



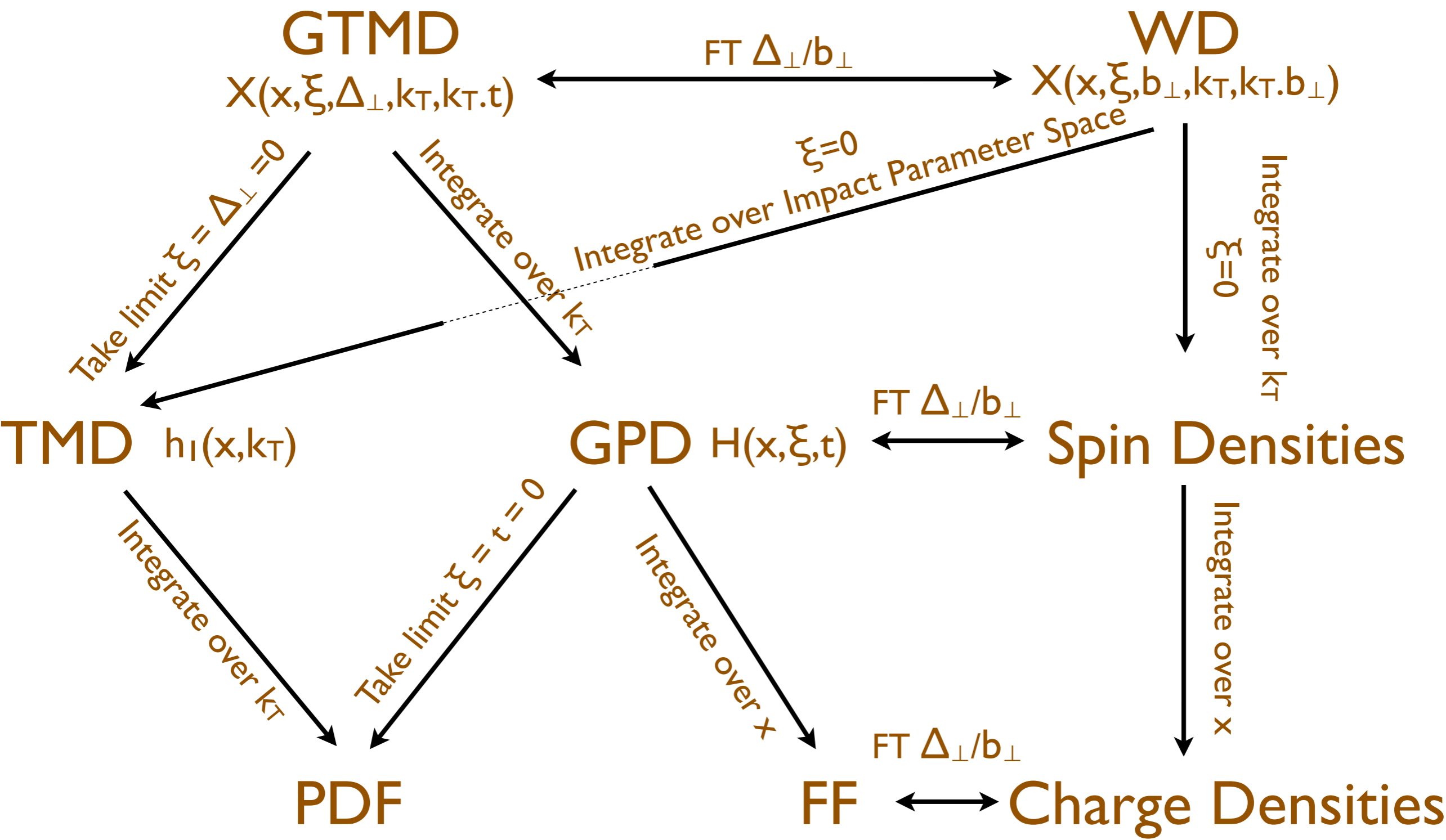
University
of Glasgow

Physics Updates from HERMES

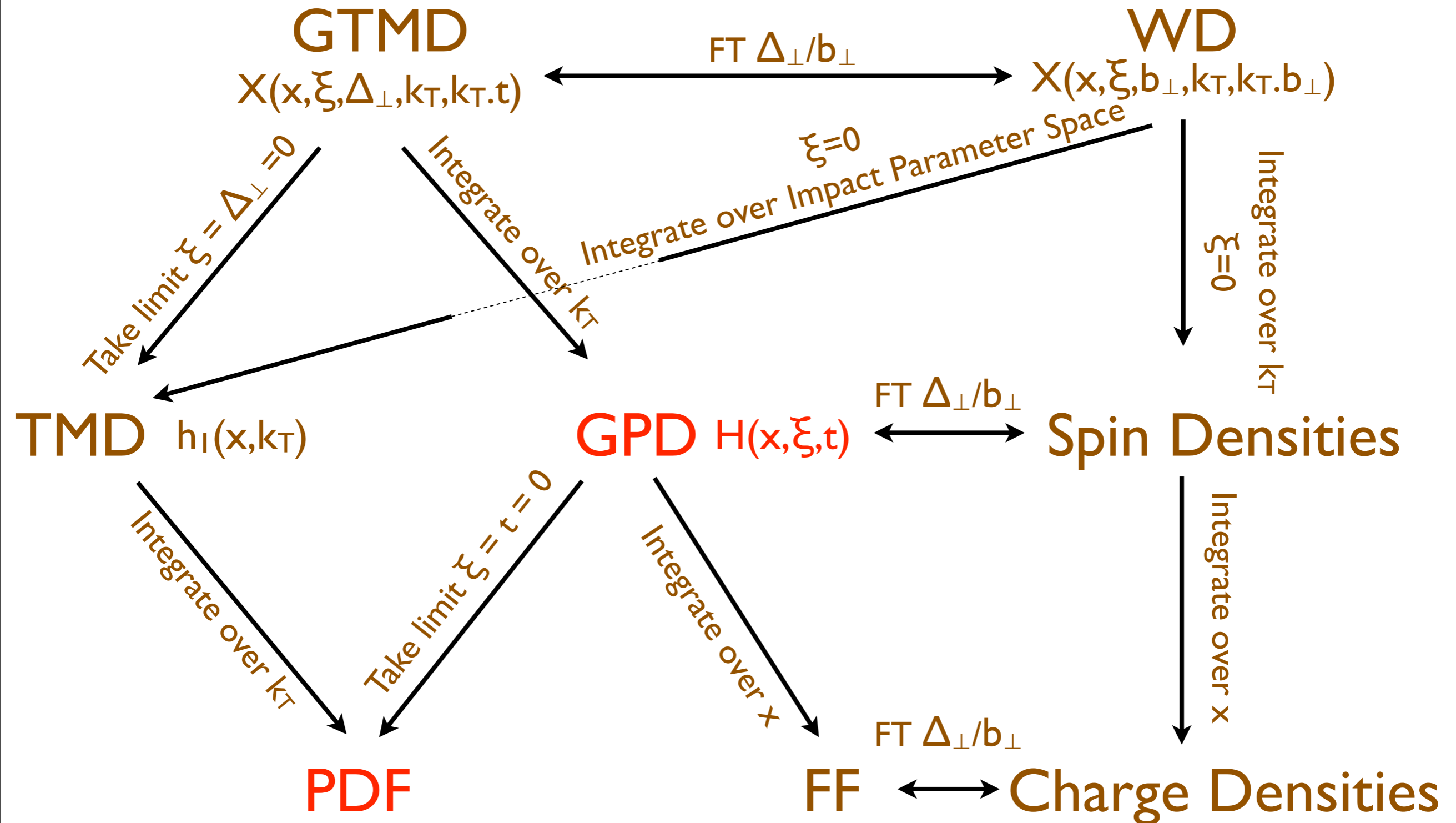
M. MURRAY, UNIVERSITY OF GLASGOW
Diffraction 2012



Distribution Graph



Distribution Graph



PDFs & Inclusive Physics

F_1 , F_2 and g_1 all comparatively well-known

New HERMES data on A_2 and g_2 available - also
measured at CERN and SLAC

$$g_2(x, Q^2) = g_s^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

PDFs & Inclusive Physics

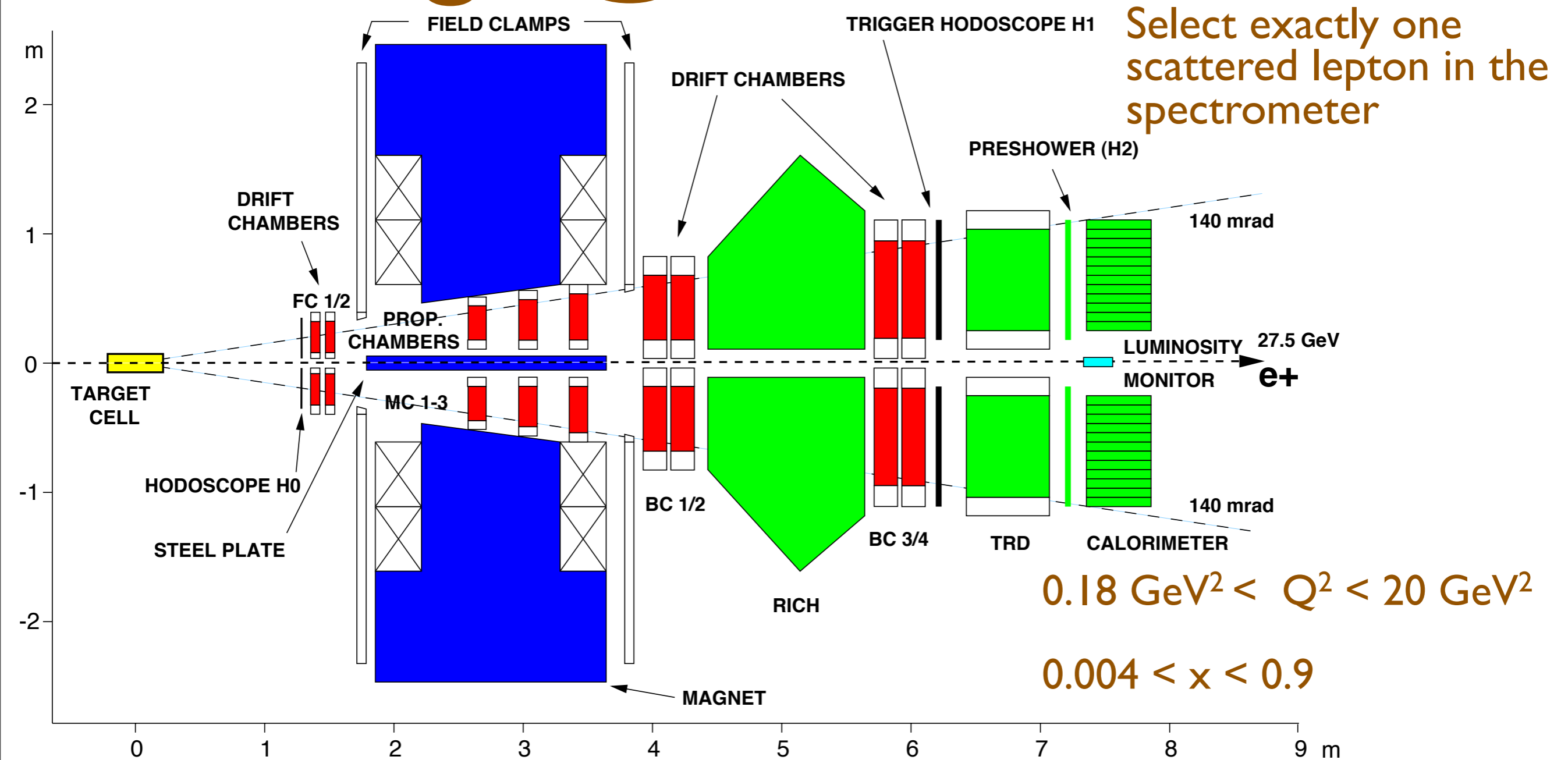
F_1 , F_2 and g_1 all comparatively well-known

New HERMES data on A_2 and g_2 available - also measured at CERN and SLAC

$$g_2(x, Q^2) = g_s^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

“pure” twist-3; related to quark-gluon correlations

g_2 @ HERMES



$$A_{LT}(x, Q^2, \phi, h_e) = h_e \frac{N^{h_e \uparrow}(x, Q^2, \phi) \mathcal{L}^{h_e \downarrow} - N^{h_e \downarrow}(x, Q^2, \phi) \mathcal{L}^{h_e \uparrow}}{N^{h_e \uparrow}(x, Q^2, \phi) \mathcal{L}_p^{h_e \downarrow} + N^{h_e \downarrow}(x, Q^2, \phi) \mathcal{L}_p^{h_e \uparrow}}$$

A_2, g_2 Extraction Procedure

Unfold in bins of (x, Q^2, ϕ)

$$A_{LT}(x, Q^2, \phi, h_\ell) = h_\ell \frac{N^{h_\ell \uparrow}(x, Q^2, \phi) \mathcal{L}^{h_\ell \downarrow} - N^{h_\ell \downarrow}(x, Q^2, \phi) \mathcal{L}^{h_\ell \uparrow}}{N^{h_\ell \uparrow}(x, Q^2, \phi) \mathcal{L}_p^{h_\ell \downarrow} + N^{h_\ell \downarrow}(x, Q^2, \phi) \mathcal{L}_p^{h_\ell \uparrow}}$$

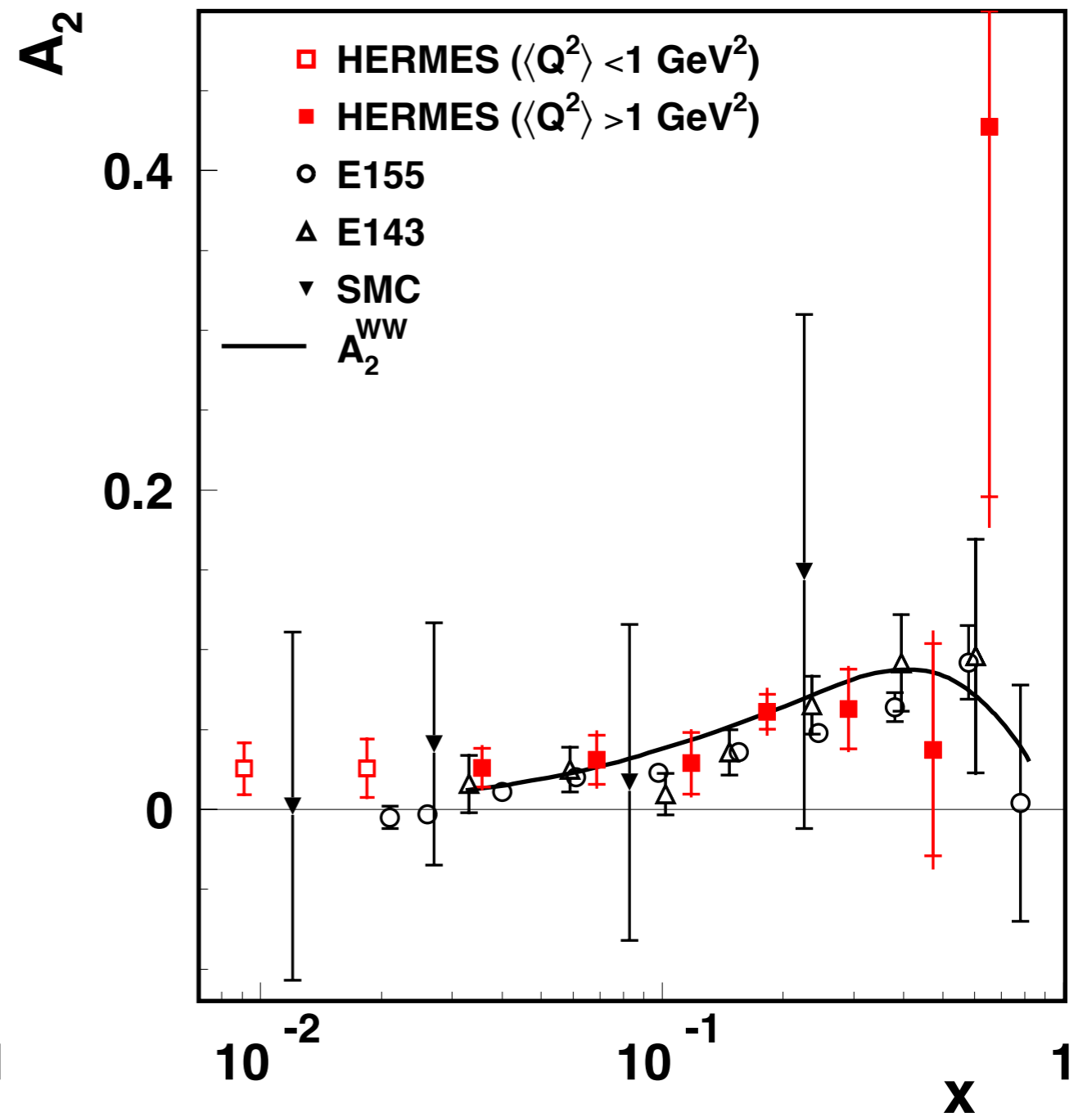
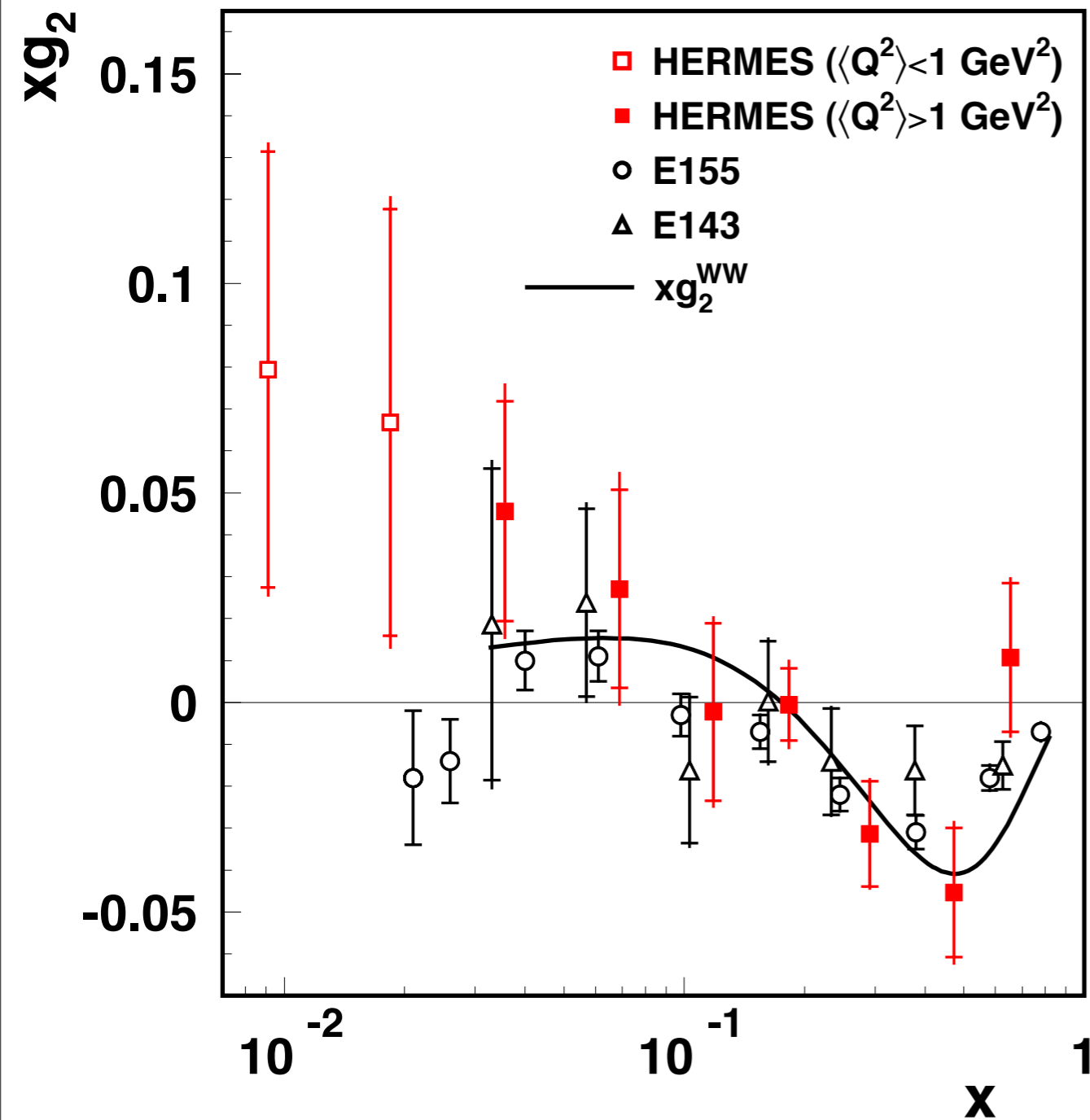
Fit result with functional form:

