

# Flavor Physics: open problems & recent developments

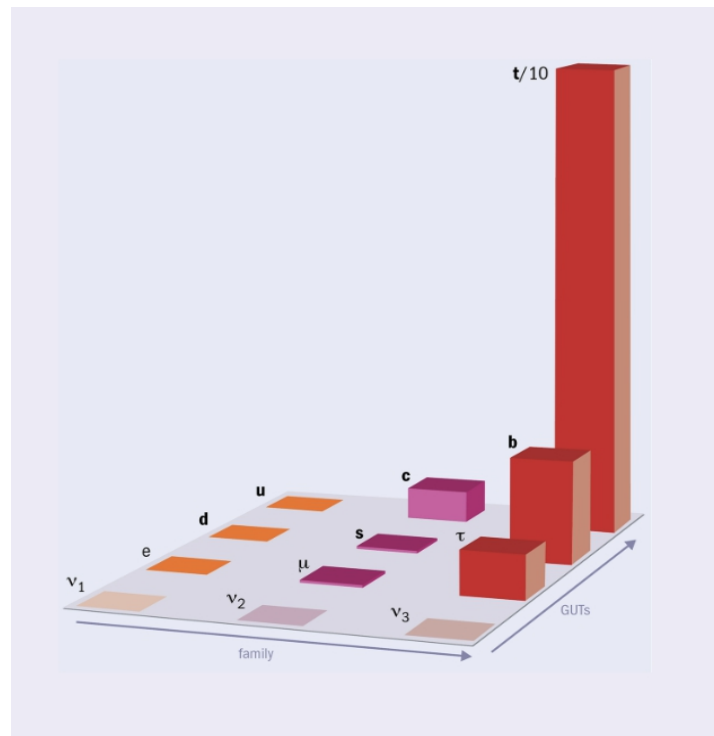
Gino Isidori  
[ *University of Zürich* ]



University of  
Zurich<sup>UZH</sup>

- ▶ Introduction: the flavor puzzle
- ▶ The flavor structure of the SMEFT
- ▶ Flavor non universality & flavor deconstruction
- ▶ A look to data & future prospects
- ▶ Conclusions

## The flavor puzzle



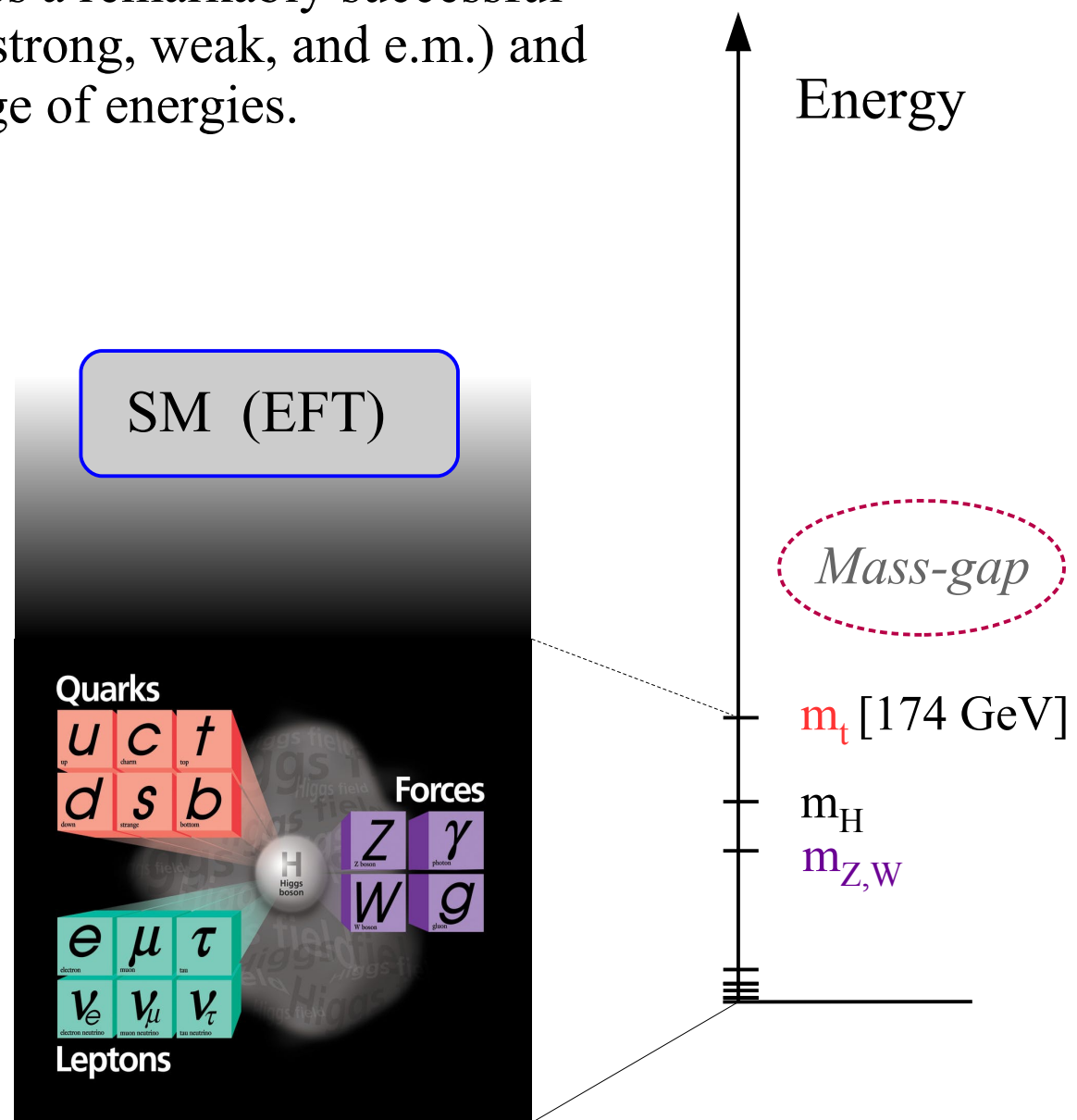
## ► Introduction

The Standard Model (SM) provides a remarkably successful description of *fundamental forces* (strong, weak, and e.m.) and *matter constituents*, over a wide range of energies.

However, as for any QFT, it is natural to consider the SM as an Effective Field Theory, i.e. the low energy limit of a more complete theory with more degrees of freedom

$$\mathcal{L}_{\text{SM-EFT}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \dots$$

We identified the *long-range* properties of this EFT



## ► Introduction

There are several reasons why we think the SM must be extended at high energies:

Electroweak hierarchy problem

Flavor puzzle

U(1) charges

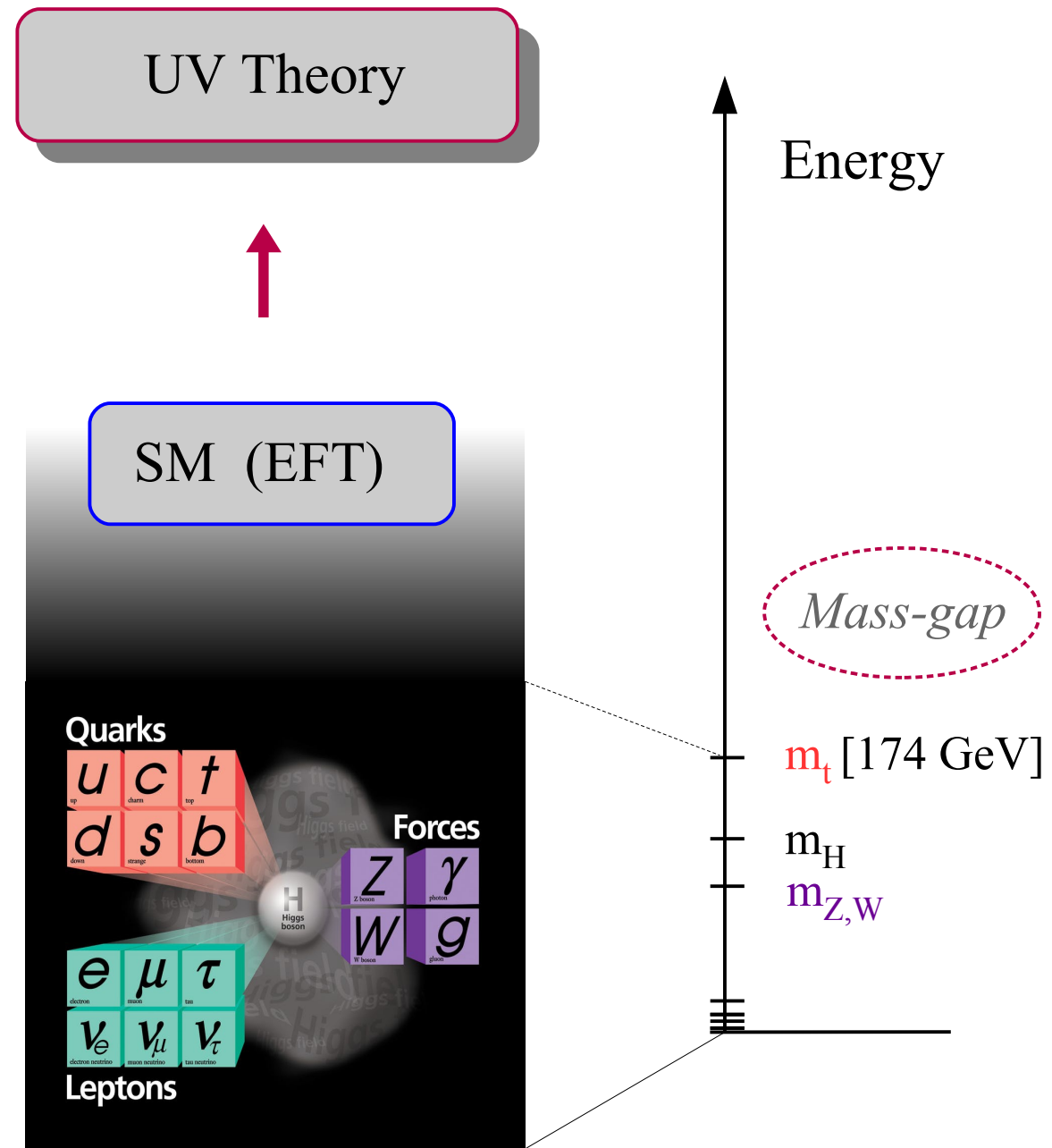
Neutrino masses

Dark-matter

Dark-energy

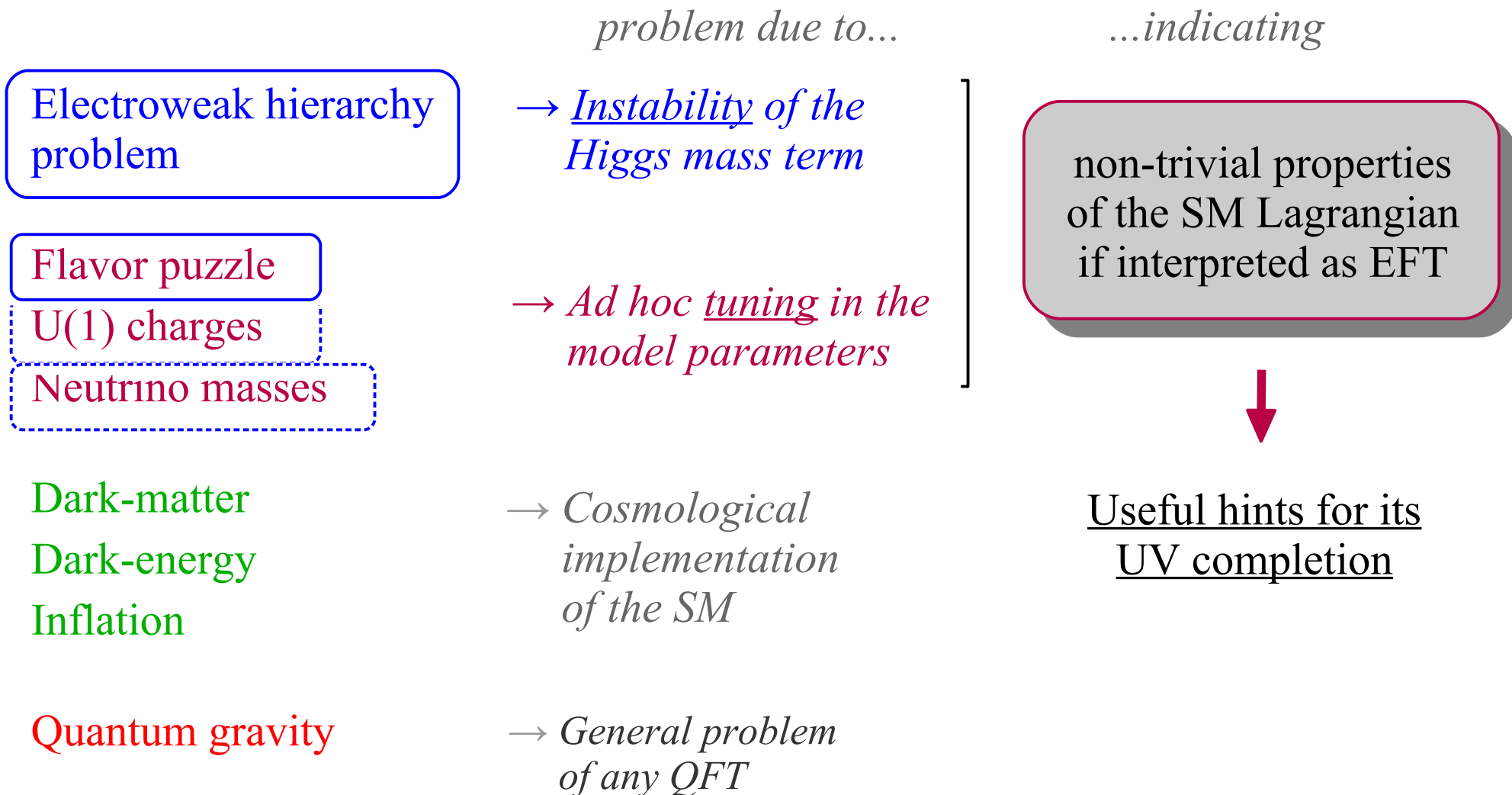
Inflation

Quantum gravity



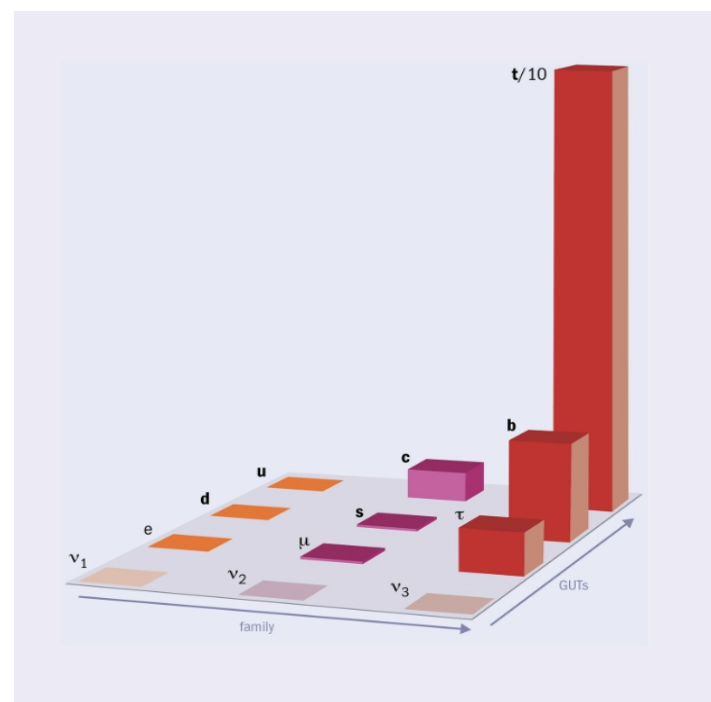
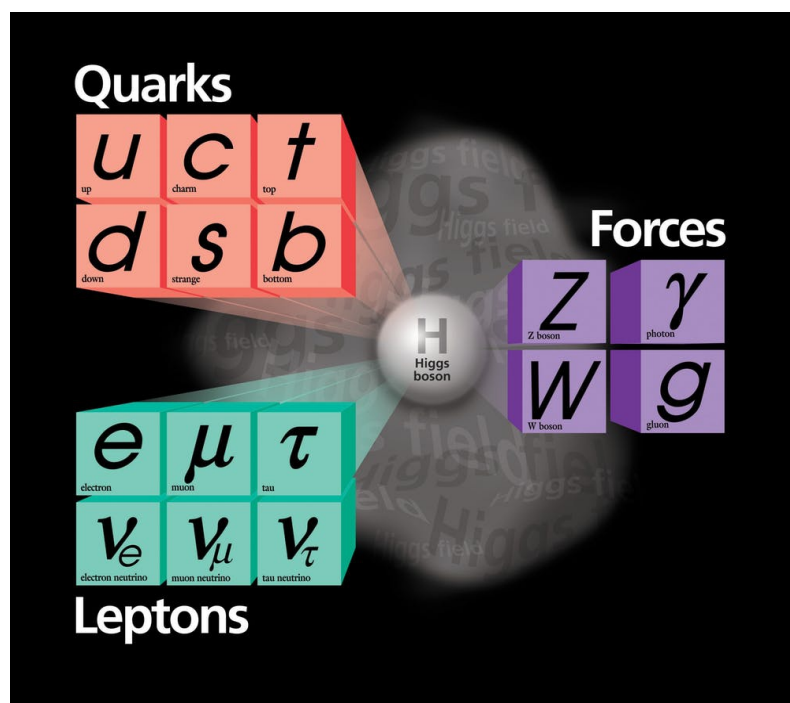
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## ► Introduction: the flavor puzzle

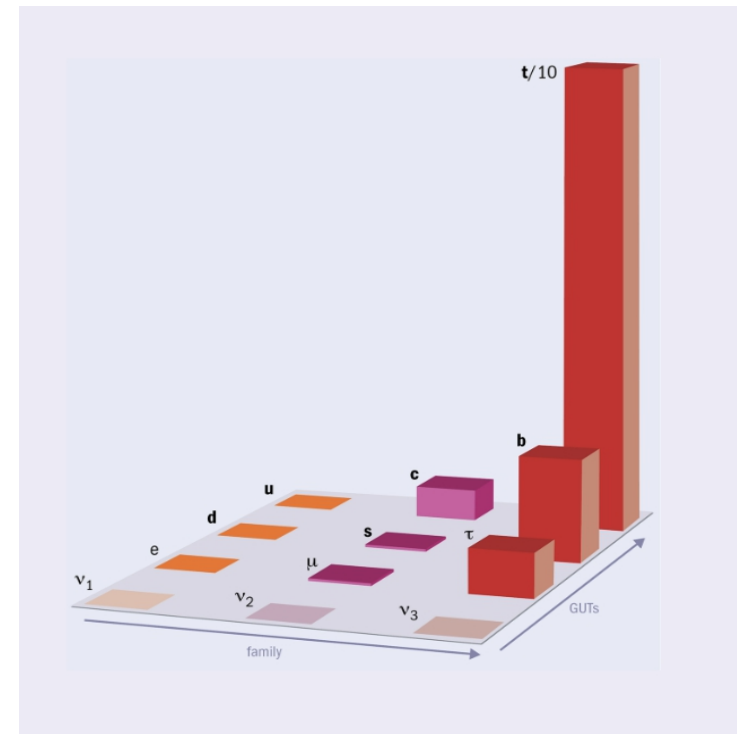
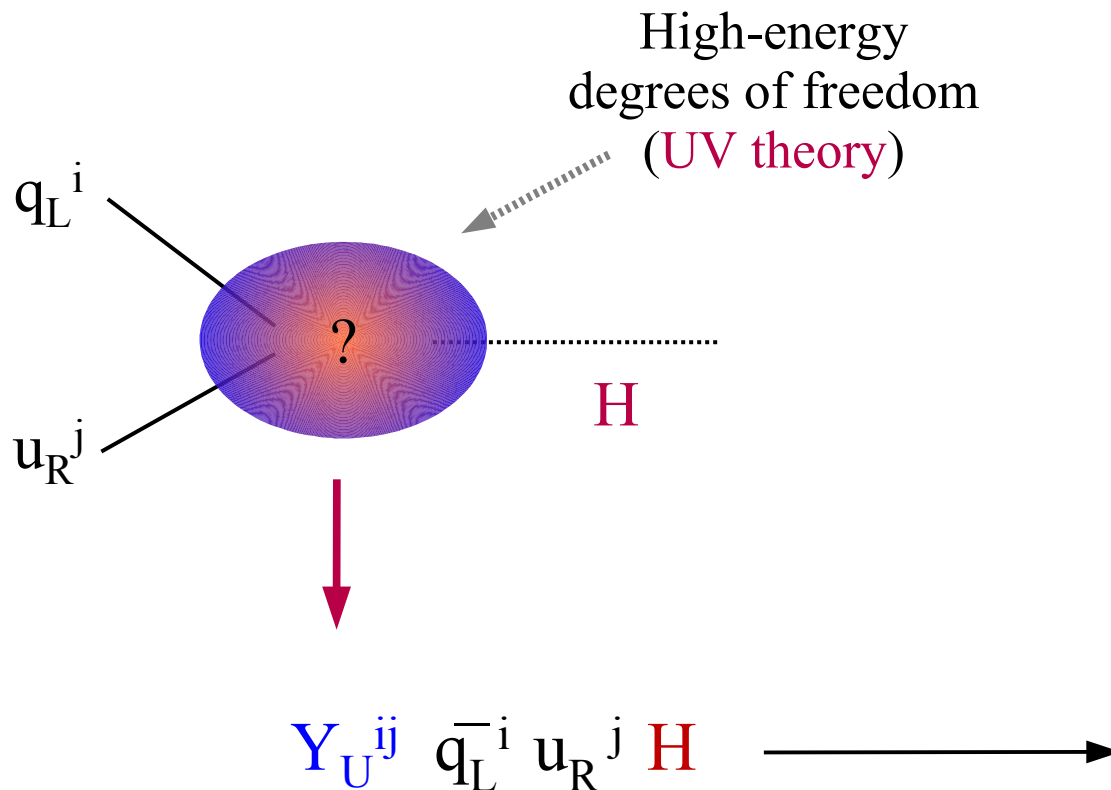
The “flavor puzzle” is an old problem, that emerged well before the Standard Model (SM) was conceived [ *“Who ordered the muon ?”* - I. Rabi (~1950) ].



