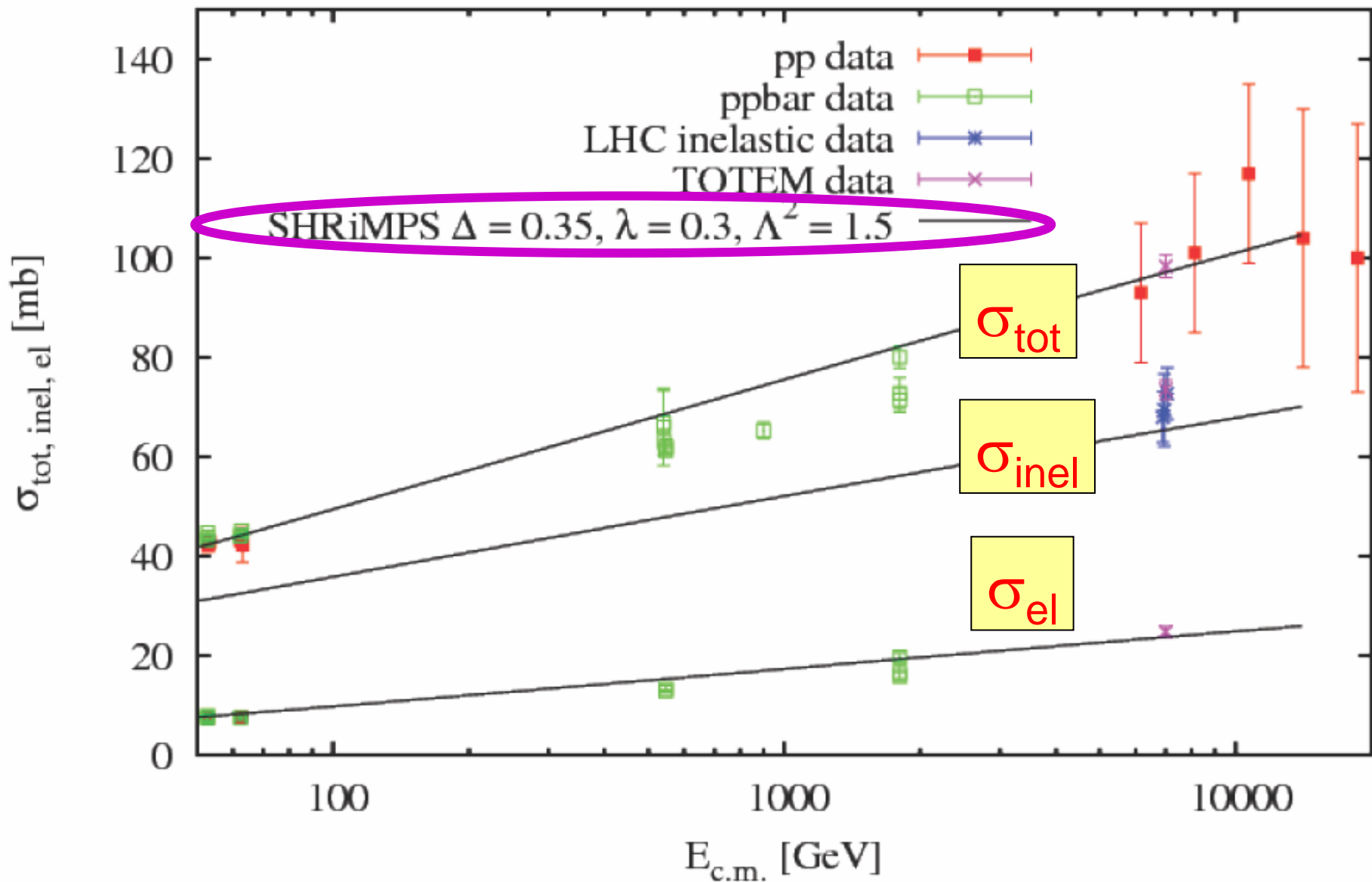
iterate until
kinematically forbidden

Special properties of “SHRiMPS” MC continued

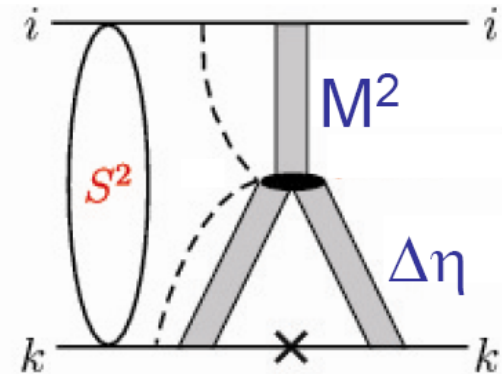
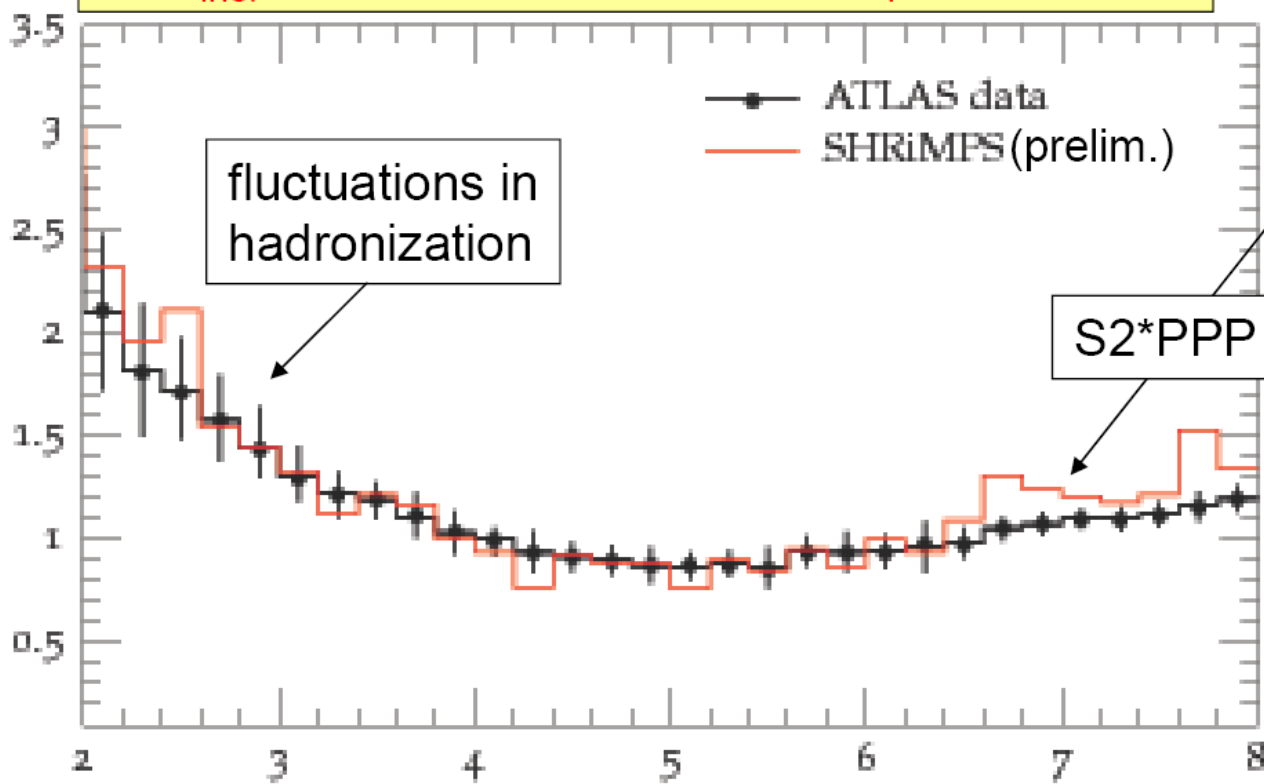
- If we have **hard interaction** then multiplicity of secondaries enhanced by strong gluon emission during DGLAP evolution up to hard scale. Inelastic scattering of these additional partons produces **new secondaries, which modify structure of underlying event**
- Finally implement parton shower, plus hadronization, plus hadron decays, plus QED
- At present, tuning MC to particle production at LHC, mainly rapidity gaps, minimum bias, and underlying event. Need correct energy dependence 900 GeV \rightarrow 7 TeV, and interface of parton shower with hard m.e.,.....