$$A_{LT}(x, Q^2, \phi, h_\ell) = -A_T(x, Q^2) \cos \phi$$

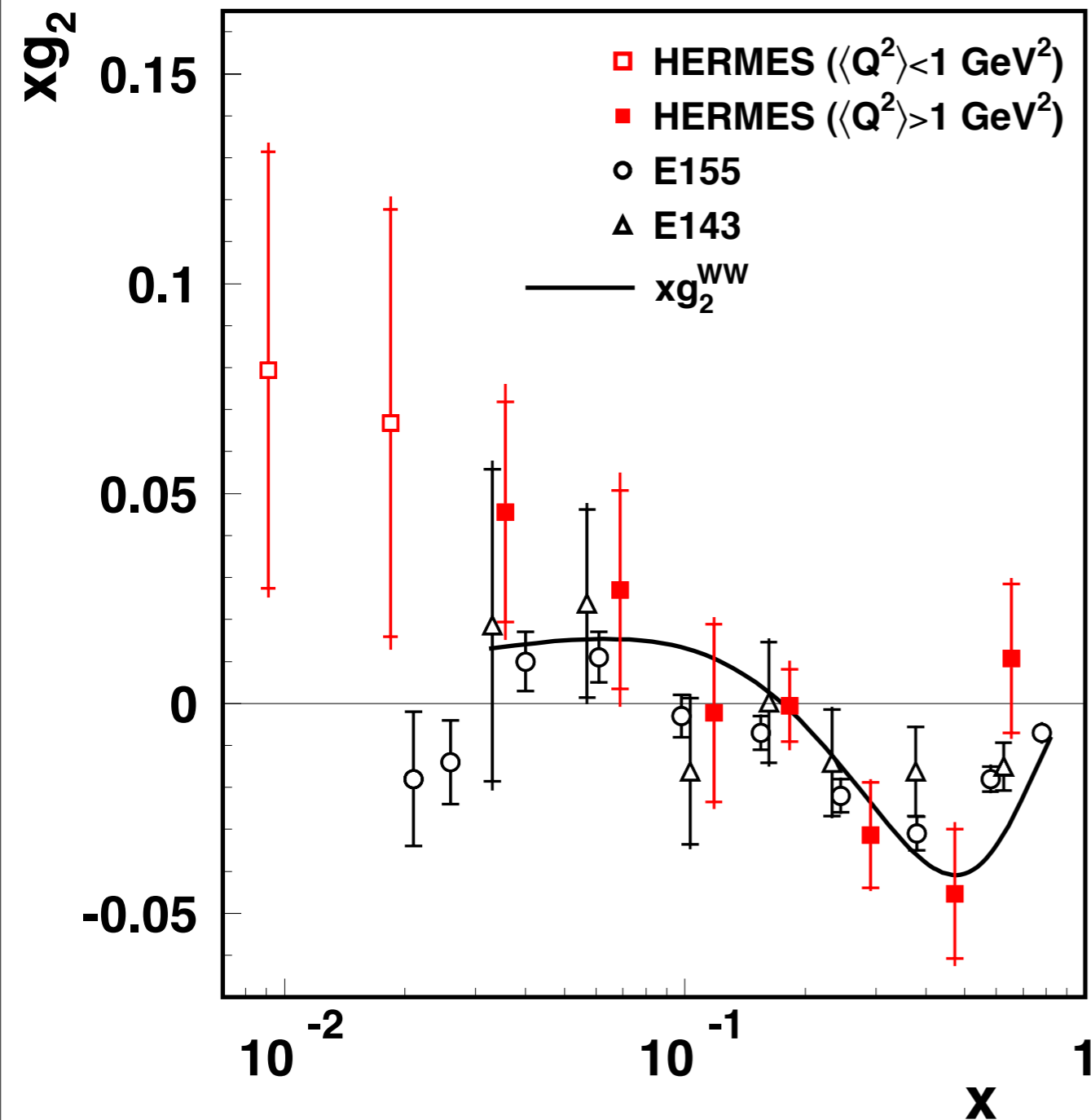
Calculate g_2 and A_2 using:

$$g_2 = \frac{F_1}{\gamma(1 + \gamma\xi)} \left(\frac{A_T}{d} - (\gamma - \xi) \frac{g_1}{F_1} \right) \quad A_2 = \frac{1}{1 + \gamma\xi} \left(\frac{A_T}{d} + \xi(1 + \gamma^2) \frac{g_1}{F_1} \right)$$

A_2, g_2 Extraction Procedure



A_2, g_2 Extraction Procedure



Results are separated into $< 1 \text{ GeV}^2$ and $> 1 \text{ GeV}^2$ series

Compared to SLAC E143 and E155 experiments

Also shown against a theoretical prediction from

[E155 Coll., P.L. Anthony et al., Phys. Lett. B 493, 19\(2000\).](#)

Consistent with CB sum-rule

A_2, g_2 Extraction Procedure

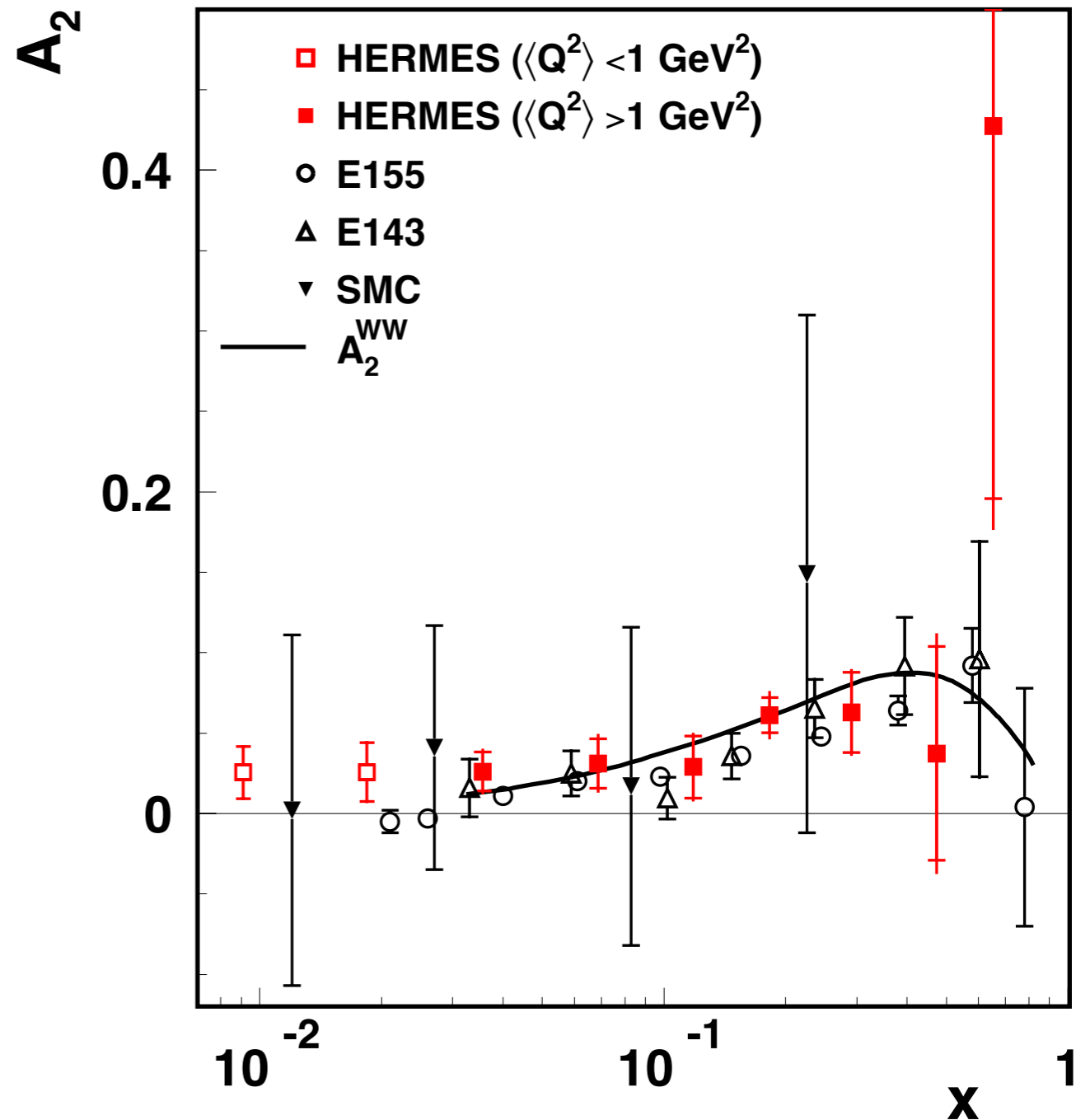
Results are separated into < 1 GeV^2 and $> 1 \text{GeV}^2$ series

Compared to SLAC E143 and E155 experiments and SMC

Also shown against a theoretical prediction from

[E155 Coll., P.L. Anthony et al., Phys. Lett. B 493, 19\(2000\).](#)

Statistical precision not enough to determine non-WW behaviour



g_2, A_2 Conclusions

A_2 and g_2 have been extracted at HERMES from the A_{LT} inclusive asymmetry.

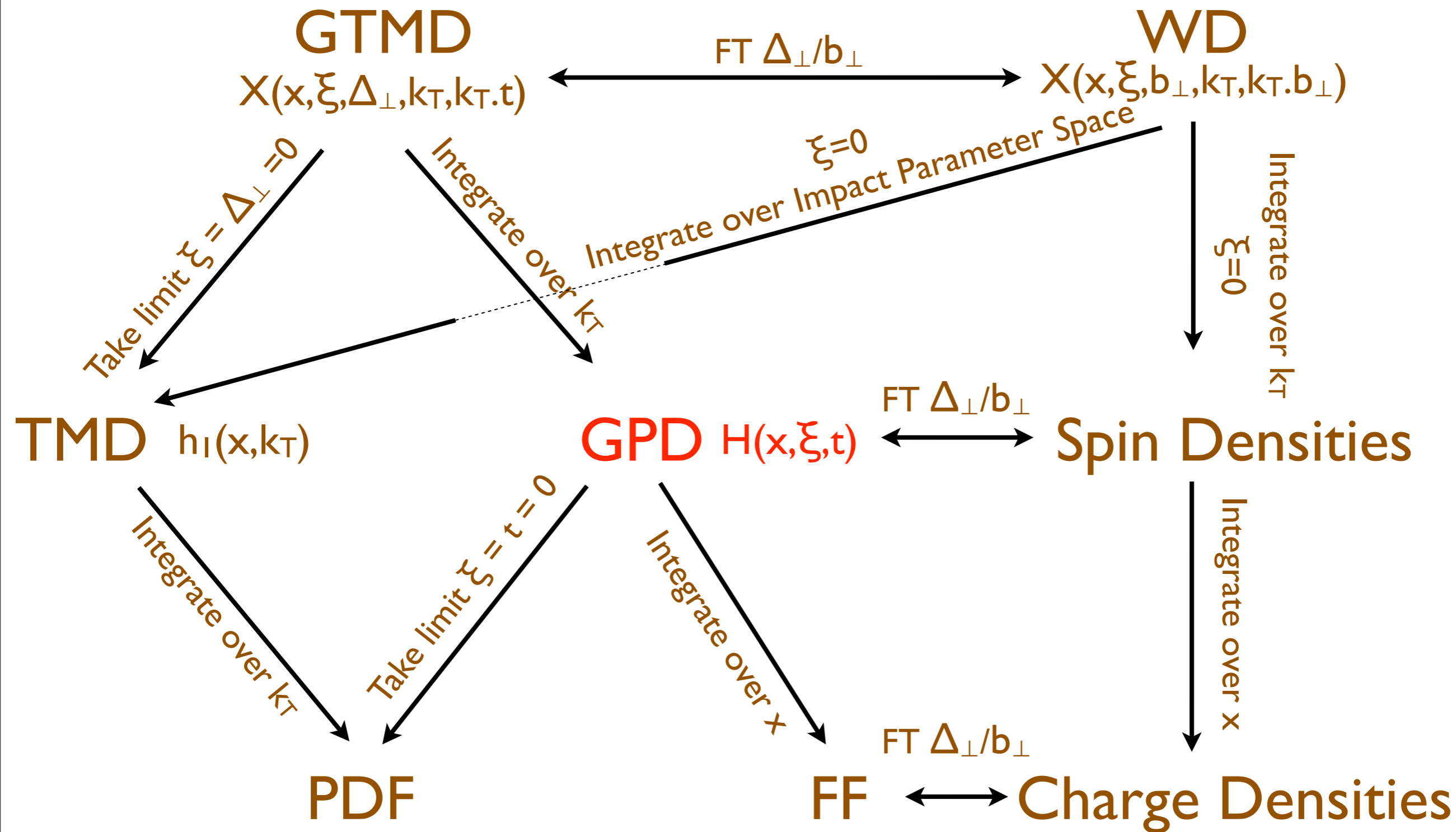
The results confirm the Burkhardt-Cottingham sum-rule for g_2 and are consistent with SLAC and CERN data (A_2 only).

Sit alongside measurements of F_2 and g_1 as contributions from HERMES to inclusive structure functions

F_2 : [A. Airapetian et al, JHEP 05 \(2011\) 126](#)

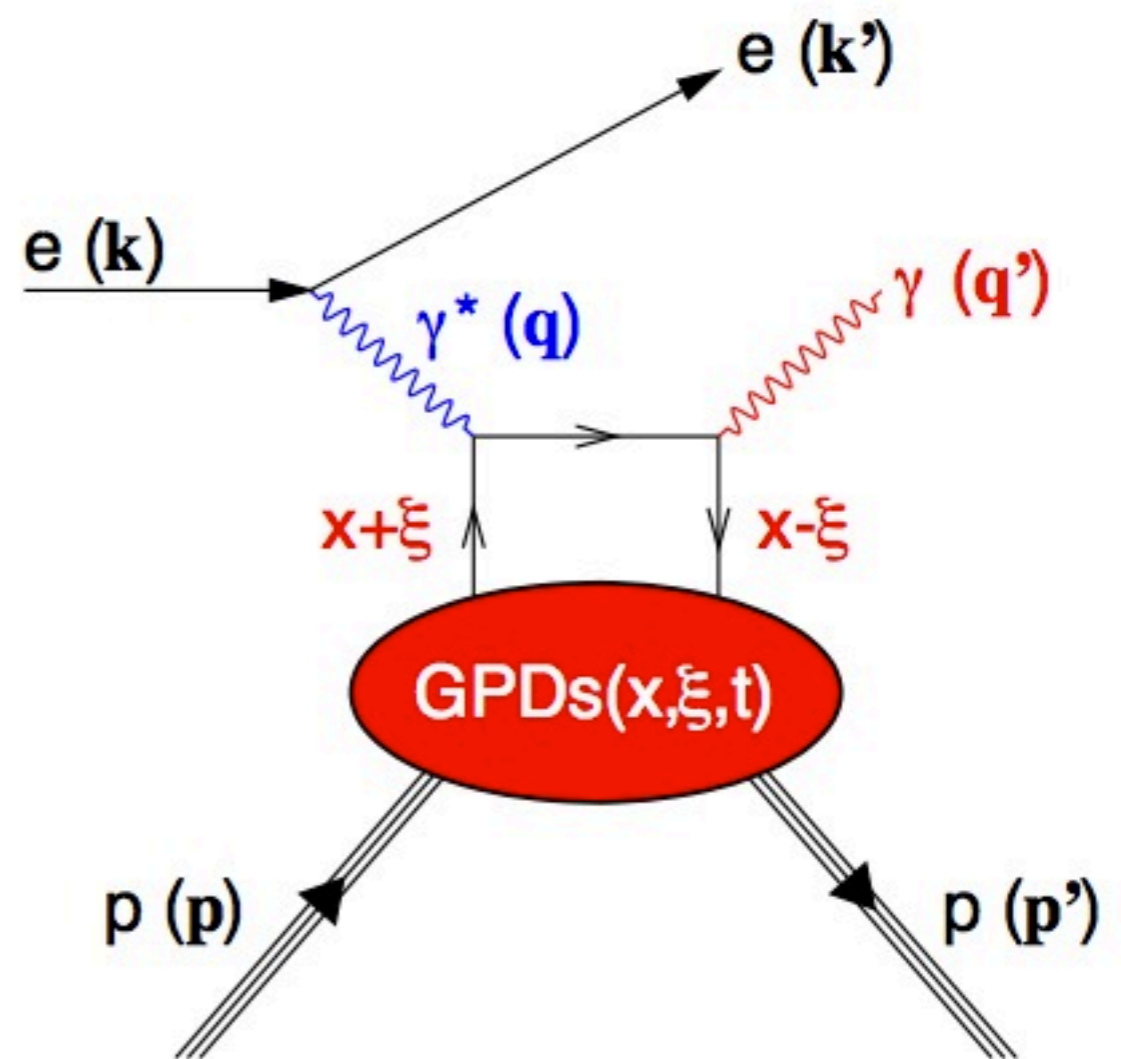
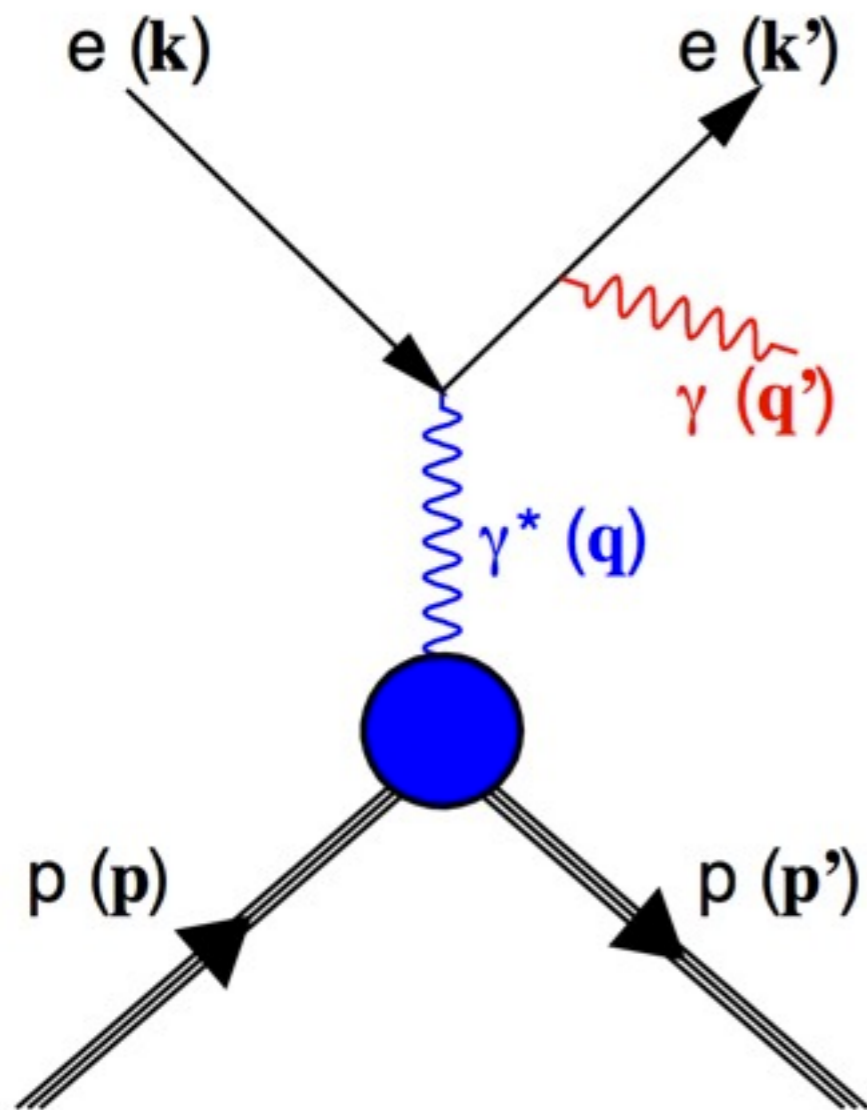
g_1 : [A. Airapetian et al, Phys. Rev. D 75 \(2007\) 012007](#)