- Why do we have 3 copies of fermions with identical gauge quantum numbers
- What determines their highly non-trivial mass matrices?

► Introduction: the flavor puzzle

Fermion masses are the results of the Yukawa interactions → [Inescapable link between Higgs and flavor](#), whose origin can be addressed only **beyond the SM...**

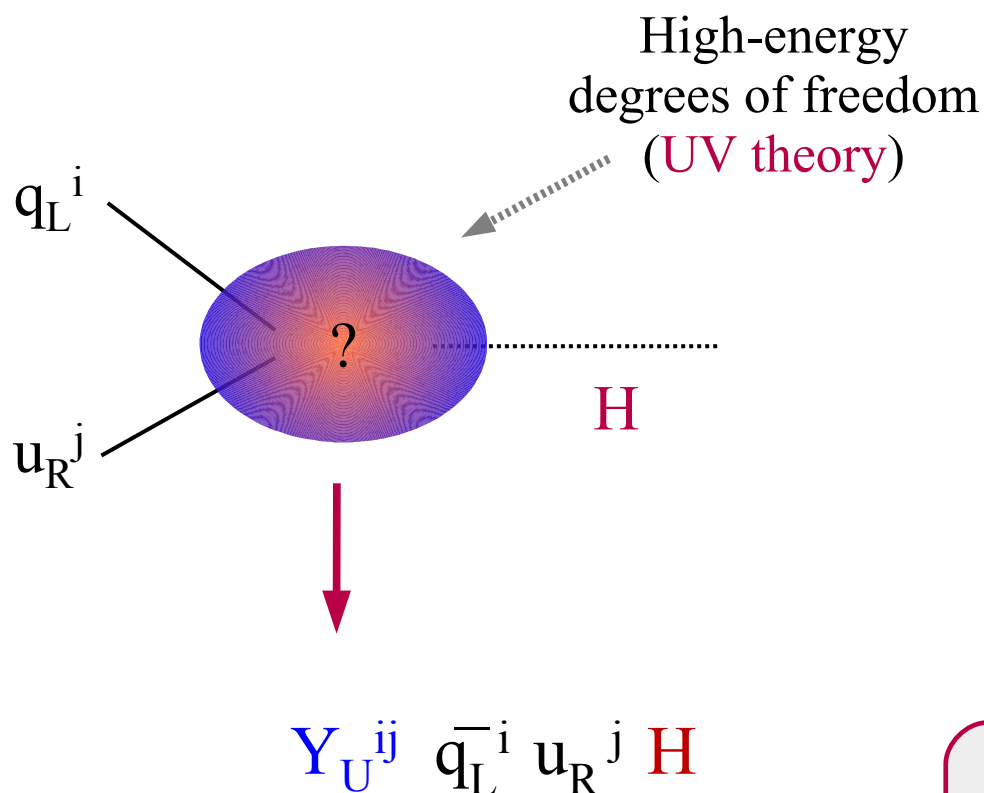


“Message from the UV”  
that we need to “decode”



## ► Introduction: the flavor puzzle

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Contrary to gauge interactions, which are “protected” by gauge symmetries [*universal couplings controlled only by field charges*] the Yukawa interactions provide more information about UV dynamics

The UV dynamics responsible for the Higgs-fermion couplings is not universal among generations

## ► Introduction: the flavor puzzle

Fermion masses are the results of the Yukawa interactions → [Inescapable link between Higgs and flavor](#), whose origin can be addressed only **beyond the SM**...

E.g.:

$$Y_U \sim \begin{pmatrix} \square & \square & 0.003 \\ & & 0.04 \\ & & \boxed{1} \end{pmatrix}$$

$< 0.01$

$Y_U$  in the basis  
where  $Y_D$  is diagonal

$$y_u = \frac{\sqrt{2} m_u}{\langle H \rangle} \approx 10^{-5}$$

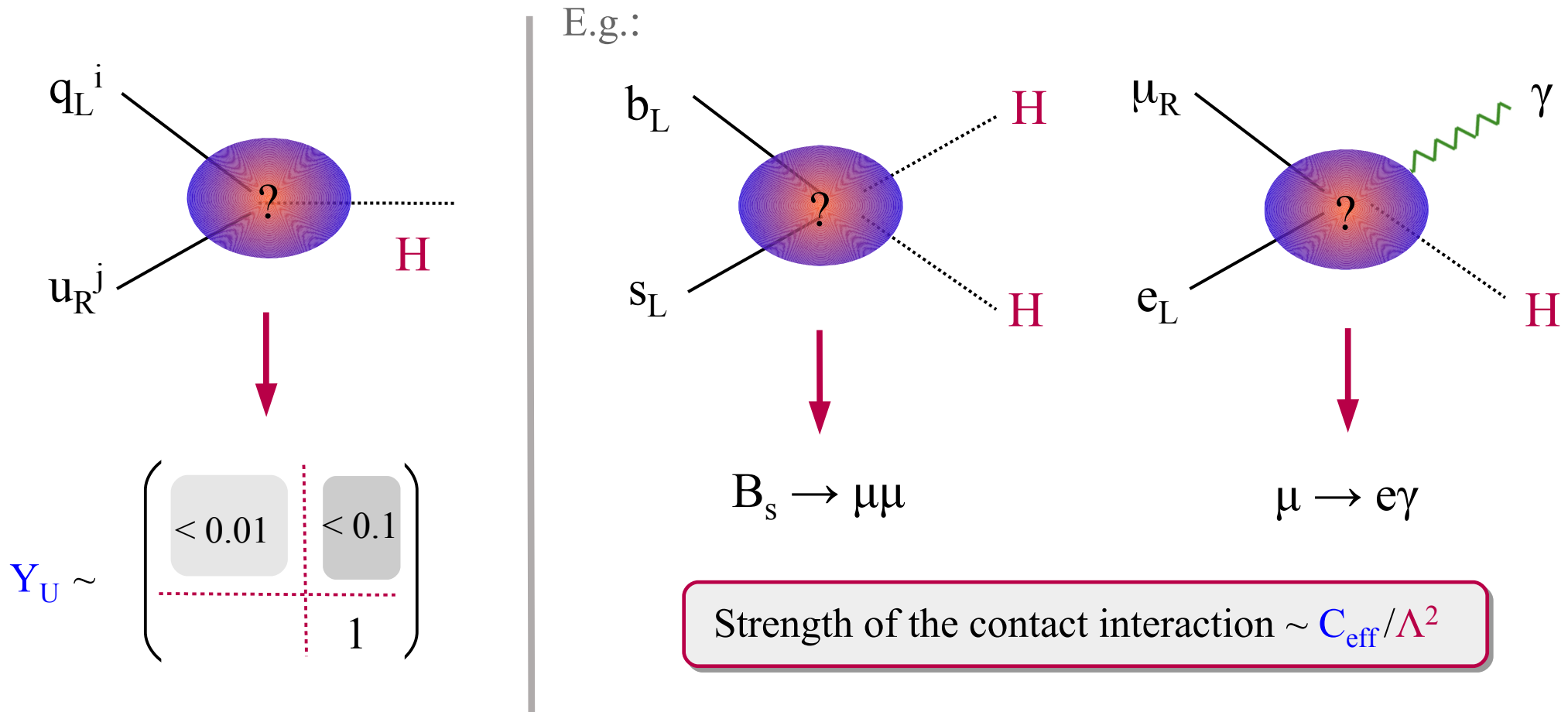
$$y_t = \frac{\sqrt{2} m_t}{\langle H \rangle}$$

The UV dynamics responsible for the Higgs-fermion couplings is not universal among generations & exhibits a strongly hierarchical pattern



## ► Introduction: the flavor puzzle

On general grounds, the flavor non-universal dynamics in the UV should give rise to new effective flavor-violating contact interactions

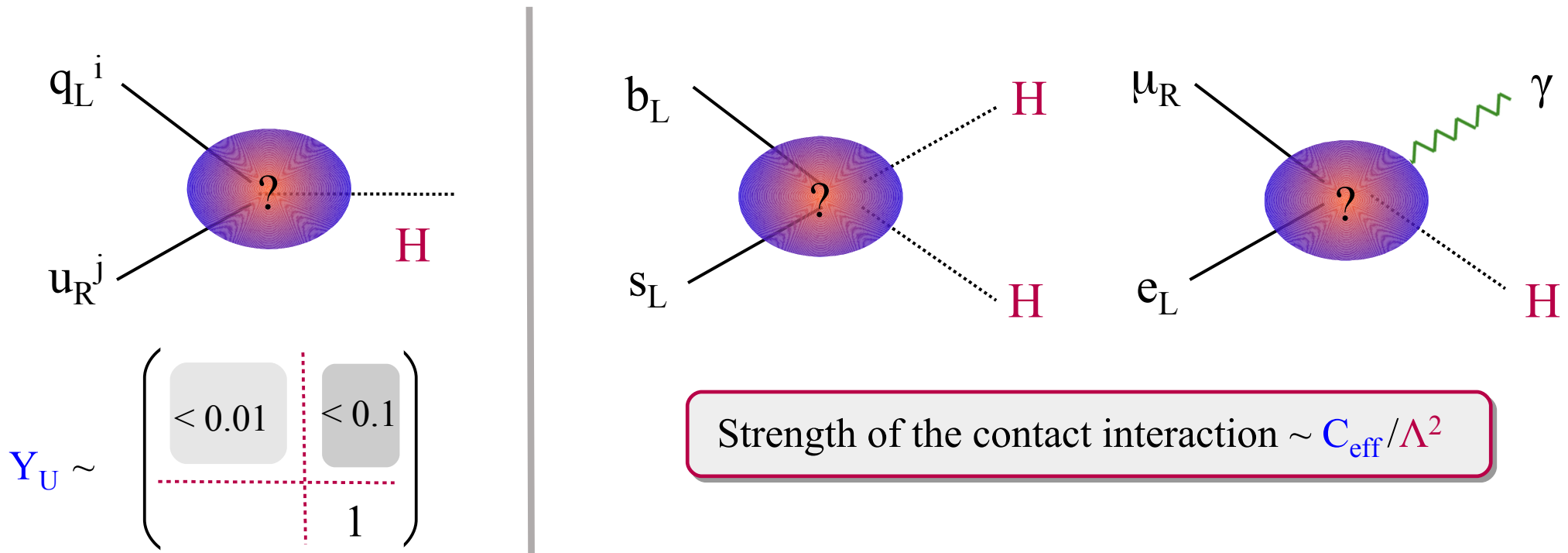


At present we do not see any of these effects.

This is often expressed in terms of bounds on  $\Lambda$  exceeding the **TeV scale** by several orders of magnitude

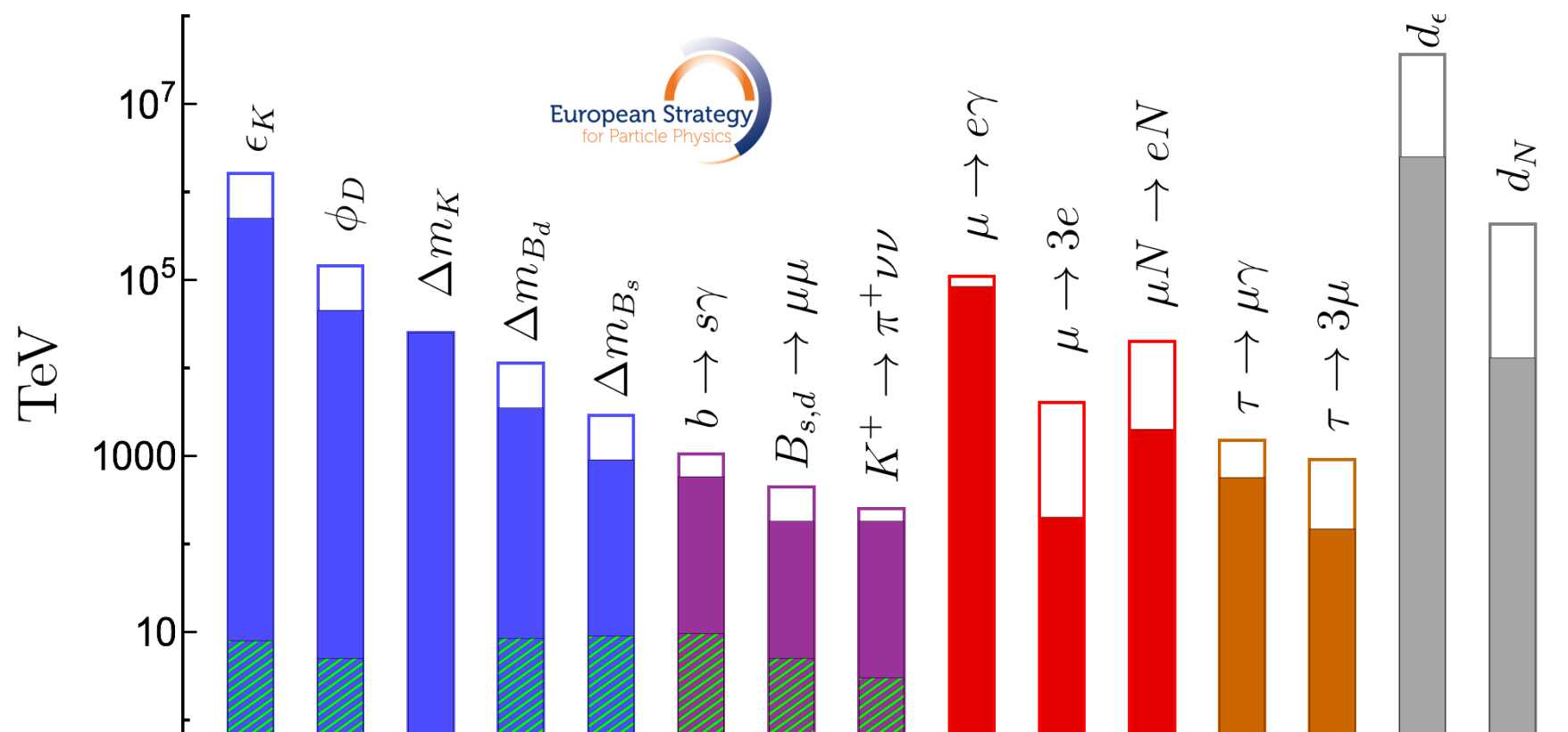
## ► Introduction: the flavor puzzle

*The two faces of the flavor puzzle:*



- *Which is the origin of this peculiar structure?*
- *Why we do not see other signs of flavor non-universality?*
- *Is there a connection between these two problems?*

## The flavor structure of the SMEFT



► The flavor structure of the SMEFT

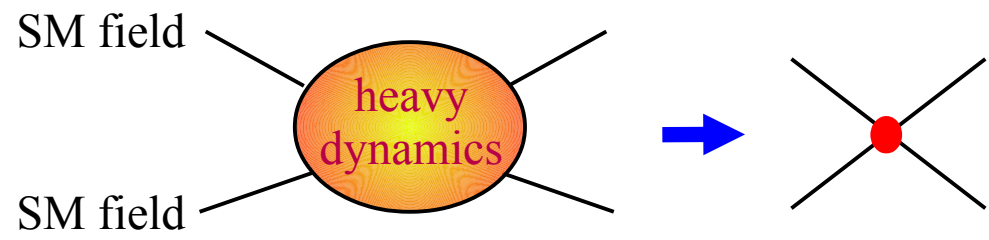
$$\mathcal{L}_{\text{SM-EFT}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \sum_{d,i} \frac{c_i^{[d]}}{\Lambda^{d-4}} \mathcal{O}_i^{d \geq 5}$$

Interactions surviving @ large distances  
(operators with  $d \leq 4$ )

Long-range forces  
of the SM particles  
+  
ground state (Higgs)

Local contact interactions  
(operators with  $d > 4$ )

“Remnant” of the heavy  
dynamics at low energies

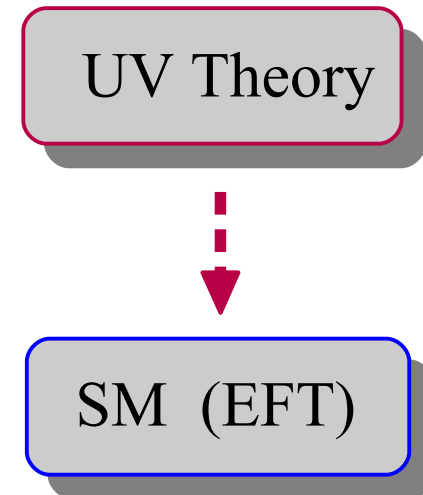


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What is the cut-off scale  $\Lambda$  of the SMEFT?