some preliminary plots \rightarrow

Total, inelastic and elastic cross section at various energies



$(d\sigma_{inel}/d(\Delta\eta))$ for particles with $p_T > 200$ MeV

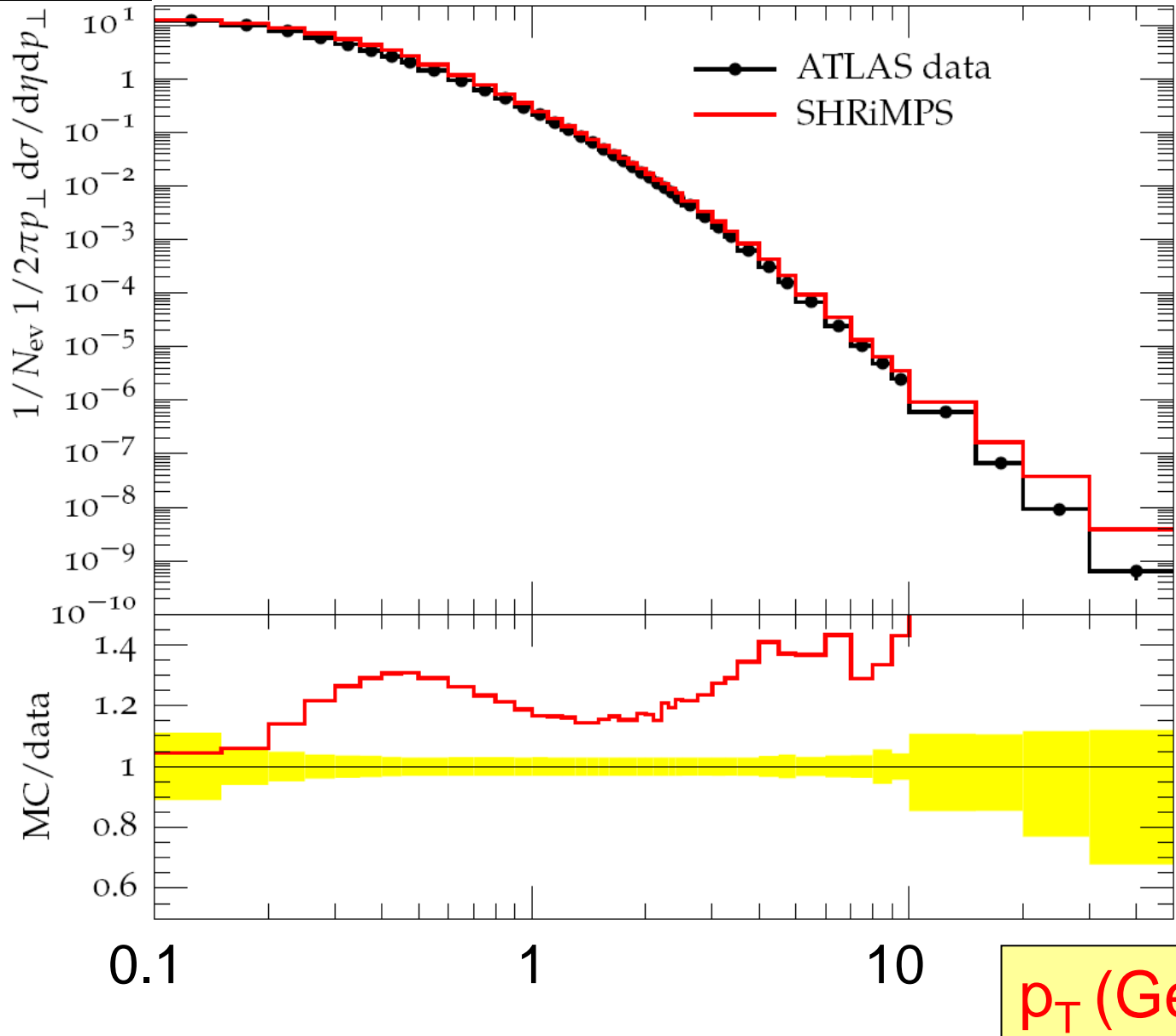