Distribution Graph

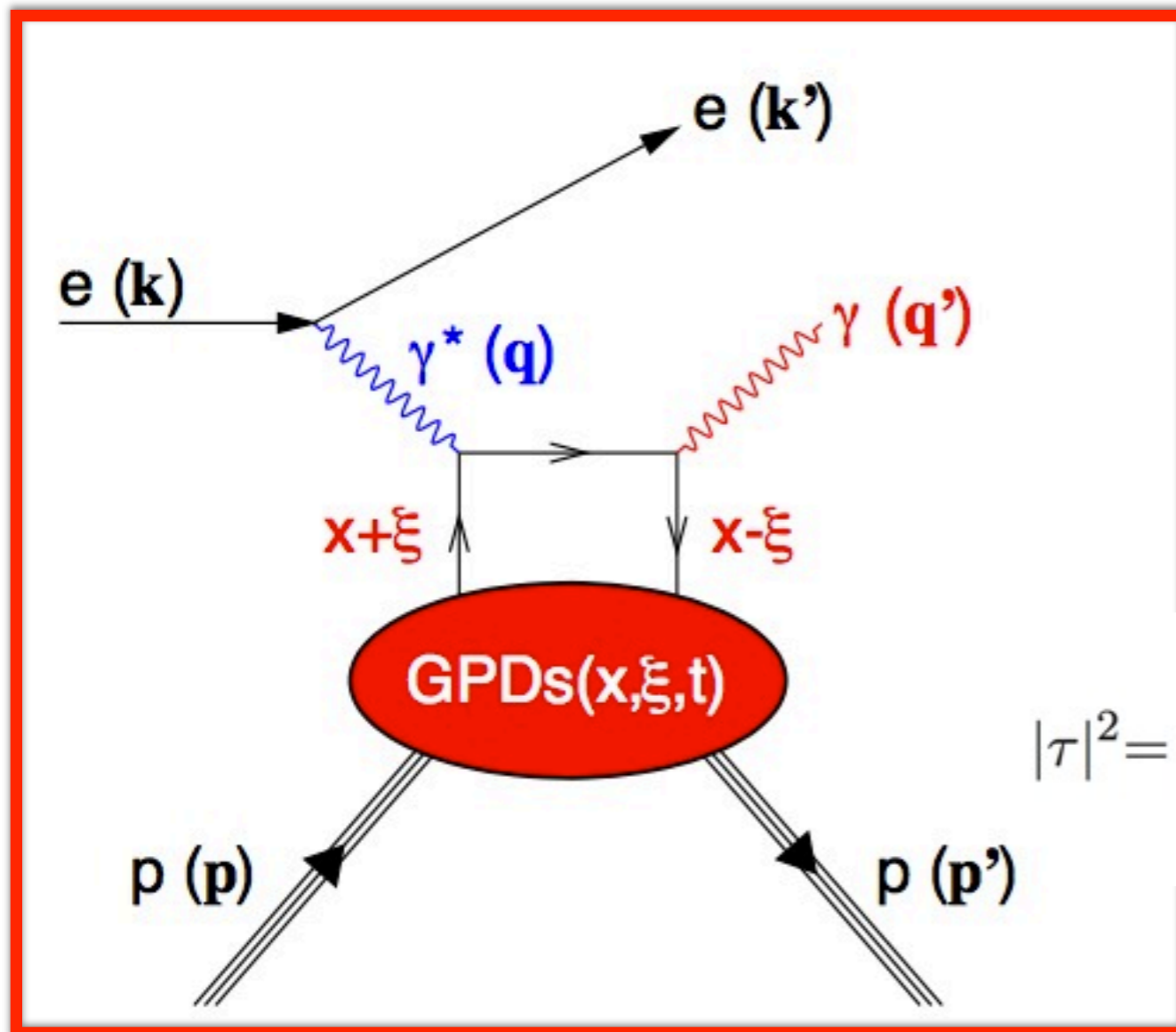


Deeply Virtual Compton Scattering

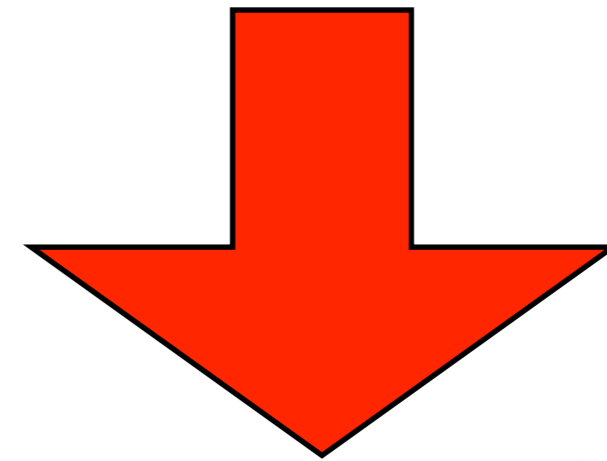
$$e p \rightarrow e p \gamma$$



Deeply Virtual Compton Scattering



$$\frac{d\sigma}{dx_B dQ^2 d|t| d\phi} = \frac{x_B e^6 |\tau|^2}{32(2\pi)^4 Q^4 \sqrt{1 + \epsilon^2}}$$



$$|\tau|^2 = |\tau_{\text{BH}}|^2 + |\tau_{\text{DVCS}}|^2 + \overbrace{\tau_{\text{BH}}\tau_{\text{DVCS}}^* + \tau_{\text{BH}}^*\tau_{\text{DVCS}}}^{\mathcal{I}}$$

DVCS @ HERMES

$$\mathcal{A}_C(\phi) \equiv \frac{d\sigma^+(\phi) - d\sigma^-(\phi)}{d\sigma^+(\phi) + d\sigma^-(\phi)} \quad \tilde{\propto} \quad \text{Re}(\mathcal{H})$$

$$\mathcal{A}_{\text{LU}}^{\text{I}}(\phi) \equiv \frac{(d\sigma(\phi)^{+\rightarrow} - d\sigma(\phi)^{+\leftarrow}) - (d\sigma(\phi)^{-\rightarrow} - d\sigma(\phi)^{-\leftarrow})}{(d\sigma(\phi)^{+\rightarrow} + d\sigma(\phi)^{+\leftarrow}) + (d\sigma(\phi)^{-\rightarrow} + d\sigma(\phi)^{-\leftarrow})} \quad \tilde{\propto} \quad \text{Im}(\mathcal{H})$$

$$\mathcal{A}_{\text{LU}}^{\text{DVCS}}(\phi) \equiv \frac{(d\sigma(\phi)^{+\rightarrow} + d\sigma(\phi)^{-\rightarrow}) - (d\sigma(\phi)^{+\leftarrow} + d\sigma(\phi)^{-\leftarrow})}{(d\sigma(\phi)^{+\rightarrow} + d\sigma(\phi)^{-\rightarrow}) + (d\sigma(\phi)^{+\leftarrow} + d\sigma(\phi)^{-\leftarrow})} \quad \tilde{\propto} \quad \text{Im}[\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*]$$

$$\mathcal{A}_{\text{UT}}^{\text{I}}(\phi, \phi_S) \equiv \frac{d\sigma^+(\phi, \phi_S) - d\sigma^+(\phi, \phi_S + \pi) - d\sigma^-(\phi, \phi_S) + d\sigma^-(\phi, \phi_S + \pi)}{d\sigma^+(\phi, \phi_S) + d\sigma^+(\phi, \phi_S + \pi) + d\sigma^-(\phi, \phi_S) + d\sigma^-(\phi, \phi_S + \pi)} \quad \tilde{\propto} \quad \text{Im}(\mathcal{E})$$

$$\mathcal{A}_{\text{UT}}^{\text{DVCS}}(\phi, \phi_S) \equiv \frac{d\sigma^+(\phi, \phi_S) - d\sigma^+(\phi, \phi_S + \pi) + d\sigma^-(\phi, \phi_S) - d\sigma^-(\phi, \phi_S + \pi)}{d\sigma^+(\phi, \phi_S) + d\sigma^+(\phi, \phi_S + \pi) + d\sigma^-(\phi, \phi_S) + d\sigma^-(\phi, \phi_S + \pi)} \quad \tilde{\propto} \quad \text{Im}(\mathcal{E})$$

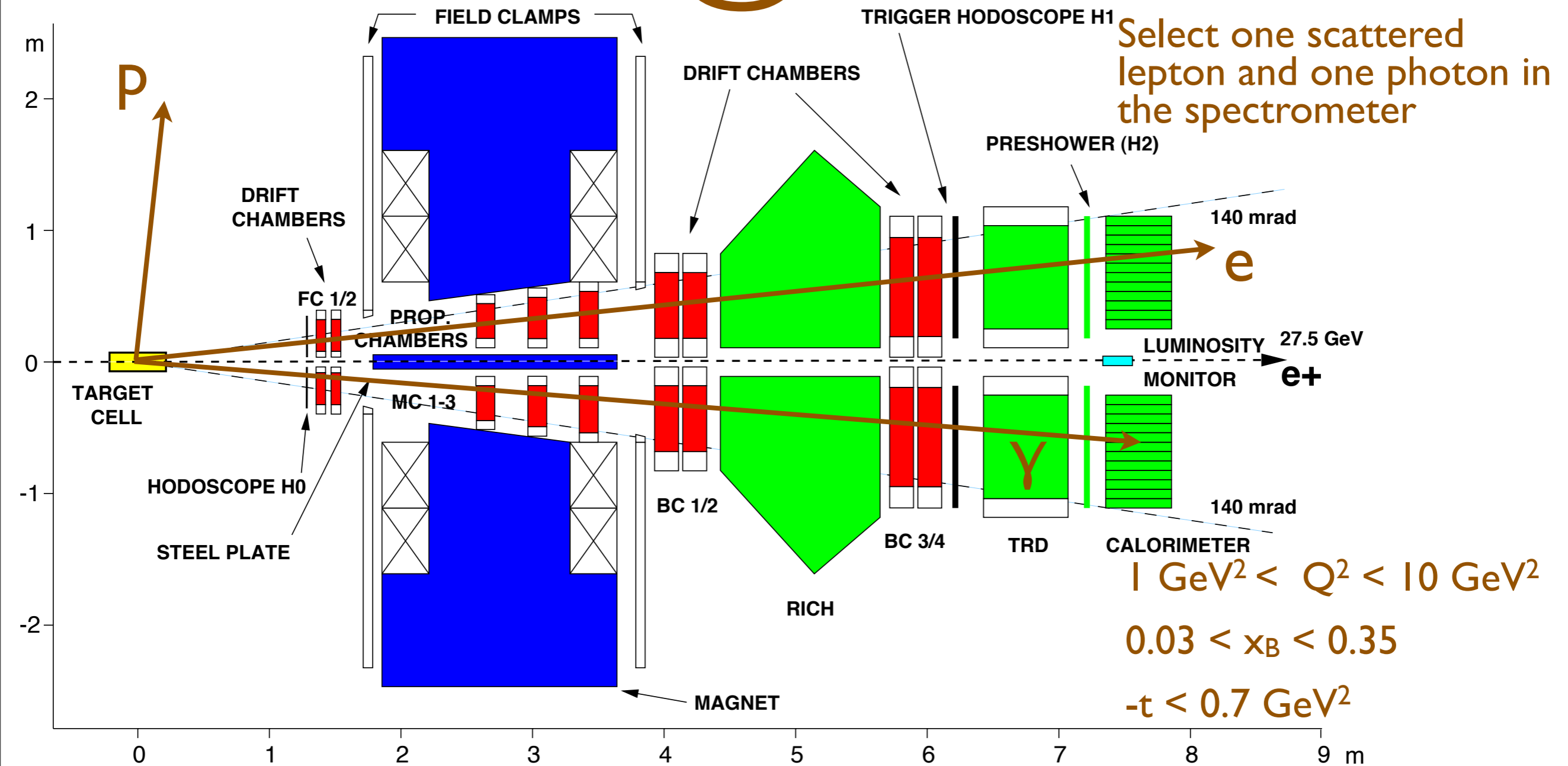
$$\mathcal{A}_{\text{LT}}^{\text{BH+DVCS}}(\phi, \phi_S) \equiv \frac{1}{8d\sigma_{\text{UU}}} \left[(d\vec{\sigma}^{+\uparrow} - d\vec{\sigma}^{+\downarrow} - d\vec{\sigma}^{-\uparrow} + d\vec{\sigma}^{-\downarrow}) + (d\vec{\sigma}^{-\uparrow} - d\vec{\sigma}^{-\downarrow} - d\vec{\sigma}^{+\uparrow} + d\vec{\sigma}^{+\downarrow}) \right] \quad \tilde{\propto} \quad \text{Re}(\mathcal{H} + \mathcal{E})$$

$$\mathcal{A}_{\text{LT}}^{\text{I}}(\phi, \phi_S) \equiv \frac{1}{8d\sigma_{\text{UU}}} \left[(d\vec{\sigma}^{+\uparrow} - d\vec{\sigma}^{+\downarrow} - d\vec{\sigma}^{-\uparrow} + d\vec{\sigma}^{-\downarrow}) - (d\vec{\sigma}^{-\uparrow} - d\vec{\sigma}^{-\downarrow} - d\vec{\sigma}^{+\uparrow} + d\vec{\sigma}^{+\downarrow}) \right] \quad \tilde{\propto} \quad \text{Re}(\mathcal{H})$$

$$\mathcal{A}_{\text{UL}}(\phi) \equiv \frac{[\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] - [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}{[\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] + [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]} \quad \tilde{\propto} \quad \text{Im}(\tilde{\mathcal{H}})$$

$$\mathcal{A}_{\text{LL}}(\phi) \equiv \frac{[\sigma^{\rightarrow\Rightarrow}(\phi) + \sigma^{\leftarrow\leftarrow}(\phi)] - [\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}{[\sigma^{\rightarrow\Rightarrow}(\phi) + \sigma^{\leftarrow\leftarrow}(\phi)] + [\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]} \quad \tilde{\propto} \quad \text{Re}(\tilde{\mathcal{H}})$$

DVCS @ HERMES

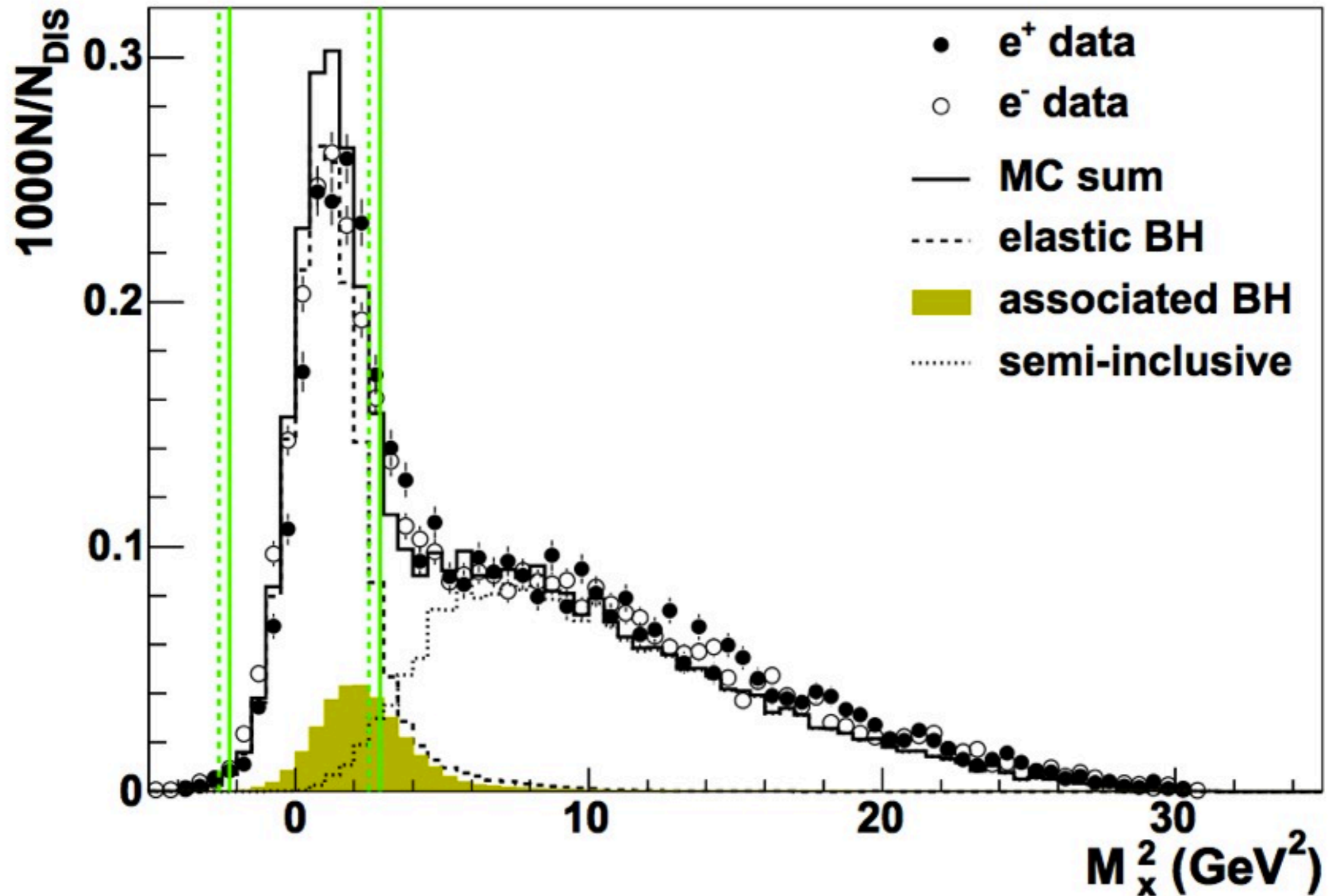


$$\langle Q^2 \rangle \cong 2.4 \text{ GeV}^2$$

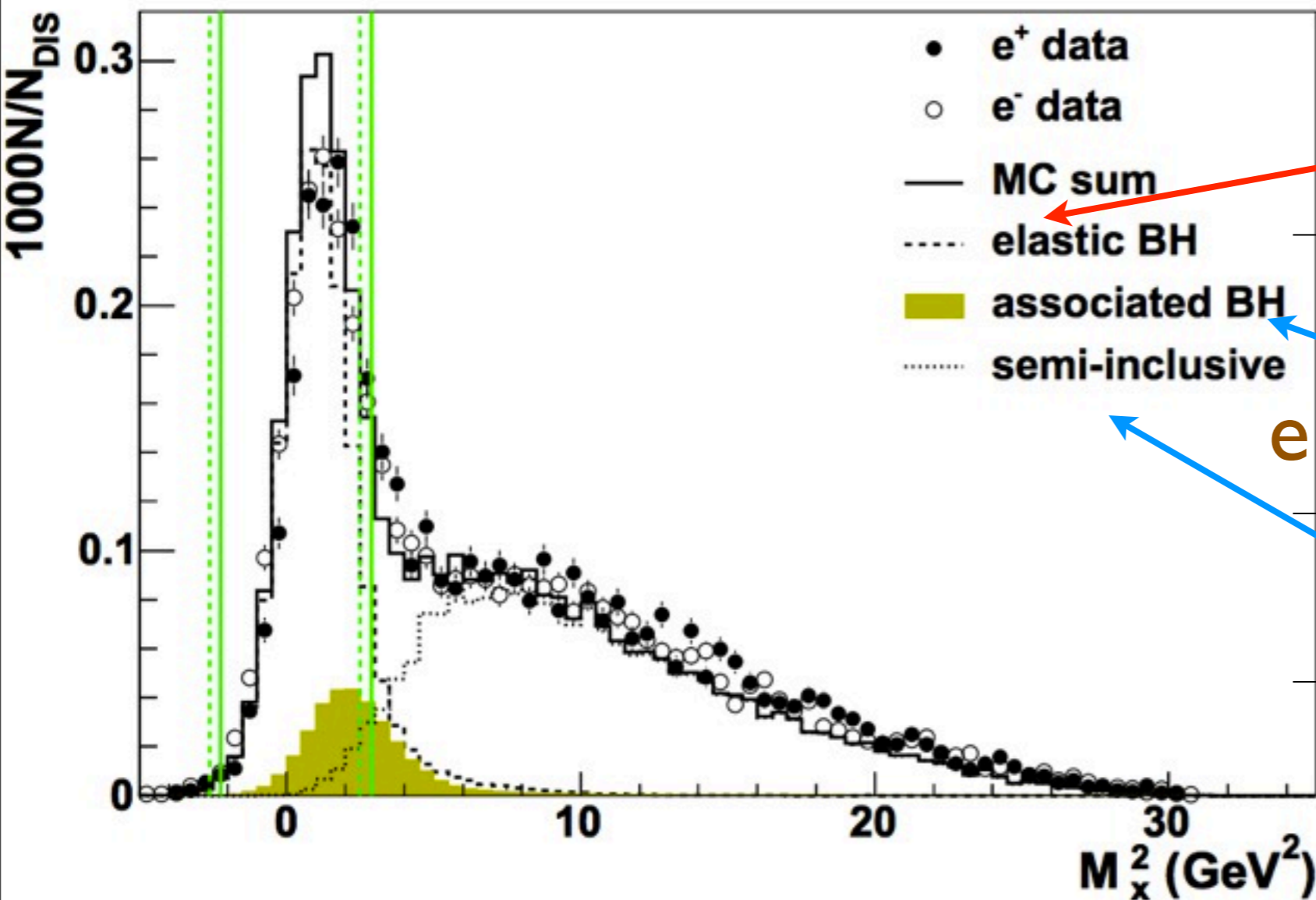
$$\langle x_B \rangle \cong 0.1$$

$$\langle -t \rangle \cong 0.1 \text{ GeV}^2$$

DVCS @ HERMES



DVCS @ HERMES



Wanted Signal

BH from Δ , e.g.