A useful (*but somewhat vague*) indication follows from the **Electroweak hierarchy problem** ( $\leftrightarrow$  *instability of the Higgs mass under quantum corrections*):



$$\text{---} \bullet \text{---} + \text{---} \bigcirc \text{NP} \text{---} \rightarrow m_H^2 \Big|_{\text{Phys}}$$

$\Delta m_H^2 \sim \Lambda^2$

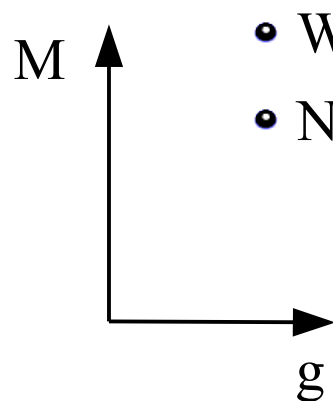
➔ (some) **New Physics** (coupled at least to H & t) in the TeV domain

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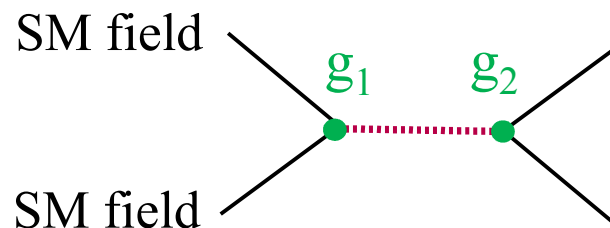
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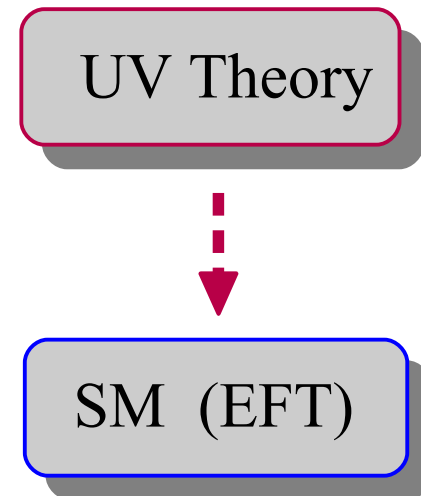
A closer look to this question reveals more “layers”



- What is the mass scale of the new d.o.f. ?
- New dynamics weakly or strongly coupled ?



$$\frac{g_1 g_2}{M^2}$$

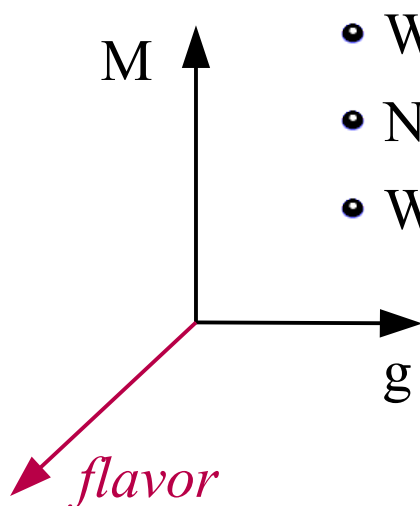
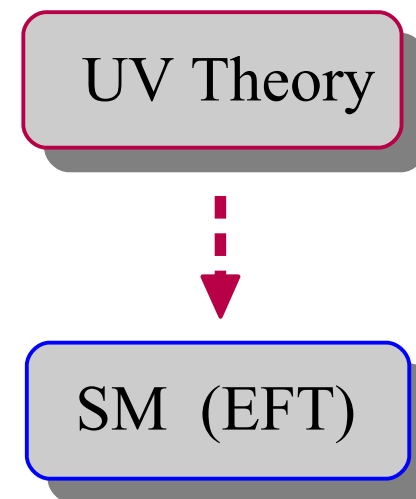


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- What is the mass scale of the new d.o.f. ?
- New dynamics weakly or strongly coupled ?
- What is the flavor structure?

SM (Yukawa) sector → flavor is highly non trivial

- No flavor symmetry → 2499 free couplings in the SMEFT @ d=6
- Exact  $U(3)^5$  → 47

Jenkins, Manohar, Trott '14

► The flavor structure of the SMEFT

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Large flavor symmetry

Three identical replica of  
the basic fermion family  
[ $U(3)^5$  symmetry]

$$\psi^i \rightarrow U^{ij} \psi^j$$

$$\Psi = Q_L, u_R, d_R, L_L, e_R$$

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[U(3)<sup>5</sup> symmetry]

$$y_{ij} \psi_L^i \psi_R^j H \rightarrow m_{ij} \psi_L^i \psi_R^j$$

“Peculiar” breaking structure

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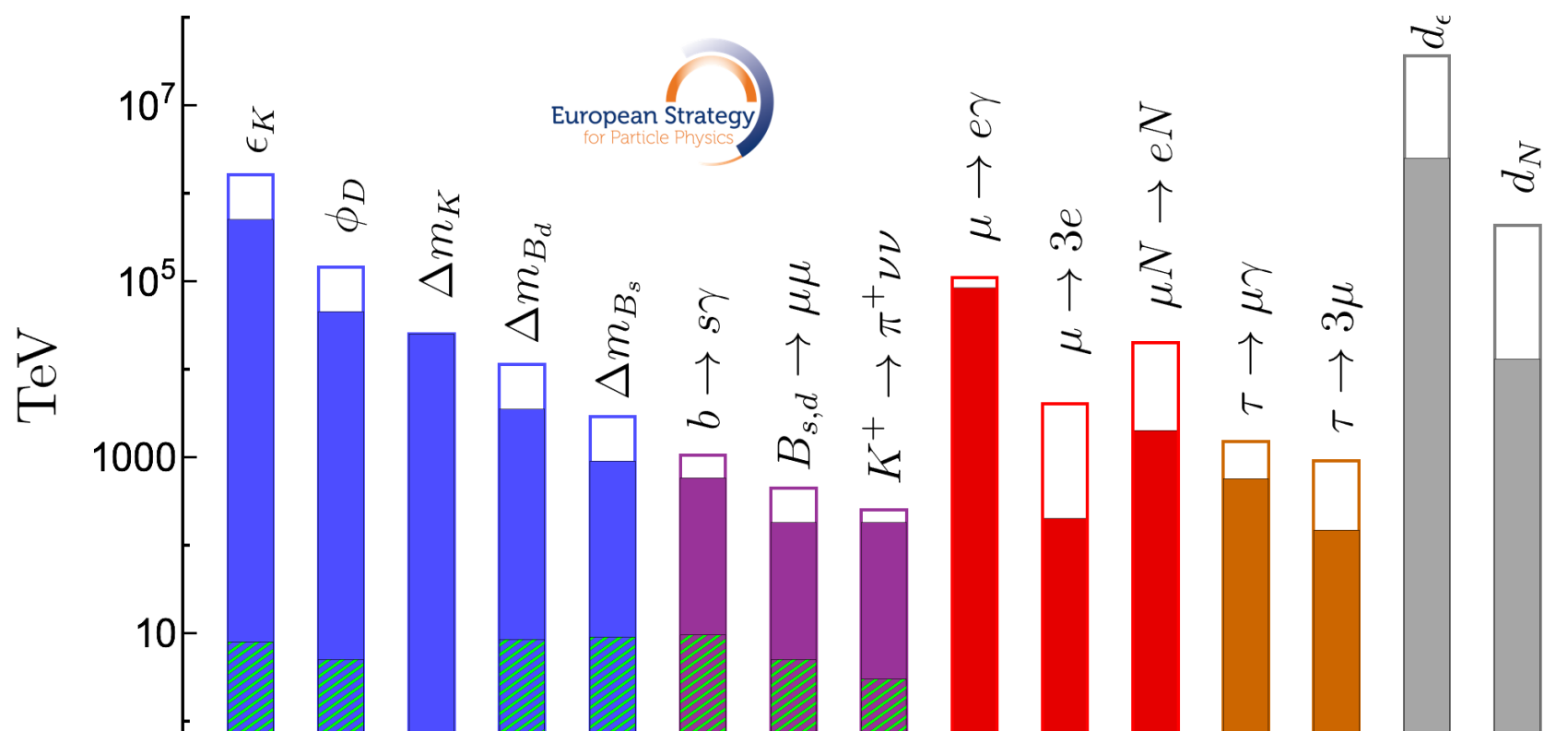
Exact & approximate (*accidental* ?) symmetries

How are these symmetries broken in the whole SMEFT ?

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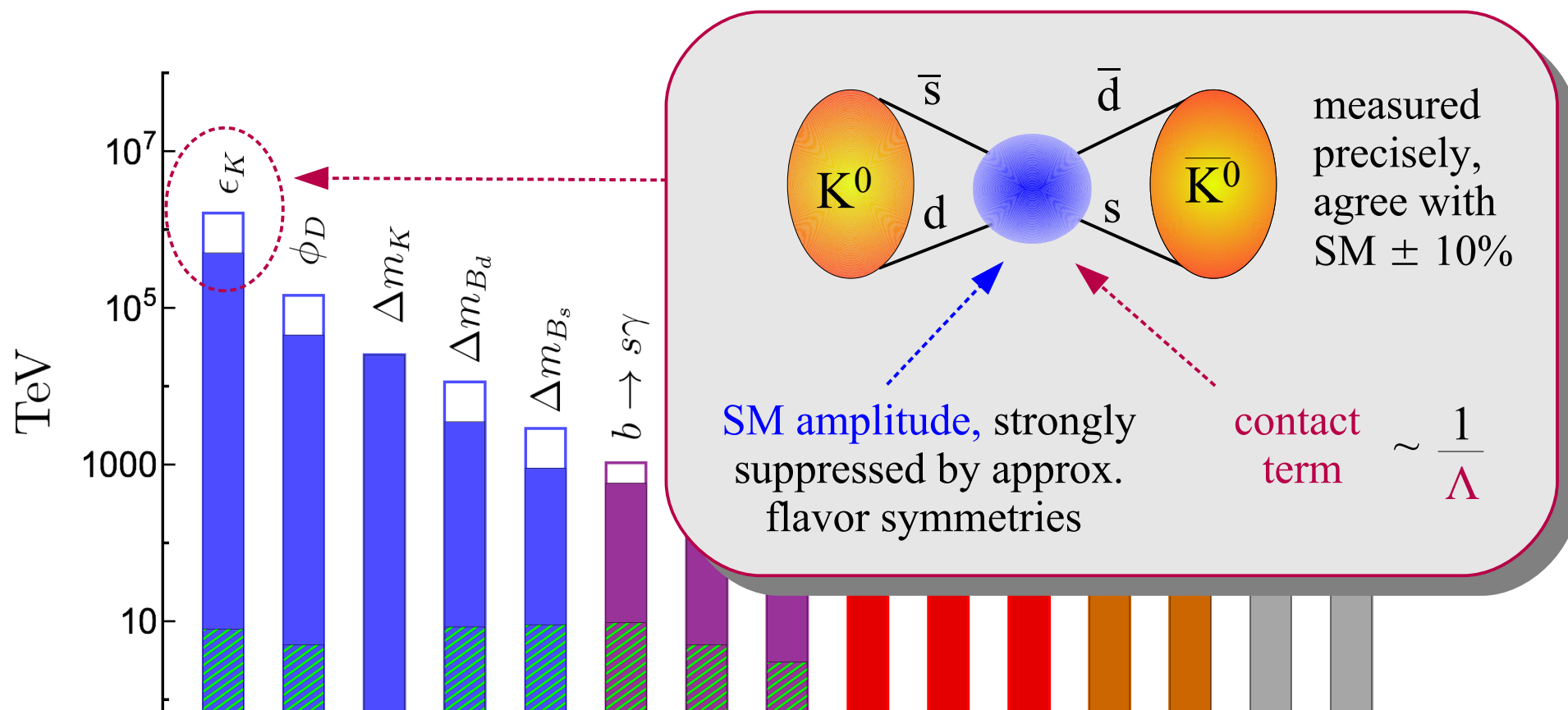
In principle, we could expect many violations of the accidental symmetries from heavy dynamics. However, no clear deviations observed so far → stringent bounds on the *effective scale* of possible new flavor violating interactions



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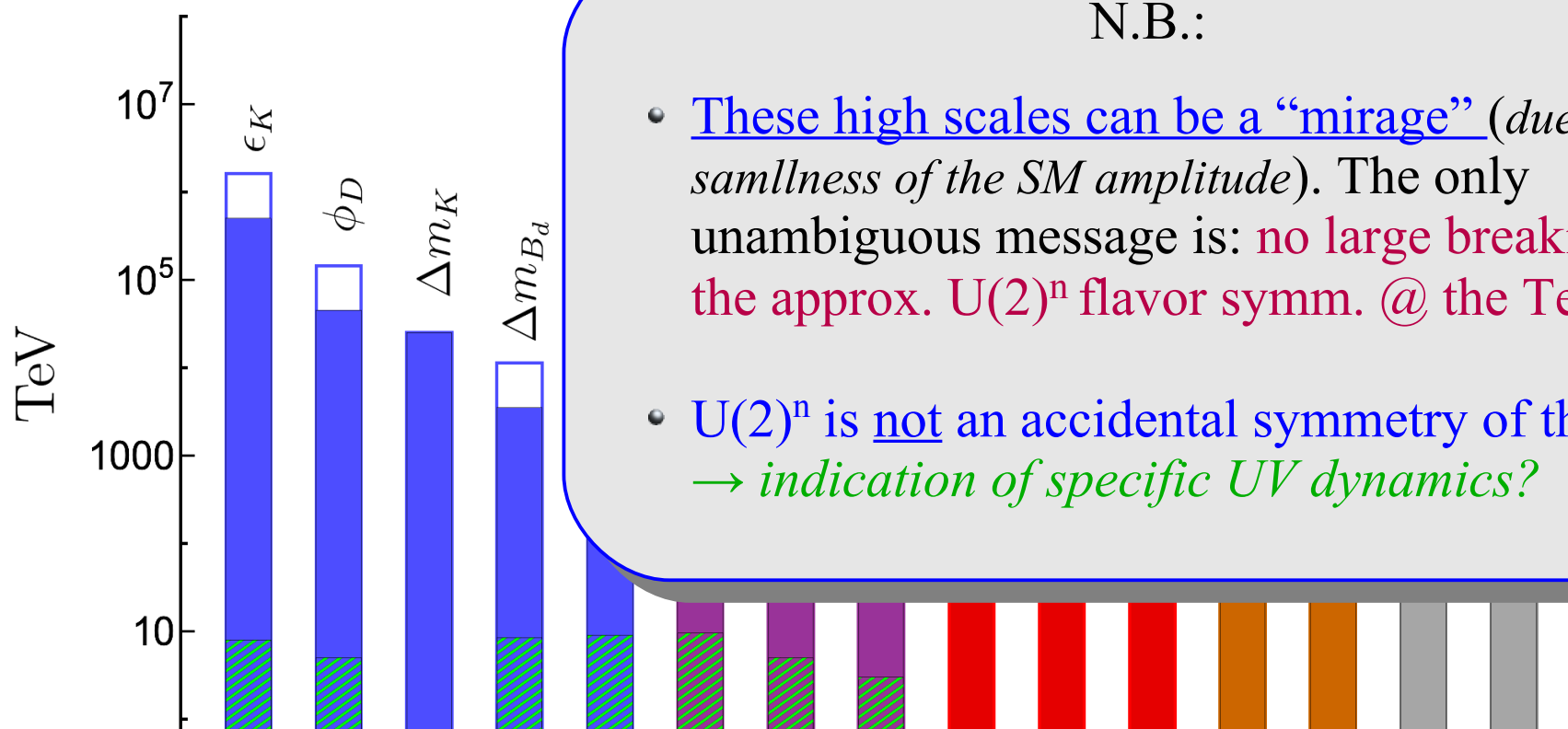
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Flavor-degeneracy:  
 $U(3)^5$  symmetry

Yukawa couplings:  
 $U(3)^5 \rightarrow \sim U(2)^n$   
*peculiar breaking of  
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**Stringent bounds**  
 on additional symmetry  
 breaking terms

3 classes of “interpretations”



- “Simplicity”: the scale of NP is high ( $> 10^5$  TeV)  
 → no way to test the origin of the Y's
- “High-scale flavor dynamics” & MFV
- “Flavor deconstruction”

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No “QFT solution” to  
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“decoupling” of flavor &  
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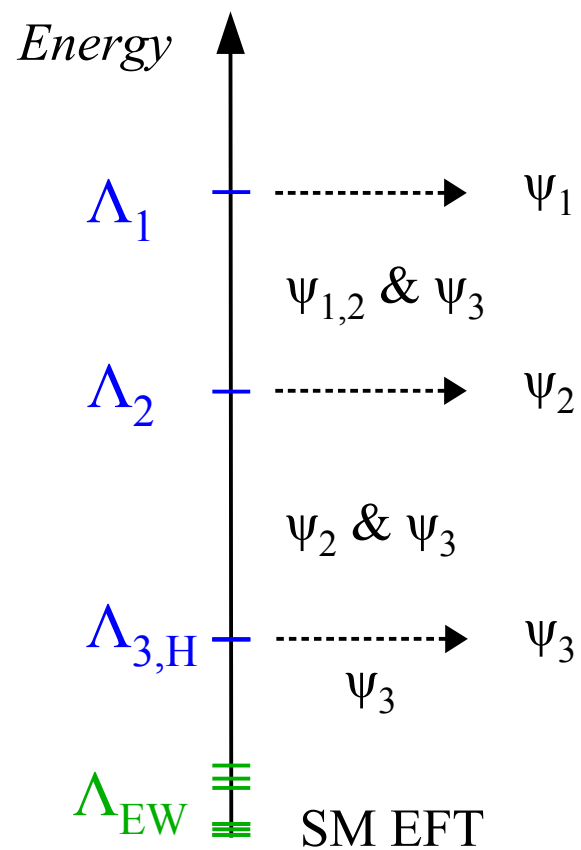
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- “Flavor deconstruction”: the scale of NP is low &  $U(2)^n$  emerges as accidental symm. from flavor non-universal gauge interactions @ nearby scales

No “QFT solution” to  
EW hierarchy problem

“decoupling” of flavor &  
EW problems

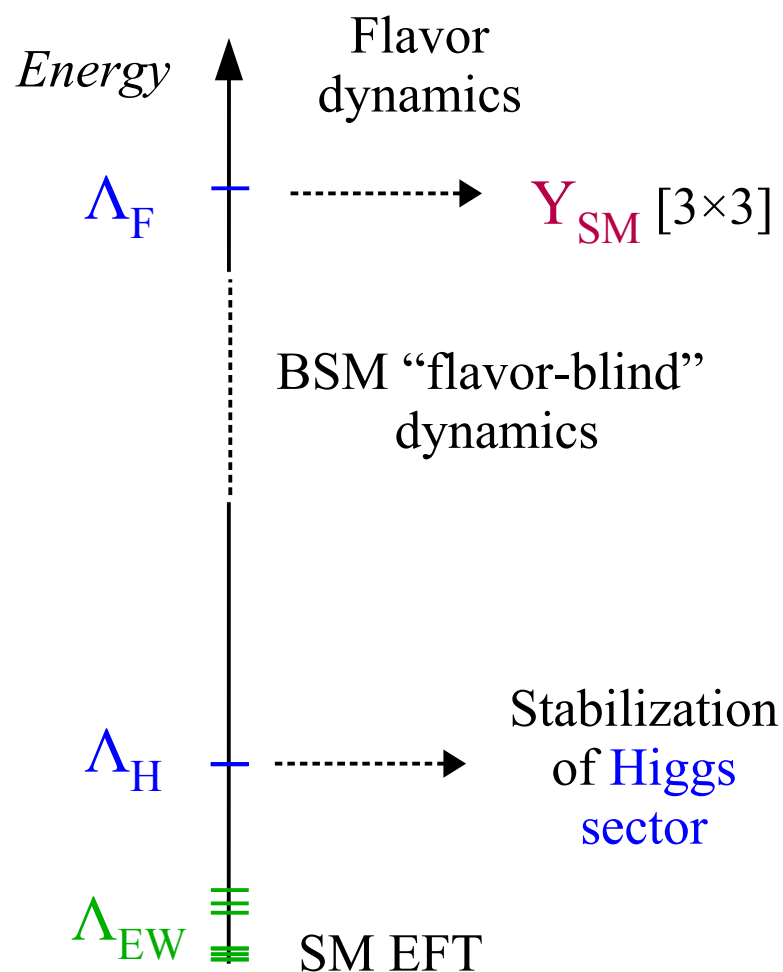
“interplay” of flavor &  
EW problems

## Flavor non-universal interactions



## ► Flavor non-universal interactions

For a long time, the vast majority of model-building attempts to extend the SM was based on the *implicit* hypotheses of *flavor-universal* New Physics