$$\Delta\eta \sim \ln(s/M^2)$$

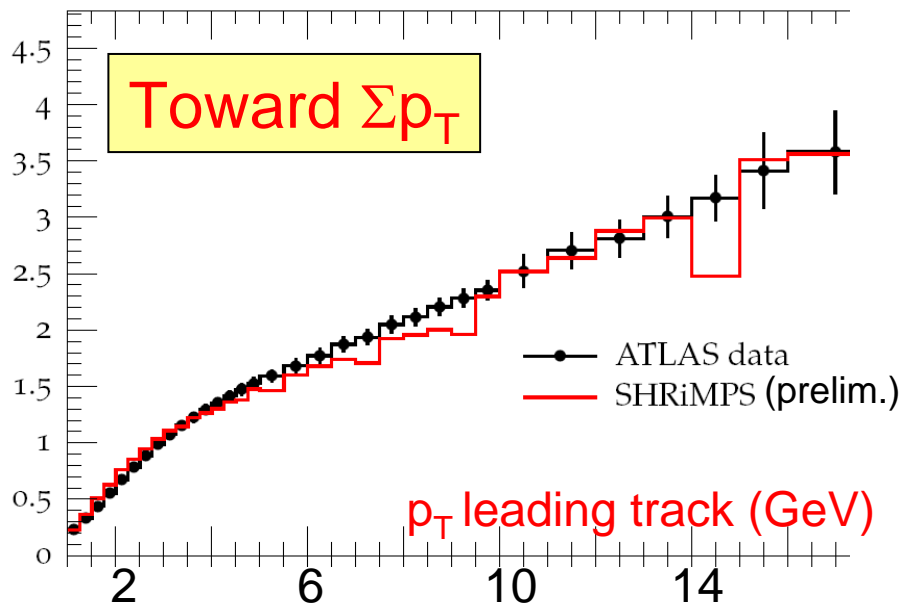
$\Delta\eta$

$d\sigma/dp_T$

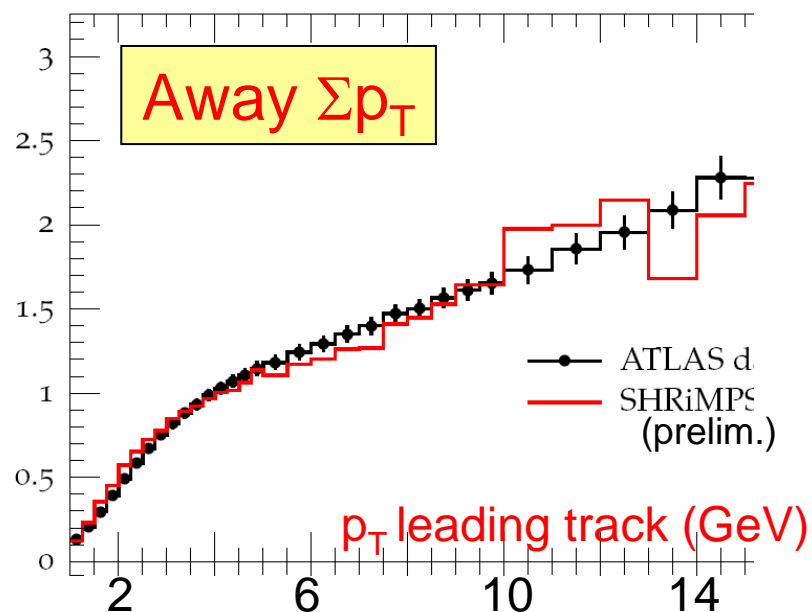
Charged particle p_{\perp} at 7 TeV, track $p_{\perp} > 100$ MeV, for $N_{\text{ch}} \geq 2$