$e \Delta \rightarrow e \gamma \Delta \rightarrow e \gamma p \pi^0$

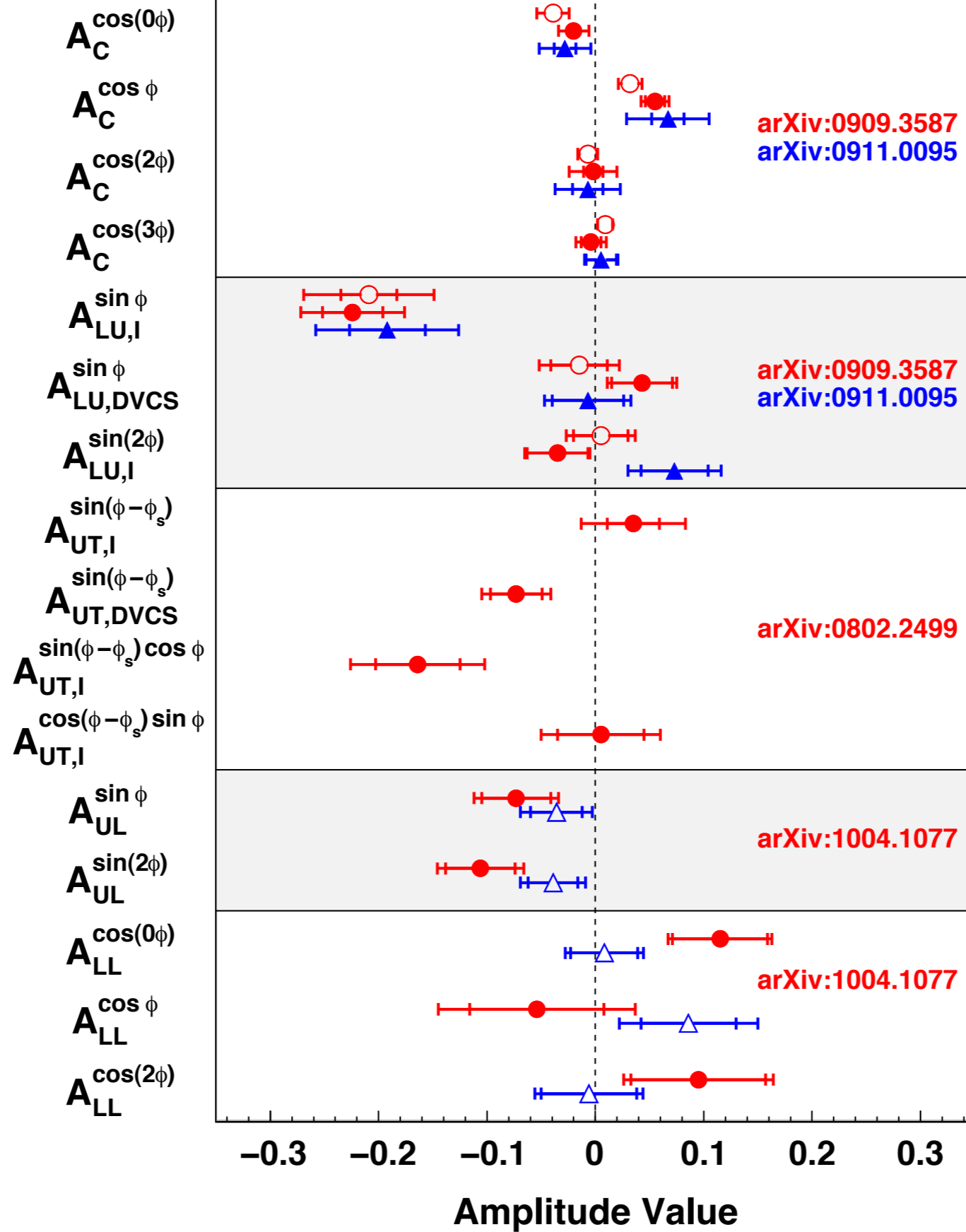
$e p \rightarrow e X \gamma$

$e p \rightarrow e p \pi^0$

D
V
C
S
@

HERMES DVCS

● Hydrogen
▲ Deuterium
○ Preliminary

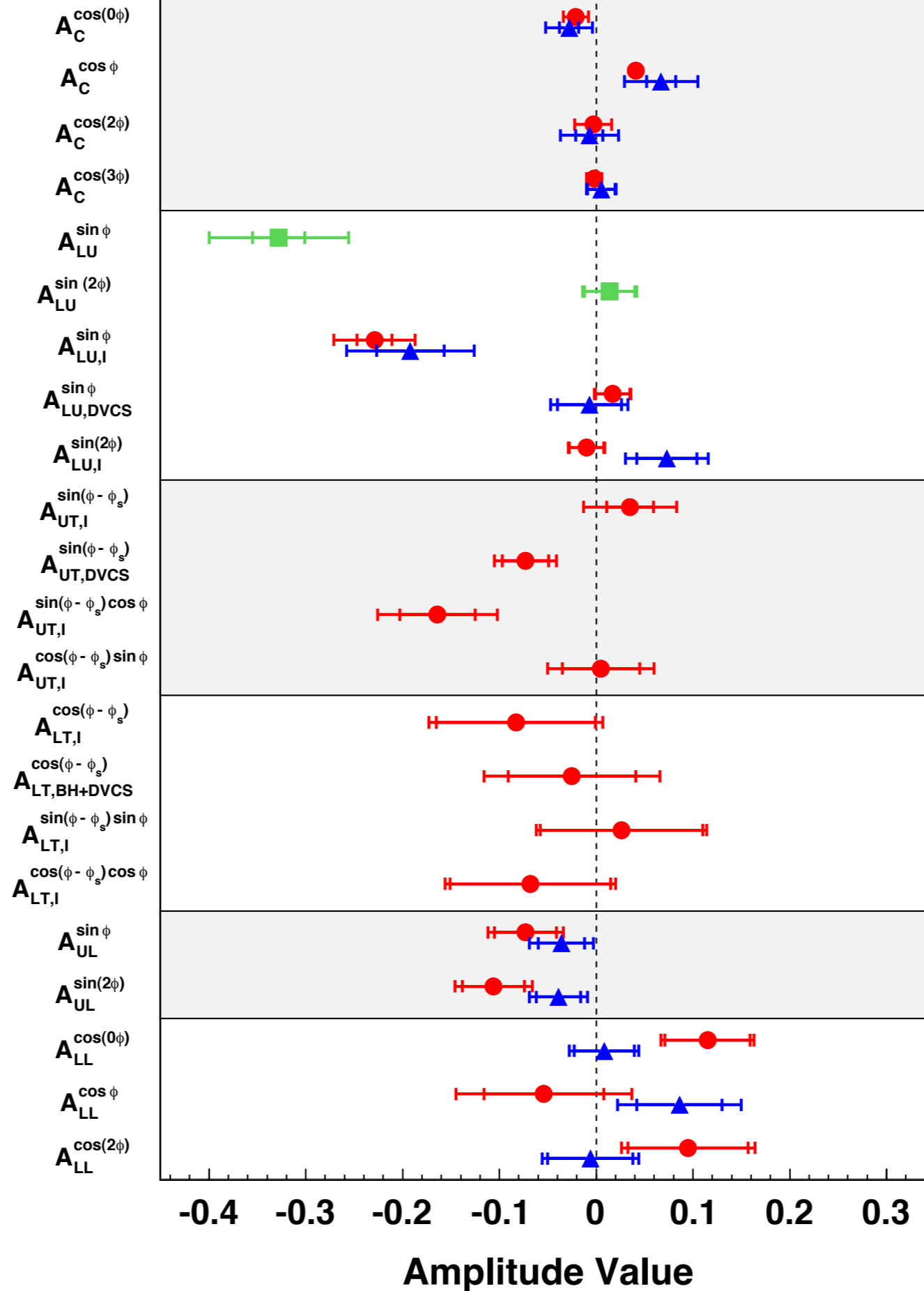


H
E
R
M
E
S

D
V
C
C
S
@

HERMES DVCS

- Hydrogen
- ▲ Deuterium
- Hydrogen Pure



H
E
R
M
E
S

Beam-Charge Asymmetries

A. Airapetian et al, JHEP 07 (2012) 032

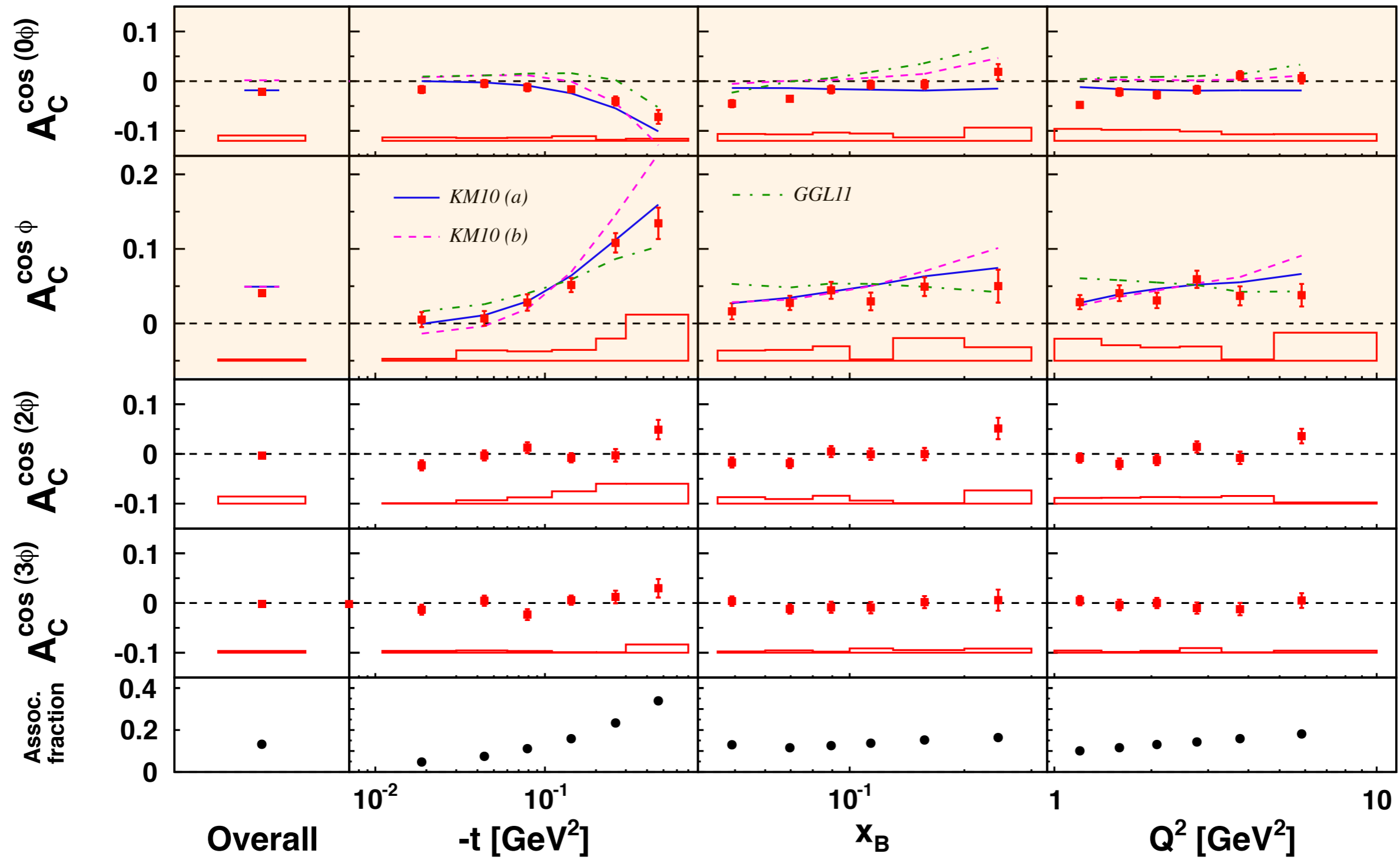
<http://arxiv.org/abs/1203.6287>

Kumerički and Müller, Nucl. Phys. **B841** (2010)

<http://arxiv.org/abs/0904.0458>

G. Goldstein, J. Hernandez and S. Liuti, Phys. Rev. **D84** (2011)

<http://arxiv.org/abs/1012.3776>

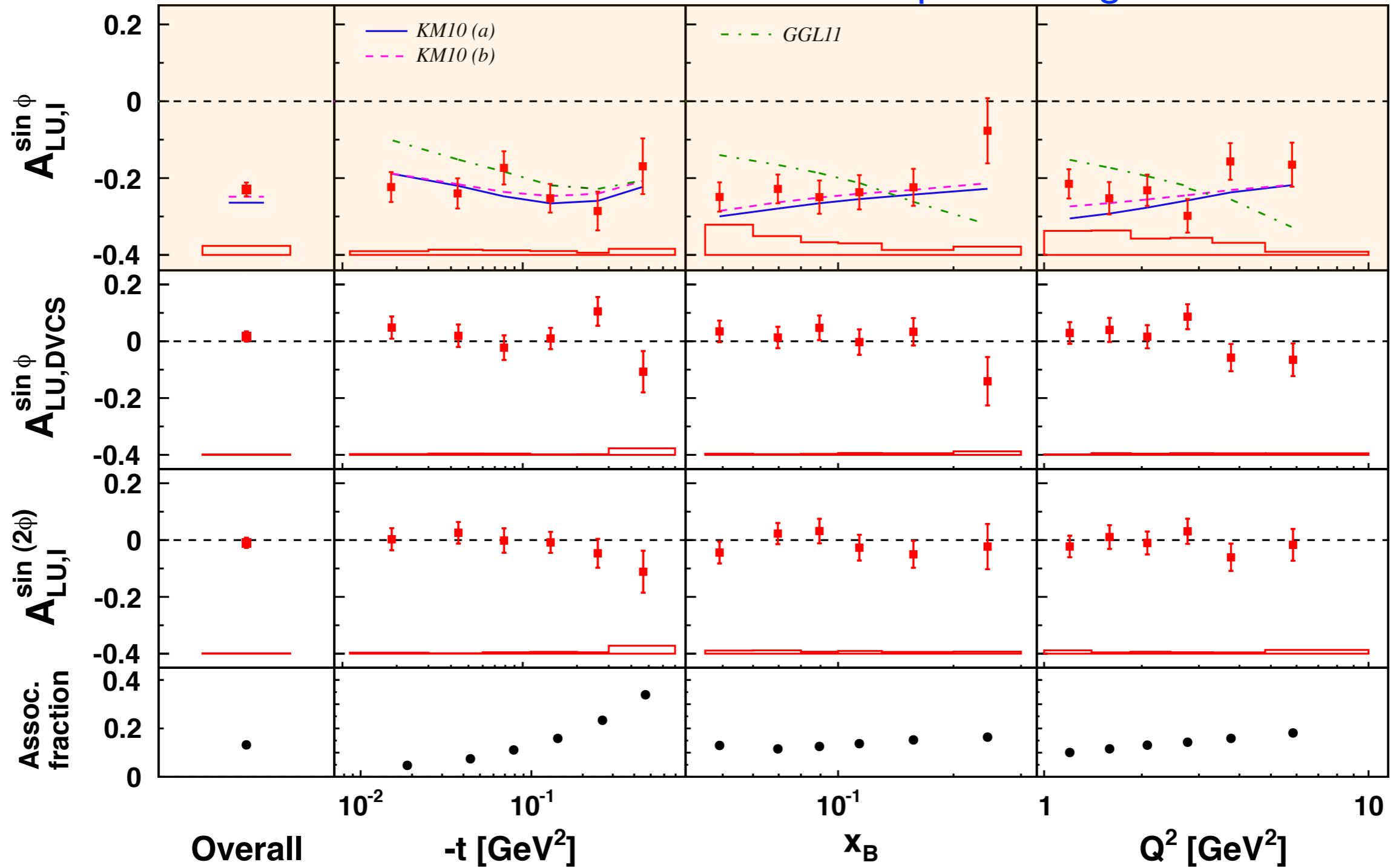


Beam Charge Asymmetries access $\text{Re}(\mathcal{H})$

Beam-Spin Asymmetries

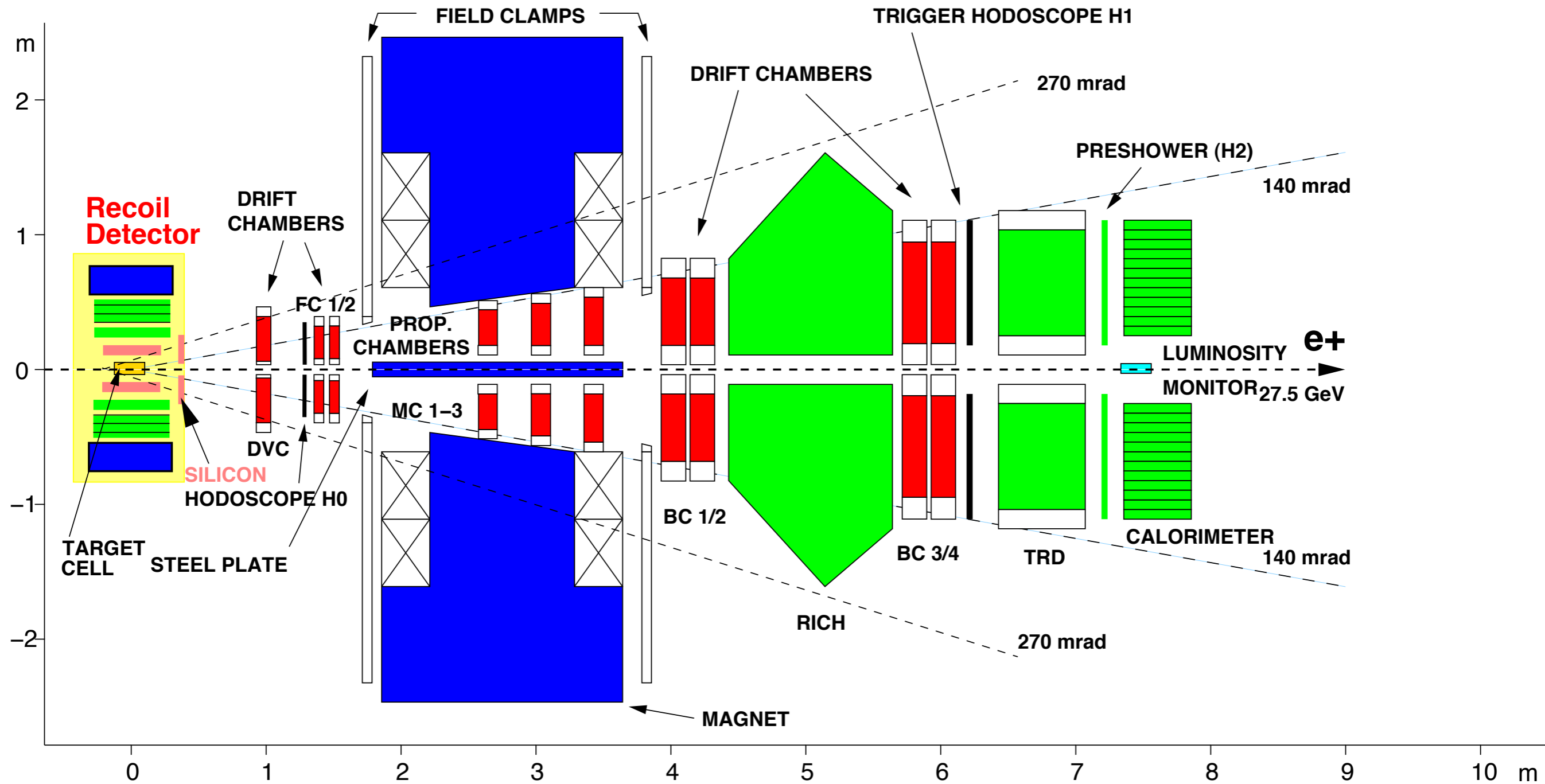
A. Airapetian et al, JHEP 07 (2012) 032

<http://arxiv.org/abs/1203.6287>



Beam Helicity Asymmetries access $\text{Im}(\mathcal{H})$

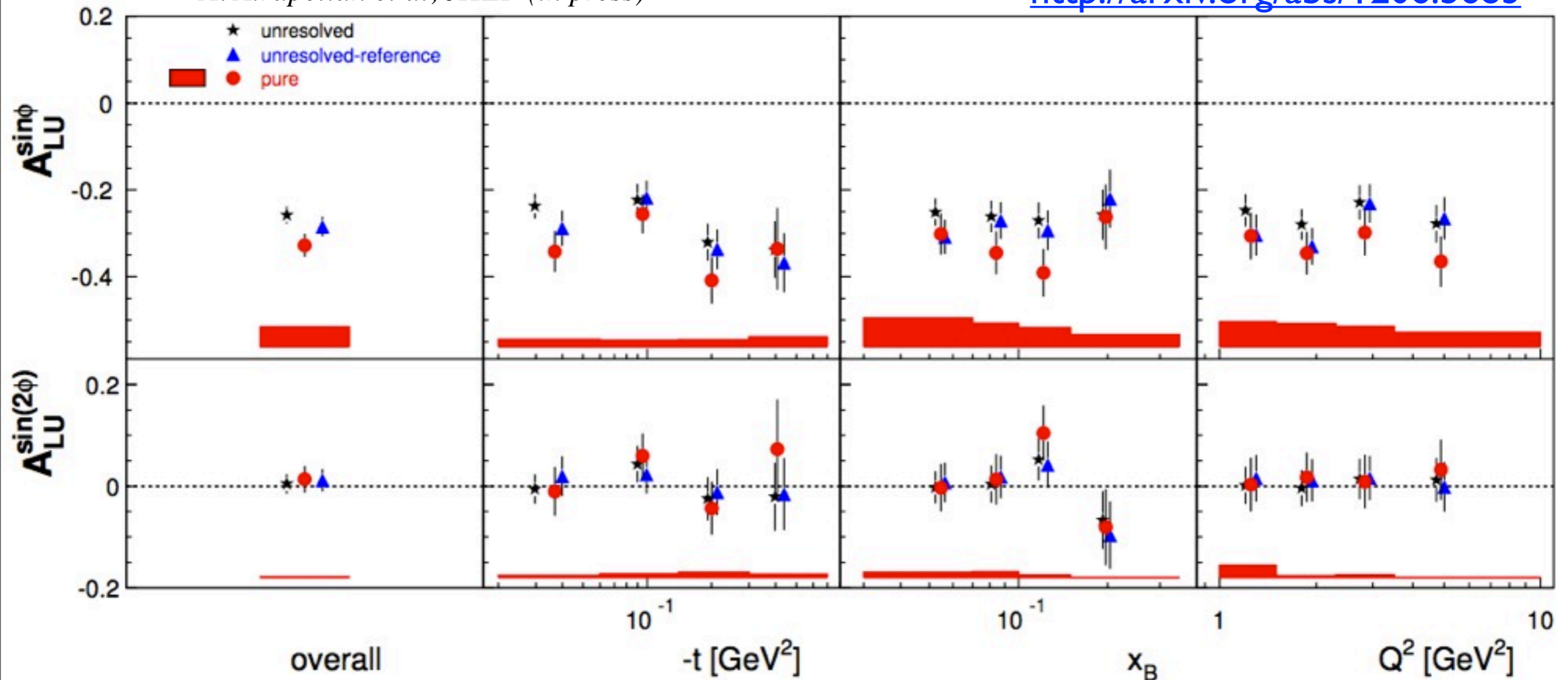
DVCS @ HERMES



Exclusive Measurement

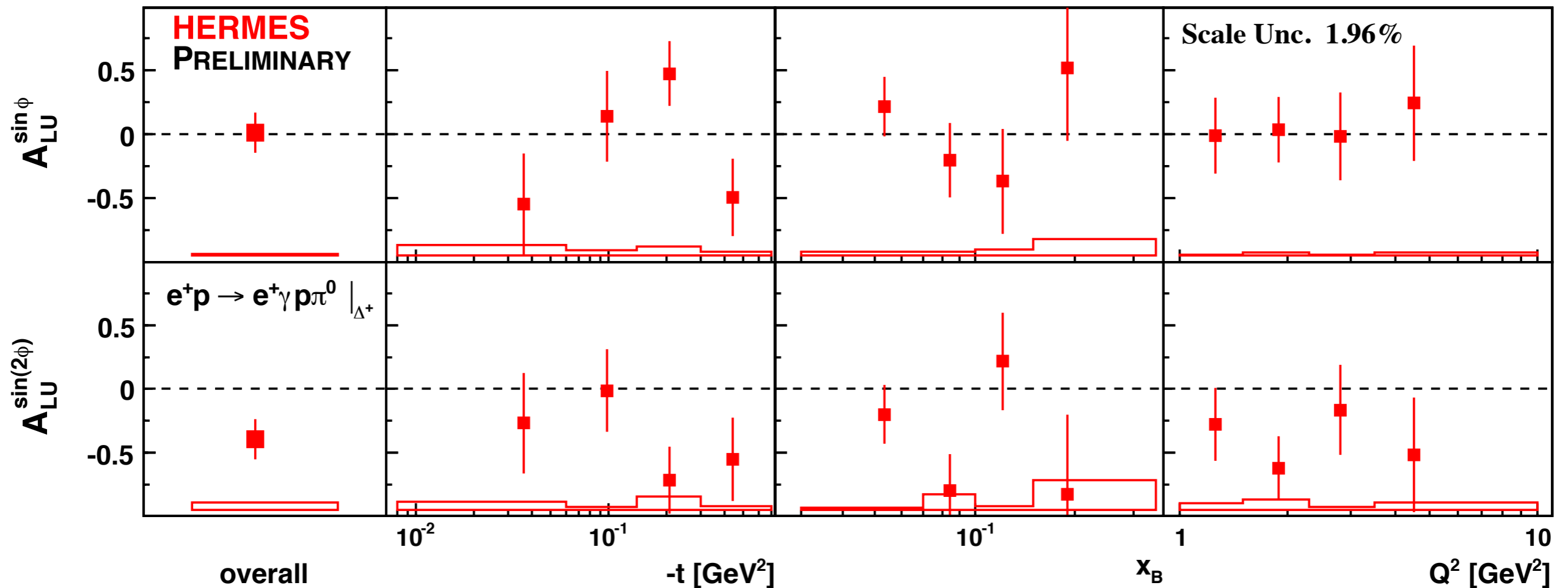
A. Airapetian et al, JHEP (in press)