- Concentrate on the **Higgs hierarchy problem**
- Postpone **the flavor problem** to higher scales



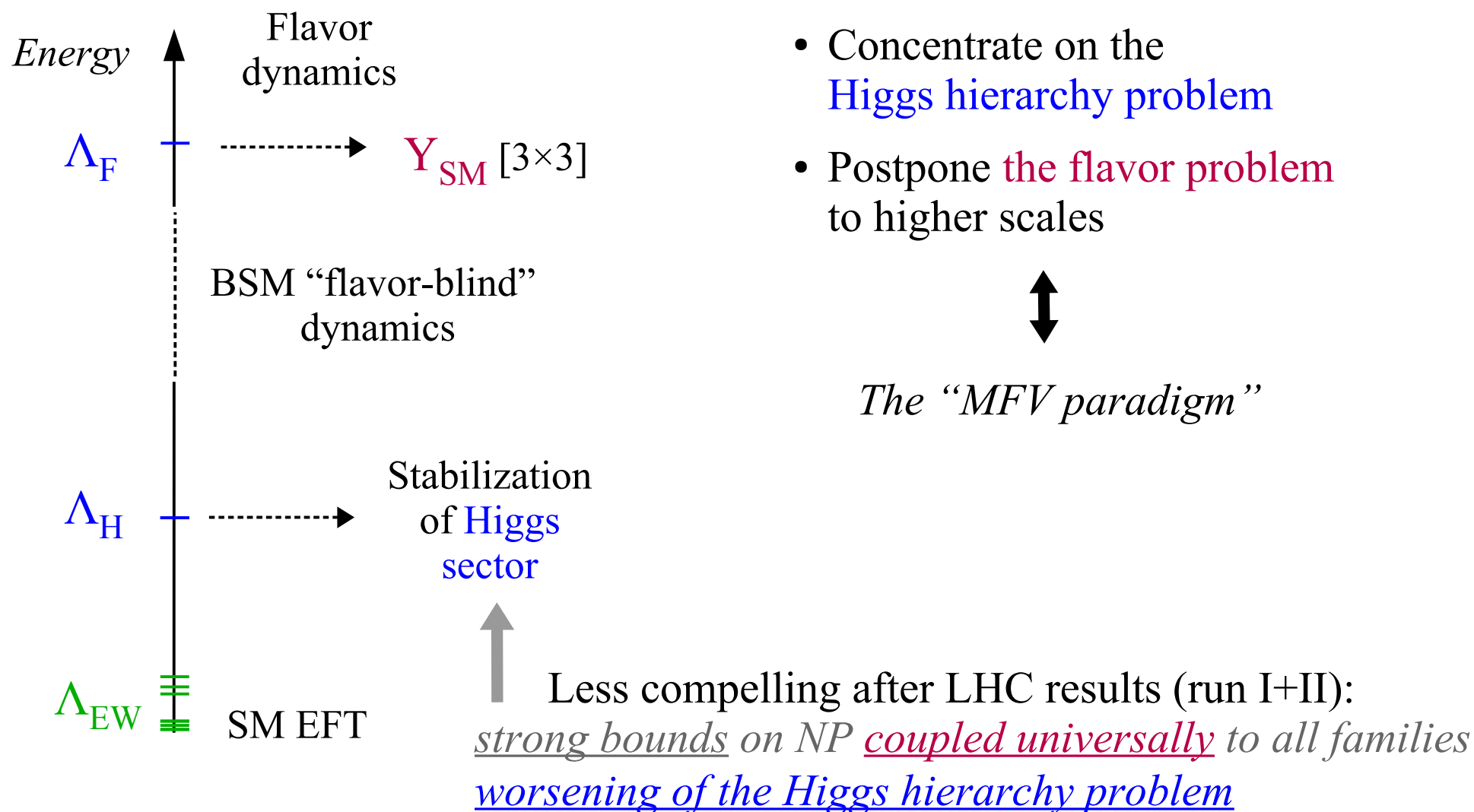
*The "MFV paradigm"*

*The Yukawa couplings are the only sources of flavor symmetry breaking accessible at low energies*

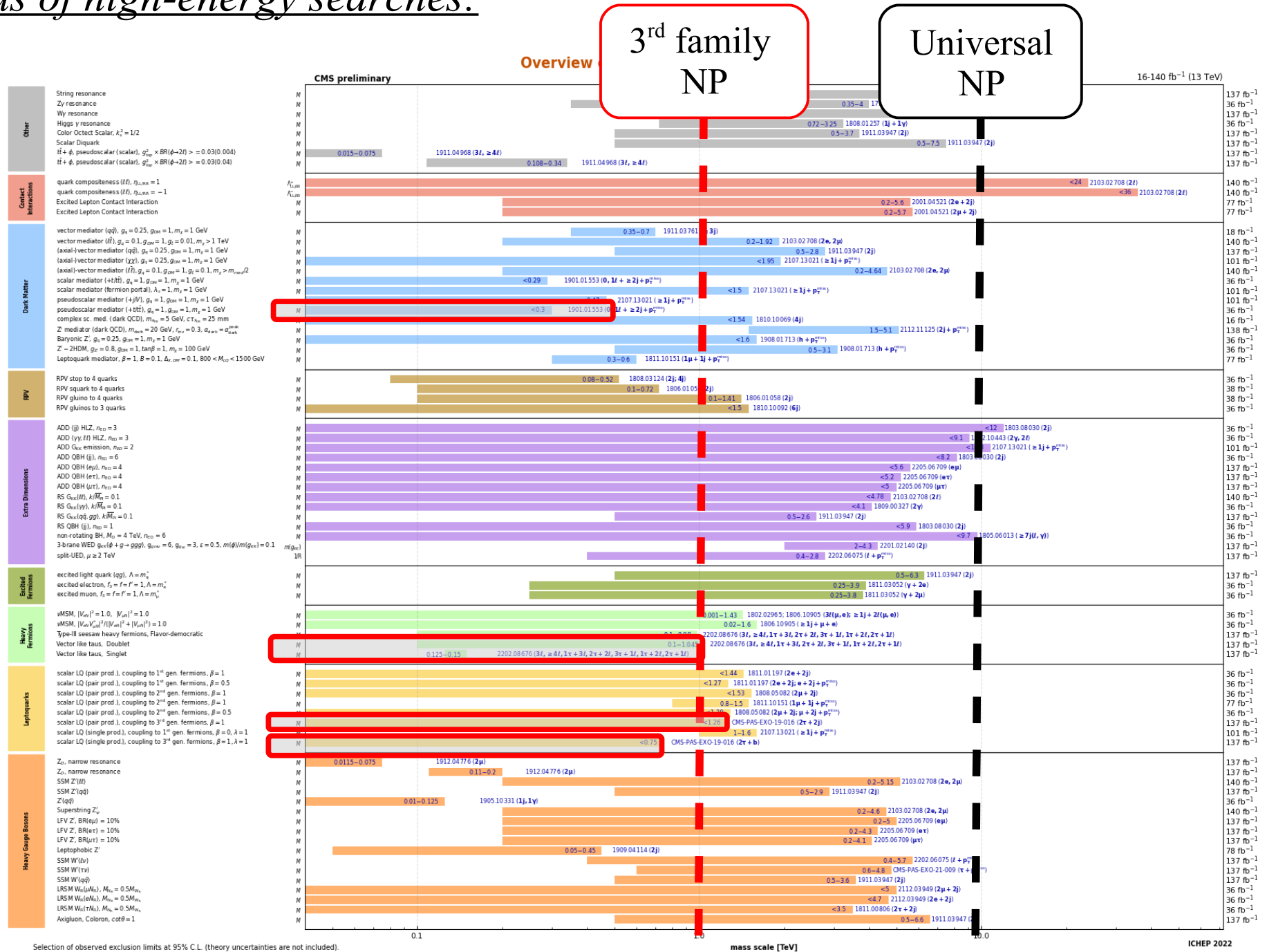
3 families = "identical copies" up to high energies

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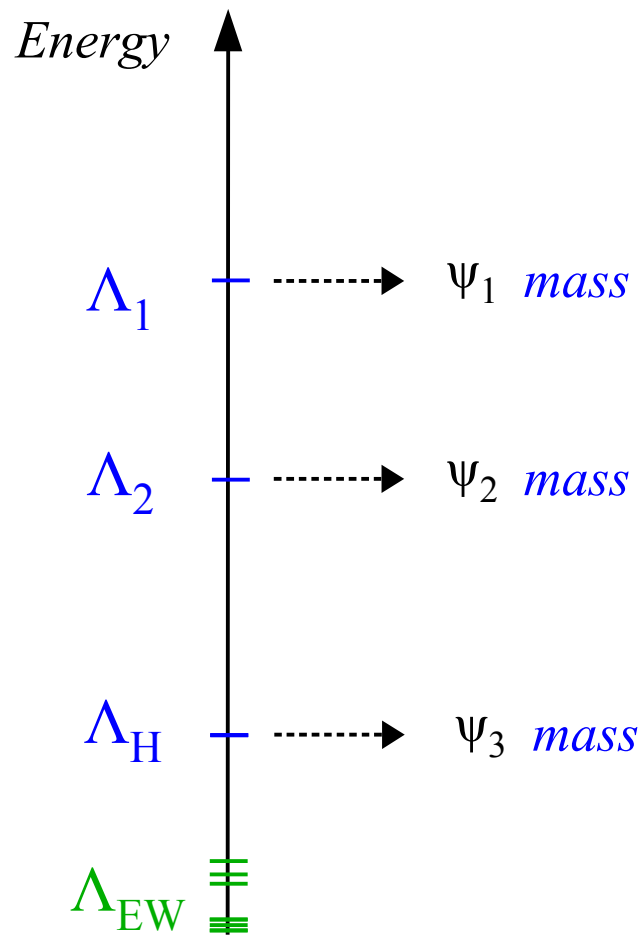


# Status of high-energy searches:



## ► Flavor non-universal interactions

A more efficient paradigm to address both flavor puzzles (I+II), & *possibly* the Higgs hierarchy, is a multi-scale UV with flavor non-universal interactions



Dvali & Shifman '00  
 Panico & Pomarol '16  
 Bordone *et al.* '17  
 Allwicher, GI, Thomsen '20  
 Barbieri '21  
 Davighi & G.I. '23  
 Barbieri & GI '24  
 ⋮

*Basic idea:*

1<sup>st</sup> & 2<sup>nd</sup> generations have small masses  
 (+ small coupling to NP) because these are  
 generated by **new dynamics at heavier scales**

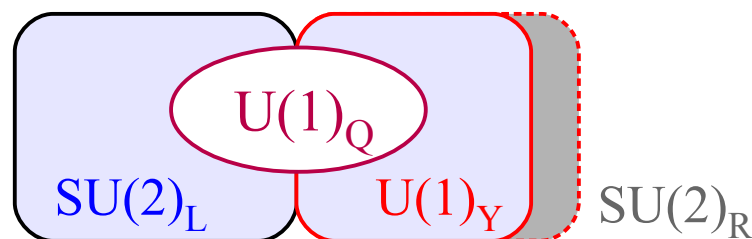


~~3 gen. = “identical copies”  
 up to high energies~~

## ► Flavor non-universal interactions

To better appreciate the change of perspective, consider the following analogy:

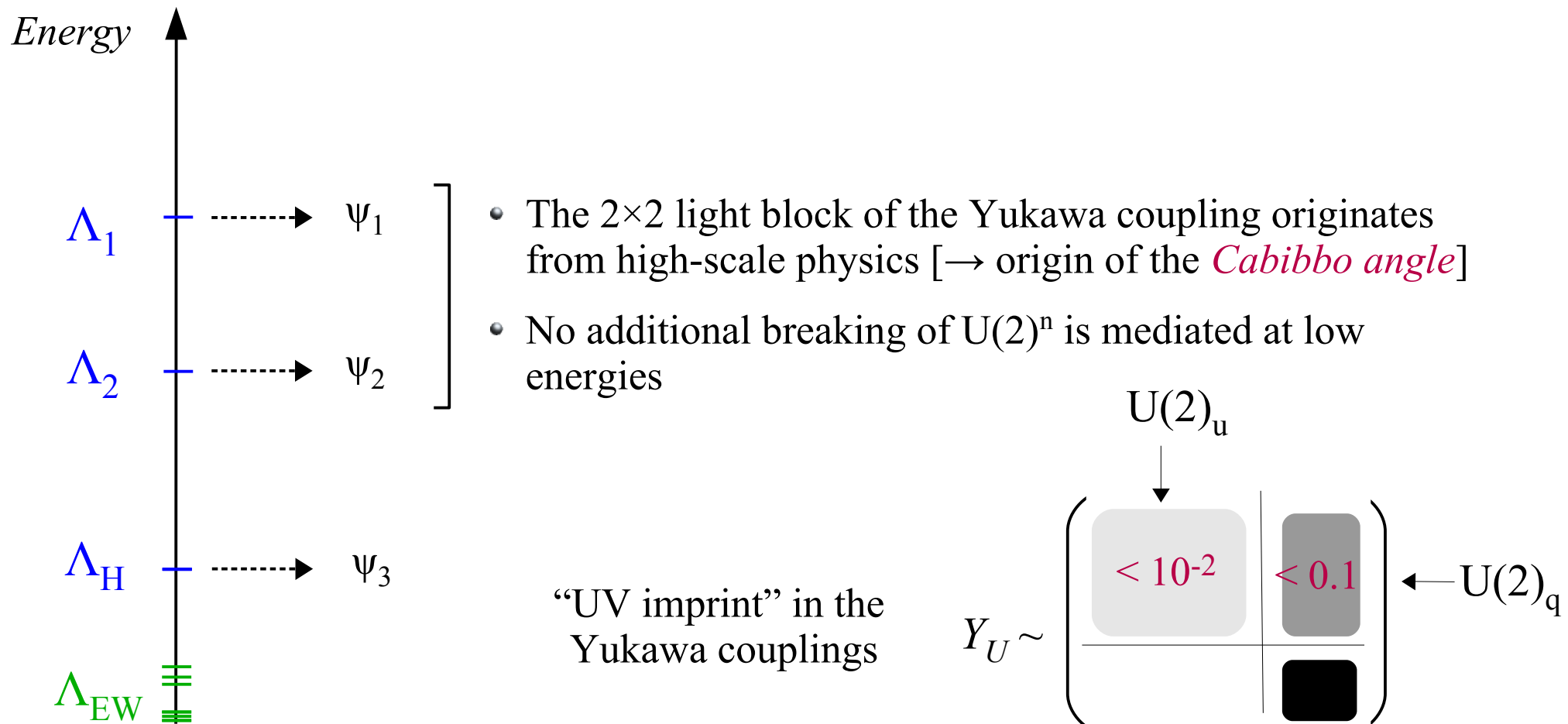
In the low-energy limit of the SM (= QED×QCD) we observe perfect universality of LH and RH gauge couplings. However, we know this is a low-energy artifact:



In a similar fashion, the flavor universality of all SM gauge interactions could be a low-energy artifact...