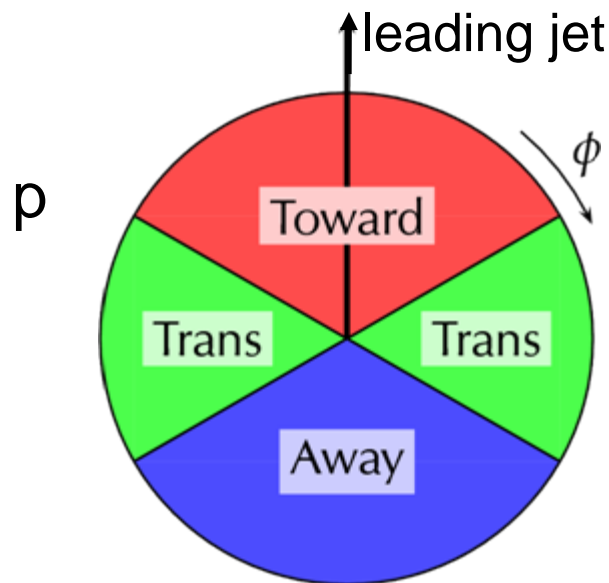
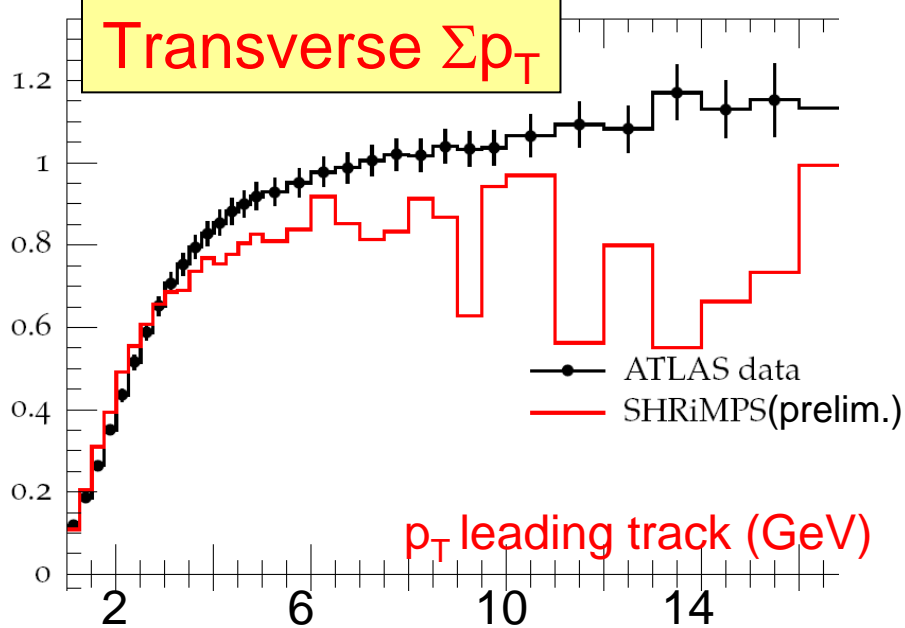
Toward Σp_{\perp} density vs. p_{\perp}^{trk1} , $\sqrt{s} = 7$ TeV



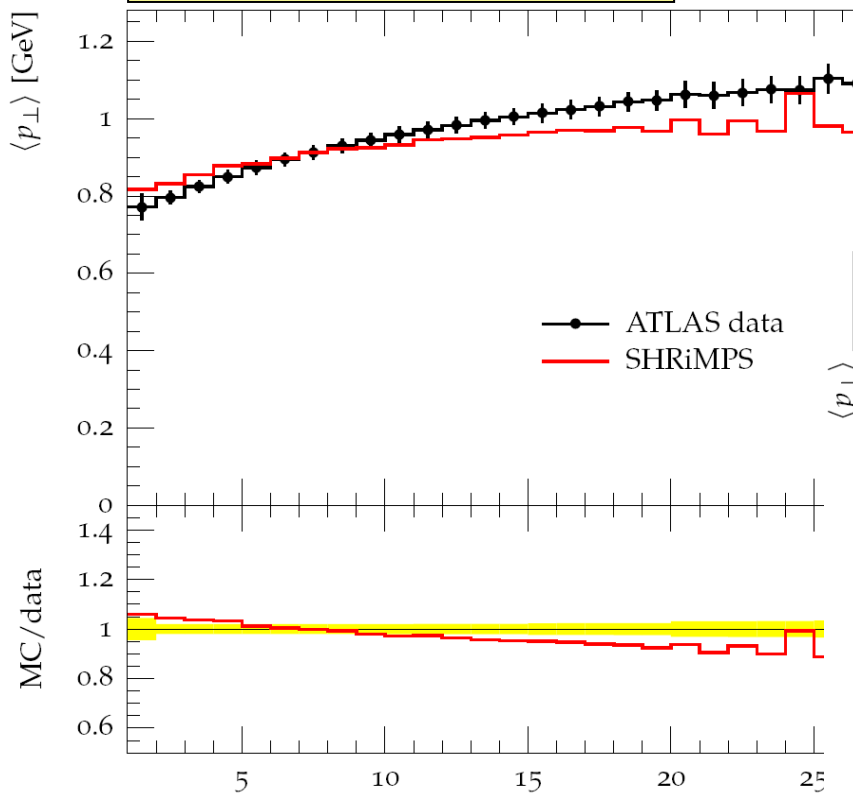
Away Σp_{\perp} density vs. p_{\perp}^{trk1} , $\sqrt{s} = 7$ TeV