<http://arxiv.org/abs/1206.5683>



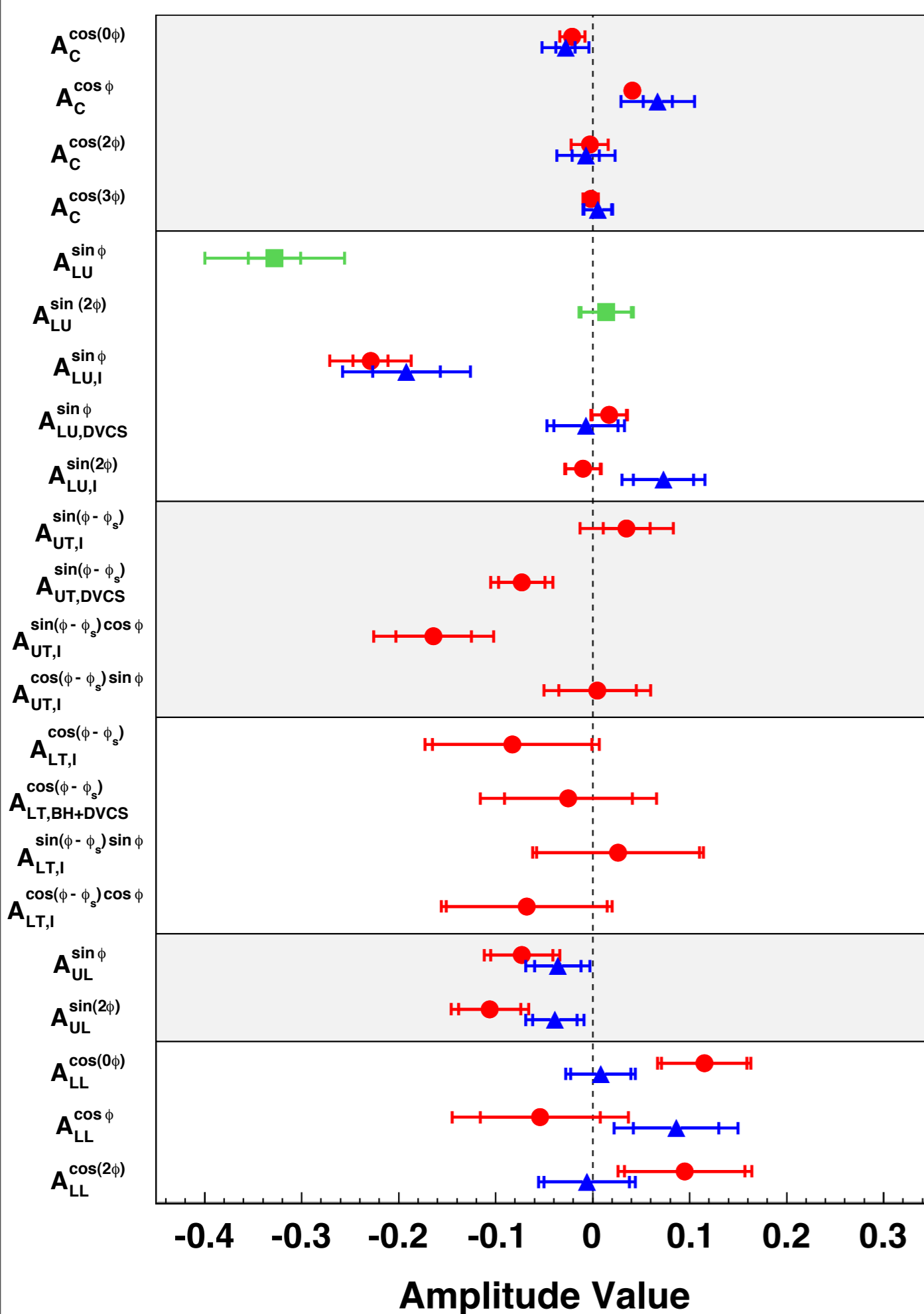
Fully reconstructed measurement of $ep \rightarrow ep\gamma$

Exclusive Measurement



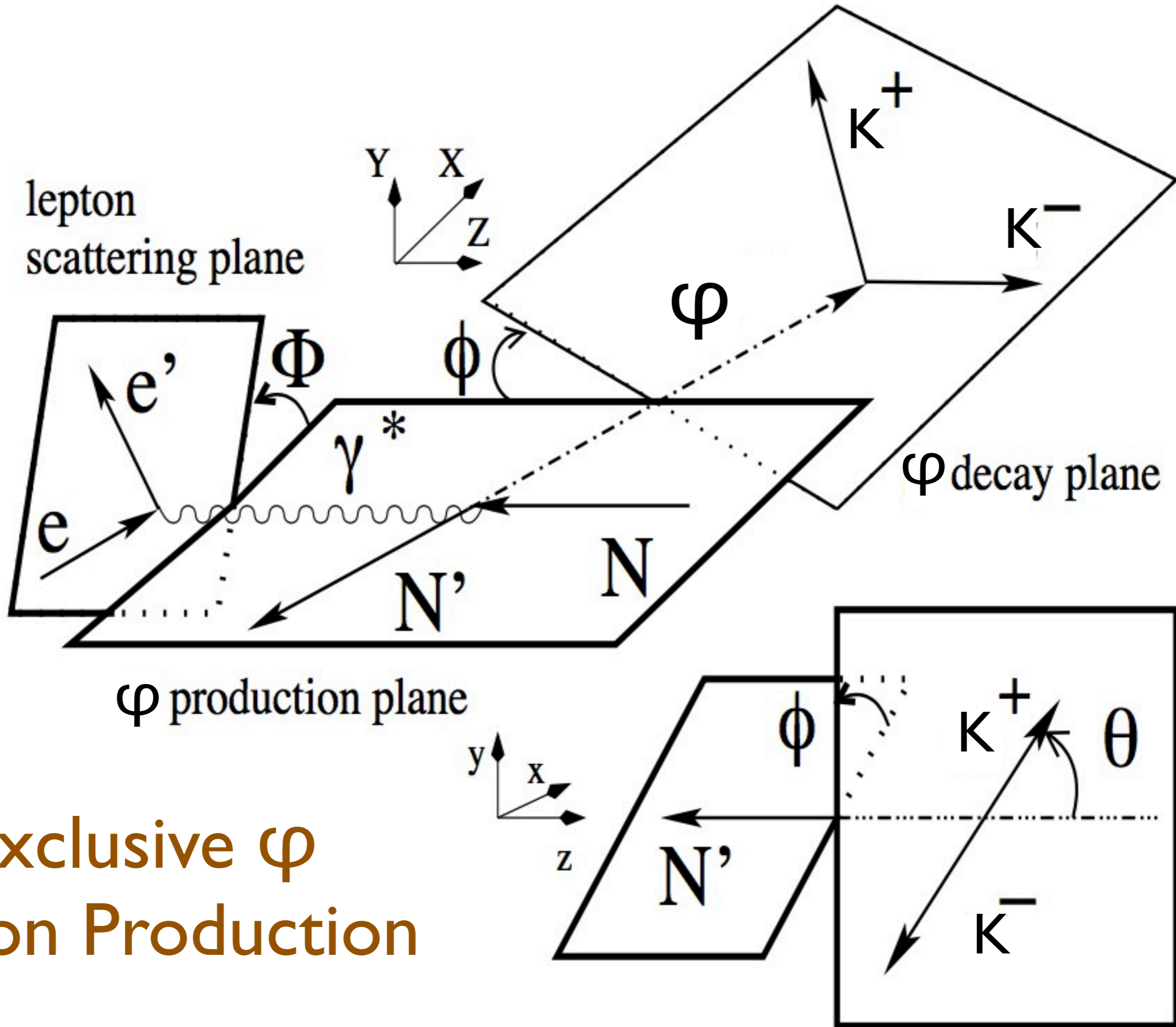
Results taken from measurement of $ep \rightarrow ep\pi^0\gamma$

(Overall 'zero' asymmetry implies that the 'associated' fraction in the non-exclusive results acts as a dilution)



- DVCS remains the leading process for access to Generalised Parton Distributions
- HERMES has the most diverse DVCS measurements of any experiment.
- Polarised target experiments are essential for the extraction of GPDs; should be seen as a fundamental experimental priority!

Exclusive ϕ Meson Production

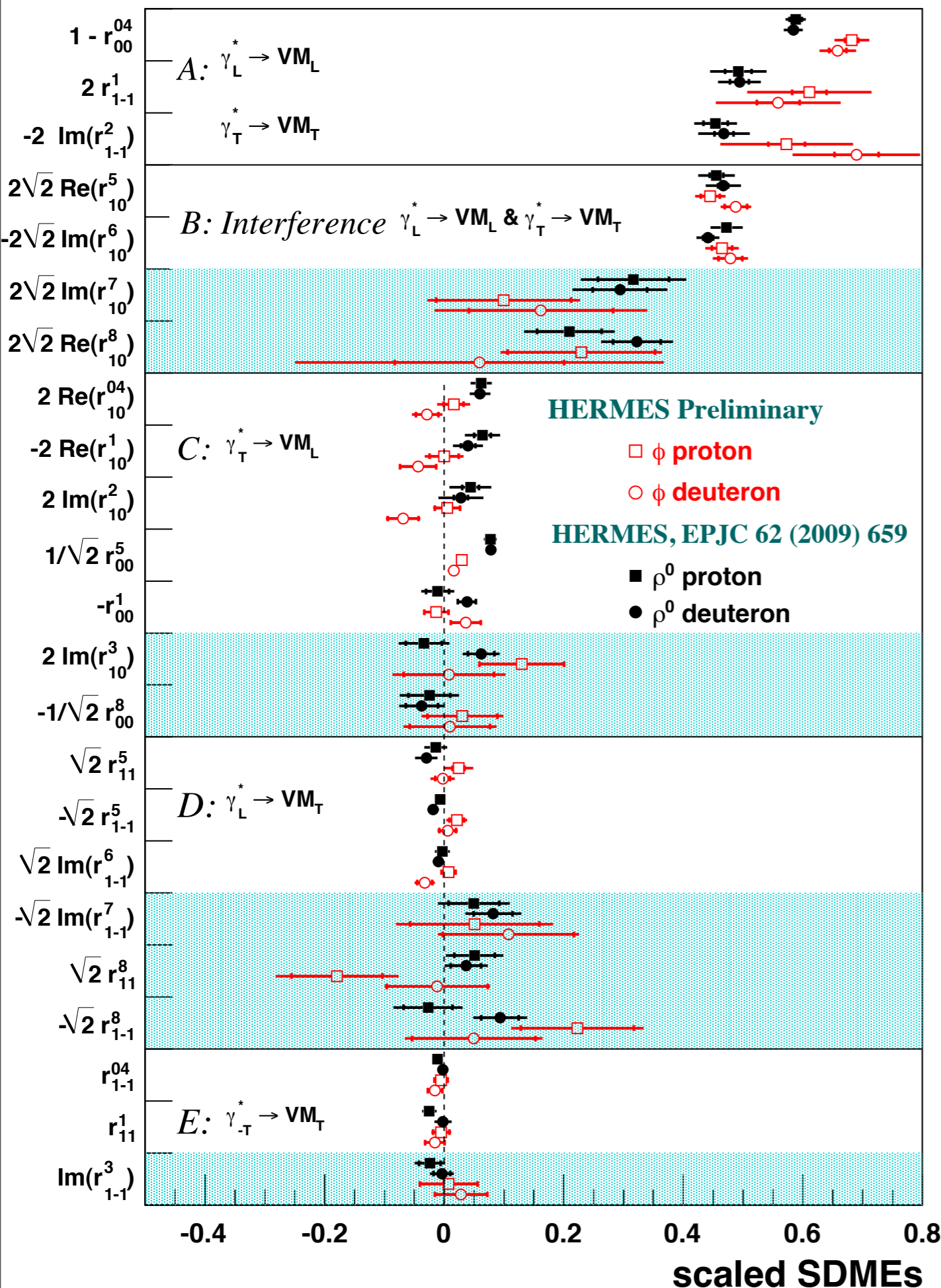


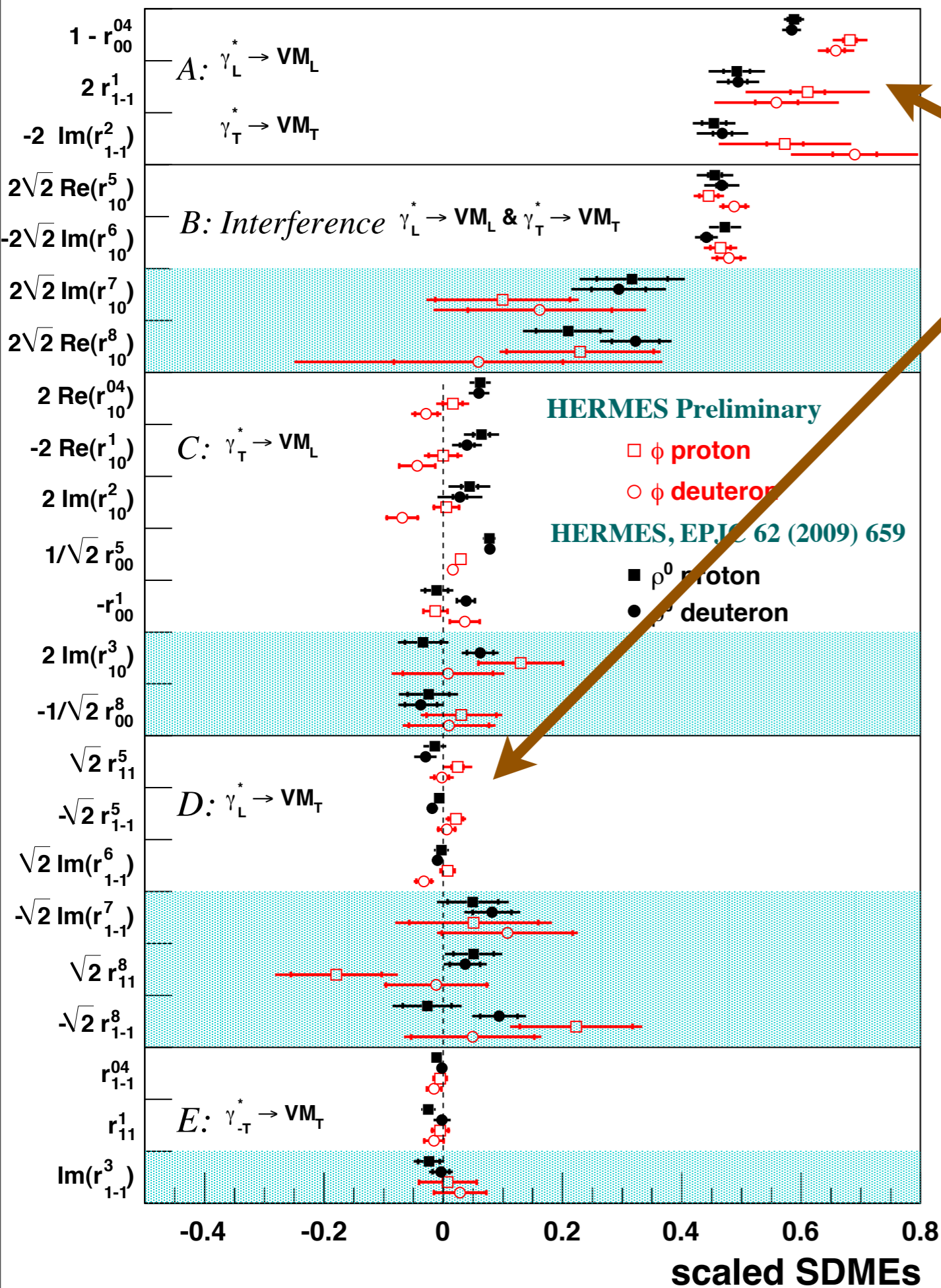
Exclusive Meson Production

Results taken from measurement of $ep \rightarrow eX\phi$.

No measured distinction between proton and deuteron data.

Leading-twist transitions are typically larger than the ρ^0 -equivalent.