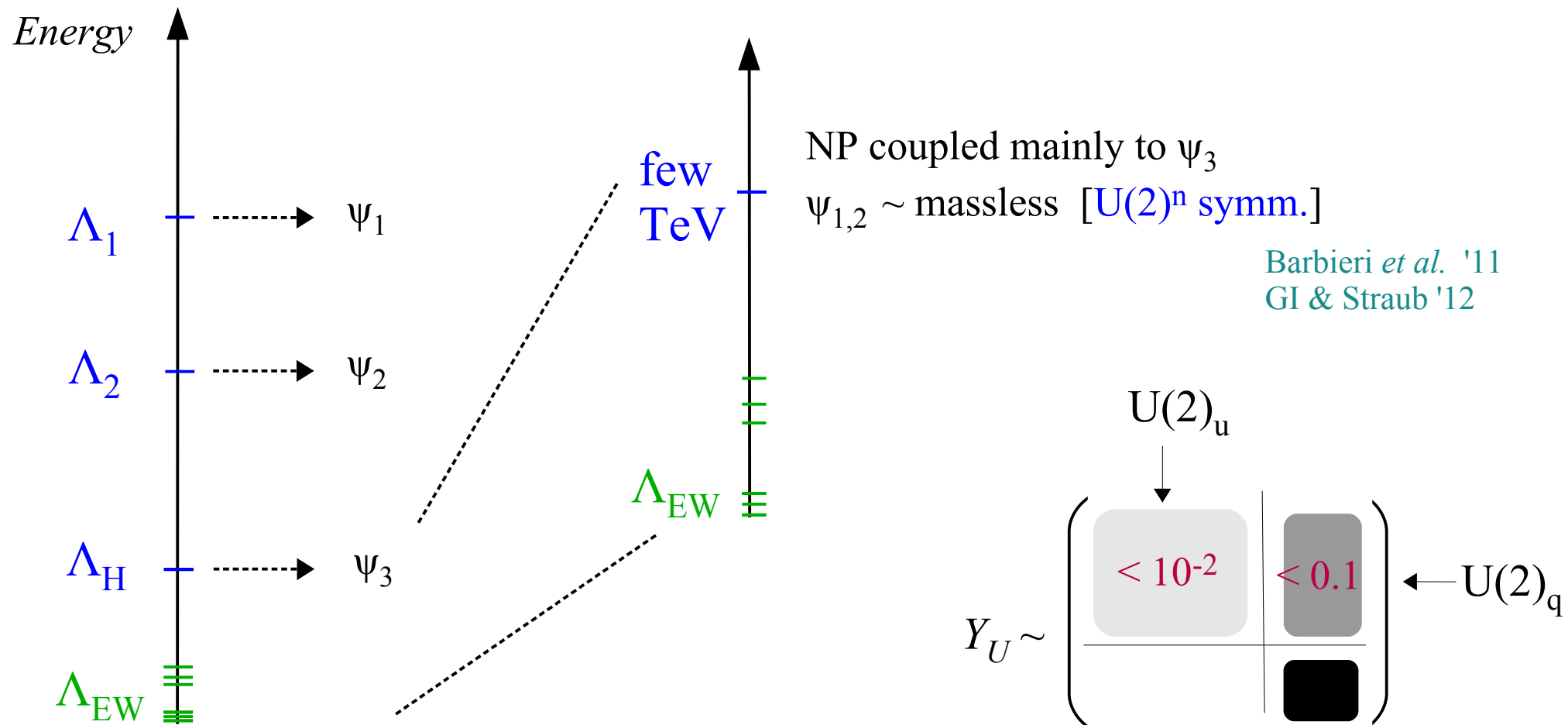
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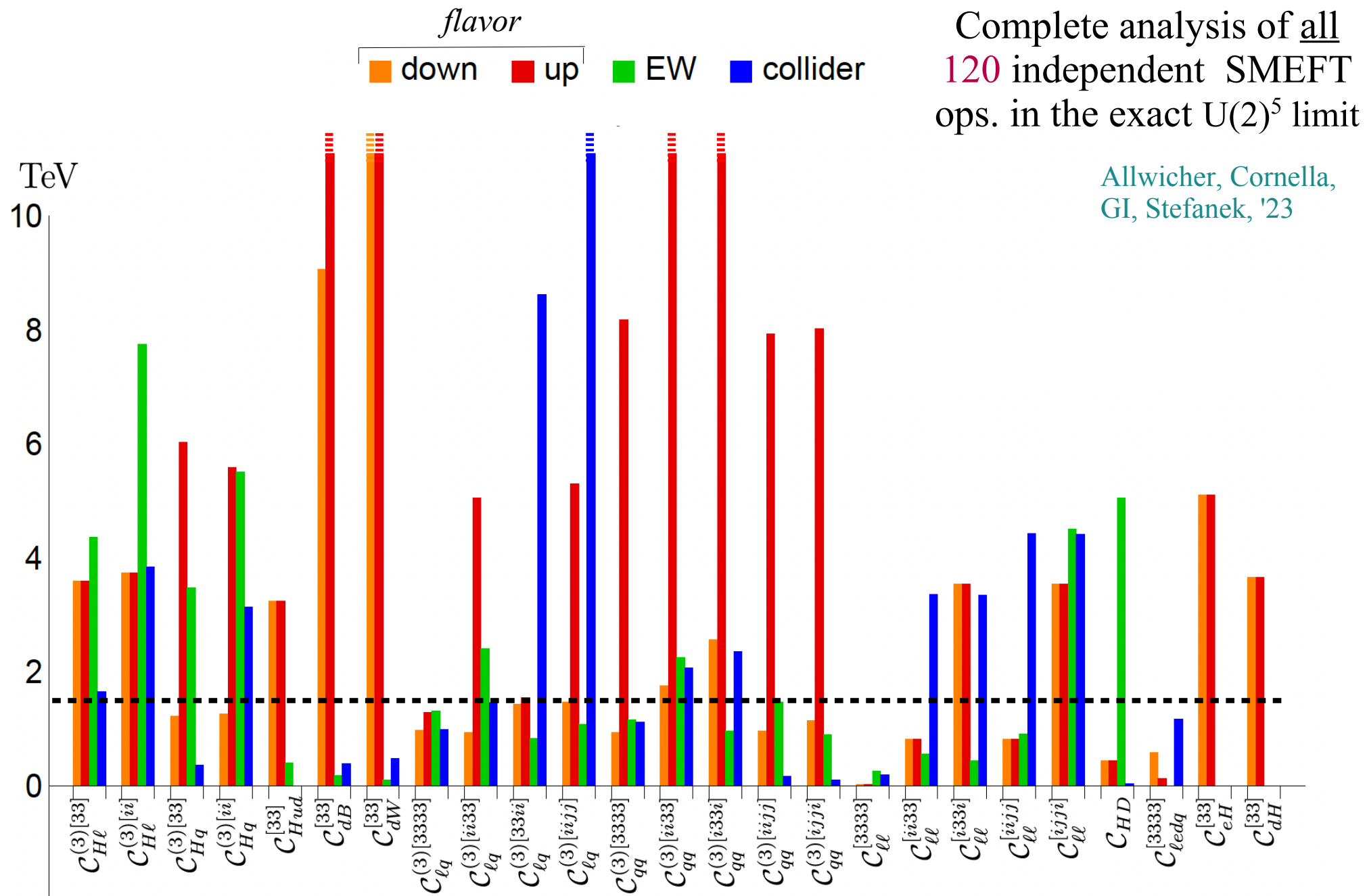
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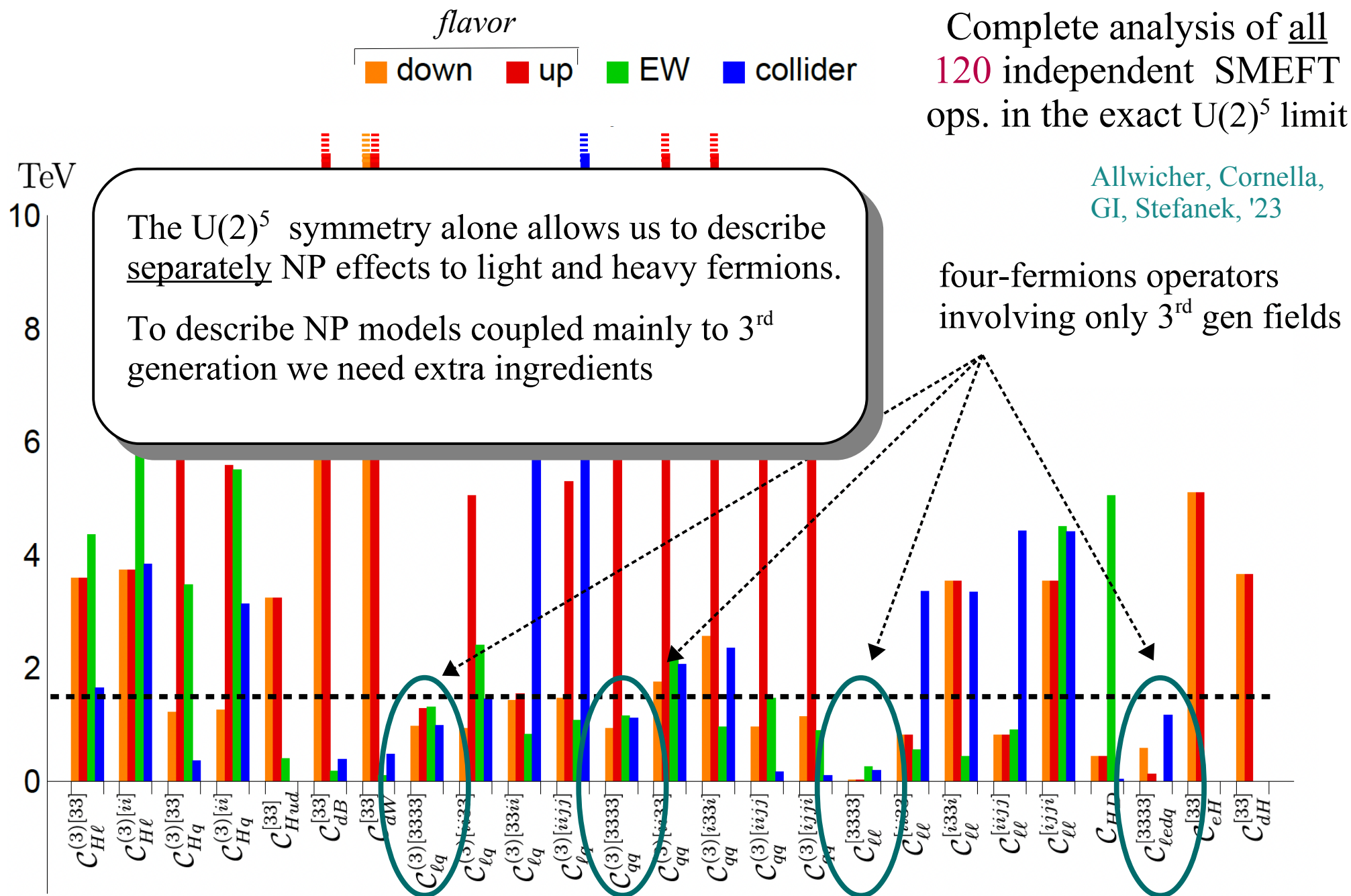


Effective organizing principle for the **flavor structure** of the **SMEFT**

→ [I] SMEFT bounds in the  $U(2)^5$  symmetric limit



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flavor

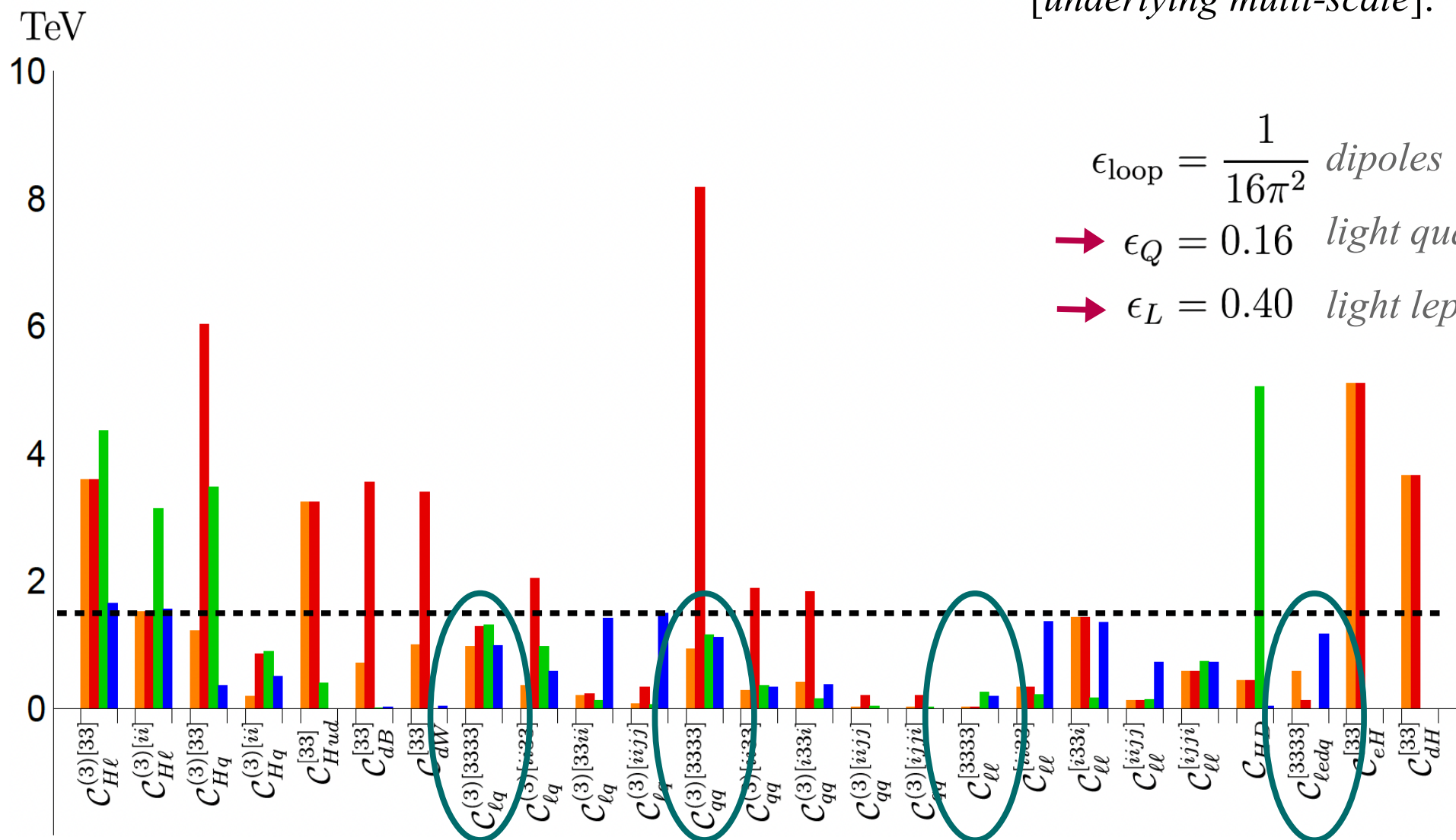


Dynamical suppression factors  
[underlying multi-scale]:

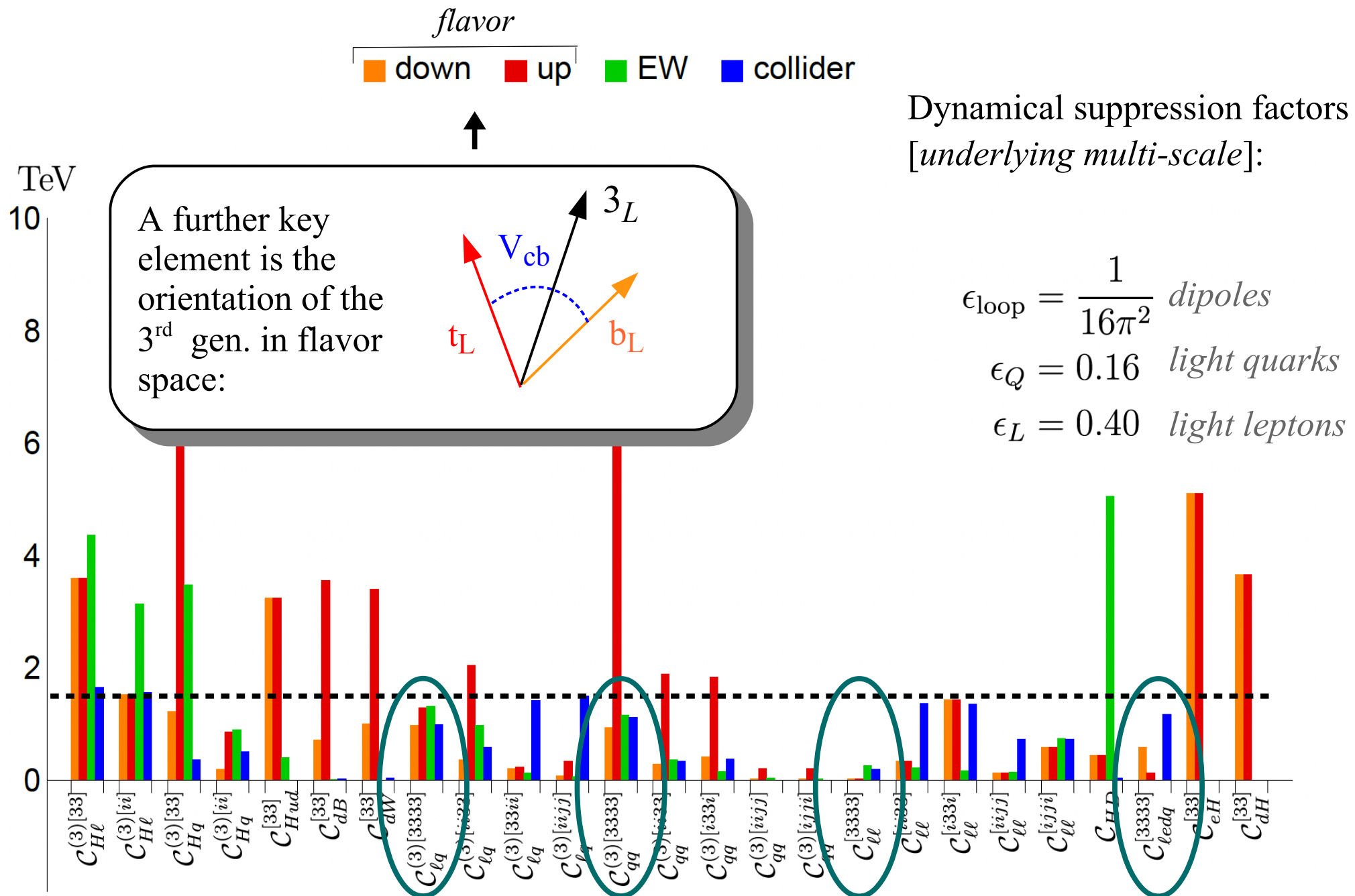
$$\epsilon_{\text{loop}} = \frac{1}{16\pi^2} \text{ dipoles}$$

$$\rightarrow \epsilon_Q = 0.16 \text{ light quarks}$$

$$\rightarrow \epsilon_L = 0.40 \text{ light leptons}$$

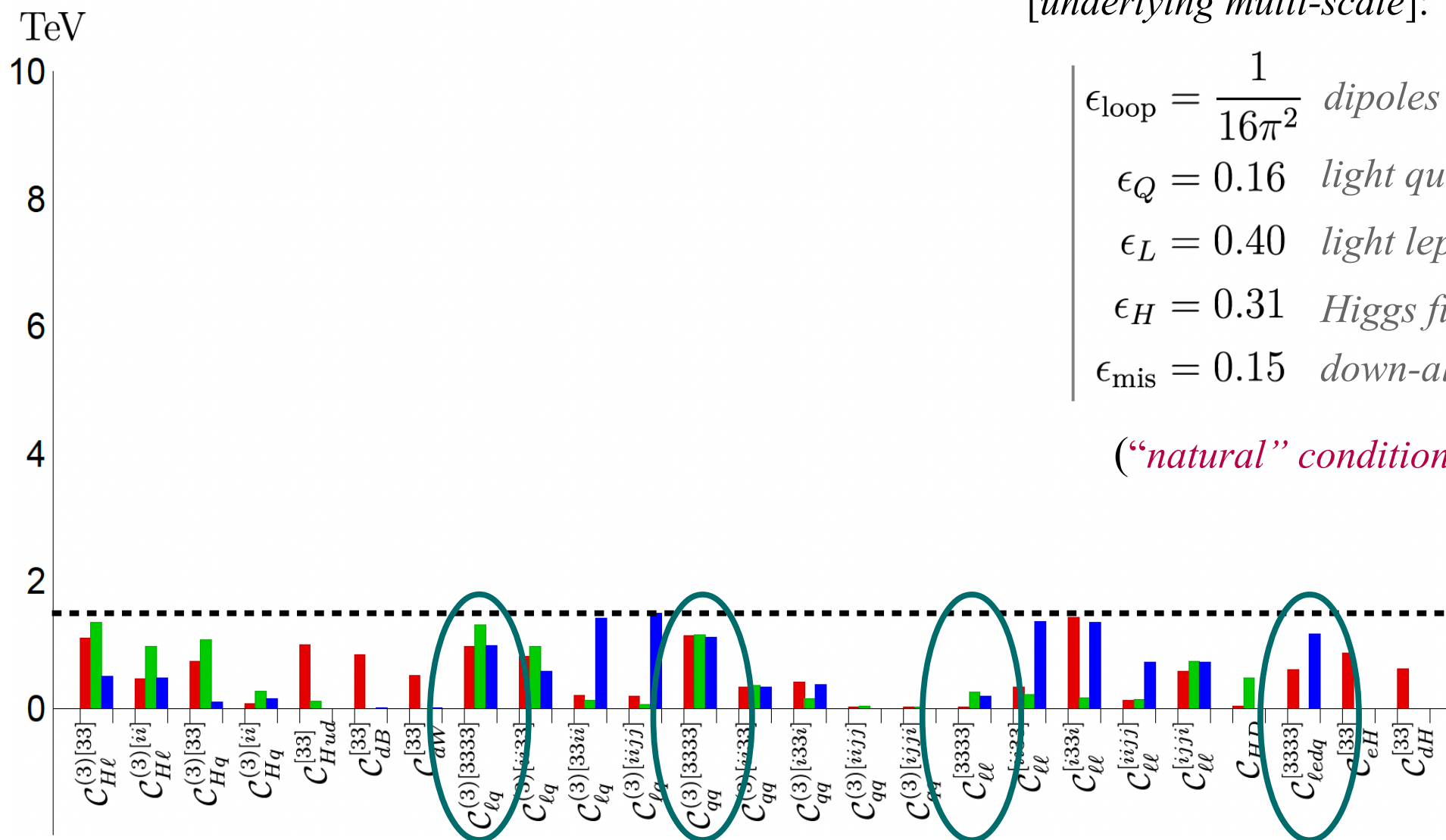


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flavor EW collider



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$$\epsilon_H = 0.31 \text{ Higgs fields}$$

$$\epsilon_{\text{mis}} = 0.15 \text{ down-align.}$$

(“natural” conditions)