Transverse Σp_{\perp} density vs. p_{\perp}^{trk1} , $\sqrt{s} = 7$ TeV

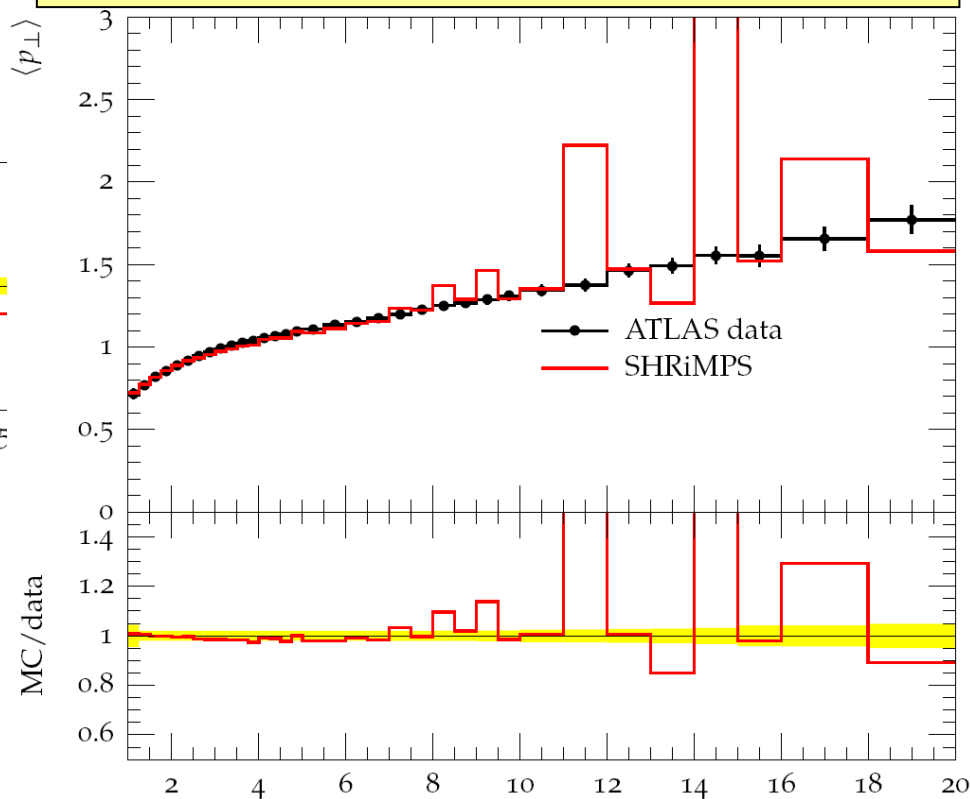


Away $\langle p_T \rangle \propto N_{ch}$



N_{ch}

Away $\langle p_T \rangle \propto p_T(\text{leading track})$



$p_T(\text{leading track})$ GeV

Conclusion

High-energy soft pp interactions may be described by the continuation of QCD/BFKL-like Pomeron into the low k_T domain, where it suffers increasingly from multi-Pomeron absorptive corrections, which automatically provides low k_T effective cutoff

Such a model forms the basis of an “all purpose” Monte Carlo -- SHRiMPS

3-ch eikonal
description of
elastic pp data