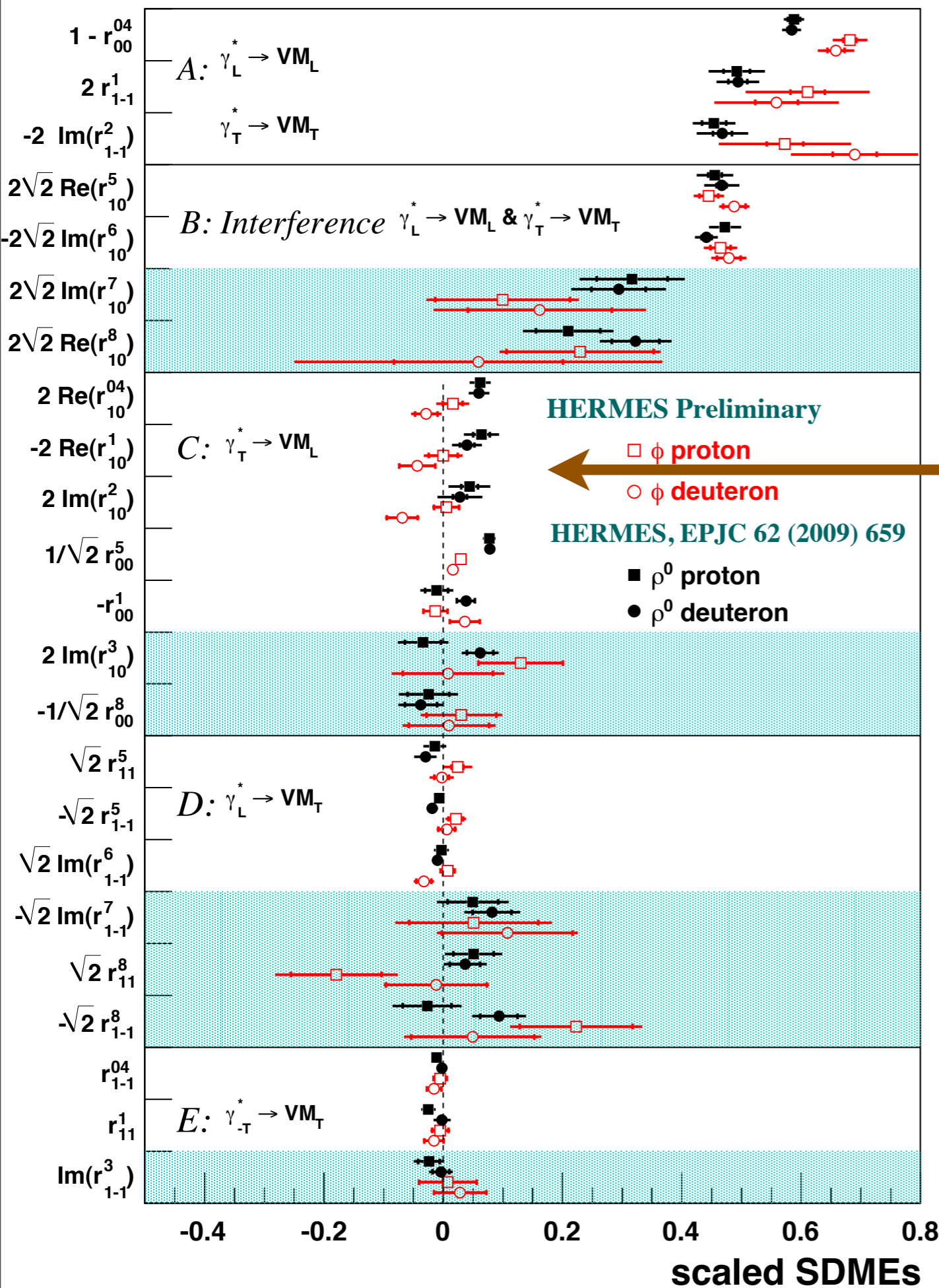




Longitudinal photons
 mostly produce
 longitudinal mesons

Longitudinal photons
mostly produce
longitudinal mesons

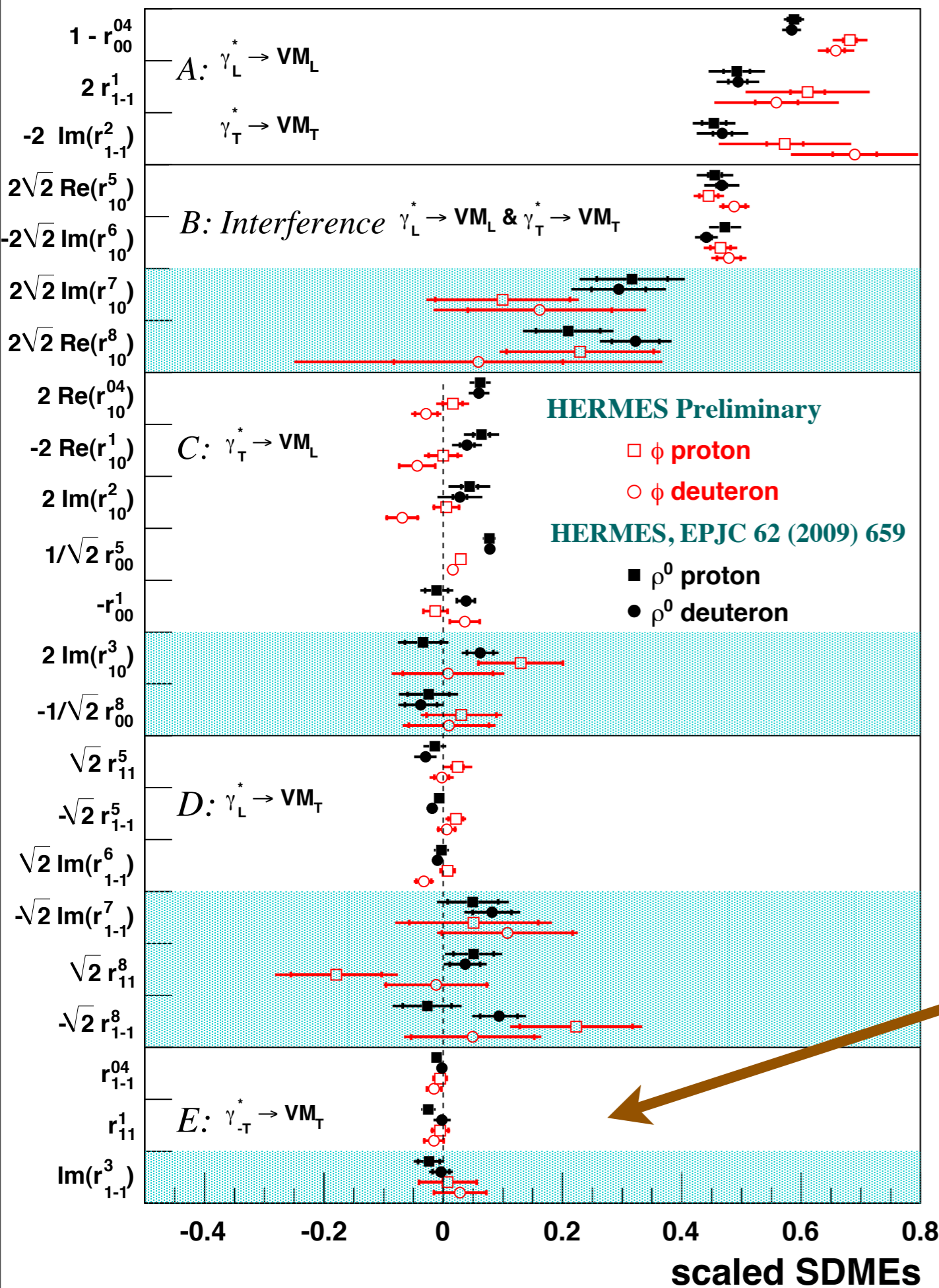
Some small indication
that transverse photons
can produce longitudinal
mesons



Longitudinal photons
mostly produce
longitudinal mesons

Some small indication
that transverse photons
can produce longitudinal
mesons

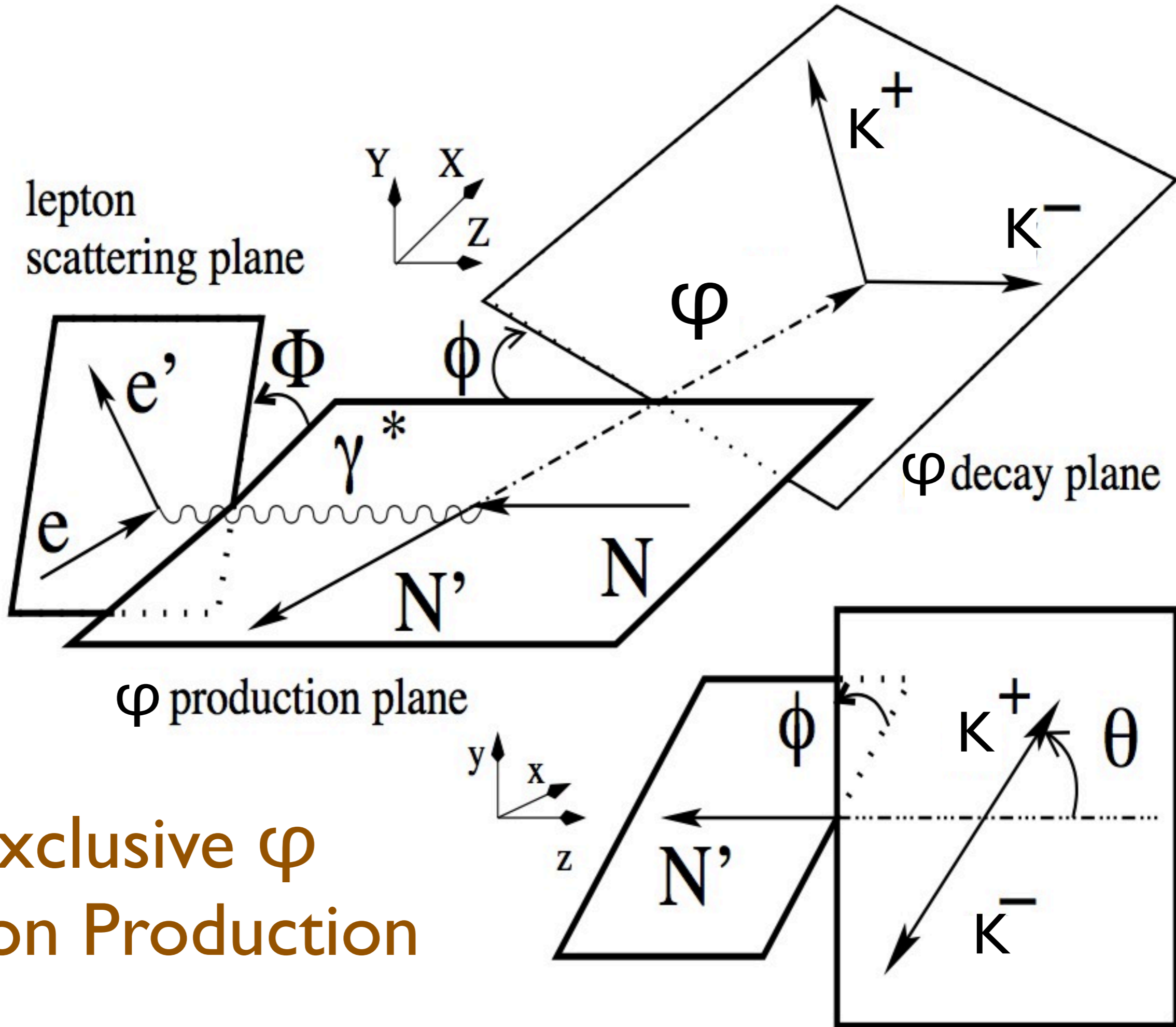
Zero indication of two
units of angular
momentum change
(-T γ makes +T ϕ)



Physics Update from HERMES

- **New inclusive measurement of g_2 and A_2** released, compatible with SLAC and SMC.
- **HERMES** has very **diverse** DVCS measurements available - programme almost complete.
- **Exclusive meson results** also available; φ SDMEs seem mostly to **match ρ^0** SDMEs.

Exclusive φ Meson Production



Angular Distribution

$$W^{U+L}(\Phi, \phi, \cos\theta) = W^{UU}(\Phi, \phi, \cos\theta) + W^{LU}(\Phi, \phi, \cos\theta)$$

For unpolarized target and beam:

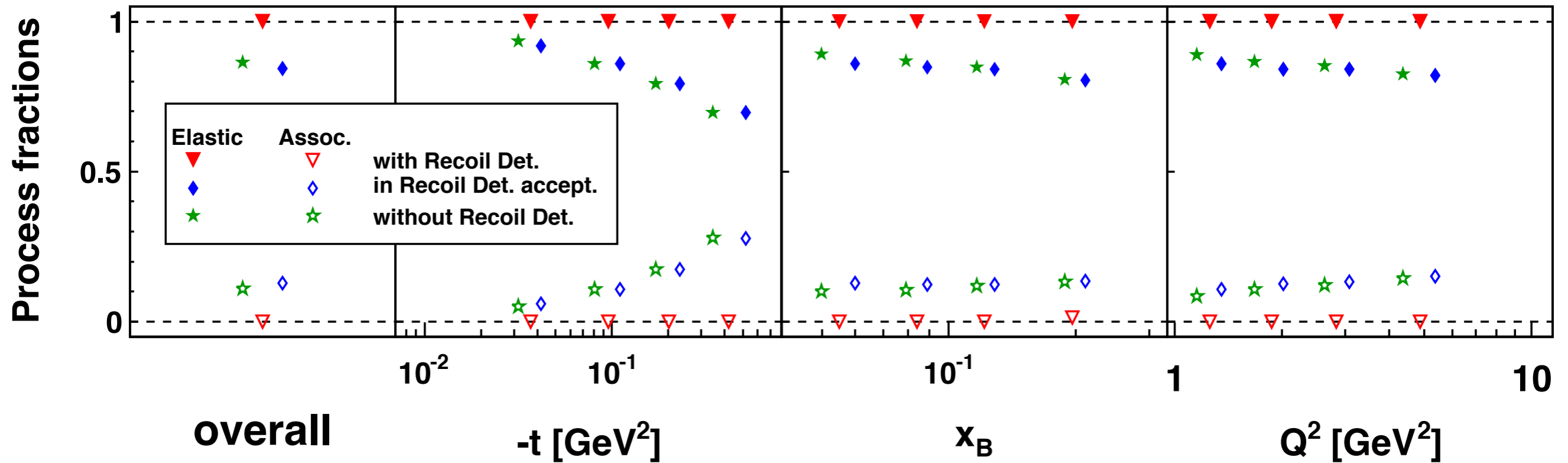
$$\begin{aligned} W^{UU}(\Phi, \phi, \cos\theta) = & \frac{3}{8\pi^2} \left[\frac{1}{2} (1 - r_{00}^{04}) + \frac{1}{2} (3r_{00}^{04} - 1) \cos^2\theta - \sqrt{2} \operatorname{Re}\{r_{10}^{04}\} \sin 2\theta \cos\phi - r_{1-1}^{04} \sin^2\theta \cos 2\phi \right. \\ & - \varepsilon \cos 2\Phi \left(r_{11}^1 \sin^2\theta + r_{00}^1 \cos^2\theta - \sqrt{2} \operatorname{Re}\{r_{10}^1\} \sin 2\theta \cos\phi - r_{1-1}^1 \sin^2\theta \cos 2\phi \right) \\ & - \varepsilon \sin 2\Phi \left(\sqrt{2} \operatorname{Im}\{r_{10}^2\} \sin 2\theta \sin\phi + \operatorname{Im}\{r_{1-1}^2\} \sin^2\theta \sin 2\phi \right) \\ & + \sqrt{2\varepsilon(1+\varepsilon)} \cos\Phi \left(r_{11}^5 \sin^2\theta + r_{00}^5 \cos^2\theta - \sqrt{2} \operatorname{Re}\{r_{10}^5\} \sin 2\theta \cos\phi - r_{1-1}^5 \sin^2\theta \cos 2\phi \right) \\ & \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\Phi \left(\sqrt{2} \operatorname{Im}\{r_{10}^6\} \sin 2\theta \sin\phi + \operatorname{Im}\{r_{1-1}^6\} \sin^2\theta \sin 2\phi \right) \right] \end{aligned}$$

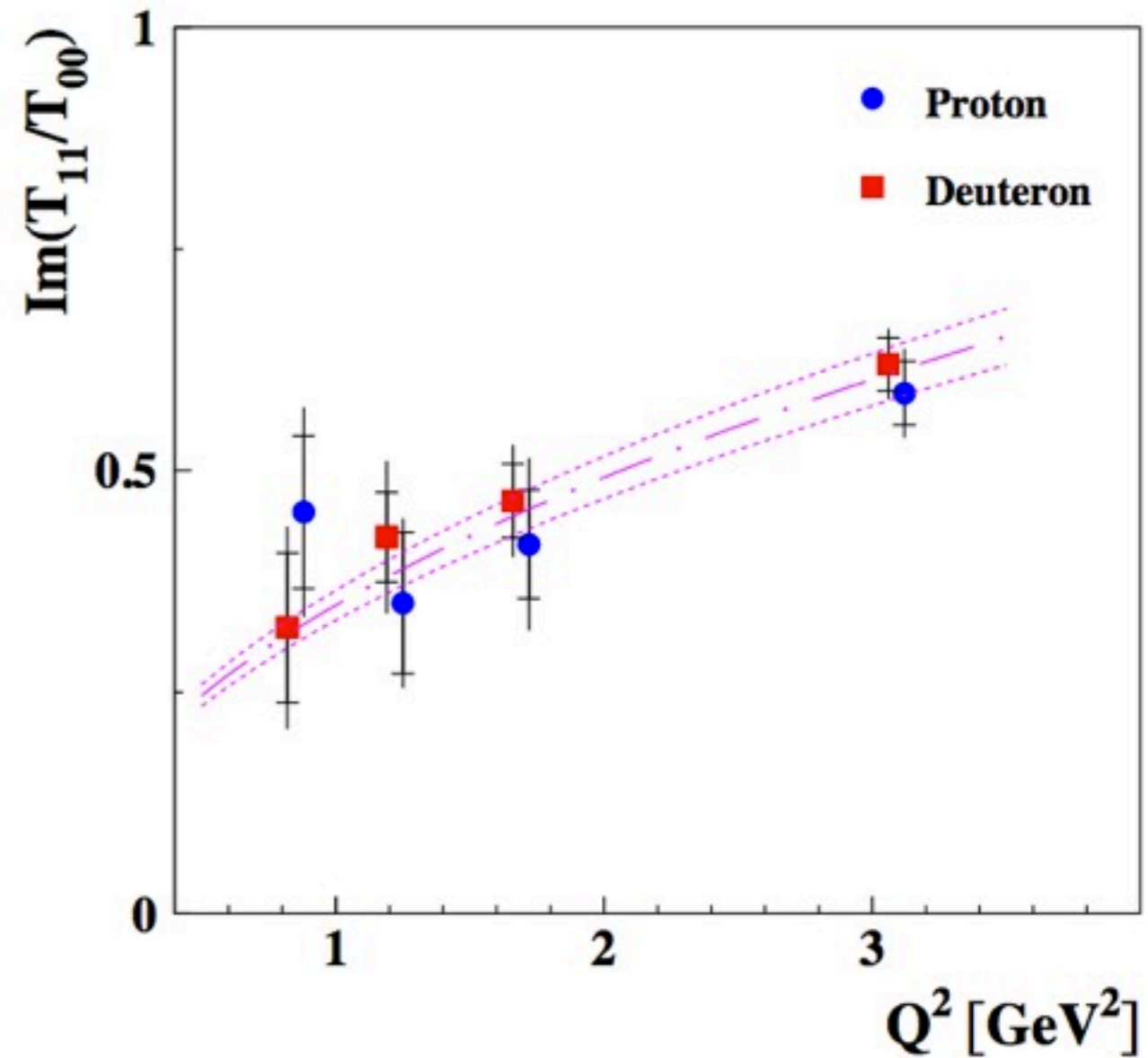
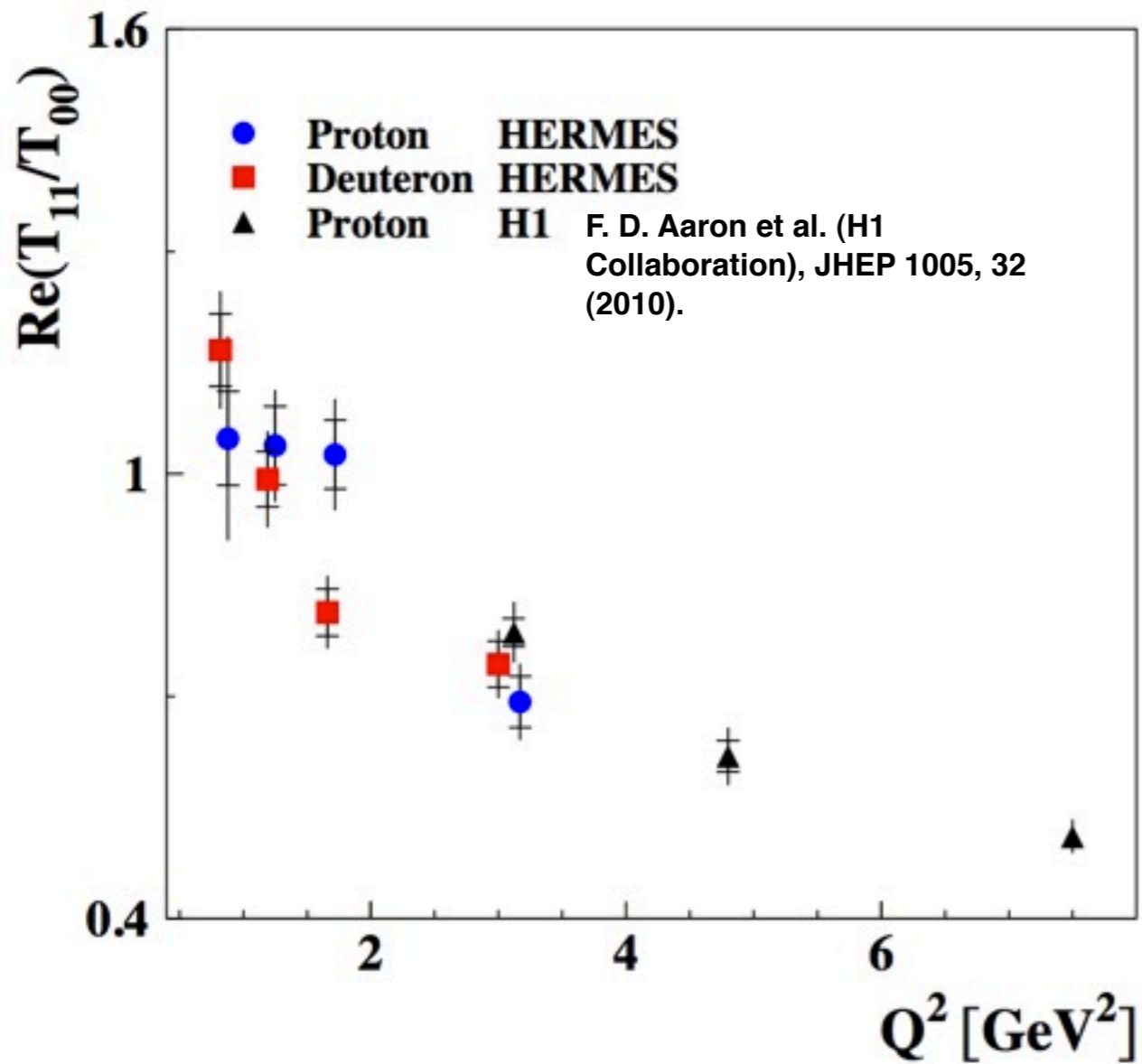
For unpolarized target and longitudinally polarized beam:

$$\begin{aligned} W^{LU}(\Phi, \phi, \cos\theta) = & \frac{3}{8\pi^2} P_{\text{Beam}} \left[\sqrt{1-\varepsilon^2} \left(\sqrt{2} \operatorname{Im}\{r_{10}^3\} \sin 2\theta \sin\phi + \operatorname{Im}\{r_{1-1}^3\} \sin^2\theta \sin 2\phi \right) \right. \\ & + \sqrt{2\varepsilon(1-\varepsilon)} \cos\Phi \left(\sqrt{2} \operatorname{Im}\{r_{10}^7\} \sin 2\theta \sin\phi + \operatorname{Im}\{r_{1-1}^7\} \sin^2\theta \sin 2\phi \right) \\ & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \sin\Phi \left(r_{11}^8 \sin^2\theta + r_{00}^8 \cos^2\theta - \sqrt{2} \operatorname{Re}\{r_{10}^8\} \sin 2\theta \cos\phi - r_{1-1}^8 \sin^2\theta \cos 2\phi \right) \right] \end{aligned}$$