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■ flavor ■ EW ■ collider

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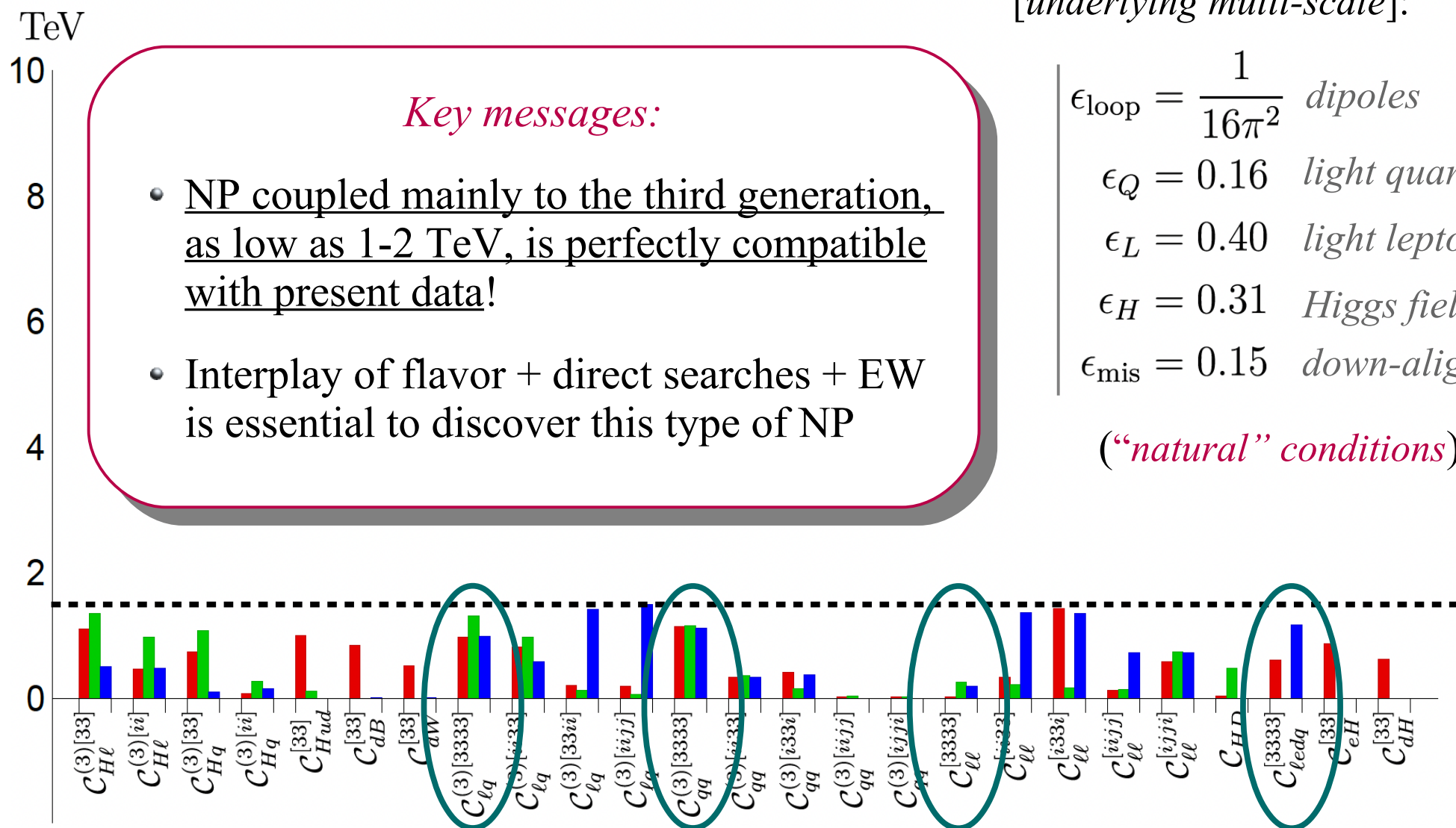
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*Key messages:*

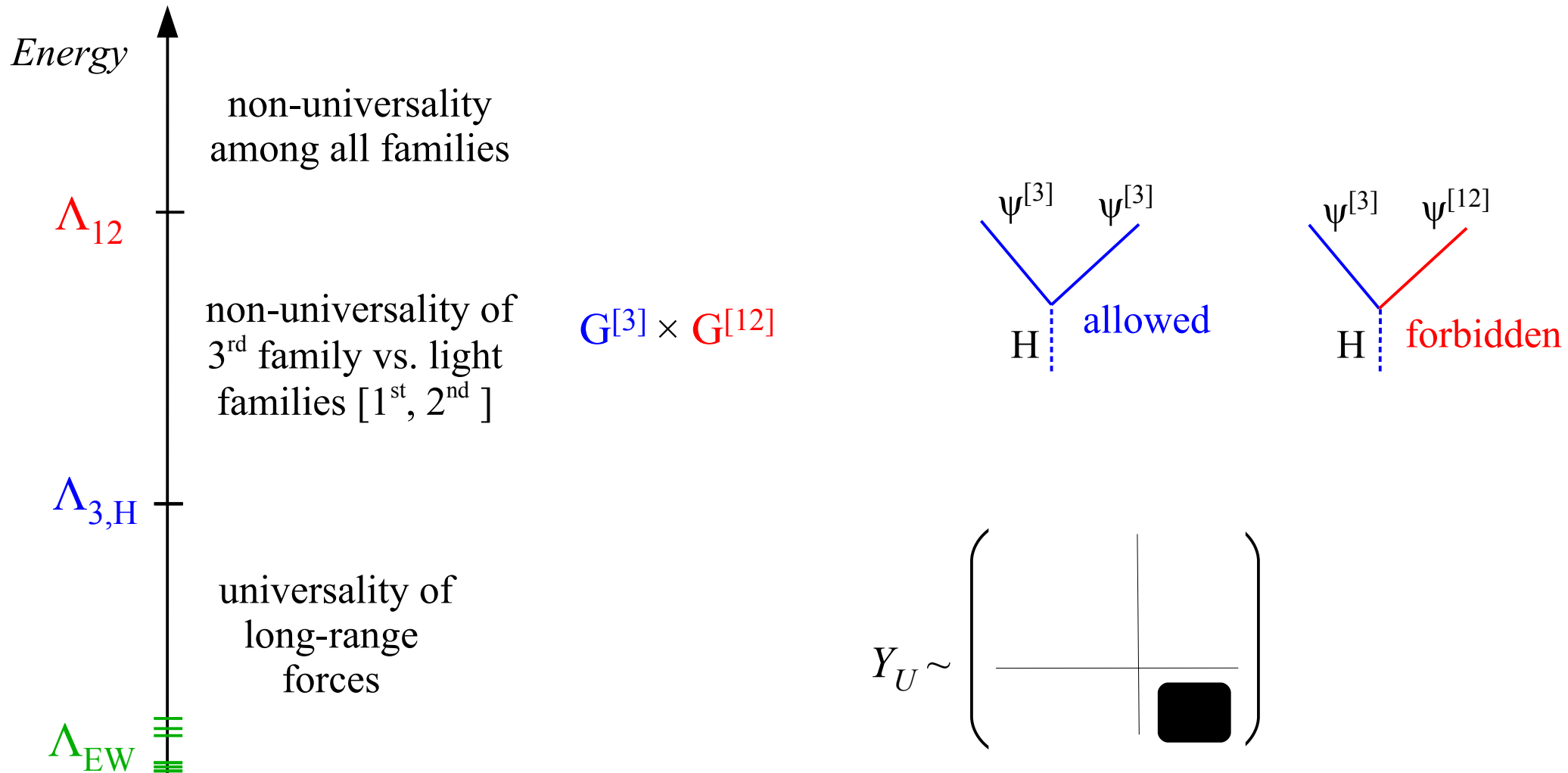
- NP coupled mainly to the third generation, as low as 1-2 TeV, is perfectly compatible with present data!
- Interplay of flavor + direct searches + EW is essential to discover this type of NP



(“natural” conditions)

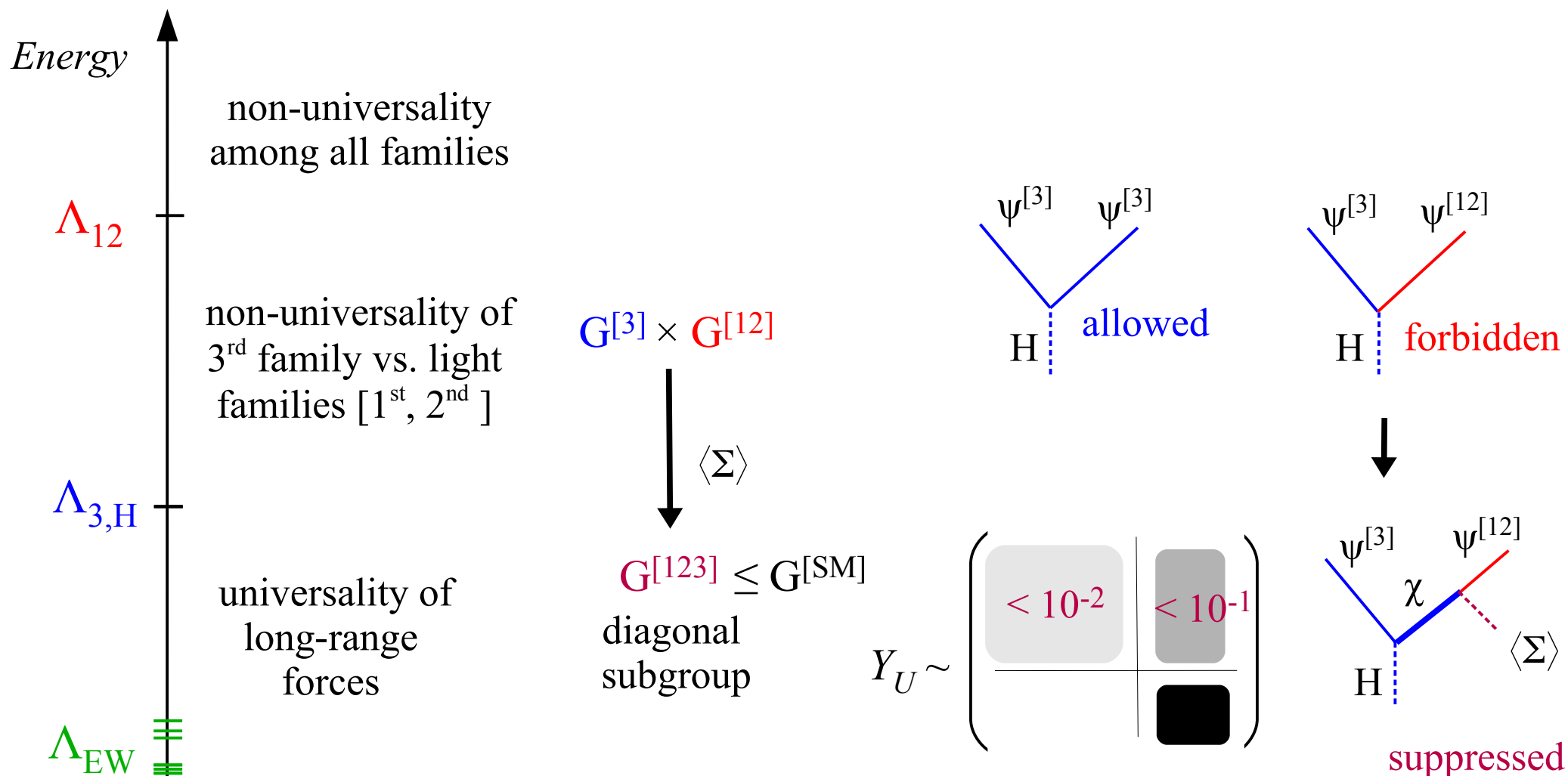
→ [II] Flavor deconstruction

Going beyond the EFT approach, a consistent way to construct a multi-scale theory with flavor non-universal interactions is via a “flavor deconstruction” of the SM gauge symmetries:



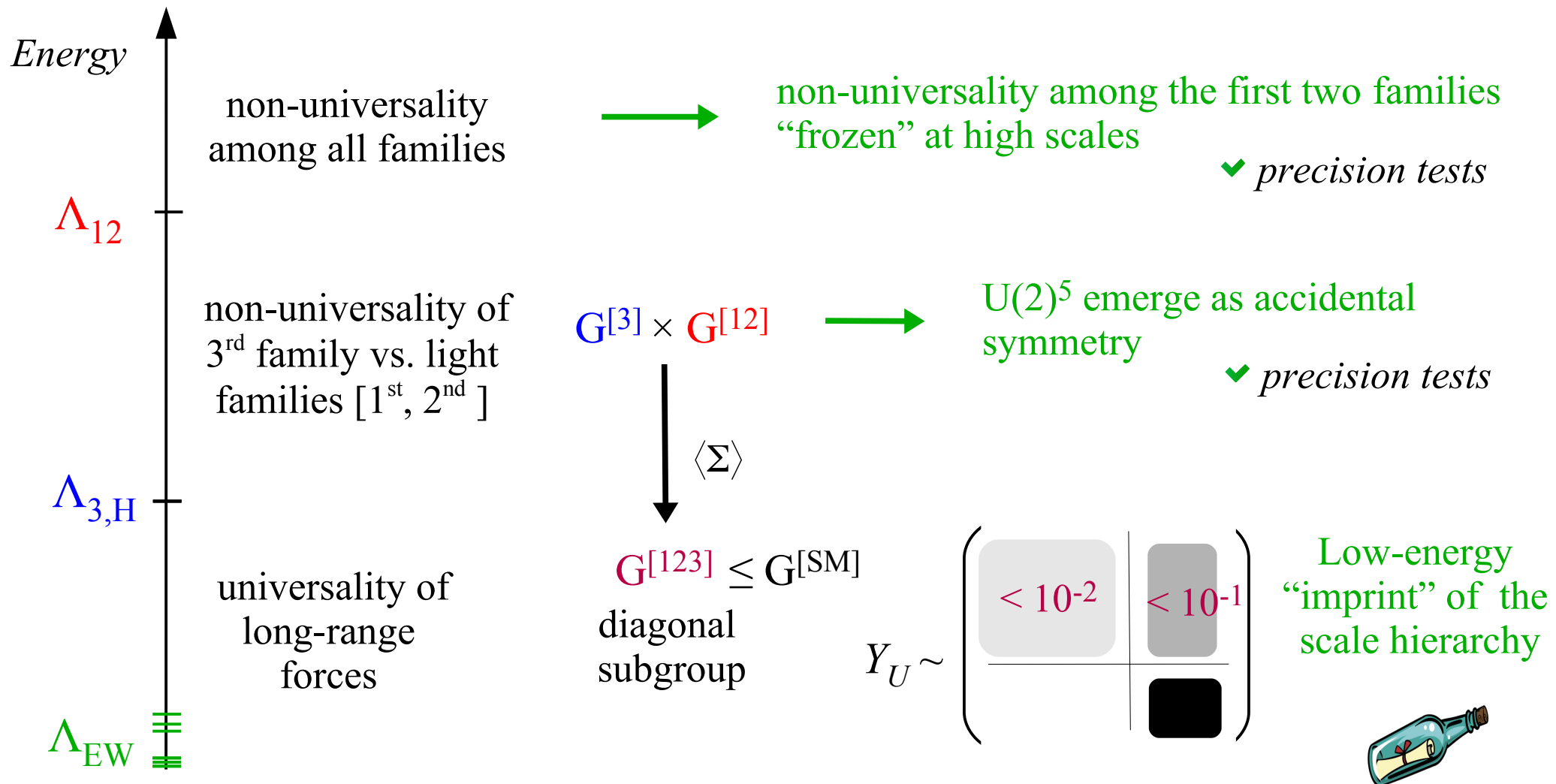
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Going beyond the EFT approach, a consistent way to construct a multi-scale theory with flavor non-universal interactions is via a “flavor deconstruction” of the SM gauge symmetries:



→ [II] Flavor deconstruction

Going beyond the EFT approach, a consistent way to construct a multi-scale theory with flavor non-universal interactions is via a “flavor deconstruction” of the SM gauge symmetries:



→ [III] Flavor deconstruction & quark-lepton unification

All possible options have been classified (*not many consistent choices*).

A particularly interesting one is allowing quark-lepton unification a la Pati-Salam for the 3<sup>rd</sup> gen

$$\begin{array}{c} \text{SU}(4)^{[3]} \times \text{SU}(3)^{[12]} \times G_{\text{EW}} \\ \downarrow \\ \text{SU}(3) \times \text{SU}(2)_L \times \text{U}(1)_Y \end{array}$$

Fermions  
in SU(4):

$$\begin{bmatrix} Q^\alpha \\ Q^\beta \\ Q^\gamma \\ L \end{bmatrix}$$

Main Pati-Salam idea:  
Lepton number as “the 4<sup>th</sup> color”

✓ Explain charge quantization

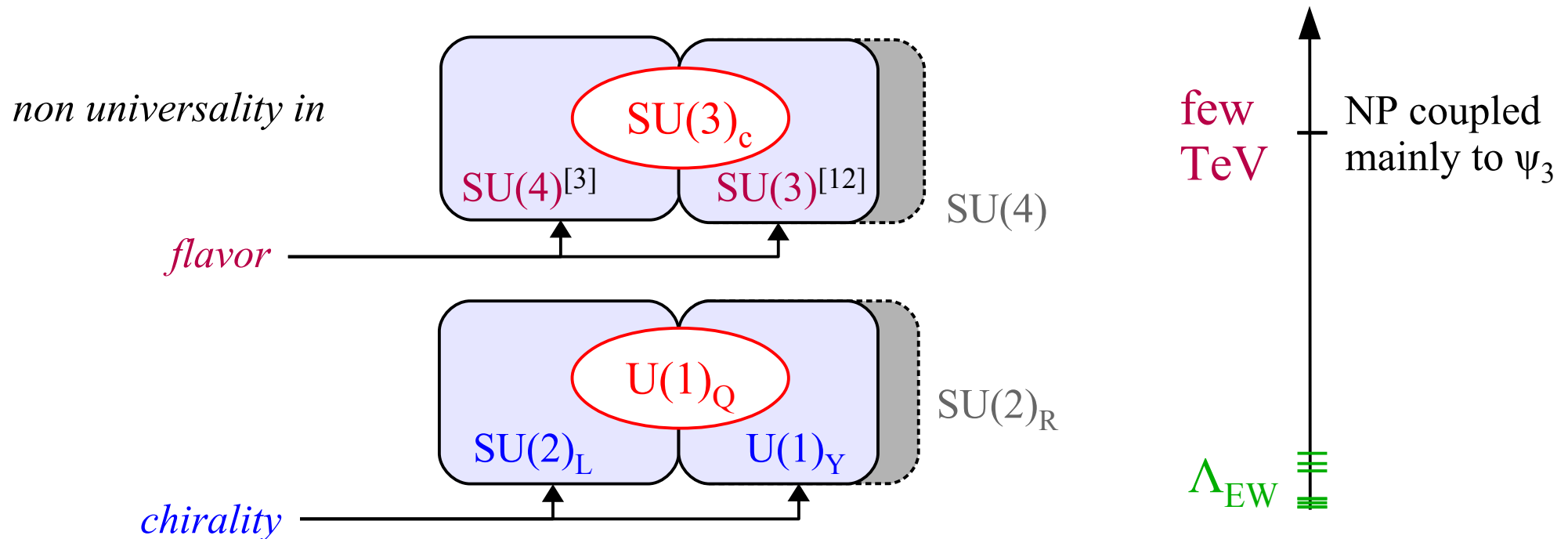
$$\text{SU}(4) \sim \left[ \begin{array}{c|c} \text{SU}(3)_C & 0 \\ \hline 0 & 0 \end{array} \right] + \left[ \begin{array}{c|c} 0 & \text{LQ} \\ \hline \text{LQ} & \end{array} \right] + \left[ \begin{array}{c|c} \frac{1}{3} & 0 \\ \hline 0 & -1 \end{array} \right] \quad \text{B-L generator}$$

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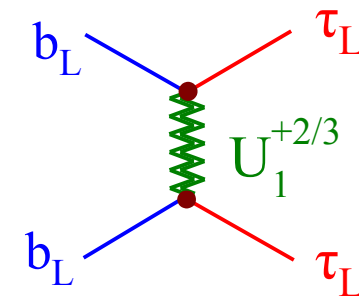
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$$\begin{array}{ccc} \text{SU}(4)^{[3]} \times \text{SU}(3)^{[12]} \times G_{\text{EW}} & & \\ \downarrow & \longrightarrow & \\ \text{SU}(3) \times \text{SU}(2)_L \times \text{U}(1)_Y & & \end{array}$$

vector  
leptoquark



✓ Explain charge quantization

✓ Might explain some existing tensions in B-physics data

$$\text{SU}(4) \sim \left[ \begin{array}{c|c} \text{SU}(3)_C & 0 \\ \hline 0 & 0 \end{array} \right] + \left[ \begin{array}{c|c} 0 & \text{LQ} \\ \hline \text{LQ} & \end{array} \right] + \left[ \begin{array}{c|c} \frac{1}{3} & 0 \\ \hline 0 & -1 \end{array} \right] \quad \begin{array}{l} \text{B-L} \\ \text{generator} \end{array}$$

→ [IV] Flavor deconstruction & partial compositeness

The FD hypothesis alone does not address the EW hierarchy problem.