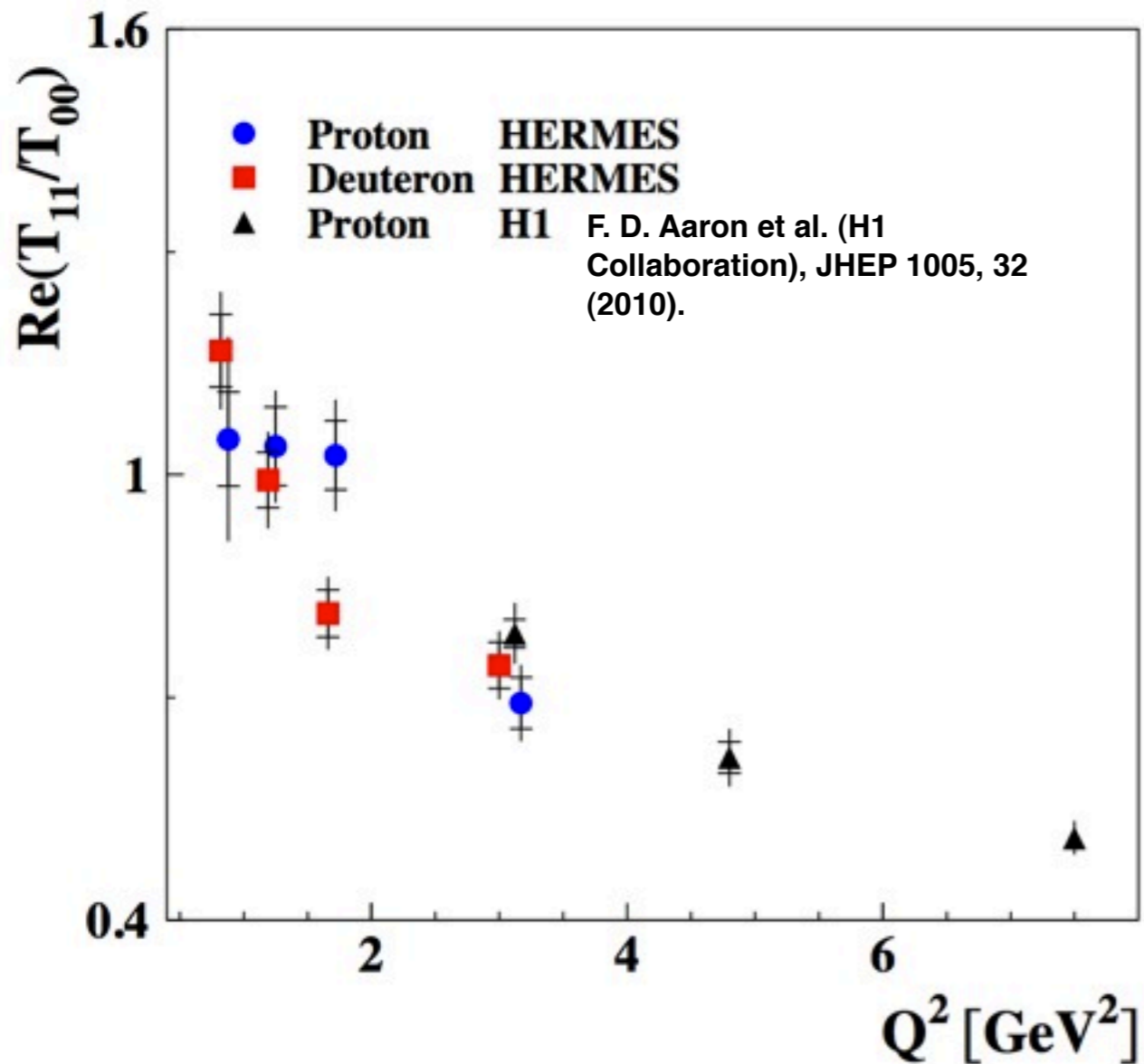
$$\varepsilon = \frac{1 - y - y^2 \frac{Q^2}{4\nu^2}}{1 - y + \frac{1}{4} y^2 \left(\frac{Q^2}{\nu^2} + 2 \right)} \quad \text{the ratio of virtual photon fluxes for longitudinal and transverse polarization}$$

Exclusive Measurement

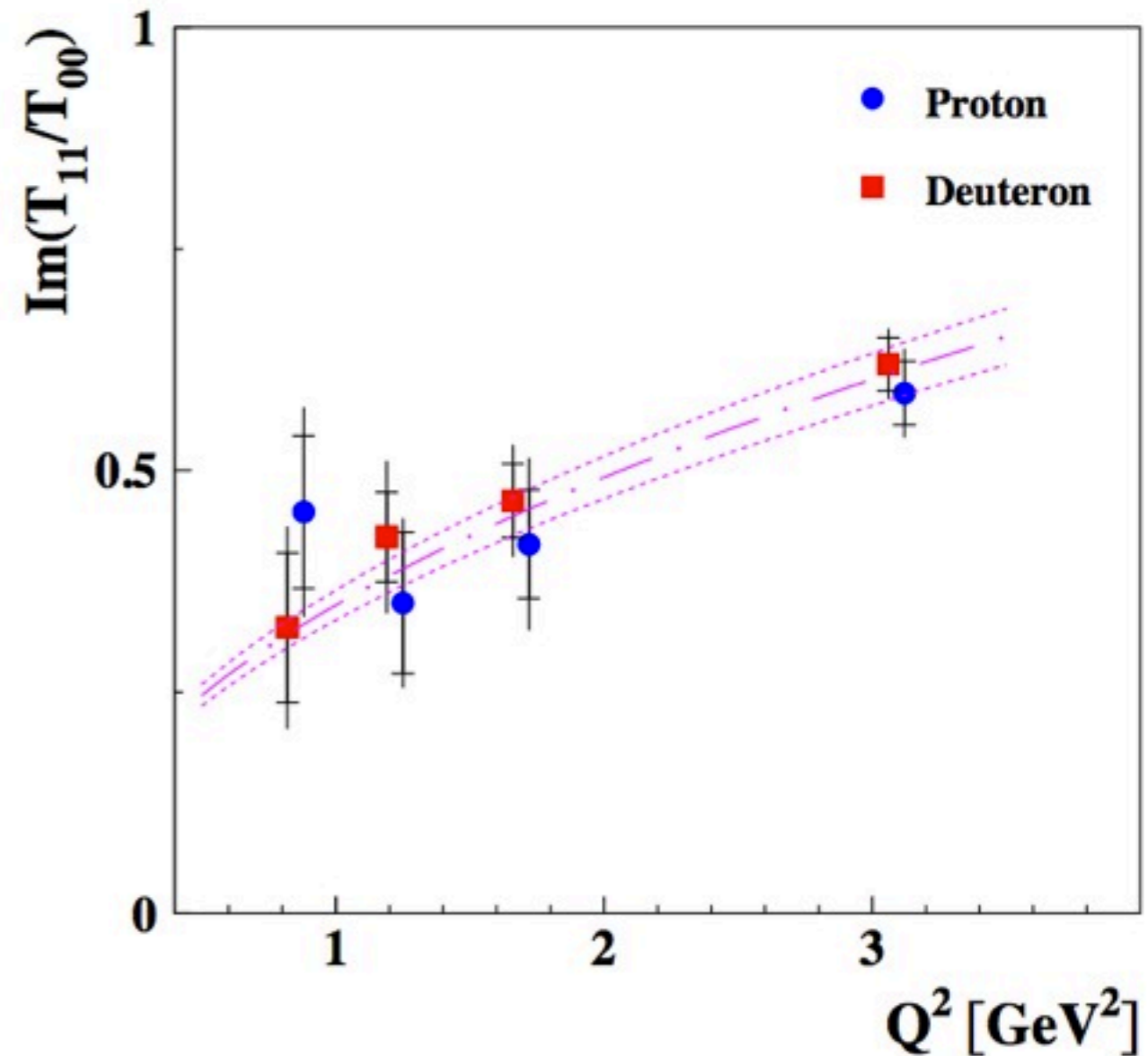




Kinematic Dependence of t_{11}



Real Part follows a/Q
 with $a = 1.11 \pm 0.03 \text{ GeV}$
 as expected!



Imaginary Part follows bQ
 with $b = 0.34 \pm 0.02 \text{ GeV}^{-1}$
 (fit has no basis in theory)

Phase Differences of HARs

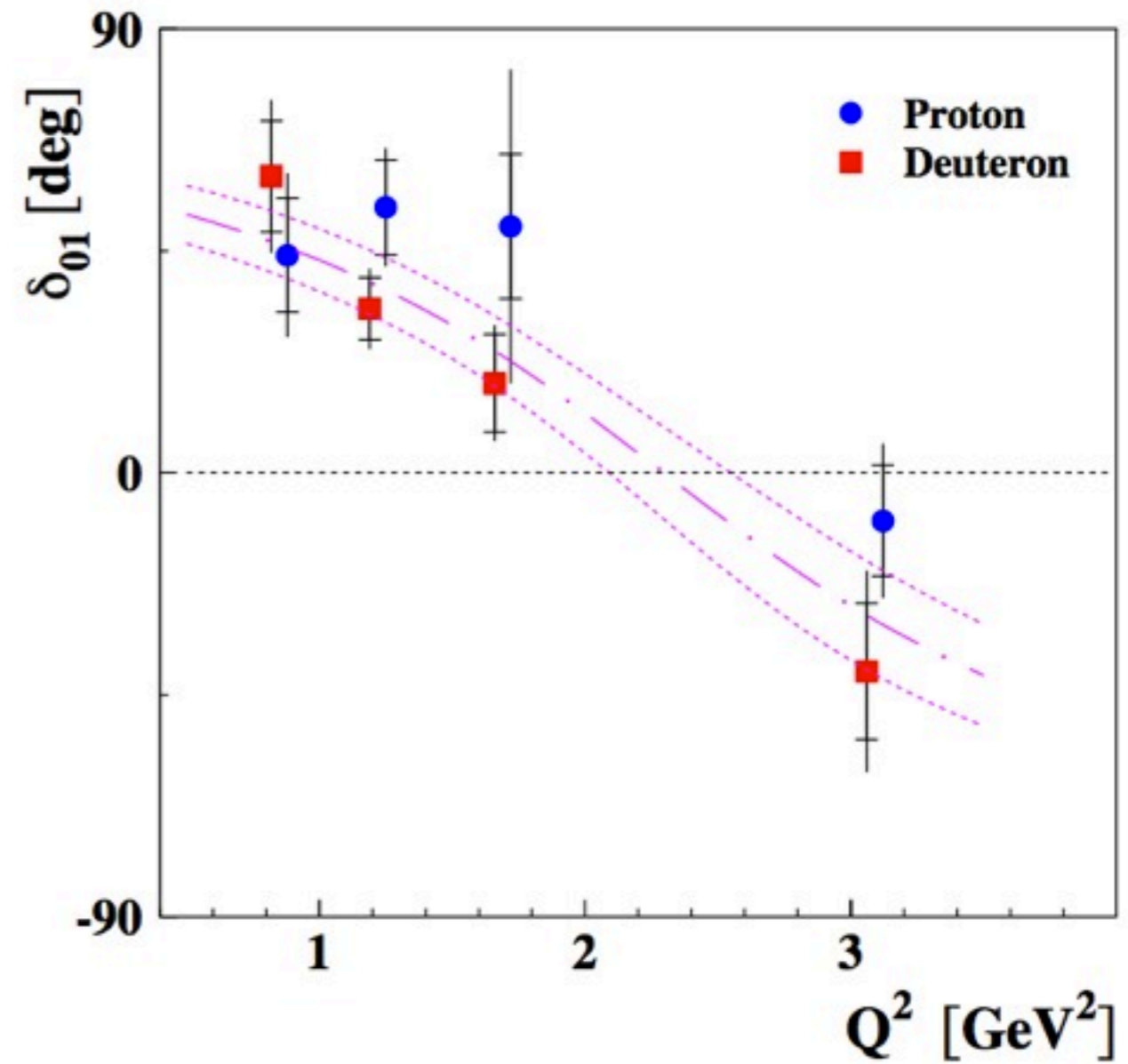
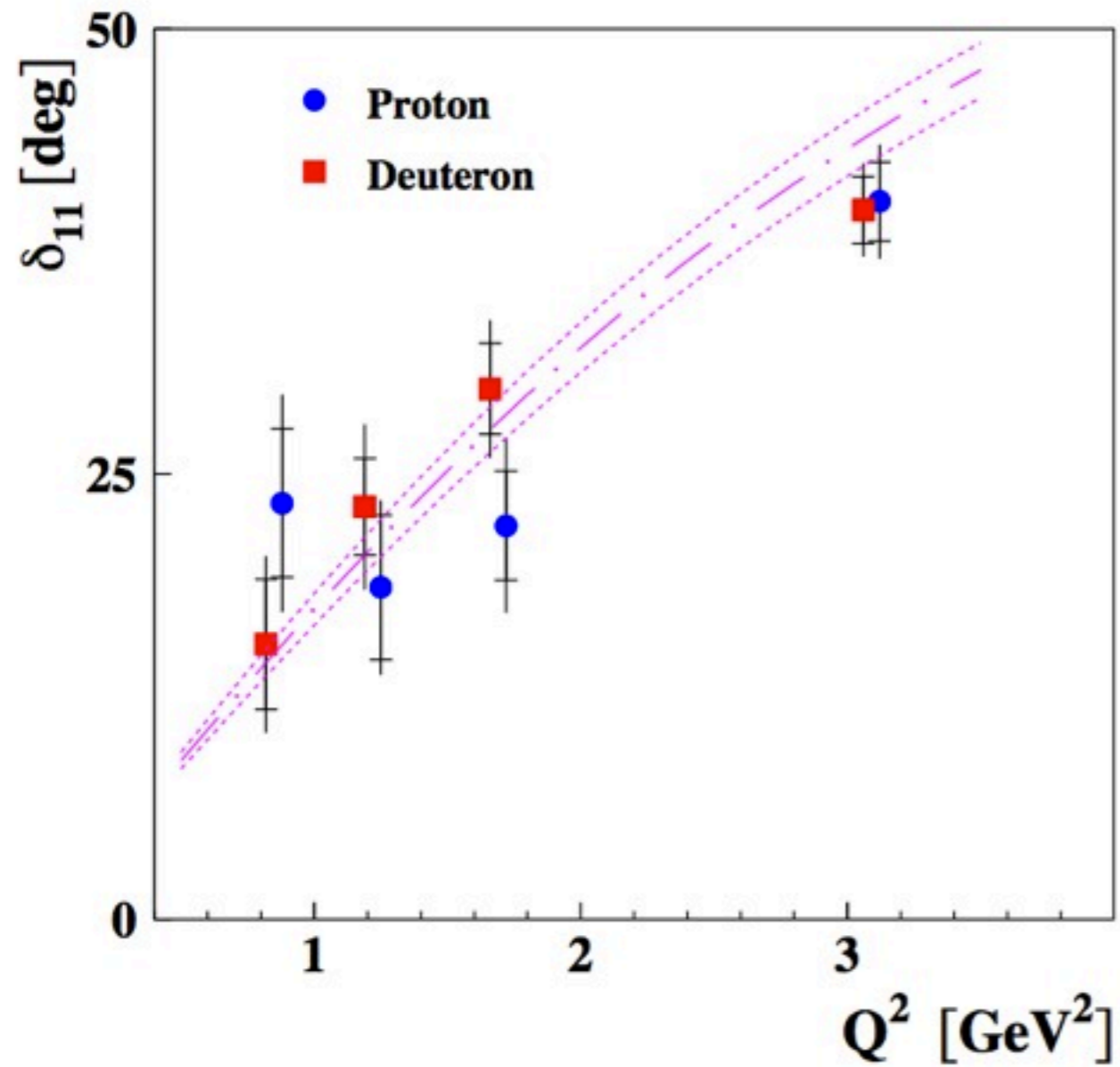
- **GPD** model predicts **small phase difference** for $\tan(\delta_{11}) = \text{Im}(t_{11})/\text{Re}(t_{11})$

[S. V. Goloskokov and P. Kroll, Eur. Phys. J. C 53, 367 \(2008\)](#)

- t_{01} is expected to be the **largest SCHC-violating amplitude** and δ_{01} should be **constant**

[D. Yu. Ivanov and R. Kirschner, Phys. Rev. D 58, 114026 \(1998\)](#)

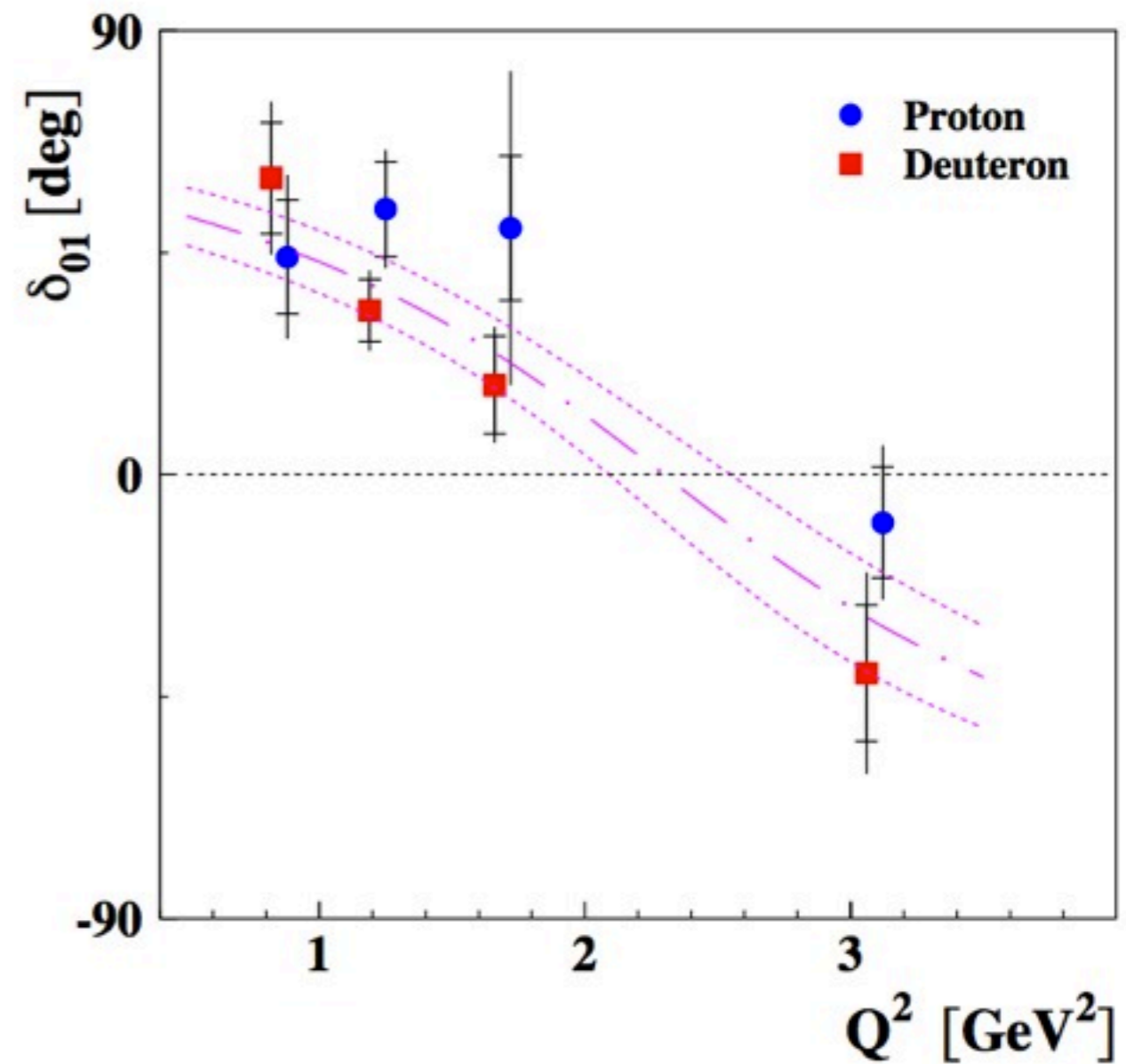
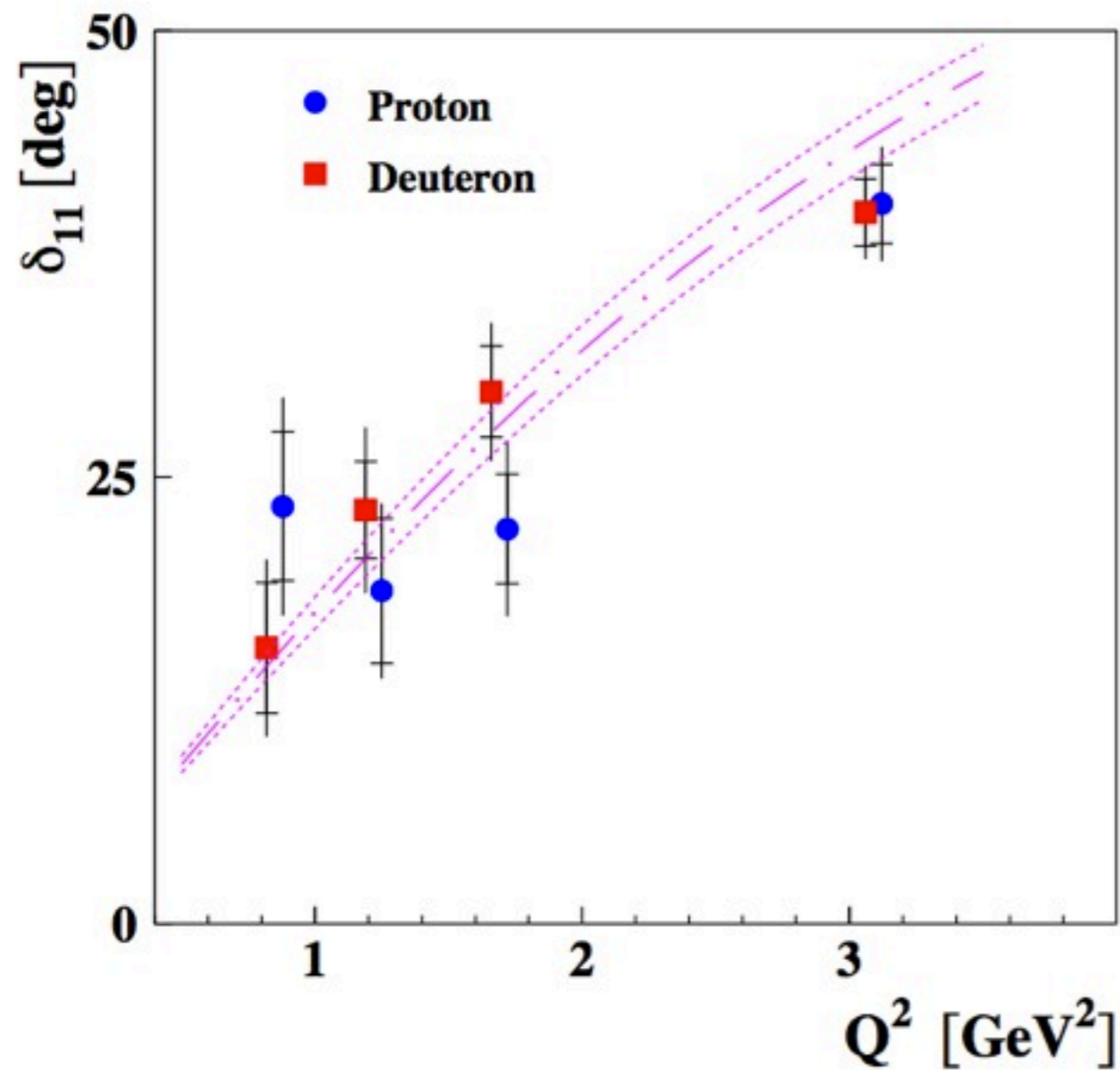
Phase Difference of HARs



Large value **contradicts**
GPD-based models

Should be a **constant**

(Neither $\text{Re}(t_{01})$ nor $\text{Im}(t_{01})$ follow theoretical dependence predictions!!!)