However, it provides some basic ingredients (“*light*” new states coupled to Higgs and 3<sup>rd</sup> gen.) which can help to address it → interesting specific advantage when merged to Higgs compositeness

Fuentes-Martin & Stangl '20  
Fuentes-Martin *et al.* '22

→ [IV] Flavor deconstruction & partial compositeness

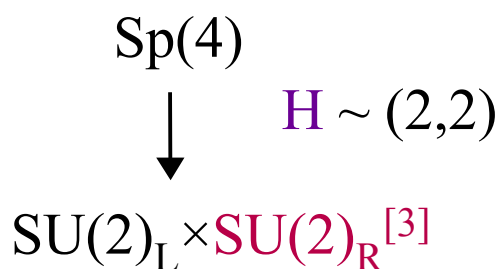
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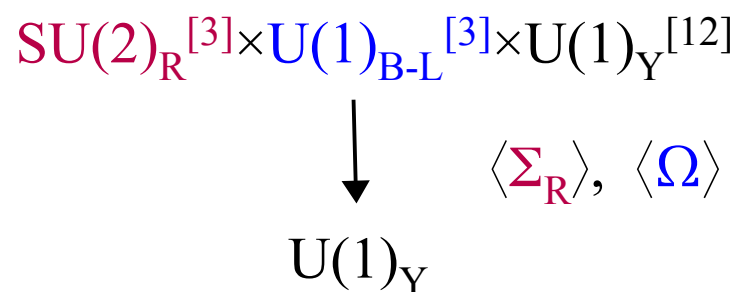
Fuentes-Martin & Stangl '20  
Fuentes-Martin et al. '22

Toy example [Covone, GI, Pesut '24]:

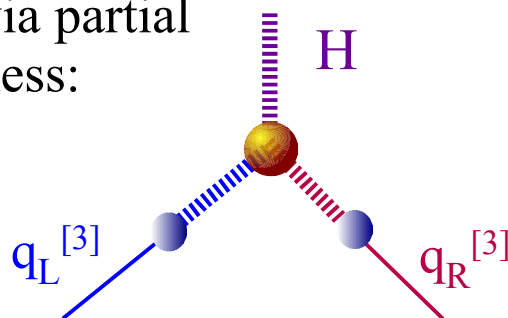
*Minimal Composite Higgs:*



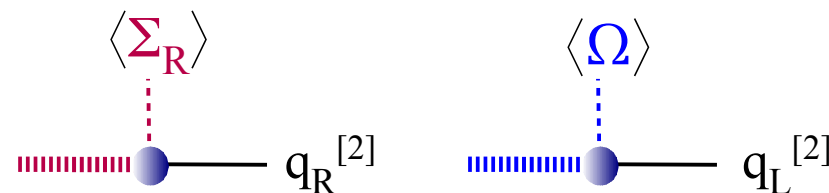
*Flavor Deconstruction of Hypercharge:*



3<sup>rd</sup> family Yukawa couplings via partial compositeness:



Light Yukawas suppressed → higher-dim ops. generated by heavy (elementary) dynamics (e.g. VL fermions) above the compositeness scale:



→ [IV] Flavor deconstruction & partial compositeness

An interesting feature emerges when computing explicitly the radiatively-induced Higgs potential:

$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2 \left( \frac{h}{2F} \right) + c_2 \sin^4 \left( \frac{h}{2F} \right)$$

$$\frac{c_1}{F^4} \Big|_{\text{phys.}} = \frac{m_h^2}{F^2} \lesssim 0.06$$

Unavoidable  $\sim 1\%$  tuning

[given current bounds on  $F$  ( $\gtrsim 0.7$  TeV)]

- O(1) model-independent fermion contribution (controlled by  $y_t$ )

- Possible cancellation from the gauge sector (unambiguous opposite sign) if

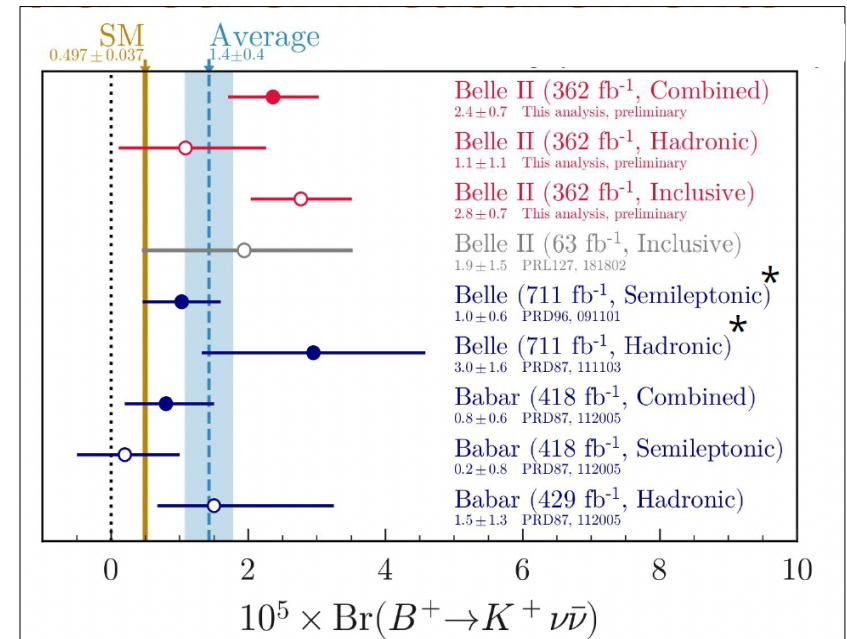
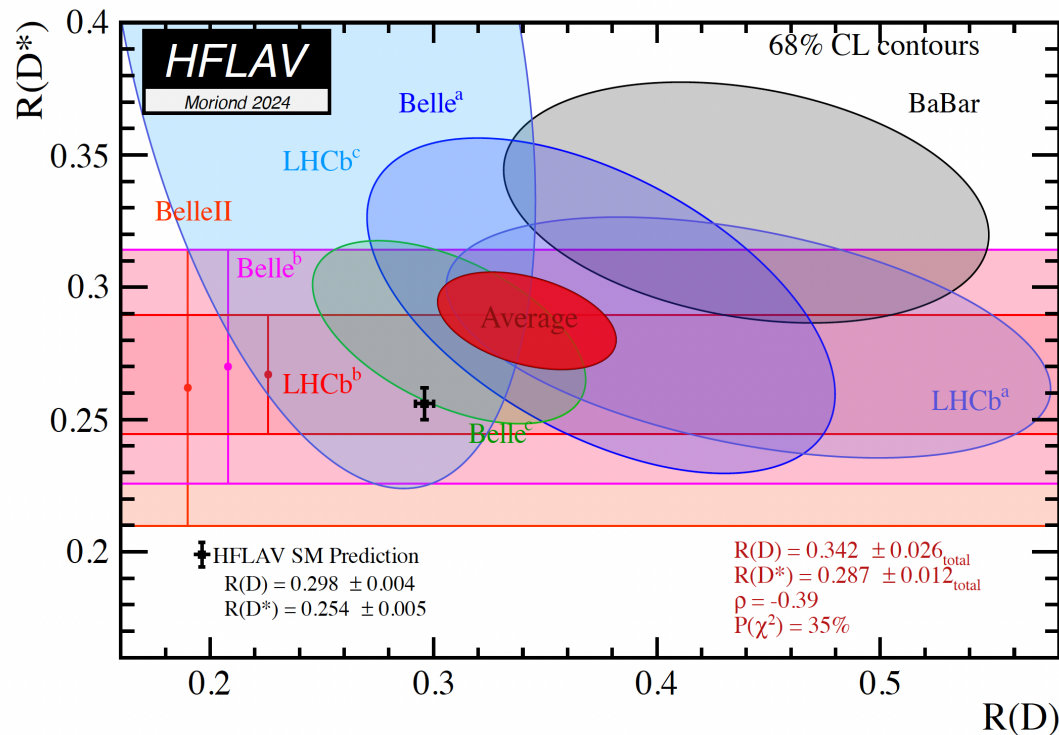
$$g_R^{[3]} \gtrsim 1 \quad \& \quad g_R \langle \Sigma_R \rangle < M_\rho$$

$$\frac{c_2}{F^4} \Big|_{\text{phys.}} = \frac{2m_h^2}{v^2} \approx \frac{1}{2}$$

requires light top partners  
(as in ordinary CH)

*The cancellation occurs in the “natural” parameter space dictated by the flavor-deconstruction hypothesis*

# A brief look to current data & future prospects



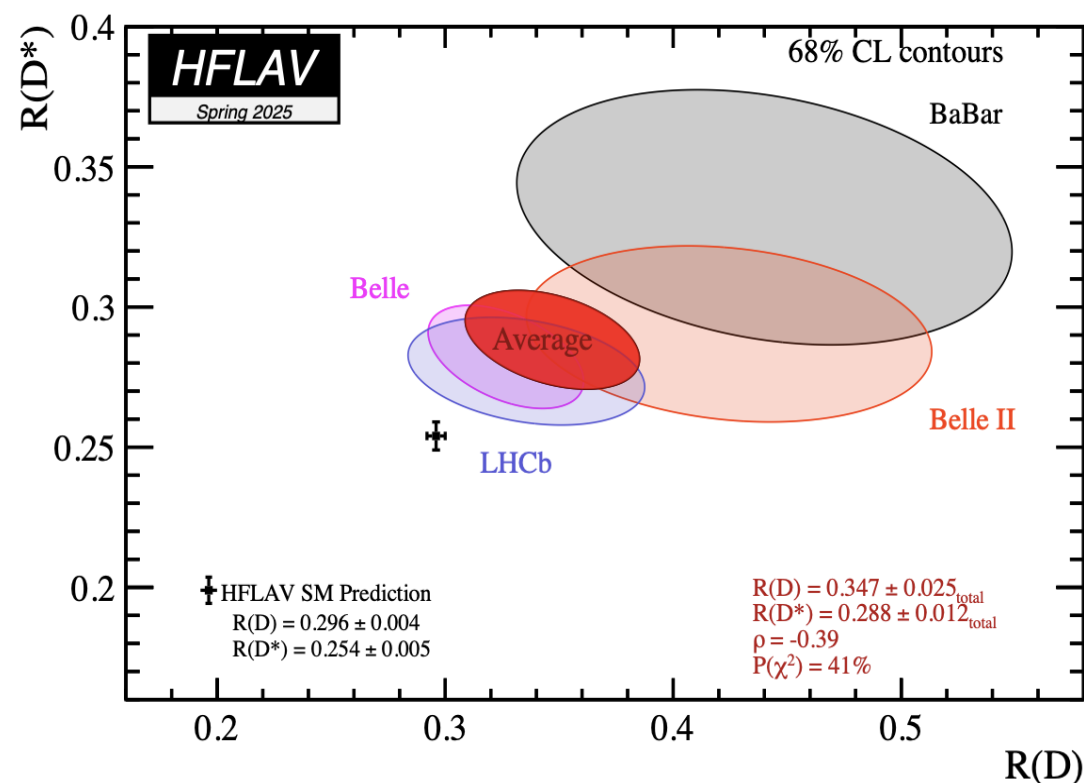
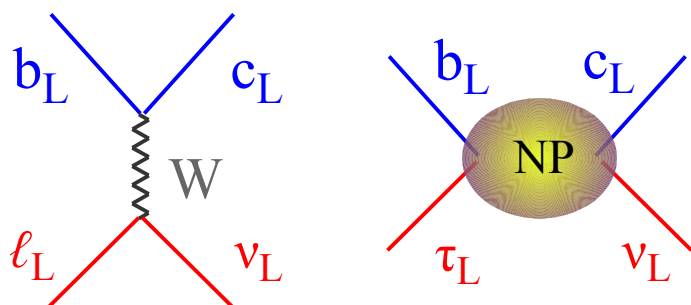
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The idea of flavor non-universal interactions – with a 1<sup>st</sup> layer of new physics already at the TeV scale – has several interesting implications for various **low-energy measurements** (with different degree of model-dependence)

E.g.: I) Lepton universality violations in  $b \rightarrow c \tau \nu$  decays

$$R(X)^{\tau/l} = \frac{\Gamma(B \rightarrow X \tau \nu)}{\Gamma(B \rightarrow X l \nu)}$$

$X = D$  or  $D^*$



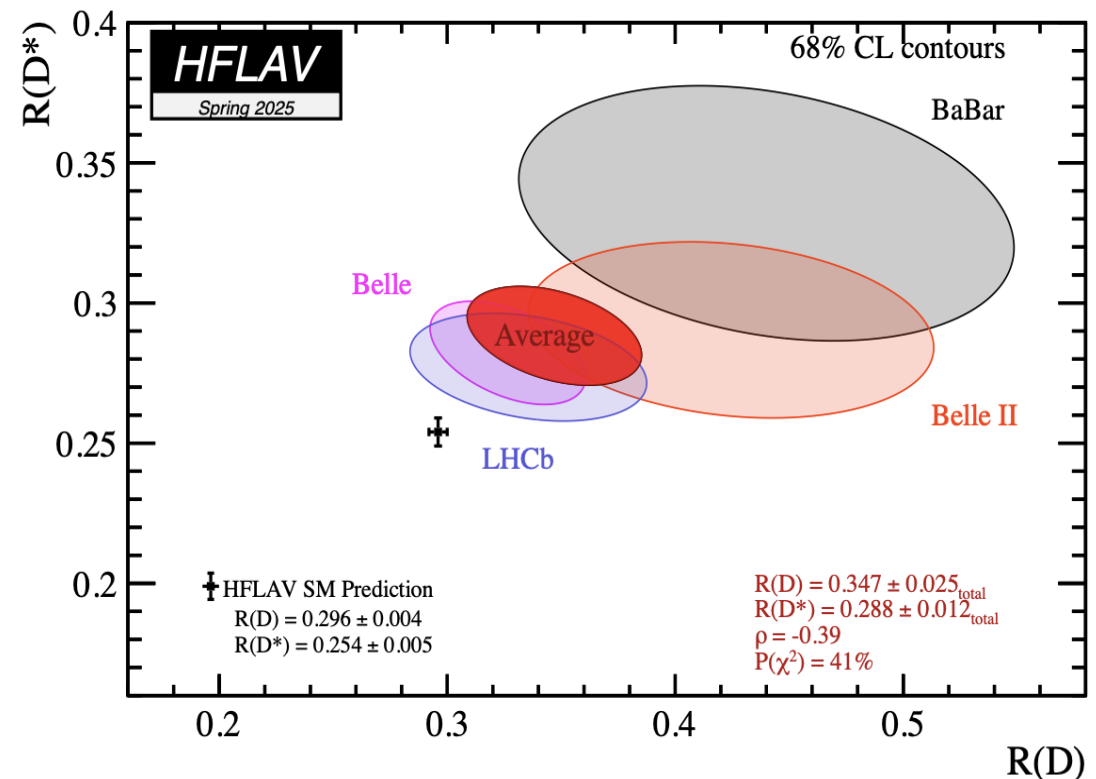
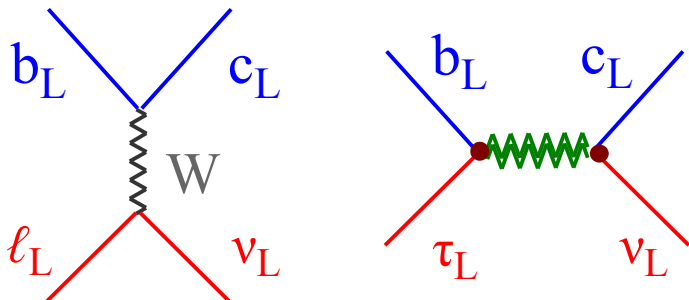
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→ The vector  $LQ$  of 3<sup>rd</sup> gen. quark-lepton unification is an ideal candidate to describe current data

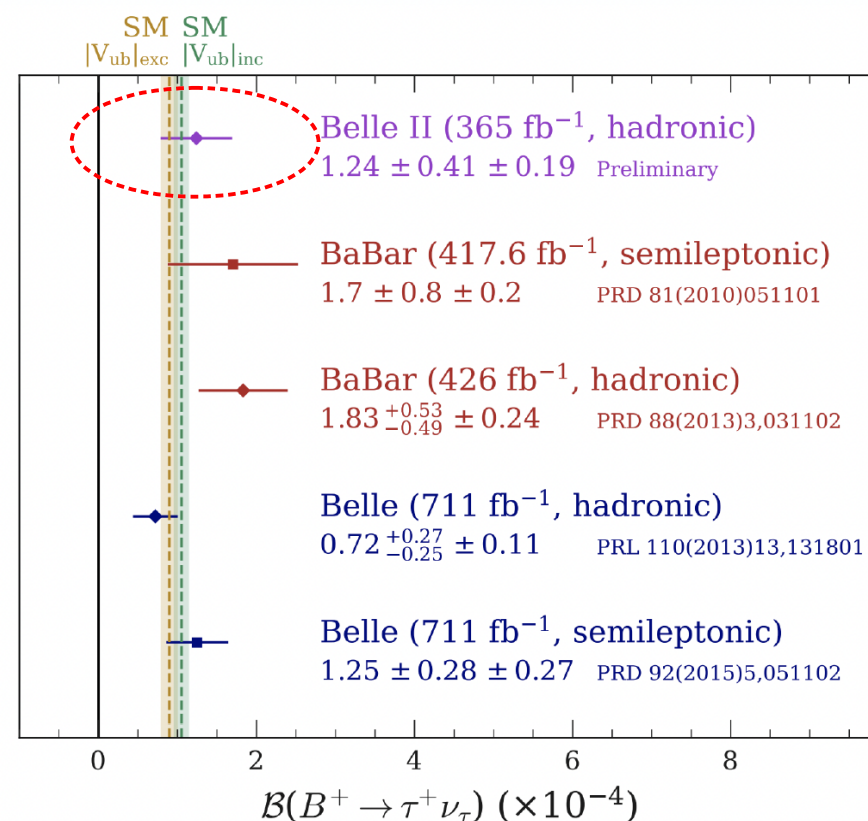
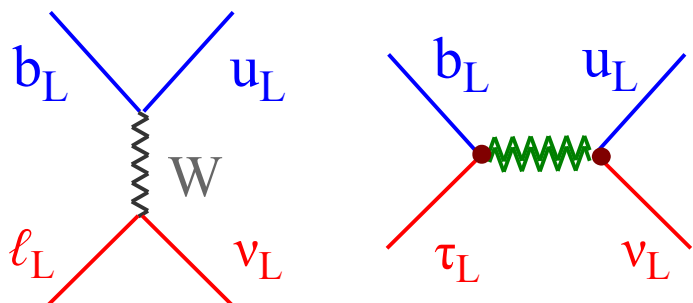
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E.g.: I) Lepton universality violations in  $b \rightarrow c \tau \nu \dots$  &  $b \rightarrow u \tau \nu$  decays

$b \rightarrow u \tau \nu$  decays are more suppressed ( $|V_{ub}| \ll |V_{cb}|$ ) but could allow an extremely clean LFU test via purely leptonic modes

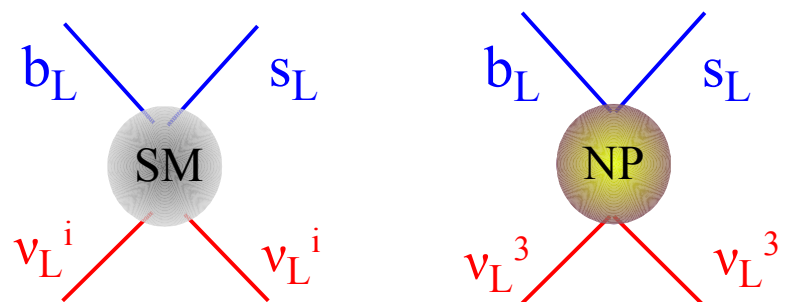
$$R_u^{\tau/\mu} = \frac{\Gamma(B \rightarrow \tau \nu)}{\Gamma(B \rightarrow \mu \nu)}$$



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The idea of flavor non-universal interactions – with a 1<sup>st</sup> layer of new physics already at the TeV scale – has several interesting implications for various **low-energy measurements** (with different degree of model-dependence)

E.g.: II) Deviations from SM in  $b \rightarrow s \nu \bar{\nu}$  rates [ 3<sup>rd</sup> gen.  $\nu$  in the final state ]

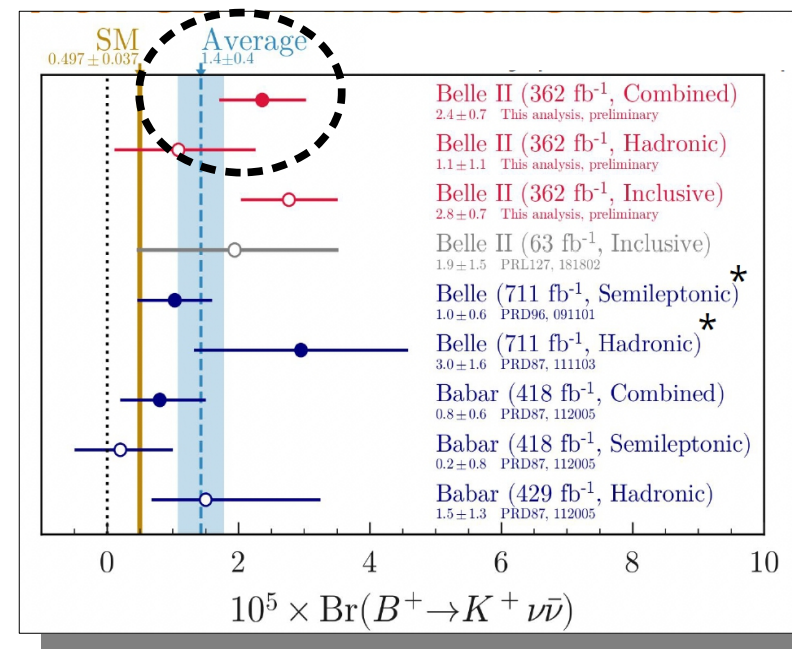


identical for all neutrino species

relevant only for 3<sup>rd</sup> gen. neutrinos

Unambiguous prediction of O(50%) enhancement of  $B(B \rightarrow K \nu \bar{\nu})$  in the model with vector LQ – given excess in R(D).

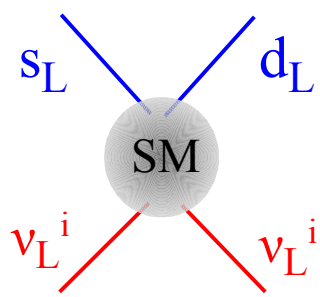
Belle-II '2023



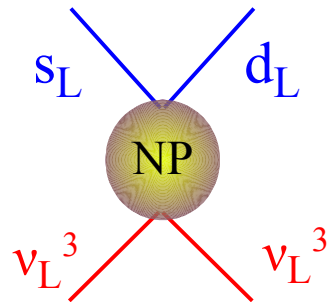
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The idea of flavor non-universal interactions – with a 1<sup>st</sup> layer of new physics already at the TeV scale – has several interesting implications for various **low-energy measurements** (with different degree of model-dependence)

E.g.: II) Deviations from SM in  $b \rightarrow svv$  rates... &  $s \rightarrow dvv$  rates



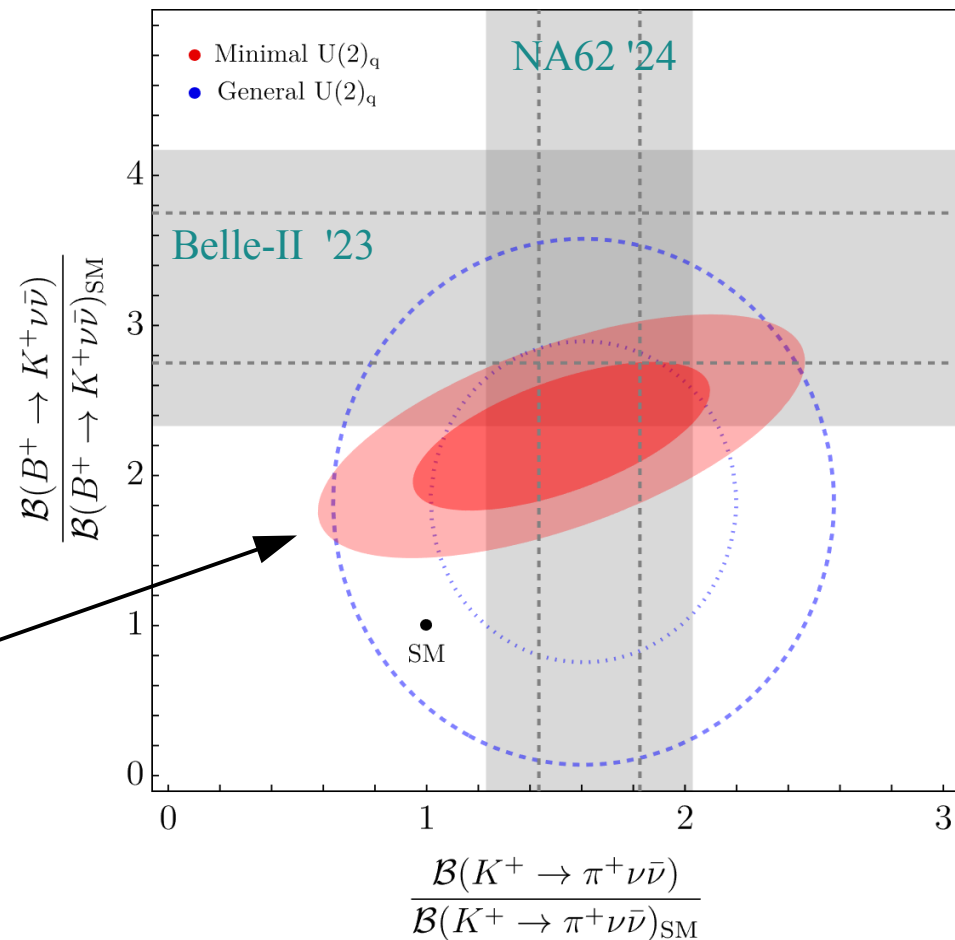
identical for all neutrino species



relevant only for 3<sup>rd</sup> gen. neutrinos

Close correlation of non-standard effects in  $B(B \rightarrow Kv\bar{v})$  &  $B(K \rightarrow \pi v\bar{v})$

Bordone *et al.* '24



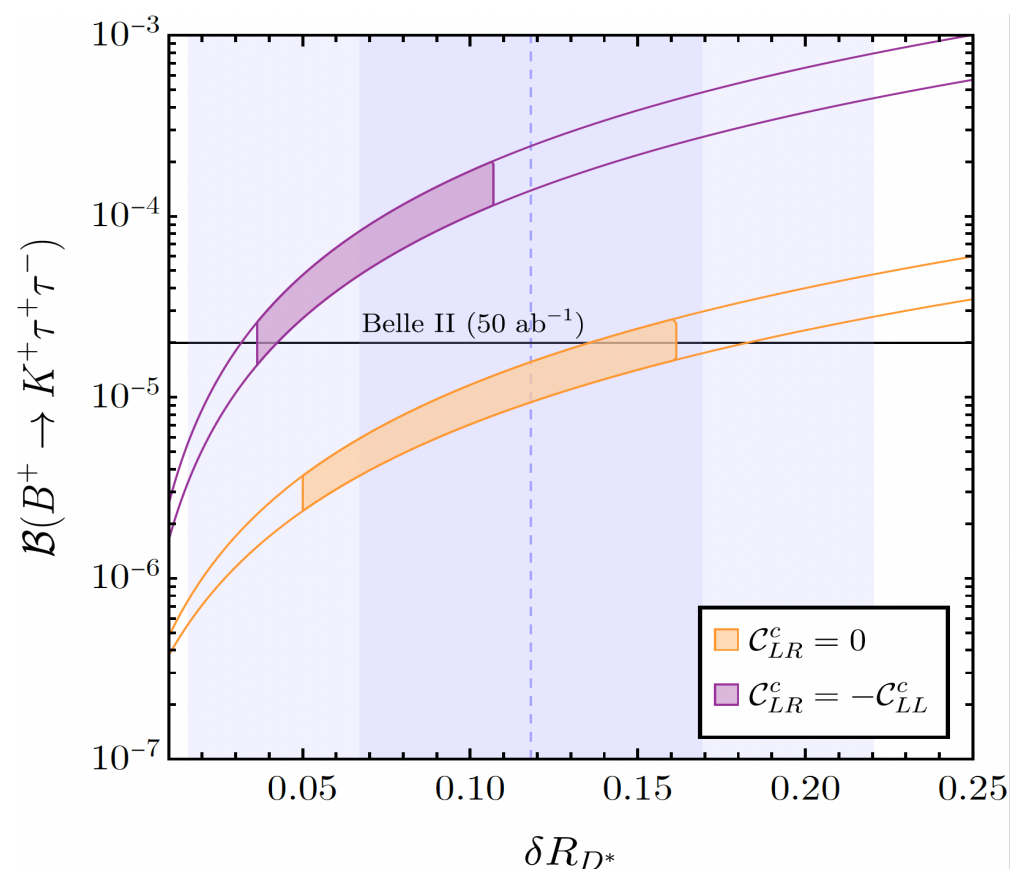
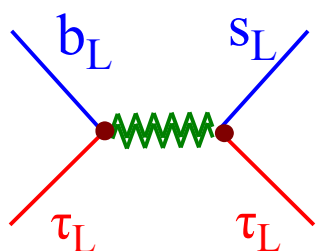
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E.g.: III) Potential large enhancement of  $b \rightarrow s \tau \tau$  rates

$b \rightarrow s \tau \tau$  are affected by the LQ already at the tree-level (contrary  $b \rightarrow s \nu \nu$ )

- huge effect compared to SM in  $b \rightarrow s \tau \tau$  (consistent with data)



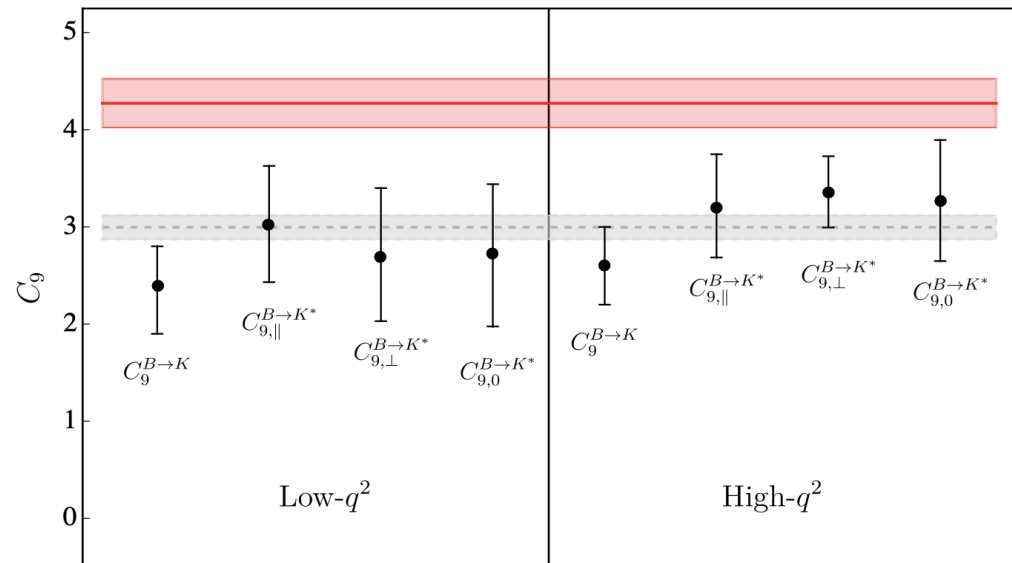
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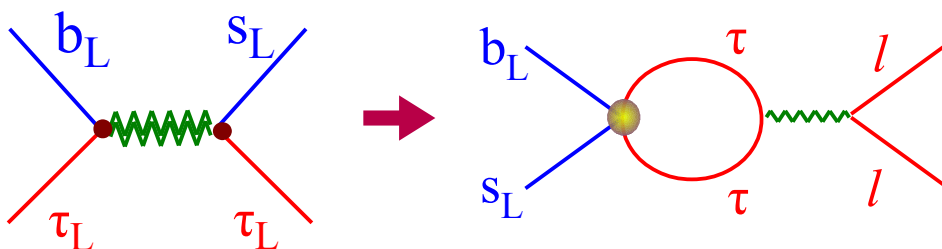
E.g.: III) Potential large enhancement of  $b \rightarrow s\tau\tau$  rates

$b \rightarrow s\tau\tau$  are affected by the LQ already at the tree-level (contrary  $b \rightarrow svv$ )

- huge effect compared to SM in  $b \rightarrow s\tau\tau$  (consistent with data)
- O(10%) effect in  $b \rightarrow sl\bar{l}$  @ 1-loop could explain persisting tension(s) in  $B \rightarrow K^{(*)}\mu\mu$  &  $B_s \rightarrow \phi\mu\mu$



Bordone, GI, Macler, Tinari'24  
GI, Tinari, Polonski, '24-'25

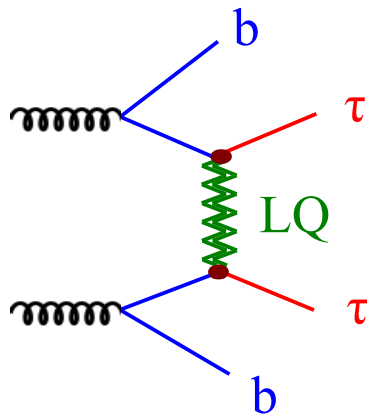


Unambiguous prediction of **magnitude**  
& **sign** consistent with excess in R(D)

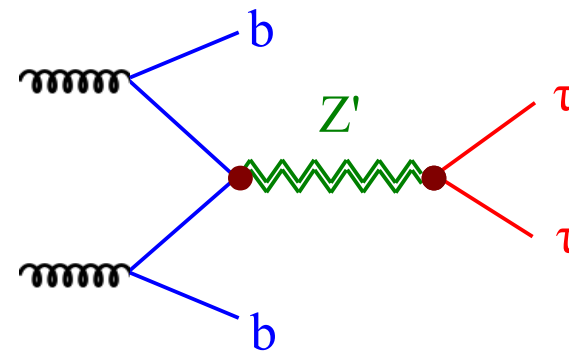
► A brief look to current data & future prospects

The idea of flavor non-universal interactions – with a 1<sup>st</sup> layer of new physics already at the TeV scale – has several interesting implications for various low-energy measurements & **high-energy measurements**

E.g.: IV)  $pp \rightarrow \tau\bar{\tau}$  (+ b-jets)



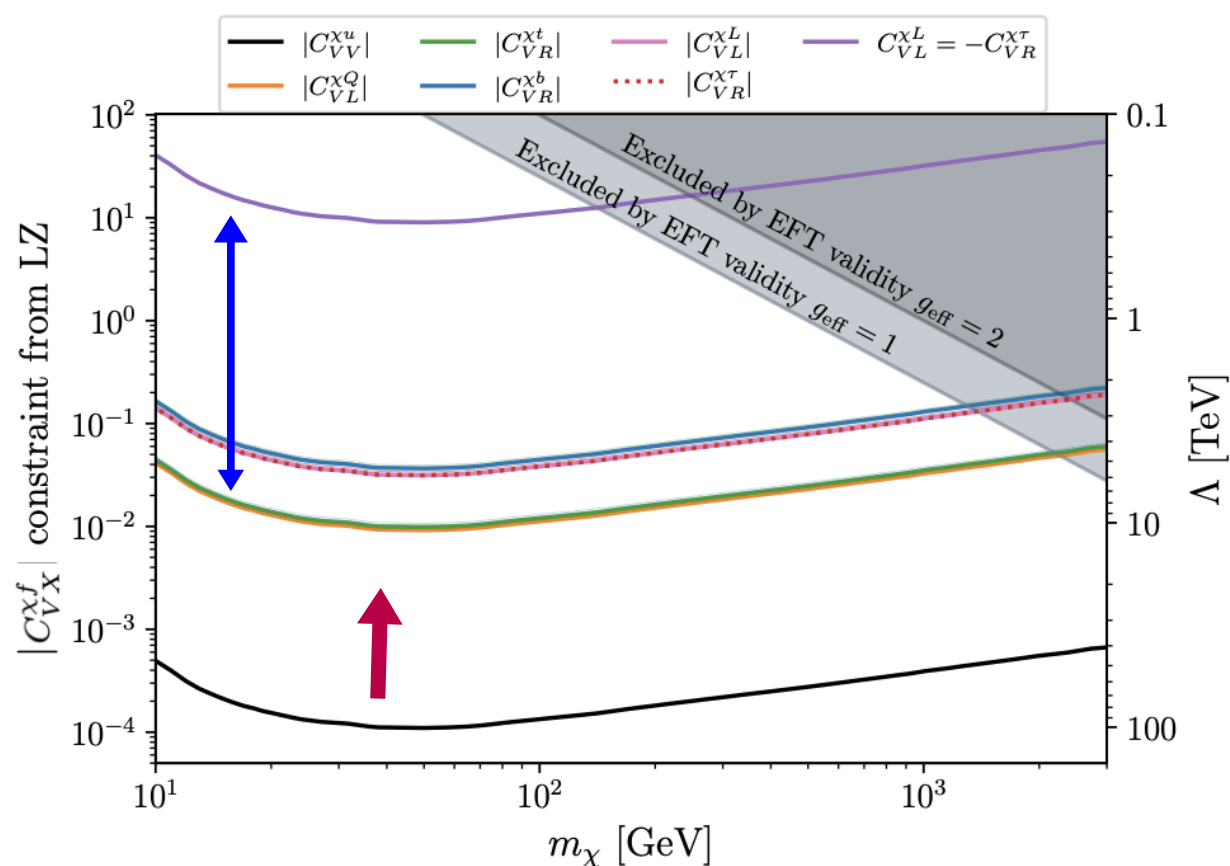
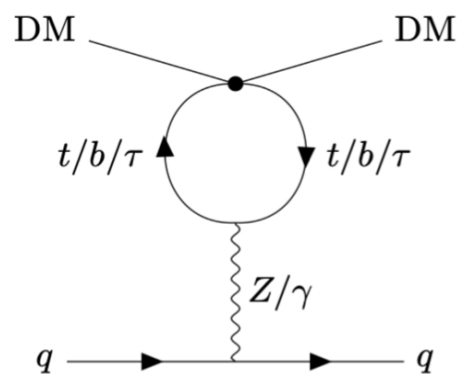
and / or



## ► A brief look to current data & future prospects

The idea of flavor non-universal interactions – with a 1<sup>st</sup> layer of new physics already at the TeV scale – has several interesting implications for various low-energy measurements & high-energy measurements, but could even have **implications for DM searches**

DM coupled mainly to 3<sup>rd</sup> gen. leads to weaker bounds from direct searches [1-loop suppression] → TeV-scale WIMP-type back in business



## Conclusions

- Flavor physics represents one the most intriguing aspects of the SM and, at the same time, a great opportunity to investigate the nature of physics beyond the SM.
- The idea of a *multi-scale construction at the origin of the flavor hierarchies* has several appealing aspects: general paradigm to consistently interpret the structure of the SMEFT, in connection with present data.
- Still a lot of work to be done on the model-building side to construct realistic models addressing both the origin of the flavor and the EW hierarchy problem...
- If these ideas corrects, new non-standard effects should emerge soon both at low and at high energies → very interesting opportunities for near-future exp. in flavor physics (**Belle-II, NA62, LHCb**) & at high energies (**LHC run-3**)