N.B: Fits have no basis in theory

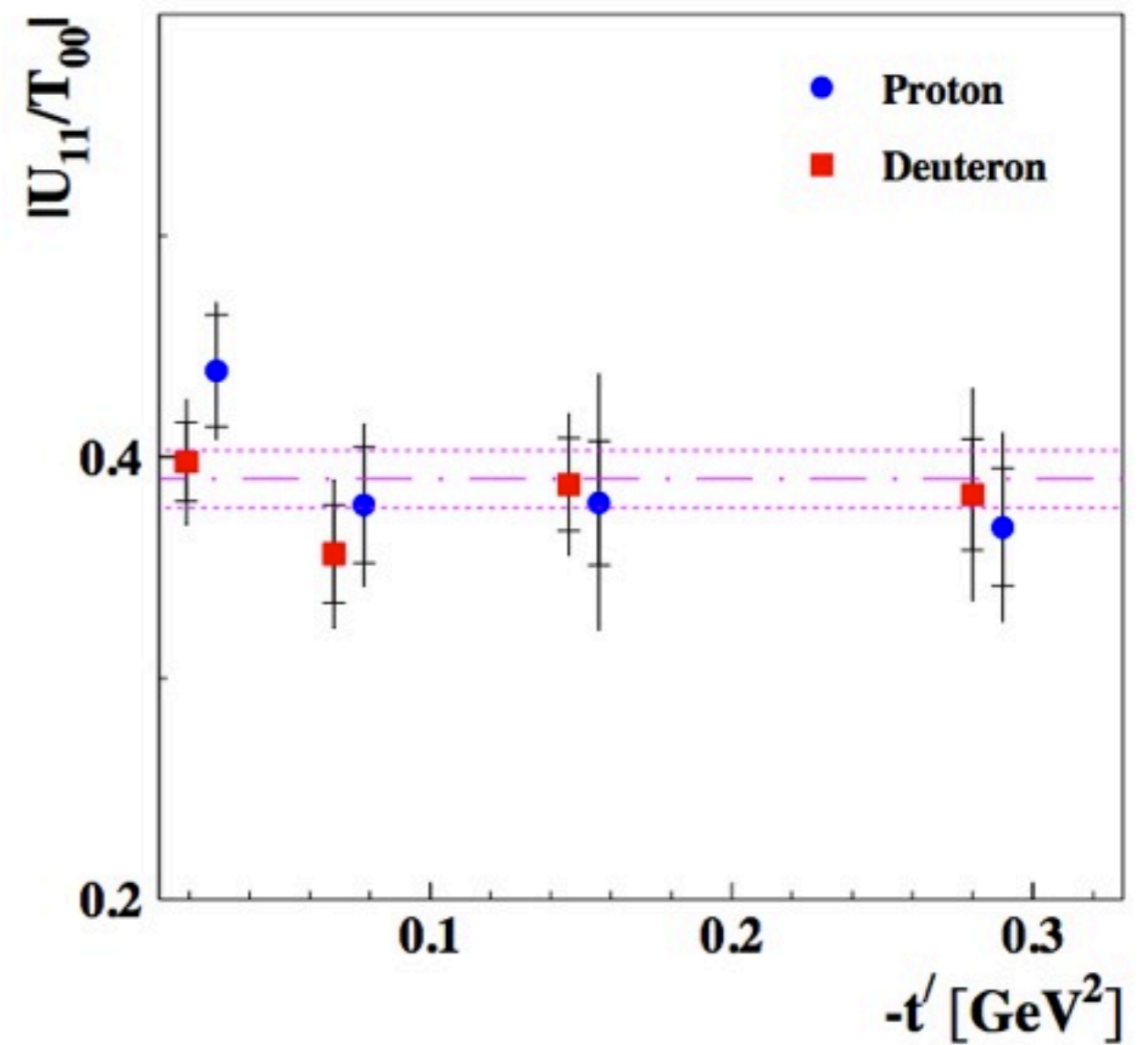
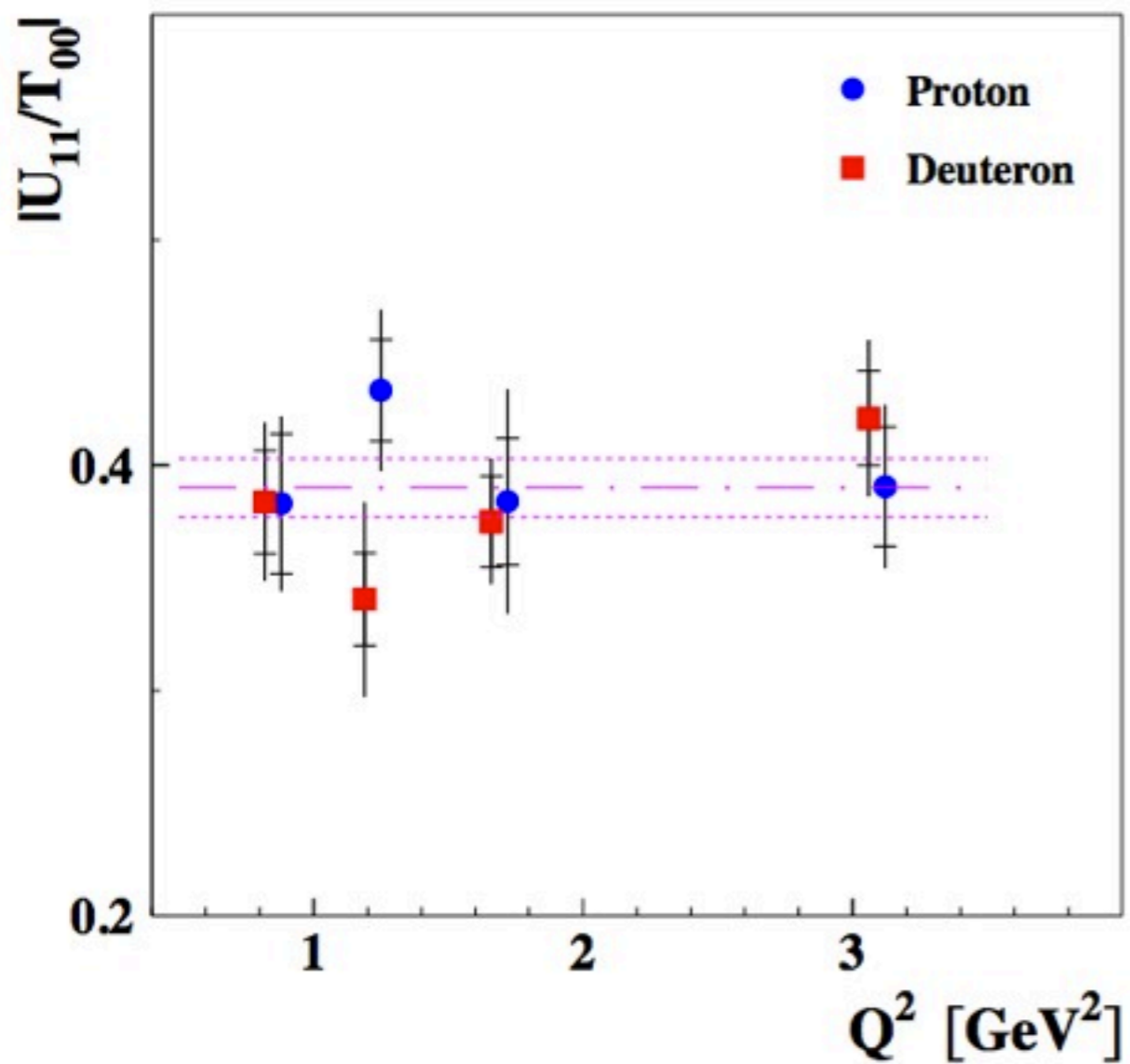
Helicity Amplitude Hierarchy

Behaviour of UPE

$$|T_{00}|^2 \approx |T_{11}|^2 \gg |U_{11}|^2 > |T_{01}|^2 \gg |T_{10}|^2 \dots$$

- $u_{11} = |U_{11}|/|T_{00}|$ should be small ($u_{11} \approx 0.2$) but **visible** (only) for ρ^0 at HERMES!
- May naively expect a $1/Q$ dependence in u_{11}
- UPE is one-pion exchange \Rightarrow may also see some influence of the **pion-pole at small t ?**

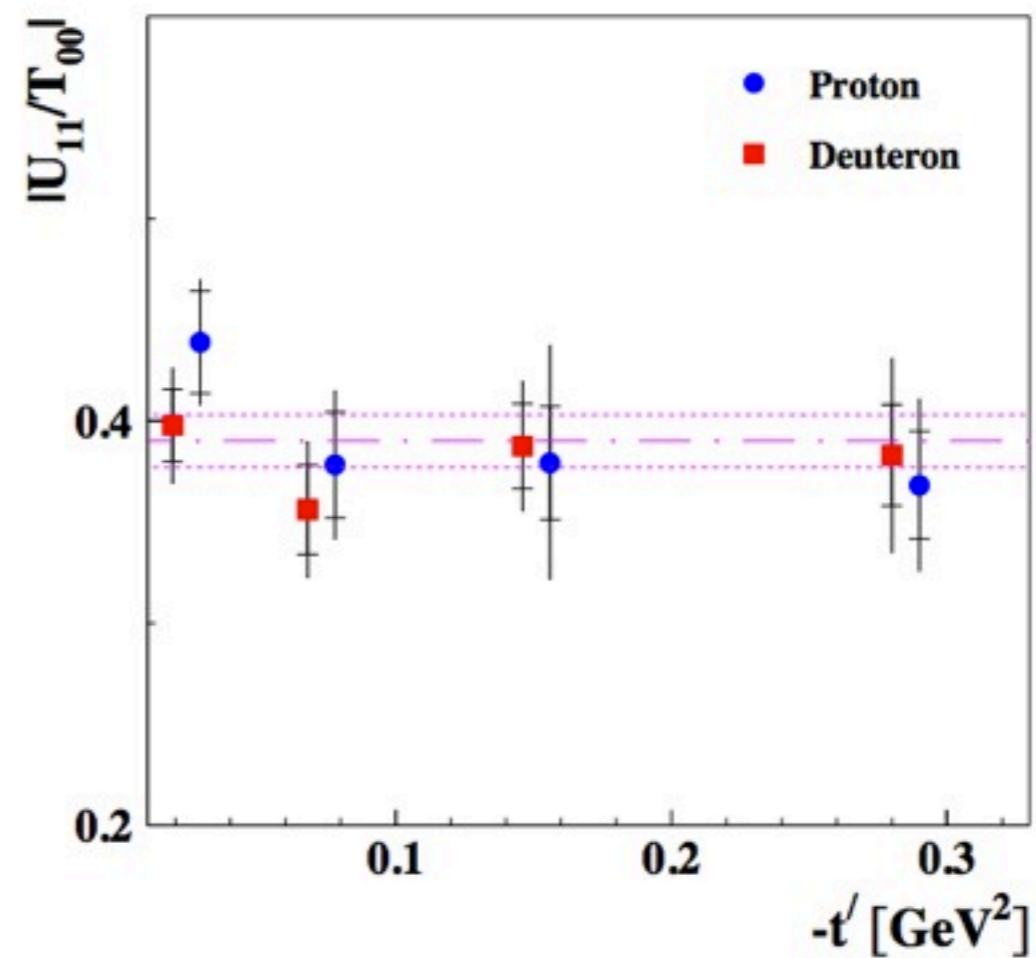
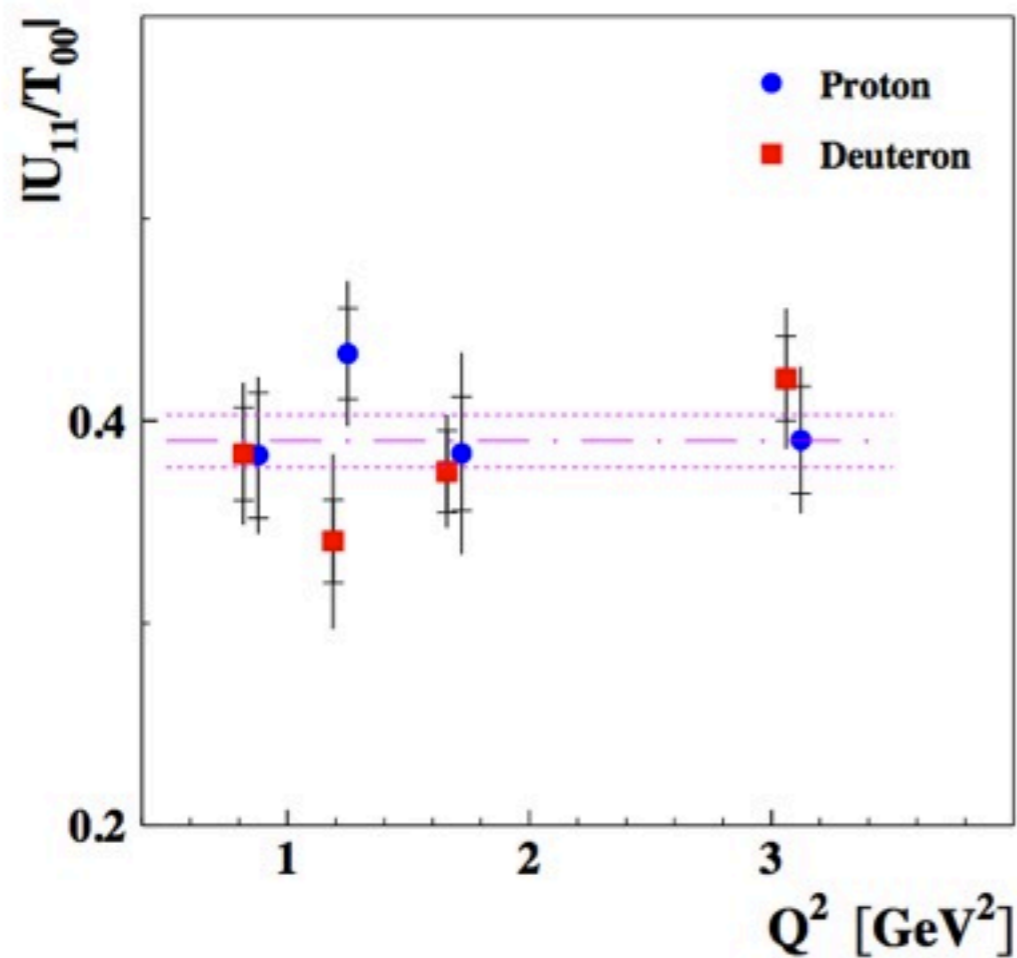
Unnatural Parity Exchange



Unnatural Parity Exchange

No dependence on Q^2 !

No dependence on t' !



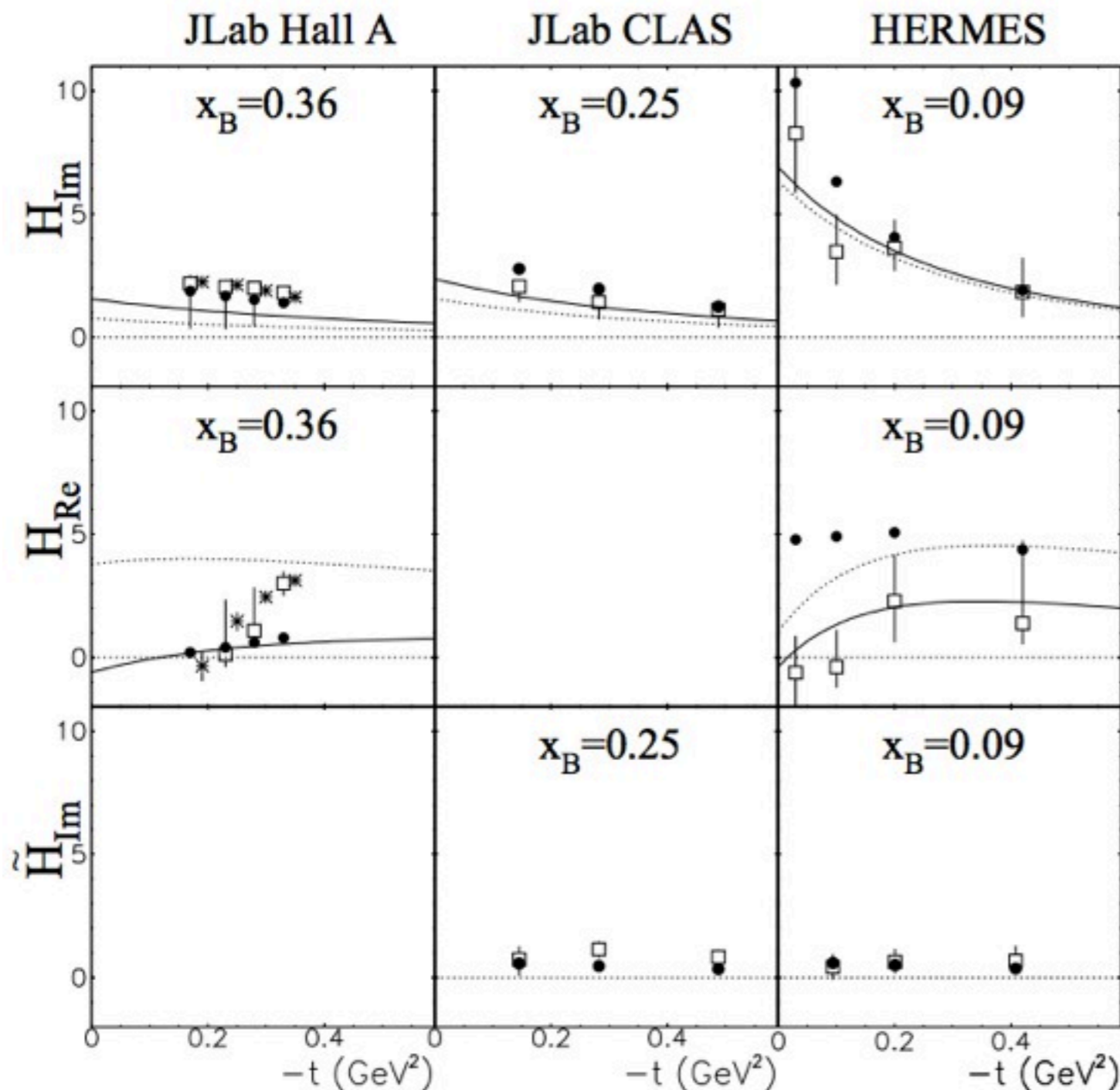
Existence established to 20σ (integrated extraction)

Magnitude of U_{11} is $2.5x$ smaller than T_{00}

Unnatural Parity Exchange

- No dependence on Q^2 may be because **HERMES is far from the asymptotic region ?**
- No dependence on t'
 - ➔ **Too far** from pion-pole ?
 - ➔ U_{11} **not dominated** by one-pion exchange ?
 - ➔ An underlying **dependence of T_{00} on t' ?**

GPD Extraction



Even for H, **VGG** model GPDs are shown **not to be consistent with experimental measurements** when CFFs are extracted from data.

<http://arxiv.org/abs/1011.4195>

Guidal, *ICHEP Procs.* (2010)

<http://arxiv.org/abs/0904.1648>

H. Moutarde, *Phys. Rev. D79* (2009)

<http://arxiv.org/abs/0904.0458>

Kumerički and Müller, *Nucl. Phys.* **B841** (2010)